

The EW (weak) corrections for DY Z->ll comparing/benchmarking: Powheg_ew; wt^{EW} (IBA approach); MCSANC; DYTURBO

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with contributions from **F. Piccinini** (Powheg_ew),
S. Bondarenko&L.Kalinovskaya (MCSANC), **A. Armbruster** (DYTURBO)

- Strategy for comparison
- EW schemes
- Results: EW LO, NLO, NLO+HO

Some slides repeated from „LHCC Precision WG meeting”, 25.09.2018

Strategy for comparison: be pragmatic

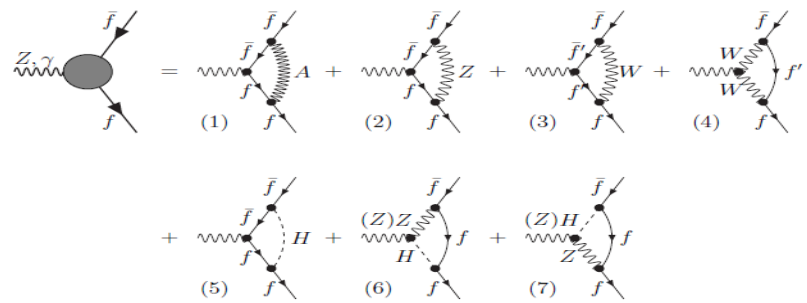
- Scope:
 - **Genuine EW and lineshape corrections** to Drell-Yan production at NLO QCD.
 - Three EW LO schemes chosen to allow for straightforward interpretation of results. We tuned EW LO parameters, otherwise out-of-the-box.
 - The highest available corrections in a given approach used.
 - QED FRS/ISR not included here.
- Observables:
 - Lineshape (cross-section) and forward-backward asymmetry A_{FB} in the full phase-space.
 - **Compared ratios or absolute differences** between different EW LO schemes and/or between NLO, NLO+HO predictions within each EW scheme and same MC generator. Allows to **minimize sensitivity to QCD details**.
- Goals:
 - Check if reweighting with **wt^{EW} (TauSpinner)** works for NLO QCD MC's. Compared distributions at EW LO (**DYTURBO, Powheg_ew**).
 - Establish how consistent are predictions between different EW schemes with EW NLO corrections (**Powheg_ew, MCSANC**).
 - Establish how consistent are EW NLO+HO corrections of **Dizet 6.21 form-factors** implemented in **wt^{EW}** and those of **Powheg_ew**.

Genuine EW and lineshape corrections

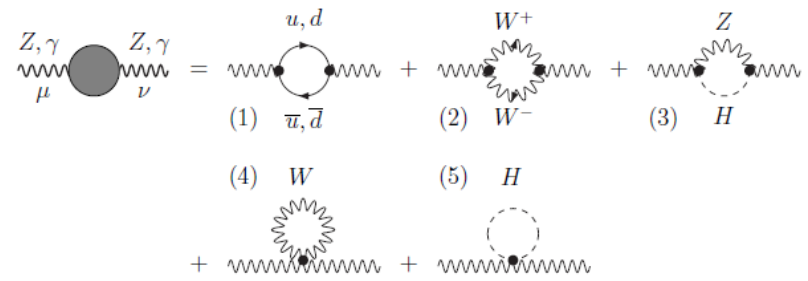
Gauge-invariant set of diagrams.

For Improved Born Approximation (IBA) approach calculated as form-factor corrections to couplings, propagators and masses.

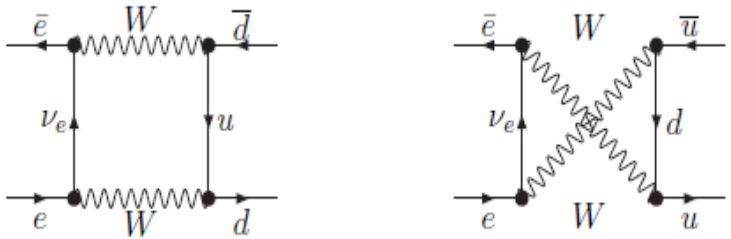
Zff and γ ff vertices



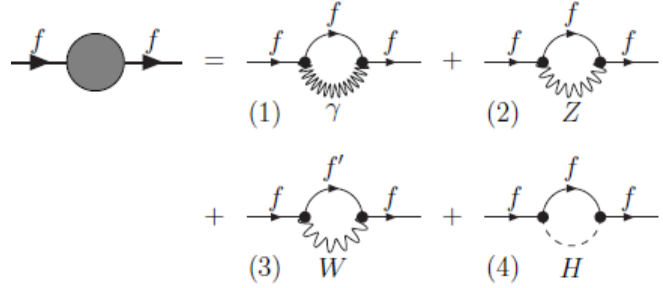
Bosonic self-energies



WW, ZZ boxes (shown only WW diagrams)



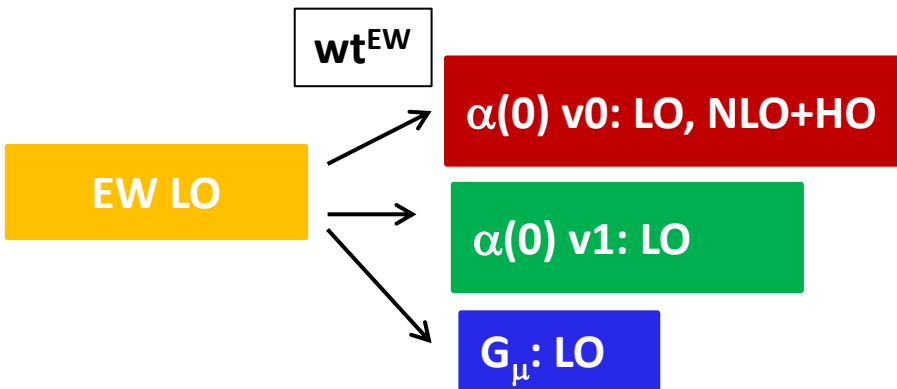
Fermionic self-energies



What we have so far

PowhegZj: QCD NLO for Z+j

wt^{EW} : TauSpinner + Dizet 6.21



Powheg_ew: QCD LO for Z

$\alpha(0) v0: LO$

$\alpha(0) v1: LO, NLO, NLO+HO$

$G_\mu: LO, NLO, NLO+HO$

DYTURBO: QCD LO, NLO for Z

$\alpha(0) v0: LO$

$\alpha(0) v1: LO$

$G_\mu: LO$

MCSANC: QCD LO for Z

$\alpha(0) v1: LO, NLO$

$G_\mu: LO, NLO$

EW corrections with event weight wt^{EW}

Reweighting possible because of Drell-Yan factorisation properties,
Mirkes et al. arXiv:9406381.

Method follows technique developed for **TauSpinner** program (for LHC!),
arXiv:1201.0117; arXiv1802.05459

Define per event electroweak weight $wt^{EW} = \sigma_{\text{Born}}^{\text{new}} / \sigma_{\text{Born}}^{\text{old}}$

$$wt^{EW} = \frac{d\sigma_{\text{Born}+EW}(x_1, x_2, \hat{s}, \cos\theta, s_W^2)}{d\sigma_{\text{Born}}(x_1, x_2, \hat{s}, \cos\theta, s_W^2)}$$

$$d\sigma_{\text{Born}}(x_1, x_2, \hat{s}, \cos\theta^*, s_W^2) = \sum_{q_f, \bar{q}_f} [f^{q_f}(x_1, \dots) f^{\bar{q}_f}(x_2, \dots) d\sigma_{\text{Born}}^{q_f \bar{q}_f}(\hat{s}, \cos\theta, s_W^2) \\ + f^{q_f}(x_2, \dots) f^{\bar{q}_f}(x_1, \dots) d\sigma_{\text{Born}}^{\bar{q}_f q_f}(\hat{s}, -\cos\theta, s_W^2)]$$

$x_1, x_2, \cos\theta$ (symmetrised)
calculated using 4-momenta
of outgoing leptons;
asymmetry in sign of $\cos\theta$
from weighted average
over PDFs

Allows to reweight MC event generated between different EW LO
scheme and to **Improved Born Approximation** in EW scheme used
for form-factors calculation. See more in **arXiv:1808.08616**

Spin amplitude: EW Improved Born (IBA)

$$\mathcal{A}^{Born+EW} = \frac{\alpha}{s} \left\{ \begin{aligned} & [\bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu u] \cdot (q_\ell \cdot q_f) \Gamma_{V\Pi} \chi_\gamma(s) \\ & + [\bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu u \cdot (v_\ell \cdot v_f \cdot vv_{\ell f}) + \bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu \gamma^5 u \cdot (v_\ell \cdot a_f) \\ & + \bar{u}\gamma^\mu \gamma^5 v g_{\mu\nu} \bar{v}\gamma^\nu u \cdot (a_\ell \cdot v_f) + \bar{u}\gamma^\mu \gamma^5 v g_{\mu\nu} \bar{v}\gamma^\nu \gamma^5 u \cdot (a_\ell \cdot a_f)] \cdot Z_{V\Pi} \chi_Z(s) \end{aligned} \right\}$$

$$\chi_\gamma(s) = 1$$

$$\chi_Z(s) = \frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} \cdot \frac{s}{s - M_Z^2 + i \cdot \Gamma_Z \cdot M_Z}$$

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot s_W^2) \cdot K_\ell(s, t) / \Delta$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot s_W^2) \cdot K_f(s, t) / \Delta$$

$$a_\ell = (2 \cdot T_3^\ell) / \Delta$$

$$a_f = (2 \cdot T_3^f) / \Delta$$

$$\Delta = \sqrt{16 \cdot s_W^2 \cdot (1 - s_W^2)}$$

$$Z_{V\Pi} = \rho_{e,f}(s, t)$$

$$\Gamma_{V\Pi} = \frac{1}{2 - (1 + \Pi_{\gamma\gamma}(s))}$$

**EW form-factors , functions of (s,t)=(m_{ll}, cosθ)
Calculated with Dizet 6.21 library.**

Vacuum polarisation corrections, used low-energy experiment input.

Warning: problem for analytic continuation

$$vv_{\ell f} = \frac{1}{v_\ell \cdot v_f} \left[(2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot s_W^2 \cdot K_f(s, t) (2 \cdot T_3^\ell) - 4 \cdot q_f \cdot s_W^2 \cdot K_\ell(s, t) (2 \cdot T_3^f) + (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2) K_{\ell f}(s, t) \right] \frac{1}{\Delta^2}$$

EW LO schemes

SM fundamental relation used to calculate EW parameters in different EW LO schemes, on-mass-shell definition.

