

Phase structure and real-time dynamics of the massive Thirring model in 1+1 dimensions using tensor-network methods

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Lattice 2019

Wuhan

18/06/2019

The 1+1 dimensional Thirring model and its duality to the sine-Gordon model

$$S_{\text{Th}} [\psi, \bar{\psi}] = \int d^2x \left[\bar{\psi} i \gamma^\mu \partial_\mu \psi - m_0 \bar{\psi} \psi - \frac{g}{2} (\bar{\psi} \gamma_\mu \psi)^2 \right]$$



strong-weak duality $g \leftrightarrow t$

Thirring	sine-Gordon
$\bar{\psi} \gamma_\mu \psi$	$\frac{1}{2\pi} \epsilon_{\mu\nu} \partial_\nu \phi$
$\bar{\psi} \psi$	$\frac{\Lambda}{\pi} \cos \phi$

$$S_{\text{SG}} [\phi] = \int d^2x \left[\frac{1}{2} \partial_\mu \phi(x) \partial^\mu \phi(x) + \frac{\alpha_0}{t} \cos \left(\sqrt{t} \phi(x) \right) \right]$$

$$\xrightarrow{\phi \rightarrow \phi / \sqrt{t}} \frac{1}{t} \int d^2x \left[\frac{1}{2} \partial_\mu \phi(x) \partial^\mu \phi(x) + \alpha_0 \cos (\phi(x)) \right]$$

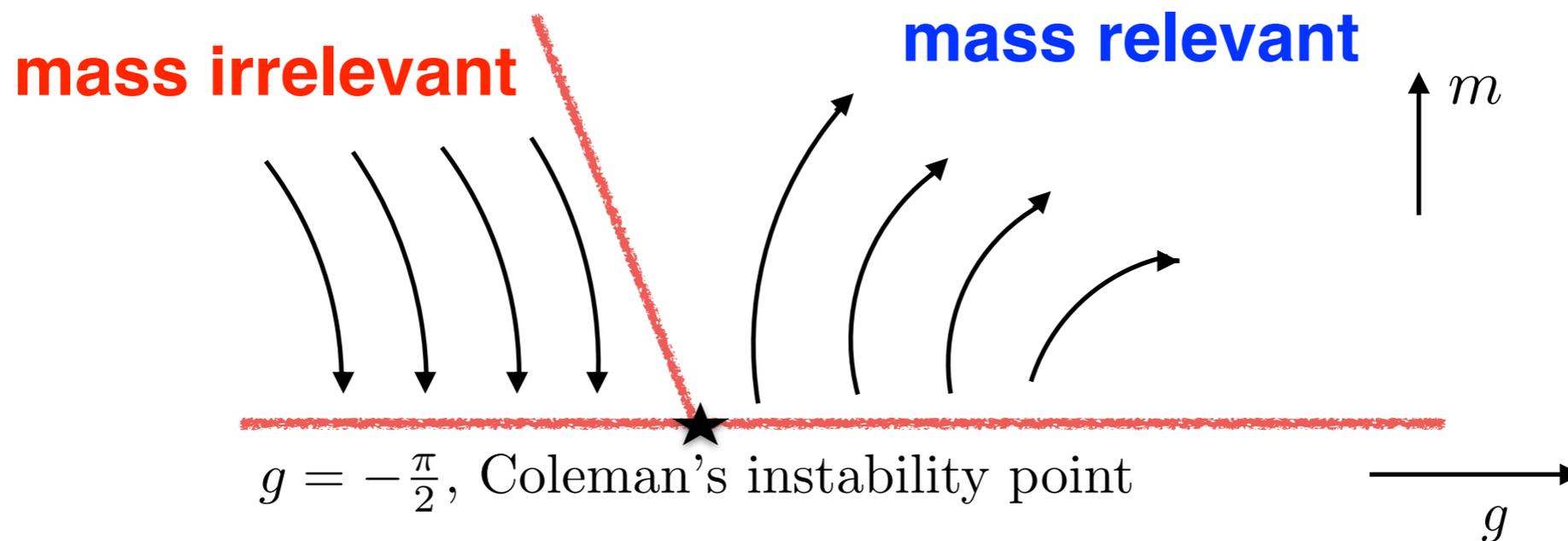
Works in the zero-charge sector

RG flows of the Thirring model

$$\beta_g \equiv \mu \frac{dg}{d\mu} = -64\pi \frac{m^2}{\Lambda^2},$$

$$\beta_m \equiv \mu \frac{dm}{d\mu} = \frac{-2(g + \frac{\pi}{2})}{g + \pi} m - \frac{256\pi^3}{(g + \pi)^2 \Lambda^2} m^3$$

★ Massless Thirring model is a conformal field theory



Operator formalism and the Hamiltonian

- Operator formalism of the Thirring model Hamiltonian

C.R. Hagen, 1967

$$H_{\text{Th}} = \int dx \left[-i\bar{\psi}\gamma^1\partial_1\psi + m_0\bar{\psi}\psi + \frac{g}{4}(\bar{\psi}\gamma^0\psi)^2 - \frac{g}{4}\left(1 + \frac{2g}{\pi}\right)^{-1}(\bar{\psi}\gamma^1\psi)^2 \right]$$

- Staggering, J-W transformation ($S_j^\pm = S_j^x \pm iS_j^y$):

J. Kogut and L. Susskind, 1975; A. Luther, 1976

$$\bar{H}_{XXZ} = \nu(g) \left[-\frac{1}{2} \sum_n^{N-2} (S_n^+ S_{n+1}^- + S_{n+1}^+ S_n^-) + a\tilde{m}_0 \sum_n^{N-1} (-1)^n \left(S_n^z + \frac{1}{2} \right) + \Delta(g) \sum_n^{N-1} \left(S_n^z + \frac{1}{2} \right) \left(S_{n+1}^z + \frac{1}{2} \right) \right]$$

$$\nu(g) = \frac{2\gamma}{\pi \sin(\gamma)}, \quad \tilde{m}_0 = \frac{m_0}{\nu(g)}, \quad \Delta(g) = \cos(\gamma), \quad \text{with } \gamma = \frac{\pi - g}{2}$$

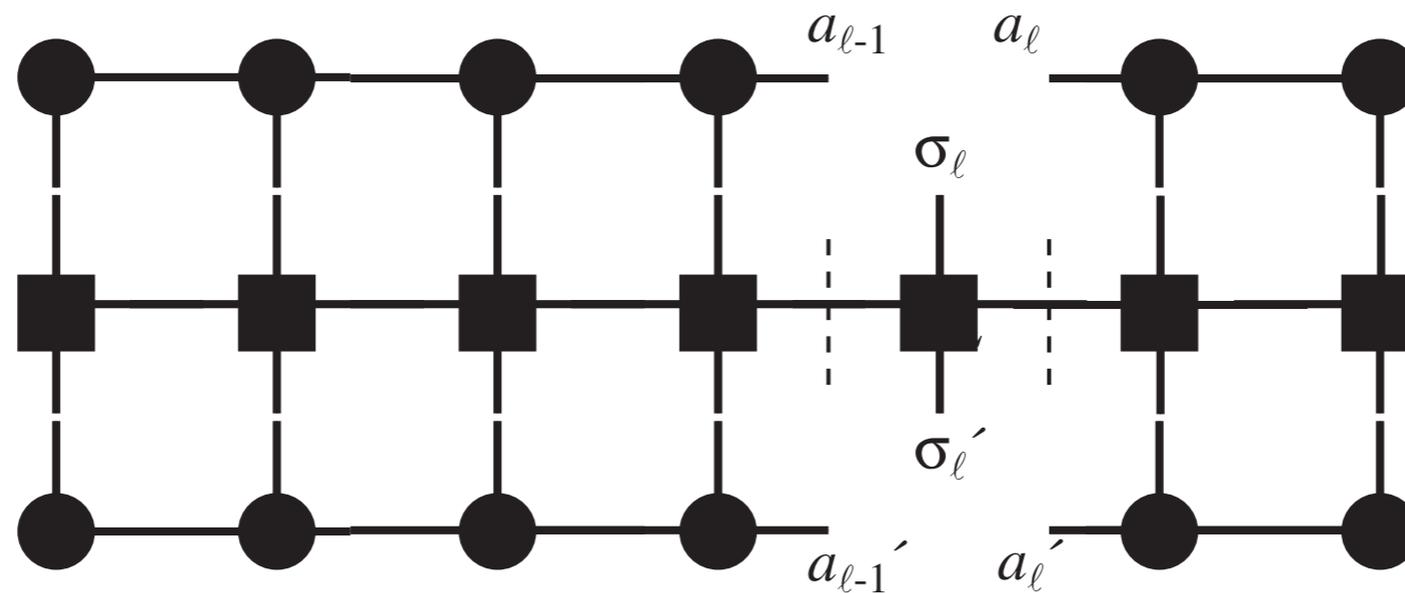
$$\bar{H}_{\text{sim}} = \frac{\bar{H}_{XXZ}}{\nu(g)} + \lambda \left(\sum_{n=0}^{N-1} S_n^z - S_{\text{target}} \right)^2$$

projected to a sector of total spin

JW-trans of the total fermion number,
Bosonise to topological index in the SG theory.

