

Livio Bianchi * Università & INFN Torino

QGP in small systems at LHC: overview

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26

During the last 5 days we've been talking about small systems very extensively...

...but this is the topic I've been working on in the last 5 years...

...mostly thanks to Rene's money...



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6





2

26















26

... are ordinary business for experimentalists and theorists are also aware about!

- Geometrical bias
- Multiplicity estimation in small systems
- ...

... if that is not the case, let's discuss (and check backup)



3

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Results in small systems





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Measurement of **relative abundances of** different particle **species** Light hadrons (composed by *u* and *d*) abundantly produced in elementary collisions, but **strange hadrons suppressed**!

What happens in heavy-ion collisions?

- **1982 (Rafelski, Muller)**: **Strangeness enhancement** relative to elementary collisions proposed as smoking gun for **QGP formation**
- **Statistical Hadronization Model** (SHM): reproduce particle yields in HI by means of a Hadron-Resonance Gas in thermal equilibrium





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Experimental evidence: strange hadrons more abundant in HI than in pp(p-Be)

MIND THE REFERENCE: strangeness production depends on \sqrt{s} in small systems!

DOES IT? \downarrow





Strangeness enhancement in small collision systems (pp and p-Pb)

The larger the content in strangeness of the hadron, the steeper the increase is





A tiny droplet of the early universe?

2407 (1 by Senit Charley rticles seen by the ALICE experiment hint at the mation of quark-gluon plasma during protonton collisions.



"Many people think that protons are too light to produce this extremely hot and dense plasma" says Livio Bianchi, a postdoc at the University of Houston who worked on this analysis. "But these new results are making us question this assumption."

Scientists at the LHC and at the US Department of Energy's Brookhaven National Laboratory's Relativistic Heavy Ion Collider, or RHIC, have previously created quark-gluon plasma in goldgold and lead-lead collisions.

In the quark gluon plasma, mid-sized quarks—such as strange quarks—freely roam and eventually bond into bigger, composite particles (similar to the way quartz crystale grow within molten granite rocks as they slowly cool). These hadrons are ejected as the plasma fizzles out and serve as a telltale signature of their soupy origin. ALICE researchers noticed numerous protonproton collisions emitting strange hadrons at an elevated rate.

"In proton collisions that produced many particles, reserved, by the approximation of the product of the server of the server Bellwied Carofessor at the University of Houston." And infutional product of the University of Houston." And infution we saw an even bigger gap between the predicted number and our experimental results when we examined particles containing two or three strange quarks."

From a theoretical perspective, a proliferation of strange hadrons is not enough to definitively confirm the existence of

LOOK, I'M IN THE PRESS!! 7 ____ 26



ALI-PREL-134498

Ratio of yields to $(\pi^++\pi^-)$ \Box $2K_{S}^{0}$ HIGH MULTIPLI $\Lambda + \overline{\Lambda}$ (×2) 曲曲₫ Å $\Xi^{-}+\overline{\Xi}^{+}$ (×6) SATURATES 10⁻² $\Omega^{-}+\overline{\Omega}^{+}$ (x16) ALICE pp, $\sqrt{s} = 7$ TeV Nat. Phys. 13 (2017) 535-539 p-Pb, $\sqrt{s_{_{
m NN}}}$ = 5.02 TeV PLB 728 (2014) 25-38 Preliminary Pb-Pb, $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ 10⁻³ 10^{2} 10^{3} 10 $\left< \mathrm{d}\textit{N}_{\mathrm{ch}} \! / \! \mathrm{d} \eta \right>_{\mid \eta \mid < \ 0.5}$

Strangeness enhancement in small collision systems (pp and p-Pb)

The larger the content in strangeness of the hadron, the steeper the increase is



Multiplicity drives strangeness enhancement across different collision systems/energies



...and this is true for all "soft" particles



Iramatically

fails

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- **<u>PYTHIA</u>** (Lund string model):
 - Hadrons come from string (κ = 1 GeV/fm) breaking. s/u fit on data
 - At high energies need MPI to describe multiplicity and re-connection of colour strings to describe <p_T > VS multiplicity
 - Recently introduced:
 - Colour ropes: packing of strings increase κ
 - p-A and A-A environments
- **<u>DIPSY</u>** (Dipole evolution in Impact Parameter Space and rapiditY)
 - Evolution of initial state and collision described in impact parameter space.
 - Strings which overlap in impact parameter space form ropes



CAVEAT: ropes favor baryons wrt mesons. No flavour preference!





