

# Determination of the endpoint of the first order deconfinement phase transition in the heavy quark region of QCD



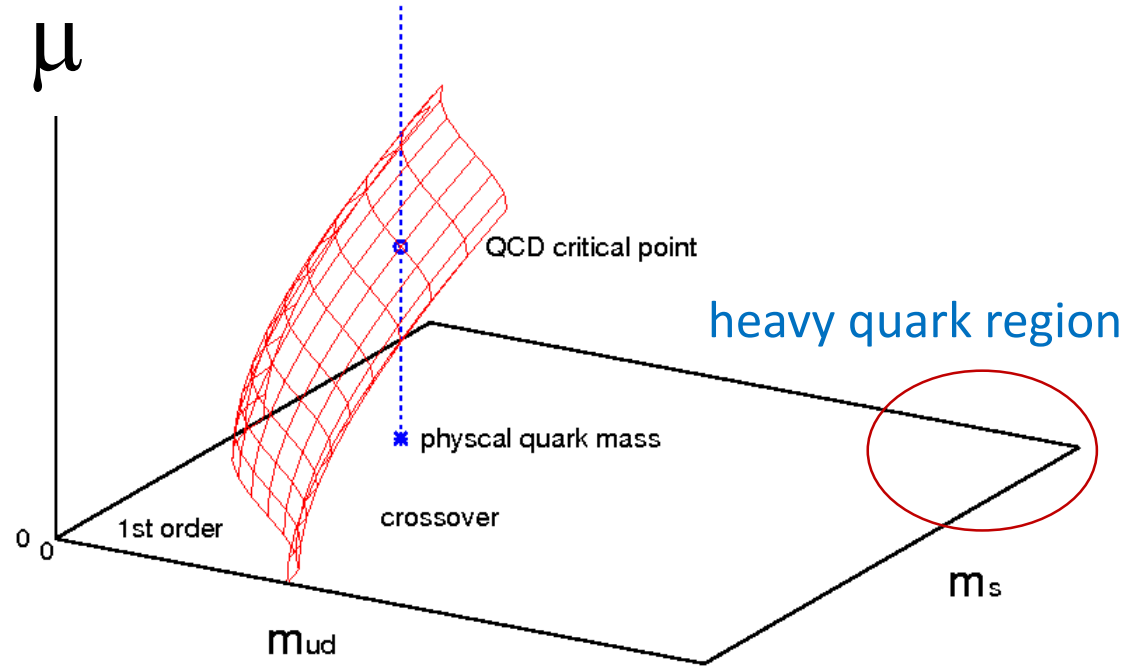
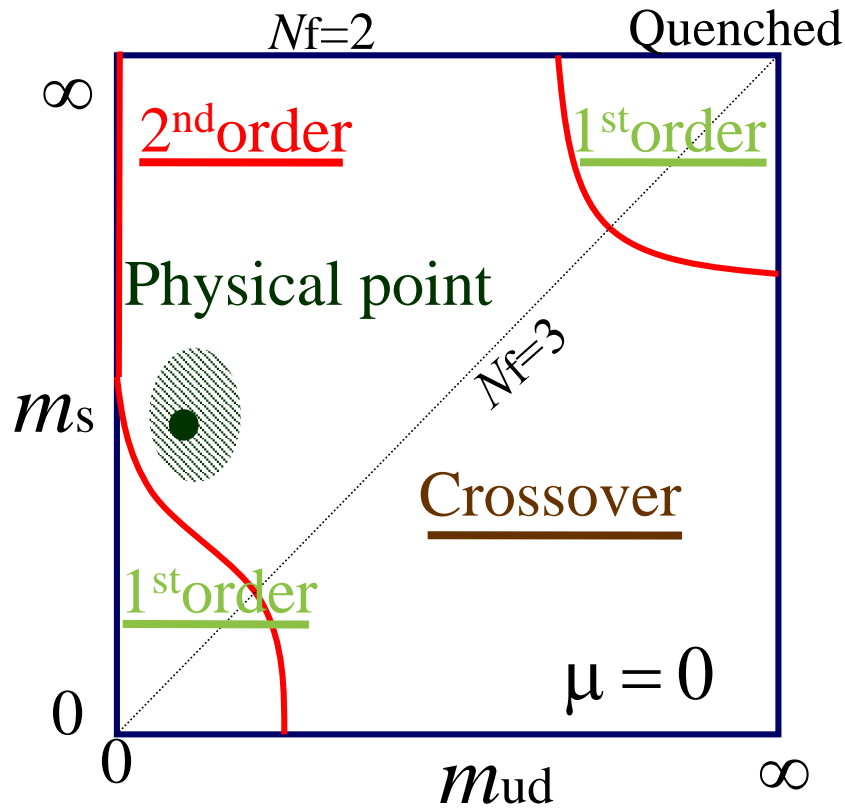
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WHOT-QCD Collaboration

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Lattice 2019, CCNU, Wuhan, June 17-22, 2019

# Quark Mass dependence of QCD phase transition



- The determination of the boundary of 1<sup>st</sup> order region: important.
- On the line of physical mass, the crossover at low density  
➔ 1<sup>st</sup> order transition at high density (?)
- We study the boundary in the heavy quark region.

# Polyakov loop distribution at $\beta_c$ in the complex plane

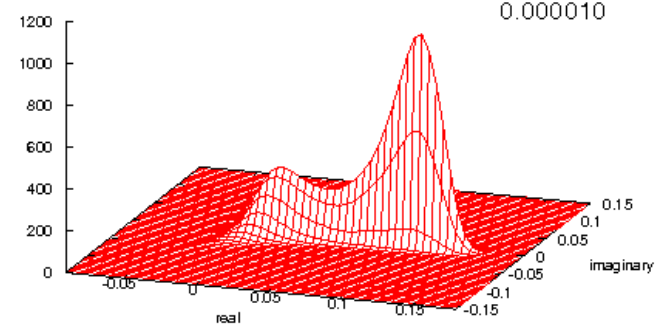
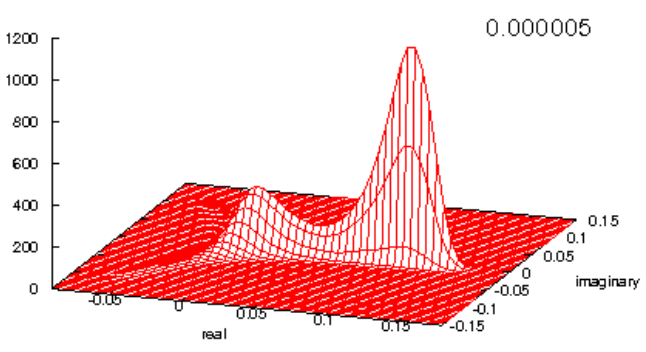
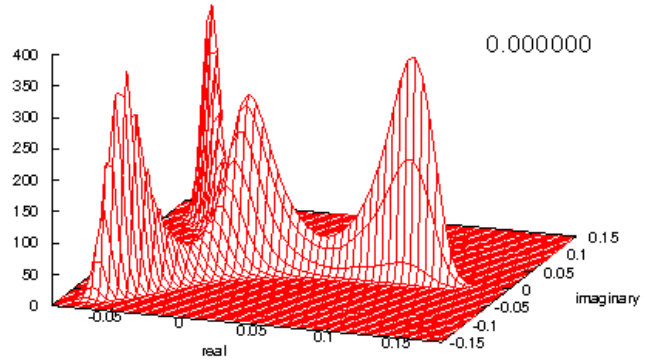
(2-flavor,  $24^3 \times 4$  lattice, Phys.Rev.D89, 034507(2014))

Quenched QCD

$K^4 = 0.0$  **Z(3) symmetric**

$K^4 = 5.0 \times 10^{-6}$

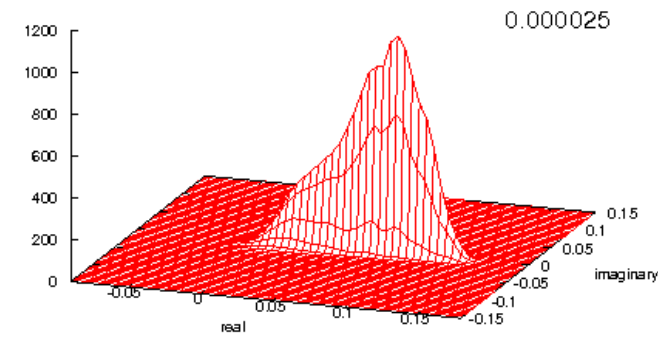
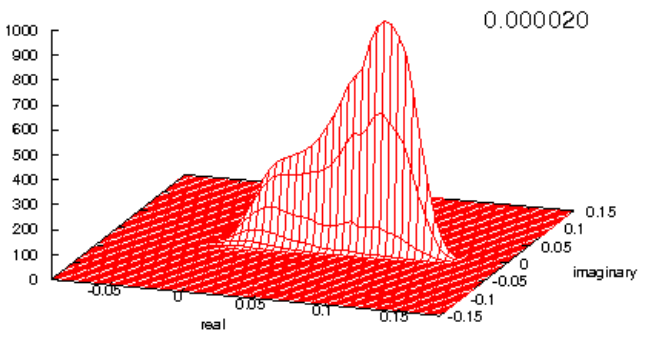
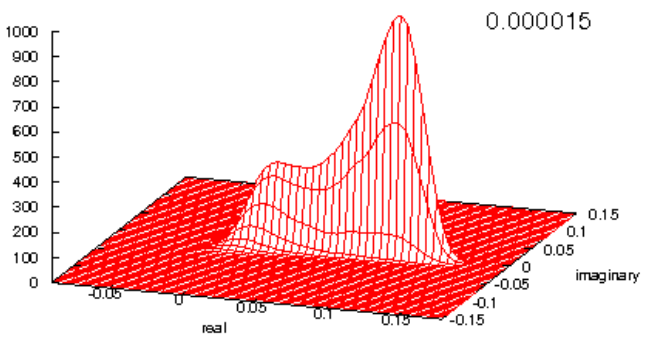
$K^4 = 1.0 \times 10^{-5}$



