

# An exploratory study of heavy-light semileptonics using distillation

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*of* EDINBURGH



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- 1 Introduction
  - Motivation
  - heavy-light semileptonic decays
- 2 Distillation (in Grid)
- 3 Analysis
  - First glance at data
  - Comparison to  $Z_2$  noise with sequential solves
- 4 Conclusions & Outlook

- Semi-leptonic decays give access to CKM matrix elements, e.g.  $|V_{cd}|$  from  $D \rightarrow \pi \ell \nu$ .
- current tension in lepton flavour universality detected in  $R(D^{(*)}) = \frac{\mathcal{B}(B \rightarrow D^{(*)} \tau \nu_\tau)}{\mathcal{B}(B \rightarrow D^{(*)} \ell \nu_\ell)}$ 
  - $\Rightarrow$  clear first-principles determination needed
- Interesting processes suffer from bad signal-to-noise ratios
  - $\Rightarrow$  advanced methods needed
- Testing ground for recently implemented distillation code in Grid (& Hadrons)



[ [github.com/paboyle/Grid](https://github.com/paboyle/Grid) ]

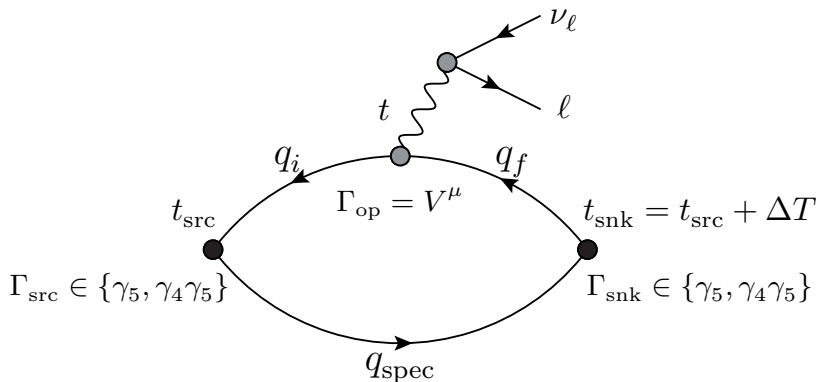
## Hadrons

[ [github.com/paboyle/Grid/tree/develop/Hadrons](https://github.com/paboyle/Grid/tree/develop/Hadrons) ]

Related RBC/UKQCD charm-to-bottom programme talks:

- Semileptonic  $B$  decays with RHQ  $b$  quarks [Mon 16:50, Ryan Hill]
- Neutral meson mixing and related observables in the  $D_{(s)}$  and  $B_{(s)}$  meson systems [Tue 15:40, Tobias Tsang]
- Semileptonic form factors for exclusive  $B_s \rightarrow K \ell \nu$  and  $B_s \rightarrow D_s \ell \nu$  decays [Tue, Poster, Oliver Witzel]

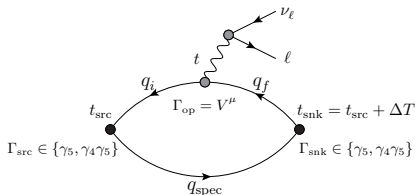
# heavy-light semileptonic decays



for  $D(B) \rightarrow \pi$ :  $q_{\text{spec}} = l$ ,  $q_i = h_1, h_2$ ,  $q_f = l$

# heavy-light semileptonic decays

Pointlike weak operator  
Three-point functions  $C_3^{PP}$  have the form:



$$\tilde{C}_3^{PP}(t, \Delta T) \propto \left( e^{E_i t} - e^{E_f(\Delta T - t)} \right).$$

we want to map out a large momentum transfer, so  $D(p_i = 0) \rightarrow \pi(p_f)$  are best suited.

We study different values of source-sink separations  $\Delta T$  and different momenta  $\mathbf{p}_f$ .

