

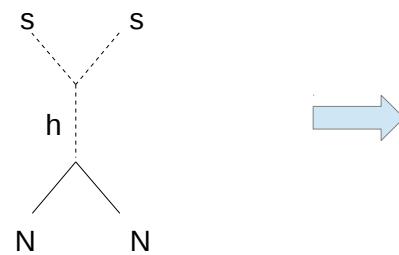
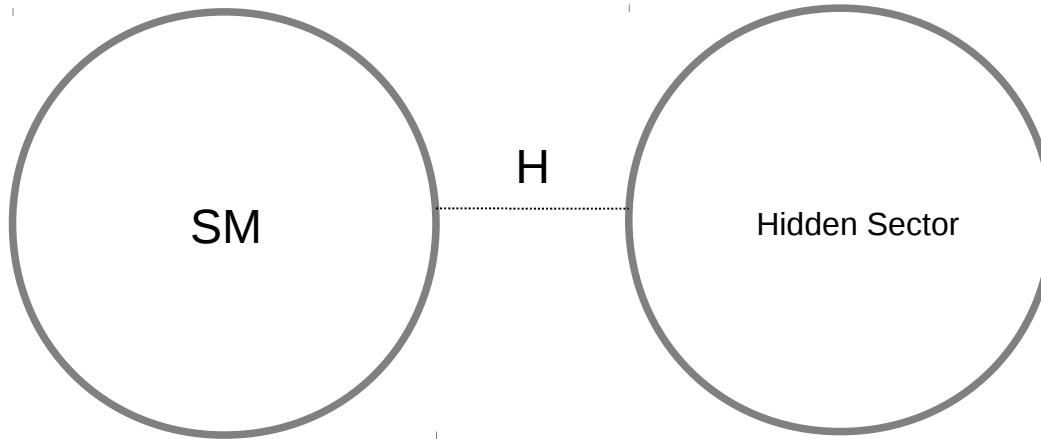
# Pseudo-Goldstone Dark Matter

through the Higgs portal

Oleg Lebedev



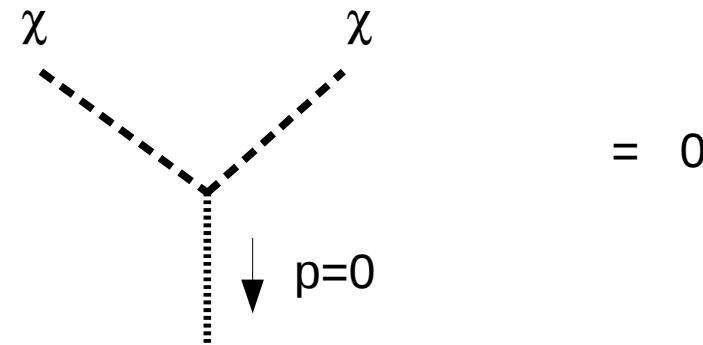
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*strong constraints on a WIMP*

## Main Goldstone feature:

derivative couplings     $\partial\chi \dots$



$$= 0$$

Perfect for  $\chi$  – dark matter:

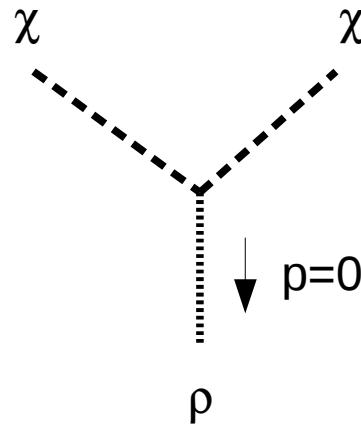
- no direct detection
- efficient annihilation ( $p \neq 0$ )

Massive "Goldstone" case:

$$L = \partial_\mu S \partial^\mu S^* - m^2 S S^* + (\mu^2 S^2 + h.c.) - V(S S^*)$$

$$S = \rho e^{i\chi}$$

( $\chi$  = *pseudo-Goldstone boson*  
with mass  $\sim \mu$ )



$$\sim \chi \rho (\square + 4 \mu^2) \chi = 0 \quad (\text{EOM})$$

**Still works!** (*NB: only for an  $S^2$  soft term*)

# Pseudo-Goldstone dark matter

**SM + a complex scalar  $S$** , require softly broken U(1) symmetry:

$$V = V_0 + V_{\text{soft}} ,$$

$$V_0 = -\frac{\mu_H^2}{2} |H|^2 - \frac{\mu_S^2}{2} |S|^2 + \frac{\lambda_H}{2} |H|^4 + \lambda_{HS} |H|^2 |S|^2 + \frac{\lambda_S}{2} |S|^4 ,$$

$$V_{\text{soft}} = -\frac{\mu'_S^2}{4} S^2 + \text{h.c.}$$

All parameters are real  $\rightarrow \langle S \rangle = \text{real} , S \rightarrow S^* \text{ symmetry}$

