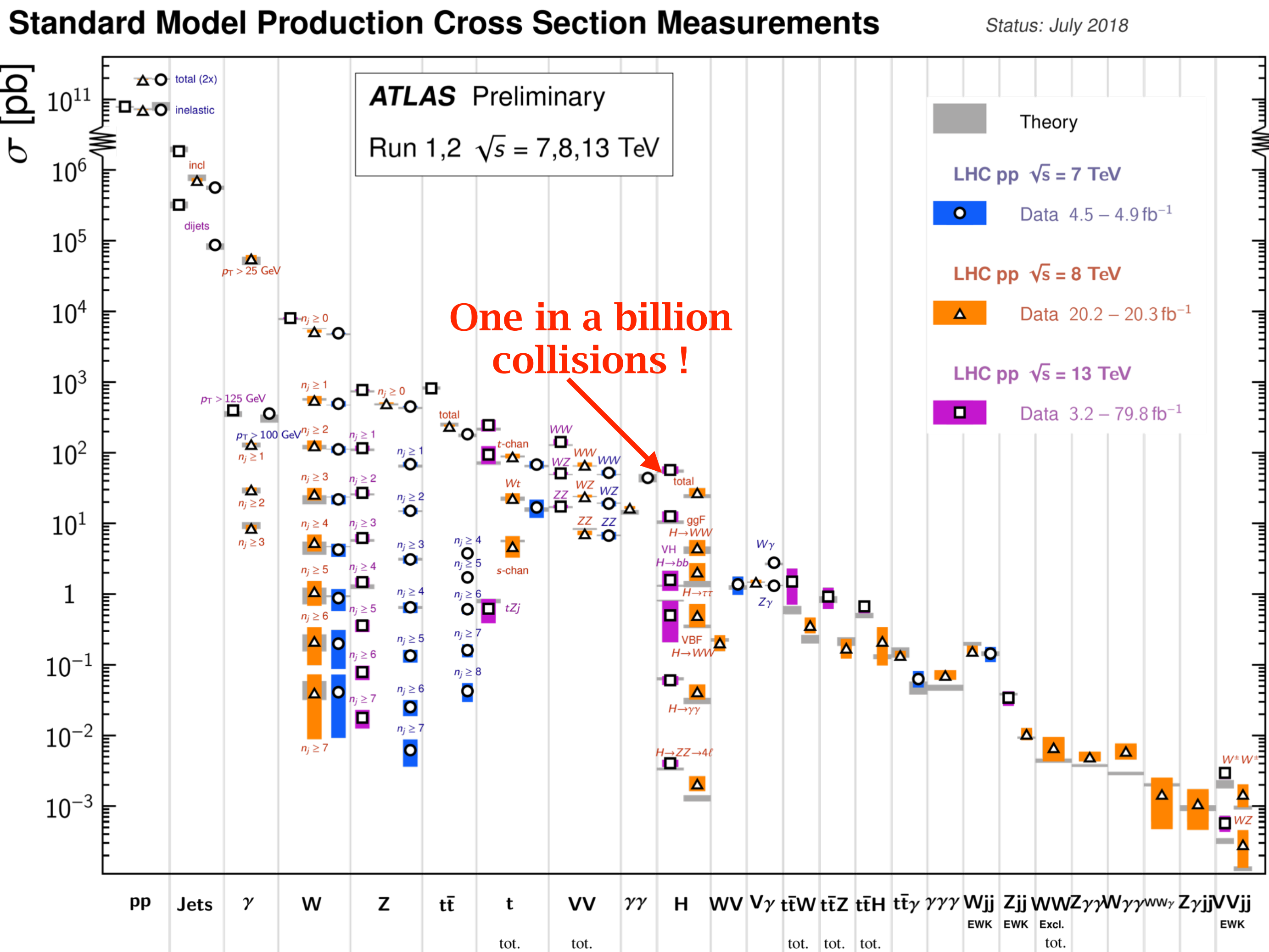


Resummation at the LHC: recent progress and future challenges

Pier Francesco Monni
CERN

Standard Model at the LHC

Some aspects are less precisely established experimentally, or not at all. Need to go differential...



$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{\Psi} \not{D} \Psi + h.c.$$

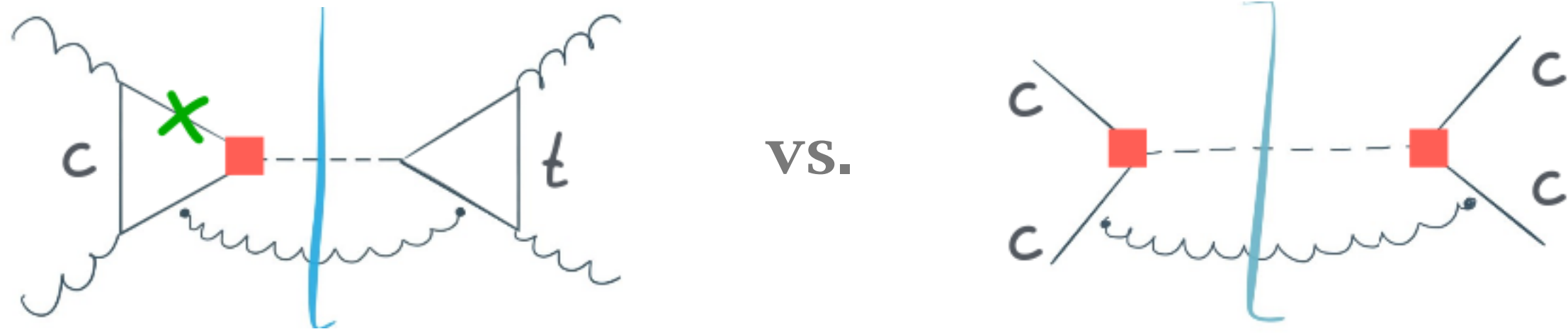
$$+ \bar{\Psi}_i y_{ij} \Psi_j \phi + h.c.$$

$$+ |D_\mu \phi|^2 - V(\phi)$$

- ▶ New physics can be encoded in small deviations in kinematic distributions

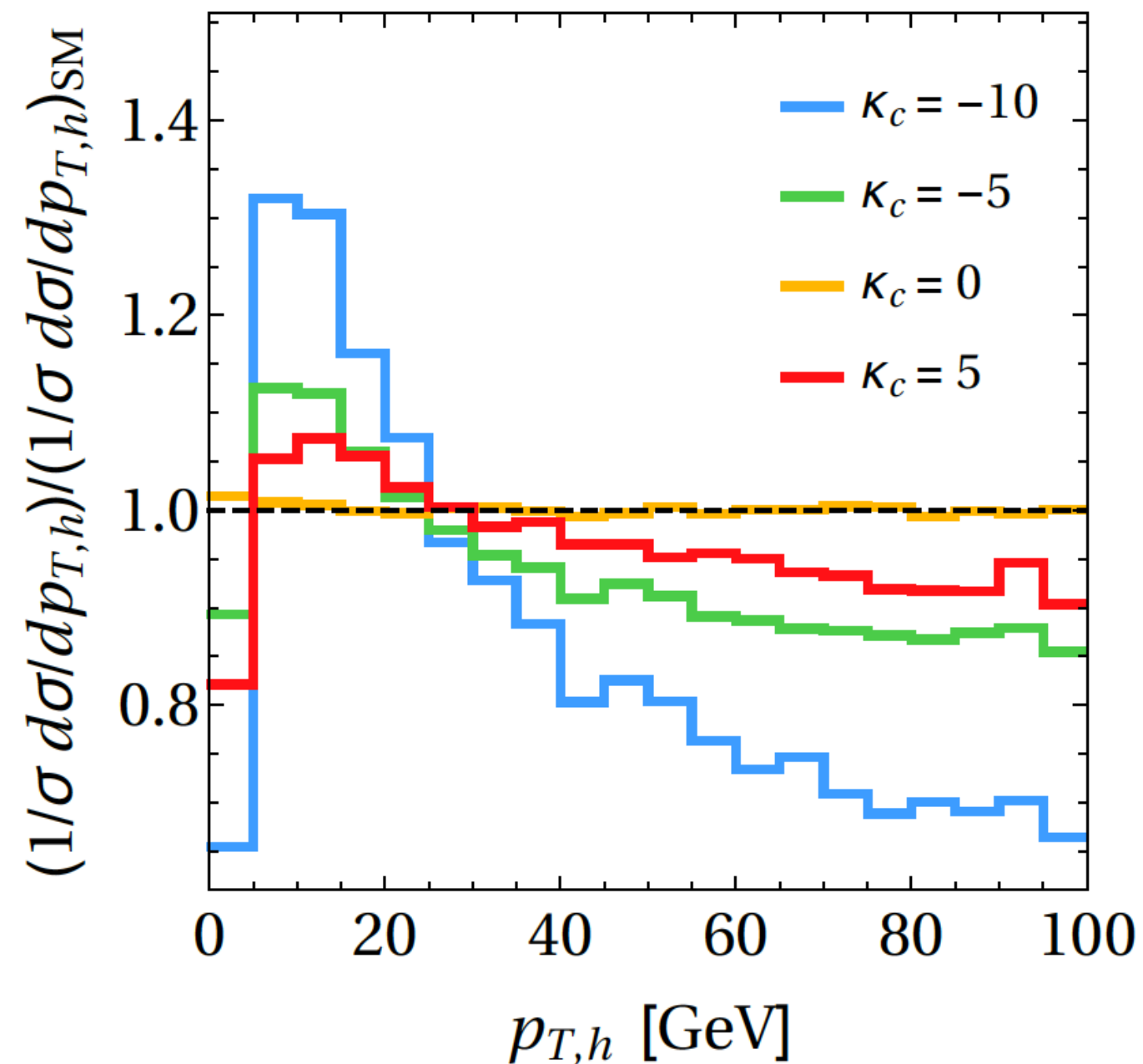
e.g. Yukawa couplings

- ▶ Interplay between production modes leads to sensitivity to light-quark couplings at small $p_{T,H}$



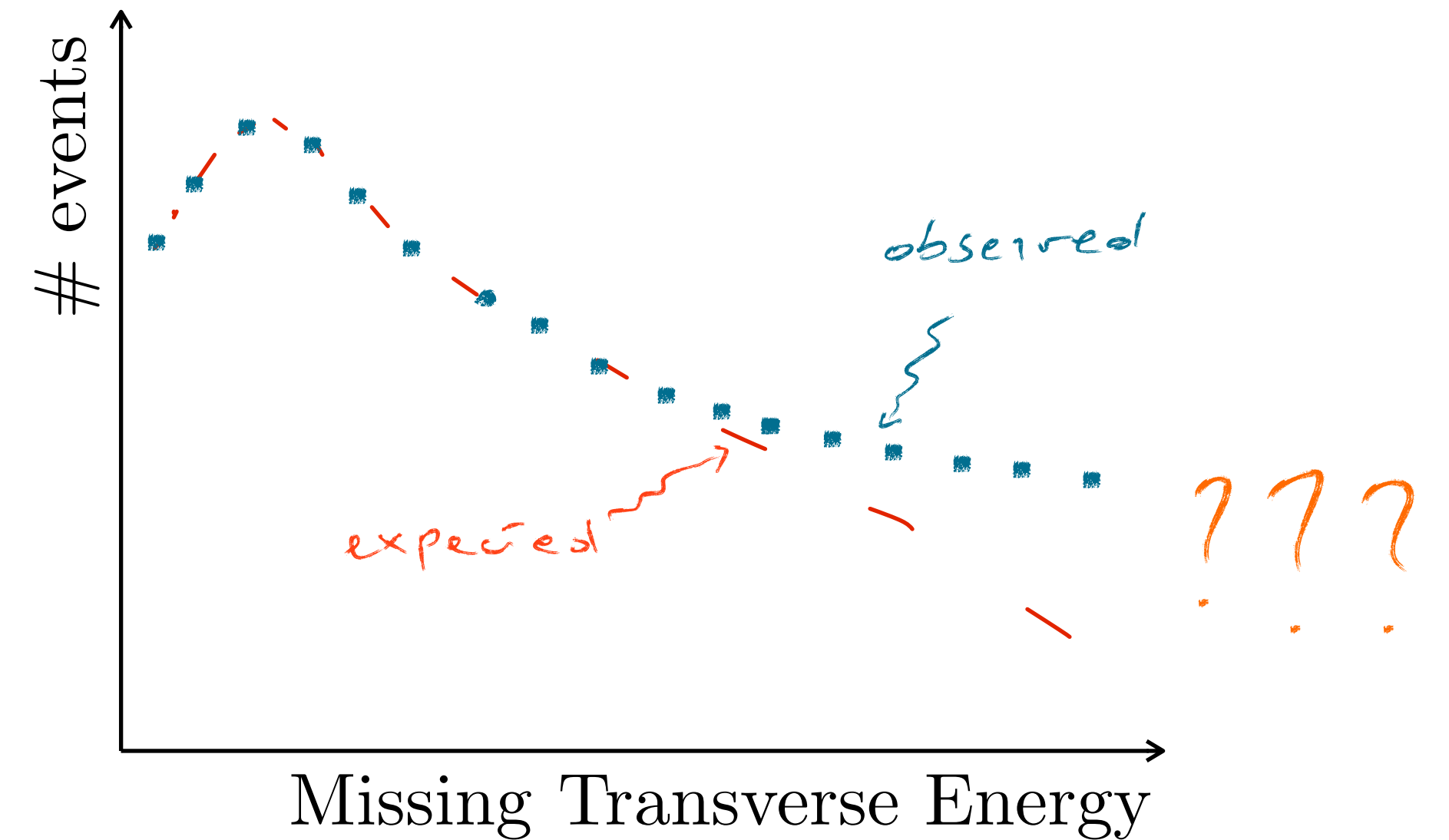
[Bishara, Haisch, PM, Re '16;
see also Soreq, Zhu, Zupan '16]

$$y_f^{\text{SM}} = \sqrt{2} \frac{m_f}{v} \kappa_f$$



e.g. NP searches in boosted regimes

- ▶ Regimes with large momentum transfer indirectly sensitive to heavy new physics states - look for (mild) effects on kinematics
- ▶ Precise comparison between Theory and Data is necessary: **mismodelling or actual signal?**

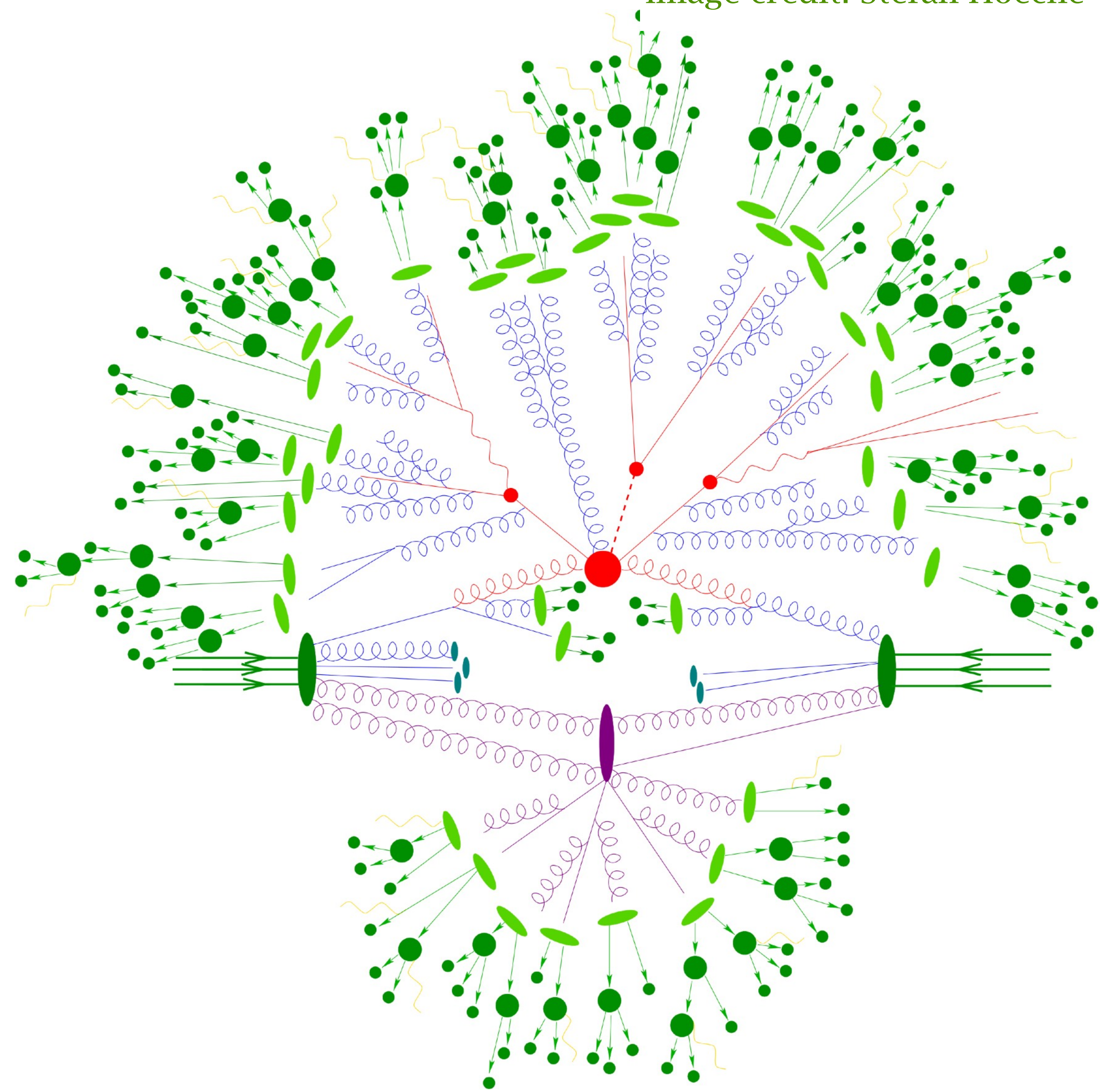


▶ Demanding a more exclusive description of final states comes with complications

☑ **Hard scattering** provides a good description of sufficiently hard final states (IRC safety)

$$\sigma \sim \sigma_{\text{BORN}} \left(1 + \alpha_s + \alpha_s^2 + \alpha_s^3 + \dots \right)$$

Image credit: Stefan Hoeche



▶ Demanding a more exclusive description of final states comes with complications

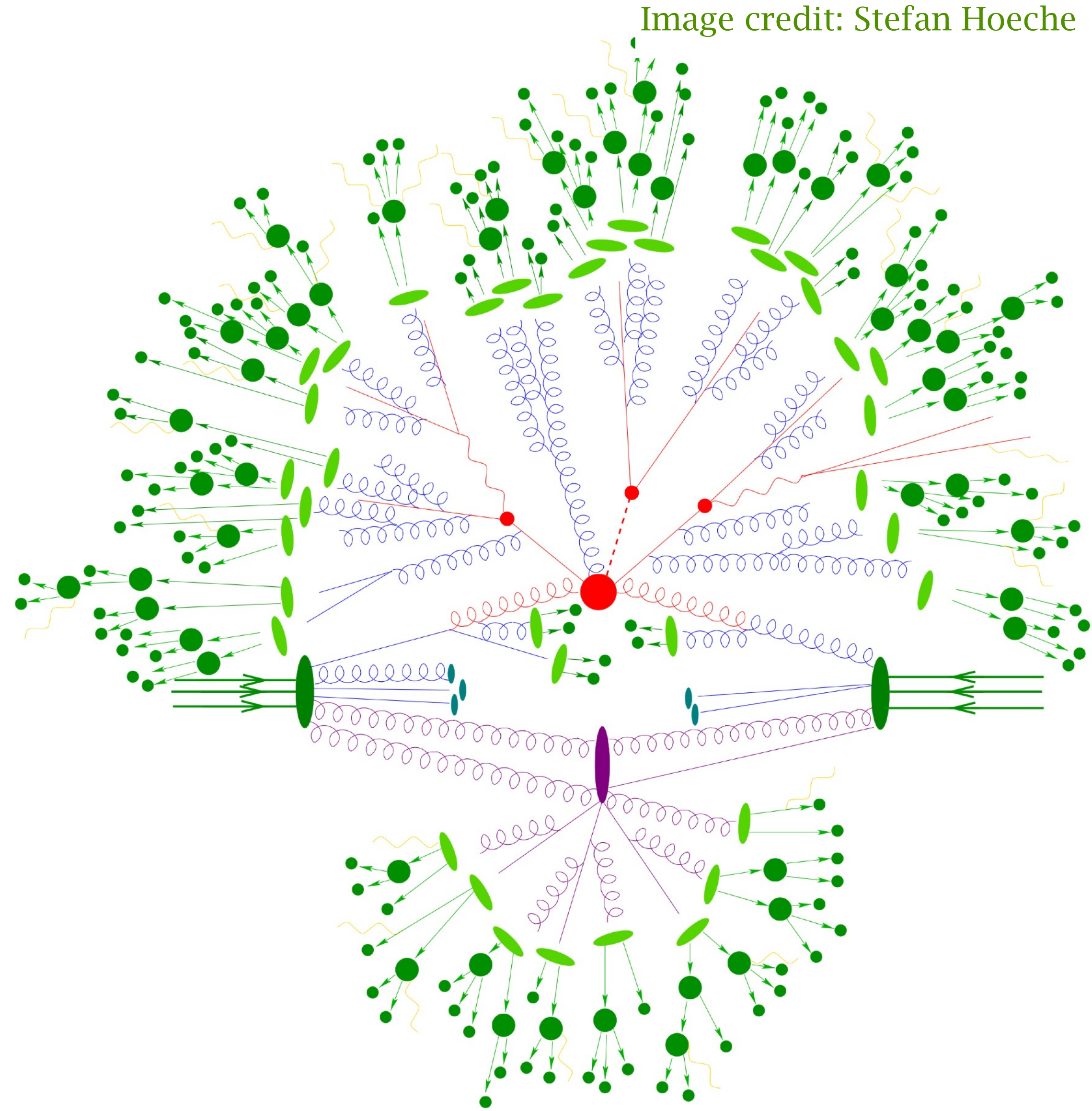
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$$\sigma \sim \sigma_{\text{BORN}} \left(1 + \alpha_s + \alpha_s^2 + \alpha_s^3 + \dots \right)$$

☑ **Multiscale dynamics** becomes dominant whenever a large gap between scales is present

- ▶ observables sensitive to soft or collinear radiation (differential distributions, jet vetoes, ...)
- ▶ reaction depends on disparate scales (production at threshold, masses, ...)

$$\sigma \sim \sigma_{\text{BORN}} \left(1 - \alpha_s L^2 + \frac{1}{2} \alpha_s^2 L^4 - \frac{1}{6} \alpha_s^3 L^6 + \dots \right)$$



▶ Demanding a more exclusive description of final states comes with complications

