

# Transplanckian axion monodromy and the swampland



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G. Buratti, J. Calderón & A. Uranga,  
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Planck '19

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-  Compatibility of gauge field theory with quantum gravity

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## Swampland vs Landscape



Credit: F. Marchesano

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## 📌 Swampland vs Landscape



## 📌 Applications to model building?

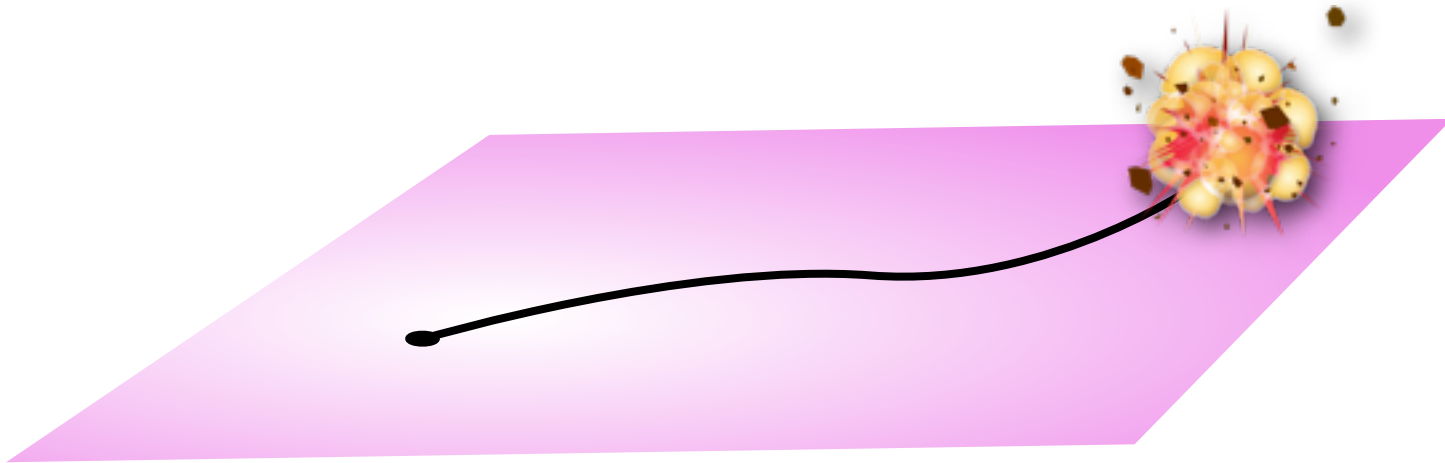
# Swampland distance conjecture

 Slogan: “Transplanckian trips in field space not allowed”

 More precise


[Ooguri, Vafa; Palti;...]

Adiabatic motion in moduli space over transplanckian (geodesic) distances leads to breakdown of effective field theory



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Adiabatic motion in moduli space over transplanckian (geodesic) distances leads to breakdown of effective field theory

- Appearance of infinite tower of light modes
- Other disasters

 Deep implications for inflation: unobservable tensor modes?

 Must deal with potential and backreaction!



# Axions



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**Axions:** Periodic scalars with (perturbative) shift symmetry

$$\phi \rightarrow \phi + \lambda$$

broken to discrete periodicity by:

- Non-perturbative effects  $\Rightarrow$  natural inflation
- Monodromic effects  $\Rightarrow$  axion monodromy

# Natural inflation

- 📌 Single axion with 1-instanton generated potential  
Freese, Frieman, Olinto

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- In deep trouble with WGC swampland constraints
  - Single field  
Arkani-Hamed, Motl, Nicolis, Vafa
  - Multifield  
Rudelius; Montero, AU, Valenzuela; ...
  - Work & proposals go on, but...