**Im  $S$  = Dark Matter**

States:

$$S = \frac{1}{\sqrt{2}}(v_s + s + i\chi)$$

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + h \end{pmatrix}$$

CP-even mass eigenstates =  $\mathbf{h}_1, \mathbf{h}_2$  (mixtures of  $\mathbf{s}$  and  $\mathbf{h}$ ) :

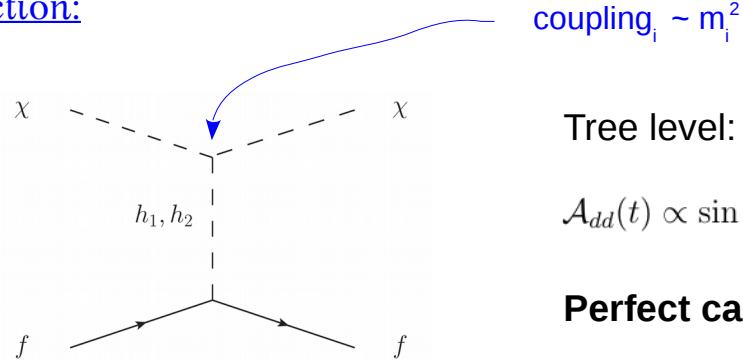
$$O^T \mathcal{M}^2 O = \begin{pmatrix} m_{h_1}^2 & 0 \\ 0 & m_{h_2}^2 \end{pmatrix} \quad O = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}$$

$$m_{h_1, h_2}^2 = \frac{1}{2} \left( \lambda_H v^2 + \lambda_S v_s^2 \mp \frac{\lambda_s v_s^2 - \lambda_H v^2}{\cos 2\theta} \right)$$

CP-odd mass eigenstate = dark matter  $\chi$  :

$$m_\chi^2 = \mu'_S^2$$

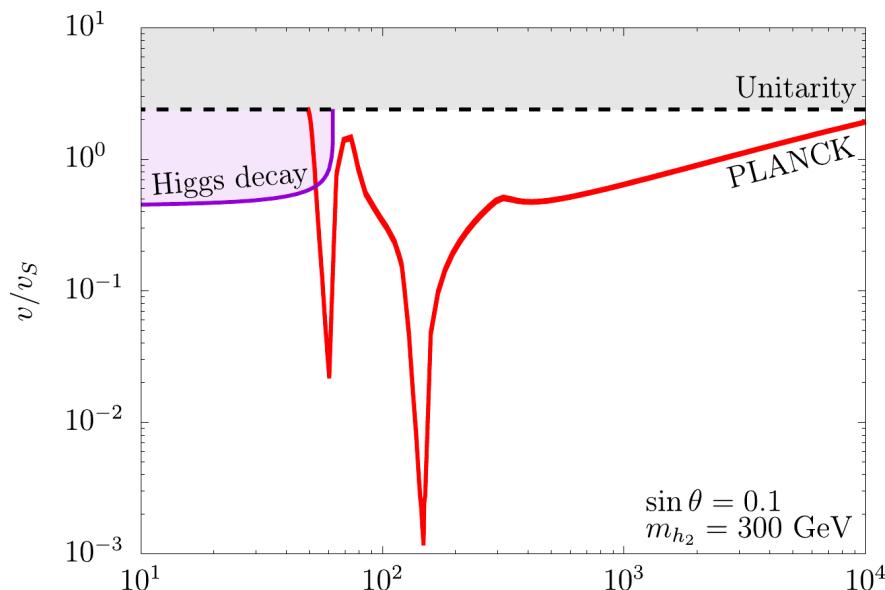
Direct detection:



Tree level:

$$\mathcal{A}_{dd}(t) \propto \sin \theta \cos \theta \left( \frac{m_2^2}{t - m_2^2} - \frac{m_1^2}{t - m_1^2} \right) \rightarrow 0$$

**Perfect cancellation for any parameter choice !**



Direct detection = loop-suppressed

Annihilation = unsuppressed



**Excellent WIMP**

(from 60 GeV to 10 TeV)