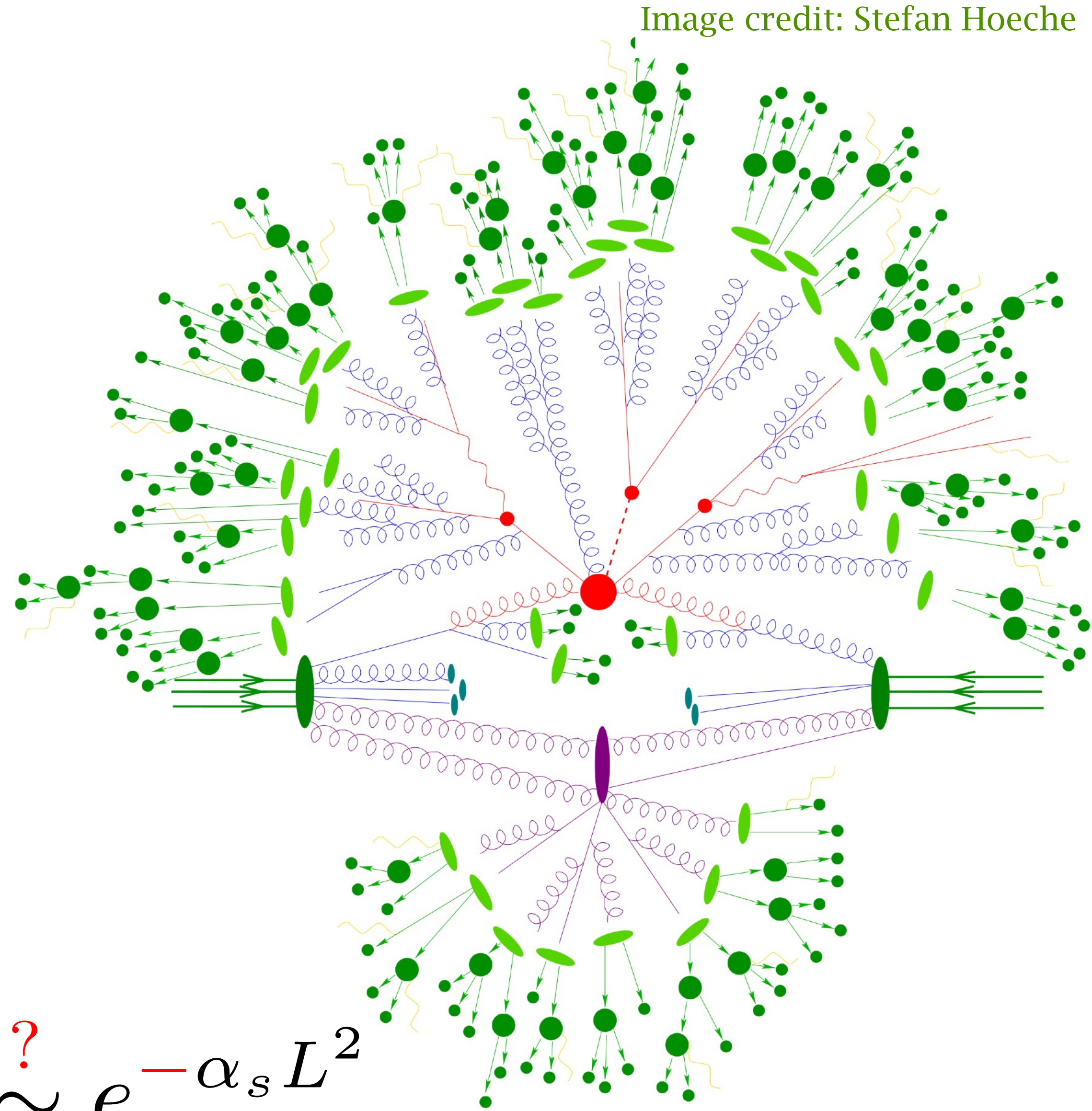
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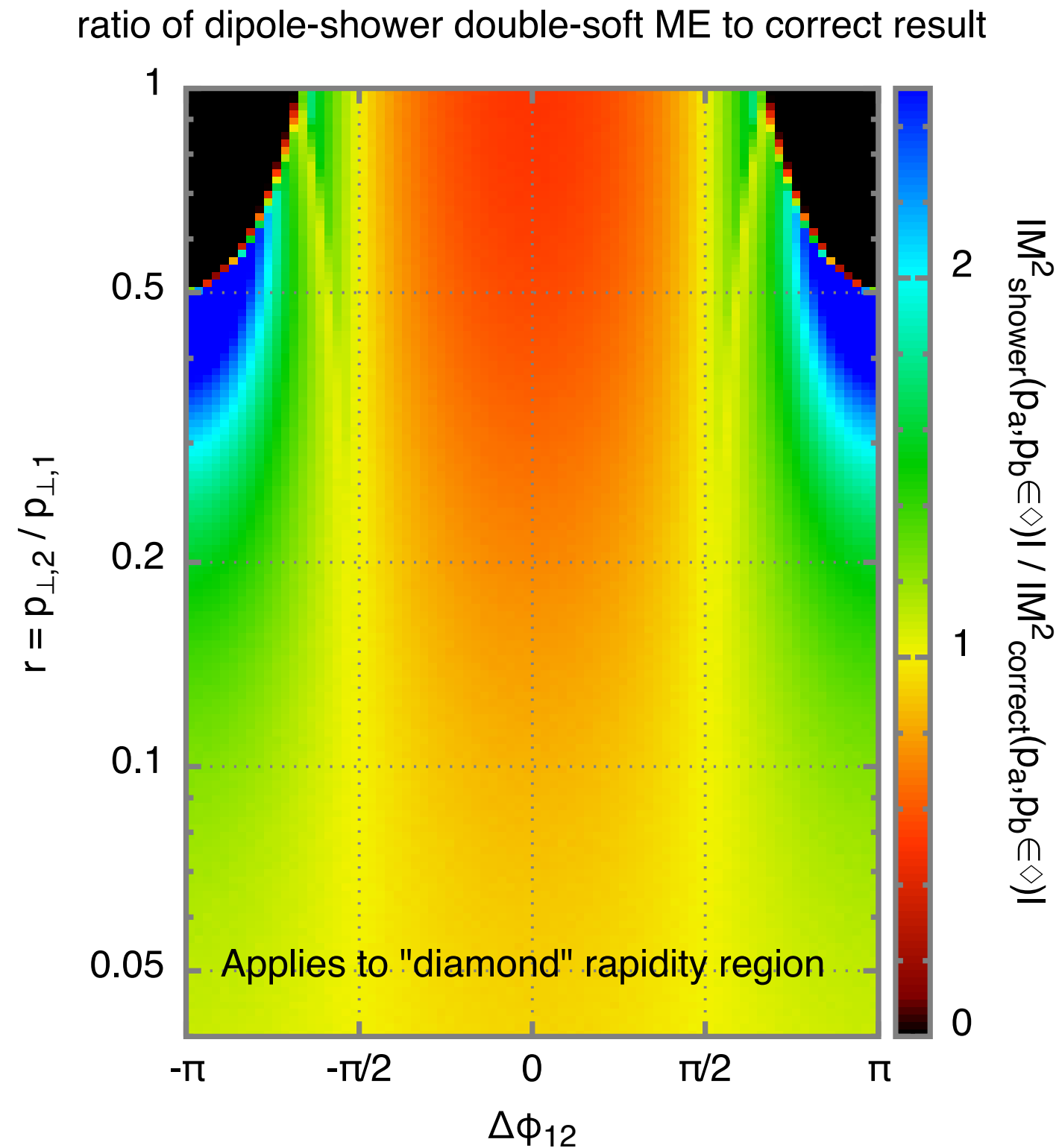


Resummation at the LHC

understanding of multi-scale dynamics plays an important role in several aspects of LHC physics

Design of more accurate event generators & parton shower algorithms

Accurate predictions in infrared sensitive kinematic regimes (differential distributions, jets & their substructure)

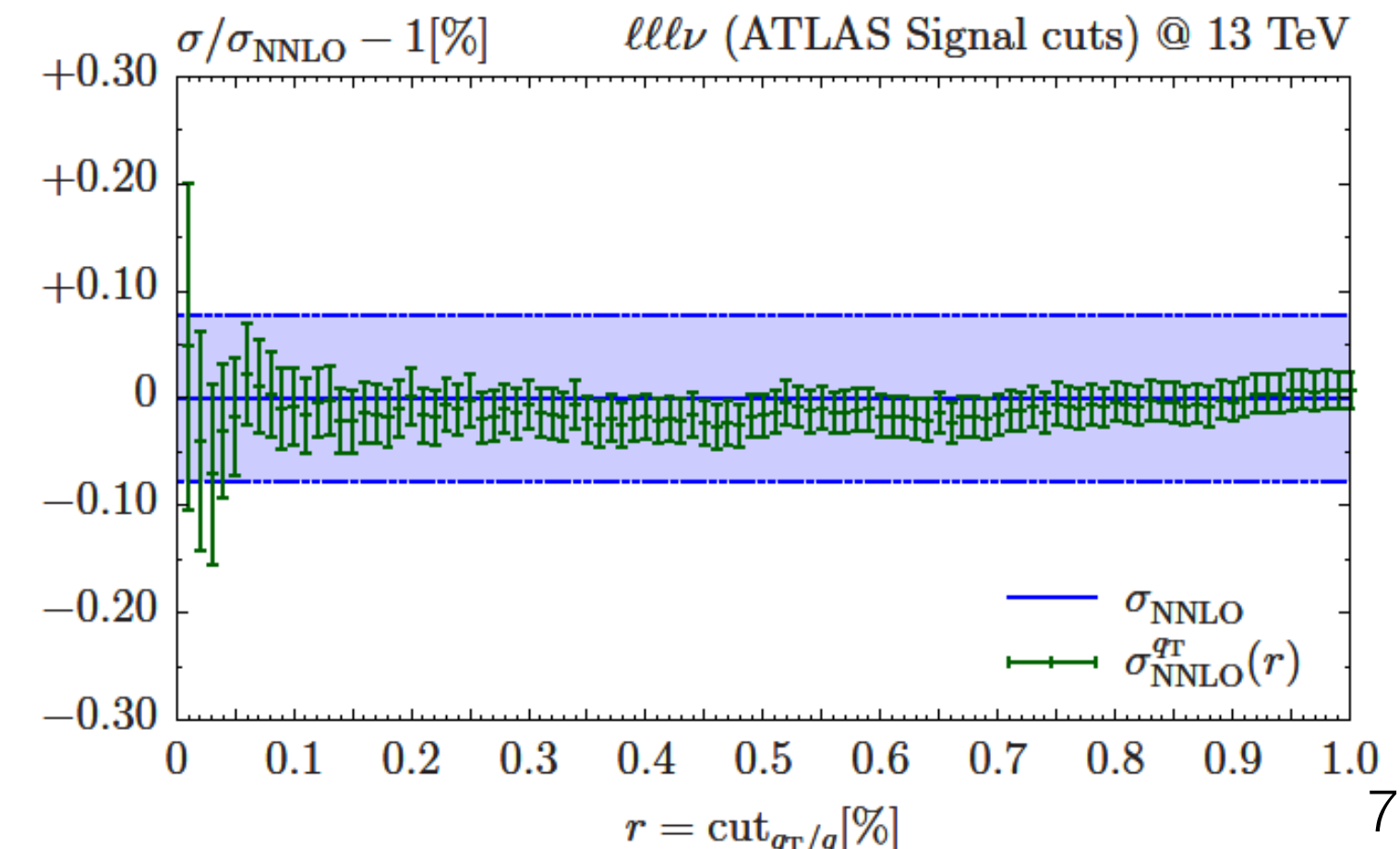
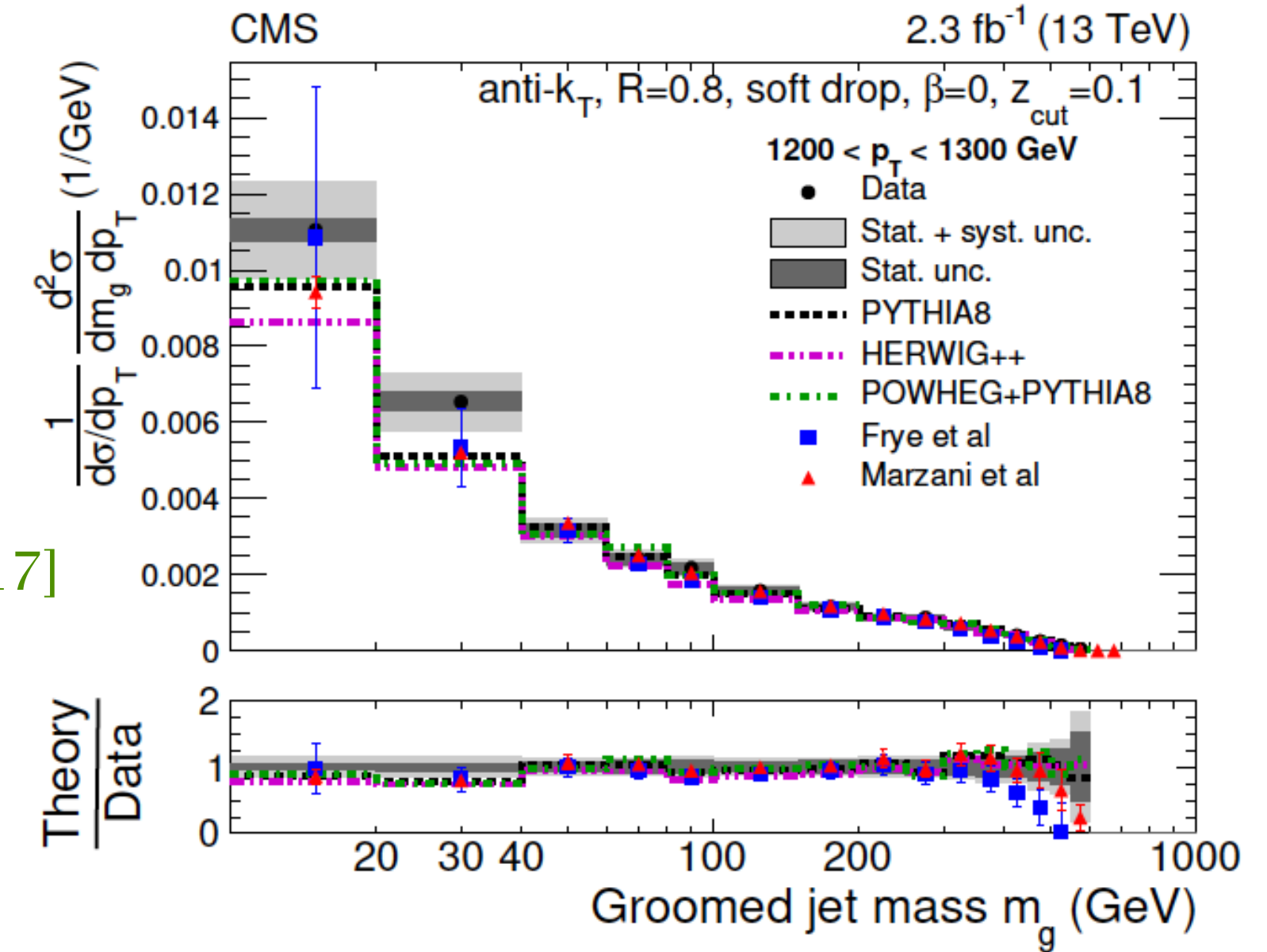
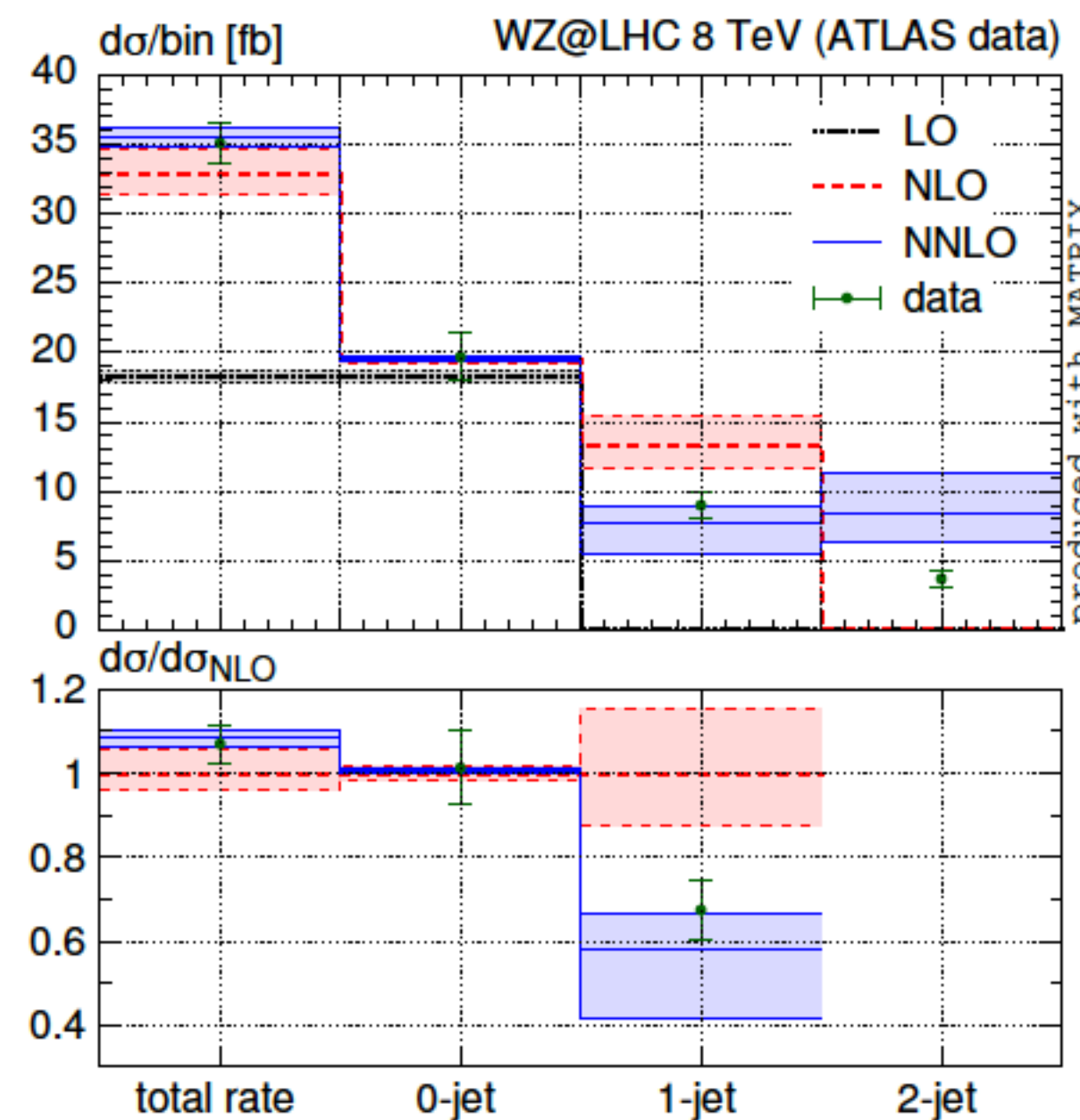


[Dasgupta, Dreyer, Hamilton, PM, Salam '18]

Control over infrared structure of the theory (extraction of parton densities, subtraction methods for higher order computations)

[>> talks by D. Walker & R. Röntsch]

[Grazzini, Kallweit, Rathlev, Wiesemann '17]



Understanding observables with jets

- ▶ Higgs + 0-jet cross section known with very high precision (3-4% uncertainty)

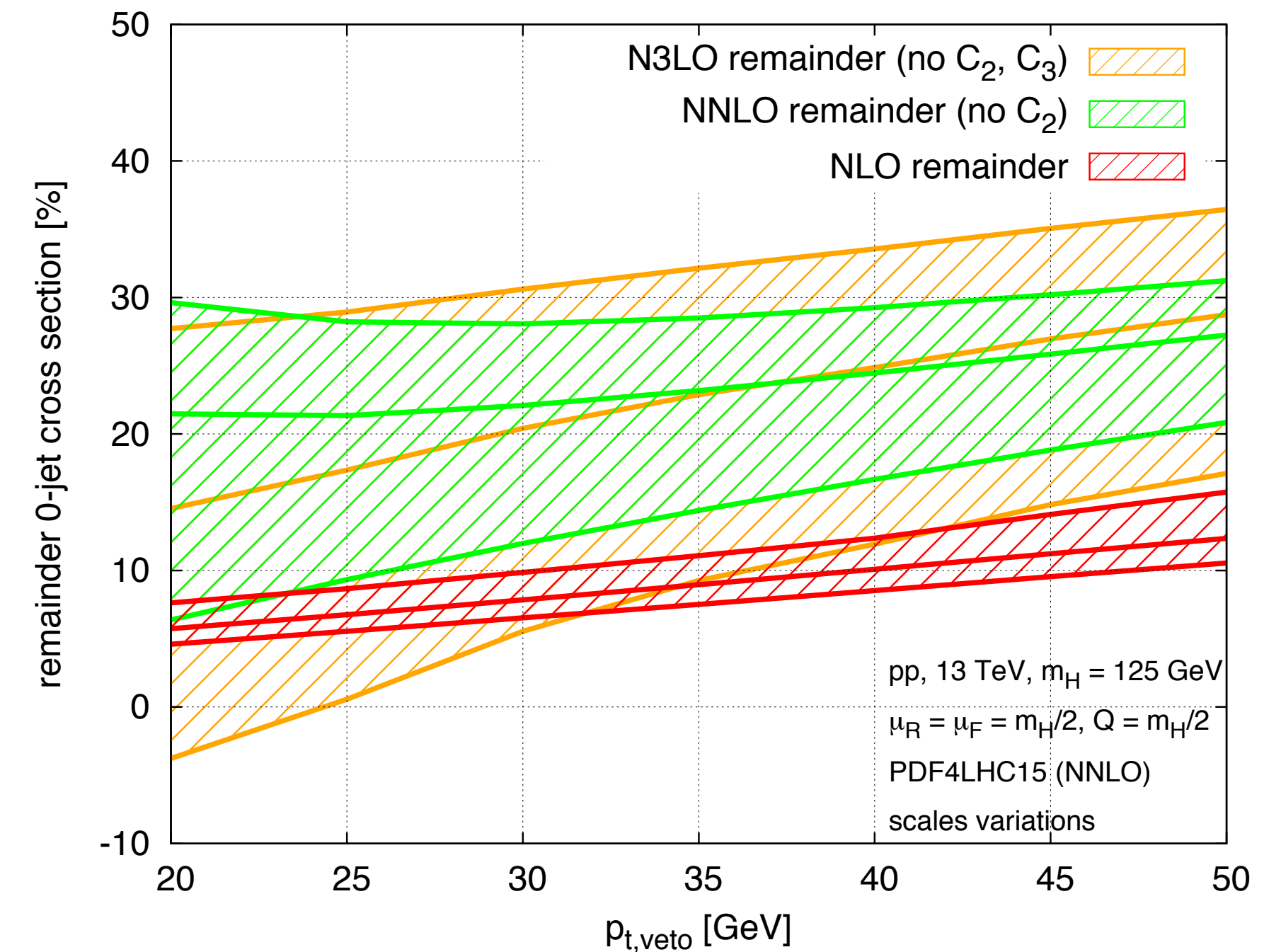
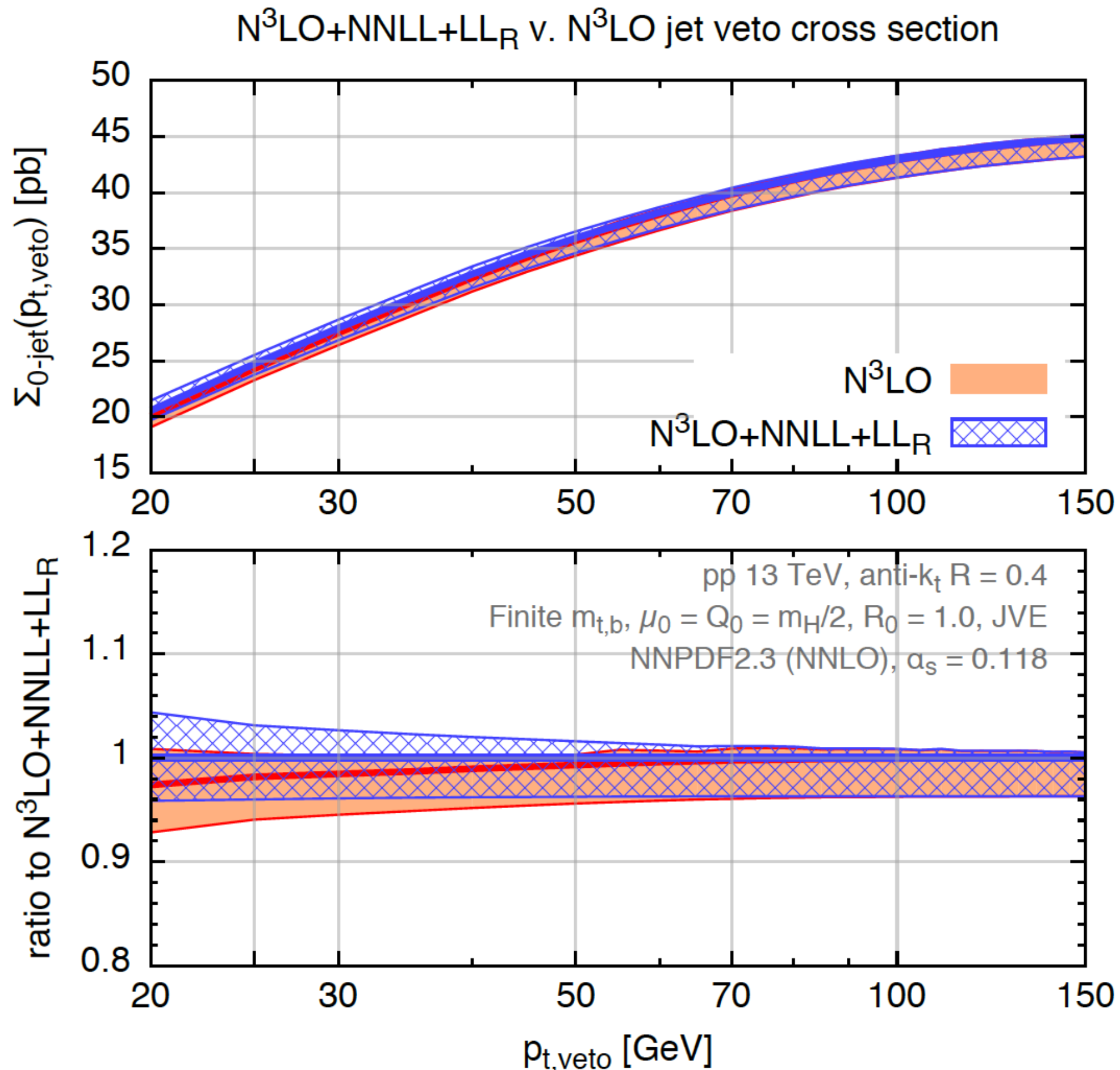
[Banfi, Caola, Dreyer, PM, Salam, Zanderighi, Dulat '15]

see also: [Banfi, PM, Salam, Zanderighi '12]

[Becher, Neubert, Rothen '13]

[Stewart, Tackmann, Walsh, Zuberi '13]

- ▶ ~ 70-90% of the perturbative series at 30 GeV is made of logarithms. Resummation provides the bulk of higher order corrections (beyond N³LO)



- ▶ Possible future improvements from resummation of subleading-power corrections. First steps in this area recently

[Moult, Stewart, Vita '18]

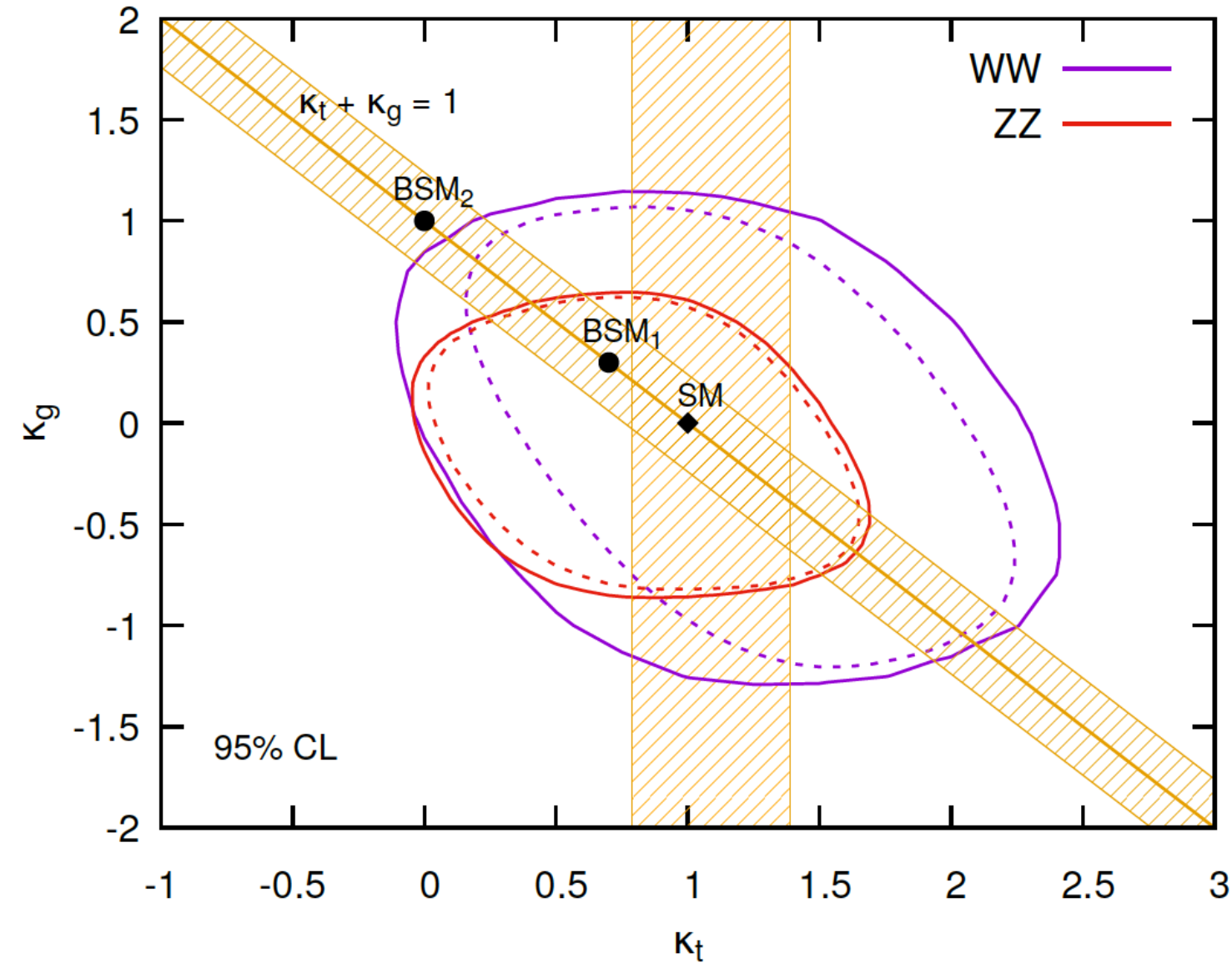
[van Beekveld, Beenakker, Basu, Laenen, Misra, Motylinski '19]

