

Yang Mills short distance potential and perturbation theory

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Wuhan, Lattice 2019, June 17



Motivation

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- ▶ α_s world average is dominated by lattice determinations
FLAG 2019:

$$\alpha_{\overline{\text{MS}}}^{(5)}(M_Z) = 0.11823(81)$$

3.5% precision on

$$\Lambda_{\overline{\text{MS}}}^{(3)} = 343(12)$$

with phenomenology determinations (PDG 16/18)

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- ▶ perturbative potential enters predictions of e+e-
cross-section across t-tbar `threshold' → top quark mass at linear collider

α_s from static potential

► Advantages

- 4-loop PT (4-loop beta-function) available, (normally only 3-loop)
- **relatively** simple observable (no fermions ...)

Peter 97; Schröder 99
Anzai, Kiyo, Sumino, 10
Smirnov, Sminov, Steinhauser, 10
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Kniehl, Penin, 99
Brambilla, Garcia i Tormo, Soto, Vairo, 07,09

► Disadvantages

- infrared divergences, starting at 4-loop, resummation with pNRQCD techniques

$$\alpha^{4+k} [\log(\alpha)]^k, \quad k = 0, 1, \dots$$

or US - scale $\alpha(1/r), \alpha(\mu_{\text{US}}), \quad \mu_{\text{US}} \sim \frac{C_A \alpha}{2r}$

- T(loop) $\rightarrow \infty$ limit, noise
- Discretisation errors at small r , window problem

Pure Yang Mills

- ▶ Precision, study of systematic effects
- ▶ Comparisons: two recent determinations of $\Lambda_{\overline{\text{MS}}}$

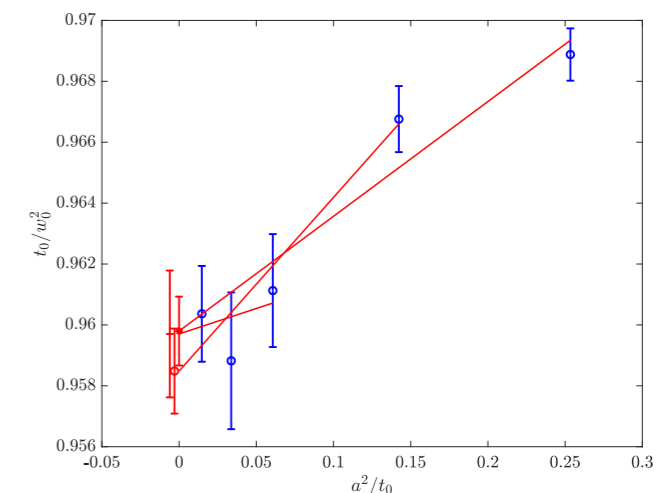
- FlowQCD: $w_0 \Lambda_{\overline{\text{MS}}} = 0.2154(12)$
converted
(by us): $\sqrt{8t_0} \Lambda_{\overline{\text{MS}}} = 0.5968(33)$

from boosted coupling,
contin. extrapolation $\beta \rightarrow \infty$ with a^2
perturbative uncertainty?

- Dalla Brida and Ramos, 2019: $\sqrt{8t_0} \Lambda_{\overline{\text{MS}}} = 0.6227(98)$

NP step scaling down to $\alpha_{\text{SF}} < 0.1$

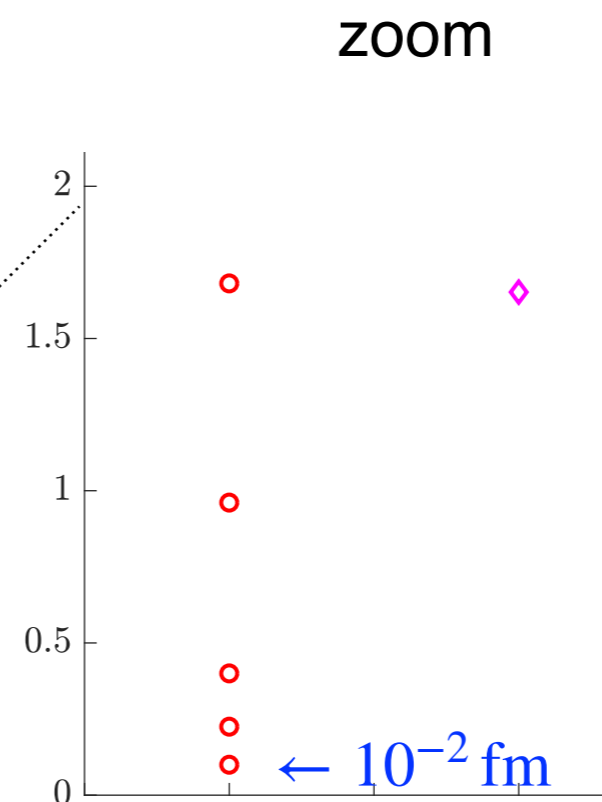
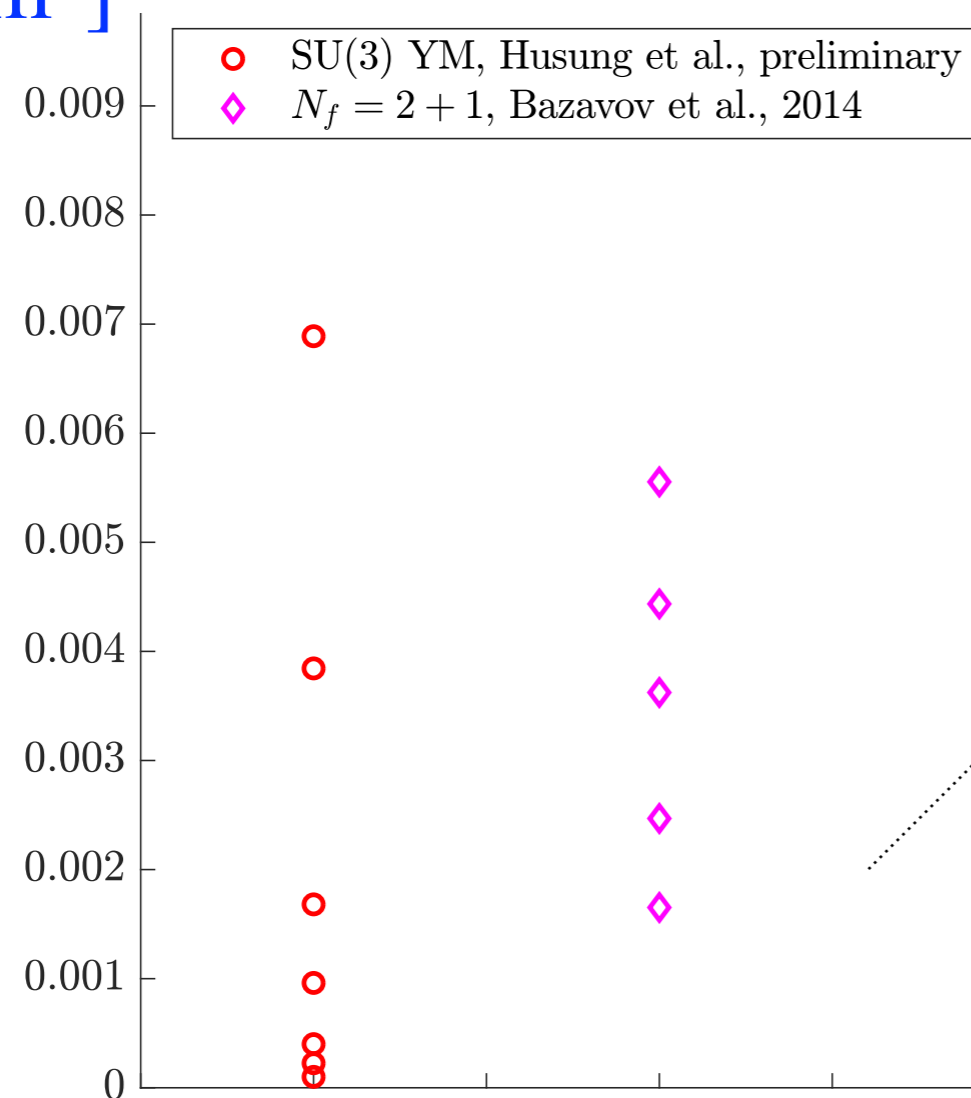
conversion is very precise



Scales, lattices

- ▶ $L/a=32, \dots, 192, L=2\text{fm}$ (big enough in YM), open BC (no topology freezing)

$a^2[\text{fm}^2]$



Strategy to get to small r (see also arXiv:1711.01860)

- ▶ basic scale from t_0 :

$$\alpha_{qq}(\mu, a^2\mu^2), \quad \mu = 1/r = (x\sqrt{8t_0})^{-1}$$

on ensembles with $a > 0.02$ fm

- ▶ Then step scaling functions

$$\Sigma(u, a/r) = \bar{g}_{qq}^2(sr) \Big|_{\bar{g}_{qq}^2(r)=u}$$

with $s = 3/4$ including $a = \{1.0, 1.4, 2.0\} \times 10^{-2}$ fm

Strategy to get to small r (see also arXiv:1711.01860)

- ▶ basic scale from t_0 :

$$\alpha_{\text{qq}}(\mu, a^2\mu^2), \quad \mu = 1/r = (x\sqrt{8t_0})^{-1}$$
$$0.25 \leq x \leq 0.4$$

on ensembles with $a > 0.02$ fm

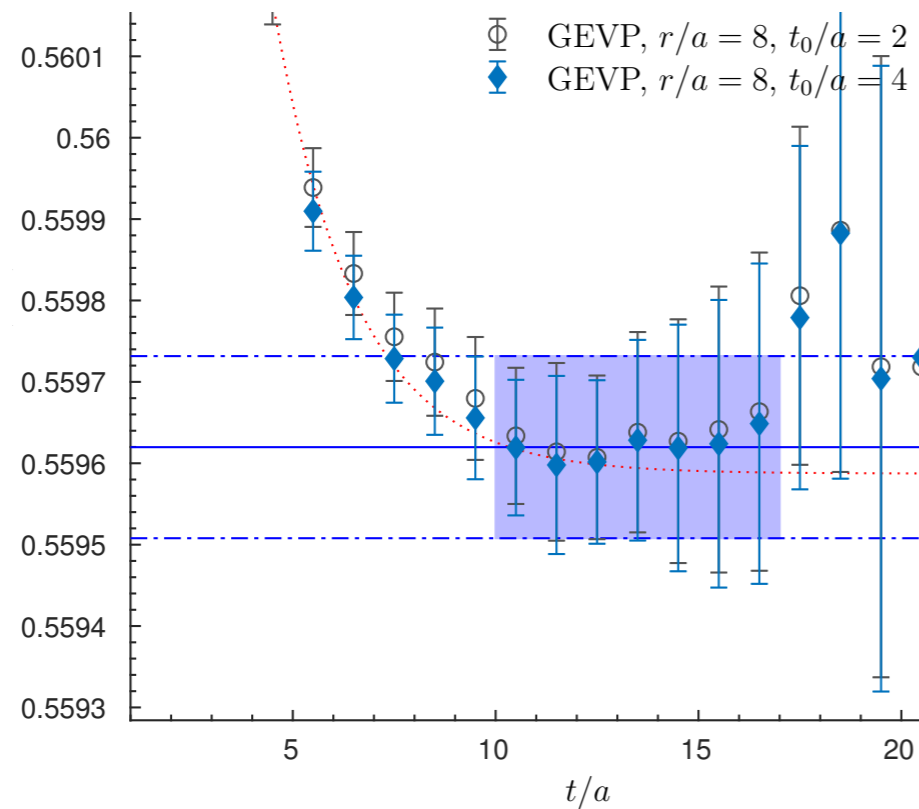
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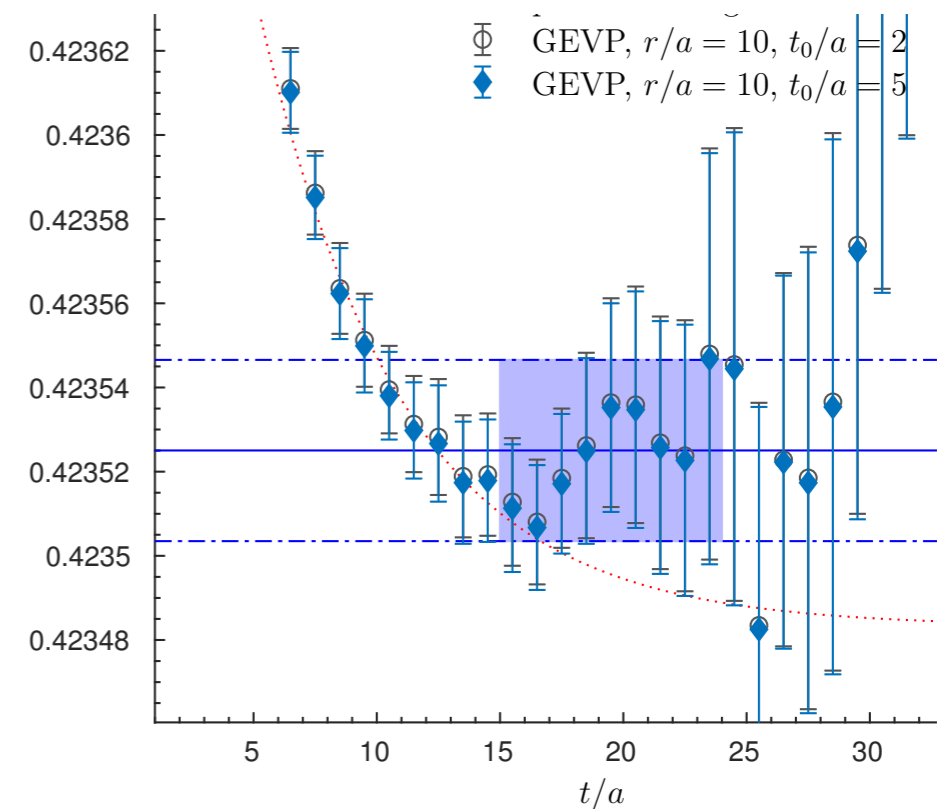
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Plateaux

- ▶ 1-link integral
- ▶ GEVP with
3 trial wave-functions
 $a=0.03$ fm $r=0.24$ fm



- ▶ 2 trial wave-functions
 $a=0.01$ fm $r=0.10$ fm



- ▶ Tree-level improved force, r_I such that

$$F_{\text{TI}}(r_I) = \frac{1}{a}[V(r+a) - V(r)] = \frac{4}{3} \frac{\alpha}{r_I^2} + \mathcal{O}(\alpha^2)$$

- ▶ Tree-level improved force with a^2 improved derivative:

$$F_n(r_n) = \frac{1}{a}[V(r+a) - V(r)] = \frac{1}{a}[V(r_n + a/2) - V(r_n - a/2)], \quad r_n = r + a/2$$

$$F_{\text{impr}}(r_n) = \frac{13}{12}F_n(r_n) - \frac{1}{24}[F_n(r_n + a) + F_n(r_n - a)]$$

$$(\tilde{r}_I)^{-2} = \frac{13}{12}(r_I(r))^{-2} - \frac{1}{24}[(r_I(r+a))^{-2} + (r_I(r-a))^{-2}]$$

$$\alpha_{\text{qq}}(1/r, a) = \alpha_{\text{qq}}(1/r, 0) \left\{ 1 + (a/r)^2 [\alpha(1/a)]^{-\hat{\gamma}_0} A_1(r) [1 + \mathcal{O}([\alpha(1/a)]^{-\Delta_{\hat{\gamma}}})] \right\} + \mathcal{O}(a^4)$$