EPOS:

- Hard scattering: parton "ladders" + CGC-inspired saturation scale
- At time τ_0 (before hadronization) strings divided into fluid (CORE) and escaping (CORONA) according to momenta and local density
 - CORONA: strings can hadronize as in the Lund approach
 - CORE: from time τ_0 evolves as a viscous <u>hydrodynamic</u> system. Hadronization happens statistically at a common T_H
- After hadronization \rightarrow afterburner (e.g. UrQMD)

Good job with version 3 of the generator! Hints to the need of "core" part in pp collisions?



K. Werner, http://www.ectstar.eu/files/talks/Klaus.pdf



26



High precision data from the LHC suggest that the production of strangeness is driven by the final-state multiplicity of the collision

Independence on the collision energy

Can we extend this observation to lower energies? High multiplicity STAR results superimpose to ALICE's points

Can we infer something looking at the trend at lower multiplicity?

Would be interesting to complement with smaller systems results @RHIC!!



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GIVE ME 2 F-O TEMP!!





Selected results





According to the hydro picture, in HI the QGP is expected to develop:

- Radial flow (important in central collisions):
 - Common expansion velocity of partons
 - Translates into p_{T} spectra modification
 - Baryon/meson anomaly



 p_{T} spectrum gets harder as the collision gets more central

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Common $\beta \rightarrow \text{larger } p$ boost to higher-mass particles (p=m $\gamma\beta$)



According to the hydro picture, in HI the QGP is expected to develop:

- Radial flow (important in central collisions):
 - Common expansion velocity of partons
 - Translates into p_{T} spectra modification
 - Baryon/meson anomaly
- Anisotropic flow (important in semi-peripheral collisions):
 - Initial spatial anisotropy translates into final momentum anisotropy (pressure gradients)
 - Measured through angular anisotropies in the momentum distribution

$$E \frac{d^3 N}{dp^3} \approx \frac{1}{2\pi} \frac{d^2 N}{p_T dp_T d\eta} \left[1 + 2 \sum_{n=1}^{\infty} v_n \cos[n(\phi - \Psi_n)] \right]$$



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baryon/meson (high/low mass) ratio: from HI...

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Increase at mid- p_{T} in all centrality classes.

Peak shifting towards the right when going more central (higher radial boost in central collisions?)

> Evolution can be described by hydro at $low-p_T$

13 ____ 26

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> Evolution can be described by hydro at low- p_{T}

NOTE: it is not a strangeness nor baryon/meson –related effect!

10

 $\dot{p}_{\rm T}$ (GeV/c)

baryon/meson (high/low mass) ratio: from HI...

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Peak shifting towards the right when going more central (higher radial boost in central collisions?)

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Hint for a similar evolution in pp from STAR?



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Same pattern in the Λ/K_{S}^{0} measured in small systems, with different magnitude...

> ...but... MIND THE MULTIPLICITY SPAN!

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ALICE, arXiv:1807.11321 PLB 728 (2014) 25-38 PRL 111 (2013) 222301 °S 2.2 N / 2 N / V ALICE Preliminary pp Vs = 7 TeV ALICE p-Pb Vs_{NN} = 5.02 TeV ALICE Preliminary V0M Class I, $\langle dN_{ch}/d\eta \rangle = 21.3$ Pb-Pb Vs_{NN} = 5.02 TeV **60-80%**, (dN_μ/dη) = 9.8 VOM Class X, $\langle dN_{n} \rangle / d\eta \rangle = 2.3$ (VOA Mult. Classes - Pb side) (VOM Multiplicity Classes) **0.8 0.6**⊢ **0.4**⊧ **0.2**⊢ 10 10 *p*_{_} (GeV/*c*) ALI-PREI -135238 0°¥ ∀ √ 1.4 1.6 2.50 < p_ < 2.90 GeV/c $6.50 < p_{_{
m T}} < 8.00 ~{
m GeV}/c$ ALICE Preliminary pp vs = 7 Te pp: pp: p-Pb: $6.00 < p_{-} < 8.00 \text{ GeV}/c$ p-Pb: $2.42 < p_{\perp} < 3.20 \text{ GeV}/c$ ALICE p-Pb Vs_{NN} = 5.02 TeV Pb-Pb: $2.40 < p_{-} < 3.00 \text{ GeV}/c$ Pb-Pb: $6.50 < p_{\perp} < 8.00 \text{ GeV}/c$ ALICE PL Pb Vs_{NN} = 2.75 TeV $0.60 < p_{_{T}} < 0.80 \text{ GeV}/c$ 0.8 0.6 0.4 0.2 10^{3} 10² 10² 10 10^{2} 10^{3} 10^{3} 10 10 $\left< \mathrm{d} \textit{N}_{\mathrm{ch}} / \mathrm{d} \eta \right>_{|\eta| < 0.5}$

