

The hadronic vacuum polarization of the muon from four-flavor lattice QCD

Fermilab Lattice, HPQCD, and MILC Collaborations

C. T. H. Davies, C. DeTar, A. X. El-Khadra, E. Gámiz, Steven
Gottlieb, D. Hatton, A. S. Kronfeld, J. Laiho, G. P. Lepage, Yuzhi
Liu, P. B. Mackenzie, C. McNeile, E. T. Neil, T. Primer, J. N.
Simone, D. Toussaint, R. S. Van de Water, and A. Vaquero

Methodology

$$a_\mu^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dQ^2 K_E(Q^2) \widehat{\Pi}(Q^2) \quad \widehat{\Pi}(Q^2) \equiv \Pi(Q^2) - \Pi(0)$$

Taylor coefficients

$$\widehat{\Pi}(\omega^2) = \frac{4\pi^2}{\omega^2} \int_0^\infty dt G(t) \left[\omega^2 t^2 - 4 \sin^2 \left(\frac{\omega t}{2} \right) \right] \quad G_{2n} \equiv a^4 \sum_t \sum_{\mathbf{x}} \sum_f t^{2n} Z_V^2 q_f^2 \langle j_i^f(\mathbf{x}, t) j_i^f(0) \rangle$$

$$G(t) = \frac{1}{3} \int d\mathbf{x} \sum_f q_f^2 Z_V^2 \langle j_i^f(\mathbf{x}, t) j_i^f(0,0) \rangle$$

$$\widehat{\Pi}(Q^2) = \sum_{j=1}^{\infty} Q^{2j} \Pi_j$$

- Calculate $G_{2n} \rightarrow \Pi_{n-1}$
- Use Padé to extend $\widehat{\Pi}(Q^2)$ to small Q^2
- Integrate to get $a_\mu^{\text{HVP,LO}}$

Lattice parameters

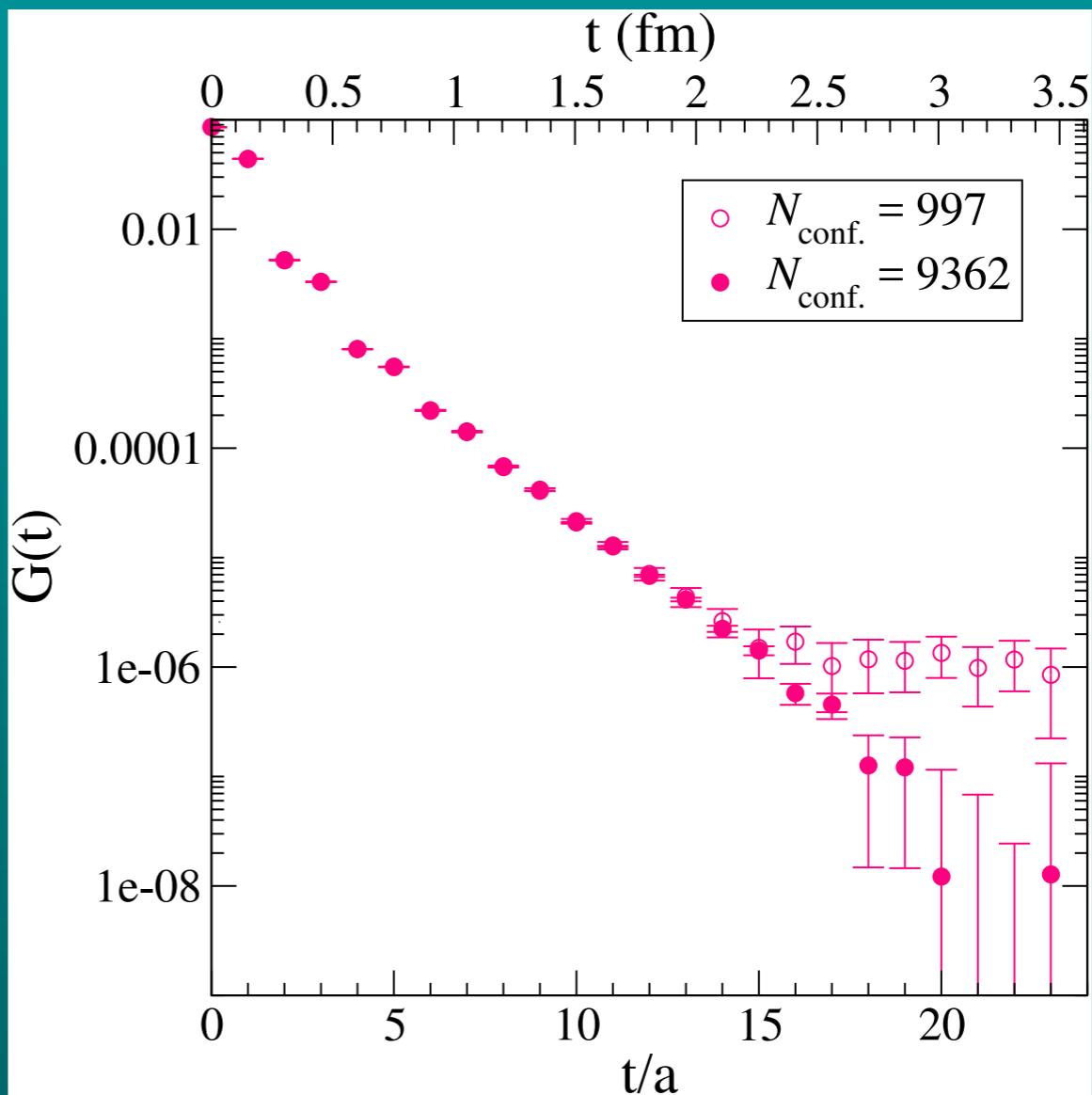
- Ensembles
- 2+1+1 sea

| $\approx a(\text{fm})$ | $M_\pi(\text{MeV})$ | size | N_{config} |
|------------------------|---------------------|-------------------|---------------------|
| 0.15 | 133 | $32^3 \times 48$ | 997 |
| 0.15 | 135 | $32^3 \times 48$ | 9362 |
| 0.12 | 133 | $48^3 \times 64$ | 998 |
| 0.09 | 128 | $64^3 \times 96$ | 1557 |
| 0.06 | 134 | $96^3 \times 192$ | 1230 |

- Valence: here only $m_u = m_d = m_\ell$ so only $a_\mu^{\ell\ell}$
- Lattice spacing a/w_0 is known to 5 per mille.