# Distillation

- small  $\Delta T$ : cannot isolate ground state
  - large  $\Delta T$ : bad signal-to-noise ratio
- ⇒ need good smearing technique and advanced numerical method
- Distillation with (stochastic) LapH might help with both [arxiv:0905.2160]  
[arxiv:1104.3870]

# Distillation

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[arxiv:1104.3870]

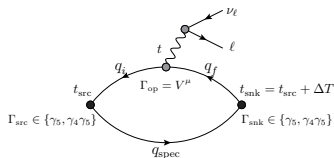
idea: construct a smearing matrix from  $N_{\text{vec}}$  low modes of the 3D lattice Laplacian (tuneable parameter):

$$S_{xy}(t) = \sum_{k=1}^{N_{\text{vec}}} V_k(x, t) V_k^\dagger(y, t),$$

- expensive only once, assembly of correlation functions (momenta,  $\Gamma$ -structure) is the last step of the computation
- Smeared propagators can be projected into smaller subspace
  - ⇒ Can be re-used in other projects



We want to compute:



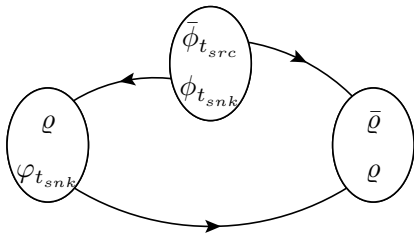
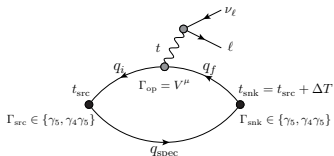
$$C_3 = \langle \Gamma_{\text{snk}} D^{-1}(t_{\text{snk}}, t) \Gamma_{\text{op}} D^{-1}(t, t_{\text{src}}) \Gamma_{\text{src}} D^{-1}(t_{\text{src}}, t_{\text{snk}}) \rangle.$$

The central semileptonic insertion must be unsmeared and we want to smear with  $S(t) = V(t)V^\dagger(t)$  (spatial sums implicit) only at the  $\Gamma_{\text{src}}, \Gamma_{\text{snk}}$  insertions:

$$C_3 = \langle \Gamma_{\text{snk}} V(t_{\text{snk}}) V^\dagger(t_{\text{snk}}) D^{-1}(t_{\text{snk}}, t) \Gamma_{\text{op}} D^{-1}(t, t_{\text{src}}) V(t_{\text{src}}) V^\dagger(t_{\text{src}}) \Gamma_{\text{src}} V(t_{\text{src}}) V^\dagger(t_{\text{src}}) D^{-1}(t_{\text{src}}, t_{\text{snk}}) V(t_{\text{snk}}) V^\dagger(t_{\text{snk}}) \rangle.$$

# Distillation

using  $\gamma_5$  hermiticity we can invert some of the quark lines and can write, using meson fields



- $\phi_{t'}(x, t) = D_{x,t,x',t'}^{-1} V_x(t')$  are the unsmeared sinks, i.e. a solve on a source with support on timeslice  $t'$ . [arxiv:1403.5575]
- Constructed from sources on two timeslices,  $t_{src}$  and  $t_{snk}$ .
- $\phi$  cannot be projected into a smaller subspace (i.e. into a perambulator object).

# Lattice setup

- RBC-UKQCD's 2+1 flavour domain wall fermions
- feasibility study on  $L^3 \cdot T = 24^3 \cdot 64$  lattice,  $m_\pi \approx 340\text{MeV}$
- one light quark ( $am_l = 0.005, M_5 = 1.8, L_s = 16$ )
- two different heavy-quark masses with  $am_h = 0.58$  and  $am_h = 0.64$  using a stout-smearred action ( $\rho = 0.1, N = 3$ ) with  $M_5 = 1.0, L_s = 12$  and Moebius-scale = 2 [arxiv:1812.08791]