Same pattern in the Λ/K^0_s measured in small systems, with different magnitude... ...but...

MIND THE MULTIPLICITY SPAN!

In order to make proper comparison, one can select $p_{\rm T}$ ranges and look at multiplicity dependence

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26

14



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MIND THE MULTIPLICITY SPAN!

In order to make proper comparison, one can select $p_{\rm T}$ ranges and look at multiplicity dependence

Clear continuity among different systems! Is the underlying mechanism the same here? Need to compare p_{T} spectra to hydro

ALI-PREL-110566



Blast wave - simplified hydro model:

- Assumes common particle expansion with β_{T} and T_{kin}
- If assumption ok: fit (e.g.) π,K,p → predict p_T shape of other particles
- Assumption ~ok for all collision systems
- pp and p-Pb: similar T_{kin} - β_T progression







ALI-PREL-122516

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- pp and p-Pb: similar T_{kin} - β_T progression
- Considering corresponding multiplicity: less "violent" expansion in A-A, but T_{kin} common for all systems



 $\langle \beta_{-} \rangle$







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$v_n \neq 0$ observed in HI at RHIC and LHC



What about small collision systems?

Global hydro fits to several bulk observables start appearing:




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 $v_2 \neq 0$ observed in all collision systems

NOTE: contribution of non-flow not easy to estimate in pp (and p-Pb)

...but does this make sense at all? Can hydro develop in so small systems?

Naïve expectation: need "large enough" and "live long enough" medium to reach thermal equilibrium and apply hydro (several interactions needed)

• $R > \lambda$

 $>\lambda/v$

MEAN FREE PATH



17



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Too restrictive: hydro can be applied far from thermalization!

W. Li, arXiv: 1704.03576





First theoretical calculations involving hydro expansion of a single fluid in all collisional systems start appearing.

26

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Selected results





Partonic energy loss

- High-p_T partons produced in the early stages of the collisions (τ<<1fm)
- Loose energy in the medium through:
 - elastic scattering
 - Induced gluon radiation (dominant at high-p_T)
- Simple prediction (dead-cone effect):

 $\Delta E_g > \Delta E_{light-quark} > \Delta E_{heavy-quark}$









Hard probes: going "smaller"

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20 ____ 26



Hard probes: going "smaller"

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20 ____ 26_





-0

ALICE Preliminary, Pb-Pb 15_{NN} = 5.02 TeV

0-5%

π⁺ + π⁻

Hard probes: going "smaller"

5-10%

10-20%

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20

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Difficult to define an R_{AA} in pp...!!!

Let's concentrate on p-Pb









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LHCb, arXiv:1707.02750



Quarkonia

22

26

 the original idea: quarkonium production suppressed via color screening in the QGP

T.Matsui and H.Satz, Phys.Lett.B178 (1986) 416





Quarkonia

22

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the original idea:
quarkonium production suppressed
via color screening in the QGP

T.Matsui and H.Satz, Phys.Lett.B178 (1986) 416



• (re)combination

enhanced quarkonium production via (re)combination during QGP phase or at hadronization

Central AA	SPS	RHIC	LHC	LHC
collisions	20 GeV	200 GeV	2.76TeV	5.02TeV
N _{ccbar} /event	~0.2	~10	~85	~115

P. Braun-Muzinger, J. Stachel, PLB 490(2000) 196 R. Thews et al, Phys.Rev.C63:054905(2001)



26



No F.S. suppression for J/ ψ in p-Pb collisions

Quarkonia melting in small systems?

----- 2.03 < y_{cms} < 3.53

-4.46 <y_ms< -2.96

12

14

 $\langle N$

16

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6

8

10

No F.S. suppression for J/ ψ in p-Pb collisions

...but ratio $\psi(2S)/J/\psi$ significantly lower than 1 at large N_{coll}!!