Correlator example

- 0.15 fm analysis
- Higher statistics has obviously improved the signal in the tail



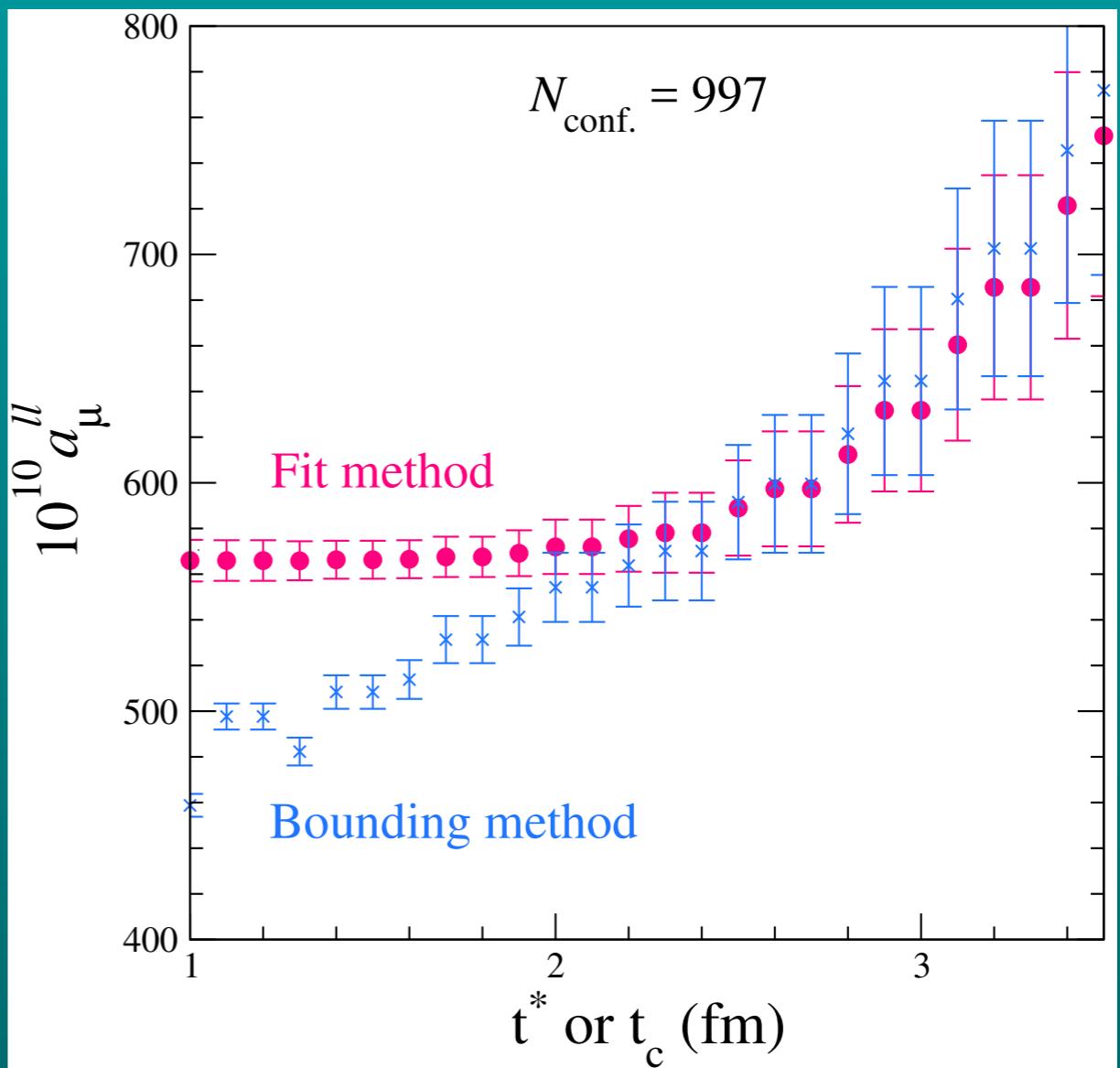
Dealing with noisy data at large t

- Bounding method: 2 pion contribution for $t > t_c$
[RBC/UKQCD, BMW, Mainz/CLS (improved)]
- Hybrid R: Replace with contribution from
$$R = \sigma(e^+e^- \rightarrow \text{hadrons})/\sigma(e^+e^- \rightarrow \mu^+\mu^-)$$

[RBC/UKQCD]
- (Here) Fit method: Fit low t data and replace
high $t > t^*$ values with an extrapolation