[Bahjat-Abbas, Bonocore, Damste, Laenen, Magnea, Vernazza, White '19]

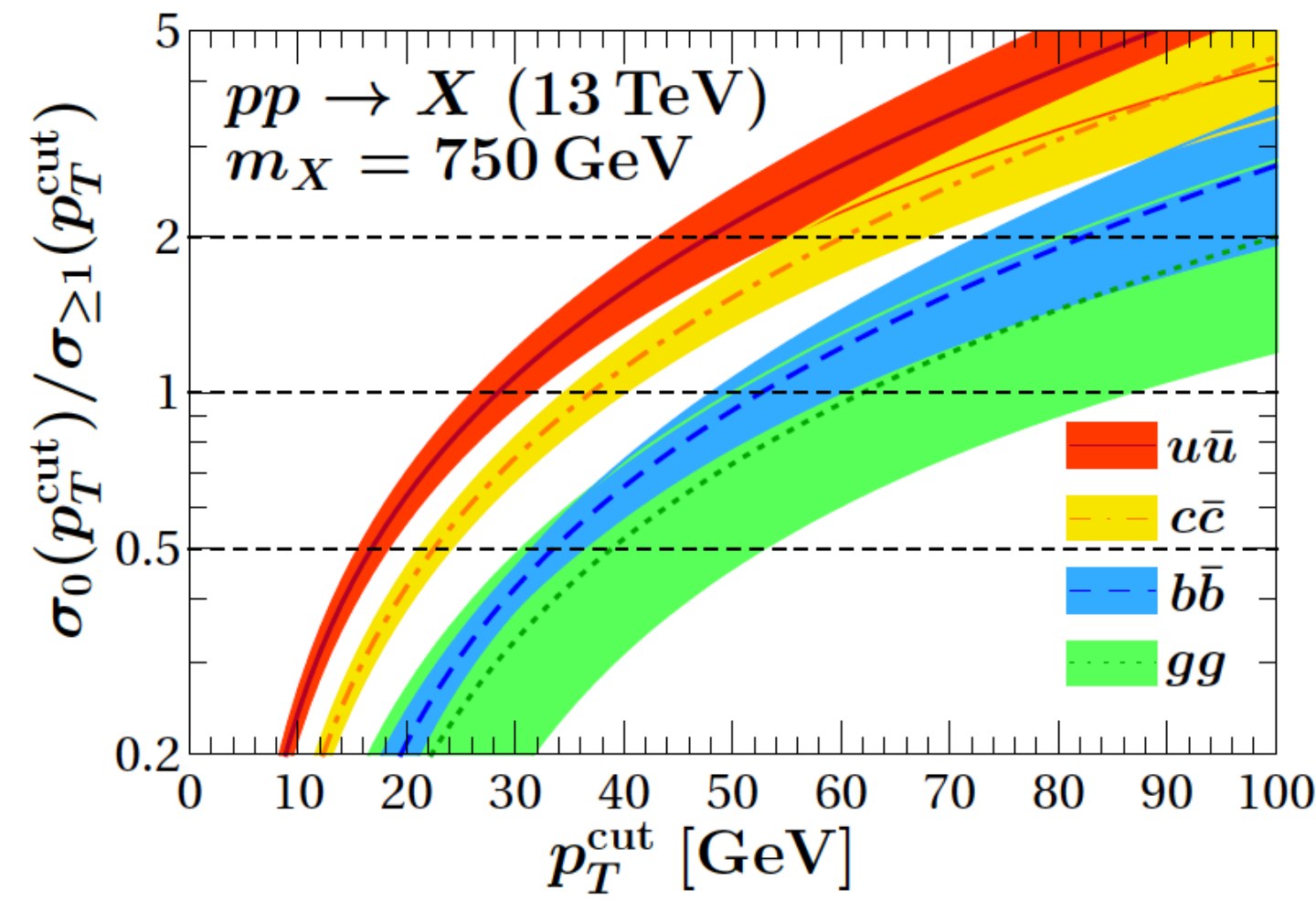
Jet vetoes: recent developments

☑ Exclusion of BSM models

[Arpino, Banfi, Jaeger, Kauer '19]

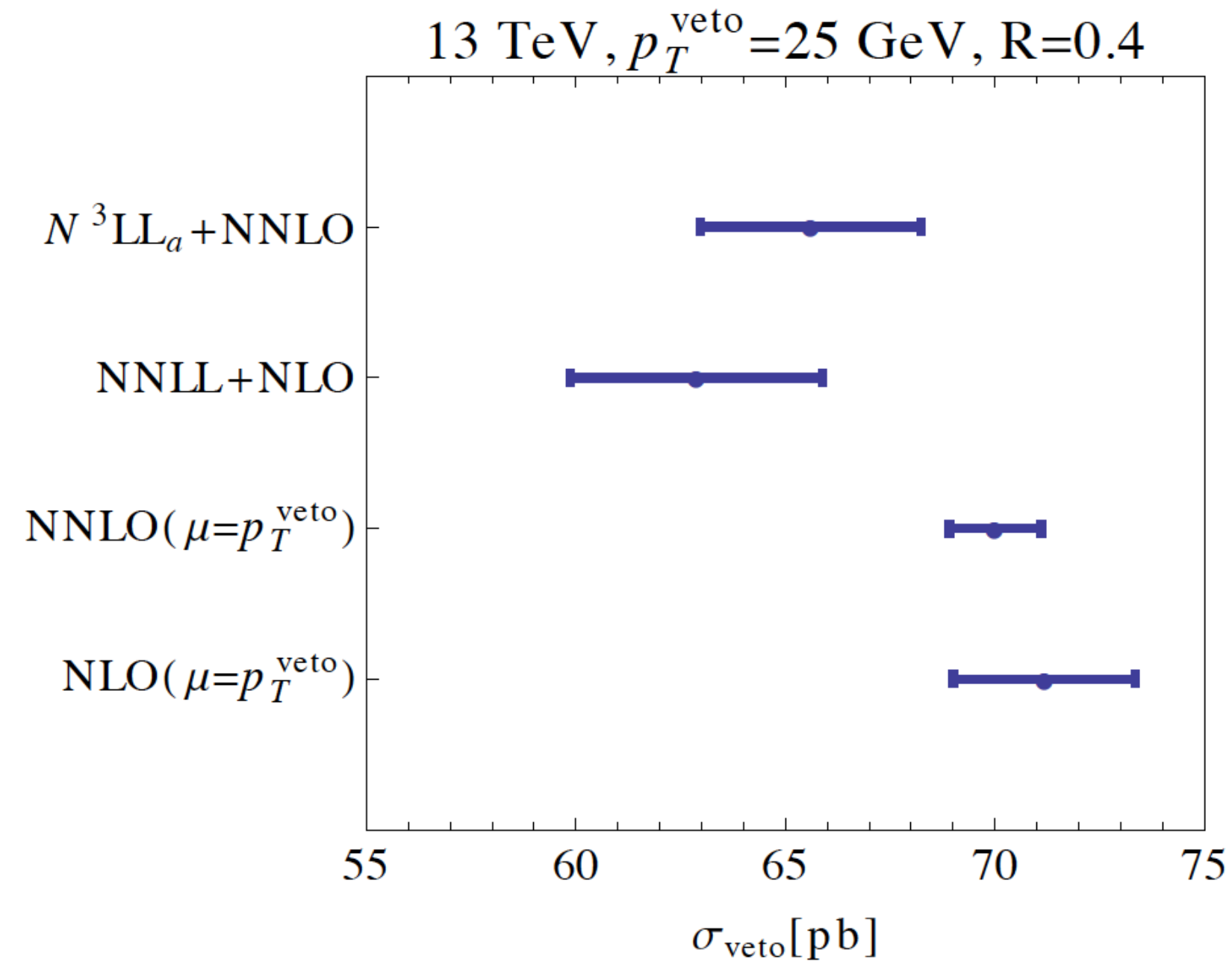


[Ebert, Liebler et al. '16]



☑ Extension to complex final states, e.g. WW

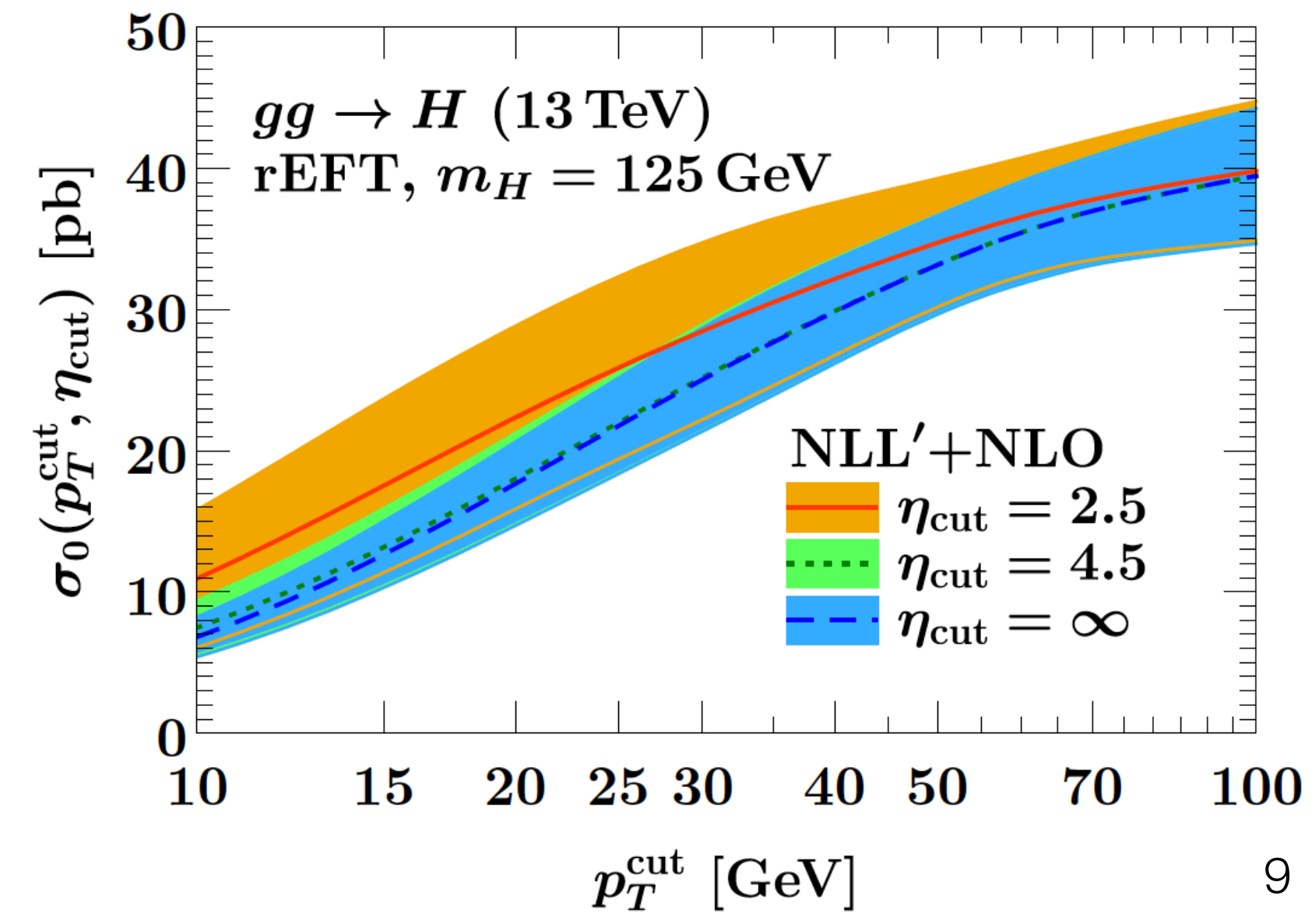
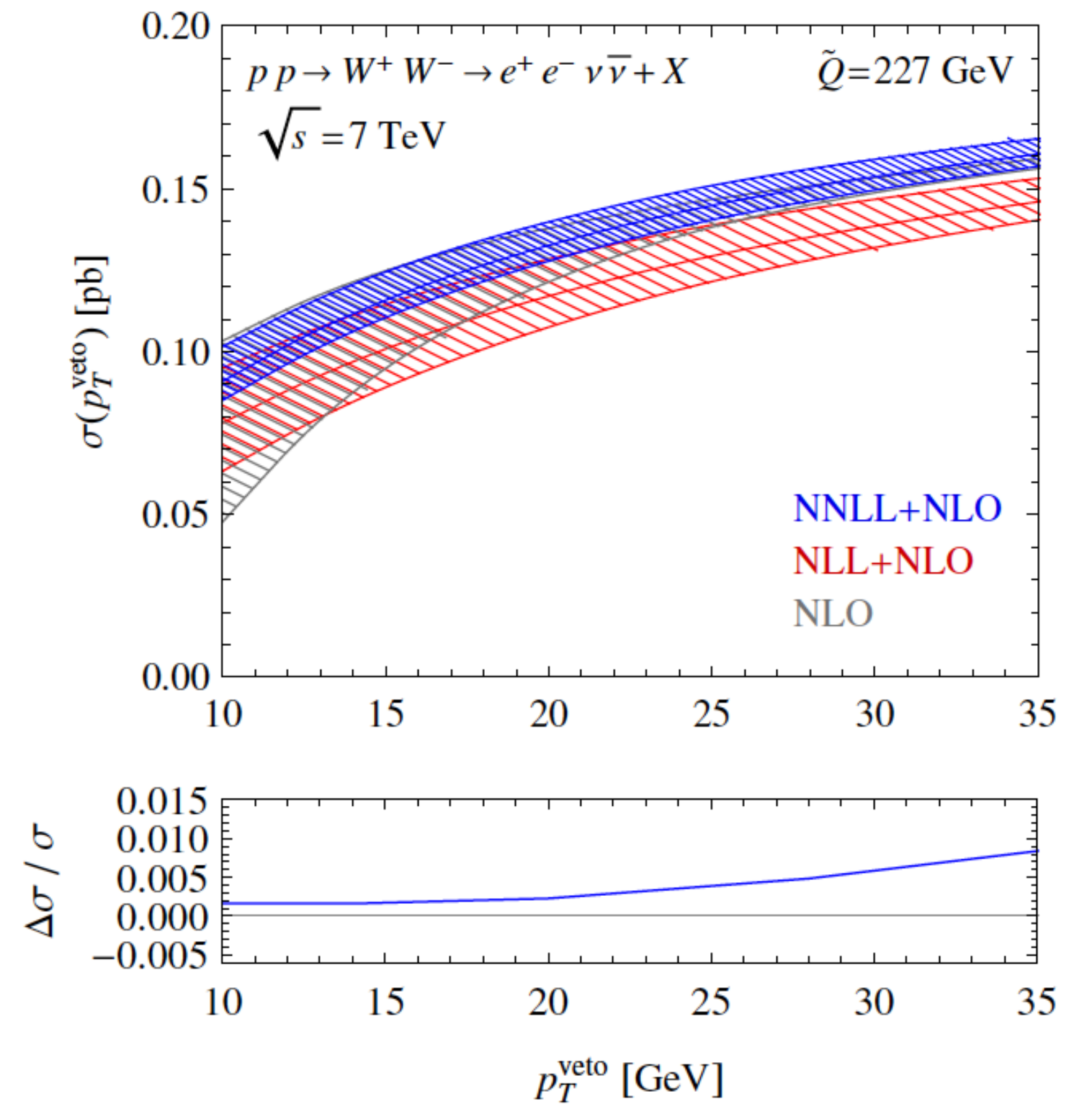
[Dawson, Jaiswal, Li, Ramani, Zeng '16]



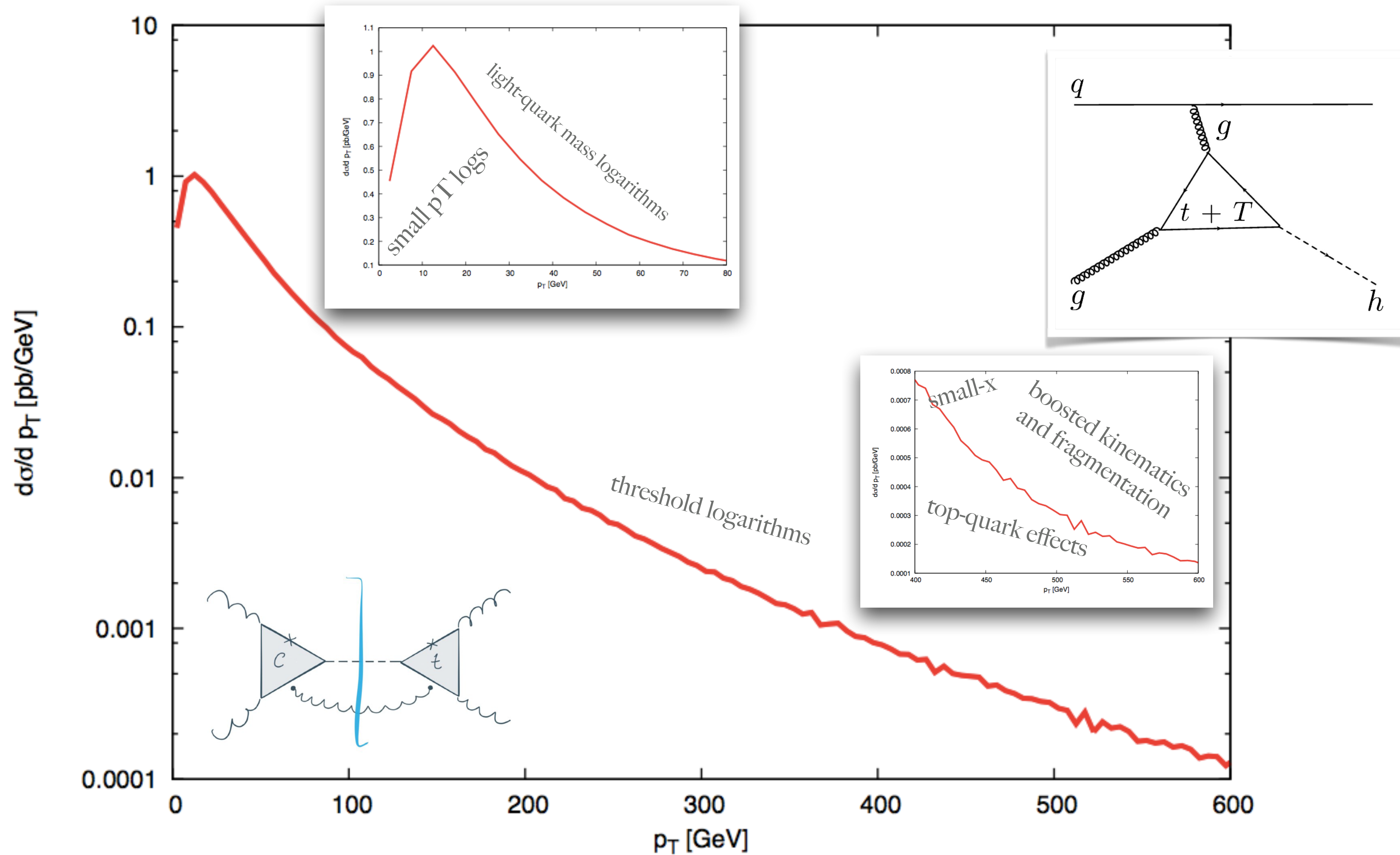
☑ Impact of cuts on leading jet's rapidity

[Michel, Pietrulewicz, Tackmann '18]

[Becher, Frederix, Neubert, Rothen '15]



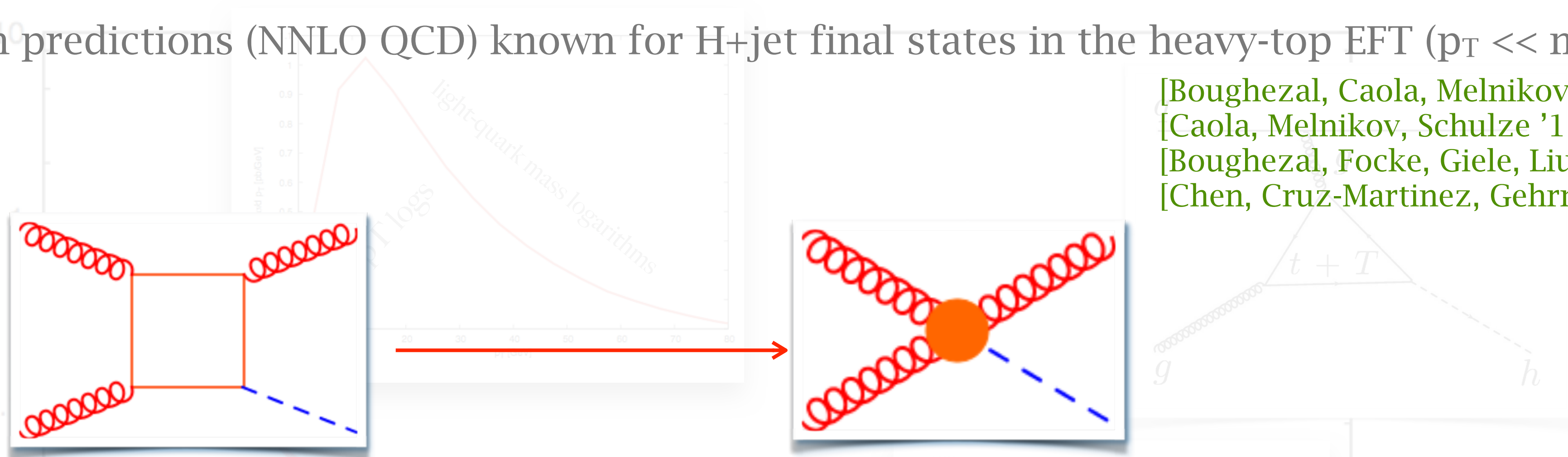
Higgs p_T distribution



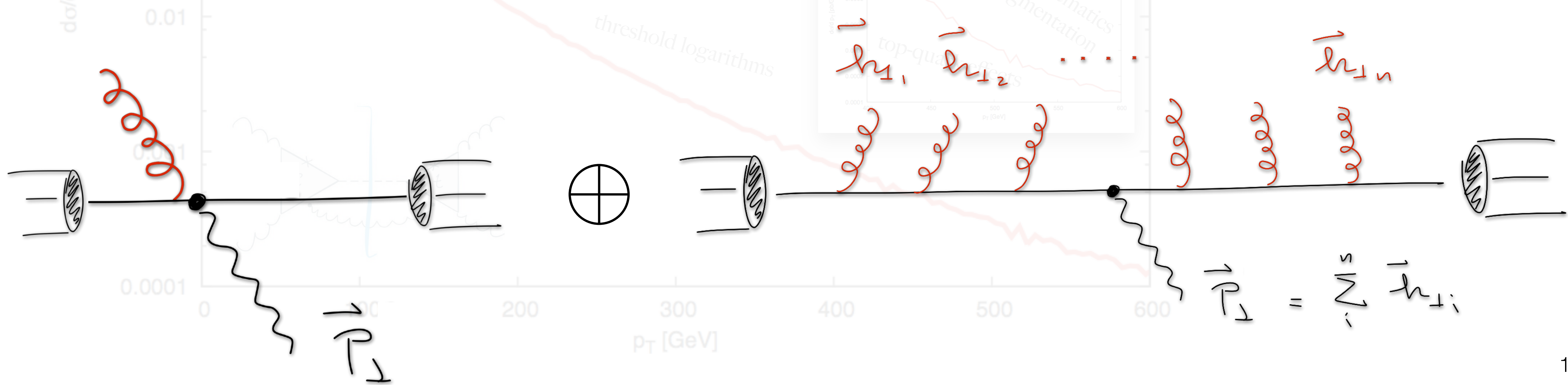
Higgs p_T distribution

- High precision predictions (NNLO QCD) known for H+jet final states in the heavy-top EFT ($p_T \ll m_{\text{top}}$)

[Boughezal, Caola, Melnikov, Petriello, Schulze '15]
 [Caola, Melnikov, Schulze '15]
 [Boughezal, Focke, Giele, Liu, Petriello '15]
 [Chen, Cruz-Martinez, Gehrmann, Glover, Jaquier '16]

