

Nucleon scalar charge with overlap fermions

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in collaboration with

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- Introduction.
- Numerical details.
- Results.
- Summary and outlook.

- Fundamental property of hadron structure.
- Probe new physics beyond standard model.
- Not well known experimentally.
- Need accurate lattice determination.

- Gauge configurations generated by the RBC/UKQCD collaboration with 2+1 flavor domain-wall fermion.

Ensemble ID	$L^3 \times T$	a^{-1} (GeV)	m_{π}^{sea} (MeV)	N_{conf}
24I	$24^3 \times 64$	1.747	337	203
32I	$32^3 \times 64$	2.310	302	309
48I	$48^3 \times 96$	1.730	139	81
32ID	$32^3 \times 64$	1.3709	171	200
32Ifine	$32^3 \times 64$	3.148	371	459

Overlap fermions

- Overlap fermion action for the valence quark.
- Chiral symmetry at finite lattice spacing via the Ginsparg-Wilson relation.
- Has better control of the systematic error, but more costly.

The overlap Dirac operator is

$$D(m) = \rho D_{ov}(\rho) + m \left(1 - \frac{D_{ov}(\rho)}{2}\right),$$

where $D_{ov} = 1 + \gamma_5 \epsilon (\gamma_5 D_w(\rho))$, we use the parameter $\rho = -1.5$.

The effective propagator is

$$D_{eff}^{-1} = \frac{1}{D_c + m},$$

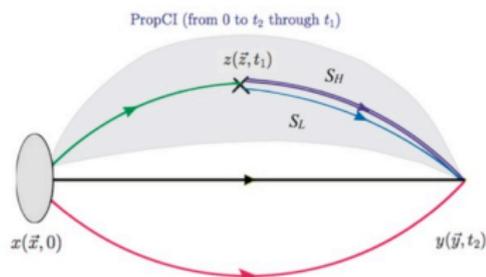
where $D_c = \frac{\rho D_{ov}}{1 - D_{ov}/2}$ is exactly chiral, i.e. $\{\gamma_5, D_c\} = 0$.

Computational strategies

- Low-mode substitution(LMS): Use the low lying eigenvalues and eigenvectors of D_c to speed up the inversion and separate the propagators into its low-mode and high-mode parts.

$$S(y, x) = S_L(y, x) + S_H(y, x), \quad S_L(y, x) = \sum_{|\lambda_i| < \epsilon} \frac{1}{\lambda_i + m} |i\rangle_y \langle i|_x$$

- Z_3 noise grid source: Tie the sources of the three quark propagators stochastically to each point so that one can have multi-to-all propagator from one inversion.
- Stochastic sandwich method with LMS: all-to-all low mode + noise many-to-all high mode



- To extract the scalar charge, we compute the two- and three-point functions:

$$C_{2pt}(t_2) = \sum_{\vec{x}} \Gamma_{\alpha\beta}^e \langle 0 | \chi_\alpha(t_2, \vec{x}) \bar{\chi}_\beta(0, \vec{0}) | 0 \rangle,$$
$$C_{3pt}^{u,d}(t_2, t) = \sum_{\vec{x}, \vec{x}'} \Gamma_{\alpha\beta}^e \langle 0 | \chi_\alpha(t_2, \vec{x}) \mathcal{O}^{u,d}(t, \vec{x}') \chi_\beta(0, 0) | 0 \rangle,$$

where χ is the proton interpolating operator $\chi_\alpha = \epsilon^{abc} [u^{aT} C \gamma_5 d^b] u^c$.

- Only the connected insertion is computed for isovector scalar charge.

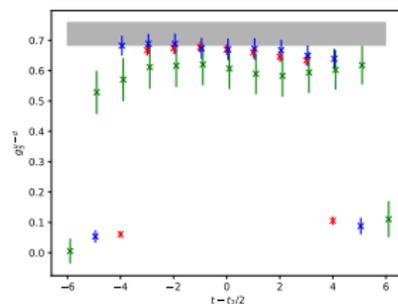
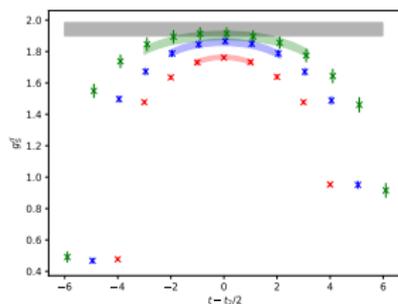
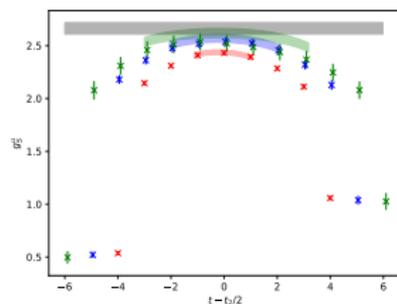
- For each ensemble, the correlation function are computed for several values of valence quark mass (use multi-mass algorithm) and source-sink separation.

