

Searches for new resonances in final states comprising leptons using the ATLAS detector

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June 4, 2019

On behalf of ATLAS Collaboration

*31st Rencontres de Blois,
June 2-7, 2019, Blois, Loire Valley, France*

Covered in this talk:

- Searches for TeV scale resonances in final states with at least one charged lepton using LHC Run-II pp collisions at $\sqrt{s} = 13$ TeV

Not covered:

- SUSY-inspired searches
- $X \rightarrow$ invisible recoiling against an SM state (covered in DM talks)
- Non-resonant signatures
- BSM decays of SM particles ($Z \rightarrow e\mu$, etc.)

Motivation: new heavy resonances appear in a variety of SM extensions.

Resonance searches are not sensitive to *all* details of the particular model.
 An amplitude is dominated by a diagram with an \sim on-shell resonance
 \Rightarrow need to know only its quantum numbers \Rightarrow Set limits in a parameter space of a **simplified model** \Rightarrow facilitate **reinterpretations**

- Spin 0:

- ▶ Extended Higgs sector: 2HDM
- ▶ Warped extra dimensions (Randall–Sundrum model, RS)
- ▶ superpartners of SM fermions

- Spin 1:

- ▶ Extended gauge group \Rightarrow HVT ...
- ▶ DM mediators
- ▶ superpartners of SM fermions

- Spin 2:

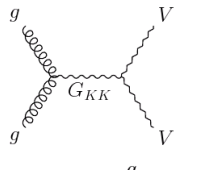
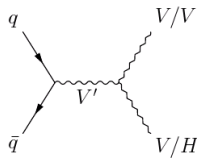
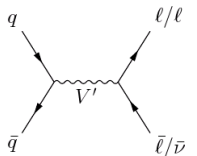
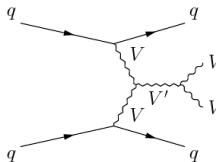
- ▶ Randall–Sundrum G_{KK}

- Spin 1/2:

- ▶ Vector-like fermions, Majorana neutrinos, excited fermions (technicolor models), superpartners of SM bosons

Resonant production:

- gg fusion (ggF) ● VV fusion (VBF) ● $q\bar{q}$ ● associated with 3rd generation fermions



• V_μ^a : $I = 1, Y = 0$

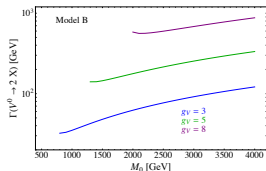
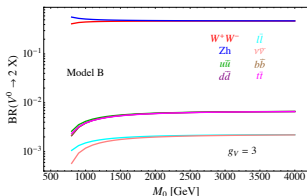
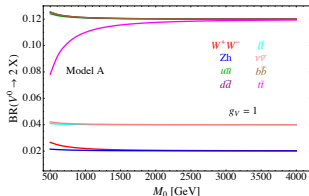
- ▶ Physical W, Z have an admixture of $V^{1,2}, V^0$
- ▶ γ remains massless with SM couplings

$$\begin{aligned} \mathcal{L}_{\text{HVT}} = & -\frac{1}{4} D_{[\mu} V_{\nu]}^a D^{[\mu} V^{\nu]} a + \frac{m_V^2}{2} V_\mu^a V^\mu a \\ & + i g_V c_H V_\mu^a H^\dagger \tau^a \overleftrightarrow{D}^\mu H + \frac{g^2}{g_V} c_F V_\mu^a J_F^{\mu a} \\ & + \frac{g_V}{2} c_{VVV} \epsilon_{abc} V_\mu^a V_\nu^b D^{[\mu} V^{\nu]} c \\ & + g_V^2 c_{VVHH} V_\mu^a V^\mu a H^\dagger H \\ & - \frac{g}{2} c_{VWW} \epsilon_{abc} W^{\mu\nu a} V_\mu^b V_\nu^c, \end{aligned}$$

• HVT realisations: $\triangleright V' \rightarrow VV/VH, \ell\ell/\ell\nu$

- ▶ (A) Weak coupling: $g_V^2 c_H \ll g^2 c_F \Rightarrow W', Z'$
- ▶ (B) Strong coupling: $g_V^2 c_H \gg g^2 c_F \Rightarrow$ composite Higgs
- ▶ Fermiophobic: $c_F = 0 \Rightarrow$ only VBF production
- ▶ Parameter phase space beyond the above is also probed.

where D_μ are $SU(2)_L$ covariant derivatives and $J_F^{\mu a} = \sum_f \bar{f}_L \gamma^\mu \tau^a f_L$ is the fermion current.



Benchmark models: Randall–Sundrum (RS1) warped extra dimensions

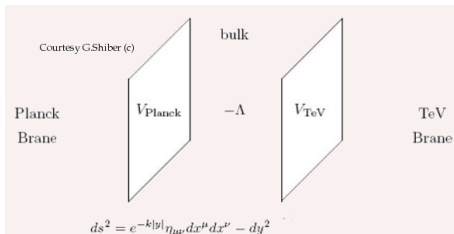
$$ds^2 = e^{-2kr_c\phi} \eta_{\mu\nu} dx^\mu dx^\nu - r_c^2 d\phi^2$$

$$0 \leq \phi \leq \pi, \quad k \sim M_{Pl}, \quad kr_c \sim 10$$

Take with caution: oscillations with \sim Planck frequency at the Planck brane are red-shifted down to the EW scale while going from Planck to the EW brane

A variant of the RS1 model as a benchmark [[Phys. Rev. D 76, 036006](#)]:

- Visible and hidden 3-branes at $\phi = 0$ and $\phi = \pi$
- The EW scale $\propto \bar{M}_{Pl} e^{-kr_c\pi} \Rightarrow$ no need for large r_c
- SM fields allowed to propagate in the bulk. Fermion localization explains flavours: light fermions are near the Planck brane, while 3rd generation and the Higgs are at the TeV brane.
- Zero mode bulk metric excitation $\Leftrightarrow m = 0$ graviton couples to $T^{\mu\nu} \propto 1/\bar{M}_{Pl}^2$
- Higher bulk metric excitations are $\sim O(\text{TeV})$ and localized near the TeV brane.
- Bulk G_{KK} couples to SM gauge bosons' and 3rd generation fermions' $T^{\mu\nu} \propto e^{2kr_c\pi} / \bar{M}_{Pl}^2 \sim 1/\text{TeV}^2$
- G_{KK} couplings to 1,2-generation fermions are suppressed.



$G_{KK} \rightarrow VV$

Final states and observables

- $X \rightarrow ll, lv, \nu\nu$

- $X \rightarrow tb \rightarrow lX$

- \rightarrow light $q\bar{q}, gg$

- $\rightarrow t\bar{t} \rightarrow lX$

(see N.Nishu's talk)

- $X \rightarrow VV, VH, HH$

- $\rightarrow V\gamma, H\gamma$

- $\blacktriangleright V \rightarrow ll'$

- $\blacktriangleright H \rightarrow bb, \tau\tau, VV^*$

Observables:

- $m(X)$

- m_T (if ν 's)

- ...

Combination of multiple channels in a context of a benchmark model to increase sensitivity: $\{VV, VH, ll'\} \rightarrow \{\{VV \& VH\}, ll'\} \rightarrow \{VV \& VH \& ll'\}$

ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: March 2019

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 139) \text{ fb}^{-1}$$

$$\sqrt{s} = 8, 13 \text{ TeV}$$

Model	Γ, γ	Jets ^b	$E_{\text{miss}}^{\text{min}}$	$ \beta \alpha(\beta^{-1})$	Limit	Reference		
Extra dimensions	ADD $G_{\mu\nu} = g^2$	$8 \mu\text{eV}$	$1-4j$	Yes	36.1	36.1	1713 0301	
	ADD non-resonant $\gamma\gamma$	$2j$	-	-	36.7	36.7	1707 0414	
	ADD GBH	-	$2j$	-	-	37.0	1703 0127	
	ADD BH High Σ_{pr}	$\geq 1.6 \mu\text{eV}$	$\geq 2j$	-	-	3.2	1608 0286	
	ADD BH multi-jet	$\geq 1.6 \mu\text{eV}$	$\geq 3j$	-	-	3.5	1512 0595	
	RS1 $G_{\mu\nu} = \gamma\gamma$	$2j$	-	-	-	36.7	1707 0417	
	Bulk RS $G_{\mu\nu} \rightarrow WW/ZZ$	multi-channel	-	-	-	36.1	1608 0286	
	Bulk RS $G_{\mu\nu} \rightarrow WW/ZZ \rightarrow \text{pppp}$	$0 \mu\text{eV}$	$2,3j$	-	-	139	1608 0286	
	Bulk RS $\beta_{\text{loop}} \rightarrow tt$	$1 \mu\text{eV}$	$\geq 1b, \geq 1LQj$	Yes	36.1	36.1	1604 1025	
	2UED/1RP	$1 \mu\text{eV}$	$\geq 2b, \geq 1j$	Yes	36.1	36.1	1603 0874	
Gauge bosons	SSM $Z' \rightarrow tt$	$2 \mu\text{eV}$	-	-	139	1603 0874		
	SSM $Z' \rightarrow \tau\tau$	$2 \mu\text{eV}$	-	-	26.1	1703 0742		
	Leptoquark $Z' \rightarrow Mb$	-	$2b$	-	-	36.1	1605 0550	
	Leptoquark $Z' \rightarrow tt$	$1 \mu\text{eV}$	$\geq 1b, \geq 1LQj$	Yes	36.1	36.1	1604 1025	
	SSM $W' \rightarrow \nu\nu$	$1 \mu\text{eV}$	$\geq 1j$	-	-	75.0	1601 0602	
	SSM $W' \rightarrow \tau\nu$	$1 \mu\text{eV}$	-	Yes	36.1	36.1	1601 0602	
	HVT $V' \rightarrow WW \rightarrow \text{pppp}$ model B	$0 \mu\text{eV}$	$2,3j$	-	-	139	1603 0874	
	HVT $V' \rightarrow WW/ZZ$ model B	multi-channel	-	-	-	36.1	1713 0316	
	LRSM $W'_2 \rightarrow tb$	multi-channel	-	-	-	36.1	1607 1475	
	CI	CI pppp	-	$2j$	-	-	37.0	1703 0127
CI pppp		$2 \mu\text{eV}$	-	-	-	36.1	1707 0424	
CI tttt		$2.1 \mu\text{eV}$	$\geq 1b, \geq 1j$	-	-	36.1	1501 0205	
DM	Axis-vector mediator (Disc DM)	$0 \mu\text{eV}$	$1-4j$	Yes	36.1	1603 0874		
	Colored scalar mediator (Disc DM)	$0 \mu\text{eV}$	$1-4j$	Yes	36.1	1603 0874		
	V/χ_{SPT} (Disc DM)	$0 \mu\text{eV}$	$1,3,4,6j$	Yes	3.2	36.1	1608 0286	
LO	Scalar reson. $\phi \rightarrow \tau\tau$ (Disc DM)	$0.1 \mu\text{eV}$	$1b, \geq 1j$	Yes	36.1	1603 0874		
	Scalar LQ 1^{st} gen	$1.2 \mu\text{eV}$	$\geq 2j$	Yes	36.1	1603 0874		
	Scalar LQ 2^{nd} gen	$1.2 \mu\text{eV}$	$\geq 2j$	Yes	36.1	1603 0874		
Heavy quarks	Scalar LQ 3^{rd} gen	$2 \mu\text{eV}$	$2b$	-	-	36.1	1603 0874	
	Scalar LQ 3^{rd} gen	$0.1 \mu\text{eV}$	$2b$	Yes	36.1	1603 0874		
	$W/Q \rightarrow T \rightarrow W/Z \rightarrow X$	multi-channel	-	-	-	36.1	1603 0874	
	$W/Q \rightarrow W/Z \rightarrow X$	multi-channel	-	-	-	36.1	1603 0874	
	$W/Q \rightarrow T \rightarrow F_{1,2} \rightarrow X$	$2BSM(\geq 2 \mu\text{eV} \geq 1b, \geq 1j)$	Yes	36.1	1603 0874	1603 0874		
	$W/Q \rightarrow W/Z \rightarrow X$	$1 \mu\text{eV}$	$\geq 1b, \geq 1j$	Yes	36.1	1603 0874		
	$W/Q \rightarrow Hb \rightarrow X$	$0.4 \mu\text{eV}$	$\geq 1b, \geq 1j$	Yes	20.9	20.9	1603 0874	
	$W/Q \rightarrow Wq \rightarrow X$	$1 \mu\text{eV}$	$\geq 4j$	Yes	26.3	26.3	1603 0874	
	Exotic fermions	Excited quark $q^* \rightarrow qq$	-	$2j$	-	-	139	1603 0874
		Excited quark $q^* \rightarrow q\gamma$	$1 \mu\text{eV}$	-	-	-	36.7	1603 0874
Excited quark $q^* \rightarrow qg$		-	$1b, 1j$	-	-	36.1	1603 0874	
Excited lepton $l^* \rightarrow ll$		$3 \mu\text{eV}$	-	-	-	26.3	1603 0874	
Excited lepton $l^* \rightarrow l\gamma$		$3 \mu\text{eV}$	-	-	-	26.3	1603 0874	
Other	Type III Seesaw	$1 \mu\text{eV}$	$\geq 2j$	Yes	75.0	1603 0874		
	LRSM Majorana ν	$2 \mu\text{eV}$	$2j$	Yes	36.1	1603 0874		
	Higgs triplet $H^{\pm} \rightarrow tt$	$23.6 \mu\text{eV}$ (BS)	-	-	-	36.1	1603 0874	
	Higgs triplet $H^{\pm} \rightarrow \tau\tau$	$23 \mu\text{eV}$	-	-	-	26.3	1603 0874	
	Multi-charged leptons	$23 \mu\text{eV}$	-	-	-	36.1	1603 0874	
Magnetic monopoles	-	-	-	-	7.0	1603 0874		

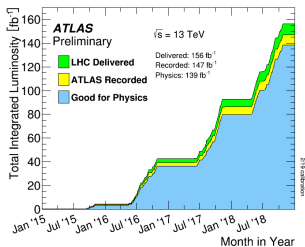
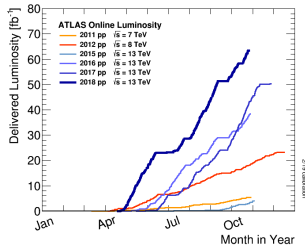
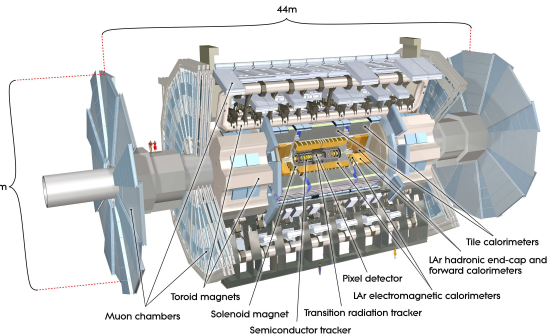
*Only a selection of the available mass limits on new states or phenomena is shown.

^bSmall-radius (large-radius) jets are denoted by the letter j (J).

The ATLAS experiment

Reported results are based on $L_{int} = 36 - 139$ /fb

Searches with a partial dataset will be extended to the full LHC Run-II statistics (*stay tuned*)



Phase space coverage:

- ▶ Tracking system: $|\eta| < 2.5$, $p_T > 0.1$ GeV
- ▶ Calorimetry: $|\eta| < 3.2$ (EM), $|\eta| < 4.9$ (HAD), $\sigma_E^{\text{jet}}/E^{\text{jet}} \sim 2\%$ for $|\eta^{\text{jet}}| < 0.8$ and $E_T \sim 1$ TeV
- ▶ Muon system: $|\eta| < 2.7$, $\sigma(p_T^\mu) \sim 2$ GeV for $20 < p_T^\mu < 300$ GeV

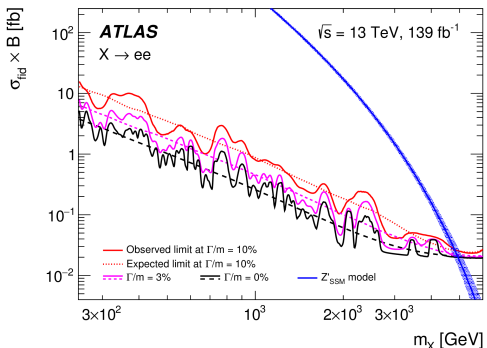
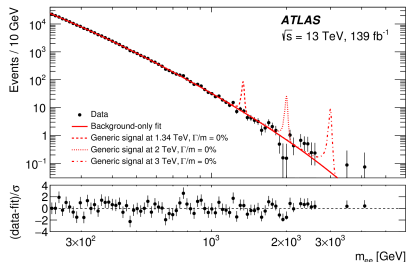
LHC pp runs at $\sqrt{s} = 13$ TeV:

- ▶ 2015: $L_{int} = 3.2$ /fb with a fully operational detector
- ▶ 2016: $L_{int} = +33$ /fb
- ▶ 2017-2018: $L_{int} \rightarrow 139$ /fb

Search for high-mass $X \rightarrow ee/\mu\mu$

arXiv:1903.06248
CERN-EP-2019-030

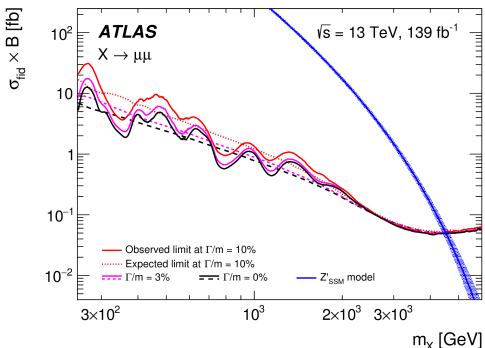
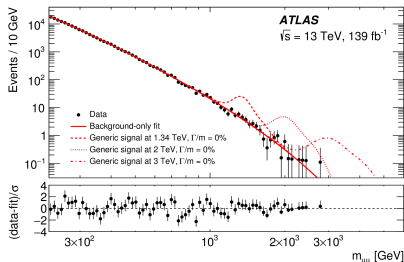
- $\sqrt{s} = 13$ TeV, $L_{int} = 139/\text{fb}$
- Fitting the $M_{\ell\ell}$ distribution by a smooth background parameterisation and $\text{BW} \oplus (\text{Gauss} + \text{CB})$ function for the signal with the given M_X, Γ_X
- **No significant deviation from the background-only hypothesis**, upper limits on $\sigma \times \mathcal{B}(X \rightarrow \ell\ell)$ are set at CL=95%
- The limits apply to X spin= 0, 1, 2
- Model-dependent $ee \oplus \mu\mu$ limits: $M_X > 4.5$ and 4.7 TeV for the E_6 $Z'_{\psi, \chi}$ bosons, $M_X > 5$ TeV for Z'_{SSM}