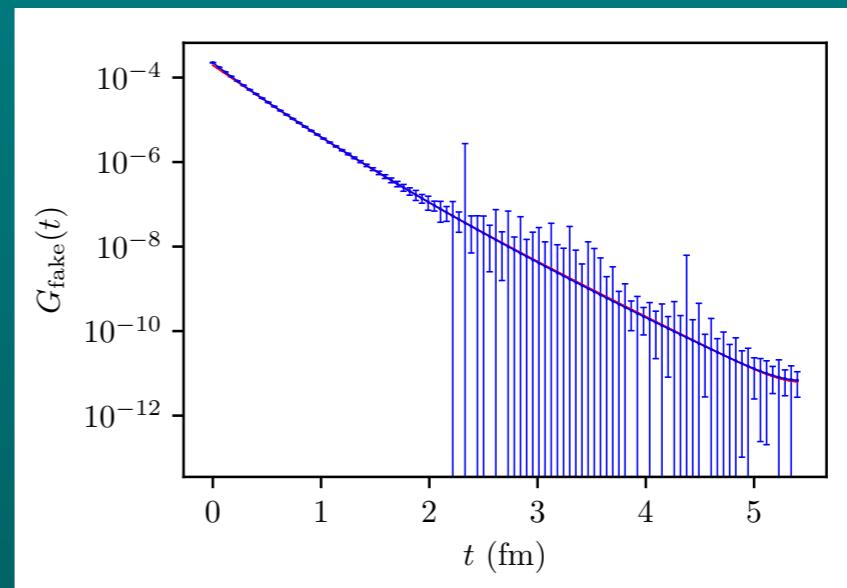
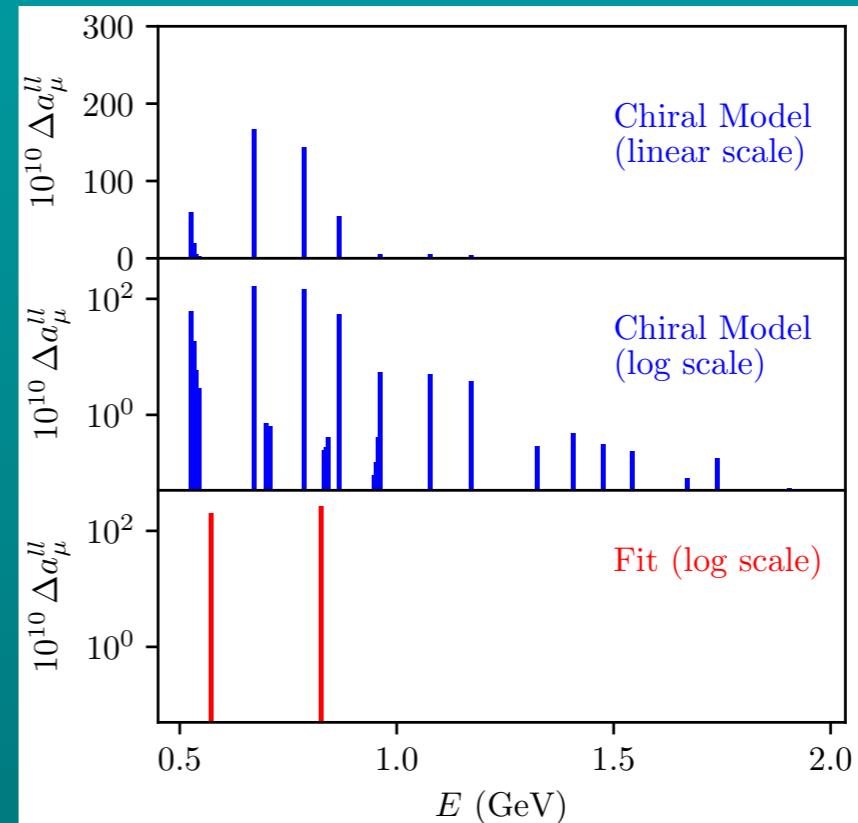
Compare bounding and fit methods