$K^4 = 1.5 \times 10^{-5}$

$K^4 = 2.0 \times 10^{-5}$

$K^4 = 2.5 \times 10^{-5}$



**critical point**

$K$ : hopping parameter  $\sim 1/(\text{mass})$

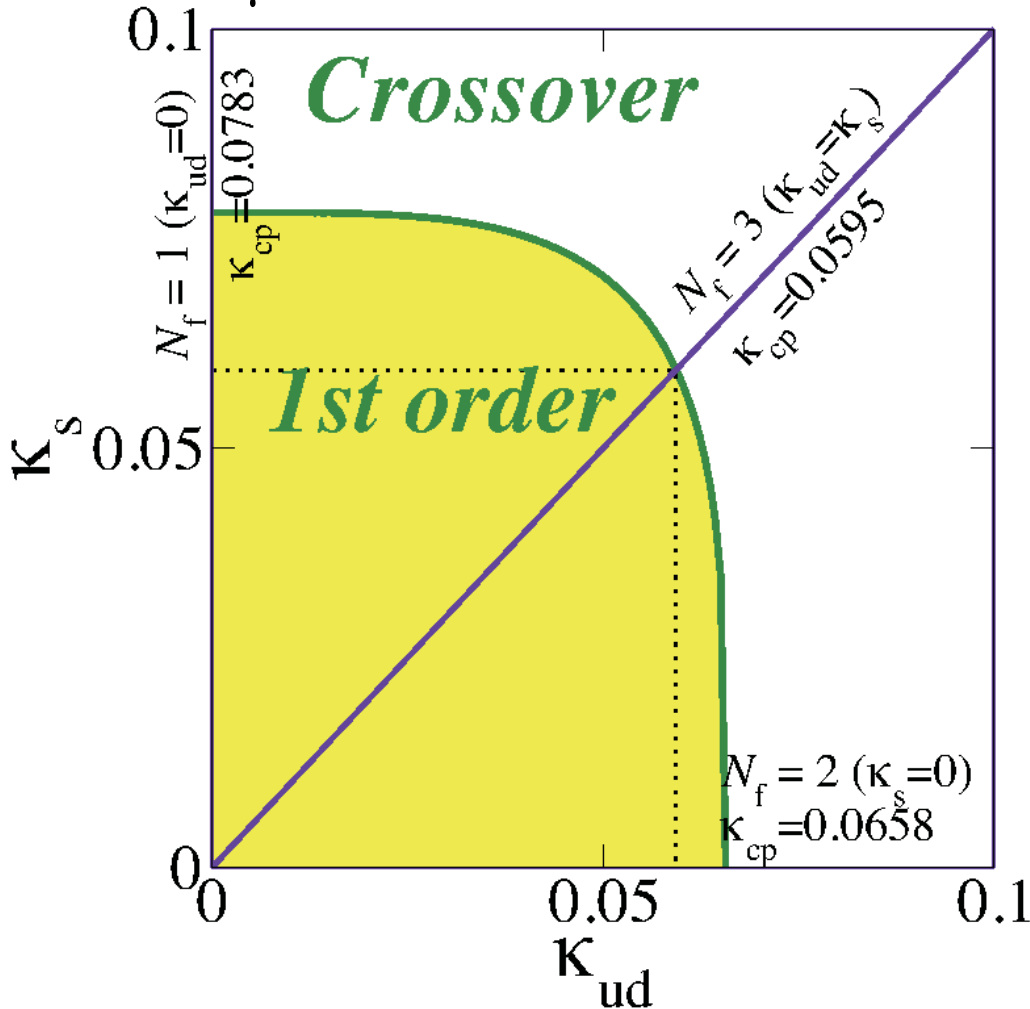
# Critical surface in the heavy quark region of (2+1)-flavor QCD

[Phys.Rev.D84, 054502(2011); Phys.Rev.D89, 034507(2014)]

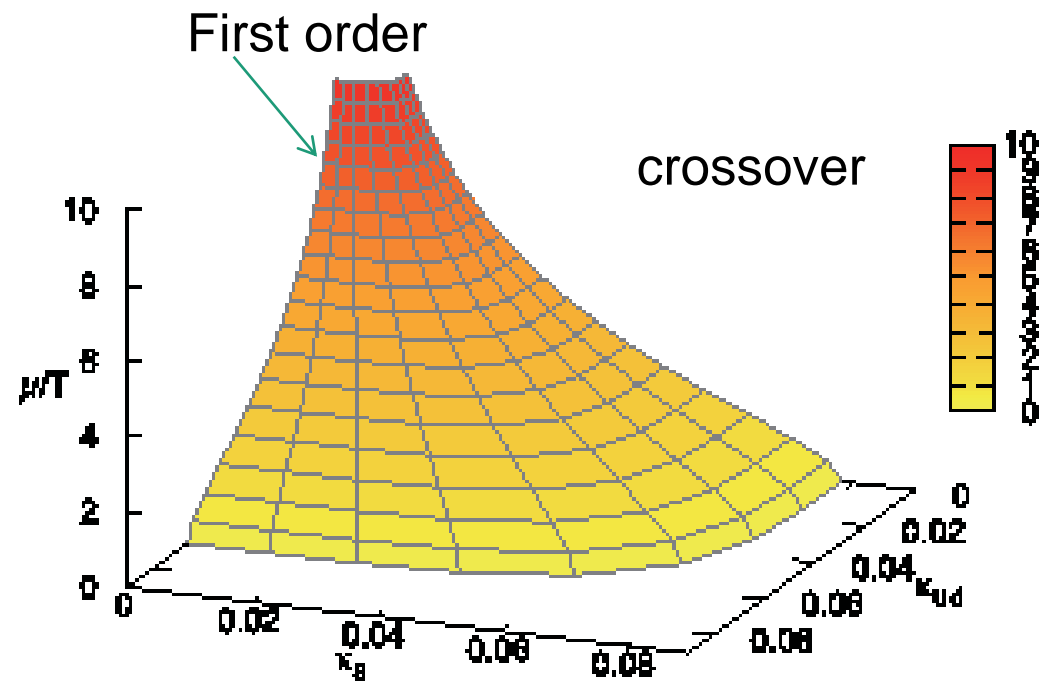
Quenched QCD simulations + reweighting

( $24^3 \times 4$  lattice)

$\mu = 0$



## Critical surface at finite density



$$\frac{m_{PS}}{T_c} \approx 16. \quad \text{at } \kappa_{cp} \text{ for 2-flavor}$$

# This talk

- Calculation of  $Kc$  on lattices of  $N_t=6$  and 8
- Limitation of this analysis by the shape of the histogram
  - Lattice spacing dependence
    - truncation error of hopping parameter expansion
  - Spatial volume dependence → Overlap problem
- New approach
  - Simulations with a Polyakov loop term
  - Finite volume scaling analysis

# Histogram method

- Probability distribution function (Histogram)

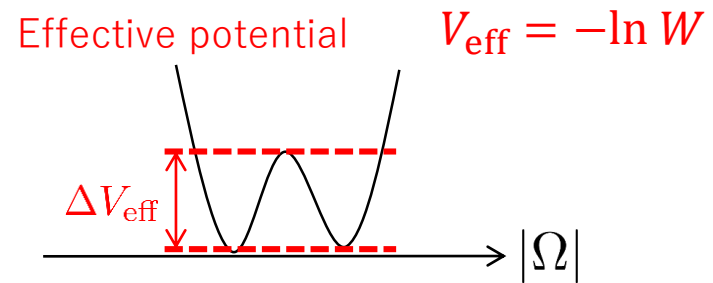
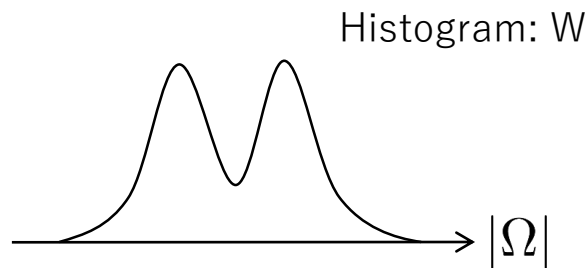
$\Omega$ : Polyakov loop (order parameter)

$$W(\Omega; \beta, K) \equiv \frac{1}{Z} \int DU \delta(\Omega - \hat{\Omega}) \prod_{f=1}^{N_f} \det M(K) e^{-S_g}$$

( $S_g$ : gauge action,  $M$ : quark matrix)

- Effective potential  $V_{\text{eff}} = -\ln W$

First order  
transition



- Critical point of  $K$ :  $\Delta V_{\text{eff}} = 0$

# Reweighting method in the heavy quark region

- Quenched QCD simulations + reweighting

$$W(\Omega; \beta, K) \equiv \frac{1}{Z} \int DU \delta(\Omega - \hat{\Omega}) \prod_{f=1}^{N_f} \det M(K) e^{-S_g} \quad \text{Histogram}$$

$$= \frac{\langle \delta(\Omega - \hat{\Omega}) \prod_f \det M(K) \rangle_{\text{quench}}}{\langle \prod_f \det M(K) \rangle_{\text{quench}}}$$

- Multi-point ( $\beta$ ) reweighting method is used.
- Hopping parameter expansion ( $K \sim 1/(ma)$ )

$$\ln(\det M(K)) = 288 N_{\text{site}} K^4 P + [768 N_{\text{site}} K^6 (3 \text{ [plaquette] } + \text{ [6-step Wilson loop] } + 6 \text{ [3-step Wilson loop] } )] + \dots$$

