

$\mathcal{N} = 1$ Supersymmetric SU(3) Gauge Theory with a Twist

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in collaboration with

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Standard model of particle physics

- describes successfully what we observe at $\lesssim \mathcal{O}(10 \text{ TeV})$

Open questions

- Higgs mass
- dark matter
- unification of forces
- ...

⇒ Not complete

⇒ More fundamental theory?

Possible solution

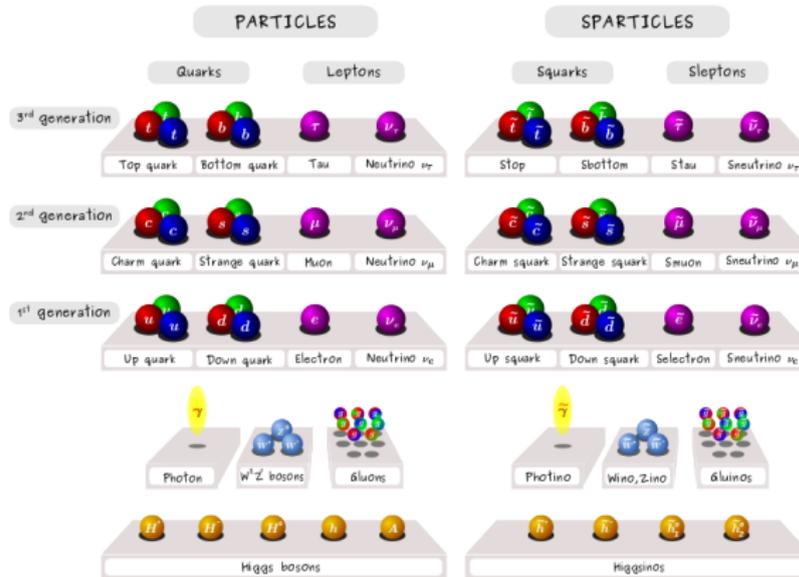
Introduce supersymmetry

Table of content

- 1 $\mathcal{N} = 1$ Super-Yang-Mills Theory (recap)
- 2 SYM on the Lattice (with a Twist)
- 3 Analytical Investigations
- 4 Numerical Investigations

Minimal Supersymmetric Standard Model

- contains a super partner for every particle



[4] <https://web2.ph.utexas.edu/~coker2/index.files/supersymmetry.htm>

$\overline{\mathcal{N}} = 1$ Super-Yang-Mills Theory

- contains a super partner for every particle

PARTICLES

SPARTICLES



⇒ important building block of Super-QCD

[4] <https://web2.ph.utexas.edu/~coker2/index.files/supersymmetry.htm>

$\overline{\mathcal{N}} = 1$ Super-Yang-Mills Theory

Fields

- Gauge boson (gluon) $A_\mu(x)$ in the adjoint representation
- Super partner (gluino) $\lambda(x)$ is Majorana fermion in the adjoint representation

On-shell Lagrange density

$$\mathcal{L}_{\text{SYM}} = \text{tr} \left(-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{i}{2} \bar{\lambda} \not{D} \lambda \right)$$

Supersymmetry: Relation between fermionic matter particles and bosonic force particles

$$\delta_\epsilon A_\mu = i \bar{\epsilon} \gamma_\mu \lambda, \quad \delta_\epsilon \lambda = i \Sigma_{\mu\nu} F^{\mu\nu} \epsilon$$

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On-shell Lagrange density

$$\mathcal{L}_{\text{SYM}} = \text{tr} \left(-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{i}{2} \bar{\lambda} \not{D} \lambda - \frac{m_g}{2} \bar{\lambda} \lambda \right)$$

Supersymmetry: Relation between fermionic matter particles and bosonic force particles

$$\delta_\epsilon A_\mu = i \bar{\epsilon} \gamma_\mu \lambda, \quad \delta_\epsilon \lambda = i \Sigma_{\mu\nu} F^{\mu\nu} \epsilon$$

Softly broken by **gluino mass term**

Symmetries

Chiral (R) symmetry breaking in $SU(3)$ SYM theory

- Global chiral $U(1)_A$ symmetry: $\lambda \mapsto e^{i\alpha\gamma_5} \lambda$
- Due to anomaly only \mathbb{Z}_6 remnant symmetry

