

# $B \rightarrow \pi l \nu$ form factors and $|V_{ub}|$ with Möbius domain wall fermions

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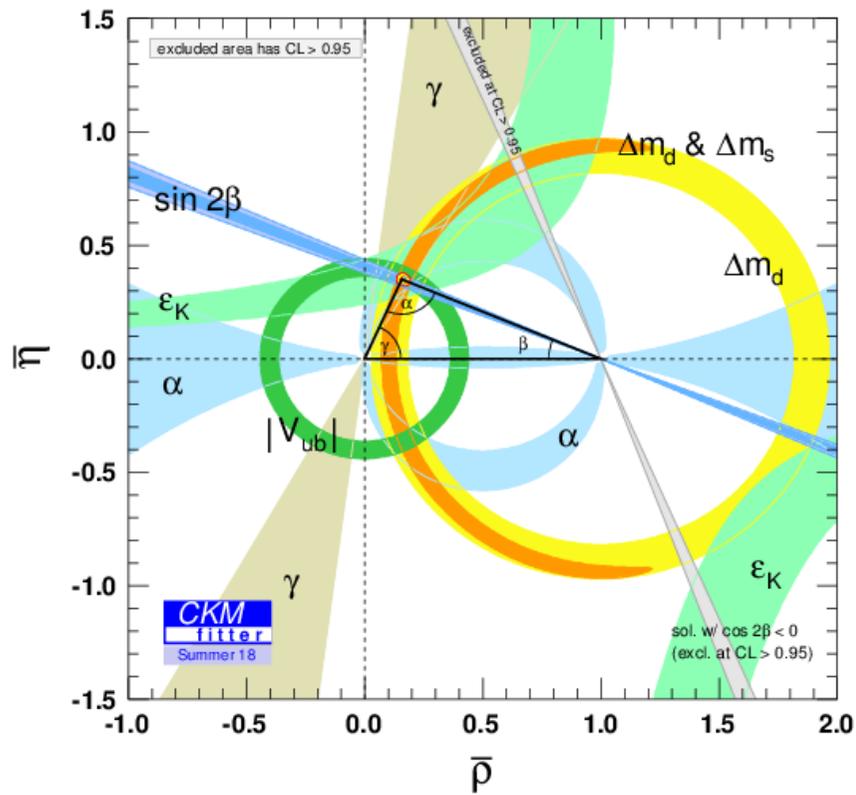


Lattice 2019, Wuhan, China

# Outline

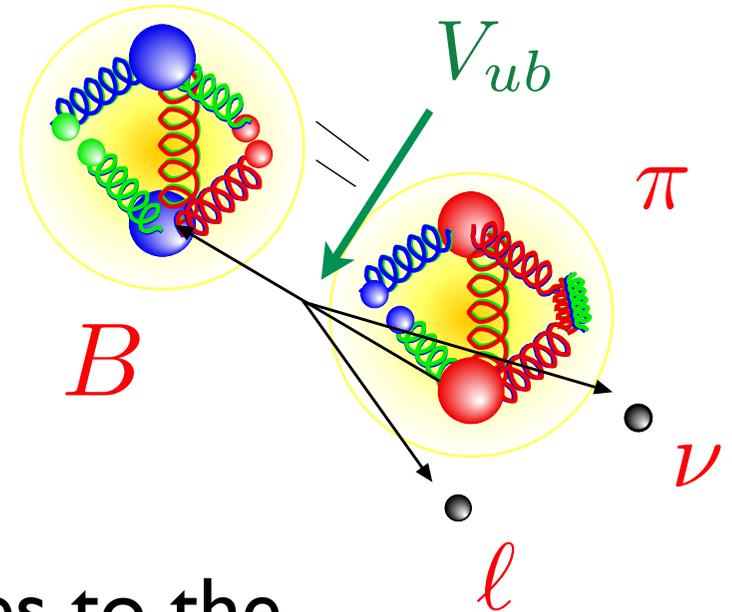
- $B$  meson semileptonic decays and CKM matrix elements
- JLQCD lattice ensembles
- Preliminary results
- Summary

# Motivation



$$\left( \begin{array}{ccc}
 V_{ud} & V_{us} & V_{ub} \\
 \pi \rightarrow \ell\nu & K \rightarrow \ell\nu & B \rightarrow \pi\ell\nu \\
 & K \rightarrow \pi\ell\nu & \\
 V_{cd} & V_{cs} & V_{cb} \\
 D \rightarrow \ell\nu & D_s \rightarrow \ell\nu & B \rightarrow D\ell\nu \\
 D \rightarrow \pi\ell\nu & D \rightarrow K\ell\nu & B \rightarrow D^*\ell\nu \\
 V_{td} & V_{ts} & V_{tb} \\
 \langle B_d | \bar{B}_s \rangle & \langle B_s | \bar{B}_s \rangle & 
 \end{array} \right)$$

# CKM matrix element $V_{ub}$



- CKM matrix element relates to the differential decay rate

$$\frac{d\Gamma(B \rightarrow \pi l \nu)}{dq^2} = \frac{G_F^2 |V_{ub}|^2}{24\pi^3} |k_\pi|^3 |f_+(q^2)|^2$$

decay rate  
(experiment)

CKM matrix  
element

theory (lattice)

$k_\pi$  = pion momentum  
 $q$  = momentum transfer

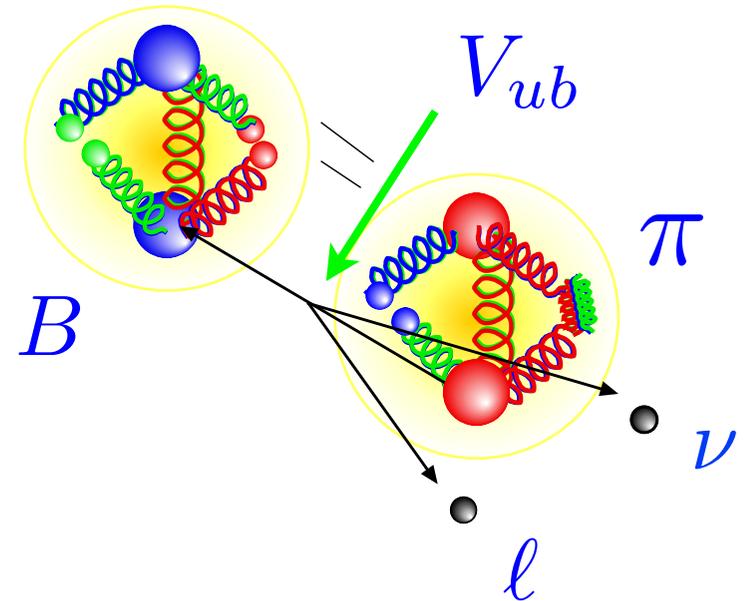
# Form factors

- For pseudoscalar to pseudoscalar decays

$$\langle \pi(k_\pi) | V^\mu | B(p_B) \rangle = f_+(q^2) \left[ (p_B + k_\pi)^\mu - \frac{M_B^2 - M_\pi^2}{q^2} q^\mu \right] + f_0(q^2) \frac{M_B^2 - M_\pi^2}{q^2} q^\mu$$

- $p_B$  and  $k_\pi$  are the  $B$  and  $\pi$  four momenta
- $q^\mu = p_B^\mu - k_\pi^\mu$  is the momentum transfer
- kinematic constraint  $f_0(0) = f_+(0)$

# Form factors



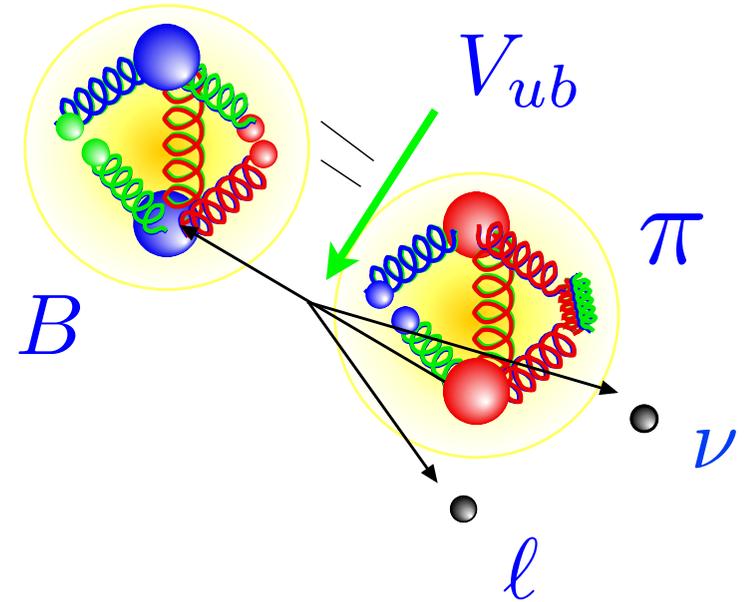
- In the context of Heavy Quark Effective Theory (HQET), a useful parametrisation is

$$\langle \pi(k_\pi) | V^\mu | B(p_B) \rangle = 2\sqrt{M_B} \left[ f_1(v \cdot k_\pi) v^\mu + f_2(v \cdot k_\pi) \frac{k_\pi^\mu}{v \cdot k_\pi} \right]$$