EW scheme: G_μ, α, M_Z α, M_W, M_Z G_μ, M_W, M_Z

Parameter	$\alpha(0) \text{ v0}$	$\alpha(0) \text{ v1}$	G_μ
M_Z	91.1876 GeV	91.1876 GeV	91.1876 GeV
Γ_Z	2.4952 GeV	2.4952 GeV	2.4952 GeV
Γ_W	2.085 GeV	2.085 GeV	2.085 GeV
α	1/137.03599	1/137.03599	1/132.23323
G_μ	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1254734 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$
M_W	80.93886 GeV	80.385 GeV	80.385 GeV
s_W^2	0.2121517	0.2228972	0.2228972
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha}$	1.0	1.0	1.0

PowhegZj
91.1876 GeV
2.4952 GeV
2.085 GeV
1/128.88859
$1.16638 \cdot 10^{-5} \text{ GeV}^{-2}$
79.958 GeV
0.2311300
1.0

$$s_W^2 = 1 - m_W^2/m_Z^2$$

$$G_\mu = \frac{\pi\alpha}{\sqrt{2}M_W^2 s_W^2}$$



MC events used for reweighting

EW schemes: $\alpha(0) \text{ v0}, \alpha(0) \text{ v1}$ – same value of α
 $G_\mu, \alpha(0) \text{ v1}$ – same value of s_W^2

EW LO schemes: pros and cons

- **EW scheme $\alpha(0)$ v0: input $\alpha(0)$, M_Z , G_μ**
 - **Pros:**
 - Precisely measured physics input, **LEP legacy EW scheme**
 - **Cons:**
 - Moderate NLO and HO corrections (few %) calculated theoretically or taken from low-energy measurements ($\alpha_{\text{had}}^{(5)}$)
- **EW scheme $\alpha(0)$ v1: input $\alpha(0)$, M_Z , M_W**
 - **Pros:**
 - Moderate NLO corrections (few %), small HO corrections (<1%)
 - **Cons:**
 - Input M_W with ± 15 MeV uncertainties, requires shifting G_μ far from its measured value.
- **EW scheme G_μ : input G_μ , M_Z , M_W**
 - **Pros:**
 - Small NLO (1%) and very small HO (0.2%) corrections
 - **Cons:**
 - Input M_W with ± 15 MeV uncertainties, requires two definitions for em coupling: $\alpha(0)$ for ISR/FSR/IFI and α_{G_μ} for matrix elements.

Goal is to see level of agreement between predictions calculated in three EW schemes, after including EW NLO+HO corrections?

EW LO schemes: details

EW schemes: come with „on-shell” or „pole” definitions!

Table 44: The EW parameters used at tree-level EW, with on-mass-shell definition (LEP convention).

Parameter	$\alpha(0) \text{ v0}$	$\alpha(0) \text{ v1}$	G_μ
M_Z	91.1876 GeV	91.1876 GeV	91.1876 GeV
Γ_Z	2.4952 GeV	2.4952 GeV	2.4952 GeV
Γ_W	2.085 GeV	2.085 GeV	2.085 GeV
α	1/137.03599	1/137.03599	1/132.23323
G_μ	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1254734 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$
M_W	80.93886 GeV	80.385 GeV	80.385 GeV
s_W^2	0.2121517	0.2228972	0.2228972
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha}$	1.0	1.0	1.0

Shift:

- -30 MeV for M_Z
- -0.00005 for s_W^2

Scaling

- 0.99906 for α

Table 45: The EW parameters used at tree-level EW, with pole definition of the Z, W masses.

Parameter	$\alpha(0) \text{ v0}$	$\alpha(0) \text{ v1}$	G_μ
M_Z	91.15348 GeV	91.15348 GeV	91.15348 GeV
Γ_Z	2.494266 GeV	2.494266	2.494266 GeV
Γ_W	2.085 GeV	2.085 GeV	2.085 GeV
α	1/137.03599	1/137.03599	1/132.3572336357709
G_μ	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.126555497 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$
M_W	80.91191 GeV	80.35797 GeV	80.35797 GeV
s_W^2	0.21208680	0.22283820939	0.22283820939
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha}$	1.0	1.0	1.0

EW LO schemes: details

Running and fixed Z-boson width in the propagator: taking into account photon-loop corrections to Γ_Z

- Fixed width $\chi_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot M_Z}$.

- Running width (LEP legacy)

$$\chi'_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot s/M_Z}$$



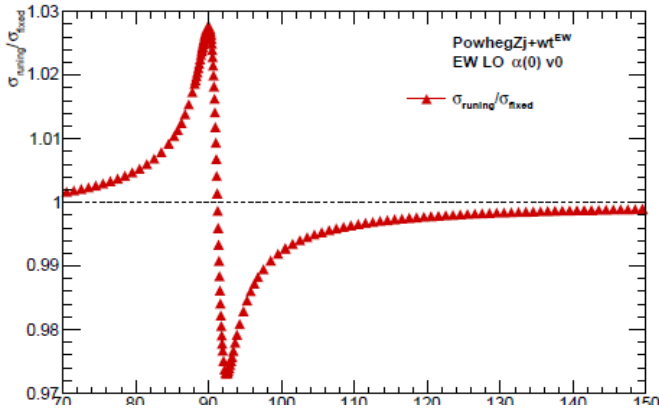
$$\begin{aligned} \chi'_Z(s) &= \frac{1}{s(1 + i \cdot \Gamma_Z/M_Z) - M_Z^2} \\ &= \frac{(1 - i \cdot \Gamma_Z/M_Z)}{s(1 + \Gamma_Z^2/M_Z^2) - M_Z^2(1 - i \cdot \Gamma_Z/M_Z)} \\ &= \frac{1}{(1 + \Gamma_Z^2/M_Z^2) \left(s - \frac{M_Z^2}{1 + \Gamma_Z^2/M_Z^2} + i \cdot \frac{\Gamma_Z M_Z}{1 + \Gamma_Z^2/M_Z^2} \right)} \\ &= N_Z \frac{1}{s - M_Z'^2 + i \Gamma_Z' M_Z'} \\ M_Z' &= \frac{M_Z}{\sqrt{1 + \Gamma_Z^2/M_Z^2}} \\ \Gamma_Z' &= \frac{\Gamma_Z}{\sqrt{1 + \Gamma_Z^2/M_Z^2}} \\ N_Z &= \frac{(1 - i \cdot \Gamma_Z/M_Z)}{(1 + \Gamma_Z^2/M_Z^2)} = \frac{(1 - i \cdot \Gamma_Z'/M_Z')}{(1 + \Gamma_Z'^2/M_Z'^2)} \end{aligned}$$

Both equivalent if redefined parameters m_Z , Γ_Z , N_Z (normalization). Change in the normalisation can (?) be absorbed into G_μ redefinition. In case of „pole” convention (last slide) it was absorbed into α .

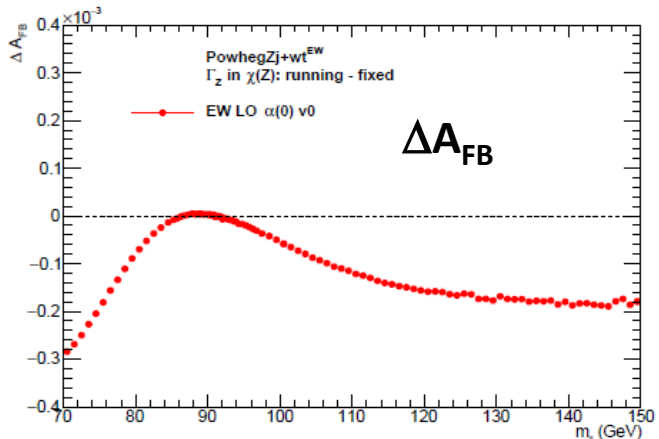
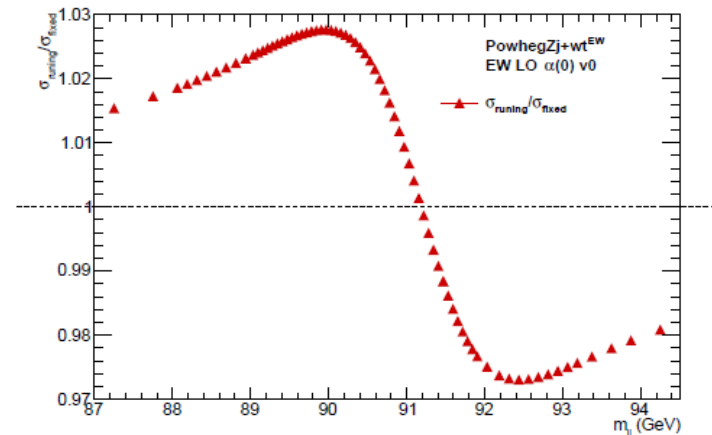
EW LO schemes: details

Nowdays MC's are using fixed-width propagators with on-shell $M_Z = 91.1876$.
 How does it affect predictions, if running- \rightarrow fixed without reparametrizing?
 $\Delta A_{FB} (m_{II}=80-100 \text{ GeV}) = 0.0005$ (thanks to D. Walker for pointing it out)

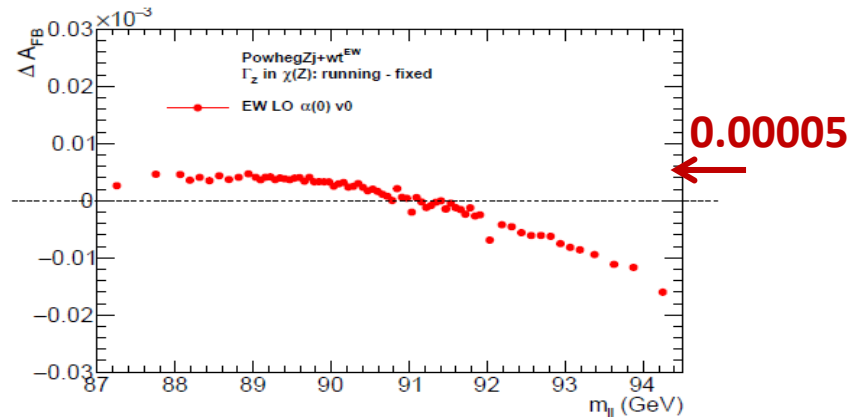
Lineshape: ratio



zoom



zoom



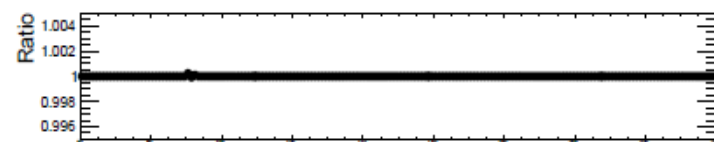
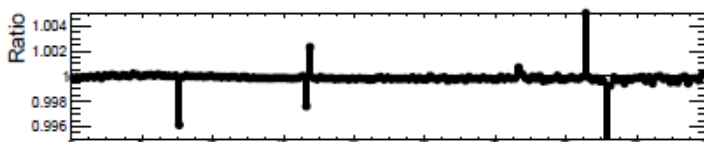
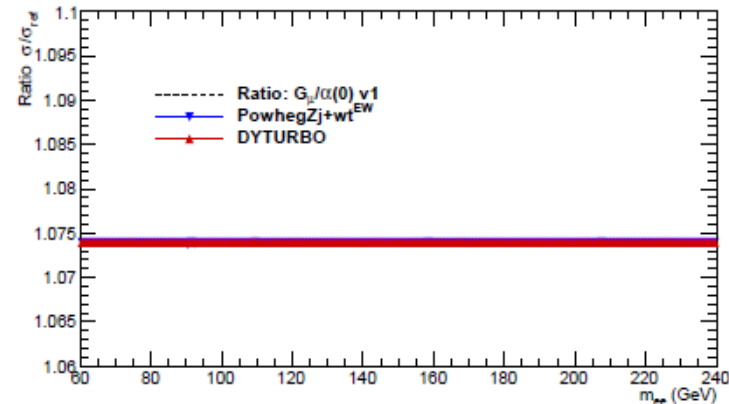
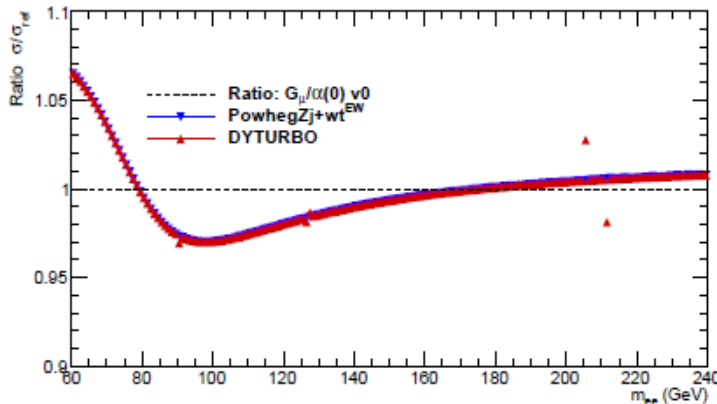
Validating reweighting with wt^{EW} : EW LO

- Ratio of differential cross-sections (lineshapes) driven by relative balance between Z and γ contributions.
- EW $\alpha(0)$ v1 and G_μ schemes chosen as such that ratio of cross-sections is equal to ratio of QED couplings squared.

