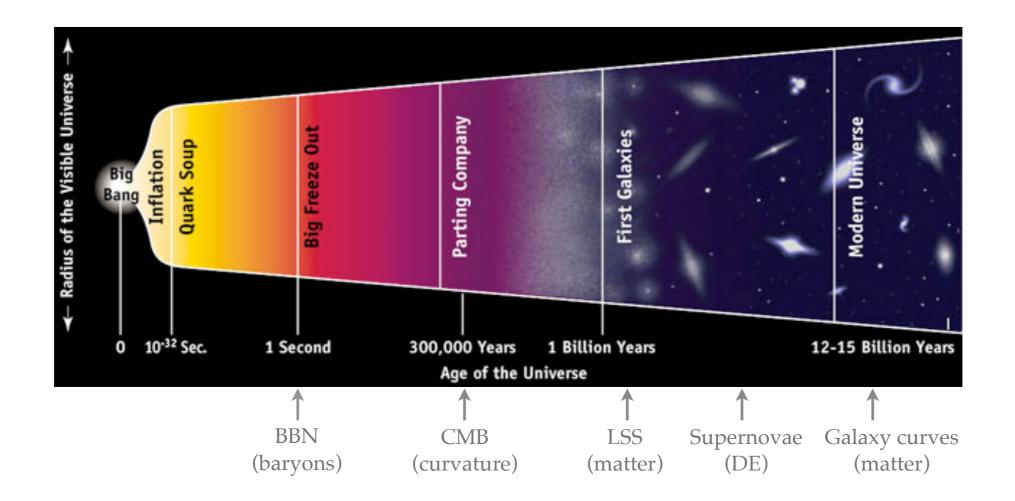


# IDENTIFICATION OF DARK MATTER PANORAMA: CHALLENGES AND PROMISE

Kathryn Zurek

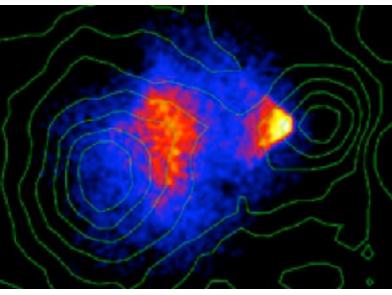
### WHY DARK MATTER? (WHY NEW PARTICLE PHYSICS?)

 The dark matter paradigm is the only successful framework for understanding the entire range of observations from the time the Universe is 1 sec old.



#### **SUPER-WEAKLY INTERACTING**

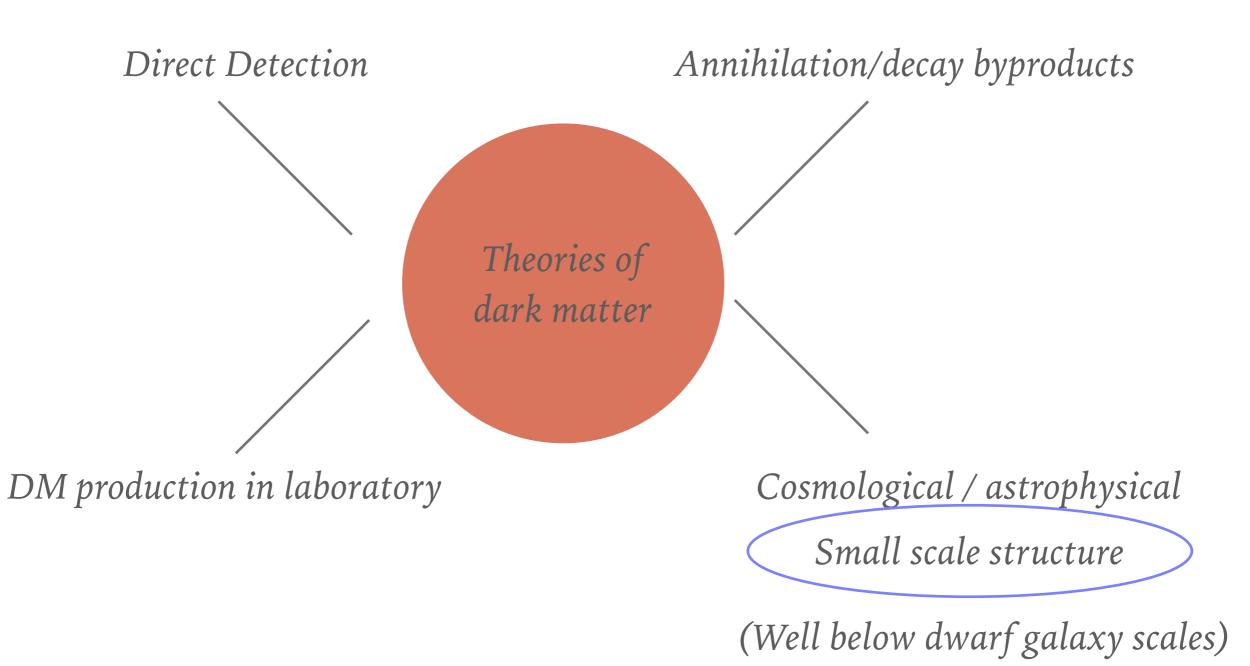
#### • Gravitational Coherence ....



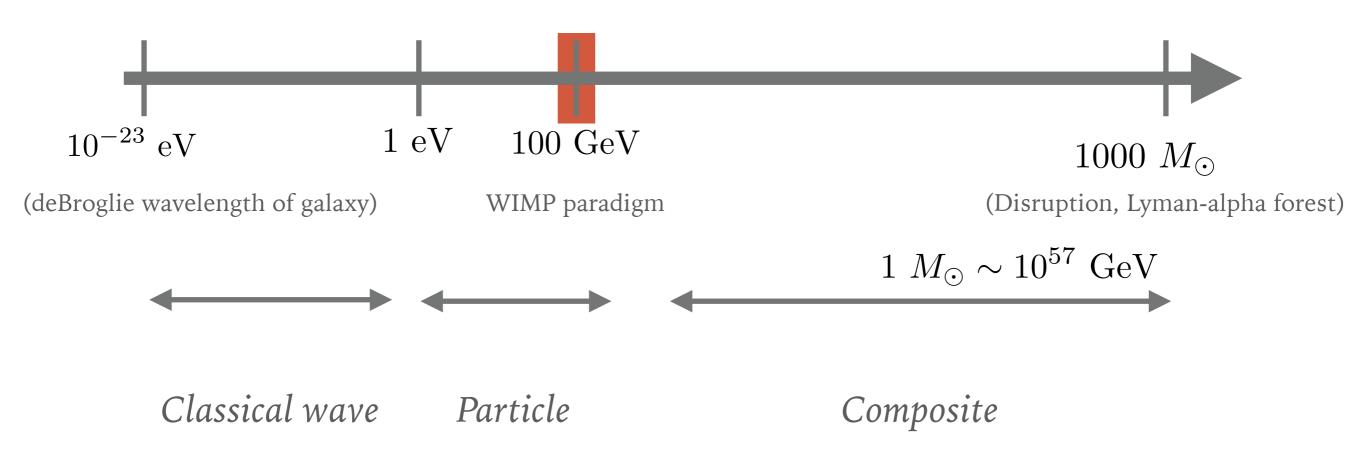
... on cosmological scales!

- Helps us learn about aggregate properties of dark matter
- Particle properties much harder
- Fundamental premise: DM has interactions other than gravitational

### MANY FACES OF DARK MATTER



#### DARK MATTER DETECTION: A FULL COURT PRESS



### **IDENTIFICATION OF DARK MATTER: CHALLENGES AND PROMISE**

 Challenge: WIMP paradigm is not dead, but it's under enormous pressure. Promise: Still opportunities for discovery in coming years

- Challenge: Vast majority of QCD axion parameter space still unprobed. Promise: ADMX has made enormous progress and there are many new ideas being developed
- Challenge: Hidden sector dark matter represents a vast and mostly unconstrained frontier. Promise: There are a suite of experiments that can probe well-motivated models in the coming 5-10 years
- Challenge: Dark matter may interact only gravitationally.
   Promise: Probes of dark matter substructure may still tell us about underlying theory

#### CHALLENGES: THE WIMP PARADIGM

- Decades of searching for the WIMP, with no evidence of new physics at the weak scale
- Production at accelerators strongly constrains colored states, not thermal electroweakinos

$$3 \times 10^{-26} \text{ cm}^3/\text{s} \sim \frac{g_{wk}^4}{4\pi (2 \text{ TeV})^2}$$

 Most promising avenue currently seems to be "minisplit" scenario, with colored super-partners well above TeV scale