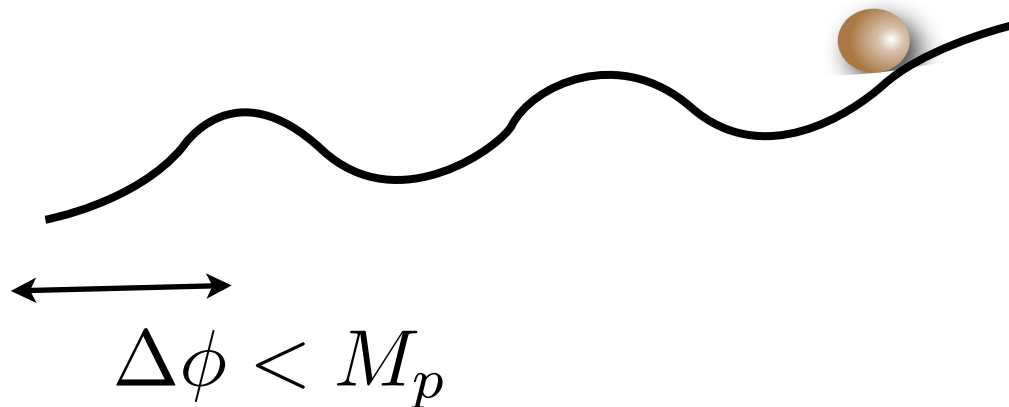
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$\Rightarrow$  higher harmonics reduce the rolling range  $< M_p$

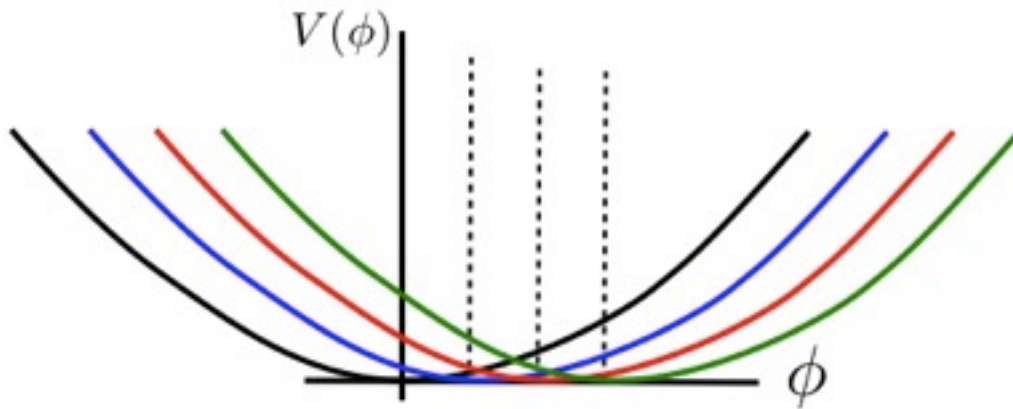




# Axion monodromy

 Alternative: Transplanckian field range with subplanckian periodicity, through multivalued potential

