

Searching for long-lived particles with displaced vertices in ATLAS at the LHC

featuring the Muon Spectrometer

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Searching for long-lived particles (LLPs)

in events with a displaced vertex (DV) and a muon in the ATLAS detector using pp-collisions from Run 2 of the LHC


ATLAS Run 1 Displaced Vertex + Muon — ATLAS-SUSY-2014-02


CMS Run 2 Displaced Jets 35.9 fb⁻¹ — CERN-EP-2018-289

CMS Run 2 Displaced Leptons 2.6 fb⁻¹ — CMS-PAS-EXO-16-022

We know Standard Model isn't full picture of the universe
but no signs of beyond Standard Model physics so far at the LHC
especially in obvious channels

	Run 1 2010-2012	Run 2 2015-2018	Run 3 2021-2023	HL-LHC 2026-
\sqrt{s}	7-8 TeV	13 TeV	14 TeV	14 TeV
Integrated Luminosity	26 fb ⁻¹	140 fb ⁻¹	150 fb ⁻¹	3000 fb ⁻¹

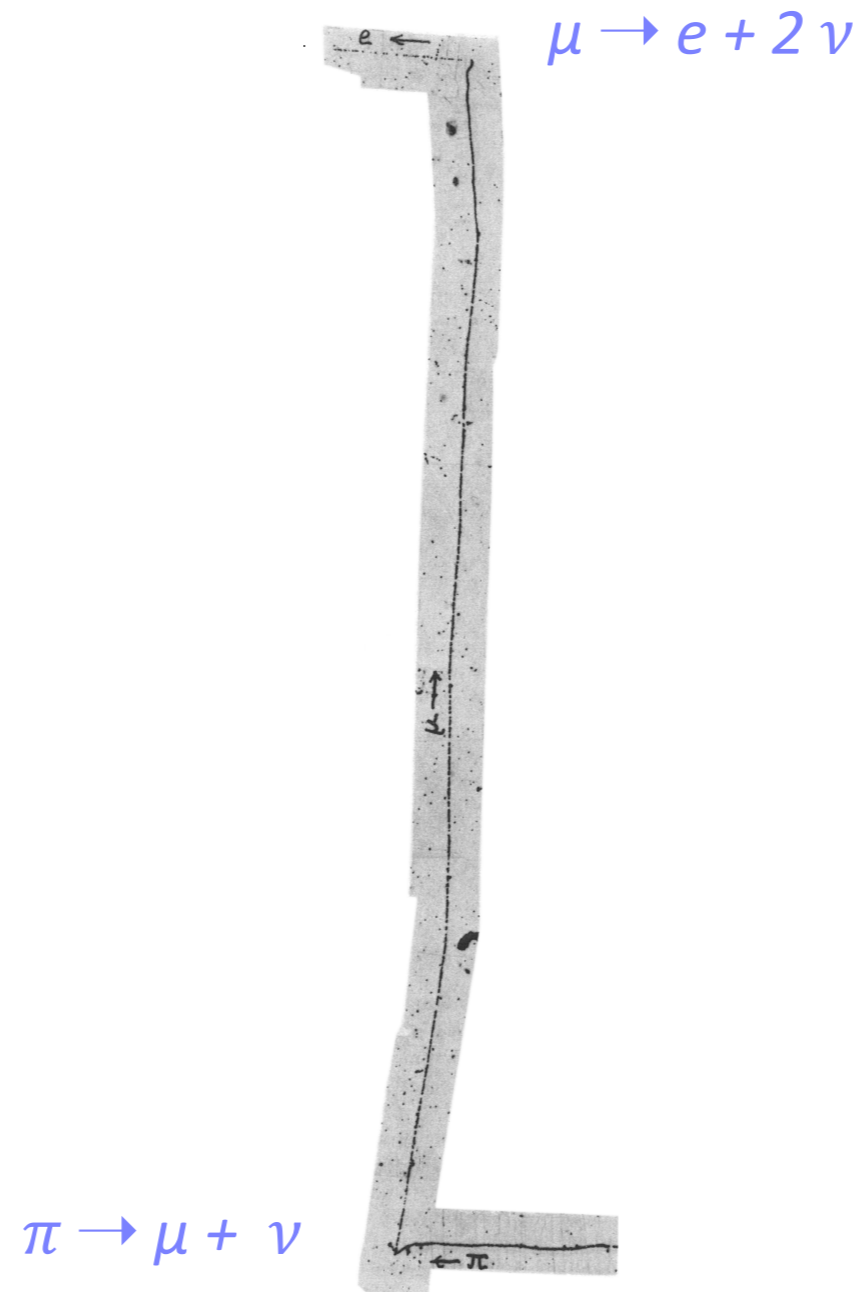
 increases in energy

 increases in dataset size

Important that we consider models with
challenging final states

Looking for long-lived particles
well motivated from physics perspective
and also exciting from an experimental perspective

Standard Model already full of long-lived particles



(1949) *Nature* **163**, 82.

Standard Model already full of long-lived particles

small couplings

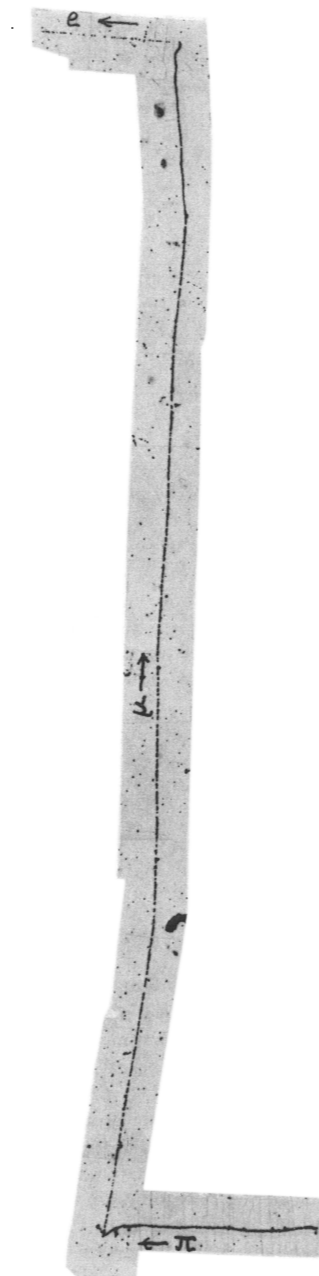
b-mesons, off-diagonal CKM, $\tau \approx \text{ps}$

high mass mediator

μ, π , via W, $\tau \approx 2 \mu\text{s}, 26 \text{ ns}$

small mass splittings

neutron, $m_n - m_p \approx 1 \text{ MeV}$, $\tau = 15 \text{ min}$



(1949) *Nature* **163**, 82.

Standard Model already full of long-lived particles

These same mechanisms come into play with new physics

small couplings

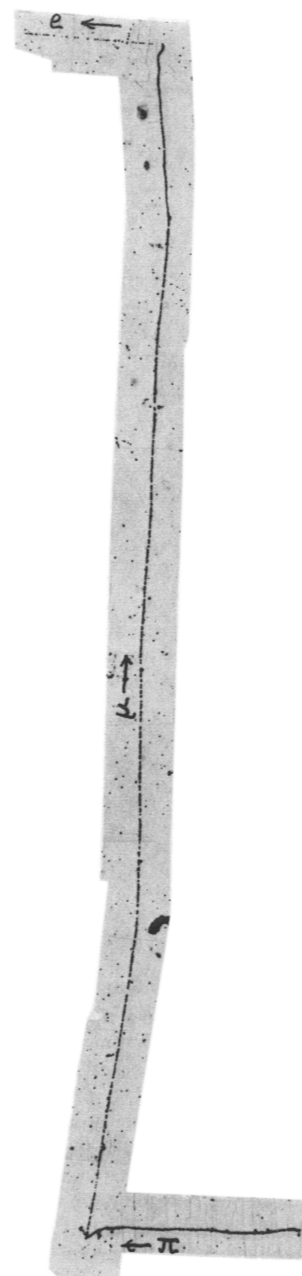
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small mass splittings

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Supersymmetry

Hidden Sectors

Dark Matter

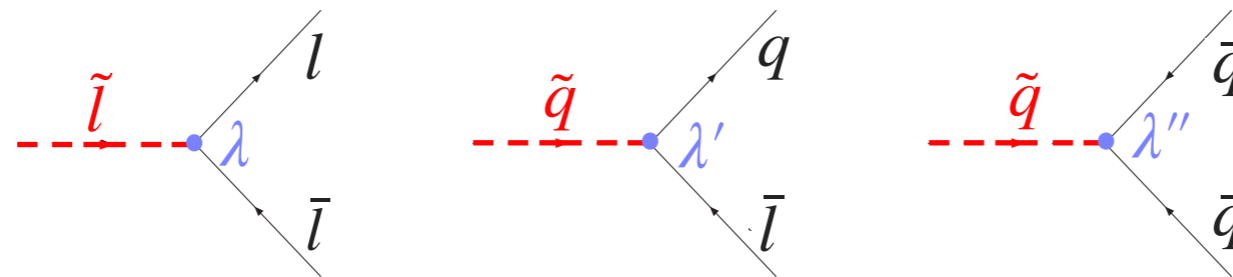
etc

(1949) *Nature* **163**, 82.

SUSY: rough idea, new relationship between fermions and bosons, fundamental Standard Model particles get a super partner

R-parity Violation: if we write down SUSY in most generic form we get the following lepton & baryon number violating couplings

$$W_{\Delta B, L} = \lambda_{ijk} L_i L_j E_k + \lambda'_{ijk} L_i Q_j D_k + \lambda''_{ijk} U_i D_j D_k + \kappa_i L_i H_u$$






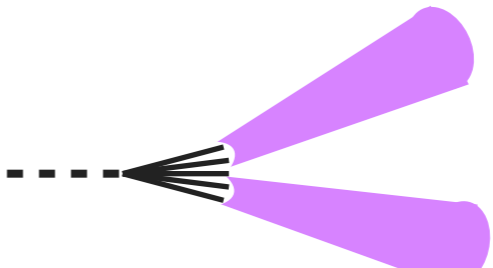
L, E = leptons
Q, D, E = quarks
H = Higgs

The only way lightest SUSY particle can decay
is to Standard Model particles via R-parity Violating couplings
we think couplings are likely small \rightarrow ps-ns lifetimes \rightarrow displaced vertices

R-parity = +1 for regular particles
R-parity = -1 for superpartners

ATLAS has several analyses looking for displaced vertices

DV+muon uniquely sensitive to λ'_{2jk}

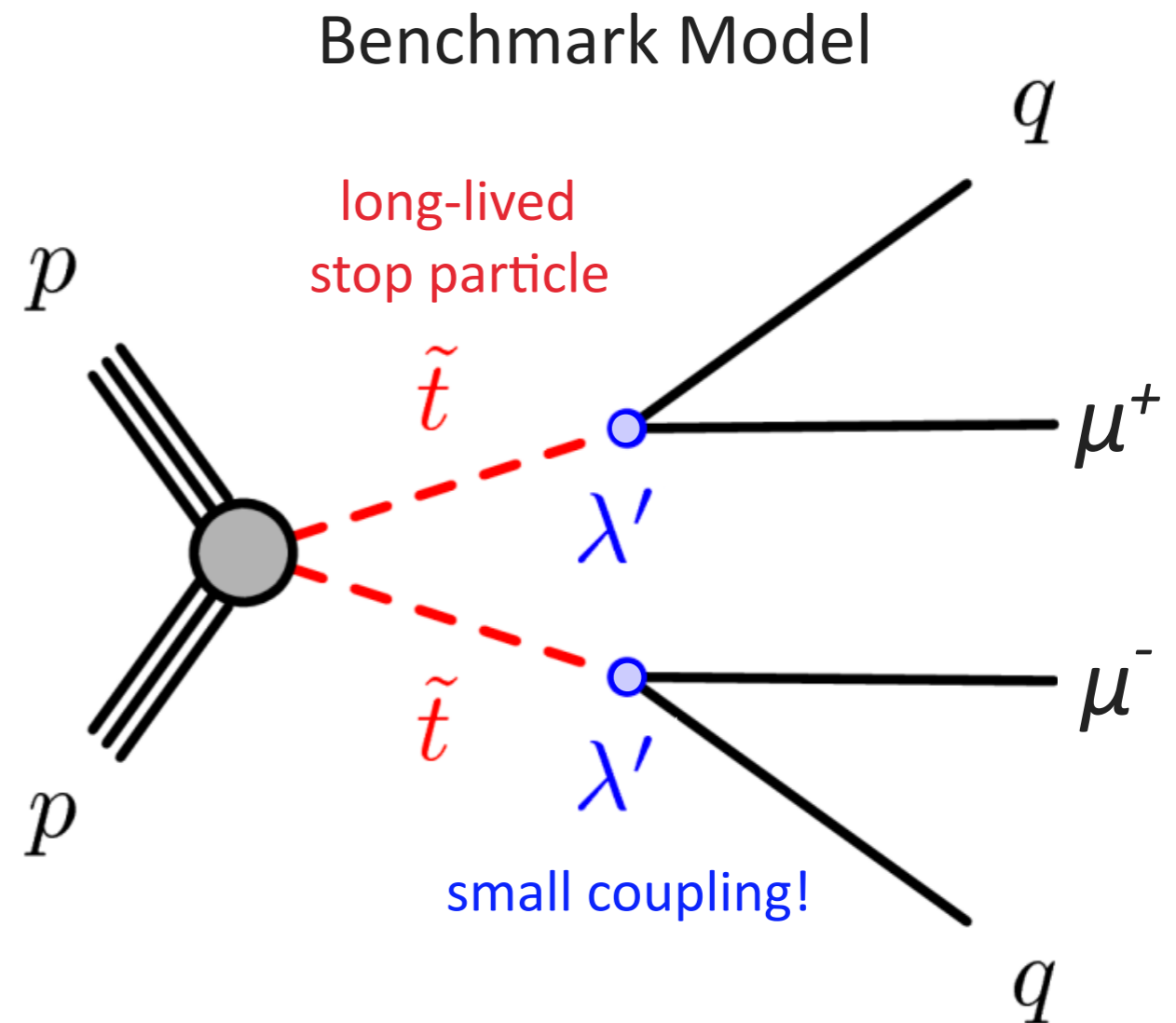
Coupling	Physics Signature	Example Detector Signature
λ_{ijk}	leptonic LLP decays $ee, \mu\mu, \mu e, \dots$	 di-lepton DV
λ'_{1jk}	semi-leptonic	 DV+electron
λ'_{2jk}	LLP decays lq, lqq	 DV+muon
λ''_{ijk}	hadronic LLP decays qq, qqq	 DV+jets

Goal: probing λ'_{2jk}

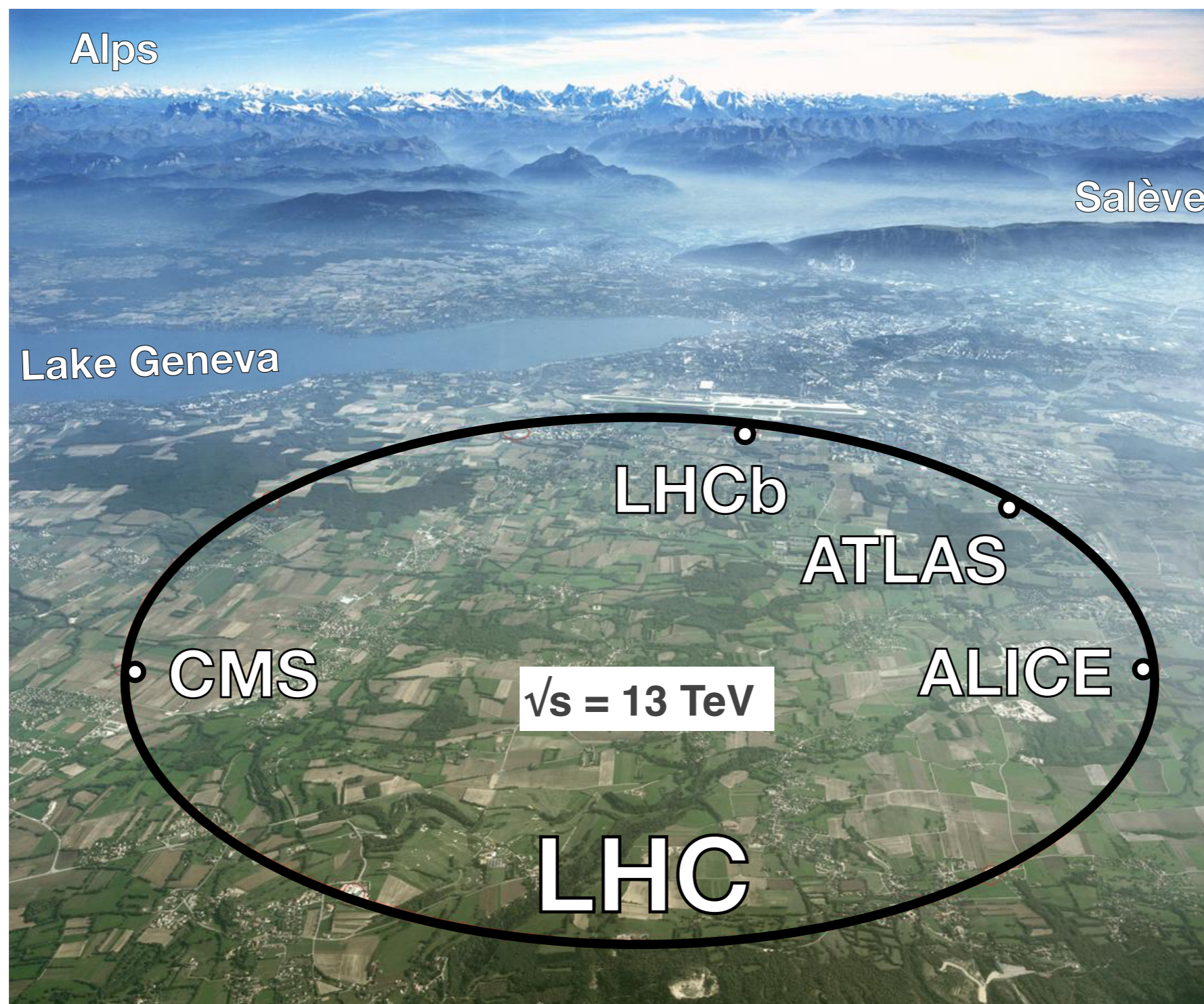
As a benchmark, consider pair production of stop particles

stop = top quark partner
lightest SUSY particle
and long-lived

stop decays to a jet and a muon
via a small λ' coupling

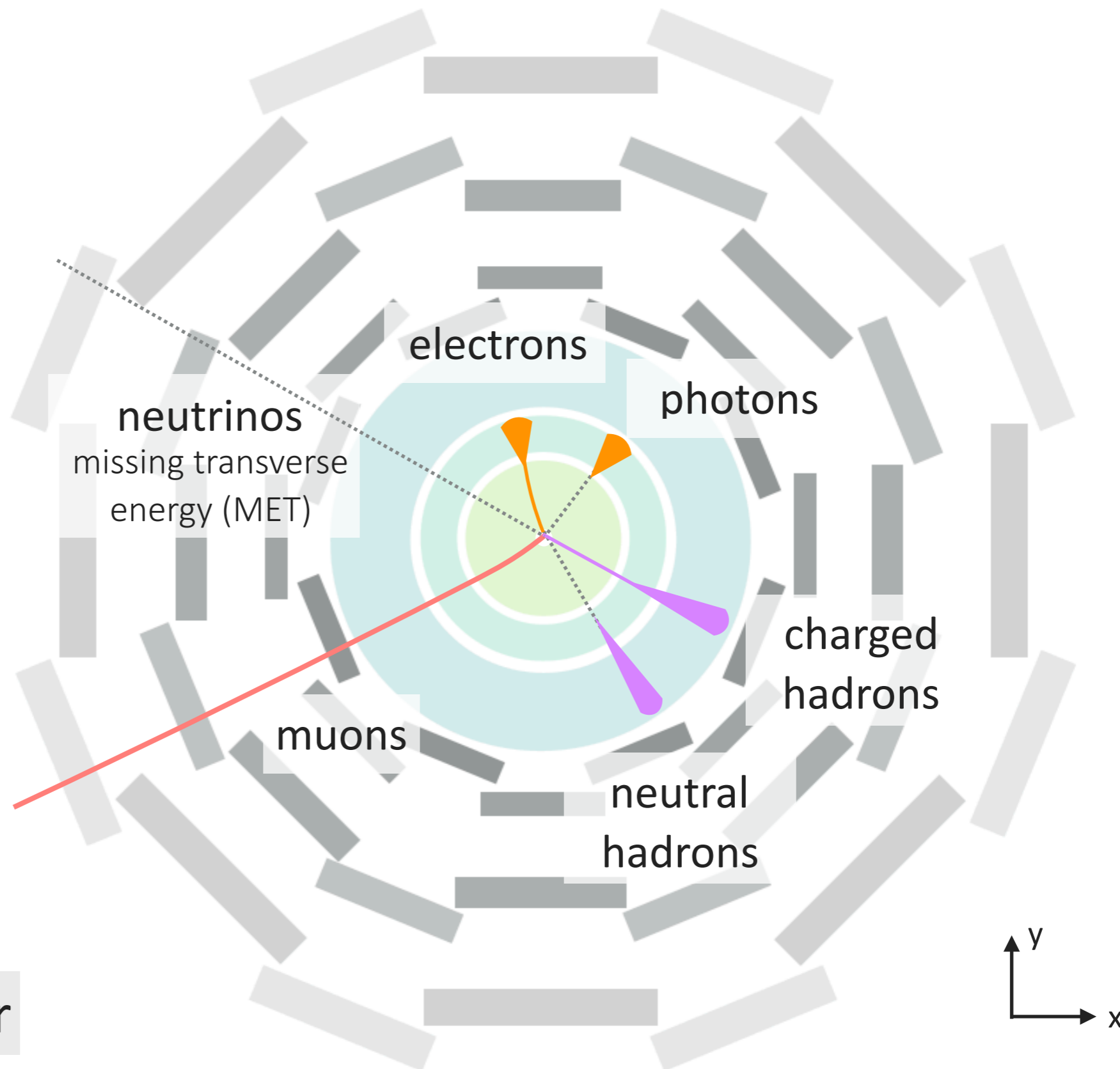


Note: stop hadronizes with Standard Model particles to form a color singlet state \rightarrow R-hadron



our analysis using 137 fb^{-1} of pp-collisions from 2016-2018
this talk uses preliminary results with 77.3 fb^{-1} from 2016-2017

4π coverage
cylindrical geometry
barrel + endcaps



Inner Detector

Electromagnetic
Calorimeter

Hadronic
Calorimeter

Muon Spectrometer

electrons

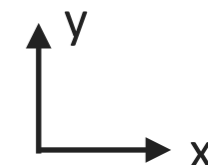
photons

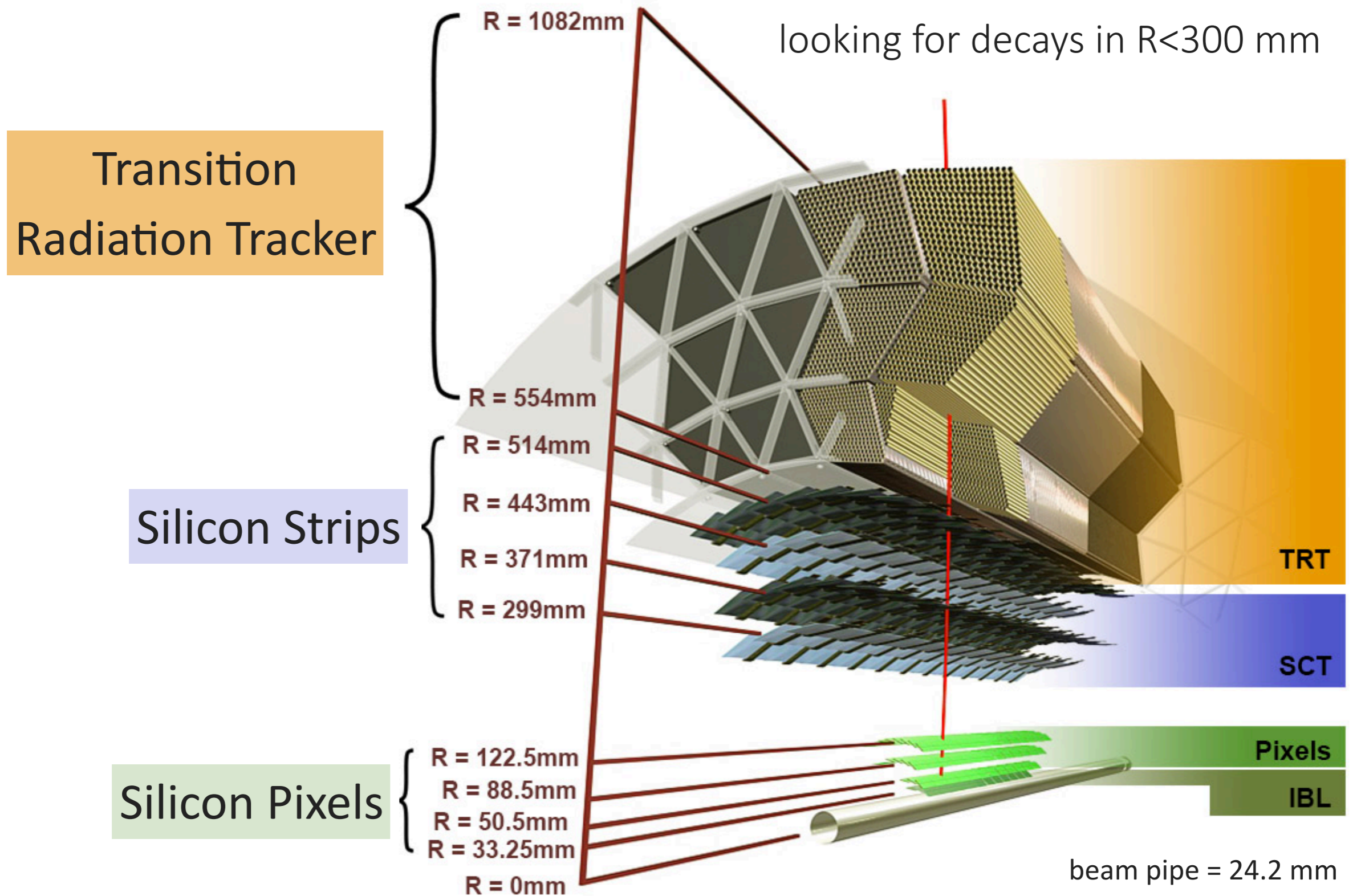
neutrinos
missing transverse
energy (MET)

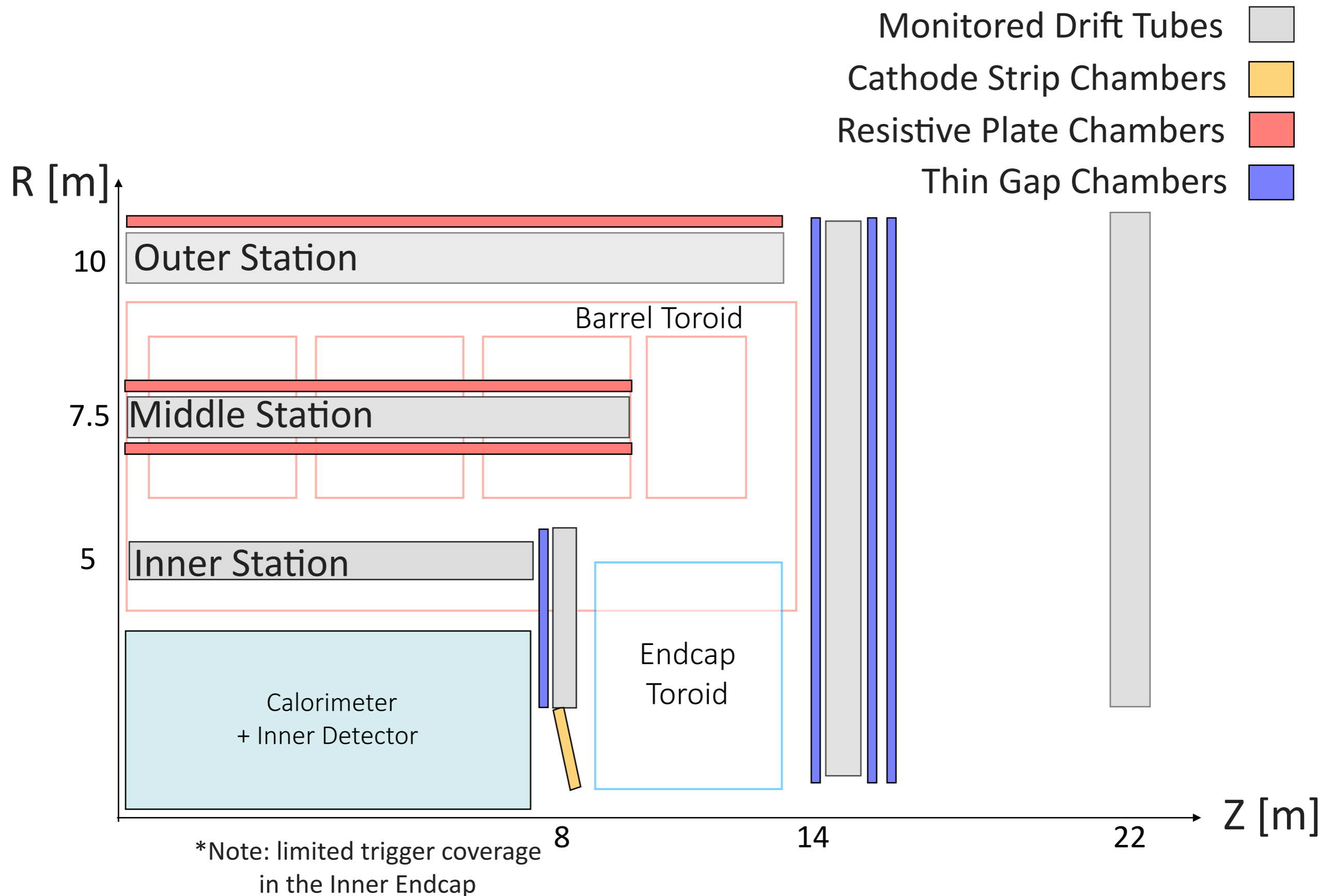
charged
hadrons

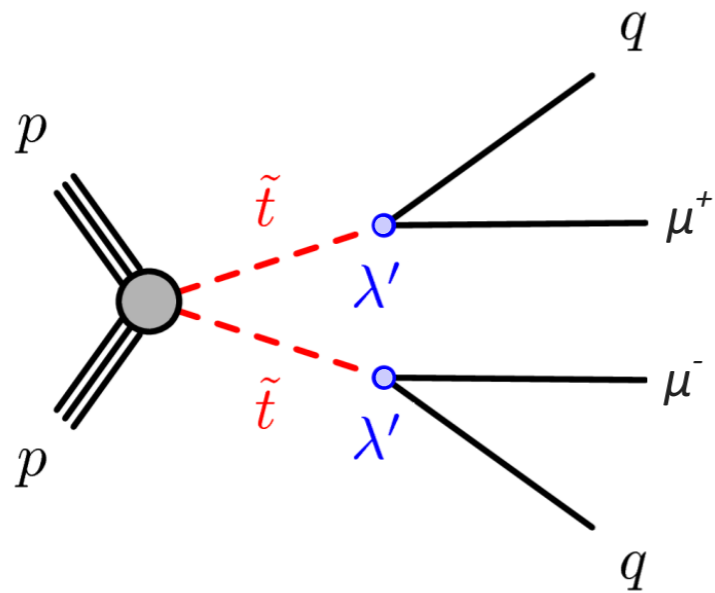
muons

neutral
hadrons



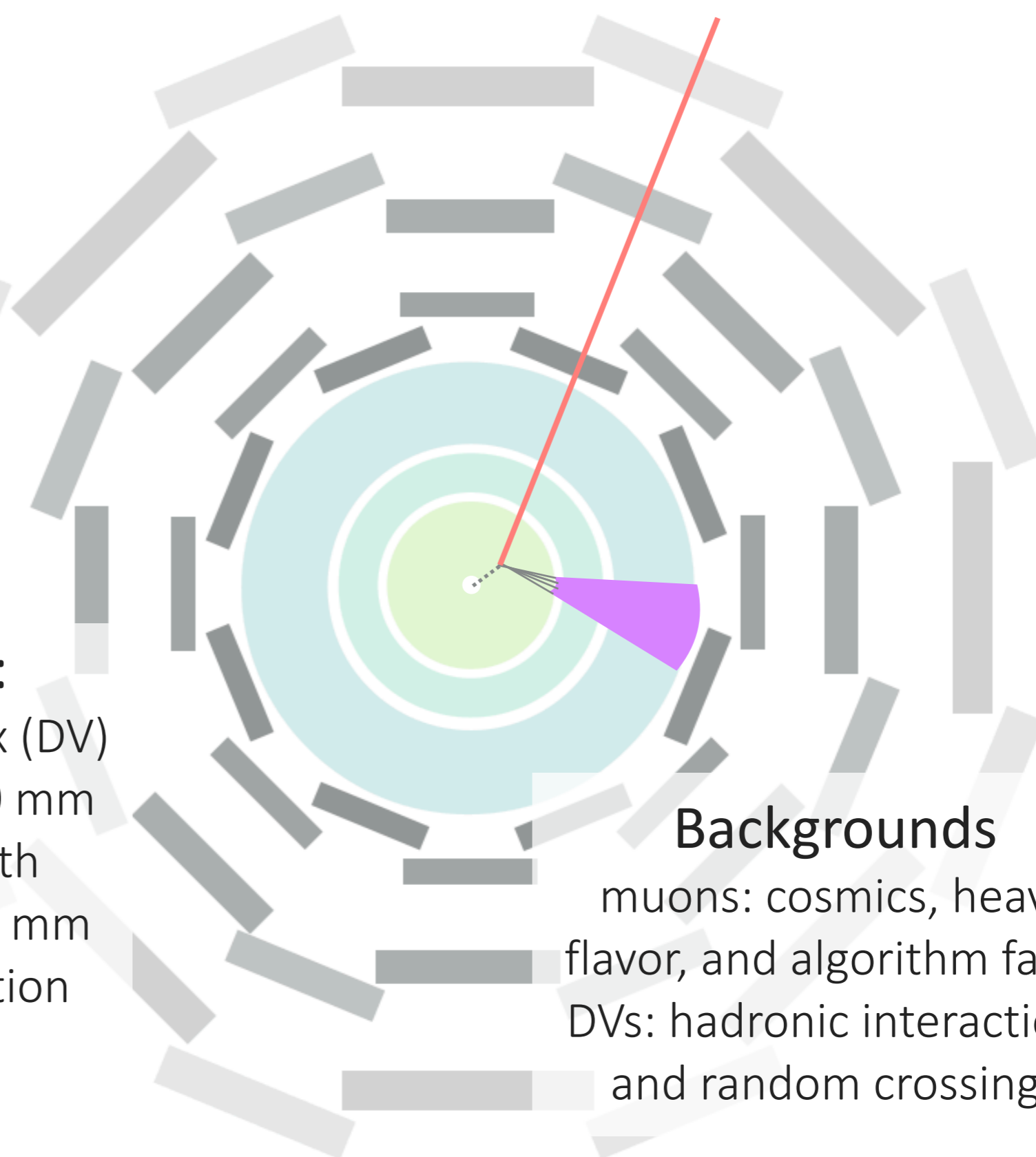






Detector Signature:

at least one Displaced Vertex (DV)
in $R < 300$ mm and $|Z| < 300$ mm
and a displaced muon with
impact parameter, $|d_0| > 2$ mm
*non-standard reconstruction



Backgrounds

muons: cosmics, heavy
flavor, and algorithm fakes
DVs: hadronic interactions
and random crossings

In general

use stop $\rightarrow \mu + \text{jet}$ as a benchmark
but remain open minded to other signals
eg. neutral LLP $\rightarrow \mu + 2 \text{ jets}$, cascade decays, etc