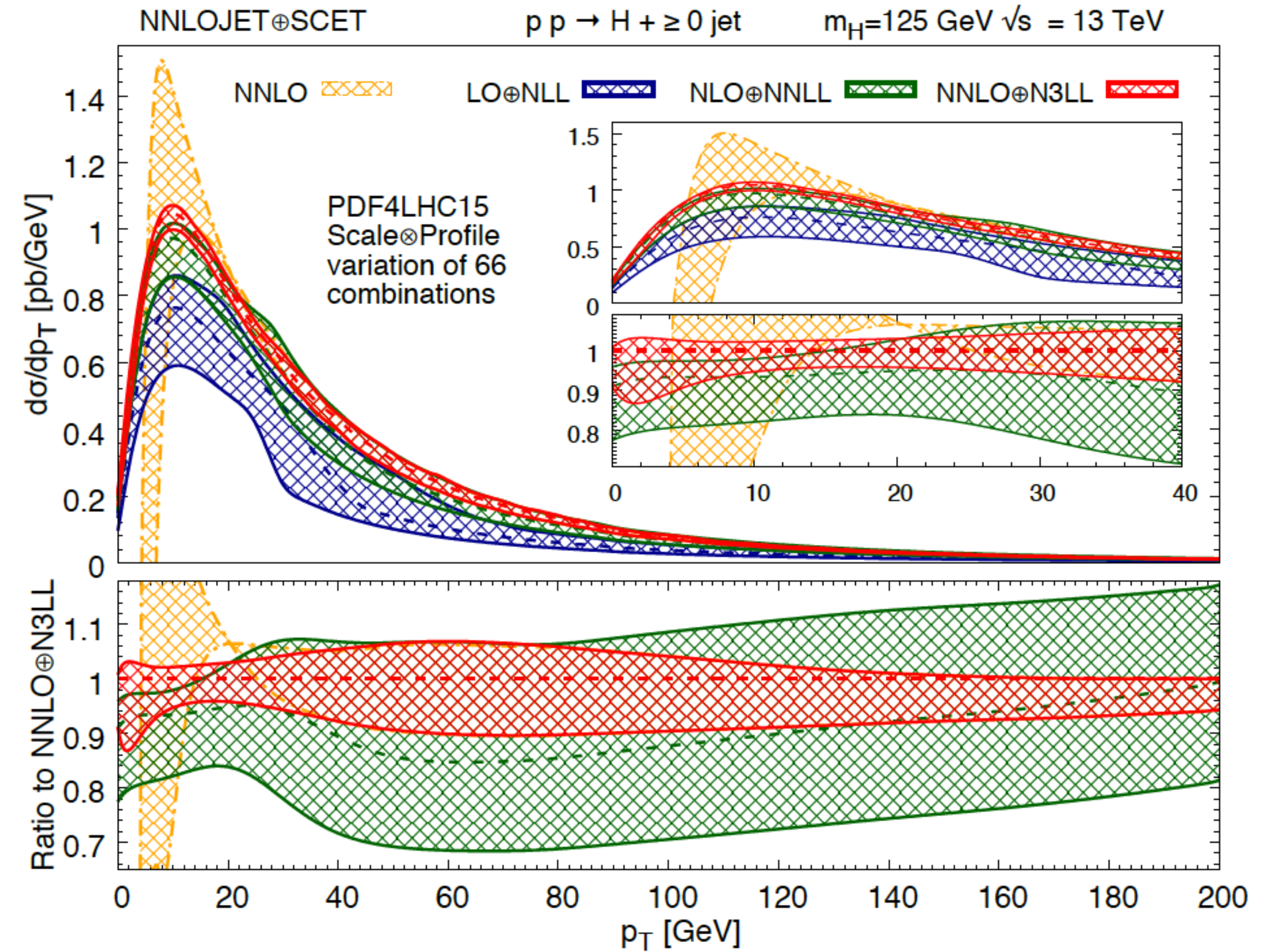
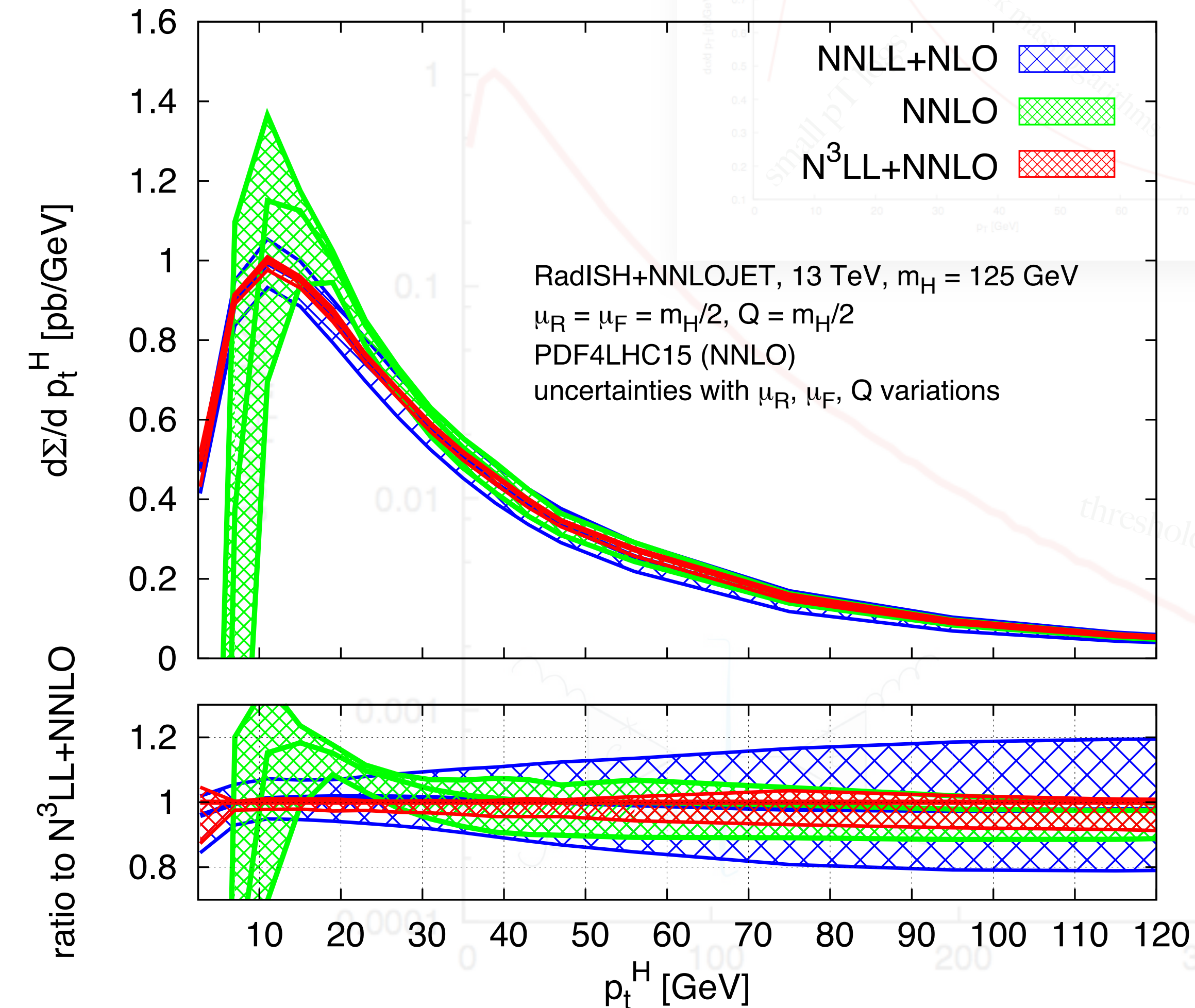


- Matching to resummation allows one to extend the prediction to $p_T \sim 0$, now available up to N³LL



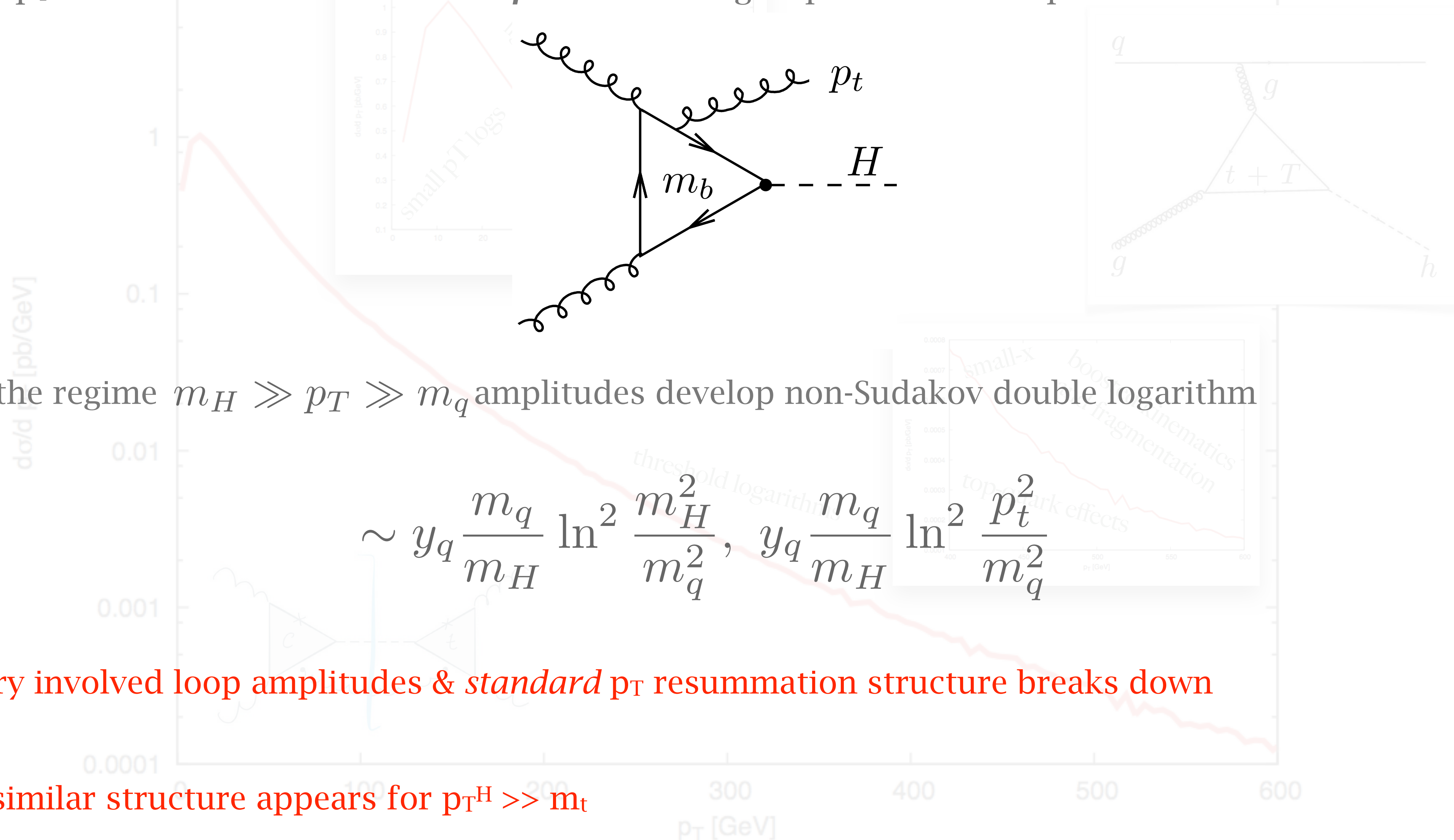
Higgs p_T distribution

- Two independent calculations with different methods: $\sim 5\%$ residual uncertainty in the spectrum
- Good agreement between different matching schemes to fixed order \rightarrow robust control over theory



Higgs p_T distribution: quark masses

- Small p_T distribution sensitive to *interference* with light-quark-mediated production



- In the regime $m_H \gg p_T \gg m_q$ amplitudes develop non-Sudakov double logarithm

- Very involved loop amplitudes & *standard* p_T resummation structure breaks down

- similar structure appears for $p_T^H \gg m_t$

Higgs p_T distribution: quark masses

- Small p_T distribution sensitive to *interference* with light-quark-mediated production

- Full all order structure still unknown, except for double logarithms in the abelian limit

$$\frac{d\sigma_{pp \rightarrow H+j}}{dp_{\perp}^2} = \frac{d\sigma_{pp \rightarrow H+j}^{(0)}}{dp_{\perp}^2} \left\{ 1 - \frac{3m_b^2}{m_H^2} L_{\text{eff}}^2 \left[1 - \frac{x_{\text{eff}}}{12} (1 - \tau^3 + \tau^4) \right. \right. \\ \left. \left. + \frac{x_{\text{eff}}^2}{48} \left(\frac{4}{15} - \tau^3 + 2\tau^4 - \frac{7\tau^5}{5} + \frac{2\tau^6}{5} \right) + \mathcal{O}(x^3) \right] + \mathcal{O}(m_b^4) \right\}$$

$$\tau = \frac{\ln(m_b^2/p_{\perp}^2)}{\ln(m_b^2/m_H^2)}$$

$$L_{\text{eff}} = \ln(m_H^2/m_b^2)$$

$$x_{\text{eff}} = \frac{\alpha_s C_F}{2\pi} L_{\text{eff}}^2$$

[Melnikov, Penin '16]

- Observed good convergence of perturbative expansion due to Yukawa suppression:

- NLO corrections likely sufficient for %-level phenomenology

[Melnikov, Tancredi, Wever '16]

[Lindert, Melnikov, Tancredi, Wever '17]

[>> C. Wever's talk]

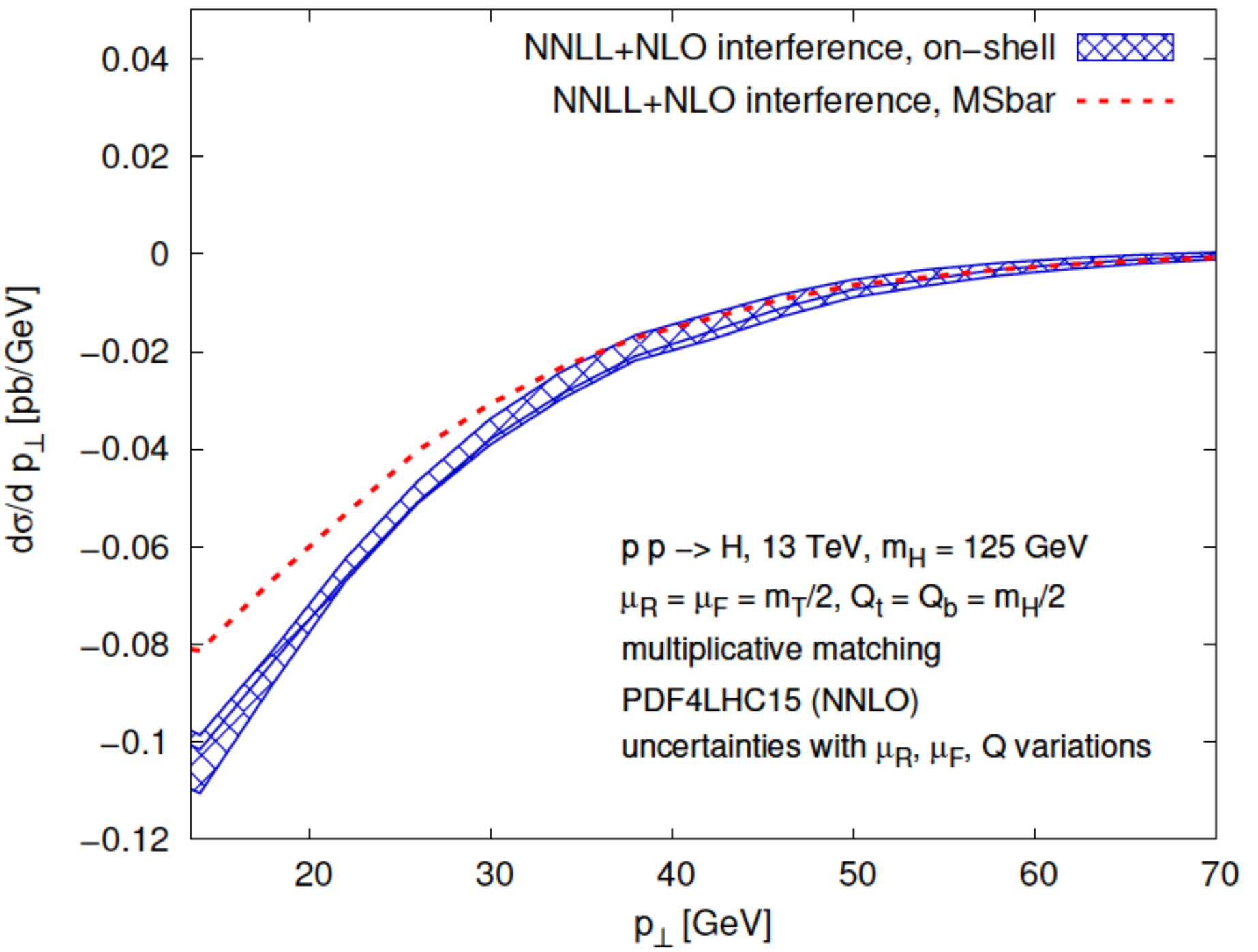
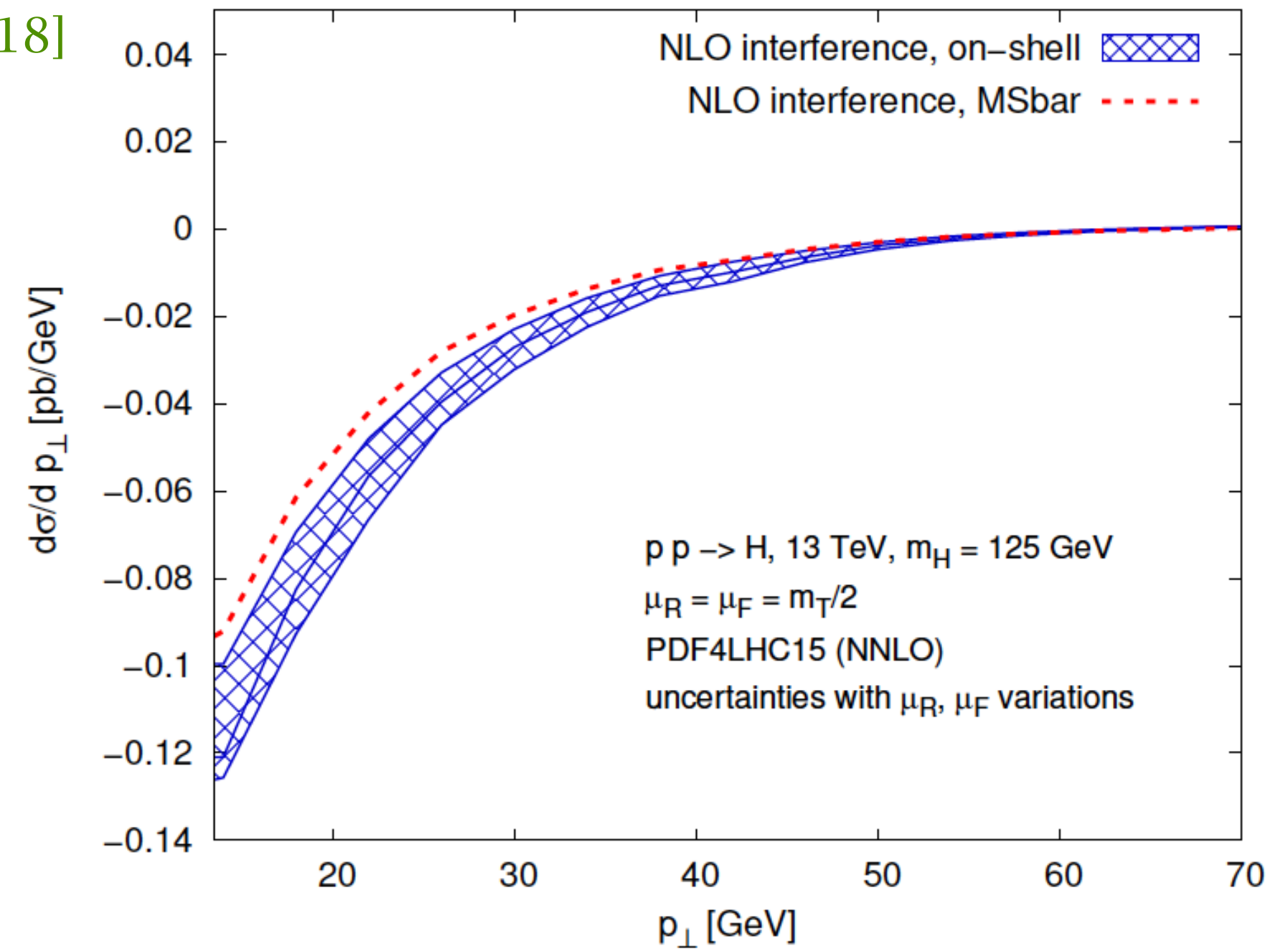
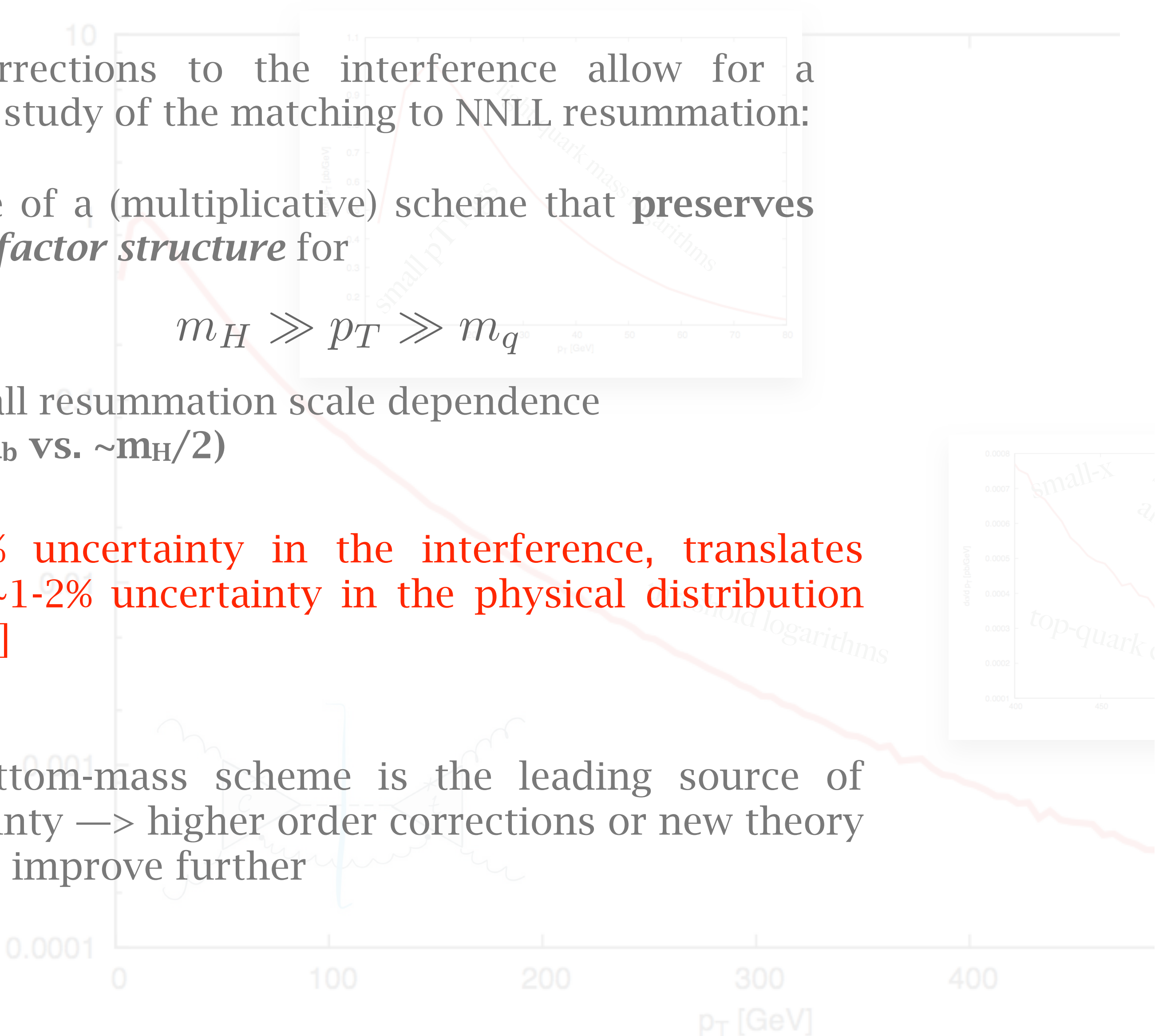
Recently full resummation of virtual corrections in inclusive $gg \rightarrow H$ production

[Liu, Penin '18]

Higgs p_T distribution: quark masses

- ▶ NLO corrections to the interference allow for a detailed study of the matching to NNLL resummation:
- ▶ choice of a (multiplicative) scheme that **preserves form factor structure** for

$$m_H \gg p_T \gg m_q$$
- ▶ small resummation scale dependence ($\sim m_b$ vs. $\sim m_H/2$)
- ▶ **$\sim 15\text{-}20\%$ uncertainty in the interference, translates into a $\sim 1\text{-}2\%$ uncertainty in the physical distribution [backup]**
- ▶ The bottom-mass scheme is the leading source of uncertainty \rightarrow higher order corrections or new theory input to improve further