Silverstein, Westphal

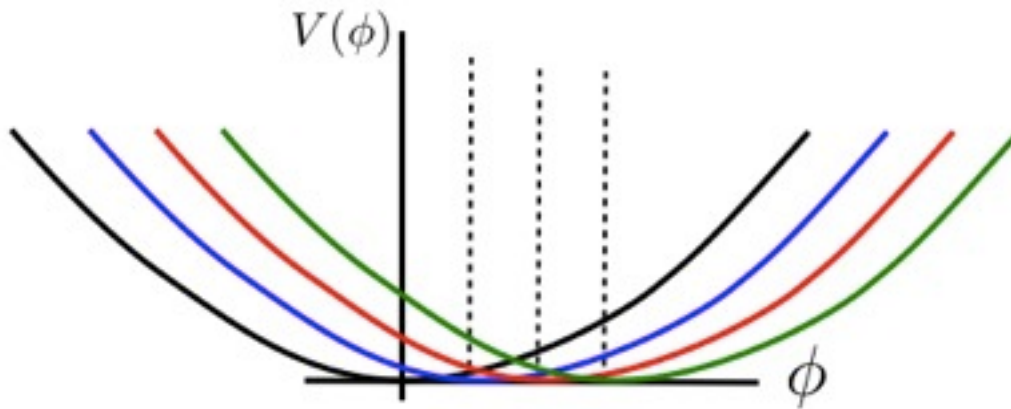


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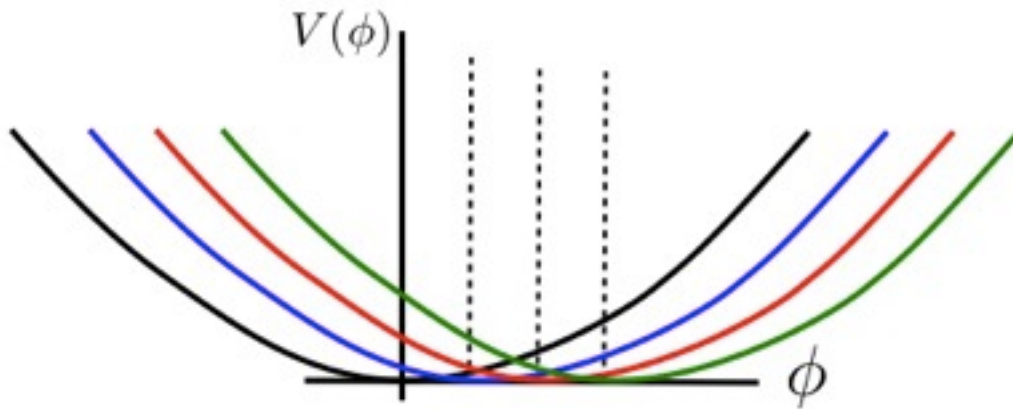
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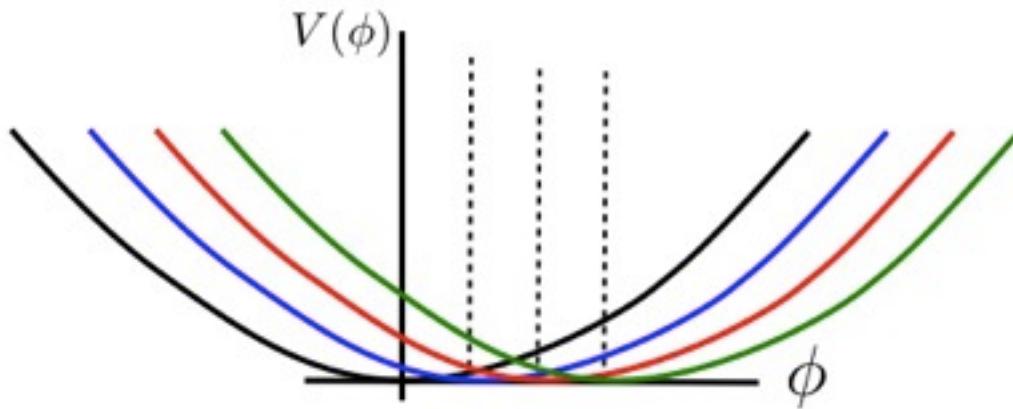
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$$|F_4|^2 + |dC_2 - nC_3|^2 \quad C_3 \rightarrow C_3 + d\Lambda_2 \quad ; \quad C_2 \rightarrow C_2 + n\Lambda_2$$

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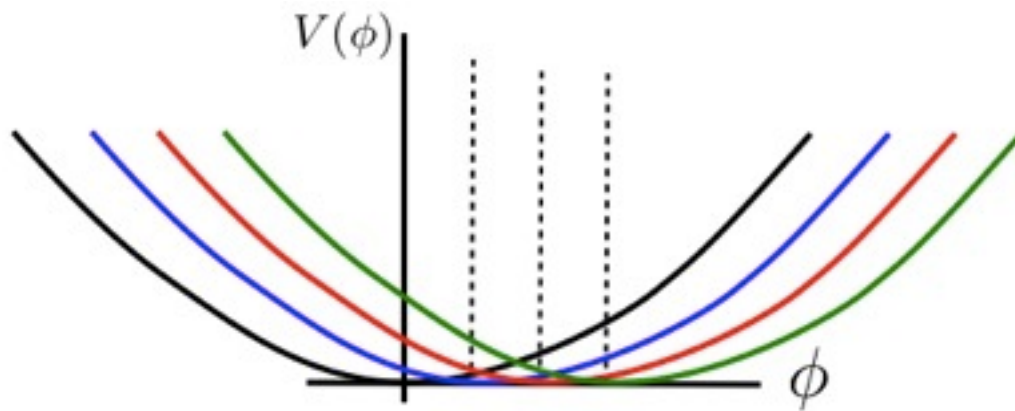
Kaloper, Sorbo+Lawrence