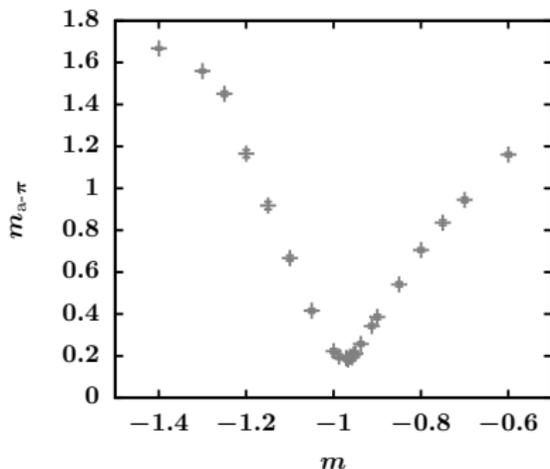
$$\lambda \mapsto e^{i\frac{2\pi n}{6}\gamma_5} \lambda \quad \text{with } n \in \{1, \dots, 6\}$$

- Spontaneously broken to \mathbb{Z}_2 symmetry in consequence of gluino condensate $\langle \bar{\lambda}\lambda \rangle \neq 0 \rightarrow 3$ different vacua

The Chiral Limit

$$D_W(x, y) = (4 + m) \delta_{x, y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} (\mathbb{1} - \gamma_\mu) \mathcal{V}_\mu(x) \delta_{x+\hat{\mu}, y}$$

with adjoint representation $[\mathcal{V}_\mu(x)]_{ab} = 2 \operatorname{tr} [\mathcal{U}_\mu^\dagger(x) T_a \mathcal{U}_\mu(x) T_b]$



Dirac Operator

$$D_W(x, y) = (4 + m) \delta_{x, y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} (\mathbb{1} - \gamma_\mu) \mathcal{V}_\mu(x) \delta_{x+\hat{\mu}, y}$$

Feature of SYM

- Particular directions of \mathbb{Z}_6 symmetry are favored by gluino condensate

Dirac Operator with Twisted Mass

$$D_W^{\text{mtw}}(x, y) = (4 + m + i\mu\gamma_5)\delta_{x,y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} (\mathbb{1} - \gamma_\mu)\mathcal{V}_\mu(x) \delta_{x+\hat{\mu},y}$$

Feature of SYM

- Particular directions of \mathbb{Z}_6 symmetry are favored by gluino condensate
- Deform lattice action by adding **parity-breaking mass** μ resembling a twisted mass
- m breaks chiral symmetry explicitly and generates a condensate $\sim \langle \bar{\lambda}\lambda \rangle$
- μ leads to a condensate $\sim \langle \bar{\lambda}\gamma_5\lambda \rangle$

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⇒ Possibility to get closer to chiral symmetry and supersymmetry at finite lattice spacing?

Dirac Operator with Twisted Mass

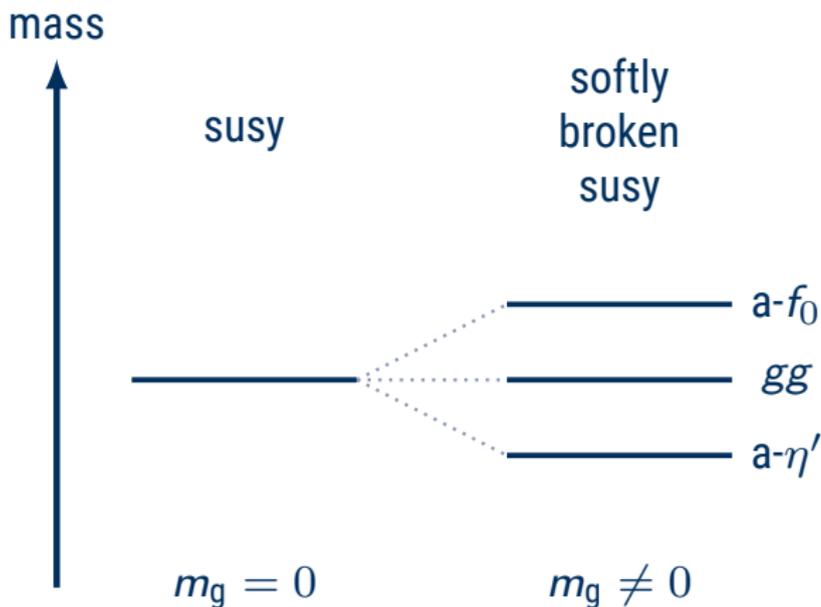
$$\begin{aligned} D_W^{\text{mtw}}(x, y) &= (4 + m + i\mu\gamma_5)\delta_{x,y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} (\mathbb{1} - \gamma_\mu)\mathcal{V}_\mu(x) \delta_{x+\hat{\mu},y} \\ &= (4 + M e^{i\alpha\gamma_5})\delta_{x,y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} (\mathbb{1} - \gamma_\mu)\mathcal{V}_\mu(x) \delta_{x+\hat{\mu},y} \end{aligned}$$