Practice of finite MPS

One step in a sweep of finite-size DMRG



- ★ Open BC
- ★ Random tensors for the smallest bond dim

Simulation details for the phase structure

- Matrix product operator for the Hamiltonian (bulk)

$$W^{[n]} = \begin{pmatrix} 1_{2 \times 2} & -\frac{1}{2}S^+ & -\frac{1}{2}S^- & 2\lambda S^z & \Delta S^z & \beta_n S^z + \alpha 1_{2 \times 2} \\ 0 & 0 & 0 & 0 & 0 & S^- \\ 0 & 0 & 0 & 0 & 0 & S^+ \\ 0 & 0 & 0 & 1 & 0 & S^z \\ 0 & 0 & 0 & 0 & 0 & S^z \\ 0 & 0 & 0 & 0 & 0 & 1_{2 \times 2} \end{pmatrix}$$

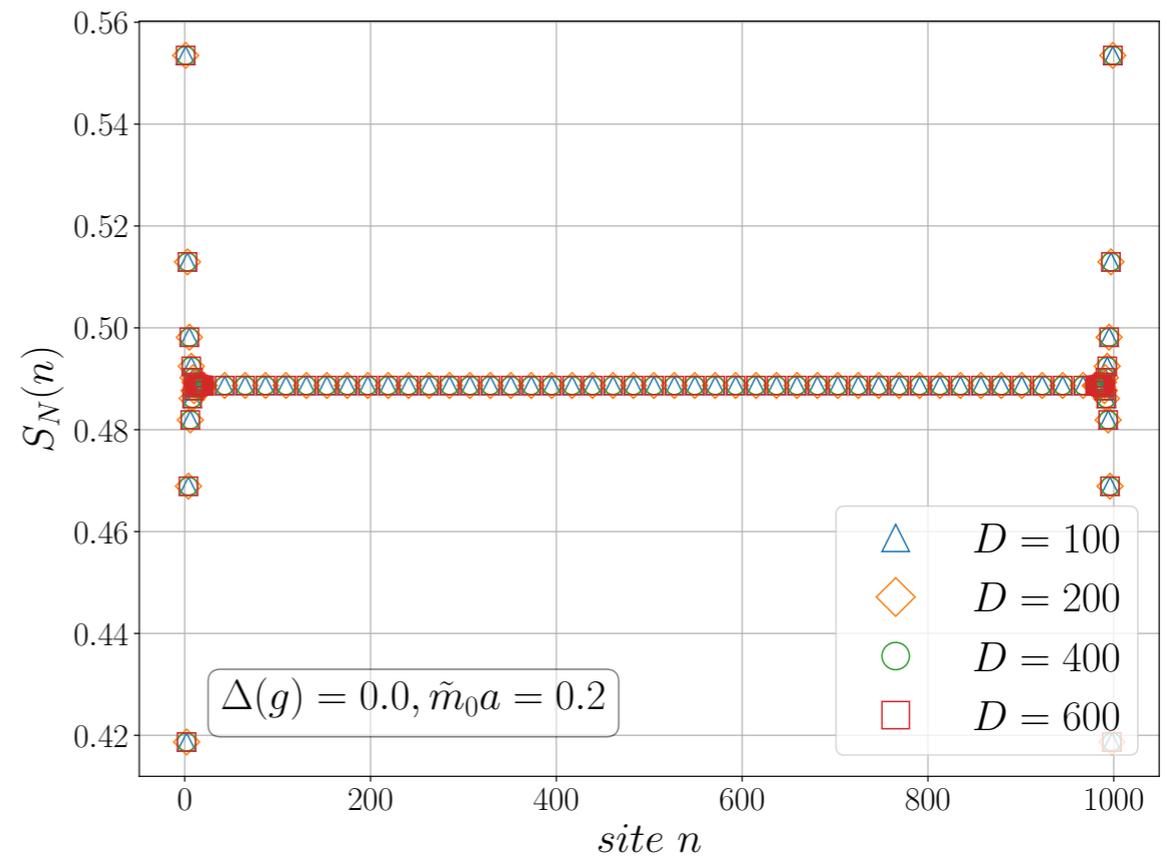
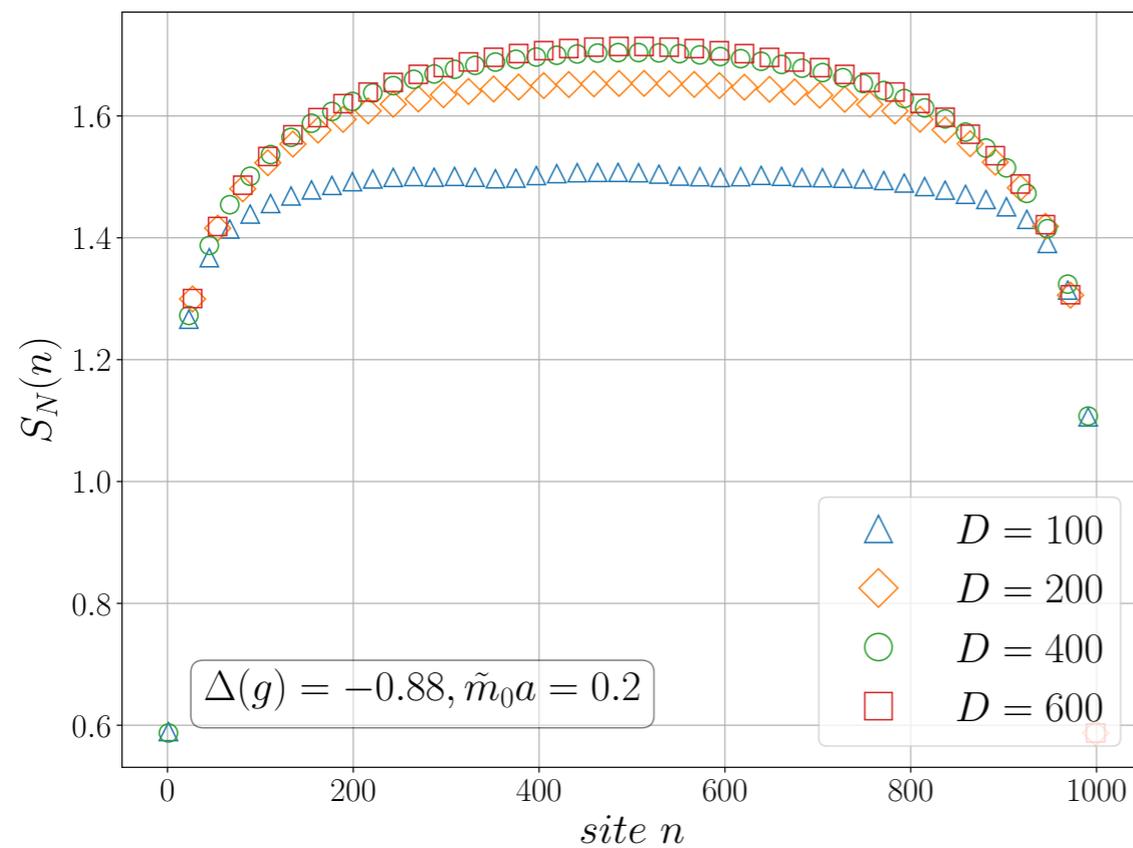
$$\beta_n = \Delta + (-1)^n \tilde{m}_0 a - 2\lambda S_{\text{target}}, \quad \alpha = \lambda \left(\frac{1}{4} + \frac{S_{\text{target}}^2}{N} \right) + \frac{\Delta}{4}$$

- Simulation parameters
 - ★ Twenty values of $\Delta(g)$, ranging from -0.9 to 1.0
 - ★ Fourteen values of $\tilde{m}_0 a$, ranging from 0 to 0.4
 - ★ Bond dimension $D = 50, 100, 200, 300, 400, 500, 600$
 - ★ System size $N = 400, 600, 800, 1000$