We know from previous versions of the analysis

we can have ~ 0 expected background
and retain excellent signal efficiency

Define two levels of selection for vertices and muons

1. preselection - loose, lets us study backgrounds
2. full selection - tight, strong background rejection

in Run 1: Muon Spectrometer Only Trigger

requires Muon Track only - agnostic to Inner Detector activity

compare this to Standard Muon Triggers, $|d_0| < 10$ mm

challenge: large background rate in endcaps, only use barrel

Trigger acceptance

defined by

$$|\eta| < 1.05$$

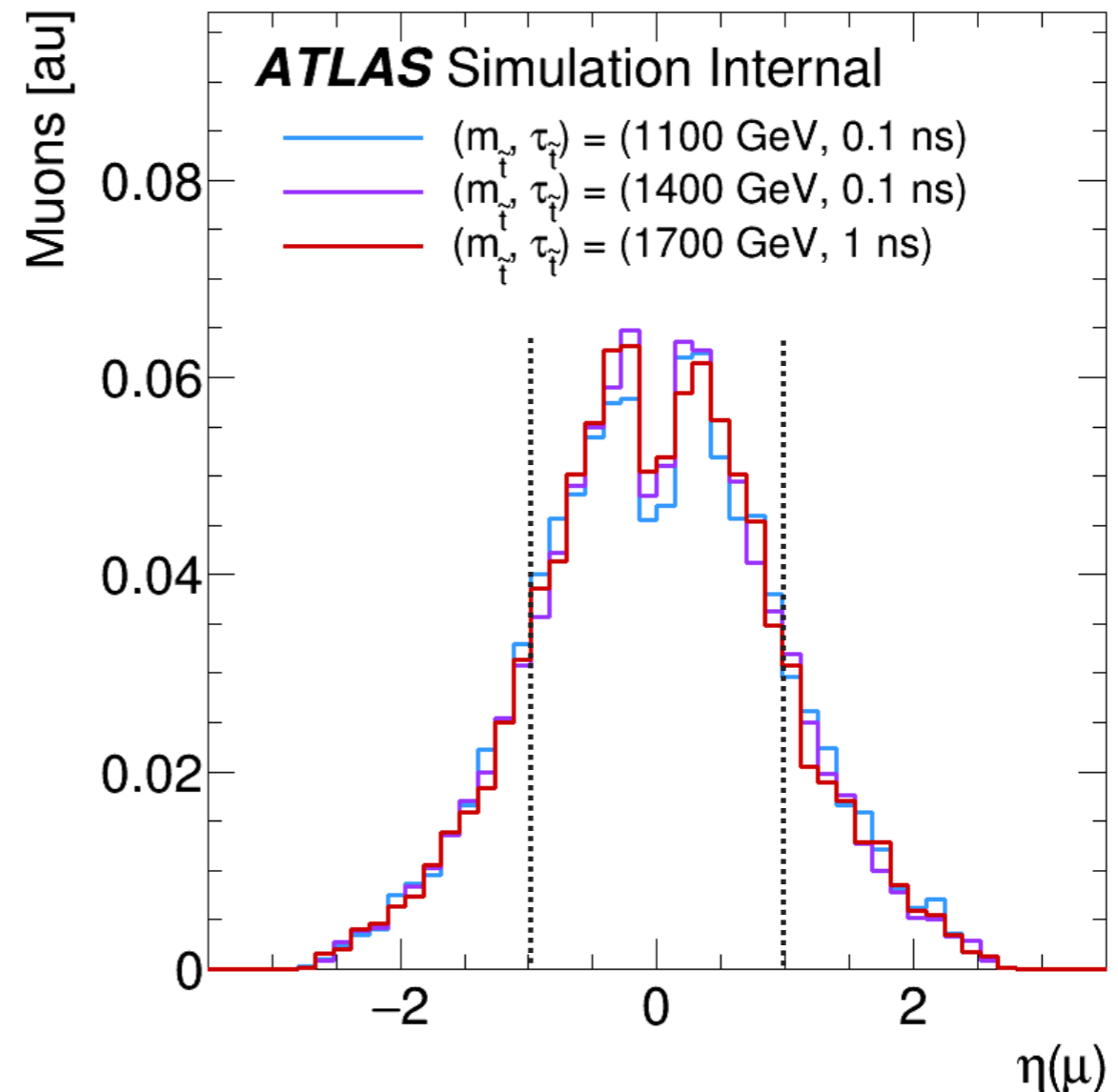
and $p_T > 62$ GeV

Trigger Efficiency

defined by

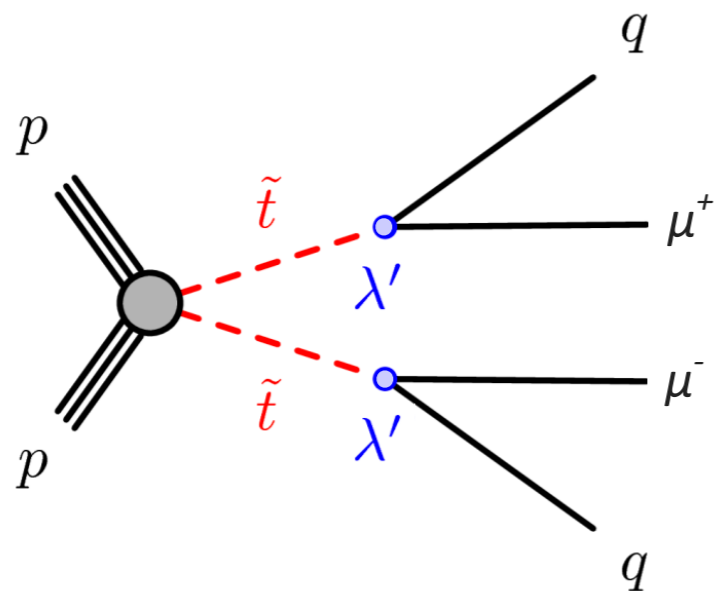
trigger detector coverage

Level 1 pointing assumptions

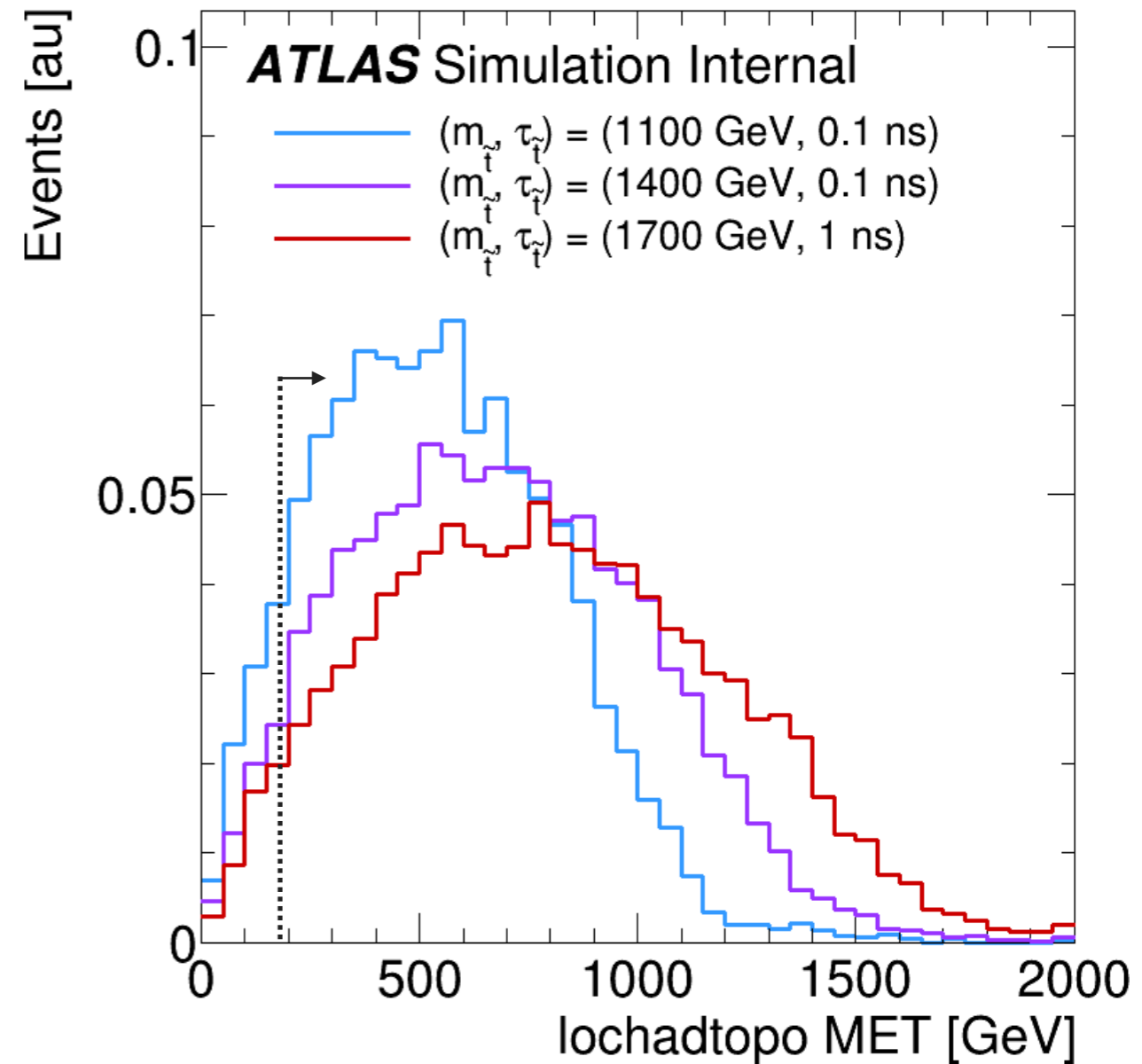


New idea: use a Missing Transverse Energy (MET) Trigger

muons ~invisible to calorimeter



Trigger is 100% efficient when calorimeter-based Missing Transverse Energy > 180 GeV



lochadtopo MET = cluster based calorimeter MET

Strategy in Run 2

MET Stream: MET Trigger if calorimeter MET > 180 GeV

Muon Stream: Muon Trigger if calorimeter MET < 180 GeV

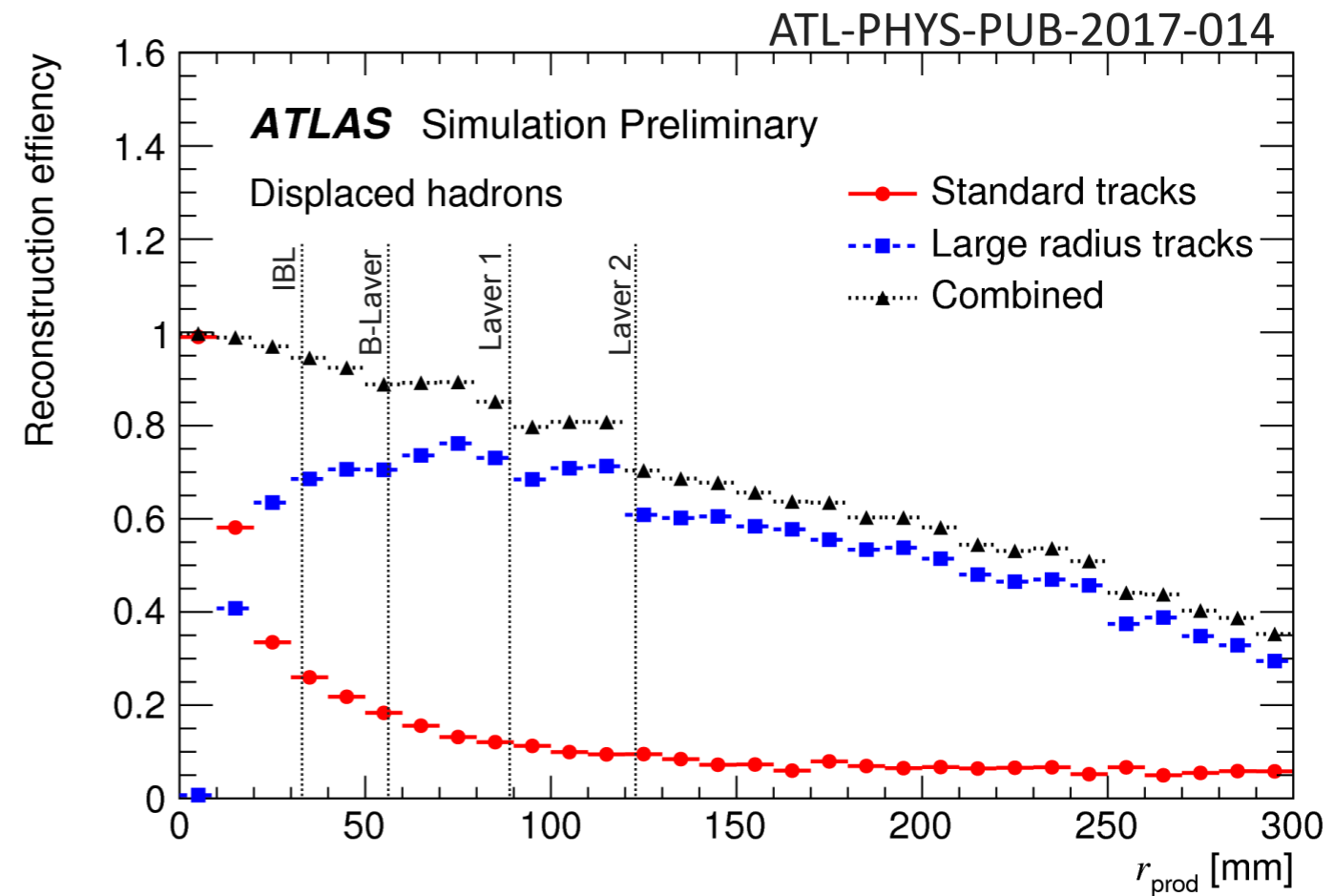
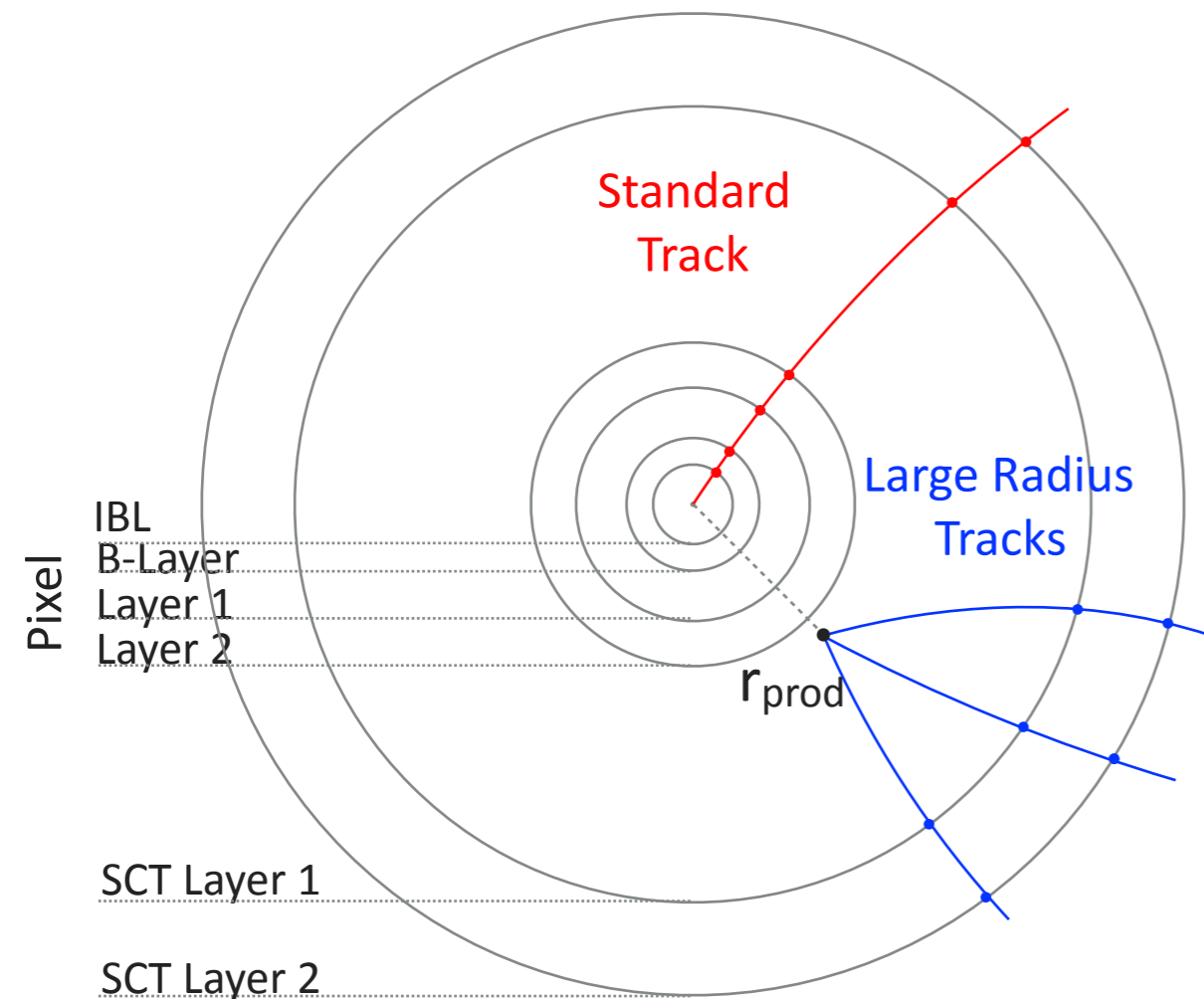
MET trigger most efficient for our model
keep Muon trigger for completeness

*new strategy improves overall signal acceptance x efficiency
for our benchmark model by $> 40\%$ with respect to Run 1*

Inner Detector Tracking

standard tracking requires $|d_0| < 10$ mm

large radius tracking is an additional pass of tracking with loosened impact parameter and hit requirements
 challenge for this analysis: fake tracks



Secondary Vertexing

forms vertices using tracks with
 $p_T > 1 \text{ GeV}$ and $|d_0| > 2 \text{ mm}$

$$R_{xy} < 300 \text{ mm}$$

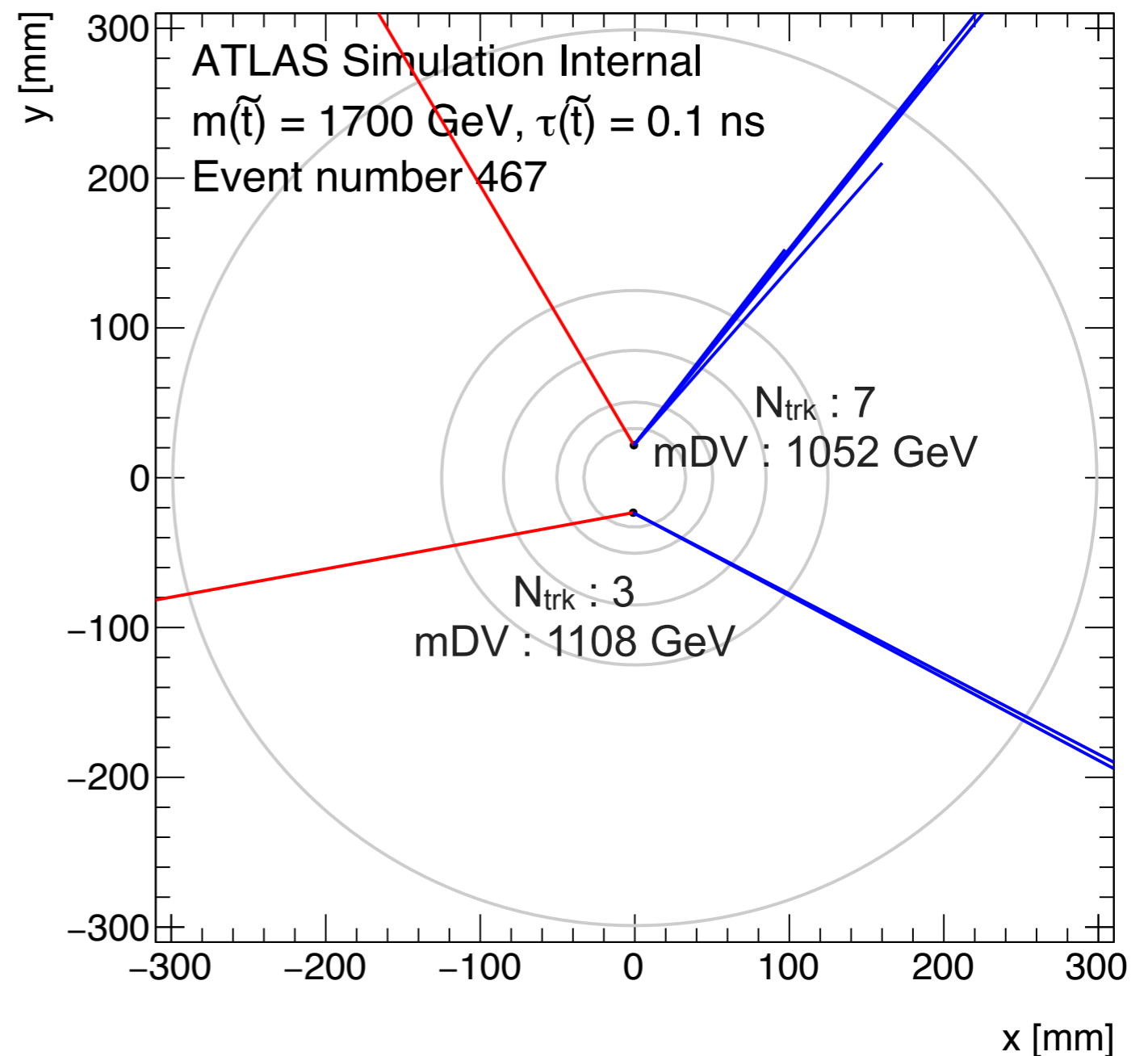
$$|Z| < 300 \text{ mm}$$

Signal Event Display

with **muons** and **tracks**
associated to displaced vertices

Challenge:

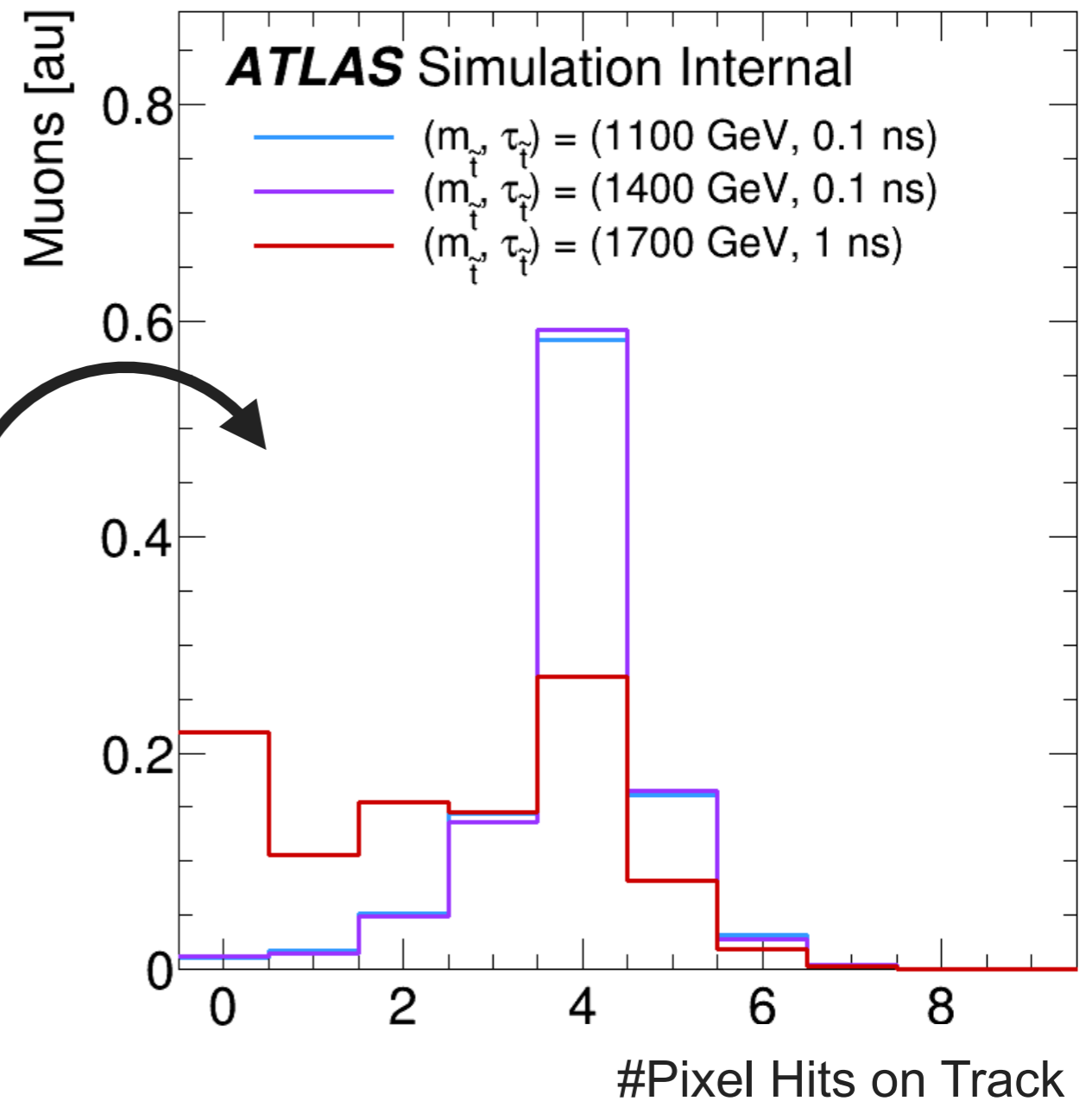
vertexing efficiency losses
at large radii



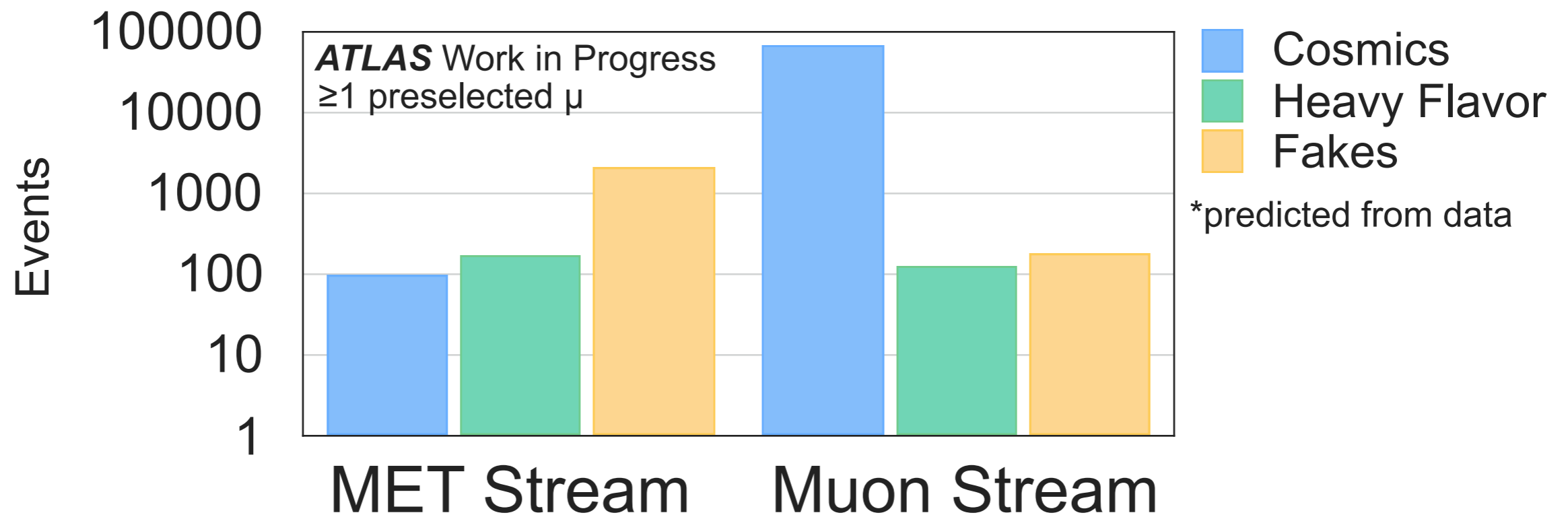
Preselected muons

combined muon, “medium” quality
with relaxed hit requirements

$|\eta| < 2.5$
 $p_T > 25 \text{ GeV}$
 $|d_0| > 2.0 \text{ mm}$



After muon preselection, events in data are dominated by algorithm fakes, muons from heavy flavor decays, cosmics

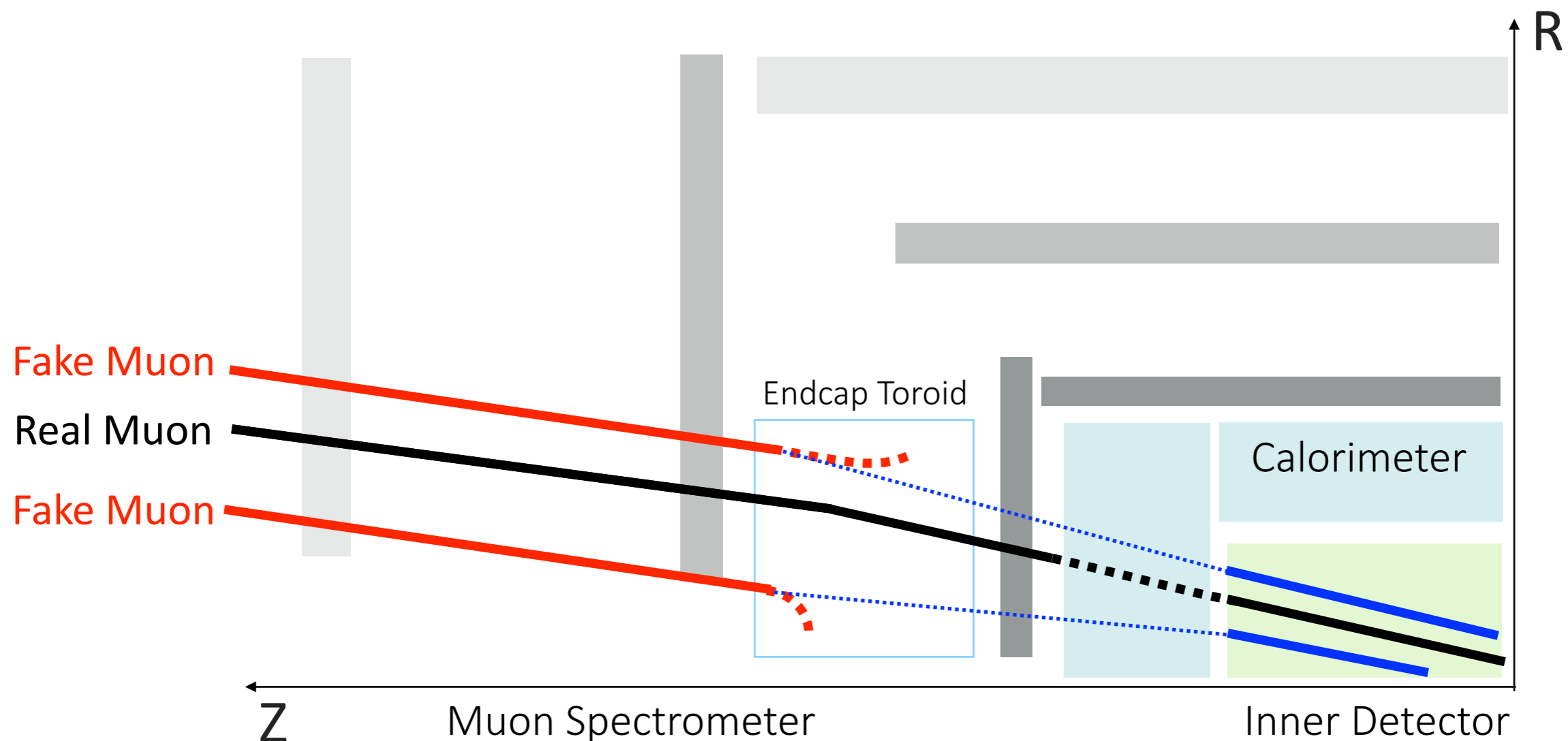


→ Design a dedicated veto for each background
with the idea that inverting each veto gives you
a control region pure in particular background

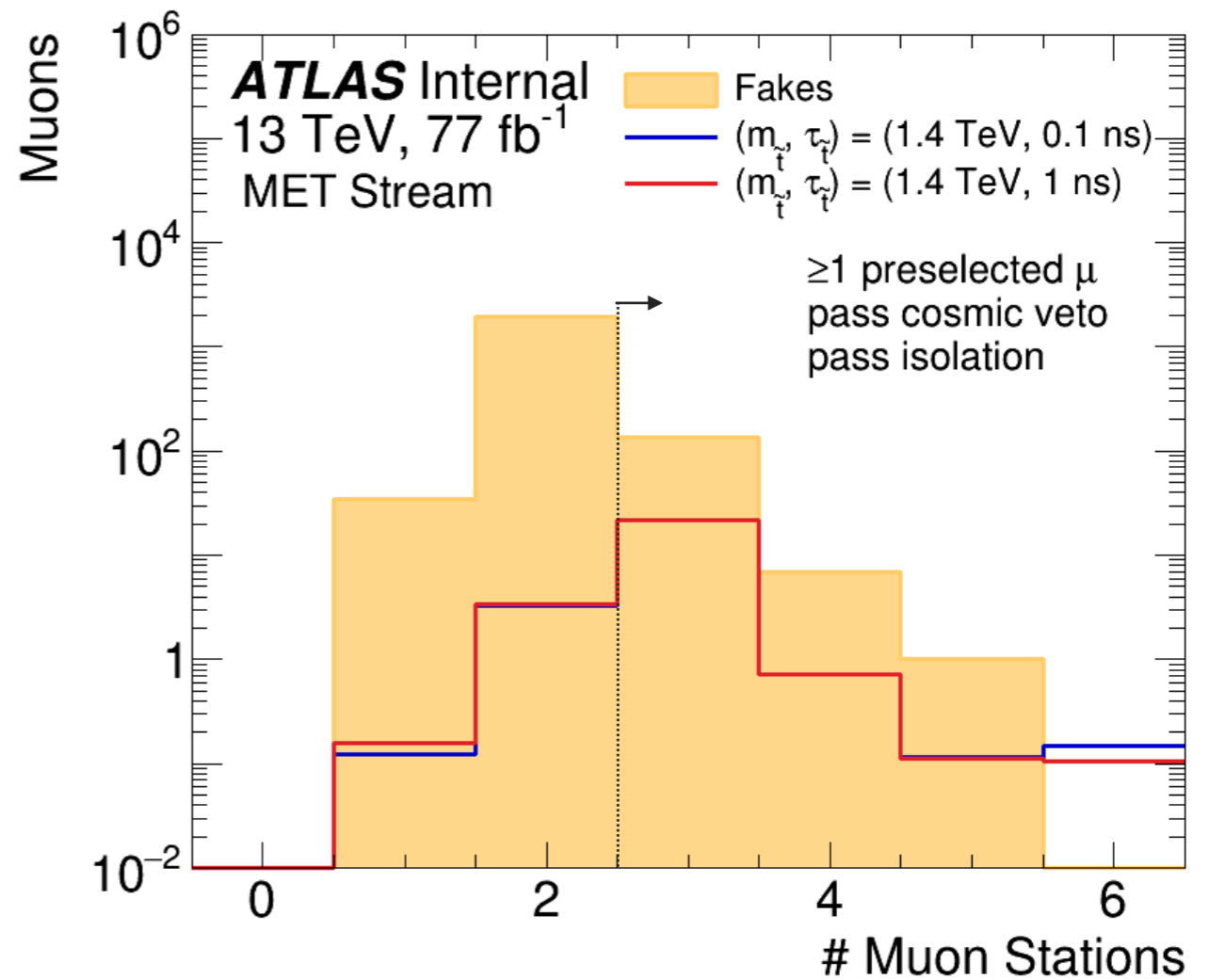
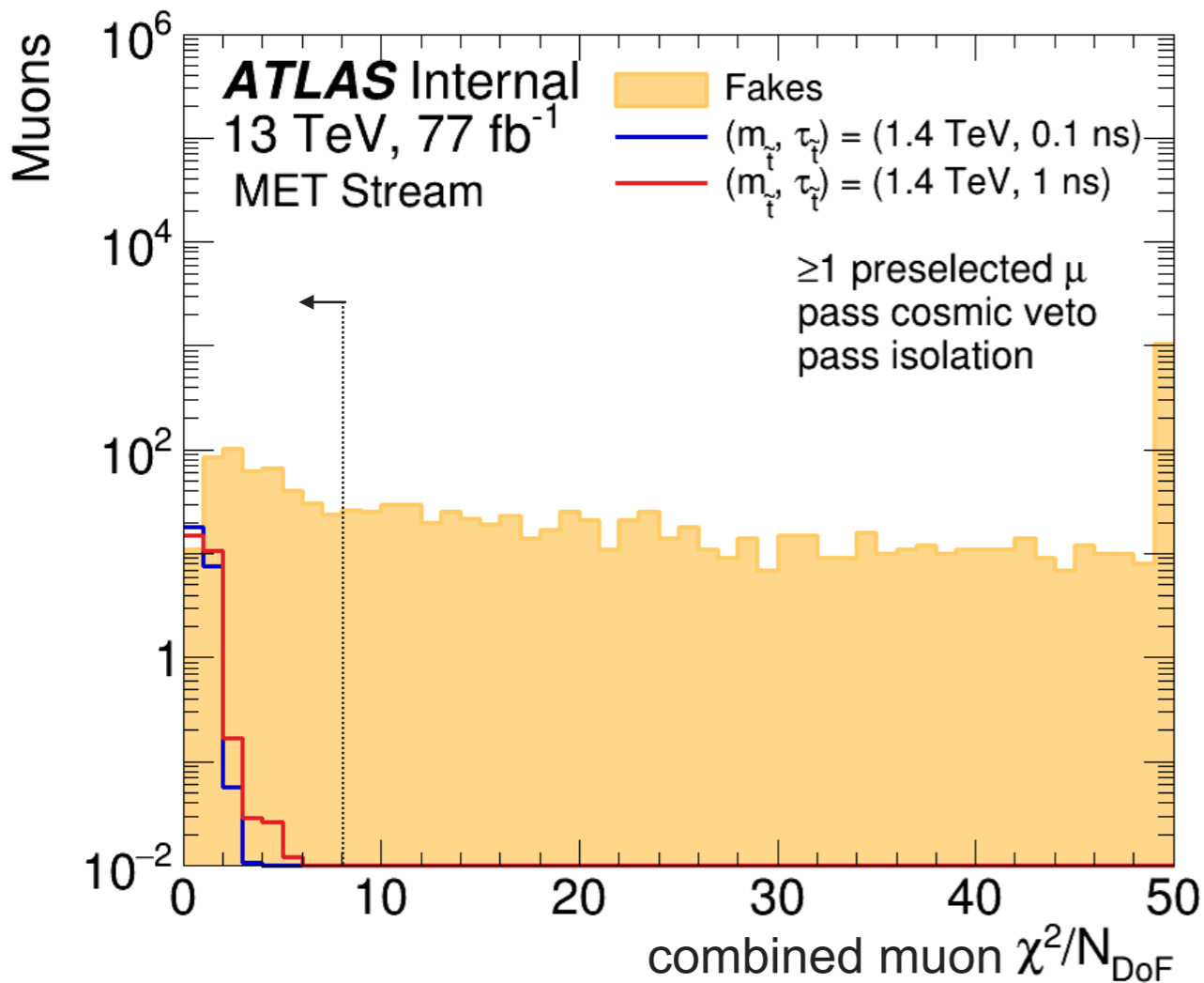
Largest background in MET triggered events

mostly in endcaps, a unique background to analyses using large radius tracking

fake **muon spectrometer track** matched to fake **inner detector track**



Can reject > 99% of fake muons by requiring muon has a segment in all 3 muon spectrometer stations and good quality of fit