Feature of SYM

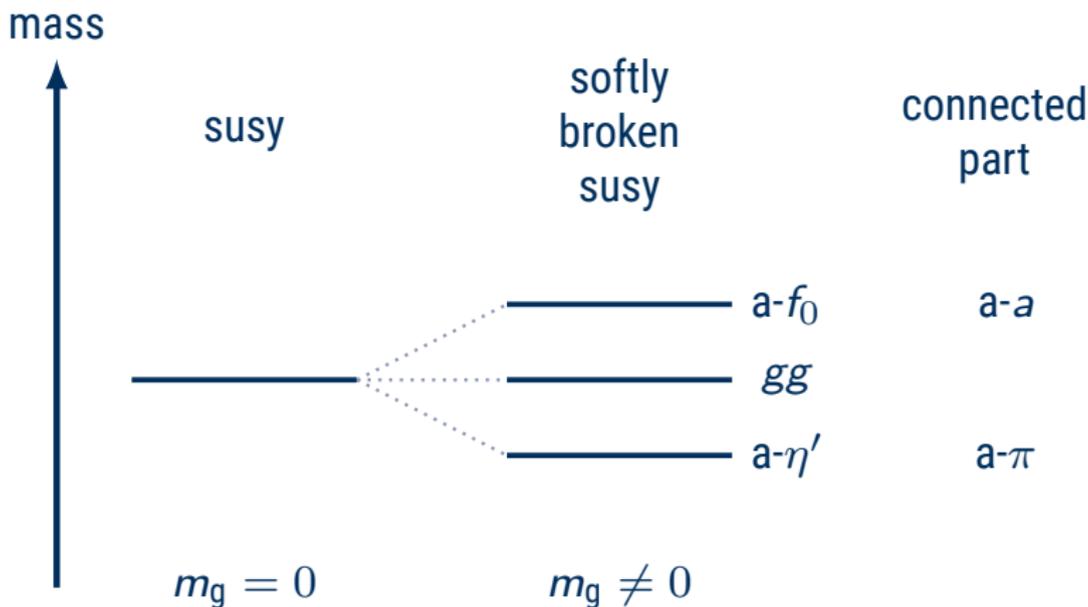
- Particular directions of \mathbb{Z}_6 symmetry are favored by gluino condensate
- Deform lattice action by adding **parity-breaking mass μ** resembling a twisted mass
- $m = M \cos(\alpha)$
- $\mu = M \sin(\alpha)$

⇒ Possibility to get closer to chiral symmetry and supersymmetry at finite lattice spacing?

Veneziano-Yankielowicz Multiplet



Veneziano-Yankielowicz Multiplet



Dirac Operator with Twisted Mass

Numerical Investigations

Spectroscopy

- for different angles $\alpha = \arctan(\mu/m)$
- of connected (unphysical) mesons with good signal-to-noise-ratio
- of full (physical) mesons including disconnected contributions
- of gluino-gluon

Dirac Operator with Twisted Mass

Numerical Investigations

Spectroscopy

- for different angles $\alpha = \arctan(\mu/m)$
- of connected (unphysical) mesons with good signal-to-noise-ratio
- of full (physical) mesons including disconnected contributions
- of gluino-gluon

Analytical Investigations

Calculation of

- fermionic expectation values of lattice interpolators
- eigenvalues of the free Dirac operator

Fermionic Expectation Values

$$\langle O \rangle = \langle \langle O \rangle_F \rangle_G = \frac{1}{Z} \int \mathcal{D}\mathcal{U} e^{-S_G[\mathcal{U}]} \mathcal{D}\lambda e^{-S_F[\lambda, \mathcal{U}]} O[\lambda, \mathcal{U}]$$

Twist

$$\lambda \mapsto e^{i\alpha\gamma_5} \lambda$$

$$\bar{\lambda} \mapsto \bar{\lambda} e^{i\alpha\gamma_5}$$

$$\begin{aligned} \bar{\lambda}\lambda &\mapsto \bar{\lambda} e^{2i\alpha\gamma_5} \lambda = \bar{\lambda}\lambda \cos(2\alpha) + i\bar{\lambda}\gamma_5\lambda \sin(2\alpha) \\ \bar{\lambda}\gamma_5\lambda &\mapsto \bar{\lambda}\gamma_5 e^{2i\alpha\gamma_5} \lambda = \bar{\lambda}\gamma_5\lambda \cos(2\alpha) + i\bar{\lambda}\lambda \sin(2\alpha) \end{aligned}$$