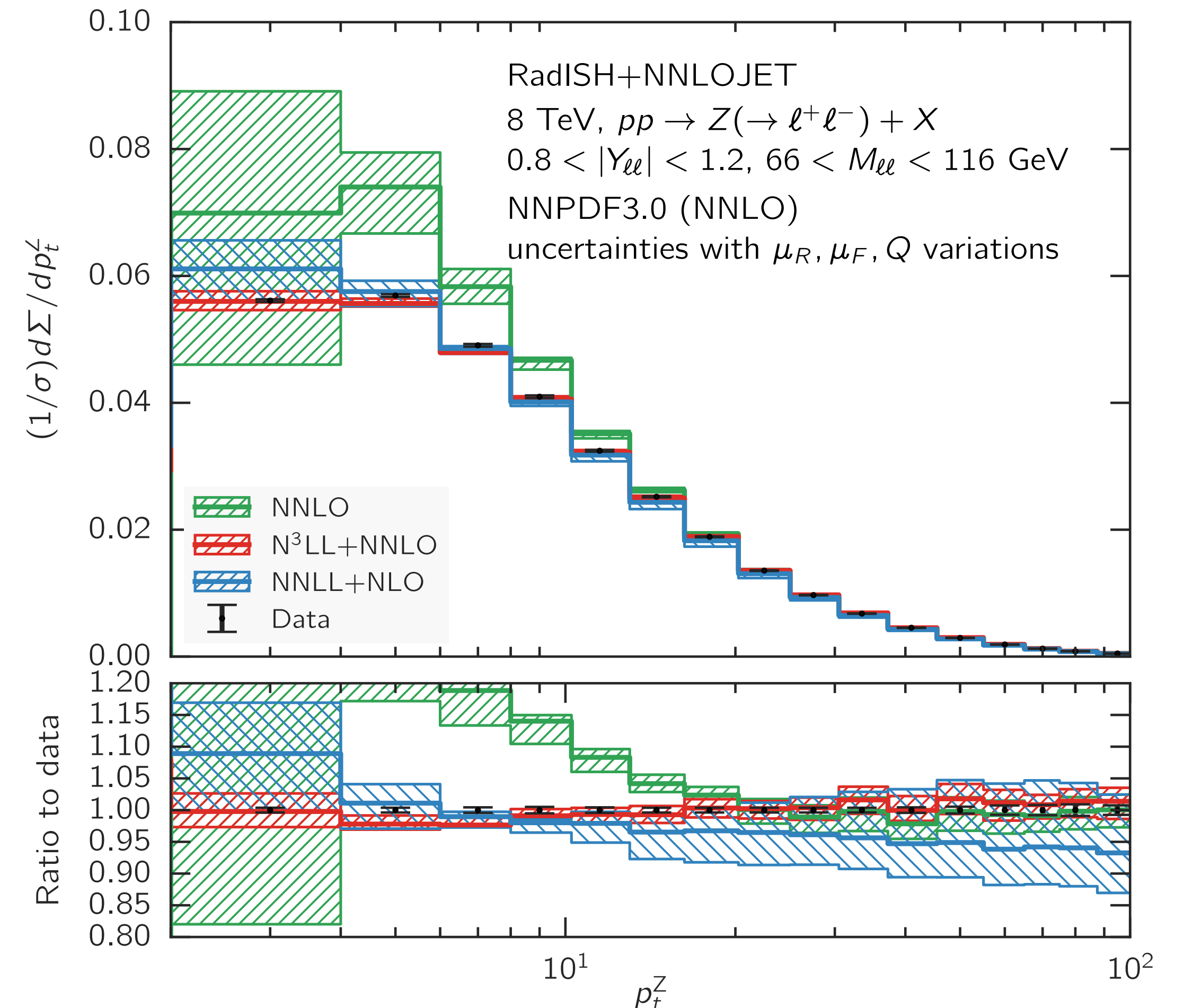


Electro-Weak physics

- ▶ This type of technology can be exploited where experimental precision is highest. E.g. the Z p_T spectrum
- ▶ Precise knowledge of the spectrum is instrumental in the extraction of SM parameters e.g. M_W , strong coupling, parton densities

[>> talks by E. Yatsenko, M. Chiesa, & N. Vranjes]

- ▶ Data and fiducial cuts from [ATLAS 1512.02192]
- ▶ **Scale uncertainties below the 5% level**
 - ▶ Similar findings for the ϕ_η^* angular observable [backup]
- ▶ Below this level of precision many corrections play a role, some of which are of non-perturbative nature

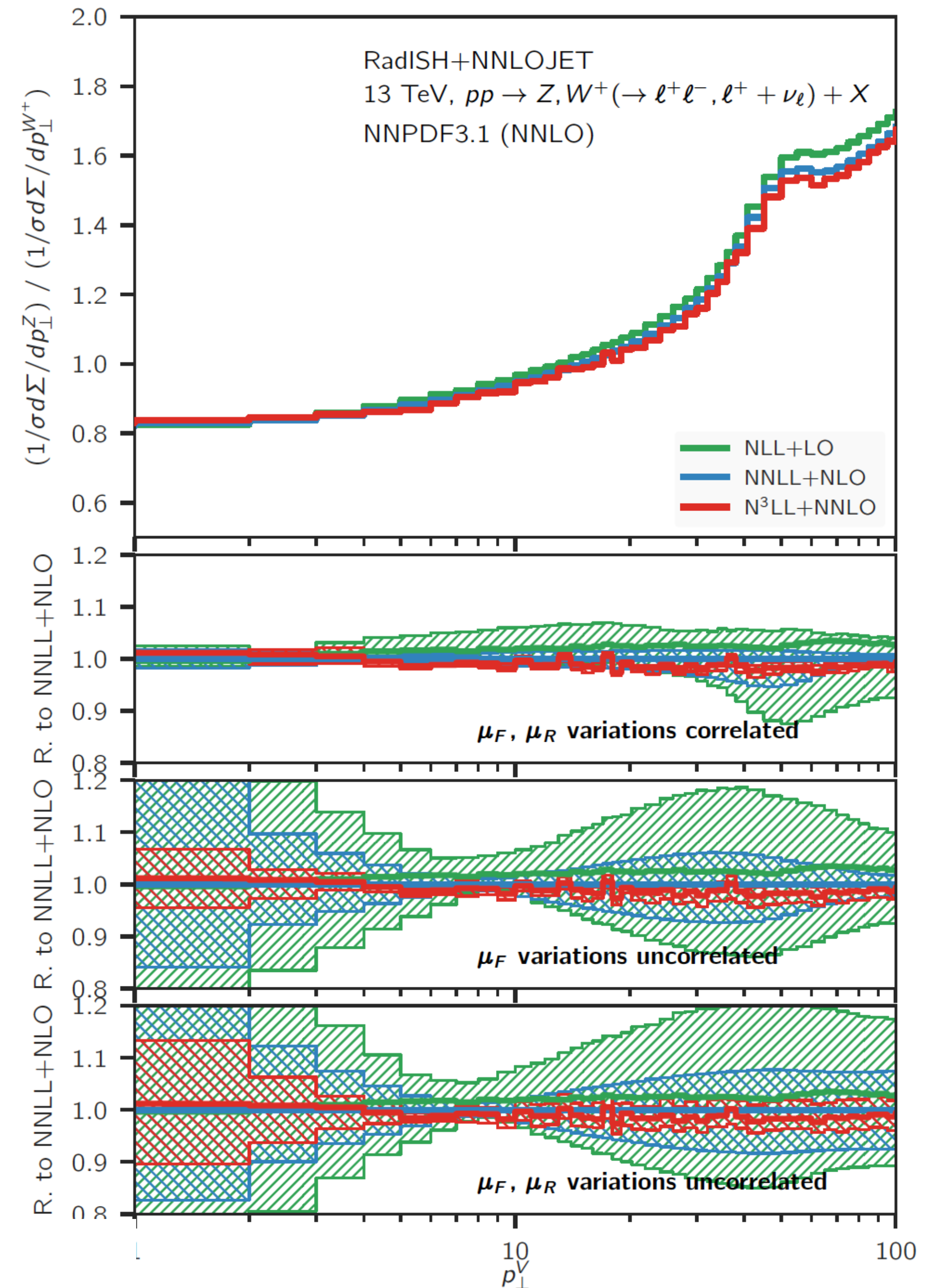


Electro-Weak physics

- Experimental data is sub-percent accurate, we must do better. Rely on data-driven approaches whenever possible

$$\frac{1}{\sigma^W} \frac{d\sigma^W}{dp_\perp} \simeq \frac{1}{\sigma_{\text{data}}^Z} \frac{d\sigma_{\text{data}}^Z}{dp_\perp} \frac{\frac{1}{\sigma_{\text{theory}}^W} \frac{d\sigma_{\text{theory}}^W}{dp_\perp}}{\frac{1}{\sigma_{\text{theory}}^Z} \frac{d\sigma_{\text{theory}}^Z}{dp_\perp}}$$

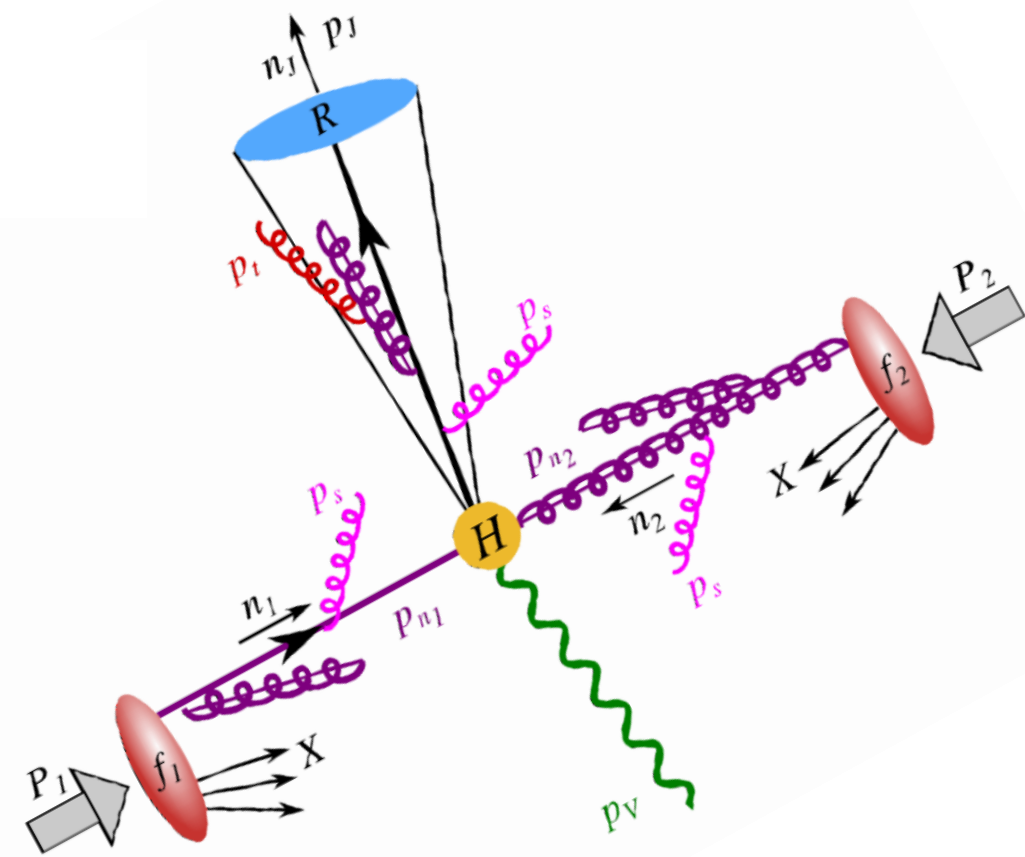
- excellent perturbative convergence for the ratio of distributions !
- % residual uncertainty at N³LL+NNLO (massless QCD + quark thresholds in PDFs)
- Study of other sources of correlation necessary



Going more differential

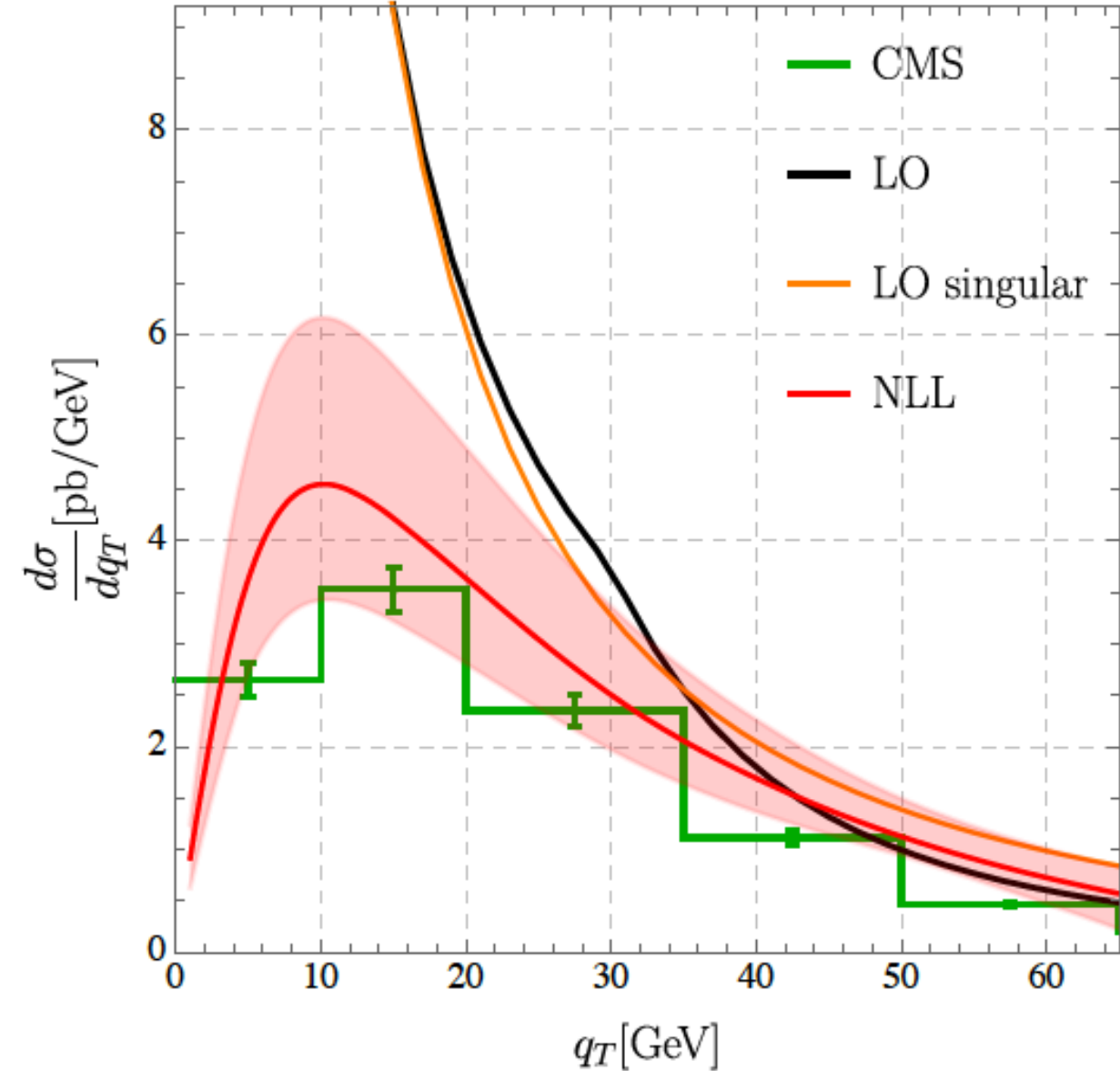
- First steps towards a better understanding of more exclusive and non-global observables at hadron colliders

e.g. Boson-jet correlations
[Chien, Shao, Wu '19]

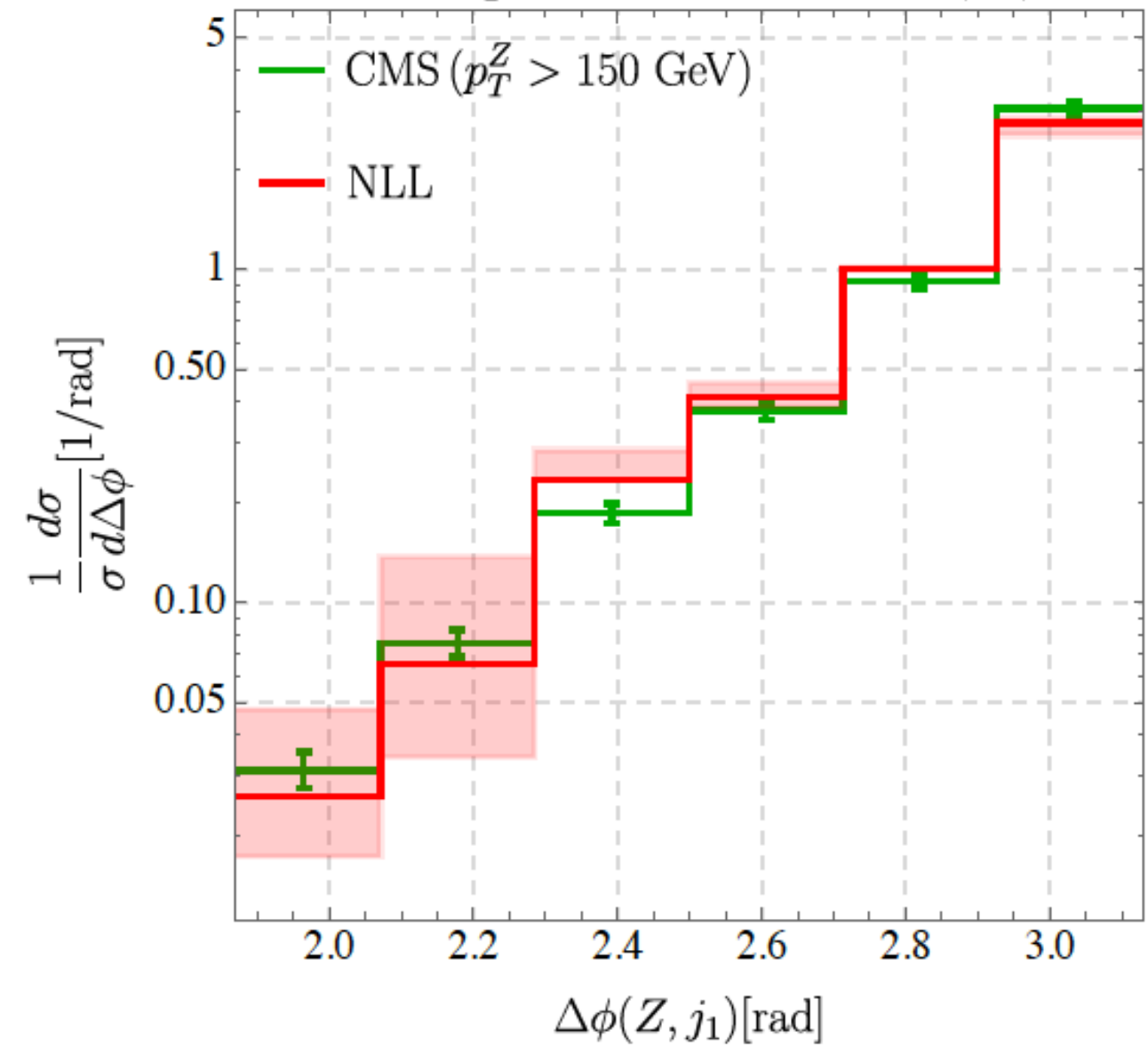


$$\vec{q}_T = \vec{p}_T^J + \vec{p}_T^V$$

LHC13TeV, $p_T^J > 30$ GeV, $R = 0.4$, $|\eta_J| < 2.4$



LHC7TeV, $p_T^J > 50$ GeV, $R = 0.5$, $|\eta_J| < 2.5$

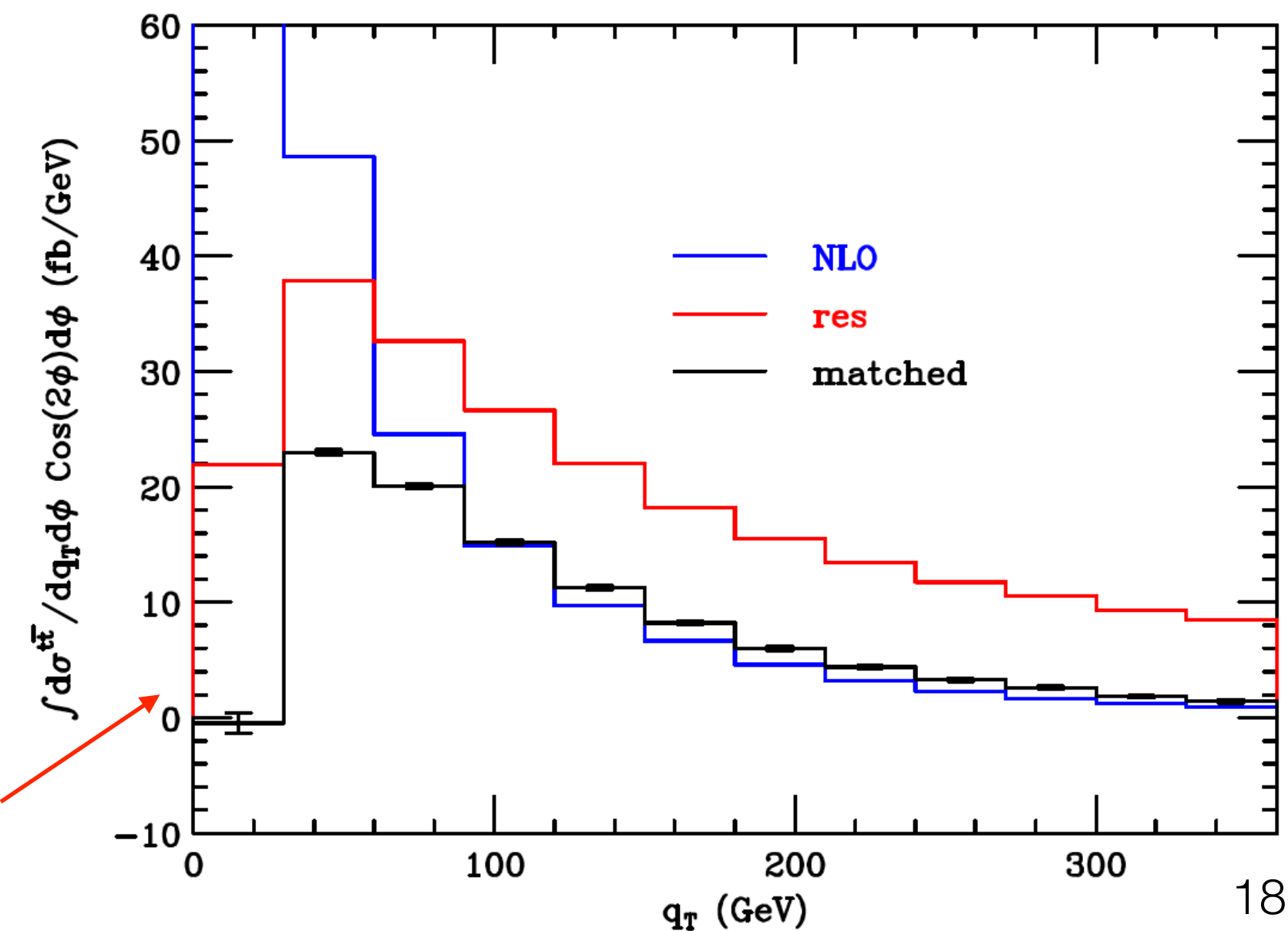


e.g. azimuthal correlations in the production of pairs of heavy particles
[Catani, Grazzini, Sargsyan '17]

$$\frac{d\sigma^{NLO}}{dM^2 d^2\mathbf{q}_T} \propto \delta^{(2)}(\mathbf{q}_T) + \alpha_S \left\{ \left(a_2 \left[\frac{1}{q_T^2} \ln \left(\frac{M^2}{q_T^2} \right) \right]_+ + a_1 \left[\frac{1}{q_T^2} \right]_+ + a_0 \delta^{(2)}(\mathbf{q}_T) + \frac{a_{\text{corr}}(\hat{\mathbf{q}}_T)}{q_T^2} \right) + \dots \right\}, \langle a_{\text{corr}}(\hat{\mathbf{q}}_T) \rangle_{\text{av.}} = 0$$

resummation at small p_T instrumental to make correlations finite. e.g. $t\bar{t}$ production

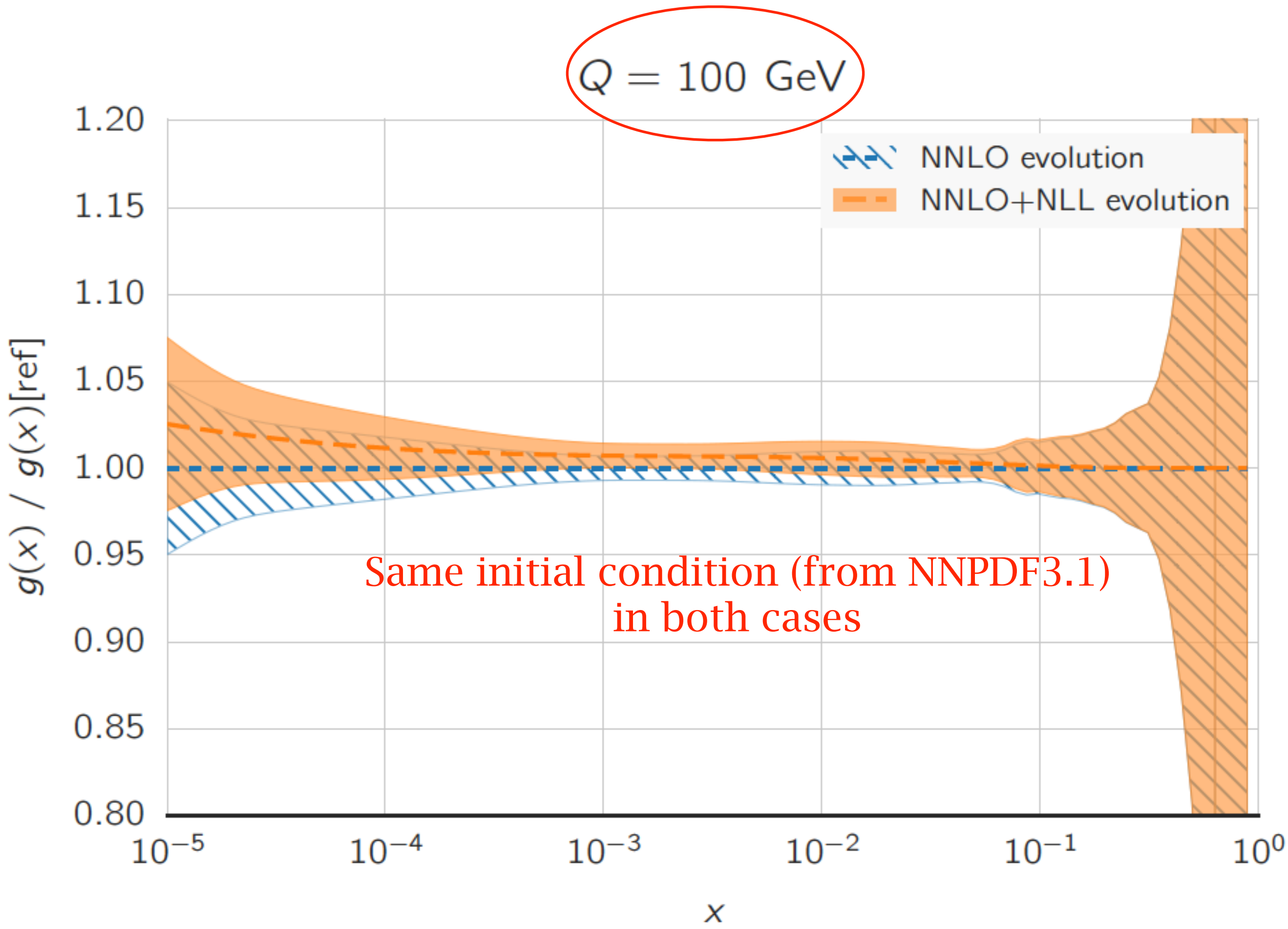
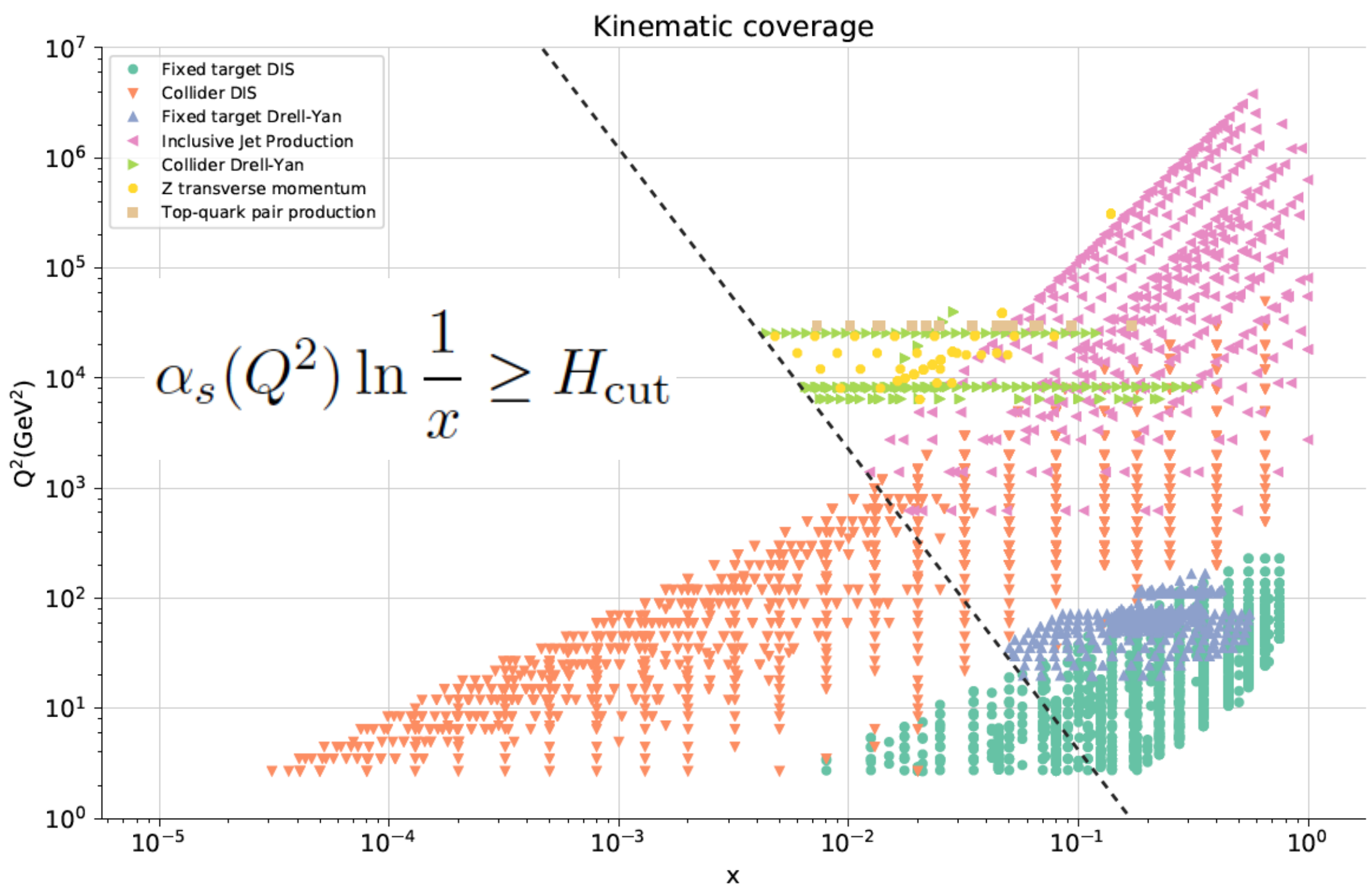
$$\frac{d\sigma_n^{(\text{res})}}{dM^2 dq_T^2} \propto q_T^n$$