-  $v^\mu = p_B^\mu / M_B$  is the heavy quark velocity

- pion energy  $E_\pi = v \cdot k_\pi = \frac{M_B^2 + M_\pi^2 - q^2}{2M_B}$

# Form factors

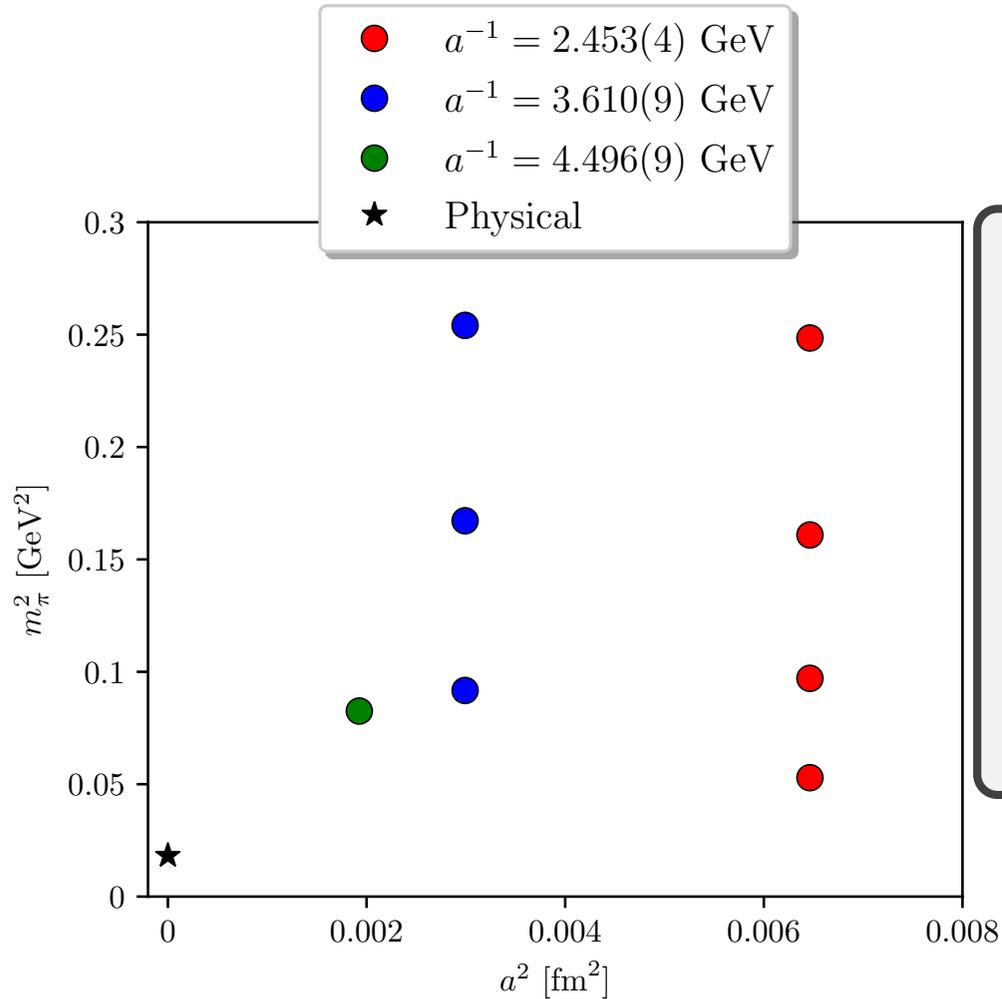


- These two form factor definitions are related by

$$f_+(q^2) = \sqrt{M_B} \left[ \frac{f_2(v \cdot k_\pi)}{v \cdot k_\pi} + \frac{f_1(v \cdot k_\pi)}{M_B} \right]$$

$$f_0(q^2) = \frac{2}{\sqrt{M_B}} \frac{M_B^2}{M_B^2 - M_\pi^2} \left[ f_1(v \cdot k_\pi) + f_2(v \cdot k_\pi) - \frac{v \cdot k_\pi}{M_B} \left( f_1(v \cdot k_\pi) + \frac{M_\pi^2}{(v \cdot k_\pi)^2} f_2(v \cdot k_\pi) \right) \right]$$

# Lattice configurations



Möbius Domain Wall Fermions

$$N_f = 2 + 1$$

$$m_\pi \approx 225, 300, 400, 500 \text{ MeV}$$

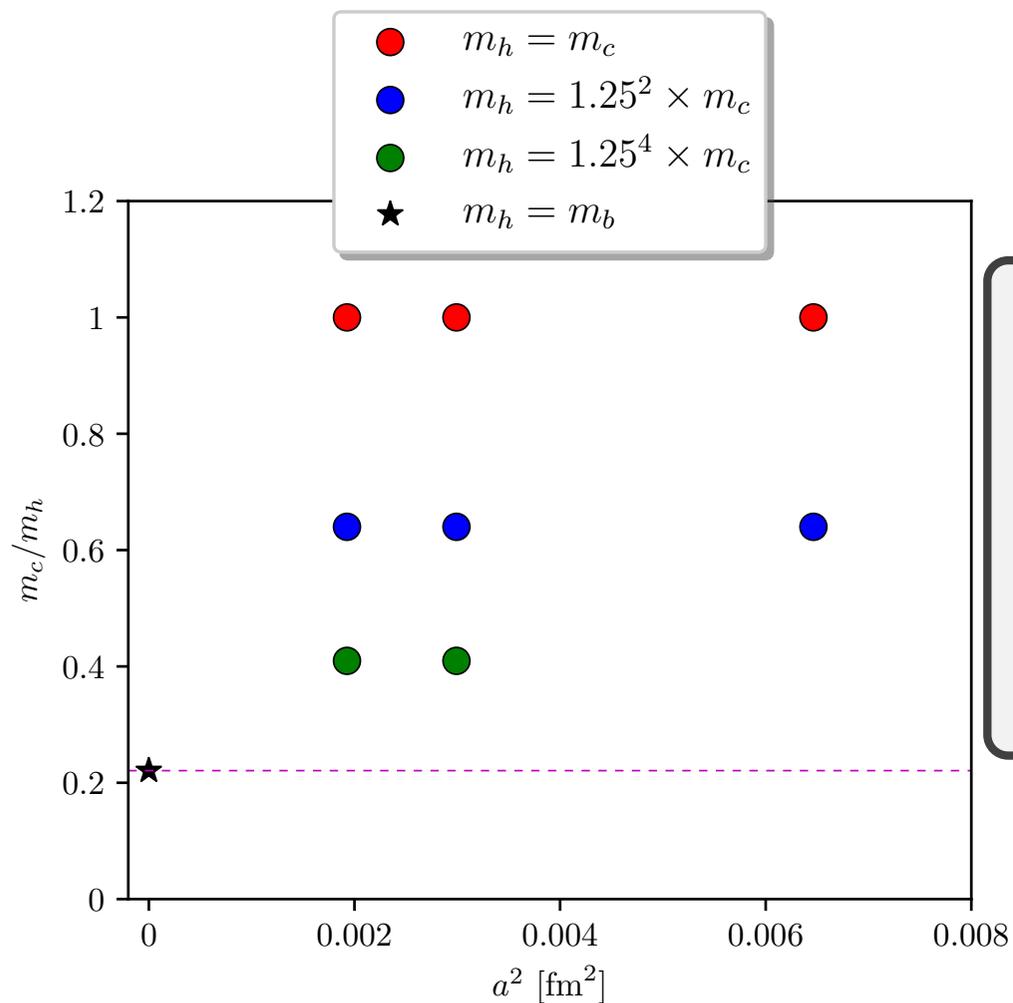
$$\beta = 4.17, 4.35, 4.47$$

$$a \approx 0.080, 0.055, 0.044 \text{ fm}$$

$$a^{-1} \approx 2.453, 3.610, 4.496 \text{ GeV}$$

$$L^3 \times T = 32^3 \times 64, 48^3 \times 96, 64^3 \times 128$$

# Valence quarks



Möbius Domain Wall Fermions

$m_l$  same values as sea quarks

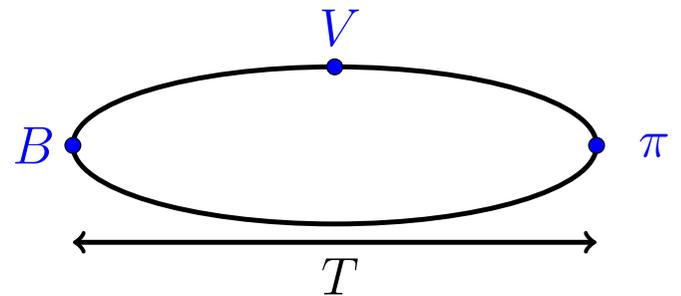
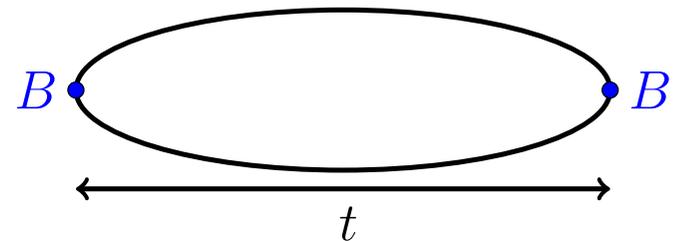
$m_h \in \{m_c, 1.25^2 m_c, 1.25^4 m_c\}$

