

# Charm CP & the lattice

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HET-BNL

*Sorry, NO lattice data yet;  
just some initial  
ideas at this  
stage*

Lattice 2019

Wuhan, China 06/21/19

# Outline

- Recently, LHCb exciting obs of  $\Delta ACP$  *EW Moniod March 2019*
- Naturalness reasoning strongly suggests other BSM-CP phase(s) must exist. *To discern Computations essential*
- CPV represents an imp. avenue for searching new pheno but precise calculations are highly desirable
- Crude estimate of  $\Delta ACP$  *→ 1905.00907*
- Propose a specific mechanism for understanding  $\Delta ACP$
- This mechanism points to new lattice difficulties for addressing  $\Delta ACP$
- Suggest mode for precision testing the SM-CKM
- Lattice feasibility
- Summary + Outlook

$$\Delta A_{CP} \equiv A_{CP}(K^+ K^-) - A_{CP}(\pi^+ \pi^-)$$

## Results



LHCb-PAPER-2019-006

$$\Delta A_{CP}^{\pi\text{-tagged}} = [-18.2 \pm 3.2 \text{ (stat.)} \pm 0.9 \text{ (syst.)}] \times 10^{-4}$$

$$\Delta A_{CP}^{\mu\text{-tagged}} = [-9 \pm 8 \text{ (stat.)} \pm 5 \text{ (syst.)}] \times 10^{-4}$$

- Compatible with **previous** LHCb results and the **WA**
- **Combination** with LHCb Run 1 gives:

*# is ~ 6 times smaller than 2012 LHCb*

$$\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}$$

CP violation observed at **5.3σ**!!

# CP violation history

**BNL-COLUMBIA**

1956  
**Parity violation**  
T. D. Lee,  
C. N. Yang and  
C. S. Wu *et al.*

1964  
**Strange particles:  
CP violation in K  
meson decays**  
J. W. Cronin,  
V. L. Fitch *et al.*

**BNL**  
**Indirect CP**  
 $\sim 10^{-3}$

2001  
**Beauty particle  
CP violation in  
meson decays**  
BaBar and Belle  
collaborations

**Direct  $\sim 10^{-6}$**   
**CERN/FNAL**

1963  
**Cabibbo Mixing**  
N. Cabibbo

1973  
**The CKM matrix**  
M. Kobayashi and  
T. Maskawa

2019  
**Charm particle  
CP violation  
meson decay**  
LHCb collabo

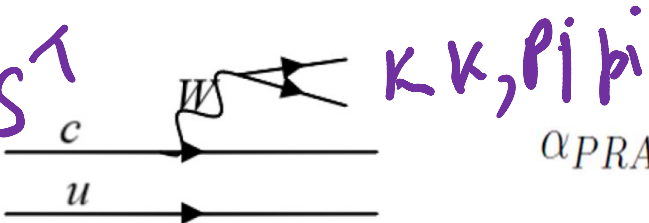


# Alas, we cannot compute!

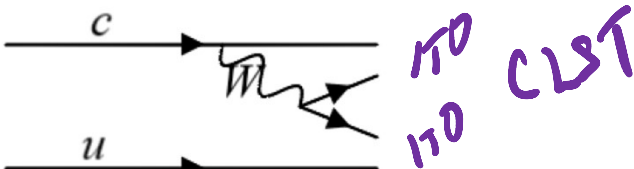
- Sadly calculations of Direct CP remain a challenge for theory; we are not able to calculate with any precision [**eps' an exception as a result of over 3 decades of concerted lattice effort**]. As a rule, we can just give crude (QCD) model-dependent estimates..
- Essential ingredient is CP-conserving scattering phases & usually complicated non-perturbative effect on weak decays....difficult for the lattice in general

# BASICS of DIRECT CP ASy

CBS T

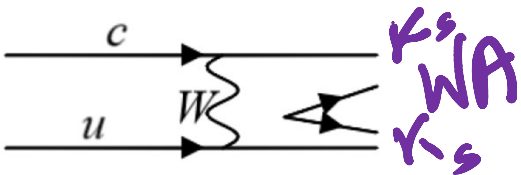


$$\alpha_{PRA} = \frac{B[\bar{I} \rightarrow \bar{F}] - B[\bar{I} \rightarrow \bar{F}]}{B[I \rightarrow F] + B[\bar{I} \rightarrow \bar{F}]}$$



$$A_{hh} = \lambda T_{hh} + \lambda^5 P_{hh} e^{i[\delta_{st} + \eta_{wk}]} \quad \text{due CPT}$$