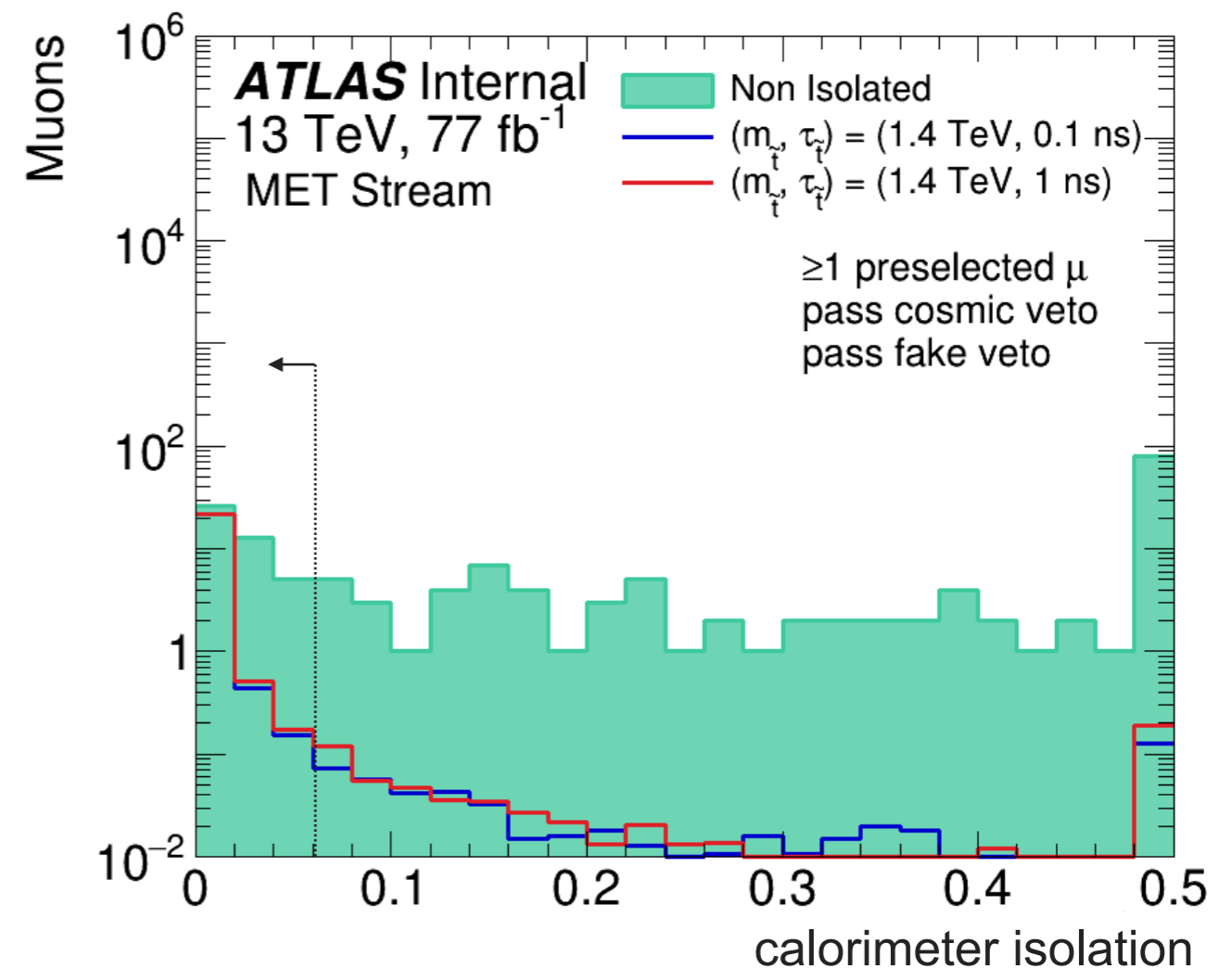
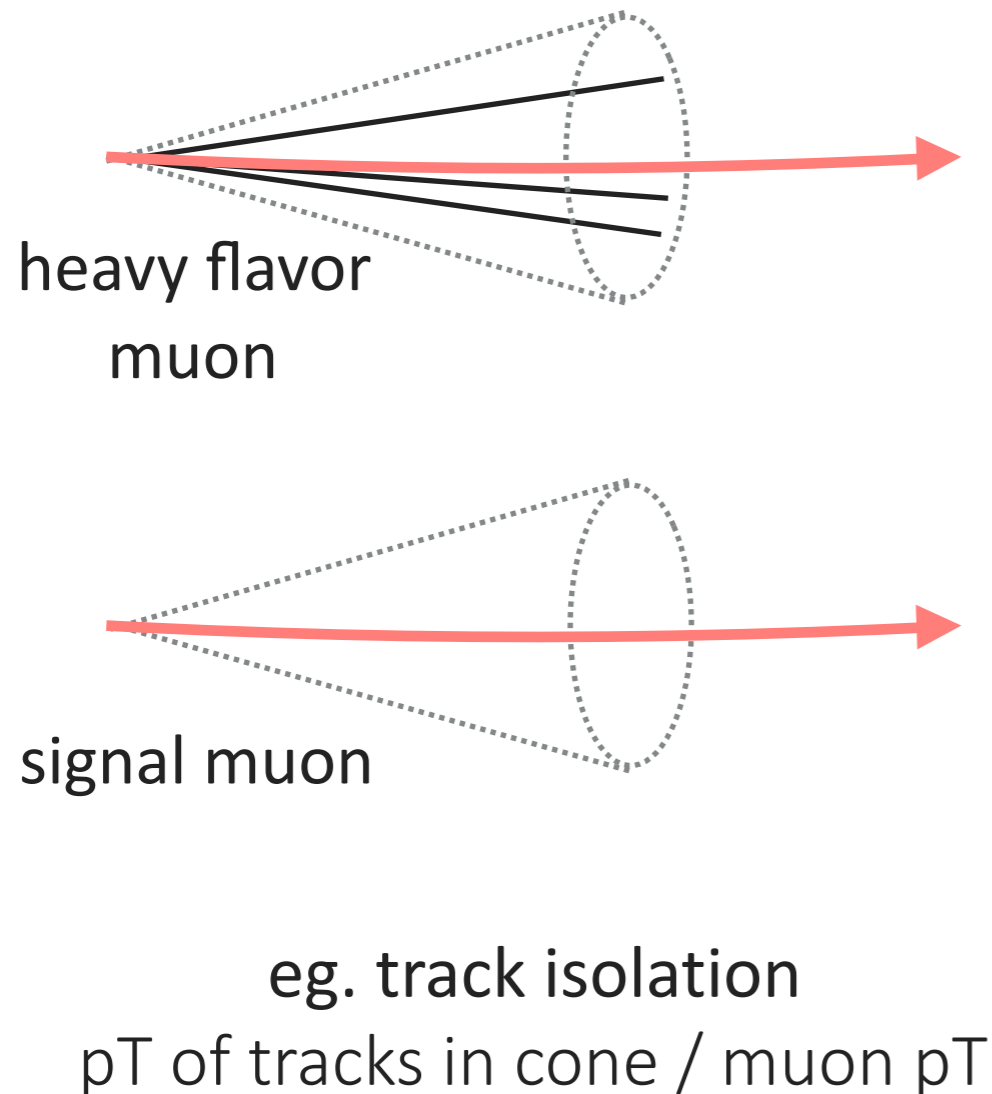


Muons from heavy flavor decays do pass $|d_0| > 2$ mm

these muons tend to be produced inside jets

by requiring muons pass track and calorimeter isolation requirements

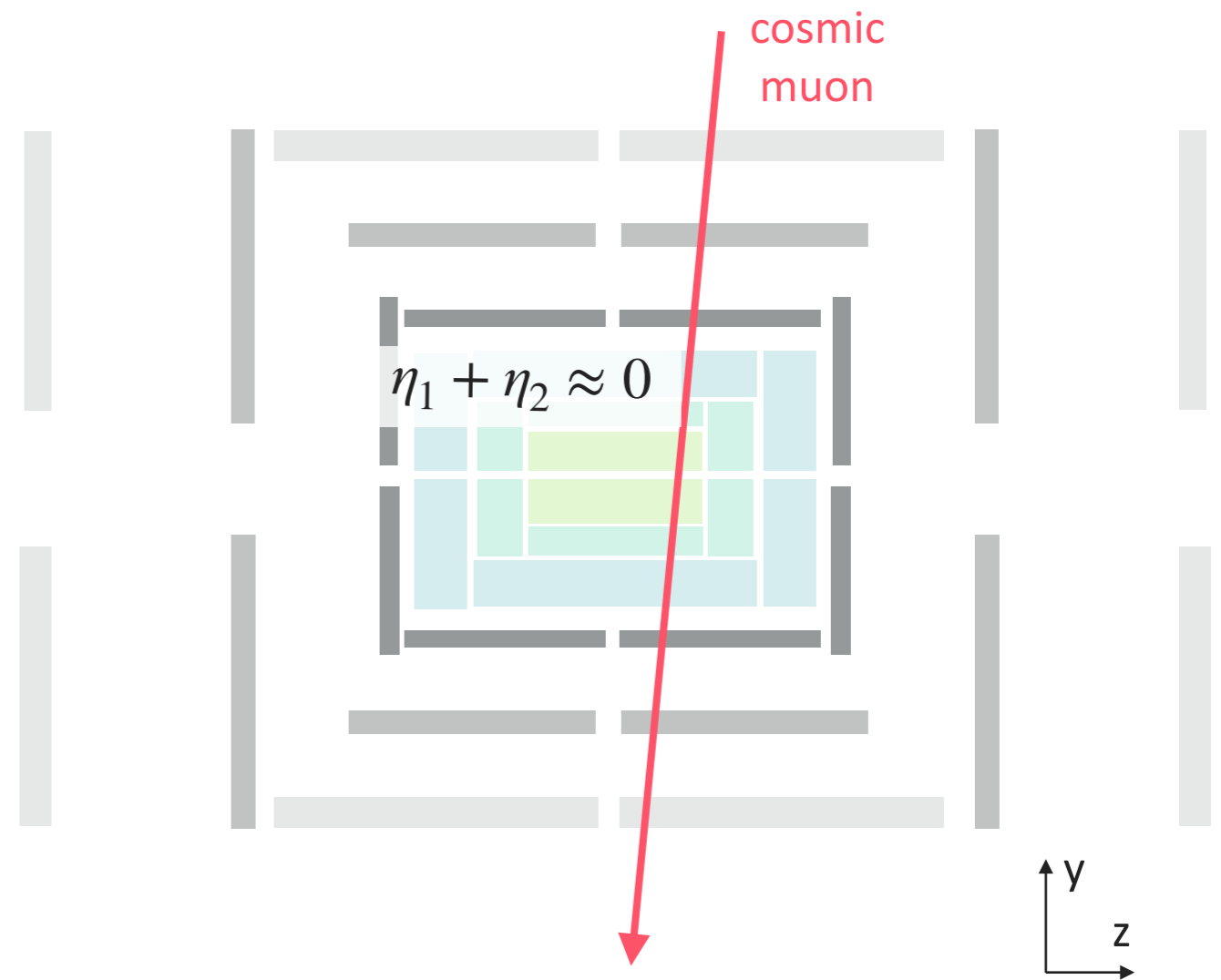
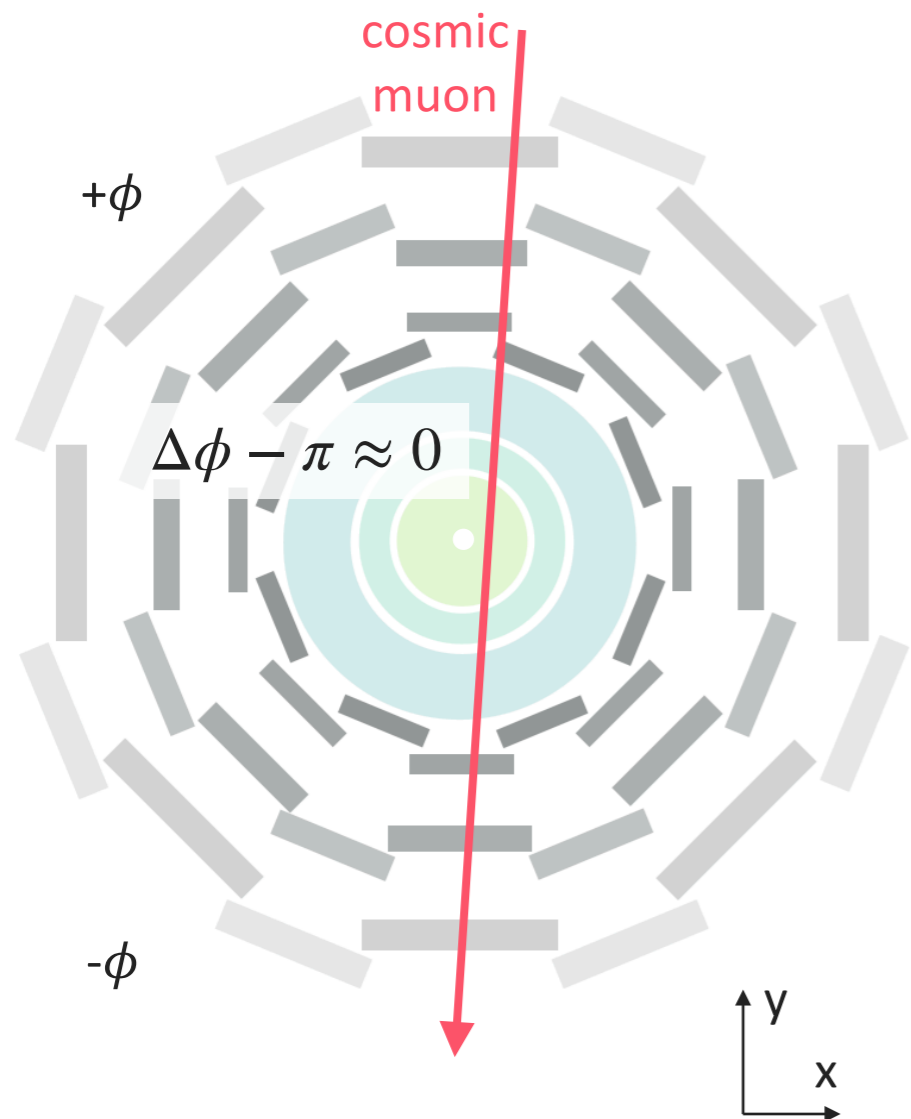
we can reject $\sim 85\%$ of heavy flavor muons and keep $\sim 99\%$ of signal muons



Dominant background in muon triggered events - difficult to reject

Original Idea: Run 1 analysis vetoed back-to-back muons

$$\Delta R_{cosmic} = \sqrt{(\eta_1 + \eta_2)^2 + (\Delta\phi - \pi)^2} \approx 0$$

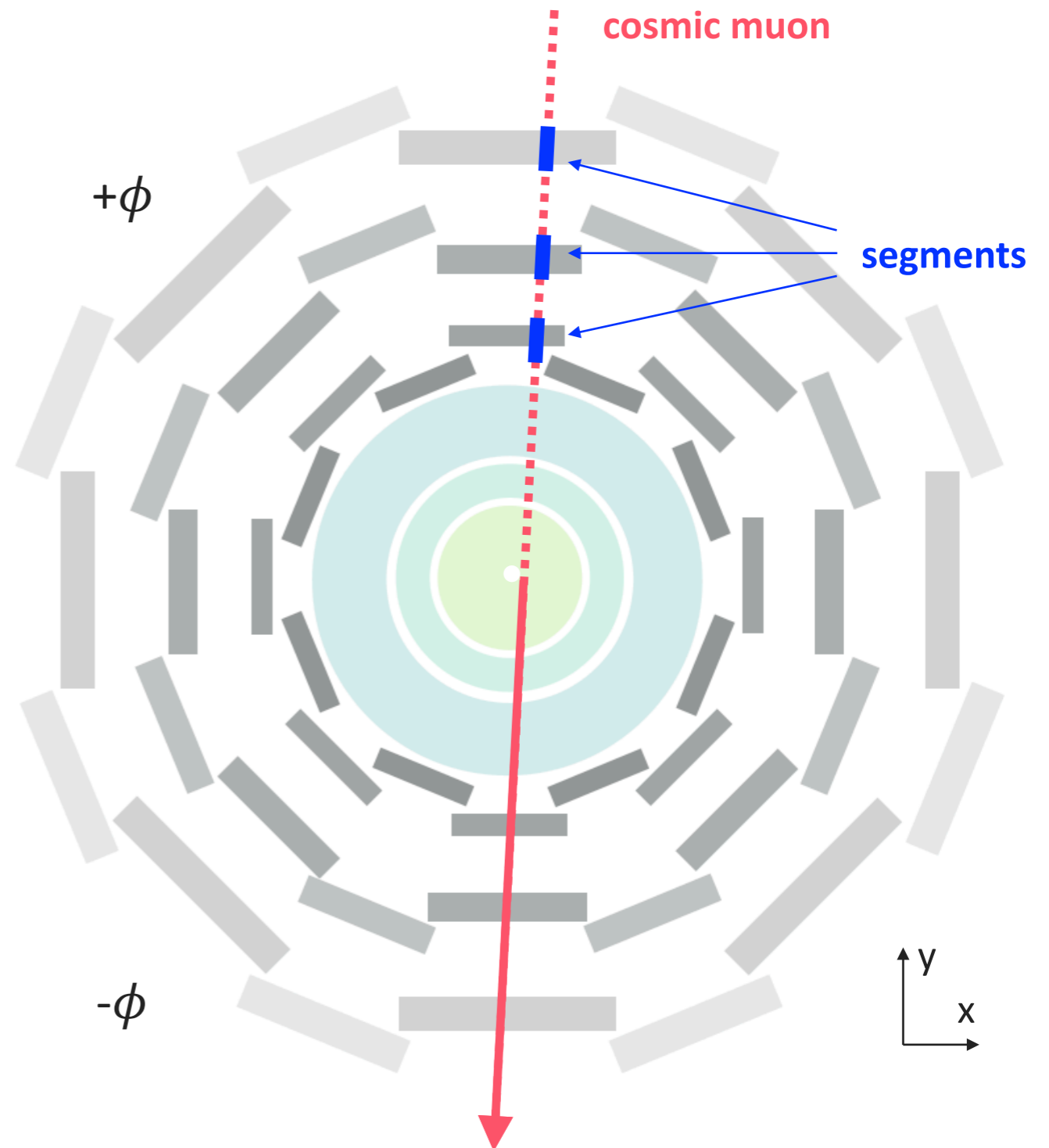


Challenge: Run 1 veto only
rejects 65% of cosmics!

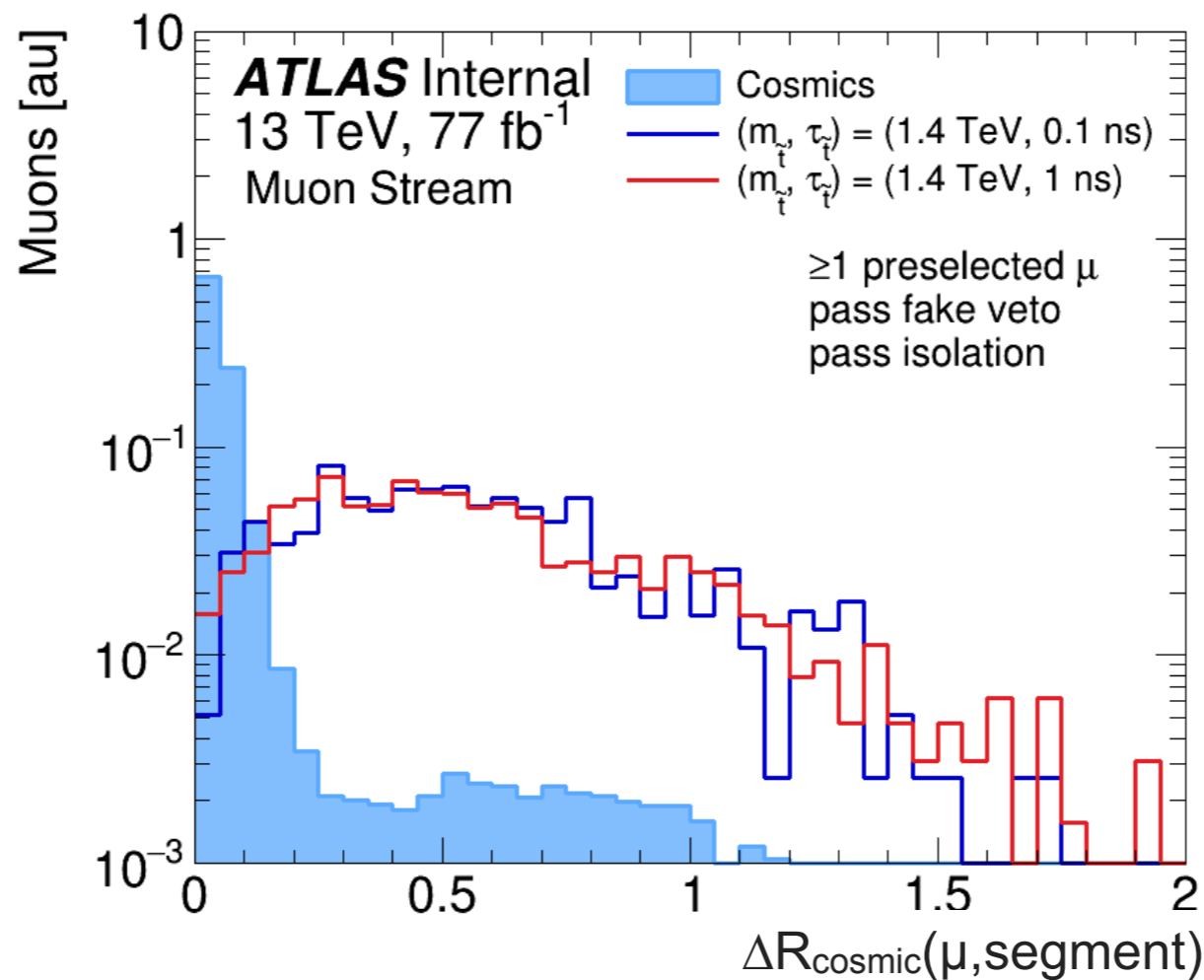
often don't reconstruct $+\phi$ muon,
because we're reconstructing
opposite actual direction of muon,
out of time with respect to collisions
by $\sim 30-70$ ns

But we do reconstruct muon
segments in $+\phi$!

timing difference for hits
within a muon chamber
 $\sim \mathcal{O}(1$ ns)



New idea: veto muons if back-to-back with a segment
rejects 95% of cosmics and keep 95% of signal

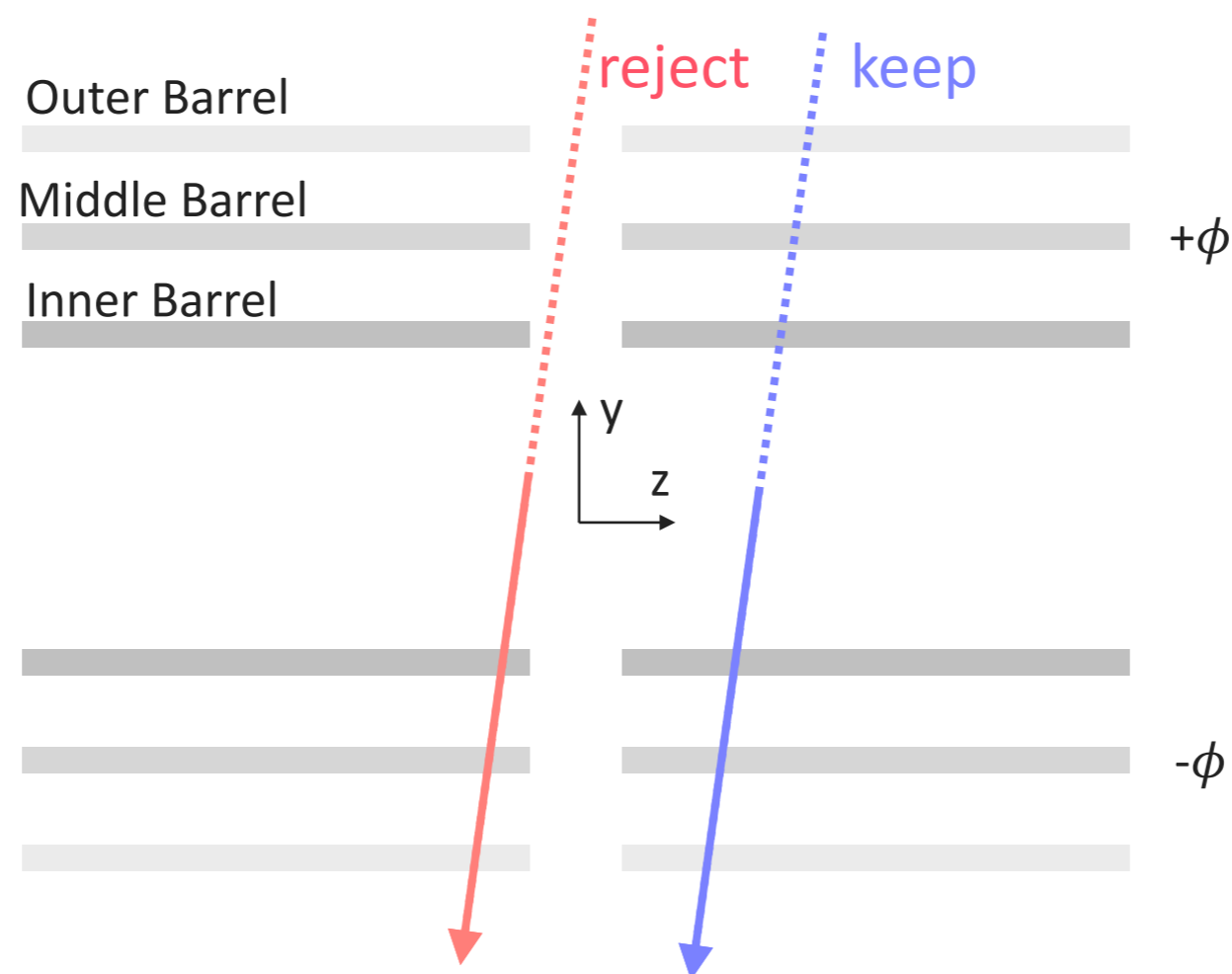


*also account for differences in segment eta-phi resolution (drift tubes point in phi)
and make a geometric correction for muon displacement

But we're not done yet!

cosmics which pass the segment veto form hot spots near $\eta \sim 0$

→ additionally require muons to point backwards in η - ϕ to regions with muon detector coverage



full cosmic veto rejects $>99\%$ of cosmics and keeps 95% of signal
allows us to loosen our displaced vertex selection with respect to run 1

Preselection

Fiducial Volume

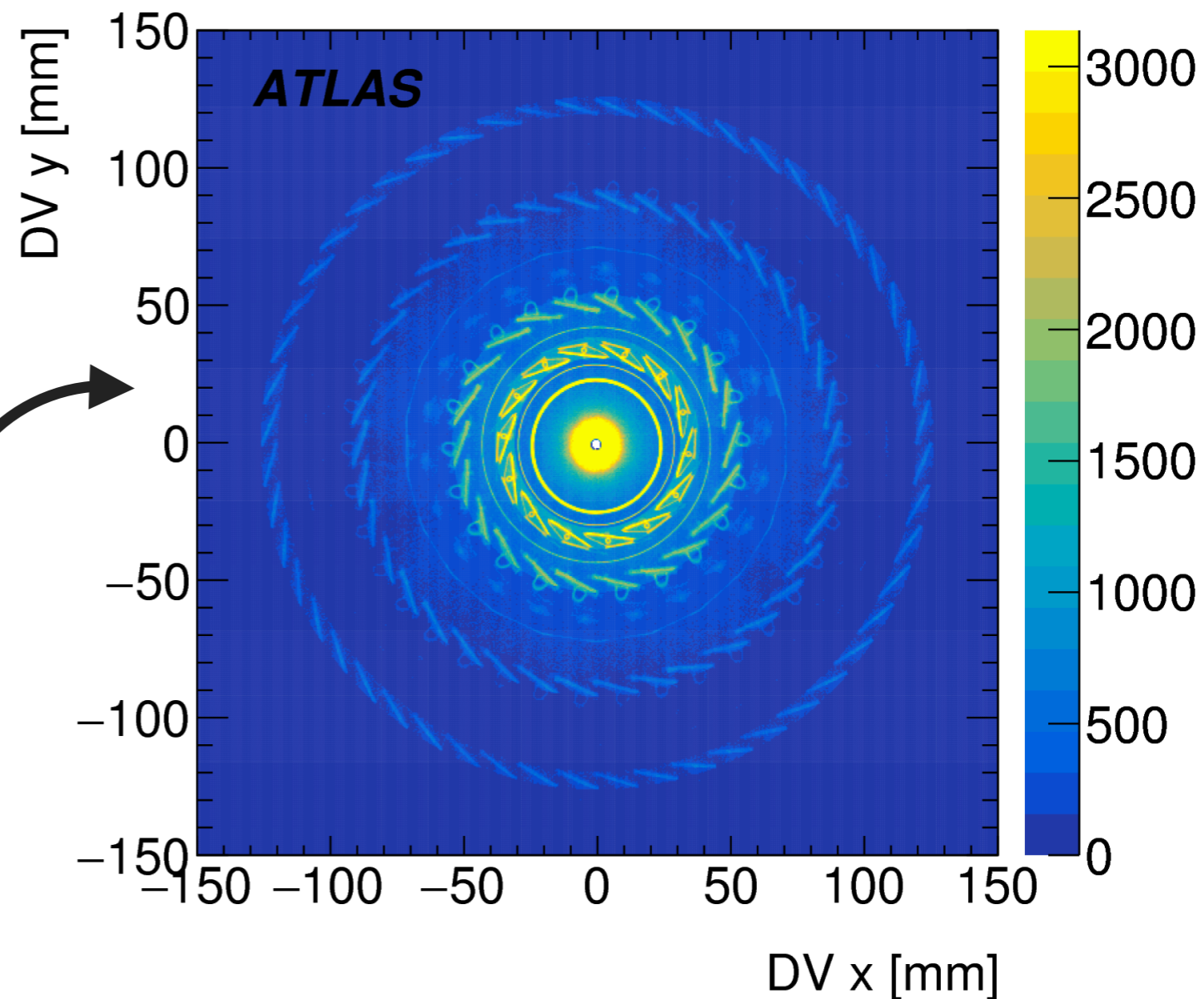
$$R_{xy} < 300 \text{ mm} \quad |z| < 300 \text{ mm}$$

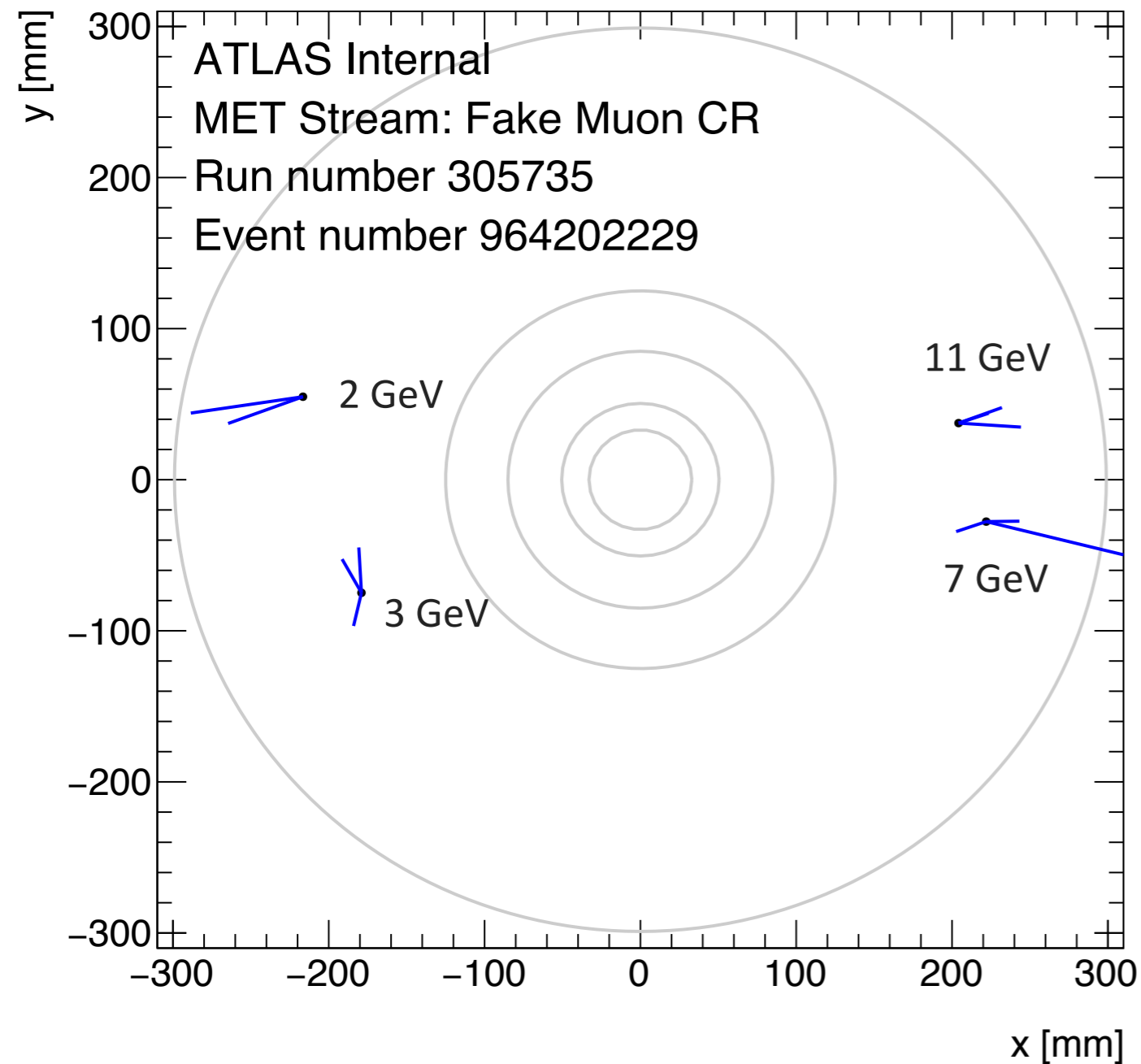
Distance from displaced vertex to any primary vertex, in transverse plane $> 4 \text{ mm}$

$$\text{Quality: } \chi^2/N_{\text{DoF}} < 5$$

Material Veto:
reject vertices with (r, z, ϕ)
in regions of material

$\sqrt{s}=13 \text{ TeV}$, $L=32.8 \text{ fb}^{-1}$, All Reconstructed Vertices





After preselection most DVs
from random track crossings
tracks with low pT form low mass and
low track multiplicity vertices