○  $m_{H_l} \approx 1.95, 2.55, 3.40$  GeV

○ Fully relativistic formalism

# 2pt and 3pt correlators

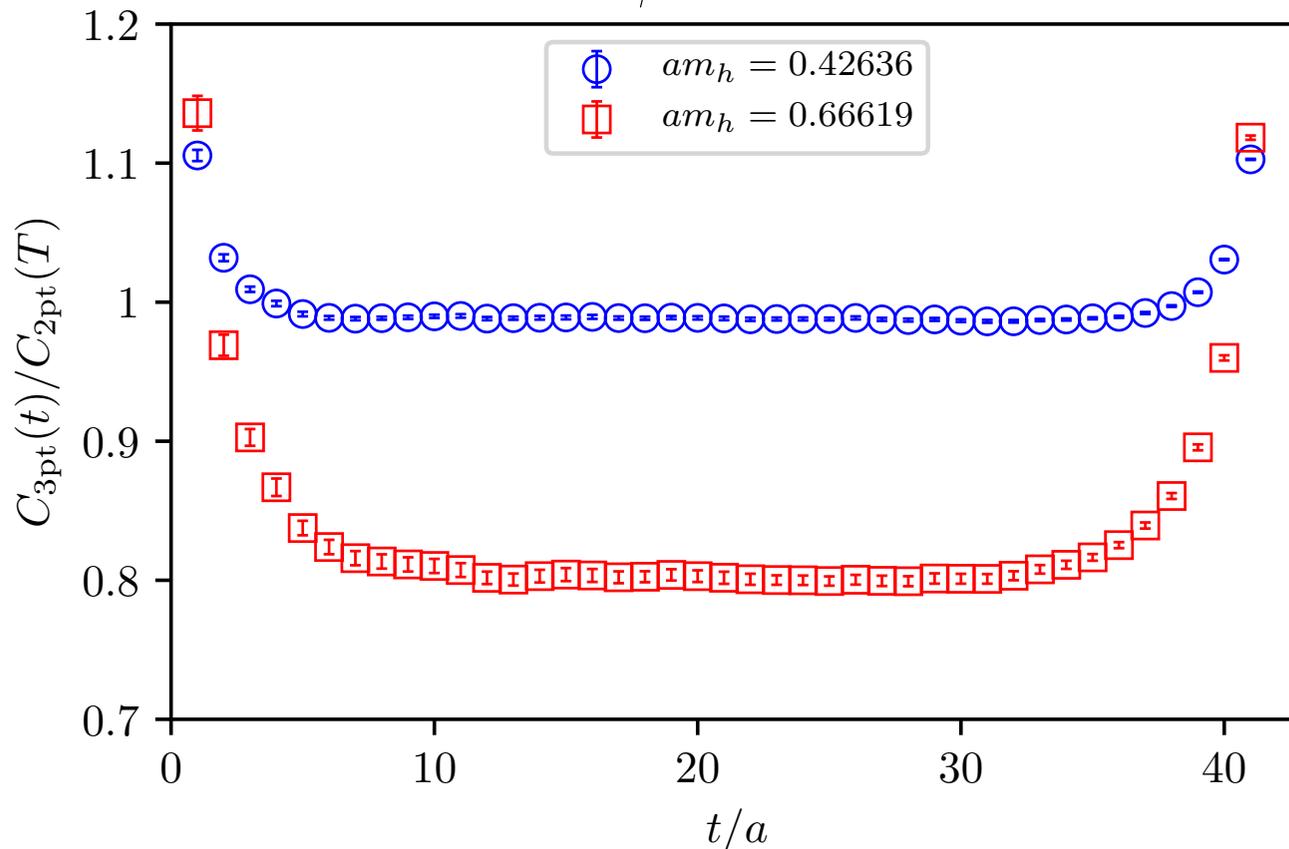
- Smeared sources with  $\mathbb{Z}_2$  noise
- Local and smeared sinks
- 1, 2 or 4 time sources
- $B$  and  $\pi$  mesons separated by time  $T$
- Vector operator  $V^\mu$  inserted at time  $0 < t < T$



# Renormalisation

- Renormalise with  $Z_{V,bl} = \sqrt{Z_{V,bb}Z_{V,ll}}$
- Require  $C_{3\text{pt}}(t)/C_{2\text{pt}}(T) = Z_{V,bb}^{-1}$
- In the continuum limit  $Z_{V,bb} = V_{V,ll}$

$$\beta = 4.35$$



# Extrapolating to $a \rightarrow 0$ ,

$$m_h \rightarrow m_b, m_l \rightarrow m_{u,d}$$

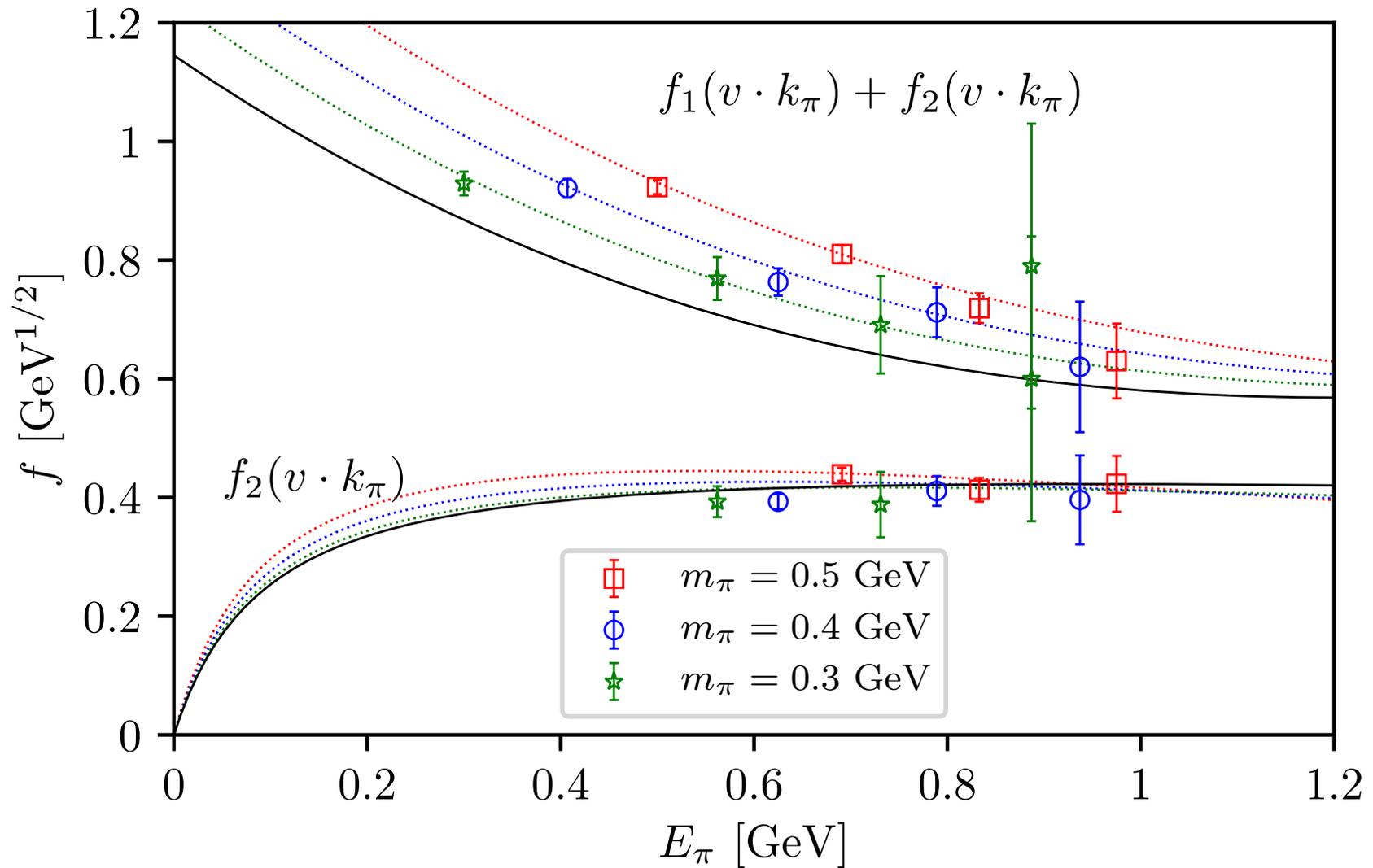
$$f_1(v \cdot k_\pi) + f_2(v \cdot k_\pi) = D_0 \left( 1 + \sum_{n=1}^3 D_{E_\pi^n} E_\pi^n + D_{M_\pi} M_\pi^2 \right. \\ \left. + \chi \log + D_{E_\pi M_\pi^2} E_\pi M_\pi^2 + \frac{D_{m_h}}{m_h} \right) (1 + D_{a^2} a^2)$$

$$f_2(v \cdot k_\pi) = C_0 \frac{E_\pi}{E_\pi + (M_{B^*} - M_B)} \left( 1 + C_{E_\pi} E_\pi + C_{M_\pi} M_\pi^2 \right) \\ + \chi \log + C_{E_\pi M_\pi^2} E_\pi M_\pi^2 + \frac{C_{m_h}}{m_h} \right) (1 + C_{a^2} a^2)$$

