

Higgs Parity, Strong CP, and Dark Matter

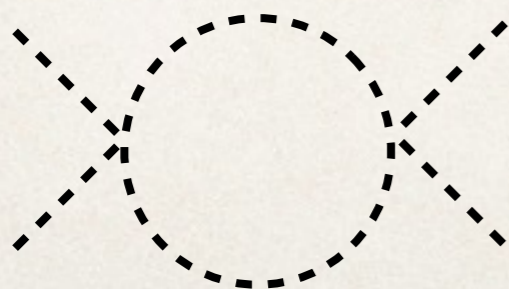
David Dunsky, Lawrence Hall, Keisuke Harigaya

arXiv:1902.07726

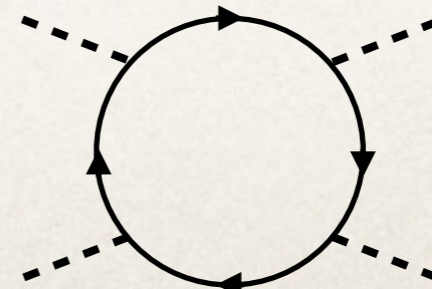
The Higgs Quartic

$$V(H) = -m^2 |H|^2 + \lambda |H|^4$$

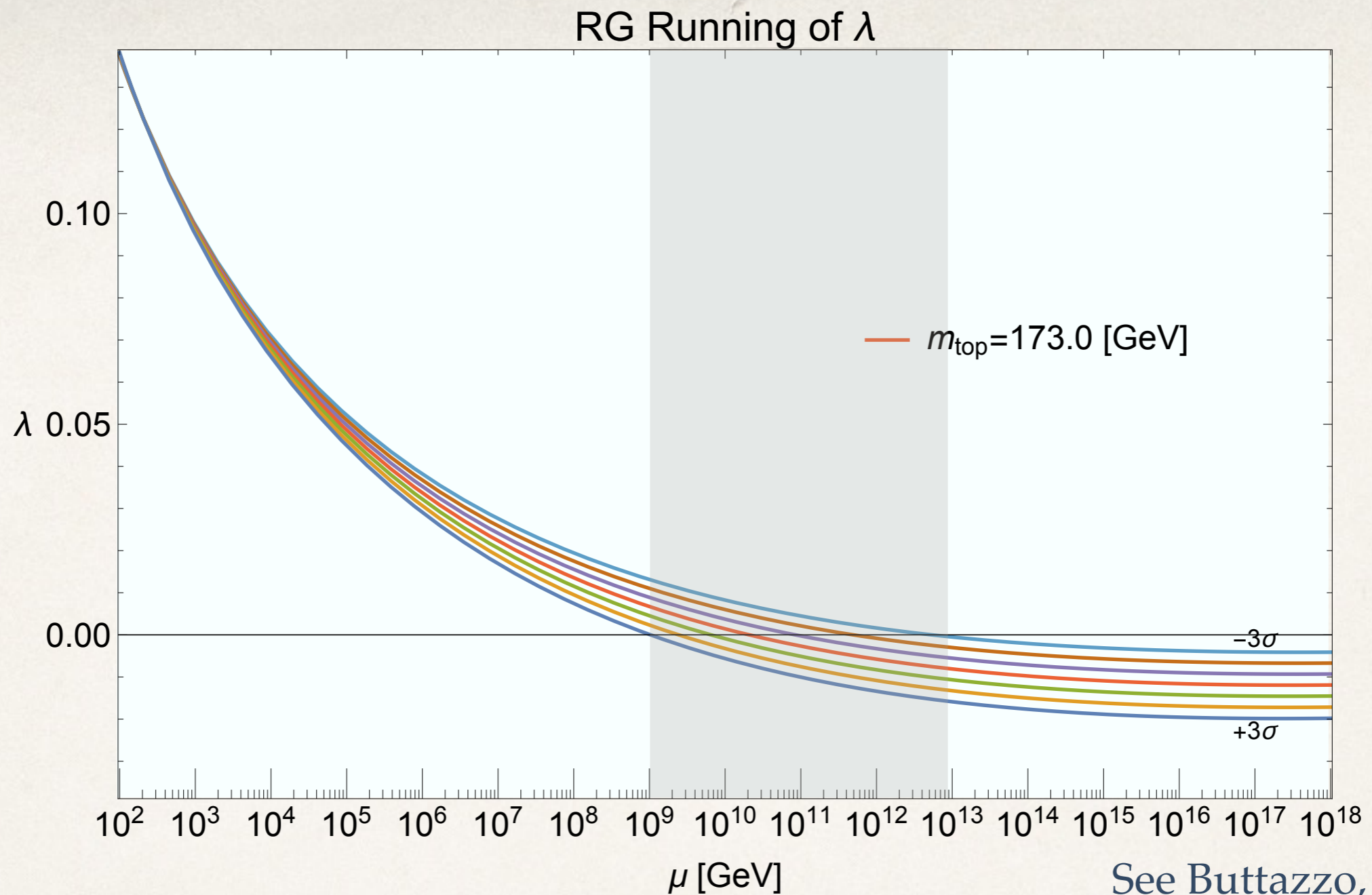
- 1967: $\langle H \rangle = v$ Weinberg
- 2012: m, λ ATLAS, CMS
- Measured value λ of appears special



λ ↑



λ ↓



- Why is $\lambda \sim .01$ above 10^9 GeV?
- Why λ crosses 0 between $10^9 - 10^{13}$?
- Vanishing of λ hint of new physics?

