

---

# What can(-not) multilepton searches tell us about SUSY?

---

Víctor Martín Lozano

Universität Bonn



Bethe Center for Theoretical Physics



Bethe Center for  
Theoretical Physics



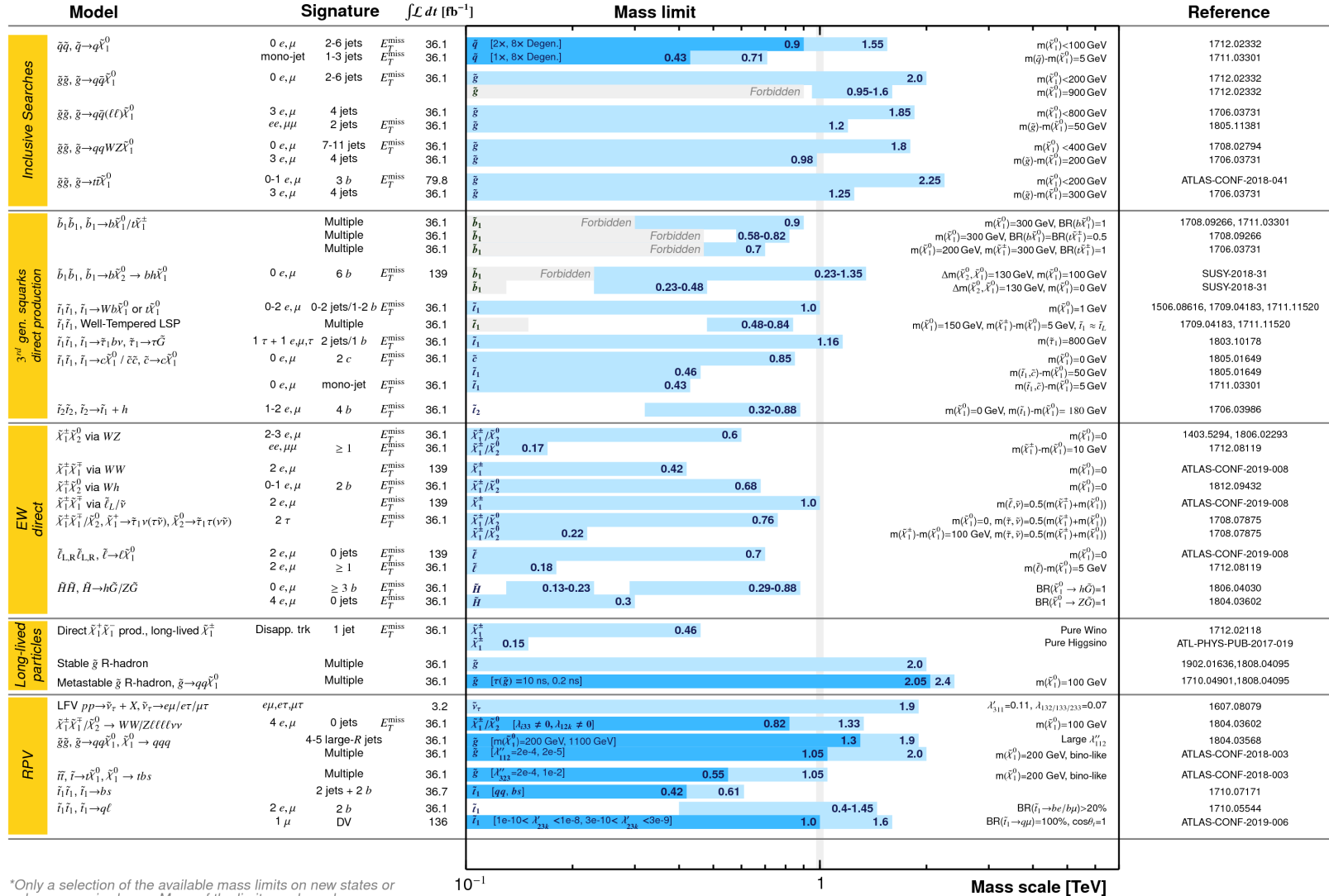
# Introduction

## ATLAS SUSY Searches\* - 95% CL Lower Limits

March 2019

ATLAS Preliminary

$\sqrt{s} = 13$  TeV



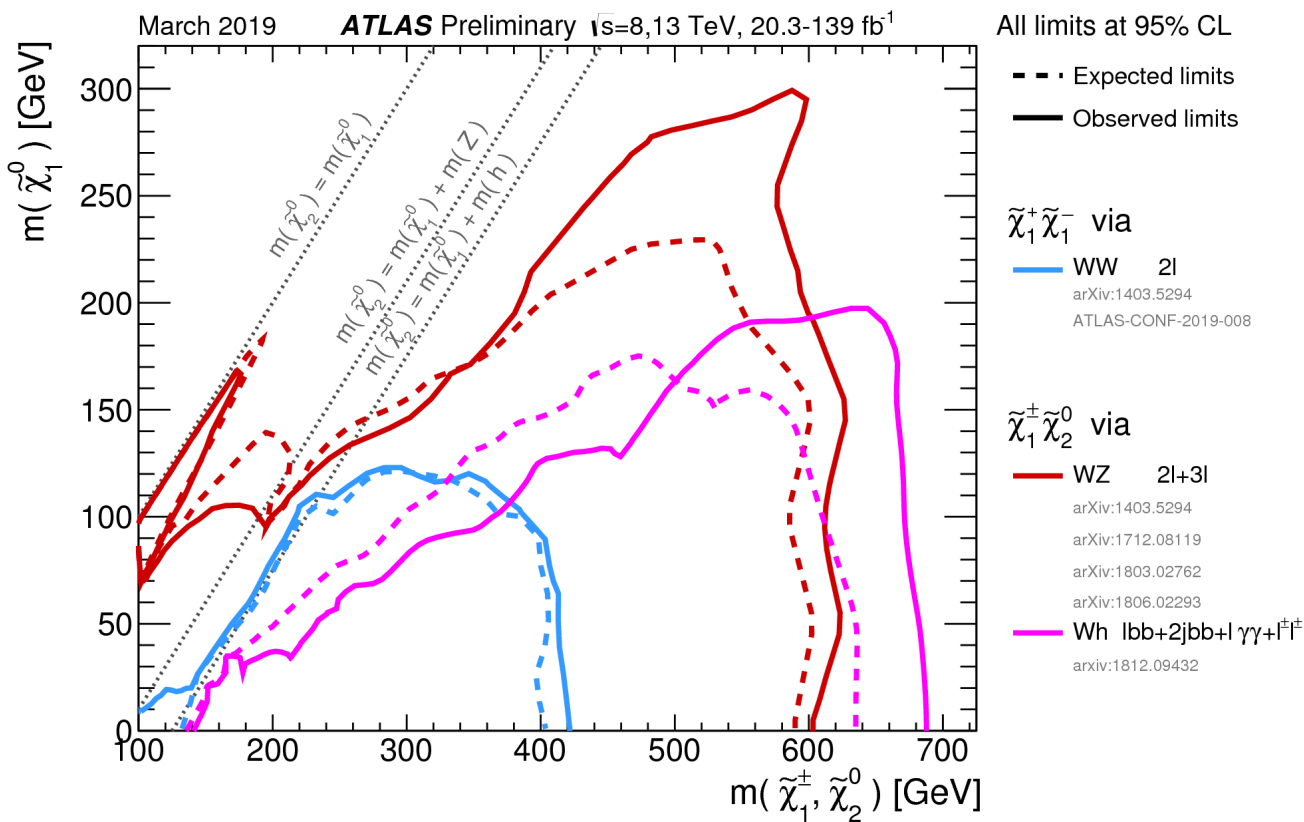
\*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10<sup>-1</sup> 1 Mass scale [TeV]