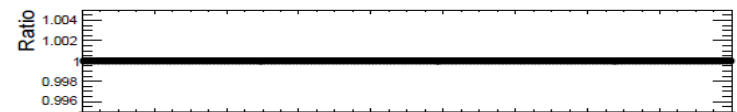
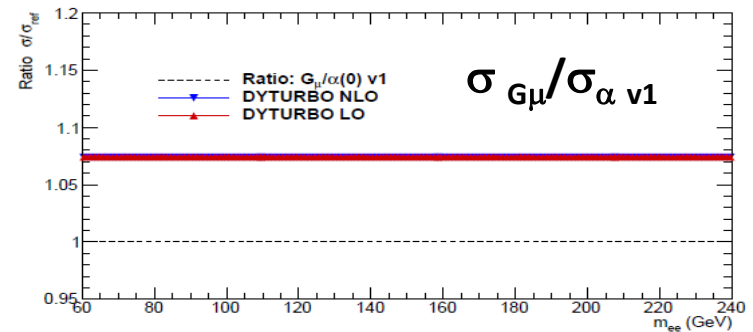
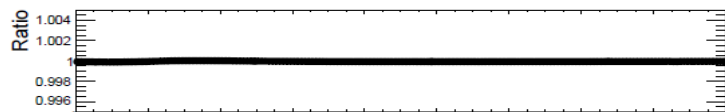
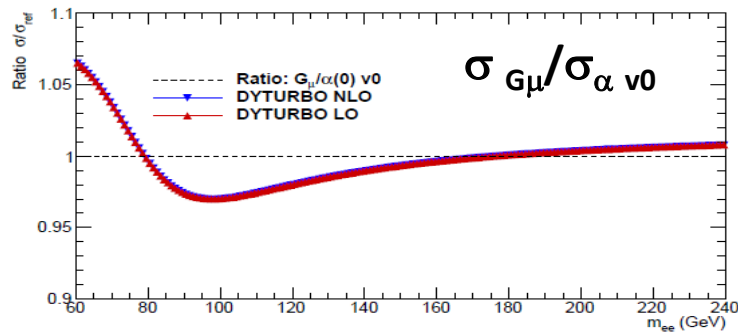
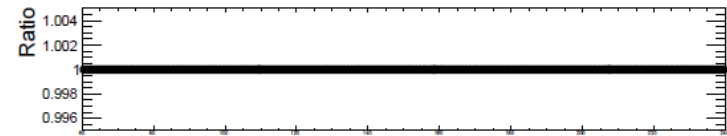
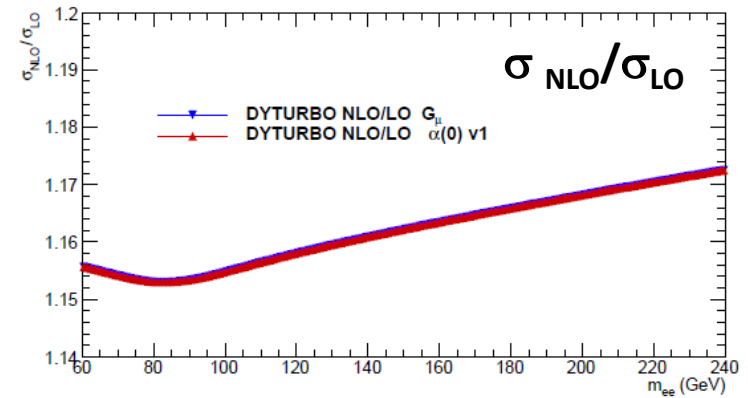
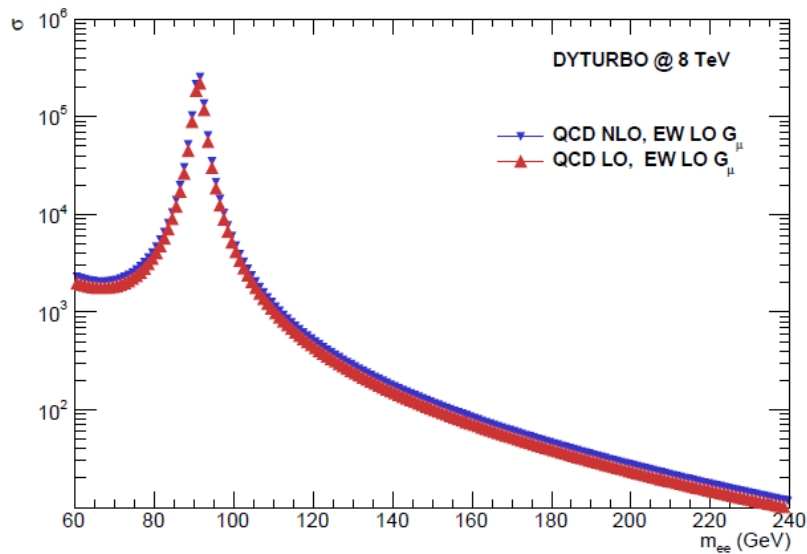
Benchmark for wt^{EW} reweighting

σ/σ_{ref}

—▼— PowhegZj+ wt^{EW}
—▲— DYTURBO

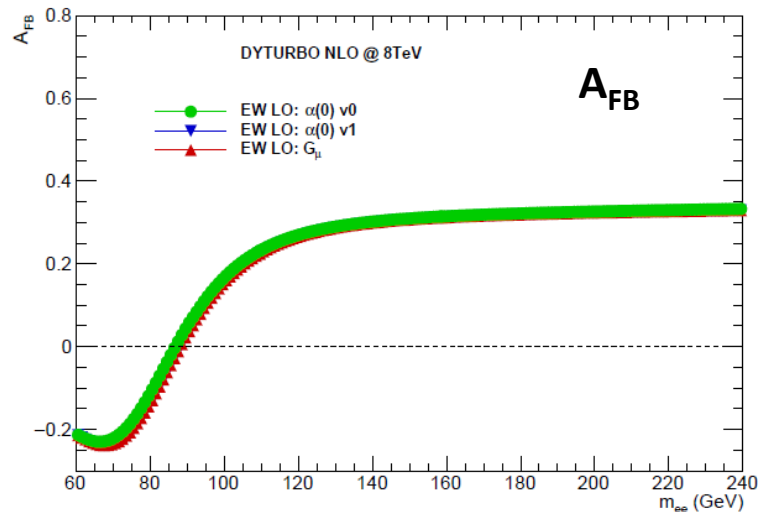


DYTURBO: QCD LO, NLO; EW LO

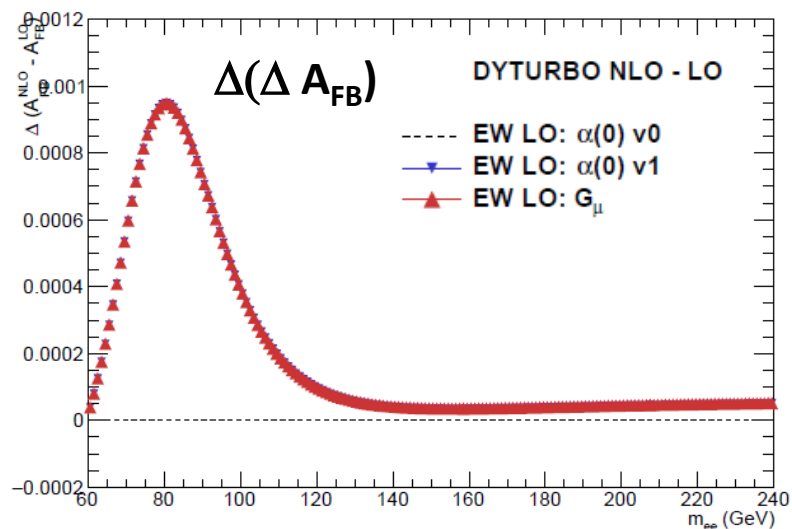
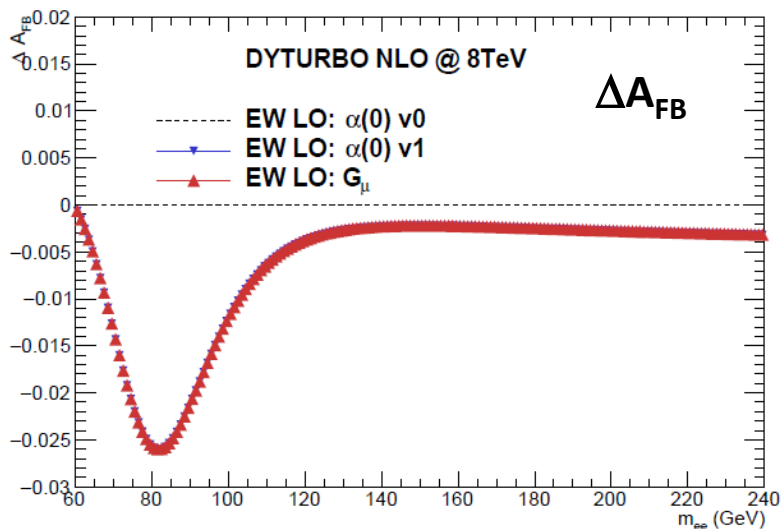
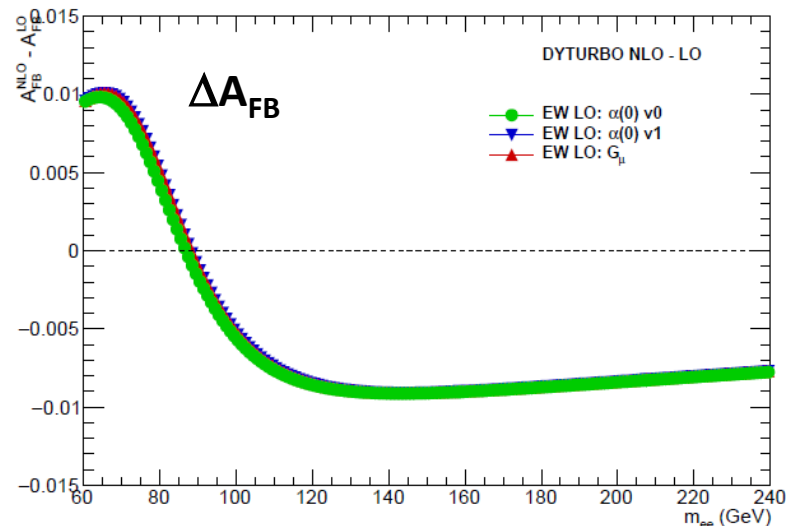