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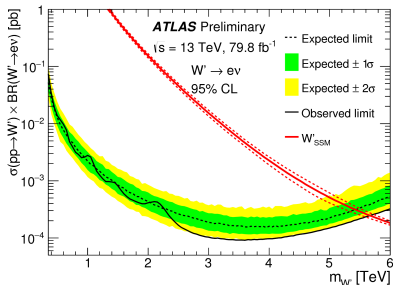
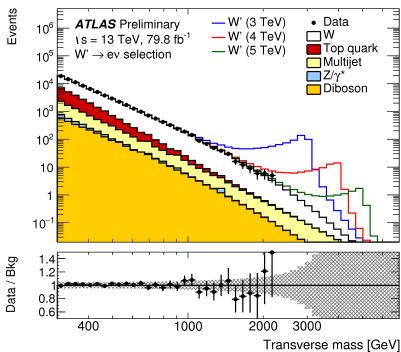
Search for $W' \rightarrow e\nu/\mu\nu$

ATLAS-CONF-2018-017

- $\sqrt{s} = 13$ TeV,
 $L_{int} = 79.8/\text{fb}$
partial LHC Run 2 data,
to be extended to the full
2015–2018 dataset,
 $L_{int} = 140/\text{fb}$
- The discriminating variable:
$$M_T(\ell\nu) = \sqrt{2p_T^\ell E_T^{\text{miss}}(1 - \cos\phi_{\ell\nu})}$$
- **No significant excess above the SM background**, upper limits on $\sigma \times \mathcal{B}(W' \rightarrow \ell\nu)$ $\ell = e, \mu$ are set at CL=95%
- Assuming SM couplings for the W' : $M_{W'} > 5.6$ TeV at CL=95%

More on lepton non-universality:

▶ $M_{W'} > 3.7$ TeV in $W'(\tau\nu)$



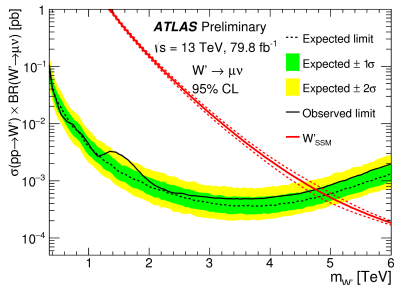
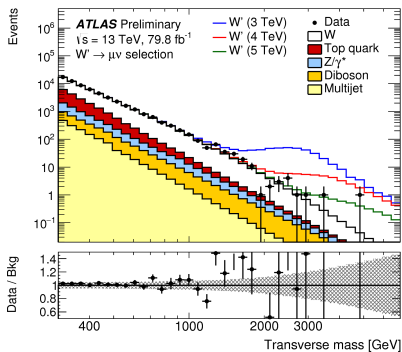
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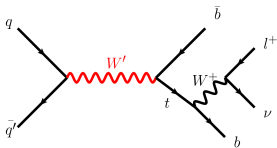
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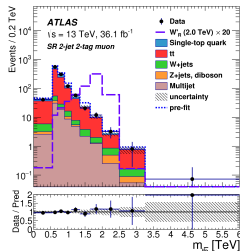
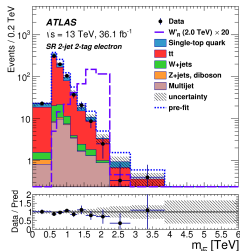
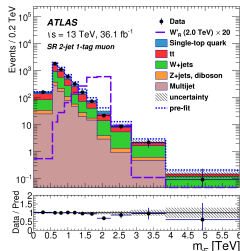
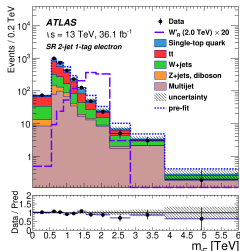
▶ $M_{W'} > 3.7$ TeV in $W'(\tau\nu)$

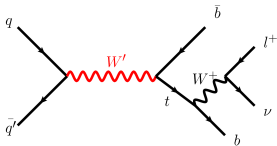




- $\sqrt{s} = 13 \text{ TeV}$, $L_{int} = 36.1/\text{fb}$
- Motivation: $W' \rightarrow t\bar{b} \rightarrow W(\ell\nu)b\bar{b}$
- $\Rightarrow \ell^\pm + \geq 1b\text{-jet} + \geq 1 \text{ jets} + E_T^{miss}$, $\ell = e, \mu$
 - ▶ $\ell^\pm, \vec{E}_T^{miss}, m_W$ mass constraint $\Rightarrow \vec{p}_\nu$
 - ▶ $\min |m(W, \text{jet}) - 172.5 \text{ GeV}| \Rightarrow m_{top}^{reco}$
- Signal regions are categorised by numbers of jets, b -tags and $\ell = e, \mu$:

m_{tb} in 2-jet and 1,2**b**-tag SR \rightarrow





- $\sqrt{s} = 13 \text{ TeV}$, $L_{int} = 36.1/\text{fb}$
- Motivation: $W' \rightarrow t\bar{b} \rightarrow W(\ell\nu)b\bar{b}$
- $\Rightarrow \ell^\pm + \geq 1b\text{-jet} + \geq 1 \text{ jets} + E_T^{miss}$,

$\ell = e, \mu$

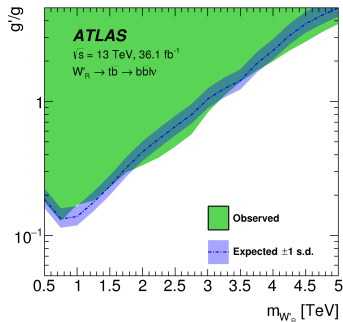
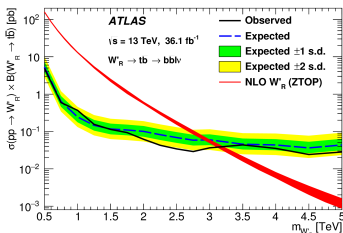
- ▶ $\ell^\pm, E_T^{miss}, m_W$ mass constraint $\Rightarrow \vec{p}_\nu$
- ▶ $\min |m(W, \text{jet}) - 172.5 \text{ GeV}| \Rightarrow m_{top}^{reco}$

- CL=95% limits on $\sigma(pp \rightarrow W'_R X) \times B(W'_R \rightarrow t\bar{b})$ for $g'/g = 1$ and on g'/g vs $m_{W'_R} \rightarrow$

- $g'/g > 0.13$ is excluded for $M_{W'_R} = 1 \text{ TeV}$

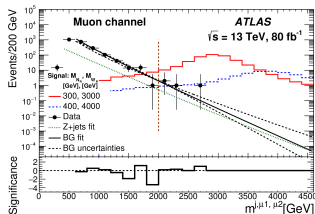
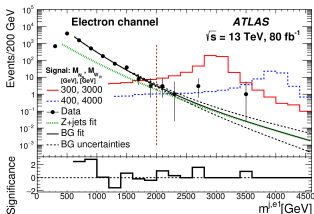
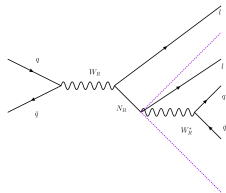
- Combination with the $W'_R \rightarrow t\bar{b} \rightarrow b\bar{b}q\bar{q}'$ search [Phys. Lett. B 781 (2018) 327] improves the $\sigma \cdot B$ limit by 35% at $M_{W'_R} > 1 \text{ TeV}$.

The $M_{W'_R}$ limit in the ZTOP model is improved from 3.15 to 3.25 TeV.



Search for W_R and a heavy neutrino

arXiv:1904.12679



• $\sqrt{s} = 13$ TeV, $L_{int} = 80/\text{fb}$

• Left-Right Symmetric Model:

- ▶ Restores P -symmetry at a high scale
- ▶ SM-singlet heavy neutrinos N_R
- ▶ Heavy gauge boson W_R coupled to right-handed fermion currents

• $m_{N_R}/m_{W_R} \leq 0.1$:

$$W_R \rightarrow N_R \ell^\pm \quad (\ell = e, \mu),$$

boosted $N_R \rightarrow \ell^\pm + \text{large-}R \text{ jet}$

- ▶ Complementary to previous searches with $m_{N_R} > m_{W_R}$
- ▶ Fully reconstructed N_R

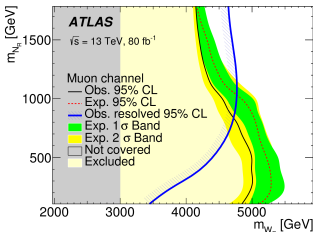
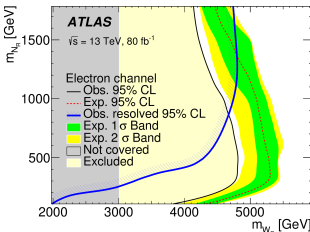
• $m_{W_R} > 3.8 - 5$ TeV for

$m_{N_R} = 0.1 - 1.8$ TeV at CL=95%

• \uparrow Signal regions: $m_{W_R, reco} > 2$ TeV, same ℓ flavour

• Control regions: $m_{W_R, reco} < 2$ TeV, same ℓ flavour.