#### **CHALLENGES: THE WIMP PARADIGM**

#### ATLAS SUSY Searches\* - 95% CL Lower Limits

#### ATLAS Preliminary

May 2020 Model Signature ∫L dt [fb<sup>-1</sup>] Mass limit Reference 0 e. µ § [10x Degen.] 2-6 jets 1.9 m(2)<400 GeV ATLAS-CONF-2019-040  $\bar{q}\bar{q}, \bar{q} \rightarrow q\bar{t}_{1}^{0}$  $E_{Liss}^{miss}$ 139 1-3 jets 0.43 mono-iet Dr. By D 0.71 36.1 m(j)m(t1)=5 GeV 1711.03301 22.2-9999 0 0.00 2-6 jets Emiss 129 m(N\_1)-0:04V ATLAS-COME-2019-040 2.25 1.15-1.95 m(2)=1000 GeV Forbidden ATLAS-CONF-2019-040 m(1)-800 GeV 3 c. H 4 jets 36.1 1.85 1706.03731  $\tilde{\chi}\tilde{\chi}, \tilde{\chi} \rightarrow q\tilde{q}(\ell\ell)\tilde{k}_{1}^{0}$ ٥ð 2 jets  $E_T^{miss}$ ee. µµ 36.1 1.2 1805.11381 m(2)-n(2)=50 GeV 0 e. µ 7-11 jets ATLAS-CONF-2020-002  $E_T^{miss}$ 139 1.97 m(F) <600 GeV  $\hat{\chi}\hat{\chi}, \hat{\chi} \rightarrow qqWZ\hat{t}_1$ SS e. µ 6 jets 139 1.15 m(¿)-m(ž)=200 GeV 1909.08457 0-1 e.µ 3b79.8 m(2)<200 GeV ATLAS-CONF-2018-041 22.2-MT  $E_T^{mi}$ 2.25 SS c. µ 6 jets 139 1.25 m(i)-m(i)=300 GeV 1909.08457 Multiple  $m(\hat{t}_1^0)=300 \text{ GeV}, BP(h\hat{t}_1^0)=1$ 1708.09266, 1711.00301  $\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b \tilde{k}_1^0 / d_1^*$ 36.1 Fabidden 0.9 0.74 m(t)=200 GeV. m(t)=300 GeV, BR(t)=1 Multiple 139 Forbidden 1909.08457  $\bar{b}_1\bar{b}_1, \bar{b}_1 \rightarrow b\bar{\ell}_2^0 \rightarrow bh\bar{\ell}_1^0$ 0 e. µ 6.6  $E_T^{miss}$ 139 Б. Forbidde 0.23-1.35 An(\$2,\$1)=130 GeV, m(\$1)=100 GeV 1908.03122 0.23-0.48 3m(\$2, \$1)=130 GeV, m(\$1)=0 GeV 1905.03122 0-1 e.µ ≥ 1 jet 1.25 m(2))=1 GeV ATLAS-CONF-2020-003, 2004.14060  $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow t \tilde{t}_1^0$ Emiss 139  $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow W b \tilde{t}_1^0$ 1 c. µ 3 jets/1 b  $E_T^{miss}$ 139 0.44-0.59 m(R<sup>2</sup>)=400 GeV ATLAS-CONF-2019-017 1 τ + 1 e,μ,τ 2 jets/1 b m(71)=800 GeV  $\tilde{l}_1\tilde{l}_1, \tilde{l}_1 \rightarrow \tilde{\tau}_1 br, \tilde{\tau}_1 \rightarrow \tau \tilde{G}$  $E_T^{miss}$ 36.1 1.16 1808.10178  $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow c \tilde{\chi}_1^0 / \tilde{c} \tilde{c}, \tilde{c} \rightarrow c \tilde{\chi}_1^0$ 36.1 0.85 m(tr)=0 GeV 1805.01649 0 e. µ 2.c  $E_T^{min}$ 1805.01649 0.46 m(i, z)-m(i)=50 GeV m(i, z)m(i)=5 GeV  $0 c. \mu$ morno-jet 36.1 1711.03301 E.,  $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}^0_2, \tilde{\chi}^0_2 \rightarrow Z/h\tilde{\chi}^0_1$ 1-2 e. µ 1-4.bFrais 139 0.067-1.18 m(2)-500 GeV SUS7-2018-09  $\tilde{l}_2\tilde{l}_2, \tilde{l}_2 \rightarrow \tilde{l}_1 + Z$ 3 c. H 1b $E_T^{min}$ 139 Forbidden 0.86 m(7)=360 GeV, m(7)-m(2)=40 GeV SUS7-2018-09  $\hat{x}_{1}^{\pm}\hat{x}_{2}^{0}$  via WZ 0.64 3 e. µ Eniss Eniss 139  $m(\hat{t}_{1}^{0})=0$ ATLAS-CONF-2020-015 ≥ 1 jet 0.205 m(t))m(t))=5 GeV ee. µµ 139  $\hat{x}^{h}/\hat{x}$ 1911.12606  $\hat{x}_{1}^{\pm}\hat{x}_{1}^{\mp}$  via WW  $2 e. \mu$ 139 0.42 m(t))=0 1908.08215 Eniss  $\hat{x}_i^*$ 0-1 e. µ 2 b/2 y  $\hat{\chi}_1^{\pm}/\hat{\chi}_2^{0}$ 2004.10854.1909.09226  $\hat{x}_{1}^{\pm}\hat{x}_{2}^{0}$  via Wh Emin 139 Forbidden 0.74 n(x1)=70 GeV  $\hat{\chi}_{1}^{*}\hat{\chi}_{1}^{*}$  via  $\hat{\ell}_{L}/\hat{\nu}$ 2 c. #  $m(\tilde{\ell}, \tilde{r}) = 0.5m(\tilde{t}_1^+) + m(\tilde{t}_1^0))$ 1908.08215  $E_{i}^{m}$ 139 1.0 0.16-0.3 0.12-0.39 1911.06660  $\hat{\tau}\hat{\tau}, \hat{\tau} \rightarrow t\hat{k}_{1}^{0}$ 2τ  $E_T^{min}$ 139 PL-PRIJ  $m(\hat{t}_{1}^{0})=0$ 2 c. µ 0 jets ET. 139 0.7  $m(\hat{t}_{1}^{0})=0$ 1908.08215  $l_{LR}l_{LR}, l \rightarrow t \hat{\chi}_1^0$ 10.111 > 1 jet 1:30 0.396 m(7) m(%)=10 GeV 1011.12006  $\hat{H}\hat{H}, \hat{H} \rightarrow h\hat{G}/Z\hat{G}$ 0 e. µ  $E_T^{miss}$  $E_T^{miss}$ 0.13-0.23 0.290.88  $BR(\tilde{t}_1^0 \rightarrow h\tilde{G})=1$ 1805.04030  $\geq 3b$ 36.1 ĥ 4 c. µ 0 jets 36.1 0.3  $B8(\tilde{x}_1^0 \rightarrow Z\tilde{G})=1$ 1804.03602 Disapp. trk 1 jet  $E_T^{miss}$ 0.46 1712.02118 Direct  $\hat{x}_1^* \hat{x}_1^-$  proc., long-lived  $\hat{x}_1^*$ 36.1 Pure Wino 0.15 £. Pure higgsing ATL-PHYSPUB-2017-019 Stable ¿ R-hadron Multiple 36.1 2.0 1902.01005.1000.04095 Multiple 2.05 2.4 1710.04901,1808.04095 36.1 Metastable g R-hadron, g→qq E<sub>1</sub> m (c)=100 GeV  $\hat{\chi}_{1}^{\pm}\hat{\chi}_{1}^{\mp}/\hat{\chi}_{1}^{0}$ ,  $\hat{\chi}_{1}^{\pm}\rightarrow 2\ell \rightarrow \ell\ell\ell$ 1.05 Pure Wino ATLAS-CONF-2020-009 30.0 139 0.625 LFV  $pp \rightarrow \bar{\nu}_{\tau} + X_{\tau}\bar{\nu}_{\tau} \rightarrow e\mu/e\tau/\mu\tau$ CH.CT.HT 3.2 1.9 £311=0.11, £12/110/203=0.07 1607.08079  $\hat{\chi}_{1}^{\pm}\hat{\chi}_{1}^{\mp}/\hat{\chi}_{2}^{0} \rightarrow WWZUUUrv$ 4 e. µ 0 jets  $E_T^{miss}$ 1.33 36.1 0.82 m(2)=100 GeV 1804.03602 Sec. # 0. Jun # 0  $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq \tilde{k}_1^0, \tilde{k}_1^0 \rightarrow qq q$ 4-5 large-R jets 1801.03568 36.1 Large 2" 1.3 1.0 2.0 Multiple 36.1 1.05 m(R1)=200 GeV, bino-like ATLAS-CONF-2018-003 ď. Multiple 0.55 1.05  $\tilde{n}, \tilde{i} \rightarrow t \tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow bs$ 36.1 m(R)=200 GeV, bino-like ATLAS-CONF-2018-003  $\vec{H}, \vec{I} \rightarrow b\hat{\chi}_{1}^{*}, \hat{\chi}_{1}^{*} \rightarrow bbs$  $\geq 4b$ Fortidden 0.95 m (t)=500 GeV ATLAS-CONF-2020-016 139  $\bar{l}_1\bar{l}_1, \bar{l}_1 \rightarrow bs$ 2 jets + 2 b 36.7 0.42 1713.07171 0.61  $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$ 2 e. µ 2b36.1 0.4-1.45 BR(7, -+br/hpt)>20% 1713.05544  $1 \mu$ DV 1395 1.6 BH(/\_-saul=100%, 0056=1 2008.11956

1

Mass scale [TeV]

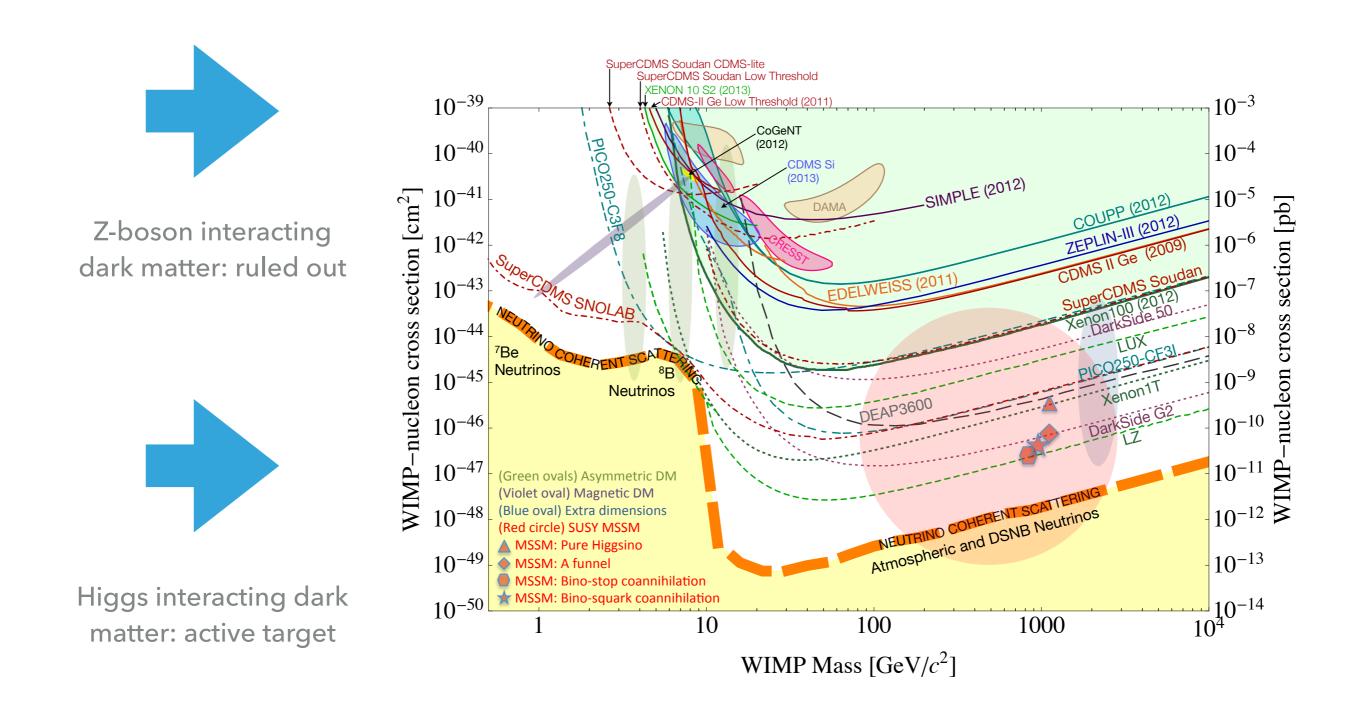
"Only a selection of the available mass limits on new states or phénomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

 $10^{-1}$ 

 $\sqrt{s} = 13 \text{ TeV}$ 

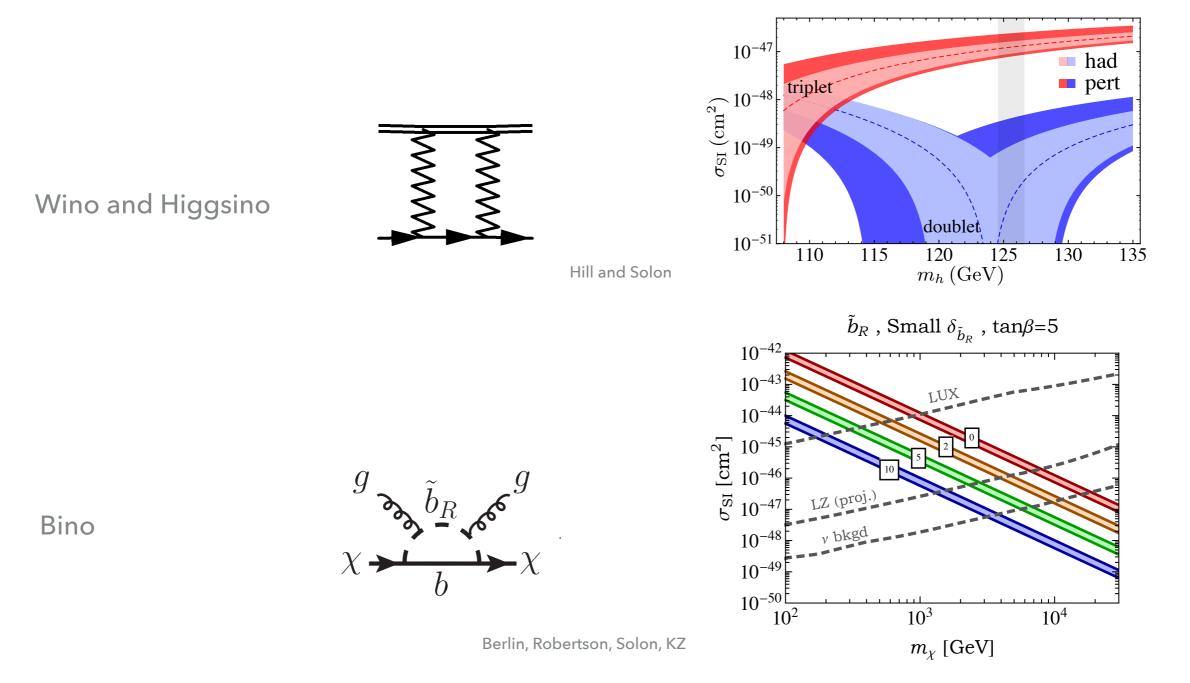
. . . . . . . .