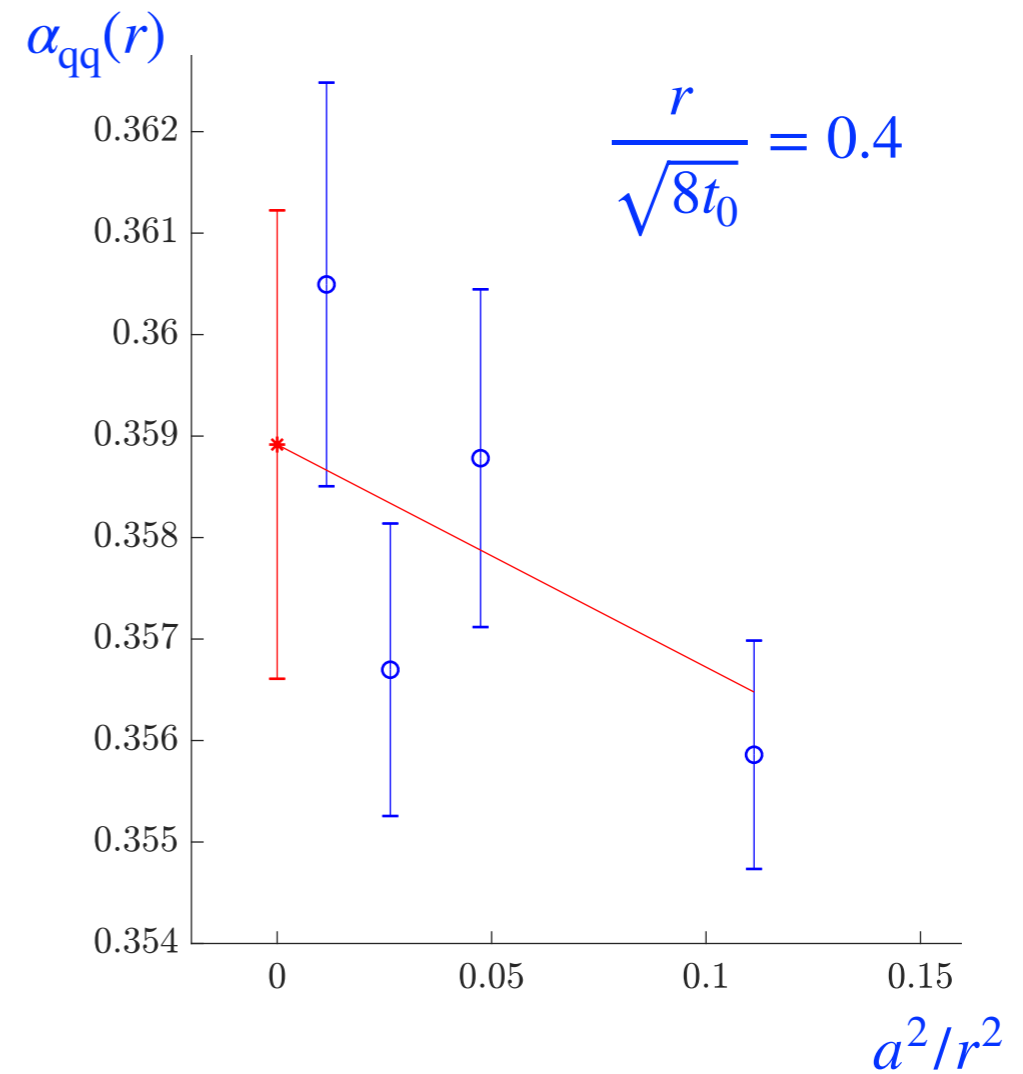
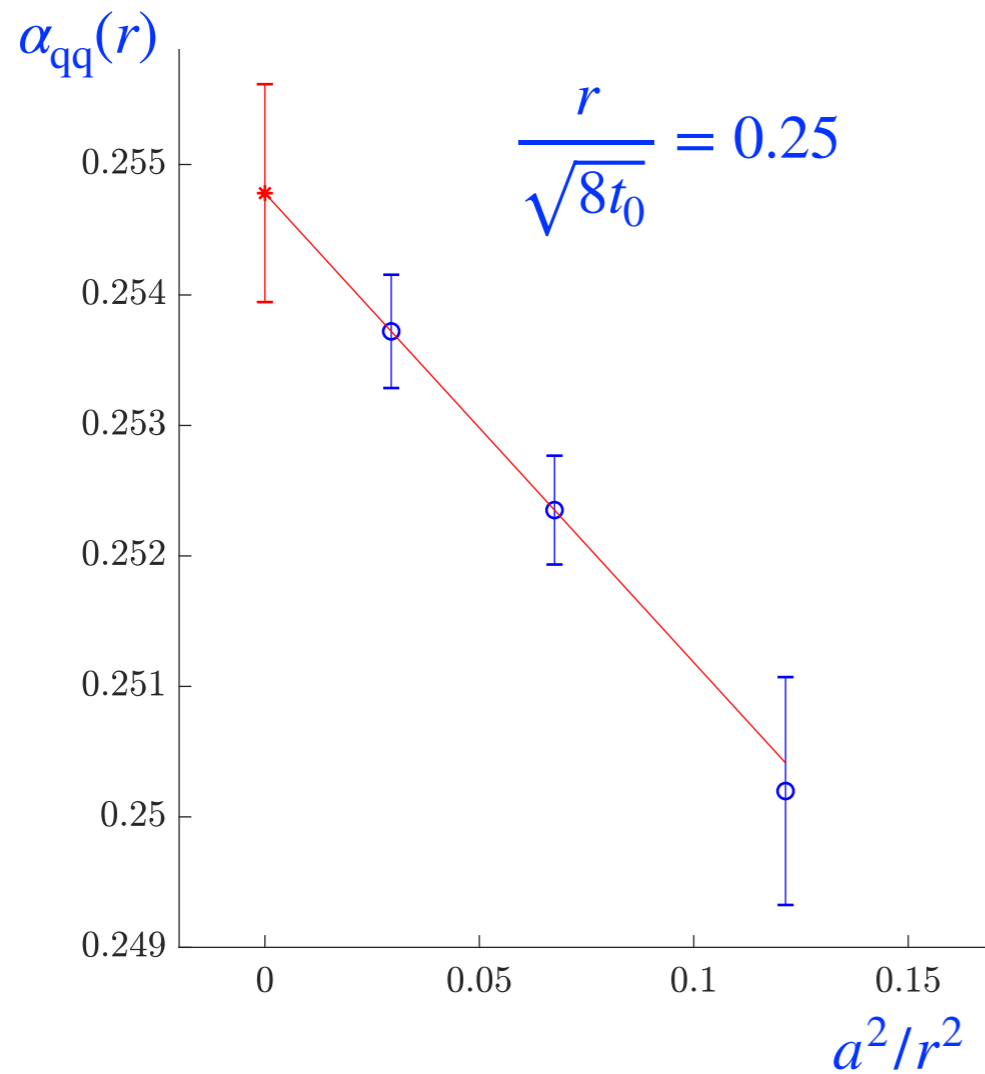
$$A_i(r) = \bar{A}_i \alpha(1/r)^{1+\hat{\gamma}_i} [1 + \mathcal{O}(\alpha(1/r))].$$

$$-\hat{\gamma}_0 = 7/11 = 0.686$$

smallest eigenvalue of 1-loop anomalous dimension matrix
of Symanzik EFT d=6 operator basis [new, talk by N. Husung]

Continuum limits

- ▶ Large r region ($r > 0.1\text{fm}$)



- ▶ Gradient flow: log-corrections to a^2 not yet known.

Continuum limits

- ▶ Small r region ($r < 0.1 \text{ fm}$), step scaling functions

(1) select $\alpha_i \in \{0.205, 0.215, \dots, 0.5\}$ (qq-scheme)

Continuum limits

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(1) select $\alpha_i \in \{0.205, 0.215, \dots, 0.5\}$ (qq-scheme)

(2) fit $\Sigma(\alpha_i, s, a/r) = \sigma_i + (a/sr)^2 [\alpha(1/a)]^{7/11} \rho_i$

$\alpha(1/a)$ from $\alpha_{\text{qq}}(1/(2.5a))$ by 4-loop running

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(3) fit slopes to

$$\rho_i = \rho(\alpha_i), \quad \rho(\alpha) = \alpha^{(2-7/11)} \times [\rho^{(0)} + \rho^{(1)}\alpha + \rho^{(2)}\alpha^2]$$

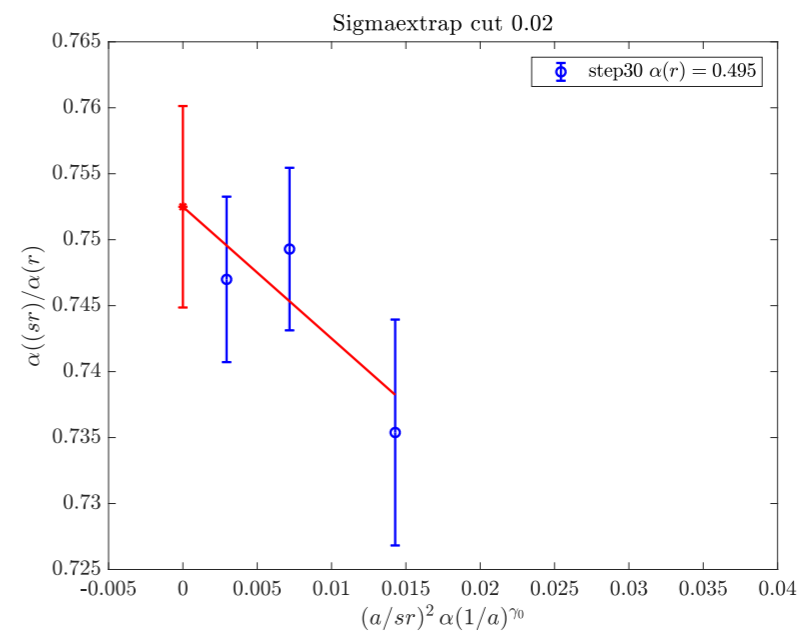
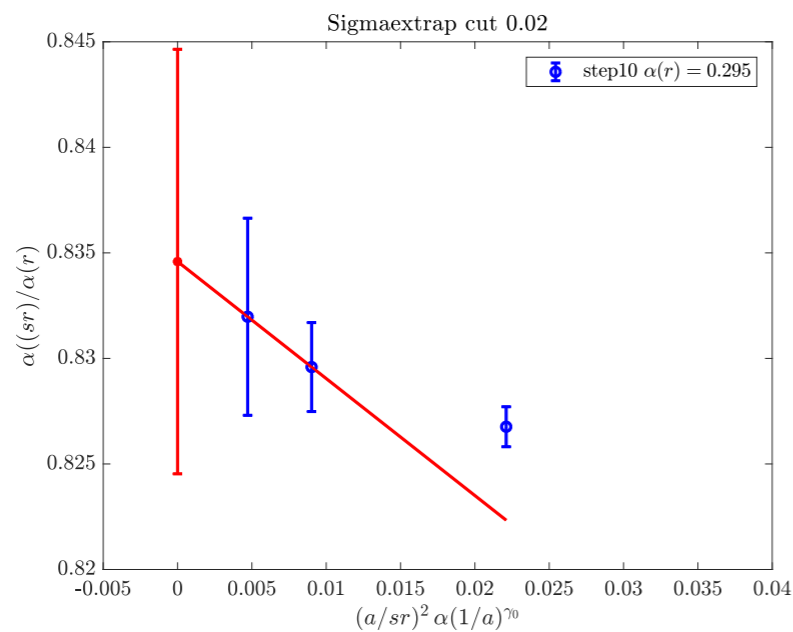
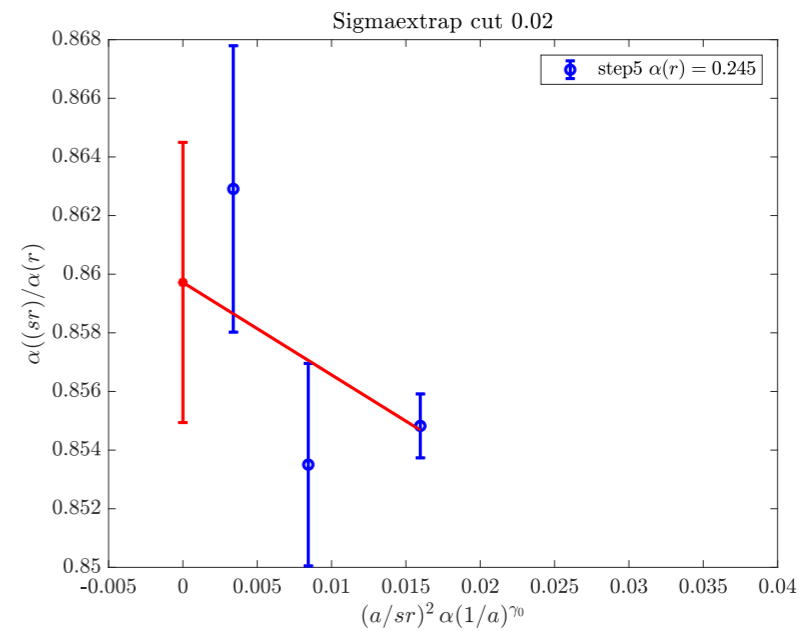
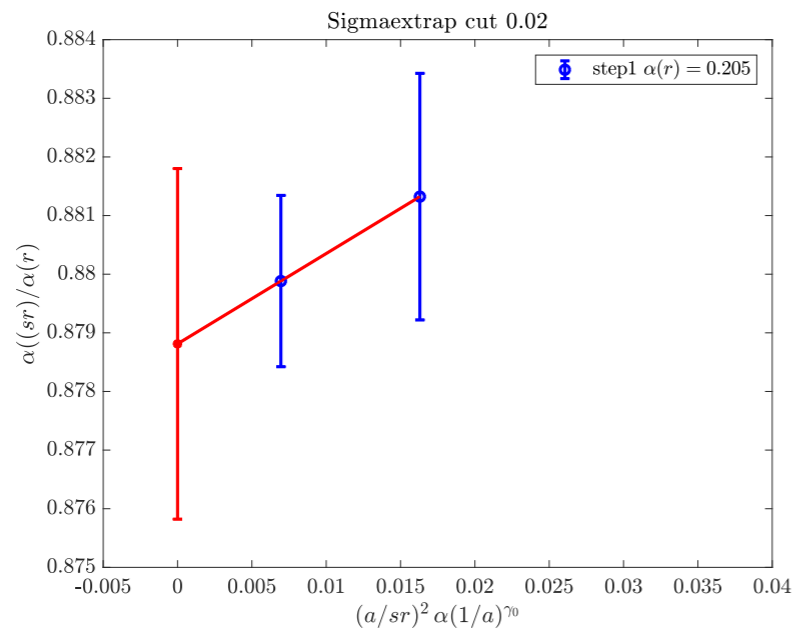
(4) use fitted slope function $\rho(\alpha)$ in

$$\Sigma(\alpha, s, a/r) = \sigma(\alpha, s) + (a/sr)^2 [\alpha(1/a)]^{7/11} \rho(\alpha)$$

for all α in range

Continuum limits

(2) fit $\Sigma(\alpha_i, s, a/r) = \sigma_i + (a/sr)^2 [\alpha(1/a)]^{7/11} \rho_i$