Entanglement entropy (Lattice 2018)

Calabrese-Cardy scaling and the central charge

$$S_N(n) = \frac{c}{6} \ln \left[\frac{N}{\pi} \sin \left(\frac{\pi n}{N} \right) \right] + k$$



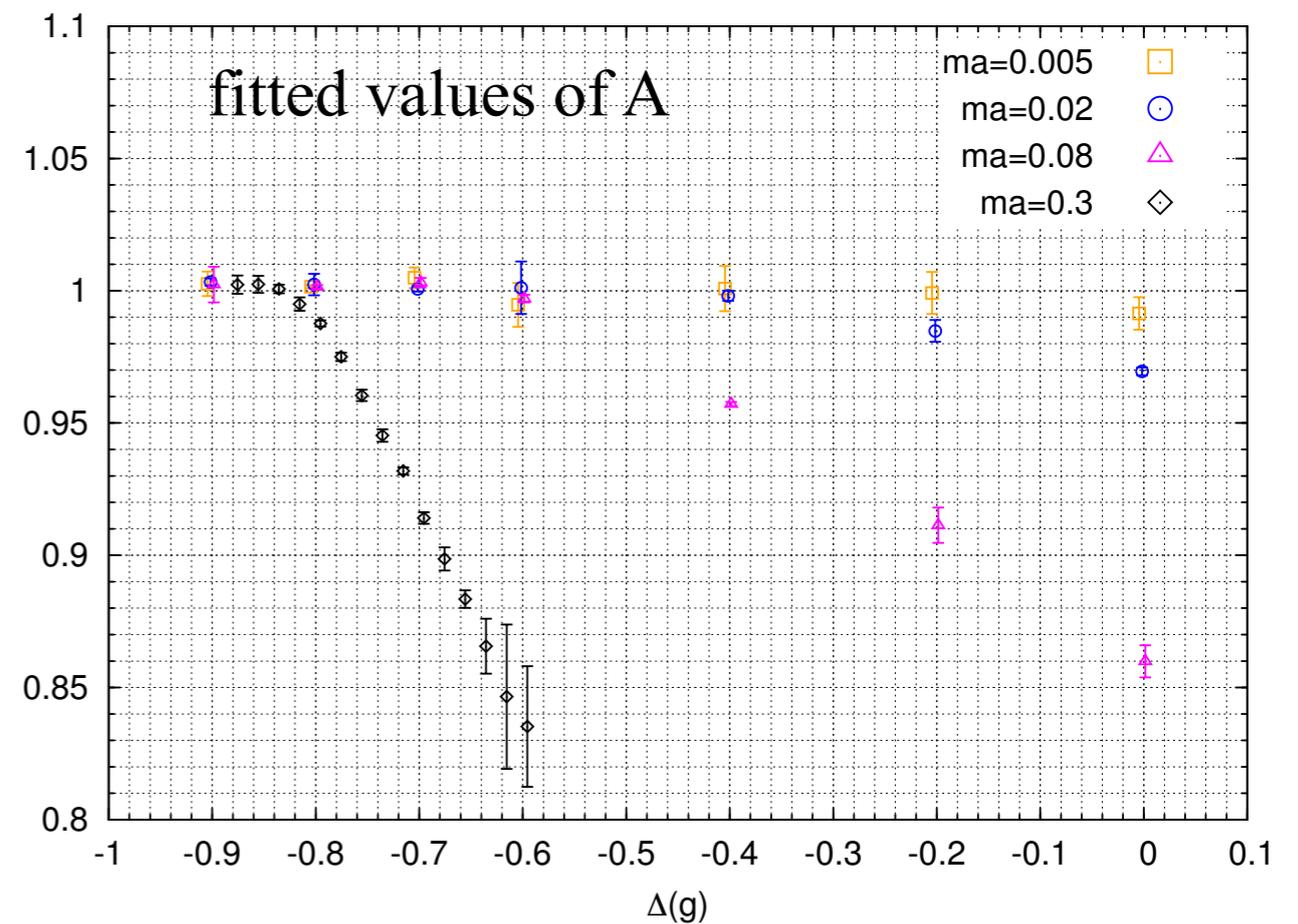
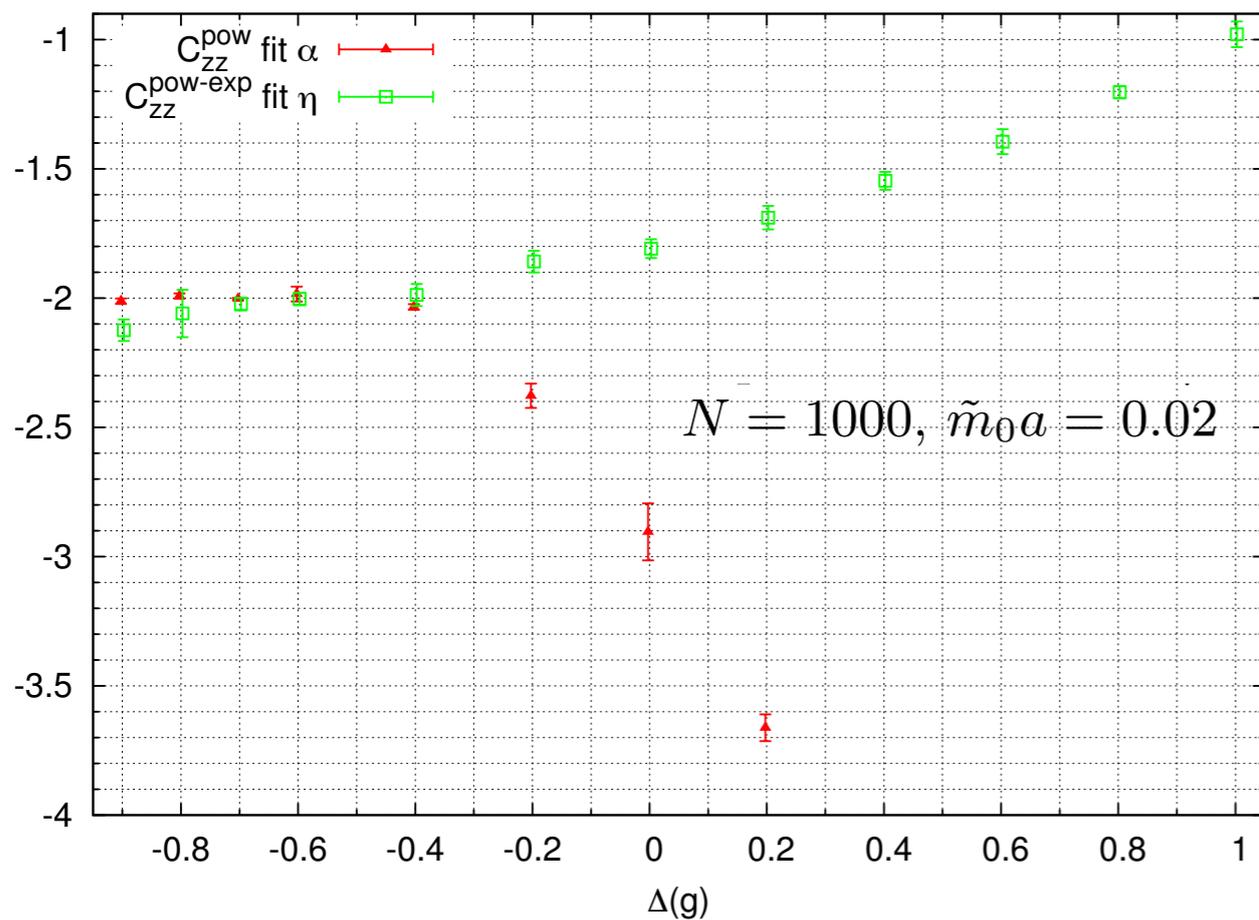
- ★ Scaling observed at $\Delta(g) \lesssim -0.7$ for $\tilde{m}_0 a \neq 0$, and for all values of $\Delta(g)$ at $\tilde{m}_0 a = 0$
- ★ In the critical phase, $c = 1$

Density-density correlators

$$C_{zz}(x) = \langle \bar{\psi}\psi(x_0 + x)\bar{\psi}\psi(x_0) \rangle_{\text{conn}} \xrightarrow{\text{JW trans}} \frac{1}{N_x} \sum_n S^z(n)S^z(n+x) - \frac{1}{N_0} \sum_n S^z(n) \sum_n S^z(n+1)$$

try fitting to

$$C_{zz}^{\text{pow}}(x) = \beta x^\alpha \text{ and } C_{zz}^{\text{pow-exp}}(x) = Bx^\eta A^x$$