Final Selection

displaced vertices must have

1. at least 3 tracks
- &
2. mass > 15 GeV

~10%-20% improvement in signal
efficiency with respect to Run 1
selection

Selection	MET Stream	Muon Stream
Cosmic Muons	<0.008	$1.02 \pm 0.06 \pm 0.07$
Heavy Flavor	$0.14 \pm 0.14 \pm 0.04$	$0.28 \pm 0.20 \pm 0.07$
Fake Muons	$0.12 \pm 0.02 \pm 0.12$	$0.01 \pm 0.01 \pm 0.01$
Total Expected Background	$0.26 \pm 0.14 \pm 0.13$	$1.31 \pm 0.21 \pm 0.10$
Expected Signal (1.4 TeV, 0.1 ns)	12.3 ± 0.06	0.45 ± 0.13
Expected Signal (1.4 TeV, 1 ns)	4.7 ± 0.41	0.21 ± 0.08

**work in progress, especially systematic uncertainties & MC statistics*

Reminder: MET stream drives sensitivity to our benchmark model but we keep the Muon Stream to retain sensitivity to other signals

Signal

15-30% Cross Section

~10% Tracking + Vertexing Efficiency

~10% Pileup Reweighting

~2% Muon Trigger

~2% Muon Identification

~2% Luminosity

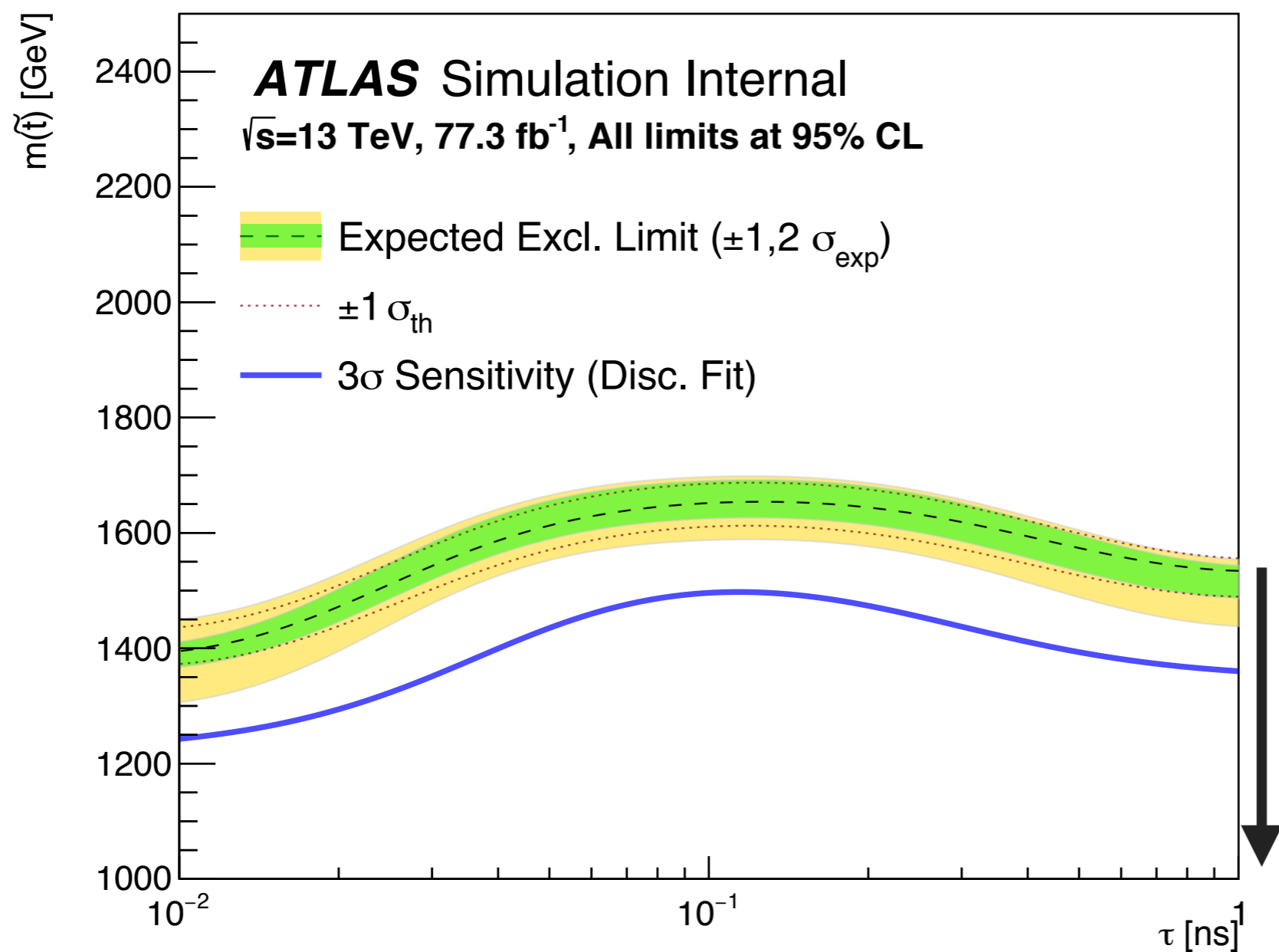
Backgrounds

~50% uncertainty on total background in MET Stream

Heavy Flavor: statistics in control regions

Fakes: statistics, and possible correlation between vertex properties & fake muons

Expected sensitivity for different stop masses and lifetimes with 2016-2017 data

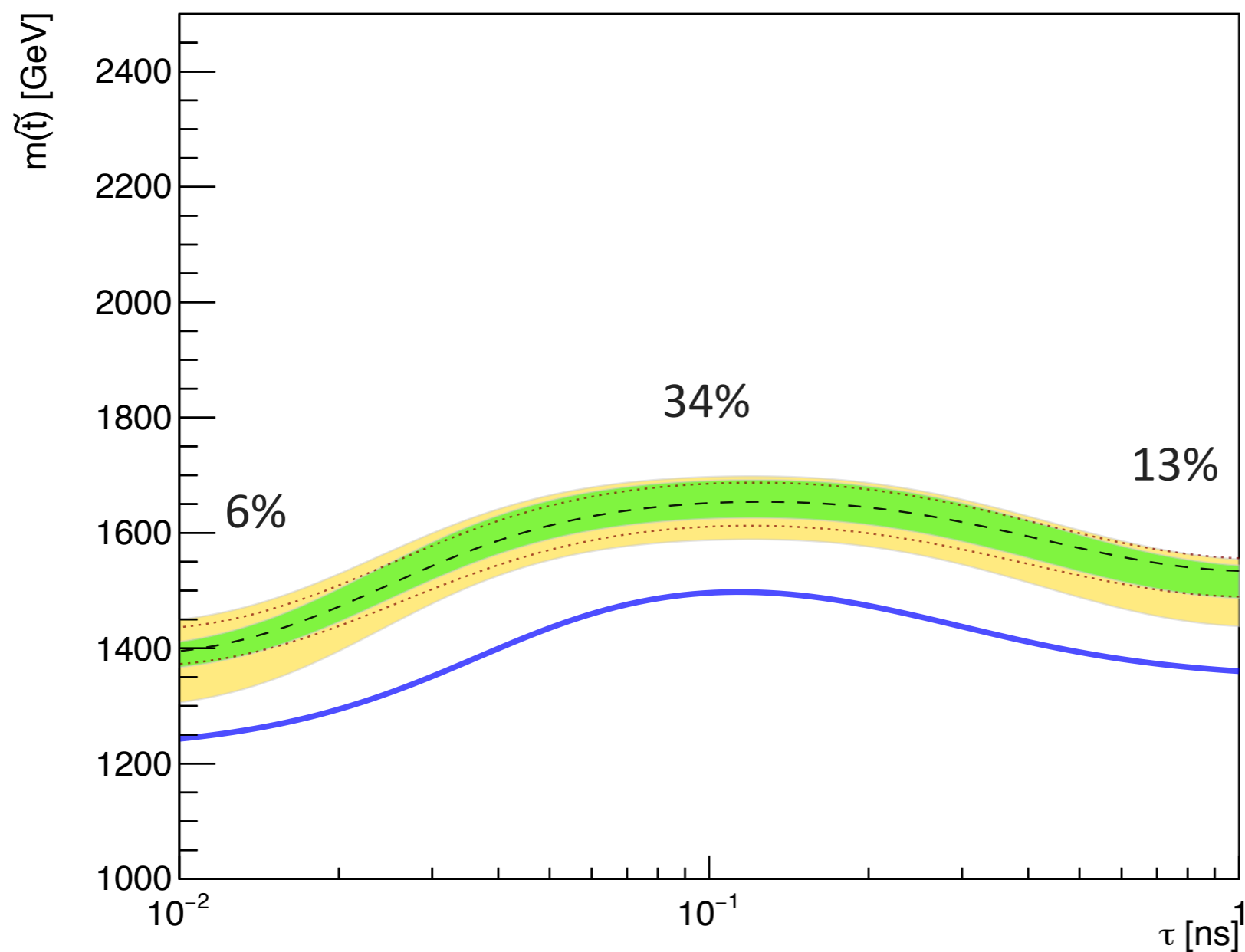


For reference:
 Converting # selected signal events to sensitivity lines
 3 events \rightarrow exclusion
 5 events $\rightarrow 3\sigma$

Increasing XS
 can exclude/discover masses
 and lifetimes below curve

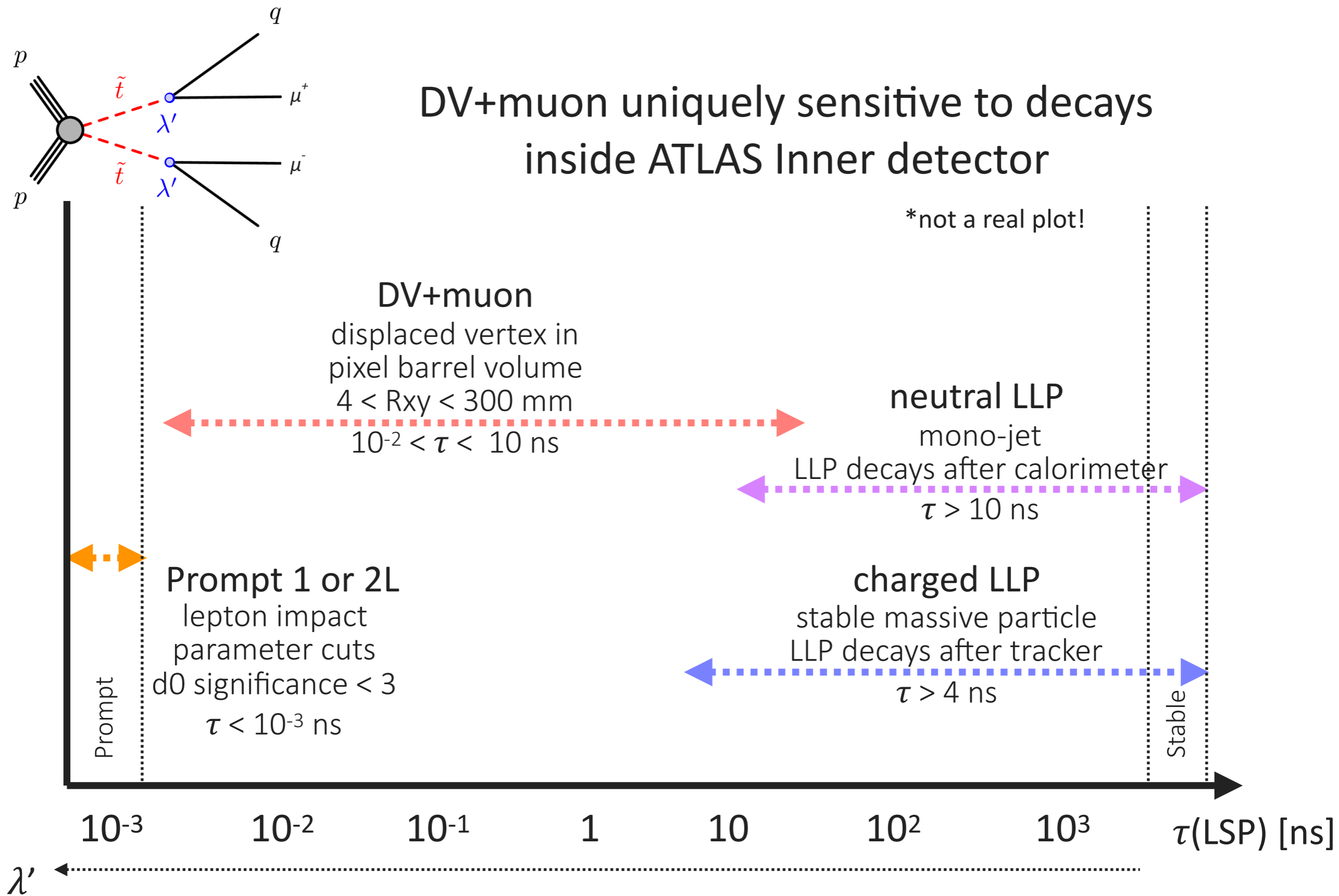
Sensitivity Shaped by Signal Acceptance x Efficiency

acceptance primarily related to signal lifetime



small lifetimes
 $|d_0| > 2$ mm
DV-PV distance

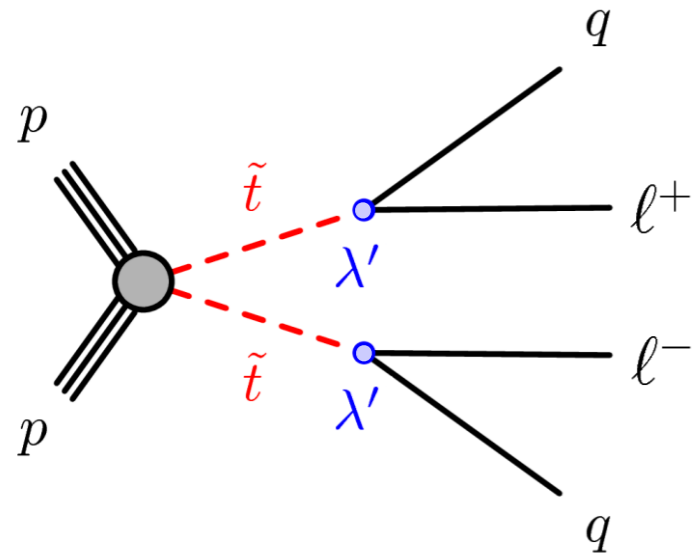
large lifetimes
Fiducial volume
 $R_{xy} < 300$ mm
and reduced
vertexing efficiency



Improvements to current analysis

try to loosen $|d_0|$ and distance(DV-PV) requirements
 reduce large radius tracking fake rate at higher pileup
 and improve vertexing efficiency at large radii
 Run 3: planned improvements to muon trigger

	Run 1	Run 2	Run 3	HL-LHC
	2010-2012	2015-2018	2021-2023	2026-
\sqrt{s}	7-8 TeV	13 TeV	14 TeV	14 TeV
Integrated Luminosity	26 fb ⁻¹	140 fb ⁻¹	150 fb ⁻¹	3000 fb ⁻¹



Other interesting long-lived particle searches

direct detection of charged LLPs using pixel dE/dx

targets slightly longer lifetimes, $\tau > 4$ ns

interesting in Run 3 because of new trigger opportunities

	Run 1 2010-2012	Run 2 2015-2018	★ Run 3 2021-2023	HL-LHC 2026-
\sqrt{s}	7-8 TeV	13 TeV	14 TeV	14 TeV
Integrated Luminosity	26 fb ⁻¹	140 fb ⁻¹	150 fb ⁻¹	3000 fb ⁻¹

ATLAS Muon Spectrometer

Challenges at high instantaneous luminosity

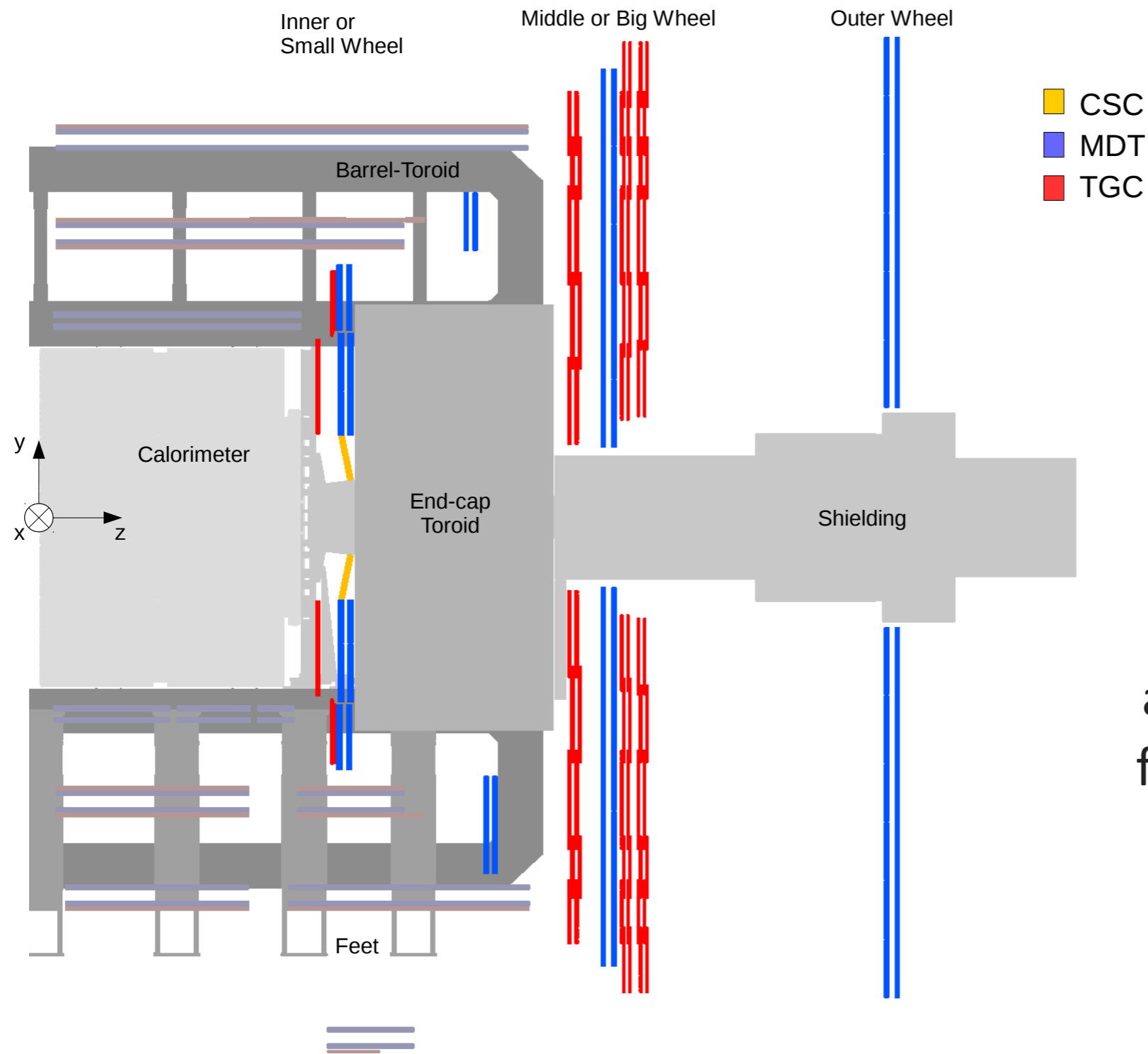
Looking forward to physics at High Luminosity LHC

	Run 1 2010-2012	Run 2 2015-2018	Run 3 2021-2023	HL-LHC 2026-
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Looking forward to physics at High Luminosity LHC

	Run 1 2010-2012	Run 2 2015-2018	Run 3 2021-2023	HL-LHC 2026-
\sqrt{s}	7-8 TeV	13 TeV	14 TeV	14 TeV
Integrated Luminosity	26 fb ⁻¹	140 fb ⁻¹	150 fb ⁻¹	3000 fb ⁻¹
Max. Inst. Luminosity	$8 \cdot 10^{33} \text{ cm}^{-2}\text{s}^{-1}$	$2 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$	$2 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$	$7 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$

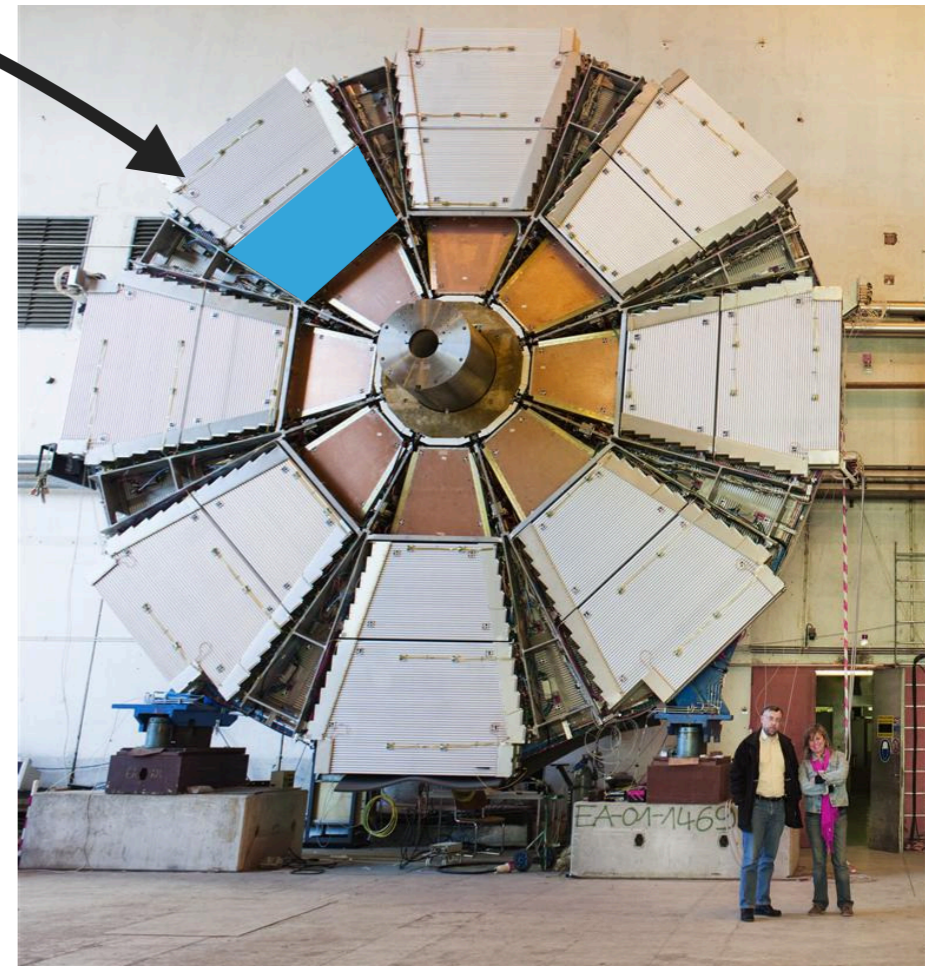
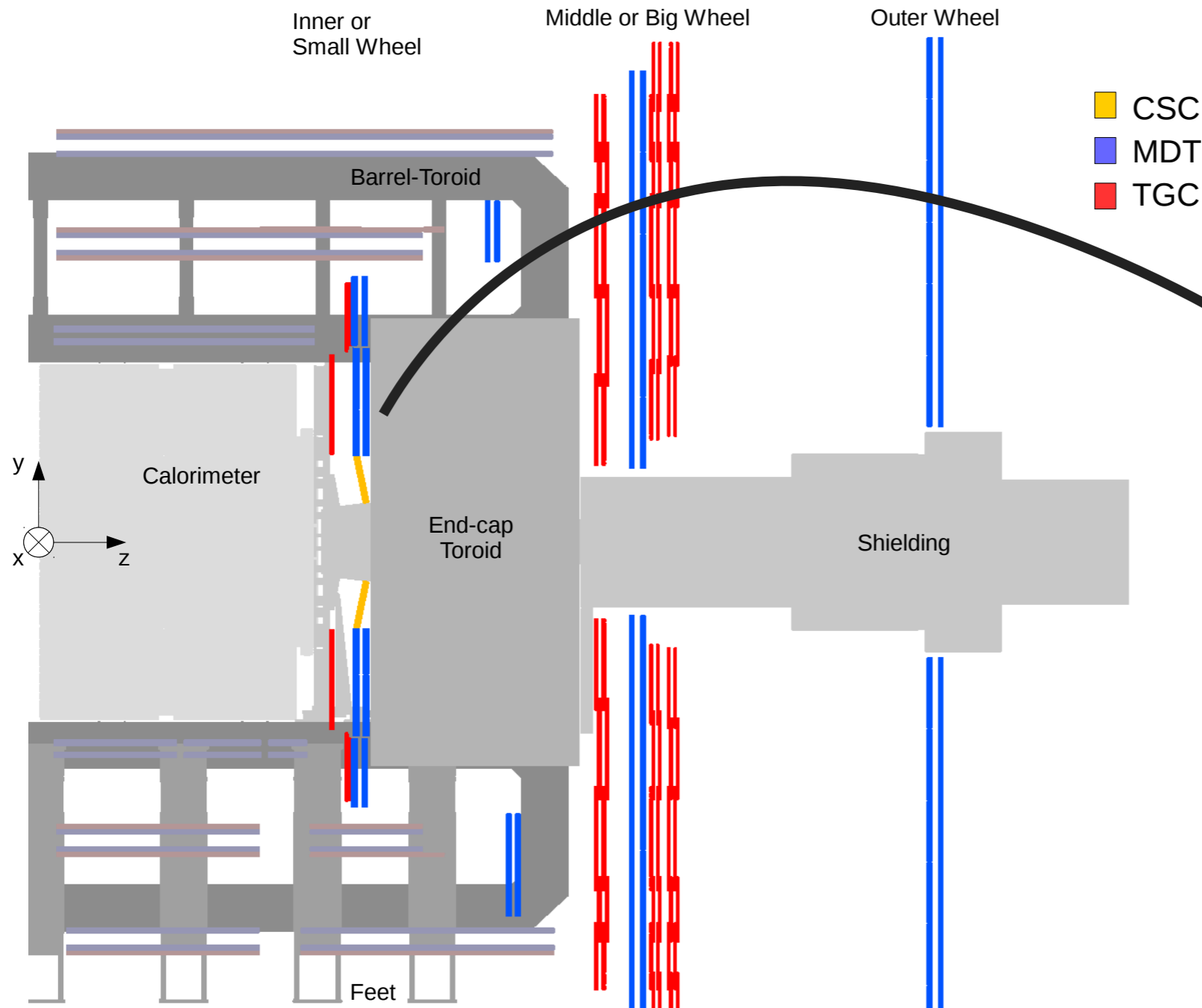
LHC instantaneous luminosity will be x7 higher than design values
need to make sure ATLAS maintains excellent performance



Cathode Strip Chambers
and **Monitored Drift Tubes**
for precision muon tracking

Thin Gap Chambers for
muon triggering

Reminder: Muon Endcap Geometry

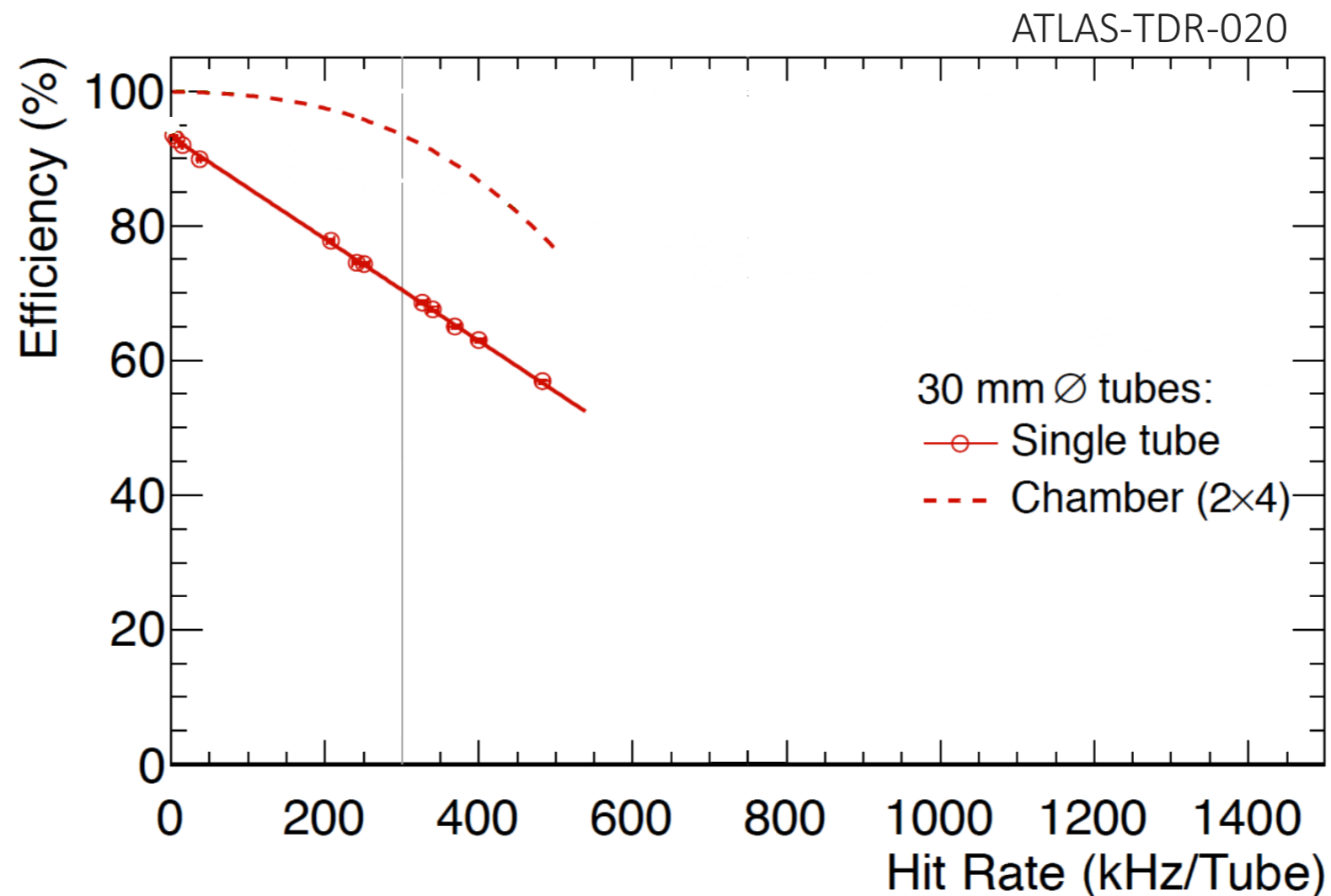


Focus on Small Wheel, and ask the question
Can the Monitored Drift Tubes survive at High Luminosity LHC?