Vanishing of Higgs Quartic by a Z_2

- Consider a Z_2 $SU(2) \leftrightarrow SU(2)'$
 $H \leftrightarrow H'$
 $(2, 1) \quad (1, 2)$

Vanishing of Higgs Quartic by a Z_2

- Consider a Z_2 $SU(2) \leftrightarrow SU(2)'$
 $H \leftrightarrow H'$
 $(2, 1) \quad (1, 2)$

- Most general potential

$$V(H, H') = -m^2(|H|^2 + |H'|^2) + \frac{\lambda}{2}(|H|^2 + |H'|^2)^2 + \lambda'|H|^2|H'|^2$$

Vanishing of Higgs Quartic by a Z_2

- Consider a Z_2 $SU(2) \leftrightarrow SU(2)'$
 $H \leftrightarrow H'$
 $(2, 1) \quad (1, 2)$

- Most general potential

$$V(H, H') = -m^2(|H|^2 + |H'|^2) + \frac{\lambda}{2}(|H|^2 + |H'|^2)^2 + \lambda'|H|^2|H'|^2$$

$$\langle H' \rangle^2 \equiv v'^2 = \frac{m^2}{\lambda} \gg v^2$$

↑
(174 GeV)²

Vanishing of Higgs Quartic by a Z_2

- Consider a Z_2 $SU(2) \leftrightarrow SU(2)'$
 $H \leftrightarrow H'$
 $(2, 1) \quad (1, 2)$

- Most general potential

$$V(H, H') = -m^2(|H|^2 + |H'|^2) + \frac{\lambda}{2}(|H|^2 + |H'|^2)^2 + \lambda'|H|^2|H'|^2$$

Integrate out H'

$$\langle H' \rangle^2 \equiv v'^2 = \frac{m^2}{\lambda} \gg v^2$$

$$V_{\text{LE}}(H) = \lambda' v'^2 |H|^2 - \lambda' \left(1 + \frac{\lambda'}{2\lambda} \right) |H|^4$$

$$(174 \text{ GeV})^2$$

Vanishing of Higgs Quartic by a Z_2

- Consider a Z_2

$$SU(2) \leftrightarrow SU(2)'$$

$$\begin{matrix} H & \leftrightarrow & H' \\ (2, 1) & & (1, 2) \end{matrix}$$

- Most general potential

$$V(H, H') = -m^2(|H|^2 + |H'|^2) + \frac{\lambda}{2}(|H|^2 + |H'|^2)^2 + \lambda'|H|^2|H'|^2$$

Integrate out H'

$$\langle H' \rangle^2 \equiv v'^2 = \frac{m^2}{\lambda} \gg v^2$$

$$V_{\text{LE}}(H) = \lambda' v'^2 |H|^2 - \lambda' \left(1 + \frac{\lambda'}{2\lambda} \right) |H|^4$$

$$(174 \text{ GeV})^2$$

$$-m_{\text{SM}}^2$$

Vanishing of Higgs Quartic by a Z_2

- Consider a Z_2

$$SU(2) \leftrightarrow SU(2)'$$

$$\begin{matrix} H & \leftrightarrow & H' \\ (2, 1) & & (1, 2) \end{matrix}$$

- Most general potential

$$V(H, H') = -m^2(|H|^2 + |H'|^2) + \frac{\lambda}{2}(|H|^2 + |H'|^2)^2 + \lambda'|H|^2|H'|^2$$

Integrate out H'