→ PRA



$$\alpha_{hh} = 2\lambda^4 P_{hh} \sin(\delta_{st}) \sin(\gamma) / T_{hh}$$



interferences with CBS T

↓  
↓ Br

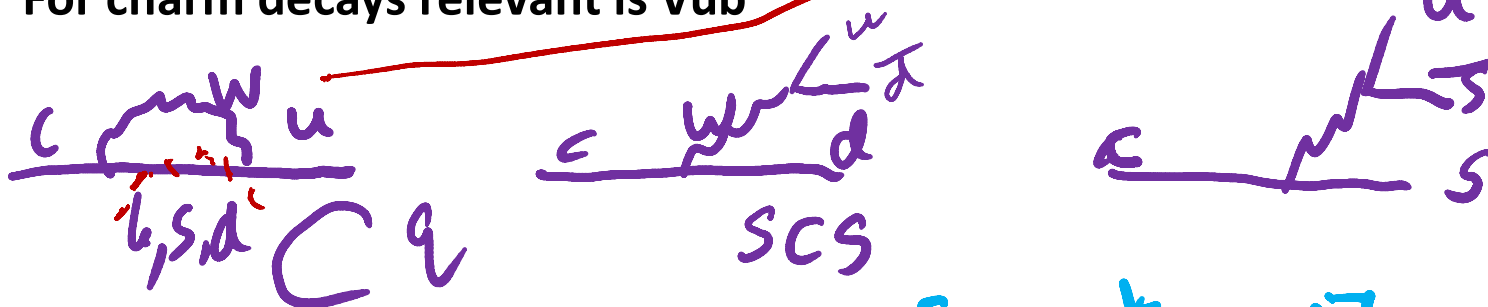
$V_{cb} V_{ub}^*$  CP-odd phase

# CRUDE ESTIMATE

ASCKM 2018

# SM expectation...DCP

- Dir CP..... See Bander, Silverman + AS, PRL 1979 for DCP when  $m_q \gg \lambda_{\text{QCD}}$ ...anticipate large corrections for charm from s-quark [K-decays] → non-perturbative
- Key points: Penguin-Tree interference; SCS modes.....Hall mark of BSS'79
- Need suitable simple changes
- SM-dominant CKM phase either in  $V_{ub}$  or in  $V_{td}$  →  $m_c^2 > 4[m_s^2, m_d^2]$
- For charm decays relevant is  $V_{ub}$



$4CLA \rightarrow \text{DCP} \sim \frac{\text{Im}[V_{cb}V_{ub}^*]}{\Lambda^2} \sim \lambda^4 \sim 10^{-3}$   
 $\rightarrow K^+K^-, \pi^+\pi^-$

maybe enhanced a lot by CLS  
 e.g.  $K_S K_S$

$\sim N \times 10^{-3} \lesssim 10^{-2} !!$   
 ↑ Very DIFFICULT to reach in SM

# Lattice difficulties for $\Delta ACP$

- $m_D \sim 1.865$  GeV;  $Br [D \Rightarrow K^+ K^-] \sim 0.4\%$

**much larger Br to multiparticle states...difficult on the lattice....see Hansen and Sharpe, PRD'12.**

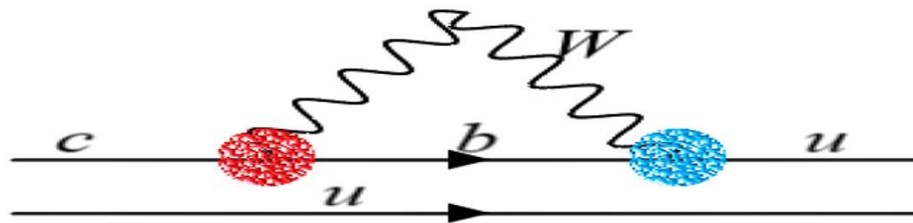
- **However, strictly speaking  $\Delta ACP$  requires only difference between KK and ppi states.**
- **Therefore, optimistically, one may hope for appreciable cancellation between multiparticle states; i.e U-spin or SU(3) may provide a basis for such a cancellation.**
- **But unlikely to be viable as will be explained.**

# Interplay with resonances

- Charm region has many resonances
- Resonances that have appreciable Br to  $KK$  &  $\pi\pi$  are relevant
- In particular ( $\rho$ s) scalar resonances
- These can enter the weak decay dynamics via the penguin  $H_{\text{eff}}$

# b-penguin

- b-penguin fig





# Penguin Hamiltonian

- Since  $mb \gg \Lambda_{\text{QCD}}$ , we can integrate it out along with  $W$  and get the familiar,  $H_{\text{eff}}$ :
- $H_{\text{eff}} = \sum C_i O_i$

$$O_3 = (\bar{c}_\alpha u_\alpha)_{V-A} [(\bar{u}_\beta u_\beta)_{V-A} + (\bar{d}_\beta d_\beta)_{V-A} + (\bar{s}_\beta s_\beta)_{V-A} + (\bar{c}_\beta c_\beta)_{V-A}] ,$$

$$O_4 = (\bar{c}_\alpha u_\beta)_{V-A} [(\bar{u}_\beta u_\alpha)_{V-A} + (\bar{d}_\beta d_\alpha)_{V-A} + (\bar{s}_\beta s_\alpha)_{V-A} + (\bar{c}_\beta c_\alpha)_{V-A}] ,$$

$$O_5 = (\bar{c}_\alpha u_\alpha)_{V-A} [(\bar{u}_\beta u_\beta)_{V+A} + (\bar{d}_\beta d_\beta)_{V+A} + (\bar{s}_\beta s_\beta)_{V+A} + (\bar{c}_\beta c_\beta)_{V+A}] ,$$