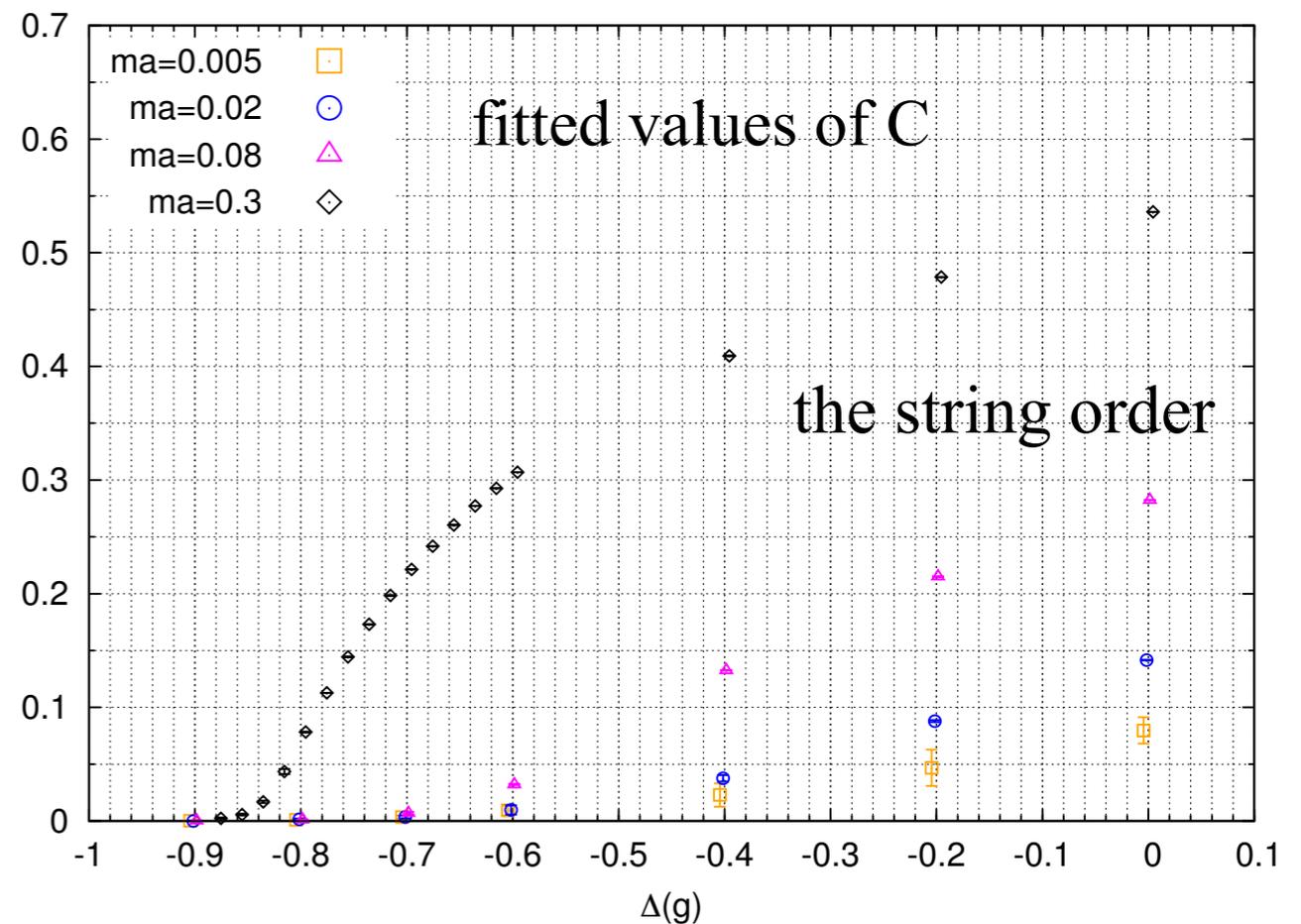
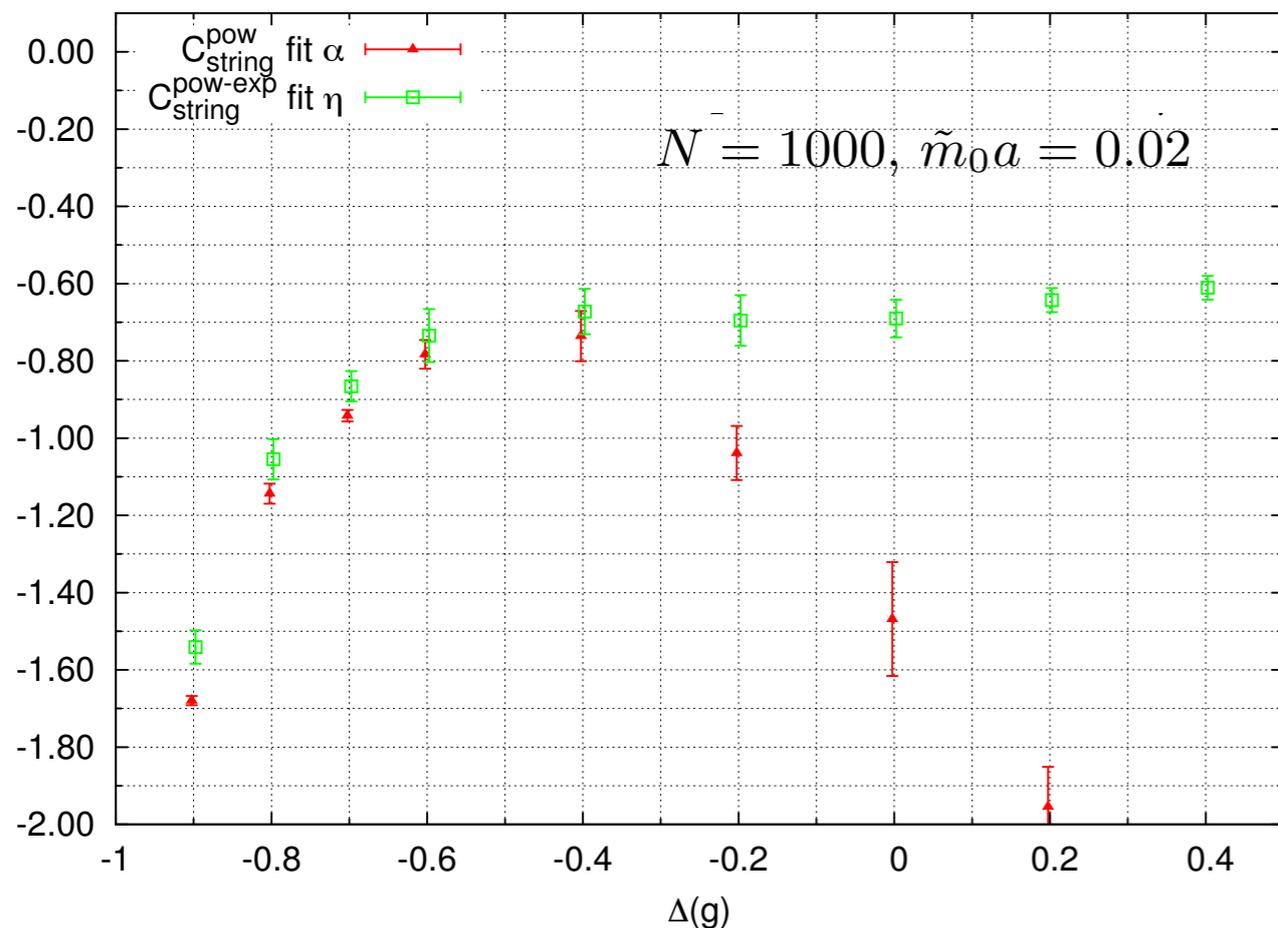
★ Evidence for a critical phase