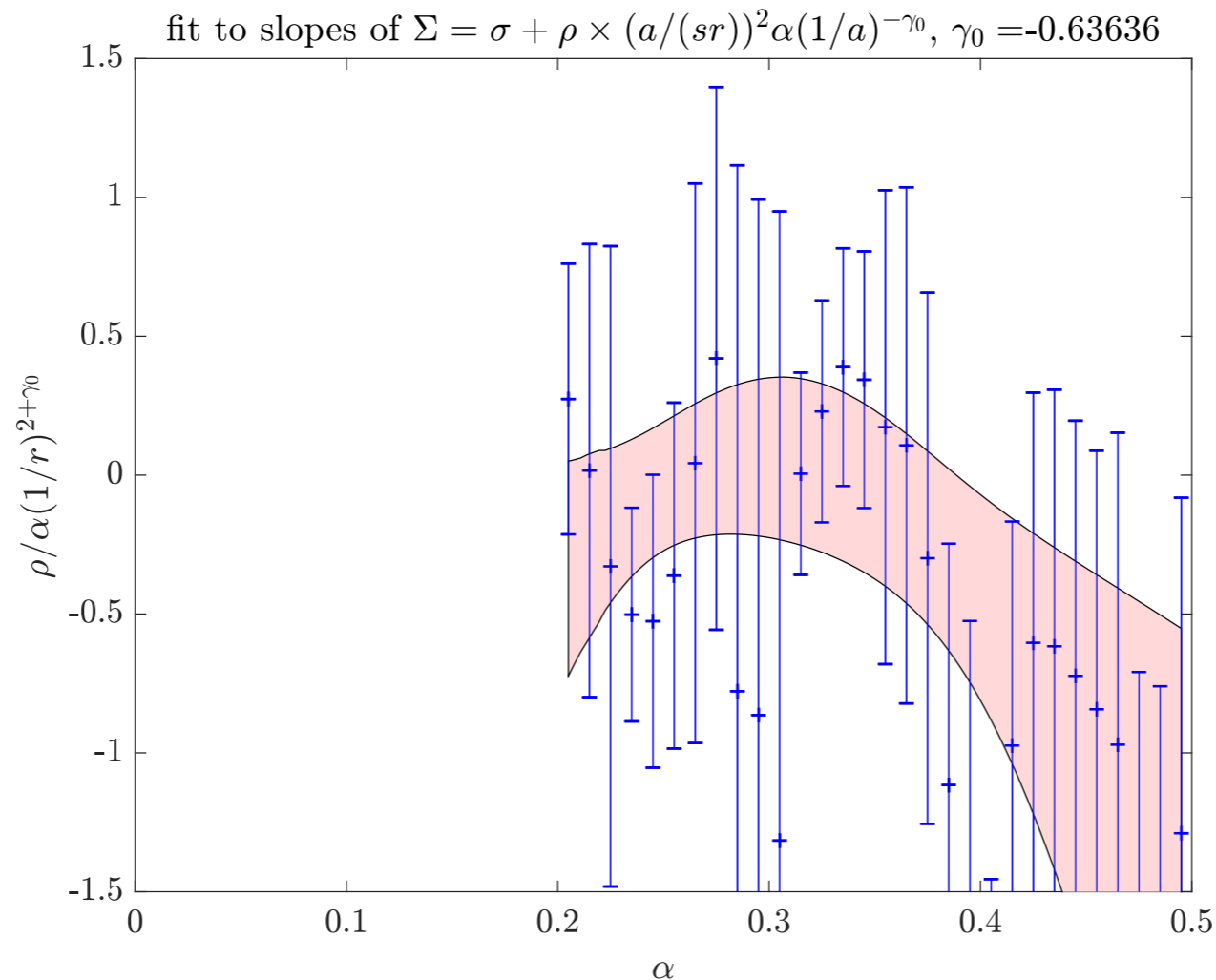
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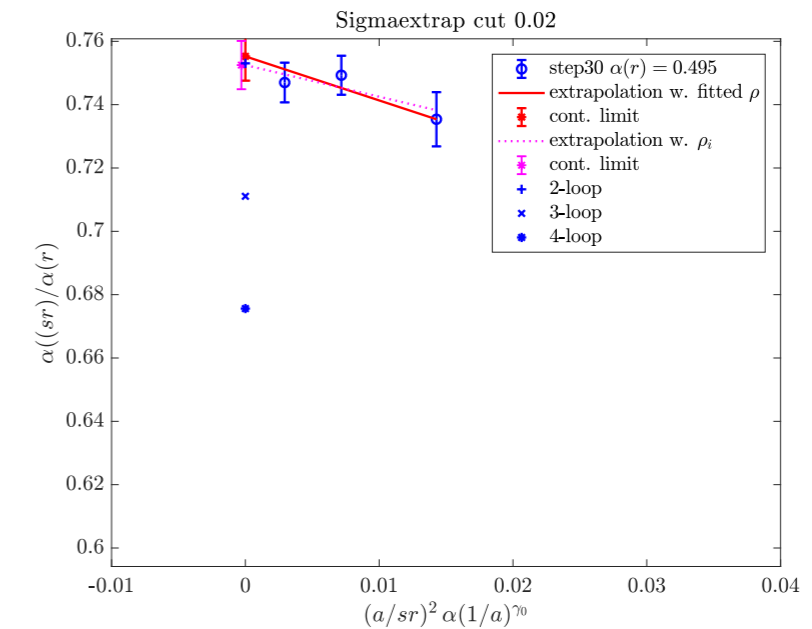
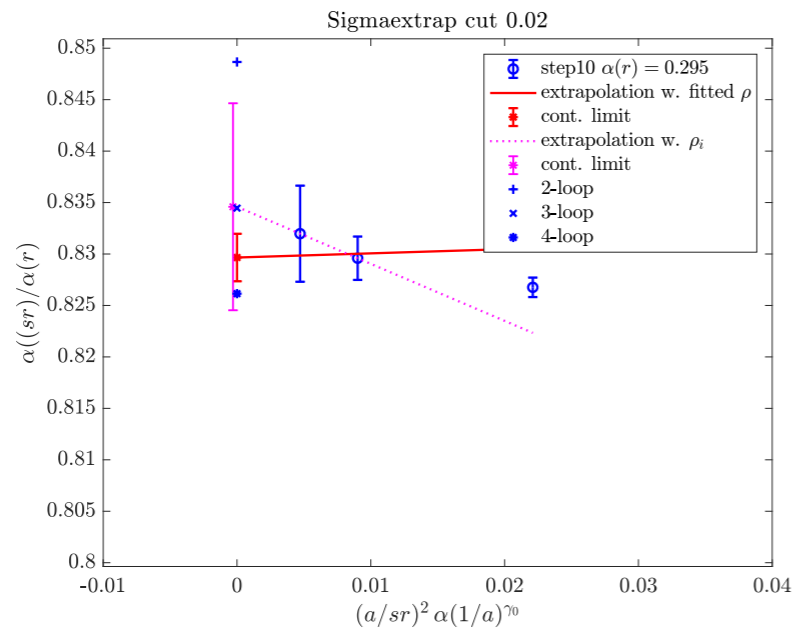
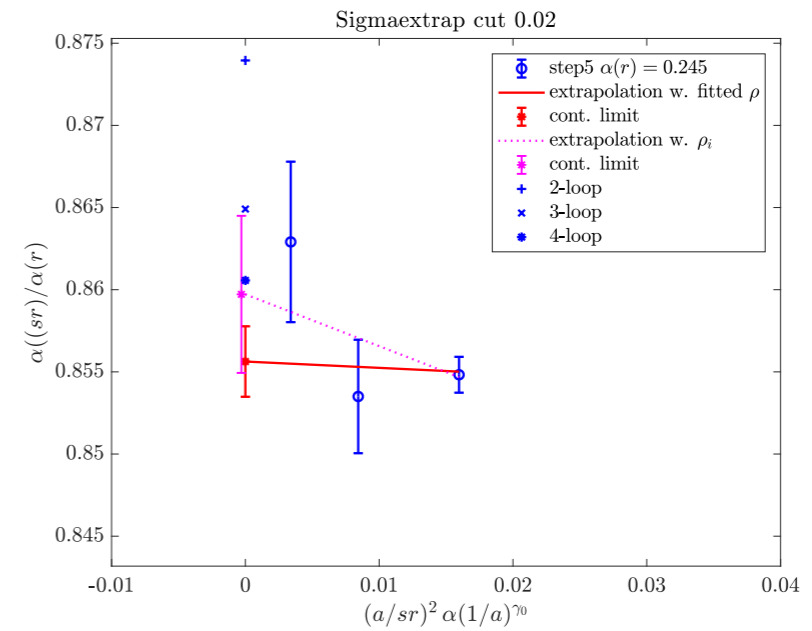
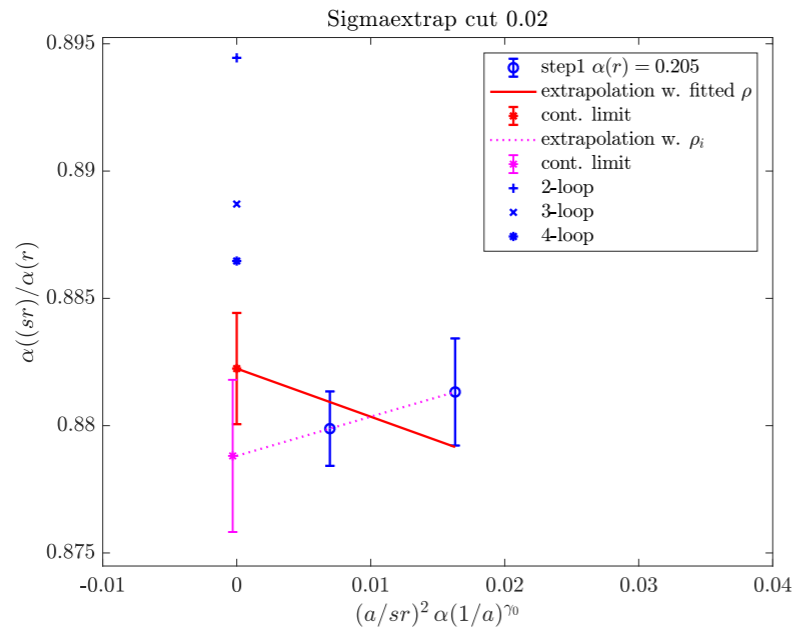
- fit $\Sigma(\alpha_i, s, a/r) = \sigma_i + (a/sr)^2 [\alpha(1/a)]^{7/11} \rho_i$ $\alpha(1/a)$ from $\alpha_{\text{qq}}(1/(2.5a))$ by 4-loop running

(3) fit slopes to $\rho_i = \rho(\alpha_i)$, $\rho(\alpha) = \alpha^{(2-7/11)} \times [\rho^{(0)} + \rho^{(1)}\alpha + \rho^{(2)}\alpha^2]$