$$\langle H' \rangle^2 \equiv v'^2 = \frac{m^2}{\lambda} \gg v^2$$

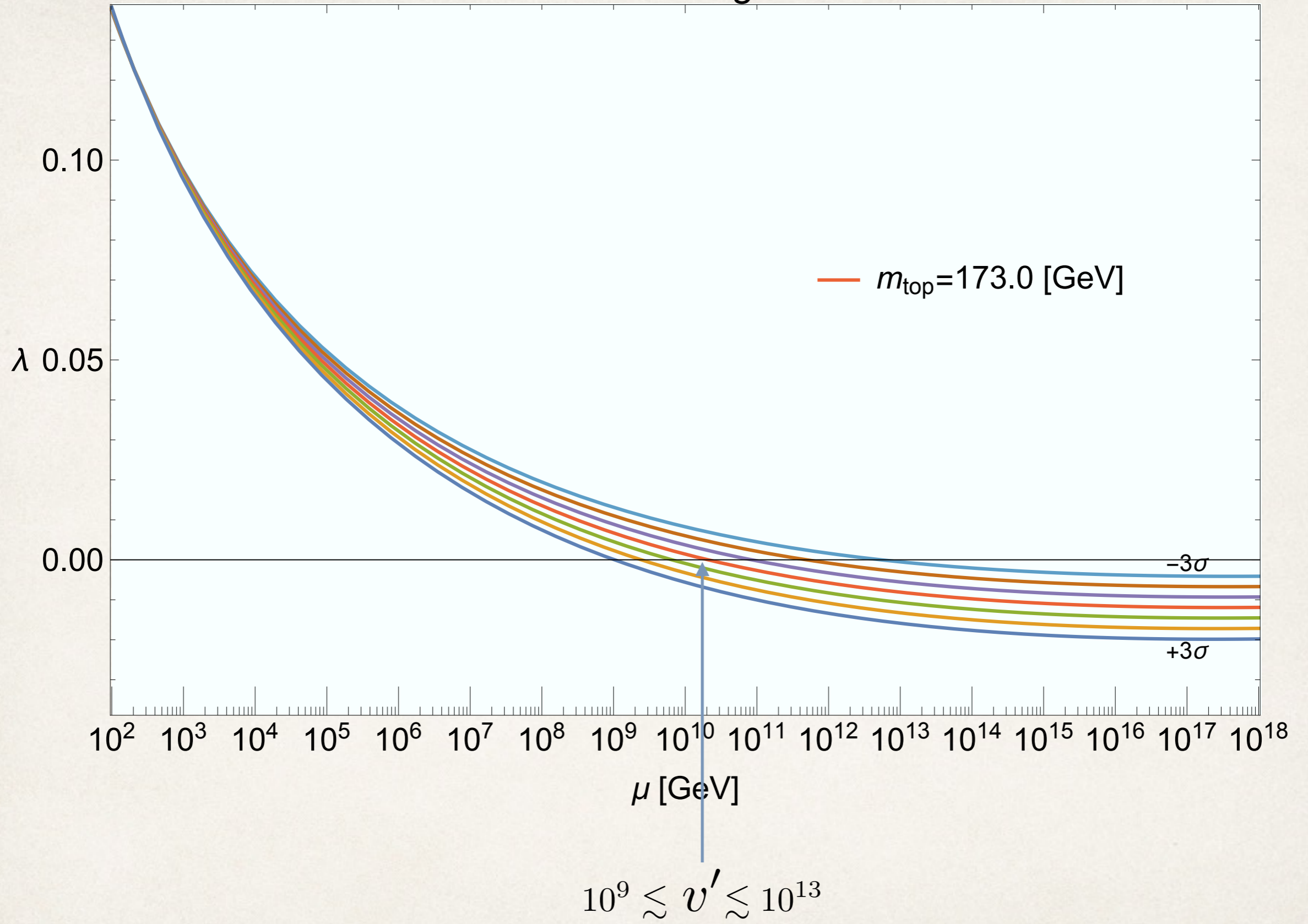
$$V_{\text{LE}}(H) = \lambda' v'^2 |H|^2 - \lambda' \left(1 + \frac{\lambda'}{2\lambda} \right) |H|^4$$

$\underbrace{\quad}_{-m_{\text{SM}}^2} \qquad \underbrace{\quad}_{\lambda_{\text{SM}}}$

$$(174 \text{ GeV})^2$$

- Requiring $v \ll v' \longrightarrow \lambda' \ll 1 \longrightarrow \lambda_{\text{SM}} \approx 0$

RG Running of λ



Fine-Tuning

- Fine tuning required, but same as SM

$$\left(\frac{v^2}{m^2}\right) \times \left(\frac{m^2}{\Lambda^2}\right) = \frac{v^2}{\Lambda^2}$$

Tuning of λ'

Tuning of v'

The Mirror Electroweak Sector

- (1) Extend Z_2 to mirror electroweak sector
- (2) Identify Z_2 with spacetime parity

$$SU(2) \times U(1) \leftrightarrow SU(2)' \times U(1)'$$

$$\vec{x} \leftrightarrow -\vec{x}$$

$$q, \bar{u}, \bar{d}, l, \bar{e} \leftrightarrow (q', \bar{u}', \bar{d}', l', \bar{e}')^\dagger$$

$$H \leftrightarrow H'$$

- Implications?

Solves Strong CP Problem

- $SU(3) \times (SU(2) \times U(1)) \times (SU(2)' \times U(1)')$ solves strong CP

Babu, Chang, Senjanovic (1991)