MDTs have a maximum drift time of 750 ns → dead time

hit efficiency v. hit rate, measured using test beam data

hit rates of 500 Hz/cm² or 300 kHz/tube result in 35% single tube efficiency loss
any higher results in drastic losses to muon resolution and efficiency



In ATLAS hit rate depends on

luminosity, # bunches

position in Rxy, Z

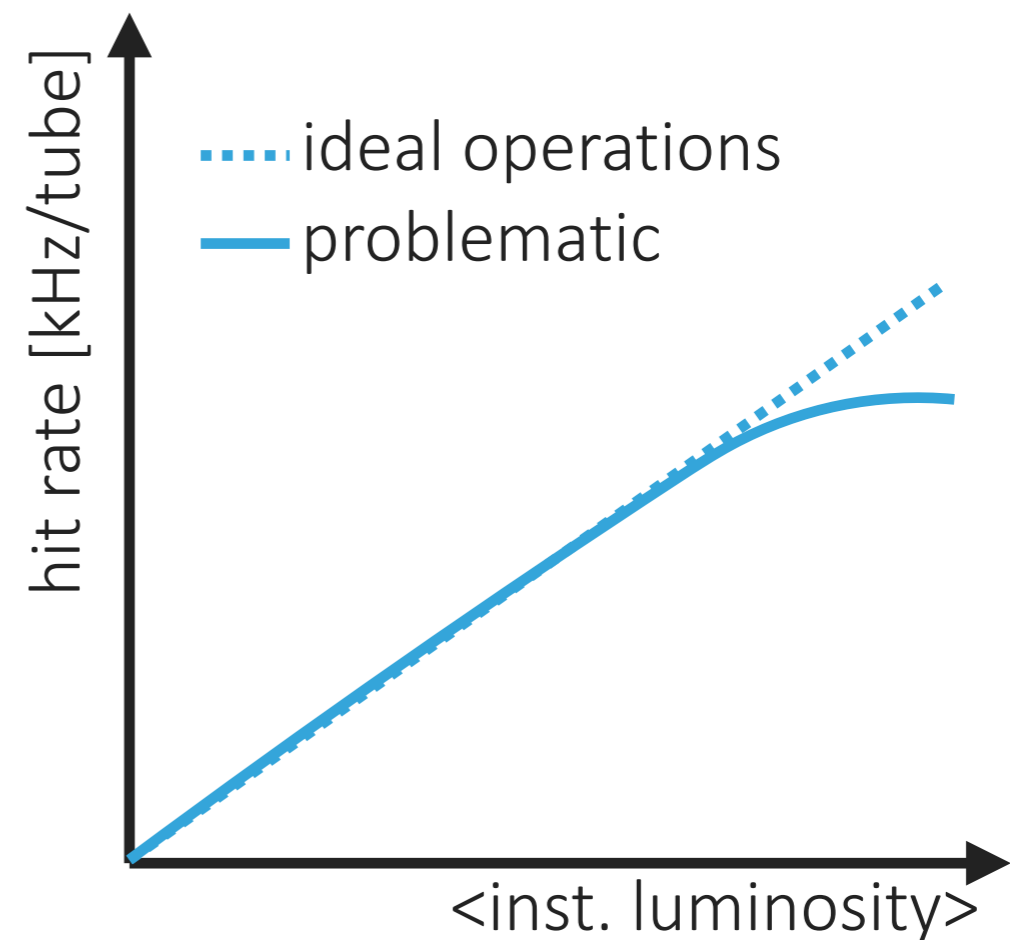
geometry, shielding

Most interested in
hit rates v. luminosity

Why we make this plot

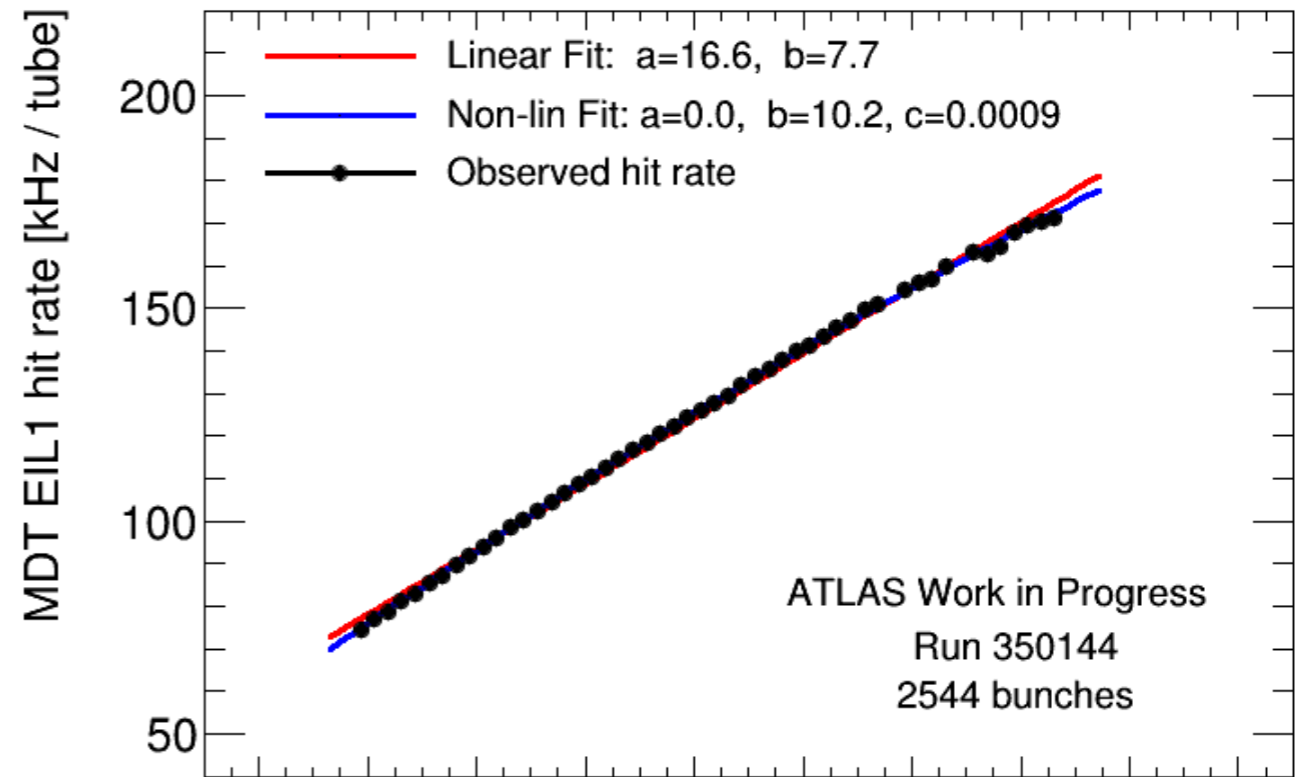
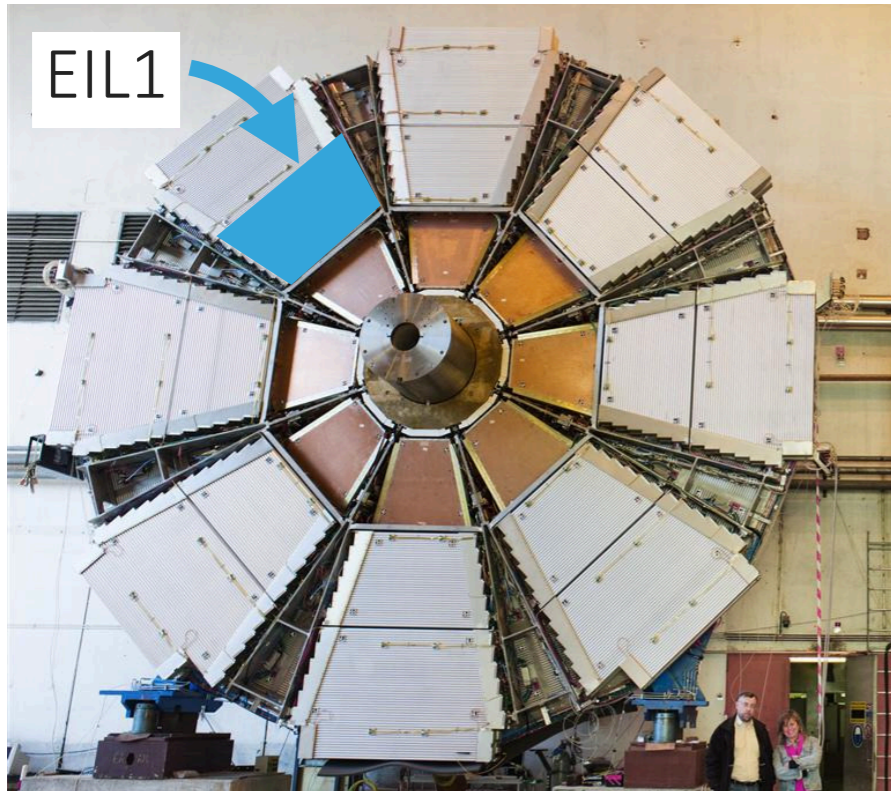
1. linear relationship is indicative of good operations
2. extrapolate to higher luminosity, will our detectors work at HL-LHC?

$$\text{hit rate} = \frac{\text{\#hits/event}}{\text{lifetime}}$$



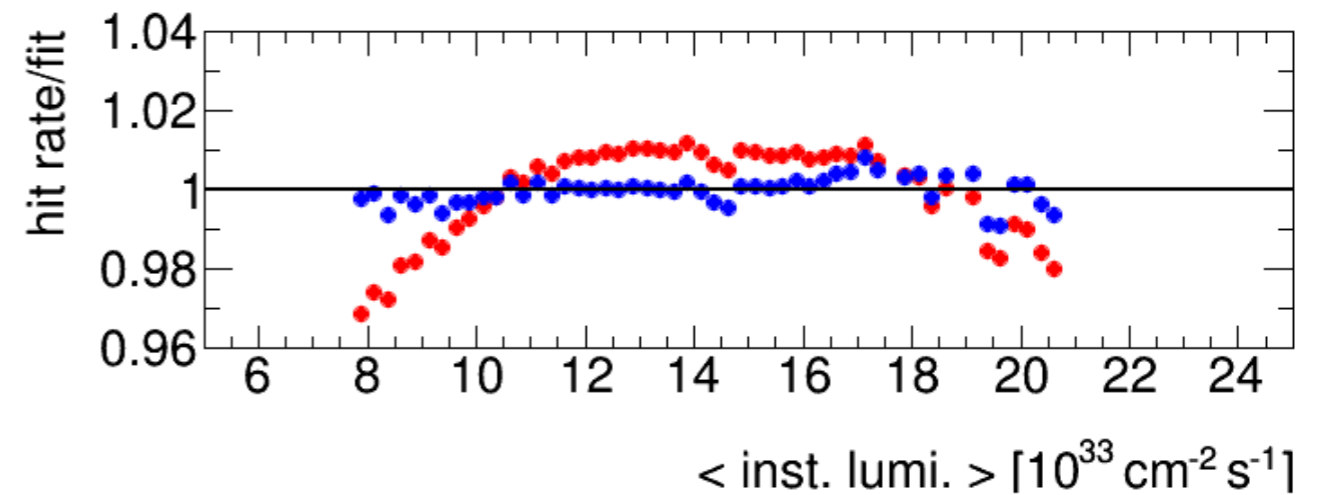
*not a real plot!

2018: non-linear hit rates observed in Inner Endcap!

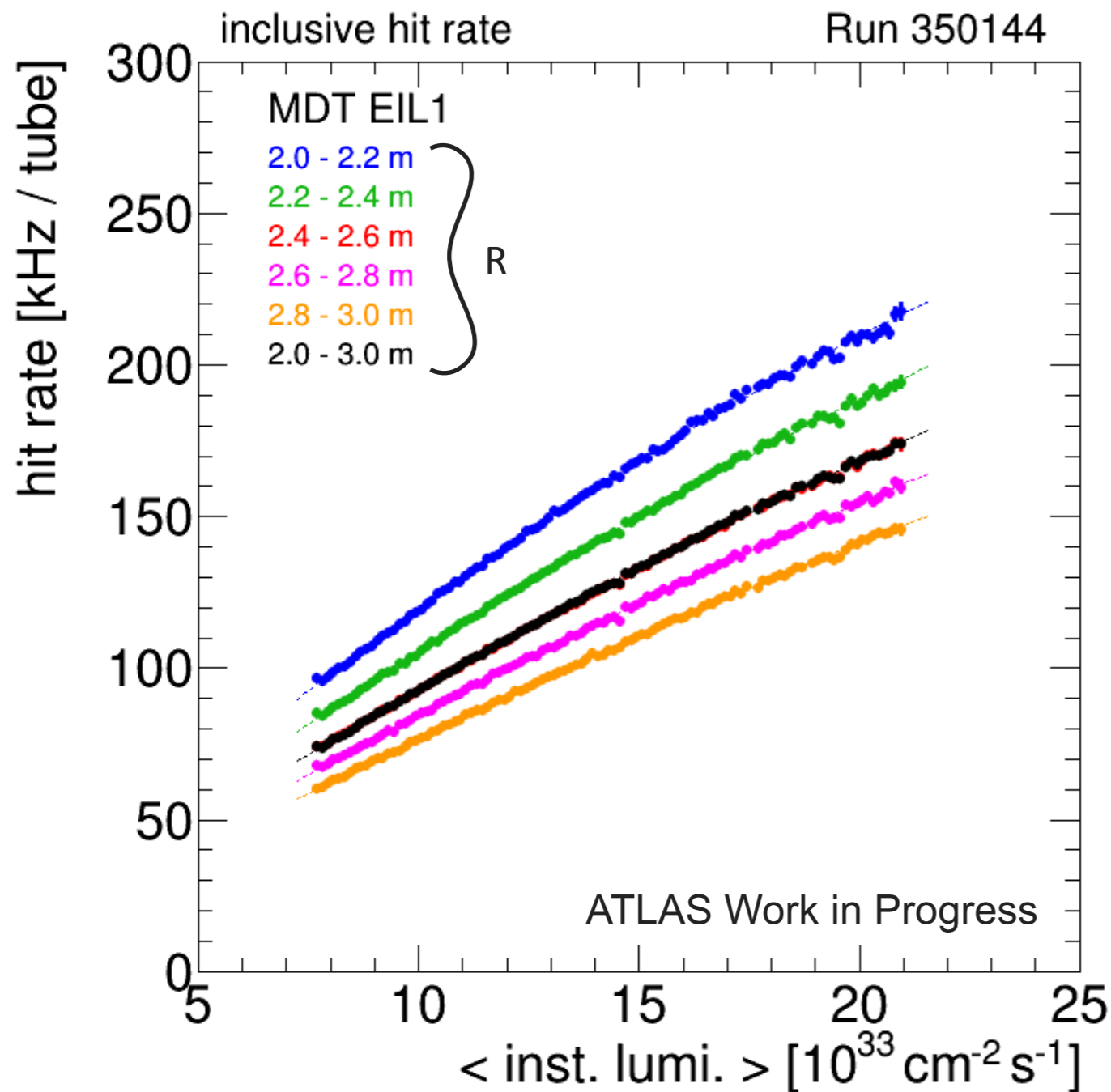


non-linearity consistent with effects from dead time \rightarrow new model

$$\text{observed hit rate} \sim A + B \cdot \mathcal{L} - C \cdot B^2 \cdot \mathcal{L}$$



1. Characterizing performance in Run 2



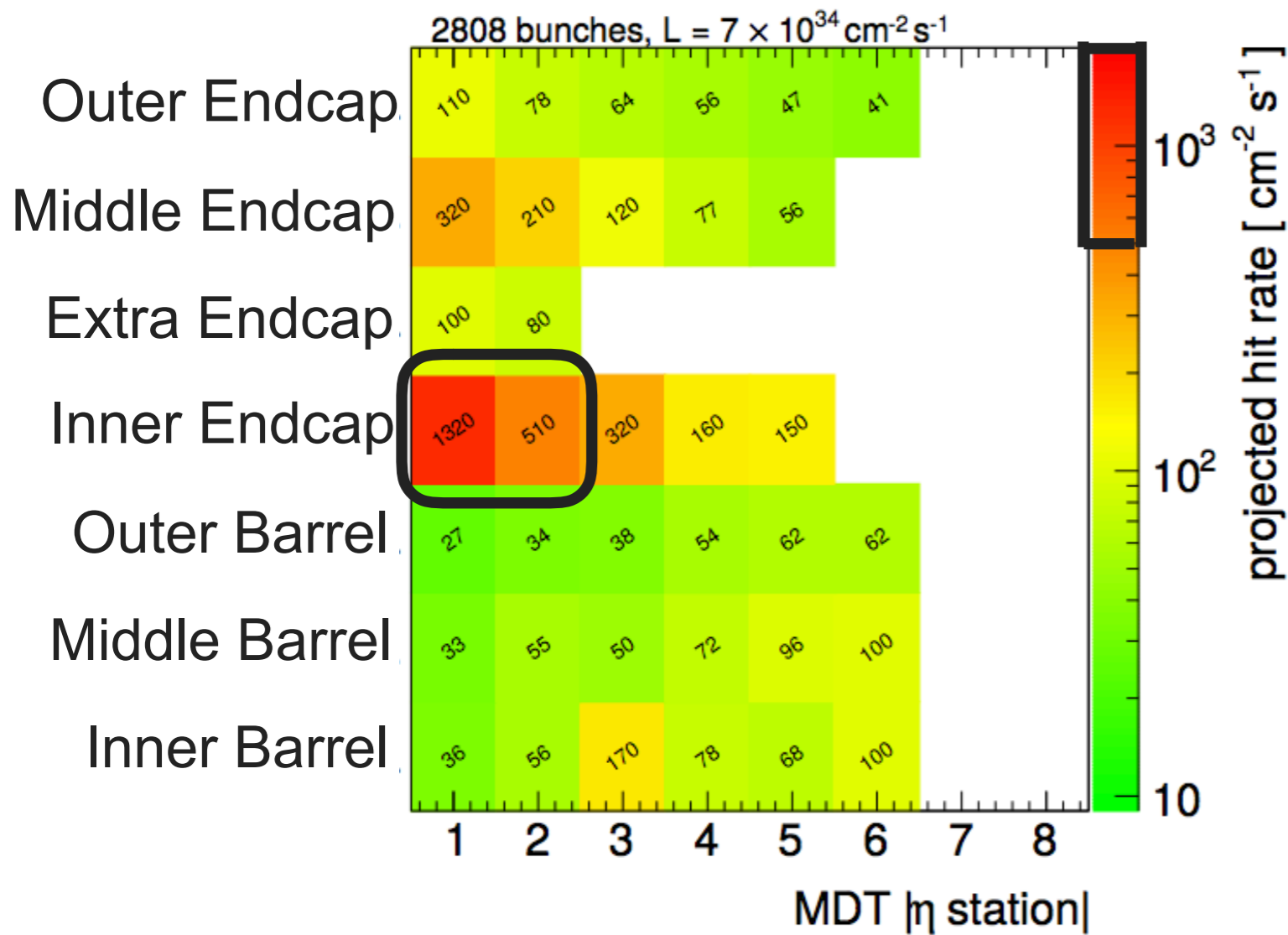
From data we can measure “C”
8.6% efficiency loss
for each 100 kHz/tube
increase in delivered hit rate

In the hottest tubes of MDT
corresponds to $\sim 20\%$ loss in single hit
efficiency at $2.0 \cdot 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

This result was expected!
and is ok for Run 3*

2. Making Projections to HL-LHC

MDT Large Sector Hit Rate Projections at HL-LHC
using extrapolation from 2015 data
Phase II Muon TDR - ATLAS-TDR-026



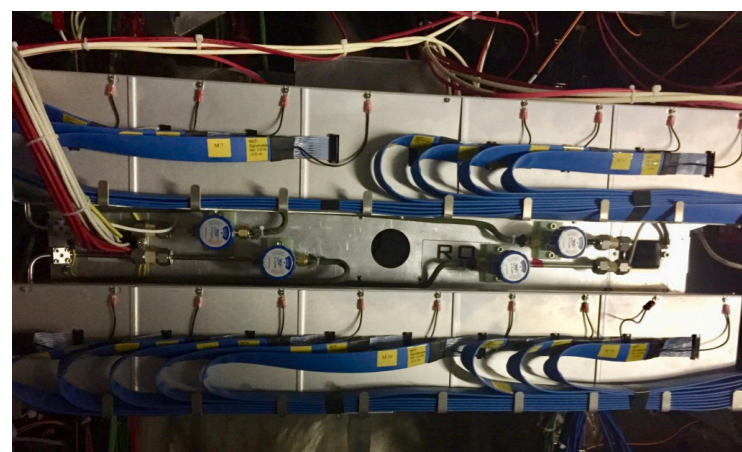
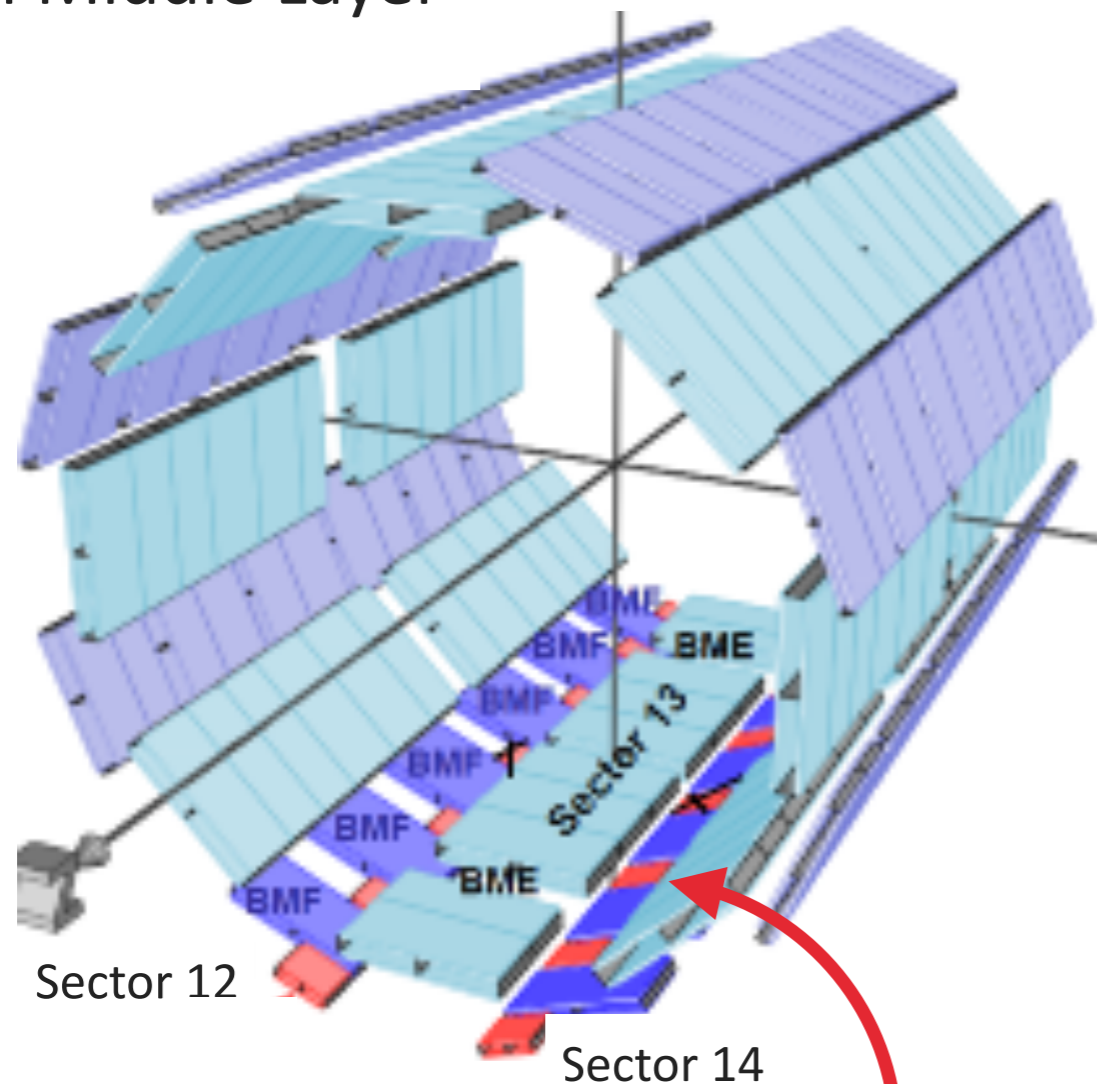
Hit rates exceed maximum
allowed 500 Hz/cm^2

Need to upgrade the Small
Wheel of Inner Endcap!

Consider two options
for our “New Small Wheel”

1. small-MDTs
2. MicroMegs

Barrel Middle Layer

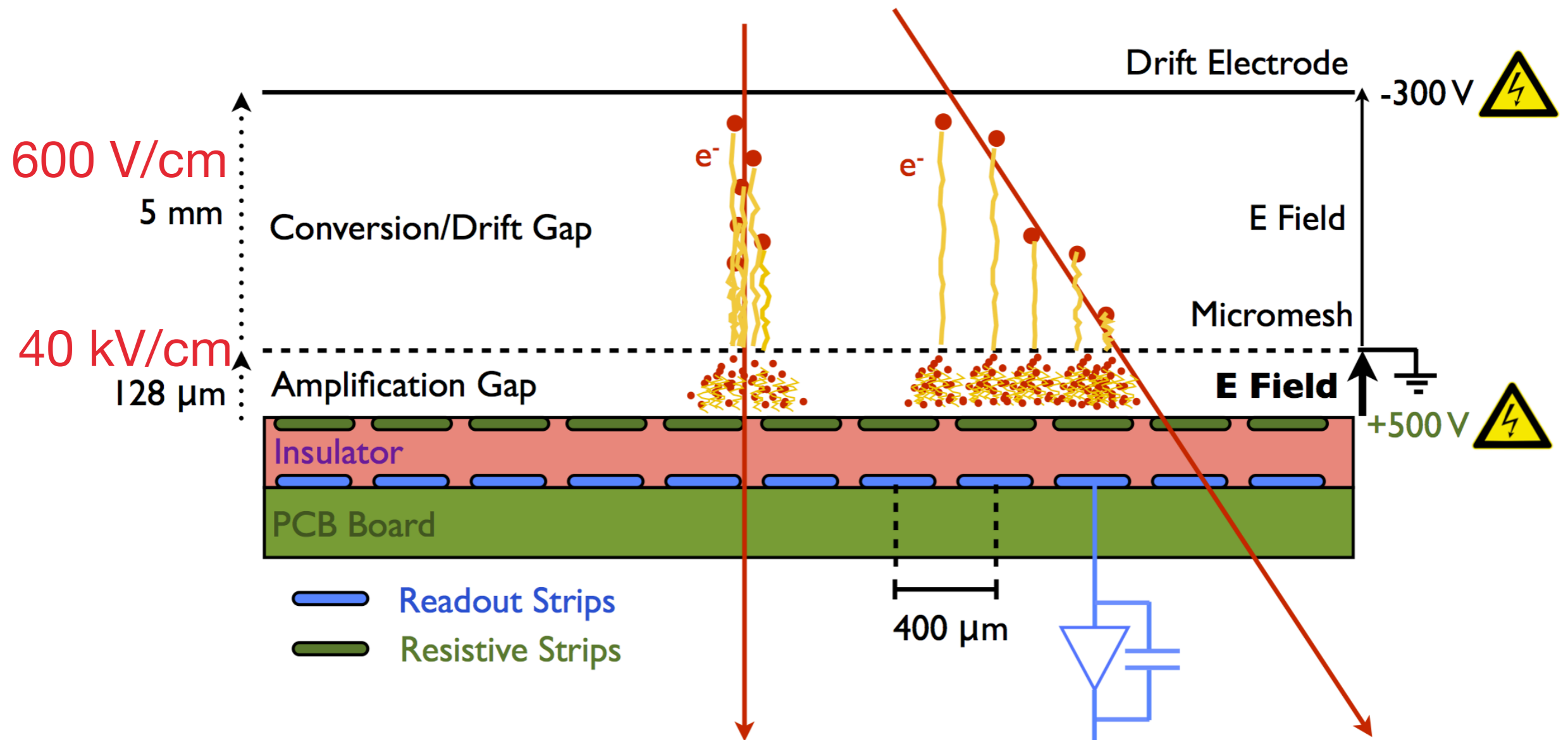


BMG2C14

Proven technology:
several small MDT chambers
already installed in ATLAS

Why they could be a solution
tubes 1/2 radius of regular MDTs
~8x better rate capability
max rate = 4 kHz/cm²

At the HL-LHC
could replace Small Wheel MDTs
but not the entire Small Wheel



1. Small amplification gap

Fast movement and absorption of ions
→ 200 ns dead time/strip

2. Small strip pitch 425-450 μm

Maximum Rate

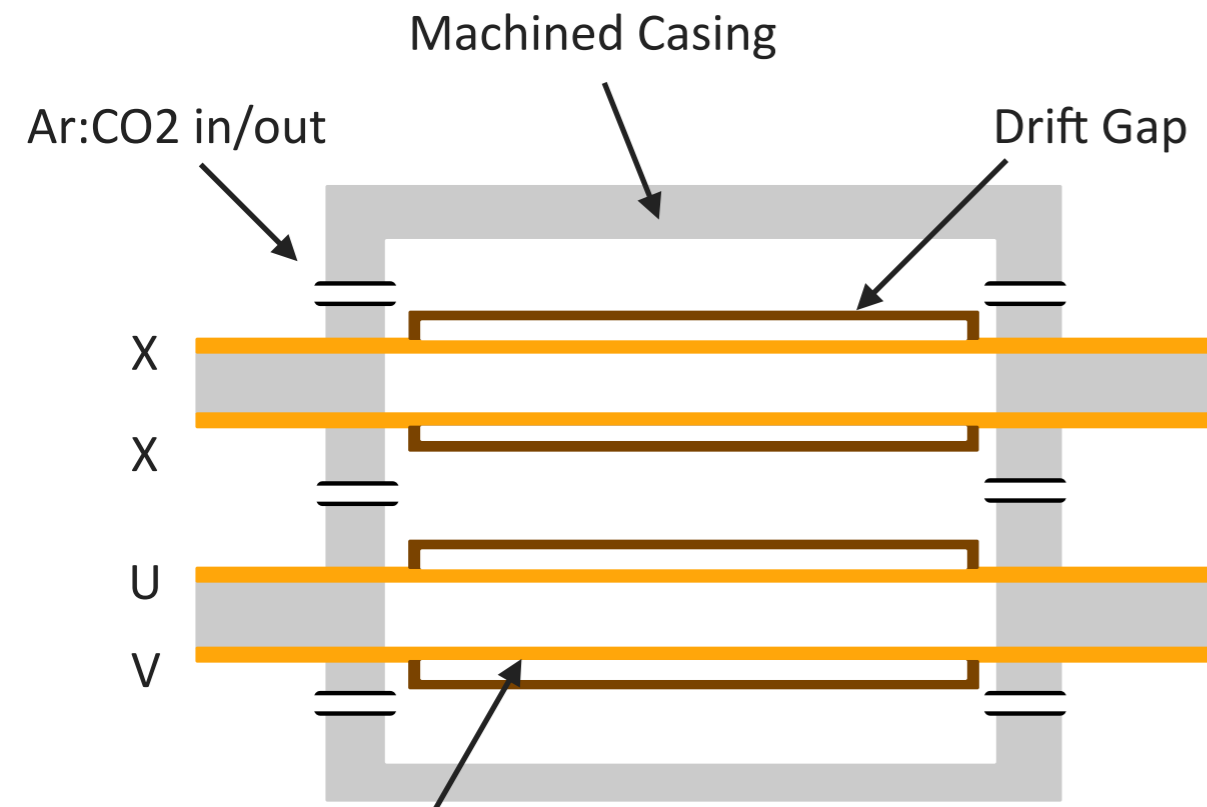
$$\sim 1 / (\text{dead time} * \text{strip area})$$

$$\sim 1 \text{ MHz} / \text{cm}^2$$

can replace entire small wheel!

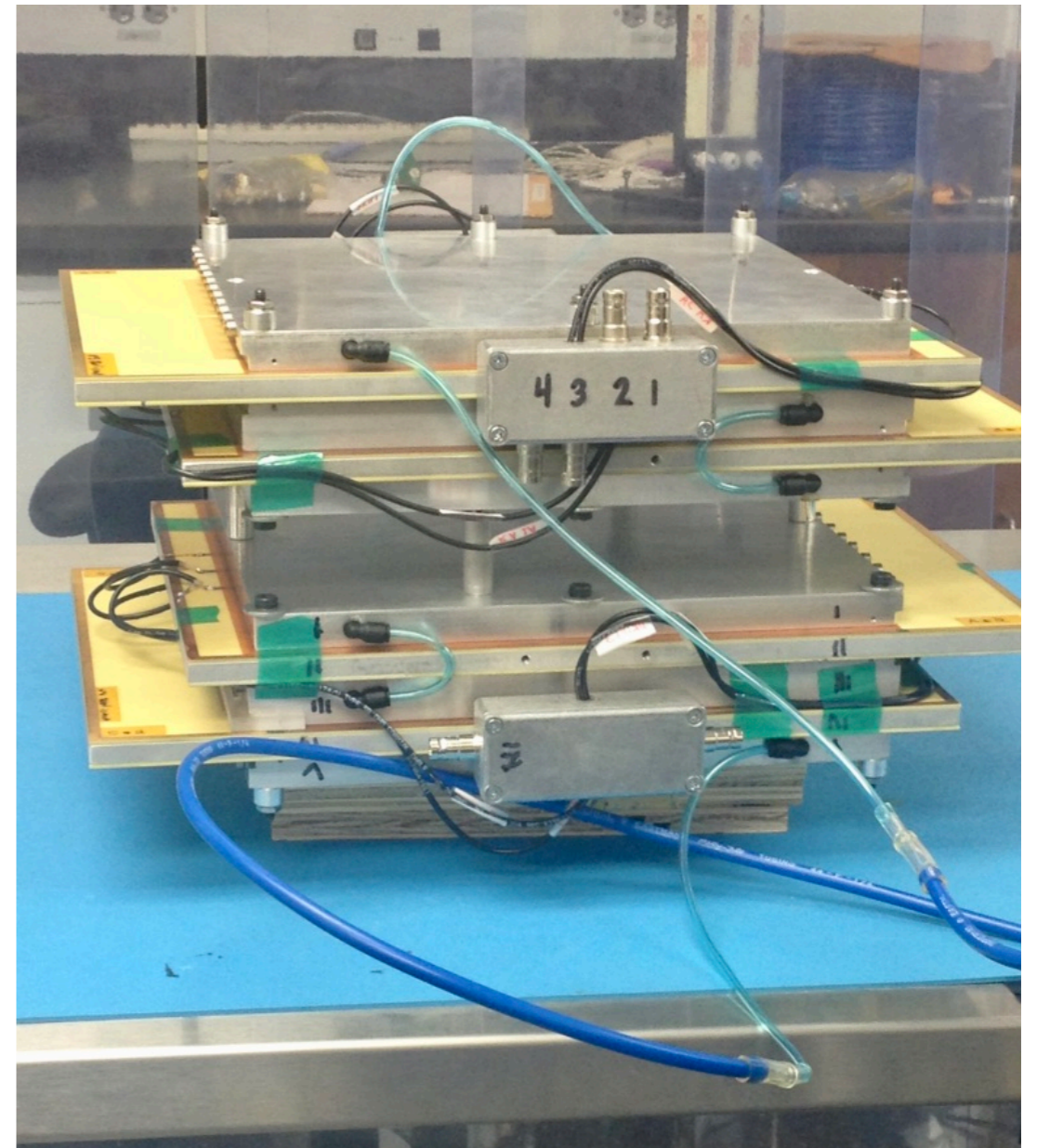
Quadruplet Design

Octuplet consists of 2 Quadruplets
20x20 cm chambers



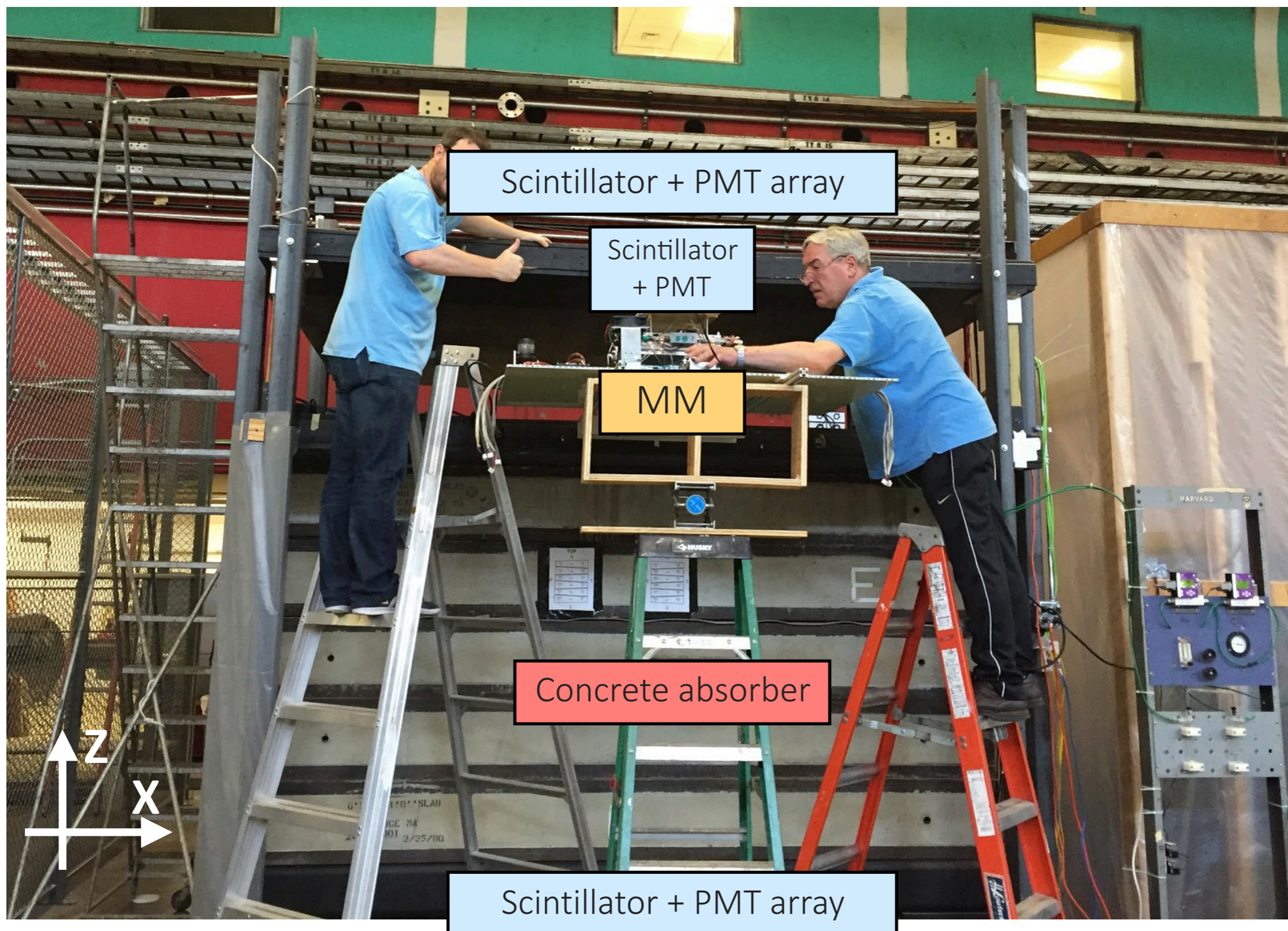
Completed PCB:
w/ Kapton, Resistive Strips,
Pillars, and Mesh