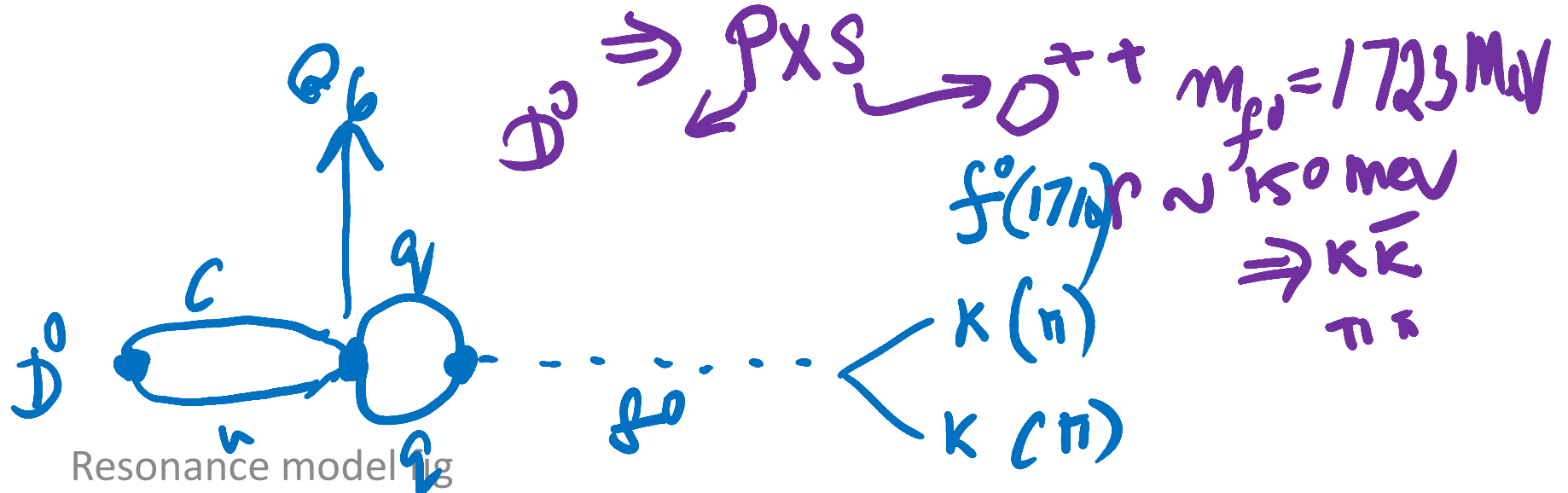
$$O_6 = (\bar{c}_\alpha u_\beta)_{V-A} [(\bar{u}_\beta u_\alpha)_{V+A} + (\bar{d}_\beta d_\alpha)_{V+A} + (\bar{s}_\beta s_\alpha)_{V+A} + (\bar{c}_\beta c_\alpha)_{V+A}] ,$$

See Abbott, Sikivie, Wise '80; Goldman Grinstein '89  
 Buchalla, Buras, Lautenbacher '96

A.S. 1905.00907

# Penguin Ops & Resonance Coupling

$D_s + D_b$  are LXR  $\Rightarrow -2[S+P] \times [S-P]$



See PDG Live

$$BR \quad f_0 \rightarrow K\bar{K} \quad \sim 40\%$$

$$\frac{(f_0 \rightarrow \pi\pi)}{(f_0 \rightarrow K\bar{K})} \sim 0.4$$

# S,P Glueballs & their SU3 breakings

Monningstern  
+  
Peardon,  
Review '99

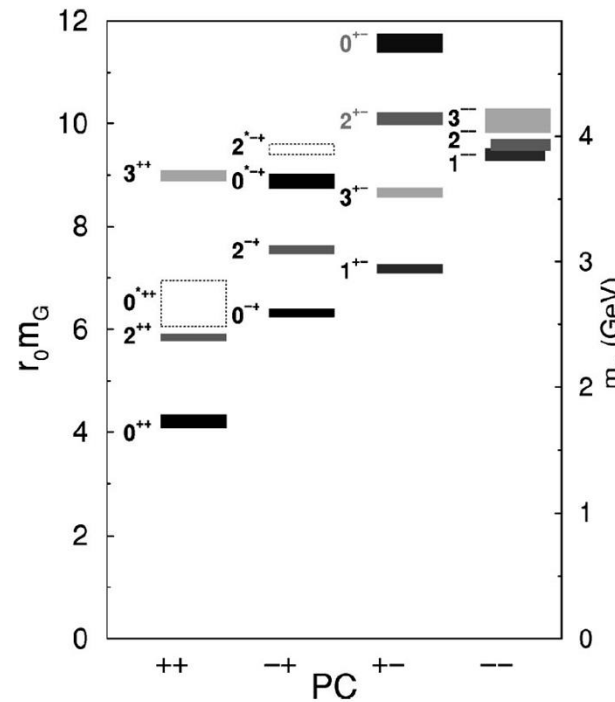


FIG. 8. The mass spectrum of glueballs in the pure SU(3) gauge theory. The masses are given in terms of the hadronic scale  $r_0$  along the left vertical axis and in terms of GeV along the right vertical axis (assuming  $r_0^{-1} = 410$  MeV). The mass uncertainties indicated by the vertical extents of the boxes do *not* include the uncertainty in setting  $r_0$ . The locations of states whose interpretation requires further study are indicated by the dashed open boxes.

# PDG gives Br's of D0

$K^+ K^-$

$$(3.97 \pm 0.07) \times 10^{-3}$$

$\pi^+ \pi^-$

$$(1.407 \pm 0.025) \times 10^{-3}$$

Ratio of  $\frac{\text{Br}(D^0 \rightarrow K^+ K^-)}{\text{Br}(D^0 \rightarrow \pi^+ \pi^-)} \approx 2.82$

Large SU3  
breaking known  
for long

Mystery for long

$(f_K/f_\pi)^2 \approx 1.7$  explains a part

# New understanding of an old issue

Cornwall + AS, PRD '84

$$\mathcal{L}_{\text{eff}} = -\bar{\psi}(G_S S + i\gamma_5 G_P P)\psi + \dots, \quad (3.2)$$

$$G_S = \frac{M_Q}{2\langle S \rangle}$$

$$\langle S \rangle \simeq 130 \text{ MeV}.$$

Couplings to quarks are proportional to quark mass  
(1.3)