Pole  
dominance

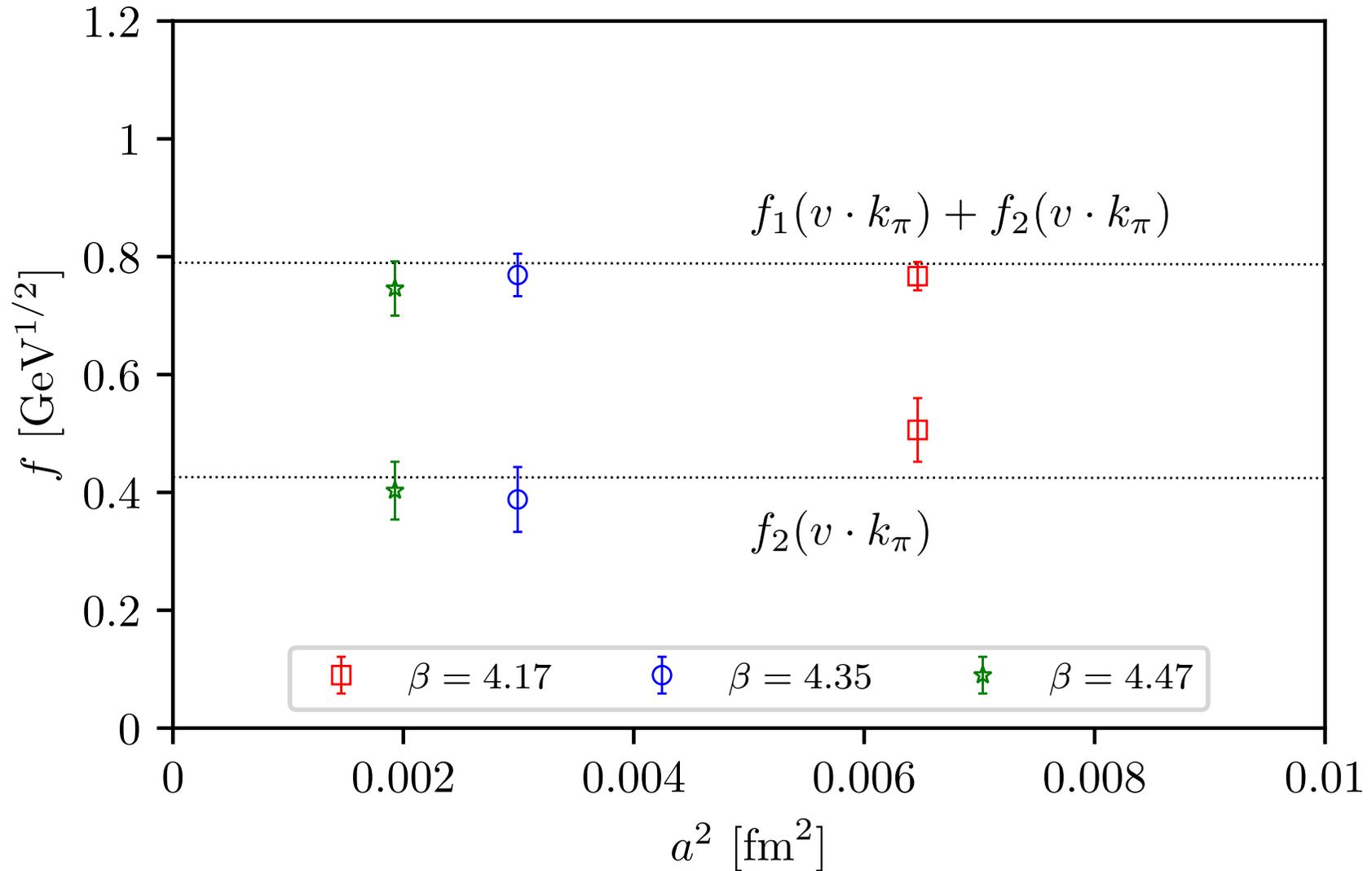
# Form factors vs $m_\pi$

$$\beta = 4.35; 1.25^2 \times m_c$$



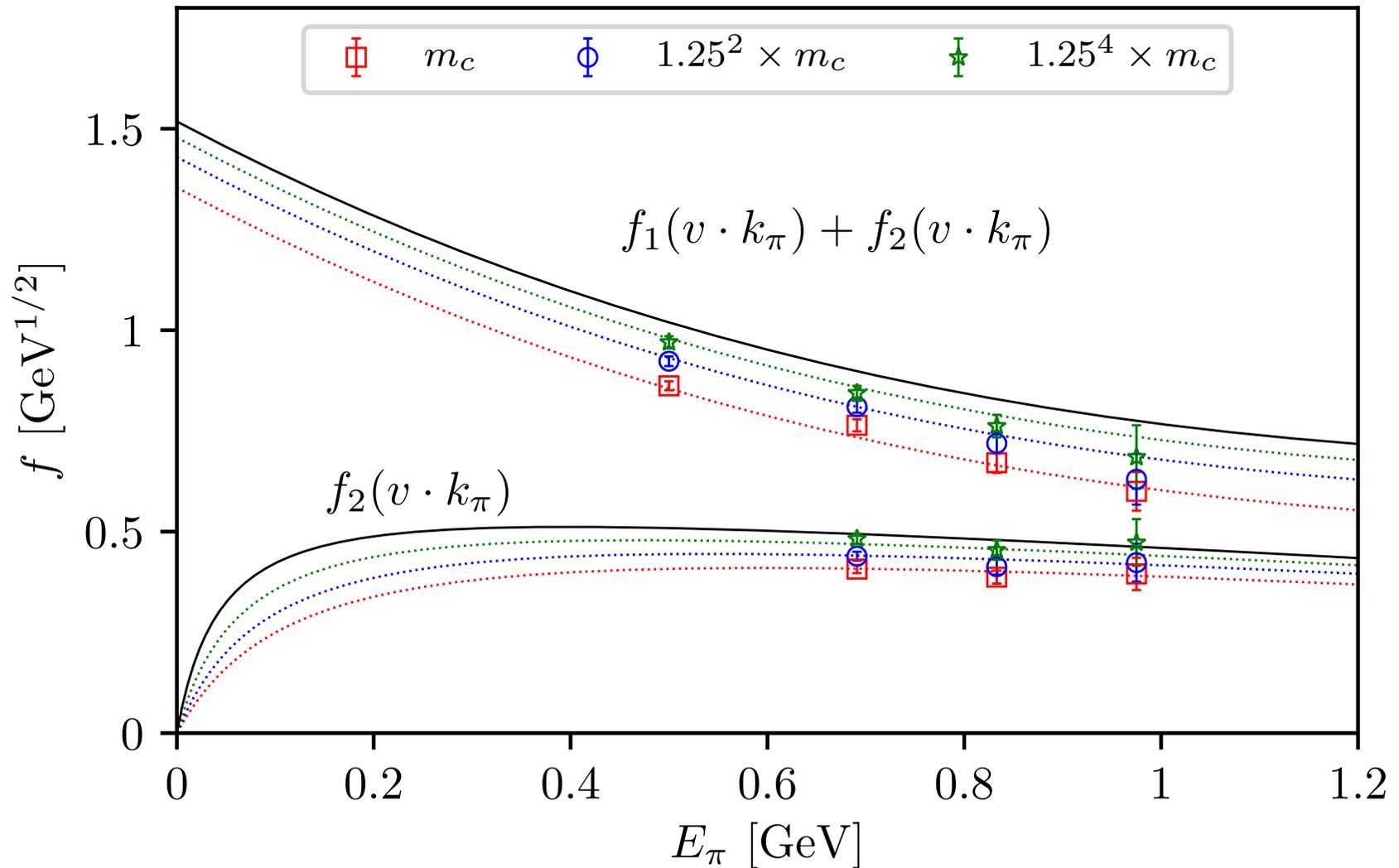
# Form factors vs $a^2$

$$|\mathbf{p}|^2 = (2\pi/L)^2; m_\pi = 0.3 \text{ GeV}; 1.25^2 \times m_c$$

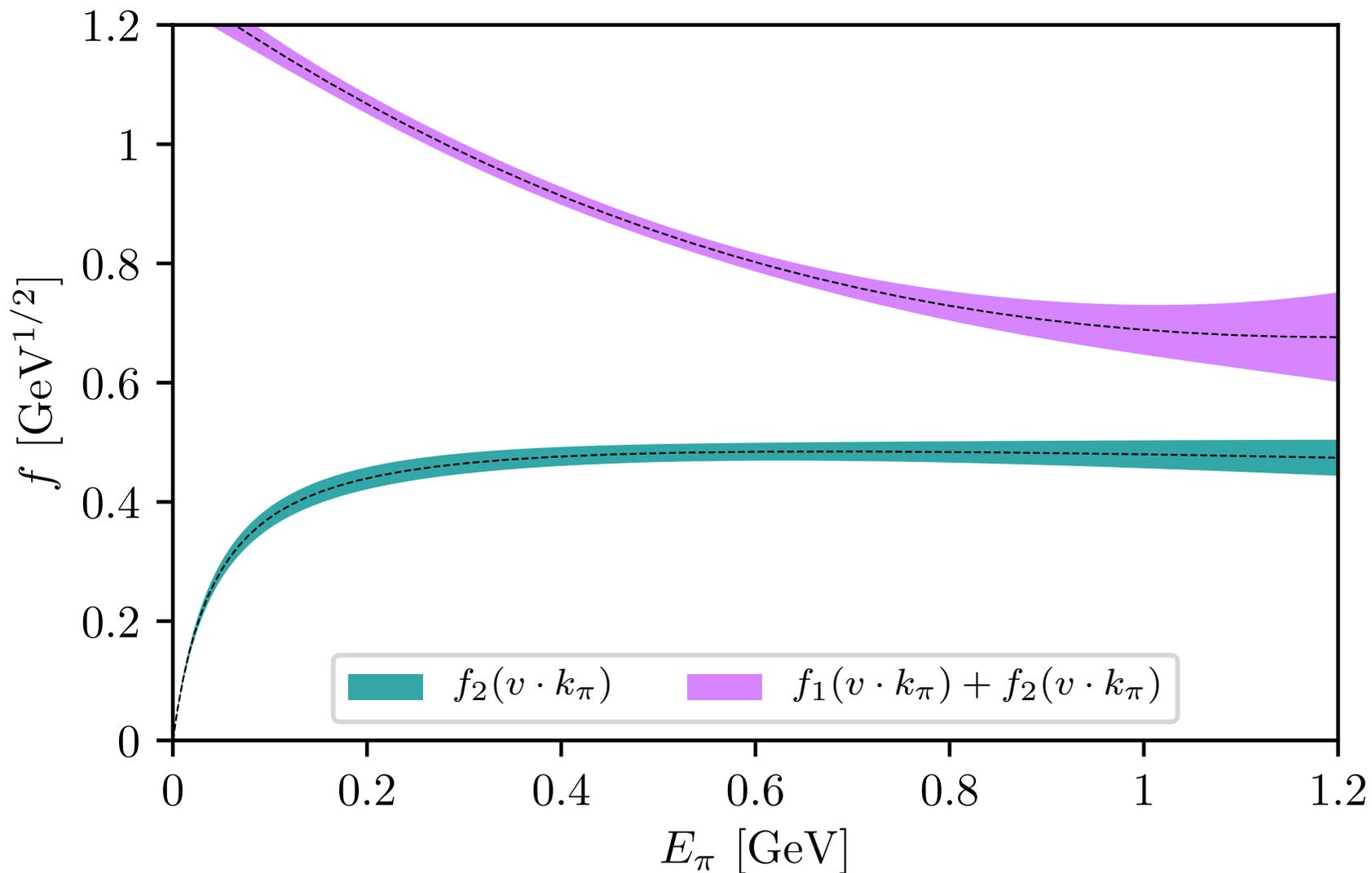


# Form factors vs $m_h$

$$\beta = 4.35; m_\pi = 0.5 \text{ GeV}$$



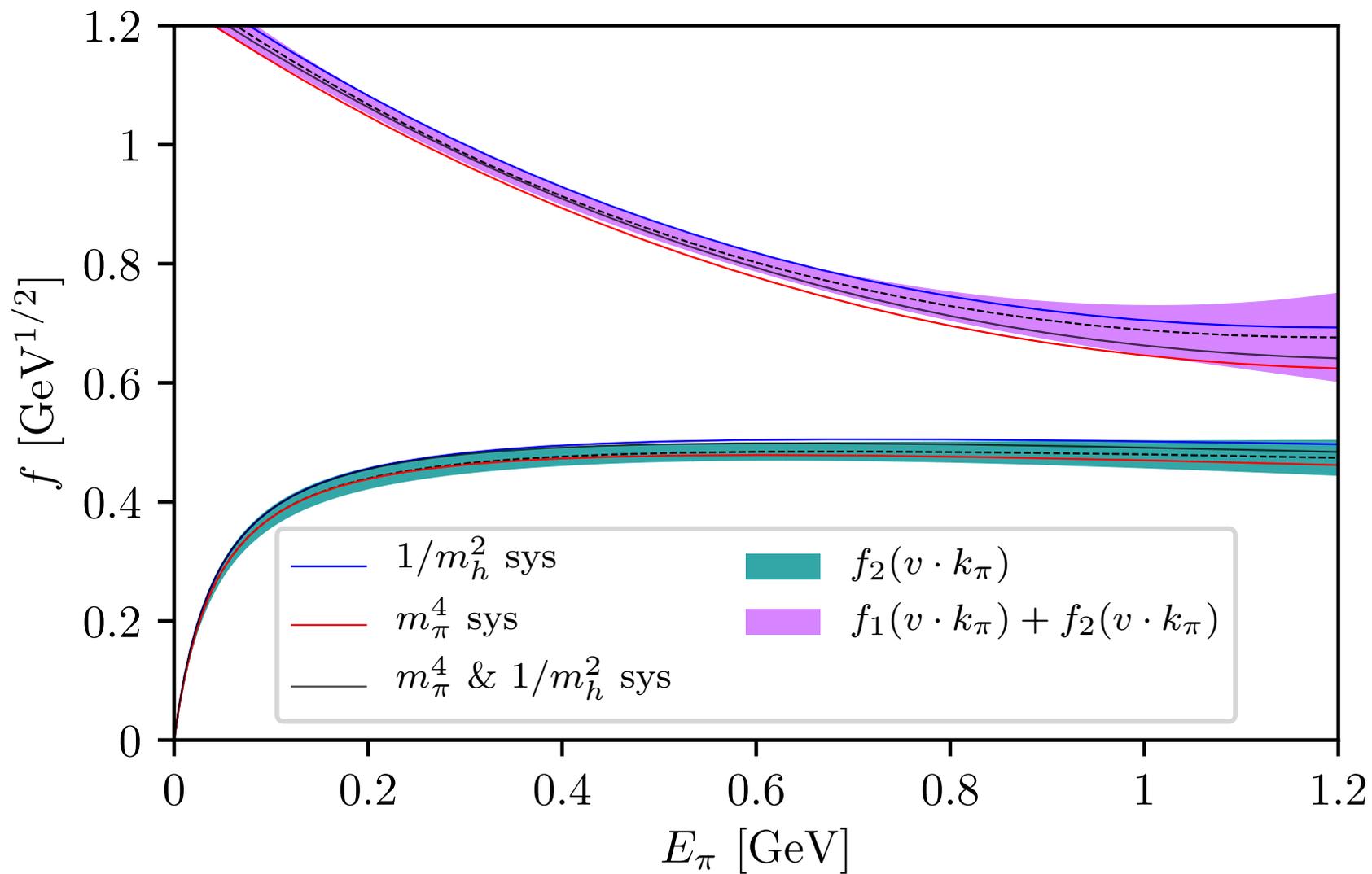
# Form factors $f_1+f_2$ and $f_2$



# Systematic effects

- Systematic effects vary across  $q^2$
- Adding higher powers of  $E_\pi$  ( $E_\pi^4$  and  $E_\pi^2$  for  $f_1+f_2$  and  $f_2$  respectively): a couple of %
- Adding  $M_\pi^4$  terms: around -5%