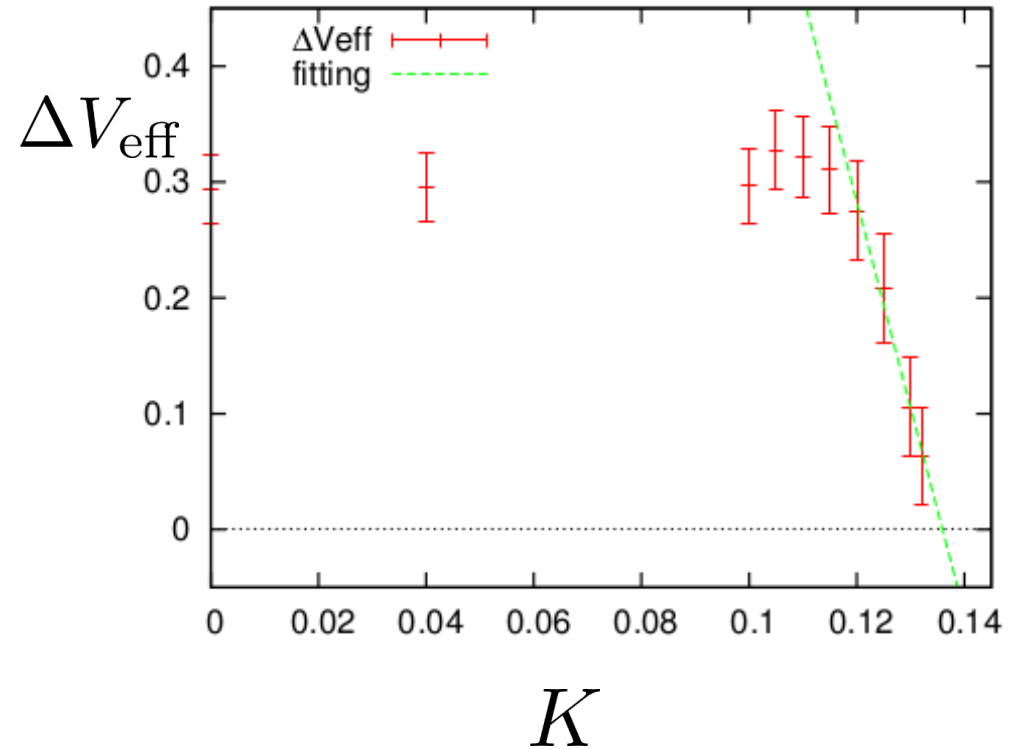
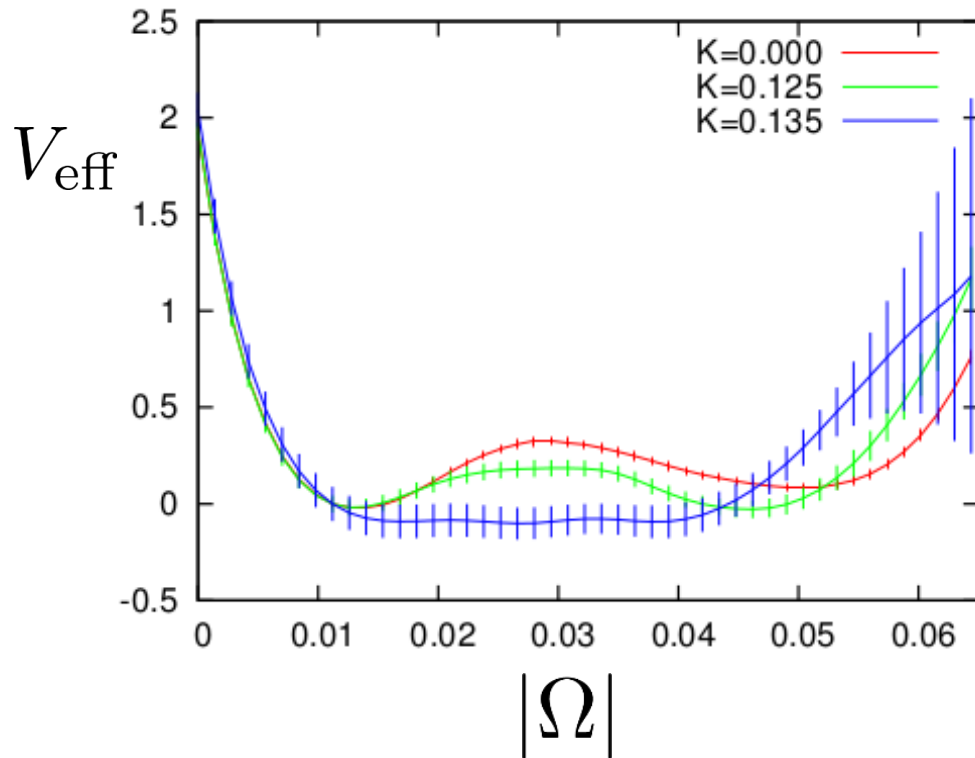
$$+ 12 \times 2^{N_t} \times N_s^3 \left[ \underbrace{K^{N_t} \text{Re}\Omega}_{\text{Polyakov loop}} + 6 N_t K^{N_t+2} \text{ [Wilson loop] } + 6 N_t K^{N_t+2} \text{ [Wilson loop] } + 3 N_t K^{N_t+2} \text{ [Wilson loop] } \right]$$

Leading term Next to leading terms

\* plaquette and 6-step Wilson loop can be absorbed into the gauge action.

# Determination of $K_c$ (leading order calculation)

$N_t = 6$  lattice ( $24^3 \times 6$ )  
Effective potential



$$K_c = 0.1359(30)$$

Lattice	Configurations
$24^3 \times 6$	676,190
$32^3 \times 6$	1,172,000
$24^3 \times 8$	342,821

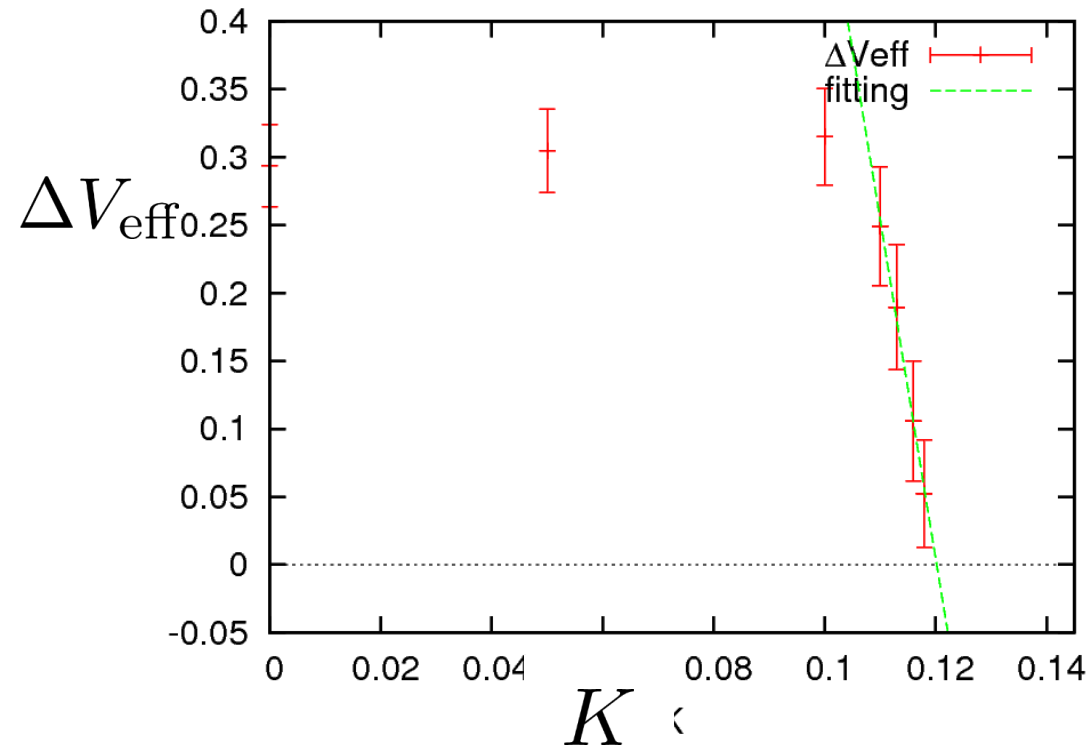
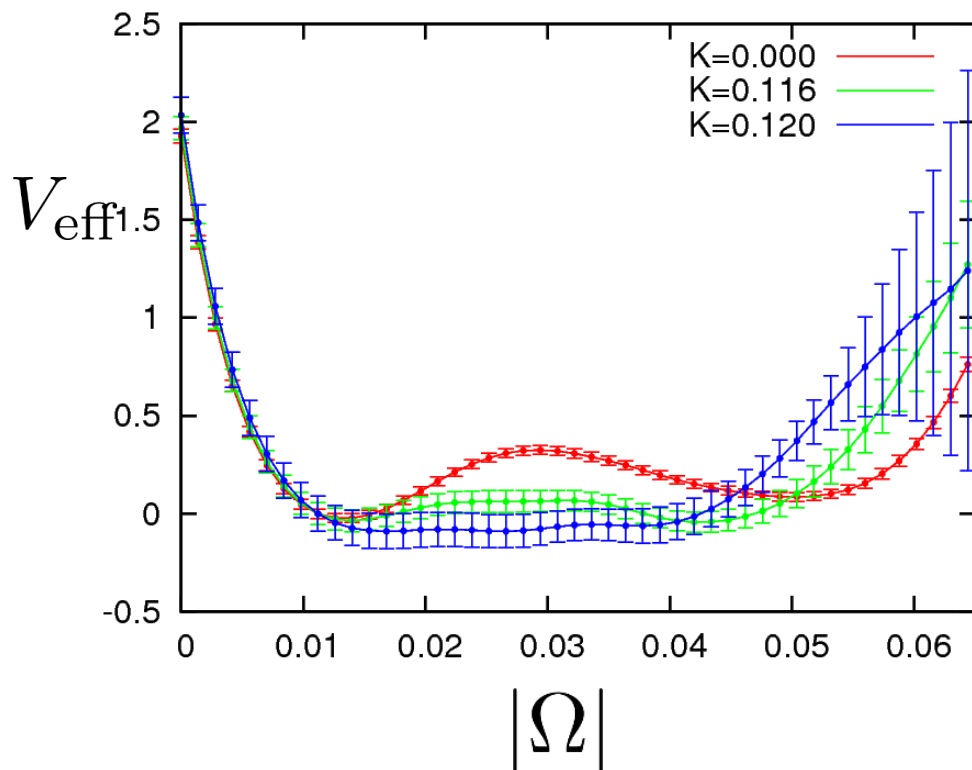


# Determination of $K_c$ (Next to leading order calculation)

To estimate the truncation error of the hopping parameter, the next to leading contribution is computed.