• Validation region: all $m_{W_R, reco}$, mixed flavour, leading μ



Combination of $X \rightarrow VV/VH, \ell\nu/\ell\ell$ searches

Phys. Rev. D 98 (2018) 052008

• $\sqrt{s} = 13$ TeV, $L_{int} = 36/\text{fb}$

• Benchmark models:

▶ HVT: [see the HVT introduction](#)

(A) Weakly coupled scenario:

$$g_H \equiv g_V c_H^{1/2} = -0.56, \quad g_f = g_q = g_\ell = 0.55$$

(B) Strong $V' VH/V' HH$ coupling: $g_H = -2.9, g_f = 0.14$

(C) Fermiophobic V' : $g_f = 0, g_H = 1$

\Rightarrow VBF V' production only

▶ Bulk RS1 with the heavy spin-2 $G_{KK} \rightarrow VV$ [see the RS introduction](#)

★ Gravity and SM fields propagate in the bulk

★ G_{KK} coupling to SM $T^{\mu\nu} \propto k/\bar{M}_{Pl}$, where k is the curvature of the extra dimension

★ G_{KK} couplings to light fermions are suppressed, $gg, WW/ZZ$ and $t\bar{t}$ preferred

▶ Empirical: heavy scalar singlet $\rightarrow VV$

• Combination of channels:

$WZ \rightarrow \ell\nu qq, \ell\nu\ell\ell, qqqq, WW \rightarrow \ell\nu\nu\nu, \ell\nu qq, qqqq, ZZ \rightarrow \ell\ell qq, \ell\ell\nu\nu, \ell\ell\ell\ell, qqqq, \nu\nu qq$

$WH \rightarrow \ell\nu bb, qqbb, ZH \rightarrow \ell\ell bb, qqbb, \nu\nu bb$

$V' \rightarrow \ell\nu, \ell\ell$

• The limits:

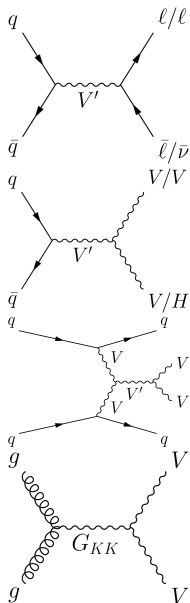
▶ HVT: $m_{V'} > 5.5$ TeV (weakly coupled) and > 4.5 TeV (strongly coupled) at CL=95%

▶ Bulk RS1 with $k/\bar{M}_{Pl} = 1$: $m_{G_{KK}} > 2.3$ TeV at CL=95% [backup](#)

▶ Heavy scalar: cross section limits for $gg \rightarrow S$ and $VV \rightarrow S$ [backup](#)

▶ Limits on V' couplings to fermions and gauge bosons, [backup](#)

better than the limits from EW fits



[▶ backup](#)

[▶ backup](#)

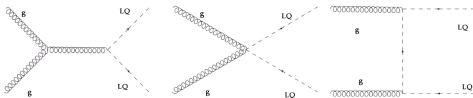
[▶ backup](#)

[▶ backup](#)

Pair production of 1st and 2nd generation scalar leptoquarks

arXiv:1902.00377
CERN-EP-2018-262

- $\sqrt{s} = 13$ TeV, $L_{int} = 36/\text{fb}$
- The benchmark model: 3 generations of scalar LQs, only same-generation couplings, $LQ_1 \rightarrow eq/\nu_e q$, $LQ_2 \rightarrow \mu q/\nu_\mu q$



- ▶ Parameters: M_{LQ} , $\beta = \text{Br}(LQ \rightarrow \ell^\pm q)$
- ▶ Production modes: $gg \rightarrow LQ\bar{L}Q$ at $M_{LQ} < 1$ TeV, $q\bar{q} \rightarrow LQ\bar{L}Q$ dominates at higher M_{LQ}

- Search channels covering all β values:

$$eejj, \mu\mu jj, e\nu jj, \mu\nu jj$$

- Variables for the BDT ordered by the discriminating power (from high to low):

$$\underline{\ell ljj}: m_{LQ}^{\min} = \min\{m(\ell j)\}, m_{\ell\ell}, \text{subleading jet } p_T, \text{subleading lepton } p_T, m_{LQ}^{\max} = \max\{m(\ell j)\}$$

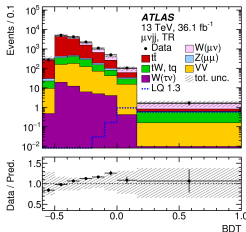
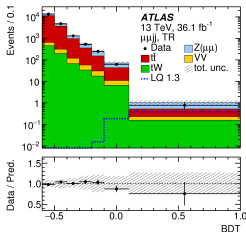
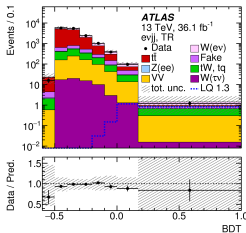
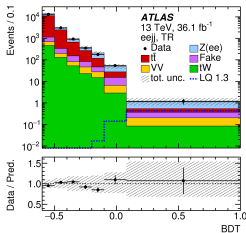
$$\underline{\ell\nu jj}: m_{LQ} = m(\ell j) \text{ and } m_{LQ}^T = m^T(\vec{E}_T^{\text{miss}}, \vec{p}_T^j) \text{ with } \ell j \text{ paired to minimize } |m_{LQ} - m_{LQ}^T|, m_T(\ell, \vec{E}_T^{\text{miss}}), E_T^{\text{miss}}, \text{subleading jet } p_T, p_T^\ell$$

- V +jets and $t\bar{t}$ backgrounds constrained in the dedicated control regions; 2×29 ℓljj and $\ell\nu jj$ signal regions (SRs) for each $\ell = e, \mu$ and for each of 29 M_{LQ} mass hypotheses. An SR is a highest BDT score bin defined individually for each M_{LQ} value. The BDT is trained for each M_{LQ} .

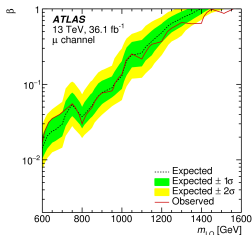
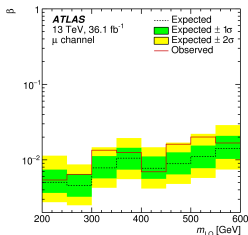
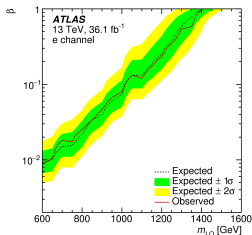
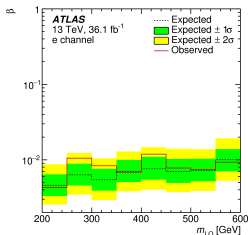
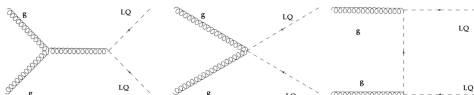
$$\underline{\ell ljj}: m_{\ell\ell} > 130 \text{ GeV}, 1 \text{ BDT bin}$$

$$\underline{\ell\nu jj}: m_T(\ell\nu) > 130 \text{ GeV}, E_T^{\text{miss}} > 150 \text{ GeV},$$

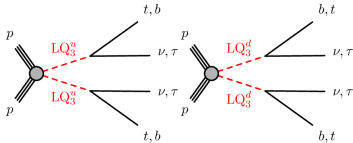
$$E_T^{\text{miss}} / \sqrt{p_T^{j1} + p_T^{j2} + p_T^\ell} > 3 \text{ GeV}^{1/2}$$



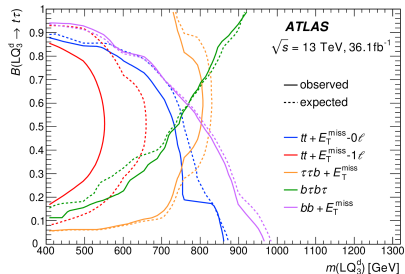
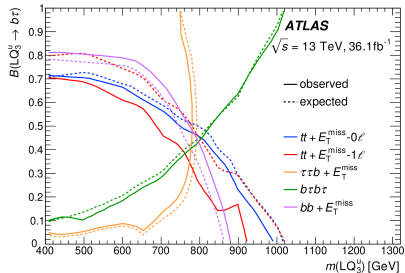
- $\sqrt{s} = 13$ TeV, $L_{int} = 36/\text{fb}$
- The benchmark model: 3 generations of scalar LQs, only same-generation couplings, $LQ_1 \rightarrow eq/\nu_e q$, $LQ_2 \rightarrow \mu q/\nu_\mu q$
 - ▶ Parameters: M_{LQ} , $\beta = \text{Br}(LQ \rightarrow \ell^\pm q)$
 - ▶ Production modes: $gg \rightarrow LQ\bar{L}Q$ at $M_{LQ} < 1$ TeV, $q\bar{q} \rightarrow LQ\bar{L}Q$ dominates at higher M_{LQ}
- Search channels covering all β values: $eejj$, $\mu\mu jj$, $e\nu jj$, $\mu\nu jj$
- Limits on LQ_1 (upper) and LQ_2 (lower) in the $\{m_{LQ}, \beta\}$ plane \rightarrow
- For $\beta = 0.5$ $m_{LQ_{1,2}} < 1.25$ TeV are excluded at CL=95%
- The paper also includes particle-level differential cross section measurements in the phase space regions not covered by "standard" SM measurements



