

# Exploring the QCD phase diagram at finite density by the complex Langevin method on a $16^3 \times 32$ lattice

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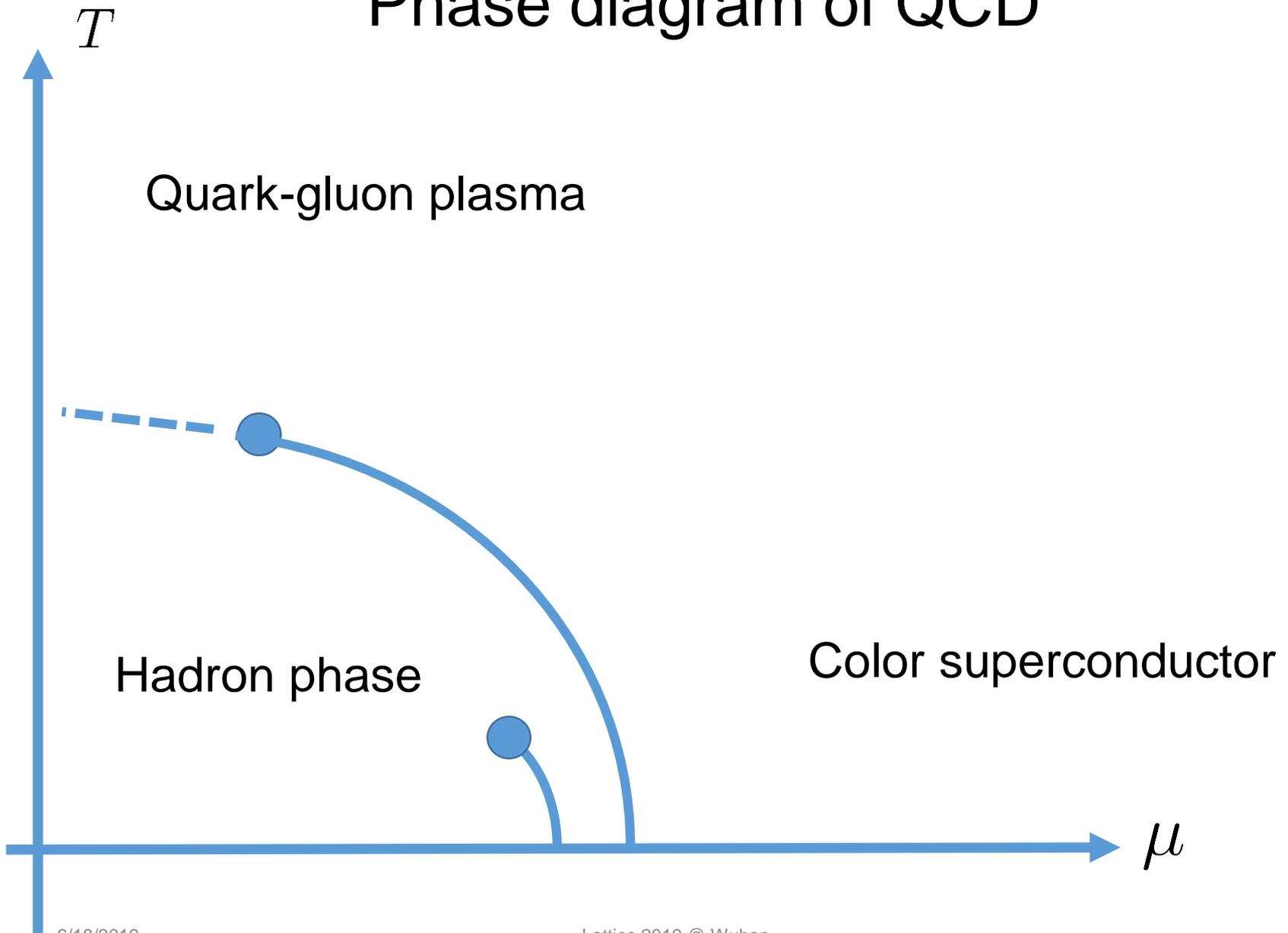
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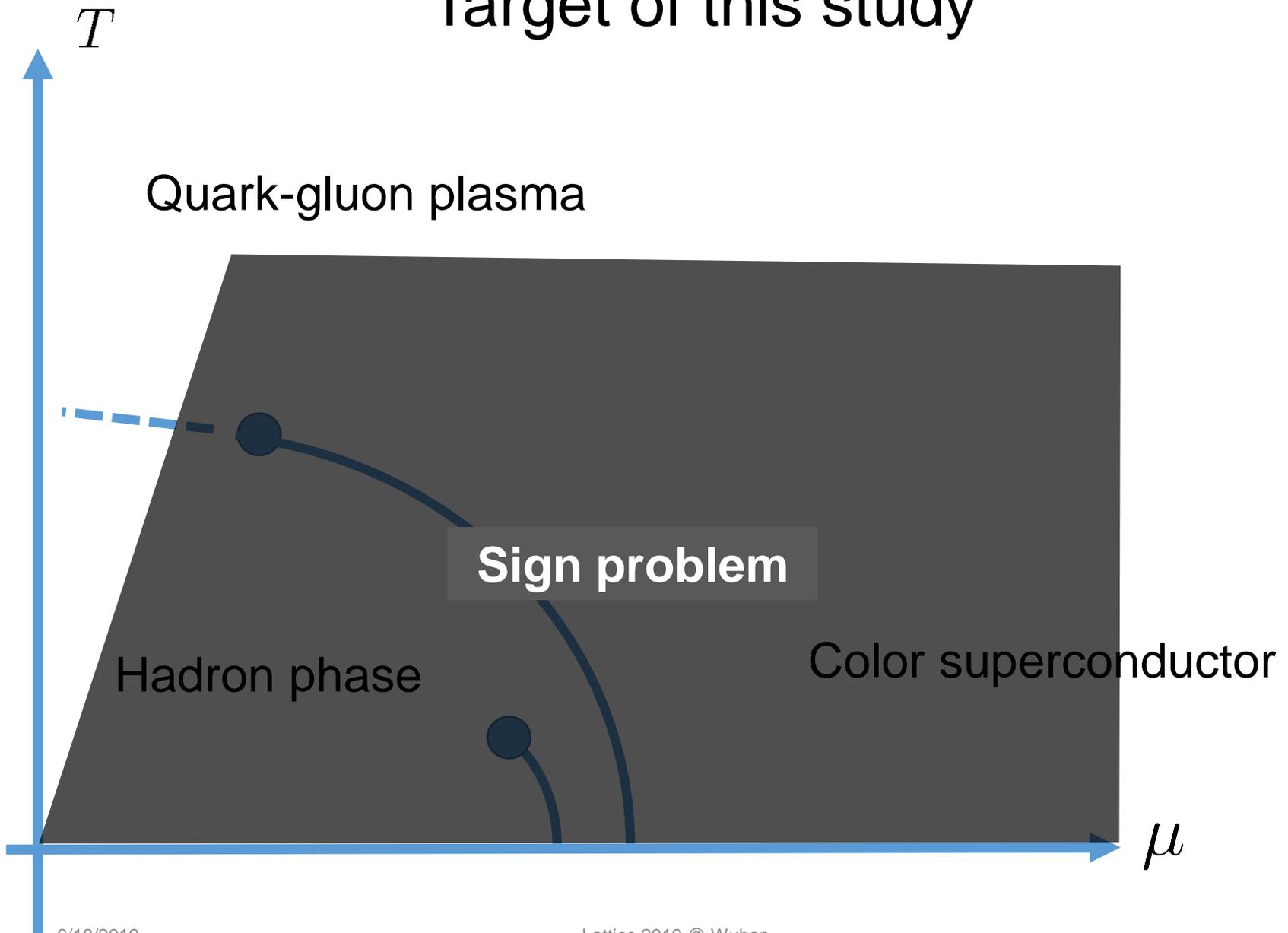
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# Phase diagram of QCD