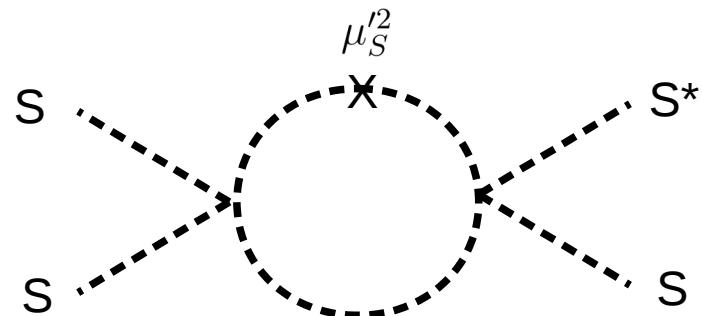
Loop corrections to the potential:

$$V_1 = \frac{\lambda'_{HS}}{2} |H|^2 S^2 + \frac{\lambda''_S}{4} |S|^2 S^2 + \frac{\lambda'_S}{4} S^4 + \text{h.c.}$$

$$\lambda'_{HS} = \frac{\lambda_{HS}\lambda_S}{32\pi^2} \ln \frac{\mu_S^2 + \mu'_S^2}{\mu_S^2 - \mu'^2_S},$$

$$\lambda''_S = \frac{\lambda_S^2}{8\pi^2} \ln \frac{\mu_S^2 + \mu'_S^2}{\mu_S^2 - \mu'^2_S},$$

$$\lambda'_S = \frac{\lambda_S^2}{64\pi^2} \left( \frac{\mu_S^2}{\mu'^2_S} \ln \frac{\mu_S^2 - \mu'^2_S}{\mu_S^2 + \mu'^2_S} + 2 \right)$$



$$\sigma \sim 10^{-49} \text{ cm}^2$$

Refined by

Azevedo, Duch, Grzadkowski, Huang, Iglicki, Santos '18

Ishiwata, Toma '18

Couplings as spurions:

$$\text{gauged } U(1) + S + \Phi : \left\{ \begin{array}{l} q_S = \text{even} \\ q_\Phi = \text{odd} \end{array} \right.$$

$$\Delta V \sim a S^2 \Phi^n + b S^4 \Phi^{2n} + \dots \quad (\text{no odd powers of } S)$$

Here

$$n \equiv -2 \frac{q_S}{q_\Phi} , \quad \epsilon \equiv \frac{\langle \Phi \rangle}{\Lambda}$$

$$\mu_S'^2 \sim \langle \Phi \rangle^2 \epsilon^{n-2} , \quad \lambda'_{HS} \sim \lambda''_S \sim \epsilon^n , \quad \lambda'_S \sim \epsilon^{2n}$$

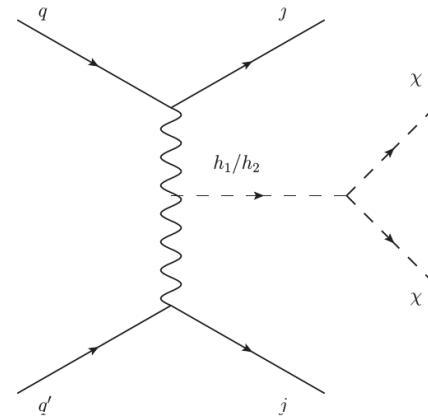

significant

small

# LHC search

**Main signature:** missing energy (via Higgs production)

VBF:

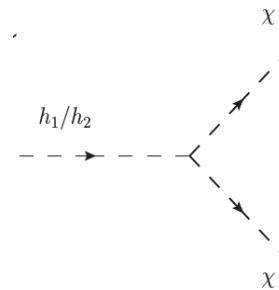


**BSM states:**  $h_2, \chi + (h_1)$

The signal is significant only for

$$m_{h_2} \text{ or } m_{h_1} > 2 m_\chi$$

invisible decay



### Constraints:

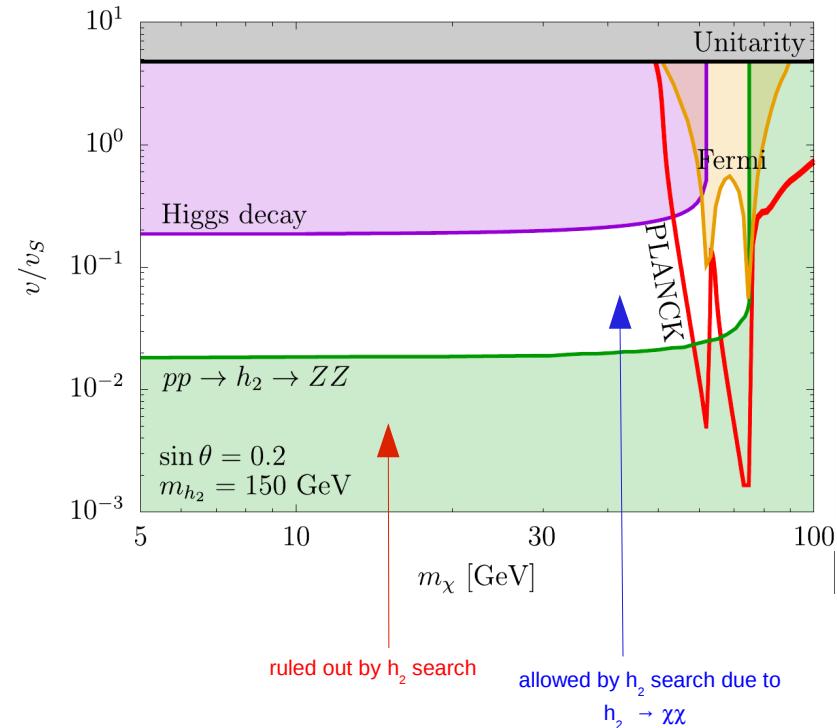
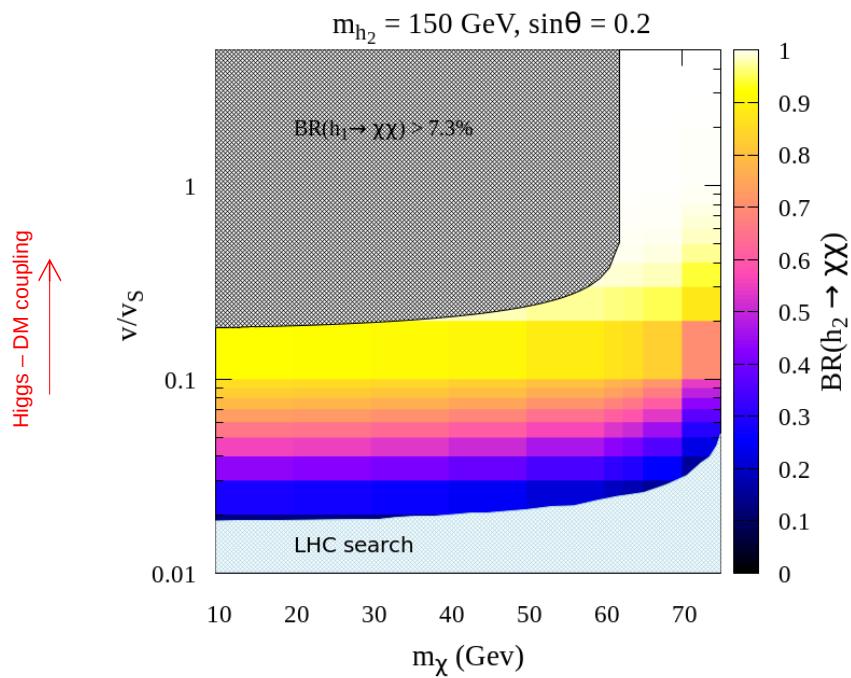
- $h_1$  Higgs signal strength  $\mu = \cos^2 \theta (1 - BR_{\text{inv}})$

→  $\theta$ -dependent bound on  $BR_{\text{inv}}$

- $h_2$  search

Example:

Use  $\mu = 1.09^{+0.11}_{-0.10}$



- light  $h_2$  allowed (e.g.  $\sim 150 \text{ GeV}$ )
- $h_2$  invisible decay very efficient