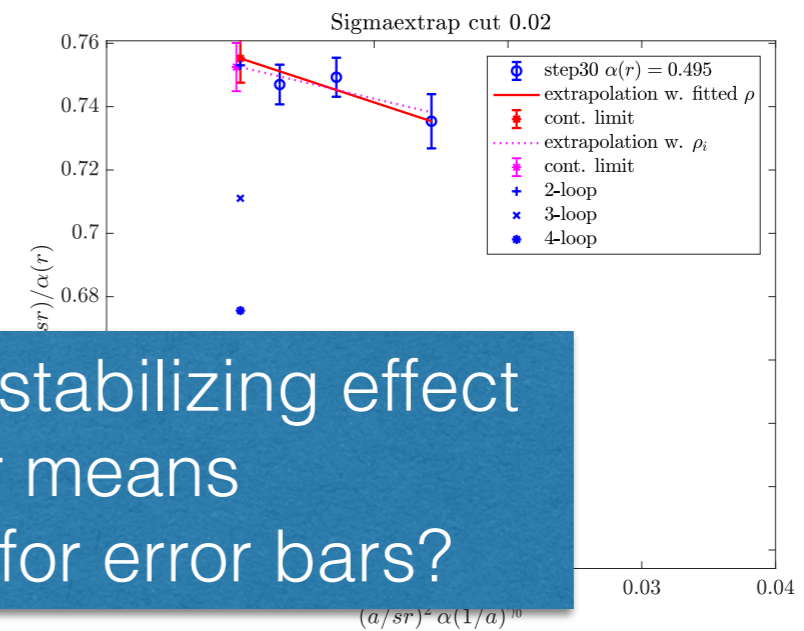
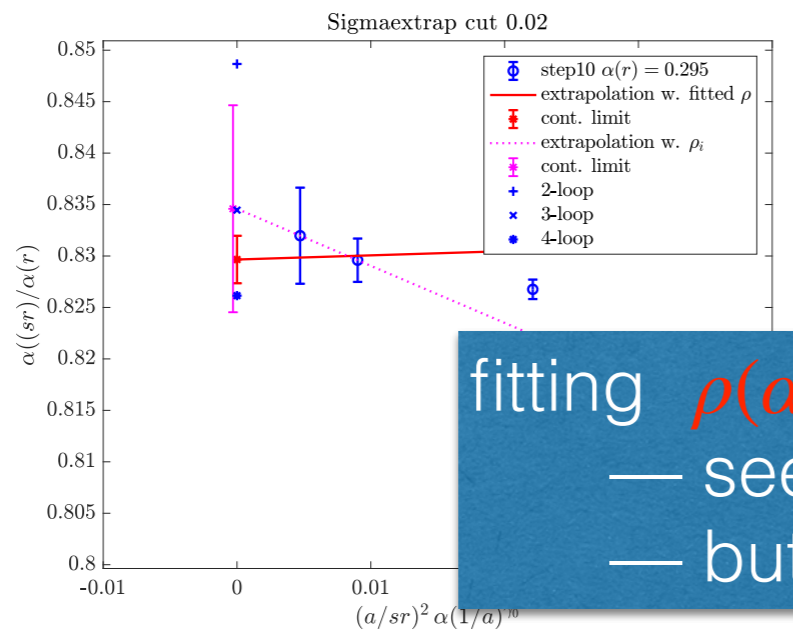
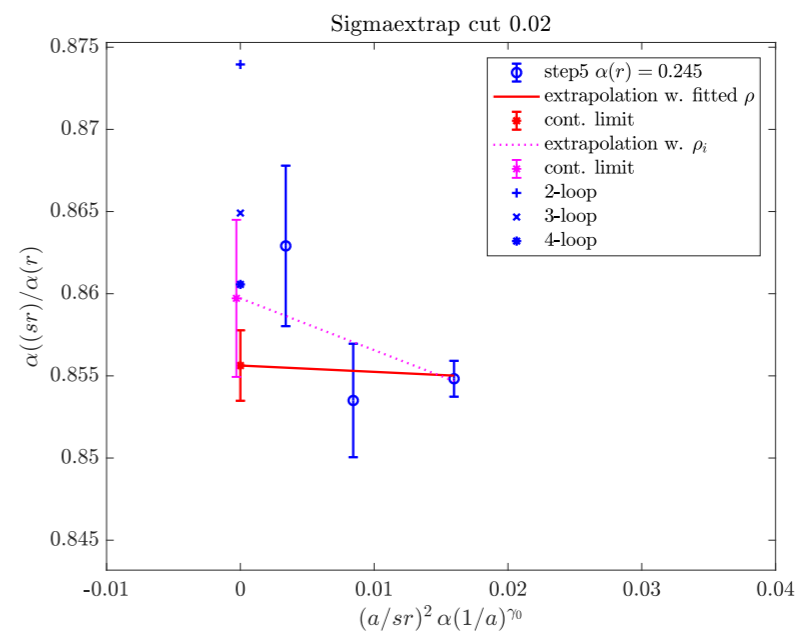
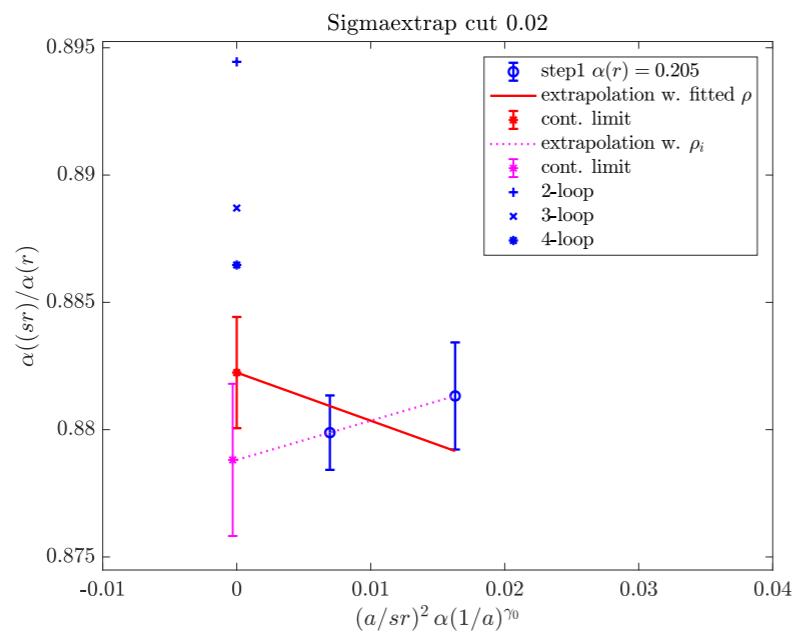
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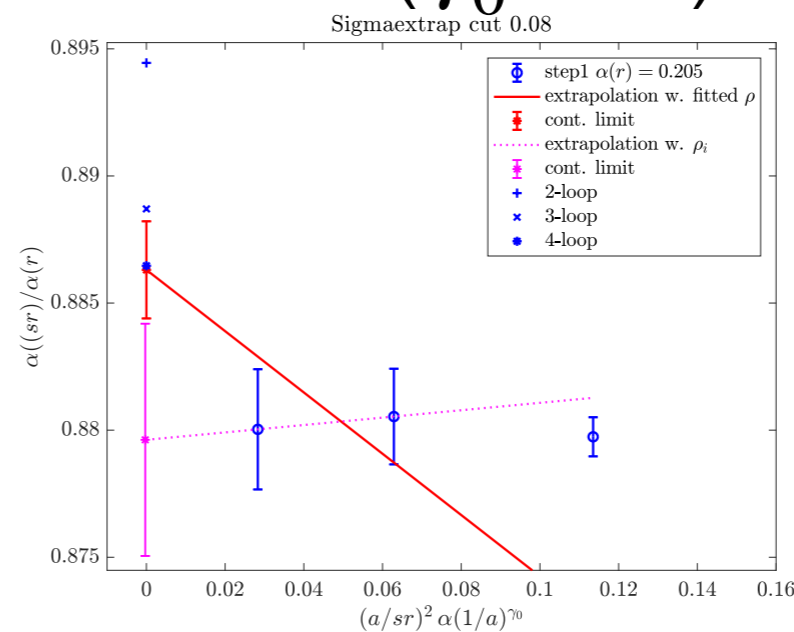
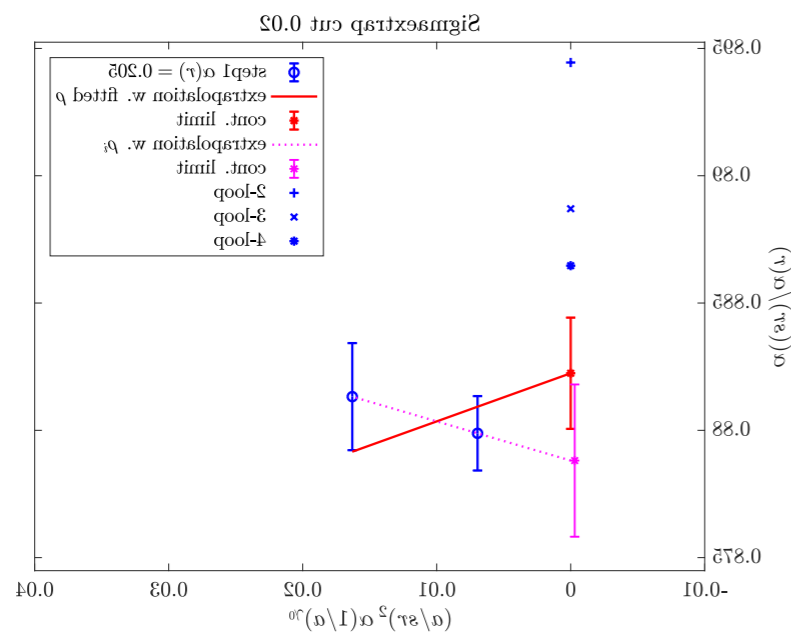


fitting $\rho(\alpha)$ first has a stabilizing effect
 — seems good for means
 — but trustworthy for error bars?

Continuum limits: compare to standard derivative

a^2 improved derivative

standard derivative
($\gamma_0 = 0$)

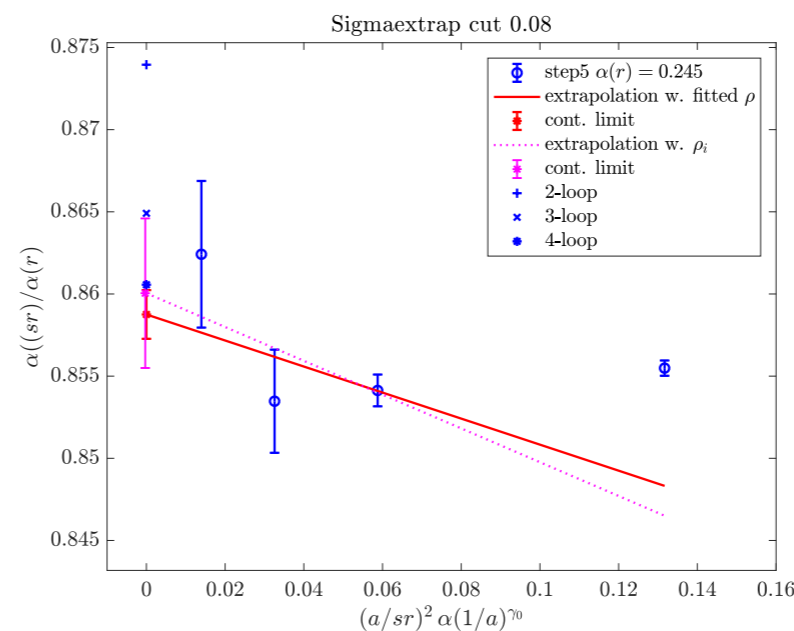
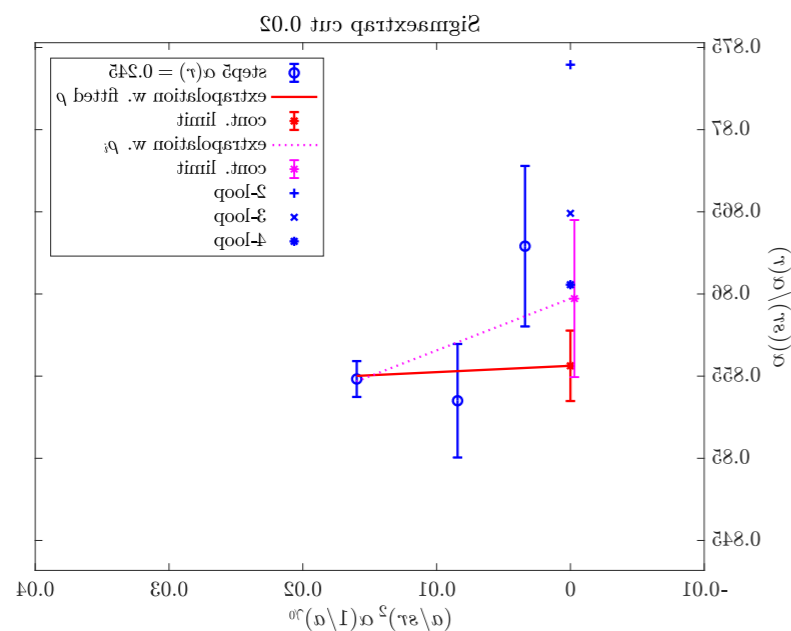


more points,
but lost
by cut

$$\left(\frac{a}{sr}\right)^2 \leq 0.08$$

$$\left(\frac{sr}{a} > 3.5\right)$$

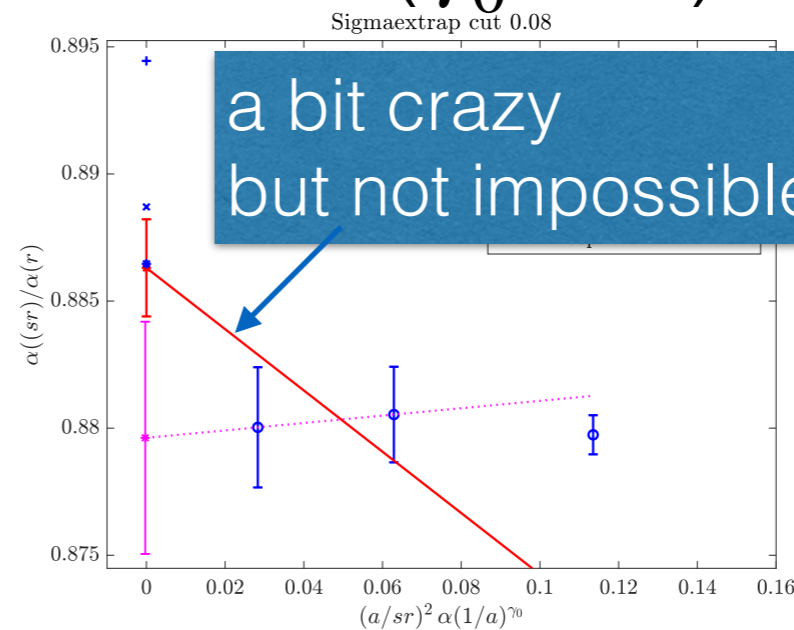
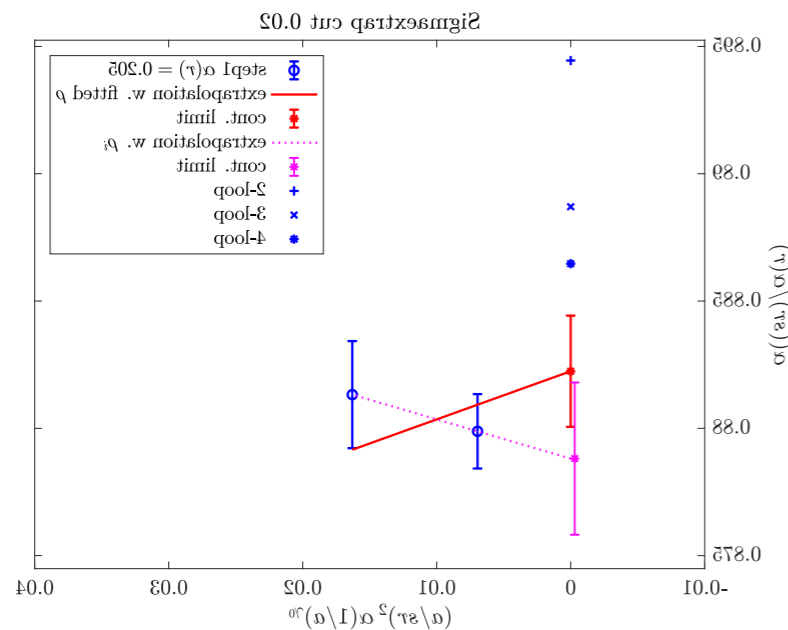
same
continuum
limits



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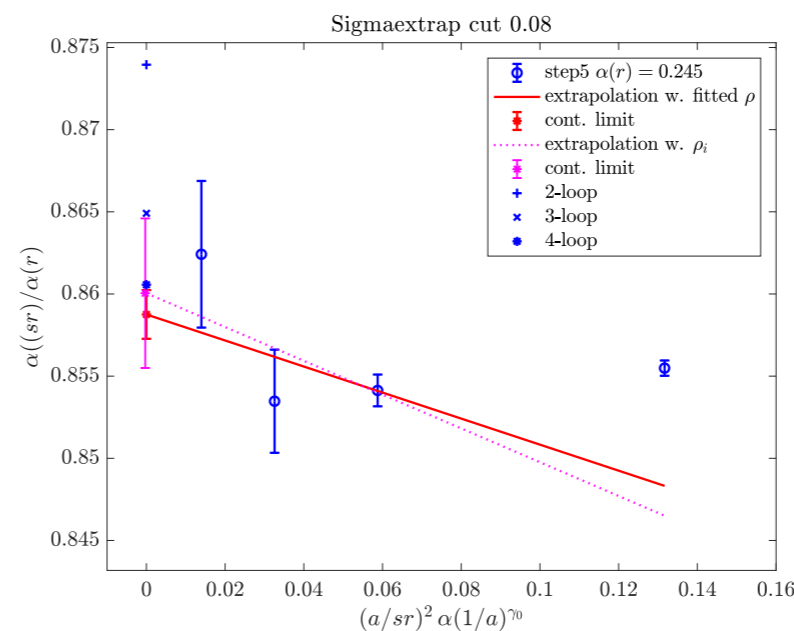
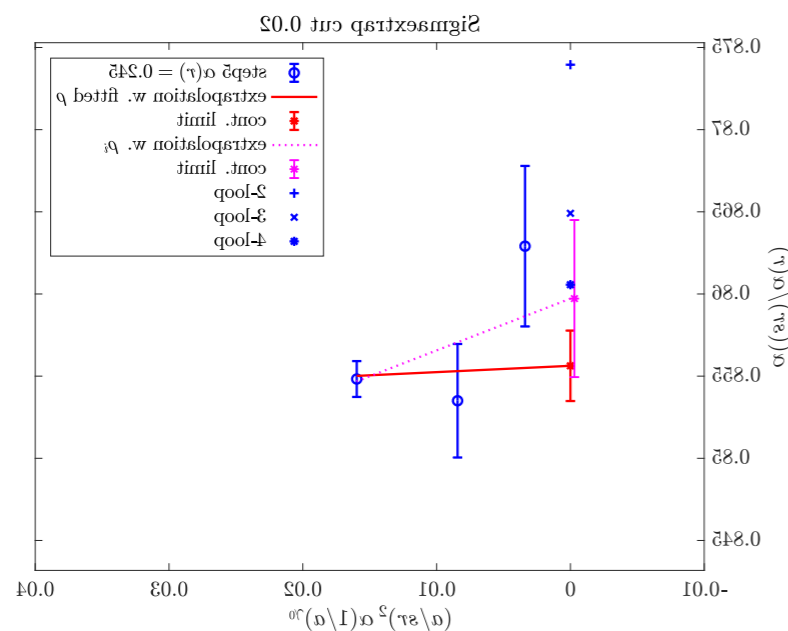
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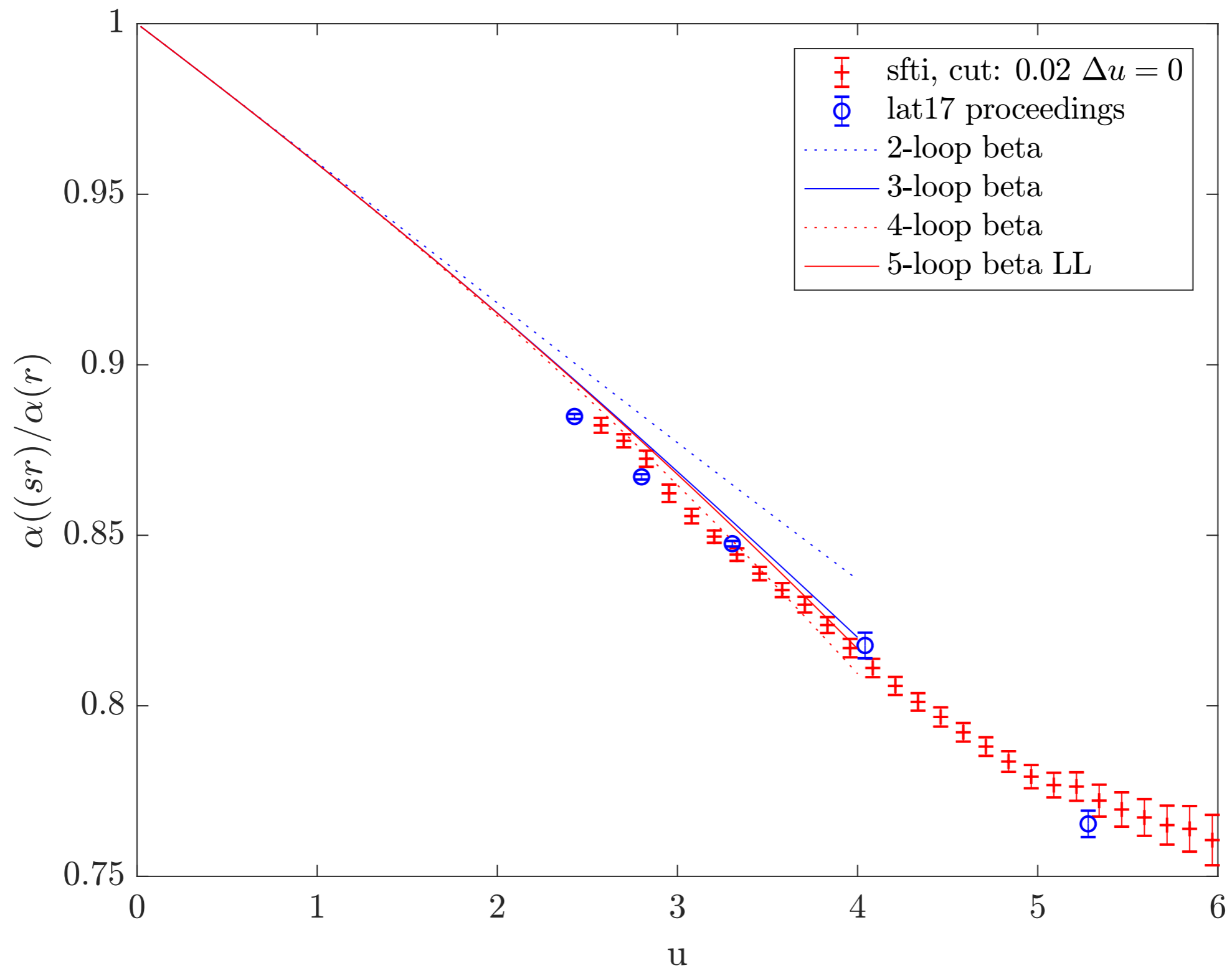
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same
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Results (from 2-stage continuum limit, improved derivative)