# Introduction

## Electroweakino sector:

Mediated through a vector boson:



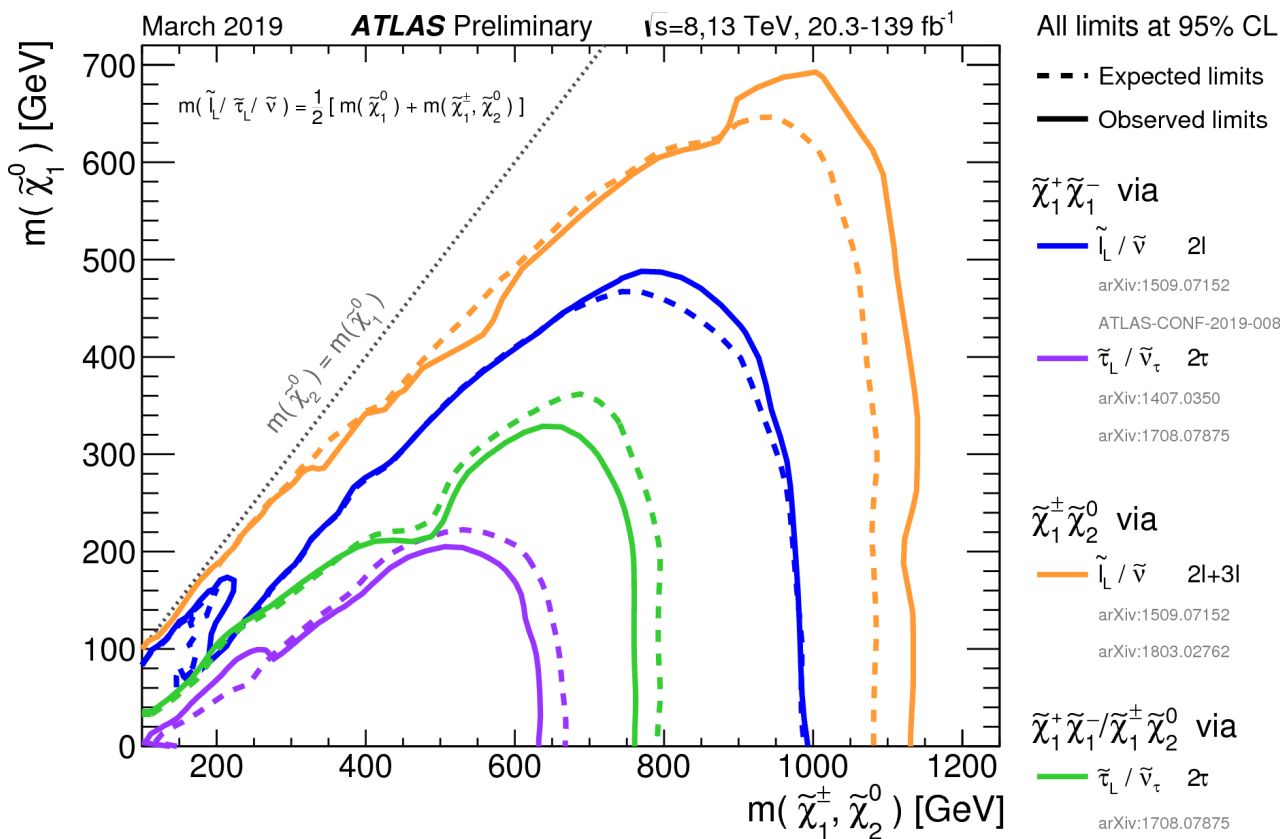
- Multilepton searches are very effective to explore the electroweakino sector

- Up to 600-700 GeV for winos and higgsinos

# Introduction

## Electroweakino sector:

Mediated through a slepton:



- Multilepton searches are also effective to explore the electroweakino sector in this case
- Up to 1.0 - 1.1 TeV for winos and higgsinos

# Introduction

---

## Electroweakino sector:

This is still true even for realistic models. Although the limits are weaker

[See for example: Arina, Chala, ML, Nardini '16]

Multilepton searches are also efficient for other SUSY models like the NMSSM or other extensions

[See for example: Domingo, Kim, ML, Ramiro, Ruiz de Austri '18]

What happens if the electroweakino sector decays differently?

- Gauge Mediation SUSY breaking: [Kim, Krauss, ML '18 ]
  - Photons
- Stealth and compressed SUSY: [Kim, Krauss, ML, Staub '18 ]
  - Soft objects (leptons and jets)
  - Photons
- R Parity Violating SUSY: [Dreiner, ML – soon ]
  - Leptons

# General Gauge Mediation Models

---

GGM models:

Messenger sector  $\rightarrow$  SUSY breaking

=====

$(\tilde{q}, \tilde{\ell}, \tilde{\nu}, \tilde{g})$

Squarks, gluino and sleptons decoupled

=====

$(\tilde{\chi}_i^0, \tilde{\chi}_j^\pm)$

Electroweakinos at EW scale:  $M_1, M_2, \mu$

NLSP decays into:  $\tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma/h/Z$

=====

$(\tilde{G})$

Bino:  $\tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma$  ( $\tilde{\chi}_1^0 \rightarrow \tilde{G}Z$ )  
 Wino:  $\tilde{\chi}_1^0 \rightarrow \tilde{G}Z$  ( $\tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma$ )  
 Higgsino:  $\tilde{\chi}_1^0 \rightarrow \tilde{G}Z/h$

$m_{\tilde{G}} = 1 \text{ eV}$ , prompt decays

# General Gauge Mediation Models

---

## GGM models:

- Squarks, gluino and sleptons decoupled
- Electroweakinos at EW scale:  $M_1$ ,  $M_2$ ,  $\mu$