- Parity is a symmetry

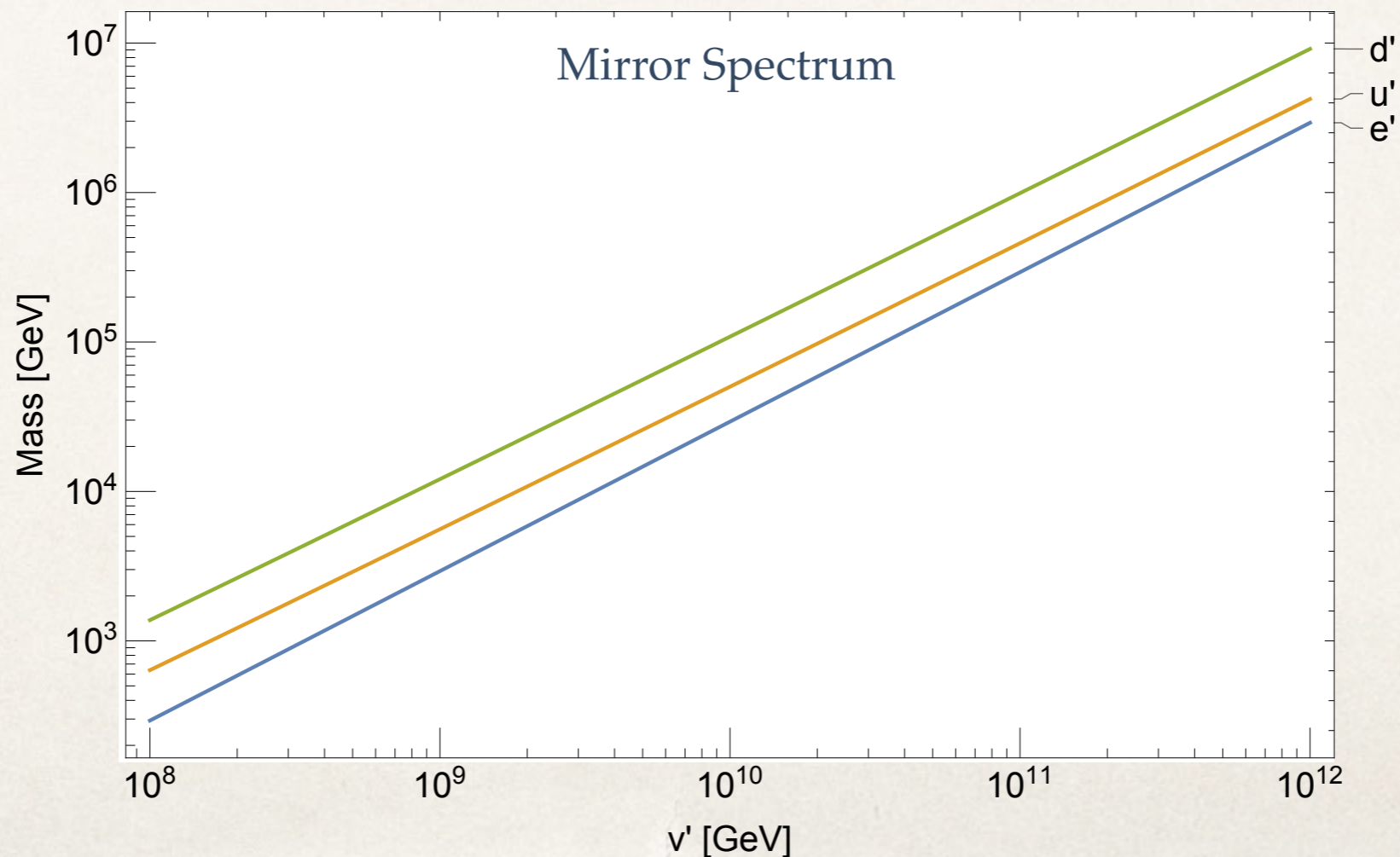
- $\frac{\theta}{32\pi^2} G\tilde{G}$ forbidden

- No contribution from Yukawa sectors

$$\arg \det \begin{pmatrix} y_{u,d} v & 0 \\ 0 & y_{u,d}^* v' \end{pmatrix} = 0$$

Mirror Dark Matter

- Natural DM candidate lightest $U(1)'_{EM}$ particle, e'
- DM mass $m_{e'} = y_{e'} v' = m_e \frac{v'}{v}$ (1-10⁴ TeV)



Signals: Neutron EDM

- Higher dimensional operators generate θ

- $$\mathcal{L}_6 = \frac{C}{M_{Pl}^2} (|H|^2 - |H'|^2) G\tilde{G}$$

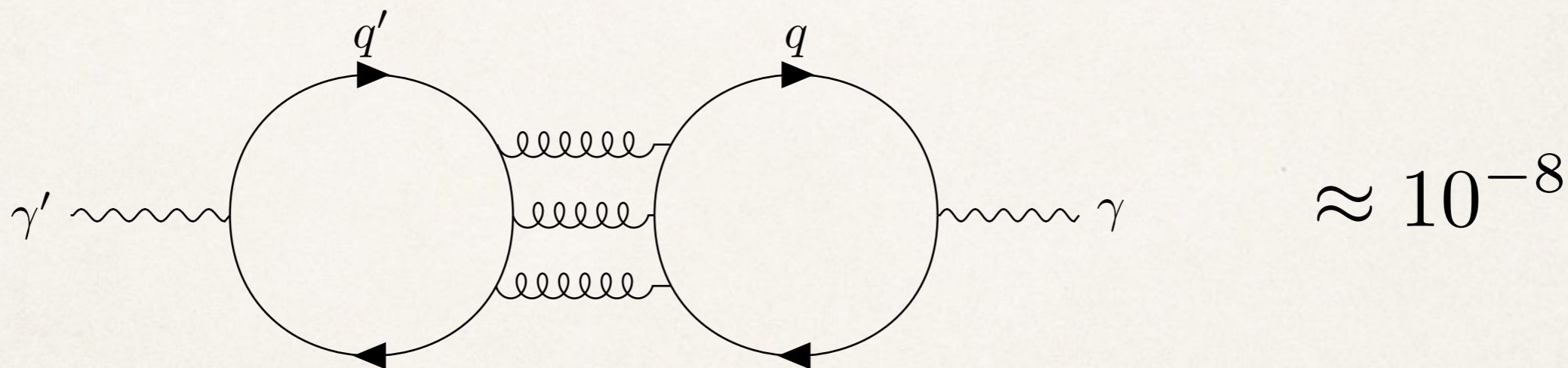
↓

$$\theta = 32\pi^2 C \left(\frac{v'}{M_{Pl}} \right)^2 = 5 \times 10^{-11} C \left(\frac{v'}{10^{12} \text{ GeV}} \right)^2$$

- Current neutron EDM limit $\theta < 10^{-10}$ Baker, et al (2006)