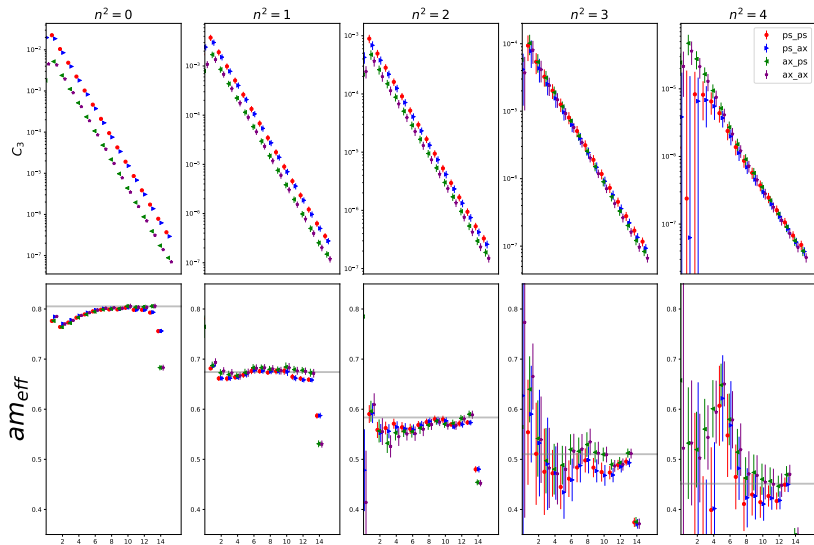
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current level of statistics:

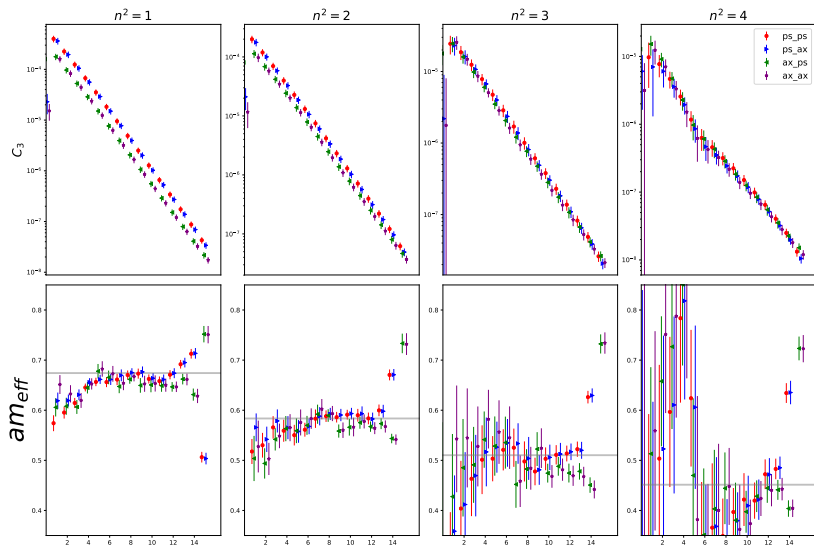
- 2 configurations
- 16 solves on each config
- 4 different  $\Delta T = 12, 16, 20, 24$
- all lattice momenta up to  $n^2 = 4$

$$V^0, \Delta T = 16, m_h = 0.58$$



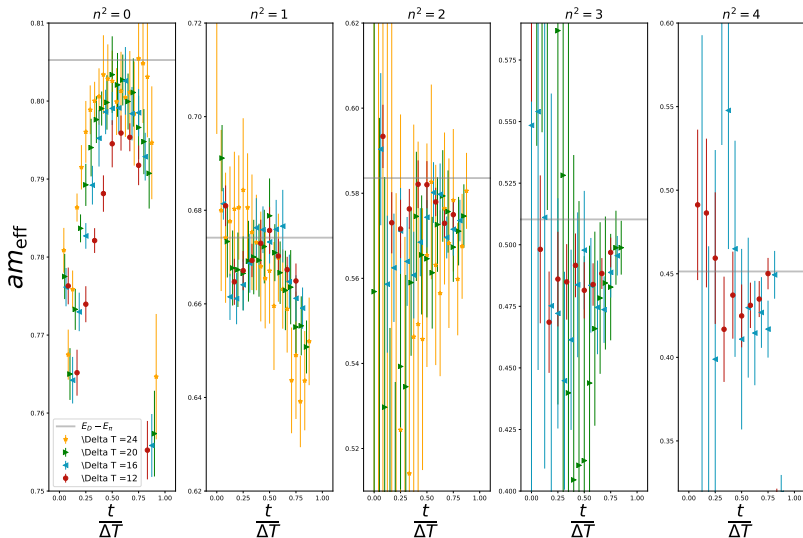
grey bands: free energies  $E_D - E_\pi(p_f)$