Scenario	$M_1$ [GeV]	$M_2$ [GeV]	$\mu$ [GeV]	Description
<b>I</b>	[100, 1000]	[100, 1000]	2000	$\mu$ decoupled
<b>II</b>	[100, 1000]	2000	[100, 1000]	$M_2$ decoupled
<b>IIb</b>	[100, 1000]	2000	[-1000, -100]	$M_2$ decoupled
<b>III</b>	50	[100, 1000]	[100, 1000]	light bino
<b>IV</b>	2000	[100, 700]	[100, 700]	heavy bino

- $m_{\tilde{G}} = 1 \text{ eV}$ , prompt decays

-Still good DM candidate + other candidates

-Long-lived decays

[Viel, Legourges, Hachnelt, Matarrese, Riotto '05]

[Feng et al. '10, Dreiner et al. '12, ATLAS '14]

# General Gauge Mediation Models

---

## GGM models:

- Squarks, gluino and sleptons decoupled
- Electroweakinos at EW scale:  $M_1$ ,  $M_2$ ,  $\mu$

Scenario	$M_1$ [GeV]	$M_2$ [GeV]	$\mu$ [GeV]	Description
<b>I</b>	[100, 1000]	[100, 1000]	2000	$\mu$ decoupled
<b>II</b>	[100, 1000]	2000	[100, 1000]	$M_2$ decoupled
<b>IIb</b>	[100, 1000]	2000	[-1000, -100]	$M_2$ decoupled
<b>III</b>	50	[100, 1000]	[100, 1000]	light bino
<b>IV</b>	2000	[100, 700]	[100, 700]	heavy bino

- $m_{\tilde{G}} = 1 \text{ eV}$ , prompt decays

-Still good DM candidate + other candidates

-Long-lived decays

[Viel, Legourges, Hachnelt, Matarrese, Riotto '05]

[Feng et al. '10, Dreiner et al. '12, ATLAS '14]



# Recasting Experimental Searches

---

Spheno 4.0.0  
(Spectrum generator)

[Porod '03 ]

[Porod, Staub '12 ]

[Porod, Staub '17 ]



Pythia 8.219  
(Monte Carlo event generator)

[Sjöstrand et al. '15 ]



CheckMATE  
(check against experimental searches)

[Drees et al. '15 ]

[Dercks et al. '17 ]

# Recasting Experimental Searches

---

How does CheckMATE work?

For a given search, it compares the number of signal events of a given model,  $S$ , with the observed limit at 95% C.L.,  $S_{\text{exp}}^{95}$

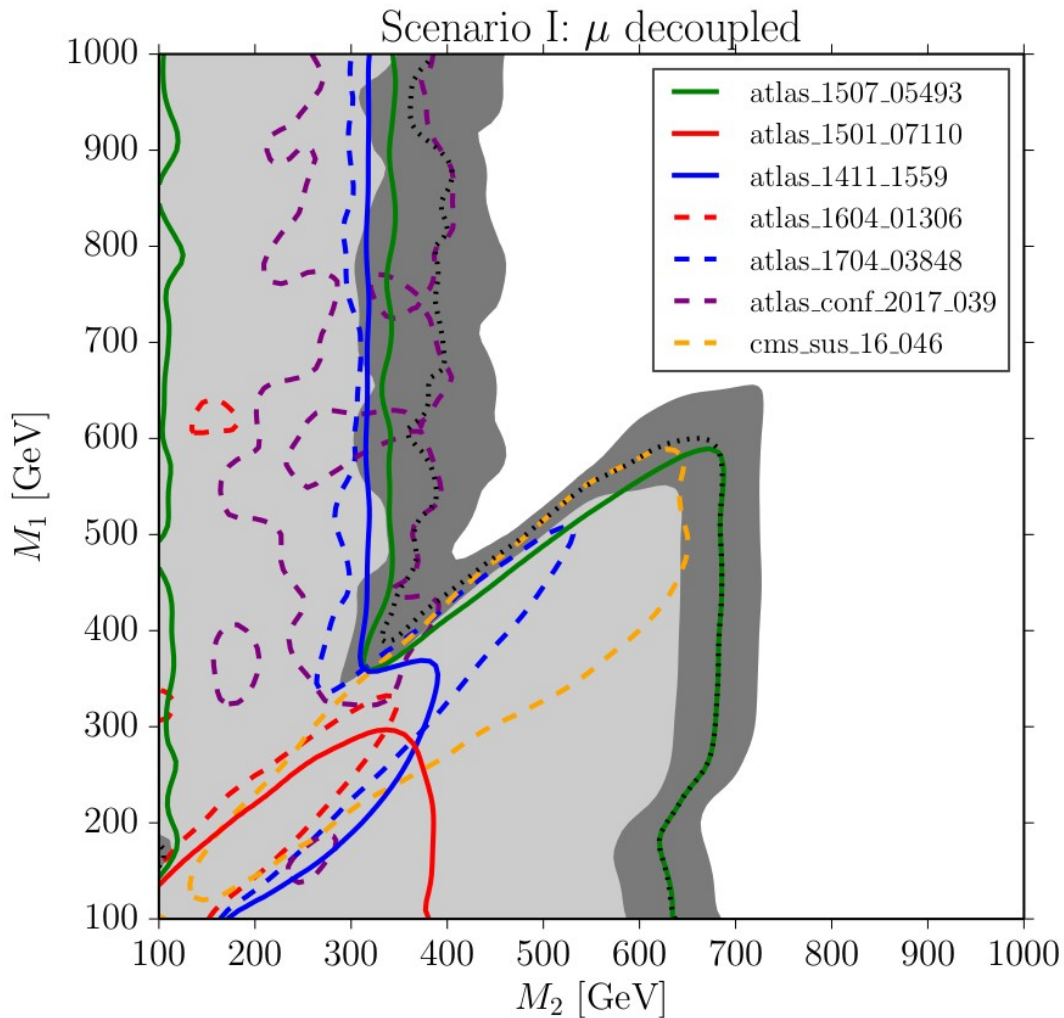
$$r = \frac{S - 1.96 \cdot \Delta S}{S_{\text{exp}}^{95}} \quad ( \Delta S = \sqrt{S} )$$