Λ - parameter

- ▶ $\Lambda_{\overline{\text{MS}}}$ from Λ_{qq} locally (α_{qq} by α_{qq})

$$\Lambda = \mu (b_0 \bar{g}(\mu)^2)^{-b_1/(2b_0^2)} e^{-1/(2b_0 \bar{g}(\mu)^2)} \exp \left\{ - \int_0^{\bar{g}(\mu)} dx \left[\frac{1}{\beta_{n\text{-loop}}(x)} + \frac{1}{b_0 x^3} - \frac{b_1}{b_0^2 x} \right] \right\} \\ \times [1 + \mathcal{O}(\bar{g}(\mu)^{2(n-1)})]$$

$$\beta_{3\text{-loop}}(g) = -g^3 [b_0 + b_1 g^2 + b_2 g^4]$$

$$\text{4-loop:} \quad +b_3 g^6 + b_{3L} g^6 \log(\alpha)$$

$$\text{4-loop LL:} \quad +b_{4L} g^8 \log(\alpha) + b_{4LL} g^8 [\log(\alpha)]^2$$

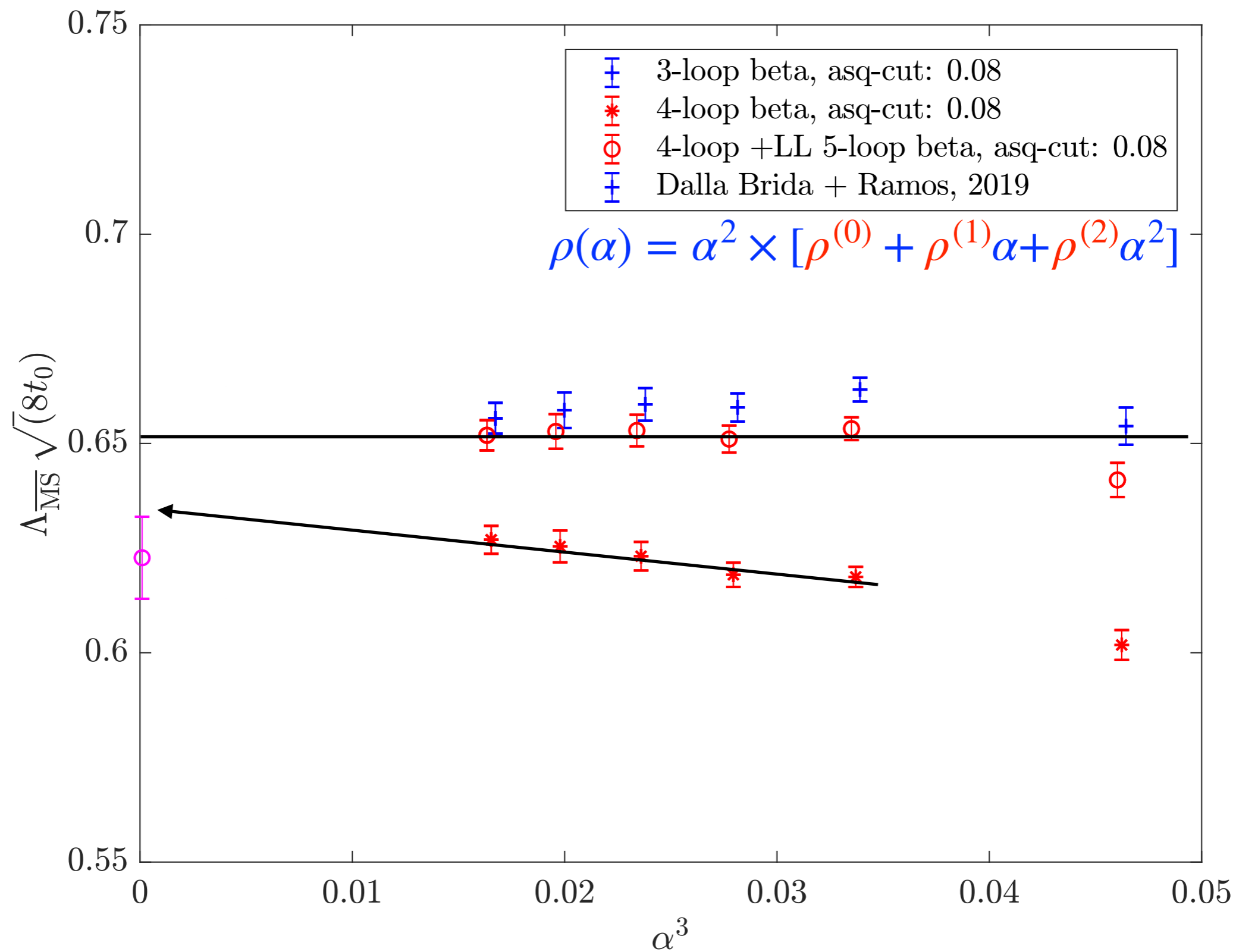
computed from

Peter 97; Schröder 99
Anzai, Kiyo, Sumino, 10

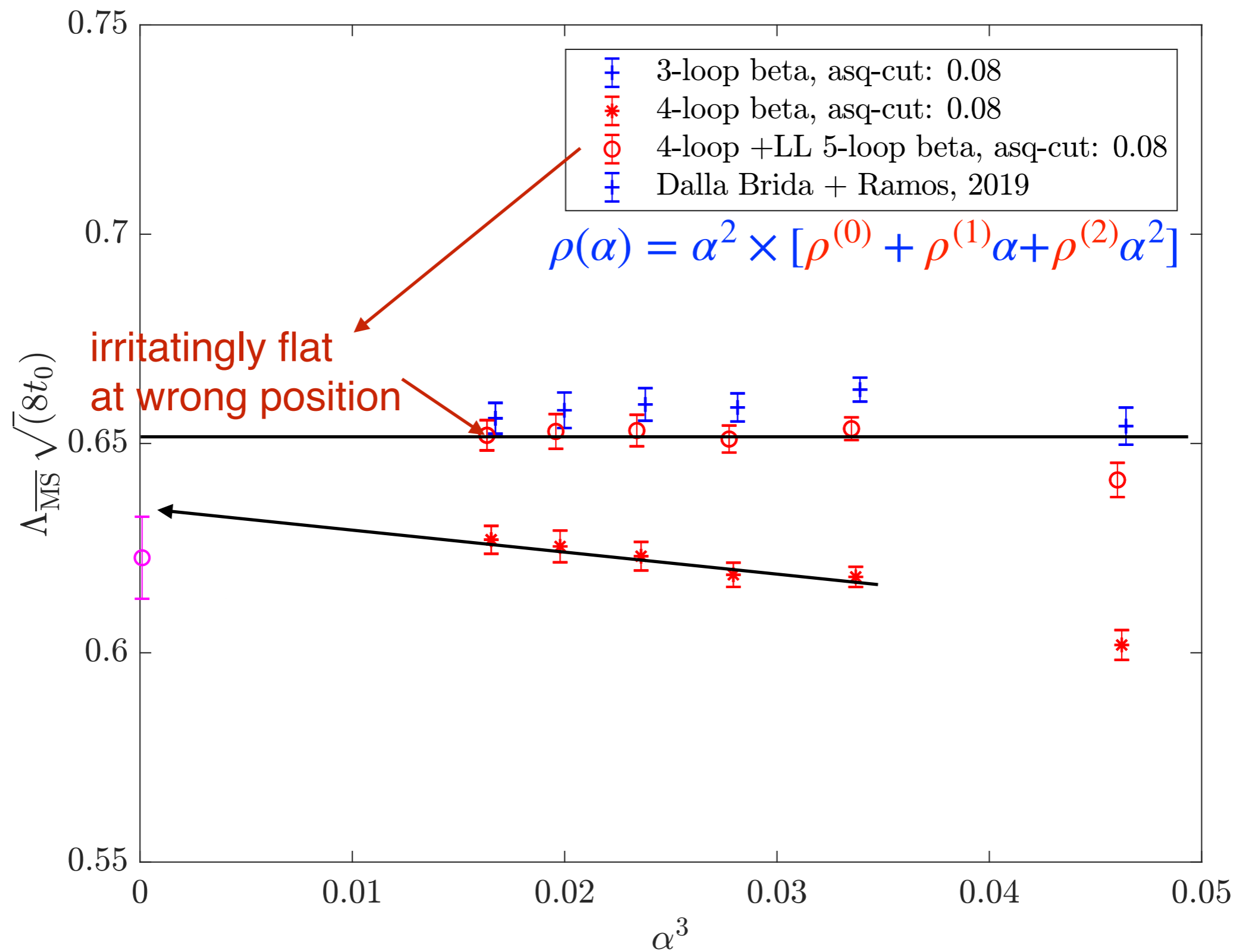
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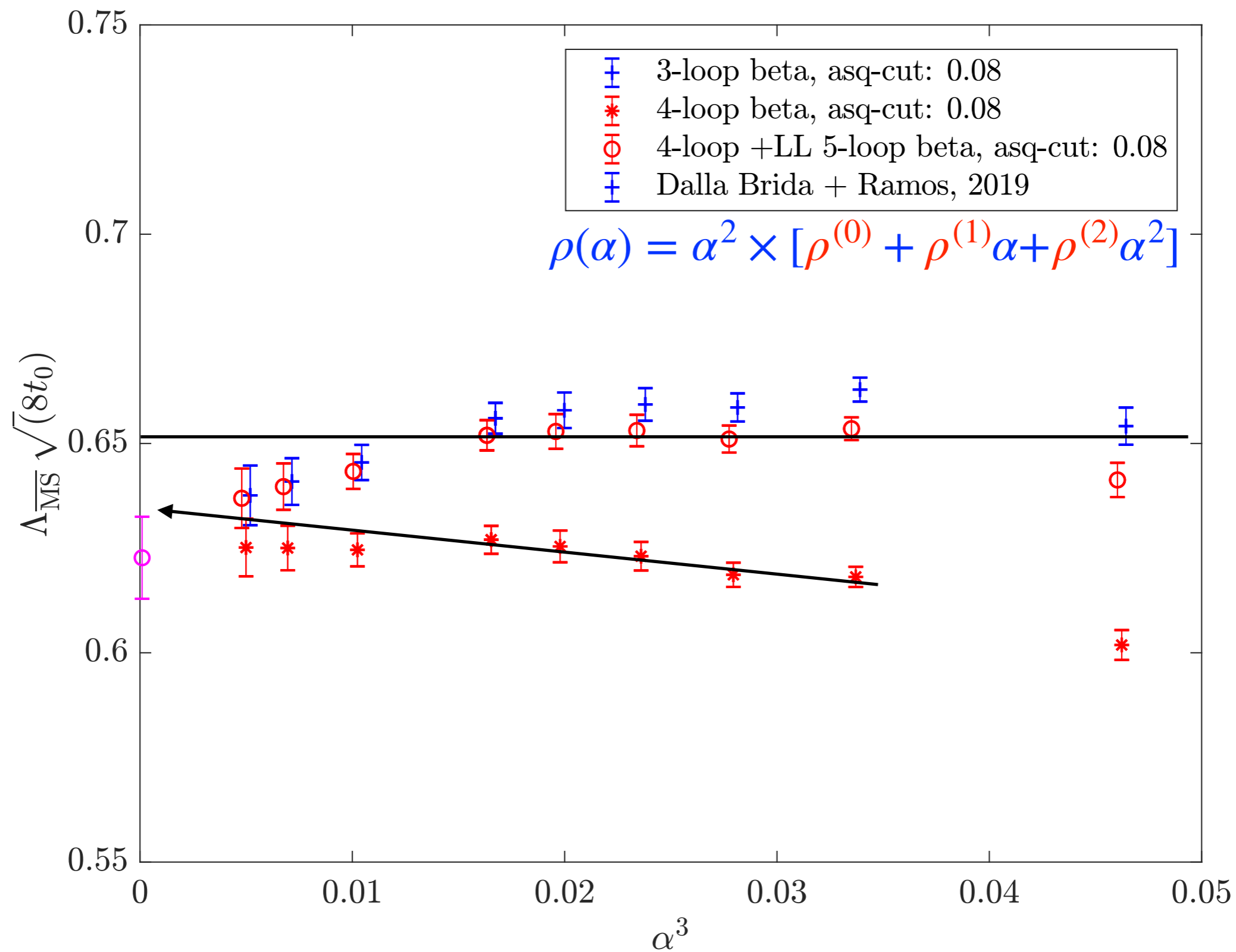
Results (from 2-stage continuum limit, standard derivative)