$$V^i, \Delta T = 16, m_h = 0.58$$



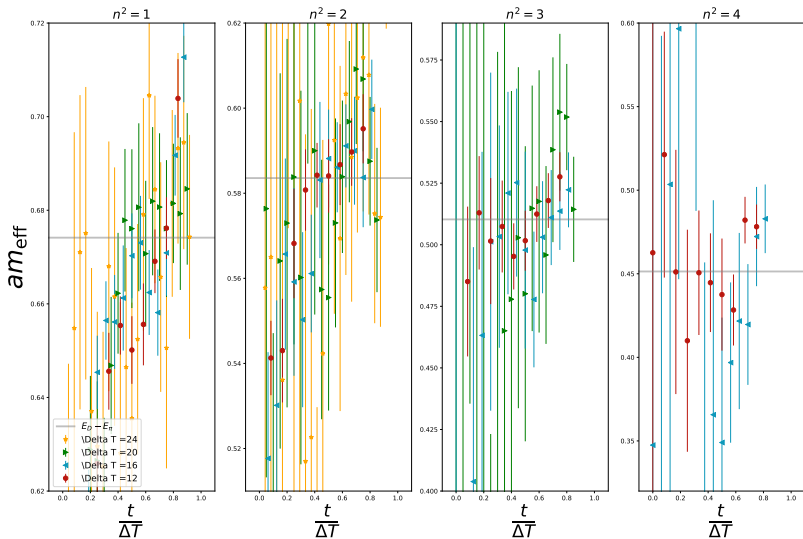
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# $V^0$ , comparison of $\Delta T$ (pp channel), $m_h = 0.58$



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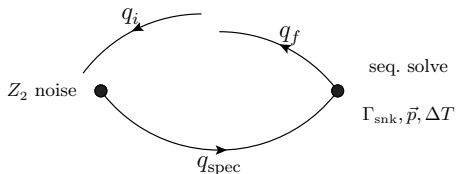


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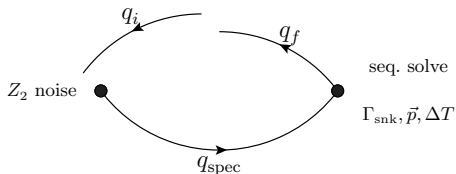
## $Z_2$ noise with sequential solves

- $Z_2$  noise at source
  - exploit  $\gamma_5$  hermiticity for  $q_i$  quark
  - compute  $q_{\text{spec}}$  quark line
  - sequential solve on  $q_{\text{spec}}$  quark line
- $\Rightarrow$  1 inversion for each  $\vec{p}$ ,  $\Gamma_{\text{snk}}$ ,  $\Delta T$



