



Real-Time-Evolution of Heavy-Quarkonium Bound States

Alexander Lehmann

Department of Mathematics and Physics University of Stavanger

> Institute for Theoretical Physics Heidelberg University

work in collaboration with A. Rothkopf

The 37th International Symposium on Lattice Field Theory – June 19th 2019 – Wuhan – PRC

Heavy Quarkonia in Heavy Ion Collisions





Heavy Quarkonia in Heavy Ion Collisions





- Heavy Quarkonia (Charmonium, Bottomonium) are well controlled experimental and theoretical probes for the quark-gluon-plasma
- Phenomenological models describe quarkonium supression via a Schrödinger equation + assumption of early formation of bound states

Heavy Quarkonia in Heavy Ion Collisions





- Heavy Quarkonia (Charmonium, Bottomonium) are well controlled experimental and theoretical probes for the quark-gluon-plasma
- Phenomenological models describe quarkonium supression via a Schrödinger equation + assumption of early formation of bound states
- Characterized by a separation of scale:

$$M_Q \gg M_Q v \gg M_Q v^2 \gg \Lambda_{\rm QCD}$$

- Very heavy states, e.g. Υ(1S), already bound Coulombically
 - M_Q ... heavy quark mass ($m_{Bottom} = 4.18(3)$ GeV [PDG 2017]) Mv
 - v ... relative velocity in centre of mass frame ($v_{Bottom}^2 \approx 0.1$)
 - Λ_{QCD} ... momentum scale below which gluons strongly interacting
- /lv ... typical momentum
- Mv² ... typical kinetic or potential energy

ALEXANDER LEHMANN The 37th

The 37th International Symposium on Lattice Field Theory – June 19th – Wuhan – PRC

Heavy Quarkonia in Early Stages in HICs





Early dynamics of heavy quarkonium in HIC largely unexplored

Rule of thumb via uncertainty relation: $\tau_{form} \sim 1/E_{bind} \approx 0.2 \dots 0.4$ fm/c

Heavy Quarkonia in Early Stages in HICs





Early dynamics of heavy quarkonium in HIC largely unexplored

Rule of thumb via uncertainty relation: $\tau_{form} \sim 1/E_{bind} \approx 0.2 \dots 0.4$ fm/c

Heavy Quarkonia in Early Stages in HICs





Early dynamics of heavy quarkonium in HIC largely unexplored

Rule of thumb via uncertainty relation: $\tau_{form} \sim 1/E_{bind} \approx 0.2 \dots 0.4$ fm/c

Can we find hints for heavy-quarkonium formation in the glasma?

Real-Time Evolution of the Gauge Fields

 Vital insight into glasma dynamics via classical statistical simulations of gauge fields in expanding geometry





Real-Time Evolution of the Gauge Fields

 Vital insight into glasma dynamics via classical statistical simulations of gauge fields in expanding geometry

In this study Hamiltonian evolution in axial gauge, formulated in spatial links and electric fields (Leapfrog) in a non-expanding box

$$\partial_t U_j(x,t) = iE_j^a U_j(x,t) \qquad \partial_t E_j^a(x,t) = -2ImTr\left\{ T^a \sum_{j \neq k} \left[U_{ij}(t,x) + U_{i(-j)}(t,x) \right] \right\} \\ U_j(x) = \exp(ia_j A_j^a T^a) \quad E_j^a = F_{0j}^a = a_0 a_j 2ImTr[T^a U_{0j}]$$

ISOQuant

ANICS .

¹ (plasma instabilities

KM (plasma instabilities)

BGLMV (const. anisotropy

Occupancy n_{Hard}

Smaller occupancy

SU(2) Lattice

•

 $\alpha_{\rm S}^{1/2}$ ¹

 $\alpha_{\rm S}^{1/3}$

 $\alpha_{\rm S}^{1/7}$

 δ_{S}

Momentum space anisotropy: Λ_L/Λ_T

University of Stavanger

Ľ

Schlichting,

avski.