$r > 1.5$  Point Excluded

$0.67 < r < 1.5$  Point Ambiguous








$r < 0.67$  Point Allowed

# Recasting Experimental Searches



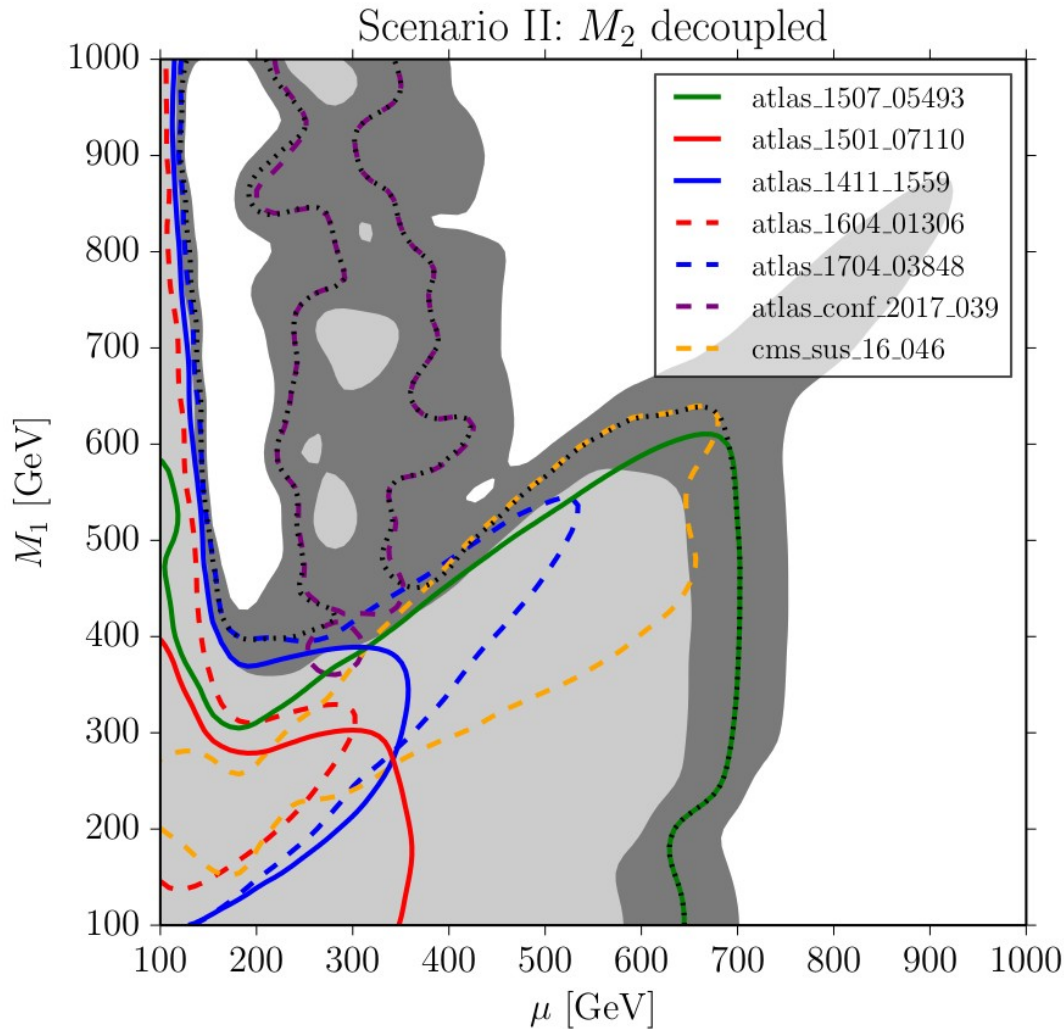
-Multilepton searches are not efficient

-Searches including photons improve the sensitivity

	$1(2)\gamma + 0(1)\ell$ + (b)-jets + $\cancel{E}_T$	(GGM)
	$h + \cancel{E}_T$	(EWK)
	$\gamma + \cancel{E}_T$	(DM)
	$\gamma + \cancel{E}_T$	(DM)
	$\gamma + \cancel{E}_T$	(DM)
	$2 - 3\ell$	(EWK)
	$\gamma + \cancel{E}_T$	(GGM)

[Kim, Krauss, ML '18]

# Recasting Experimental Searches



-Multilepton searches are efficient for only Higgsino scenarios

-Searches including photons are not efficient in only-Higgsino scenarios

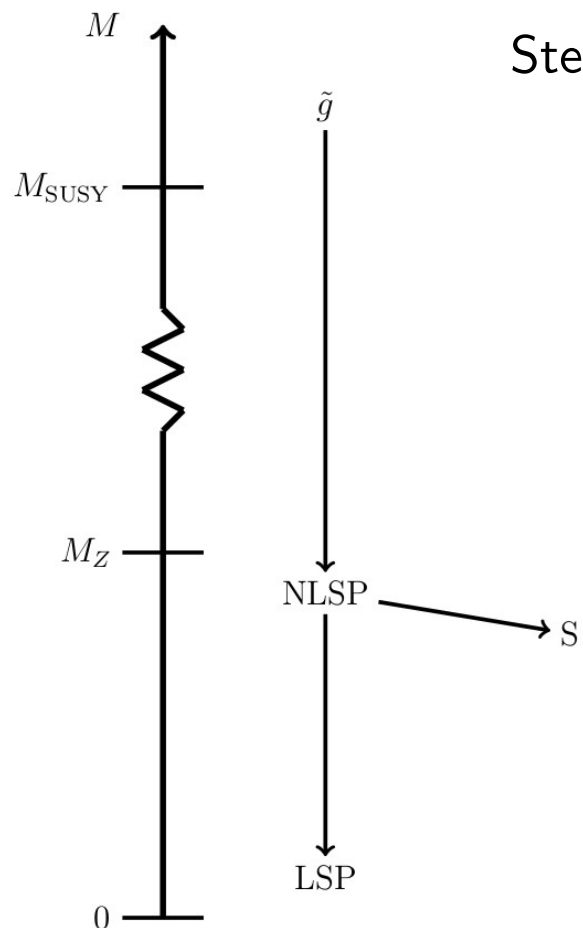
- $1(2)\gamma + 0(1)\ell$  (GGM)  
+ (b)-jets +  $\cancel{E}_T$
- $h + \cancel{E}_T$  (EWK)
- $\gamma + \cancel{E}_T$  (DM)
- $\gamma + \cancel{E}_T$  (DM)
- $\gamma + \cancel{E}_T$  (DM)
- $2 - 3\ell$  (EWK)
- $\gamma + \cancel{E}_T$  (GGM)

[Kim, Krauss, ML '18]

# Stealth SUSY

## Stealth SUSY models:

NLSP and S close in mass  $\rightarrow$  LSP small momentum ( $\cancel{E}_T$ )



Stealth SUSY spectrum hierarchy:

$$m_{\tilde{g}} \gg m_h, m_Z > m_{\tilde{\chi}_2^0} > m_S \gg m_{\tilde{\chi}_1^0}$$

