Evading the model sign problem in the PNJL model with repulsive vector-type interaction via path optimization

Akira Ohnishi¹, Yuto Mori², Kouji Kashiwa³

Yukawa Inst. for Theoretical Physics, Kyoto U.,
 Dept. Phys., Kyoto U., 3. Fukuoka Inst. Tech.

The 37th Int. Symp. on Lattice Field Theory (Lattice 2019), June 16-22, 2019, Wuhan, China







What is the nature of finite density QCD phase transition ?

- First order phase transition boundary may exist at finite Re μ_B
- **First order p.t. boundaries EXIST at \theta=Im \mu_q/T=\pi/3, \pi, 5\pi/3, ...**





Conjectured 3D phase diagram in (T, Re μ_B , θ) space

- RW & finite density transition are CONNECTED
 - → Deconfinement assisted chiral phase transition



 or These two are DISCONNECTED
 → Independent of RW (deconf.) transition TA



Which is true ? PNJL (Polyakov loop extended NJL) model should give answer, but has the sign problem at complex μ_B !

Outline

Introduction

Nature of finite density phase transition and the phase diagram in (T, Re μ_B , θ) space

- Path Optimization Method
 - Variational method & Euler-Lagrange equation for the path
 - **Example of repulsive vector-type interaction: One-site Hubbard model**
- Application to Polyakov loop extended NJL (PNJL) model with repulsive vector-type interaction at real μ
 - Average phase factor and Observables
 - Configurations on Optimized Path
- Summary and Outlook







Complexified variable methods for the sign problem

Lefchetz thimble method

Witten ('10), Cristoforetti et al. (Aurora)('12), Fujii et al. ('13), Alexandru et al. ('16)

- Complex Langevin method *Parisi ('83), Klauder ('83), Aarts et al. ('11), Nagata et al. ('16).*
- Path Optimization method Mori+('17,'18,'19), Kashiwa+('19,'19), AO+('17,'18), Alexandru+('17, '18, '18), Bursa, Kroyter ('18)
 - Integration path is variationally optimized to enhance the average phase factor.





Euler-Lagrange equation for the integral path

Maximizing APF = Minimizing phase quenched partition fn.

$$\begin{aligned} \mathcal{Z}_{pq} &= \int d^N \varphi_R \left| \det \left(\delta_{ij} + i \frac{\partial \varphi_{j,I}}{\partial \varphi_{i,R}} \right) \exp[-S(\varphi_R + i \varphi_I)] \right. \\ &= \int d^N x \left| W(x_i + i y_i, \partial_i y_j) \right| \ (\varphi = x + i y(x)) \end{aligned}$$

Stationary condition of Z_{pq} **w.r.t.** $y(x) \rightarrow Euler-Lagrange eq.$

$$\frac{\delta}{\delta y_j} \mathcal{Z}_{pq} = 0 \to \left[\frac{\partial}{\partial x_i} \frac{\partial}{\partial (\partial_i y_j)} - \frac{\partial}{\partial y_j} \right] |W(x + iy, \partial y)| = 0$$

One variable Euler-Lagrange equation

$$\ddot{y} = (1 + \dot{y}^2)^2 \left[\frac{\partial (\mathrm{Im}S)}{\partial x} + \frac{\dot{y}}{1 + \dot{y}^2} \frac{\partial (\mathrm{Re}S)}{\partial x} \right] \quad (\dot{y} = dy/dx, \\ \ddot{y} = d^2y/dx^2)$$



Example of repulsive vector-type interaction

■ One-Site Hubbard model (strong coupling limit → Hopping term=0) $S = Un_{\uparrow}n_{\downarrow} - \mu(n_{\uparrow} + n_{\downarrow}) \quad (n_i = \psi_i^{\dagger}\psi_i)$

Path integral representation *Tanizaki*, *Hidaka*, *Hayata* ('16)

$$\mathcal{Z} = \sqrt{\frac{\beta U}{2\pi}} \int d\varphi \left[1 + \exp(\beta U(i\varphi + \mu/U + 1/2))\right]^2 \exp\left[-\beta U\varphi^2/2\right]$$
Complex !

Cancellation among multi-thimbles,
and # of thimbles increases with $\beta = 1/T$

$$\stackrel{\Theta}{=}$$

$$\stackrel{0.6}{0.58}_{0.54}_{0.52}_{0.54}_{0.52}_{0.5}_{0.48}_{0.52}_{0.2}_{0.2}_{0.4}_{0.6}_{0.6}_{0.8}_{$$

Example of repulsive vector-type interaction





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PNJL model

Polyakov-loop extended Nambu–Jona-Lasinio model

$$\mathcal{L}_E = \bar{q} \left(\mathcal{D}(\Phi, \bar{\Phi}) + m_0) q - G \left[(\bar{q}q)^2 + (\bar{q}i\gamma_5 \boldsymbol{\tau}q)^2 \right] + G_v (\bar{q}\gamma_\mu q)^2 + \mathcal{V}_g(\Phi, \bar{\Phi})$$

Hubbard-Stratonovich transformation

$$\mathcal{L}_{\text{eff}} = \bar{q} \left[\mathcal{D} + m_0 \right] q - 2G \left[\bar{q}\sigma q + \bar{q}i\gamma_5 \boldsymbol{\pi} \cdot \boldsymbol{\tau} q \right] + \mathcal{V}_g(\Phi, \bar{\Phi}) + G(\sigma^2 + \boldsymbol{\pi}^2) + 2iG_v \omega_4 \bar{q}\gamma_4 q + G_v \omega_4^2 \right]$$

Model sign problem arises from Polyakov loop & Vector field \rightarrow **Ansatz !**

 $\omega_4 = -i\rho_q$

- CK symmetry ansatz $ImA_4^3 = 0, ReA_4^8 = 0$ Nishimura, Ogilvie, Pangeni ('14, '15)
- Vector field (MF)

Are these ansatz justified ?