# Target of this study



# QCD at finite density

Partition function:

$$Z = \int dU \det M[U, \mu] e^{-S_g[U]}$$

$$S_g = -\frac{\beta}{6} \sum_x \sum_{\mu < \nu} \text{tr}(U_{x,\mu\nu} + U_{x,\mu\nu}^{-1})$$

$$M = \gamma_\mu D_\mu + m + \mu \gamma_4$$

The origin of the sign problem:

$$\det(M(\mu)^\dagger) = \det M(-\mu) \quad \dots \text{complex number !}$$

# Complex Langevin Method (CLM)

Complexify the link variable:

$$U_{x\mu} \in SU(3) \rightarrow \mathcal{U}_{x\mu} \in SL(3, \mathbb{C})$$

[Parisi '83], [Klauder '84]  
 [Aarts, Seiler, Stamatescu '09]  
 [Aarts, James, Seiler, Stamatescu '11]  
 [Sexty '14] [Fodor, Katz, Sexty, Torok '15]  
 [Sinclair, Kogut '16]  
 [Nishimura, Shimasaki '15]  
 [Nagata, Nishimura, Shimasaki '15]

Solve the Langevin equation:

$$\mathcal{U}_{x\mu}(t + \epsilon) = \exp \left[ i \left( \underbrace{-\epsilon \mathcal{D}_{x\mu} S[\mathcal{U}]}_{\text{Drift term}} + \underbrace{\sqrt{\epsilon} \eta_{x\mu}(t)}_{\text{Gaussian noise}} \right) \right] \mathcal{U}_{x\mu}(t)$$

Drift term

Gaussian noise

Calculate observables:

$$\langle O(\mathcal{U}) \rangle_{\text{CLM}} = \lim_{T \rightarrow \infty} \frac{1}{T} \int_{t_0}^{t_0+T} dt O(\mathcal{U}) P_{\text{eq}}(\mathcal{U})$$

$\langle \dots \rangle_{\text{CLM}}$

Noise average

# Validity of the CLM 1/2

## (1) Excursion problem

Complexified link variables deviate from SU(3) subspace

## (2) Singular drift problem

The drift term becomes large due to the zero eigenvalues of the Dirac operator

$$v_{x,\nu} = D_{x,\nu} S[\mathcal{U}] \propto \text{tr} \left( M[\mathcal{U}]^{-1} \mathcal{D}_{x,\nu} M[\mathcal{U}] \right)$$

Drift term

$$M = \gamma_\mu D_\mu + m + \mu \gamma_4$$

If one of these happens, CLM  $\neq$  path integral quantization.

# Validity of the CLM 2/2

[Nagata, Nishimura, Shimasaki '15]

We can judge the validity by histograms of the drift term.

Norm of the drift term: 
$$v = \sqrt{\frac{1}{3} \max_{x,\nu} \text{tr}(v_{x,\nu}^\dagger v_{x,\nu})}$$

The histogram shows exponential fall-off: CLM = path integral

The histogram shows power-law fall-off: CLM  $\neq$  path integral

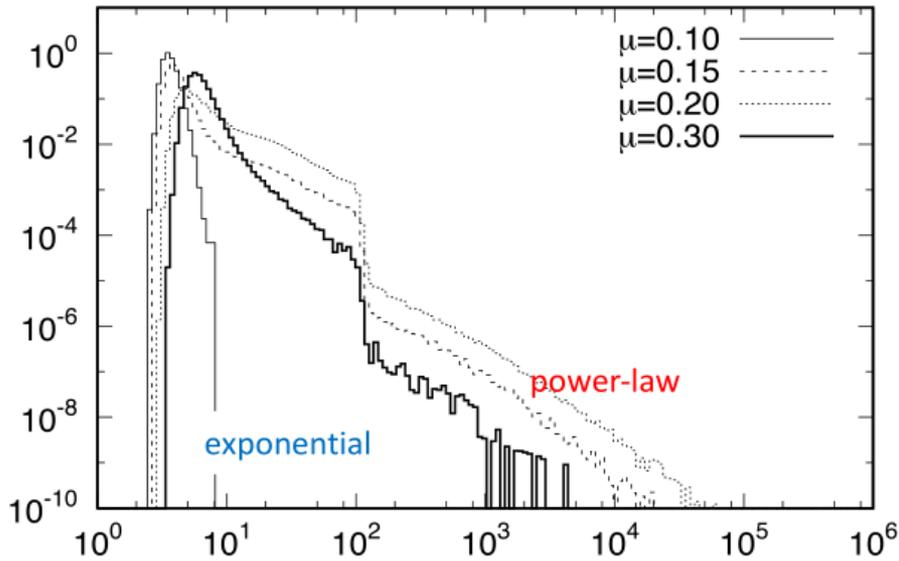
# Lattice setup

- ◆  $N_f = 4$  (staggered fermions),
- ◆ Lattice size:  $8^3 \times 16$ ,  $16^3 \times 32$
- ◆  $\beta = 5.7$  ( $a \sim 0.045$  fm),
- ◆  $ma = 0.01$ ,  $\mu a = 0.1 - 0.5$ .

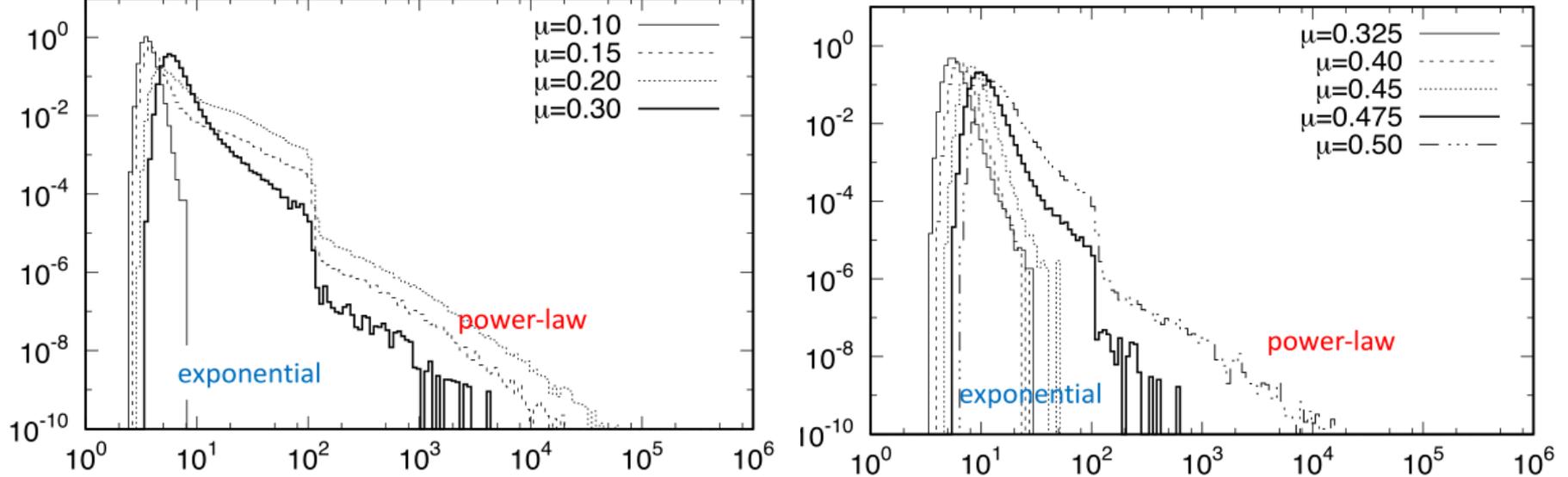
c.f)  $N_f=2$  is studied in [Kogut, Sinclair '19]  
(next talk)