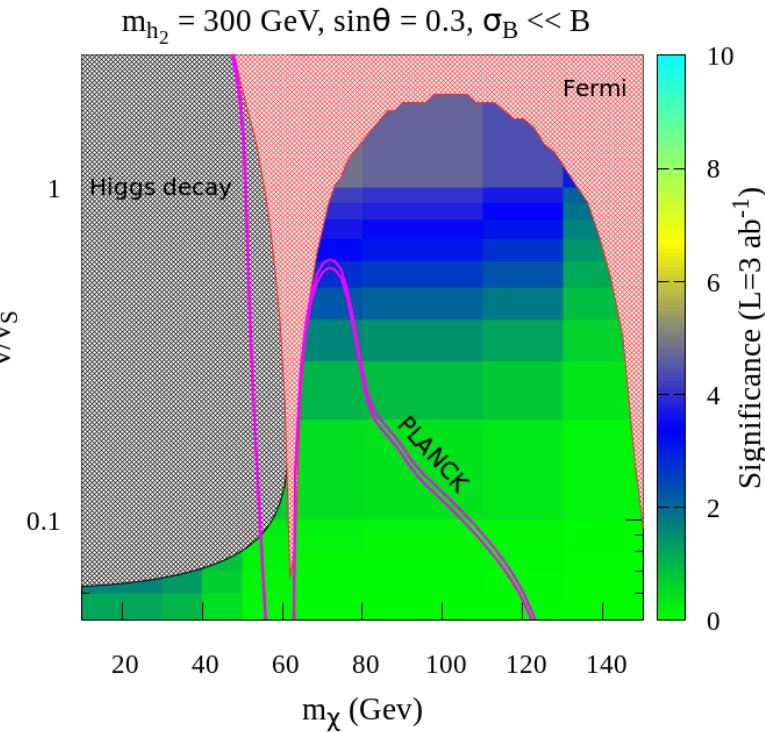
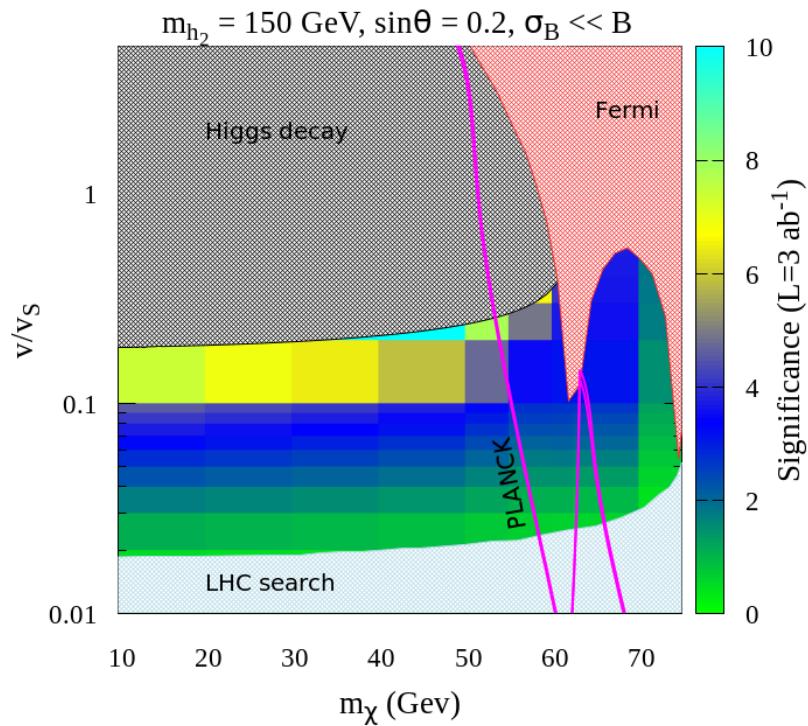
## Signal selection (CMS, 13 TeV):

Observables	Requirements
Leading (trailing) jet	$p_T > 80(40)$ GeV, $ \eta  < 4.7$
$\cancel{E}_T$	$> 250$ GeV
$\Delta\phi(\cancel{E}_T, \vec{p}_T^{\text{jet}})$	$> 0.5$
$ \Delta\phi_{jj} $	$< 1.5$
$\eta_{j1} \cdot \eta_{j2}$	$< 0$
$ \Delta\eta_{jj} $	$> 4$
$ m_{jj} $	$> 1300$ GeV
Leptons	$N_{\mu,e} = 0$ , $p_T > 10$ GeV, $ \eta  < 2.4(2.5)$
$\tau$ leptons	$N_{\tau_h} = 0$ , $p_T > 10$ GeV, $ \eta  < 2.3$
Photons	$N_\gamma = 0$ , $p_T > 15$ GeV, $ \eta  < 2.5$
$b$ -jets	$N_b = 0$ , $p_T > 20$ GeV, $ \eta  < 2.4$

Background:      Z + jets,    W + jets, ...

$$\mathcal{S} = \frac{S}{\sqrt{S+B+\sigma_B^2}}$$

Discovery potential at  $3 \text{ ab}^{-1}$ :



good prospects for non-thermal dark matter

*Further studies:*

- non–Abelian generalization      Karamitros '19
- gravitational waves                Kannike, Raidal '19

# Conclusion

- very simple DM model:  $SM + \text{complex singlet}$
- natural direct detection suppression:  $pseudo\text{--}Goldstone\ nature$
- light DM  $< O(100)$  GeV can be probed at HL-LHC in  $VBF + \cancel{E}_T$