Signals: Kinetic Mixing

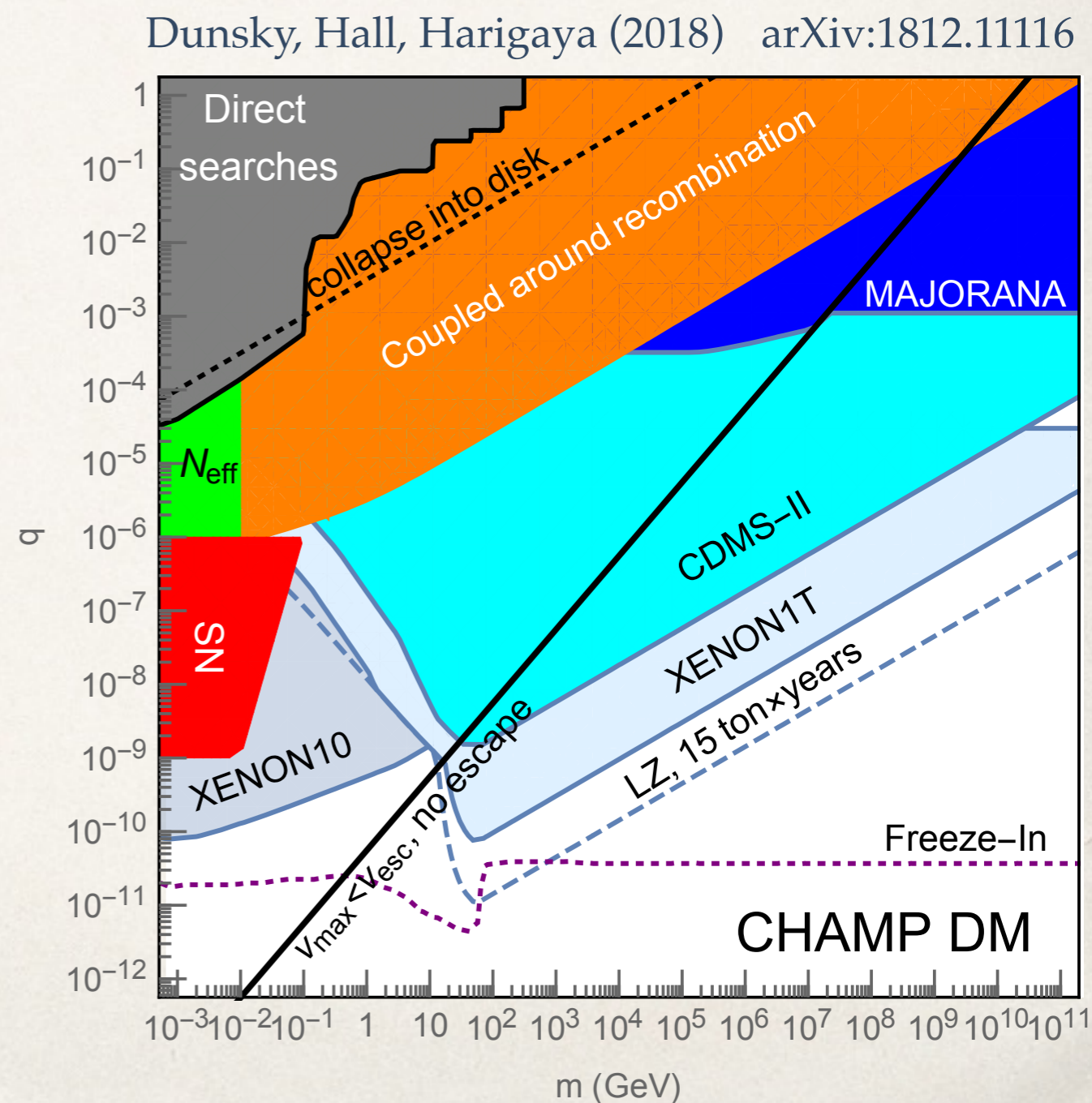
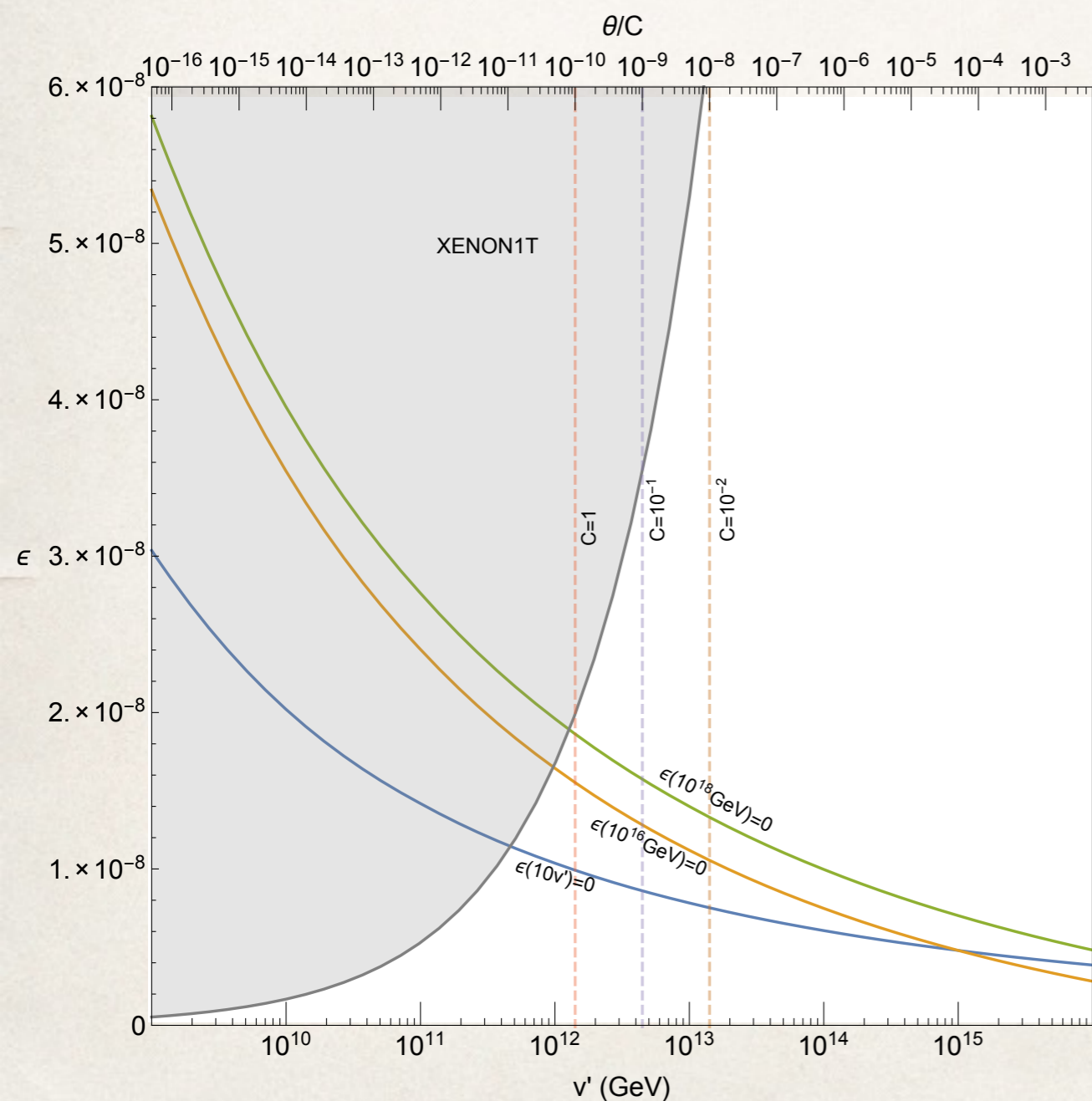
- $\mathcal{L} \supset -\frac{\epsilon_B}{2} B^{\mu\nu} B'_{\mu\nu}$ allowed
- Generated by