# $Z_2$ noise with sequential solves

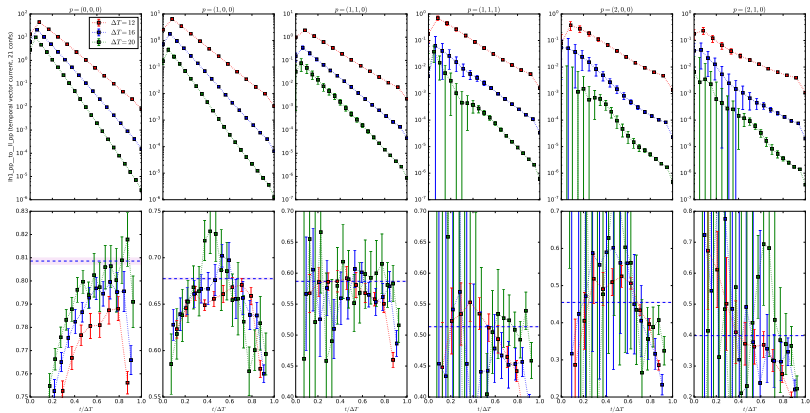
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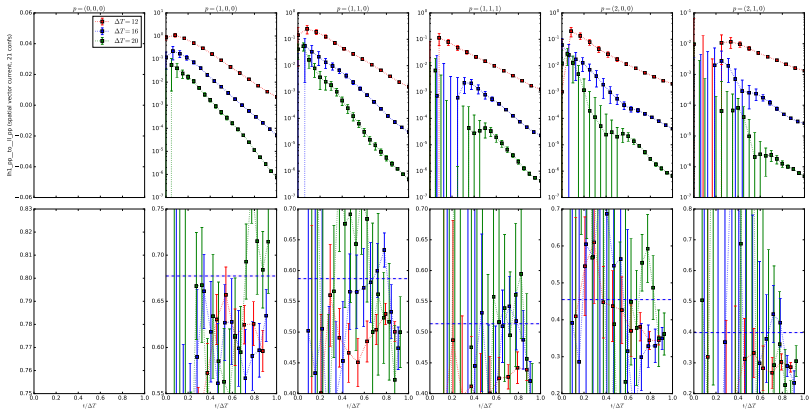
current level of statistics:

- 21 configurations
- 2 solves on each config
- 3 different  $\Delta T = 12, 16, 20$
- one lattice momentum each up to  $n^2 = 5$

# $V^0$ , comparison of $\Delta T$ (pp channel)



# $V^i$ , comparison of $\Delta T$ (pp channel)



# cost of production per configuration

Presented here:  $q_{\text{spec}} = l$ ,  $q_f = l$ ,  $q_i = h_{1,2,\dots}$

	Distillation	$Z_2$ seq.
#Inv / conf / $t_{\text{src}}$	$N_{\text{vec}} \times 4$	$N_{\Delta T} \times N_{\vec{p}} \times N_{\Gamma_{\text{snk}}}$
total #Inv	7680	1008

⇒  $\approx$  factor 8 in inversion cost for Distillation, plus non-negligible cost for meson fields (another factor 2-3)

- $\approx$  factor 2-6 error reduction in Distillation, depending on  $\mathbf{p}$  and  $\Gamma_{\text{op}}$

⇒ no clear winner

Possible setup we are interested in:

$D \rightarrow \pi$ ,  $D \rightarrow K$ ,  $D_s \rightarrow K$ ,  $D_{(s)} \rightarrow D'_{(s)}$ :

$q_{\text{spec}} = l, s$ ,  $q_f = l, s, h$ ,  $q_i = h_{1,2,\dots}$

	Distillation	$Z_2$ seq.
#Inv / conf / $t_{\text{src}}$	$N_{\text{vec}} \times 4$	$N_{\Delta T} \times N_{\vec{p}} \times N_{\Gamma_{\text{snk}}} \times \#\{q_{\text{spec}}, q_f\}$
total #Inv	unchanged	increased by factor $\approx 5$

⇒ Cost of Distillation might pay off

# Conclusions & Outlook

## Conclusions:

- computed 2pt-functions and 3pt-functions to study heavy-light semileptonics using the newly implemented distillaton code in Grid and Hadrons
- Distillation is expensive, but has several advantages:
  - different momenta and  $\Gamma_{\text{snk}}$  are free
  - smearing automatically implemented
  - perambulators (inversions) can be re-used for other projects

## Outlook:

- we want to study  $D \rightarrow \pi$ ,  $D \rightarrow K$ ,  $D_s \rightarrow K$ ,  $D_{(s)} \rightarrow D'_{(s)}$  with larger statistics on the  $24^3$  ensemble
- Some interesting RBC-UKQCD ensembles in production at the moment (2.8GeV at the physical point) [Thu, 14:00, Robert Mawhinney]

## Longer-term goal:

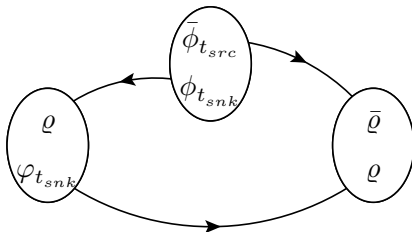
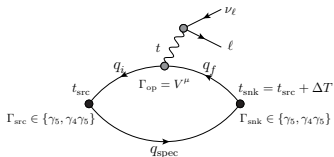
- use ensembles with smaller lattice spacing to extrapolate to physical  $B$

Thank you!