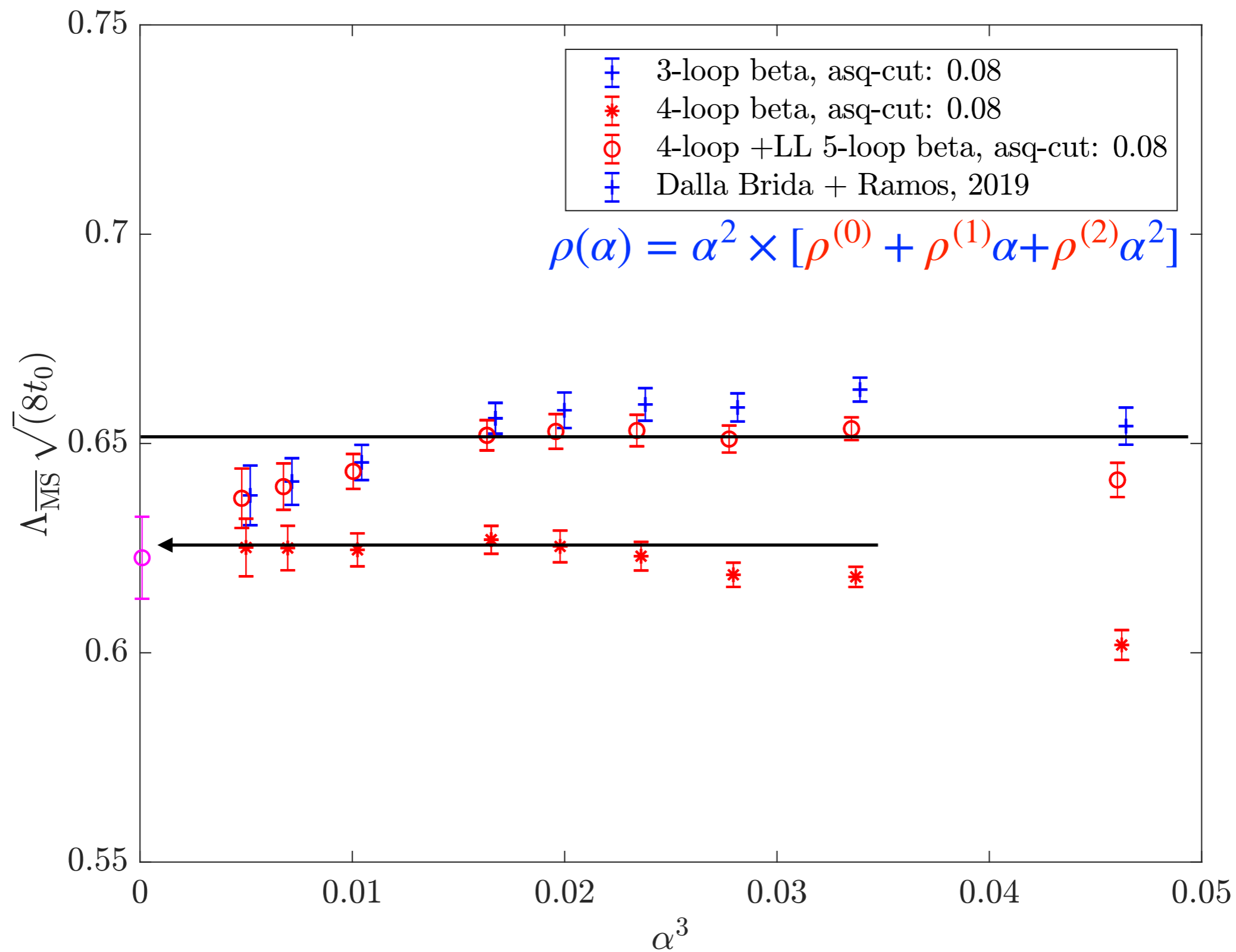
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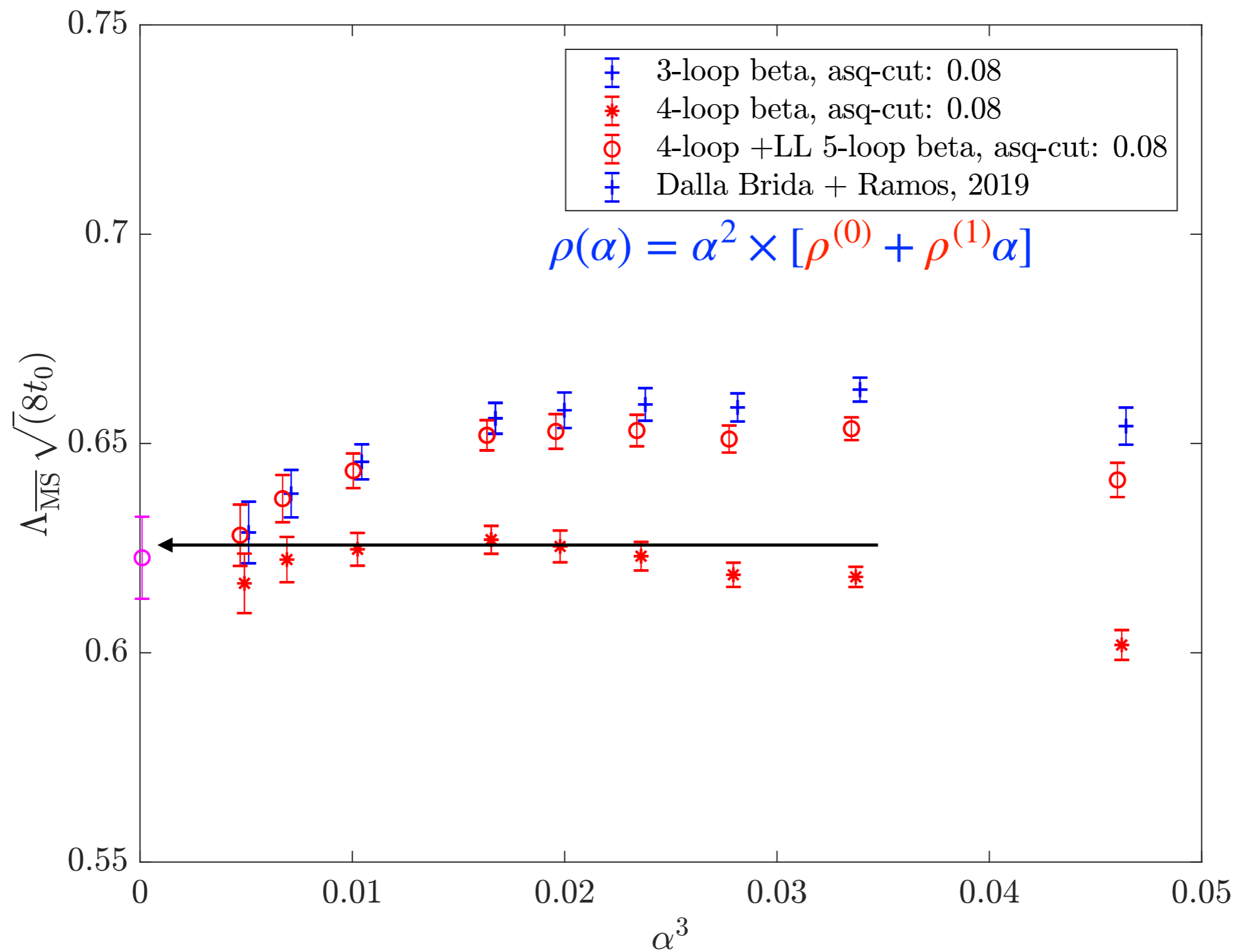
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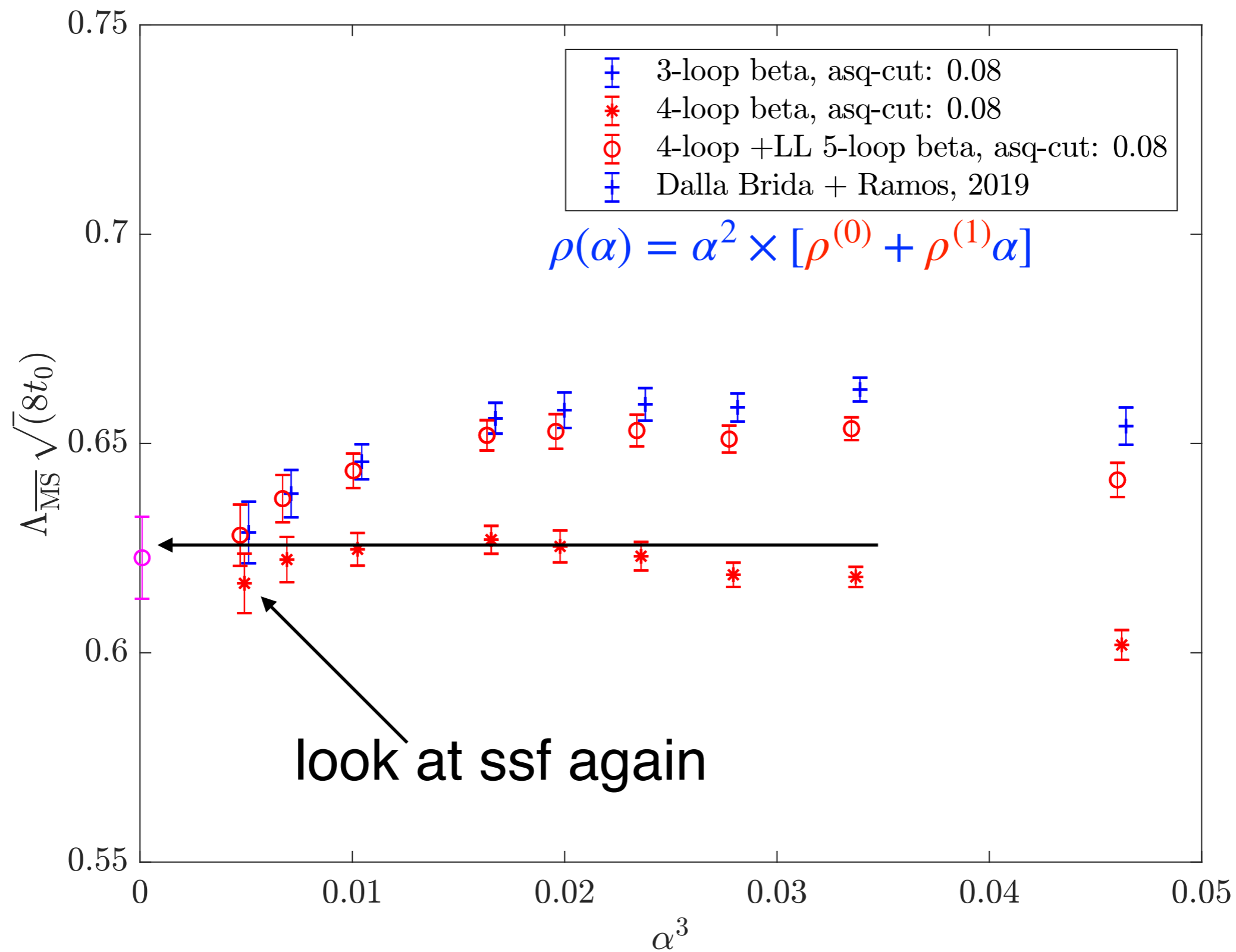
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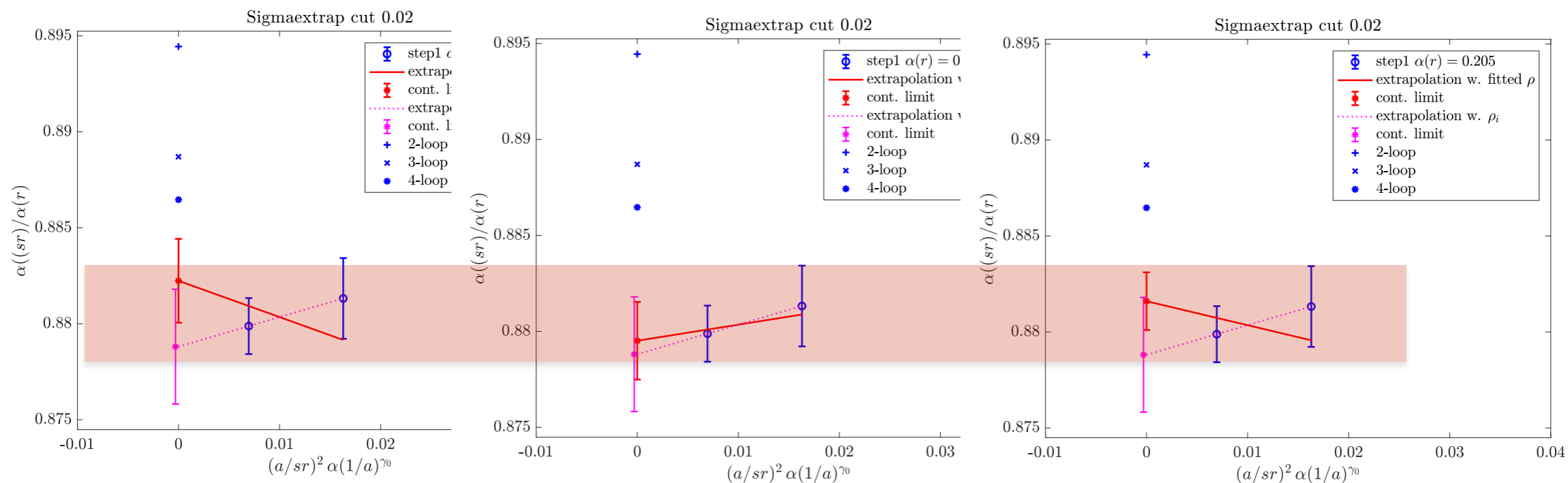
Results (from 2-stage continuum limit, standard derivative)



look at ssf again

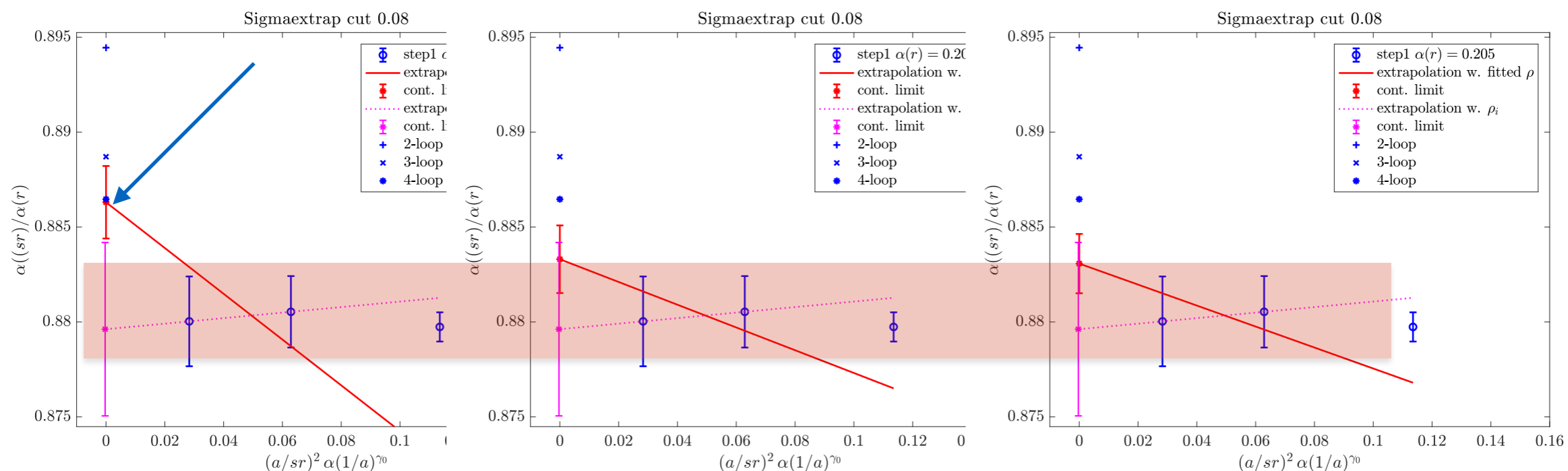
Continuum limits: compare to standard derivative