Fermionic Expectation Values

$$\langle O \rangle = \langle \langle O \rangle_F \rangle_G = \frac{1}{Z} \int \mathcal{D}\mathcal{U} e^{-S_G[\mathcal{U}]} \mathcal{D}\lambda e^{-S_F[\lambda, \mathcal{U}]} O[\lambda, \mathcal{U}]$$

Twist

$$\lambda \mapsto e^{i\alpha\gamma_5} \lambda$$

$$\bar{\lambda} \mapsto \bar{\lambda} e^{i\alpha\gamma_5}$$

$$\begin{pmatrix} \bar{\lambda} \lambda \\ i\bar{\lambda}\gamma_5\lambda \end{pmatrix} \mapsto \begin{pmatrix} \cos(2\alpha) & \sin(2\alpha) \\ -\sin(2\alpha) & \cos(2\alpha) \end{pmatrix} \begin{pmatrix} \bar{\lambda} \lambda \\ i\bar{\lambda}\gamma_5\lambda \end{pmatrix}$$

Fermionic Expectation Values

General Observables

$$O_{\text{gen}}(x) = a \bar{\lambda}_x \lambda_x + b \bar{\lambda}_x \gamma_5 \lambda_x$$

Mesonic Observables

Results

Fermionic Expectation Values

General Observables

Mesonic Observables

	a	b	\mapsto	a	b
$O_{a-f_0}(x) = \bar{\lambda}_x \lambda_x$	1	0		$\cos(2\alpha)$	$i \sin(2\alpha)$
$O_{a-\eta'}(x) = \bar{\lambda}_x \gamma_5 \lambda_x$	0	1		$i \sin(2\alpha)$	$\cos(2\alpha)$

Results

Fermionic Expectation Values

General Observables

Mesonic Observables

Results

- $\langle O_{a-f_0}(n) \bar{O}_{a-f_0}(m) \rangle_F(0^\circ) = \langle O_{a-\eta'}(n) \bar{O}_{a-\eta'}(m) \rangle_F(90^\circ)$
 $\langle O_{a-\eta'}(n) \bar{O}_{a-\eta'}(m) \rangle_F(0^\circ) = \langle O_{a-f_0}(n) \bar{O}_{a-f_0}(m) \rangle_F(90^\circ)$
- $\langle O_{a-\eta'}(n) \bar{O}_{a-\eta'}(m) \rangle_F(45^\circ) = \langle O_{a-f_0}(n) \bar{O}_{a-f_0}(m) \rangle_F(45^\circ)$

Eigenvalues

Goals

- Effect of twist angle α
- $\mathcal{O}(a)$ improvement

Dirac Operator with Mass Twist

$$D_W^{\text{tw}}(x, y) = (4 + M e^{i\alpha\gamma_5})\delta_{x,y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} \left(\mathbb{1} - \gamma_\mu \right) \mathcal{V}_\mu(x) \delta_{x+\hat{\mu},y}$$

Eigenvalues

Goals

- Effect of twist angle α
- $\mathcal{O}(a)$ improvement \Rightarrow angle φ

Dirac Operator with Double Twist

$$D_W^{\text{dtw}}(x, y) = (4 e^{i\varphi\gamma_5} + M e^{i\alpha\gamma_5}) \delta_{x,y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} \left(e^{i\varphi\gamma_5} - \gamma_\mu \right) \mathcal{V}_\mu(x) \delta_{x+\hat{\mu},y}$$

- Idea based on G. Bergner, T. Kästner, S. Uhlmann, A. Wipf, Annals Phys. 323 (2009)
- General action found in G. Immirzi, K. Yoshida, Nucl. Phys. B210 (1982)