#### **PROMISE: DETECTING HIGGS INTERACTING DARK MATTER**



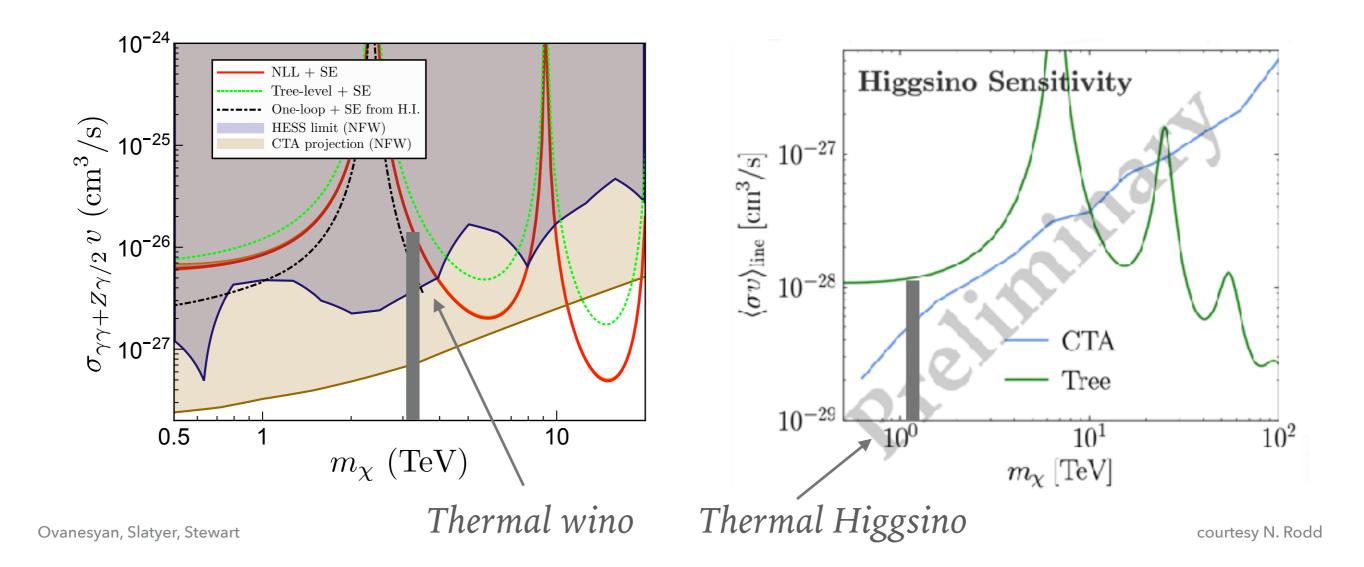
### **PROMISE: ELECKTROWEAKINOS VIA 1-LOOP PROCESSES**

Challenge: even LZ will struggle to reach the wino

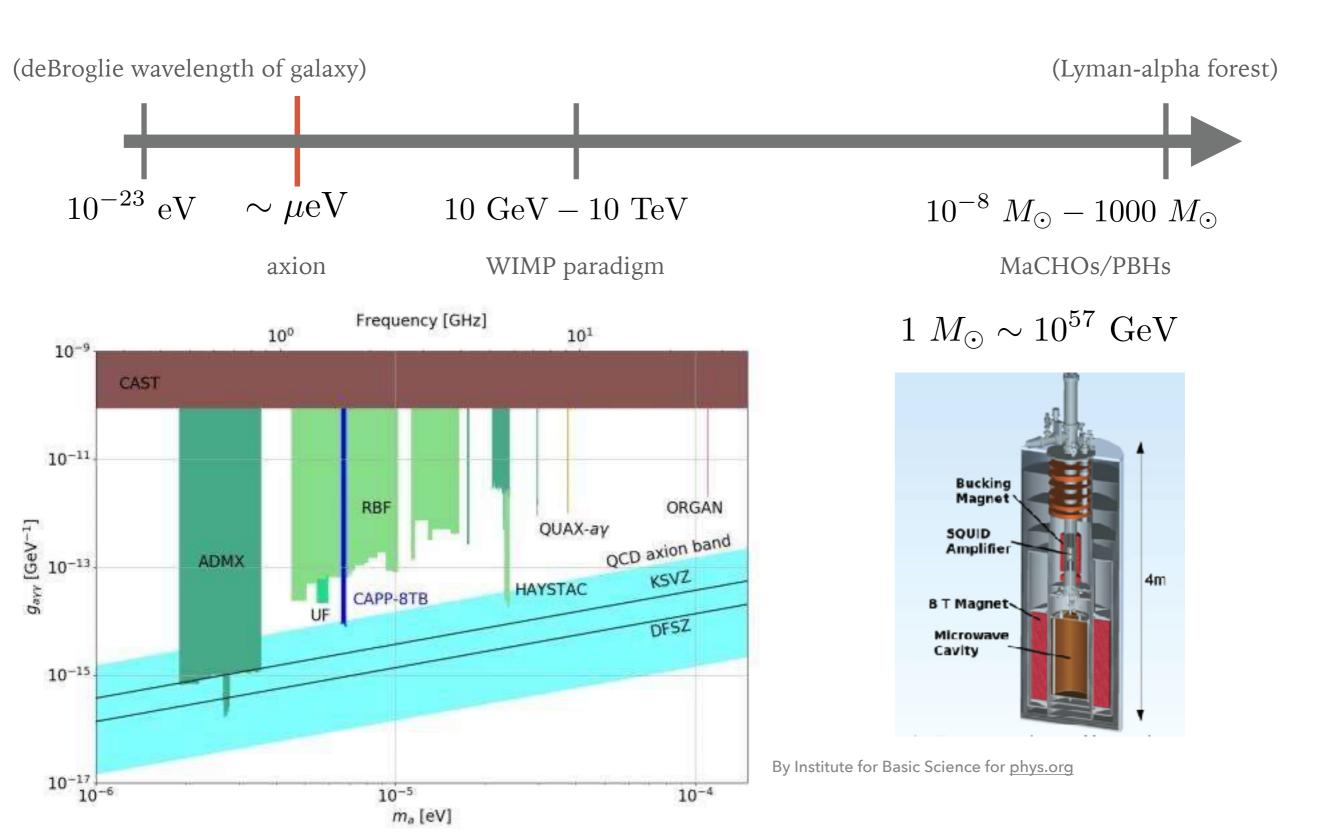


#### PROMISE: INDIRECT DETECTION STILL HAS REACH ON THERMAL ELECTROWEAKINOS

Currently some sensitivity to winos and possible future sensitivity with CTA to Higgsinos

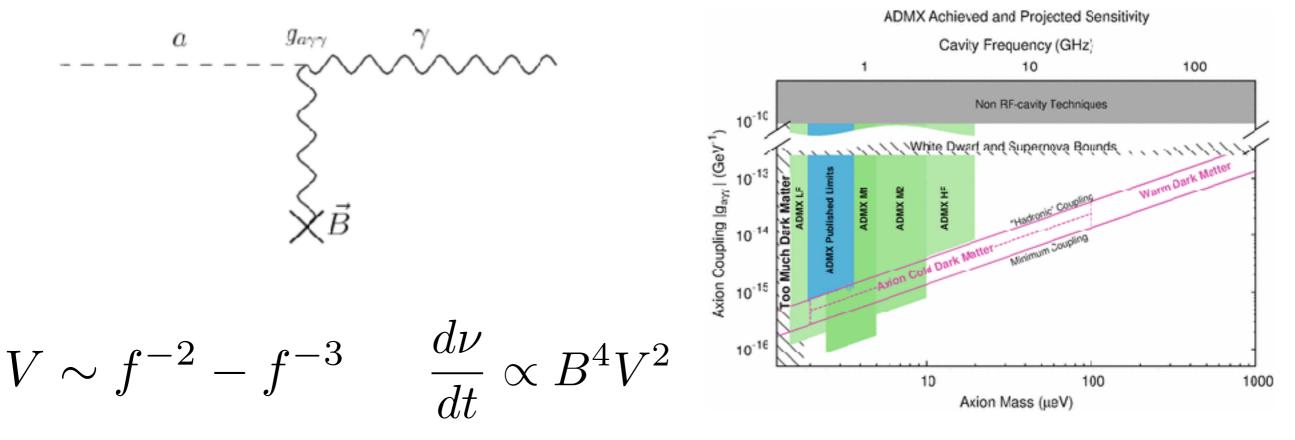


#### PROMISE: AXION DETECTION HAS SENSITIVITY TO THE QCD AXION!



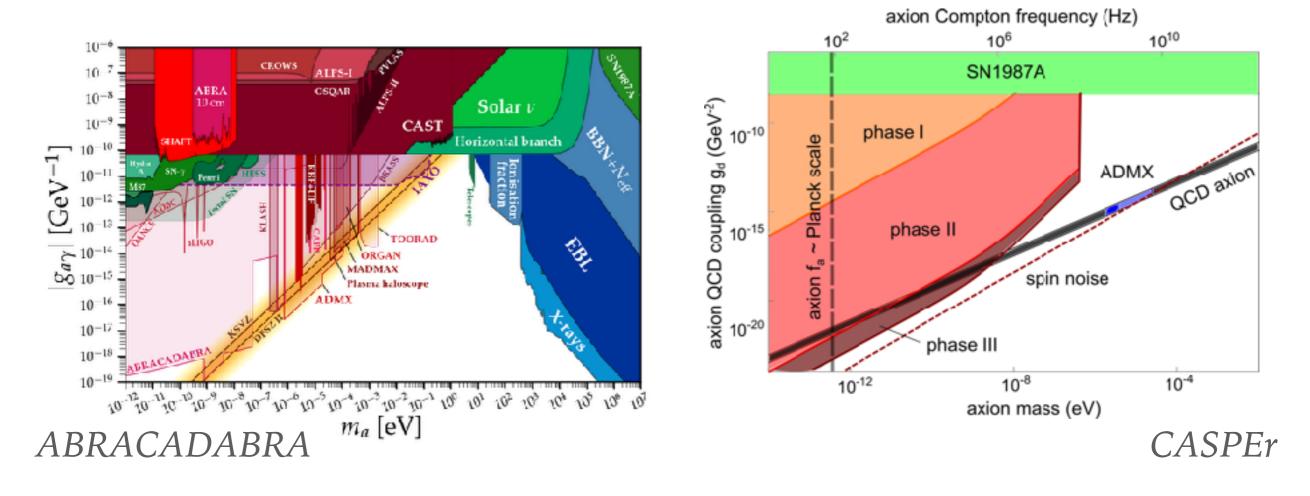
#### CHALLENGE: VERY LIMITED FREQUENCY RANGE IN MASS

- The classic, decades-old Sikivie proposal is now achieving sensitivity to the QCD axion
- Challenge: scaling to higher frequency and covering the whole range originally motivated by misalignment will be very challenging because of matching axion mass to cavity size