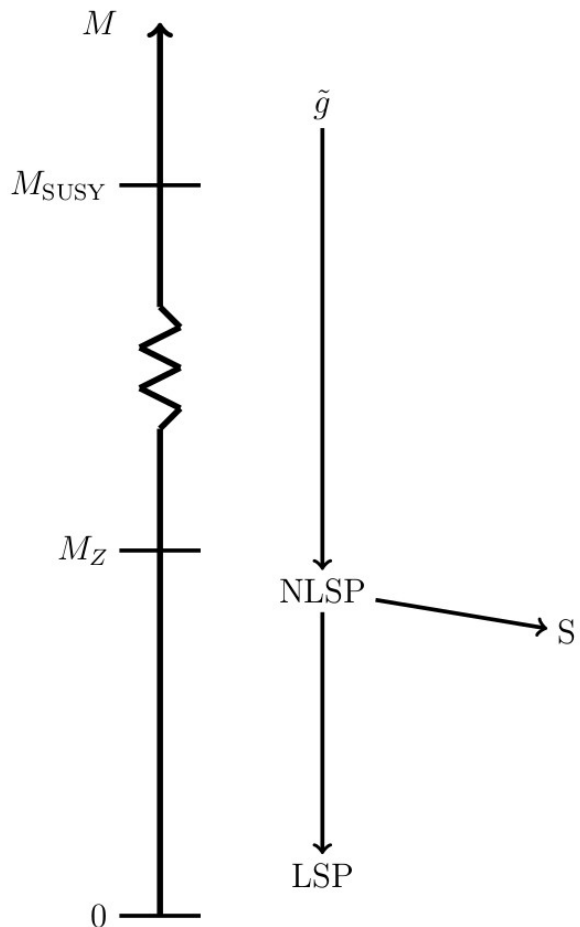
$$m_{\tilde{\chi}_2^0} \simeq m_S + m_{\tilde{\chi}_1^0} + (0.5 - 1) \text{ GeV}$$

Particle	Mass [GeV]
$\tilde{g}$	1100 – 2000
$\tilde{\chi}_2^0$	89
$\tilde{\chi}_1^0$	5
S	83

# Stealth SUSY

## Stealth SUSY models:

NLSP and S close in mass  $\rightarrow$  LSP small momentum ( $\cancel{E}_T$ )



$$pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_{3,4}^0 \rightarrow W^\pm \chi_2^0 Z/h \chi_2^0$$

$$pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_1^\mp \rightarrow W^\pm \chi_2^0 W^\pm \chi_2^0$$

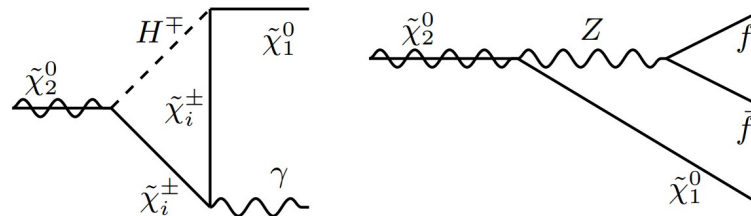
$$pp \rightarrow \tilde{\chi}_{3,4}^0 \tilde{\chi}_{3,4}^0 \rightarrow Z/h \chi_2^0 Z/h \chi_2^0$$

NLSP decays into:

$$\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 S$$

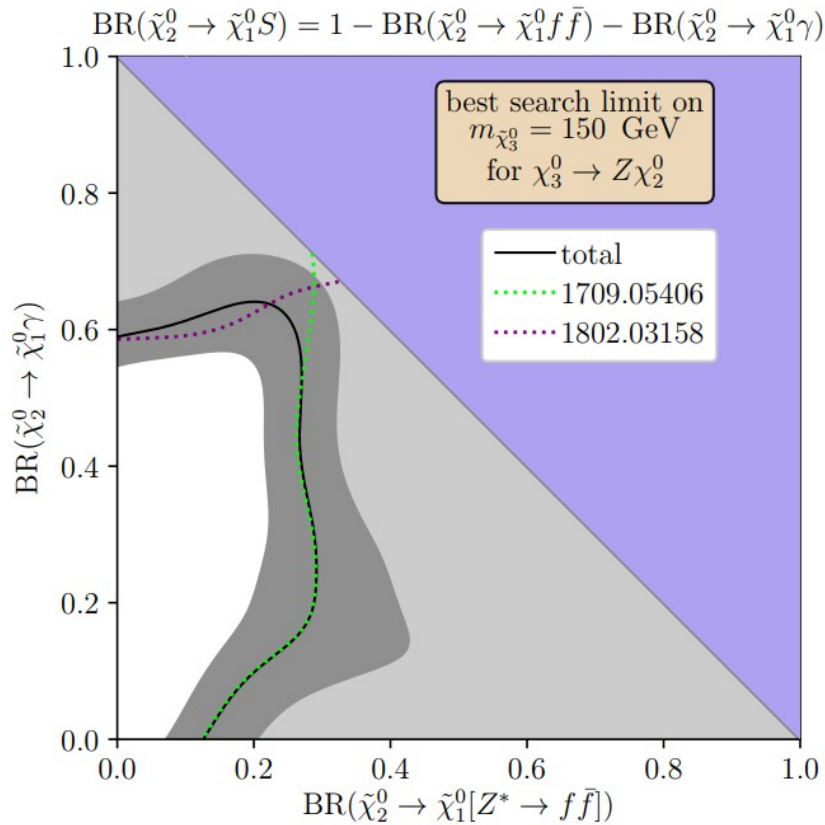
$$\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma$$

$$\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 Z^* \rightarrow \tilde{\chi}_1^0 f \bar{f}$$