DYTURBO: QCD LO, NLO

QCD NLO



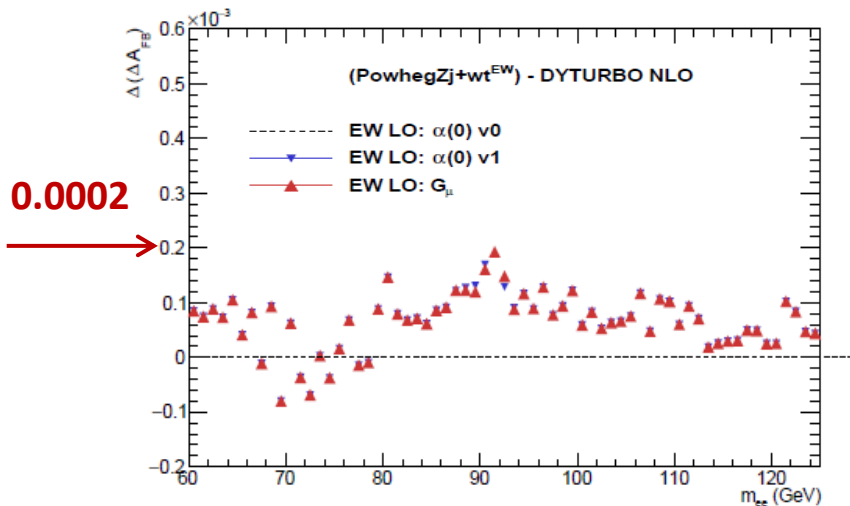
QCD NLO - LO



Validating reweighting with wt^{EW} : EW LO

ΔA_{FB} : driven by s^2_W value (same for $\alpha(0)$ v1 and G_μ schemes)

Benchmark for wt^{EW} reweighting



Double difference:

ΔA_{FB} (DYTURBO) - ΔA_{FB} (PowhegZj+wt^{EW})

$\alpha(0)$ v1 - $\alpha(0)$ v0

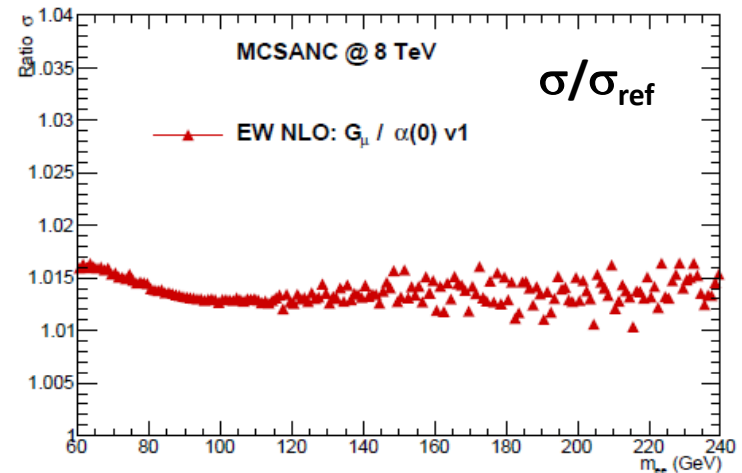
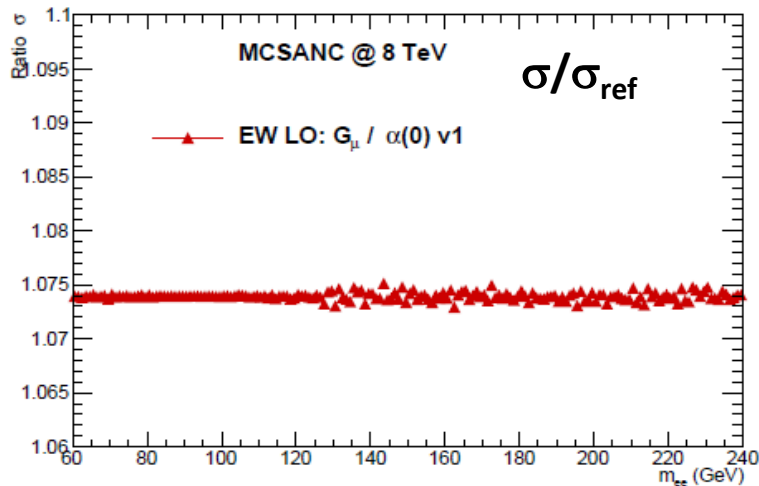
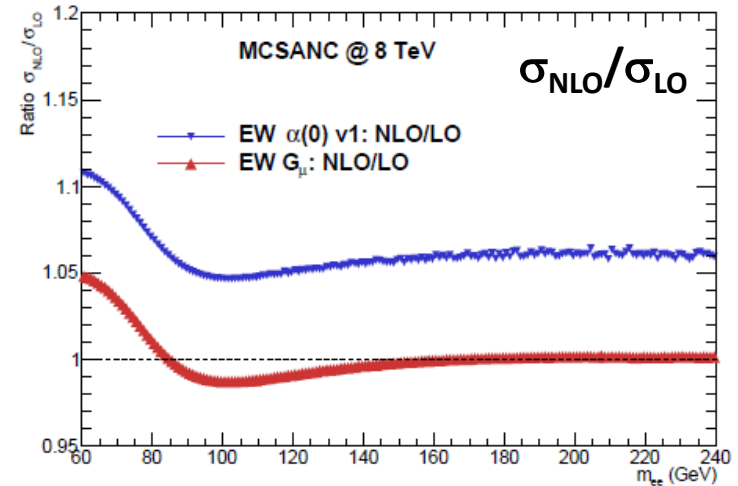
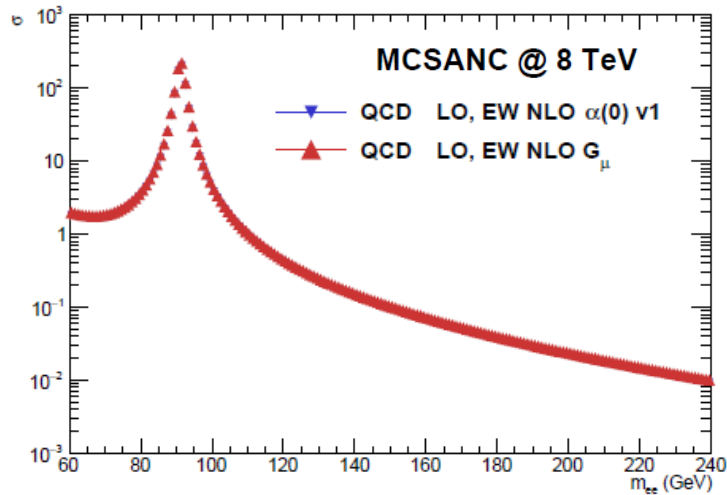
G_μ - $\alpha(0)$ v0

Agreement on $\Delta(\Delta A_{FB})$ within ± 0.0002

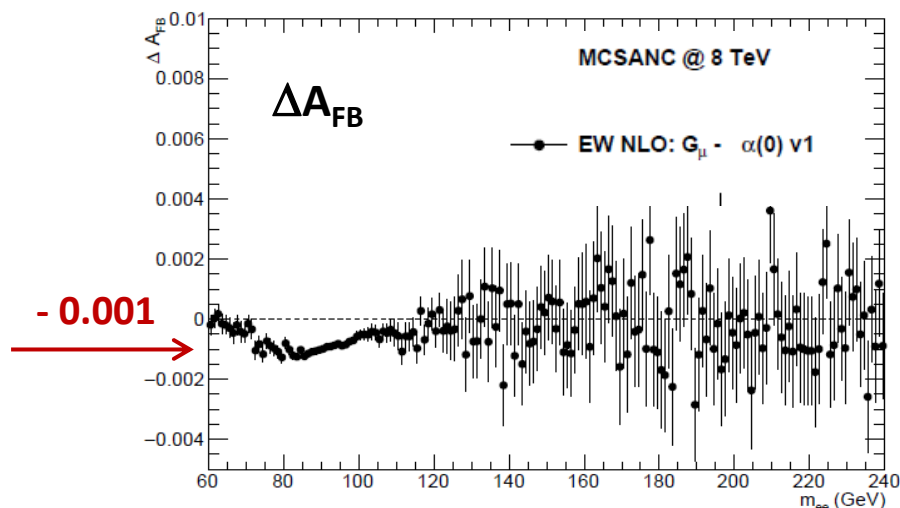
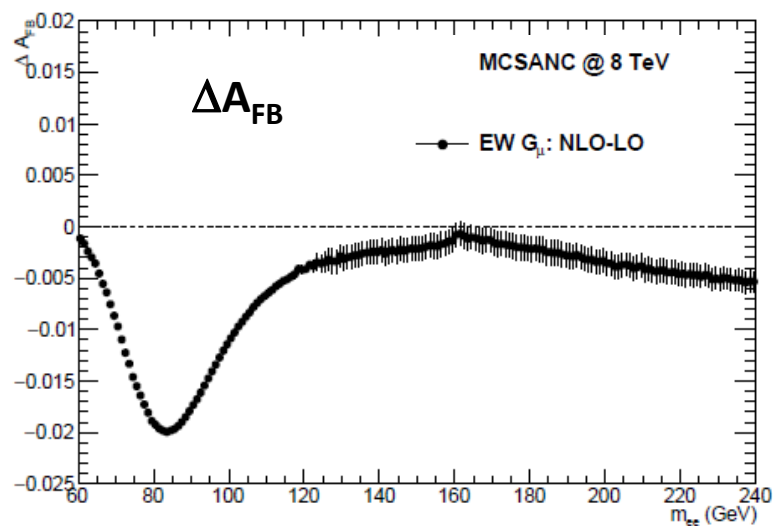
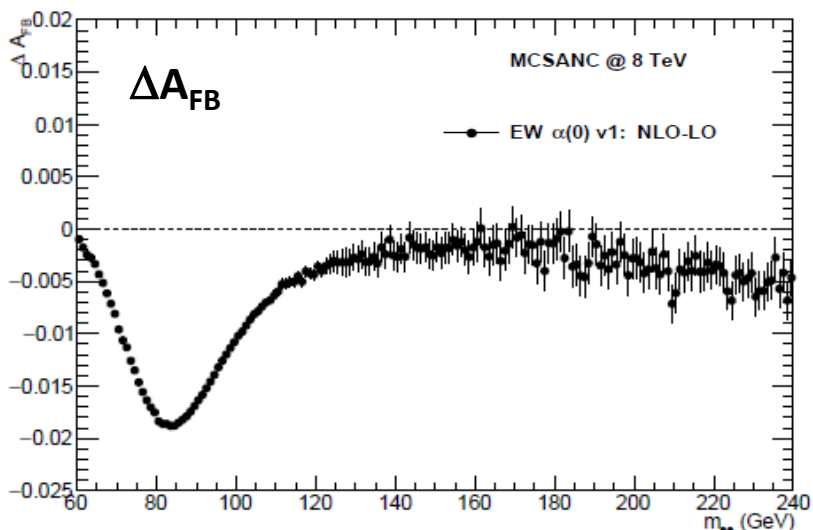
15

Should redo it with much finer binning around Z-pole to better estimate precision.