Bogusl

Venugopal

Higher anisotropy

Real-Time Evolution of the Gauge Fields

 Vital insight into glasma dynamics via classical statistical simulations of gauge fields in expanding geometry

In this study Hamiltonian evolution in axial gauge, formulated in spatial links and electric fields (Leapfrog) in a non-expanding box

$$\partial_t U_j(x,t) = iE_j^a U_j(x,t) \qquad \partial_t E_j^a(x,t) = -2\mathrm{Im}\mathrm{Tr}\left\{T^a \sum_{j \neq k} \left[U_{ij}(t,x) + U_{i(-j)}(t,x)\right]\right\}$$
$$U_j(x) = \exp(ia_j A_j^a T^a) \quad E_j^a = F_{0j}^a = a_0 a_j 2\mathrm{Im}\mathrm{Tr}[T^a U_{0j}]$$

Initial conditions drawn from a statistical ensemble





The real-time NRQCD setup





The real-time NRQCD setup





The real-time NRQCD setup





The 37th International Symposium on Lattice Field Theory – June 19th – Wuhan – PRC

Real-time Lattice NRQCD



- Effective **non-relativistic formulation** of heavy quarks from systematic expansion of QCD action in quark velocity v for 2-component **pauli spinors** ψ , χ
- Hamiltonian to order O(v³) with leading order Wilson coefficients c_i=1

$$H^{\psi} = -\frac{\overline{D}^2}{2M} - c_1 \frac{g}{2M} \vec{\sigma} \cdot \vec{B} - c_2 \frac{g}{8M^2} \vec{D} \cdot \vec{E} - c_3 \frac{ig}{8M^2} \vec{\sigma} \cdot \left(\vec{D} \times \vec{E} - \vec{E} \times \vec{D}\right)$$

$$D_i \psi(x) = (U_{i,x} \psi_{x+\hat{i}} - U_{i,x-\hat{i}}^{\dagger} \psi_{x-\hat{i}})/2a_i + O(a_i^2)$$

Real-time Lattice NRQCD



- Effective **non-relativistic formulation** of heavy quarks from systematic expansion of QCD action in quark velocity v for 2-component **pauli spinors** ψ , χ
- Hamiltonian to order O(v³) with leading order Wilson coefficients c_i=1

$$H^{\psi} = -\frac{\vec{D}^2}{2M} - c_1 \frac{g}{2M} \vec{\sigma} \cdot \vec{B} - c_2 \frac{g}{8M^2} \vec{D} \cdot \vec{E} - c_3 \frac{ig}{8M^2} \vec{\sigma} \cdot \left(\vec{D} \times \vec{E} - \vec{E} \times \vec{D}\right)$$

$$D_i \psi(x) = (U_{i,x} \psi_{x+\hat{i}} - U_{i,x-\hat{i}}^{\dagger} \psi_{x-\hat{i}})/2a_i + O(a_i^2)$$

Real-time quarkonium current **correlator** $D^>$ from heavy quark propagator G

 $D_{V}^{>}(t_{2},t_{1}) = \langle J_{NRQCD}^{i}(t_{2}) J_{i,NRQCD}^{+}(t_{1}) \rangle \qquad D_{V}^{>}(t) = \int DU \ G^{\psi} \sigma^{i} G^{\chi +} \sigma_{i} \ e^{iS[U]}$

Real-time Lattice NRQCD



 $G^{c}[U]_{x_{2},x_{1}} = 1_{2N_{c} \times 2N_{c}} \delta_{x_{2}.x_{1}}$

- Effective **non-relativistic formulation** of heavy quarks from systematic expansion of QCD action in quark velocity v for 2-component **pauli spinors** ψ , χ
- Hamiltonian to order O(v³) with leading order Wilson coefficients c_i=1

$$H^{\psi} = -\frac{\vec{D}^2}{2M} - c_1 \frac{g}{2M} \vec{\sigma} \cdot \vec{B} - c_2 \frac{g}{8M^2} \vec{D} \cdot \vec{E} - c_3 \frac{ig}{8M^2} \vec{\sigma} \cdot \left(\vec{D} \times \vec{E} - \vec{E} \times \vec{D}\right)$$

$${}_{D_i \psi(x) = (U_{i,x} \psi_{x+\hat{i}} - U_{i,x-\hat{i}}^{\dagger} \psi_{x-\hat{i}})/2a_i + O(a_i^2)}$$

Real-time quarkonium current **correlator** $D^>$ from heavy quark propagator G

$$D_{V}^{>}(t_{2}, t_{1}) = \langle J_{NRQCD}^{i}(t_{2}) J_{i,NRQCD}^{+}(t_{1}) \rangle \qquad D_{V}^{>}(t) = \int DU \, G^{\psi} \sigma^{i} G^{\chi +} \sigma_{i} \, e^{iS[U]}$$

• Heavy quark equation of motion: $G^{\psi}[U]^{t+\Delta t}_{x_2,x_1} = \exp[-i\Delta t H^{\psi}[U]] \cdot G^{\psi}[U]^{t}_{x_2,x_1}$