- Agree for $t > 2.3$ fm
- Fit method stable for $1 < t^* < 2$ fm



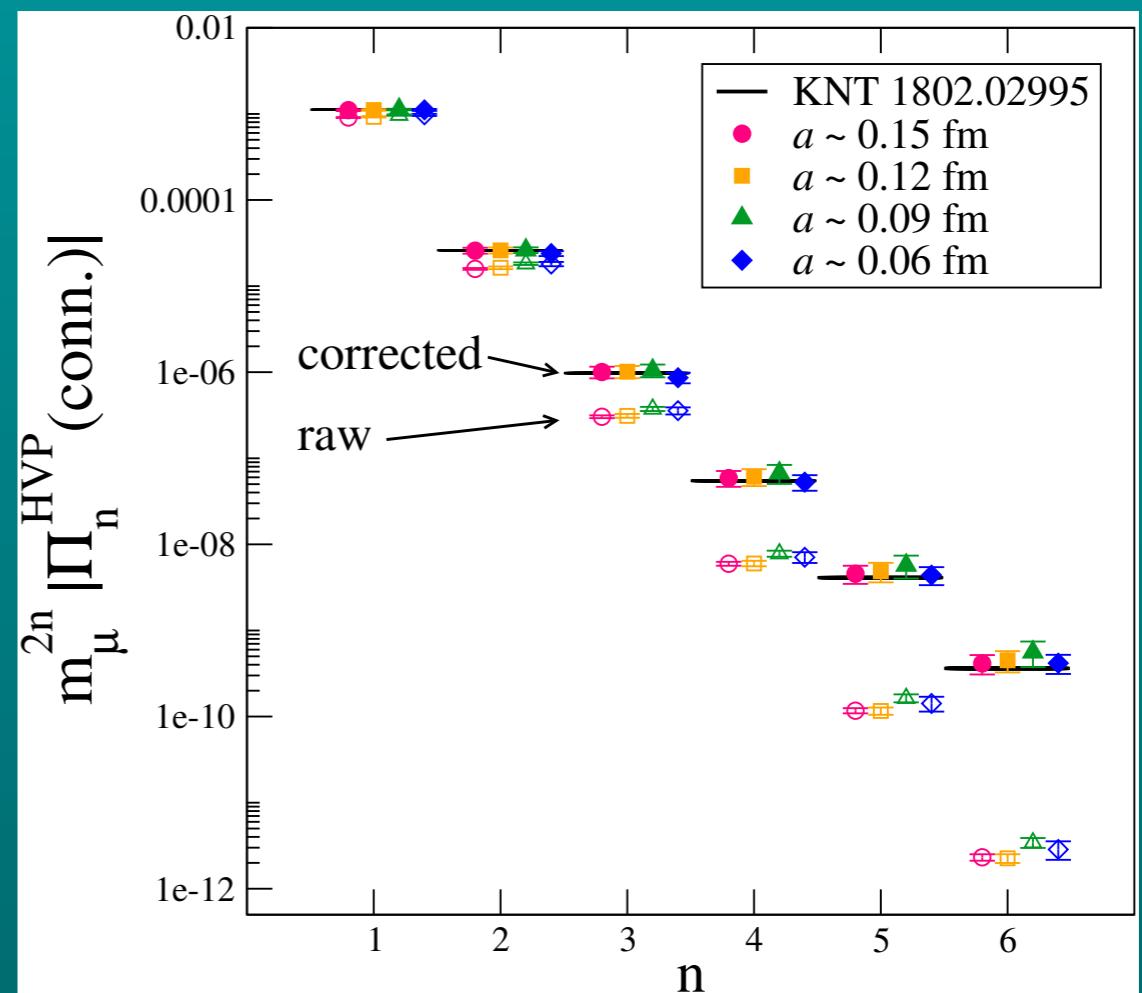
Synthetic data test of fit method

- Generate synthetic data set based on a chiral model with rho plus staggered 2 pion states
- Fit with only two states (red curve) agrees perfectly with model (blue curve)



Corrections

- Use staggered chiral model to correct for
 - Finite volume effects
 - Discretization due to taste-splitting
 - Quark mass tuning
- Subtract lattice $\pi - \pi$ contributions and replace with physical $\pi^0 - \pi^0$ contributions.
- Nice agreement with phenomenological R method (black lines)



n^{th} Taylor coefficient

Light-light extrapolation

$$a_\mu^{ll}(\text{latt.}) = a_\mu^{ll}(\text{conn.}) \left(1 + c_s \sum_{f=l,l,s,c} \frac{\delta m_f}{\Lambda} + c_{a^2} \frac{(a\Lambda)^2}{\pi^2} \right)$$