# Histograms of the drift term

$8^3 \times 16$  lattice



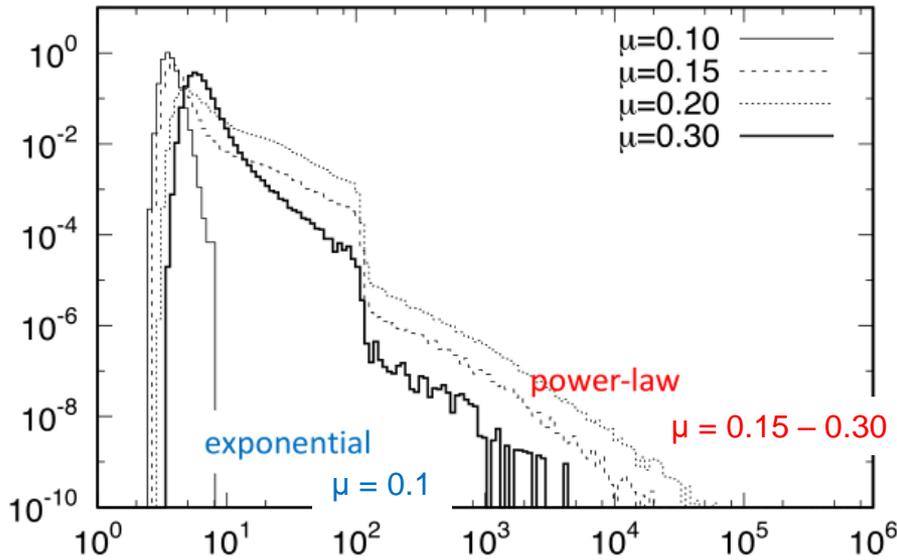
[ST, Ito, Matsufuru, Nishimura, Shimasaki, Tsuchiya '18]



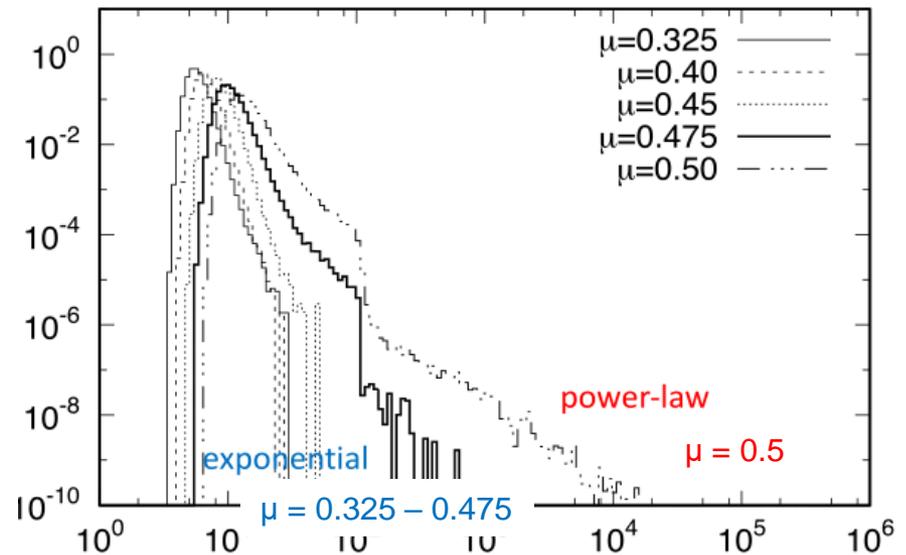
$$\mu a = 0.1 - 0.5$$

# Histograms of the drift term

$8^3 \times 16$  lattice



[ST, Ito, Matsufuru, Nishimura, Shimasaki, Tsuchiya '18]