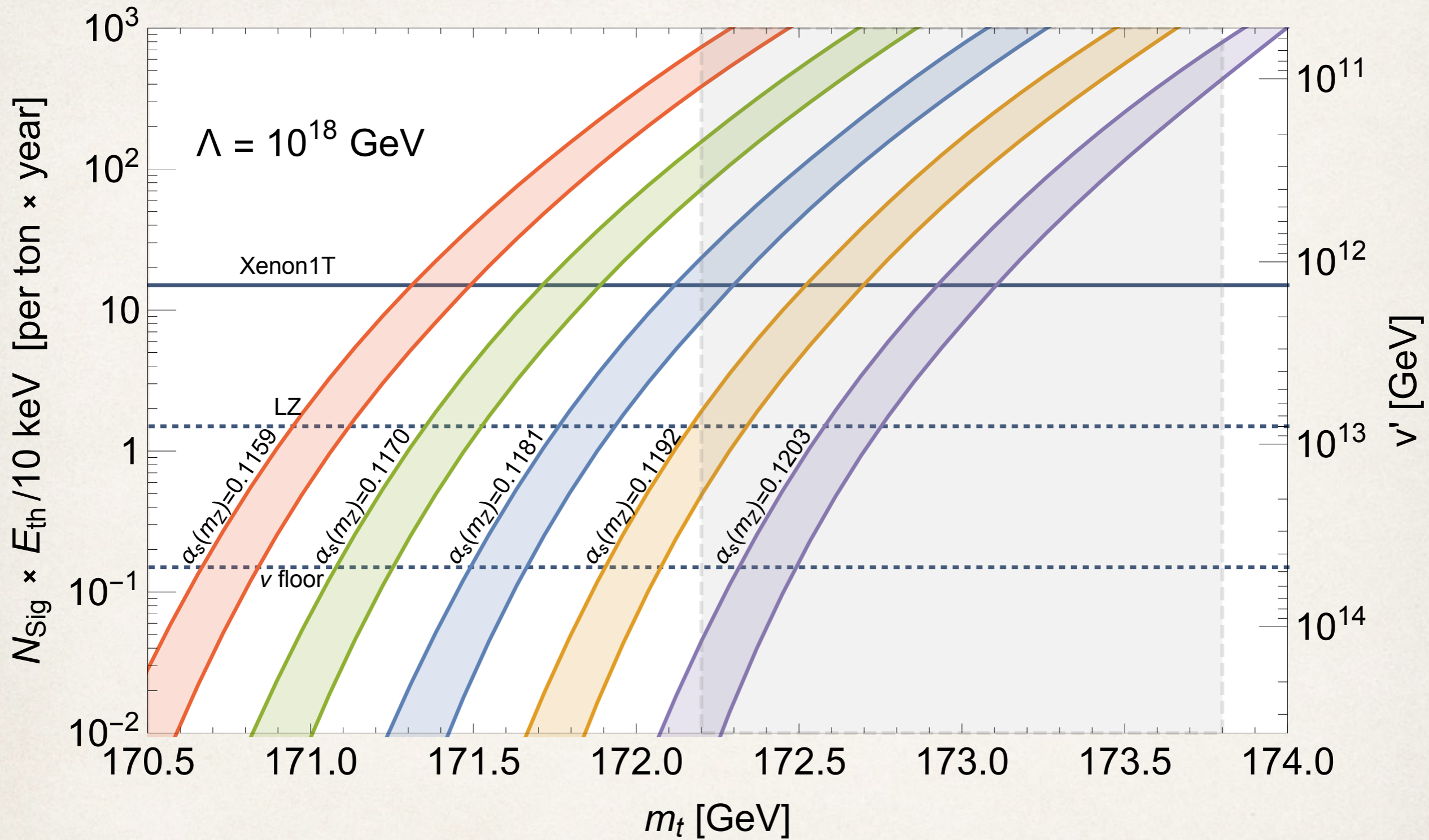
(Calculable from 4-loop QCD beta function)