S, P are low lying gluonia  
 $\delta^{++}$  &  $\delta^{-+}$

# **f0 decay to KK & pi pi**

- **PDG live gives**
- **$\text{Br}[f_0 \Rightarrow \pi^+ + \pi^-] / \text{Br}[f_0 \Rightarrow K^+ + K^-] \sim 0.4$**

**enhancing the SU3 breaking decays of the D0**

# Imp implication of Resonances for CP

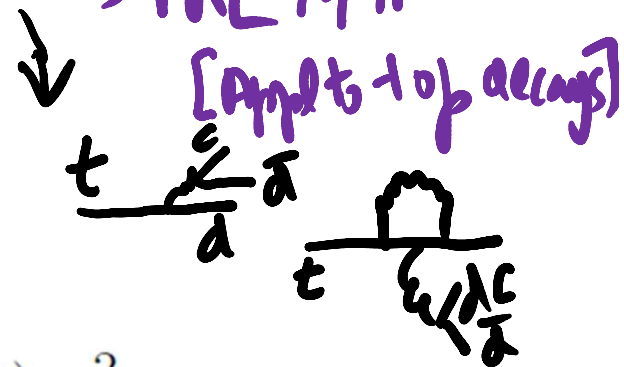
**WHEN DEALING WITH FINAL STATES THAT ARE DOMINATED BY RESONANCES, THEN THE WIDTH OF THE RESONANCE PROVIDES THE “STRONG”-CP-EVEN PHASE NEEDED FOR COMPUTING PRA...1<sup>ST</sup> ILLUSTRATED BY EILAM, HEWETT + AS, PRL'91**



Resonance idea

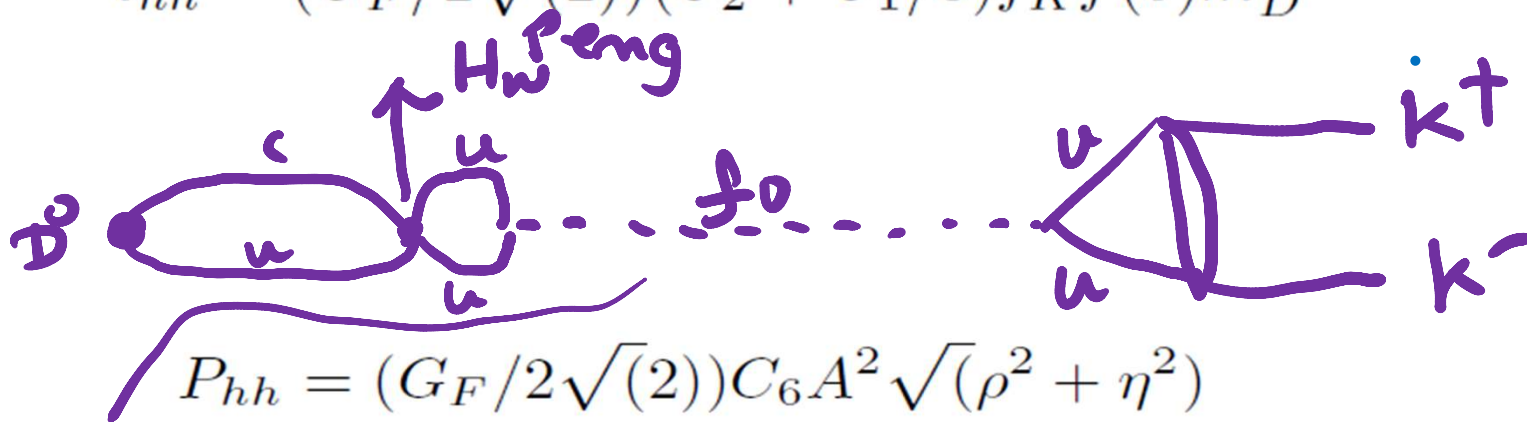
For basic idea see Eidelam, Hewett + AS PRL 1991

See also ATW + AS ZP 1994



For  $D^0 \rightarrow KK$

$$t_{hh} = (G_F/2\sqrt{(2)})(C_2 + C_1/3)f_K f(0)m_D^2$$



$$P_{hh} = (G_F/2\sqrt{(2)})C_6 A^2 \sqrt{(\rho^2 + \eta^2)}$$

$$[\frac{m_D^2}{f_D} - \frac{m_f^2}{f_f} + i\eta_f \frac{r_f}{f_f}]$$

$$(f_D m_D / m_c) K_f m_f^2 \quad (7)$$

$$K_f \sim 3.7$$

Nice paper  
Calo H-N Li, C-D Lu + Fu-S Yu PRD 86, 2012

$$\alpha_{K^+K^-} \approx 5.5 \times 10^{-4}$$

$$\alpha_{\pi^+\pi^-} = -1.06 * \alpha_{K^+K^-}; \Delta A \approx 12 \times 10^{-4}$$

## Interplay of CPT & resonant CP

Because of the CPT constraint that life-time of particle must equal to that of its antiparticle, we must have,

$$\sum_X \Delta\Gamma(X) = 0 \quad (9)$$

where  $X$  are the various final states that emerge from the decays of  $D^0$  and

$$\Delta\Gamma(X) = \Gamma(D^0 \rightarrow X) - \Gamma(\bar{D}^0 \rightarrow \bar{X})$$

Since at the quark level in charm decays there are only two channels,  $c \rightarrow u\bar{d}d$  and  $c \rightarrow u\bar{s}s$ , this means because of CPT we must have [20, 21],

$$\Delta\Gamma(c \rightarrow u\bar{d}d) = -\Delta\Gamma(c \rightarrow u\bar{s}s) \quad (10)$$

At the meson level, it is suggested [20] that this materializes into,

$$\Delta\Gamma(\pi^+\pi^-) = -\Delta\Gamma(K^+K^-) \quad (11)$$

Using the tree dominance, as emphasized above, one then arrives at the relation,

$$\alpha_{K^+K^-} \propto -\alpha_{\pi^+\pi^-} / \sqrt{2.8} \quad (12)$$

where, we have taken for simplicity that the two Brs differ by about  $\sqrt{2.8}$  [18].