Thank you!

# Distillation

using  $\gamma_5$  hermiticity we can invert some of the quark lines and can write, using meson fields



$$C = M_{\Gamma_{op}}(\bar{\phi}_{t_{src}}, \phi_{t_{snk}}, t) M_{\Gamma_{src}}(\varphi_{t_{snk}}, \varrho, t_{src}) M_{\Gamma_{snk}}(\bar{\varrho}, \varrho, t_{snk}),$$



$$C = M_{\Gamma_{op}}(\bar{\phi}_{t_{src}}, \phi_{t_{snk}}, t) M_{\Gamma_{src}}(\varphi_{t_{snk}}, \varrho, t_{src}) M_{\Gamma_{snk}}(\bar{\varrho}, \varrho, t_{snk}),$$

where

$$\varrho_{a\alpha}^{[n,d]}(\vec{x}, t) = \sum_{k,l,t',\beta} v_{ka}(\vec{x}; t) P_{k\alpha,l\beta}^{[d]}(t, t') \rho_{l\beta}^{[n]}(t'),$$

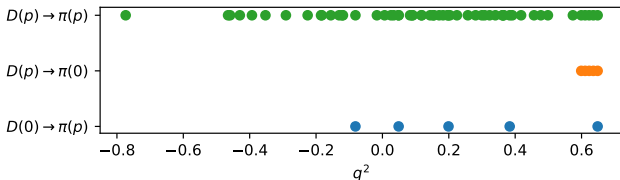
$$\phi_{t,a\alpha}^{[n,d]}(\vec{x}', t') = \sum_{a,b,\beta,t,\vec{x}} D_{a\alpha,b\beta}^{-1}(\vec{x}', t'; \vec{x}, t) \varrho_{b\beta}^{[n,d]}(\vec{x}, t),$$

$$\tau_{t,k\alpha}^{[n,d]}(t') = \sum_{a,b,\beta,t,\vec{x}',\vec{x}} v_{ka}(\vec{x}'; t') \phi_{t,a\alpha}^{[n,d]}(\vec{x}', t'),$$

$$\varphi_{t,a\alpha}^{[n,d]}(\vec{x}, t) = \sum_k v_{ka}(\vec{x}; t) \tau_{k\alpha}^{[n,d]}(t).$$

$$M_{\Gamma}^{[n_1, d_1; n_2, d_2]}(\varphi_q, \varrho_{q'}, t, \vec{p}) \\ = \sum_{\vec{x}, a, \alpha, \beta} e^{-i\vec{p} \cdot \vec{x}} (\varphi_q)_{a\alpha}^{[n_1, d_1]}(\vec{x}, t) \Gamma_{\alpha\beta}(\varrho_{q'})_{a\beta}^{[n_2, d_2]}(\vec{x}, t).$$

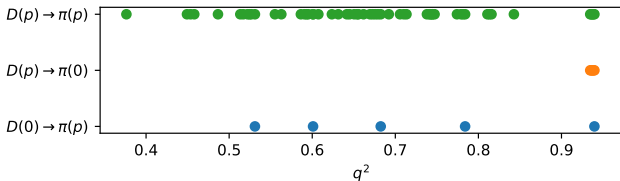
# $q^2$ range, $24^3$ ensemble



$$\begin{aligned} q^2 &= (E_D - E_\pi)^2 - (\mathbf{p}_D - \mathbf{p}_\pi)^2 \\ &= m_D^2 + \mathbf{p}_D^2 + m_\pi^2 + \mathbf{p}_\pi^2 - 2\sqrt{m_D^2 + \mathbf{p}_D^2}\sqrt{m_\pi^2 + \mathbf{p}_\pi^2} - (\mathbf{p}_D - \mathbf{p}_\pi)^2 \end{aligned}$$