a^2 improved derivative



$\gamma_0 = 7/11$

standard derivative

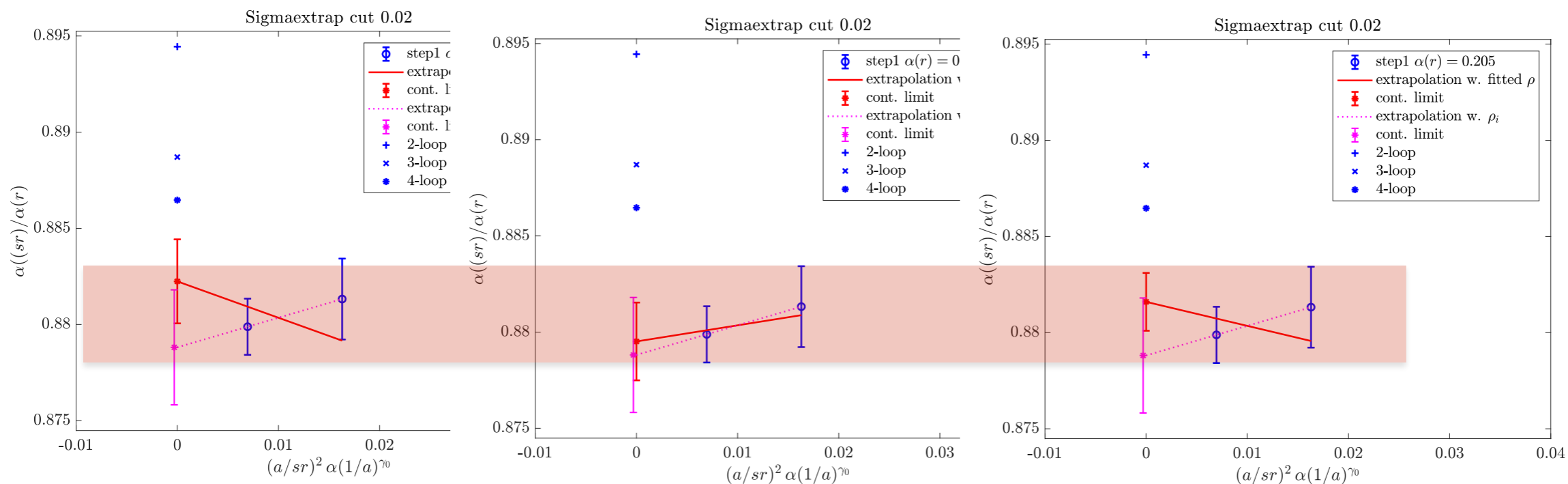


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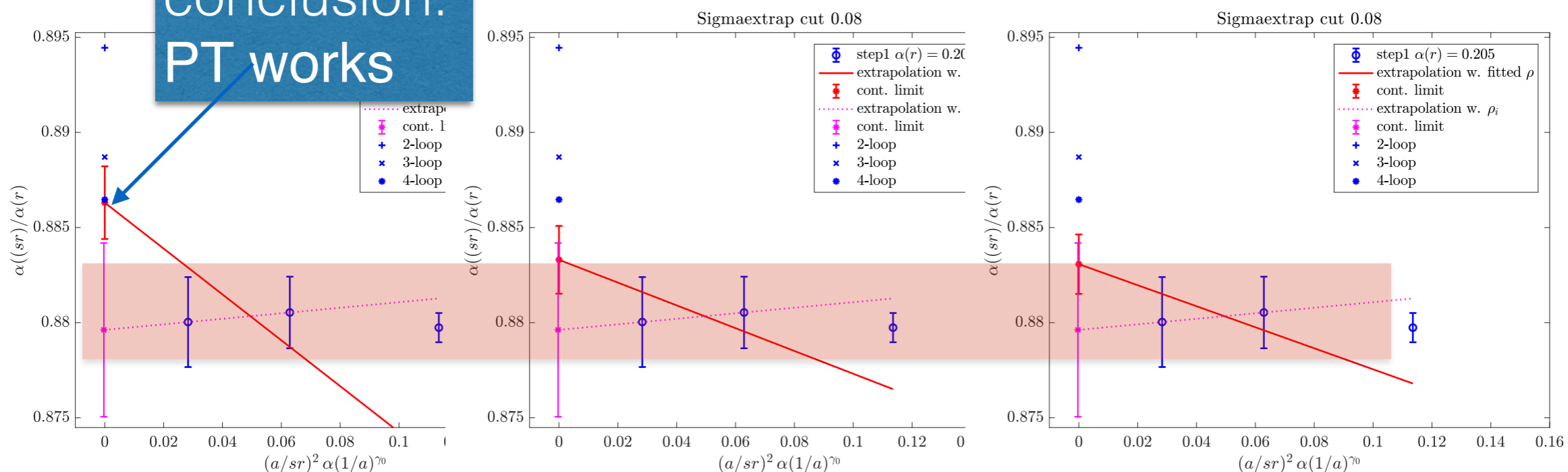
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conclusion:
PT works



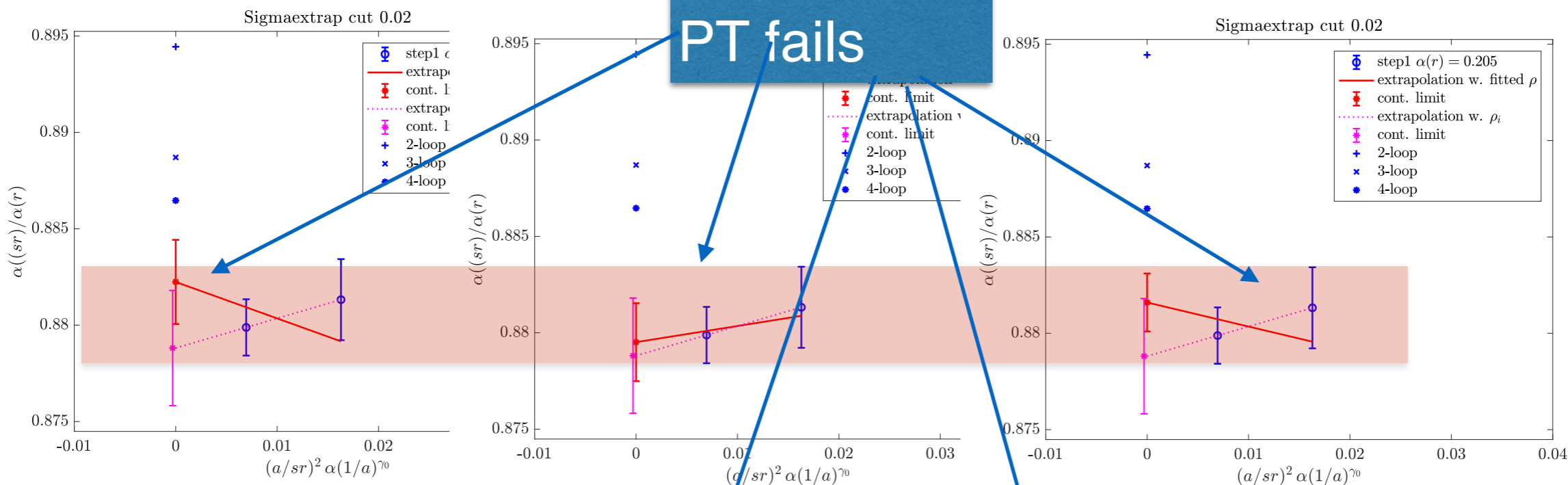
$$\gamma_0 = 0$$



Continuum limits: compare to standard derivative

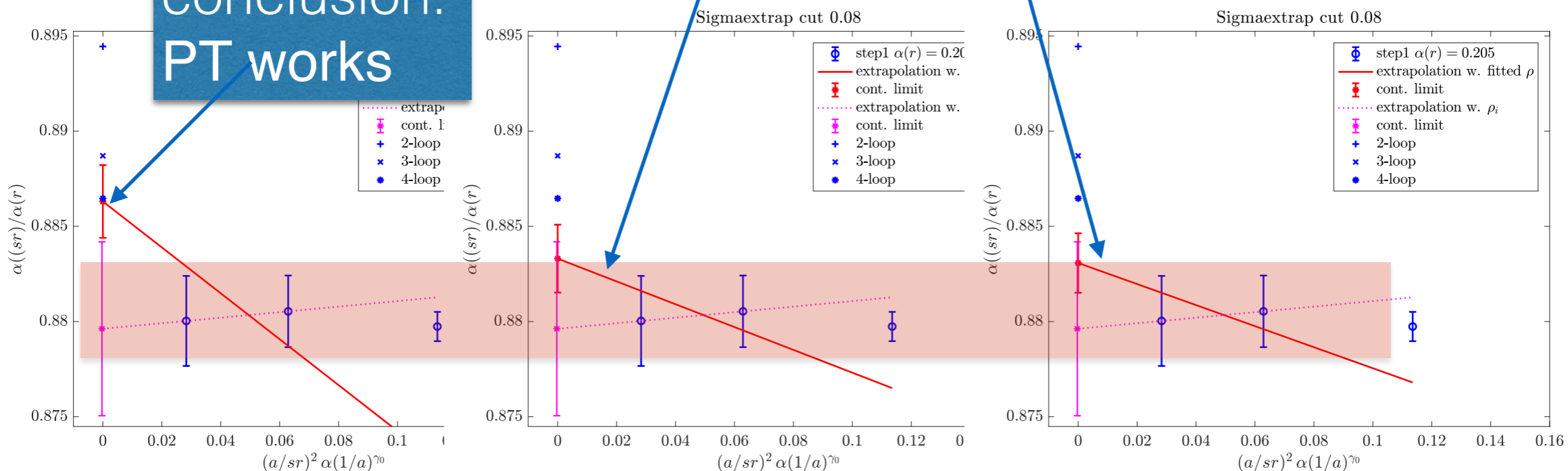
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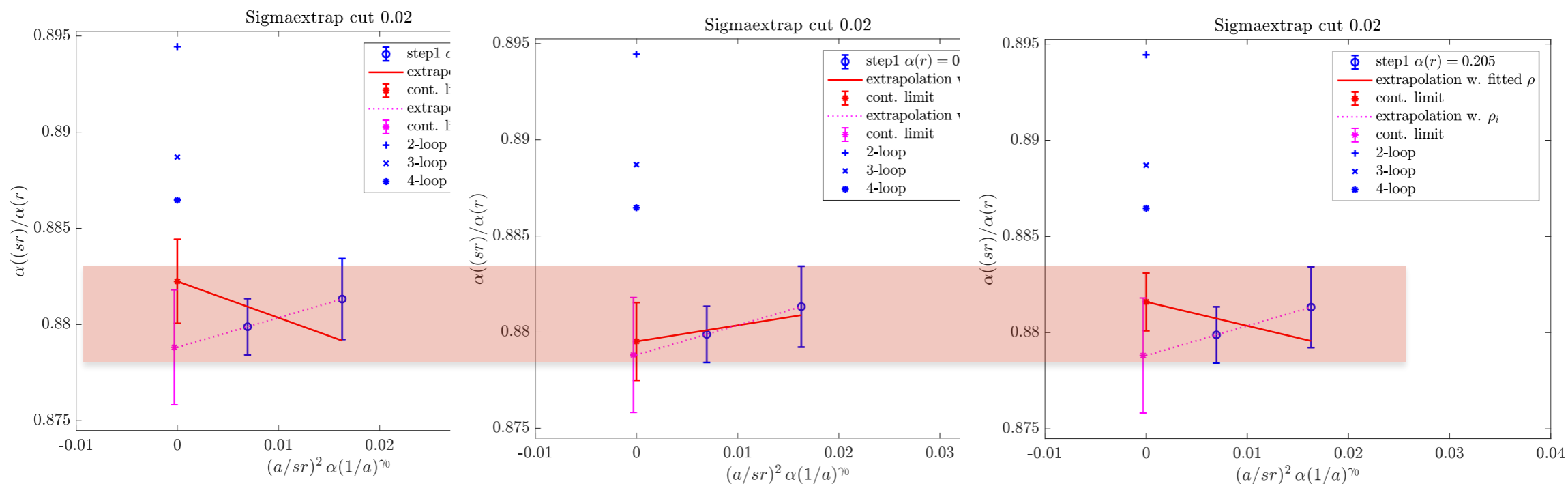


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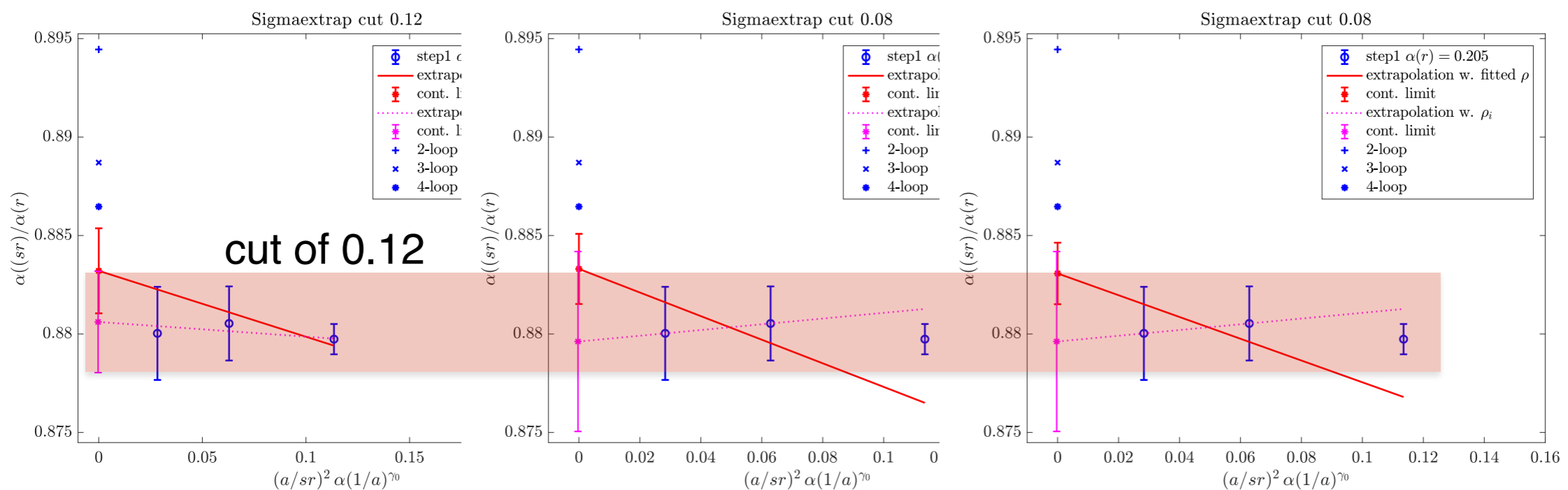