#### **PROMISE: IDEAS FOR AXION DETECTION HAS UNDERGONE A RENAISSANCE**

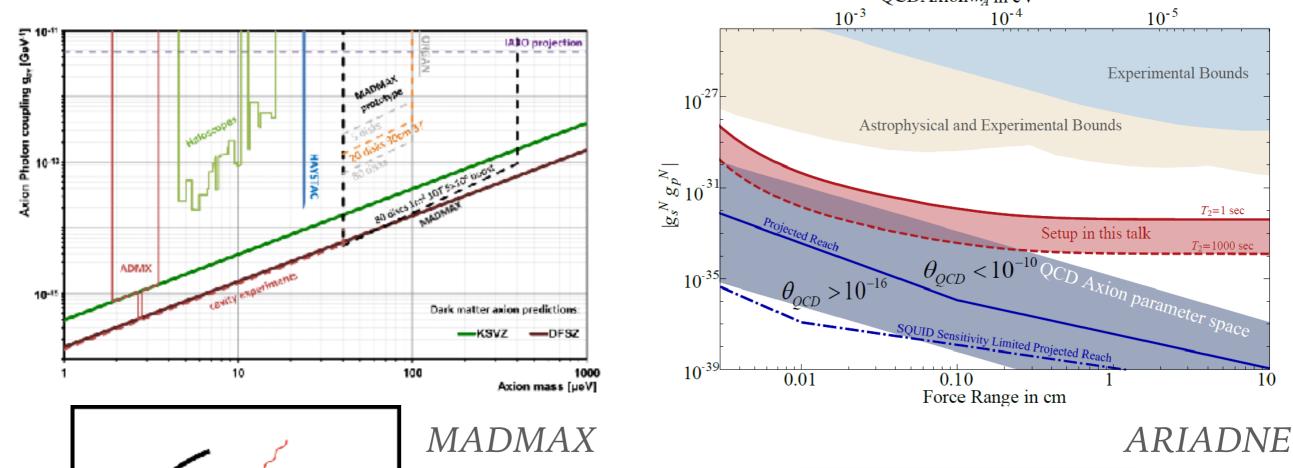
Much broader range of experiments to probe the entire mass range

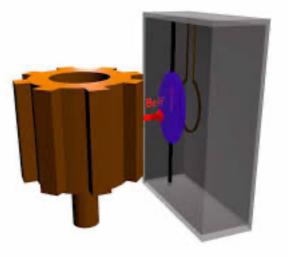


Challenge: remains to be seen how real constraints evolve

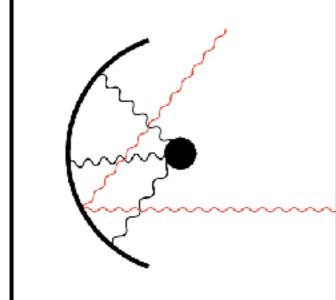
#### **PROMISE: IDEAS FOR AXION DETECTION HAS UNDERGONE A RENAISSANCE**

#### For example, dish antenna or spin-dependent force

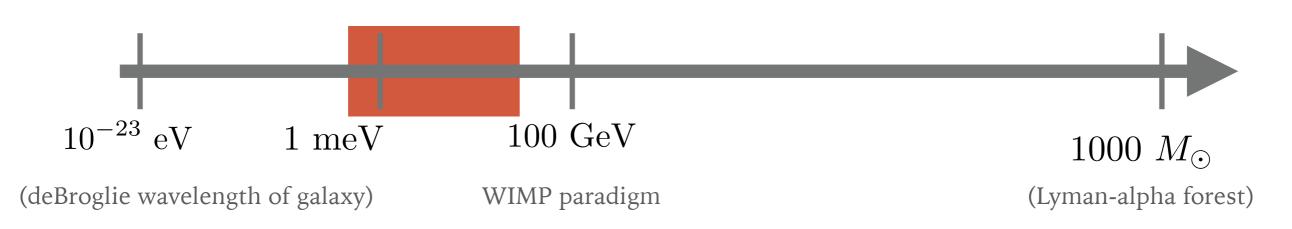




QCDAxion  $m_A$  in eV



### DARK MATTER DETECTION: A FULL COURT PRESS



- Between 1 meV and 1 eV, dark matter behaves like a classical wave, but cavity techniques will be challenging. Detection may be possible via excitation of "collective modes" = phonons
- Intermediate mass range above an eV where observation via particle interactions with SM is still highly motivated though not detectable with traditional WIMP experiments
- Arise generically in top-down constructions hidden sector/ valley paradigm

#### DARK MATTER DETECTION: A FULL COURT PRESS

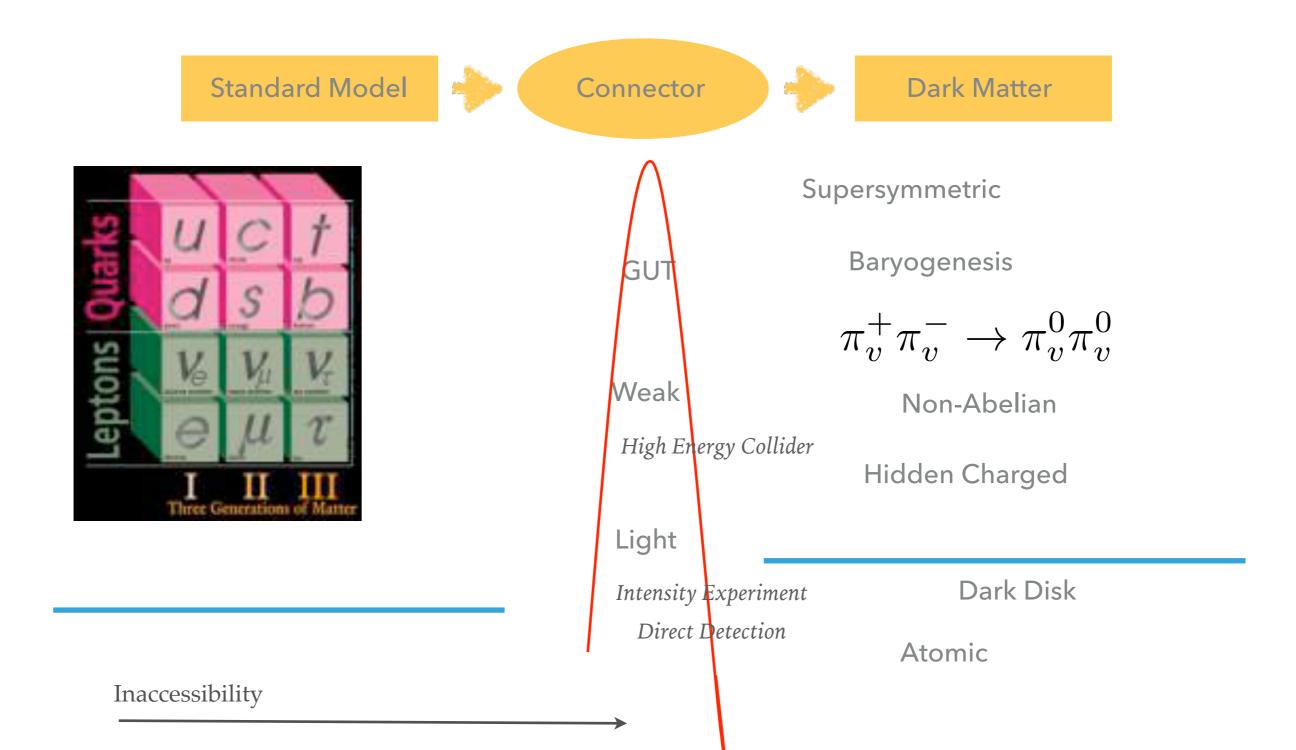


- Dark sector dynamics are complex and astrophysically relevant.  $\sigma_{str} \simeq \frac{4\pi\alpha_s^2}{M^2} \simeq 10^{-24} \text{ cm}^2 \left(\frac{1 \text{ GeV}}{M}\right)^2$
- Abundance may still be set by (thermal) population from SM sector

$$\sigma_{wk} v_{fo} \simeq \frac{g_{wk}^4 \mu_{XT}^2}{4\pi m_Z^4} \frac{c}{3} \simeq 10^{-24} \frac{\text{cm}^3}{\text{s}} \left(\frac{100 \text{ GeV}}{M}\right)^2$$

### HIDDEN SECTOR DARK MATTER

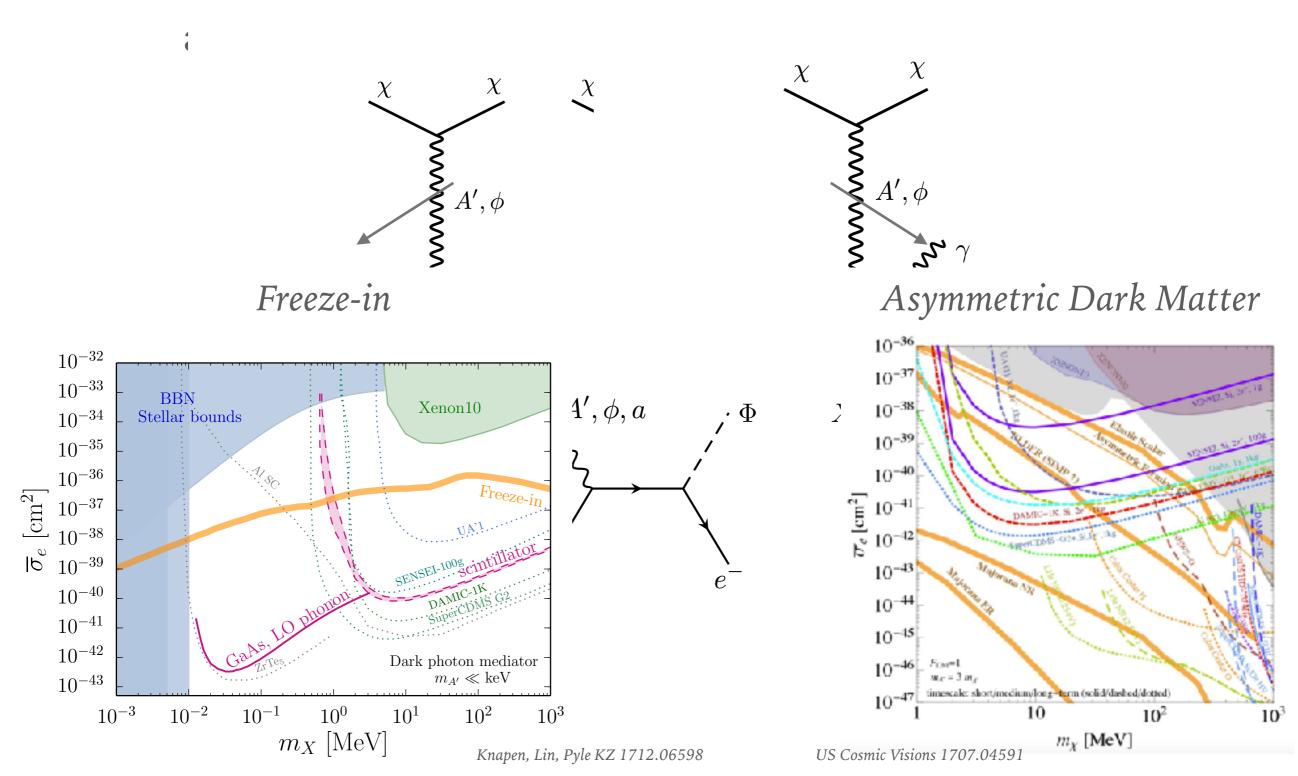
#### Scale of connector sector fixes terrestrial experiment