MCSANC: QCD LO; EW LO, NLO



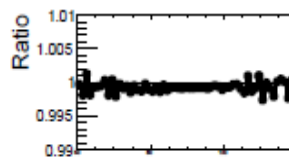
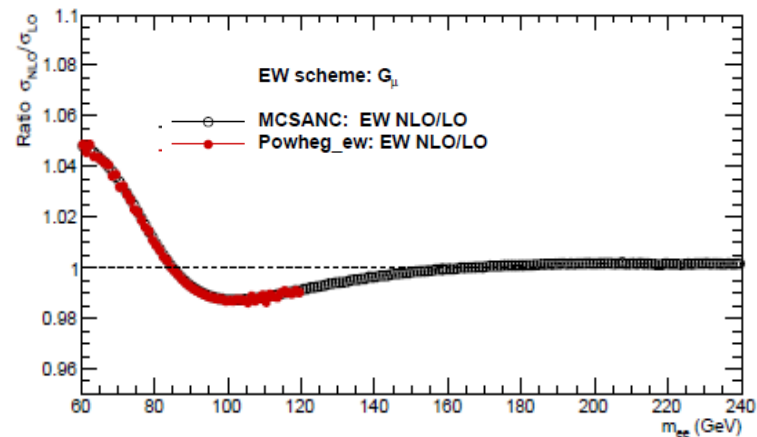
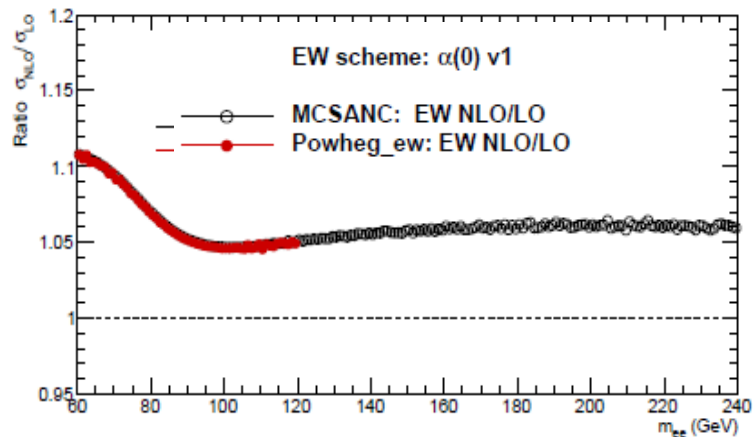
MCSANC: QCD LO, EW LO, NLO



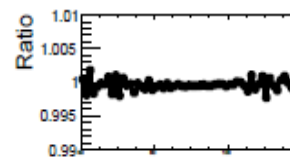
**EW NLO not enough
for precise s^2w measurement.
Needed EW HO corrections.**

MCSANC and Powheg_ew: EW LO, NLO

- Comparing ratio of cross-sections



In progress

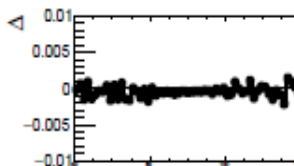
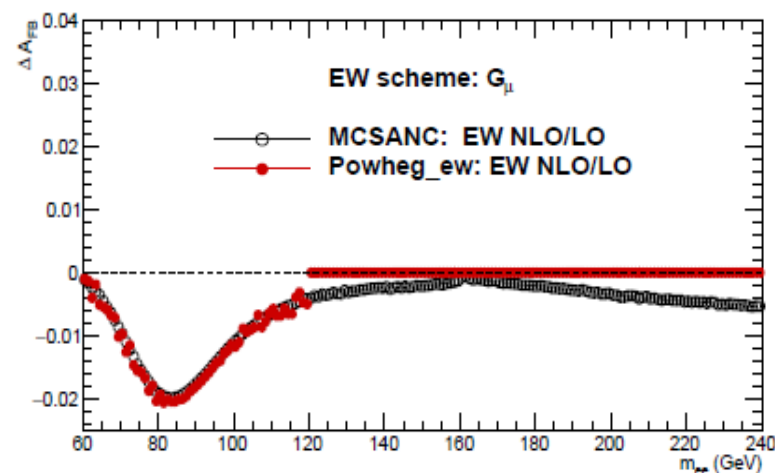
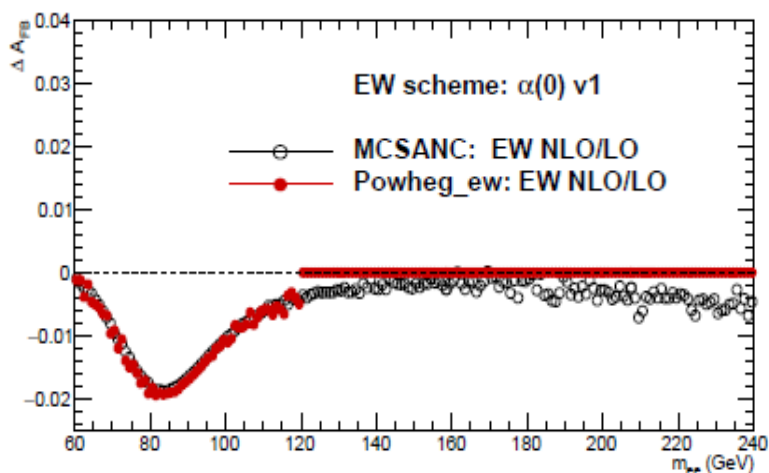


In progress

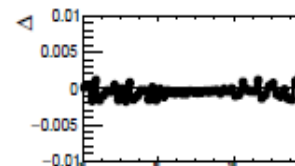
Very good agreement for EW NLO/LO corrections, shown for two EW schemes.

MCSANC and Powheg_ew: EW LO, NLO

- Comparing ΔA_{FB}



In progress

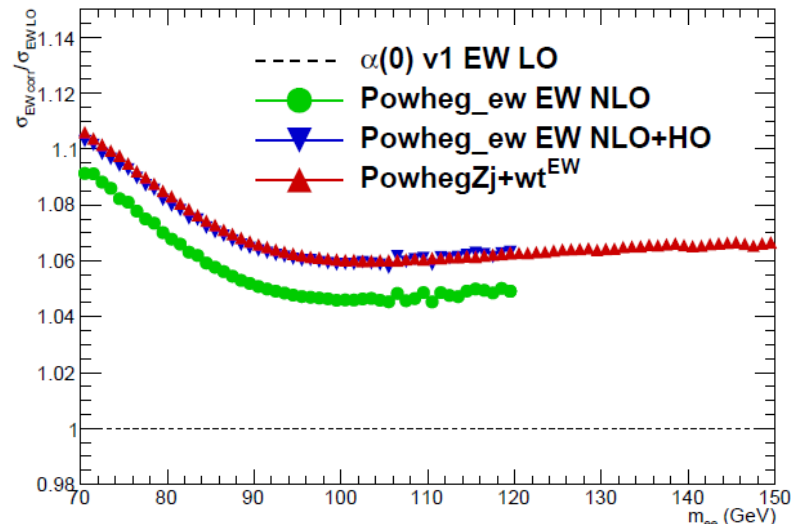
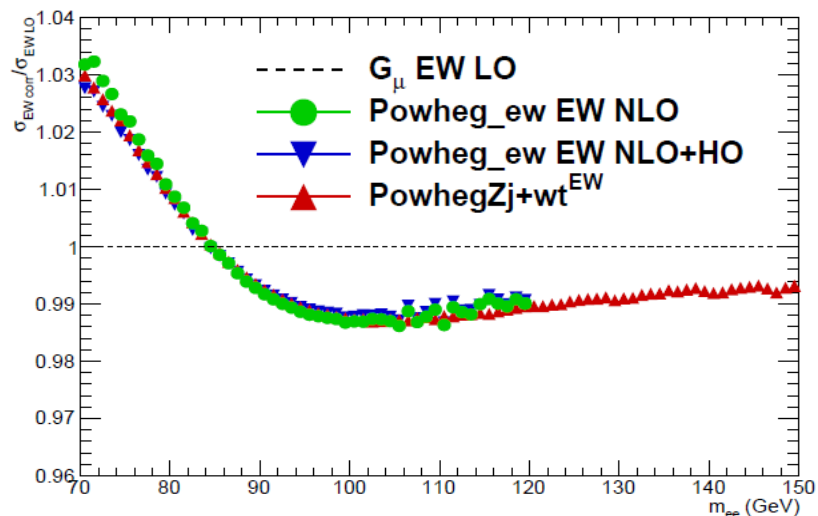


In progress

Very good agreement for EW NLO/LO corrections, shown.

Beyond EW LO: xsection NLO+HO

Slide from 25.09.2018



Powheg_ew

	EW order	$m_{ee} = 89 - 93 \text{ GeV}$
$\alpha(0) \text{ v1}$	NLO/LO	1.050350
G_μ	NLO/LO	0.991230
$\alpha(0) \text{ v1}$	NLO+HO/LO	1.063247
G_μ	NLO+HO/LO	0.991038

	EW order	$m_{ee} = 89 - 93 \text{ GeV}$
Powheg_ew	NLO+HO/LO	
$\alpha(0) \text{ v1}$		1.06325
G_μ		0.99104
PowhegZj+wt ^{EW}	NLO+HO/LO	
$\alpha(0) \text{ v0}$		0.96452
$\alpha(0) \text{ v1}$		1.06506
G_μ		0.99167

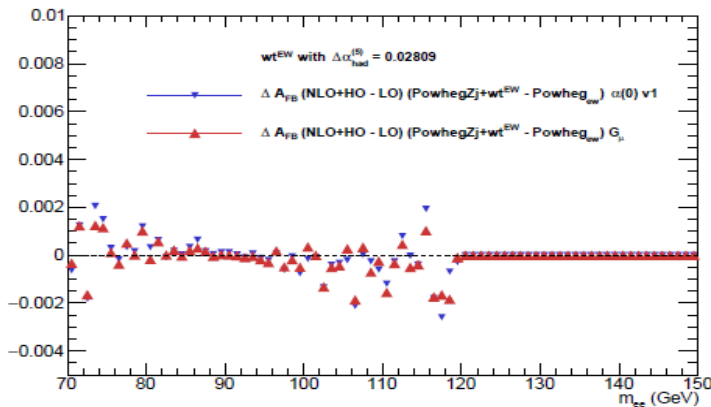
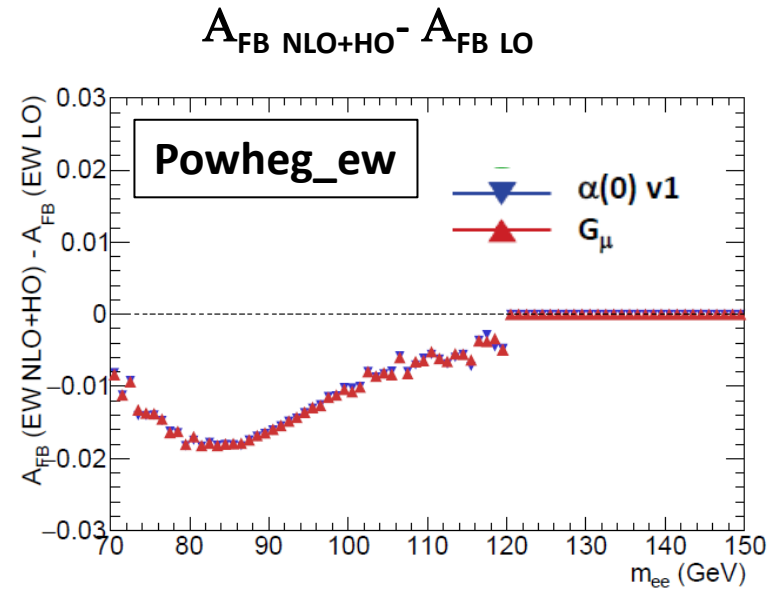
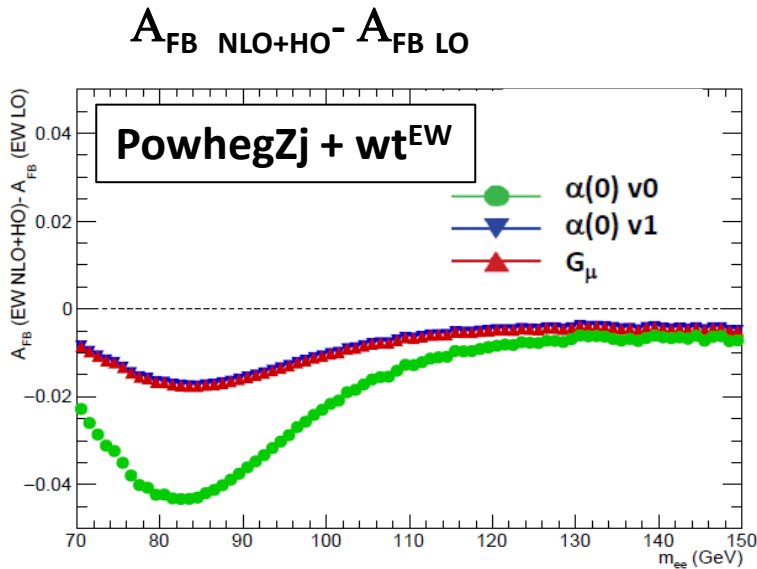
In G_μ scheme, NLO corrections < 1%