Makes sense in the "sequential suppression scenario": $\psi(2S)$ should dissociate at lower T





 $T/T_c 1/\langle r \rangle [fm^{-1}]$

Y(15)

χ_b(1P)

 $\chi_{c}(1P)$

≤T_c

J/ψ(1S) Y'(2S)

χ_b'(2P) Υ"(35)

Ψ'(25)

ALI-DER-105924

0.5

n

0

2



26





...but then, why $\Upsilon(2S)$ is suppressed in p-Pb and even pp high-multiplicity events?

Perspective:

 ψ (2S)/J/ ψ versus multiplicity in pp collisions?

Quarkonia melting in small systems?

24

26





...but then, why $\Upsilon(2S)$ is suppressed in p-Pb and even pp high-multiplicity events?

Perspective:

 ψ (2S)/J/ ψ versus multiplicity in pp collisions?



Conclusions



"small systems" path the way to a possibly deeper (microscopic) understanding of QGP phenomena :

- Final state multiplicity drives light flavours observables across systems and energies.
- Strangeness enhancement in pp collisions. In highest multiplicity, hadrochemistry \approx to the one in the QGP
- $v_2 \neq 0$ in pp and p-Pb collisions at the LHC.
- No parton energy loss observed in pp and p-A
- Intriguing (and unclear) results on quarkonium suppression in p-A (and pp!) collisions

NAMK NOU



My experience in an US research group

26

26

Huston is a great city:









My experience in an US research group

Huston is a great city:







26











The dream :

The reality :





Loano: the saddest sea place in Italy (sorry Claudia)



The dream :

The reality :













The dream :

The reality :

















Enjoy your **ies!!

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Enjoy your **ies!!

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Enjoy your **ies!!

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Multiplicity estimation in small systems

68

23



y-values:

- Measure p_T spectra of strange particles and pions in pp events characterized by different multiplicities (fwd-rapidity estimator)
- Integrate spectra extrapolating at low and high p_T with suitable functions.
- Calculate Y^{S}/Y^{π}

x-values:

- ∀ multiplicity class (fwd-rapidity estimator), count the number of primary charged particles at central rapidity and build-up dN_{ch}/dη distribution
- Take statistical average of every distribution





Multiplicity bias in small systems



Selection bias and R_{AA} in peripheral collisions at LHC



The resonances' story (I)

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71

23

Resonances are powerful tools to probe the hadronic phase after chemical freeze-out



The resonances' story (II)

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72

23

Resonances are powerful tools to probe the hadronic phase after chemical freeze-out

Lifetime [fm/c] : ρ [1.3] < K* [4.2] < Λ * [12.6] < Ξ^{0*} [21.7] < ϕ [46.2]








12

34

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Fix yield's ratio to saturation limit. Check the evolution when decreasing the volume (multiplicity)

> Qualitatively the thermal fit describes K, Λ , Ξ , Ω

Notable exception is the ϕ !

Slightly decreasing protons Hint for hadronic re-scattering? Need to evaluate degree of correlation on systematics across multiplicity!



If interested in re-scattering in the hadronic phase, more in the backup!



19

34



 $v_n ≠ 0$ observed at RHIC and LHC: means that in semi-central collisions the p_T distribution of particles is anisotropic wrt the event plane...

does this mean we have flow?



Hydrodynamic models reproduce v_n in all centralities by means of an "almost" perfect fluid: $\eta/s=0.2$

The Lund string model: basics

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75

23

- Linear confinement potential for large distances (confirmed by lattice QCD). For short distances perturbation theory holds
- Confined colour fields described as strings with tension $\kappa = 1 \text{ GeV/fm}$
- Breaking of strings (tunneling) give hadrons

 $P \propto e^{-\frac{\pi m_T^2}{\kappa}} = e^{-\frac{\pi m_q^2}{\kappa}} \cdot e^{-\frac{\pi p_{Tq}^2}{\kappa}}$

 Flavour of hadrons determined by the Gaussian mass suppression term (which mass to put? If current → less ssuppression than observed. If constituent → too much ssuppression. s/u empirical number to be tuned on data)



- In hadronic collisions multiple strings needed to describe multiplicity distribution (MPI)
- In the LC Lund model each string hadronizes separately with respect to the others
- The multiplicity increases, but not the $\langle p_T \rangle$ nor the relative flavor abundancies!
 - Multiple strings are close in space-time. Dynamical interaction not implemented in this model, but colour re-arrangement can happen: Colour Reconnection (CR)
 - Takes place after parton shower and takes into account all SU(3) permitted configurations. **Selection parameter: minimum total string length**
 - After re-arrangement of strings, hadronization takes place
 - Correctly takes into account colour re-arrangement in remnant



PYTHIA: effect of CR





CR mimics features that we traditionally attribute to collective flow, but something more is needed. Tuning?