$$\mathbf{p}_{\max}^2 = 4$$

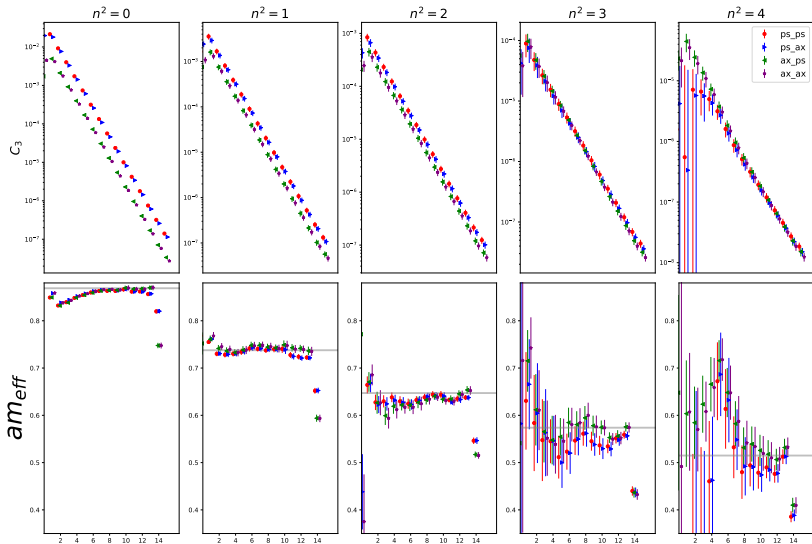
# $q^2$ range, physical meson masses



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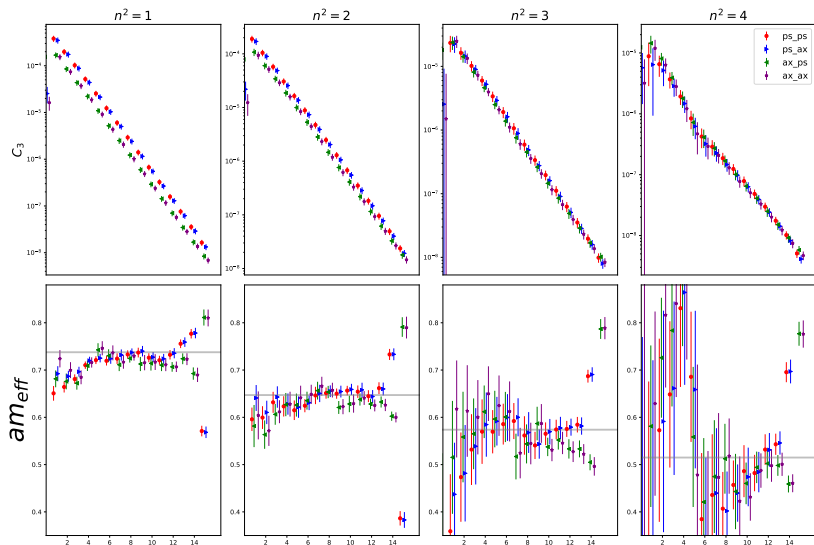
$$\mathbf{p}_{\max}^2 = 4$$