- $\Lambda = 500 \text{ MeV}$

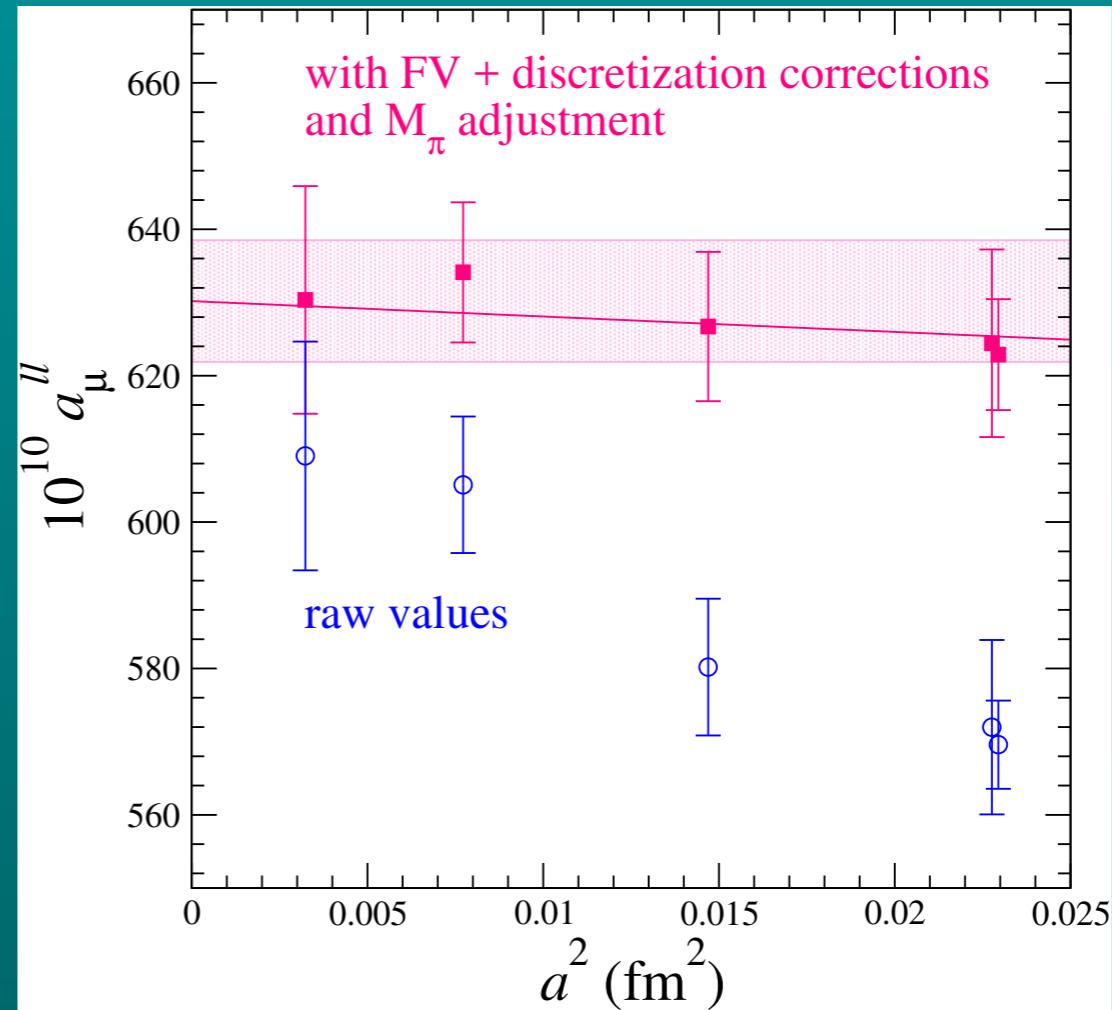
- Priors

$$c_s = 0.0(3)$$

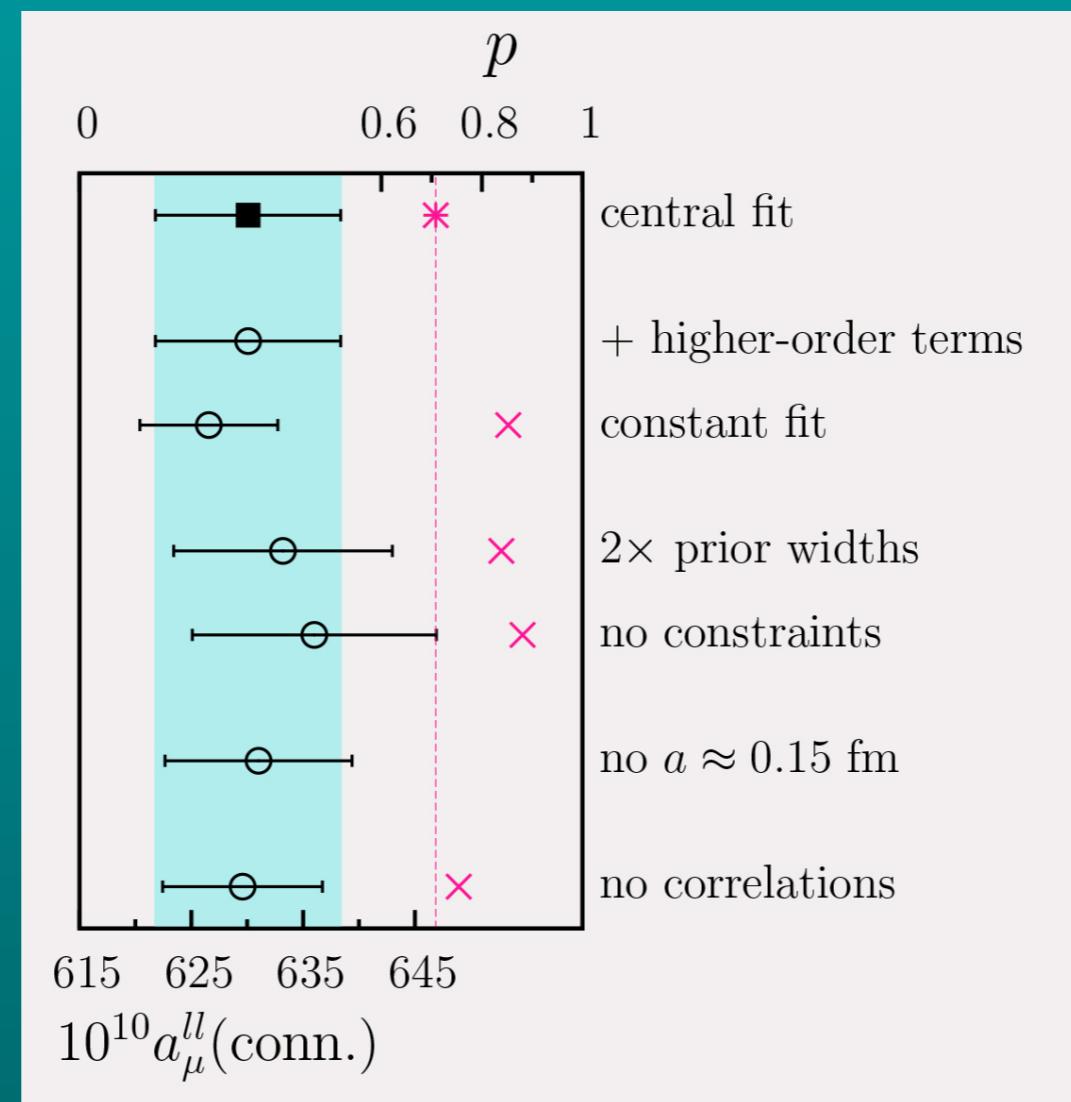
$$c_{a^2} = 0(1)$$

- Result

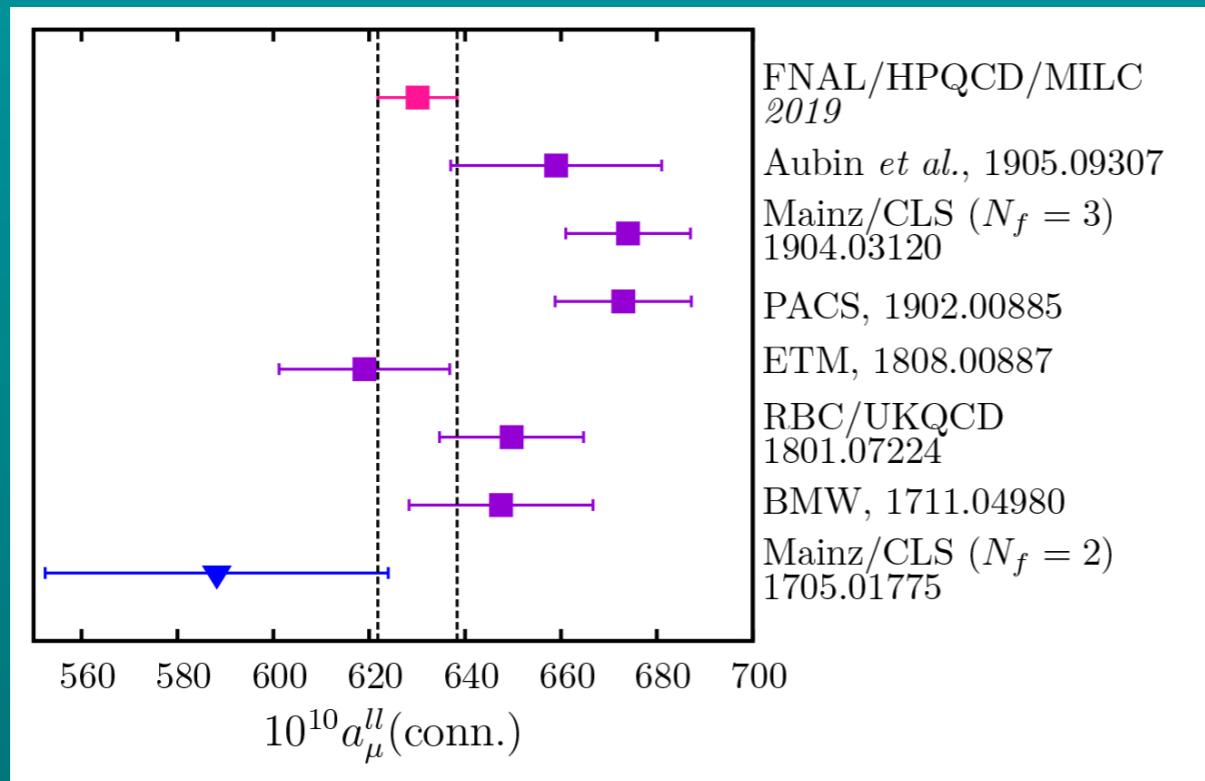
$$a_\mu^{\ell\ell}(\text{conn}) = 630.1(8.3)$$