Pair production of 3rd generation scalar leptoquarks



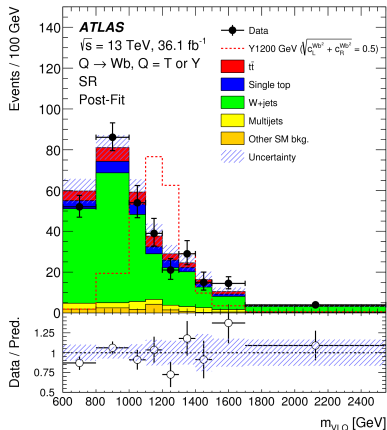
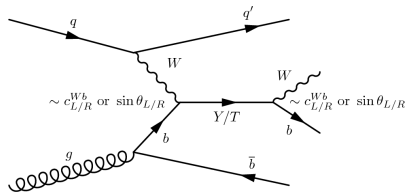
- $\sqrt{s} = 13$ TeV, $L_{int} = 36.1/\text{fb}$
- The LQ model is identical to the one in 1,2-generation LQ searches. Results presented in terms of m_{LQ} and $B(LQ \rightarrow \tau q)$.
- Multiple search channels:
 - ▶ $bb\tau\tau$ (based on $HH \rightarrow bb\tau\tau$ search)
 - ▶ $tt + E_T^{\text{miss}}$: $LQ\bar{L}\bar{Q} \rightarrow t\nu\bar{t}\nu$, reinterpretation of $t\bar{t}$ search
 - ▶ $tt + E_T^{\text{miss}}$ in 0ℓ channel: $LQ\bar{L}\bar{Q} \rightarrow t\nu\bar{t}\nu$, $LQ^d\bar{L}Q^d \rightarrow b\nu\bar{b}\nu/\tau\tau$
 - ▶ $\tau\tau b + E_T^{\text{miss}}$, reinterpretation of $t\bar{t}$ search
 - ▶ $bb + E_T^{\text{miss}}$, reinterpretation of the $b\bar{b}$ search
 - ★ 0ℓ for $LQ\bar{L}\bar{Q} \rightarrow b\nu\bar{b}\nu$ ($B = 0$)
 - ★ 1ℓ for $LQ_3^{d,u}$ with $B \neq 0$
- The main result: exclusion plots in the $\{M(LQ_3^{u,d}), B(LQ_3^{u,d} \rightarrow \tau b, \tau\tau)\}$ plane \rightarrow



Search for VLQ $\rightarrow Wb$

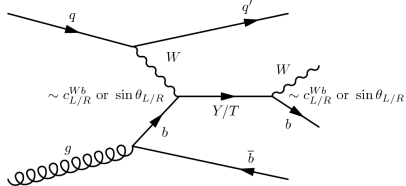
arXiv:1812.07343

- $\sqrt{s} = 13$ TeV, $L_{int} = 36/\text{fb}$
- Motivation:
 - ▶ VLQs are safe for EW precision fits: L/R isospin singlets or mass degenerate doublet
 - ▶ small Yukawa couplings \Rightarrow small contribution to $\sigma(gg \rightarrow H)$, safe for $V(H)$
 - ▶ VLQ appear in some models: extra dimensions, composite/little Higgs
- Benchmark models:
 - ▶ Renormalisable SM extension, fully parameterised by $\theta_{L,R}$ mixing angles and m_{VLQ}
 - ▶ Non-renormalisable: $c_{L,R}^{Wb}$ couplings, $B(VLQ \rightarrow Wb, Zt, Ht)$ as free parameters
- VLQ = T ($Q = +2/3$, any isospin) or Y ($Q = -4/3$, can't be an SU(2) singlet)
 - ▶ Non-negligible mixing with 3rd generation
- VLQ $\rightarrow Wb$, $W \rightarrow \ell\nu$
- Search optimised for high $m_{VLQ} \leftrightarrow$ back-to-back b -jet and ℓ^\pm
- Control regions to constrain $t\bar{t}$ and W +jets
- $m_{VLQ} \leftarrow p_z(\nu)$ from $m_{\ell\nu} = m_W$ constraint
- Interference with SM backgrounds is crucial

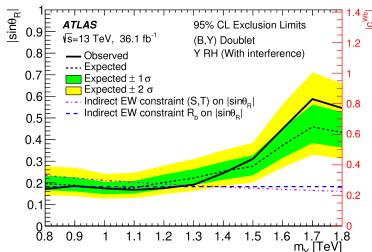
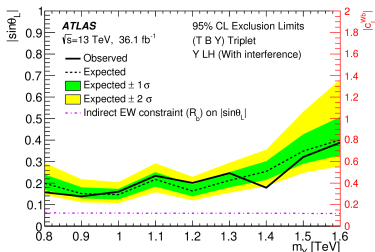
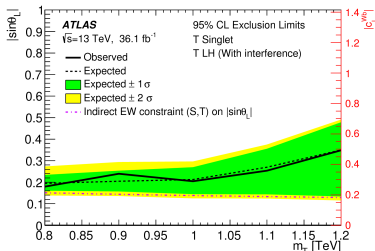


Search for VLQ $\rightarrow Wb$

arXiv:1812.07343

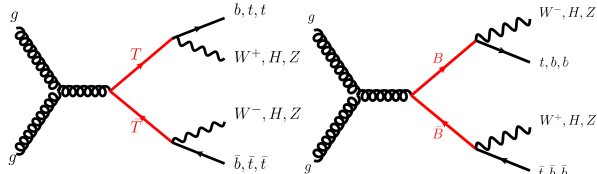


- $\sqrt{s} = 13$ TeV, $L_{int} = 36/\text{fb}$
- Benchmark models:
 - ▶ Renormalisable SM extension, fully parameterised by $\theta_{L,R}$ mixing angles and m_{VLQ}
 - ▶ Non-renormalisable: $c_{L,R}^{Wb}$ couplings, $B(VLQ \rightarrow Wb, Zt, Ht)$ as free parameters
- VLQ = T ($Q = +2/3$, any isospin) or Y ($Q = -4/3$, can't be an SU(2) singlet)
 - ▶ Non-negligible mixing with 3rd generation
- Limits on VLQ–3rd gen. mixing parameters vs m_{VLQ} for T-singlet, (T, B, Y) triplet, and (B, Y) doublet models



Combination of searches for vector-like $T\bar{T}$, $B\bar{B}$ Phys. Rev. Lett. 121 (2018) 211801

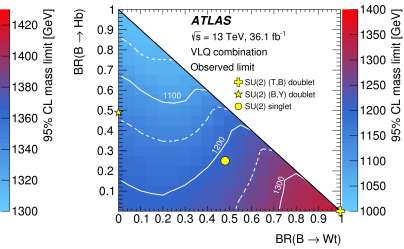
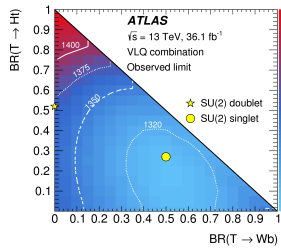
- $\sqrt{s} = 13$ TeV, $L_{int} = 36/\text{fb}$
- Combination of $T\bar{T}$ and $B\bar{B}$ searches with
 $T \rightarrow Zt/Wb/Ht$,
 $B \rightarrow Zb/Wt/Hb$
- Multiple channels:



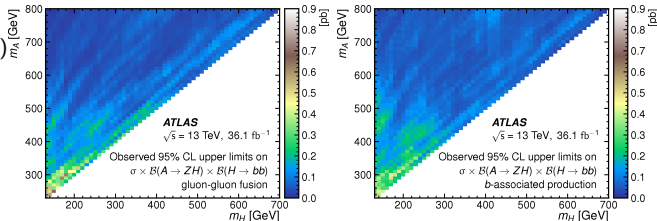
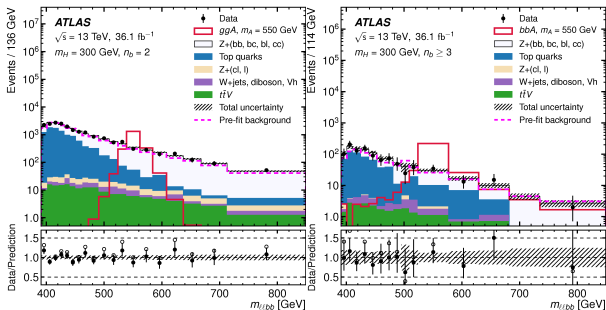
Analysis	$T\bar{T}$ decay	$B\bar{B}$ decay
$H(bb)t + X$ [16]	$HtH\bar{t}$	-
$W(\ell\nu)b + X$ [30]	$WbW\bar{b}$	-
$W(\ell\nu)t + X$ [32]	-	$WtW\bar{t}$
$Z(\nu\nu)t + X$ [33]	$ZtZ\bar{t}$	-
$Z(\ell\ell)t/b + X$ [35]	$ZtZ\bar{t}$	$ZbZ\bar{b}$
Tril./s.s. dilepton [36]	$HtH\bar{t}$	$WtW\bar{t}$
Fully hadronic [37]	$HtH\bar{t}$	$HbH\bar{b}$

- ▶ leptonic:
 - $H(bb)t + X$
 - $W(\ell\nu)b + X$,
 - $W(\ell\nu)t + X$
 - $Z(\ell\ell)t/b + X$,
 - $Z(\nu\nu)t + X$
 - $\ell^\pm\ell^\pm + \ell^\mp + X$
- ▶ fully hadronic

- Combination significantly increases sensitivity.
- Singlet T and B are excluded for $m_{T,B} < 1.31$ and 1.22 TeV. For the (T, B) doublet, $m_{T,B} < 1.37$ TeV are excluded.