Soliton (string) correlators

$$C_{\text{string}}(x) = \langle \psi^\dagger(x_0 + x)\psi(x_0) \rangle \xrightarrow{\text{JW trans}} \frac{1}{N_x} \sum_n S^+(n)S^z(n+1) \cdots S^z(n+x-1)S^-(n+x)$$

try fitting to

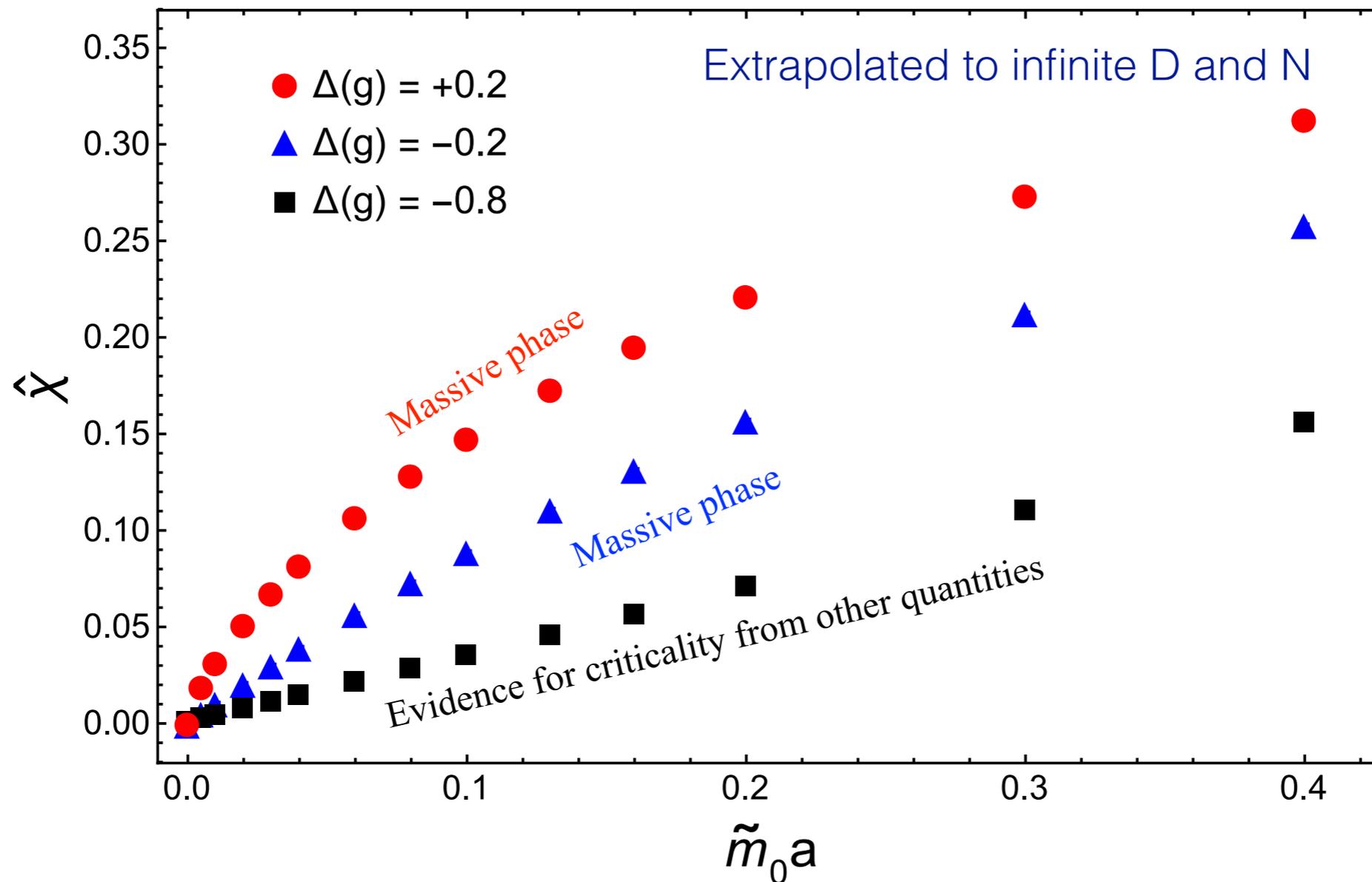
$$C_{\text{string}}^{\text{pow}}(x) = \beta x^\alpha + C \quad \text{and} \quad C_{\text{string}}^{\text{pow-exp}}(x) = Bx^\eta A^x + C$$



★ Similar behaviour in A. Evidence for a critical phase

Chiral condensate

$$\hat{\chi} = a |\langle \bar{\psi} \psi \rangle| = \frac{1}{N} \left| \sum_n (-1)^n S_n^z \right|$$



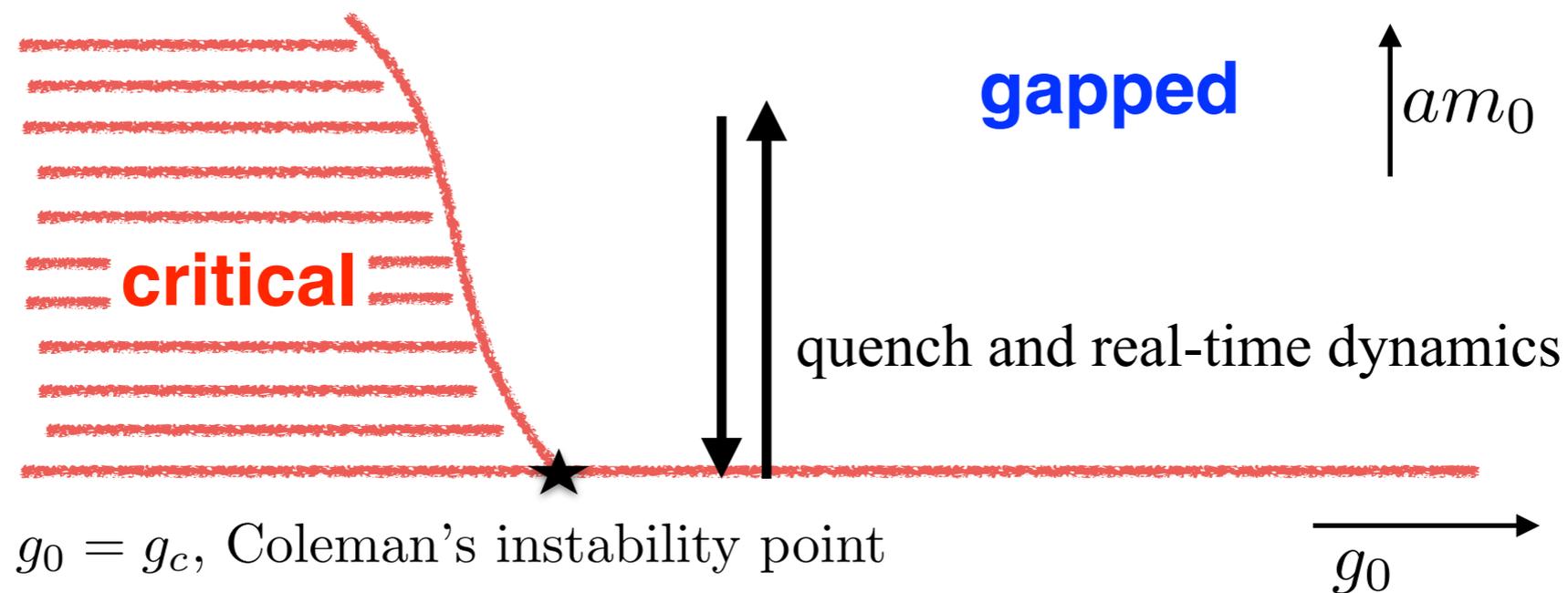
★ Chiral condensate is not an order parameter

Phase structure of the Thirring model

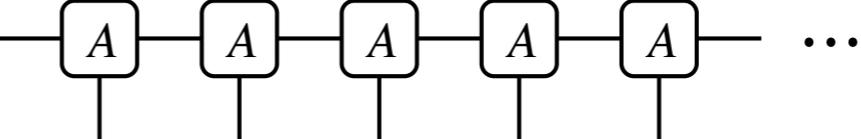
$$\beta_g \equiv \mu \frac{dg}{d\mu} = -64\pi \frac{m^2}{\Lambda^2},$$

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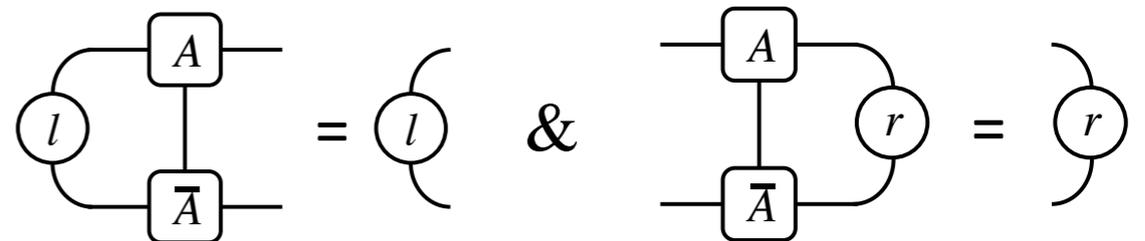
Massless Thirring model is a conformal field theory