$$V^0, \Delta T = 16, m_h = 0.64$$



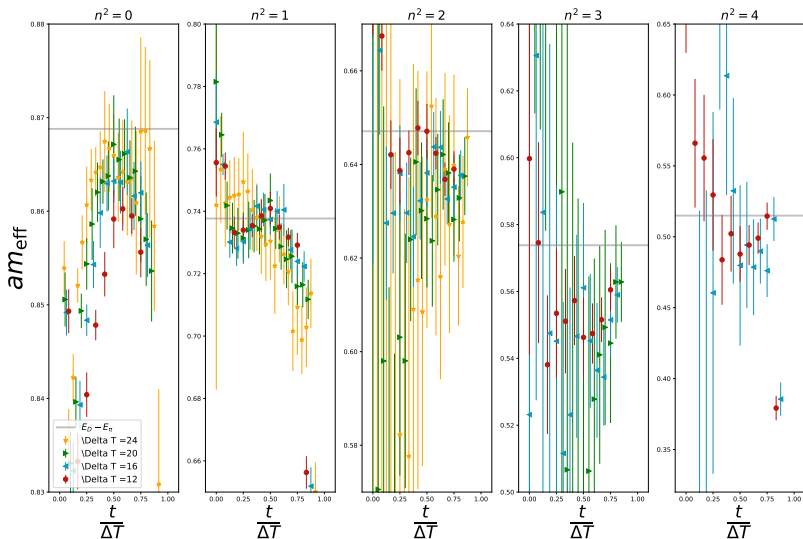
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$$V^i, \Delta T = 16, m_h = 0.64$$



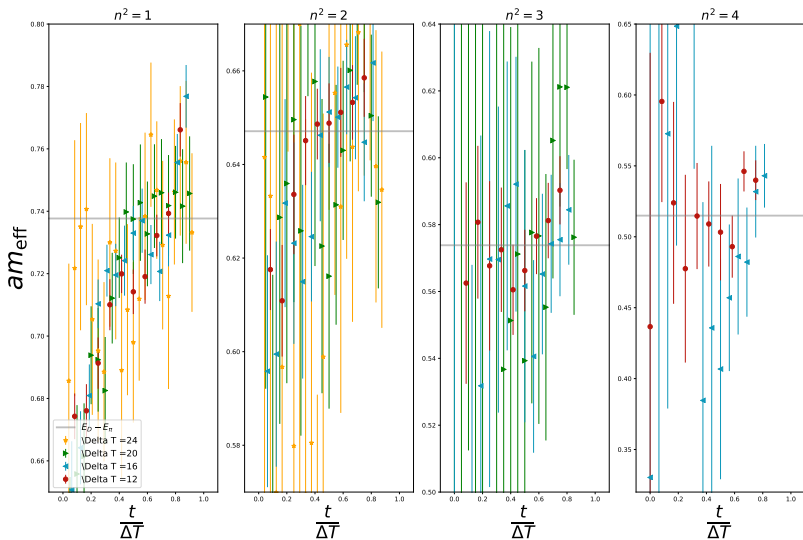
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# $V^0$ , comparison of $\Delta T$ (pp channel), $m_h = 0.64$



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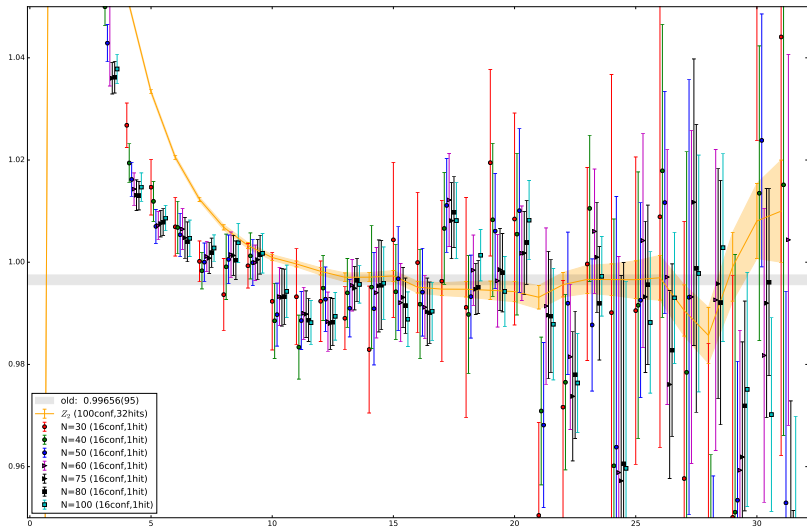
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# $N_{\text{vec}}$ comparison, pp



# $N_{\text{vec}}$ comparison, ii