Eigenvalues of Free Dirac Operator

Fermionic Kernel D

$$\gamma^\mu \partial_\mu + m - \frac{r}{2} \Delta$$

Eigenvalues $\lambda^\dagger \lambda$

$$p^2 + m^2 + amrp^2 + \mathcal{O}(a^2)$$

Eigenvalues of Free Dirac Operator

Fermionic Kernel D

$$\gamma^\mu \partial_\mu + m - \frac{r}{2} \Delta$$

$$\gamma^\mu \partial_\mu + M e^{i\alpha \gamma_5} - \frac{r}{2} \Delta$$

Eigenvalues $\lambda^\dagger \lambda$

$$p^2 + m^2 + amrp^2 + \mathcal{O}(a^2)$$

$$p^2 + M^2 + \cancel{aMrp^2 \cos(\alpha)} + \mathcal{O}(a^2)$$

0 for $\alpha=90^\circ$

Eigenvalues of Free Dirac Operator

Fermionic Kernel D	Eigenvalues $\lambda^\dagger \lambda$
$\gamma^\mu \partial_\mu + m - \frac{r}{2} \Delta$	$p^2 + m^2 + amrp^2 + \mathcal{O}(a^2)$
$\gamma^\mu \partial_\mu + M e^{i\alpha \gamma_5} - \frac{r}{2} \Delta$	$p^2 + M^2 + \cancel{aMrp^2 \cos(\alpha)} + \mathcal{O}(a^2)$ <small>0 for $\alpha=90^\circ$</small>
$\gamma^\mu \partial_\mu + m + \frac{ir}{2} \gamma_5 \Delta$	$p^2 + m^2 + \kappa a^2 + \mathcal{O}(a^4)$

- with $\kappa = -\frac{1}{3} \sum_\mu p_\mu^4 + \frac{r^2}{2} \left(\sum_\mu p_\mu^2 \right)^2$

Eigenvalues of Free Dirac Operator

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$$p^2 + m^2 + amrp^2 + \mathcal{O}(a^2)$$

$$\gamma^\mu \partial_\mu + M e^{i\alpha \gamma_5} - \frac{r}{2} \Delta$$

$$p^2 + M^2 + aMrp^2 \cos(\alpha) + \mathcal{O}(a^2)$$

$\rightarrow 0 \text{ for } \alpha=90^\circ$

$$\gamma^\mu \partial_\mu + m + \frac{ir}{2} \gamma_5 \Delta$$

$$p^2 + m^2 + \kappa a^2 + \mathcal{O}(a^4)$$

$$\gamma^\mu \partial_\mu + M e^{i\alpha \gamma_5} + \frac{r}{2} e^{i\varphi \gamma_5} \Delta$$

$$p^2 + M^2 + aMrp^2 \cos(\alpha - \varphi) + \kappa a^2 + \mathcal{O}(a^4)$$

$\rightarrow 0 \text{ for } \alpha - \varphi = 90^\circ$

- with $\kappa = -\frac{1}{3} \sum_\mu p_\mu^4 + \frac{r^2}{2} \left(\sum_\mu p_\mu^2 \right)^2$

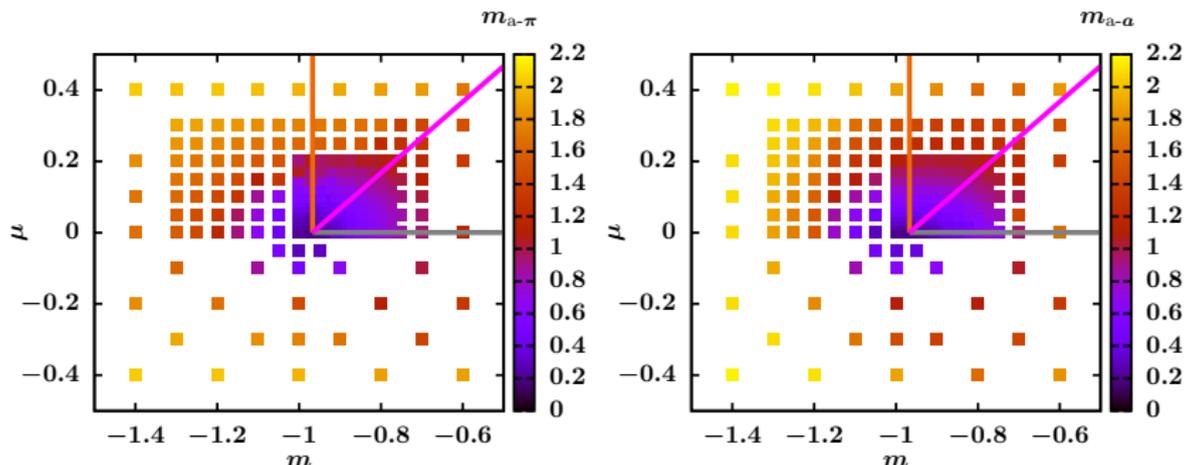
Eigenvalues of Free Dirac Operator

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$\gamma^\mu \partial_\mu + M e^{i\alpha \gamma_5} - \frac{r}{2} \Delta$	$p^2 + M^2 + aMrp^2 \cos(\alpha) + \mathcal{O}(a^2)$ <small>0 for $\alpha=90^\circ$</small>
$\gamma^\mu \partial_\mu + m + \frac{ir}{2} \gamma_5 \Delta$	$p^2 + m^2 + \kappa a^2 + \mathcal{O}(a^4)$
$\gamma^\mu \partial_\mu + M e^{i\alpha \gamma_5} + \frac{r}{2} e^{i\varphi \gamma_5} \Delta$	$p^2 + M^2 + aMrp^2 \cos(\alpha - \varphi) + \kappa a^2 + \mathcal{O}(a^3)$ <small>0 for $\alpha - \varphi = 90^\circ$</small>