Fundamental parameters: PDFs

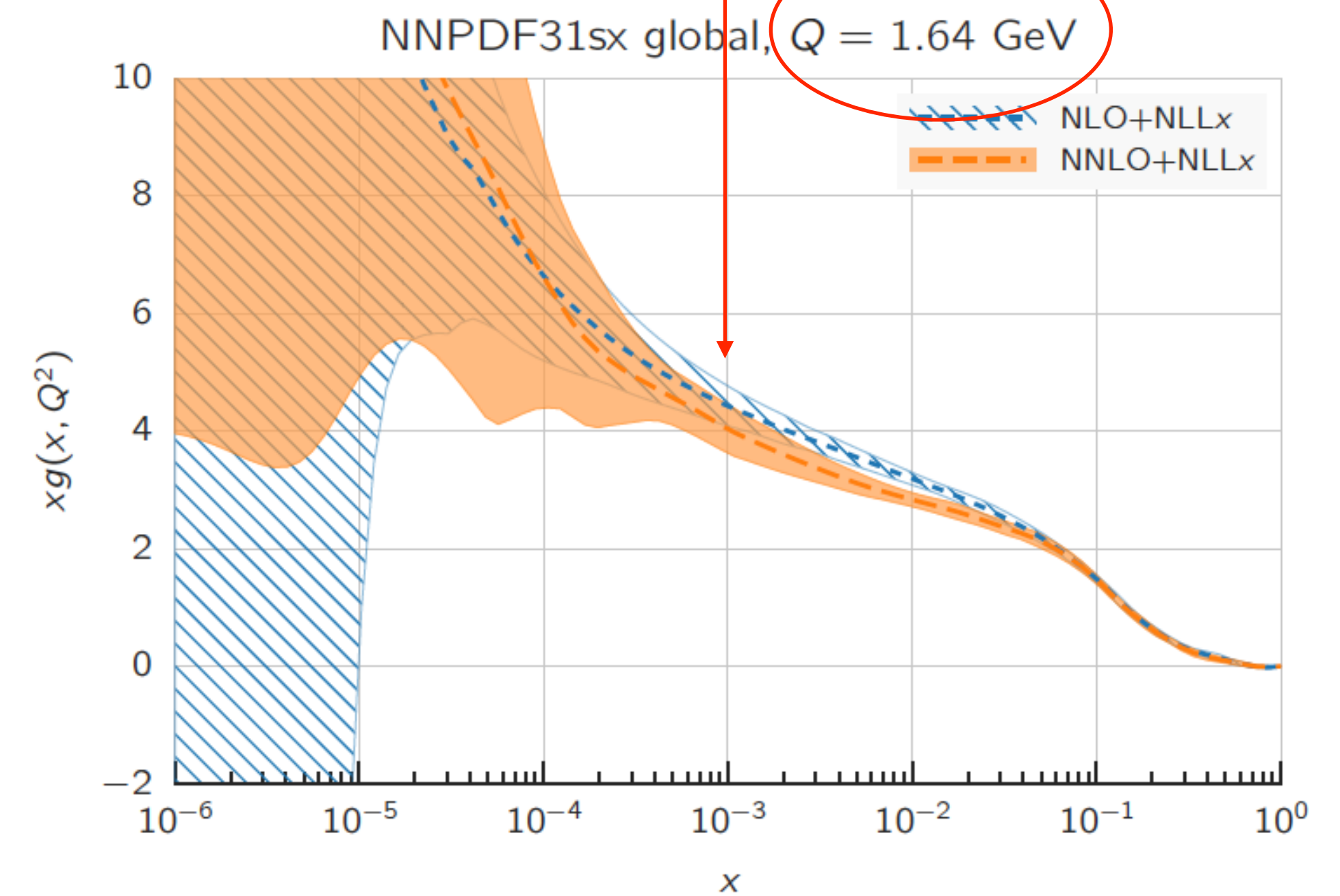
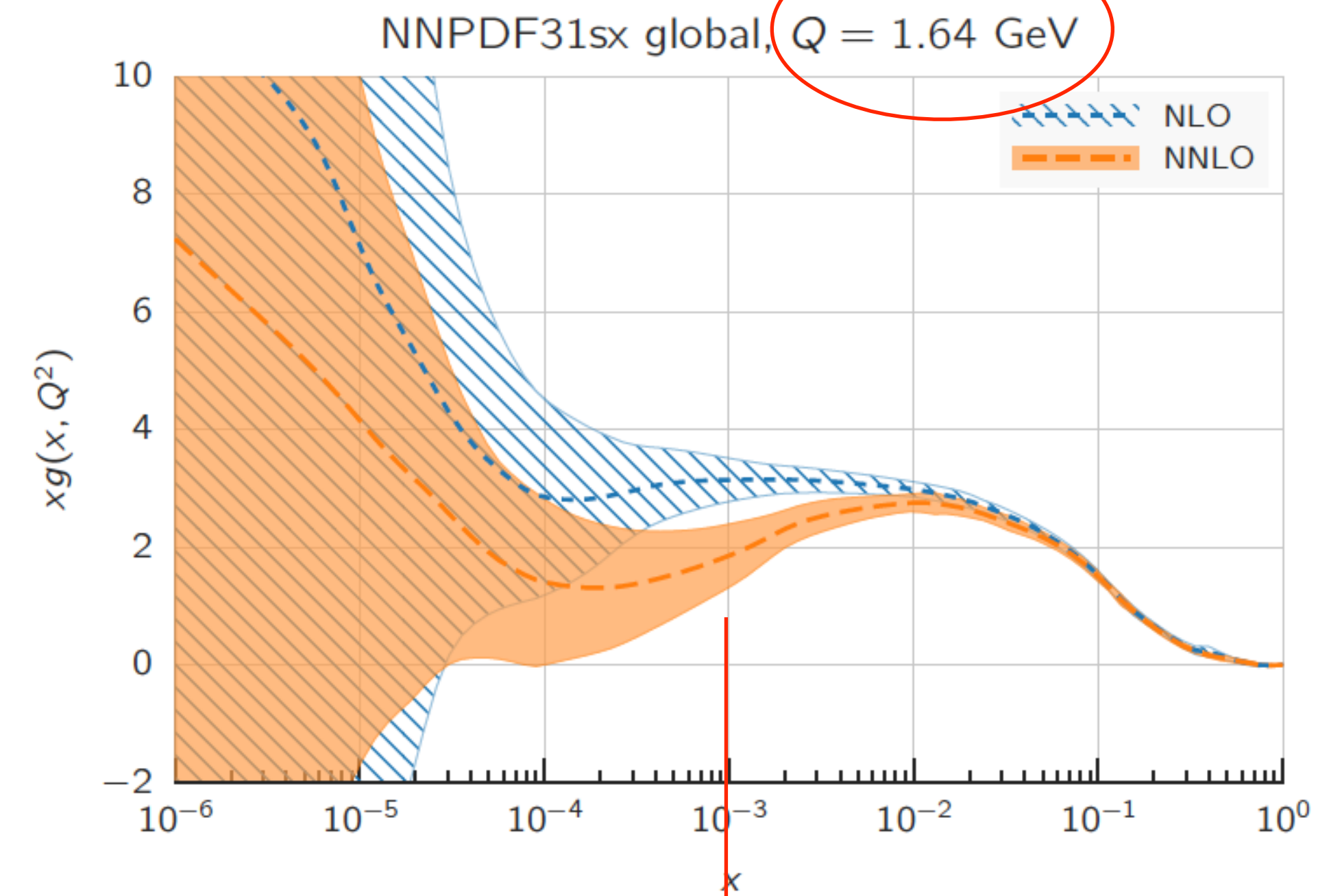
[Ball, Bertone, Bonvini, Marzani, Rojo, Rottoli '18]

- Modern PDF fits contain small-x data from HERA (mainly) and LHCb
- One may wonder how important the impact of the resummation of $\ln(x)$ is
- small-x resummation leads to moderate corrections to the PDF evolution to high scales ...



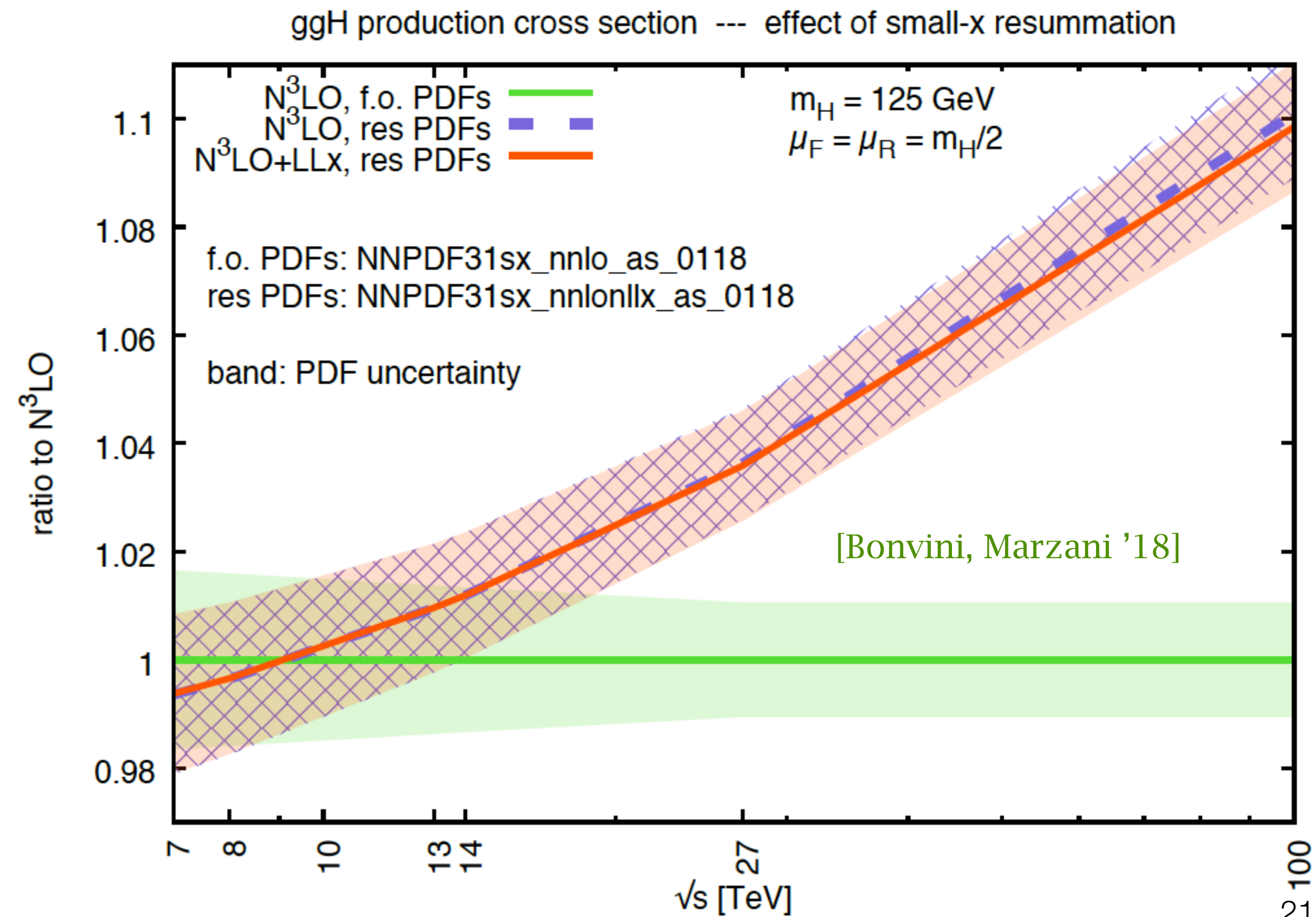
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- ▶ ... however, the initial condition at the low scale may change substantially



Higgs cross section at high energies

- ▶ Modern PDF fits contain small-x data from HERA (mainly) and LHCb
- ▶ One may wonder how important the impact of the resummation of $\ln(x)$ is
- ▶ small-x resummation leads to moderate corrections to the PDF evolution to high scales ...
- ▶ ... however, the initial condition at the low scale may change substantially
- ▶ Sizeable corrections to high-energy processes. e.g. Higgs cross section at 27-100 TeV 4-10% larger than N^3LO
- ▶ Difficult to test these conclusions at current LHC experiments (Drell-Yan at LHCb ?)