Stability of the continuum extrapolation



Comparison of light-quark-connected contribution

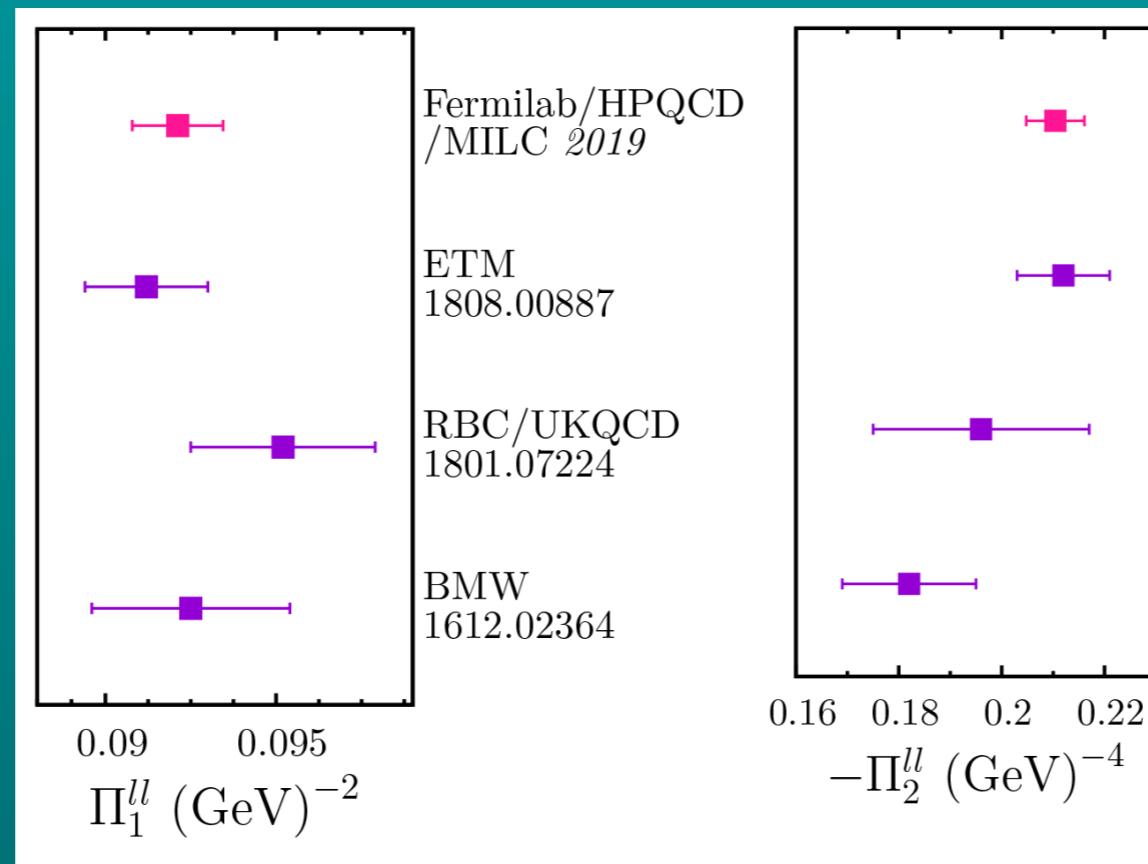


- Small error results from small error in fit-method extrapolation and finite-volume correction

Error budget - light-quark connected

| | |
|--|-------------|
| Lattice spacing uncertainty (w_0) | 0.8% |
| Monte Carlo statistics | 0.7 |
| Continuum extrapolation | 0.7 |
| Finite volume and discretization corrections | 0.3 |
| Current renormalization | 0.1 |
| Chiral interpolation | 0.1 |
| Strange sea quark mass adjustment | 0.1 |
| Pion mass uncertainty | 0.0 |
| Total | 1.3% |