- with $\kappa = -\frac{1}{3} \sum_\mu p_\mu^4 + \frac{r^2}{2} \left(\sum_\mu p_\mu^2 \right)^2$
- Mass & Wilson term orthogonal $\Rightarrow \mathcal{O}(a)$ improvement

Connected Correlators

$$D_W^{\text{tw}}(x, y) = (4 + m + i\mu\gamma_5)\delta_{x,y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} (\mathbb{1} - \gamma_\mu)\mathcal{V}_\mu(x) \delta_{x+\mu,y}$$

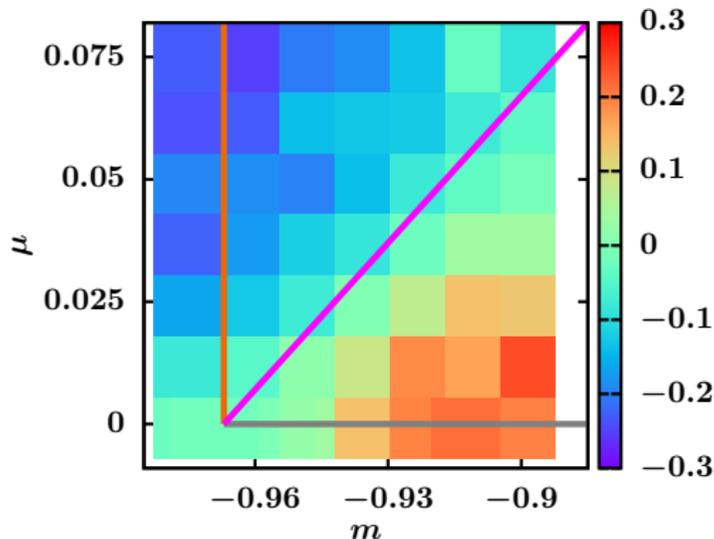


$8^3 \times 16$

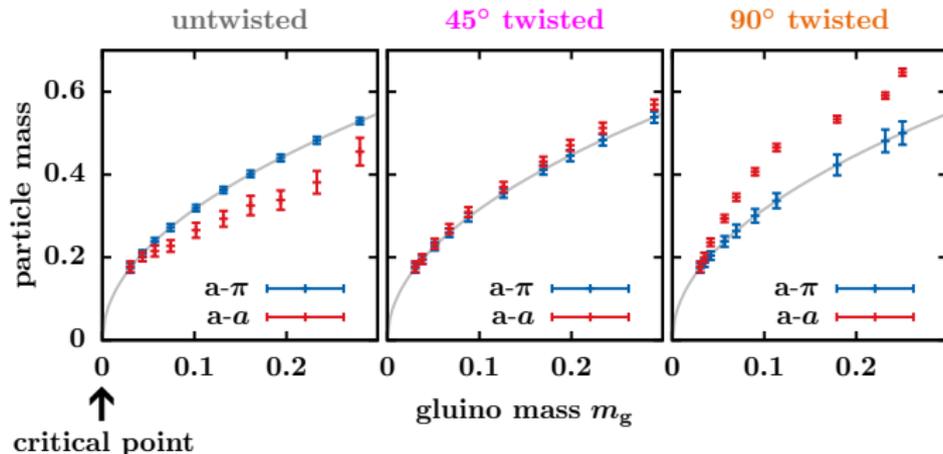
Connected Correlators

$$D_W^{\text{tw}}(x, y) = (4 + m + i\mu\gamma_5)\delta_{x,y} - \frac{1}{2} \sum_{\mu=\pm 1}^{\pm 4} (1 - \gamma_\mu) \mathcal{V}_\mu(x) \delta_{x+\mu, y}$$

$$m_{a-\pi}/m_{a-a} - 1$$



Connected Correlators



⇒ improvement of the chiral symmetry & supersymmetry at finite lattice spacing may be possible