- Adding  $1/m_h^2$  terms: around +5%
- No chiral log: <1%

# Systematic effects



# $z$ -expansion

- $q^2$  range is large, and  $z$  is often used instead:

$$z = \frac{\sqrt{t_+ - q^2} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - q^2} + \sqrt{t_+ - t_0}}$$

where  $t_+ = (M_B + M_\pi)^2$  and we choose

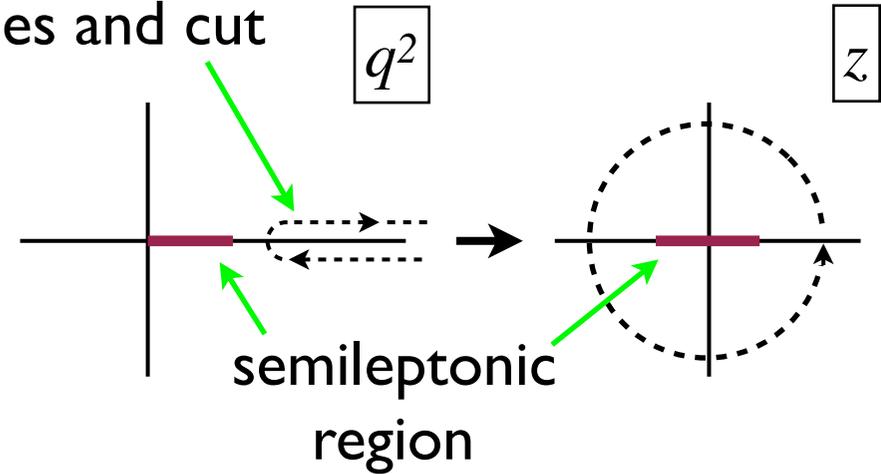
$$t_0 = (M_B + M_\pi)(\sqrt{M_B} - \sqrt{M_\pi})^2$$

- Full kinematic range is now  $-0.3 < z < 0.3$
- Fit functions for the form factors are

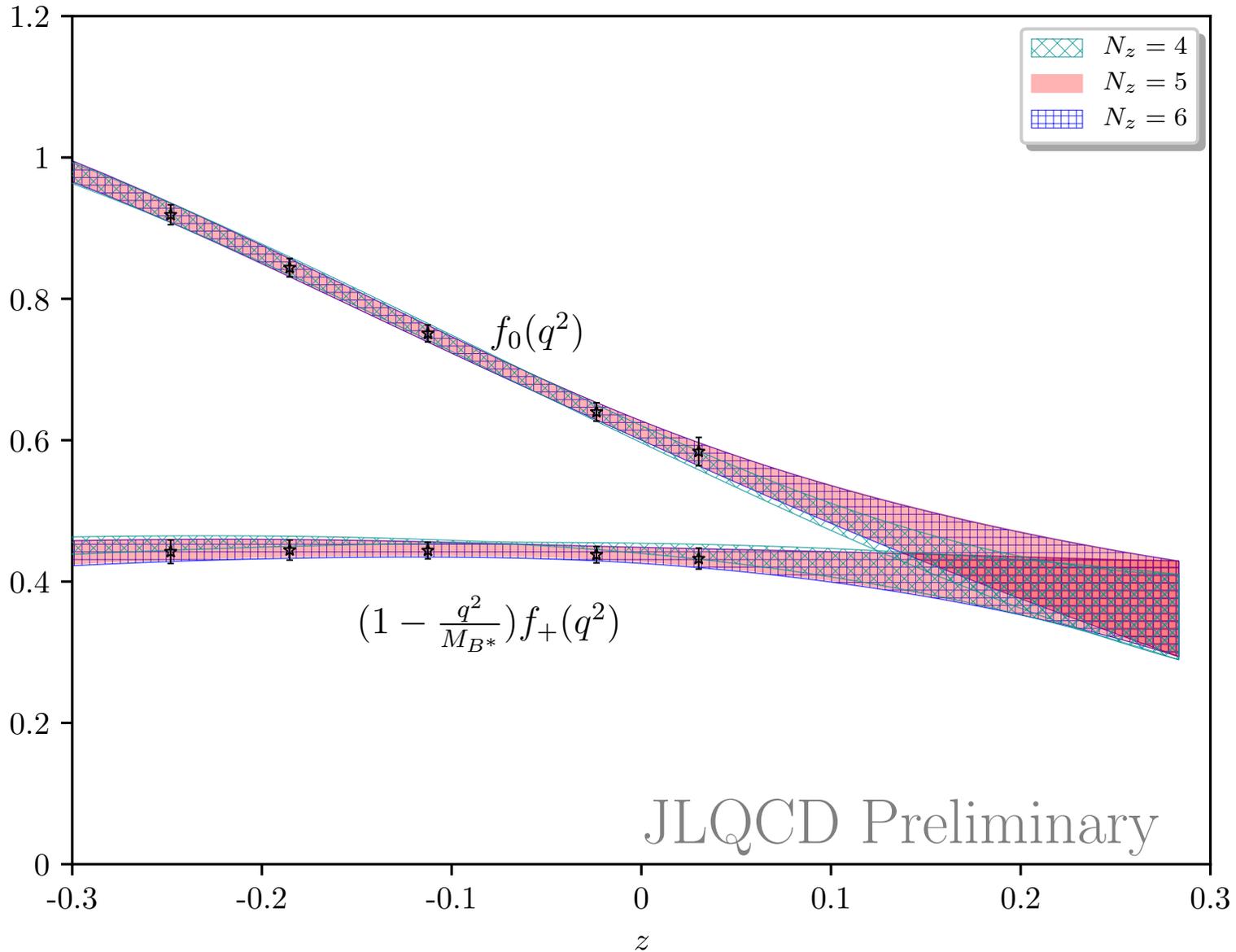
$$f_+(z) = \frac{1}{1 - q^2(z)/M_{B^*}^2} \sum_{n=0}^{N_z-1} b_n \left[ z^n - (-1)^{n-N_z} \frac{n}{N_z} z^{N_z} \right]$$