Comparison of light-light Taylor coefficients



Adjusting to get the total u-d HVP contribution

All values $\times 10^{-10}$

| | |
|---------------------------------|---------------|
| $M(\pi^0) \rightarrow M(\pi^+)$ | -4.3 |
| $\pi - \pi$ disconnected | -7.9 |
| Total $\pi - \pi$ | -12(3) |
| $\rho - \omega$ disconnected | -5(5) |
| Strong isospin breaking | 10(10) |
| Electromagnetism | 0(5) |
| Total adjustment | -7(13) |

$$\text{Net } a_\mu^{ud} = 623.1(8.3)(13.0)$$

Including the s, c, b contributions

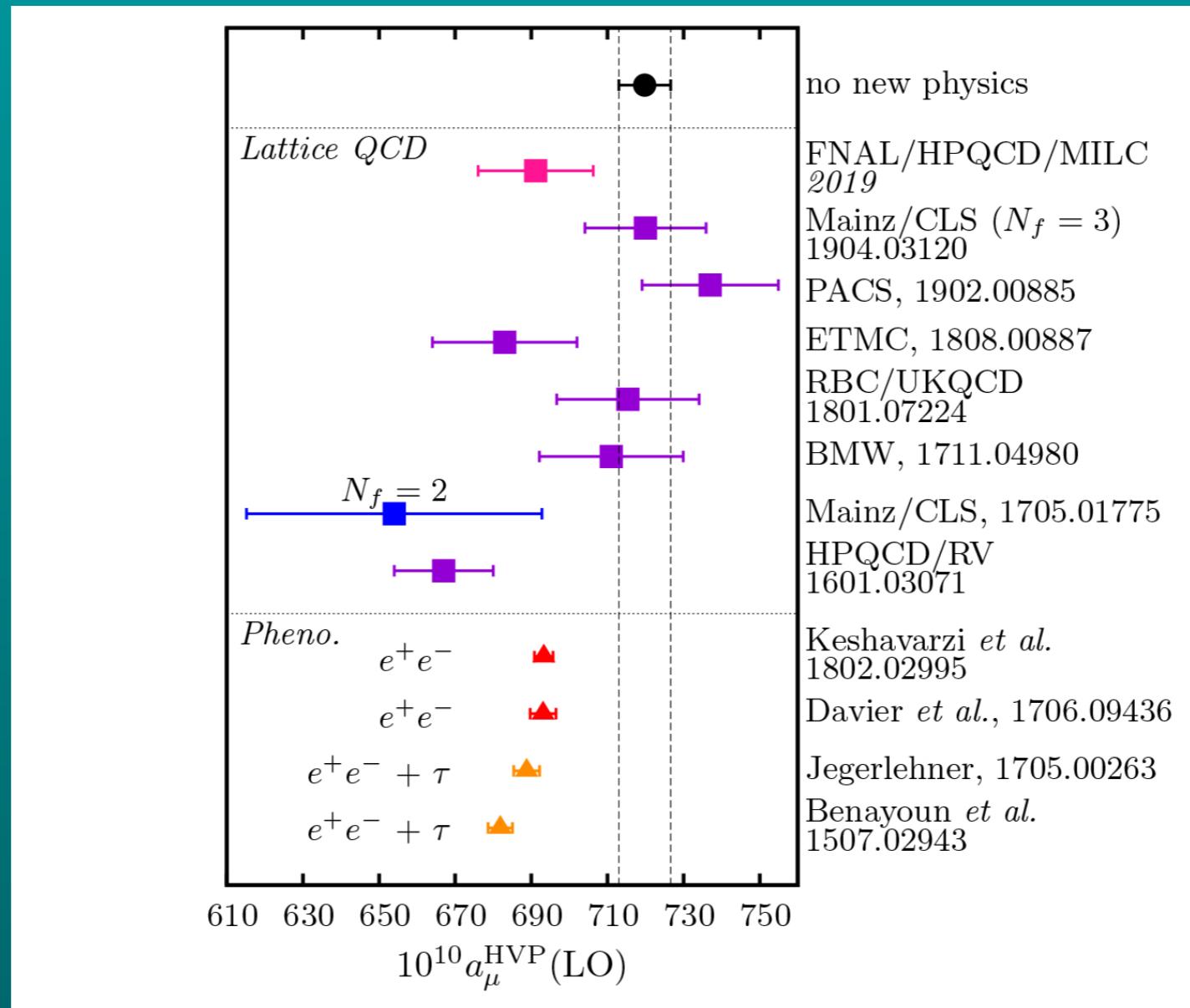
| | |
|---------|----------------|
| light | 623.1(8.3)(13) |
| strange | 53.40(60) |
| charm | 14.40(40) |
| bottom | 0.270(40) |