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Path Optimization in PNJL

Truncation of aux. field only with k=0 (Homogeneous field ansatz)

$$\mathcal{Z} = \int \prod_{\boldsymbol{k}} dz_{\boldsymbol{k}} e^{-\Gamma(z)} \simeq \int dz_{\boldsymbol{0}} e^{-\Gamma(z_{\boldsymbol{0}})}, \quad \Gamma = \beta V \mathcal{V}_{\text{eff}} = \frac{k}{T^4} \mathcal{V}_{\text{eff}}$$

z=auxiliary fields & gauge field Cristoforetti, Hell, Klein, Weise ('10)

Variables (7 dyn. + 3 dep.) $x = (\sigma, \pi^{0,+,-}, \operatorname{Re}A_3, \operatorname{Re}A_8, \operatorname{Re}\omega_4)$ $y = (\operatorname{Im}A_3, \operatorname{Im}A_8, \operatorname{Im}\omega_4)$ (Complexified)

Path Optimization

- HMC for x (H=Re S) \rightarrow 80k configs.
- Mono hidden layer neural network





Observables

- ullet ullet dependence of order parameters $\,\sigma, {
 m Im}\,\,\omega_4, \Phi, ar \Phi$
 - Rapid change around μ = 370 MeV (transition region)
 - Results agree with MF results under ansatz in the large space-time volume region
 — Supports these ansatz





Configurations

Obtained configurations after training the neural network

• Configs. are well localized at around Re $A_8 = 0$, Re $\omega_4 = 0$, Re $A_3 \neq 0$ \rightarrow Confirms CK symmetry and standard MF ansatz

 \mathcal{CK} symmetry ansatz: $(\theta_1, \theta_2, \theta_3) = (\theta - i\psi, -\theta - i\psi, 2i\psi)$ MF ansatz: $\omega_4 = i\rho_q$, Re $\omega_4 = 0$





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Summary

- QCD phase diagram at complex μ should be useful to understand the nature of QCD phase transition. Partition fn. is an holomorphic fn. of μ with some cuts, with cuts being the 1st order p.t. boundary.
- Path optimization method is a kind of Jacobian-phase improved Lefshetz thimble method, and is flexible enough to cover the multi-thimble manifold.
 - Do not care too much about # of thimbles. Care more about integrating wide enough range in the complexified field variables.
 - Euler-Lagrange eq. for the path is derived, and the variational path is confirmed to agree with the solution of EL equation in the one-site Hubbard model.
- Polyakov-loop extended Nambu–Jona-Lasino model with repulsive vector-type interaction is studied in the path optimization method.
 - CK sym. ansats (gluons) and mean field (vector field) ansatzs' have been confirmed in the path integral formulation. The latter is done for the first time.
 - Ready to study QCD phase diagram in PNJL at complex μ.



Thank you for your attention !

Collaborators: Akira Ohnishi¹, Yuto Mori², Kouji Kashiwa³ 1. Yukawa Inst. for Theoretical Physics, Kyoto U., 2. Kyoto U., 3. Fukuoka Inst. Tech.







AO (11 yrs ago)

Y. Mori (grad. stu.)

K. Kashiwa (main contributor in PNJL)

1D integral: Y. Mori, K. Kashiwa, AO, PRD 96 ('17), 111501(R) [arXiv:1705.05605] $\varphi 4 w/NN$: Y. Mori, K. Kashiwa, AO, PTEP 2018 ('18), 023B04 [arXiv:1709.03208] Lat 2017: AO, Y. Mori, K. Kashiwa, EPJ Web Conf. 175 ('18), 07043 [arXiv:1712.01088] NJL thimble: Y. Mori, K. Kashiwa, AO, PLB 781('18),698 [arXiv:1705.03646] PNJL w/NN: K. Kashiwa, Y. Mori, AO, PRD 99 ('19), 014033 [arXiv:1805.08940] PNJL w/ Vector + NN: K. Kashiwa, Y. Mori, AO, PRD 99 ('19) 114005 [arXiv:1903.03679] 0+1D QCD: Y. Mori, K. Kashiwa, AO, arXiv:1904.11140 1+1D QCD: Y. Mori, K. Kashiwa, AO, in prep., 0+1D Hubbard: AO, Mori, Kashiwa, in prep. **One-site Hubbard model**





One-site Hubbard model





One-site Hubbard model



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Extended Hubbard-Stratonovich Transformation

$$e^{\alpha AB} = \int d\varphi \, d\phi \, e^{-\alpha \left\{ (\varphi - (A+B)/2)^2 + (\phi - i(A-B)/2)^2 \right\} + \alpha AB}$$
$$= \int d\varphi \, d\phi \, e^{-\alpha \left\{ \varphi^2 - (A+B)\varphi + \phi^2 - i(A-B)\phi \right\}} \,. \tag{44}$$

Miura, Nakano, AO, Kawamoto, PRD80 (2009) 074034

QCD phase diagram