Energy

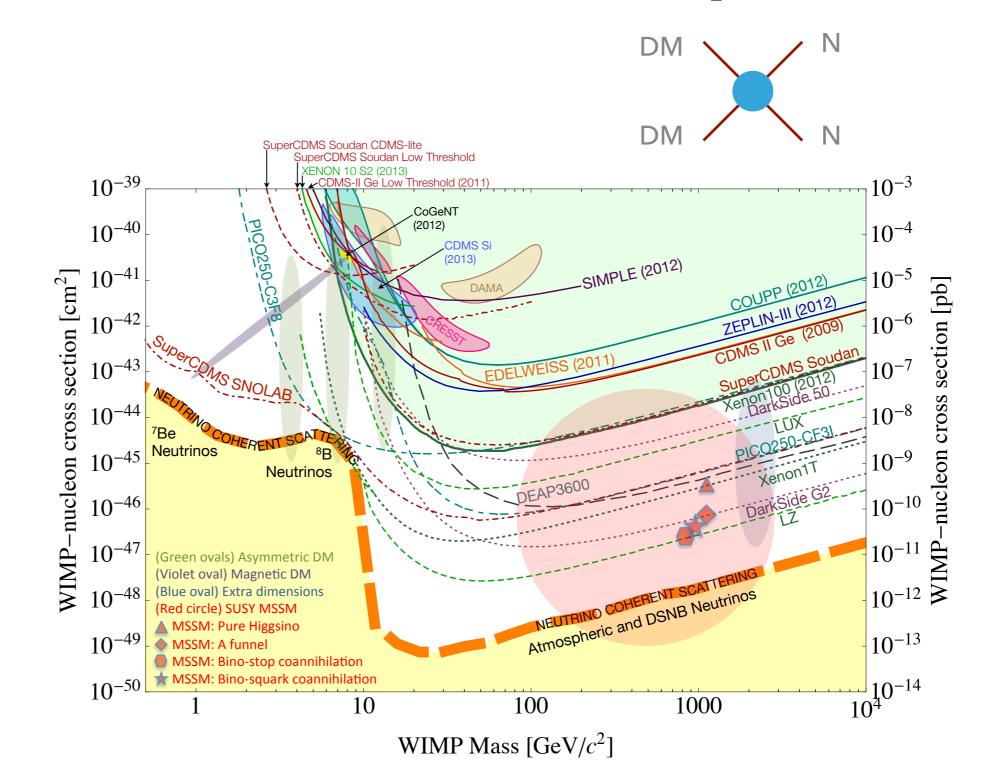
#### HIDDEN DM ABUNDANCE AS A GUIDE

#### If DM abundance is related to its coupling to the SM in



#### TOWARDS HIDDEN SECTOR DARK MATTER

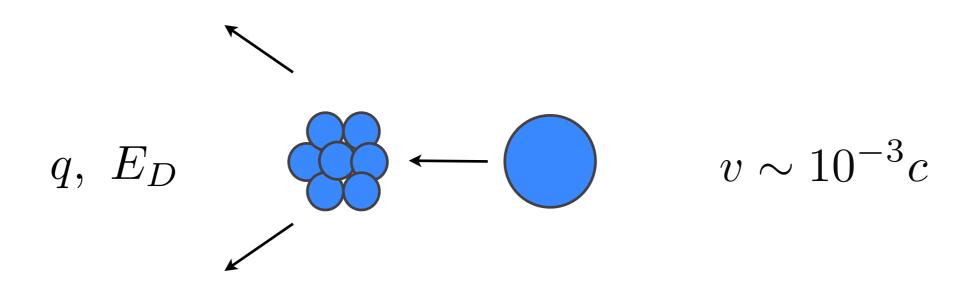
#### Developments in condensed matter make this possible



???

#### LIGHTER TARGETS FOR LIGHTER DARK MATTER

 Nuclear recoil experiments; basis of enormous progress in direct detection

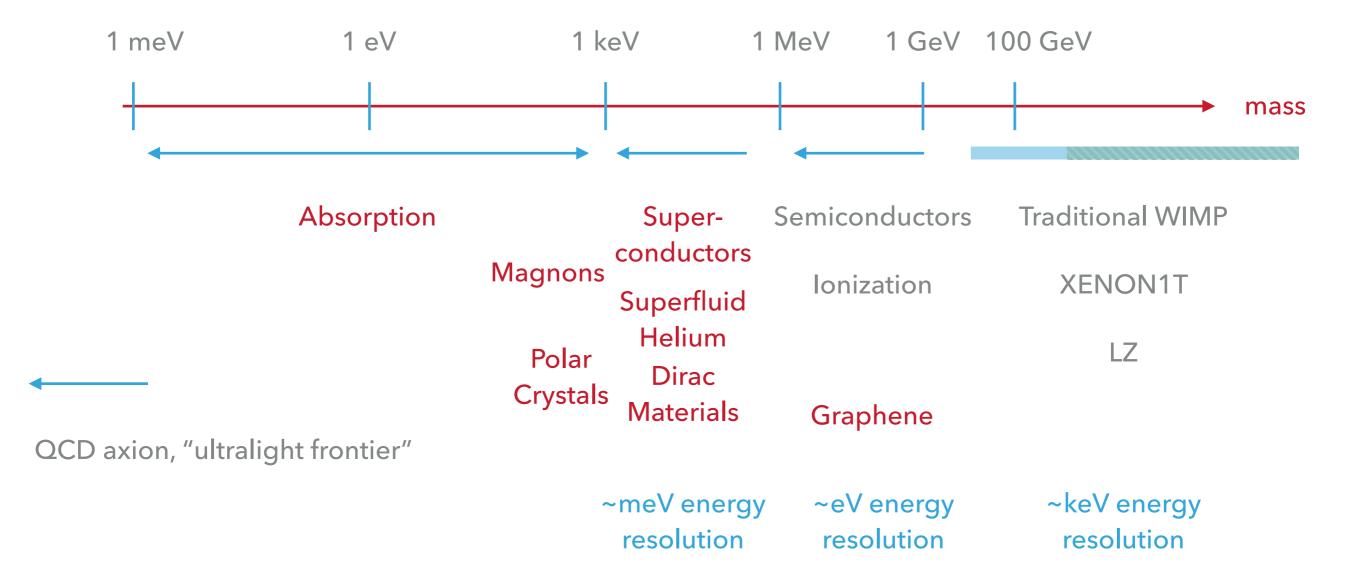


 $v \sim 300 \text{ km/s} \sim 10^{-3} c \implies E_D \sim 100 \text{ keV}$ 

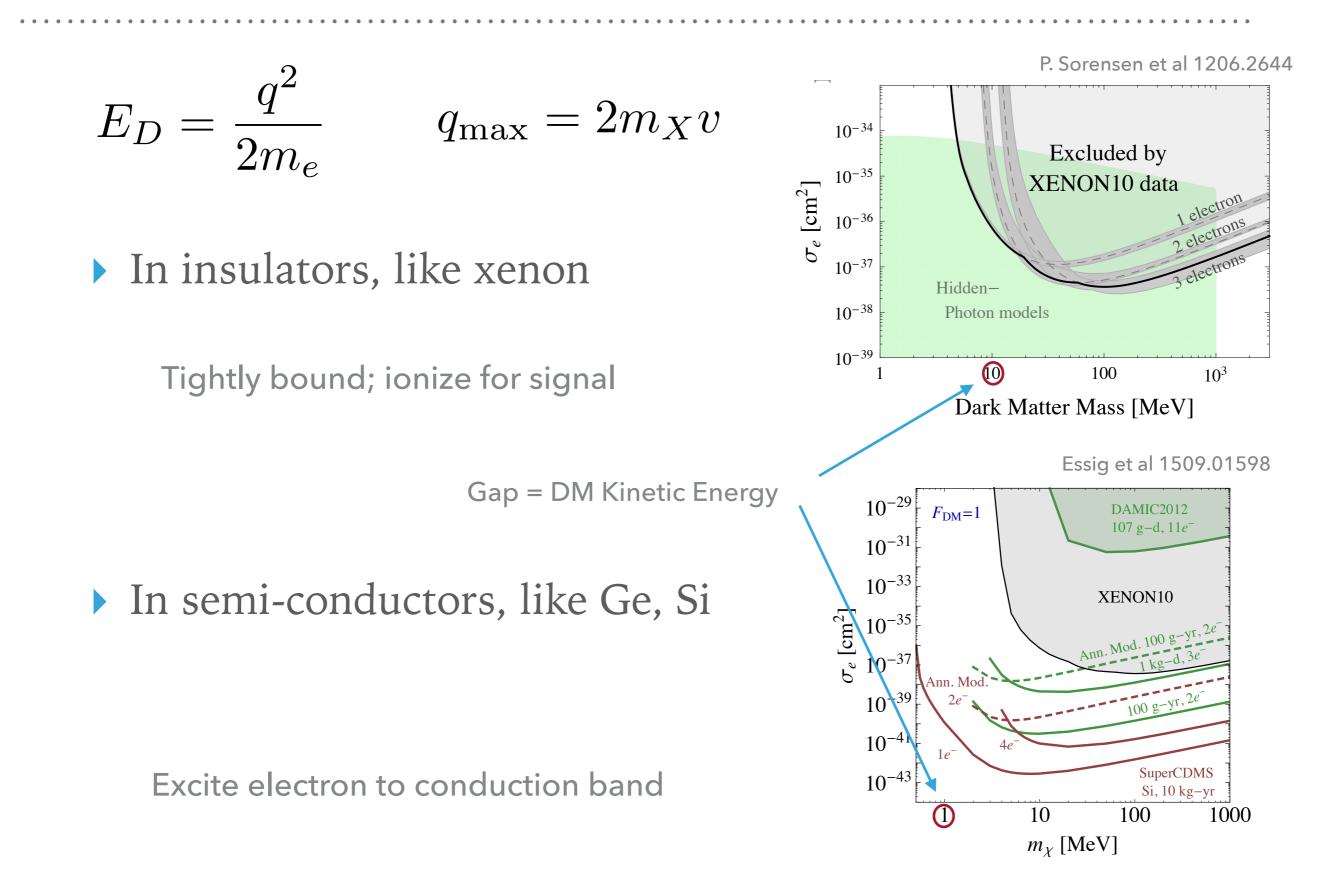
$$E_D = \frac{q^2}{2m_N} \qquad \qquad q_{\max} = 2m_X v$$