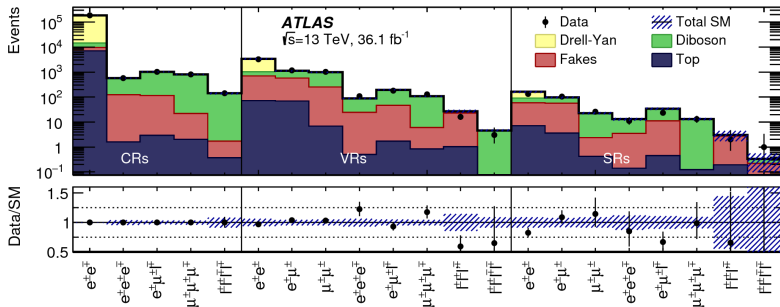
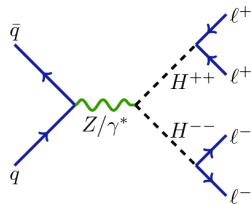
- $\sqrt{s} = 13$ TeV, $L_{int} = 36.1/\text{fb}$
- Motivation: 2HDM with $m_A > m_H > m_h$, $A \rightarrow ZH \rightarrow \ell^+\ell^- b\bar{b}$,
- Signature: $\ell^+\ell^- (\ell = e, \mu) + \geq 2b$ -jets
 $2b$ -jets for $gg \rightarrow A$ production
 $\geq 3b$ -jets for $gg \rightarrow b\bar{b}A$ production
- Scan $m_{\ell\ell b\bar{b}}$ distributions for different m_{bb} windows
- CL=95% upper limits on $\sigma(pp \rightarrow AX) \times B(A \rightarrow ZH) \times B(H \rightarrow b\bar{b})$ for ggF and b -associated A production as functions of $(m_A, m_H) \rightarrow$
- (m_A, m_H) exclusion regions vs $\tan\beta$ for Type I,II, lepton specific and flipped 2HDM [▶ see backup](#)



Search for $H^{++}H^{--}$ with same-charge leptons

Eur. Phys. J. C 78 (2018) 199

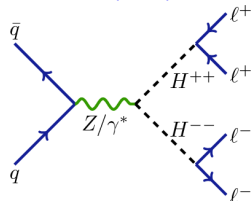
- $\sqrt{s} = 13$ TeV, $L_{int} = 36.1/\text{fb}$
- Motivation: left-right symmetric models, Higgs triplet models, little Higgs, ...
- Signature: $\ell^{\pm}\ell^{\pm}\ell'^{\mp}\ell'^{\mp}$, $\ell, \ell' = e, \mu$, b -jet veto, $Z(\ell^+\ell^-)$ veto
- SM backgrounds constrained in control regions
- Analysis regions are categorised by lepton multiplicity:
 - ▶ 2 or 3 leptons: one same-charge lepton pair
 - ▶ 4 leptons: two same charge lepton pairs, $\sum Q_{\ell} = 0$



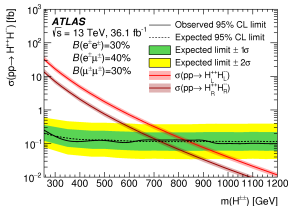
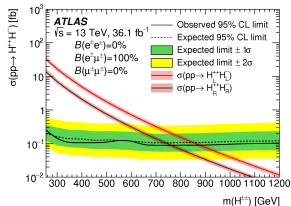
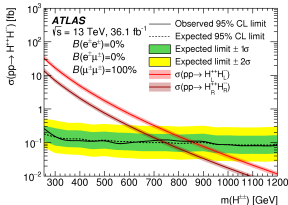
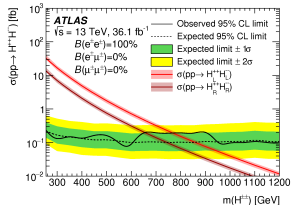
▶ see backup for region definitions and $m(\ell\ell')$ spectra

Search for $H^{++}H^{--}$ with same-charge leptons

Eur. Phys. J. C 78 (2018) 199



- $\sqrt{s} = 13$ TeV, $L_{int} = 36.1/\text{fb}$
- Motivation: left-right symmetric models, Higgs triplet models, little Higgs, ...
- Signature: $l^{\pm}l^{\pm}l'^{\mp}l'^{\mp}$, $l, l' = e, \mu$, b -jet veto, $Z(l^+l^-)$ veto
- CL=95% upper limits on $\sigma(pp \rightarrow H^{++}H^{--})$ for several $B(ee)/B(e\mu)/B(\mu\mu)$ ratio values



$m(H_L^{\pm\pm}) > 0.77$ TeV and > 0.47 TeV
 $m(H_R^{\pm\pm}) > 0.66$ TeV and > 0.32 TeV,
 for $B(H^{\pm\pm} \rightarrow l^{\pm}l^{\pm}) = 1$ and ≥ 0.1 ,
 resp.

Summary

- ATLAS searches for BSM resonances in final states with a charged lepton utilised 36 – 139/fb of LHC Run-II pp data at $\sqrt{s} = 13$ TeV
- **No significant deviations from the SM background hypothesis so far**
- Limits in terms of benchmark BSM models:
 - ▶ Heavy vector triplet: $M_{V'} \gtrsim 4.5$ TeV
 - ▶ Randall–Sundrum extra dimensions: $m_{G_{KK}} \gtrsim 2.3$ TeV
 - ▶ Vector-like quarks: $m_{VLQ} \gtrsim 1$ TeV
 - ▶ Leptoquarks:
 - ★ 1st, 2nd generation: $m_{LQ_{1,2}} > 1.25$ GeV for $\mathcal{B}(\ell^\pm q) = 0.5$
 - ★ 3rd generation: $m_{LQ_3} \gtrsim 0.8$ TeV for $\mathcal{B}(\ell^\pm q) = 0.5$
- Combined limits on several benchmark models using multiple final states:
 - ▶ HVT and RS1: $\{VV, VH, \ell\ell'\} \rightarrow \{VV \& VH\} \rightarrow \{VV \& VH \& \ell\ell'\}$
 - ▶ Vector-like quarks: $H(bb)t + X$ & $W(\ell\nu)b + X$, $W(\ell\nu)t + X$ & $Z(\ell\ell)t/b + X$, $Z(\nu\nu)t + X$ & $\ell^\pm\ell^\pm + \ell^\mp + X$
- Processing of the full LHC Run-II 139/fb dataset is in progress

Backup

Search for $W' \rightarrow \tau\nu$

PRL 120 (2018) 161802

• $\sqrt{s} = 13$ TeV, $L_{int} = 36.1/\text{fb}$

• The discriminating variable:

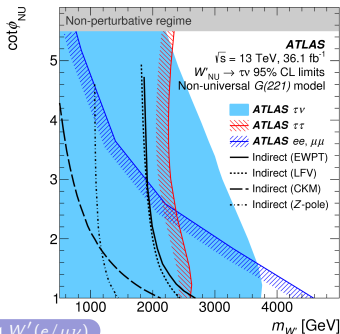
$$M_T(\tau_{had-vis}\nu) = \sqrt{2p_T^{\tau_{had-vis}} E_T^{\text{miss}} (1 - \cos \phi_{\tau\nu})}$$

• Assuming SM couplings for the W'_{SSM} :

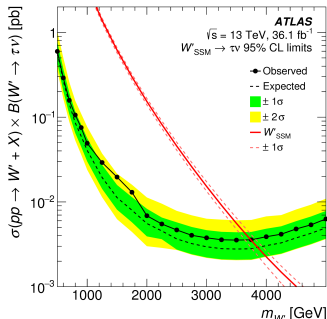
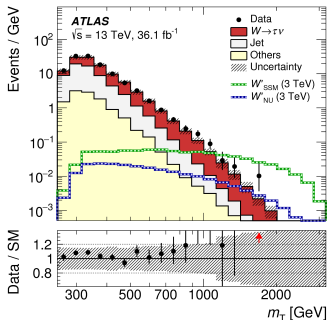
$$M_{W'_{SSM}} > 3.7 \text{ TeV at CL=95\%}$$

• \downarrow Excluded region for $G(221)$ W'_{NU} non-universally

coupled to 1st/2nd and 3rd generation fermions.
 $G(221)$ model: $SU(2)_1 \times SU(2)_2 \times U(1)_Y$ gauge group, '1' couples to 1st and 2nd gen. fermions $\propto g_1$, '2' to 3rd gen. $\propto g_2$, $\tan \phi_{NU} = g_2/g_1$. Upon breaking to $SU(2)_L \times U(1)_Y$ W' couples to 3rd gen. with $g \cot \phi_{NU}$ and to 1st and 2nd gen. with $g \tan \phi_{NU}$

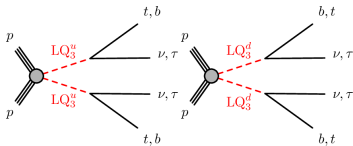


$\leftarrow W'(e/\mu\nu)$



Related topic: search for LFV $Z \rightarrow \tau e / \tau \mu$, Phys. Rev. D 98 (2018) 092010

Pair production of 3rd generation scalar leptoquarks



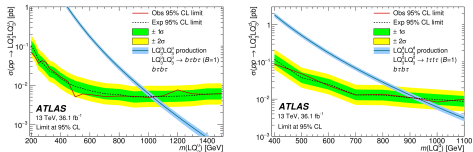
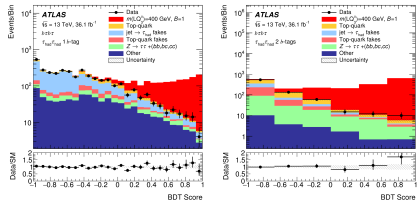
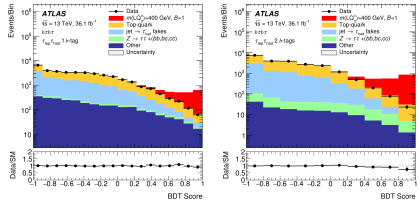
- $\sqrt{s} = 13$ TeV, $L_{int} = 36.1/\text{fb}$
- The LQ model is identical to the one in 1,2-generation LQ searches. Results presented in terms of m_{LQ} and $B(LQ \rightarrow \tau q)$.