Mesonic States

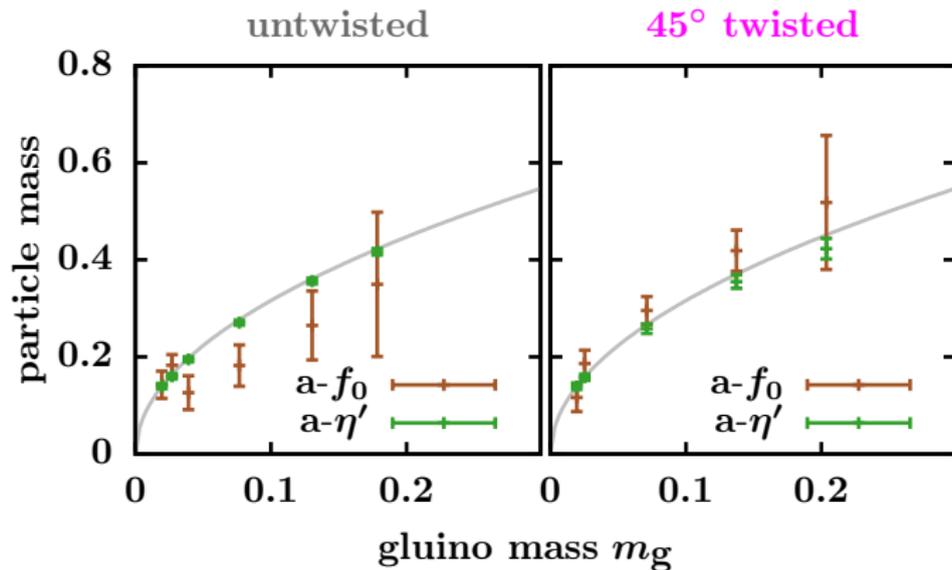


Figure: $8^3 \times 16$, preliminary

Mesonic States & Gluino-Glue

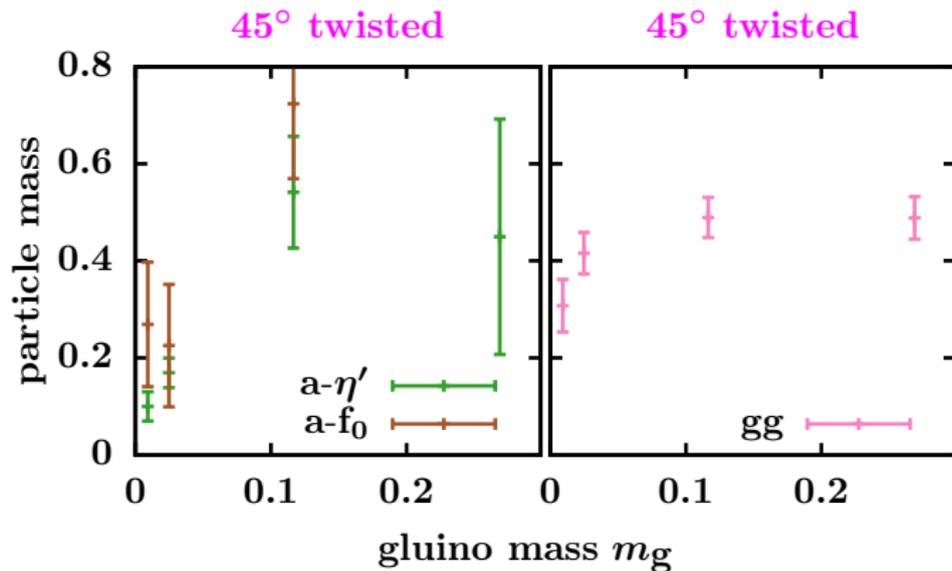


Figure: $16^3 \times 32$, preliminary

Summary

SYM on the lattice

- New analytical investigations
 - Fermionic expectation values
 - Eigenvalues of the free Dirac operator
- Lattice results
 - Fine tuning of bare gluino mass m (and μ) necessary to restore supersymmetry & chiral symmetry
 - Chiral symmetry of multiplet improved at finite lattice spacing with 45° twist
 - Wilson term twist provides an opportunity for $\mathcal{O}(a)$ improvement

Outlook

Summary

SYM on the lattice

Outlook

- Start of SuperMUC New Generation
- More statistics to confirm $m_{a-\eta'} \approx m_{a-f_0} \approx m_{gg}$ along 45° twist
⇒ improvement of the susy at finite lattice spacing may be possible
- Spectroscopy at 3 different couplings β for continuum limit

Appendix: $pq\chi pt$

SYM

- Only 1 flavor \Rightarrow partially quenched framework

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SYM

- Only 1 flavor \Rightarrow partially quenched framework

Pro

- Inclusion of anomalous chiral symmetry
- Masses of pseudoscalars, i.e. $a\text{-}\pi$, $a\text{-}\eta'$

Appendix: $pq\chi pt$

SYM

- Only 1 flavor \Rightarrow partially quenched framework

Pro

- Inclusion of anomalous chiral symmetry
- Masses of pseudoscalars, i.e. $a\text{-}\pi$, $a\text{-}\eta'$