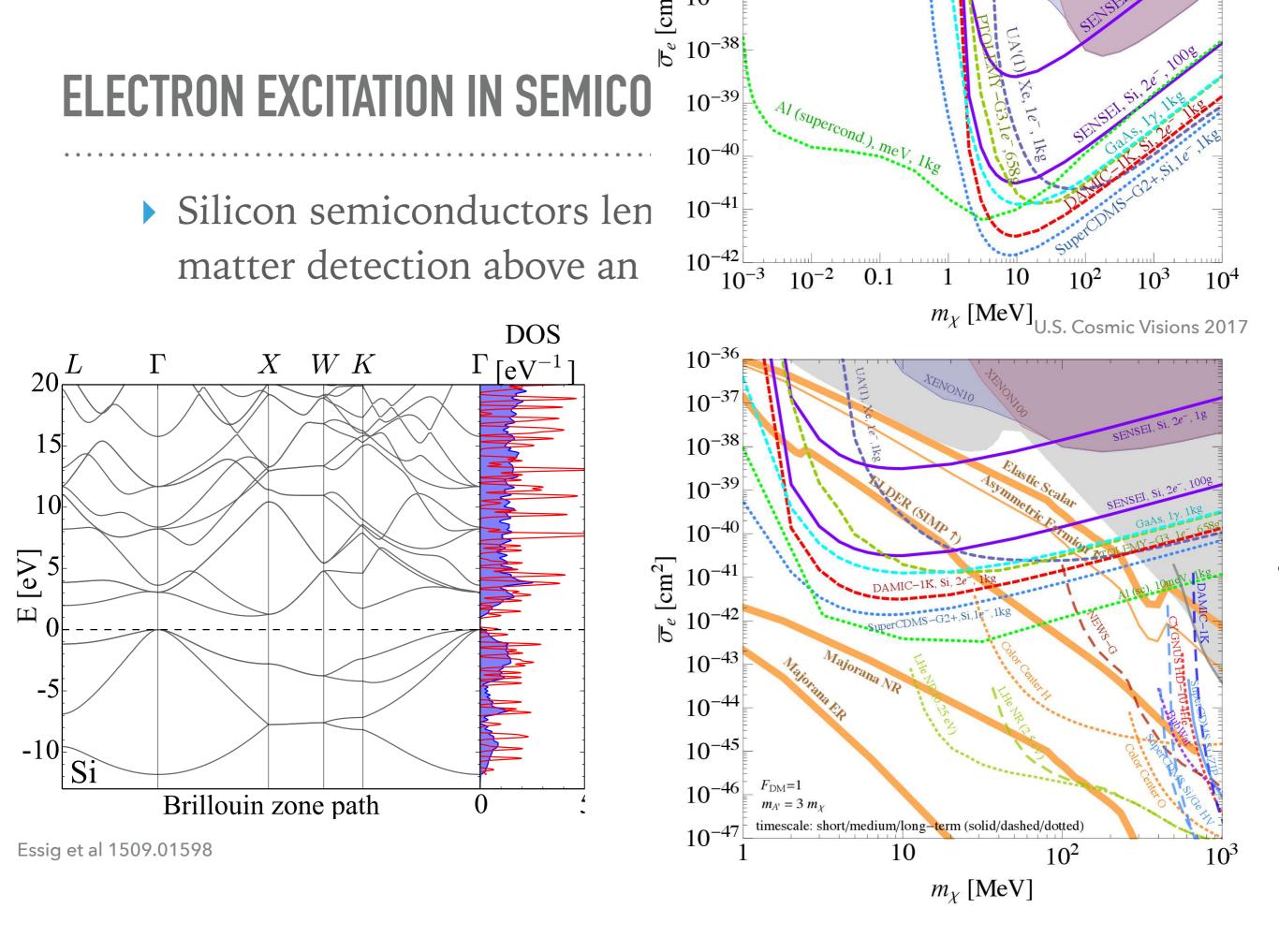
### **PROMISE: NEW DETECTION PARADIGMS FOR LIGHT DARK MATTER**

#### Experimental Panorama



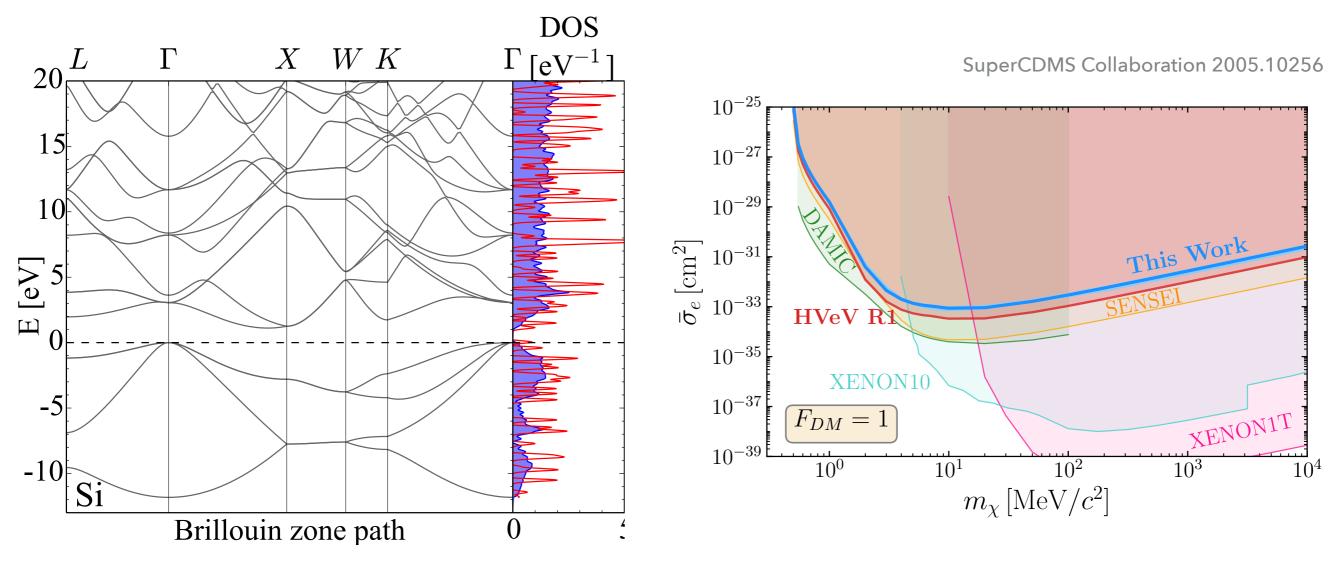
#### LIGHTER TARGETS FOR LIGHTER DARK MATTER





### ELECTRON EXCITATION IN SEMICONDUCTORS AND NOBLE LIQUIDS

 Silicon semiconductors lend themselves well to light dark matter detection above an MeV



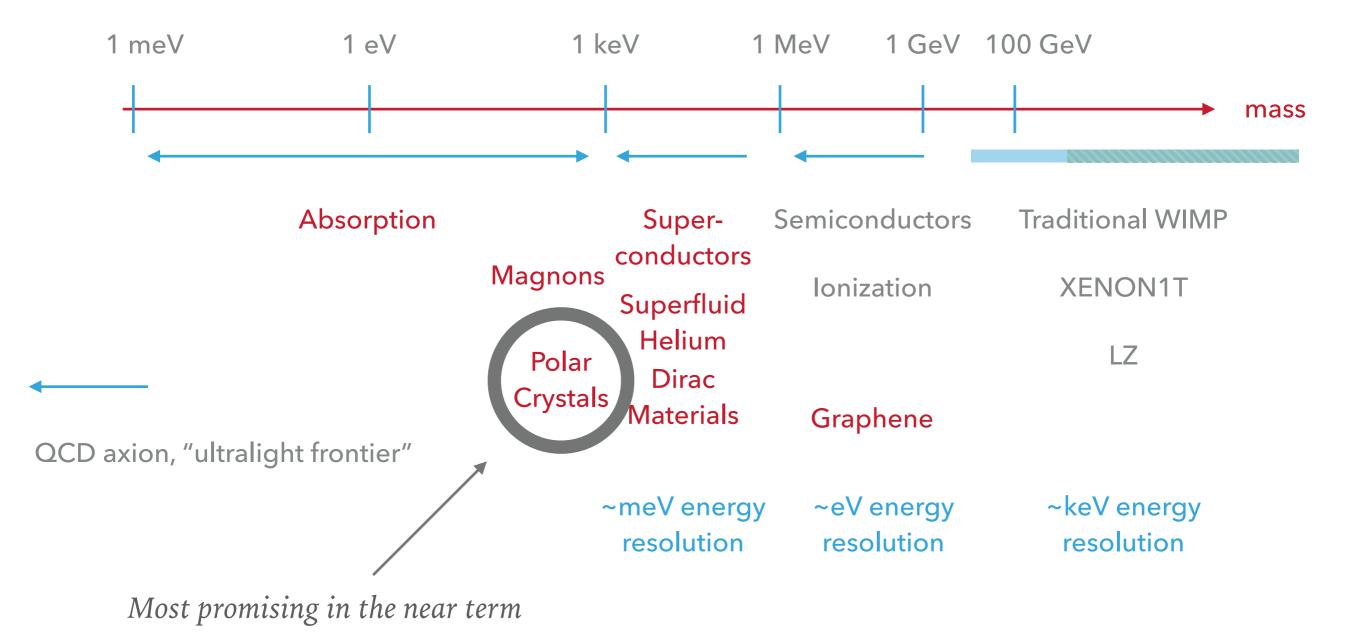
Essig et al 1509.01598

### **EXCITING COLLECTIVE MODES**

- Once DM drops below an MeV, its deBroglie wavelength is longer than the inter particle spacing in typical materials
- Therefore, coupling to collective excitations in materials makes sense!
- Collective excitations = phonon modes, spin waves (magnons)
- Can be applied to just about any material
- (partial) calculations exist for superfluid helium, semiconductors, superconductors, polar materials
- Details depend on
  - 1) nature of collective modes in target material
  - 2) nature of DM couplings to target

### **PROMISE: NEW DETECTION PARADIGMS FOR LIGHT DARK MATTER**

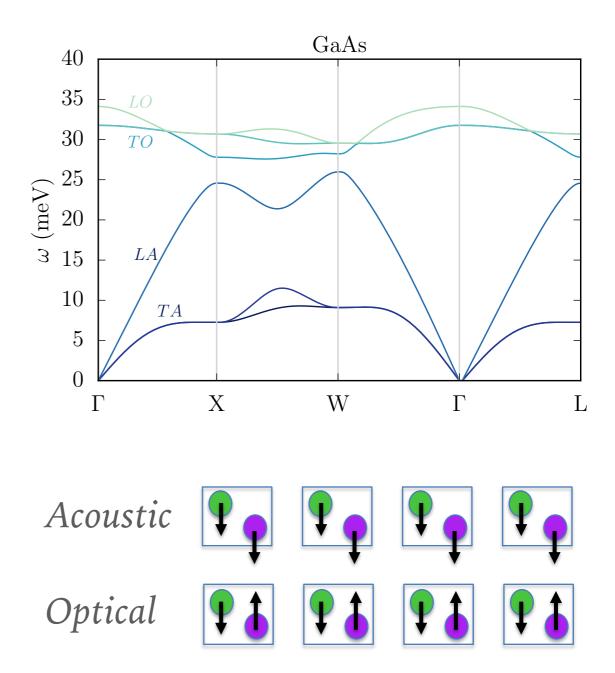
#### Experimental Panorama



## NATURE OF COLLECTIVE MODES

- Number of collective modes:
   3 x number of ions in unit cell
- 3 of those modes describe in phase oscillation — acoustic phonons — and have a translation symmetry implying gapless dispersion
- When these gapped modes result from oscillations of more than one type of ion, it sets up an oscillating dipole
- Polar Materials

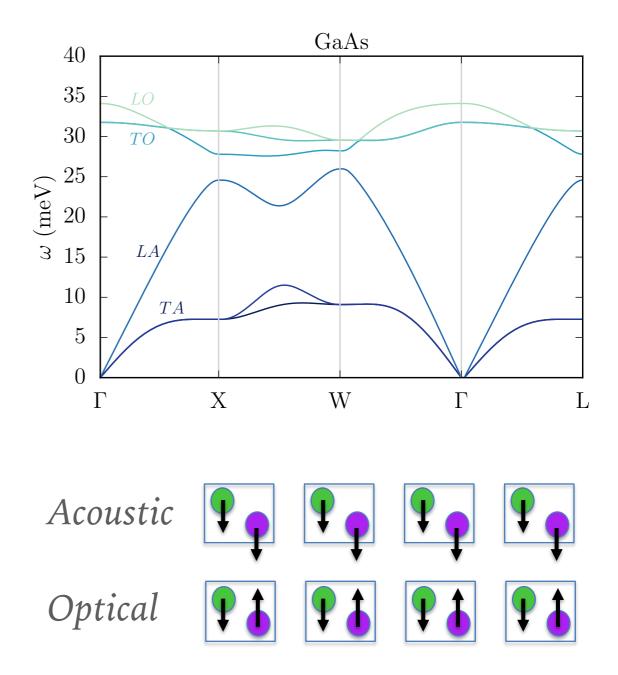




## SCATTERING ON COLLECTIVE MODES IN POLAR MATERIALS

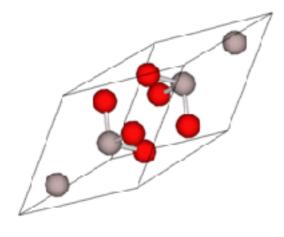
- Gapped optical modes are also ideal for sub-MeV dark matter scattering
  - i) energy spectrum of modes matches dark matter kinetic energy
  - ii) When these gapped modes result from oscillations of more than one type of ion, we have an oscillating dipole





