

Faculty of Physics



Euclidean correlation functions of the topological charge density

Lukas Mazur

Collaborators: Olaf Kaczmarek, Hai-Tao Shu, Luis Altenkort

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Motivation

- The classical QCD vacuum is not unique. QCD possesses infinitely many classical vacua which can be classified by their topological properties.
- The change of the topological charge is given by the sphaleron rate.
- It is related to the fermion number non-conservation through axial anomaly.
- Perturbative calculations exist but are not reliable in the relevant coupling region.
- Non-perturbative calculations required from lattice QCD.

Sphaleron rate

The sphaleron rate is defined in Minkowski time by the correlation function: [Moore, Tassler, 2010]

$$\Gamma_{sphal} = \int d^4x \left\langle \frac{g^2}{32\pi^2} F^a_{\mu\nu} \tilde{F}^a_{\mu\nu}(x) \frac{g^2}{32\pi^2} F^a_{\mu\nu} \tilde{F}^a_{\mu\nu}(0) \right\rangle$$

 Euclidean correlation function has same spectral function and is defined by

$$\begin{split} G_q(\tau,\vec{p}) &= \int d^3x \ e^{i\vec{p}\vec{x}} \langle q(\vec{x},\tau)q(\vec{0},0) \rangle \\ G_q(\tau,\vec{p}) &= \int_0^\infty \frac{d\omega}{\pi} \rho_q(\omega,\vec{p}) \frac{\cosh\omega\left(\frac{1}{2T}-\tau\right)}{\sinh\frac{\omega}{2T}} \\ \Gamma_{sphal} &= \lim_{\omega \to 0} \frac{2T\rho_q(\omega,0)}{\omega}. \end{split}$$

Extracting ρ_q is ill-posed inversion problem. See talk of Anna-Lena Kruse and Hiroshi Ohno.

Topological charge

Gluonic definition of the topological charge density:

$$q(x) = rac{g^2}{32\pi^2} \epsilon_{\mu
u
ho\sigma} \mathrm{tr} \left\{ F_{\mu
u}(x) F_{
ho\sigma}(x)
ight\}.$$

We use Highly-improved field-strength tensor on the lattice. [Bilson-Thompson,Leinweber,Williams,2002]

Topolocial charge:

$$Q=\int d^4x\,q(x).$$

Topological susceptibility:

$$\chi_{ ext{top}} = \lim_{V o \infty} rac{\langle Q^2
angle}{V} = \int d^4 x \, \langle q(x) q(0)
angle.$$

The gradient flow

(

- All gluonic definitions valid only on smooth configurations.
- Introduce extra coordinate t (flowtime) and define a t-dependent gauge field B_µ(x, t) [Lüscher,2014]

$$egin{aligned} &rac{d}{dt}B_\mu(x,t)=D_
u}G_{
u\mu}(x,t),\ &D_\mu=\partial_\mu+[B_\mu(x,t),\ \cdot\]\,,\ &G_{\mu
u}(x,t)=\partial_\mu B_
u(x,t)-\partial_
u}B_\mu(x,t)+[B_\mu(x,t),B_
u(x,t)] \end{aligned}$$

- with the initial condition: $B_{\mu}(x, t)|_{t=0} = A_{\mu}(x)$.
- The flow smoothens the fields over a region of radius $\sqrt{8t}$.
- We use Symanzik improved gradient flow on the lattice. [Ramos, Sint,2015]

Setup

- SU(3) pure gauge configurations were generated using the quenched approximation.
- One sweep consists of one heat bath and four overrelaxation steps.
- Temperature $T/T_c = 1.5$.
- All configurations are separated by 500 sweeps.

Ns	$N_{ au}$	β	<i>a</i> [fm]	#conf
64	16	6.873	0.026045	10000
80	20	7.035	0.021342	10000
96	24	7.192	0.017650	7892

At what flowtime should be start?

q only valid on sufficiently smoothed configurations.

• χ_t shows when UV-fluctuations are smoothed out.



Is topological tunneling sufficient?

▶ Q gets frozen with finer lattices \Rightarrow Does that affect q(x)?





• At $\tau T = 0 \Rightarrow G(\tau) > 0$ and at $\tau T > 0 \Rightarrow G(\tau) < 0$.

Small cut-off effects are seen.

Is the $G(\tau)$ stable under flow?

Expectation:

- UV-part of correlation function $0 \le \tau \le \sqrt{8t}/N_{\tau}$ gets destroyed.
- ► Too much smoothing destroys the IR-part as well.

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A plateau should be seen in the flowtime dependence of the correlation function $G_{\tau}(t)$.



- Low flowtime area $< 0.15 N_{\tau}$ affected by UV fluctuation.
- Signal in high flowtime area $> 0.3N_{\tau}$ gets destroyed.
- Plateaus are seen in between.



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Conclusion

- Gradient flow reveals $\langle q(x)q(0)\rangle$ correlation function.
- Correlation function shows small cut-off effects.
- Flow smoothens out $\approx 2\sqrt{8t}$ of the correlation function. \Rightarrow Large N_{τ} are needed.

Outlook:

- Continuum limit of correlation function. This might allow us to use even more distances.
- Compare with perturbation theory.
- Obtain correlation function from full QCD (HISQ).