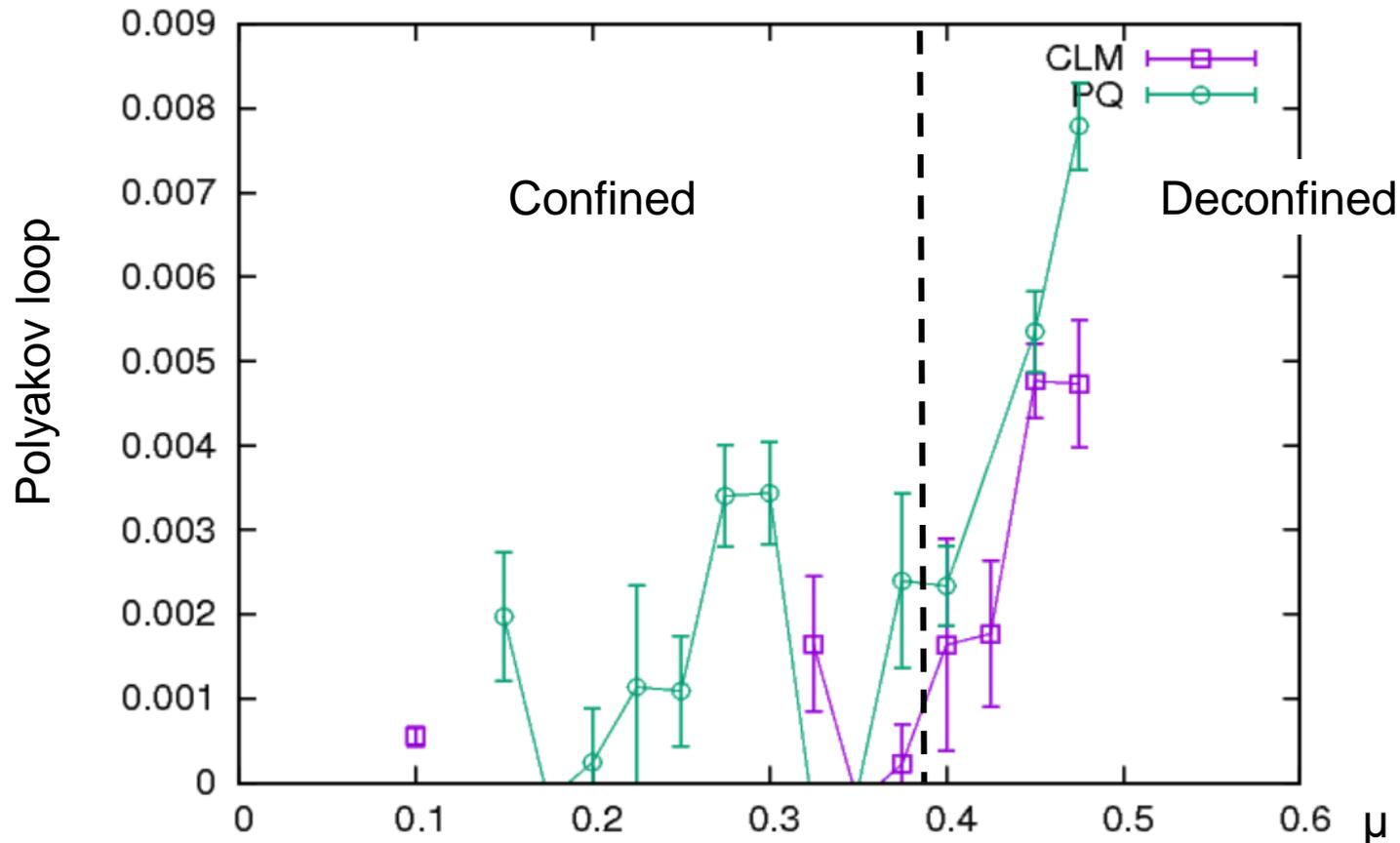


$$\mu a = \boxed{0.1}, \boxed{0.15-0.30}, \boxed{0.325-0.475} - 0.5$$

CLM is justified

# Results on $8^3 \times 16$ lattice 1/2

[ST, Ito, Matsufuru, Nishimura, Shimasaki, Tsuchiya' 18]

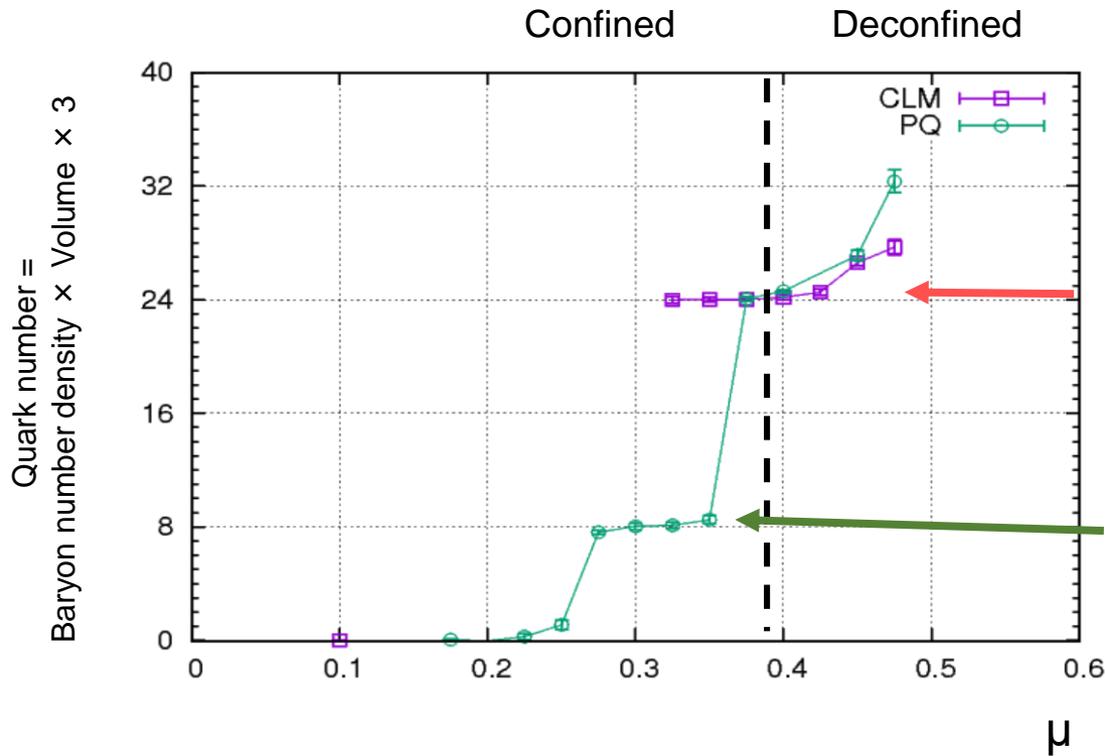


PQ = Phase Quenched  
CLM: only reliable data are shown

Reliable data can be obtained  
in the high density region

# Results on $8^3 \times 16$ lattice 2/2

[ST, Ito, Matsufuru, Nishimura, Shimasaki, Tsuchiya' 18]



CLM predicts a plateau structure in the confined phase, where the quark number is 24.

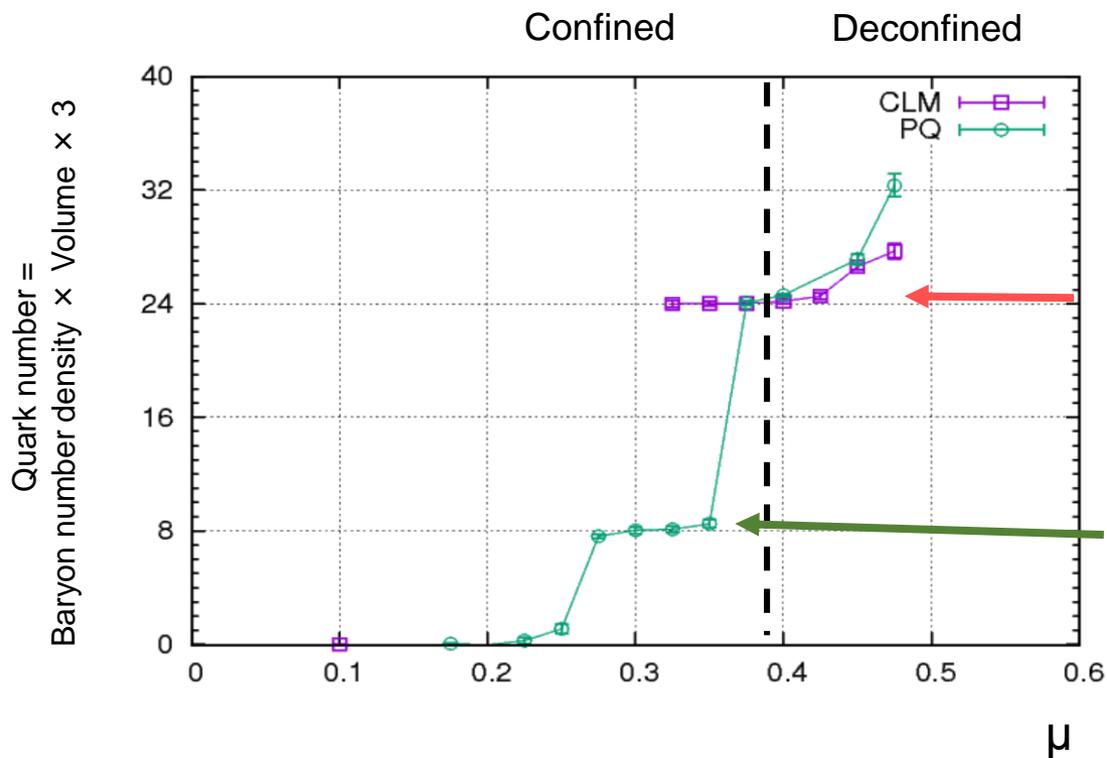


8 baryons are packed in a box.