Uniform MPS and real-time evolution

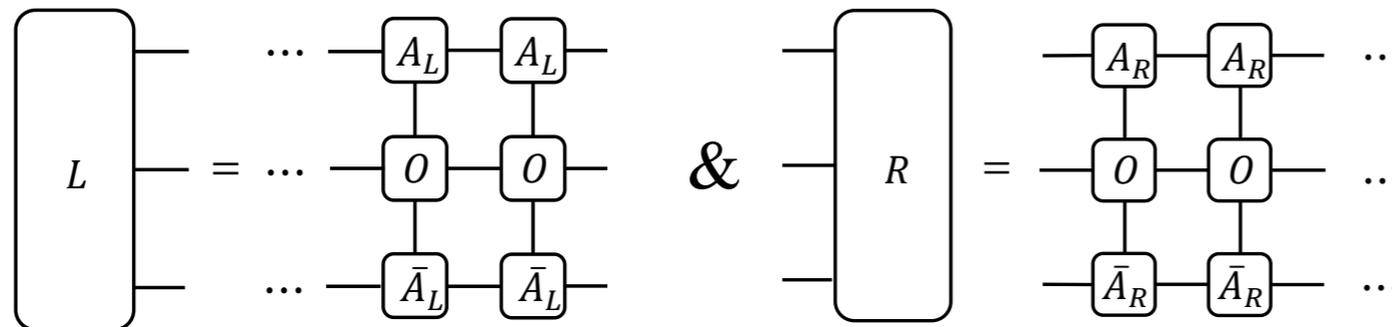
★ Translational invariance in MPS \dots  \dots

★ Finding the infinite BC for amplitudes
(largest eigenvalue normalised to be 1)



H.N. Phien, G. Vidal and I.P. McCulloch, Phys. Rev. B86, 2012

★ Similar (more complicated) procedure in the variation search for the ground state



...V. Zauner-Stauber *et al*, Phys. Rev. B97, 2018

★ Real-time evolution *via* time-dependent variational principle

➔ Key: projection to MPS in $i \frac{d}{dt} |\Psi(A(t))\rangle = P_{|\Psi(A)\rangle} \hat{H} |\Psi(A(t))\rangle$

J. Haegeman *et al*, Phys. Rev. Lett.107, 2011

Dynamical quantum phase transition

- ★ “Quenching” : Sudden change of coupling strength in time evolution

$$H(g_1)|0_1\rangle = E_0^{(1)}|0_1\rangle \quad \text{and} \quad |\psi(t)\rangle = e^{-iH(g_2)t}|0_1\rangle$$

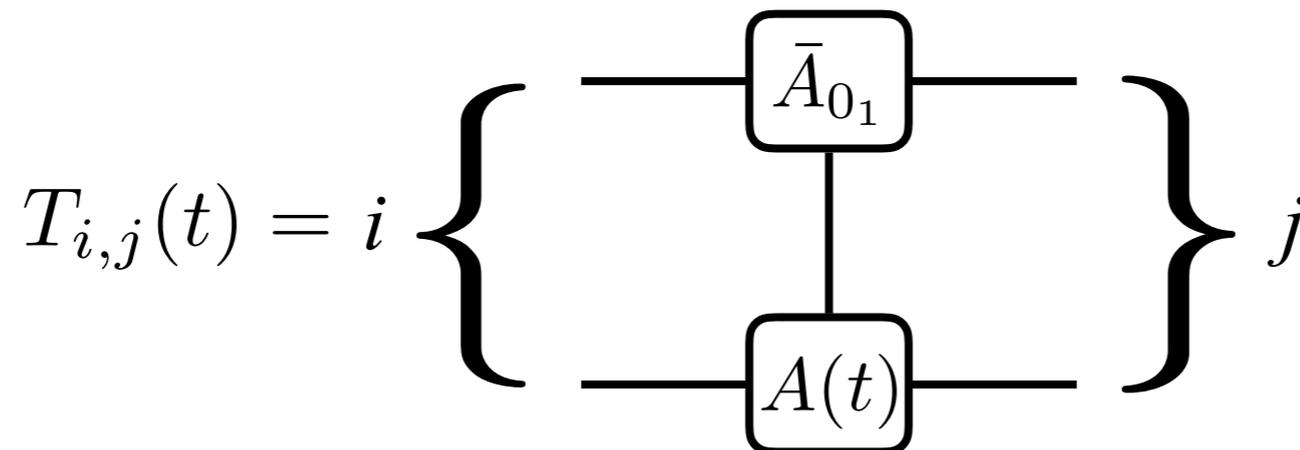
- ★ Questions: Any singular behaviour? Related to equilibrium PT?

- ★ The Loschmidt echo and the return rate

$$L(t) = \langle 0_1 | e^{-iH(g_2)t} | 0_1 \rangle \quad \& \quad g(t) = - \lim_{N \rightarrow \infty} \frac{1}{N} \ln L(t)$$

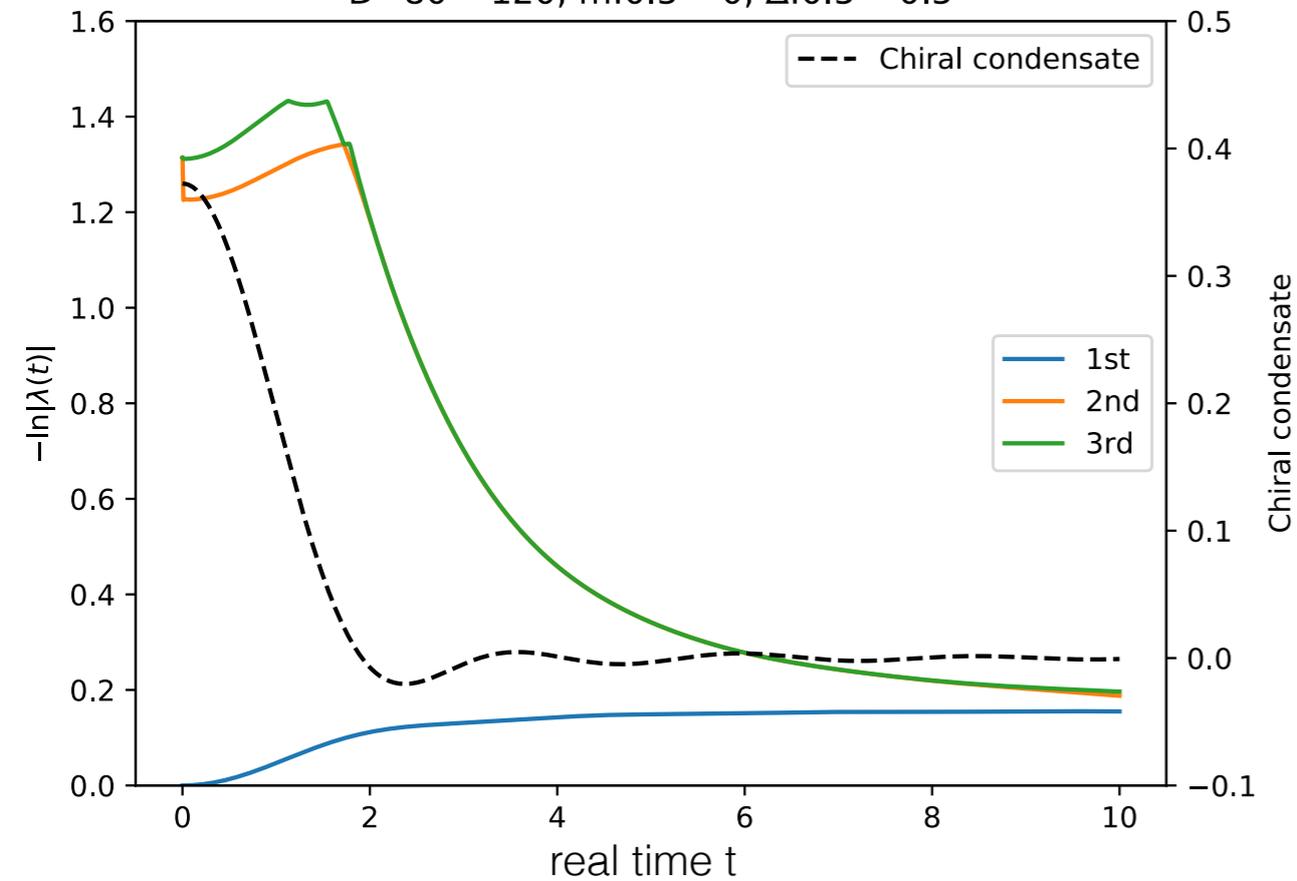
➔ *c.f.*, the partition function and the free energy