However, f-dominance of the penguin amplitude, in the  $h h$  channel that is our central focus, appearing in the numerator of (4) modifies this expectation appreciably. This is because, it appears that [5].

$$Br(f_0 \rightarrow \pi\pi) / Br(f_0 \rightarrow KK) \approx 0.4 \quad (13)$$

• See ATWOOD+AS  
PTEP 2013

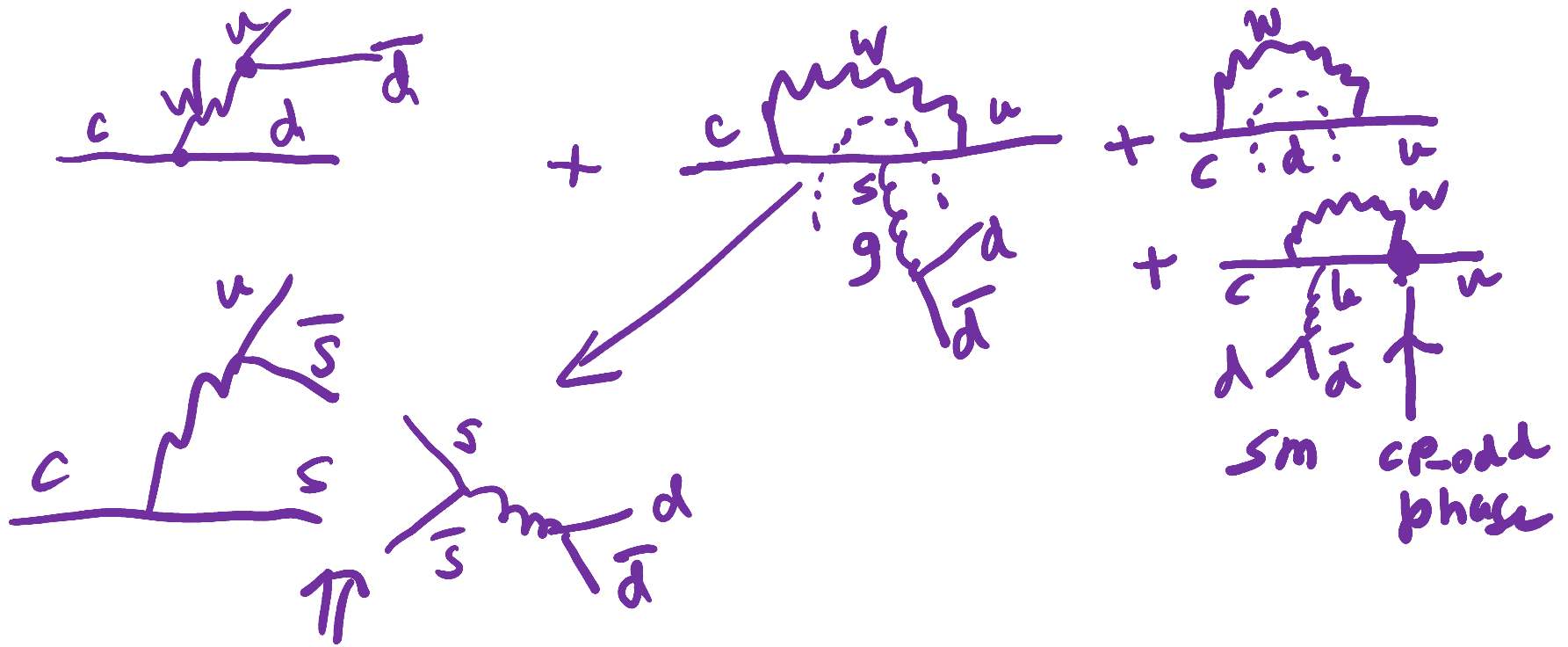
$c \rightarrow u\bar{d}d, u\bar{s}s$



In charm decays  
 $T \gg P \Rightarrow$

$\frac{d}{PR} \propto \frac{1}{f_B}$

PDGLIVE



on-shell rescattering phase  
 CP-even phase  $\Rightarrow$  Total amplitude  
 for  $c \rightarrow u d \bar{d}$  is  
 complex

central value

$$-\alpha_{\pi^+\pi^-} / \alpha_{K^+K^-} = \sqrt{(0.4 \times 2.8)} \approx 1.06$$

I  
Large uncertainty  
(PDG)

Mode	BR ( $10^{-3}$ )	Current PRA bound ( $10^{-2}$ )
$K^+K^-$	$3.97 \pm .07$	$-0.07 \pm 0.11$
$K_s K_s$	$0.17 \pm 0.012$	$-0.4 \pm 1.5$
$\pi^+\pi^-$	$1.407 \pm 0.025$	$0.13 \pm 0.14$
$\pi^0\pi^0$	$0.822 \pm 0.025$	$0.0 \pm 0.6$