Real-time Lattice NRQCD



- Effective **non-relativistic formulation** of heavy quarks from systematic expansion of QCD action in quark velocity v for 2-component **pauli spinors** ψ , χ
- Hamiltonian to order O(v³) with leading order Wilson coefficients c_i=1

$$H^{\psi} = -\frac{\vec{D}^2}{2M} - c_1 \frac{g}{2M} \vec{\sigma} \cdot \vec{B} - c_2 \frac{g}{8M^2} \vec{D} \cdot \vec{E} - c_3 \frac{ig}{8M^2} \vec{\sigma} \cdot \left(\vec{D} \times \vec{E} - \vec{E} \times \vec{D}\right)$$

$${}_{D_i \psi(x) = (U_{i,x} \psi_{x+\hat{i}} - U_{i,x-\hat{i}}^{\dagger} \psi_{x-\hat{i}})/2a_i + O(a_i^2)}$$

Real-time quarkonium current **correlator** $D^>$ from heavy quark propagator G

$$D_{V}^{>}(t_{2},t_{1}) = \langle J_{NRQCD}^{i}(t_{2}) J_{i,NRQCD}^{+}(t_{1}) \rangle \qquad D_{V}^{>}(t) = \int DU \, G^{\psi} \sigma^{i} G^{\chi +} \sigma_{i} \, e^{iS[U]}$$

Heavy quark equation of motion: $G^{\psi}[U]^{t+\Delta t}_{x_2,x_1} = \exp[-i\Delta t H^{\psi}[U]] \cdot G^{\psi}[U]^{t}_{x_2,x_1}$

$$G(t + a_t) = (1 - ia_t H[U(t)]) \cdot G(t)$$

Often via forward Euler: cheap but 1st order in dt, inherently unstable (Courant), range of validity of NRQCD mixed with breakdown of discretization

ALEXANDER LEHMANN

The 37th International Symposium on Lattice Field Theory – June 19th – Wuhan – PRC

Real-time Lattice NRQCD



- Effective **non-relativistic formulation** of heavy quarks from systematic expansion of QCD action in quark velocity v for 2-component **pauli spinors** ψ , χ
- Hamiltonian to order O(v³) with leading order Wilson coefficients c_i=1

$$H^{\psi} = -\frac{\vec{D}^2}{2M} - c_1 \frac{g}{2M} \vec{\sigma} \cdot \vec{B} - c_2 \frac{g}{8M^2} \vec{D} \cdot \vec{E} - c_3 \frac{ig}{8M^2} \vec{\sigma} \cdot \left(\vec{D} \times \vec{E} - \vec{E} \times \vec{D}\right)$$

$${}_{D_i \psi(x) = (U_{i,x}\psi_{x+\hat{i}} - U_{i,x-\hat{i}}^{\dagger}\psi_{x-\hat{i}})/2a_i + O(a_i^2)}$$

Real-time quarkonium current **correlator** $D^>$ from heavy quark propagator G

$$D_{V}^{>}(t_{2},t_{1}) = \langle J_{NRQCD}^{i}(t_{2}) J_{i,NRQCD}^{+}(t_{1}) \rangle \qquad D_{V}^{>}(t) = \int DU \ G^{\psi} \sigma^{i} G^{\chi +} \sigma_{i} \ e^{iS[U]}$$

 $= \text{ Heavy quark equation of motion:} \quad G^{\psi}[U]^{t+\Delta t} = \exp[-i\Delta t H^{\psi}[U]] \cdot G^{\psi}[U]^{t} \\ x_{2},x_{1} = \exp[-i\Delta t H^{\psi}[U]] \cdot G^{\psi}[U]$

$$G(t+a_t) = \left(1 + \frac{ia_t}{2}H[U(t)]\right)^{-1} \cdot \left(1 - \frac{ia_t}{2}H[U(t)]\right) \cdot G(t)$$

Optimal rational approximation of exp (Crank-Nicholson, O(dt²)): **unconditionally stable**, no mixing of range of validity. (No operator splitting via MPI PETSC)

ALEXANDER LEHMANN

The 37th International Symposium on Lattice Field Theory – June 19th – Wuhan – PRC