➔ In uMPS computed from the largest eigenvalue of the “transfer matrix”



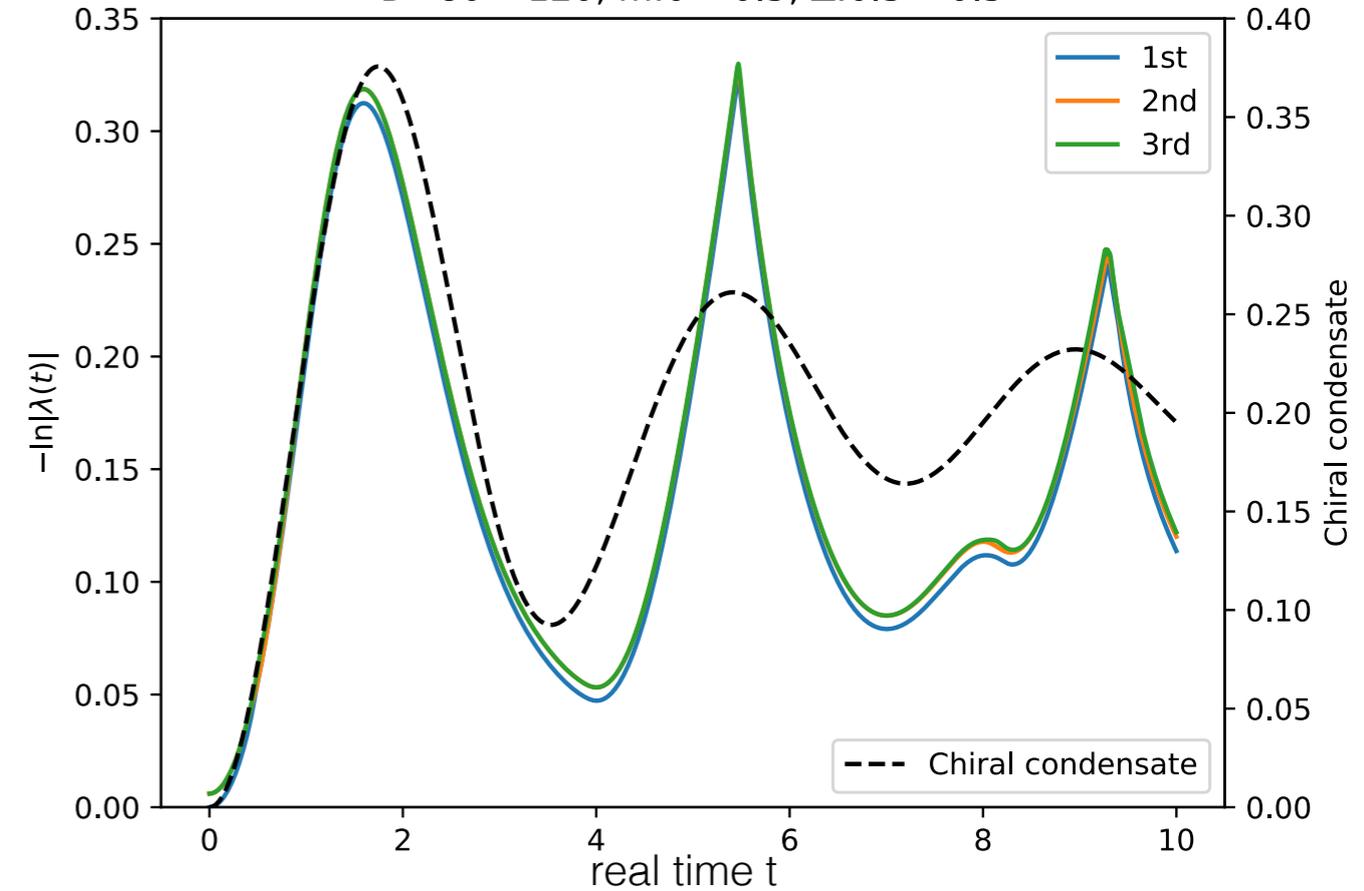
Observing DQPT

Spectrum of transfer matrix and chiral condensate
D=80 → 120, m:0.5 → 0, Δ:0.5 → 0.5



massive \longrightarrow critical

Spectrum of transfer matrix and chiral condensate
D=80 → 120, m:0 → 0.5, Δ:0.5 → 0.5