→ out dated

→ out dated

$$\alpha_{K_s K_s} / \alpha_{K^+K^-} \approx \sqrt{(23/4)} \approx 2.4$$

$$\alpha_{\pi^0\pi^0} = 1.3 \alpha_{\pi^+\pi^-}$$

II  
III

**CANDIDATES FOR PRECISION TEST:  
NEEDS TO BE EXPERIMENTALLY  
ACCESSIBLE AND BE ALSO AMENABLE  
TO PRECISE LATTICE CALCULATIONS**

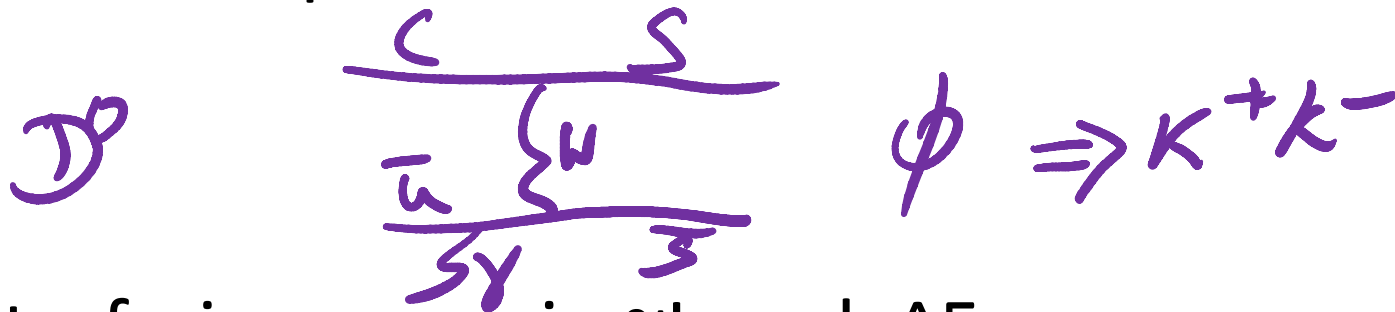
# Candidate for precision test: CPV in

**$D^0 \Rightarrow \gamma + [\phi \Rightarrow K^+ K^-]$**

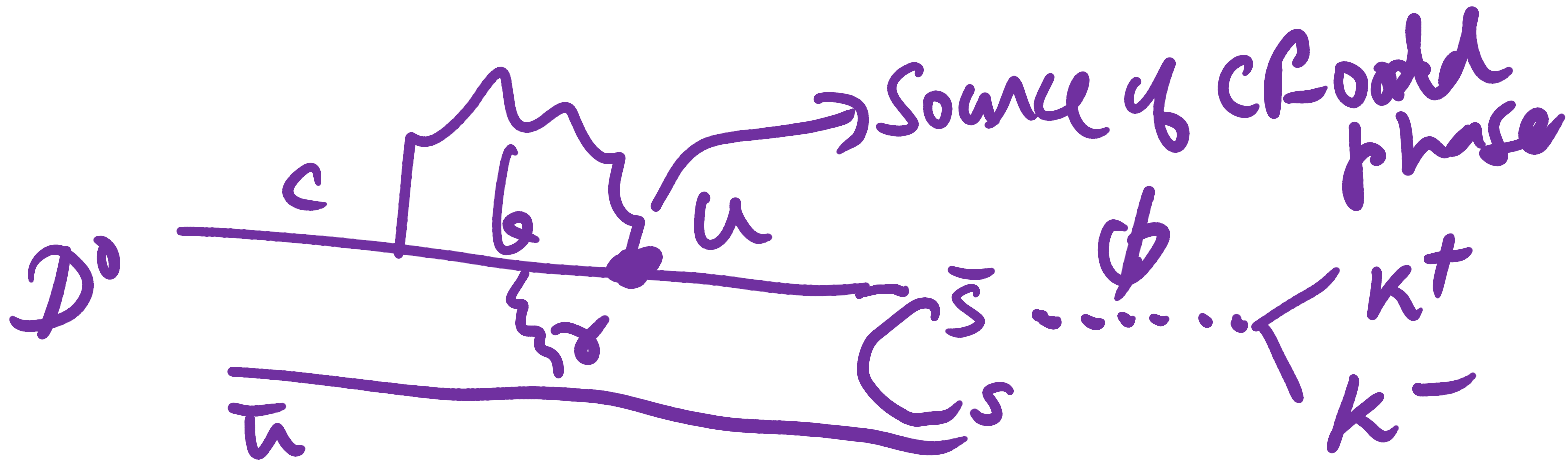
- Requires KK scattering phases for CM energy below around 1 GeV
- Differential photon energy spectrum can be used for precision test
- Integrated BR  $\sim 2.74 \text{ pm } 0.19 \times 10^{-5}$  well measured
- Belle-II expected  $\sim 10^{10}$  D's/yr ...precise measurement of differential spectrum is clearly feasible

# Production mechanisms

- Dominant production: Tree  $\sim \Lambda$



- Interfering penguin  $\sim \Lambda^5$





# Feasibility of lattice extraction of KK phases in $\sim 1$ GeV energy region

- Good reason to expect that this is doable: see works of JLAB [Briceno et al PRL 2017], RBC-UKQCD [Tainle Wang, Chris Kelly, Aaron Meyer, Mattia Brown, Dan Hoying, David Murphy et al, Lat2018 & 2019 presentations] paper(s) in preparation; see also G. Rendon et al Lat2018
- At  $\sim 1$  GeV or below, relevant to radiative D0 decays, error due to KKpipi, 4 pi contamination likely small *if needs be can focus on higher  $E_\gamma$  so as to reduce contamination from 4 p's*