On the other hand, another plateau found in the PQ simulation has different height. It corresponds to meson creation.

# Results on $8^3 \times 16$ lattice 2/2

[ST, Ito, Matsufuru, Nishimura, Shimasaki, Tsuchiya' 18]



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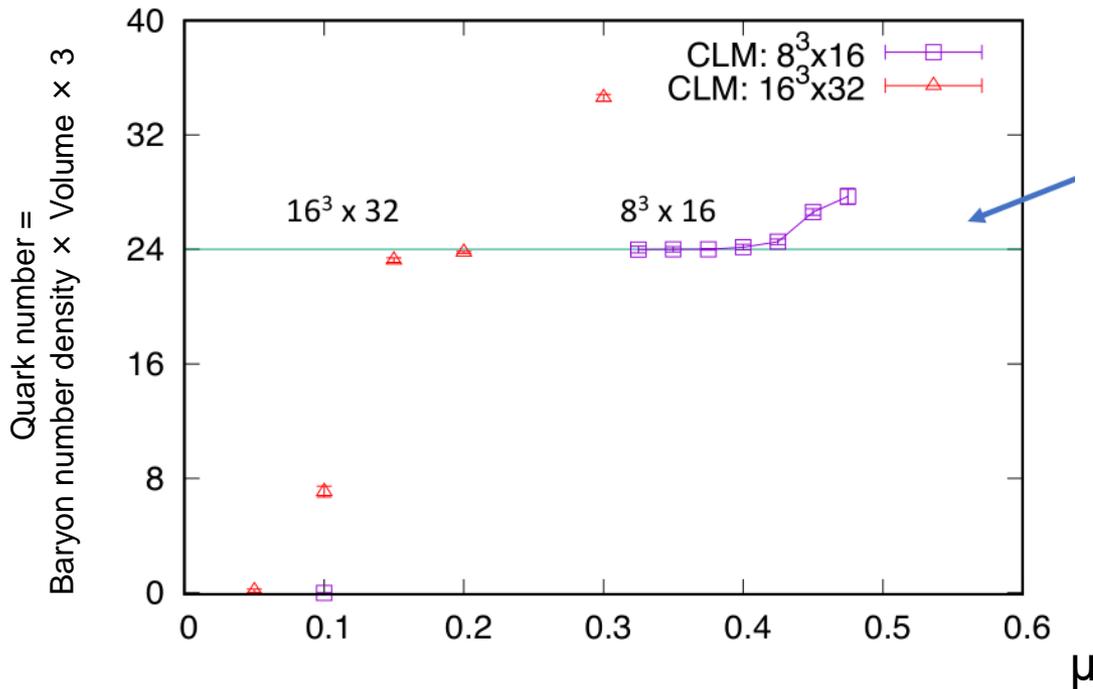


8 baryons are packed in a box.

On the other hand, another plateau found in the PQ simulation has different height. It corresponds to meson creation.

Caveat: finite volume effect is large.  $a \sim 0.045$  fm,  $L = 0.36$  fm,  
 $\mu a = 0.3$  corresponds to  $\mu \sim 1300$  MeV

# Results on $16^3 \times 32$ lattice



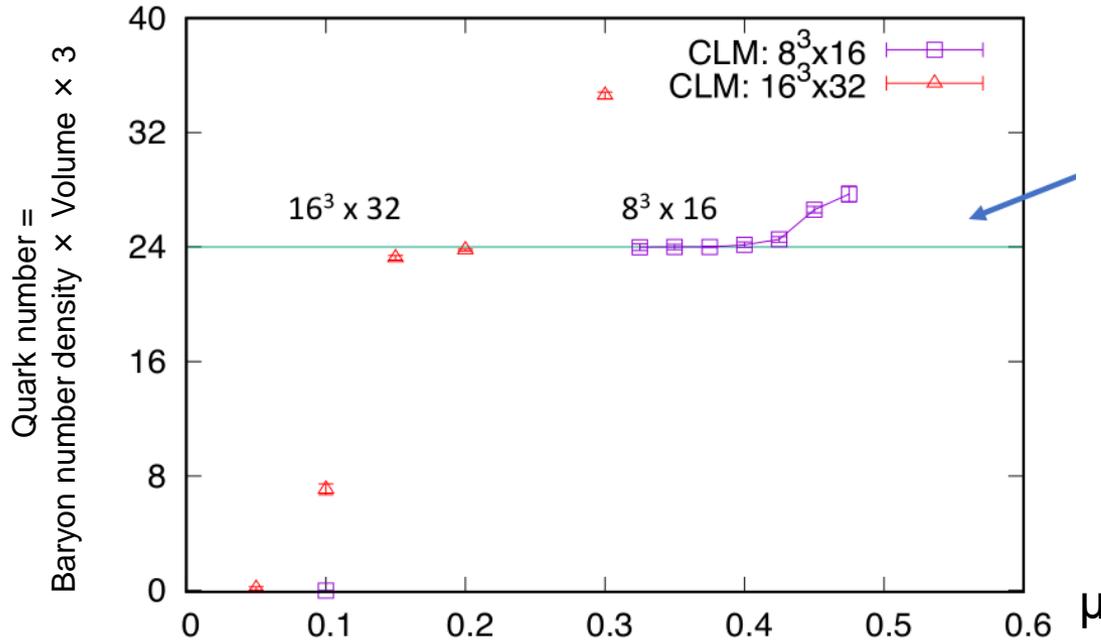
Use same  $\beta$ :  $\beta=5.7$

CLM simulation on  $16^3 \times 32$  lattice also show the plateau structure at  $N_{\text{quark}} = 24$ .

- 8 baryons are created again.
- The critical  $\mu$  is lower than the previous case.

Reduction of the finite volume effect  $\rightarrow$  creation energy per baryon decreases

# Comparison with $\rho_0$



Baryon number densities are

$$\langle n \rangle = 0.016 (8^3 \times 16),$$

$$\langle n \rangle = 0.002 (16^3 \times 32)$$

In physical unit,

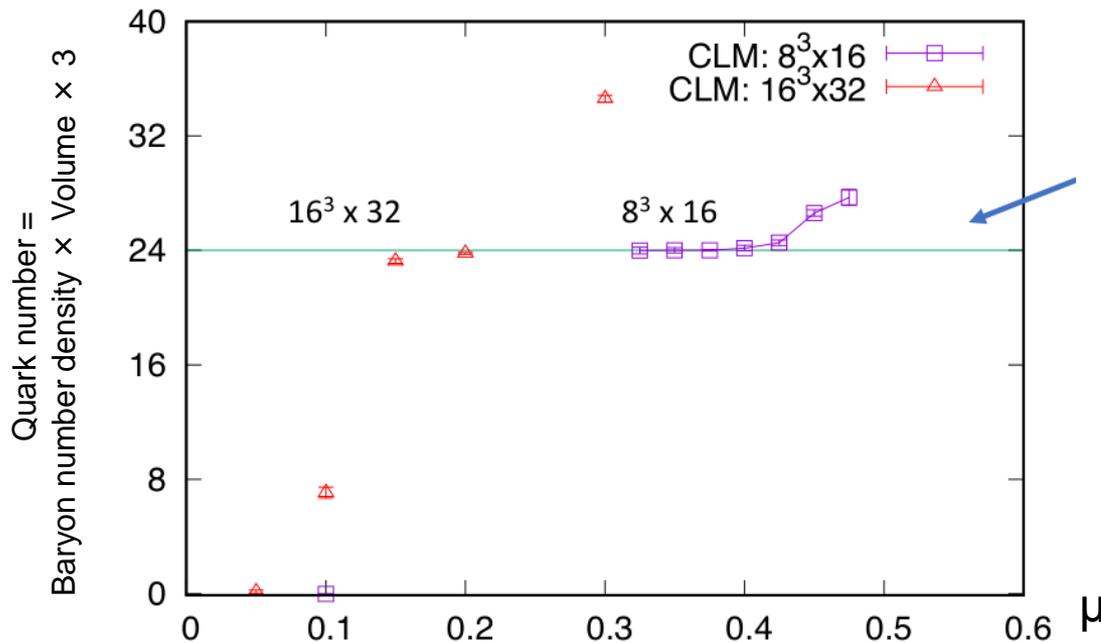
$$\langle n \rangle = 171.5 \text{ fm}^{-3} (8^3 \times 16),$$