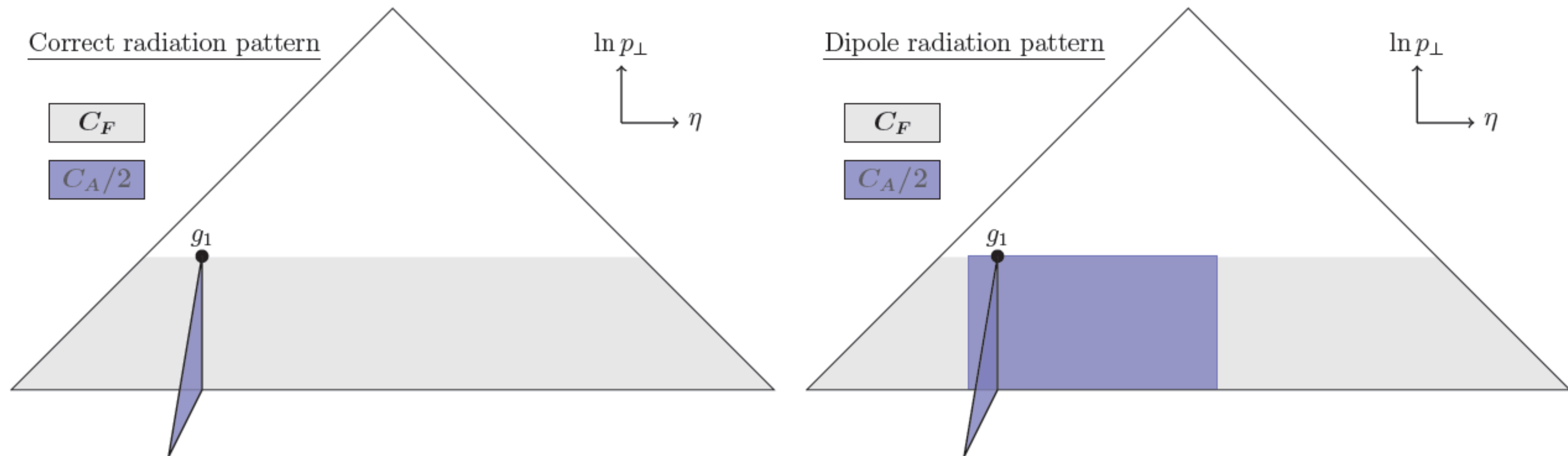


Understanding parton showers from resummations

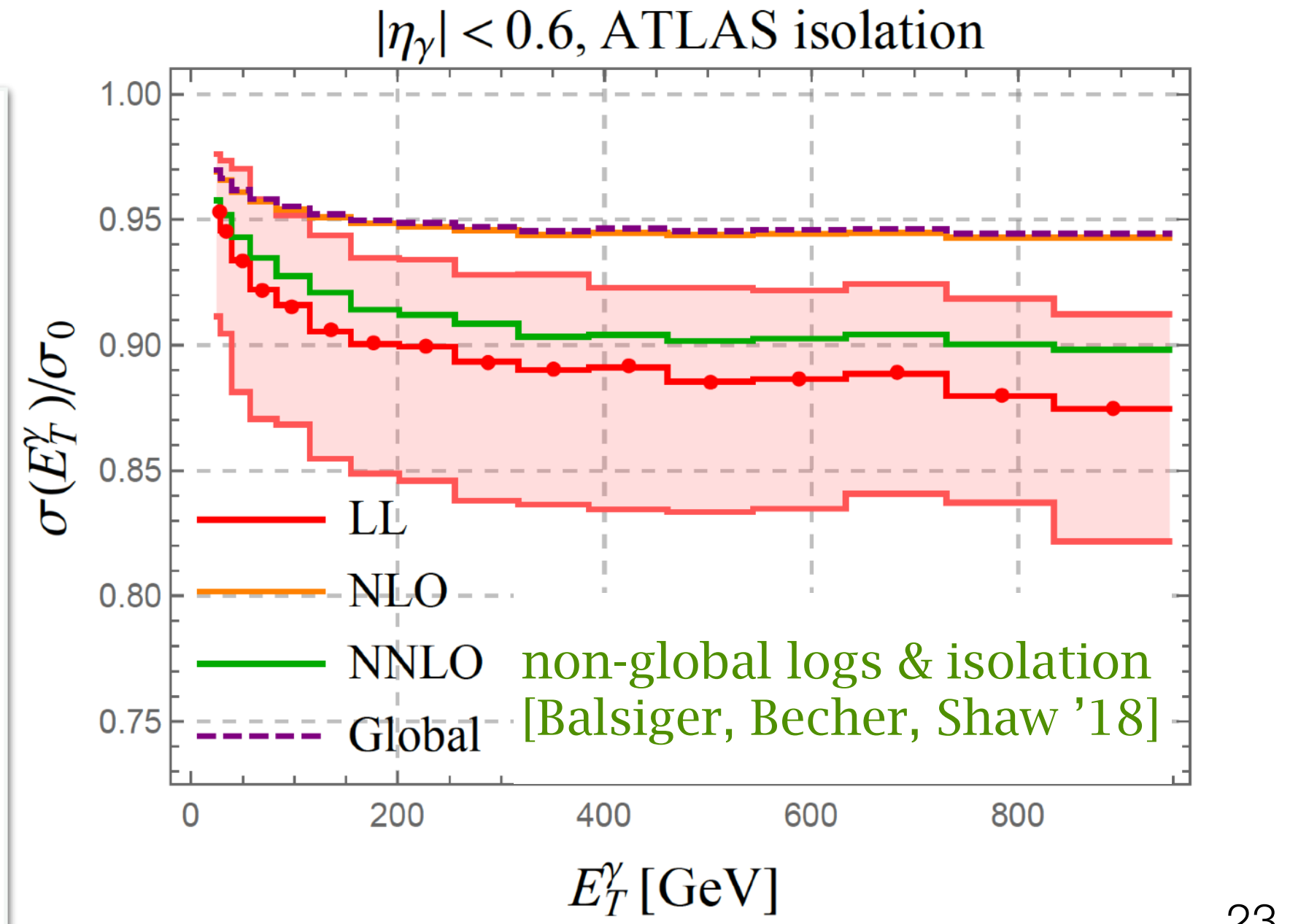
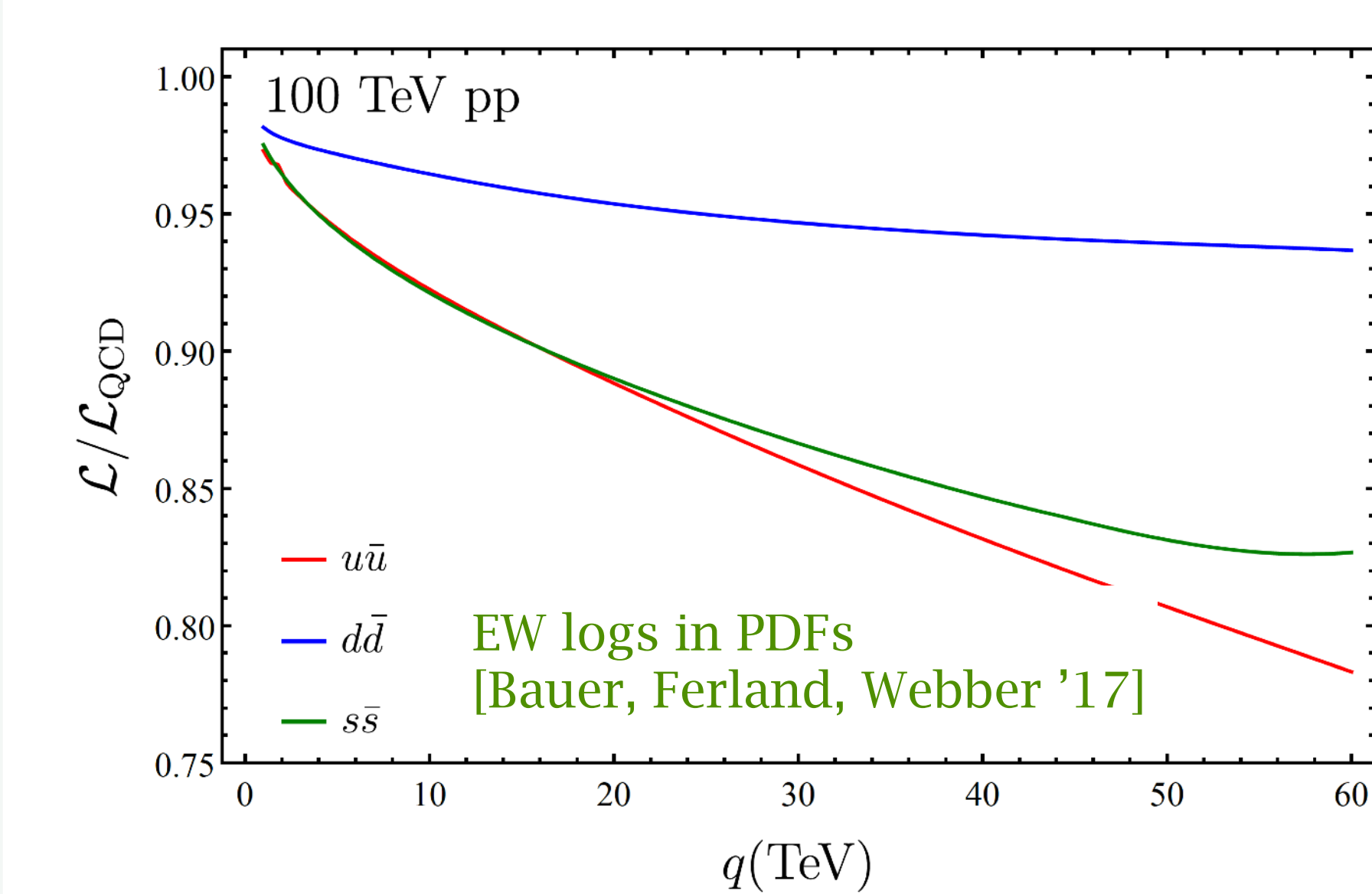
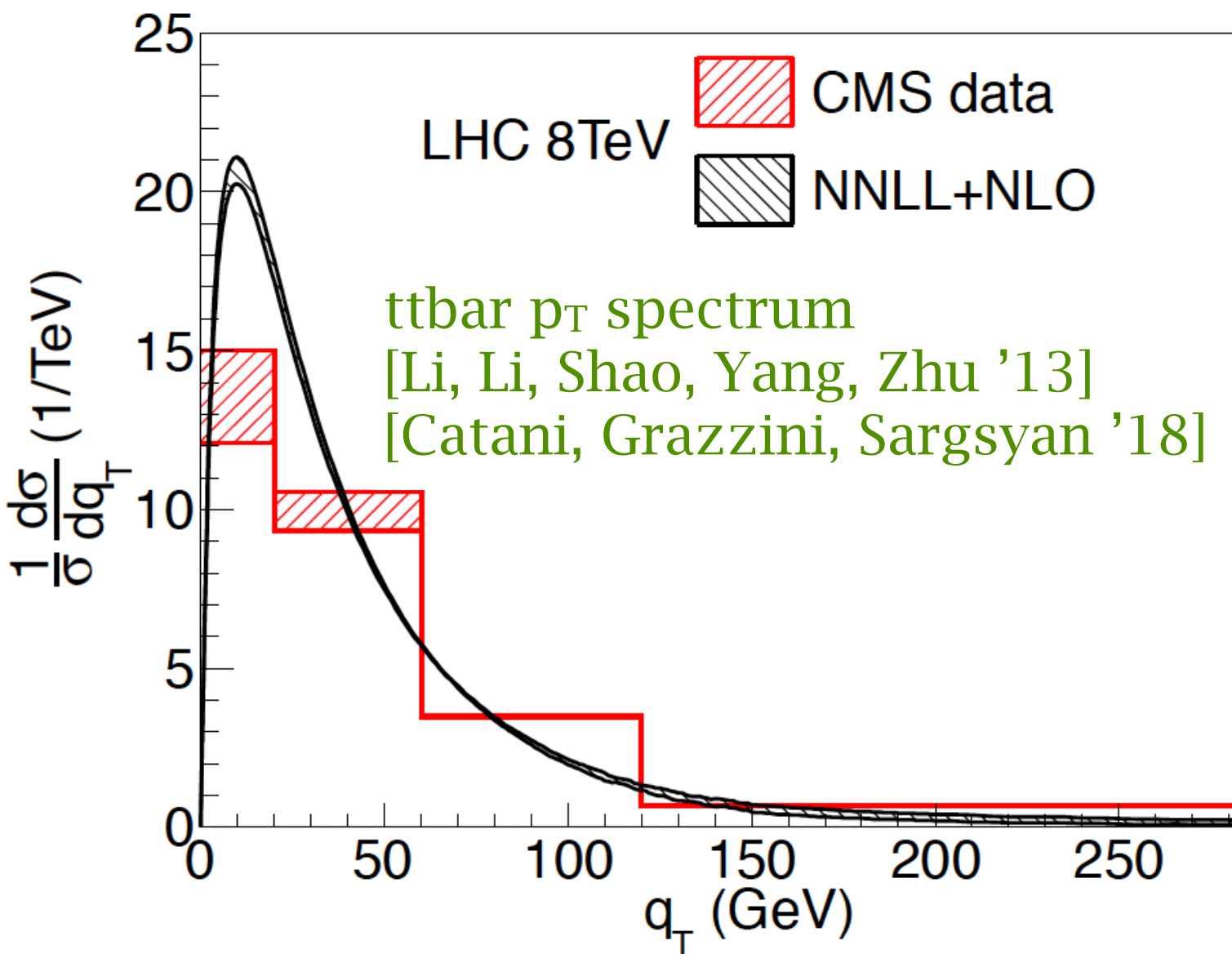
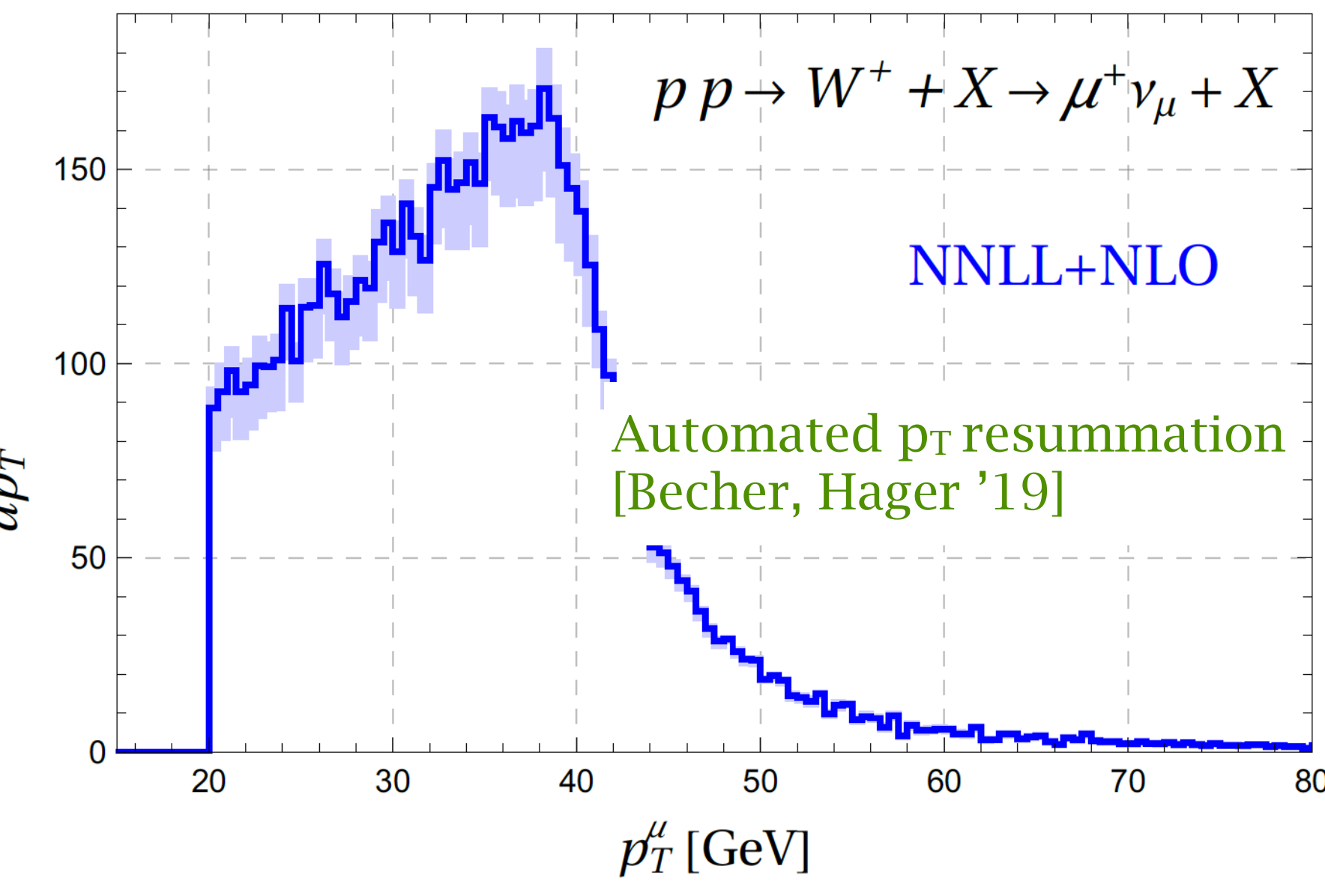
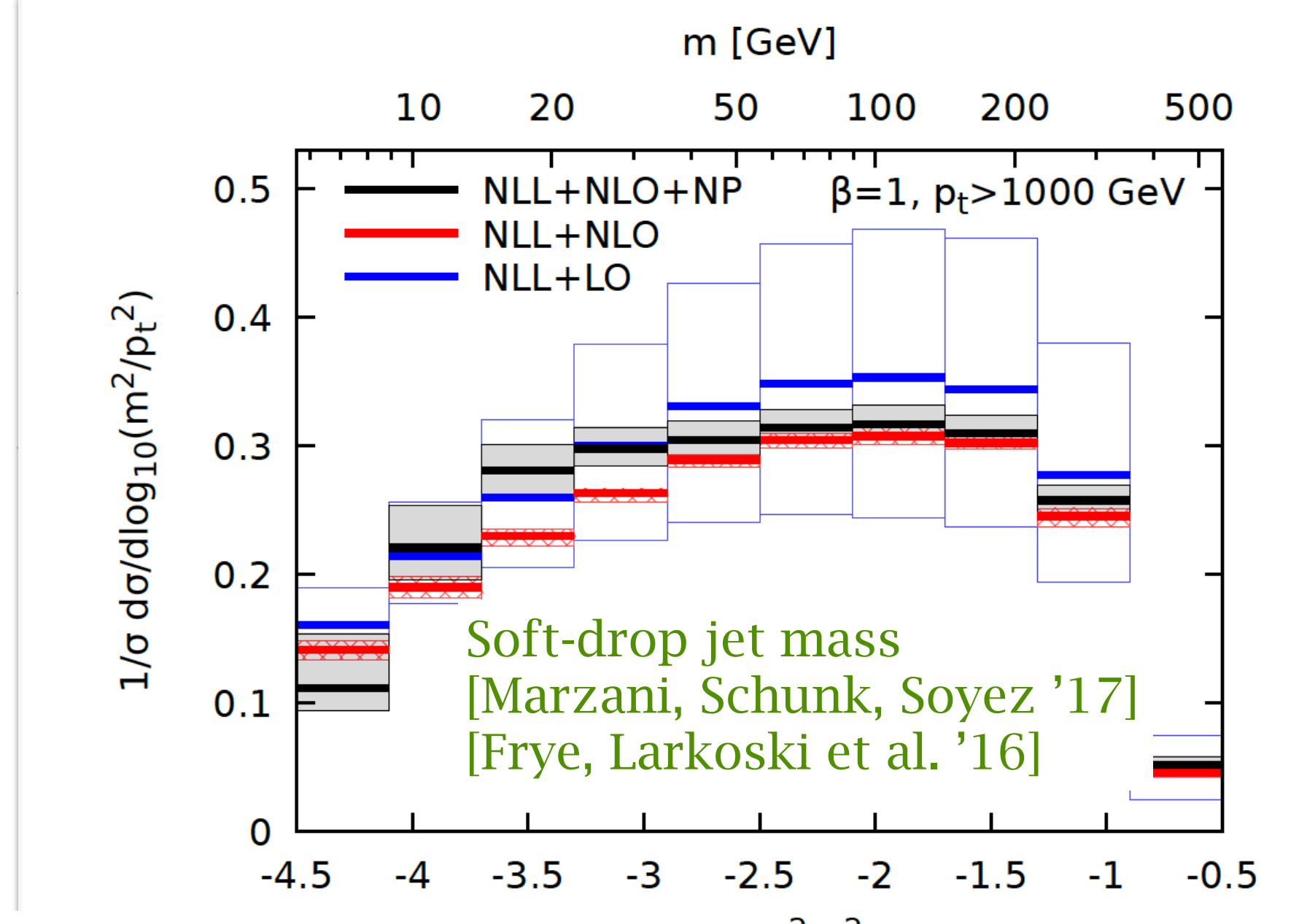
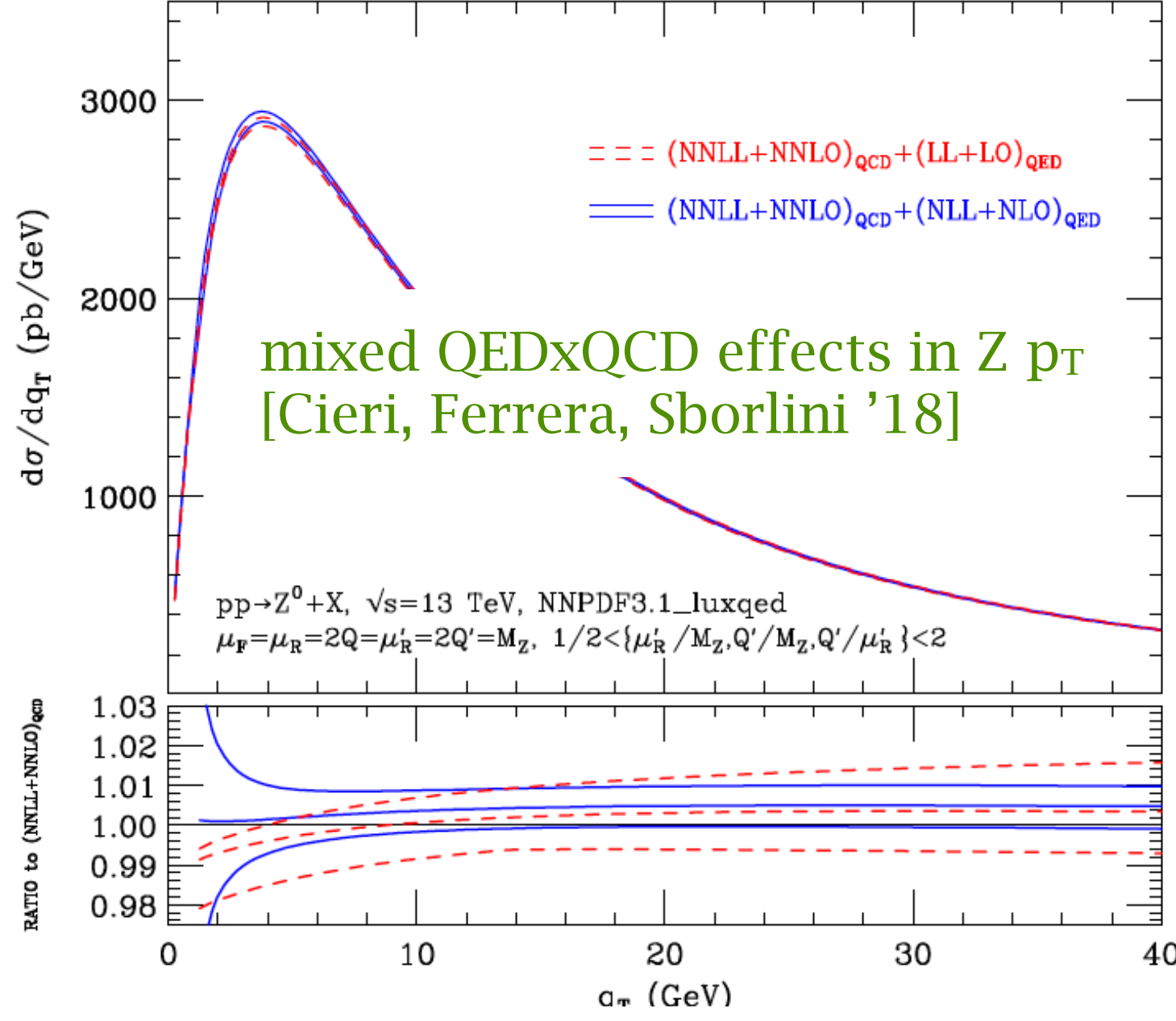
- ▶ The key to control and improve future parton shower algorithms (PS) may lie in their link to resummations
- ▶ **Resummations as limiting case of PS dynamics**
- ▶ **Assess perturbative accuracy of PS & devise new algorithms**

Study of radiation pattern unveils important constraints to go beyond LL in future designs

[Dasgupta, Dreyer, Hamilton, PM, Salam '18]
Related studies also in
[Hoang, Plaetzer, Samitz '18] [[>> A. Hoang's talk](#)]
[Bewick, Ravasio, Richardson, Seymour '19]



...just a glimpse, much more out there



Conclusions

- ▶ The future LHC programme deeply relies on precision (both for SM & BSM searches). Achieving the accuracy demanded by experiments requires **joint efforts in different areas of QCD**
- ▶ Understanding of infrared and all-order dynamics is crucial to control the theory at the few-% level in exclusive measurements
- ▶ Impressive progress in the past 10 years has led to the evolution of technology that allows us to tackle several problems of phenomenological relevance
 - ▶ precise predictions, subtractions, & event generators
 - ▶ better-behaved observables & substructure of jets
- ▶ Much more to be done on the theory side (multi-leg reactions, multi-differential/scales & exclusive measurements) ... exciting times ahead !

Thank you for listening

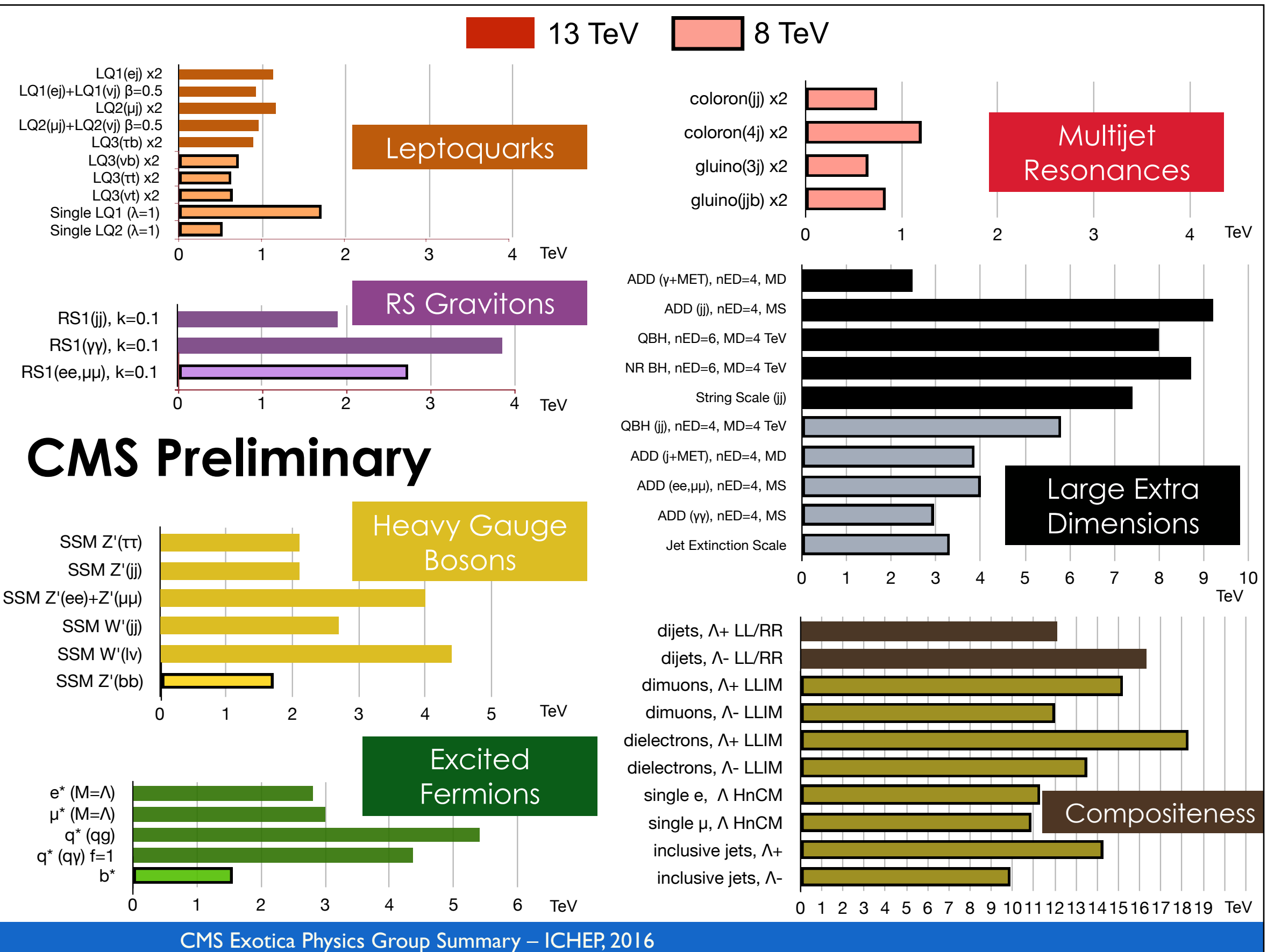
Search for new physics at the LHC

ATLAS SUSY Searches* - 95% CL Lower Limits July 2018

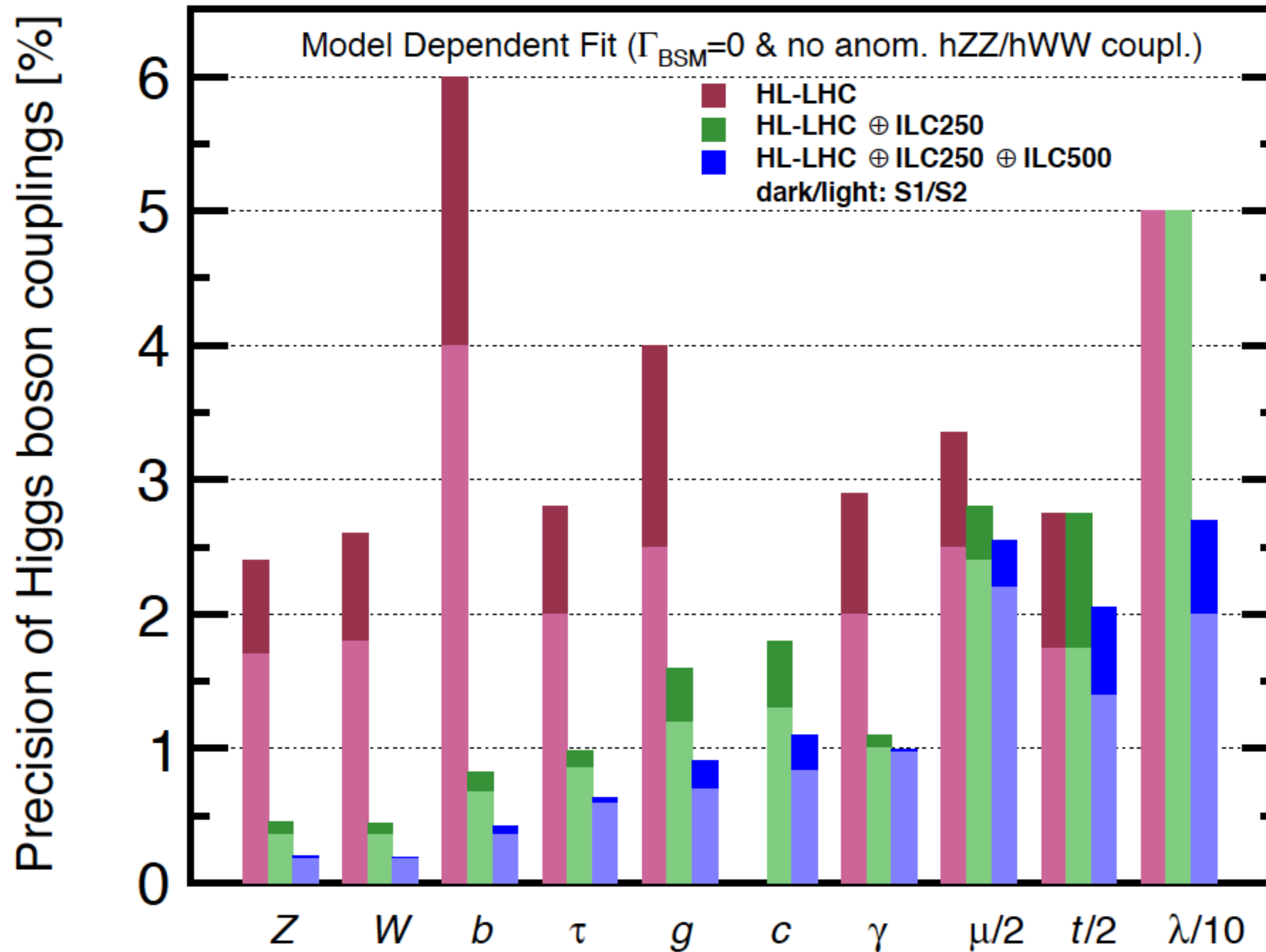
Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	$\sqrt{s} = 7, 8 \text{ TeV}$	$\sqrt{s} = 13 \text{ TeV}$	References		
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	36.1	\tilde{q} [2x, 8x Degen.]	0.9	1.55	$m(\tilde{\chi}_1^0) < 100 \text{ GeV}$	1712.023
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	1-3 jets	Yes	36.1	\tilde{q} [1x, 8x Degen.]	0.43	0.71	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1711.033
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	0	2-6 jets	Yes	36.1	\tilde{g}	Forbidden	0.95-1.6	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$	1712.023
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}WZ\tilde{\chi}_1^0$	3 e, μ	4 jets	-	36.1	\tilde{g}	Forbidden	1.85	$m(\tilde{\chi}_1^0) = 900 \text{ GeV}$	1706.037
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	e, μ	2 jets	Yes	36.1	\tilde{g}	1.2	1.8	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$	1805.113
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}WZ\tilde{\chi}_1^0$	0	7-11 jets	Yes	36.1	\tilde{g}	Forbidden	1.8	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$	1708.027
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$	0	4 jets	-	36.1	\tilde{b}_1	0.98	2.0	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200 \text{ GeV}$	1706.037
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$	0-1 e, μ	3 b	Yes	36.1	\tilde{b}_1	1.25	2.0	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$	1711.019
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$	3 e, μ	4 jets	-	36.1	\tilde{b}_1	Forbidden	0.9	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300 \text{ GeV}$	1706.037
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$	Multiple	Multiple	-	36.1	\tilde{b}_1	Forbidden	0.58-0.82	$m(\tilde{\chi}_1^0) = 300 \text{ GeV}, \text{BR}(\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0) = 1$	1708.09266, 171
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$	Multiple	Multiple	-	36.1	\tilde{b}_1	Forbidden	0.7	$m(\tilde{\chi}_1^0) = 300 \text{ GeV}, \text{BR}(\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0) = 0.5$	1708.092
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$	Multiple	Multiple	-	36.1	\tilde{b}_1	Forbidden	0.7	$m(\tilde{\chi}_1^0) = 200 \text{ GeV}, m(\tilde{\chi}_1^\pm) = 300 \text{ GeV}, \text{BR}(\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0) = 1$	1706.037
EW direct	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{\chi}_1^0$	0	2c	Yes	36.1	\tilde{t}_1	0.7	0.9	$m(\tilde{\chi}_1^0) = 60 \text{ GeV}$	1709.04183, 1711.115
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{\chi}_1^0$	0-2 e, μ	0-2 jets/1-2 b	Yes	36.1	\tilde{t}_1	Forbidden	1.0	$m(\tilde{\chi}_1^0) = 200 \text{ GeV}$	1709.04183, 1711.115
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{\chi}_1^0$	Multiple	Multiple	-	36.1	\tilde{t}_1	Forbidden	0.4-0.9	$m(\tilde{\chi}_1^0) = 150 \text{ GeV}, m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}, \tilde{t}_1 \approx \tilde{t}_L$	1506.08616, 1709.041
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{\chi}_1^0$	Multiple	Multiple	-	36.1	\tilde{t}_1	Forbidden	0.6-0.8	$m(\tilde{\chi}_1^0) = 300 \text{ GeV}, m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}, \tilde{t}_1 \approx \tilde{t}_L$	1709.04183, 171
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{\chi}_1^0$	Multiple	Multiple	-	36.1	\tilde{t}_1	Forbidden	0.48-0.84	$m(\tilde{\chi}_1^0) = 150 \text{ GeV}, m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}, \tilde{t}_1 \approx \tilde{t}_L$	1709.04183, 171
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{\chi}_1^0$	0	2c	Yes	36.1	\tilde{t}_1	0.46	0.85	$m(\tilde{\chi}_1^0) = 0 \text{ GeV}$	1805.016
Long-lived particles	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0/\tilde{\chi}_1^\pm, \tilde{c} \rightarrow c\tilde{\chi}_1^0$	0	mono-jet	Yes	36.1	\tilde{t}_1	0.43	0.5	$m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 50 \text{ GeV}$	1805.016
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0/\tilde{\chi}_1^\pm, \tilde{c} \rightarrow c\tilde{\chi}_1^0$	0	mono-jet	Yes	36.1	\tilde{t}_1	0.43	0.5	$m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1711.033
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$	1-2 e, μ	4 b	Yes	36.1	\tilde{t}_2	0.32-0.88	0.32-0.88	$m(\tilde{\chi}_1^0) = 0 \text{ GeV}, m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 180 \text{ GeV}$	1706.039
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via WZ	2-3 e, μ	-	Yes	36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.6	0.6	$m(\tilde{\chi}_1^0) = 0$	1403.5294, 180
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh	$ee, \mu\mu$	≥ 1	Yes	36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.17	0.17	$m(\tilde{\chi}_1^0) = 10 \text{ GeV}$	1712.081
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh	$lll\ell\gamma\gamma/lbb$	-	Yes	20.3	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.26	0.26	$m(\tilde{\chi}_1^0) = 0$	1501.071
RPV	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm/\tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}\nu(\tau\tilde{\nu}), \tilde{\chi}_2^0 \rightarrow \tilde{\tau}\tau(\nu\tilde{\tau})$	2 τ	-	Yes	36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.22	0.76	$m(\tilde{\chi}_1^0) = 0, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	1708.078
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm/\tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}\nu(\tau\tilde{\nu}), \tilde{\chi}_2^0 \rightarrow \tilde{\tau}\tau(\nu\tilde{\tau})$	2 τ	-	Yes	36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$	0.22	0.76	$m(\tilde{\chi}_1^0) = 100 \text{ GeV}, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	1708.078
	$\tilde{t}_L, \tilde{t}_L, \tilde{t}_L \rightarrow \tilde{t}\tilde{\chi}_1^0$	2 e, μ	0	Yes	36.1	\tilde{t}	0.18	0.5	$m(\tilde{\chi}_1^0) = 0$	1803.027
	$\tilde{t}_L, \tilde{t}_L, \tilde{t}_L \rightarrow \tilde{t}\tilde{\chi}_1^0$	2 e, μ	≥ 1	Yes	36.1	\tilde{t}	0.18	0.5	$m(\tilde{t}) - m(\tilde{\chi}_1^0) = 5 \text{ GeV}$	1712.081
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0	$\geq 3b$	Yes	36.1	\tilde{H}	0.13-0.23	0.29-0.88	$\text{BR}(\tilde{H} \rightarrow h\tilde{G}) = 1$	1806.040
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	4 e, μ	0	Yes	36.1	\tilde{H}	0.3	0.29-0.88	$\text{BR}(\tilde{H} \rightarrow Z\tilde{G}) = 1$	1804.036
RPV	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	36.1	$\tilde{\chi}_1^\pm$	0.15	0.46	Pure Wino	1712.021
	Stable \tilde{g} R-hadron	SMP	-	-	3.2	\tilde{g}	1.6	1.6	Pure Higgsino	ATL-PHYS-PUB-1606.051
	Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	Multiple	-	-	32.8	\tilde{g} [$\tau(\tilde{g}) = 100 \text{ ns}, 0.2 \text{ ns}$]	1.6	2.4	$m(\tilde{\chi}_1^0) = 100 \text{ GeV}$	1710.04901, 1604.04520
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	20.3	$\tilde{\chi}_1^0$	0.44	0.44	$1 < \tau(\tilde{\chi}_1^0) < 3 \text{ ns}$, SPSB model	1409.5542
	$\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow ee\nu/\mu\nu$	displ. $ee/\mu\nu$	-	-	20.3	\tilde{g}	1.3	1.3	$6 < c\tau(\tilde{\chi}_1^0) < 1000 \text{ mm}, m(\tilde{\chi}_1^0) = 1 \text{ TeV}$	1504.05162
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu$	$e\mu, e\tau, \mu\tau$	-	-	3.2	$\tilde{\nu}_\tau$	1.9	1.9	$\lambda_{311}^{\nu} = 0.11, \lambda_{132}/\lambda_{233} = 0.07$	1607.08079
$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm/\tilde{\chi}_2^0 \rightarrow WW/Zll\ell\nu\nu$	4 e, μ	0	Yes	36.1	$\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ [$\lambda_{333} \neq 0, \lambda_{12k} \neq 0$]	0.82	1.33	$m(\tilde{\chi}_1^0) = 100 \text{ GeV}$	1804.03602	
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow q\tilde{q}q$	0	4-5 large-R jets	-	36.1	\tilde{g} [$m(\tilde{\chi}_1^0) = 200 \text{ GeV}, 1100 \text{ GeV}$]	1.3	1.9	Large λ_{112}^{ν}	1804.03568	
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow q\tilde{q}q$	Multiple	Multiple	-	36.1	\tilde{g} [$\lambda_{112}^{\nu} = 2e-4, 2e-5$]	1.05	2.0	$m(\tilde{\chi}_1^0) = 200 \text{ GeV}$, bino-like	ATLAS-CONF-2018-003	
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{b}s / \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow t\tilde{b}s$	Multiple	Multiple	-	36.1	\tilde{g} [$\lambda_{323}^{\nu} = 1, 1e-2$]	1.8	2.1	$m(\tilde{\chi}_1^0) = 200 \text{ GeV}$, bino-like	ATLAS-CONF-2018-003	
$\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow t\tilde{b}s$	Multiple	Multiple	-	36.1	\tilde{t}_1 [$\lambda_{323}^{\nu} = 2e-4, 1e-2$]	0.55	1.05	$m(\tilde{\chi}_1^0) = 200 \text{ GeV}$, bino-like	ATLAS-CONF-2018-003	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{s}$	0	2 jets + 2 b	-	36.7	\tilde{t}_1 [qq, bs]	0.42	0.61	$m(\tilde{\chi}_1^0) = 200 \text{ GeV}$, bino-like	1710.07171	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{s}$	2 e, μ	2 b	-	36.1	\tilde{t}_1	0.4-1.45	0.4-1.45	$\text{BR}(\tilde{t}_1 \rightarrow b\tilde{s}) > 20\%$	1710.05544	