van Ritbergen, Vermaseren, Larin (1997)

Signals: Kinetic Mixing



e' DM Direct Detection Rate



Summary

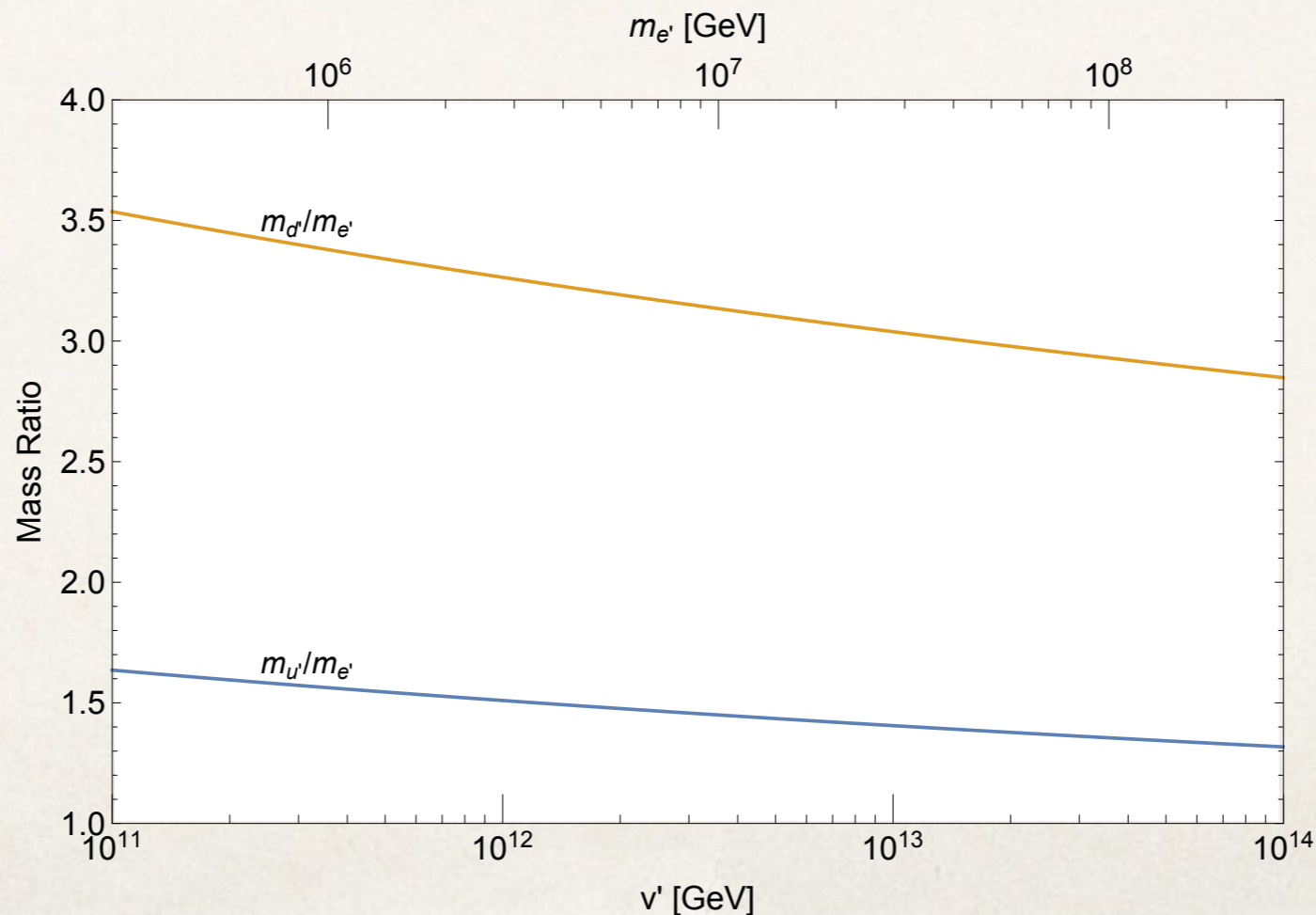
- Observed Higgs mass imply next symmetry breaking scale of nature?
- Motivated by Strong CP \longrightarrow Higgs Parity, no QCD'
- Motivated by DM \longrightarrow mirror electroweak
- Same number of parameters as SM below v'
- Future measurements of $\{m_t, m_h, \alpha_s(M_Z)\}$ will hone in on v'
- Entire parameter space will be probed by future detectors