#### KINEMATICS OF COLLECTIVE MODES (2)

#### First element to enter is the kinematics



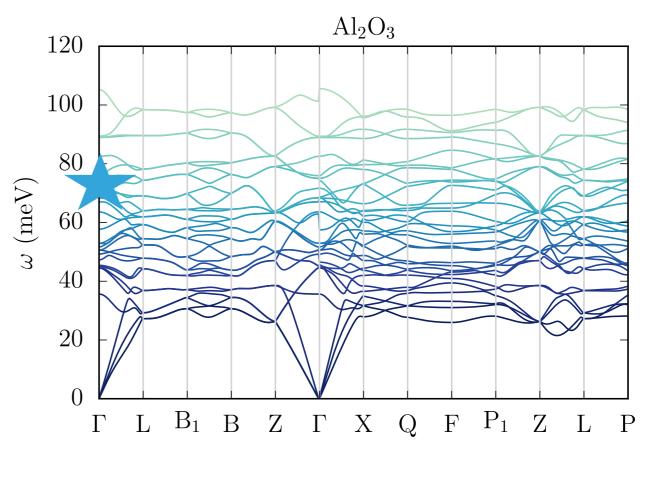
```
E_D \sim v_X q
```

 $c_s \ll v_X$ 

 $E_D \sim c_s q$ 

VS

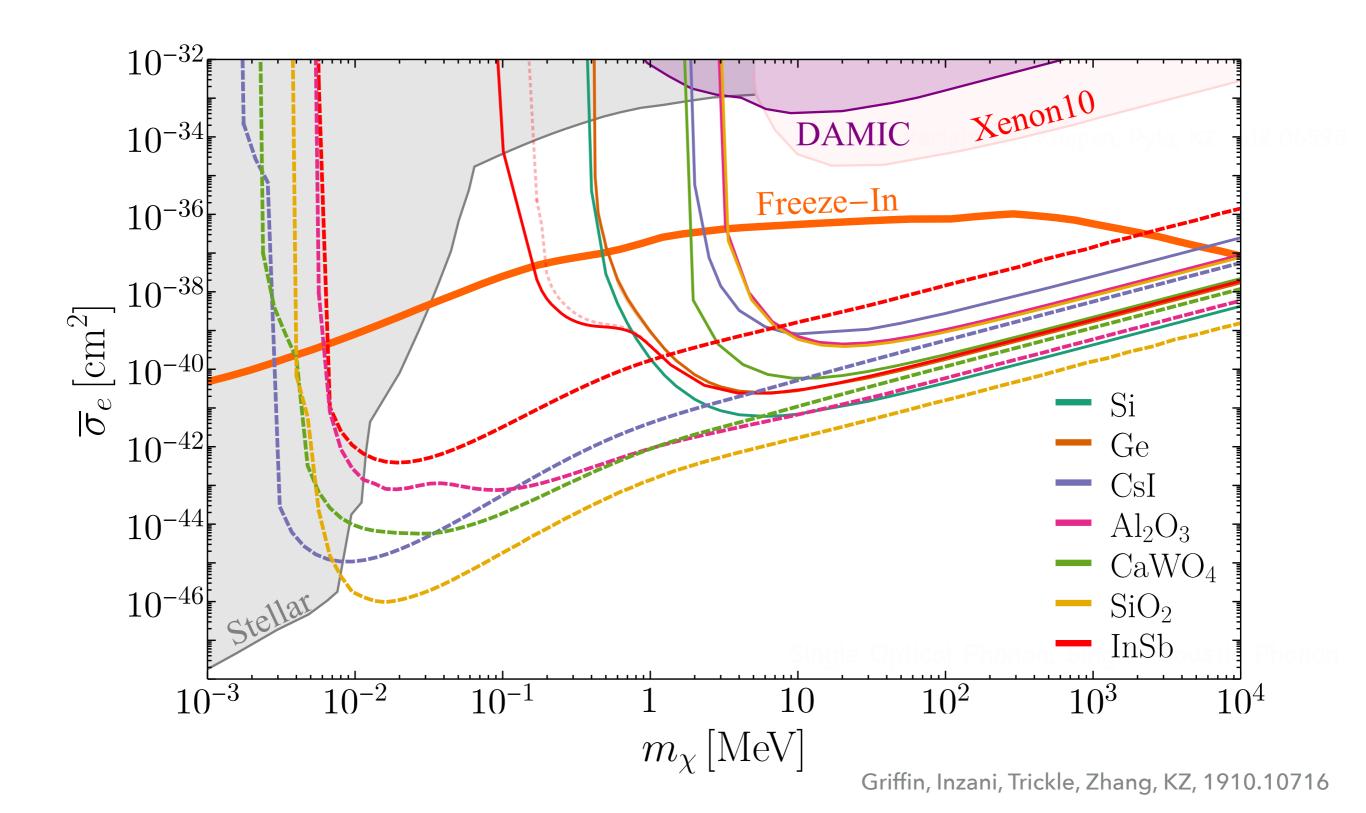
Coupling to gapped modes



Acoustic

Optical

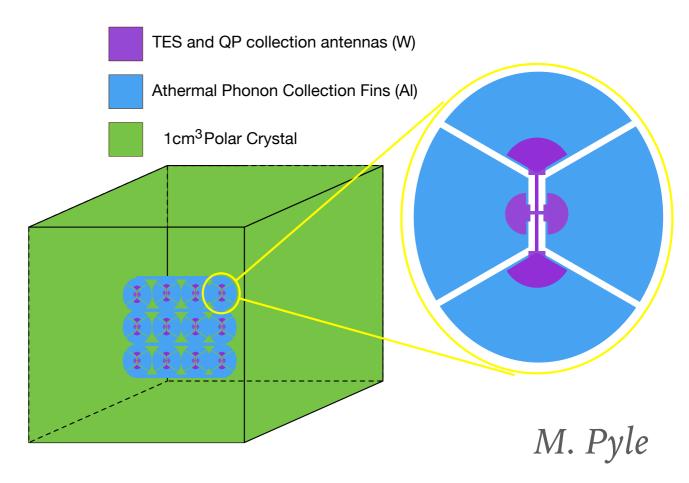
#### **OPTICAL PHONONS IN POLAR MATERIALS (2)**



### **COMMON R&D PATH**

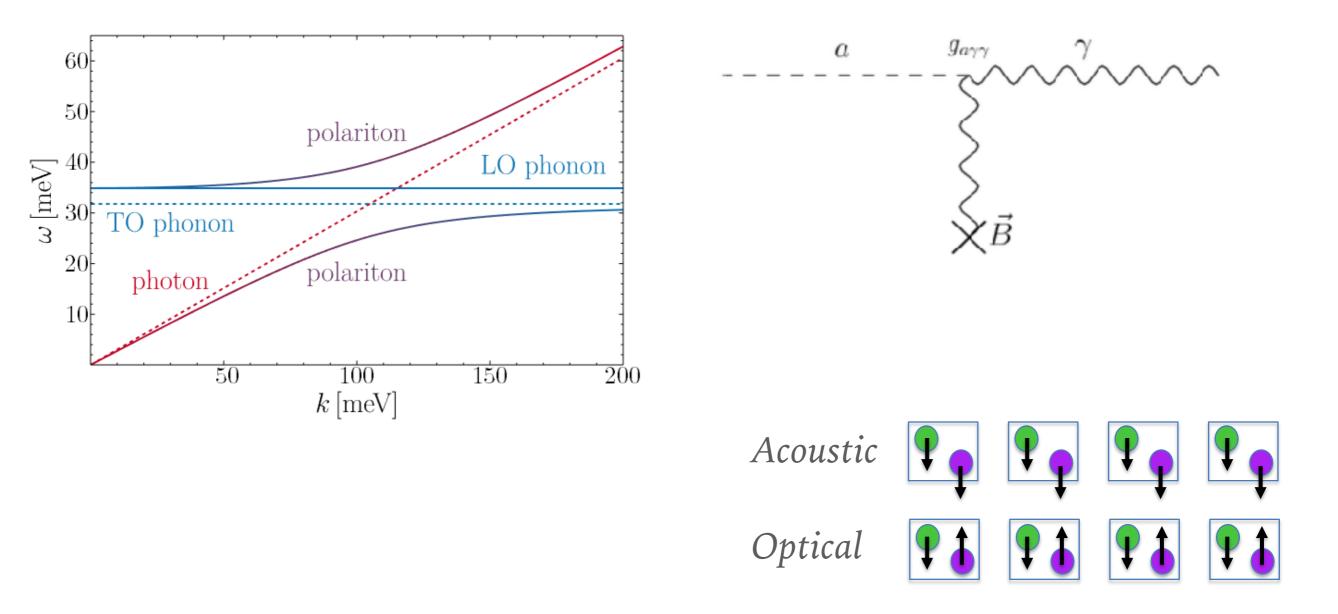
- Sensor can be coupled to multiple targets
- Zero-field read-out of phonons
- Funded for R&D by
   DoE Dark Matter New
   Initiatives
- For a polar crystal target

   Sub-eV Polar
   Interactions Cryogenic
   Experiment (SPICE)



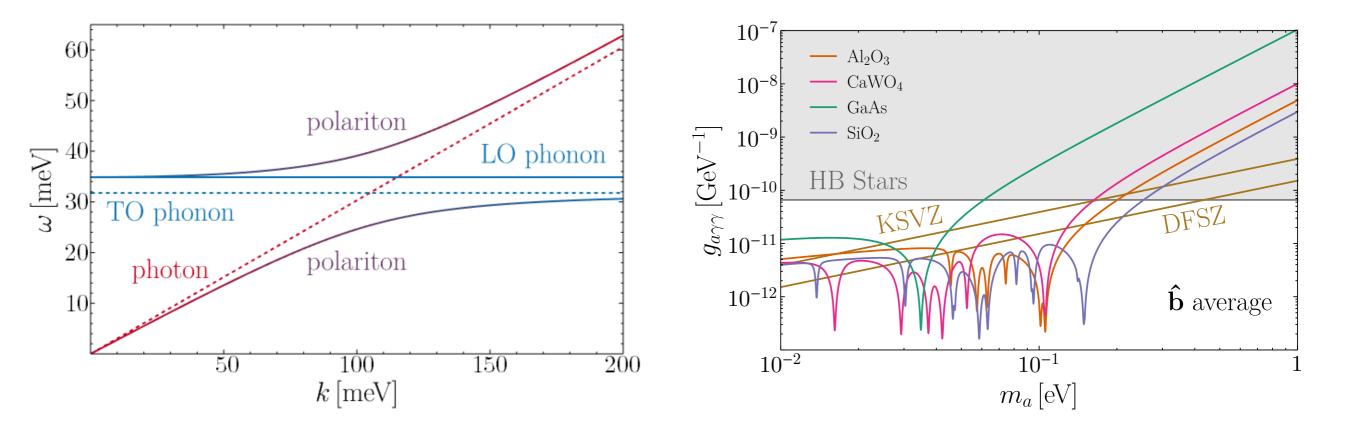
#### **AXION EXCITATION OF PHONON-POLARITONS**