Wigner Coordinates for Non-Equilibrium



No time translational invariance: need to correctly account for relative and central time coordinate in 2pt functions: $t_2 + t_1$

$$t = \frac{t_2 + t_1}{2} \ s = t_2 - t_1$$

Wigner Coordinates for Non-Equilibrium



- No time translational invariance: need to correctly account for relative and central time coordinate in 2pt functions: $t = \frac{t_2 + t_1}{2} \quad s = t_2 - t_1$
- Spectral function from Fourier transform over finite temporal extent in s

$$\rho(t, \omega, \boldsymbol{p} = 0)$$

$$= 2Im\left[\int_{0}^{s_{max}} D^{>}\left(t + \frac{s}{2}, t - \frac{s}{2}, \boldsymbol{p} = 0\right)e^{-i\omega s} ds\right] \text{ supportion for } D^{>}(t_{2}, t_{1})$$

Wigner Coordinates for Non-Equilibrium



- **No time translational invariance:** need to correctly account for relative and central time coordinate in 2pt functions: $t = \frac{t_2 + t_1}{2} \ s = t_2 - t_1$
- Spectral function from Fourier transform over **finite** temporal extent in s

$$\rho(t, \omega, \boldsymbol{p} = 0)$$

$$= 2Im \left[\int_{0}^{s_{max}} D^{>} \left(t + \frac{s}{2}, t - \frac{s}{2}, \boldsymbol{p} = 0\right) e^{-i\omega s} ds\right]$$
Spectral function has **explicit t dependence**, signaling real-time evolution of gauge fields

Spectral function has **explicit t dependence**, signaling real-time evolution of gauge fields



Free theory sanity check





Real-time correlation function is complex – finite volume effects as recurrence

Free theory sanity check





Real-time correlation function is complex – finite volume effects as recurrence

Free spectral function reproduced – reducing mass does not lead to breakdown

Free theory sanity check





- Real-time correlation function is complex finite volume effects as recurrence
- Free spectral function reproduced reducing mass does not lead to breakdown

Quarkonium in the Glasma (I)





Low enegy gluons do not significantly impact quarks at early times

Quarkonium in the Glasma (I)





Low enegy gluons do not significantly impact quarks at early times

Quarkonium in the Glasma (I)





- Low enegy gluons do not significantly impact quarks at early times
- Bulk glue effects manifest in the intermediate (s,t) time physics of heavy quarks
- At the parameters used here, **no signs for binding** into clear resonances

Quarkonium in the Glasma (II)





Reduction of singlet amplitude and broadening understood from gluon absoprtion

Octet enhancement from interaction with low enegy gluonic bulk

Understanding the absence of binding



- Consider static quarks via the **non-equilibrium real-time Wilson loop** W(t,s,x)
- Attempt to extract effective real-time potential via Wilson loop spectral function Re[V] from position of lowest lying peak, Im[V] from width, see Y.Burnier, A. Rothkopf PRD86 (2012) 051503



Understanding the absence of binding



- Consider static quarks via the **non-equilibrium real-time Wilson loop** W(t,s,x)
- Attempt to extract effective real-time potential via Wilson loop spectral function Re[V] from position of lowest lying peak, Im[V] from width, see Y.Burnier, A. Rothkopf PRD86 (2012) 051503



Understanding the absence of binding



- Consider static quarks via the non-equilibrium real-time Wilson loop W(t,s,x)
- Attempt to extract effective real-time potential via Wilson loop spectral function Re[V] from position of lowest lying peak, Im[V] from width, see Y.Burnier, A. Rothkopf PRD86 (2012) 051503



Similar to results in thermal equilibrium: **no real-part** of the potential emerges See M. Laine et.al. JHEP 0709 (2007) 066

Understanding the absence of binding



- Consider static quarks via the **non-equilibrium real-time Wilson loop** W(t,s,x)
- Attempt to extract effective real-time potential via Wilson loop spectral function Re[V] from position of lowest lying peak, Im[V] from width, see Y.Burnier, A. Rothkopf PRD86 (2012) 051503



Similar to results in thermal equilibrium: **no real-part** of the potential emerges see M. Laine et.al. JHEP 0709 (2007) 066

No indications of binding, not even Coulombic, found out of equilibrium so far

Summary



- Combination of real-time classical statistical simulations for gauge fields with novel stable lattice NRQCD solver
- Direct computation of non-equilibrium real-time quarkonium correlators and spectral functions in Wigner coordinates
- Enhancement in quarkonium colour octet channel and no signs of binding in the singlet channel
- Consistent with absence of a real-part in effective potential
- Need further study at stronger couplings to confirm absence or presence of binding

Thank you for your attention - 謝謝