U strips angled at $+ 1.5^\circ$
V strips angled at $- 1.5^\circ$





Karri Folan DiPetrillo



Scintillator + PMT array

Scintillator + PMT

MM

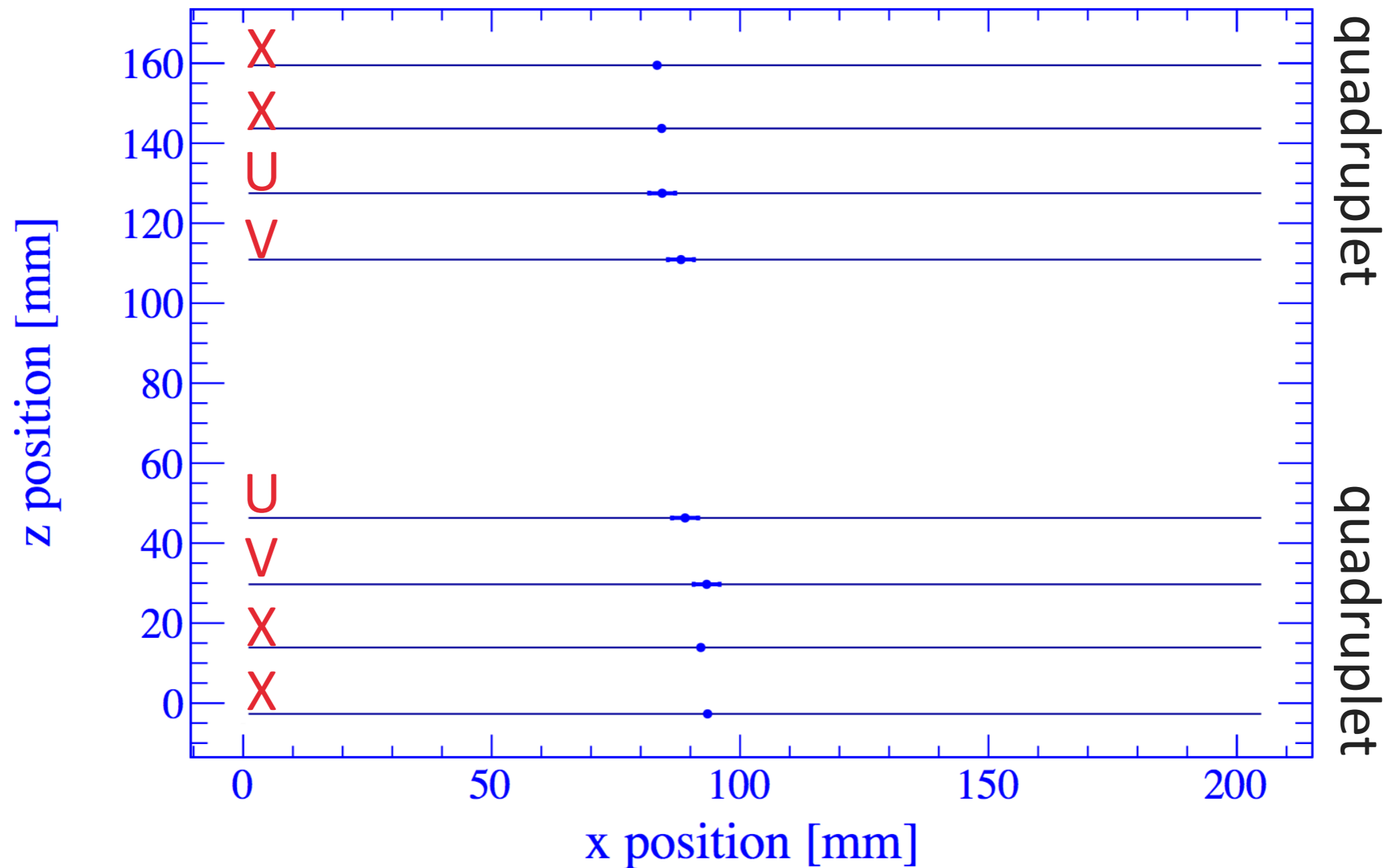
Concrete absorber

Scintillator + PMT array



Collected millions of cosmic muon events between 2016-2018

many useful studies of detector and electronics performance
with a particular focus on testing full trigger path



Building MicroMegas detectors is challenging
very sensitive to cleaning procedures, sparking, and noise

ATLAS also dealing with these challenges
working very hard on MicroMegas production
for New Small Wheel installation before 2021

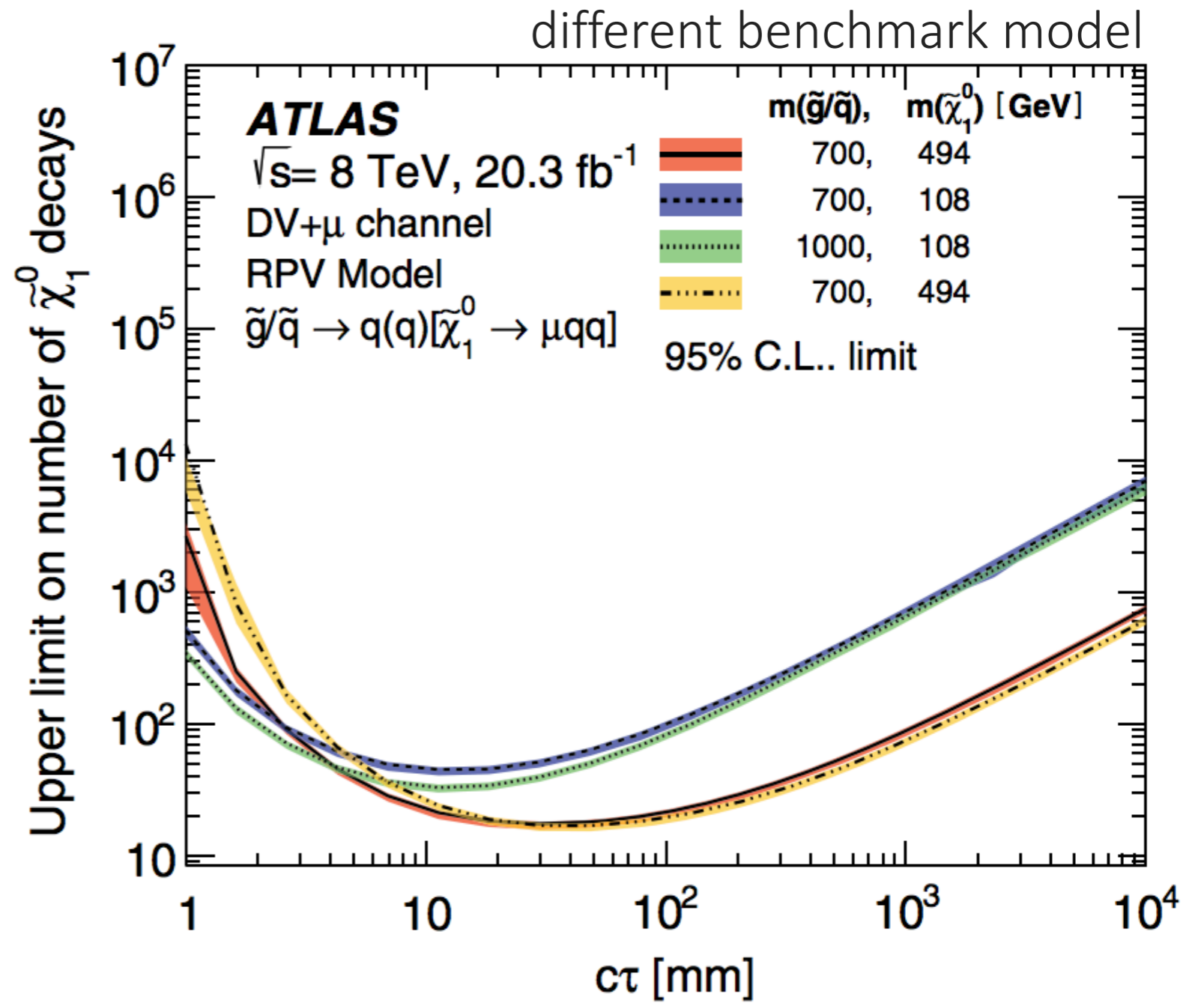
Looking forward to hunting for signs of new physics in Run 2 & 3
and preparing ATLAS detector for future data-taking

And special thanks to

DV+muon: Lawrence Lee Jr., Melissa Franklin, and SUSY DV Team

Muon Spectrometer: Alex Tuna, Siyuan Sun, Tony Tong, Ann Wang, Tomo Lazovich,
Chris Rogan, Melissa Franklin, Paolo Giromini, Tiesheng Dai, Zhen Yan, and Philipp
Fleischmann

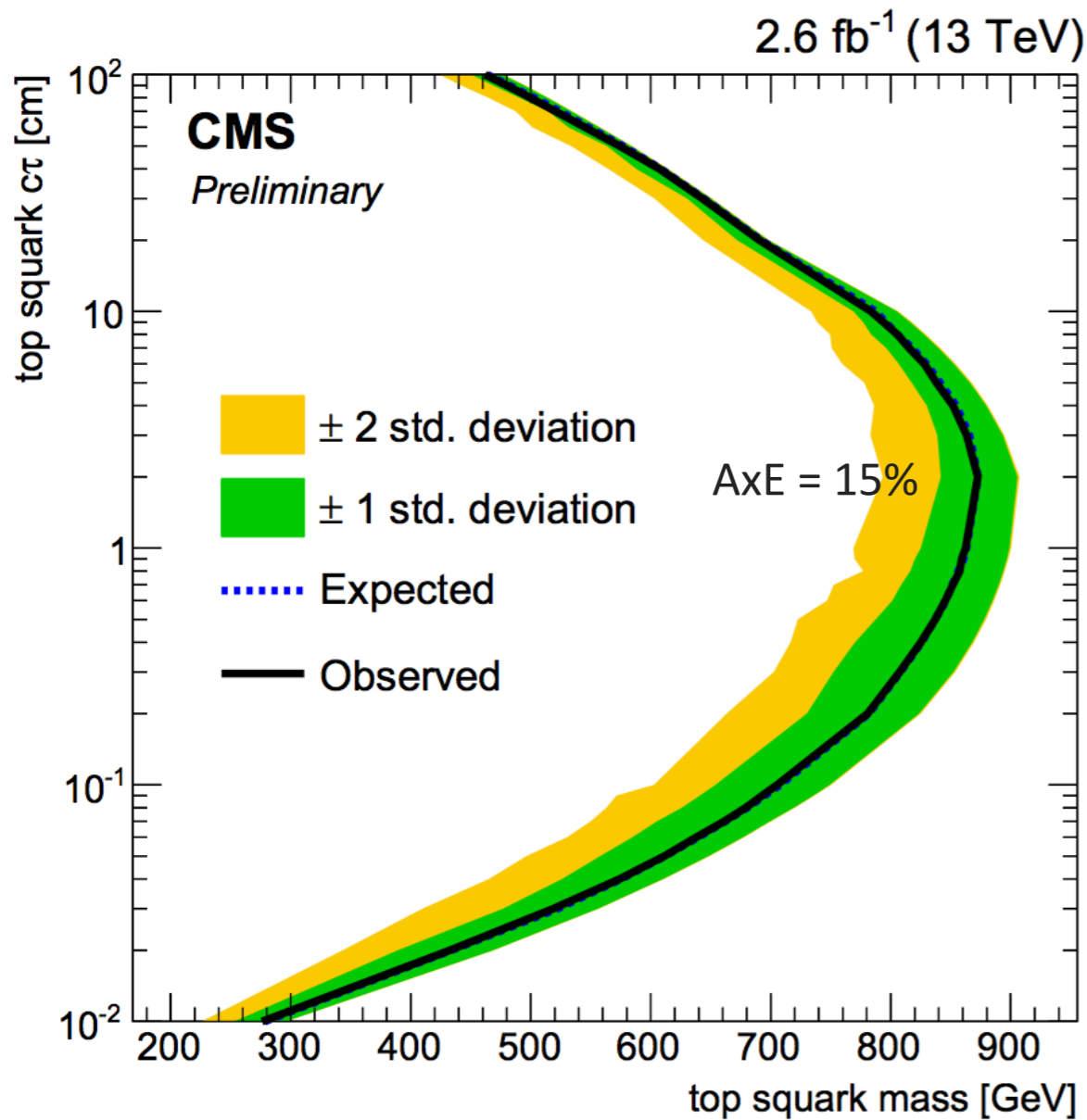
Best sensitivity
 ~20% signal acceptance x
 efficiency
 for $t \rightarrow q + \mu$ at $\tau = 0.1$ ns
 excludes stops up to 1 TeV



Displaced Leptons, $e\mu$

probes two λ' couplings

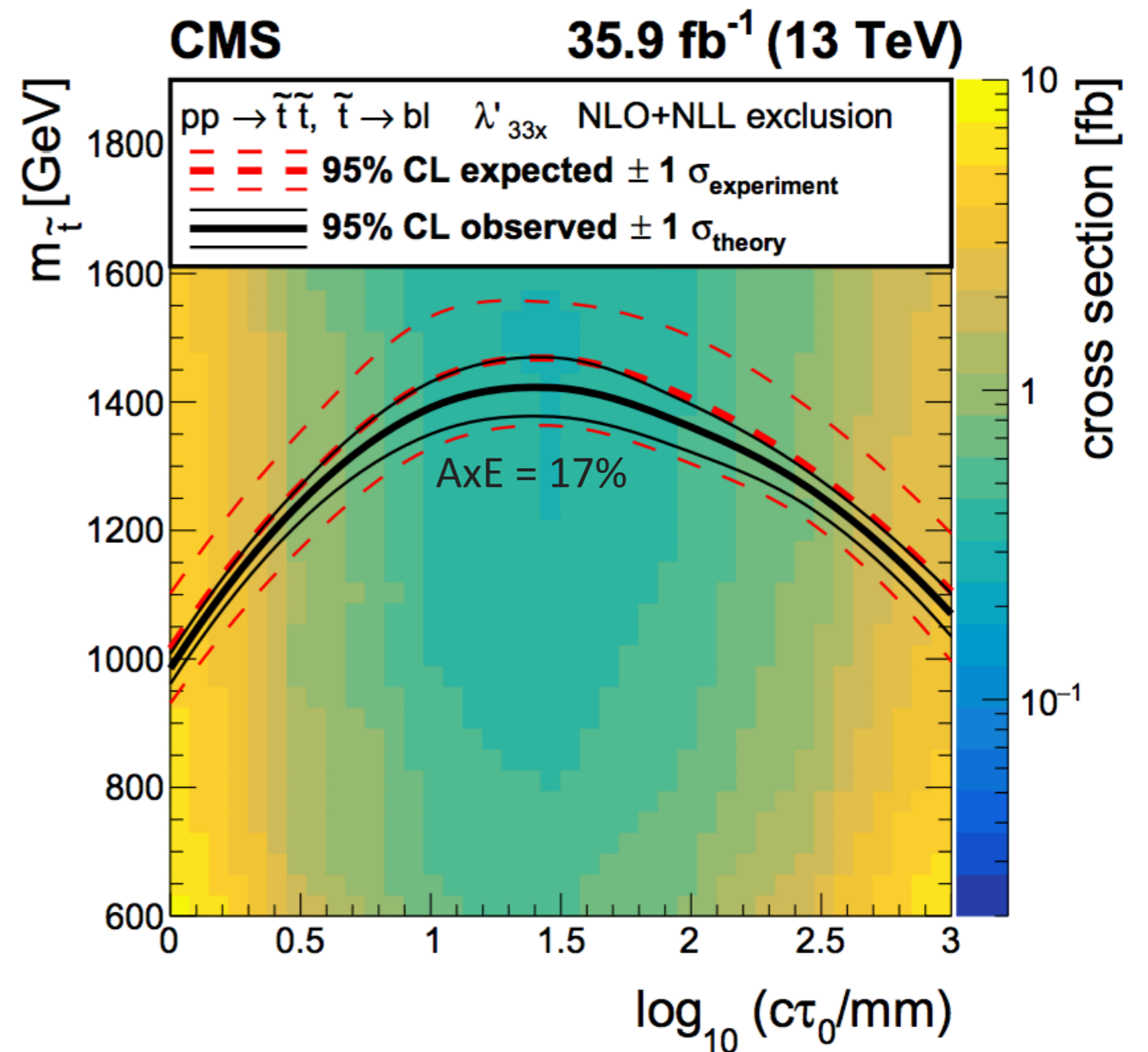
850 TeV stops at $\tau \approx 0.1$ ns, 2.6 fb^{-1}

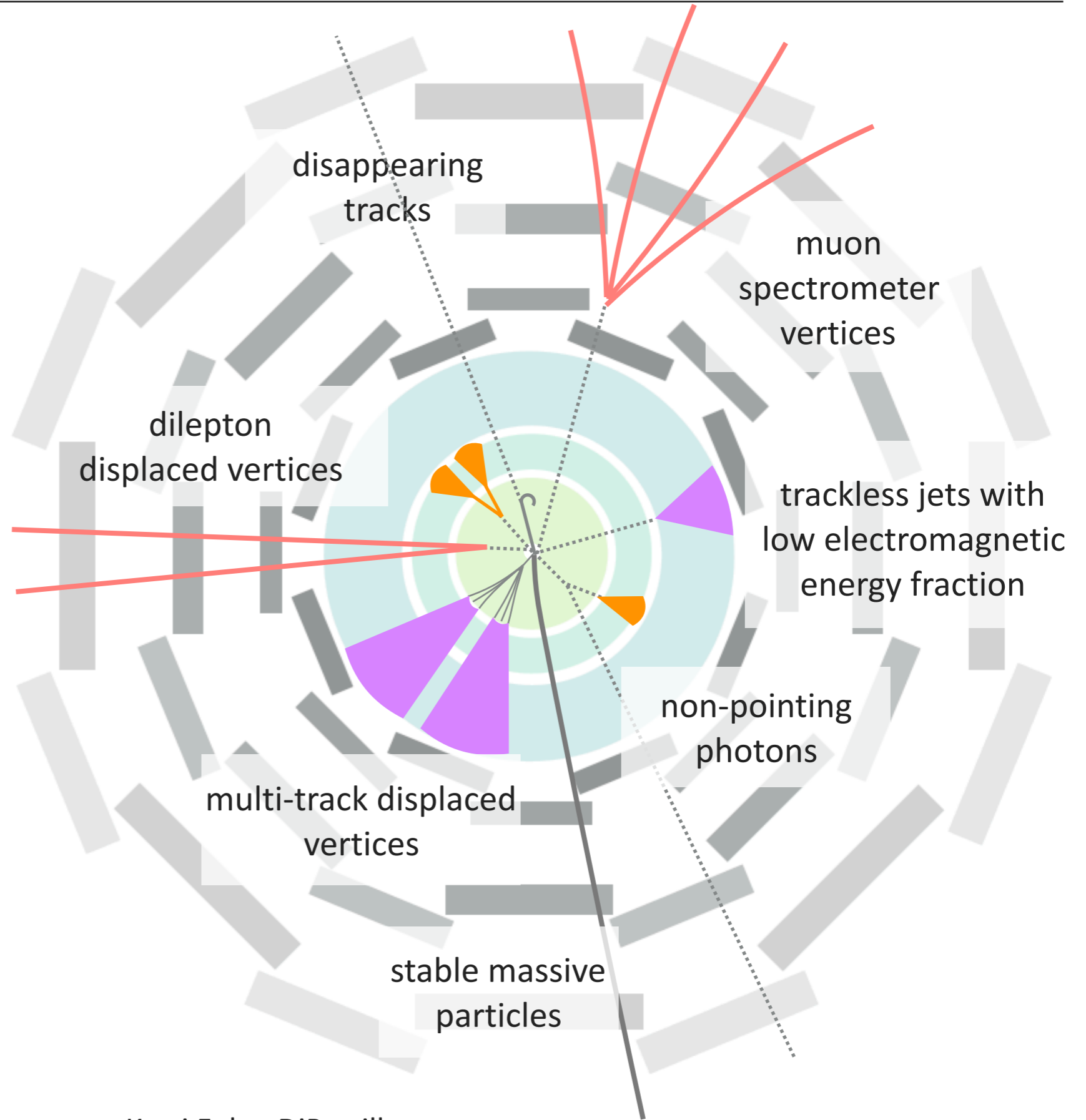


Displaced Jets

jets need not be from the same vertex

best: 1.4 TeV stops at $\tau \approx 0.1$ ns, 35.9 fb^{-1}





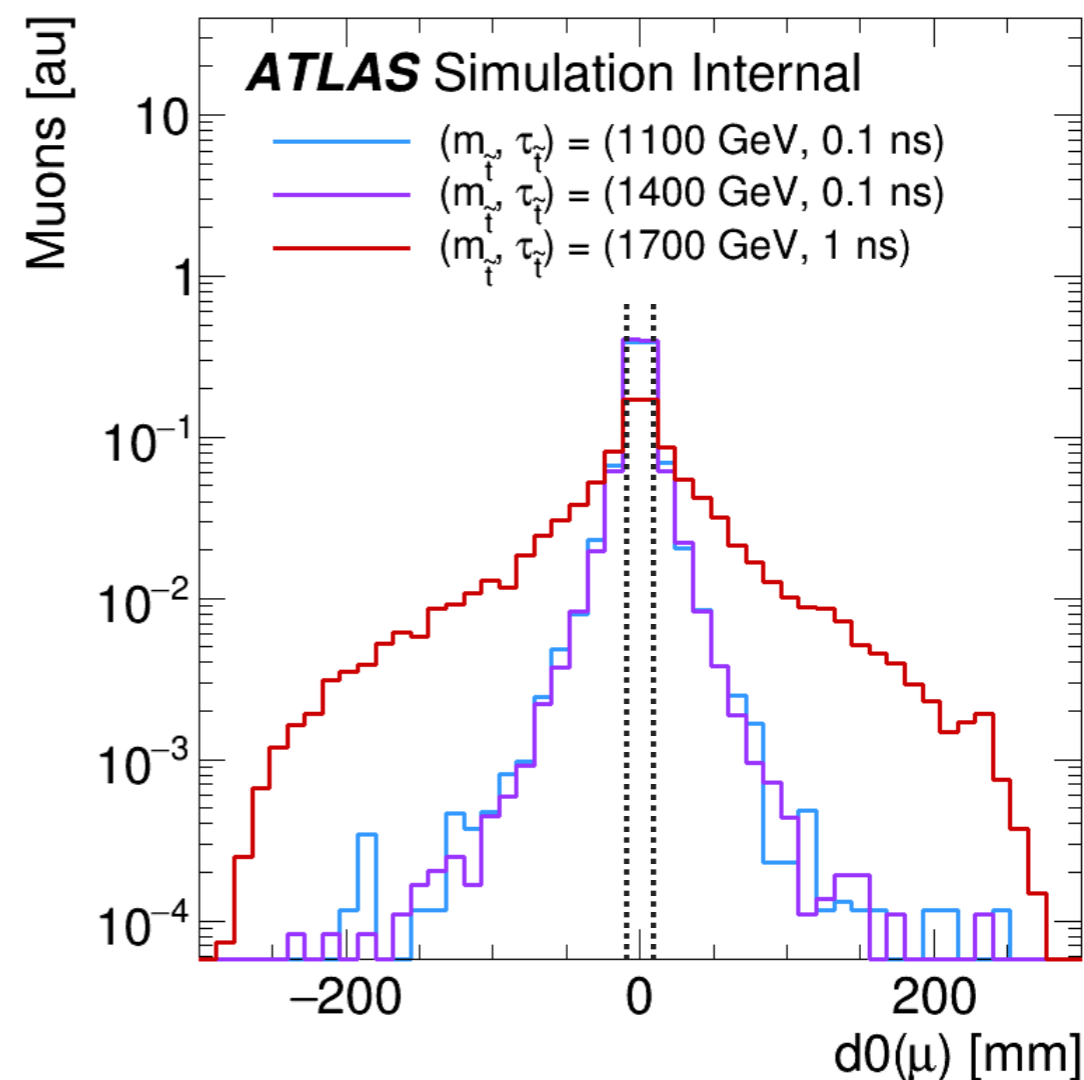
Level 1 Trigger uses Calorimeter & Muon Spectrometer information
cannot trigger on inner detector displaced vertices in ATLAS

Standard muon triggers

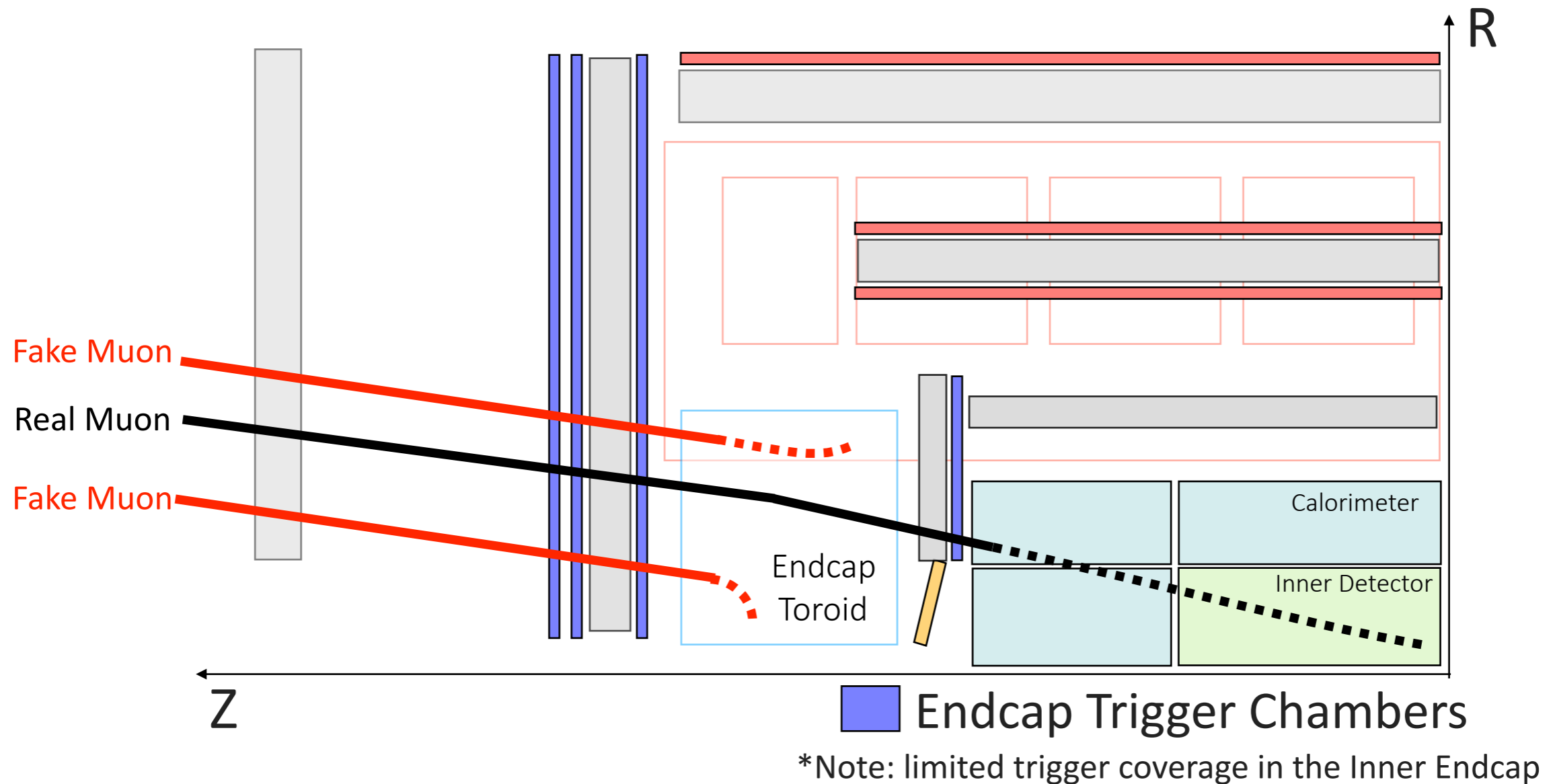
Require a Muon Track
matched to an Inner Detector Track

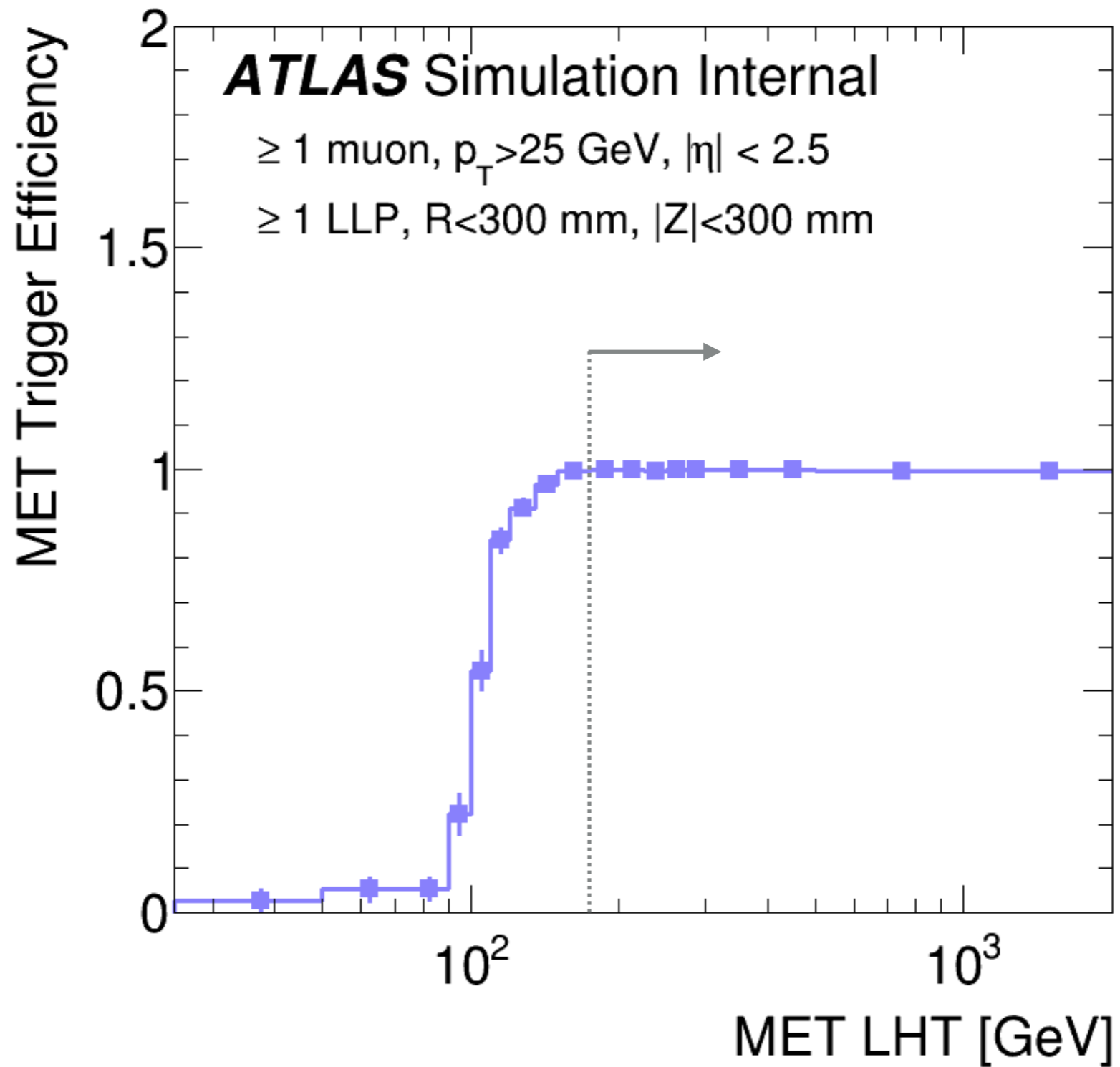
Strict impact parameter requirements
 $|d_0| < 10$ mm

Inefficient for signal!