HO corrections < 0.02%

Beyond EW LO: A_{FB} NLO+HO

- Comparing Powheg_ew and PowhegZj+wt^{EW}

Slide from 25.09.2018



- $\Delta A_{FB}^{(NLO+HO - LO)} (\text{Powheg}_{ew} - \text{PowhegZj+wt}^{EW}) \alpha(0) v_1$
- $\Delta A_{FB}^{(NLO+HO - LO)} (\text{Powheg}_{ew} - \text{PowhegZj+wt}^{EW}) G_\mu$

Excellent agreement on ΔA_{FB} !

Is it accidental?

More discussion (25.09.2018) and in SPARES slides.

Powheg_ew: EW LO, NLO, NLO+HO

Cross-section

From slides 25.09.2018

	EW order	$m_{ee} = 89 - 93 \text{ GeV}$	$m_{ee} = 80 - 100 \text{ GeV}$	$m_{ee} = 70 - 120 \text{ GeV}$
$\alpha(0) \text{ v0}$	LO	630.848722	906.156051	959.658977
$\alpha(0) \text{ v1}$	LO	571.411296	821.363274	870.729908
G_μ	LO	612.514433	880.446121	933.363827
$\alpha(0) \text{ v1}$	NLO	600.185042	863.142557	915.580114
G_μ	NLO	607.142292	873.173294	926.253246
$\alpha(0) \text{ v1}$	NLO+HO	607.551746	873.717147	926.761229
G_μ	NLO+HO	607.515354	873.655348	926.681425
$\alpha(0) \text{ v1}$	NLO/LO	1.050350	1.05087	1.05151
G_μ	NLO/LO	0.991230	0.99174	0.99238
$\alpha(0) \text{ v1}$	NLO+HO/LO	1.063247	1.063740	1.064349
G_μ	NLO+HO/LO	0.991038	0.992287	0.992840
$\alpha(0) \text{ v1} / \alpha(0) \text{ v0}$	LO	0.90578	0.906426	0.90733
$G_\mu / \alpha(0) \text{ v1}$	LO	1.07193	1.07193	1.07193
$G_\mu / \alpha(0) \text{ v1}$	NLO	1.01159	1.01162	1.01166
$G_\mu / \alpha(0) \text{ v1}$	NLO+HO	0.99994	0.99993	0.99991
$G_\mu / \alpha(0) \text{ v0}$	LO	0.97094	0.97163	0.97260

σ (pb)
 $\sigma_{\text{NLO}} / \sigma_{\text{LO}}$
 $\sigma_{\text{NLO+HO}} / \sigma_{\text{LO}}$
 Ratios between
 EW schemes
 LO, NLO, NLO+HO

Better than 0.01% agreement on σ between EW schemes at NLO+HO !

Powheg_ew: EW LO, NLO, NLO+HO

From slides 25.09.2018

	EW order	$m_{ee} = 89 - 93 \text{ GeV}$	$m_{ee} = 80 - 100 \text{ GeV}$	$m_{ee} = 70 - 120 \text{ GeV}$
$A_{FB} \alpha(0) v0$	LO	0.06691361	0.06392369	0.06253754
$A_{FB} \alpha(0) v1$	LO	0.04653886	0.04343789	0.04212883
$A_{FB} G_\mu$	LO	0.04653886	0.04343789	0.04212883
$A_{FB} \alpha(0) v1$	NLO	0.03004289	0.02690785	0.02569858
$A_{FB} G_\mu$	NLO	0.02905841	0.02592168	0.02471918
$A_{FB} \alpha(0) v1$	NLO+HO	0.03083234	0.02770533	0.02649700
$A_{FB} G_\mu$	NLO+HO	0.03090286	0.02777783	0.02656851
$\Delta A_{FB} \alpha(0) v1$	NLO-LO	-0.0164959	-0.0165300	-0.0164302
$\Delta A_{FB} G_\mu$	NLO-LO	-0.0174805	-0.0175162	-0.0174096
$\Delta A_{FB} \alpha(0) v1$	NLO+HO-LO	-0.0157065	-0.0157326	-0.0156318
$\Delta A_{FB} G_\mu$	NLO+HO-LO	-0.0156360	-0.0156596	-0.0155603

A_{FB}

$\Delta A_{FB} \text{ (NLO - LO)}$

$\Delta A_{FB} \text{ (NLO+HO - LO)}$

ΔA_{FB}	EW order	$m_{ee} = 89 - 93 \text{ GeV}$	$m_{ee} = 80 - 100 \text{ GeV}$	$m_{ee} = 70 - 120 \text{ GeV}$
$\alpha(0) v1 - \alpha(0) v0$	LO	-0.020375	-0.020486	-0.020487
$G_\mu - \alpha(0) v0$	LO	-0.020375	-0.020486	-0.0204871
$G_\mu - \alpha(0) v1$	LO	0.0	0.0	0.0
$G_\mu - \alpha(0) v1$	NLO	-0.00098	-0.00098	-0.00098
$G_\mu - \alpha(0) v1$	NLO + HO	-0.00007	-0.00007	-0.00007

ΔA_{FB} between
EW schemes at
LO, NLO, NLO+HO

Better than 0.0001 agreement on A_{FB} between EW schemes at NLO+HO !

Summary

Comparing ratios, double ratios or double-differences turned out to be a very effective strategy. No need (so far) to fine-tune the QCD part of the calculations.

Tests done so far support applicability of IBA approach for Drell-Yan productions around Z-pole at LHC. Implemented in TauSpinner framework using form-factors from Dizet 6.21 library. Documentation: arXiv:1808.08616.

Flexibility to compare different EW schemes in a systematic manner.

Very promising consistency achieved between predictions from:

- Powheg_ew (EW LO, NLO, NLO+HO)
- MCSANC (EW LO, NLO)
- DYTURBO (EW LO)
- PowhegZj + w_t^{EW} from TauSpinner+Dizet (EW LO, NLO+HO)

Ironing details

Adding predictions of NLO+HO type (MCSANC)

Extending m_{\parallel} range (Powheg_ew)

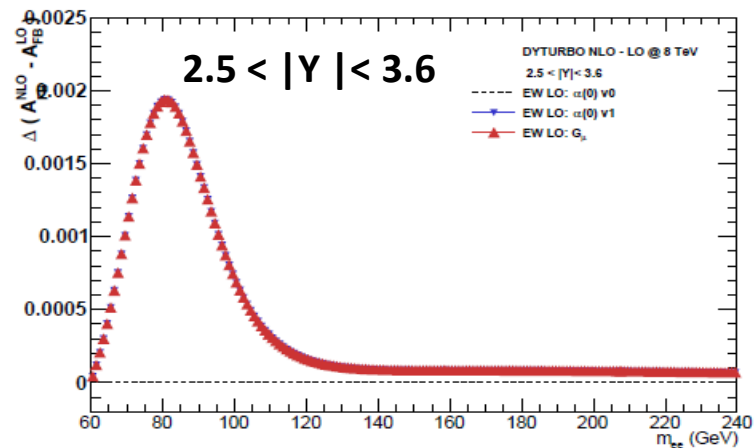
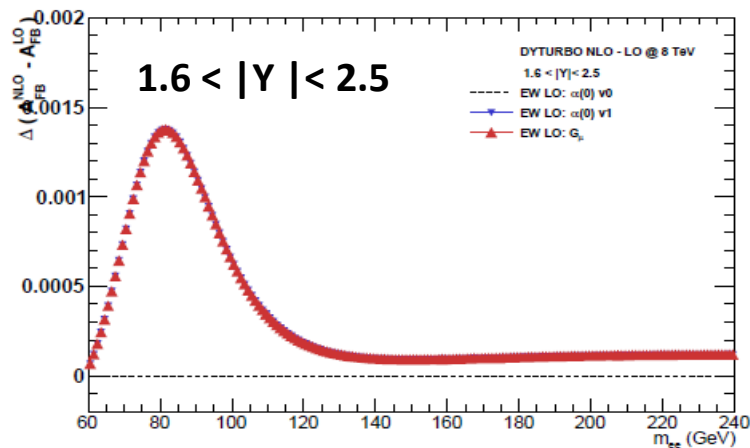
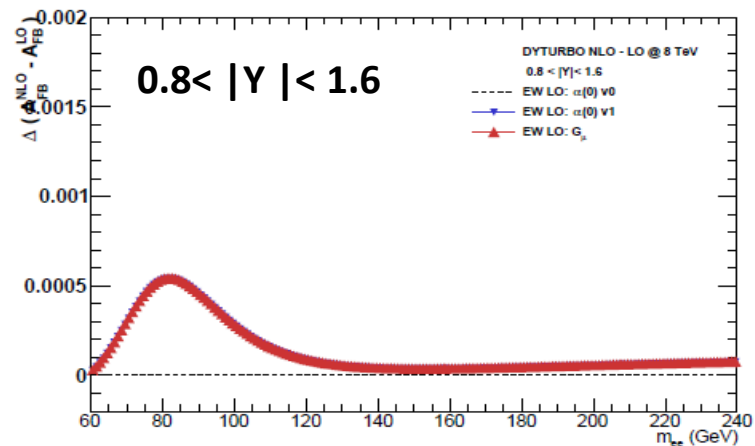
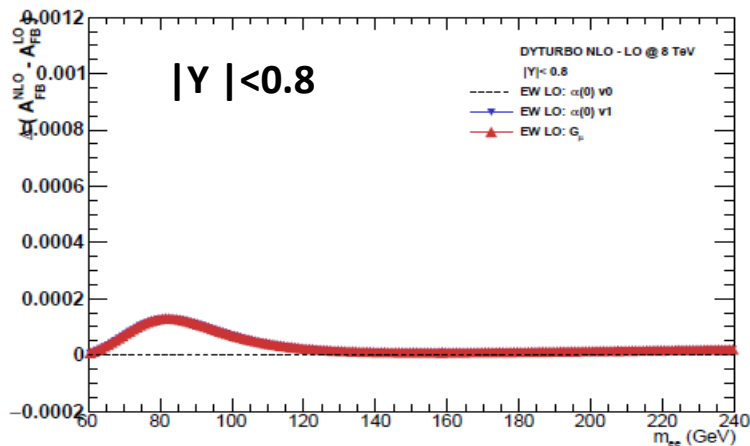
Fine-tuning w_t^{EW} (TauSpinner + Dizet)