$N_t = 6$  lattice ( $24^3 \times 6$ )

Effective potential



$$K_c = 0.1202(19)$$

Lattice	Configurations
$24^3 \times 6$	676,190
$32^3 \times 6$	1,172,000
$24^3 \times 8$	342,821

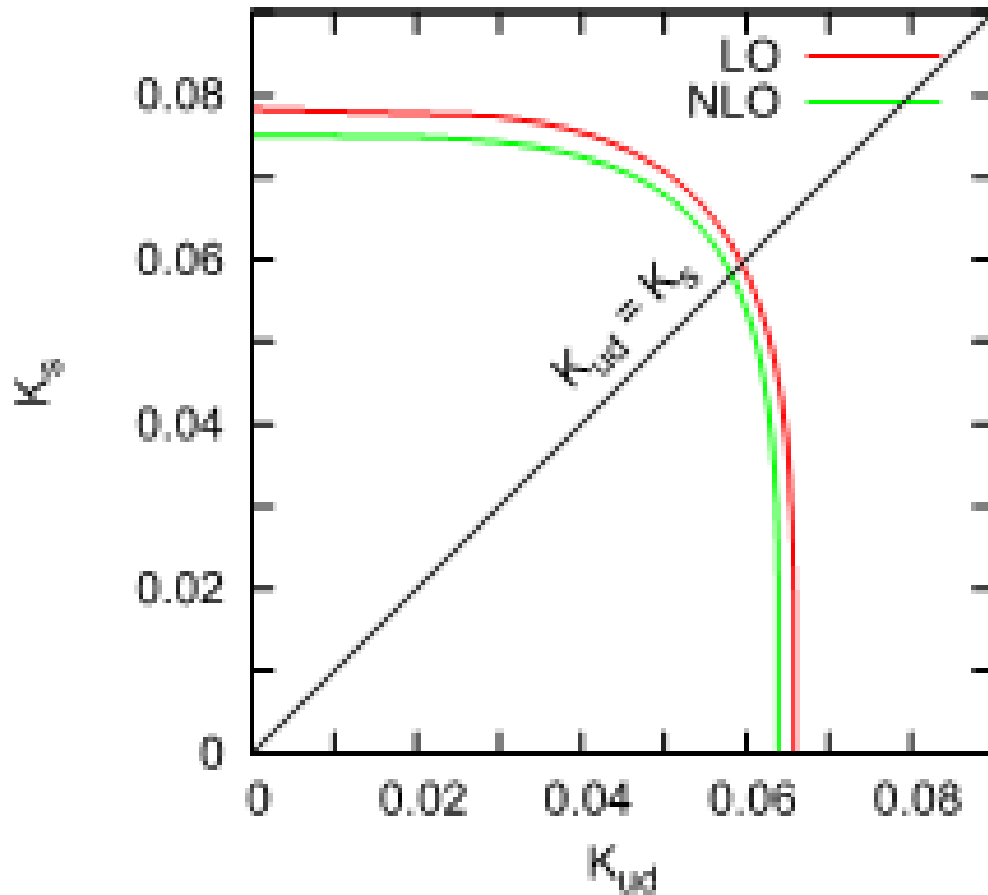
# Determination of Critical K for 2-flavor QCD

	Leading order		Next to leading order	
Lattice	$K_c$	$m_{PS}/T_c$	$K_c$	$m_{PS}/T_c$
$24^3 \times 4$	0.0658(10)	15.47(14)	0.0640(10)	15.73(14)
$24^3 \times 6$	<u>0.1359(30)</u>	<u>7.43(78)</u>	<u>0.1202(19)</u>	<u>11.15(42)</u>
$24^3 \times 8$	> 0.18 > (chiral limit)			

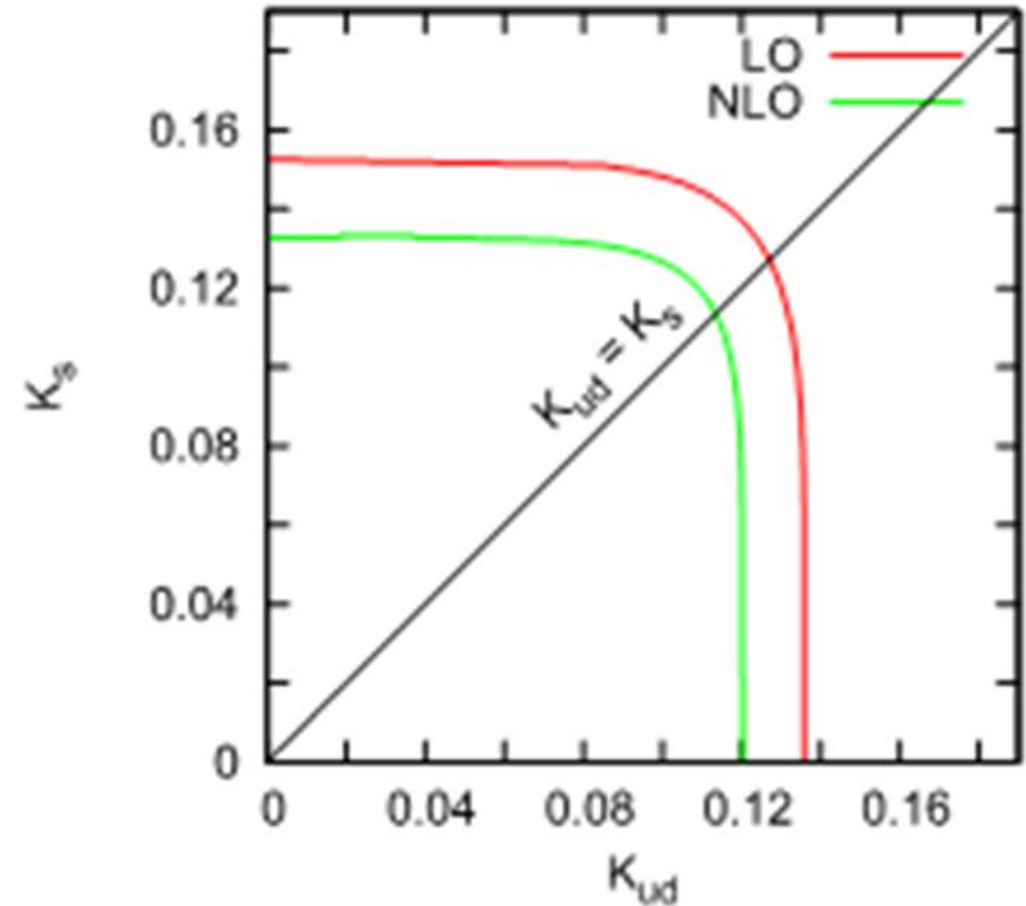
- In the previous study on a  $24^3 \times 4$  lattice, the truncation error of the hopping parameter expansion is negligible.
- The truncation error is visible for the  $24^3 \times 6$  lattice.
- Pseudo-scalar meson mass  $m_{PS}$  measured by  $T=0$  full QCD simulations at  $K_c$  for  $N_t=6$  is smaller than that for  $N_t=4$ .
- This analysis is not applicable for  $N_t=8$ .
- Higher order terms are necessary for large  $N_t$ .

# Determination of $K_c$ in 2+1 flavor QCD

$24^3 \times 4$  lattice



$24^3 \times 6$  lattice



**LO: leading order**      **NLO: next to leading order**

- Difference between the leading order and next to leading calculations are sizeable for the  $24^3 \times 6$  lattice.

# Volume dependence of Critical $K$

Lattice	$K_c$
$24^3 \times 6$	0.1359(30)
$32^3 \times 6$	0.1286(40)
$36^3 \times 6$	[overlap problem]

- $K_c$  becomes smaller as the volume increases.
- $36^3 \times 6$  lattice: overlap problem arises before  $K_c$ .
- Reweighting from quenched simulation does not work for large volume.
  - > Simulations with a Polyakov loop term.

# Simulations with the Polyakov loop term $\Omega$

- Simulations with an effective action

$$S_{\text{eff}} = -6N_s^3 \beta P + N_s^3 \lambda \text{Re}\Omega$$

- The Polyakov loop term corresponds to the leading order contribution of the hopping parameter expansion of  $\det M$ .

$$\lambda = 384 K^4 \quad (\text{for 2-flavor, } N_t=4)$$

- Heat bath algorithm is applicable.
  - computational cost is small.
- We include the next leading contribution of the hopping parameter expansion by the reweighting.
- Overlap problem can be avoided.
- Simulation parameter
  - Lattice size:  $N_t=4, N_s=32, 36, 40, 48$

# Binder cumulant

$$B_4 = \frac{\langle (\Omega - \langle \Omega \rangle)^4 \rangle}{\langle (\Omega - \langle \Omega \rangle)^2 \rangle^2}$$

- No volume dependence at the critical point  $\lambda_{cp}$
- 3D Ising universality class:  
 $B_4=1.604$  at  $\lambda_{cp}$  and  $\nu=0.63$ .