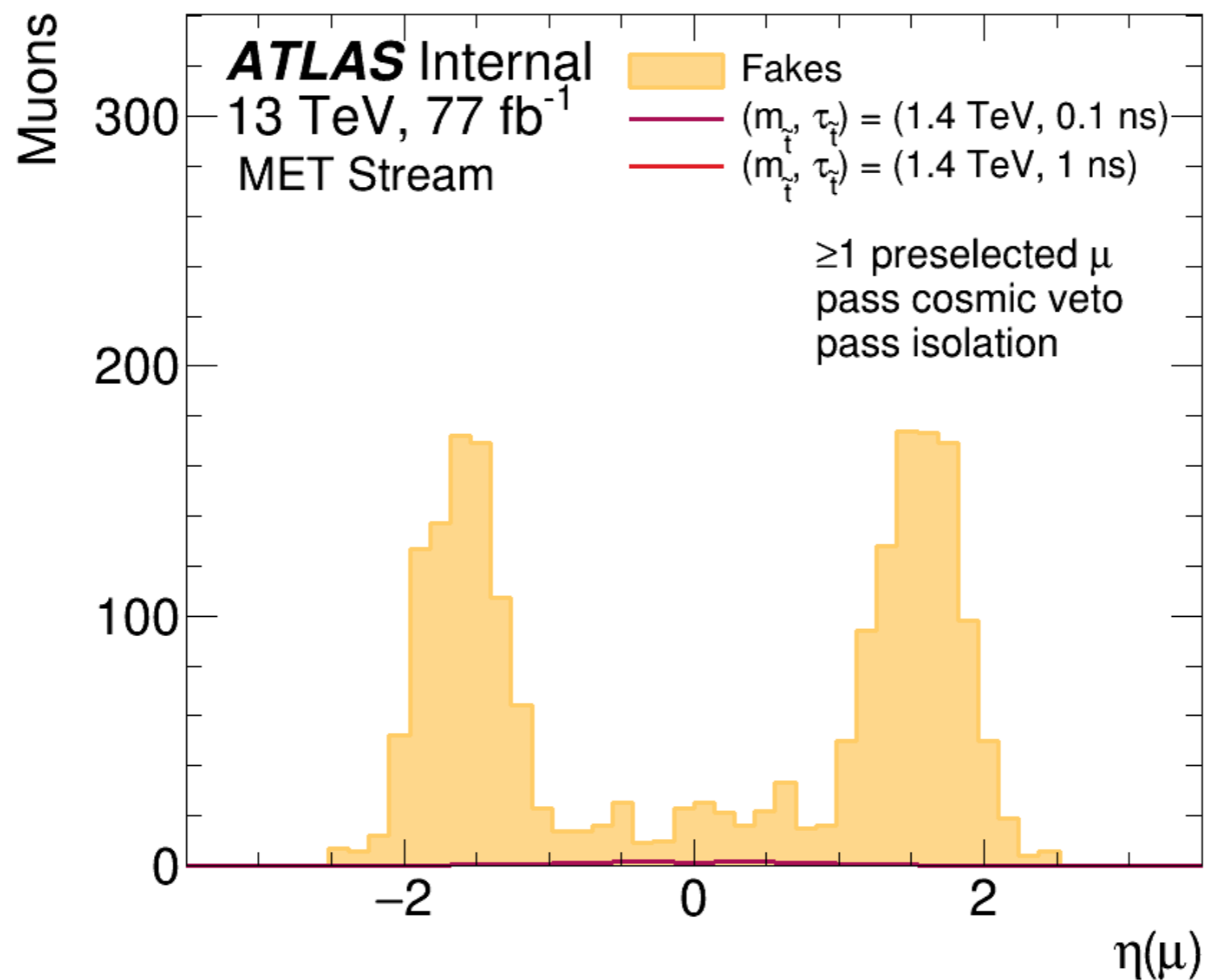
Muon Spectrometer Only Trigger Backgrounds cannot use Endcaps because of large background rate





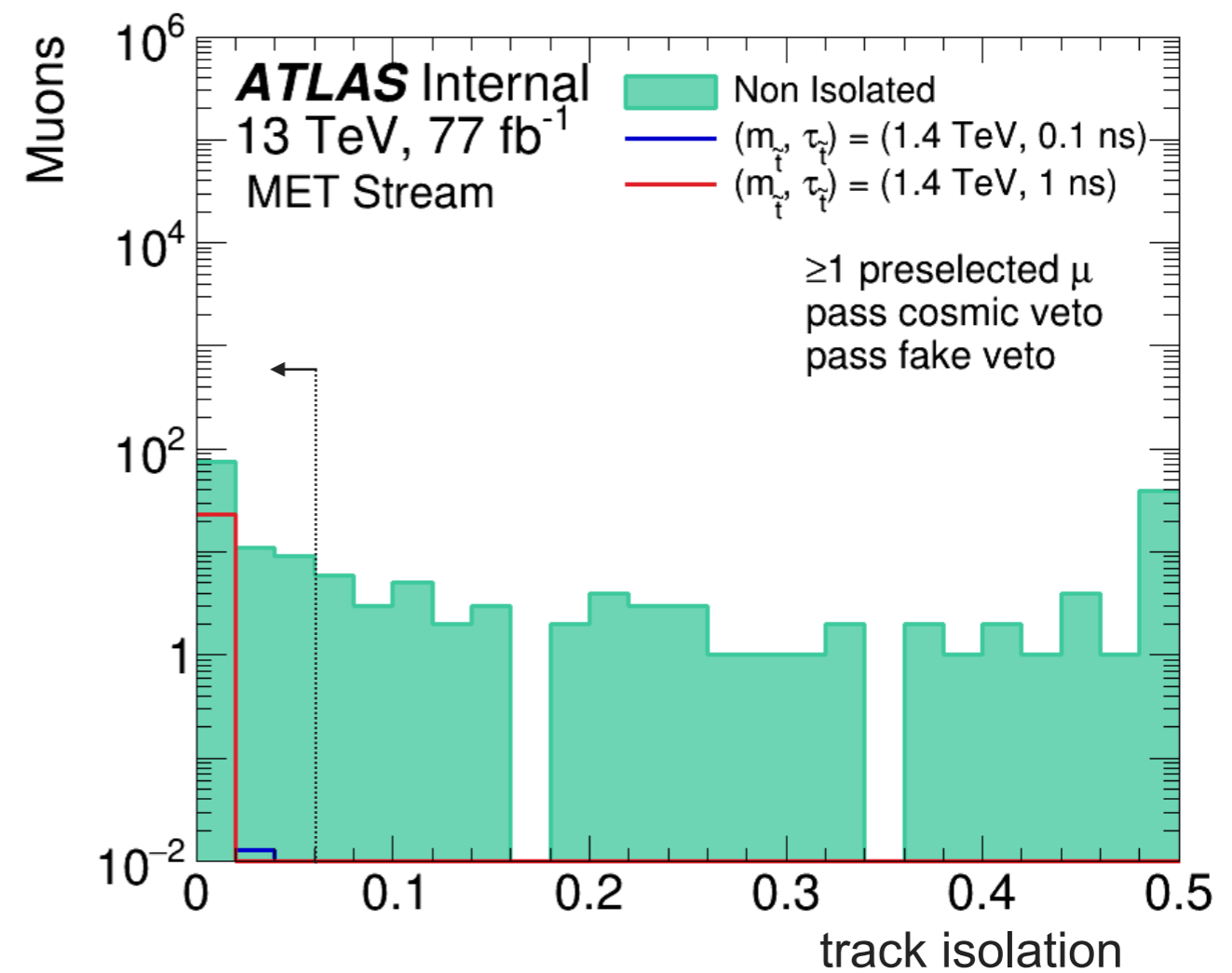
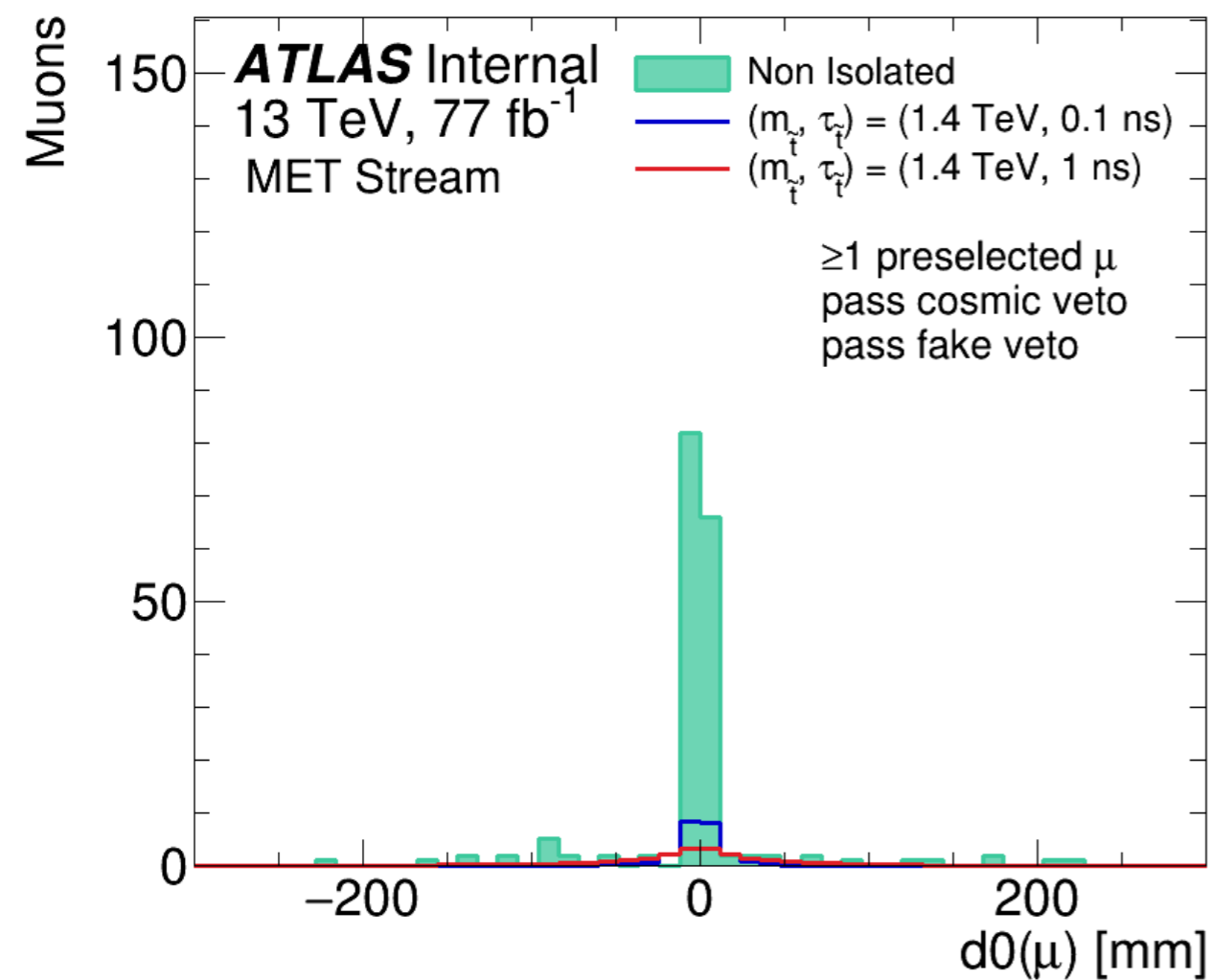
Large background in MET triggered events

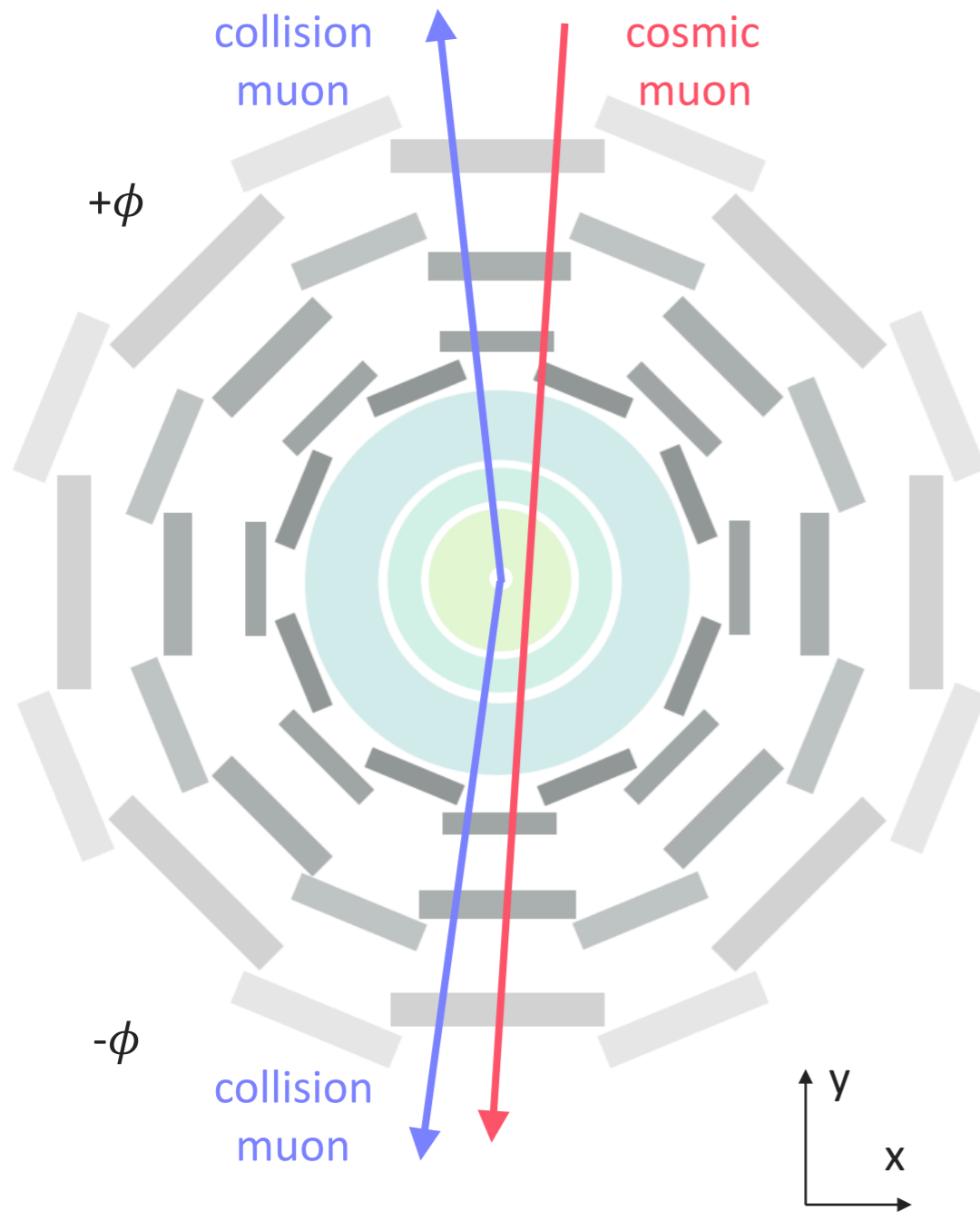
fake inner detector track matched to fake muon spectrometer track
 mostly affects endcaps



Heavy Flavor Muons

tend to have small impact parameters
also shown, track isolation



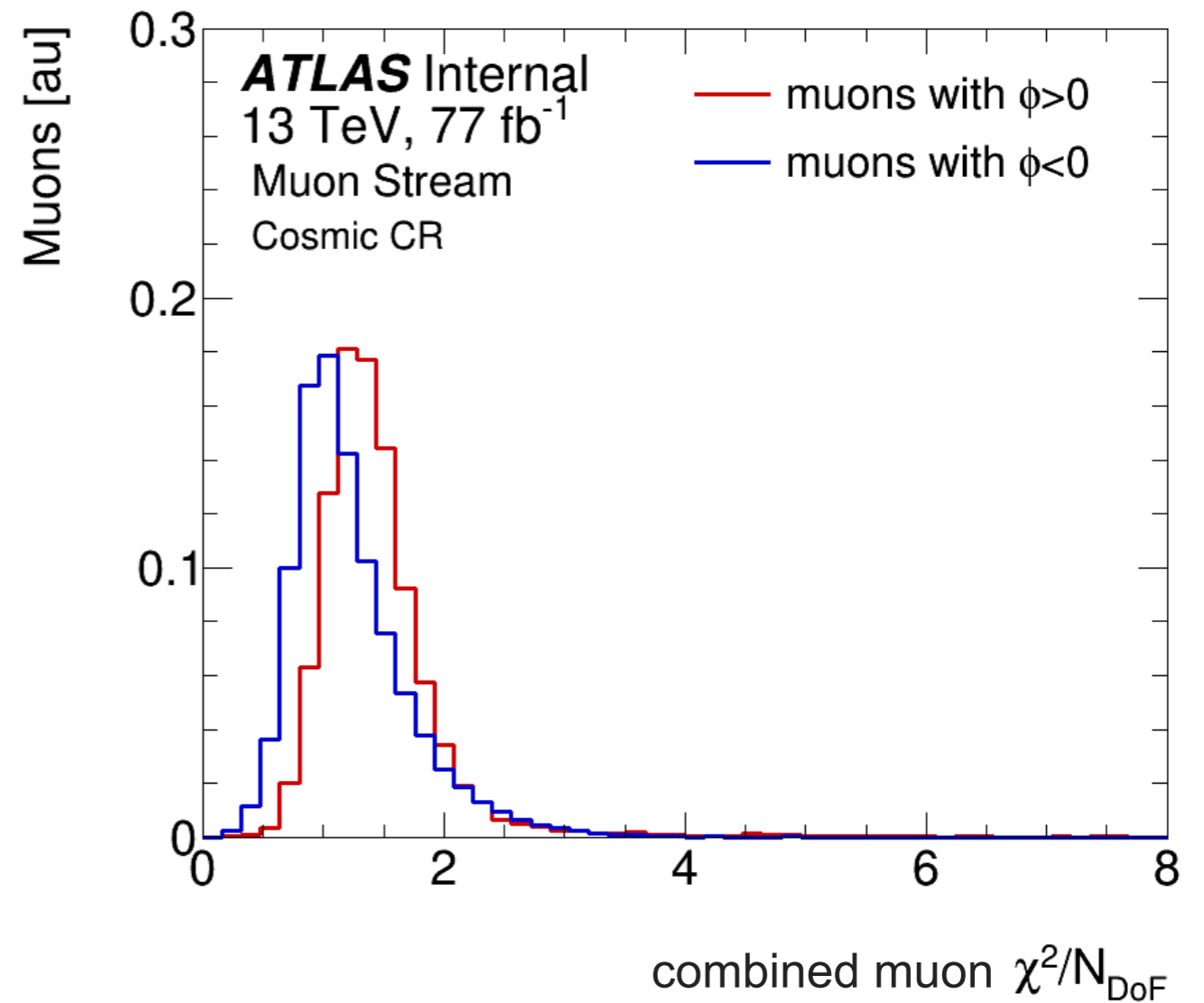
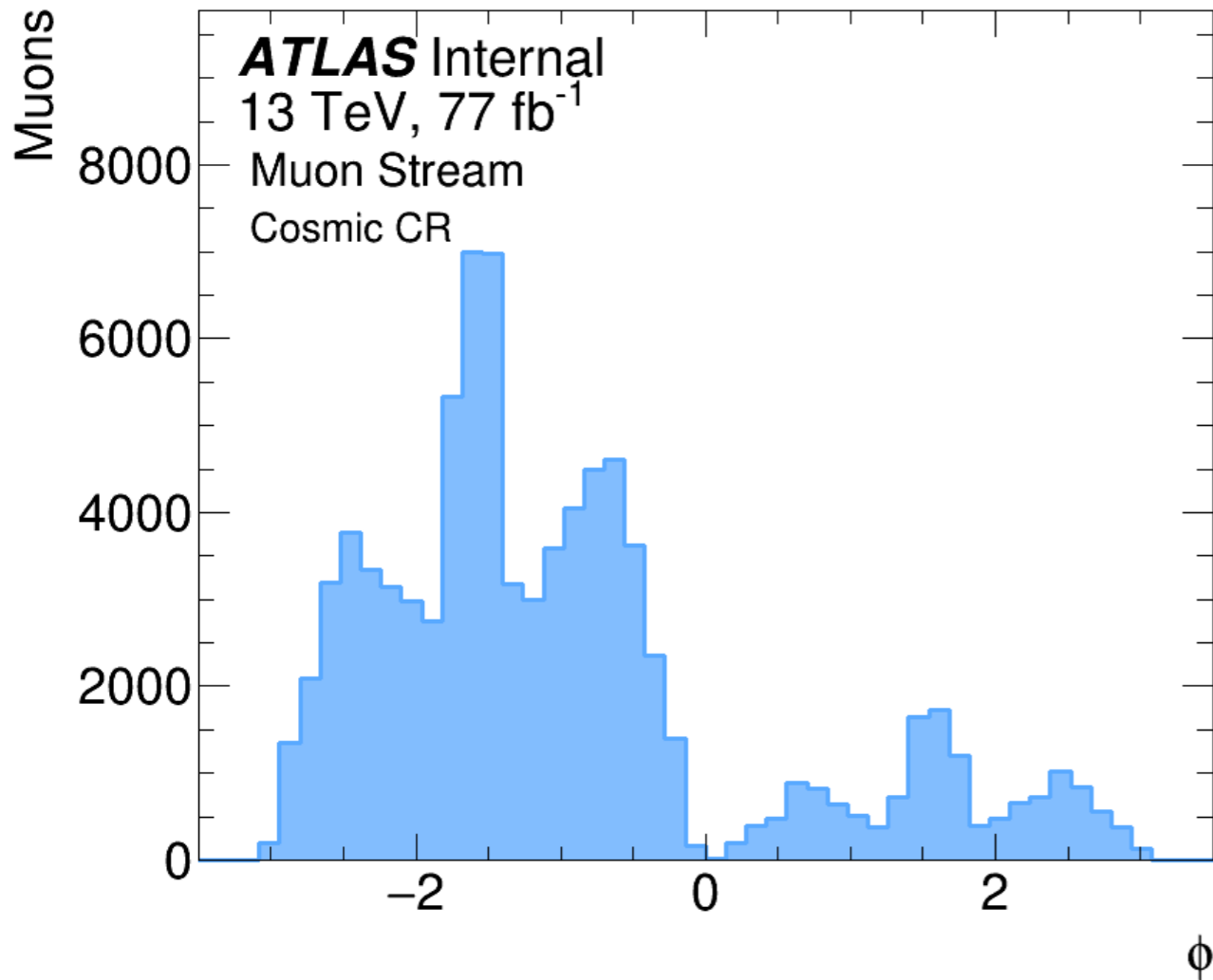


$+\phi$ muon timing difference with respect to collision muons

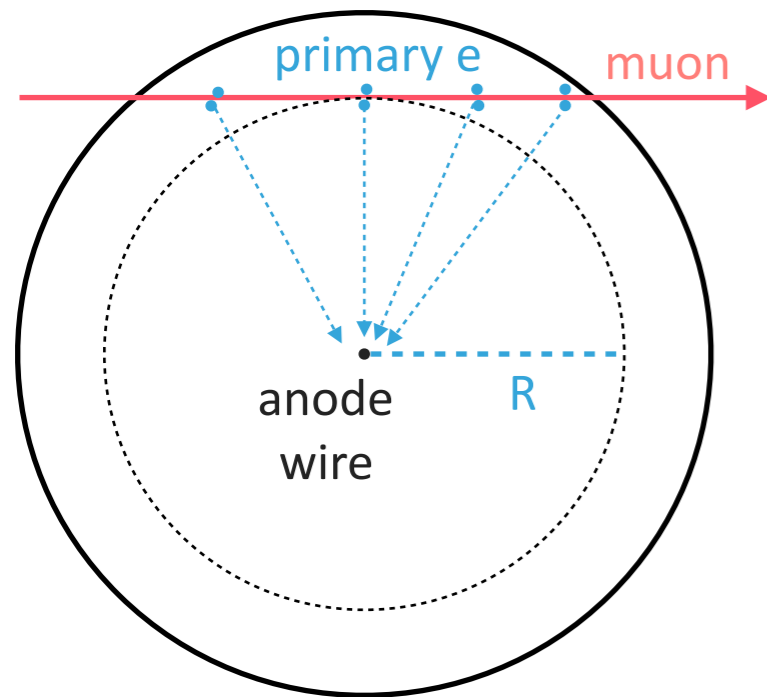
Muon Station	R [m]	Δt
Inner	5	-33 ns
Middle	7.5	-50 ns
Outer	10	-67 ns

*if $-\phi$ muon is in time with collisions

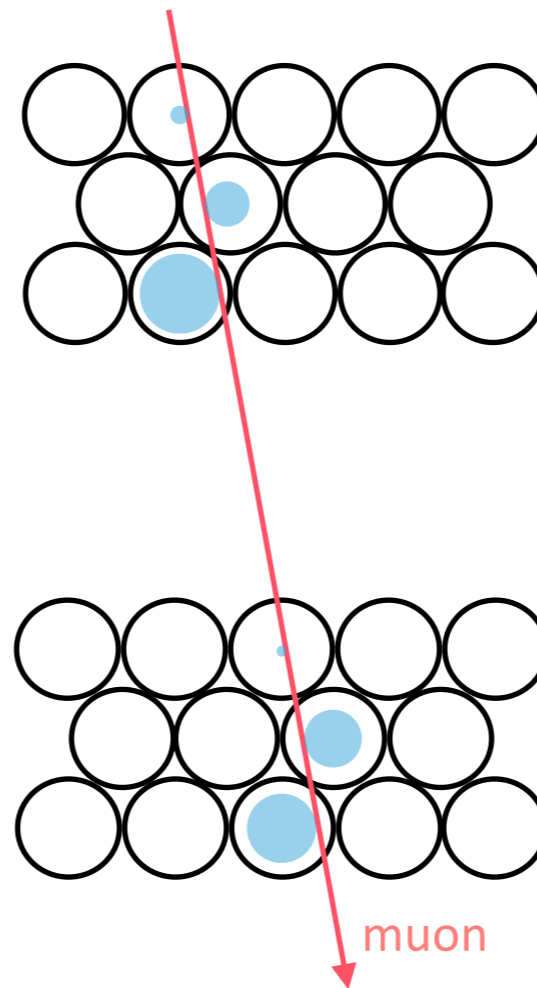
positive phi cosmic legs often aren't reconstructed as muons
when they are reconstructed, they tend to have fewer hits on track,
and poorer fit quality



why timing affects quality of fit in monitored drift tubes



cathode tube
3 cm diameter



muon passing through
MDT chamber

if timing assumptions
are incorrect
eg. time of flight

measured hit positions
are incorrect

→ poorly fit muons

muon η measured at impact parameter
segment η measured w.r.t origin

Goal: correct muon η to detector η
such that $\eta_{\text{corr}}(\mu) + \eta(\text{seg}) = 0$

Step 1: find correct $Z(R, z_0)$

$$Z_{\text{corr}} = Z_{\mu}(R) - z_0$$

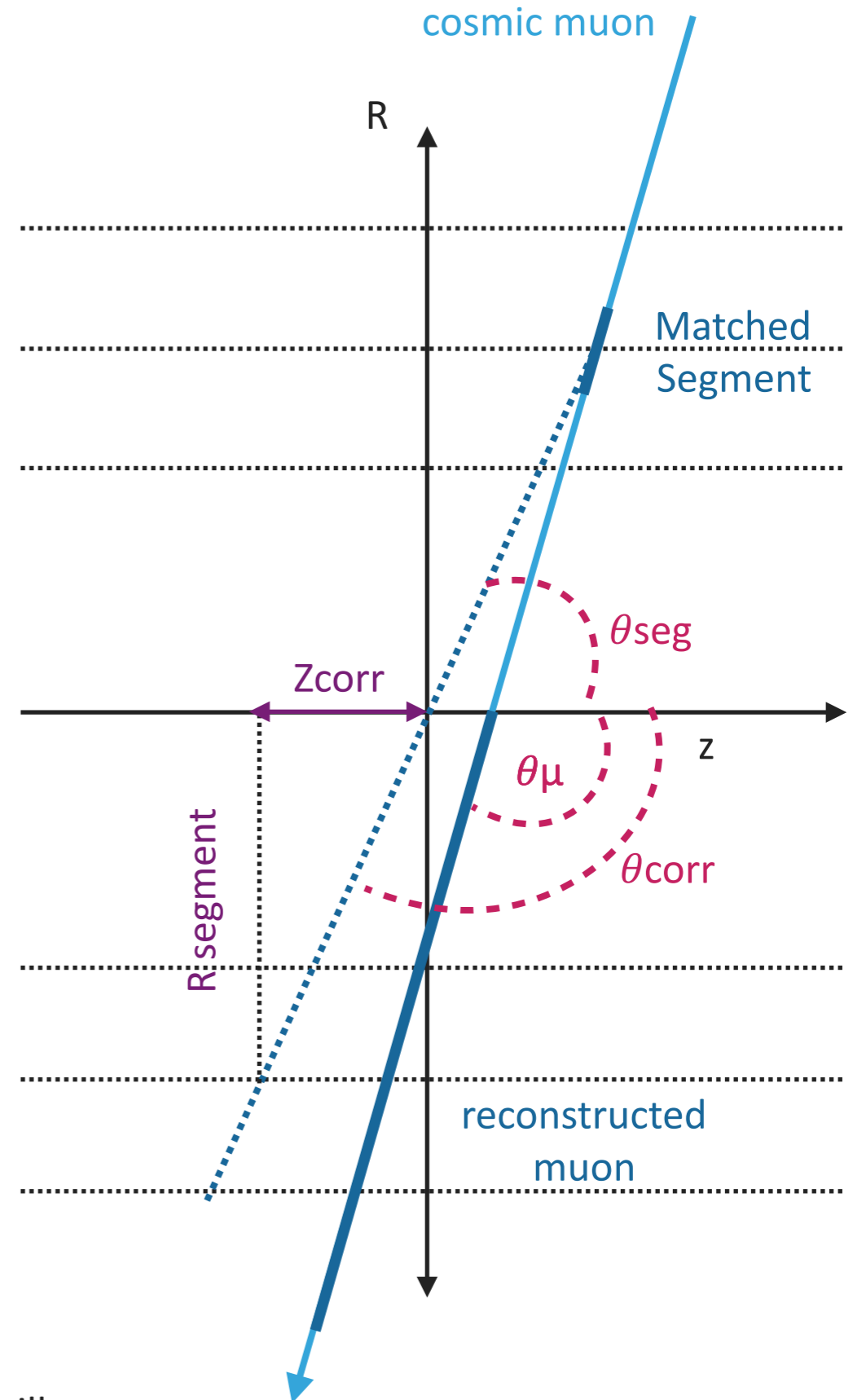
$\tan(\theta_{\mu}) = R_{\text{segment}} / z_{\text{muon}}$
gives you

$$Z_{\text{corr}} = R_{\text{segment}} / \tan(\theta_{\mu}) - z_0$$

Step 2: convert Z to $\eta(R, z_0)$

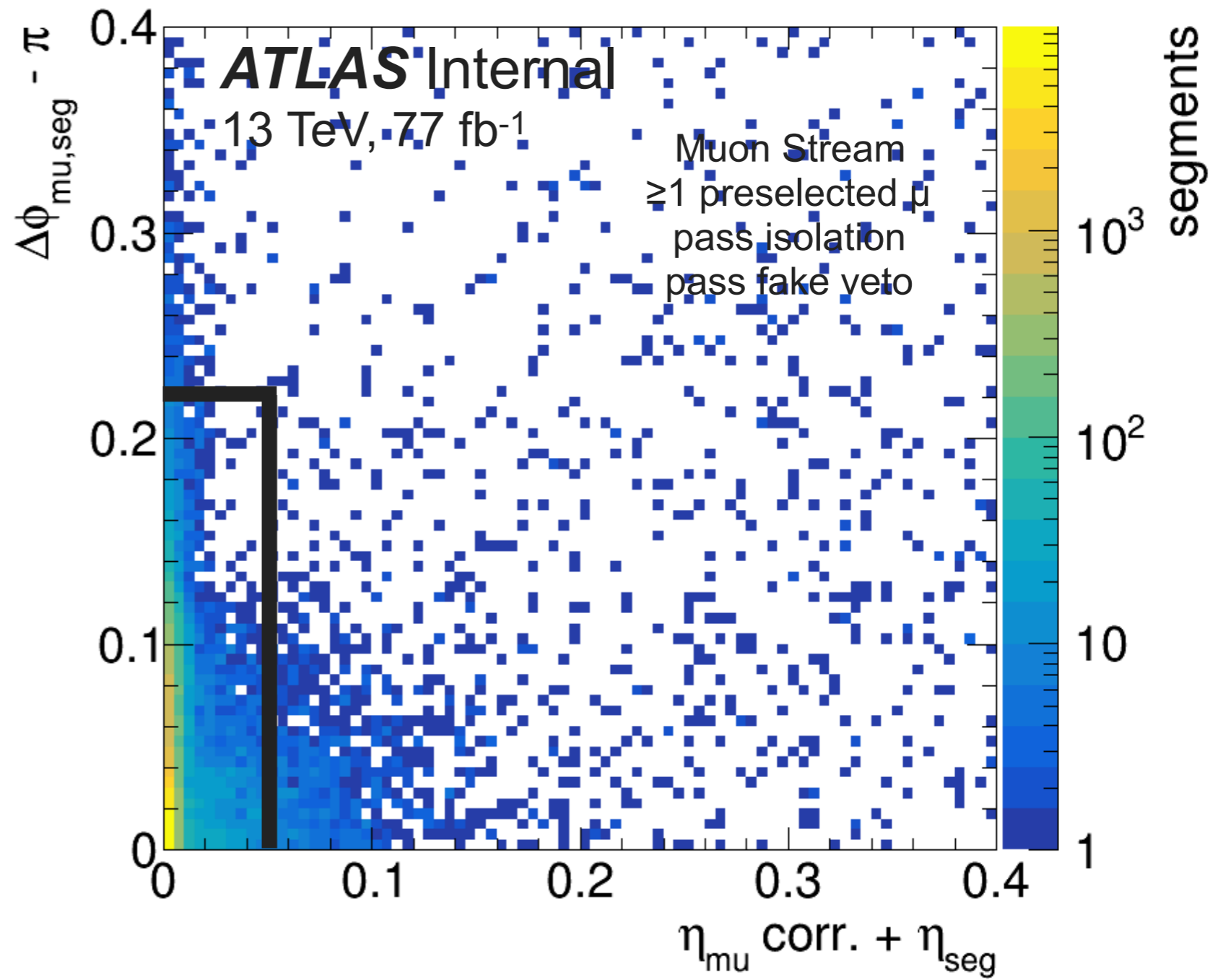
$$\theta_{\text{corr}} = \arctan(R_{\text{segment}} / Z_{\text{corr}})$$

$$\eta_{\text{corr}} = -\ln(\tan(\theta_{\text{seg}}/2))$$

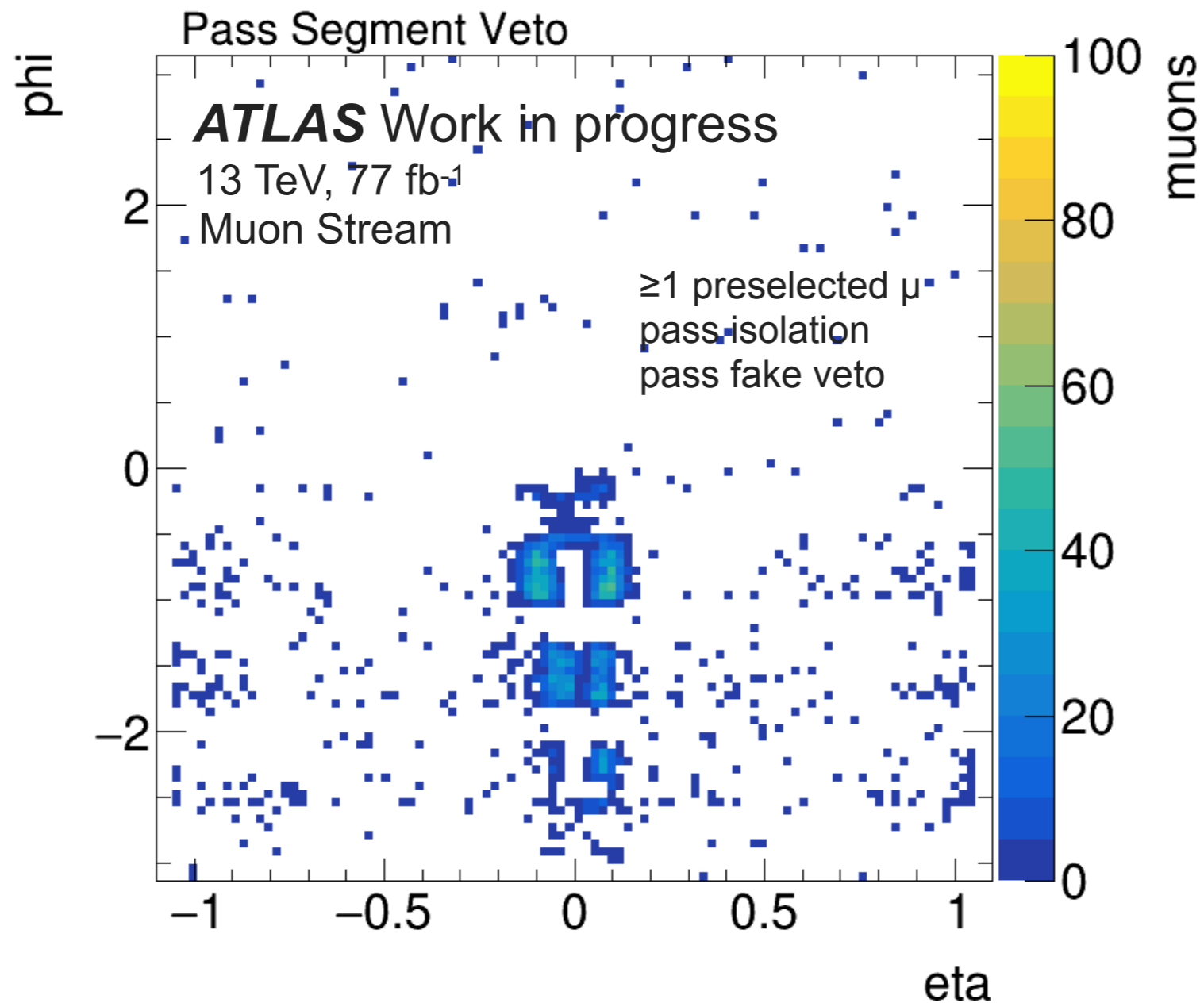


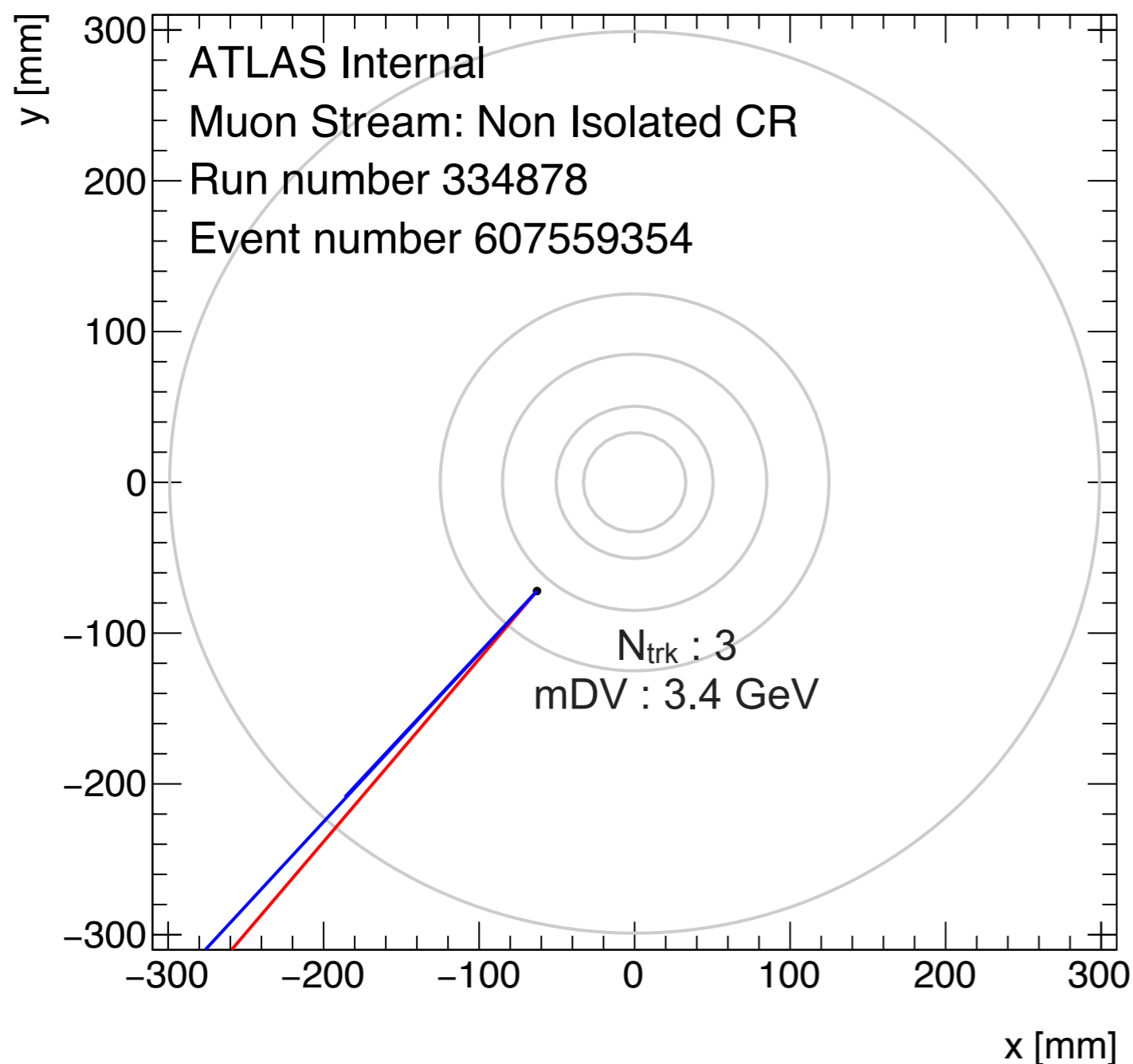
$\text{Sum } \eta_{\text{corr}} < 0.05$
 roughly size of a few tubes