$$|d\phi|^2 + \phi^2$$

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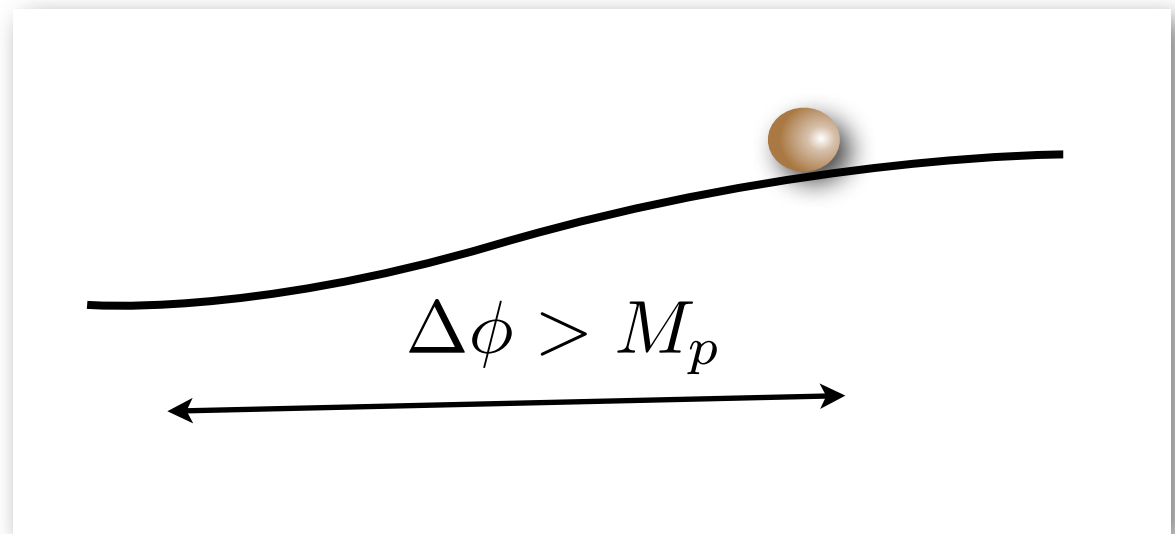
**Transplanckian vs backreaction?**

Palti et al

# Potentials and backreaction

- 📌 Adiabatic motion in field space with potential is “offshell”

Equations of motion  
don't allow for constant vev!



- 📌 In general, hard problem

What potential? What dynamics? ...?

- 📌 Address in concrete setup: fully backreacted 10d solution

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Klebanov, Tsetlyn  
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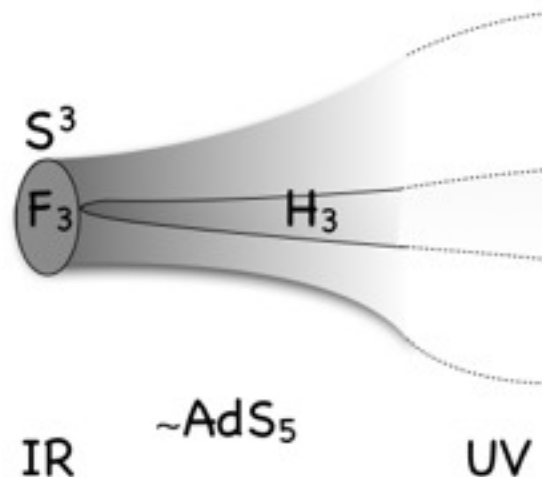
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📌 Flux compactification:

Add  $M$  units of RR 3-form flux  $F_3$  over  $S^3$

Locally  $\text{AdS}_5 \times S^2 \times S^3$ , with slow variation along  $r$



$$F_5 \text{ flux} \Rightarrow N \sim N_0 + M \log r$$

$$H_3 \Rightarrow \phi \sim M \log r$$

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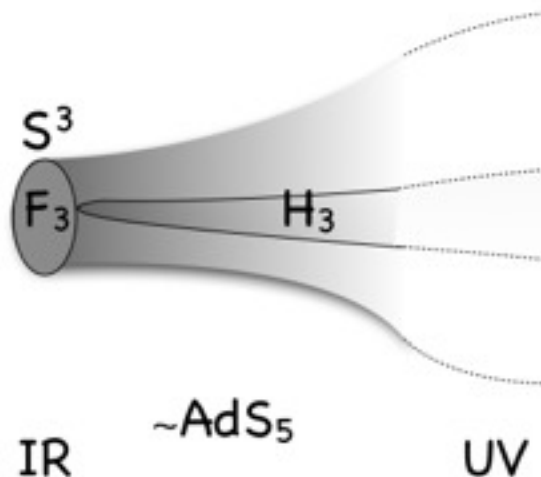
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Full 10 supergravity solution is explicitly known