## BRIEF REMARKS ON RADIATIVE CHARM DECAYS

- New class of radiative decays...interesting nearby resonances relevant for radiative transitions: [see PDGLive]:

$$\phi(1680), \Gamma \sim 150 \text{ MeV} \Rightarrow K \bar{K}^*(892)$$

$$\rho(1702), \Gamma \sim 250 \text{ MeV} \rightarrow \rho \pi \pi, \gamma \pi \dots$$

$$D^0 \rightarrow \gamma \rho(1702), \gamma \phi(1650) \quad E_\gamma \lesssim 300 \text{ MeV} \Leftarrow$$


- Best candidates for precision tests of the SM

$$D^0 \rightarrow \gamma \rho(770), \gamma \phi(1020) \quad BR \sim 10^{-4} - 10^{-5}$$

$$\sigma_{\mu(1-\gamma_5)} \times \gamma_{\mu(1+\gamma_5)} \Rightarrow \rho \times \rho$$

$\text{FT}$

# Another important resonance nearby


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**2018 Review of Particle Physics.**  
M. Tanabashi *et al.* (Particle Data Group), Phys. Rev. D **98**, 030001 (2018)

**LIGHT UNFLAVORED MESONS**  
( $S = C = B = 0$ )  
For  $I = 1$  ( $\pi, \rho, a$ ):  $u\bar{d}, (u\bar{u} - d\bar{d})/\sqrt{2}, d\bar{u}$ ;  
for  $I = 0$  ( $\eta, \eta', h, h', \omega, \phi, f, f'$ ):  $c_1(u\bar{u} + d\bar{d}) + c_2(s\bar{s})$

$\eta(1760) \quad I^G(J^{PC}) = 0^+(0^{-+})$

Seen by DM2 in the  $\rho\rho$  system (BISELLO 1989B). Structure in this region has been reported before in the same system (BALTRUSAITIS 1986B) and in the  $\omega\omega$  system (BALTRUSAITIS 1985C, BISELLO 1987).

---

$\eta(1760)$ MASS	1751 ± 15 MeV
$\eta(1760)$ WIDTH	240 ± 30 MeV

**Decay Modes**

Mode	Fraction ( $\Gamma_i / \Gamma$ )	Scale Factor/ Conf. Level	P (MeV/c)
$\Gamma_1$ $4\pi$			823
$\Gamma_2$ $2\pi^+2\pi^-$	seen		819
$\Gamma_3$ $\pi^+\pi^-2\pi^0$	seen		822
$\Gamma_4$ $\rho^0\rho^0$	seen		407
$\Gamma_5$ $\rho^+\rho^-$	seen		407
$\Gamma_6$ $2(\pi^+\pi^-\pi^0)$			733
$\Gamma_7$ $\omega\omega$	seen		393
$\Gamma_8$ $\eta'\pi^+\pi^-$	seen		572
$\Gamma_9$ $\gamma\gamma$	seen		876

$\eta(1760)$   
 $0^-$   
 Eng Asy  
 TeA  
 $4\pi, \eta'\pi\pi$   
 $50, 90 \dots$

$R_D(*)$  ANOMALY AND NEW PHYSICS IN  
 B-PENGUIN CONTRIBUTION TO CHARM  
 DECAYS [18]