$$\langle n \rangle = 21.4 \text{ fm}^{-3} (16^3 \times 32)$$

$$(\beta = 5.7, a \sim 0.045 \text{ fm})$$

$\langle n \rangle$  is much higher than the normal nuclear density  $\rho_0 = 0.17 \text{ fm}^{-3}$

# Increase the lattice size further



Baryon number densities are

$$\langle n \rangle = 0.016 \quad (8^3 \times 16),$$

$$\langle n \rangle = 0.002 \quad (16^3 \times 32)$$

In physical unit,

$$\langle n \rangle = 171.5 \text{ fm}^{-3} \quad (8^3 \times 16),$$

$$\langle n \rangle = 21.4 \text{ fm}^{-3} \quad (16^3 \times 32)$$

$$(\beta = 5.7, a \sim 0.045 \text{ fm})$$

Speculation: quark number remains to be 24 until  $\langle n \rangle$  decreases down to  $\rho_0$  (around  $N_s = 80$ )

- ◆ If the volume is sufficiently large, the quark number should be
 
$$(\text{quark num.}) = 3 \times (\text{Volume}) \times \rho_0$$

$$\sim 4.6 \times 10^{-5} \times N_s^3 \dots \leftarrow \rho_0 = 0.17 \text{ fm}^{-3}, (\text{Volume}) = (N_s a)^3$$
- ◆  $(\text{quark num.}) = 24 \Leftrightarrow N_s^3 = 80^3$

# Summary

- We have extended our our previous work on a  $8^3 \times 16$  lattice to a  $16^3 \times 32$  lattice with the same beta to reduce finite size effects.
- The plateau of the baryon number starts at much lower chemical potential. The energy cost for creating a baryon seems to be reduced considerably due to the increase of lattice size.
- The height of the plateau corresponds to 24 quarks (8 baryons), which turns out to coincide with our previous result on a  $8^3 \times 16$  lattice.
- Eventually, the height should grow linearly with the lattice volume since the nuclear density should be reproduced from the height.
- What happens if we increase the lattice size further?  
Speculation: the height of the plateau remains to be the same baryon number until the density decreases down to the nuclear density (around  $N_s=80$ , due to the chosen large  $\beta=5.7$ ).

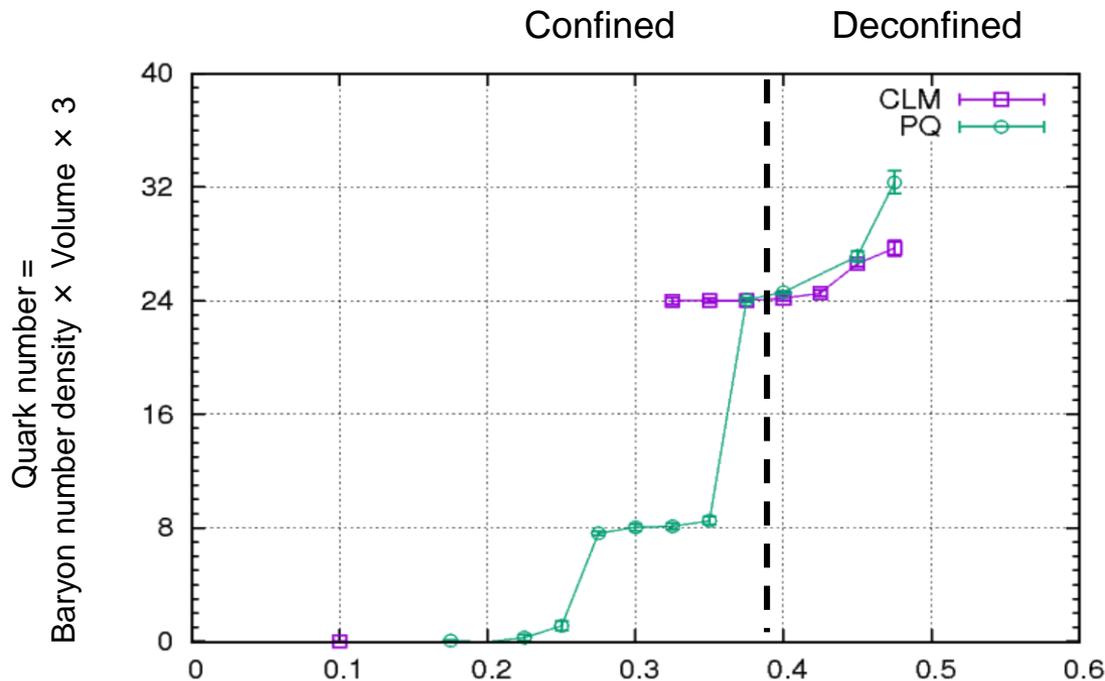
## Future plans:

Increase lattice size, use more finer lattice (larger  $\beta$ ) and study quark matter phase. For instance, can we approach to color-superconductor phase?

# Appendix

# Phase quenched simulation

[ST, Ito, Matsufuru, Nishimura, Shimasaki, Tsuchiya' 18]



In the phase quenched simulation,  $\mu$  is interpreted as the isospin chemical potential due to