**Fit result**  $Nt=4$

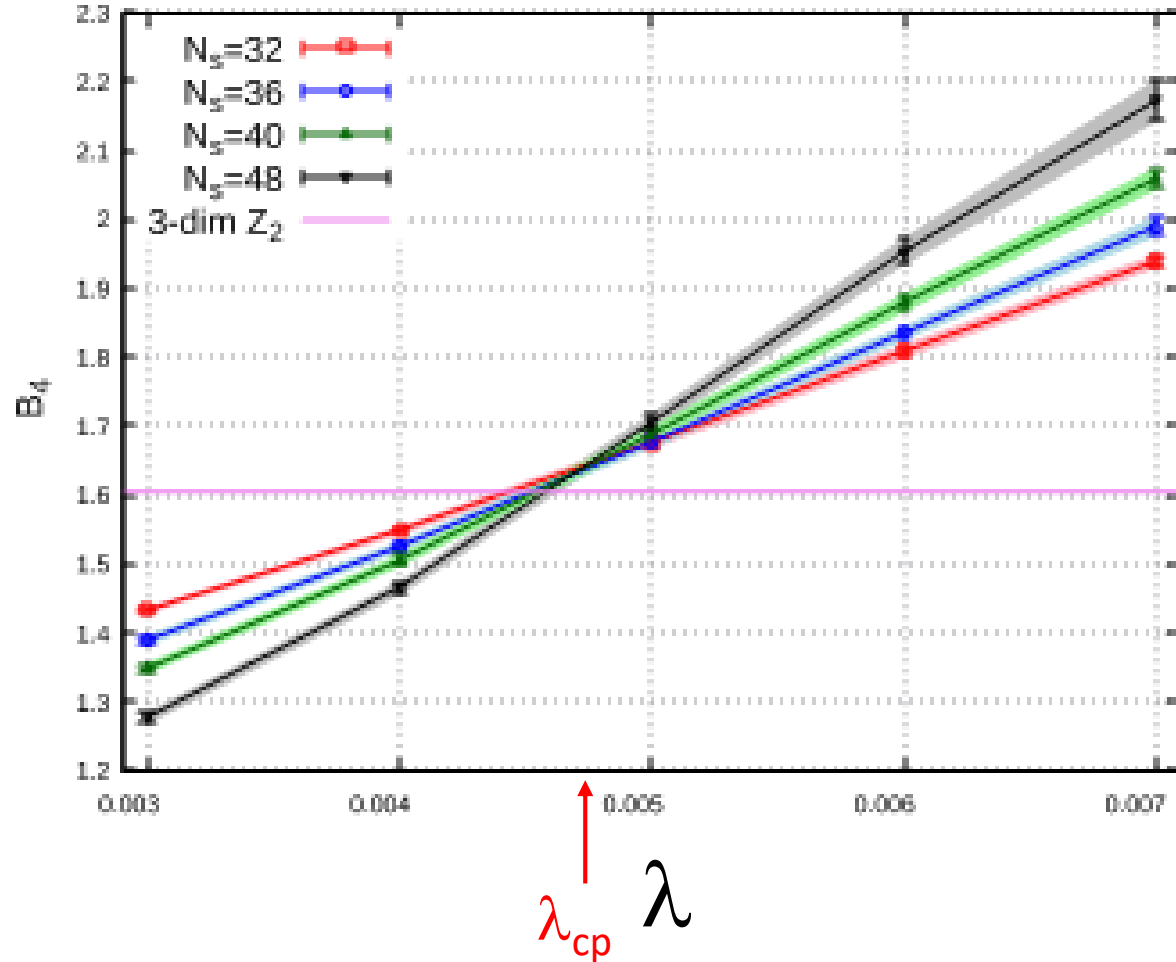
$$B_4(N_s, \lambda) \equiv B_{cp} + AN_s^{1/\nu}(\lambda - \lambda_{cp})$$

$N_s=32, 36, 40, 48$

- $\lambda_{cp}=0.004754(84)$
- $B_{cp}=1.644(13)$
- $\nu=0.65(8)$

$N_s=36, 40, 48$

- $\lambda_{cp}=0.00468(11)$
- $B_{cp}=1.630(20)$
- $\nu=0.65(11)$

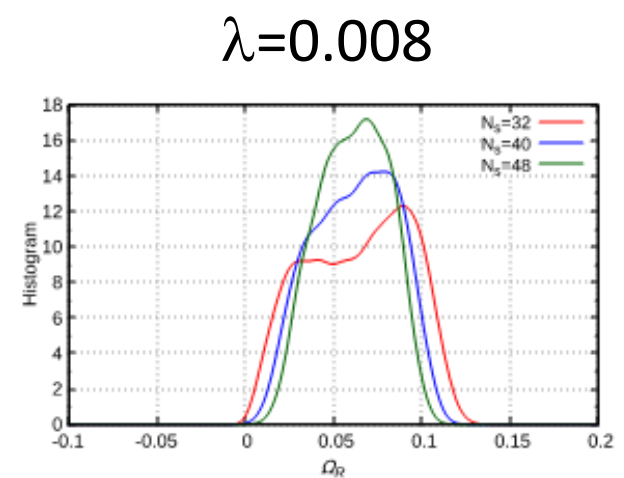
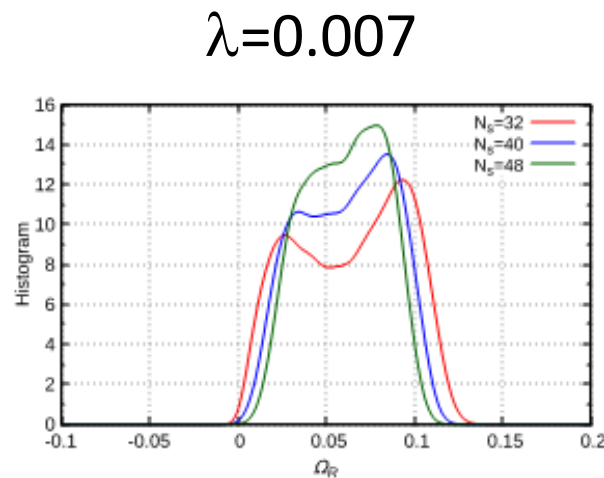
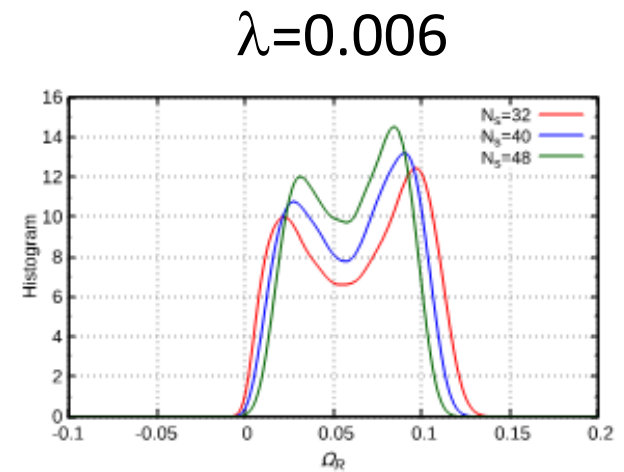
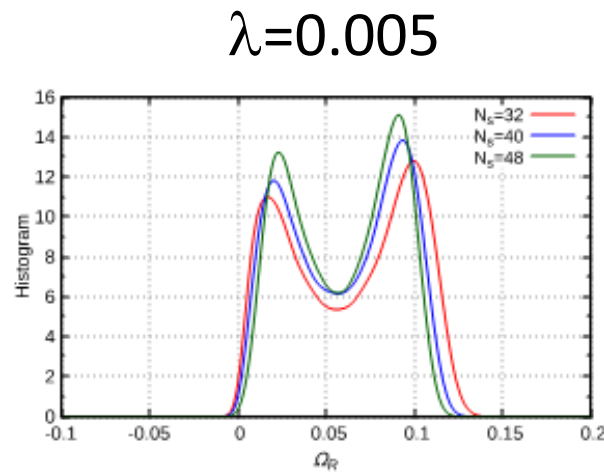
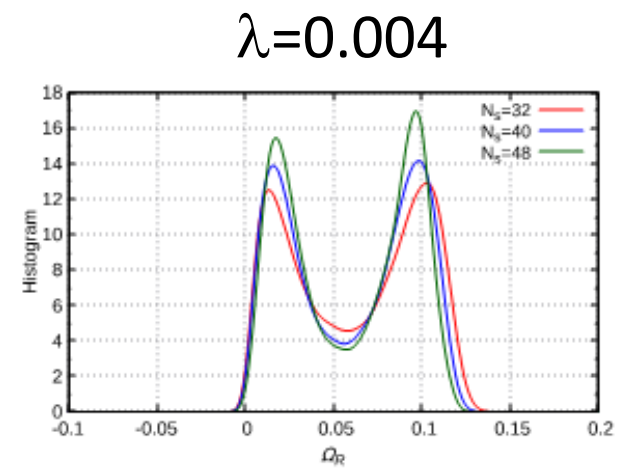
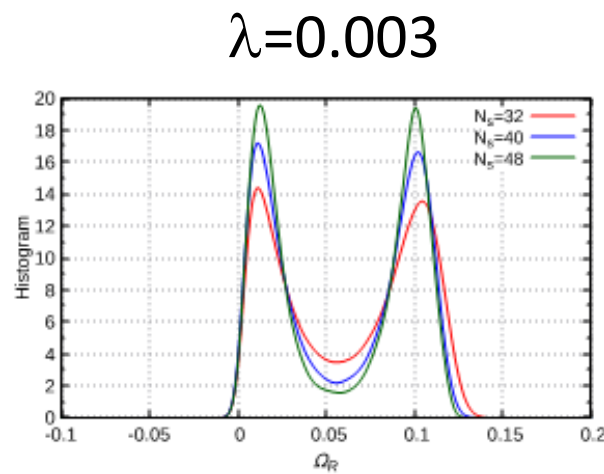