$b \rightarrow c \Rightarrow R_D(*)$

Another reason  
 for precision  
 test

New Phys

SM-CKM  
 phase

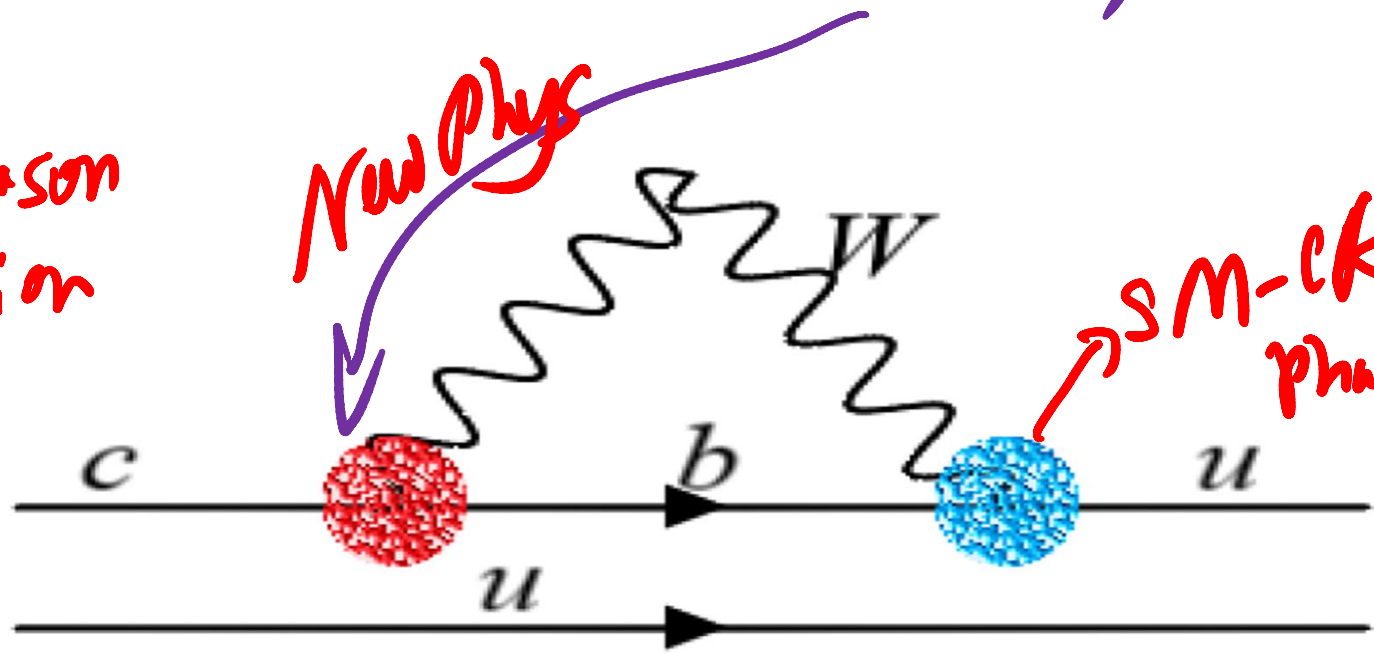


FIG. 2. b-penguin in charm decay; left c-b (red) vertex may have a CP-odd phase endowed by new physics affecting  $R_D(*)$  whereas the right b-u (blue) vertex contains the SM-CKM phase.

# Summary + outlook

- Charm region is rich with resonances
- Some of these are expected to play an important role in driving CP asymmetries, PRA, energy (3 or more particle FS), TCA (4 or more spin-less)
- New class of radiative decays with photon energy  $< \sim 300$  MeV to resonances which cascade down to several diff FS, such as KK, pi pi, eta'(eta, pi0) pi pi....
- Mixing of D0 weak decays with f0 accounts roughly for the observed size of  $\Delta ACP$  & large SU3 breaking seen for long. Resonances complicate addressing of  $\Delta ACP$  by lattice methods.
- Important target should be precise CP-tests for SM-CKM. For this purpose, it is suggested that the best candidate for lattice is CPV asymmetries (O .1%), Br  $\sim 10^{-5}$ :  $D0 \Rightarrow \gamma \phi$  with  $\phi \Rightarrow K^+ K^-$

*PRA as well as energy ( $\gamma$ ) asym are expected*

# XTRAS

## Repercussions of flavour symmetry breaking on CP violation in $D$ -meson decays

Thorsten Feldmann,<sup>a</sup> Soumitra Nandi<sup>a</sup> and Amarjit Soni<sup>b</sup>

<sup>a</sup>*Theoretische Elementarteilchenphysik, Naturwissenschaftlich Technische Fakultät, Universität Siegen, 57068 Siegen, Germany*

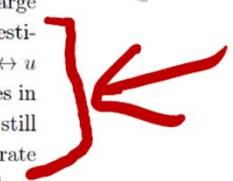
<sup>b</sup>*Physics Department, Theory Group, Brookhaven National Laboratory, Upton, NY 11973, U.S.A.*

*E-mail:* [thorsten.feldmann@uni-siegen.de](mailto:thorsten.feldmann@uni-siegen.de), [nandi@physik.uni-siegen.de](mailto:nandi@physik.uni-siegen.de), [soni@bnl.gov](mailto:soni@bnl.gov)

ABSTRACT: We investigate to what extent the recently measured value for a non-vanishing direct CP asymmetry in  $D^0 \rightarrow K^+K^-$  and  $D^0 \rightarrow \pi^+\pi^-$  decays can be accommodated in the Standard Model (SM) or extensions with a constrained flavour sector, for instance from a sequential 4<sup>th</sup> generation of quarks (4G). From the comparison with  $D^0 \rightarrow K^-\pi^+$  branching ratios, we establish large U-spin symmetry ( $d \leftrightarrow s$ ) breaking effects with large strong phases between different interfering amplitudes. On the basis of conservative estimates on amplitude ratios — which are supported by an analysis of the breaking of a  $c \leftrightarrow u$  symmetry in non-leptonic  $B^0$  decays — we find that, in the SM, direct CP asymmetries in the  $\pi^+\pi^-$  or  $K^+K^-$  modes (or in their difference) of the order of several per mille are still plausible. Due to the constraints on the new CP phases in the 4G model, only moderate effects compared to the SM estimates are possible. We suggest CP studies at LHCb as well as at (Super)B-factories of several distinctive modes, such as  $D^+ \rightarrow \bar{K}^{(*)0}\pi^+$ ,  $\phi\pi^+$  and  $D_s \rightarrow K^{(*)0}\pi^+$ ,  $\phi\pi^+(K^+)$  etc., which should shed more light on the short- and long-distance

arXiv:1111.3523 [hep-th] [inSPIRE] Lattice-2019; soni@BNL-HET

JHEP06(2012)007



*Do I: d, s + accurate calculation remain a challenge for theorists!*



**$f_0(1710)$  [ $t$ ]**

$$I^G(J^{PC}) = 0^+(0^{++})$$

Mass  $m = 1723^{+6}_{-5}$  MeV ( $S = 1.6$ )Full width  $\Gamma = 139 \pm 8$  MeV ( $S = 1.1$ )

$f_0(1710)$ DECAY MODES	Fraction ( $\Gamma_i/\Gamma$ )	$p$ (MeV/c)
$K\bar{K}$	seen	706
$\eta\eta$	seen	665
$\pi\pi$	seen	851
$\omega\omega$	seen	360



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$$\eta(1760) \quad I^G(J^{PC}) = 0^+(0^{-+})$$

Seen by DM2 in the  $\rho\rho$  system (BISELLO 1989B). Structure in this region has been reported before in the same system (BALTRUSAITIS 1986B) and in the  $\omega\omega$  system (BALTRUSAITIS 1985C, BISELLO 1987).

INSPIRE search

 $\eta(1760)$  MASS  
 $\eta(1760)$  WIDTH $1751 \pm 15$  MeV  
 $240 \pm 30$  MeV