The photon mixes with these collective modes at very low momentum transfer

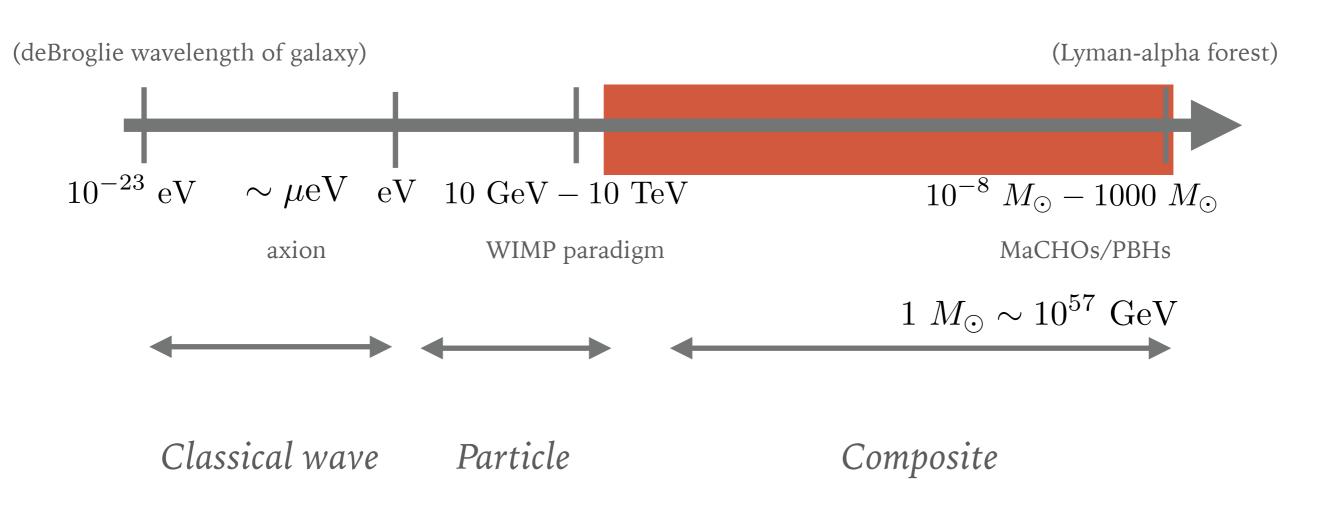


#### **AXION EXCITATION OF PHONON-POLARITONS**

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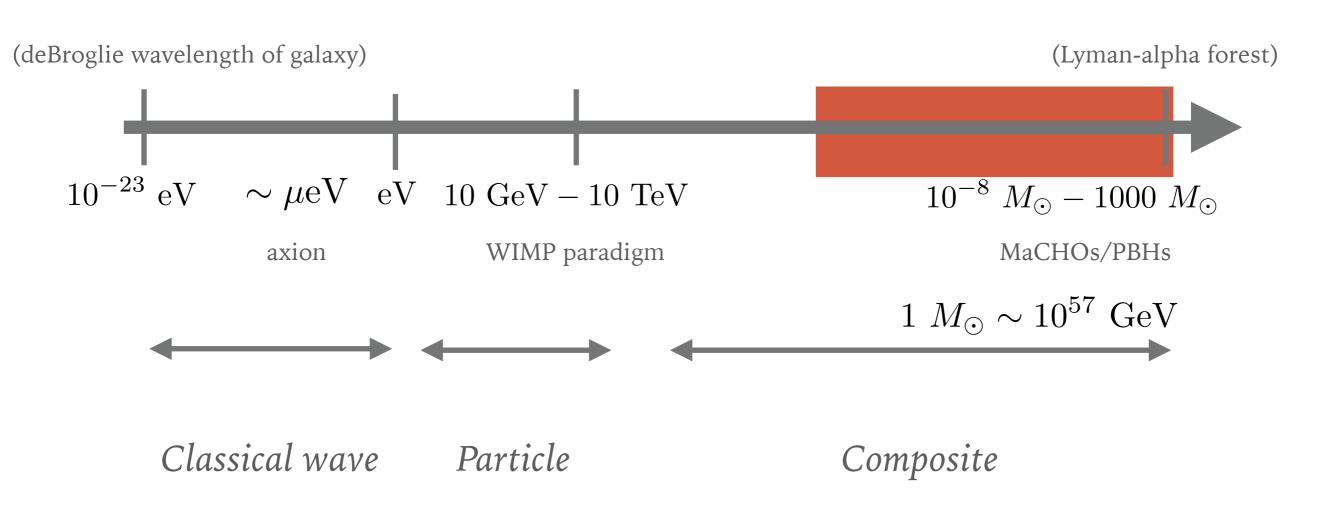


### DARK MATTER DETECTION: A FULL COURT PRESS



- Heavier dark matter: setting relic abundance through interactions with Standard Model is challenging
- At heavier masses, detection through Standard Model interactions is (generally) not motivated by abundance

### DARK MATTER DETECTION: A FULL COURT PRESS



- Look for gravitational means to detect structure
- Above  $10^{-13} M_{\odot}$  e.g. pulsar timing may be effective
- Project of the (far) future to use laboratory clocks to detect small gravitational redshift effects

### **GRAVITATIONAL EFFECTS OF DARK MATTER SUBSTRUCTURE**

• Accurate clocks and transiting objects — the time-of-arrival of a pulse is stable. Deviations can signal transiting object.

$$\phi(t) = \phi_0 + \nu t + \frac{1}{2}\dot{\nu} t^2 + \frac{1}{6}\ddot{\nu} t^3 + \dots$$

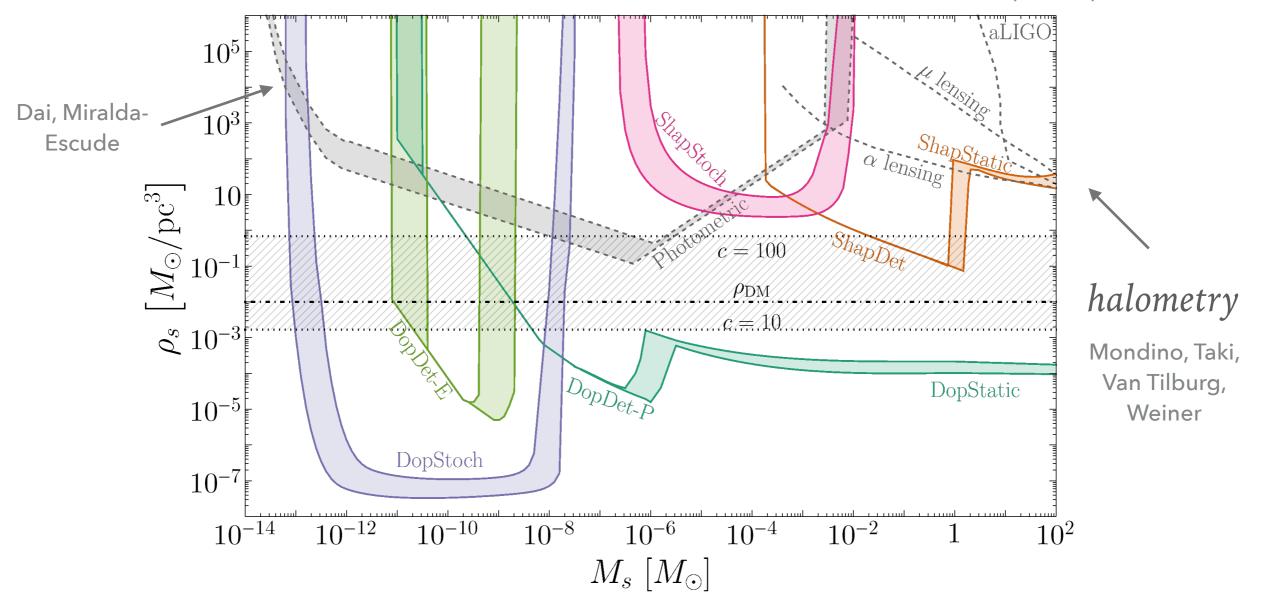
Earth

Pulsar

- Principle can be applied to many systems with accurate clocks
- Transiting clump can accelerate earth or pulsar (Doppler) or change the potential along line of sight (Shapiro)

### **GRAVITATIONAL EFFECTS OF DARK MATTER SUBSTRUCTURE**

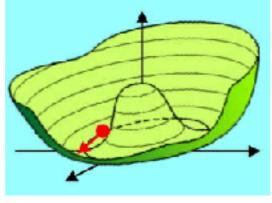
Ramani, Trickle, KZ 2005.03030



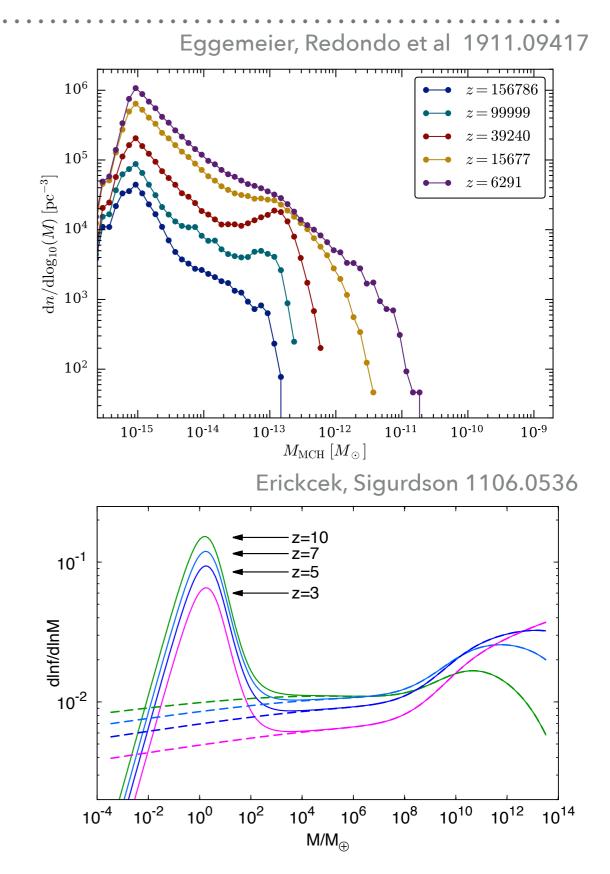
 Working with Nanograv-developed (PTA collaboration) machinery to forecast more realistic prospects for detection

### MODELS WITH ENHANCED SMALL SCALE POWER

 Axion models, where PQ symmetry breaks after inflation



 Periods of early matter domination can lead to growth of structure on small scales

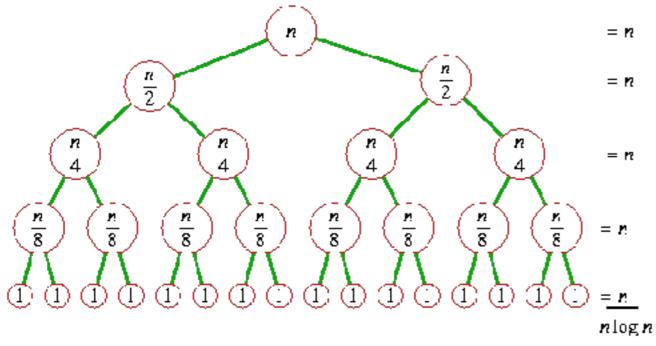


#### THE CHALLENGE

- Now is not the time for narrowing our search for Invisibles; the playing field is still wide open
- Fortunately the Identification of Dark Matter has spurred ideas for many new experimental and observational efforts, on all avenues
- Lack of evidence for WIMP has given rise to renewed interest in axion, including more robust theoretical predictions and new ideas for probes
- New ideas for hidden sector/valley dark matter and probes for such light dark matter states
- New opportunities to search for dark matter substructure over the next decade

#### THE OUTLOOK

• We are not without tools!



- The universe is dominated by invisibles!
- " WIMP or (axion)
  - How to be ready for anything? Hidden Sectors

How do I search for these new candidates?