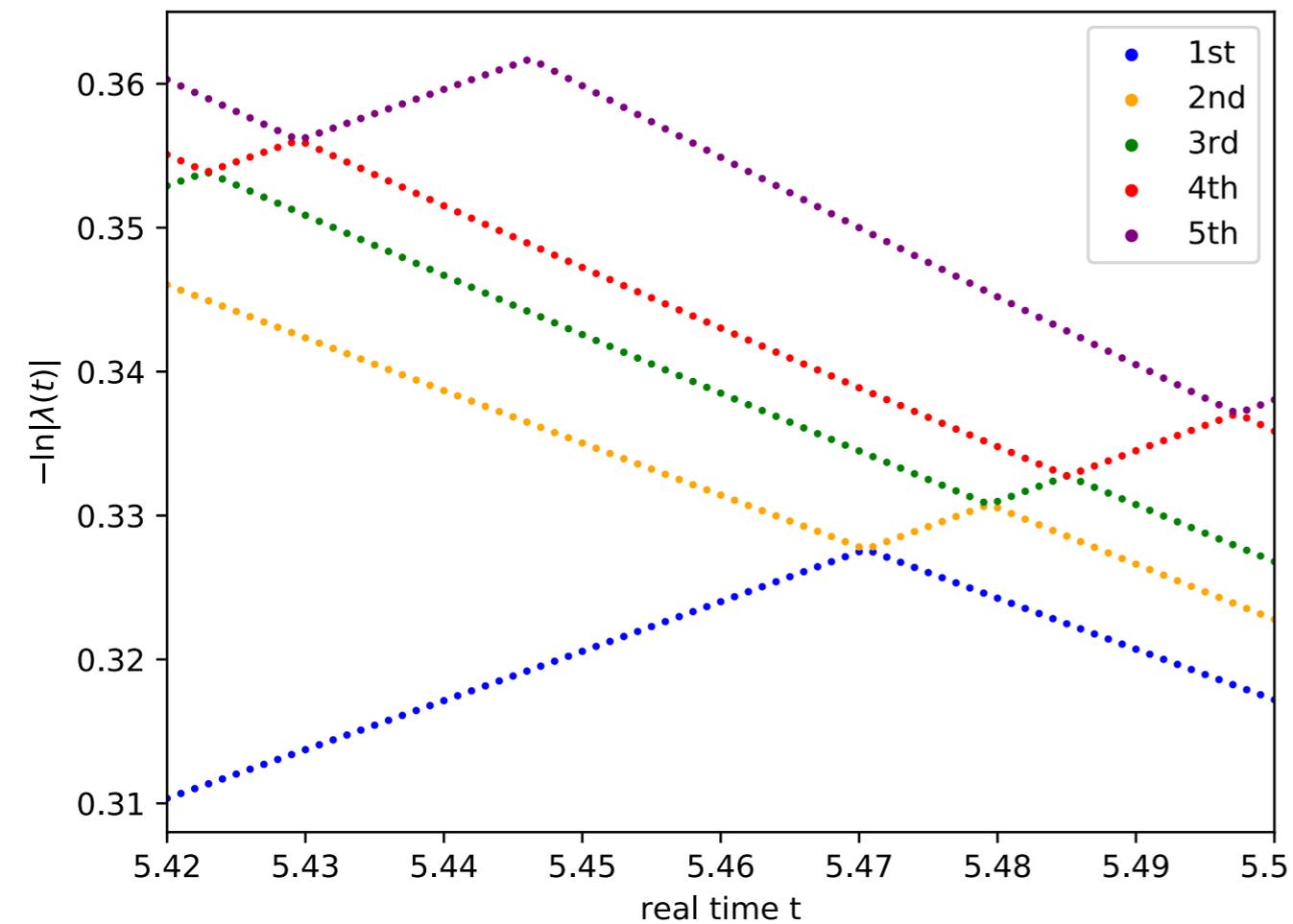


critical \longrightarrow massive

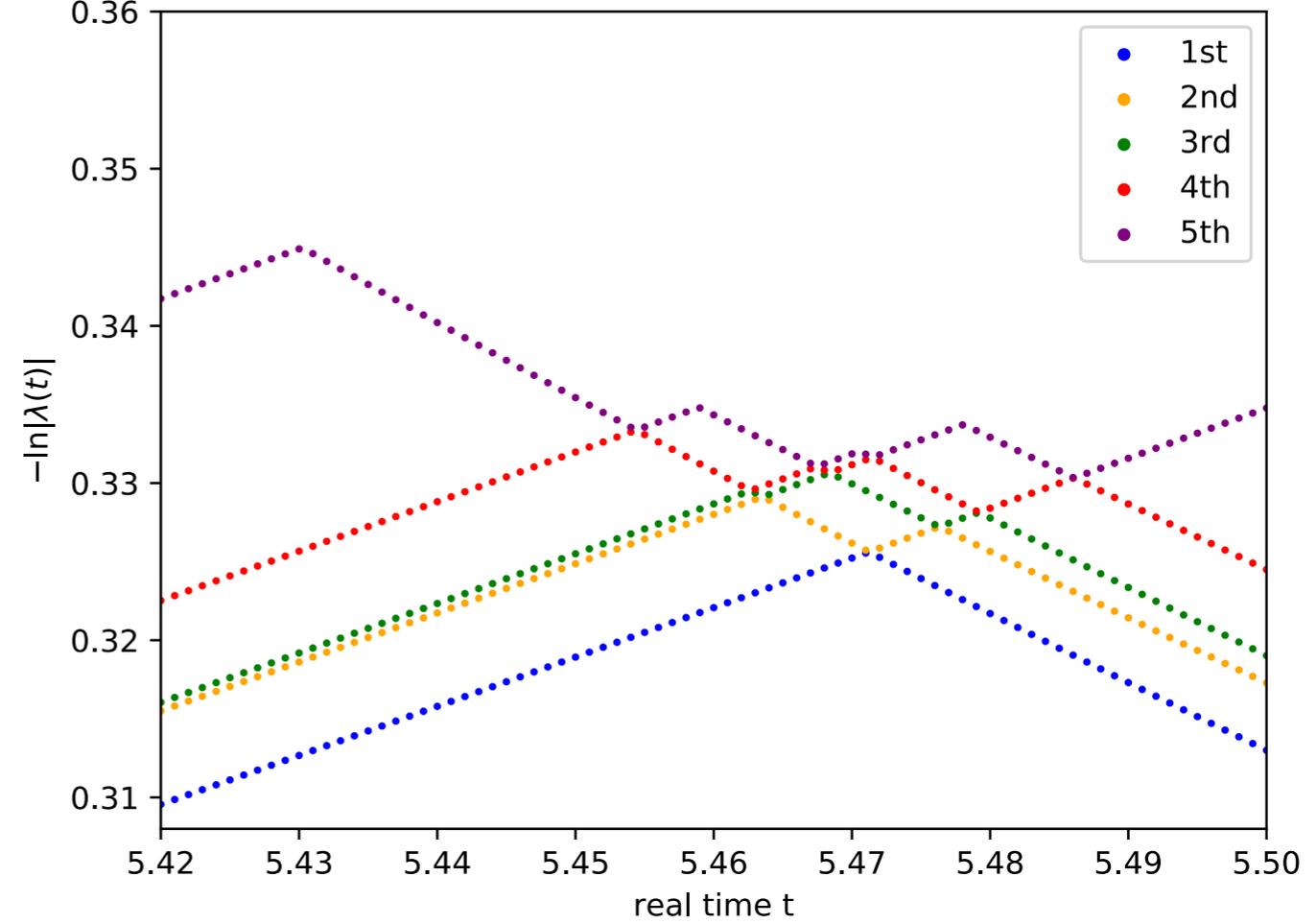
★ DQPT is a “one-way” transition...

DQPT and eigenvalue crossing

Spectrum of transfer matrix
D30 \rightarrow D45, $m:0 \rightarrow 0.5$, $\Delta:0.5 \rightarrow 0.5$



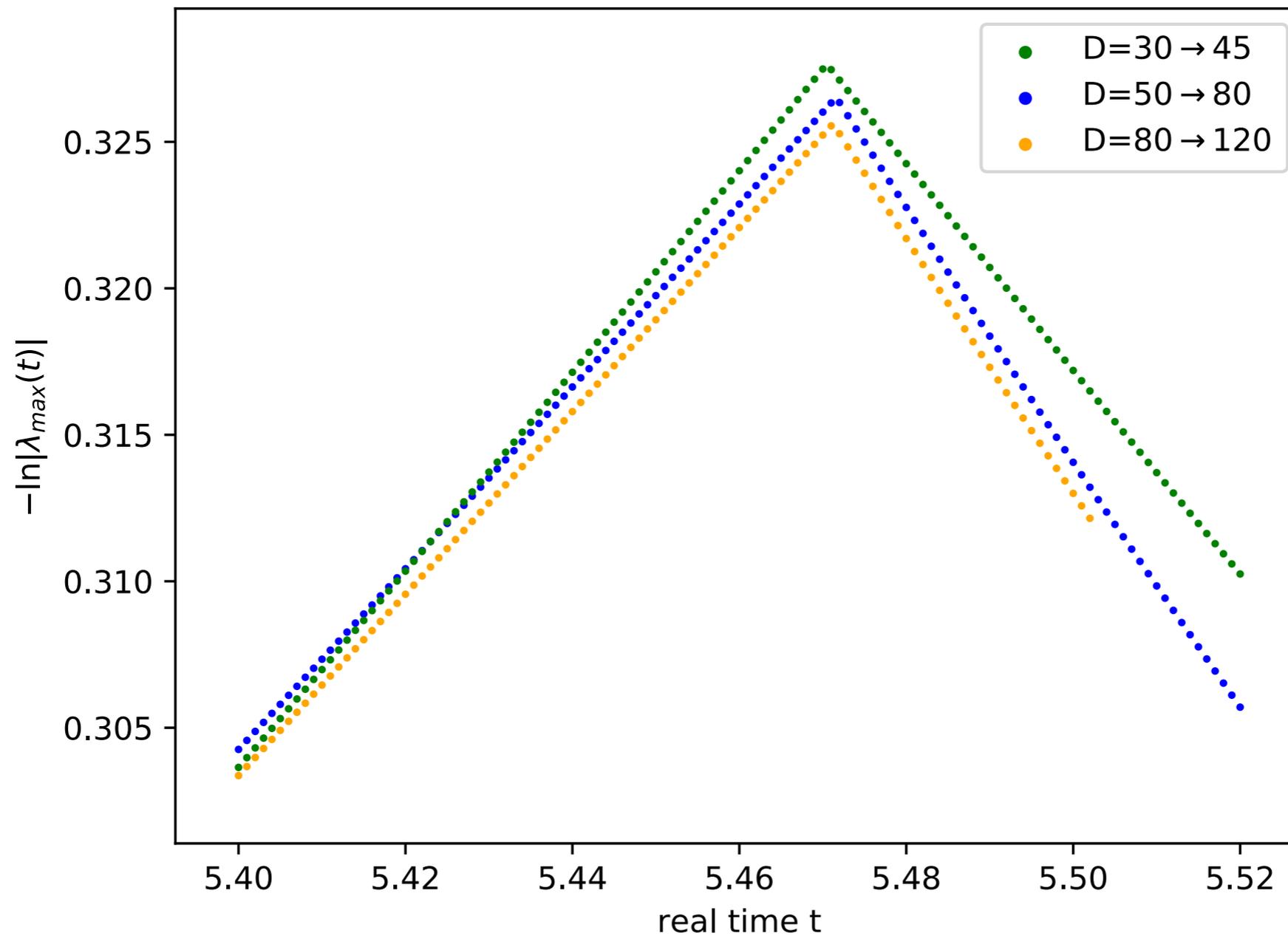
Spectrum of transfer matrix
D80 \rightarrow 120, $m:0 \rightarrow 0.5$, $\Delta:0.5 \rightarrow 0.5$



★ D-dependence in the crossing points

“Universality” in DQPT?

Return rate function, $m:0 \rightarrow 0.5$, $\Delta:0.5 \rightarrow 0.5$



Conclusion and outlook

- Concluding results for phase structure
 - ★ KT-type transition observed
- Exploratory results for real-time dynamics
 - ★ DQPT observed
 - ★ Relation to equilibrium KT phase transition?