$$\det D(-\mu) = (\det D(\mu))^*$$

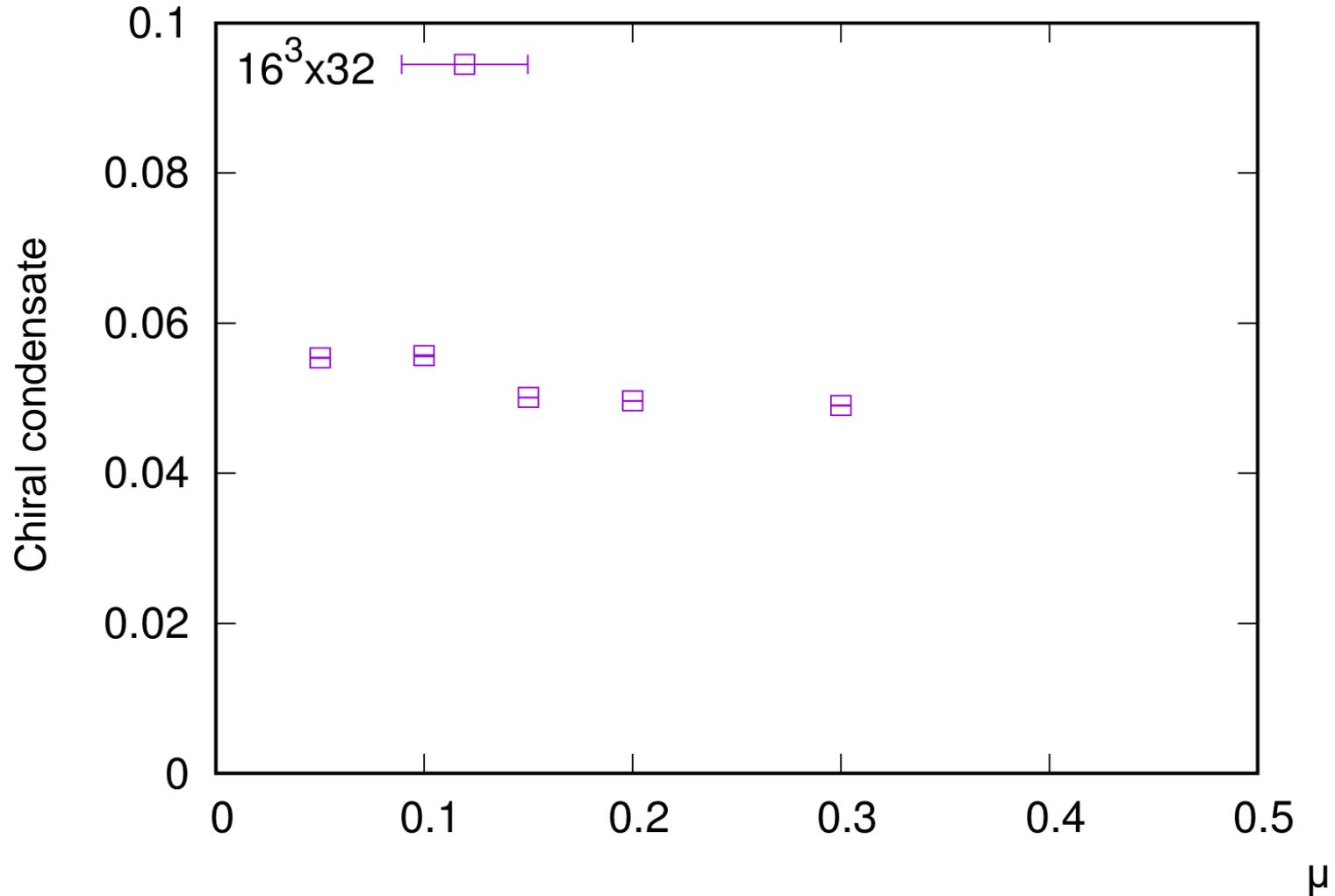
Therefore, the number density represents