Add fiducial selection

SPARES slides

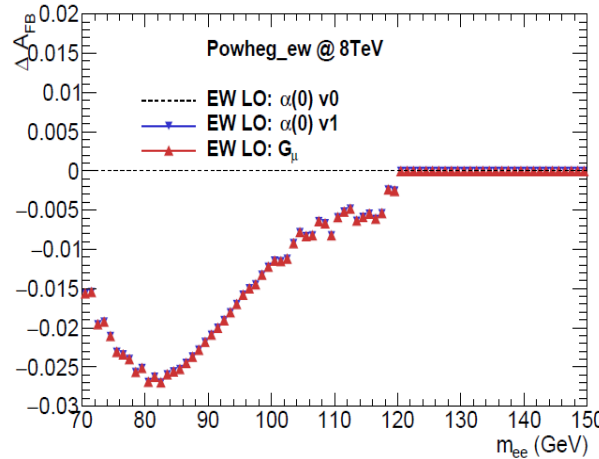
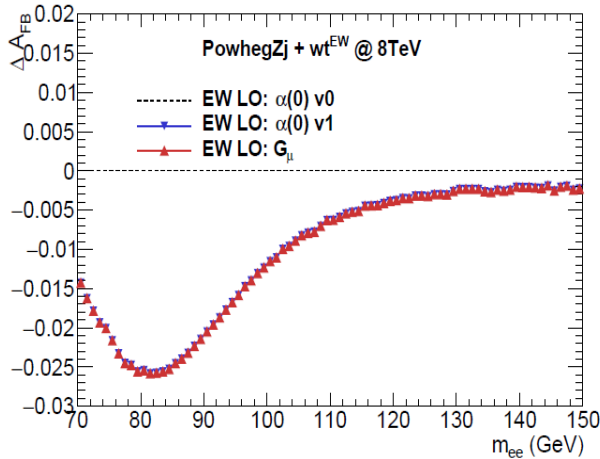
DYTURBO: QCD NLO - LO

$\Delta (A_{\text{FB}} (\text{NLO}) - A_{\text{FB}} (\text{LO}))$ between different EW LO schemes



Validating reweighting wt^{EW} : EW LO

ΔA_{FB} : driven by s^2_w value (same for $\alpha(0)$ v1 and G_μ schemes)

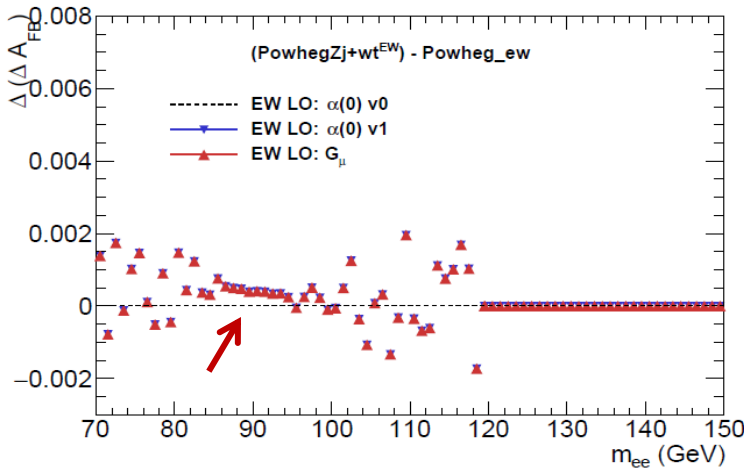


From slides 25.09.2018

$$\Delta A_{FB} = A_{FB} - A_{FB \text{ ref}}$$

$$\text{---}\blacktriangledown\text{---} \quad \alpha(0) \text{ v1} - \alpha(0) \text{ v0}$$

$$\text{---}\blacktriangle\text{---} \quad G_\mu - \alpha(0) \text{ v0}$$



Double difference:

$$\Delta A_{FB} (\text{Powheg_ew}) - \Delta A_{FB} (\text{Powheg}+wt^{EW})$$

$$\text{---}\blacktriangledown\text{---} \quad \alpha(0) \text{ v1} - \alpha(0) \text{ v0}$$

$$\text{---}\blacktriangle\text{---} \quad G_\mu - \alpha(0) \text{ v0}$$

Shift of 0.0005 at the Z-pole at EW LO, not precise enough tuning of s^2_w between „pole” and „on-mass-shell” definitions? Or the limitation of the reweighting with wt^{EW} ?

EW LO schemes in practice

- SM fundamental relations used to calculate EW parameters in EW LO schemes

$$G_\mu = \frac{\pi\alpha}{\sqrt{2}M_W^2 s_W^2} \longrightarrow \frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} = 1 \quad \Delta^2 = 16 \cdot s_W^2 \cdot (1 - s_W^2)$$

EW scheme: G_μ, α, M_Z

$\alpha(0) \nu 0$

$$d2 = \frac{\sqrt{2} \cdot 8\pi \cdot \alpha}{G_\mu \cdot M_Z^2}$$

$$\boxed{s_W^2} = (-1 + \sqrt{1 - d2/4})/2$$

EW scheme: α, M_W, M_Z

$\alpha(0) \nu 1$

$$\boxed{s_W^2} = 1 - m_W^2/m_Z^2$$

$$c_W^2 = m_W^2/m_Z^2$$

$$g2 = 4 \cdot \pi \cdot \alpha / s_W^2$$

$$\boxed{G_\mu} = \sqrt{2} \cdot g2 / 8 / m_W^2$$

EW scheme: G_μ, M_W, M_Z

G_μ

$$\boxed{s_W^2} = 1 - m_W^2/m_Z^2$$

$$c_W^2 = m_W^2/m_Z^2$$

$$g2 = 8 \cdot G_\mu \cdot m_W^2 / \sqrt{2}$$

$$\boxed{\alpha} = g2 \cdot s_W^2 / 4 / \pi$$

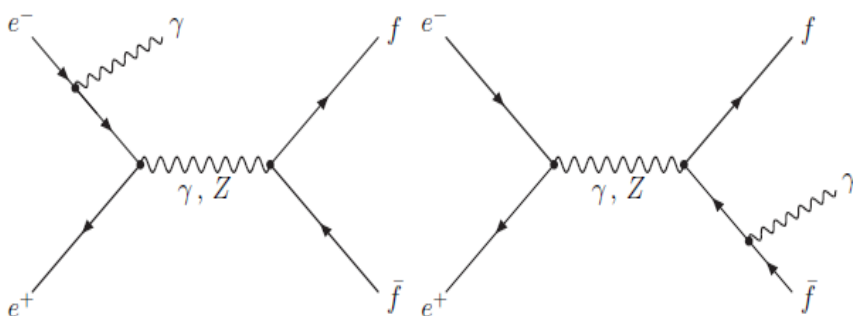
 calculated

QED (radiative) corrections

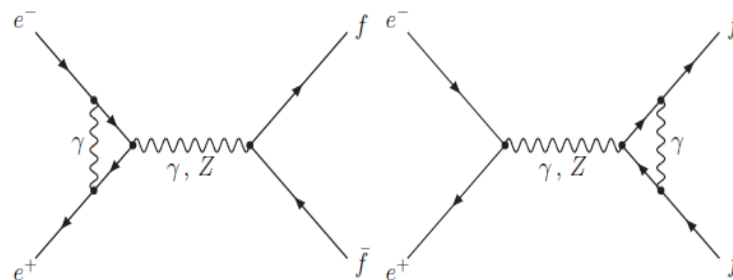
NOT discussed here.

QED FSR can be simulated by PHOTOS implemented as after-burner step on already generated event. QED ISR should be convoluted with QCD ISR. For QED initial/final state interference (see slide 42) and **talks by A. Sapronov and S. Yost.**

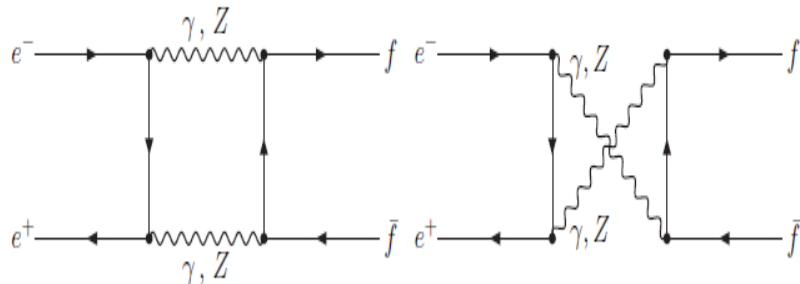
Real emission + pairs creation



Vertex corrections



$\gamma\gamma$ and γZ box diagrams



It is **QED gauge-invariant set of diagrams** (D. Bardin, hep-ph/9908433) which can be factorised out and/or convoluted with QCD corrections.

Calculated with fixed value of α_{QED}
 $\alpha_{\text{QED}} = 1./137.0359895$

Dictionary

EW LO Born (LO = lowest order):

tree-level vertex and propagator of the Z/γ^* bosons, setting of SM EW parameters defines the EW scheme.

EW effective Born:

tree-level vertex and propagator of the Z/γ^* bosons, EW couplings: $\alpha(m_Z)$, $\sin^2\theta_W(m_Z)$, m_Z , set of best measured values.

EW Improved Born Approximation (IBA):

tree-level vertex and propagator of the Z/γ^* bosons, EW couplings and propagators multiplied by form-factors dependent on the scattering angle of the lepton (choice of frame) and virtuality of Z/γ^* .

QED/EW corrections: D. Bardin et al. arXiv:9908433

separate set of „QED corrections” and „genuine EW + lineshape corrections”

Tree level

Standard Model relations:

$$G_F = \frac{\pi\alpha}{\sqrt{2}m_W^2 \sin^2\theta_W^{\text{tree}}}$$

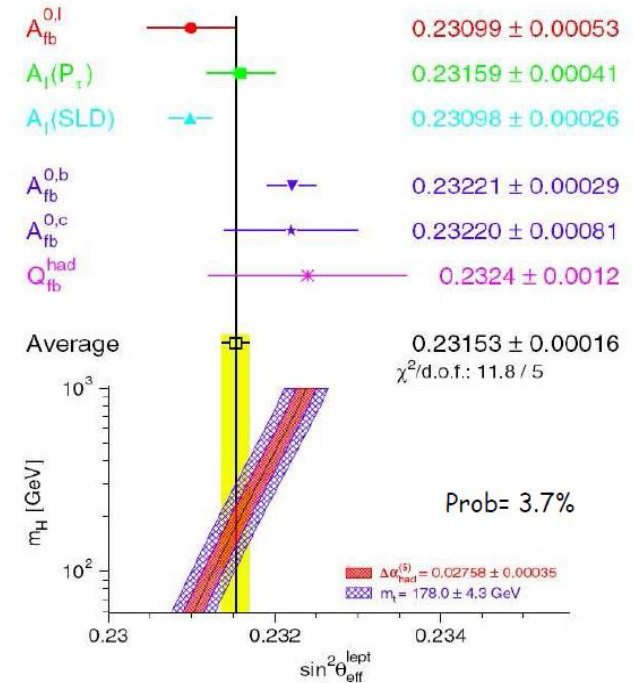
$$\rho_0 = \frac{m_W^2}{m_Z^2 \cos^2\theta_W^{\text{tree}}} = 1$$

Uncertainties on TH predictions

We have established such excellent agreement between very different: EW calculations, codes, QCD details (matrix element, PDFs,...).

Lets start now asking more detailed questions about individual corrections.