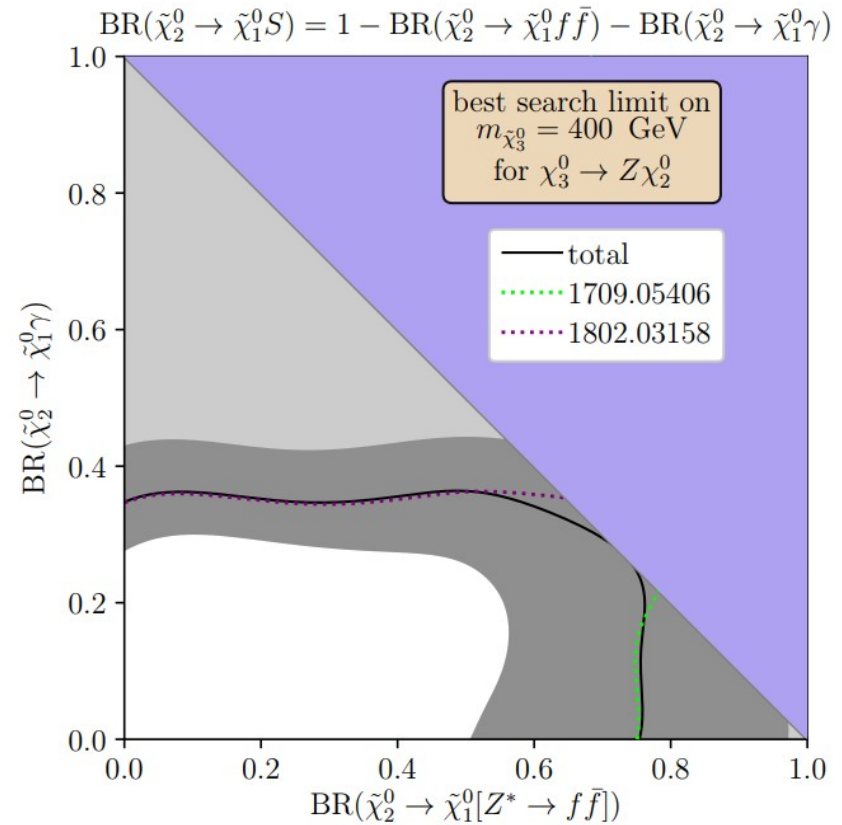


# Recasting Experimental Searches

$$\tilde{\chi}_3^0 \rightarrow Z \tilde{\chi}_2^0$$



- - - 2 - 3  $\ell$  (EWK)

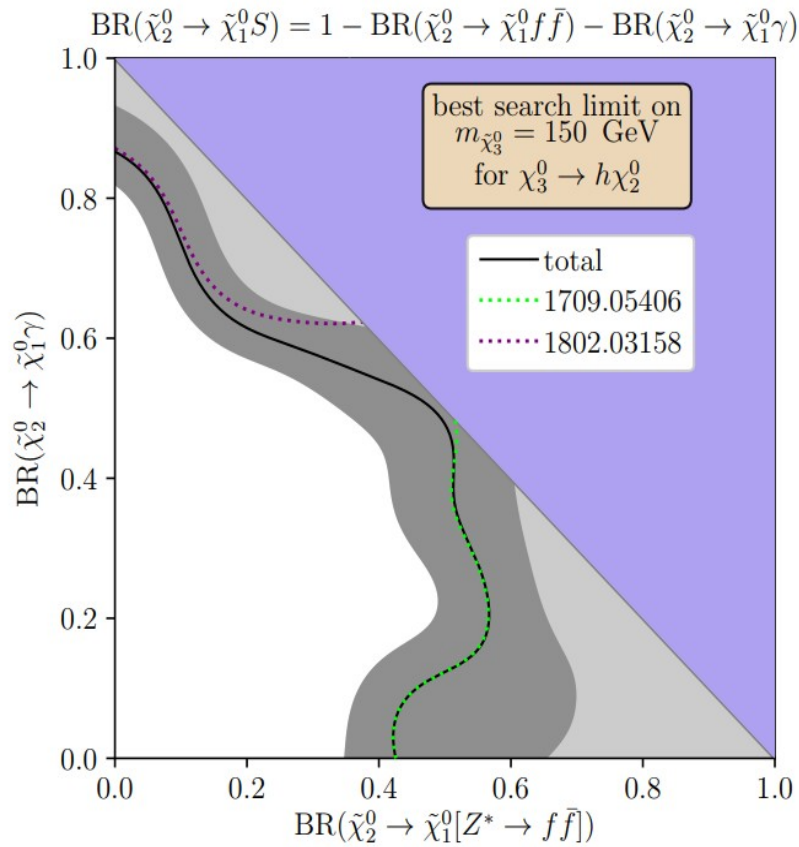


- - -  $\gamma + \cancel{E}_T$  (GMSB)

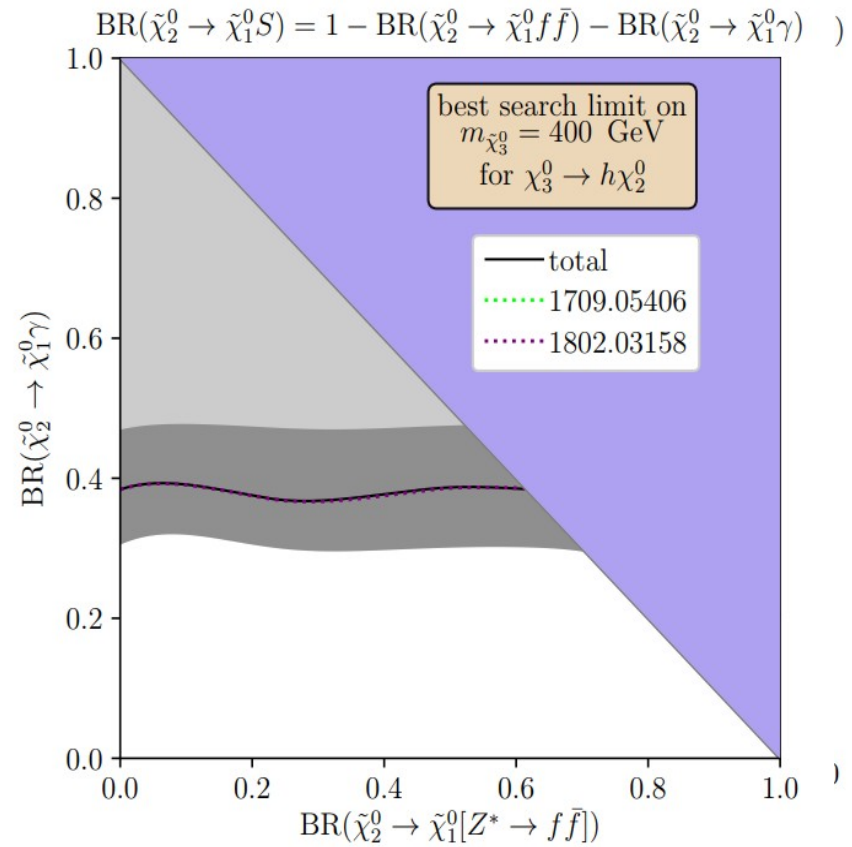
[Kim, Krauss, ML, Staub '18]

# Recasting Experimental Searches

$$\tilde{\chi}_3^0 \rightarrow h\tilde{\chi}_2^0$$



..... 2 - 3  $\ell$  (EWK)



.....  $\gamma + \cancel{E}_T$  (GMSB)

[Kim, Krauss, ML, Staub '18]



# RPV SUSY

---

RPV SUSY:

$$W = \epsilon_{ab} \left[ \frac{1}{2} \lambda_{ijk} L_i^a L_j^b \bar{E}_k + \lambda'_{ijk} L_i^a Q_j^b \bar{D}_k - \kappa_i L_i^a H_u^b \right] \\ + \frac{1}{2} \epsilon_{xyz} \lambda''_{ijk} \bar{U}_i^x \bar{D}_j^y \bar{D}_k^z .$$

R-Parity is no longer a symmetry  $\longrightarrow$  LSP can decay

The different terms violate Lepton and Barion number  $\longrightarrow$  Proton decay

New Phenomenology!!