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📌 “Roll” up to arbitrarily large distance in field space

$$G_{\phi\phi} \sim (N_0 + M\phi)^{-1} \quad \Delta = \int G_{\phi\phi}^{1/2} d\phi \sim \int \frac{d\phi}{\phi^{1/2}} \sim \phi^{1/2}$$

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 Comments:

Admittedly 5d, but 4d examples from M-theory on  $AdS_4 \times X_7$

Admittedly AdS, and susy, but tractable for a fundamental question

“Rolling”, but required from backreaction. In fact, essential see later

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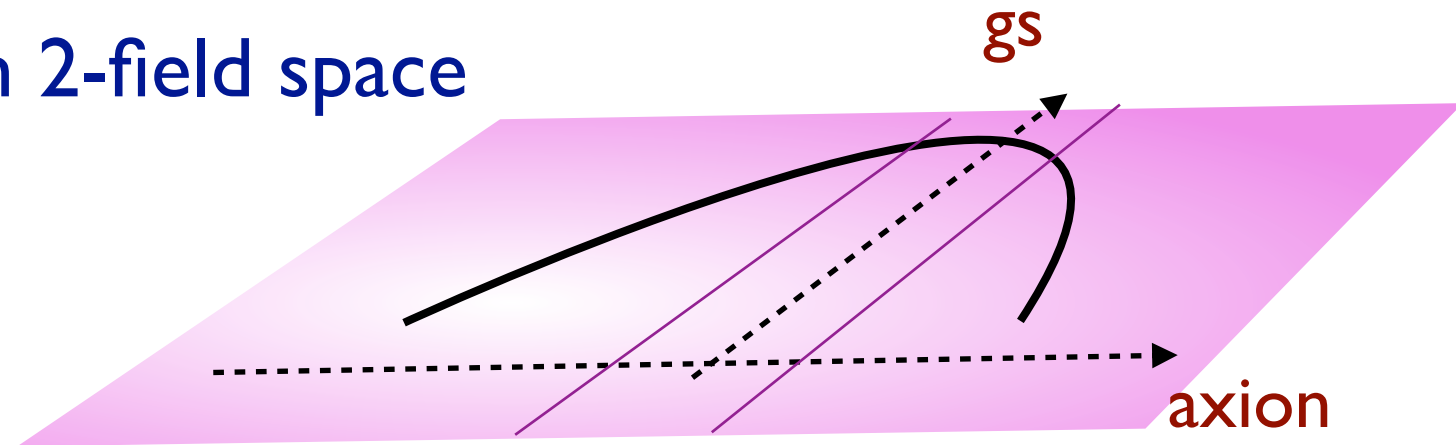
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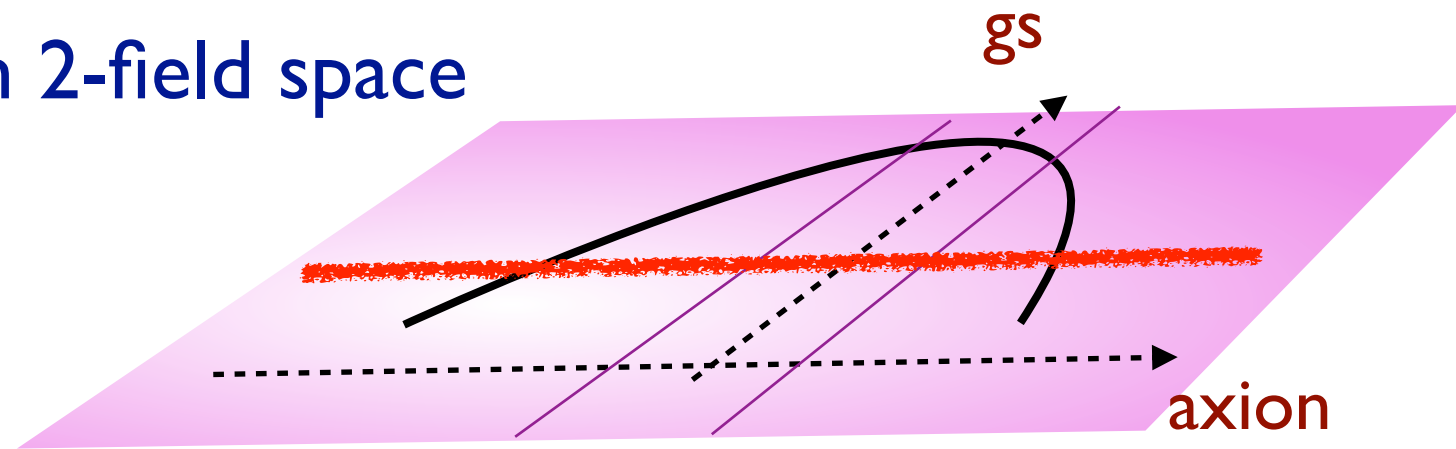
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## 📌 Geodesics in 2-field space