$$f_0(z) = \sum_{n=0}^{N_z-1} a_n z^n$$

poles and cut

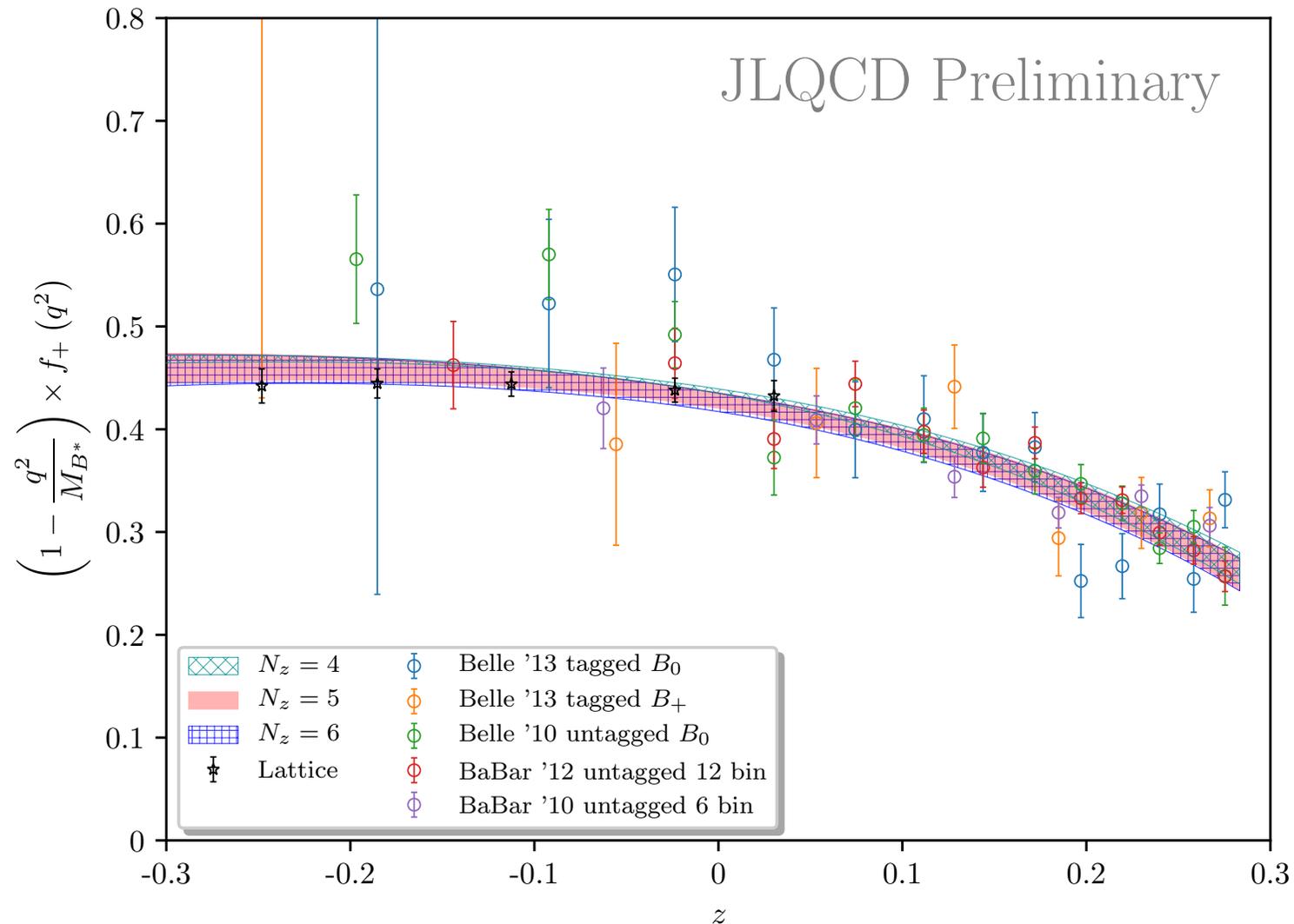


# $z$ -parameter expansion & fit

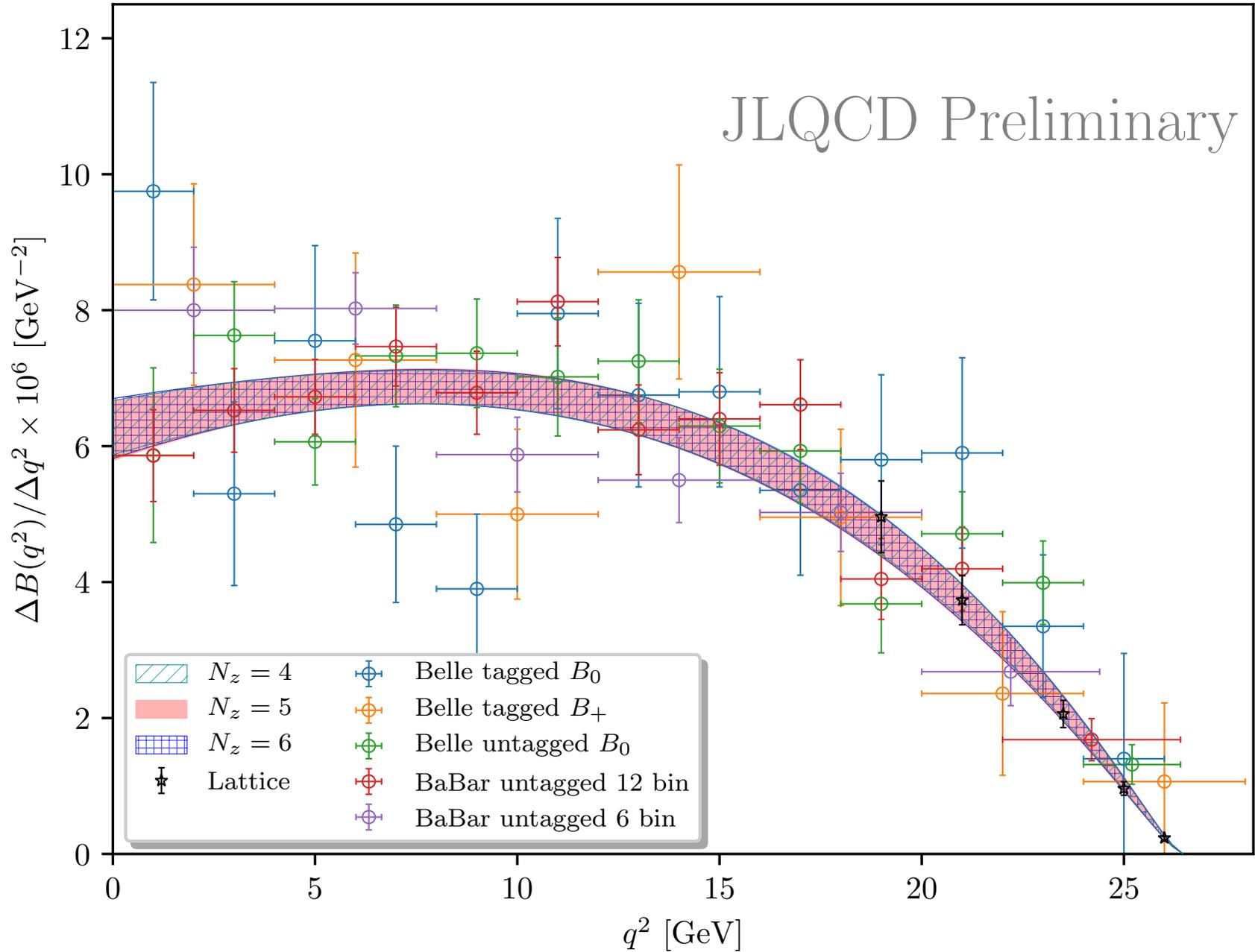


# Combined fit

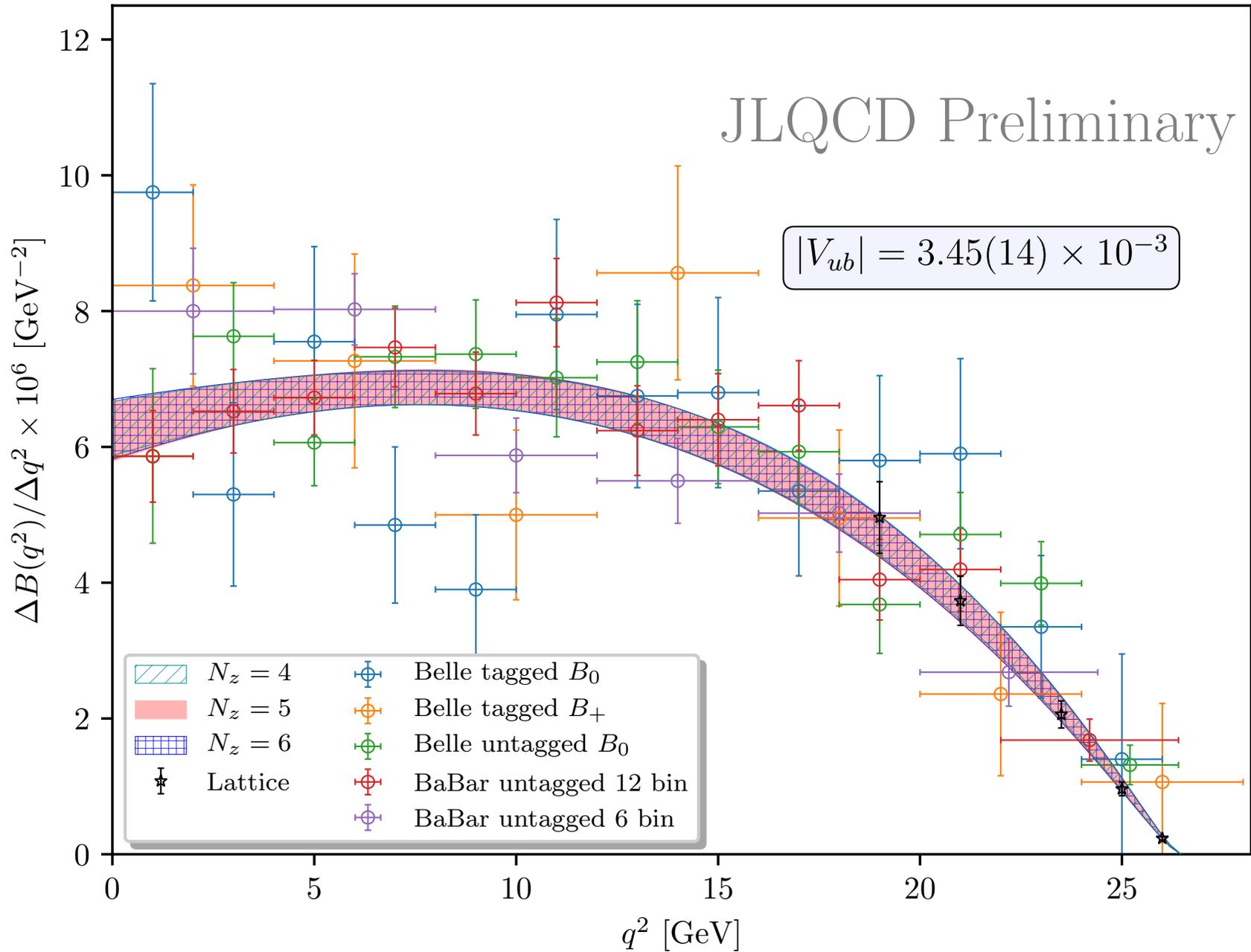
- Use  $z$ -expansion to fit both lattice and exptl. data, including  $V_{ub}$  as a fit parameter