$\lambda_{cp}=0.00475 \rightarrow Kc=0.0539$  is smaller than the result by histogram method:  $0.0640(10)$  ( $24^3 \times 4$  lattice)

# Volume dependence of the histogram

- $\lambda < \lambda_{cp}$ , The middle dent in the histogram gets deeper as the volume increases.
- $\lambda > \lambda_{cp}$ , the middle dent becomes shallower and disappears.
- $\lambda_{cp}$  is the boundary that divides one peak or two peaks in the volume infinity limit.

$$\lambda_{cp} = 0.004754(84)$$



# Finite volume Scaling of the gap

- Assuming the gap  $\Delta$  corresponds to (magnetization)/(2×volume) of 3D ising,

$$\Delta \propto N_S^{-0.519}$$

is expected at  $\lambda_{cp}$ .

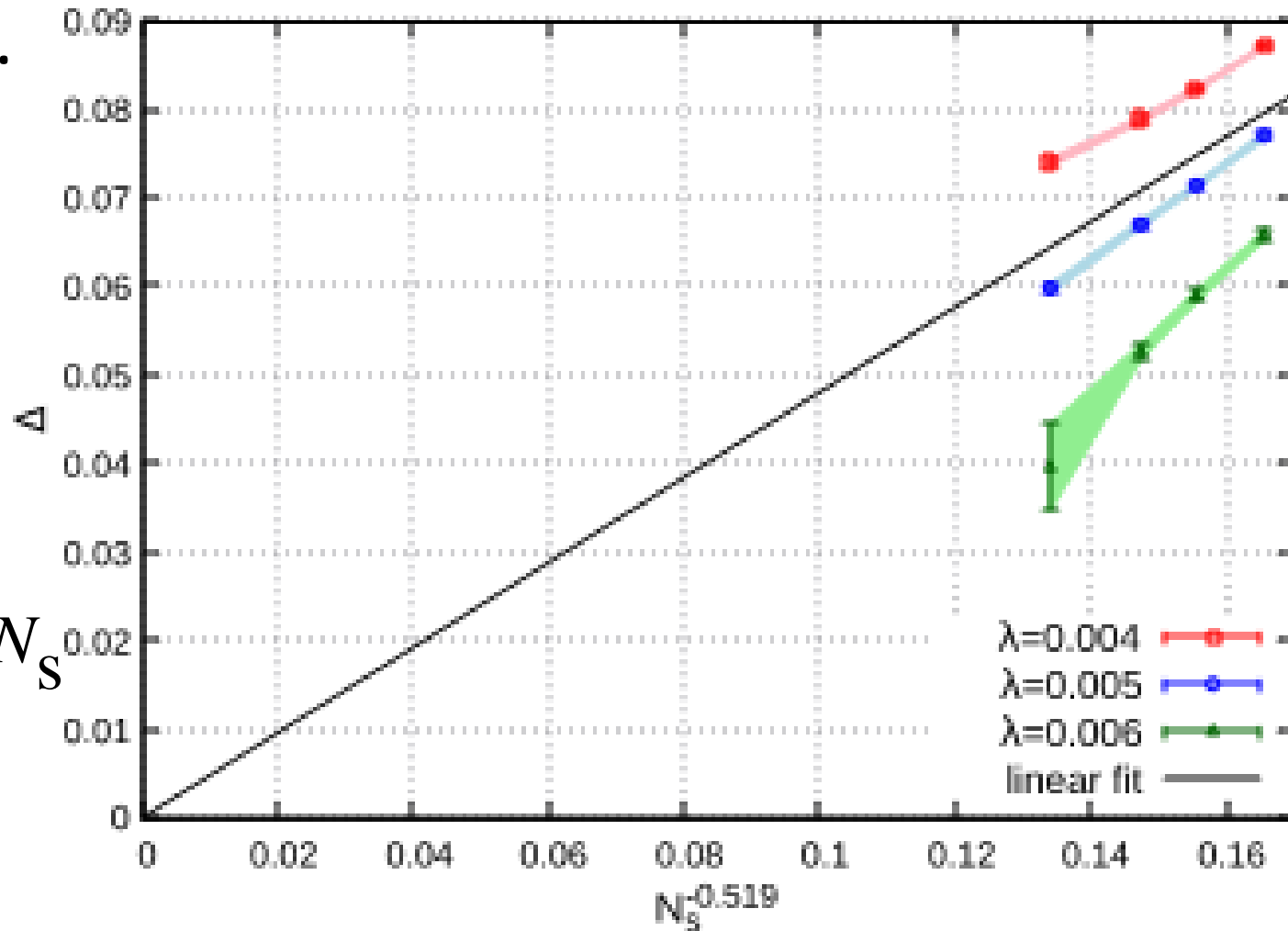
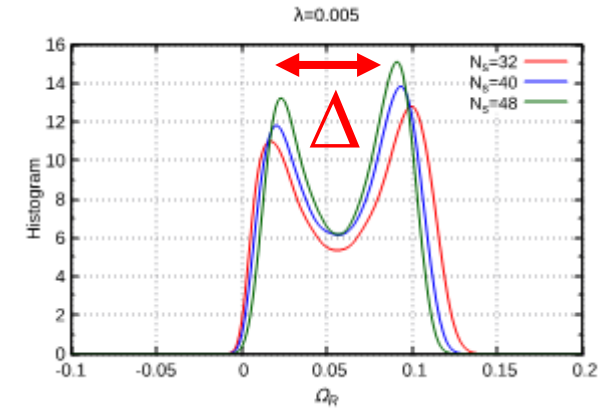
In  $N_S \rightarrow \infty$ ,

- $\lambda=0.005$ :

Linearly decrease

- $\lambda=0.006$ :

$\Delta$  vanishes at finite  $N_S$





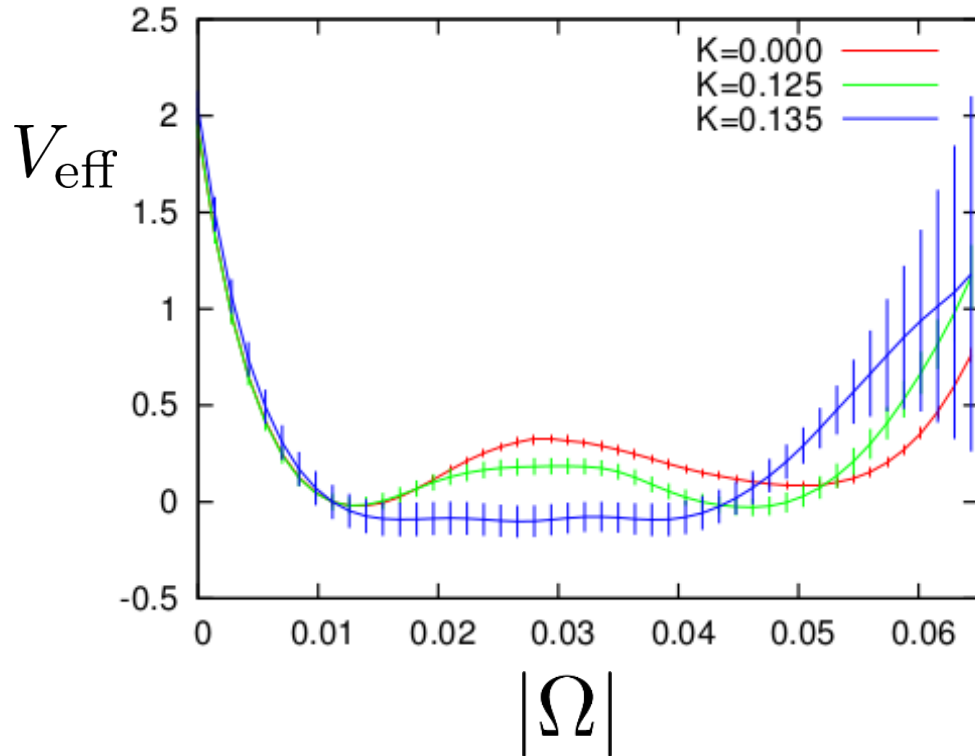
# Summary and out look

- We study the location of critical point at which the first order phase transition changes to crossover in the heavy quark region by investigating the histogram of the Polyakov loop and applying the finite-size scaling analysis.
  - simulations of quenched QCD
  - reweighting method
  - quark determinant: hopping parameter expansion
- Truncation error of the hopping parameter expansion: Visible for  $N_t \geq 6$ .  $\rightarrow$  Higher order terms are needed.
- Overlap problem arises for large volume.
- To reduce the overlap problem, we introduce an external source term of the Polyakov loop in the simulation.
  - Scaling behavior at  $\lambda_{cp}$  is consistent with 3D ising model.