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Conclusions

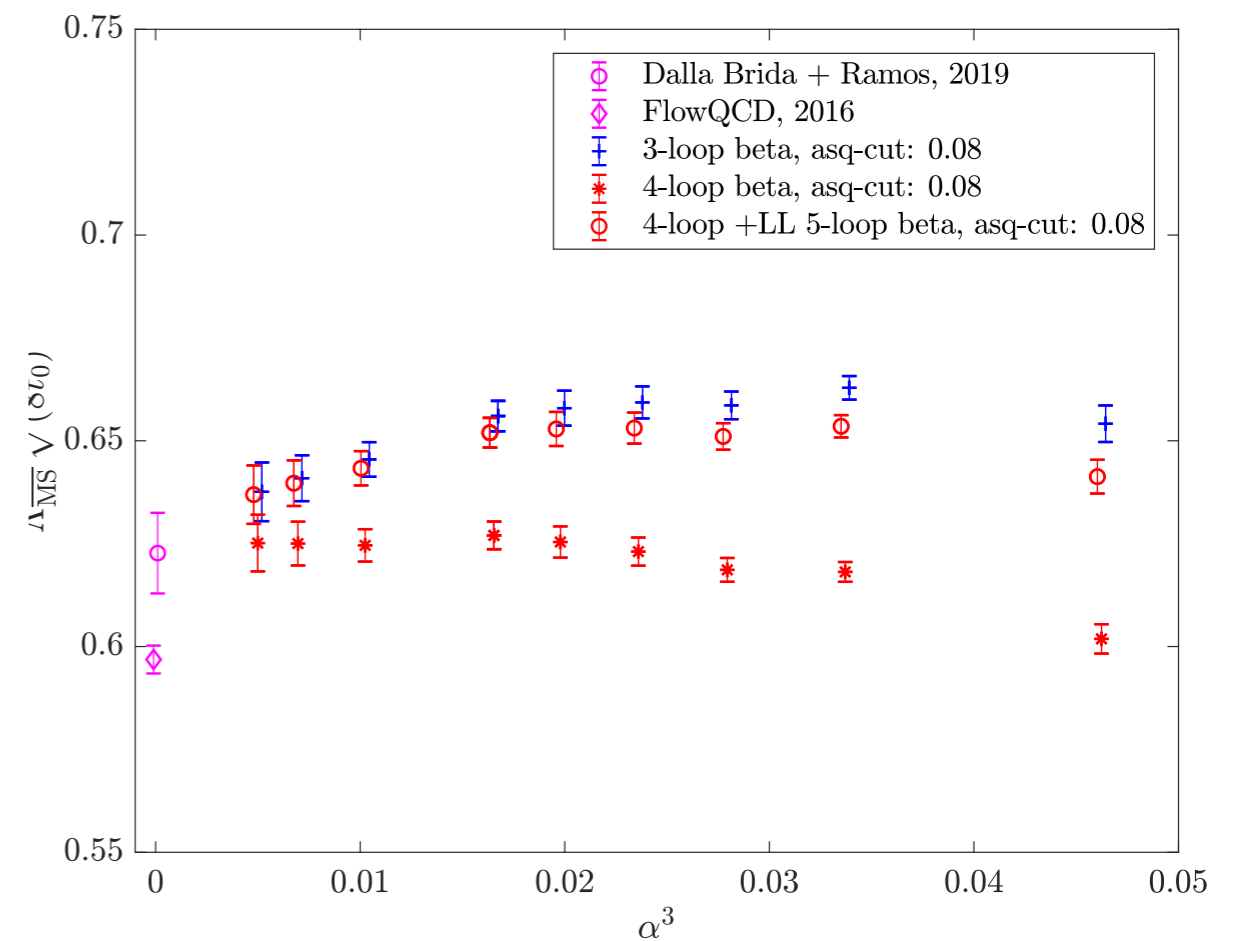
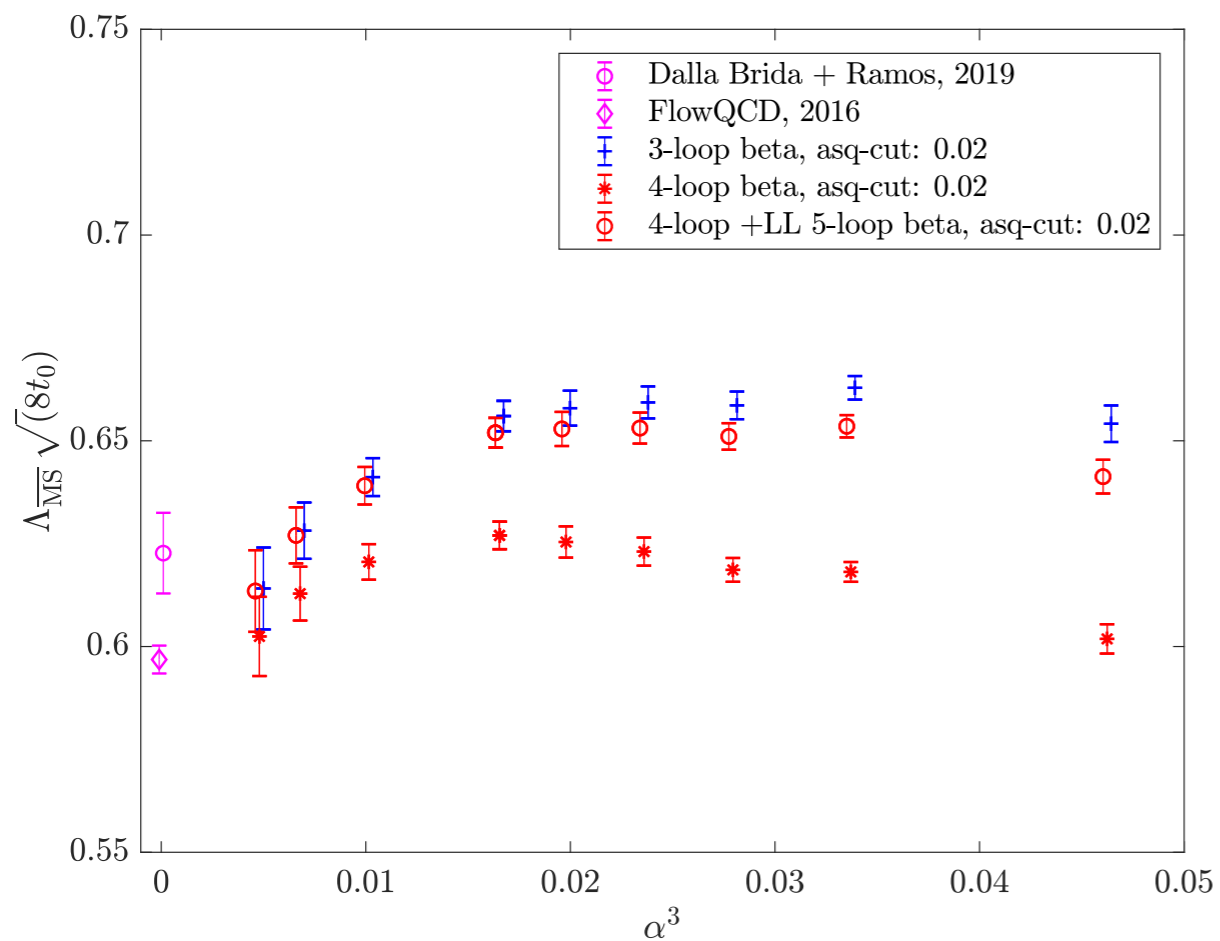
- ▶ Continuum limit with known leading first log-correction to a^2 scaling

$$a^2 [\alpha(1/a)]^{7/11}$$

- ▶ Semiquantitative agreement with perturbation theory is convincing for distances below 0.1 fm
- ▶ Precision test of PT is very difficult
 - do known US contributions apply / help for accessible α ?
 - can continuum limit be controlled sufficiently well?
 - **uncertainty in Λ very difficult to assess** even with $a=0.01$ fm lattice

Conclusions

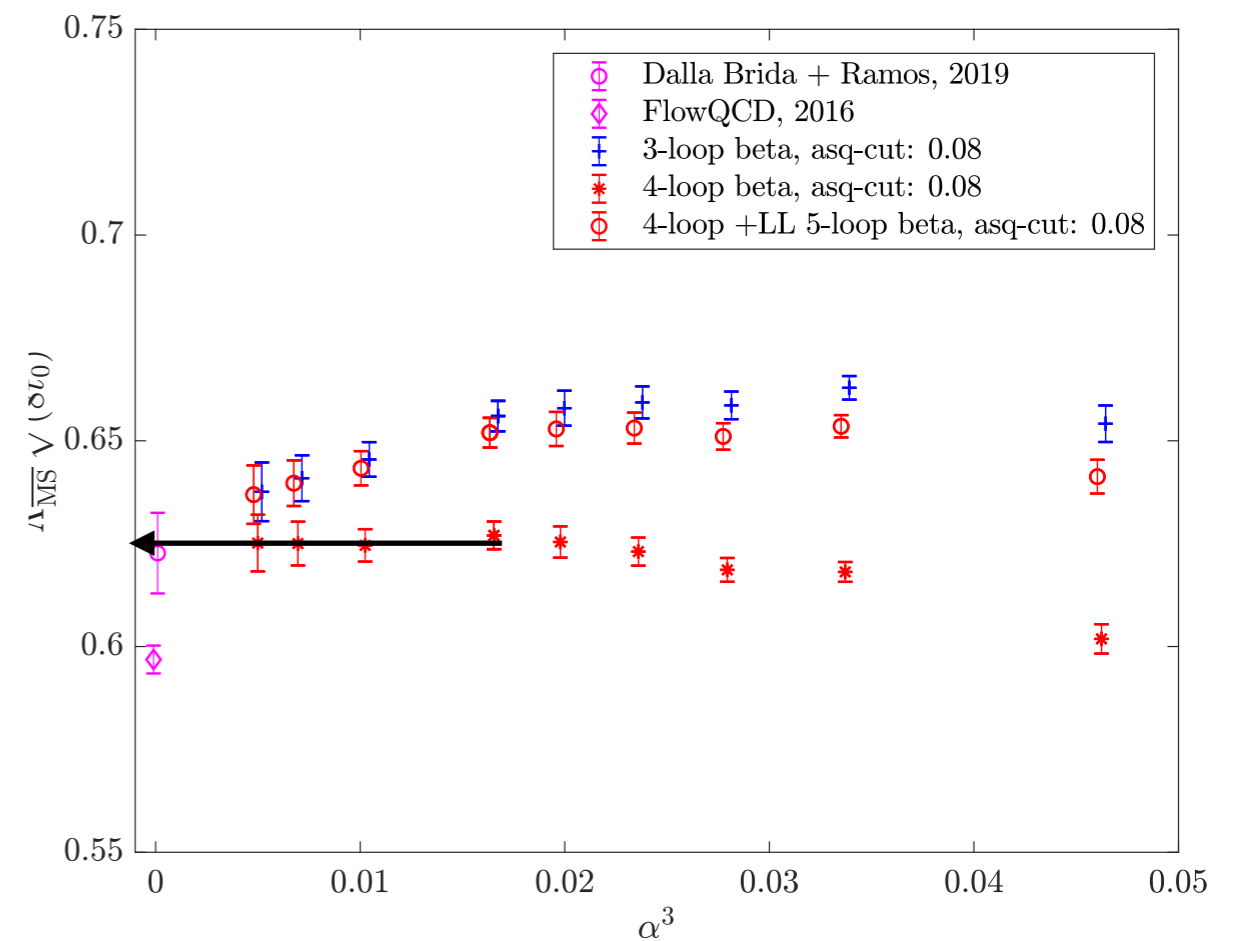
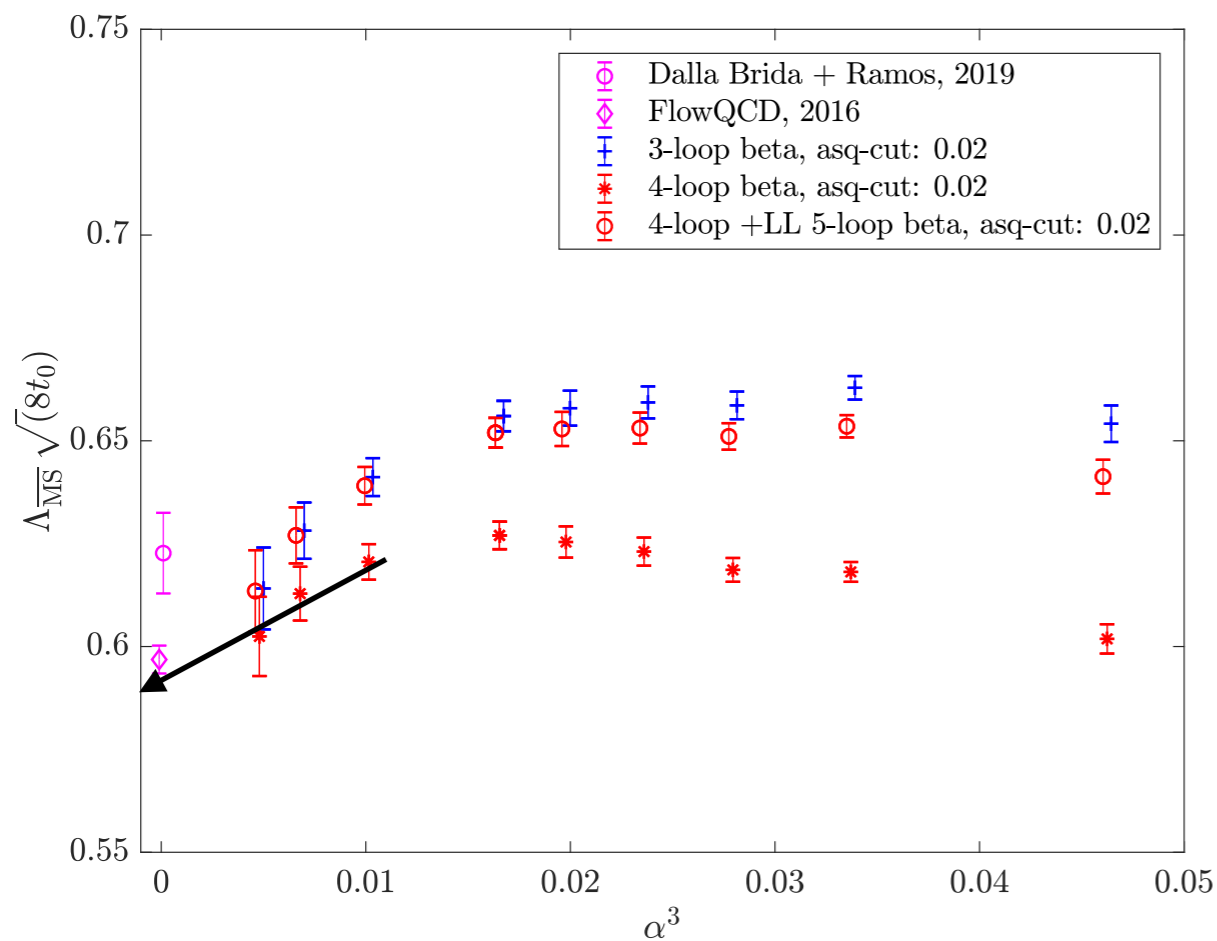
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7% difference from combination of $\alpha \rightarrow 0$ and $a \rightarrow 0$?

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Thank you