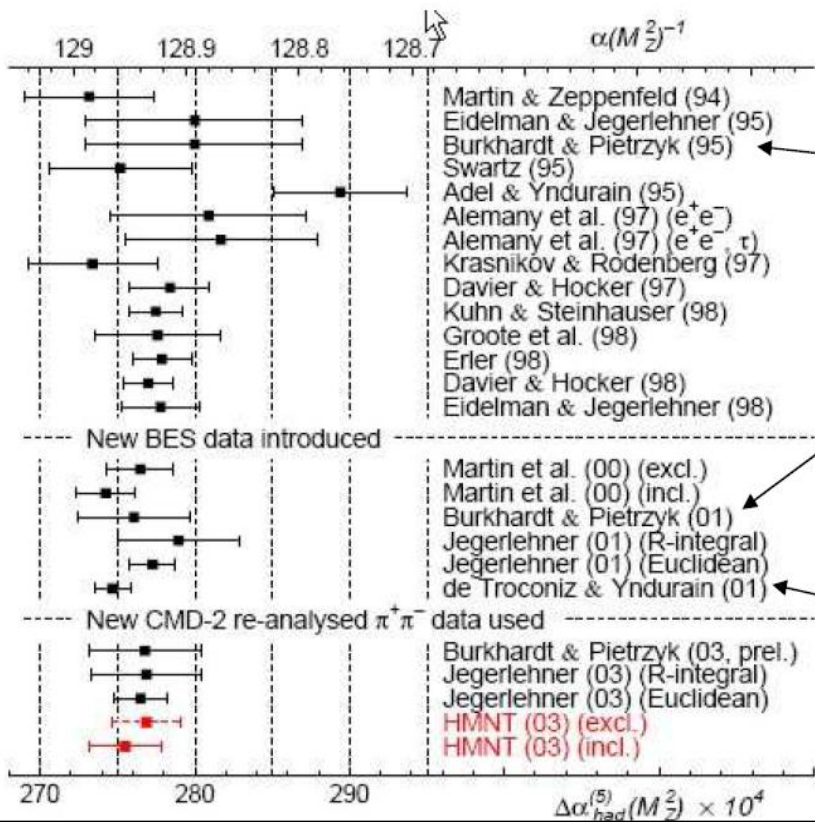
- Not all terms can/should be directly compared.
- The dominant systematics for LEP $\sin^2\theta_{\text{eff}}$ measurement, which was $\Delta\alpha^{(5)}_{\text{had}}(M_Z^2)$. Are we consistent about this correction term?
- The m_t since then known with 10 x better precision. Not an issue anymore.



New measurement (arXiv:1706.09436)
 $\Delta\alpha_{\text{had}}^{(5)}(M_Z) = 0.02753 \pm 0.00009$

Status of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

from D. Schlatter, 2007



At end of LEP $\Delta\alpha_{\text{had}}$ became limiting uncertainty in SM fits.

$$\Delta\alpha_{\text{had}}(M_Z) = 0.02804 \pm 0.00065$$

Post LEP measurements from BES and CMD-2 improvement.

$$\Delta\alpha_{\text{had}}(M_Z) = 0.02758 \pm 0.00035$$

This value is used by EWWG

Using perturbative QCD

$$\Delta\alpha_{\text{had}}(M_Z) = 0.02749 \pm 0.00012$$

Recent measurements:

arXiv:1706.09436

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z) = 0.02753 \pm 0.00009$$

arXiv:1802.02995

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z) = 0.02761 \pm 0.00011$$

$$\alpha(s) = \frac{\alpha(0)}{1 - \Delta\alpha_h^{(5)}(s) - \Delta\alpha_\ell(s) - \Delta\alpha^t(s) - \Delta\alpha^{\alpha_s}(s)}$$

$$\Delta\alpha_h^{(5)}(M_Z^2) = 0.0280398$$

$$\Delta\alpha_\ell(M_Z^2) = 0.0314976$$

$$\Delta\alpha^t(M_Z^2) = -0.585844 \cdot 10^{-4}$$

$$\Delta\alpha^{\alpha_s}(M_Z^2) = -0.103962 \cdot 10^{-4}$$

Dizet 6.21 default

LHCC Precision WG meeting, 15.11.2018

Impact of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Predictions from Dizet 6.21 library

Parameter	$\Delta\alpha_h^{(5)}(M_Z^2) = 0.0280398$	$\Delta\alpha_h^{(5)}(M_Z^2) = 0.02753$	Ratio
$\alpha(M_Z^2)$	0.00775884	0.00775463	
$1/\alpha(M_Z^2)$	128.885224	128.95522	0.99932
s_W^2	0.22351946	0.22331458	1.00092
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (electron, muon)	0.23175990	0.23157062	1.00082
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (up-quark)	0.23164930	0.23146414	1.00080
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (down-quark)	0.23152214	0.23133715	1.00080
M_W	80.35281 GeV	80.36341 GeV	1.00013
Δr	0.03694272	0.03631342	1.01733
Δr_{rem}	0.01169749	0.01170244	0.99958
ρ_{eu}	1.005408	1.005426	0.99998
K_e	1.036649	1.036770	0.99988
K_u	1.036172	1.036293	0.99988
K_{eu}	1.074146	1.074397	0.99977
ρ_{ed}	1.005894	1.005906	0.99999
K_e	1.036649	1.036699	0.99995
K_d	1.035603	1.035719	0.99989
K_{ed}	1.073556	1.073859	0.99972

shift of about -0.00020
due to corrections to M_W



← shift by +11 MeV

ATLAS measurement
 $M_W = 80370 \pm 19$ MeV

$$M_W = \frac{M_Z}{\sqrt{2}} \sqrt{1 + \sqrt{1 - \frac{4A_0^2}{M_Z^2(1 - \Delta r)}}}$$

$$\Delta r = \Delta\alpha(M_Z^2) + \Delta r_{EW}$$

$$A_0 = \sqrt{\frac{\pi\alpha(0)}{\sqrt{2}G_\mu}}$$

Impact of m_t

Parameter	$m_t = 171$ GeV	$m_t = 173$ GeV	$m_t = 175$ GeV
$\alpha(M_Z^2)$	0.00775882	0.00775884	0.00775885
$1/\alpha(M_Z^2)$	128.888558	128.885224	128.885079
s_W^2	0.22375411	0.22351946	0.22328310
$\sin^2\theta_W^{eff}(M_Z^2)$ (electron, muon)	0.23181756	0.23175990	0.23169368
$\sin^2\theta_W^{eff}(M_Z^2)$ (up-quark)	0.23171096	0.23164930	0.23169368
$\sin^2\theta_W^{eff}(M_Z^2)$ (down-quark)	0.23158377	0.23152214	0.23145996
Δr	0.03766186	0.03694272	0.03621664
Δr_{rem}	0.01165959	0.01169749	0.01173500
ρ_{eu}	1.005229	1.005408	1.005589
K_e	1.035837	1.036649	1.037467
K_u	1.035361	1.036172	1.036990
K_{eu}	1.072465	1.074146	1.075843
ρ_{ed}	1.005714	1.005894	1.006075
K_e	1.035837	1.036649	1.037467
K_d	1.034792	1.035603	1.036420
K_{ed}	1.071876	1.073556	1.075252

**± 2 GeV shift in m_t
corresponds to
 ± 0.00005 shift
in $\sin^2_{eff}{}^{lep}$**

Dizet 6.21 -> 6.42-> 6.44

AMT4 = 4 – available in Dizet 6.21

Pragmatic question: is it indeed more precise estimate to use AMT4=5 or AMT4=6?

Or better stay with well tested AMT4=4 ? What uncertainty attribute to this correction?

arXiv:1302.1395v3

Table 1: ZFITTER v.6.44beta, with the input values $\alpha_s = 0.1184$, $M_Z = 91.1876$ GeV, $M_H = 125$ GeV, $m_t = 173$ GeV. The dependence on electroweak NNLO corrections is studied for IMOMS=1 (input values are α_{em} , M_Z , G_μ). AMT4=4: with two-loop sub-leading corrections and re-summation recipe of [23-28] of [13]; AMT4=5: with fermionic two-loop corrections to M_W according to [29,30,32] of [13]; AMT4=6: with complete two-loop corrections to M_W [37] and fermionic two-loop corrections to $\sin^2 \theta_W^{\text{lept,eff}}$ [52] of [13]. IBAIKOV=0 (no α_s^4 QCD corrections) or IBAIKOV=2012 [190].

AMT4	4	5	6	Diff.	Exp. Err.
IBAIKOV=0					
$\Gamma_Z(\mu^+\mu^-)$, MeV	83.9782	83.9748	83.9807	0.0059	0.086
Γ_Z , MeV	2494.7863	2494.6019	2494.8688	0.2669	2.3
$\Gamma_W(l\nu)$, MeV	226.3185	226.2877	226.2922	0.0308	1.9
Γ_W , MeV	2090.3308	2090.0465	2090.0882	0.2843	42
M_W , GeV	80.3578	80.3541	80.3546	0.0037	0.015
$\sin^2 \theta_{\text{eff}}^{\text{lept}}$	0.231722	0.231791	0.231670	0.000121	0.00012
IBAIKOV=2012					
$\Gamma_Z(\mu^+\mu^-)$, MeV	83.9782	83.9748	83.9807	0.0059	0.086
Γ_Z , MeV	2494.5591	2494.3747	2494.6416	0.2669	2.3
$\Gamma_W(l\nu)$, MeV	226.3185	226.2877	226.2922	0.030	1.9
Γ_W , MeV	2090.1117	2089.8274	2089.8691	0.2843	42
M_W , GeV	80.3578	80.3541	80.3546	0.0037	0.015
$\sin^2 \theta_{\text{eff}}^{\text{lept}}$	0.231722	0.231791	0.231670	0.000121	0.00012



± 0.00005
around nominal
value of $\sin^2 \theta_{\text{eff}}$
with AMT4=4

Dizet 6.21 initialisation (KKMC default)

Internal flag	Default value	Optional value	Description
ibox	1	0,1	EW boxes on/off
Ihvp	1	1,2,3	Jegerlehner/Eidelman, Jegerlehner(1988), Burkhardt et al.
Iamt4	4	0,1,2,3,4	=4 the best, Degrassi/Gambino
Iqcd	3	1,2,3	approx/fast/lep1, exact/Slow!/Bardin/, exact/fast/Kniehl
Imoms	1	0,1	=1 W mass recalculated
Imass	0	0,1	=1 test only, effective quark masses
Iscre	0	0,1,2	Remainder terms
Ialem	3	1,3 or 0,2,	for 1,3 DALH5 not input
Imask	0	0,1	=0: Quark masses everywhere; =1 Phys. threshold in the ph.sp.
Iscal	0	0,1,2,3	Kniehl=1,2,3, Sirlin=4
Ibarb	2	-1,0,1,2	Barbieri???
Iftjr	1	0,1	FTJR corrections
Ifacr	0	0,1,2,3	Expansion of δ_r ; =0 none; =3 fully, unrecommend.
Ifact	0	0,1,2,3,4,5	Expansion of kappa; =0 none
Ihigs	0	0,1	Leading Higgs contribution resummation
Iafmt	1	0,1	=0 for old ZF
Iewlc	1	0,1	???
Iczak	1	0,1	Czarnecki/Kuehn corrections
Ihig2	1	0,1	Two-loop higgs corrections off,on
Iale2	3	1,2,3	Two-loop constant corrections in δ_α
Igfer	2	0,1,2	QED corrections for fermi constant
Iddzz	1	0,1	??? DD-ZZ game, internal flag