Ensemble ID	m_{π}^{sea} (MeV)	$m_{\pi}^{valence}$ (MeV)	$t_{sink} - t_{source}$
24l	337	282 ~ 389	8, 10, 12
32l	302	295 ~ 410	12, 14, 15
48l	139	134 ~ 372	8, 10, 12
32ID	171	146 ~ 325	7, 8, 9, 10, 11
32lfine	371	344 ~ 485	14, 16, 18

Two-state fit

$$R^q(t_2, t) = \frac{C_{3pt}^q(t_2, t)}{C_{2pt}(t_2)} = g_S^q + C_1^q e^{-\Delta M t_2} + C_2^q e^{-\Delta M(t_2-t)} + C_3^q e^{-\Delta M t},$$

Joint fit of $R^u(t, t_S)$ and $R^d(t, t_S)$:



Renormalization factor Z_S (Z. Liu et al. Phys. Rev. D 90, 034505 (2014)):

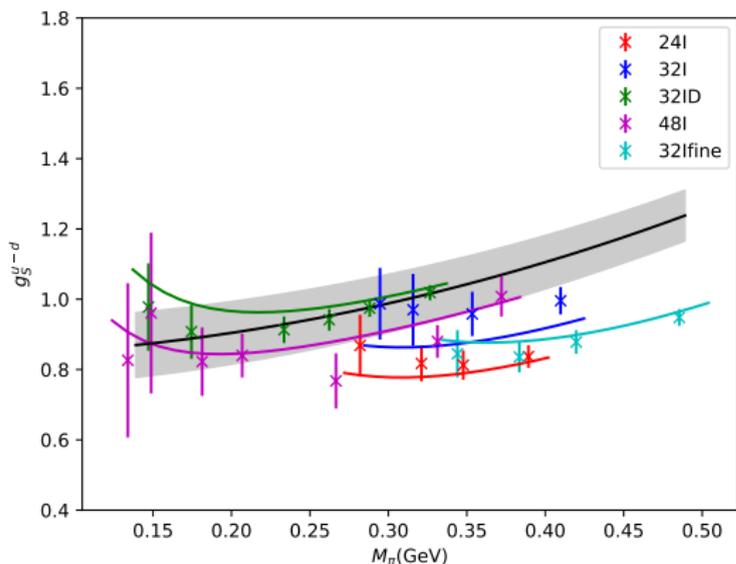
24l	32l	32lD	48l	32lfine
1.1272	1.0563	1.2355	1.1272	0.9539

Results

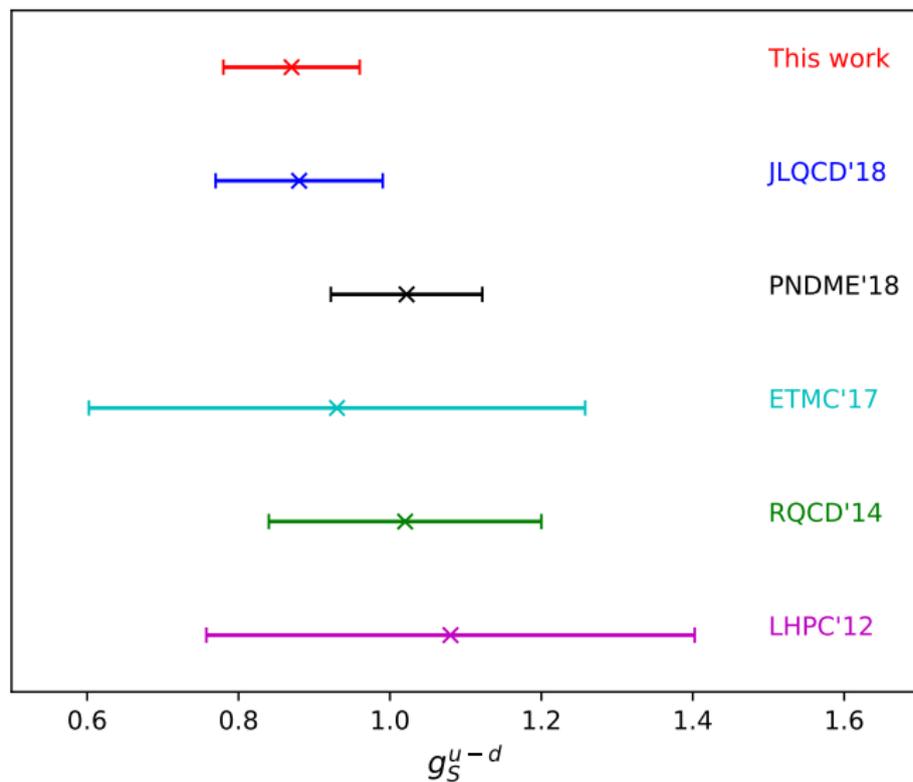
$$g_S(m_{\pi, \nu V}, m_{\pi, \nu S}, a, L) = c_0 + c_1 m_{\pi, \nu V}^2 + c_2 m_{\pi, \nu S}^{mix 2} + c_3 (c_4) a^2 + c_5 e^{-m_{\pi, \nu V} L}$$

c_0	c_1	c_2	c_3	c_4	c_5	χ^2/dof
0.88(0.10)	4.03(0.62)	-4.71(1.00)	-1.11(5.42)	3.13(3.79)	4.83(2.05)	1.06

$$g_S^{u-d} = 0.87(9)$$



Results



- More statistics on the physical ensemble.
- Systematic error needs further investigation.
- Axial, tensor charges, disconnected insertion.
- σ -terms, Form factors.