- Multiple search channels:

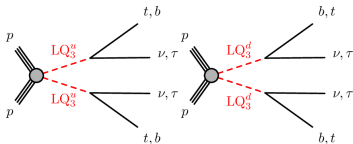
$bb\tau\tau$ channel based on

$HH \rightarrow bb\tau_{lep/had}\tau_{had}$ search:

- ▶ Event subsamples with 1 and 2 b -tagged jets
- ▶ no additional jet veto \Rightarrow sensitive to both $LQ^u L\bar{Q}^u \rightarrow b\tau b\tau$ and $LQ^d L\bar{Q}^d \rightarrow t\tau t\tau \rightarrow W(\rightarrow \text{jets})b\tau W(\rightarrow \text{jets})b\tau$
- ▶ Variables: S_T , $m_{T(\ell),jet}$, $\Delta R_{T(\ell),jet}$, $\Delta\phi(E_T^{miss}, \tau)$, $p_T^\tau \Rightarrow$ **BDT score**



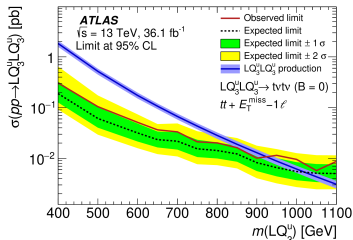
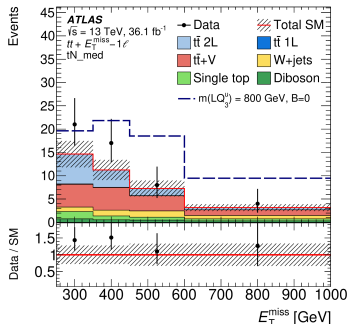
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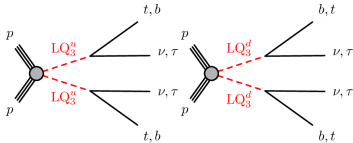
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- Multiple search channels:

$$tt + E_T^{\text{miss}} \text{ channel} \quad LQ\bar{LQ} \rightarrow t\nu\bar{t}\nu$$

- ▶ reinterpretation of $\tilde{t}\tilde{t}$ search
- ▶ $t(\ell\nu b)t(\text{all hadronic}) \rightarrow e/\mu + \geq 1 b\text{-jet} + \geq 3\text{jets} + E_T^{\text{miss}}$
- ▶ two SRs targeting $m_{LQ} \sim 600$ GeV (SR binned in E_T^{miss}) and $m_{LQ} \sim 1$ TeV (single bin)



Pair production of 3rd generation scalar leptoquarks



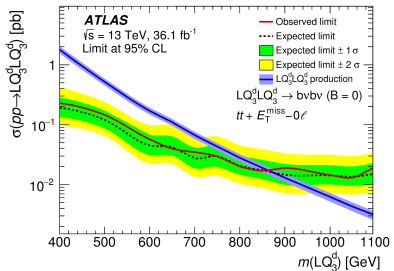
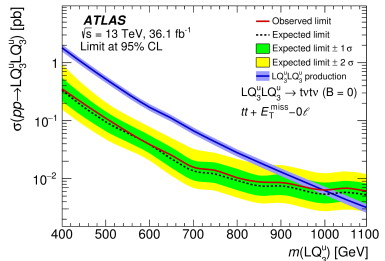
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Results presented in terms of m_{LQ} and $B(LQ \rightarrow \tau q)$.

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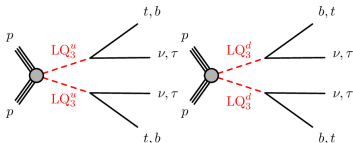
- $tt + E_T^{\text{miss}} + 0\ell$ channel $LQ\bar{L}Q \rightarrow t\nu\bar{t}\nu$,

$$LQ^d\bar{L}Q^d \rightarrow b\nu\bar{b}\nu/\tau\tau$$

- ▶ reinterpretation of $t\bar{t}$ search
- ▶ $t(\text{all had.})\bar{t}(\text{all had.}) \rightarrow \geq 1b\text{-jet} + \geq 3\text{jets} + E_T^{\text{miss}}$
- ▶ Variables: E_T^{miss} , $\Delta R(b, b)$, $m_T(b, E_T^{\text{miss}})$ for b -jet with min/max $\Delta\phi(b, E_T^{\text{miss}})$
- ▶ Three SRs:
 $m_{LQ} \sim 1$ TeV: high E_T^{miss} , boosted t 's $\rightarrow 2$ R=1.2 jets;
 $m_{LQ} \sim 0.5$ TeV: lower E_T^{miss} , lower boost of t 's $\rightarrow 2$ R=1.2 jets;
for $LQ^d \rightarrow b\nu$: ≥ 5 jets including 2 high- p_T b -jets + high E_T^{miss}



Pair production of 3rd generation scalar leptoquarks

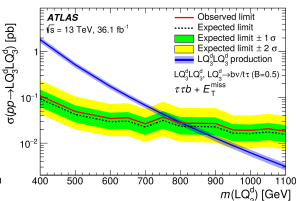
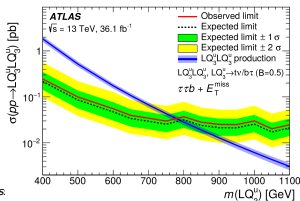
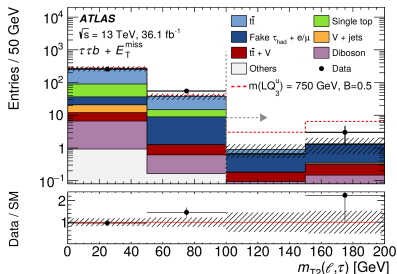


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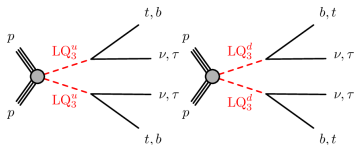
Multiple search channels:

$\tau\tau b + E_T^{miss}$ channel

- ▶ reinterpretation of $t\bar{t}$ search
- ▶ $\tau_{lep/had}^\pm \tau_{had}^\mp + \geq 1b\text{-jet} + \geq 1\text{jets} + E_T^{miss}$
- ▶ Variables: E_T^{miss} , $m_{T2}(\tau\tau, E_T^{miss})$



◀ To main slide

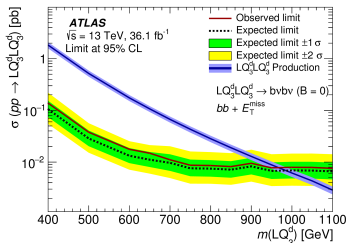
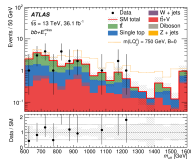
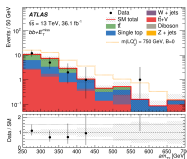
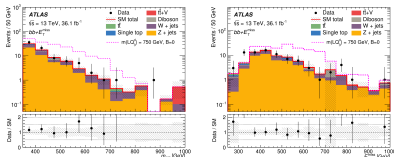