# RPV SUSY

---

RPV SUSY:

$$W = \epsilon_{ab} \left[ \frac{1}{2} \lambda_{ijk} L_i^a L_j^b \bar{E}_k + \lambda'_{ijk} L_i^a Q_j^b \bar{D}_k - \kappa_i L_i^a H_u^b \right] + \frac{1}{2} \epsilon_{xyz} \lambda''_{ijk} \bar{U}_i^x \bar{D}_j^y \bar{D}_k^z.$$

R-Parity is no longer a symmetry  $\longrightarrow$  LSP can decay

The different terms violate Lepton and Barion number  $\longrightarrow$  Proton decay

New Phenomenology!!  $pp \rightarrow \tilde{\tau}_1 \tilde{\tau}_1$

If stau is the LSP:  $\tilde{\tau}_1 \rightarrow e\nu_\mu$       Leptons + MET!!

$\tilde{\tau}_1 \rightarrow \mu\nu_e$

# RPV SUSY

Stau LSP:

$$\begin{pmatrix} \tilde{\tau}_1 \\ \tilde{\tau}_2 \end{pmatrix} = \begin{pmatrix} \cos \theta_{\tilde{\tau}} & \sin \theta_{\tilde{\tau}} \\ -\sin \theta_{\tilde{\tau}} & \cos \theta_{\tilde{\tau}} \end{pmatrix} \begin{pmatrix} \tilde{\tau}_L \\ \tilde{\tau}_R \end{pmatrix}$$

If  $\theta_{\tilde{\tau}} = 0$  then  $\tilde{\tau}_1 = \tilde{\tau}_L$ , if  $\theta_{\tilde{\tau}} = \frac{\pi}{2}$  then  $\tilde{\tau}_1 = \tilde{\tau}_R$

Benchmark Models:

$$\lambda_{ijk} L_i^a L_j^b \bar{E}_k$$

- Model I:**  $L_a L_3 \bar{E}_c$ ,  $\tilde{\tau}_1^+ \rightarrow \ell_c^+ \nu_a$ ,  
**Model II:**  $L_1 L_2 \bar{E}_3$ ,  $\tilde{\tau}_1^+ \rightarrow (e^+ \bar{\nu}_\mu, \mu^+ \bar{\nu}_e)$ ,  
**Model III:**  $L_a L_3 \bar{E}_3$ ,  $\tilde{\tau}_1^+ \rightarrow (\tau^+ \nu_a, \tau^+ \bar{\nu}_a, \ell_a^+ \nu_\tau)$

Model	Coupling	$\tilde{\tau}_1$ -Decays	Signatures
Ia	$\lambda_{a31}$	$e \nu_a$	$e^+ e^- + \cancel{E}_T$
Ib	$\lambda_{a32}$	$\mu \nu_a$	$\mu^+ \mu^- + \cancel{E}_T$
II	$\lambda_{123}$	$\mu \nu_e, e \nu_\mu$	$e^+ e^- + \cancel{E}_T$ $\mu^+ \mu^- + \cancel{E}_T$ $e^\pm \mu^\mp + \cancel{E}_T$
IIIa	$\lambda_{133}$	$e \nu_\tau, \tau \nu_e$	$e^+ e^- + \cancel{E}_T$ $\tau^+ \tau^- + \cancel{E}_T$ $e^\pm \tau^\mp + \cancel{E}_T$
IIIb	$\lambda_{233}$	$\mu \nu_\tau, \tau \nu_\mu$	$\mu^+ \mu^- + \cancel{E}_T$ $\tau^+ \tau^- + \cancel{E}_T$ $\mu^\pm \tau^\mp + \cancel{E}_T$

# Versatility of multilepton searches

Model Ia:  $\lambda_{a31} L_a L_3 \bar{E}_1$

$$\Gamma(\tilde{\tau}_1^+ \rightarrow e^+ \nu_a) = |\lambda_{a31}|^2 \cos^2 \theta_{\tilde{\tau}} \frac{(M_{\tilde{\tau}_1}^2 - m_e^2)^2}{16\pi M_{\tilde{\tau}_1}^3}$$

2-3 lepton analysis is really powerful

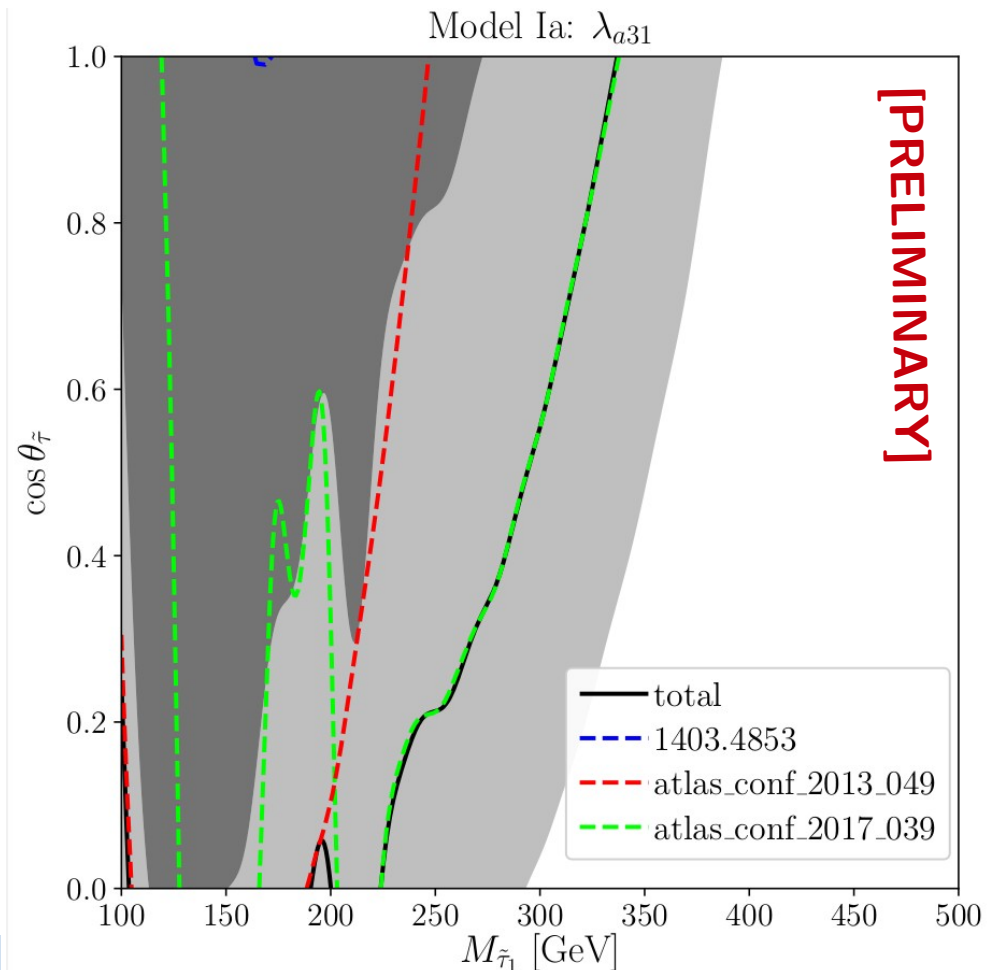
$$m_{\tilde{\tau}_1} < 225 \text{ GeV}, \quad \cos \theta_{\tilde{\tau}} \rightarrow 0$$