Contra

- No native scalars, i.e. $a\text{-}a$, $a\text{-}f_0$
- No direct access to gauge field for hybrids, i.e. gluino-gluon

Appendix: Sign of the Pfaffian

Wilson-Dirac operator $D_W = D_W^{\text{tw}}(\mu=0)$

- is γ_5 -Hermitian: $(\gamma_5 D_W)^\dagger = \gamma_5 D_W$
- is \mathcal{C} -Antisymmetric: $(\mathcal{C} D_W)^\text{T} = -\mathcal{C} D_W$
- $\det(D_W) \in \mathbb{R}^+$
- $\text{Pf}(D_W) \in \mathbb{R}$

Appendix: Sign of the Pfaffian

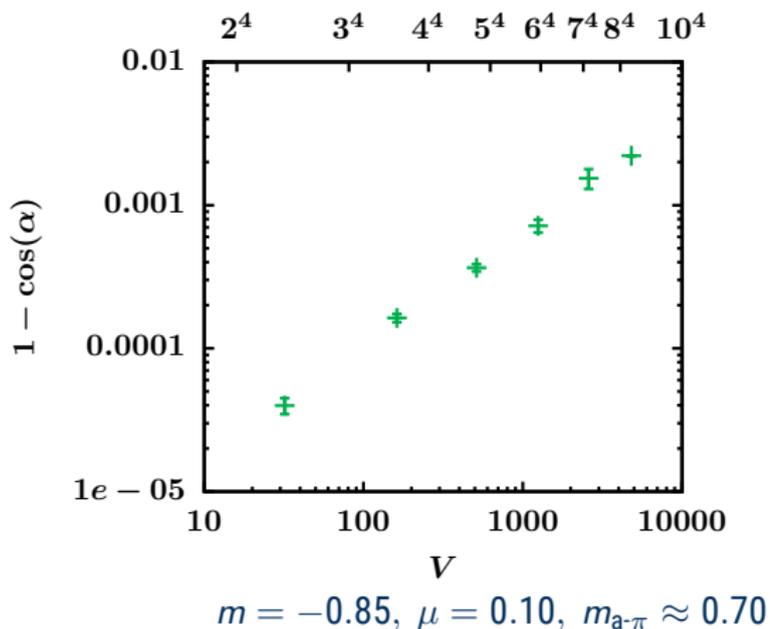
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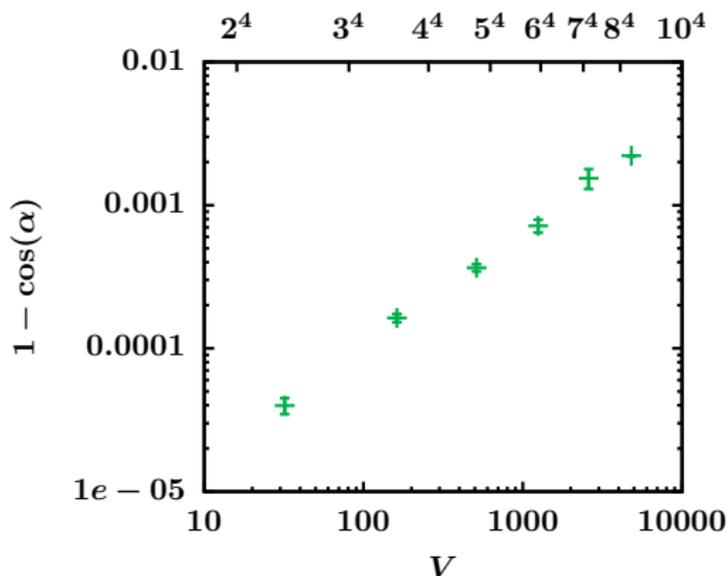
Twisted Wilson-Dirac operator $D_W^{\text{tw}}(\mu \neq 0)$

- in general $\text{Pf}(D_W^{\text{tw}}) \in \mathbb{C}$
- in continuum theory $m \rightarrow m_{\text{crit}}, \mu \rightarrow 0, a \rightarrow 0$: $\text{Pf}(D_W^{\text{tw}}) \in \mathbb{R}$
- at finite lattice spacing: phase of $\text{Pf}(D_W^{\text{tw}}) = |\text{Pf}(D_W^{\text{tw}})| \cdot e^{i\alpha}$ negligible

Appendix: Sign of the Pfaffian



Appendix: Sign of the Pfaffian



extrapolated to $16^3 \times 32$:
 $1 - \cos(\alpha) < 0.035$

$$m = -0.85, \mu = 0.10, m_{a-\pi} \approx 0.70$$