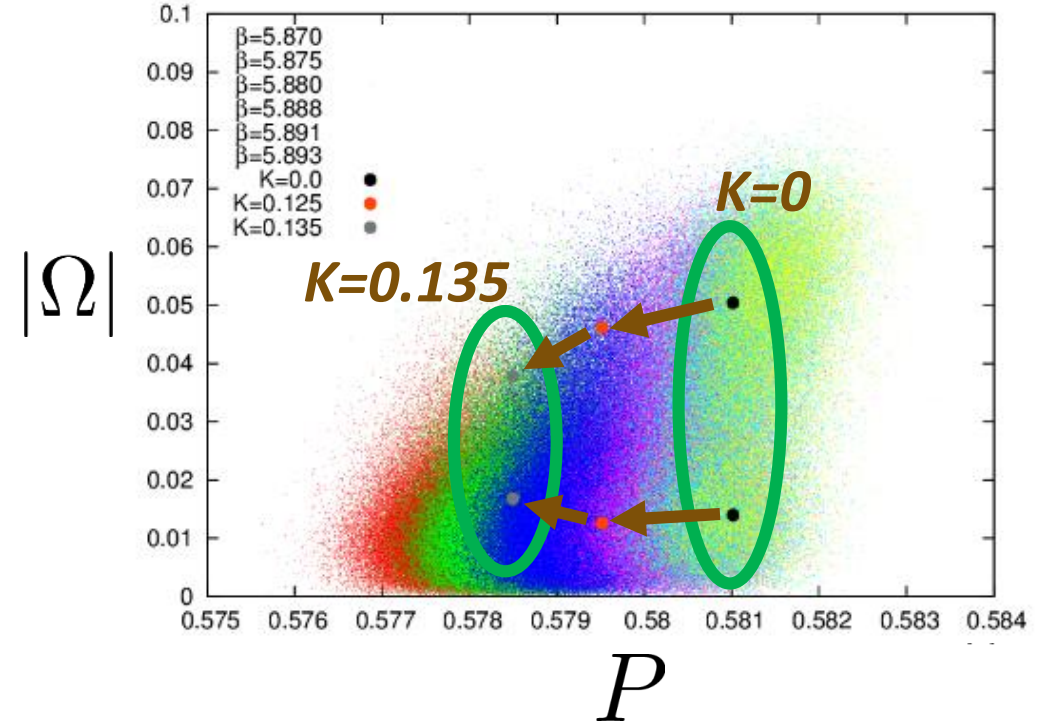
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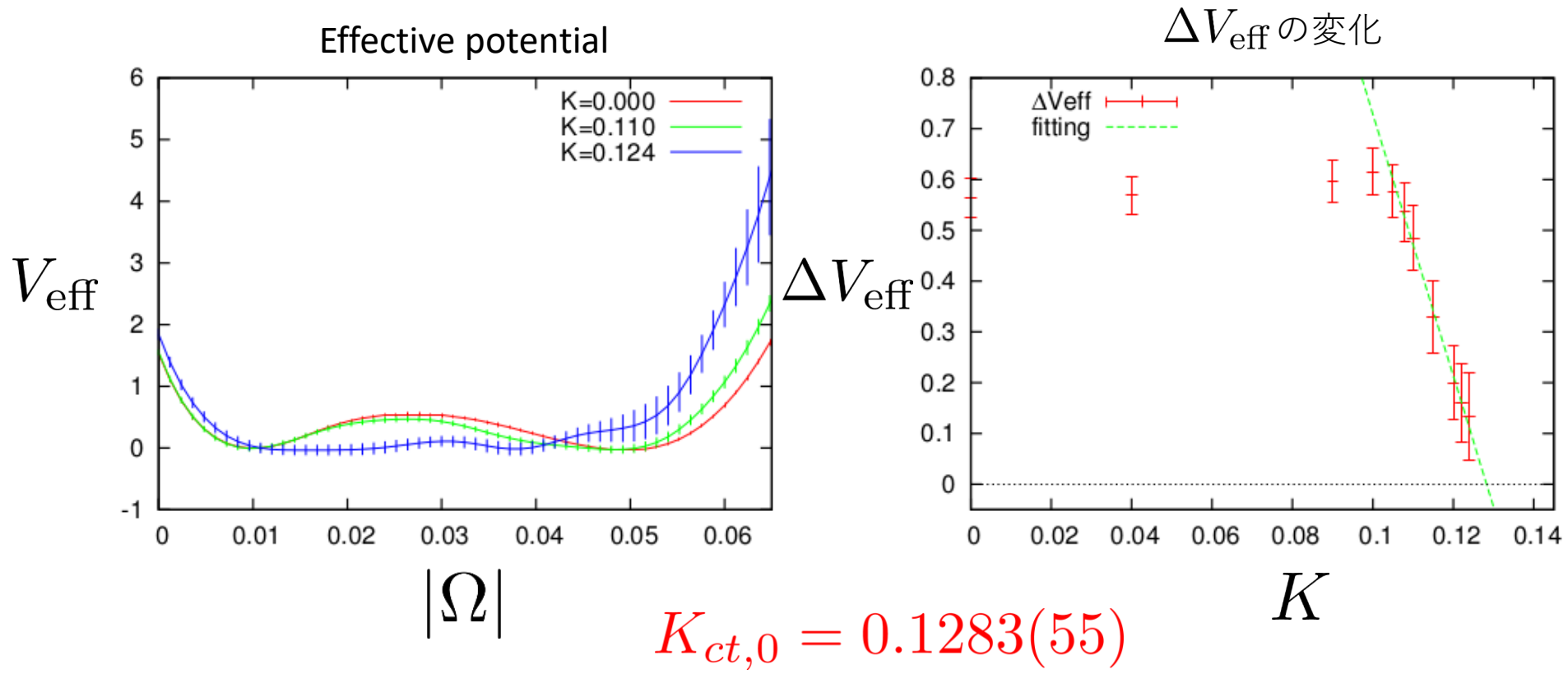
Distribution of  $P$  and  $\Omega$   
Peak positions of the histogram



Lattice	Configurations
$24^3 \times 6$	676,190
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$24^3 \times 8$	342,821

# Determination of $K_c$ (leading order calculation)

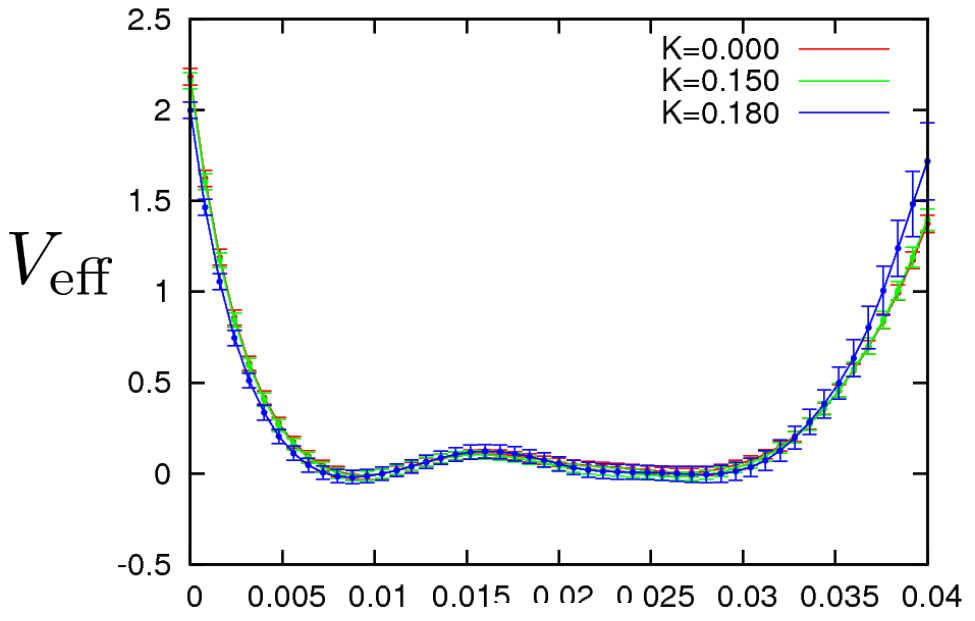
$N_t = 6$  lattice ( $32^3 \times 6$ )



# Determination of $K_c$ (leading order calculation)

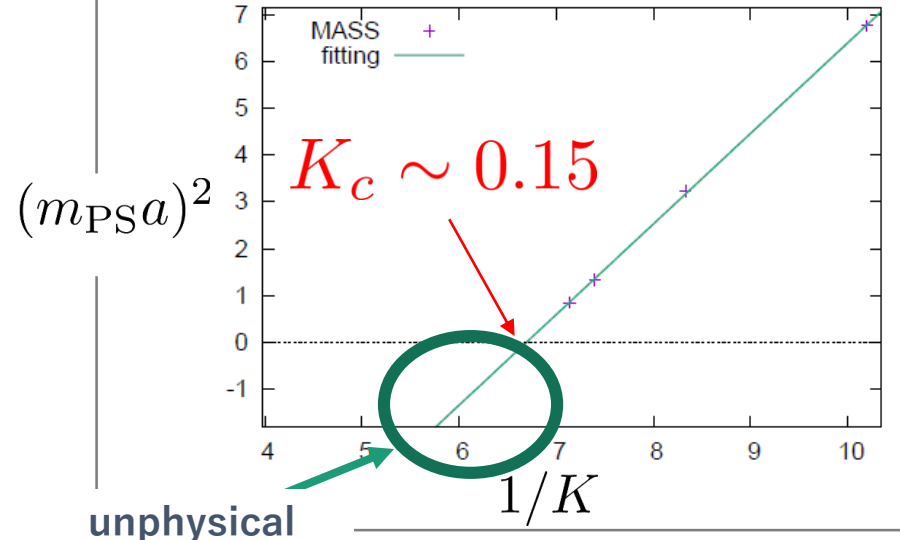
$N_t = 8$  lattice ( $24^3 \times 8$ )