- 3 main parameters tuned on data: c_{time} ($\langle p_T \rangle$), c_j (Λ/K_S^0) and p_T^{ref} ($dN_{ch}/d\eta$).
- The presence of junctions increases baryon production at intermediate p_{T} , but not sufficient to reproduce data
- Λ/K_S^0 shape (magnitude is tuned!) reproduces data up to 3 GeV/ $c \rightarrow$ problem in spectra common to baryons and mesons?



76

23

The DIPSY model: basics & ropes

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77

23

- Partonic model in impact parameter space and rapidity (Dipole evolution in Impact Parameter Space and rapiditY)
- Mueller dipole model (LL-BFKL)
- Proton/Nucleus structure built up dynamically from dipole splittings
- Builds-up initial state + collision in impact parameter space. Naturally treats saturation and MPI

0.4 0.2 0.1 0.1 -0.1 -0.1 -0.2 -0.3 -2.0 -1.0-1.5-0.50.0 0.5 1.0 1.5 2.0 V 0.8 0.6 0.4 0.2 [fm] 0 -_ 0.2 م -0.4 -0.6 -0.8 **Transverse size of strings** 10 the one visualized here!!

To the question "Which are the strings that can interact?" the DIPSY model answers following the evolution of colour strings during the whole parton shower

How do strings interact?

Stack of colour strings close in the IP-y space:

can form colour singlets or multiplets according to the summing rules of SU(3) Singlets correspond to simple re-arrangement of single strings, Multiplets correspond to **ROPES**.

Hadronizing a rope means fragmenting stringby-string with an effective string tension $\kappa > \kappa_0$

As we know from previous works, higher string tension \Rightarrow more baryons and more flavours \neq (u,d)

Before hadronizing a string a "swing" mechanism further allow colour rearrangements (in analogy with colour re-connection)

Flensburg et al. arXiv:1103.4321

EPOS: the melting pot

78

23



- Hard scattering treated with the addition of several DGLAP parton "ladders" (pomerons) + a CGC-inspired saturation scale
- Parton ladders are then considered as relativistic strings, conveniently treated in a string fragmentation approach (a-la Lund)
- At time τ_0 (well before hadronization) strings are divided into: fluid (CORE) and escaping (CORONA) according to their momenta and density of the string segments
 - **CORONA**: strings can hadronize as in the Lund approach
 - **CORE**: from the time τ_0 evolves as a viscous <u>hydrodynamic</u> system. Hadronization happens statistically at a common T_H
- After hadronization hadron-hadron rescattering can be considered, making use of an afterburner (e.g. UrQMD)

<u>NOTE</u>: parameters governing the core-only part are 6 $(\tau_0, \rho_0, \varepsilon_{FO}, \gamma_{rad}, f_{ecc}, \gamma_s)$, to be tuned on data!!

Werner, PRL 98, 152301 (2007)

EPOS: effects of Core-Corona (I)

79

23



0

- Radial flow of the core also dominates the intermediate region of the p_T spectrum
- High p_T is dominated by escaping fragmenting strings

NOTE: the exact onset of the effect depends on tuning (p_T cut-off for escaping strings)

Pierog & Karpenko & Katzy & Yatsenko & Werner, arXiv:1306.0121

10

 $p_t (GeV/c)$





Observed trends of relative particle yields **reproduced** thanks to **interplay** between **core** and **corona** (+ UrQMD)

TAKE HOME

Spectra + yields described in EPOS through evolution with multiplicity of relative importance of CORE and CORONA

<u>NOTE</u>: Does this imply QGP in small systems? NO! May or may not be.

- Relative importance of CORE/CORONA in the yields for long and short living resonances is strikingly different
- Mild Φ enhancement with multiplicity observed in EPOS