Cosmology

- Freeze-Out and Freeze-In of e' DM problematic:

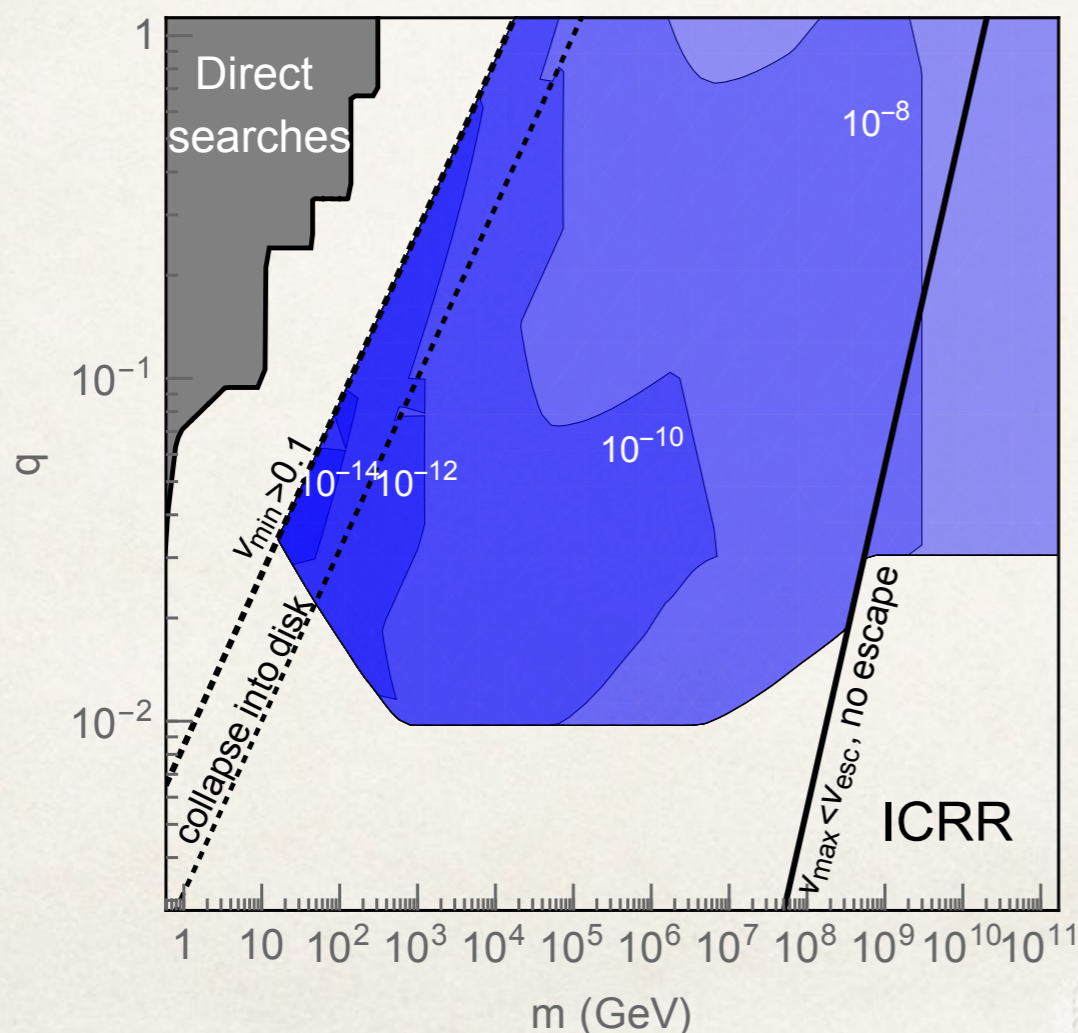
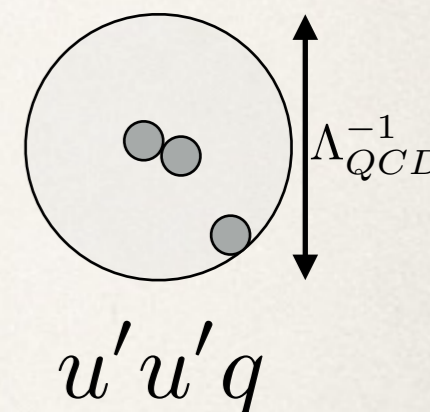
- Thermally produce $e' \rightarrow$ also produce dangerous u'

\uparrow
(also stable)



Fractionally Charged Hadrons

- Below, $T = \Lambda_{QCD}$ quarks hadronize
- Mixed hadrons fractionally charged



$$\frac{\Omega_{u'}}{\Omega_{DM}} \lesssim 10^{-8}$$

Dunsky, Hall, Harigaya (2018)
arXiv:1812.11116

Non-Thermal Production of e' DM

- Non-thermal production of e' required, via decays

- Ingredients for an abundance hierarchy:

1. $T_{RH} \lesssim \frac{m_{u'}}{40}$ \longrightarrow Prevent u' thermal production (Boltzmann suppression)

2. $B_{u'} \lesssim 10^{-8} B_{e'}$ \longrightarrow Prevent u' non-thermal production (kinematic suppression)

3. $B_{e'} \simeq \frac{eV}{T_{RH}} \frac{m_\phi}{m_{e'}}$ \longrightarrow Get correct DM abundance

- Possible via inflaton decays, ν' beta decays, etc.