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

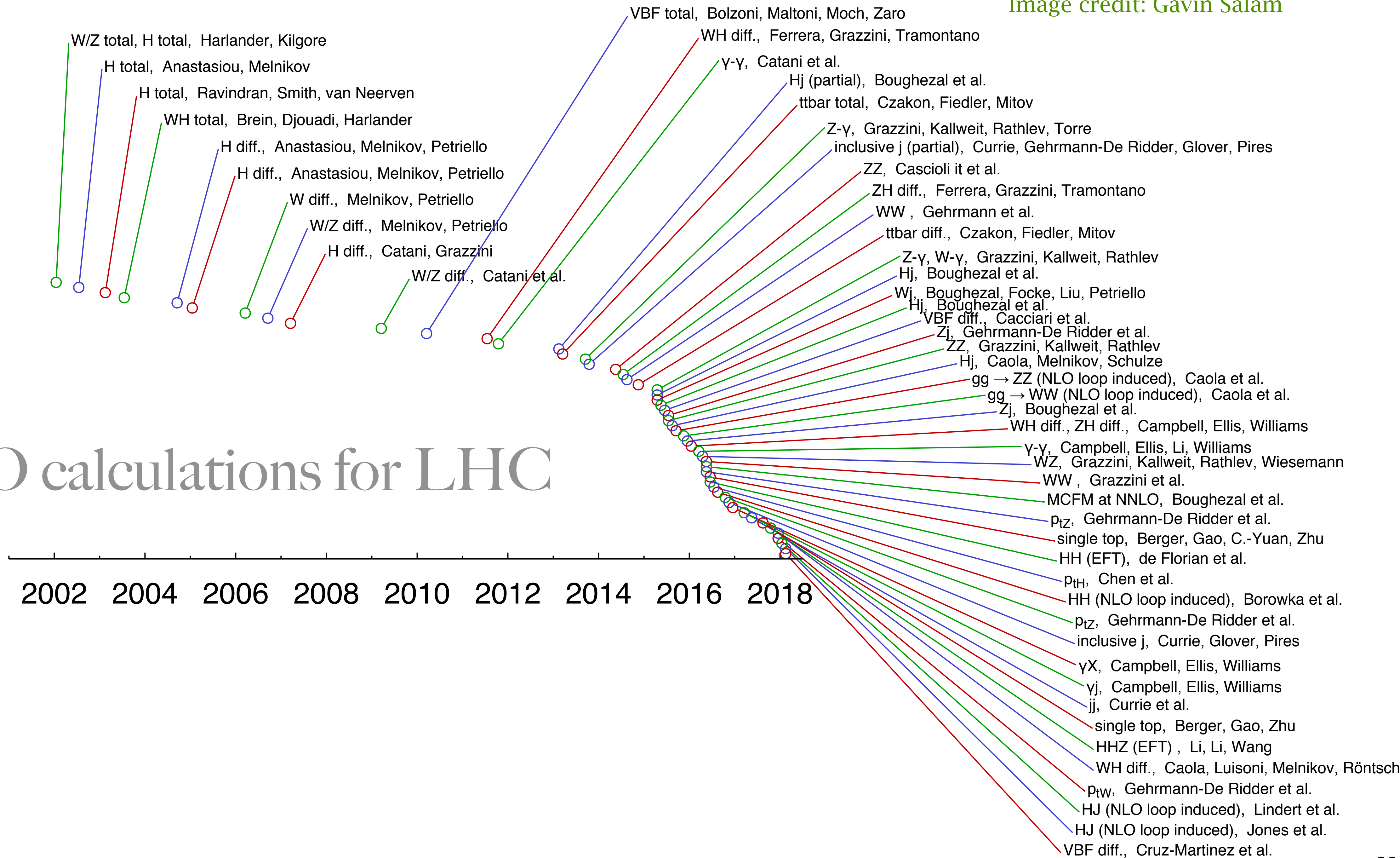
ATLAS Preliminary



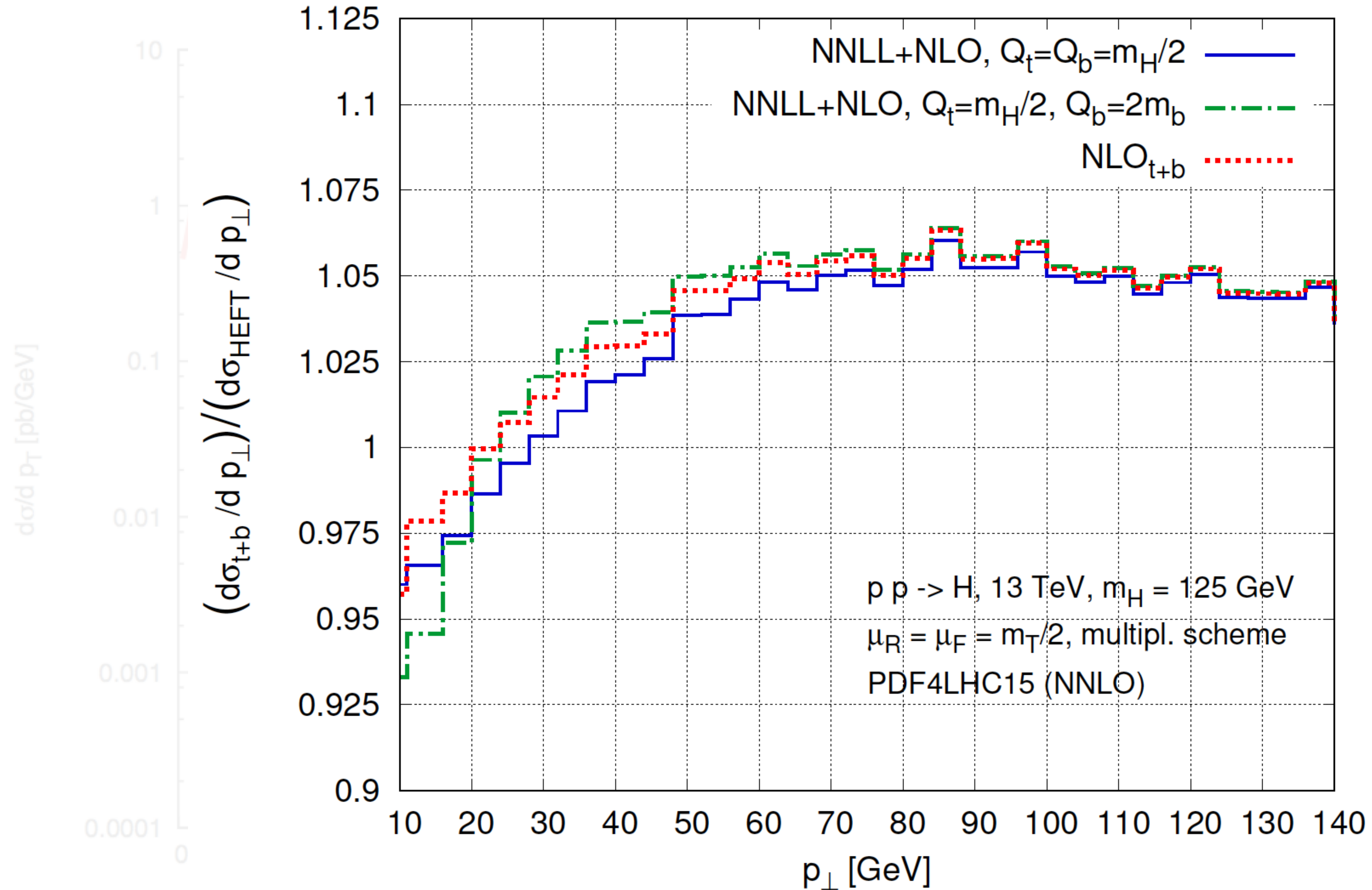
The charm Yukawa: future prospects



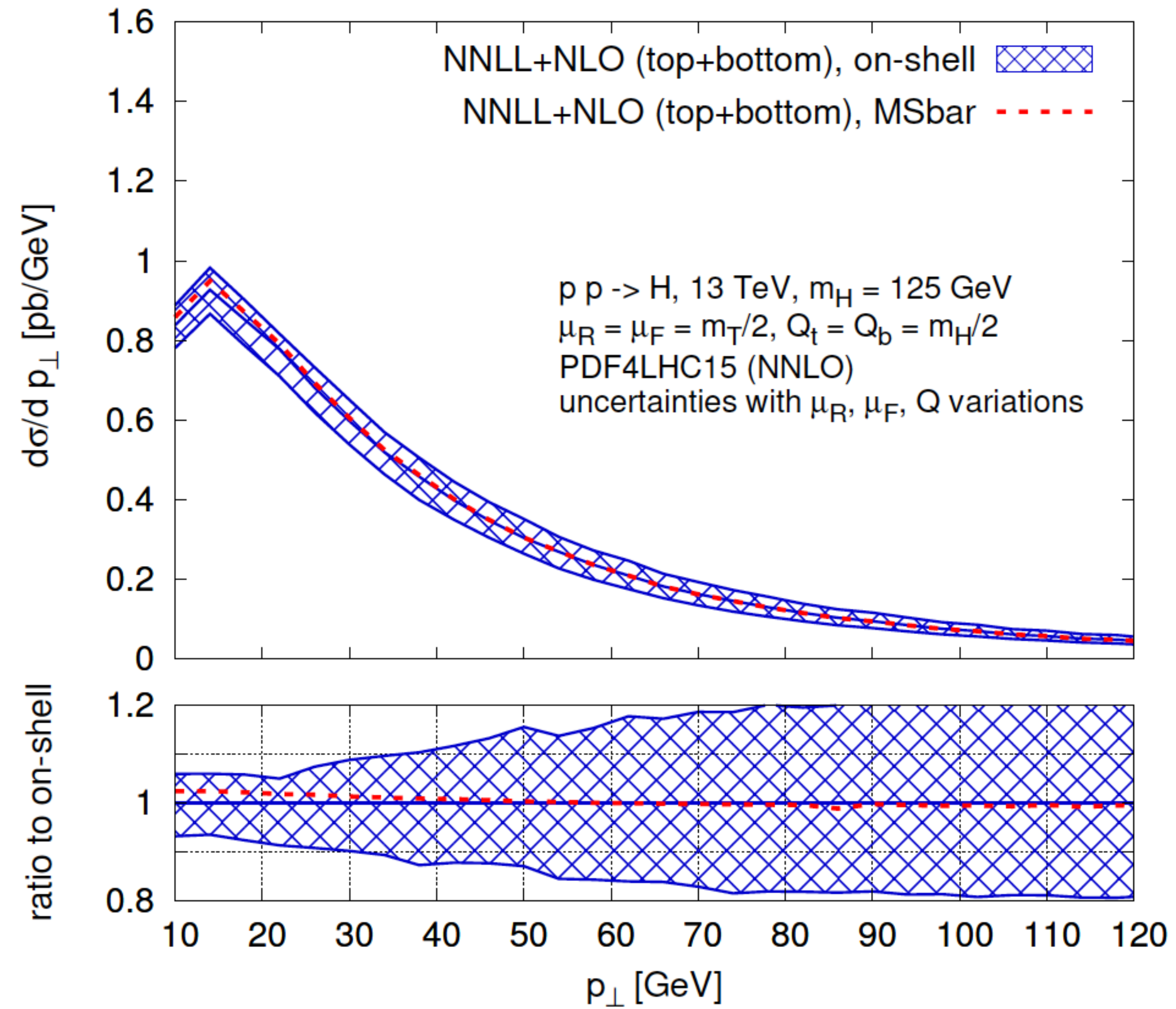
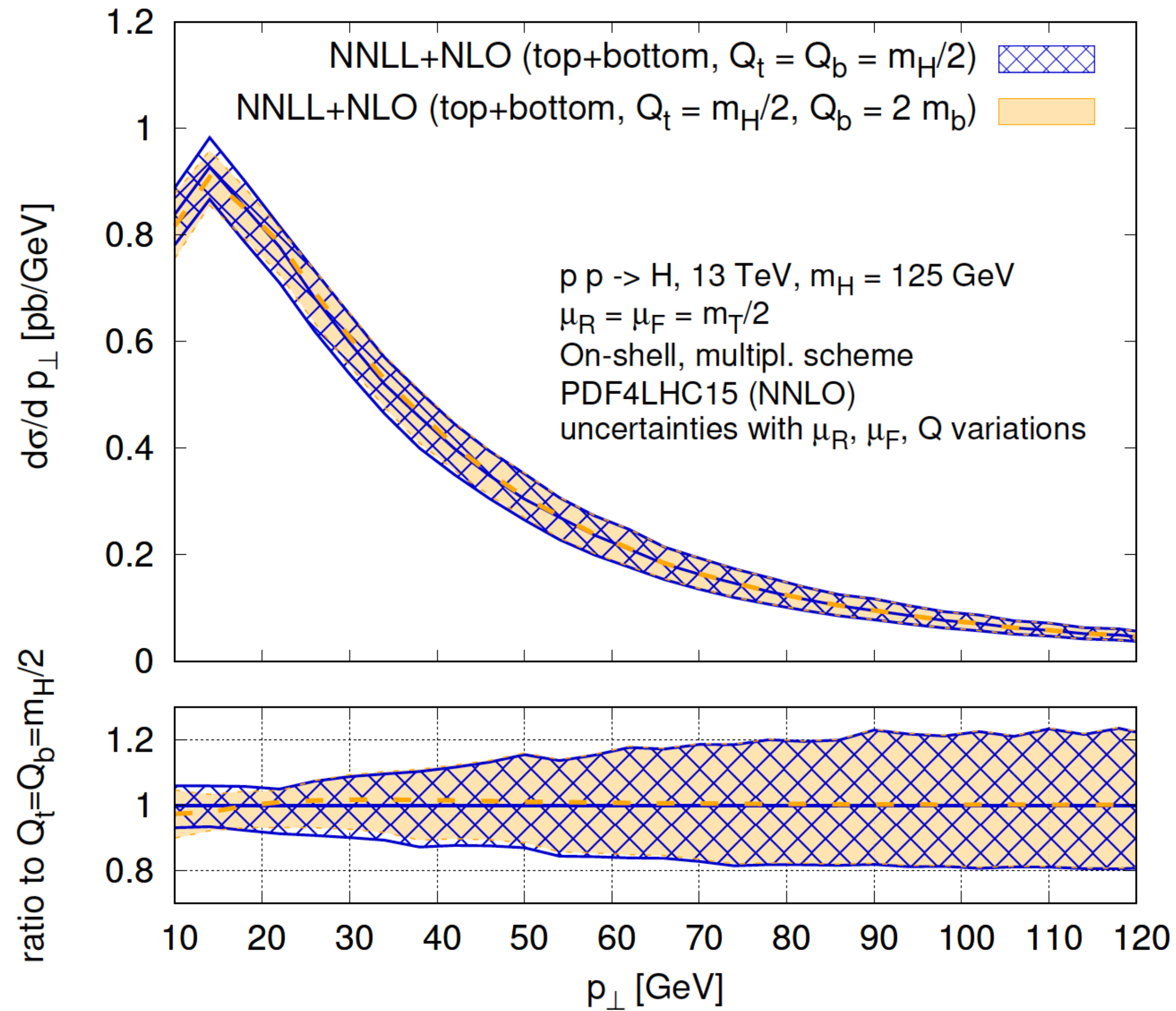
NNLO calculations for LHC



Higgs p_T distribution: quark masses



Higgs p_T distribution: quark masses



PDF uncertainty (Z @ LHC8)

▶ Data and fiducial cuts from [ATLAS 1512.02192]

▶ PDF errors at the 1% level, but difference between sets can be as large as 3.5%

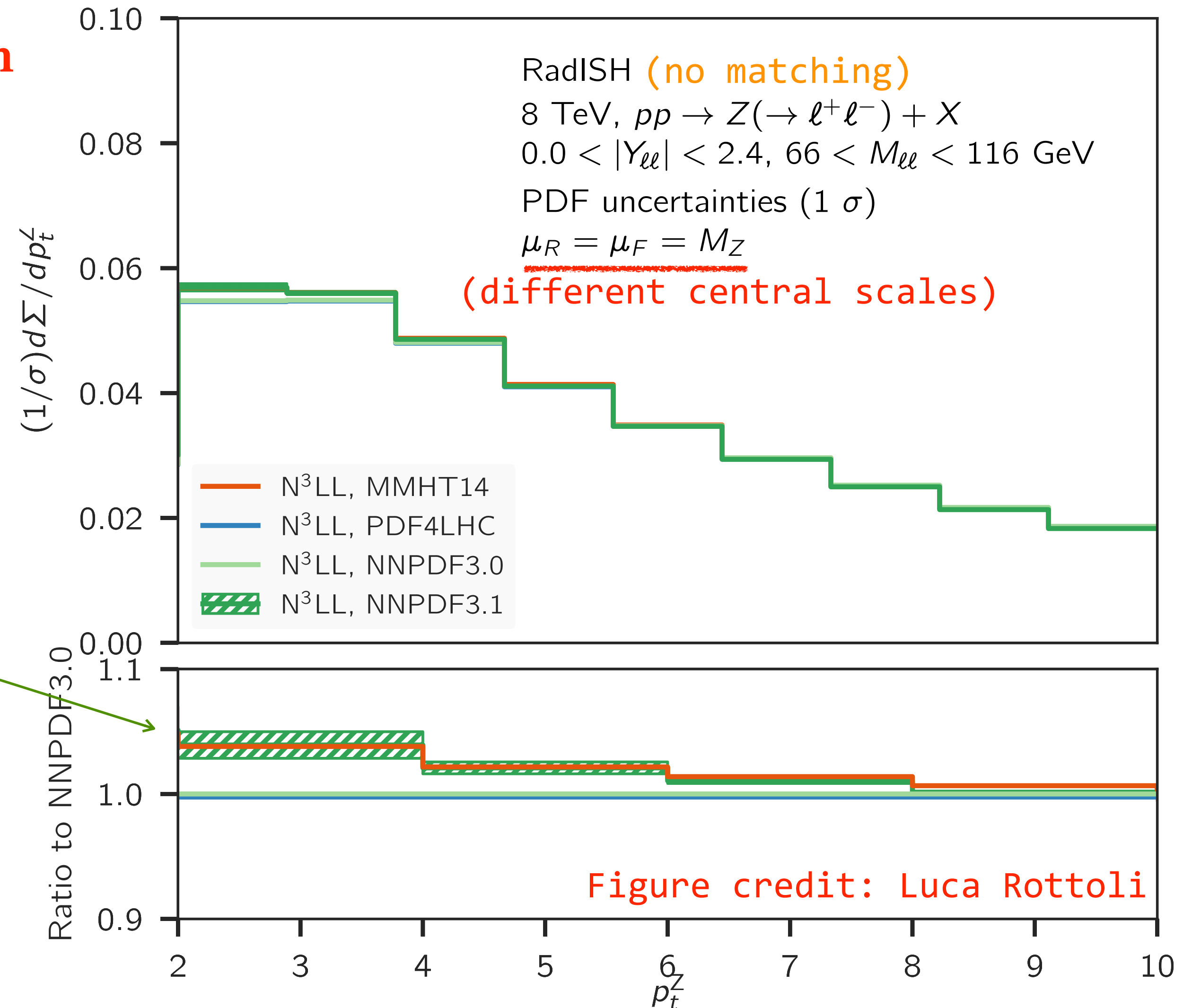
▶ Spectrum slightly harder with latest sets

▶ Theory uncertainties in PDFs become relevant

▶ see also Juan Rojo's talk (this morning)

Envelope of
NNPDF3.1 sets

$$p_t^{\ell^\pm} > 20 \text{ GeV}, \quad |\eta^{\ell^\pm}| < 2.4,$$



More predictions at LHC8 (ϕ_η^*)