$$\mu \quad \langle n \rangle = \langle (n_u - n_d) / 3 \rangle$$

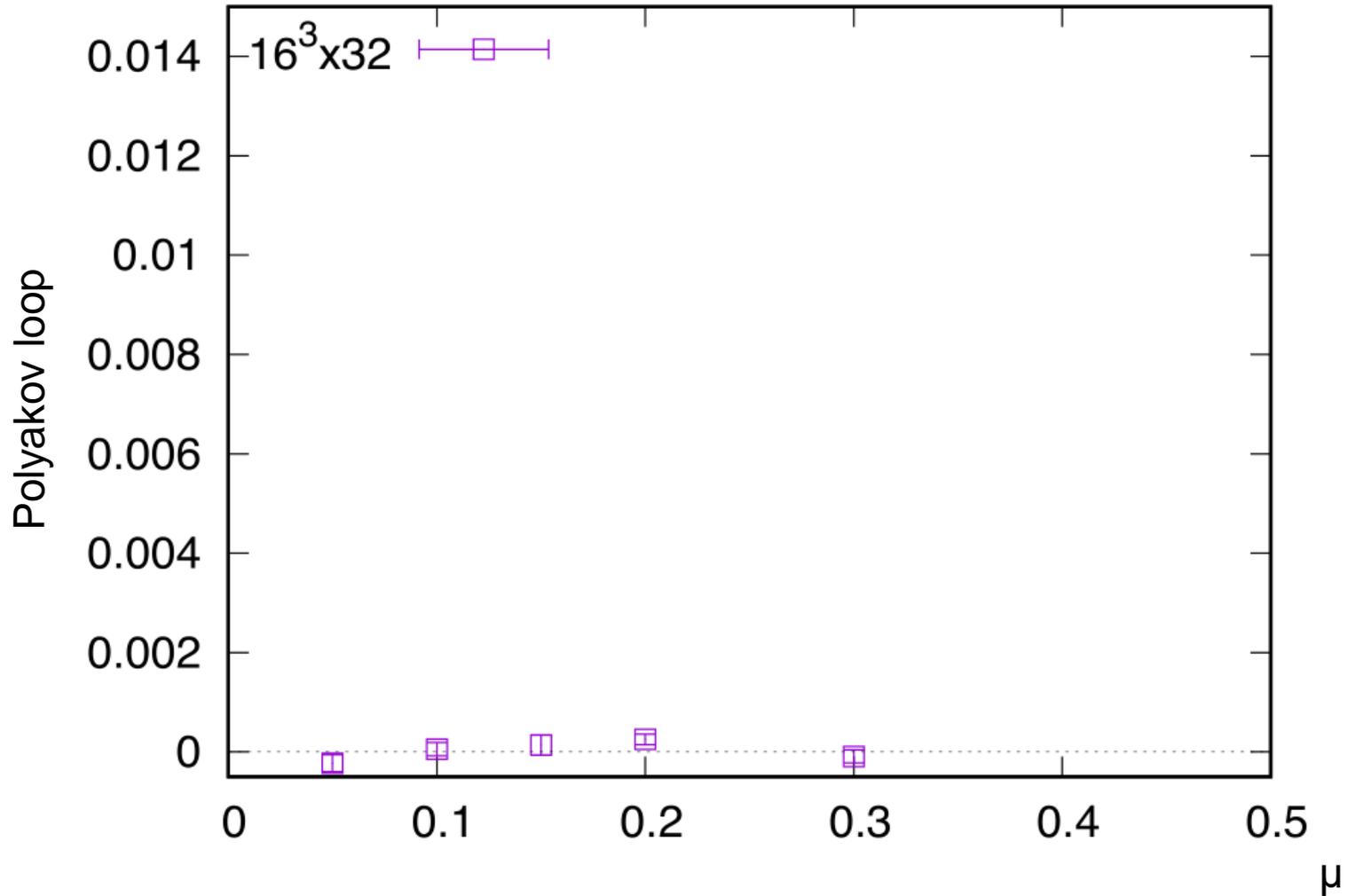
↑                      ↑

Number density of up and down quarks

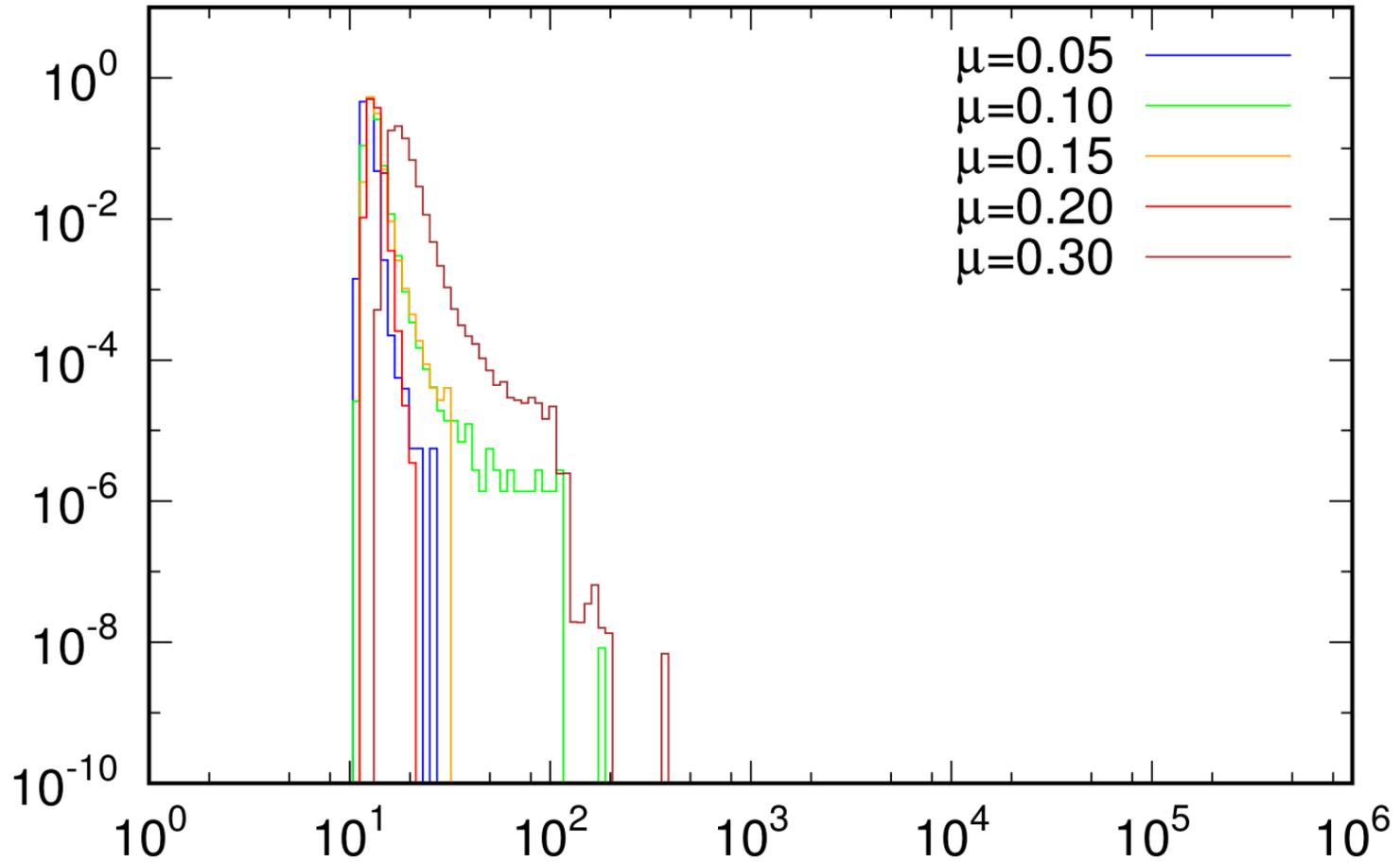
# Results on $16^3 \times 32$ lattice

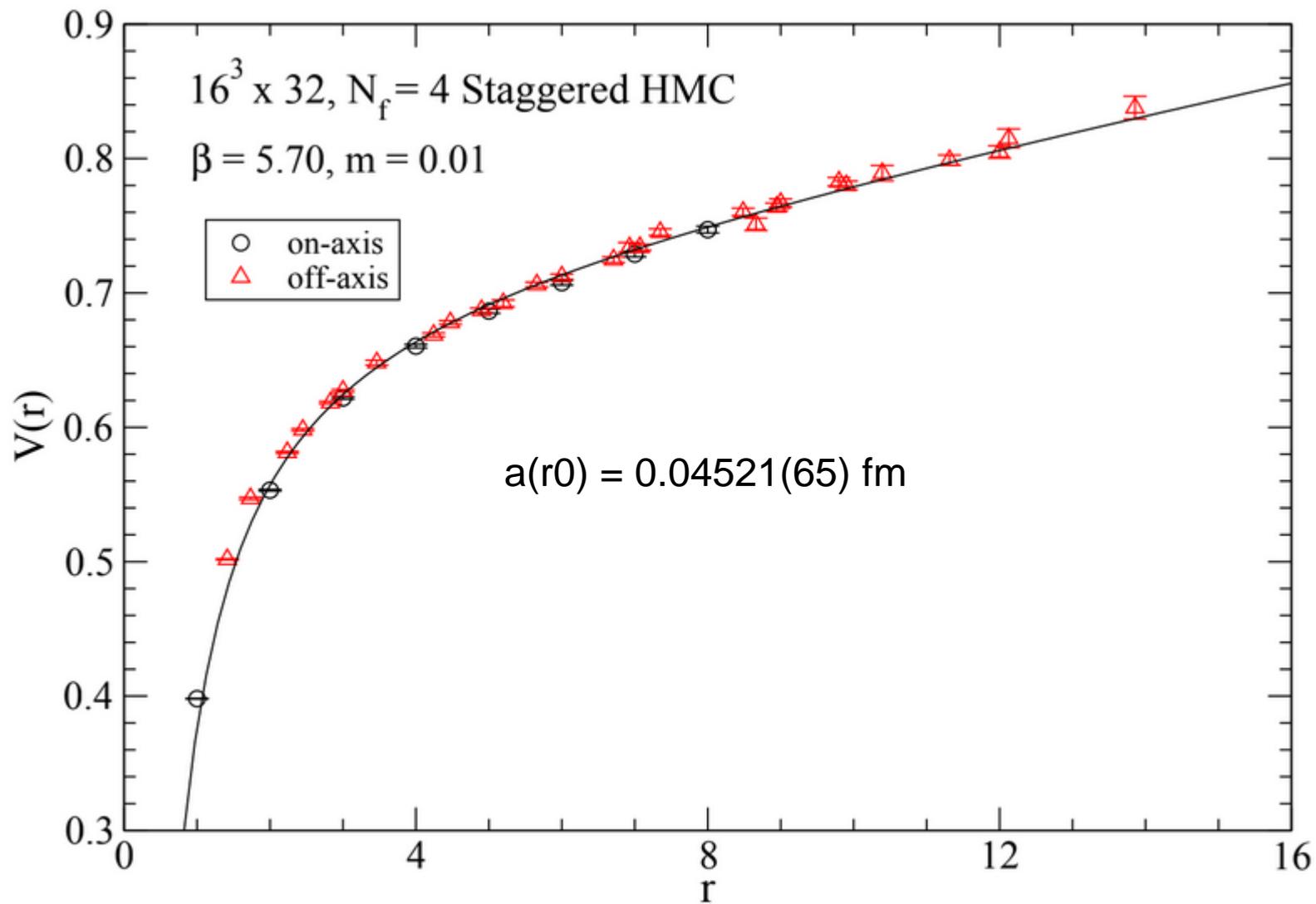


# Results on $16^3 \times 32$ lattice



# Results on $16^3 \times 32$ lattice





# Phase diagram of QCD with 4-flavor staggered fermion