Effective potential



$K_{ct,0} > 0.18$   $|\Omega|$

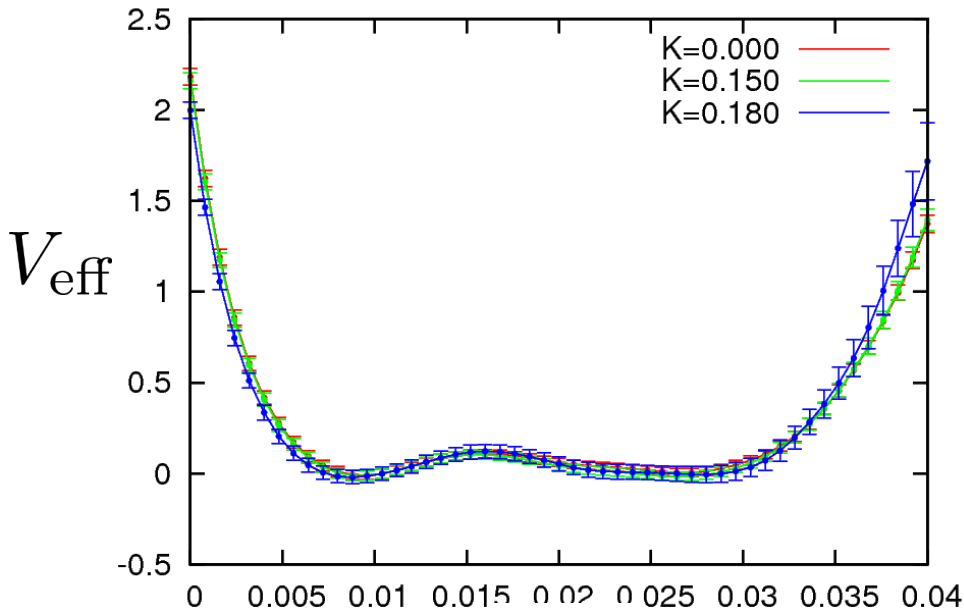
- The shape of the effective potential does not change even when  $K_c > 0.15$



# Determination of $K_c$ (leading order calculation)

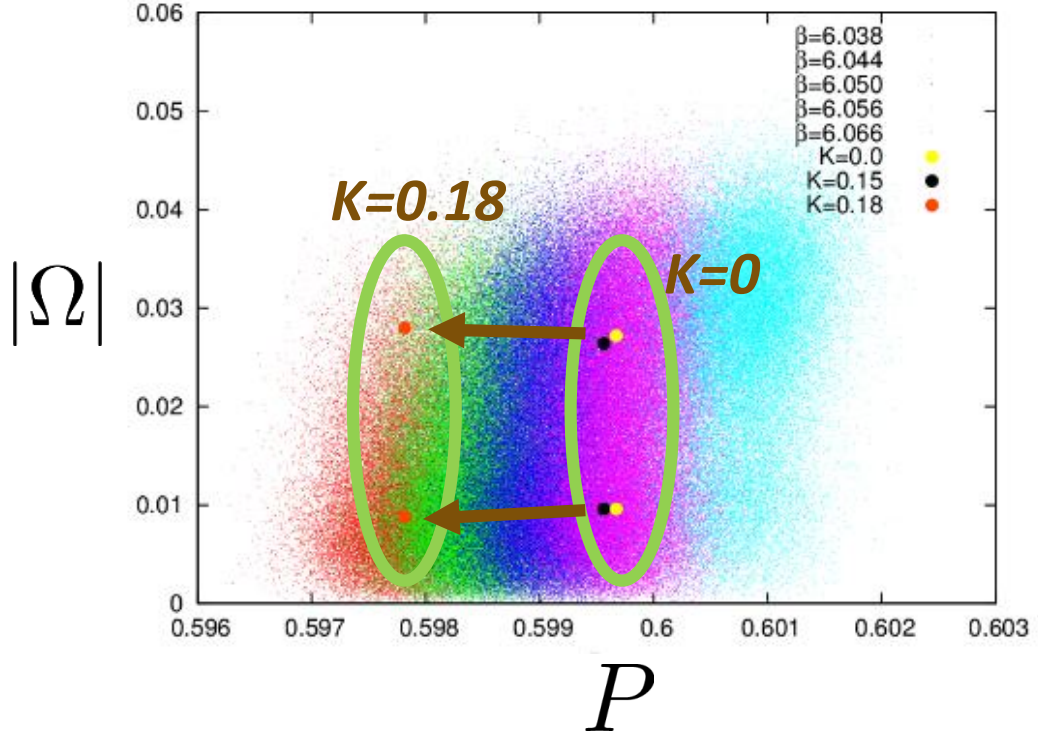
$N_t = 8$  lattice ( $24^3 \times 8$ )

Effective potential



$K_{ct,0} > 0.18 \quad |\Omega|$

Distribution of  $P$  and  $\Omega$   
Peak positions of the histogram

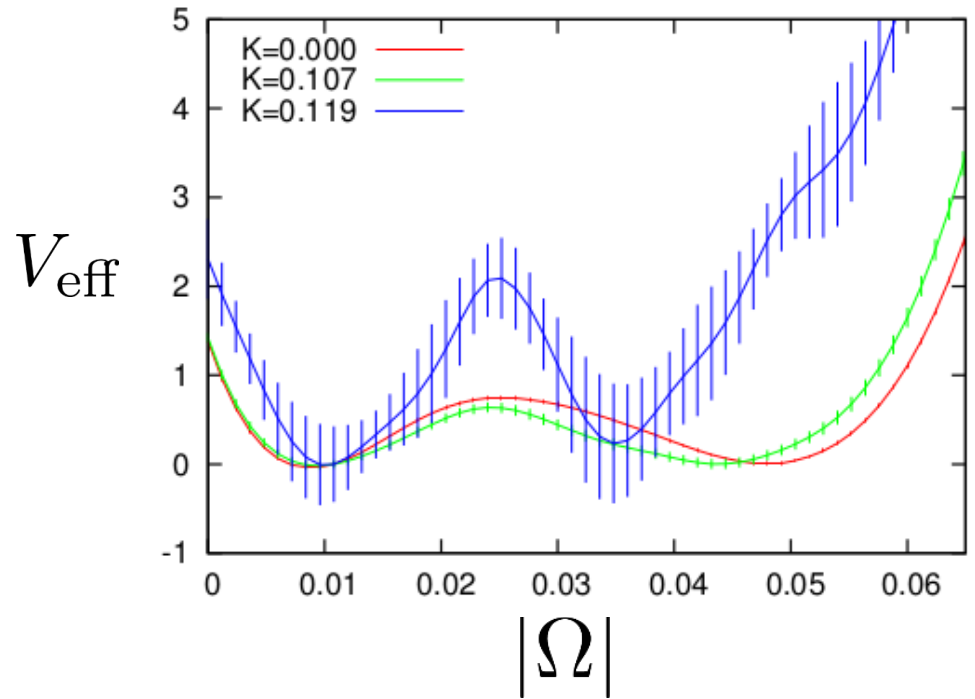


No overlap problem. Truncation error of hopping parameter expansion: large for  $N_t=8$

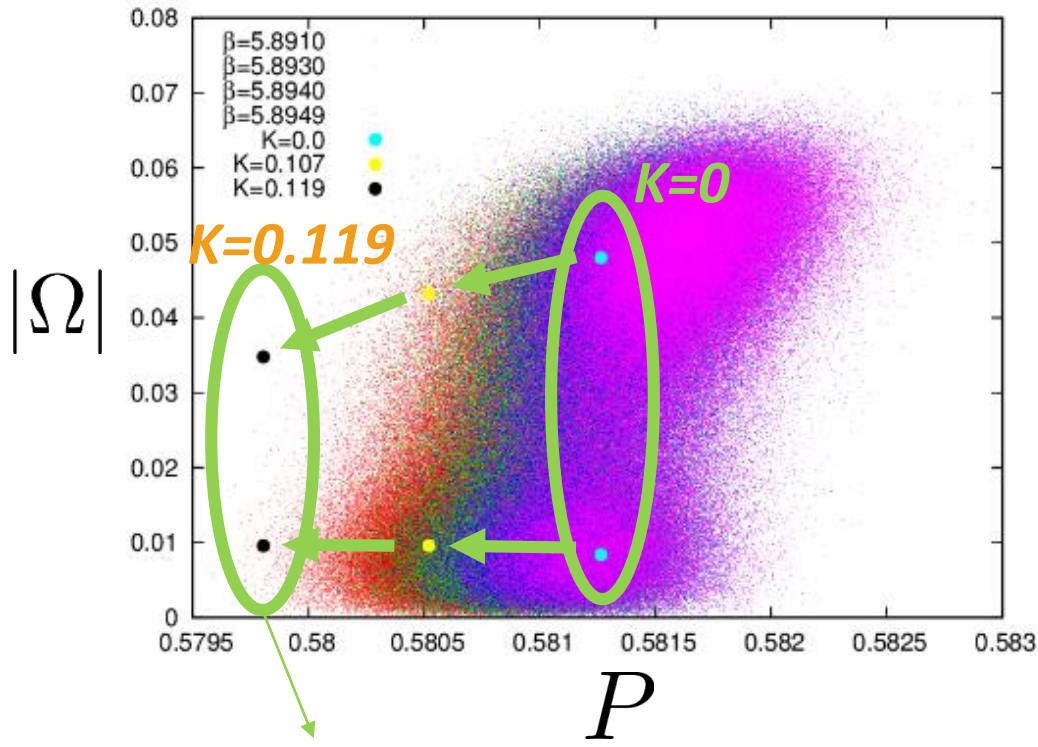
# Determination of $K_c$ (leading order calculation)

$N_t = 6$  lattice ( $36^3 \times 6$ )

Effective potential



Distribution of  $P$  and  $\Omega$   
Peak positions of the histogram



Overlap problem arises.