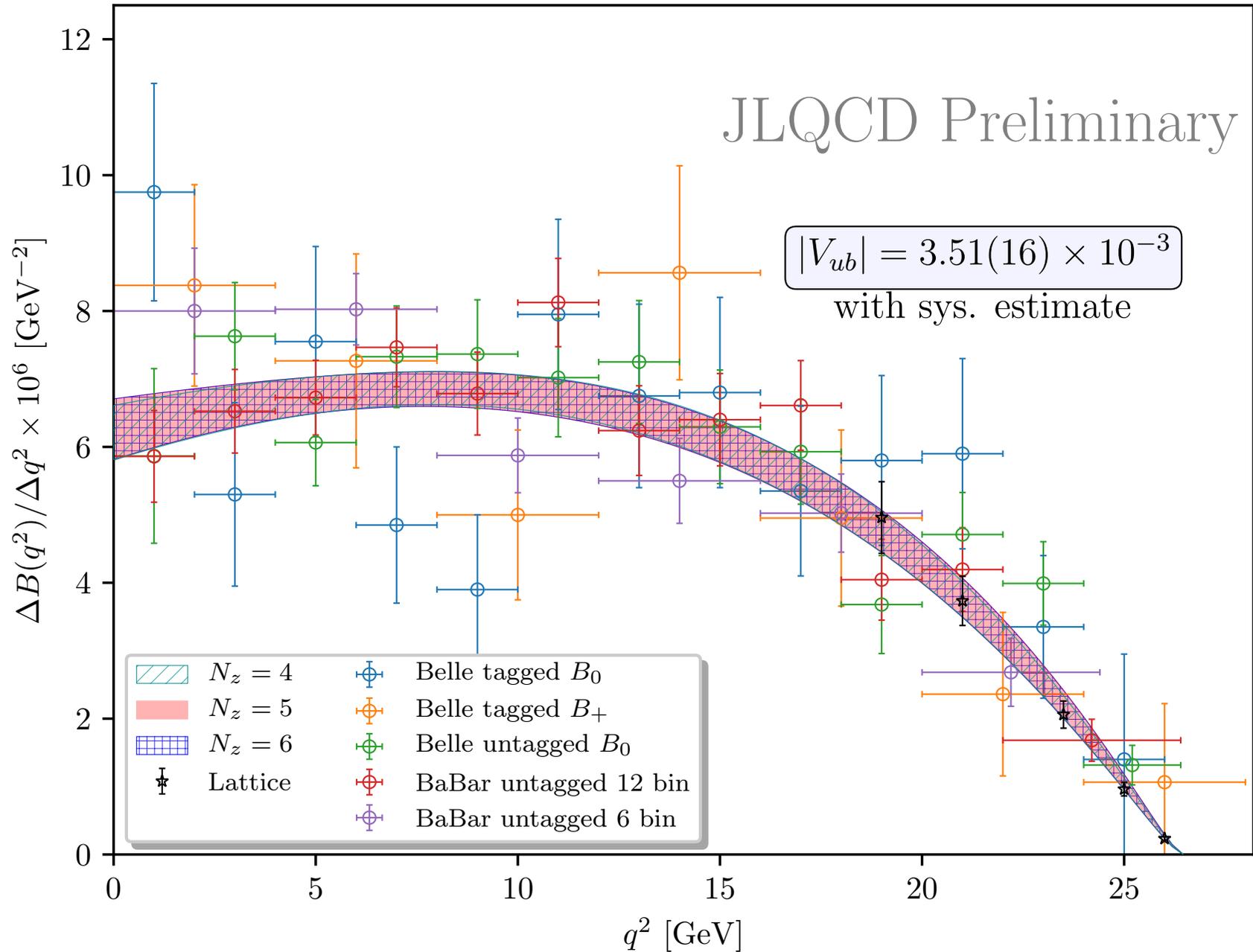
# Combined fit



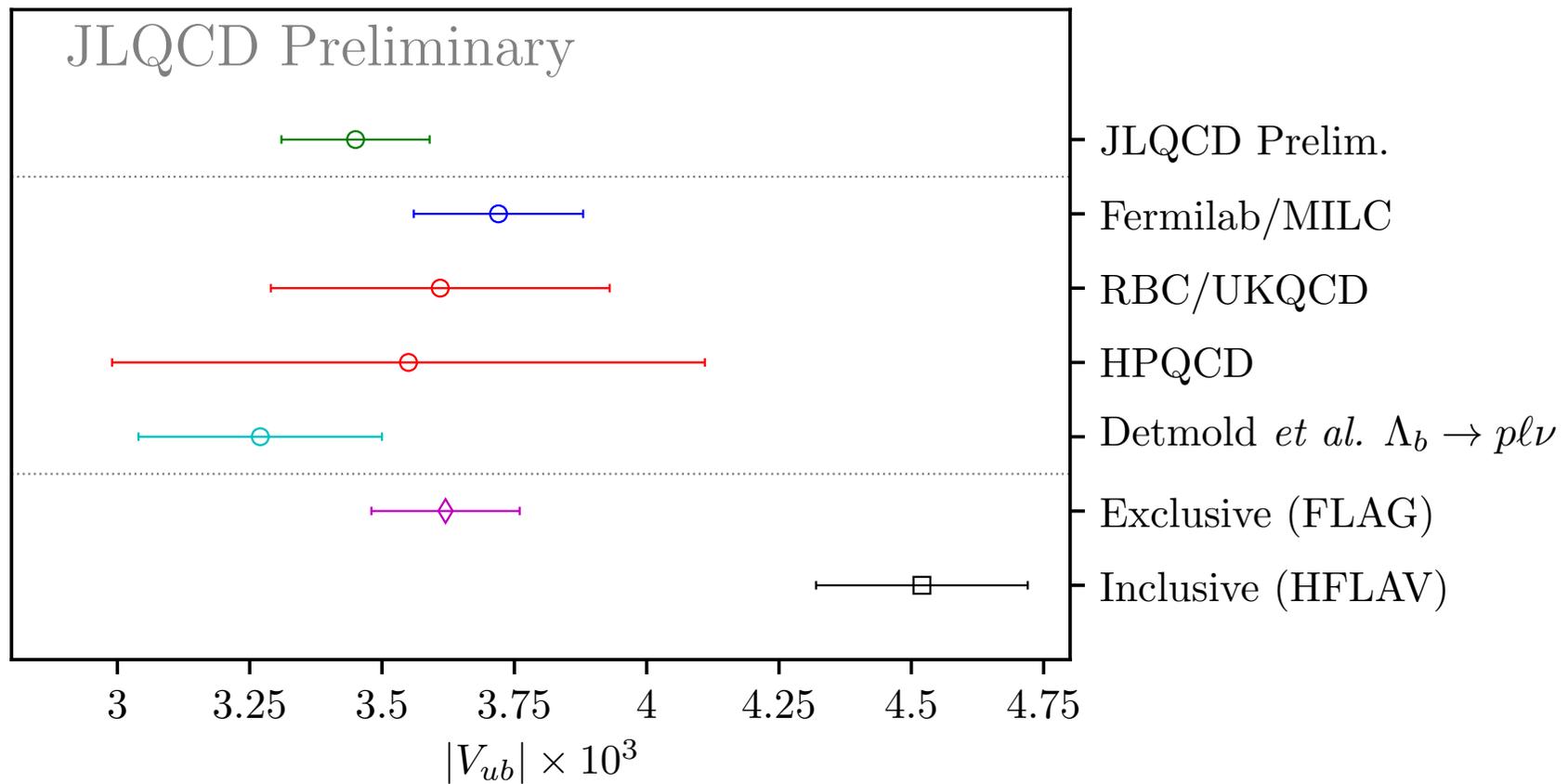
# Combined fit



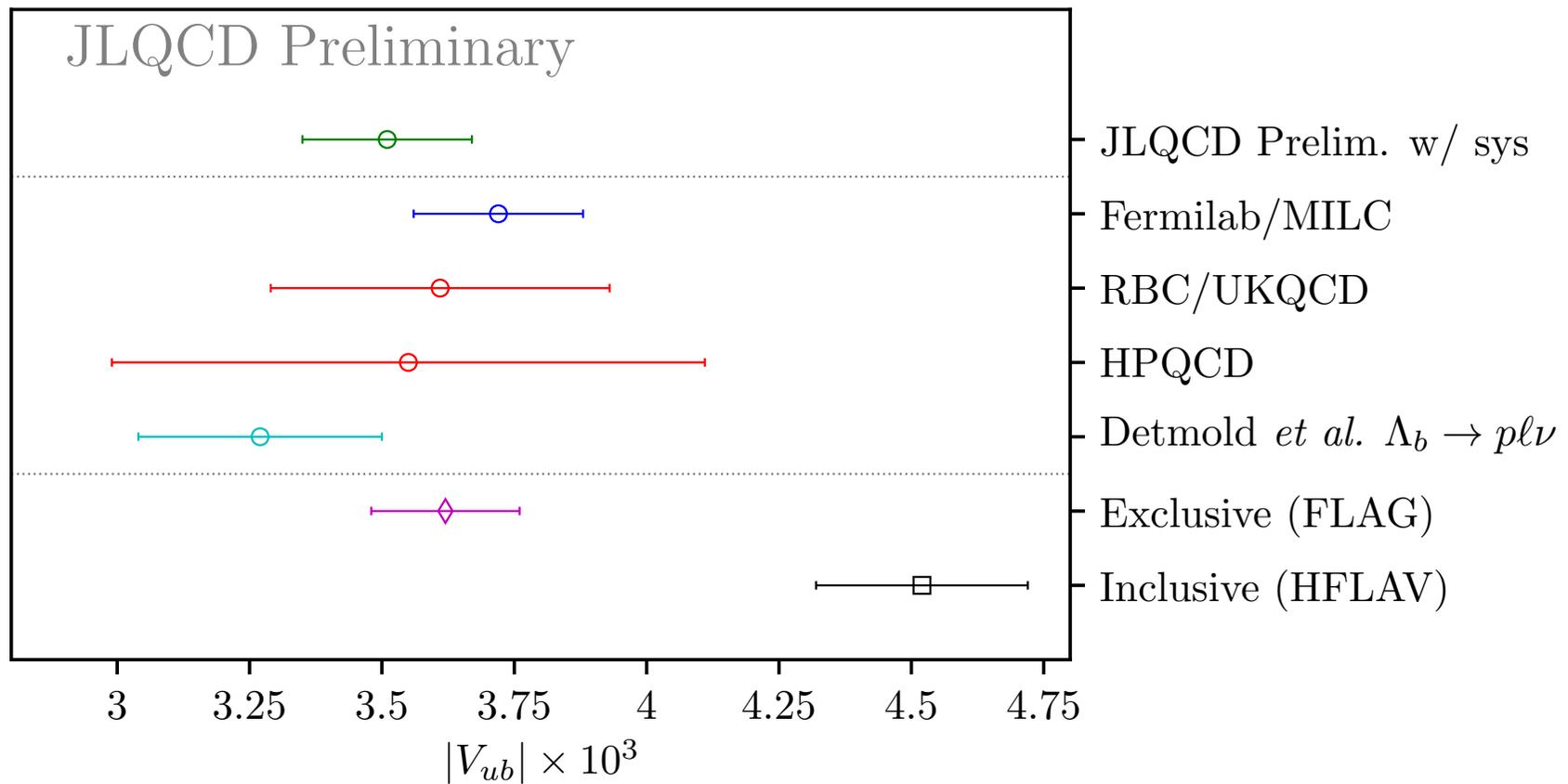
# Combined fit



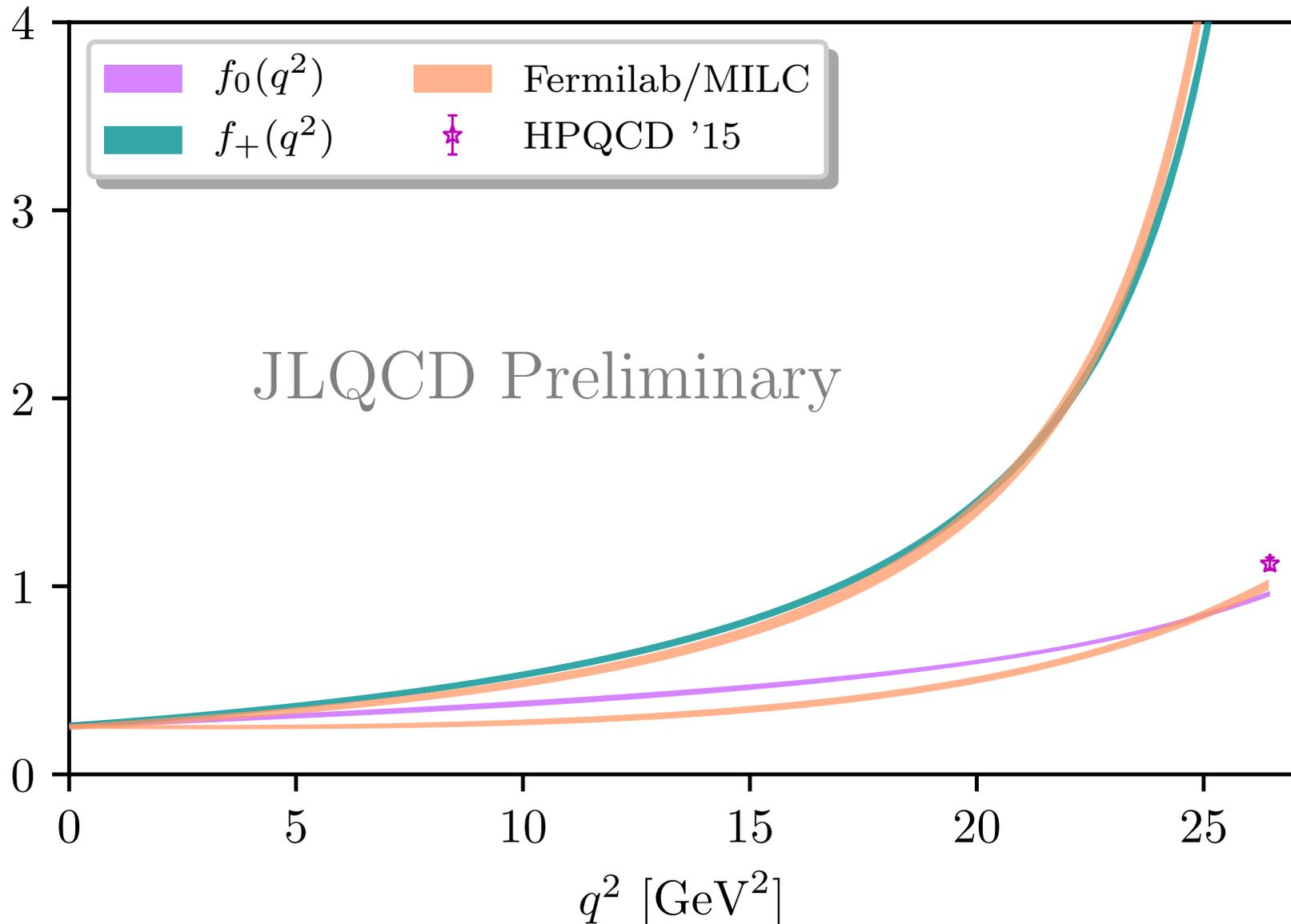
# CKM element $V_{ub}$



# CKM element $V_{ub}$



# Form factor comparison



# Summary

- $B \rightarrow \pi l \nu$  using Möbius Domain Wall fermions for all quarks
  - heavy quark treated relativistically
  - use multiple  $m_h > m_c$  and extrapolate to  $m_b$
  - use multiple light quark masses and extrapolate to physical  $M_\pi$
- Fit form factors calculated on the lattice together with experimental partial branching fractions to get  $|V_{ub}|$

# Thank you!



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