- 📌 Our path is not geodesic!



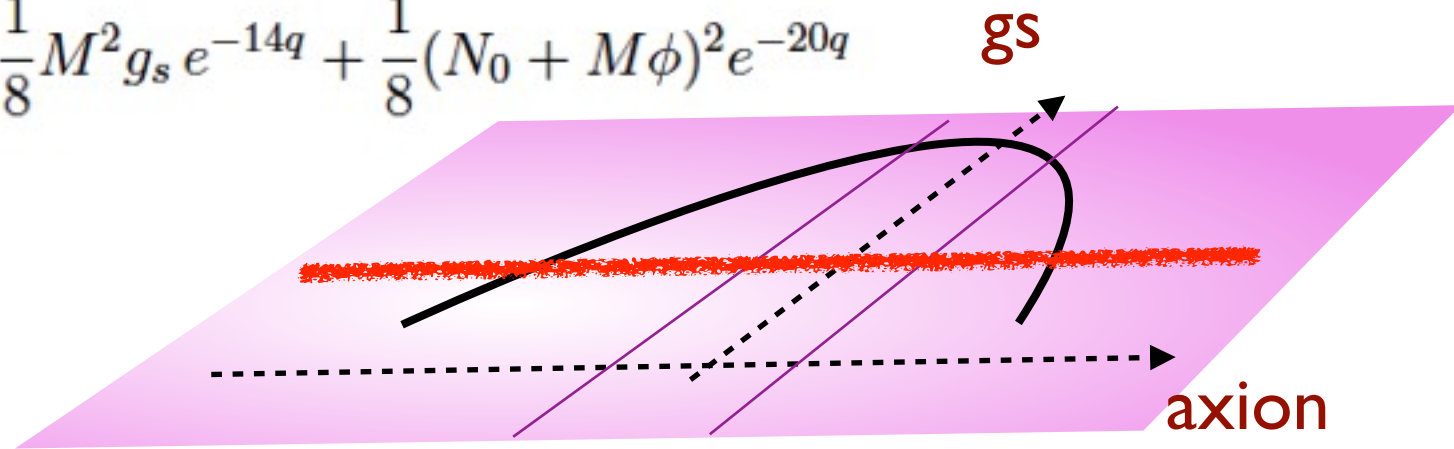
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## 📌 Constant dilaton from non-trivial balance

- Potential from flux compactification
- Effective potential from “rolling” axion

$$G_{ab}(\varphi)\partial\varphi^a\partial\varphi^b = 15(\partial q)^2 + \frac{1}{4}g_s^{-1}e^{-6q}(\partial\phi)^2,$$


$$V(\varphi) = -5e^{-8q} + \frac{1}{8}M^2g_s e^{-14q} + \frac{1}{8}(N_0 + M\phi)^2e^{-20q}$$



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 Future: Physical applications? Time dependent roll?

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# Navigating the Swampland

MADRID, 25-27 SEPTEMBER 2019



## SWAMP GONDOLIERS

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**Thank you!**