$$m_{\tilde{\tau}_1} < 325 \text{ GeV}, \quad \cos \theta_{\tilde{\tau}} \rightarrow 1$$

2 lepton analysis from 8 TeV is useful  
to constrain lower masses

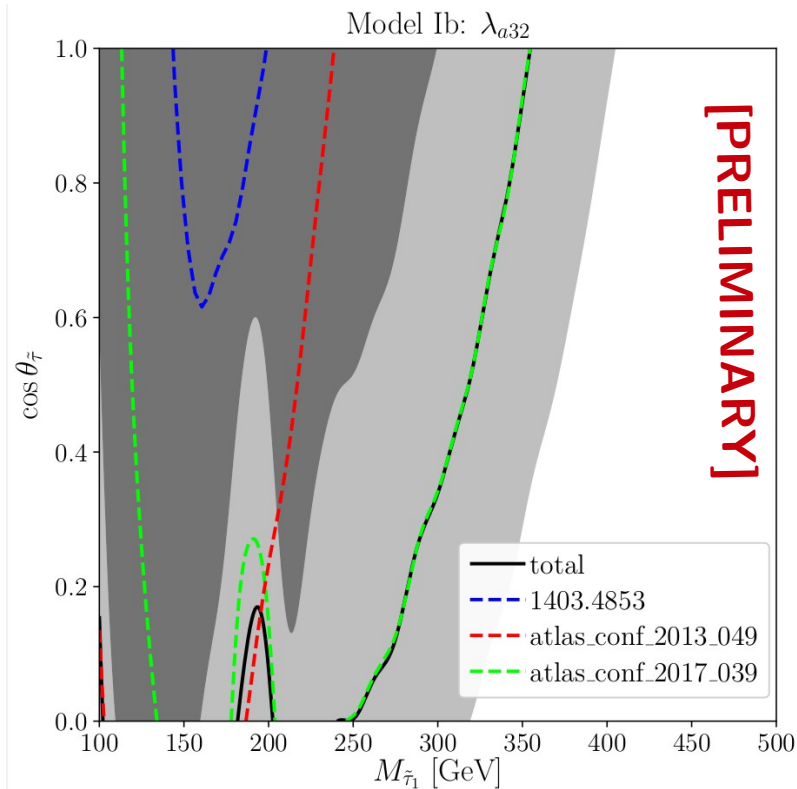
Lower than 100 GeV is still free

[Dreiner, ML – soon]



# Versatility of multilepton searches

Model Ib:  $\lambda_{a32} L_a L_3 \bar{E}_2$

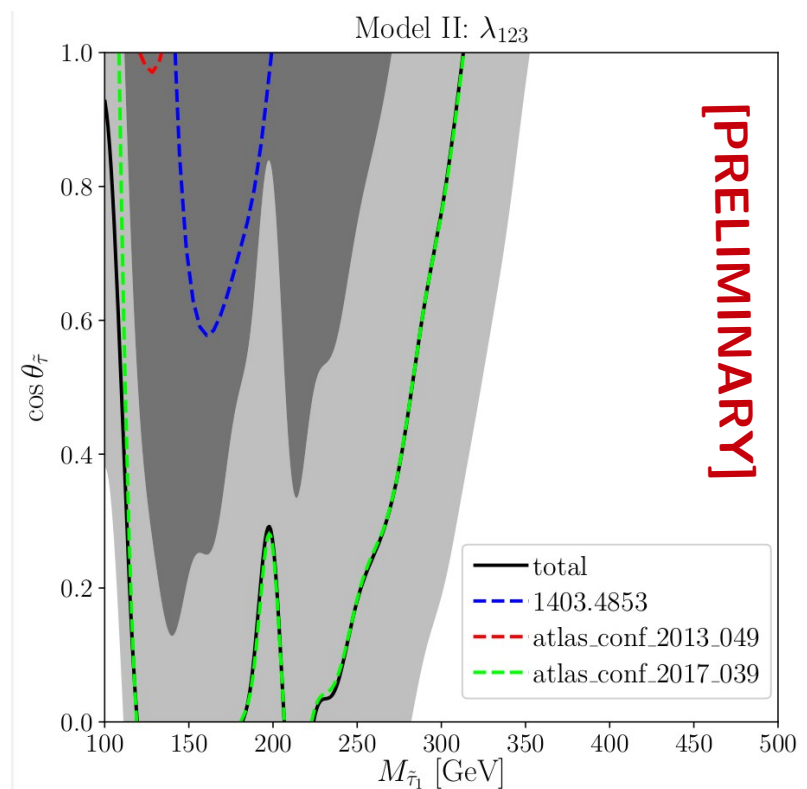


$$m_{\tilde{\tau}_1} < 250\text{GeV}, \quad \cos \theta_{\tilde{\tau}} \rightarrow 0$$

$$m_{\tilde{\tau}_1} < 340\text{GeV}, \quad \cos \theta_{\tilde{\tau}} \rightarrow 1$$

[Dreiner, ML – soon ]

Model II:  $\lambda_{123} L_1 L_2 \bar{E}_3$



$$m_{\tilde{\tau}_1} < 175(210 - 225)\text{GeV}, \quad \cos \theta_{\tilde{\tau}} \rightarrow 0$$

$$m_{\tilde{\tau}_1} < 280\text{GeV}, \quad \cos \theta_{\tilde{\tau}} \rightarrow 1$$

# Conclusions

---

- Multilepton searches are efficient in a wide variety of SUSY scenarios
- If one wants to test different scenarios and cover wider regions of the parameter space, one should make use of non-dedicated searches.
- For GGM and stealth SUSY scenarios the use of dedicated searches, multilepton, and DM-dedicated searches allows to cover vast regions of the parameter space
- However multilepton searches can be used to test other models different from the ones they were designed to (for example light staus in RPV)

---

Thank you!!