arXiv:1403.1778
arXiv:1208.2855
arXiv:1408.5768

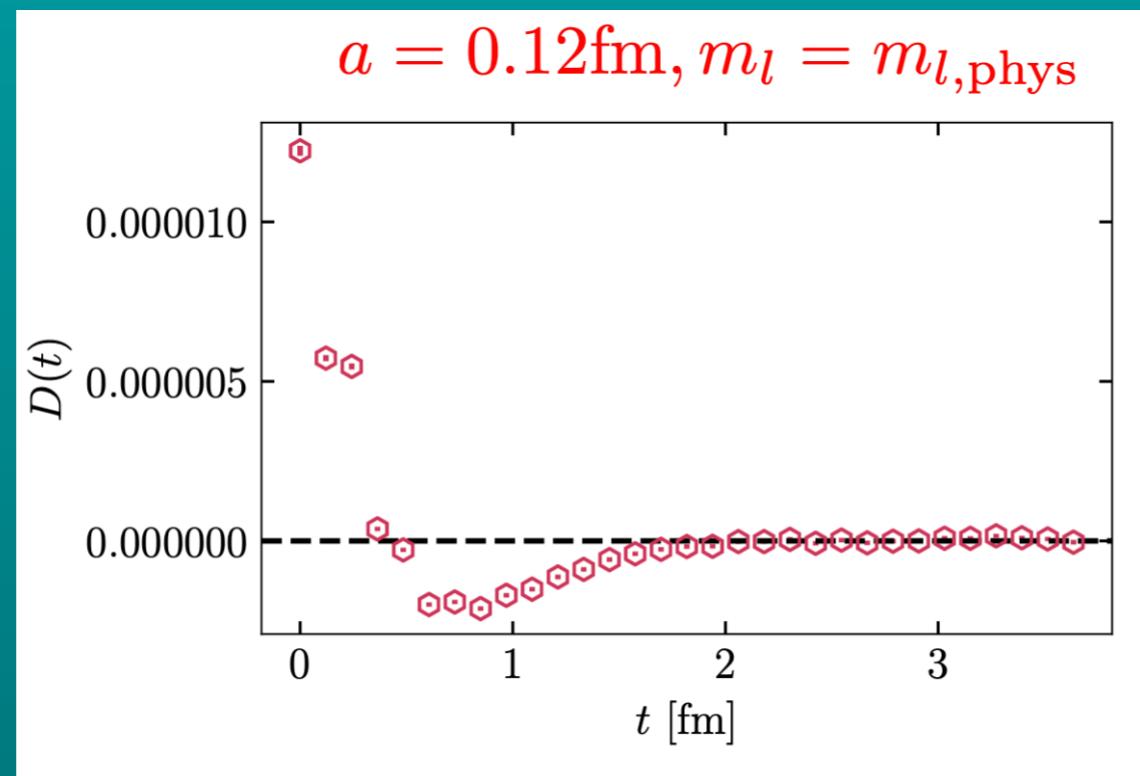
Dan Hatton Mon 4:30 PM

$$\text{Net } a_{\mu}^{\text{HVP-LO}} = 691(15)$$

Total LO HVP result



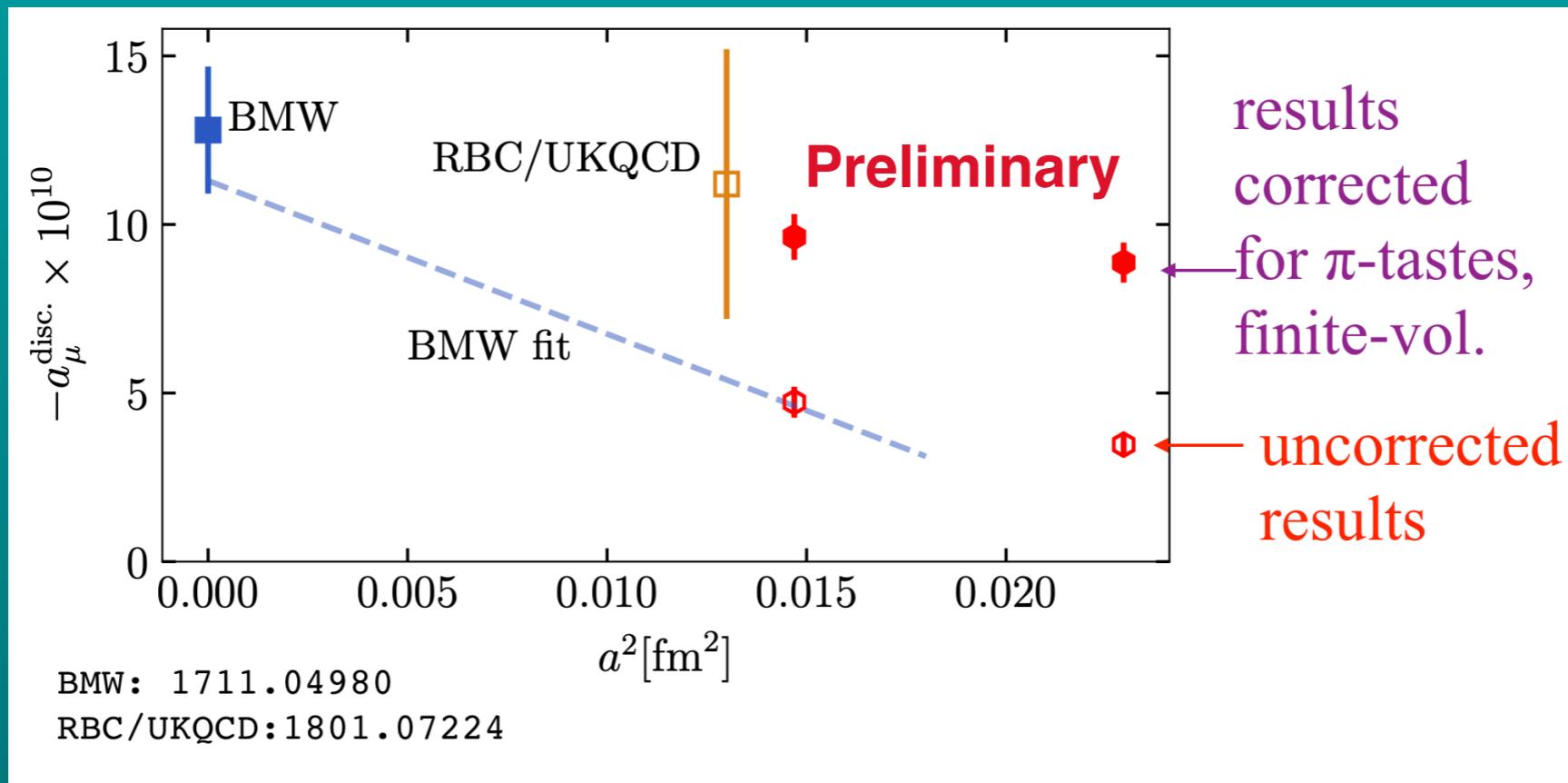
Progress with disconnected HVP



Fit to “ ω -like” states minus “ ρ -like” states

Replace data with fit for $t > 2$ fm
80% of result is from data.

Progress with disconnected HVP



See less a dependence than for BMW.
Results at 0.09 fm will allow an accurate continuum result.
We find isospin breaking effects in disconnected correlators
are relatively large.

Outlook

- Further reductions in the LO HVP require
 - More accurate lattice spacing determination
 - Better control of strong isospin breaking
 - Better control of electromagnetic effects
 - Better statistics -> better control of the continuum extrapolation
- Reaching 0.5% uncertainty is feasible over the next couple of years

Backup

Stability of fit method

- Stability against variations in the number of states and t_{\min}
- $a \sim 0.09$ fm