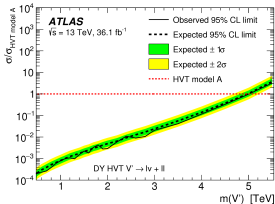
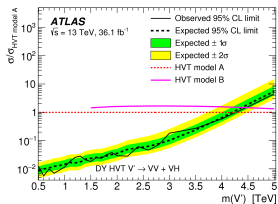
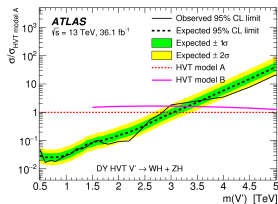
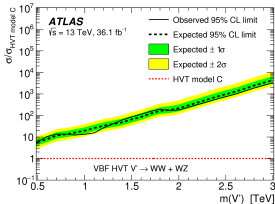
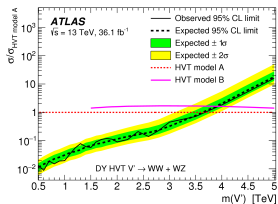
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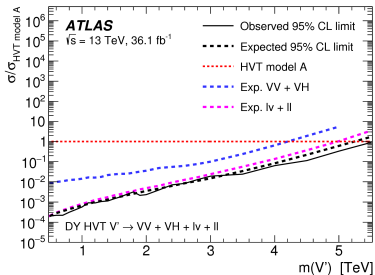
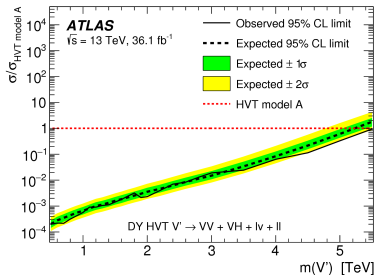
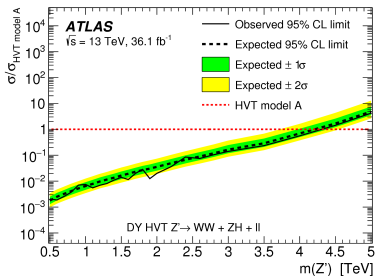
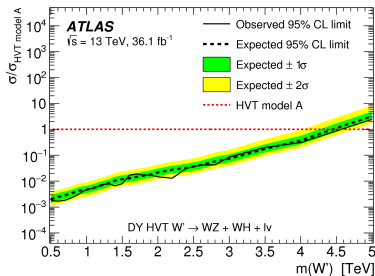
$bb + E_T^{\text{miss}}$ channel

- ▶ reinterpretation of the $b\bar{b}$ search
- ▶ $2b\text{-jets} + 0/1\ell + E_T^{\text{miss}}$,
0 ℓ for $LQ\bar{L}Q \rightarrow b\nu b\nu$ ($B = 0$),
1 ℓ for $LQ_3^{d,u}$ with $B \neq 0$
- ▶ Main discriminating variables:
cotransverse mass m_{CT} (0 ℓ), E_T^{miss} (1 ℓ)

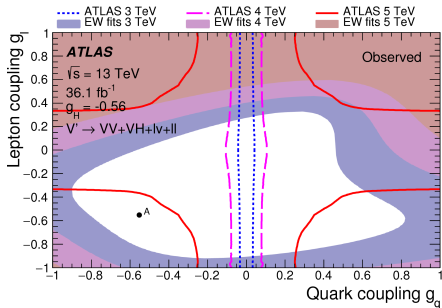
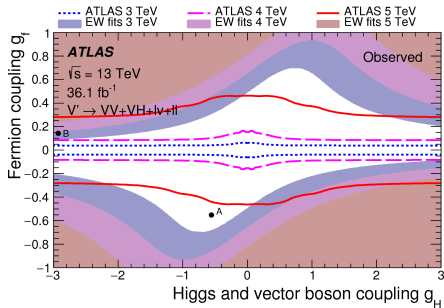




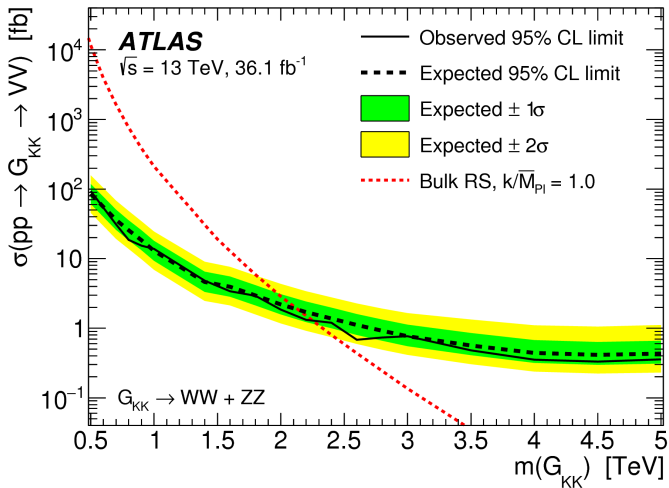
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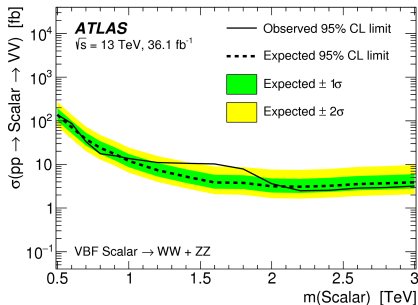
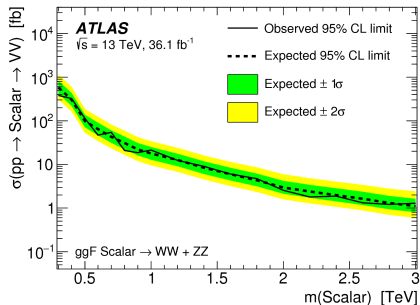
Combination of $X \rightarrow VV/VH, \ell\nu/\ell\ell$ searches: HVT couplings *Phys. Rev. D* 98 (2018) 052008



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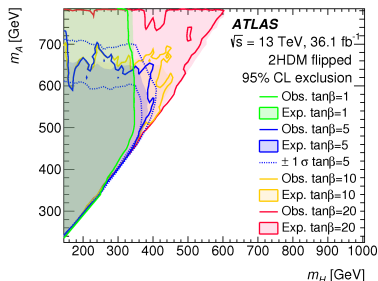
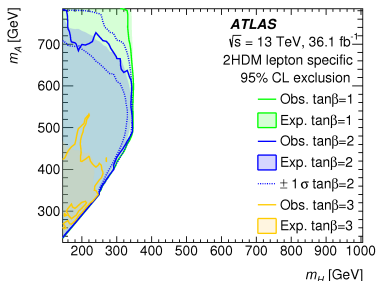
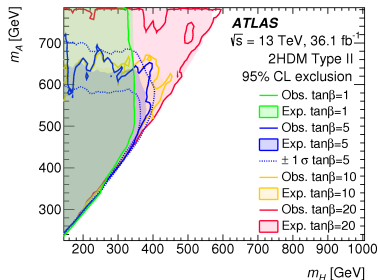
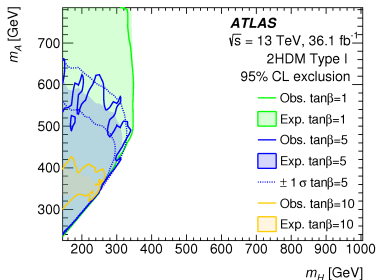
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(m_A, m_H) exclusion regions vs $\tan\beta$ for Type I,II, lepton specific and flipped 2HDM

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Region Channel	Control Regions			Validation Regions			Signal Regions		
	OCCR	DBCR	4LCR	SCVR	3LVR	4LVR	1P2L	1P3L	2P4L
Electron channel	$e^{\pm}e^{\mp}$	$e^{\pm}e^{\pm}e^{\mp}$		$e^{\pm}e^{\pm}$	$e^{\pm}e^{\pm}e^{\mp}$		$e^{\pm}e^{\pm}$	$e^{\pm}e^{\pm}e^{\mp}$	
Mixed channel	-	$e^{\pm}\mu^{\pm}\ell^{\mp}$	$\ell^{\pm}\ell^{\pm}$ $\ell^{\mp}\ell^{\mp}$	$e^{\pm}\mu^{\pm}$	$e^{\pm}\mu^{\pm}\ell^{\mp}$ $\ell^{\pm}\ell^{\pm}\ell'^{\mp}$	$\ell^{\pm}\ell^{\pm}$ $\ell^{\mp}\ell^{\mp}$	$e^{\pm}\mu^{\pm}$	$e^{\pm}\mu^{\pm}\ell^{\mp}$ $\ell^{\pm}\ell^{\pm}\ell'^{\mp}$	$\ell^{\pm}\ell^{\pm}$ $\ell^{\mp}\ell^{\mp}$
Muon channel	-	$\mu^{\pm}\mu^{\pm}\mu^{\mp}$		$\mu^{\pm}\mu^{\pm}$	$\mu^{\pm}\mu^{\pm}\mu^{\mp}$		$\mu^{\pm}\mu^{\pm}$	$\mu^{\pm}\mu^{\pm}\mu^{\mp}$	
$m(e^{\pm}e^{\pm})$ [GeV]	[130, 2000]	[90, 200]		[130, 200]	[90, 200]		[200, ∞)	[200, ∞)	
$m(\ell^{\pm}\ell^{\pm})$ [GeV]	-	[90, 200]	[60, 150]	[130, 200]	[90, 200]	[150, 200]	[200, ∞)	[200, ∞)	[200, ∞)
$m(\mu^{\pm}\mu^{\pm})$ [GeV]	-	[60, 200]		[60, 200]	[60, 200]		[200, ∞)	[200, ∞)	
b -jet veto	✓	✓	✓	✓	✓	✓	✓	✓	✓
Z veto	-	inverted	-	-	✓	-	-	✓	✓
$\Delta R(\ell^{\pm}, \ell^{\pm}) < 3.5$	-	-	-	-	-	-	✓	✓	-
$p_{\text{T}}(\ell^{\pm}\ell^{\pm}) > 100$ GeV	-	-	-	-	-	-	✓	✓	-
$\sum p_{\text{T}}(\ell) > 300$ GeV	-	-	-	-	-	-	✓	✓	-
$\Delta M/\bar{M}$ requirement	-	-	-	-	-	-	-	-	✓

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