$\Delta\phi - \pi < 0.22$
 roughly size of 1 MDT
 chamber in phi



Cosmics which pass the back-to-back segment veto
form interesting pattern in eta-phi
Near eta = 0, and in large sectors



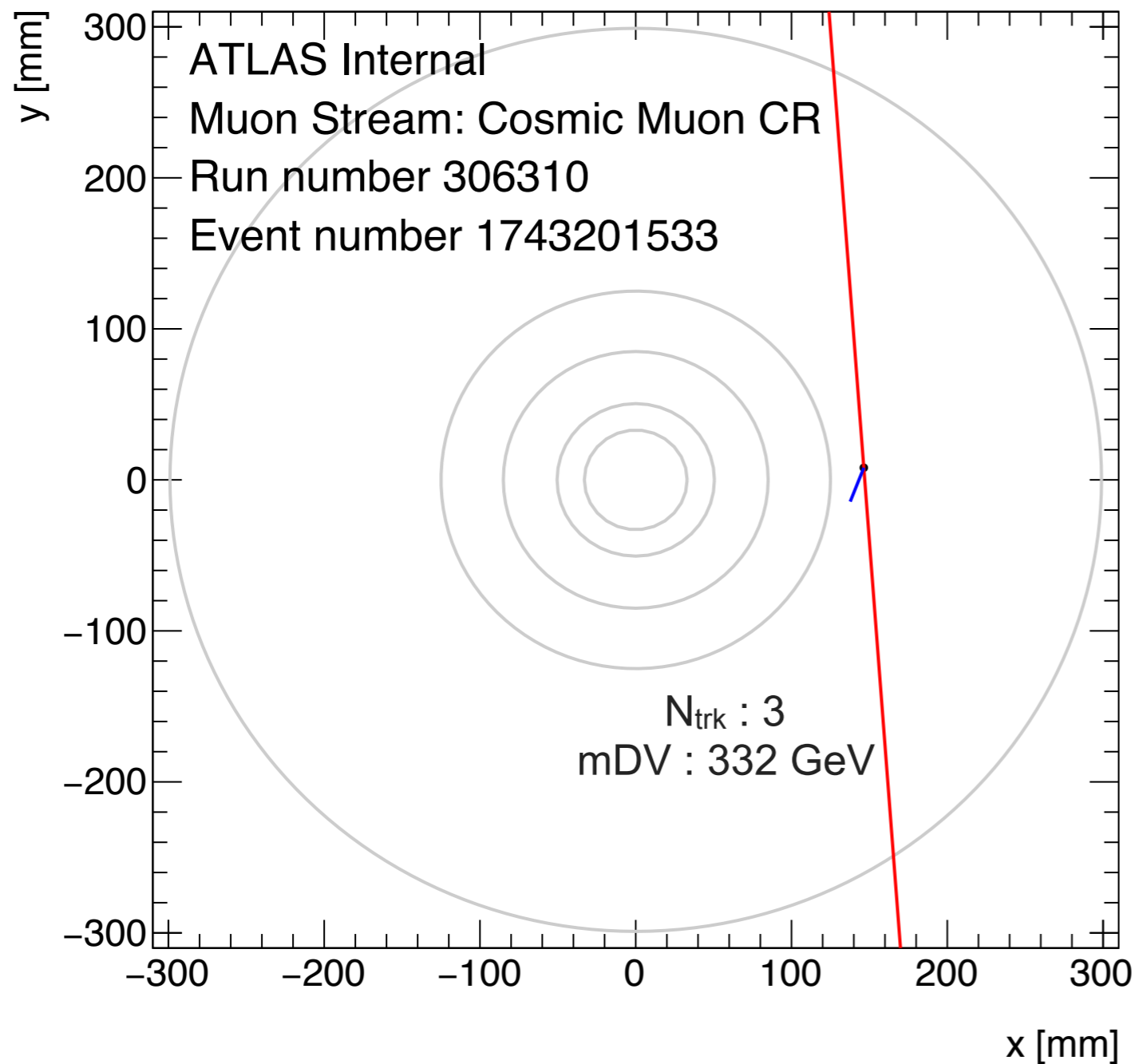


A little more rare

in non-isolated control region
material interactions from high
momentum jets

low mass vertices, but with
high momentum and highly
collimated tracks

muons sometimes reconstructed
as part of the vertex

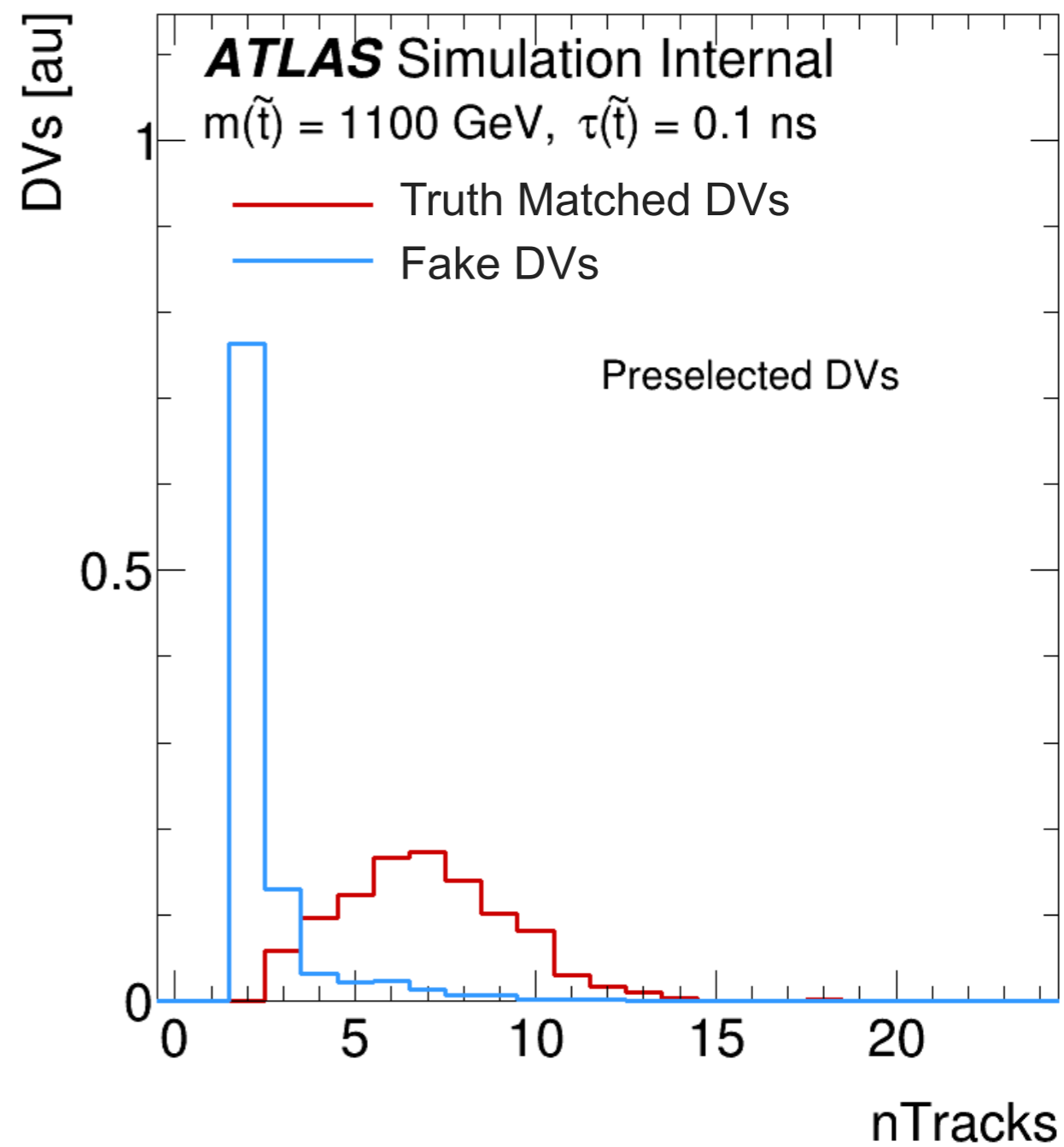
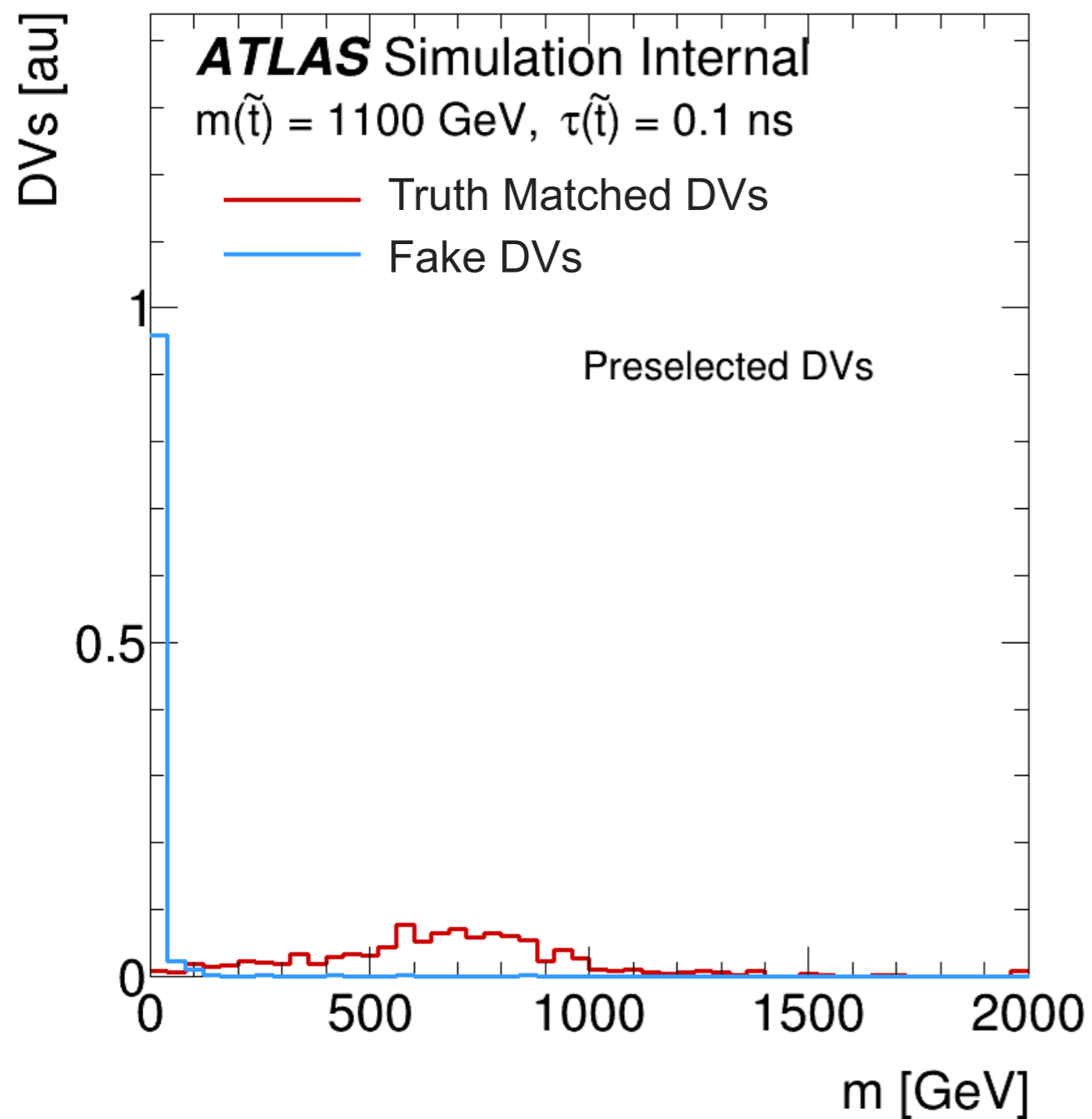


Cosmic vertices

most often two back-to-back
high momentum tracks from
cosmic

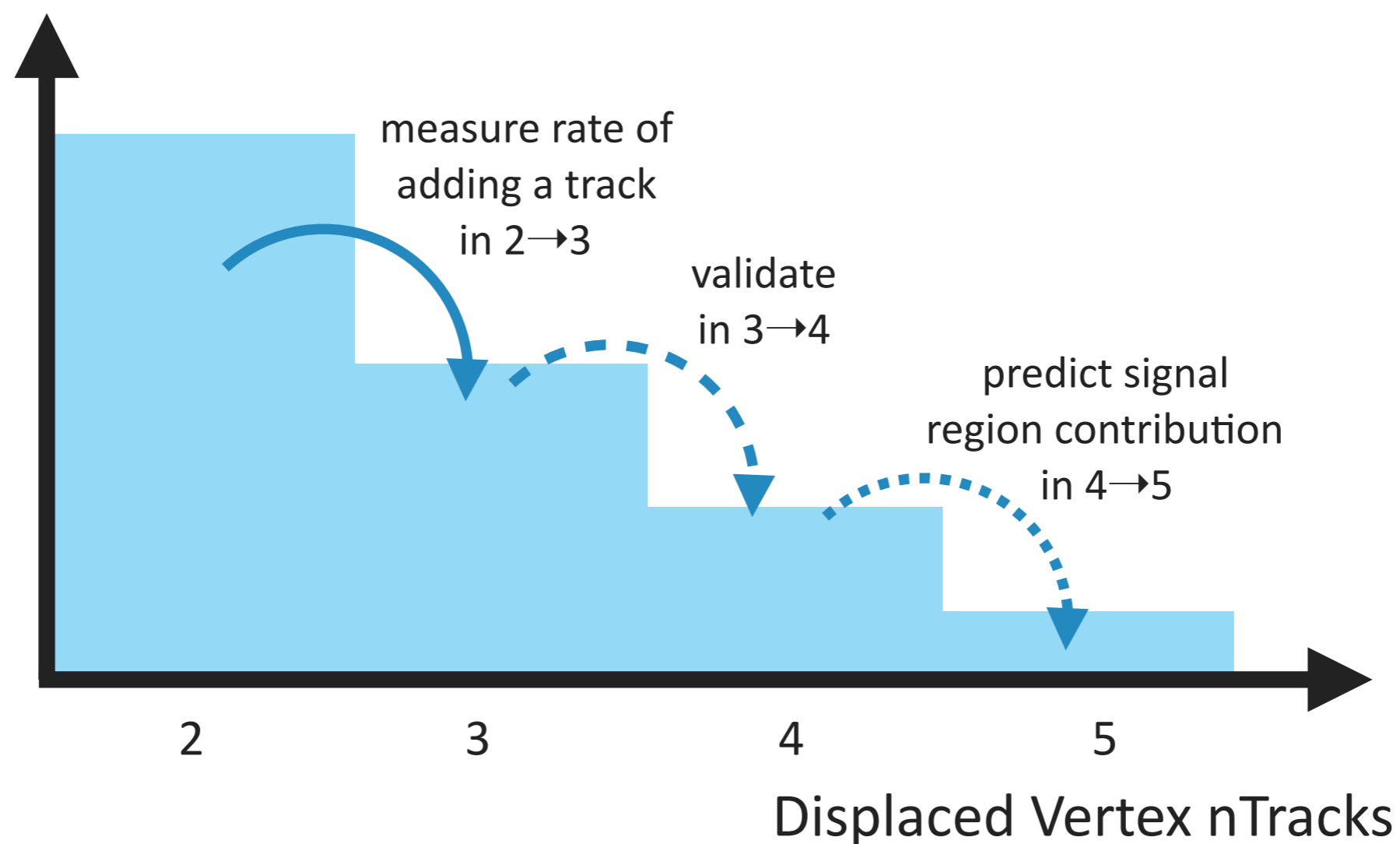
sometimes an additional
crossing 3rd track - very rare!

form very high mass
displaced vertices



In Run 1

backgrounds were estimated using displaced vertex properties
eg. random track crossings



requires cutting on DV nTracks ≥ 5

General Idea:

use both muon and displaced vertex properties of event, in “ABCD” inspired method

eg. for cosmic background at right

Very different from run 1!
Means we don't need to require DVs with ≥ 5 tracks

Displaced Vertex Properties

high mass and nTrack

low mass

low nTrack

0 preselected DVs

C → D

validation region

validation region

A → B

fail cosmic veto

pass cosmic veto

Muon Properties

Need to define two pure regions to measure transfer factors

impact parameter is best way to separate fakes and decays in flight

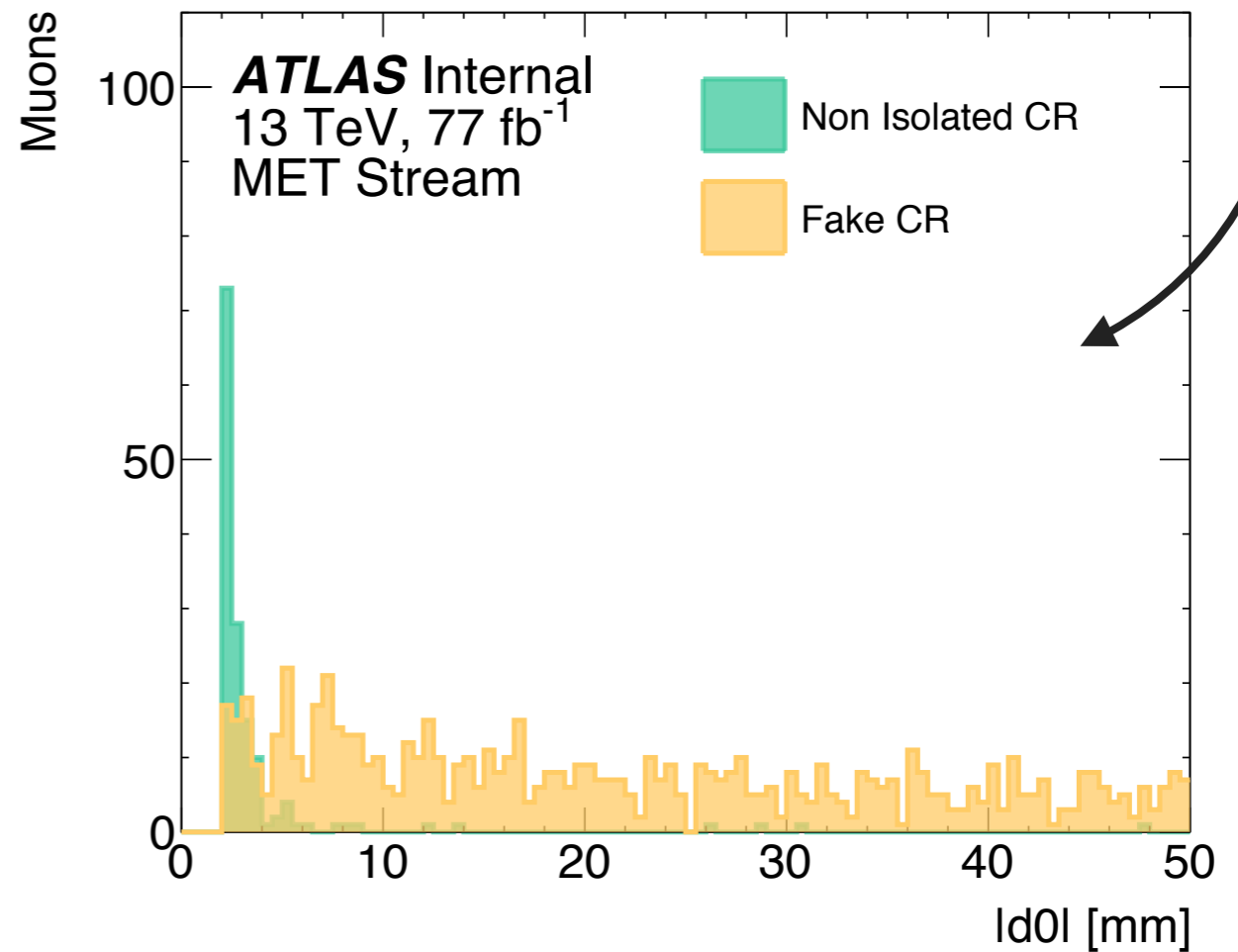
Decays in Flight

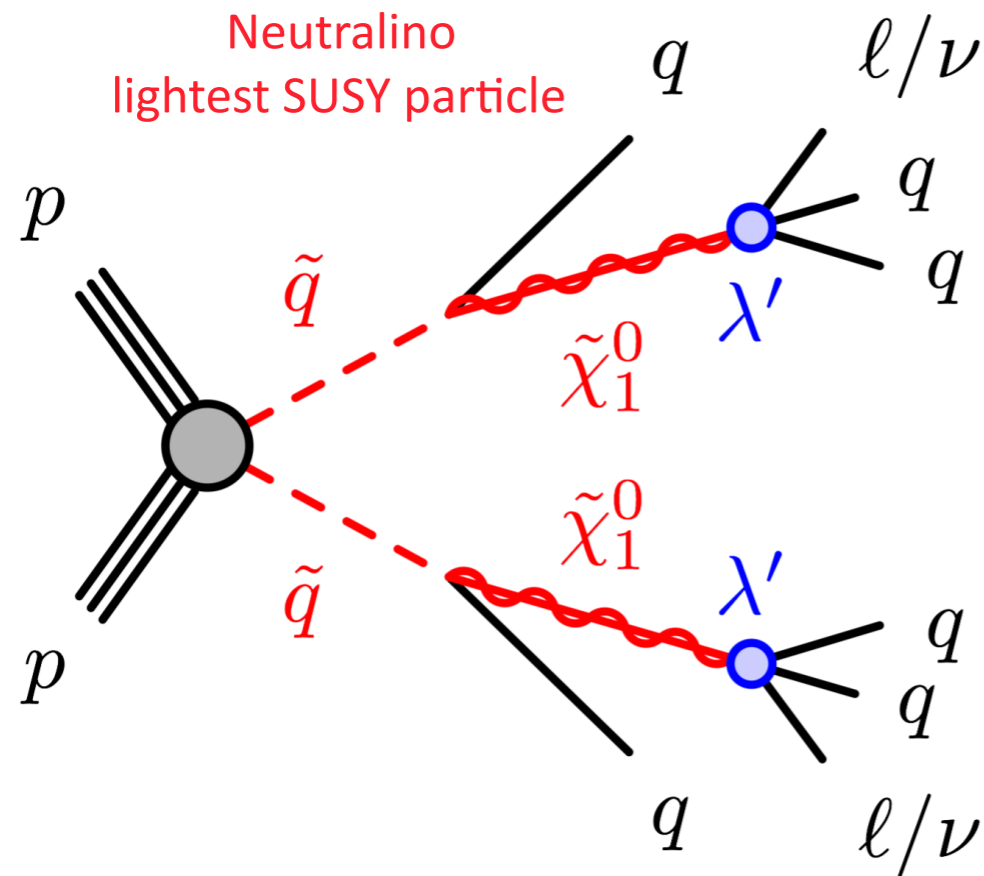
measure using muons with
 $1.5 < |d_0| < 2.0$ mm

Fake Muons

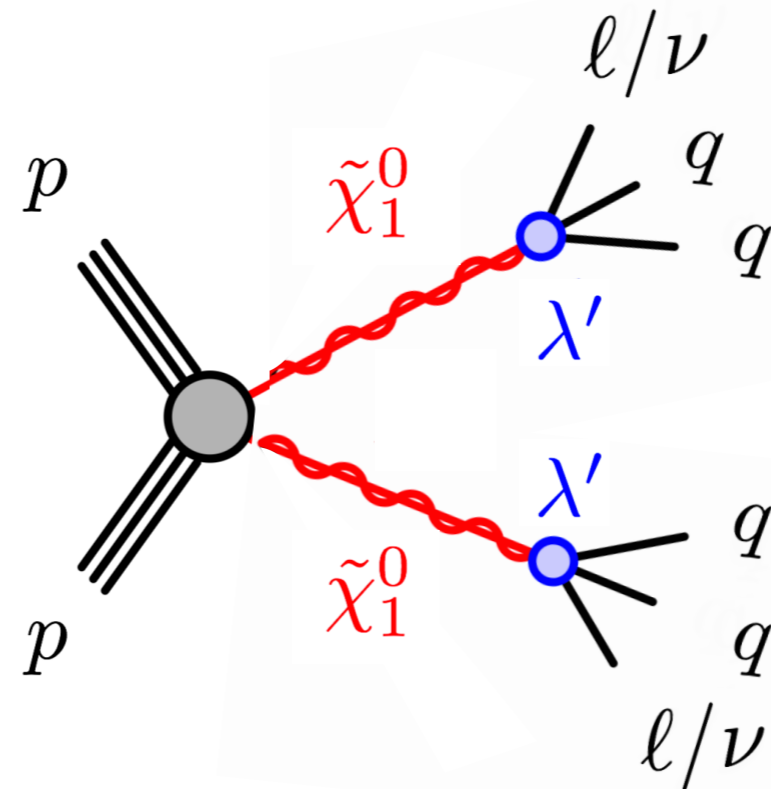
measure using muons with
 $|d_0| > 5$ mm

Challenging because we're limited
by statistics for transfer factor
numerators





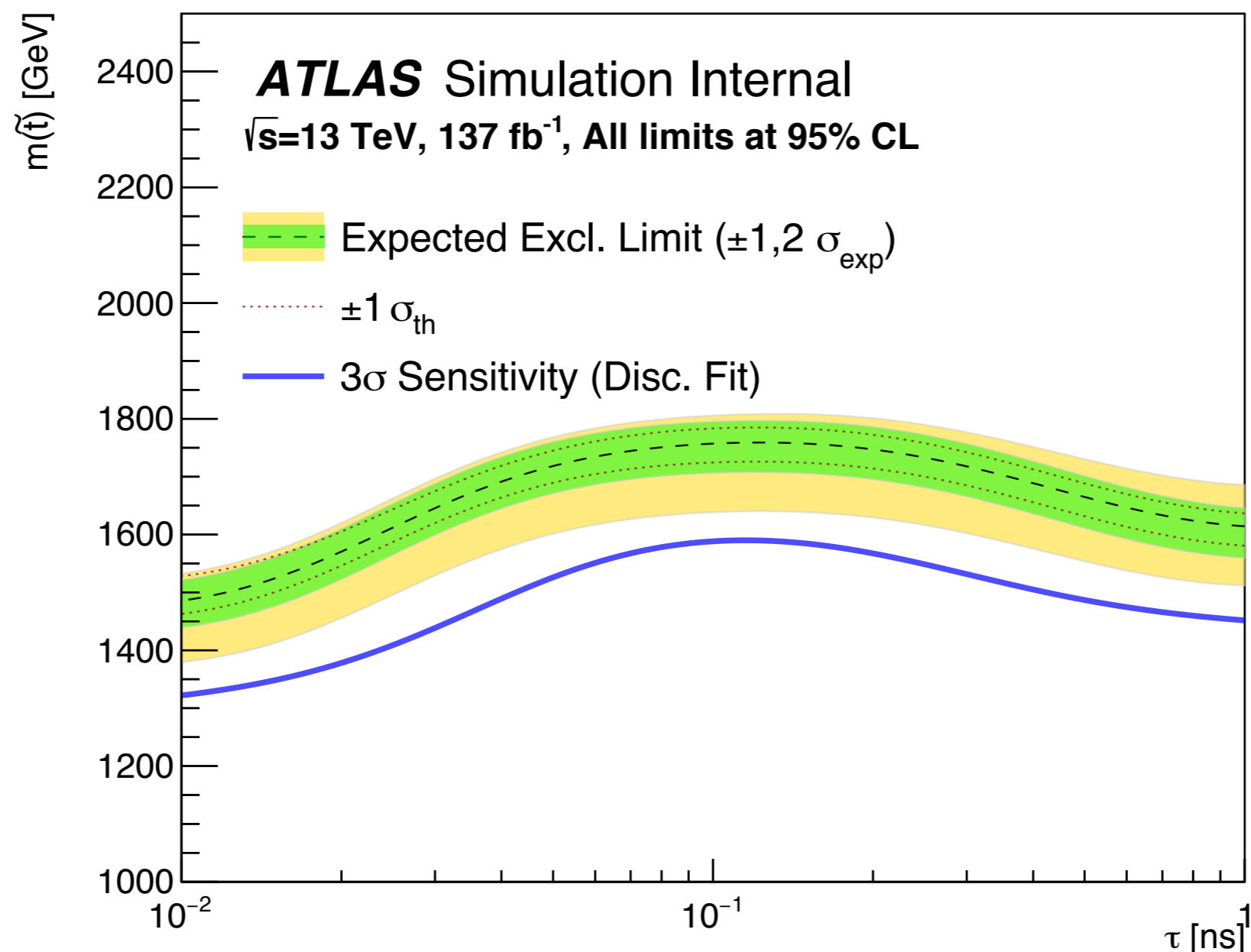
Retain sensitivity to other signals
particularly muons with lower pT



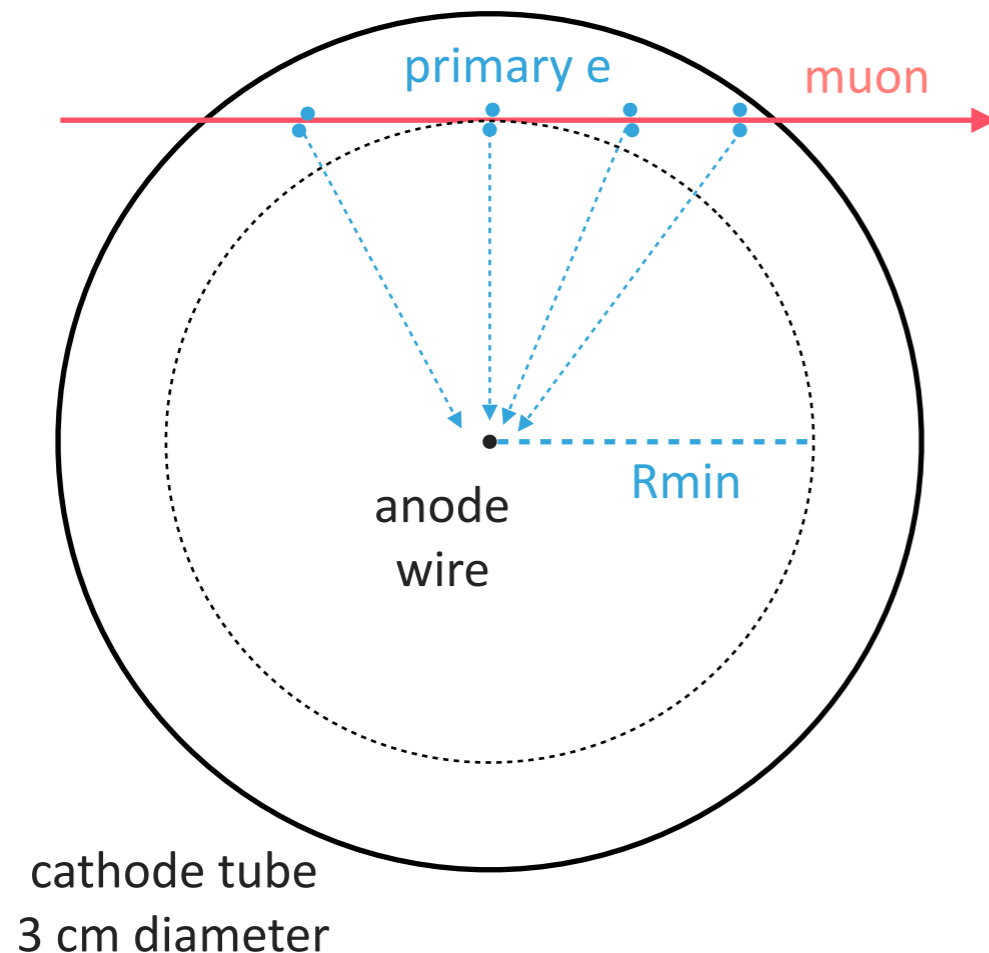
Expected sensitivity for different stop masses and lifetimes

with 2016-2017 data *+2018 data*

if we scale signal and background by luminosity



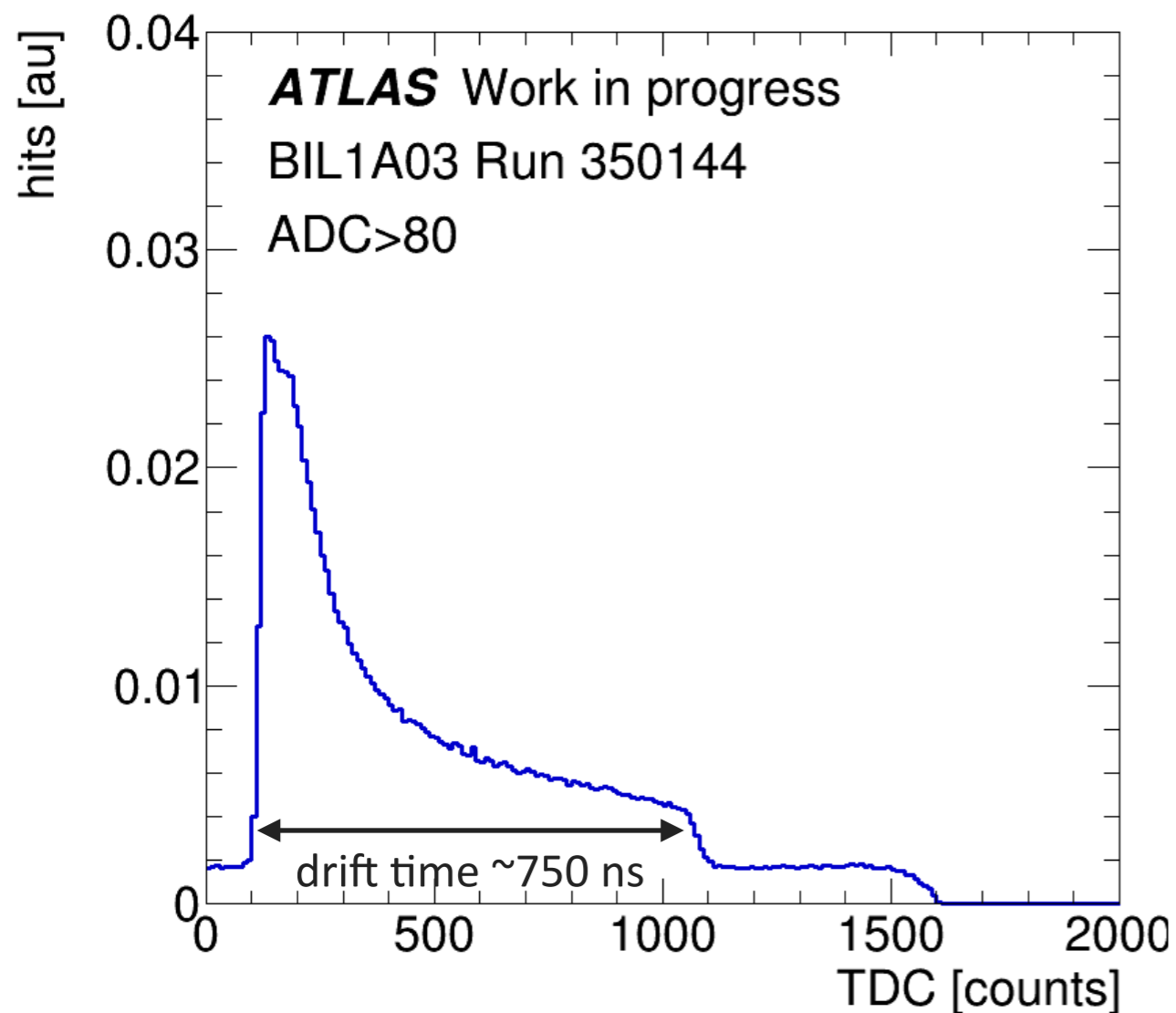
Drift Tube Reminder



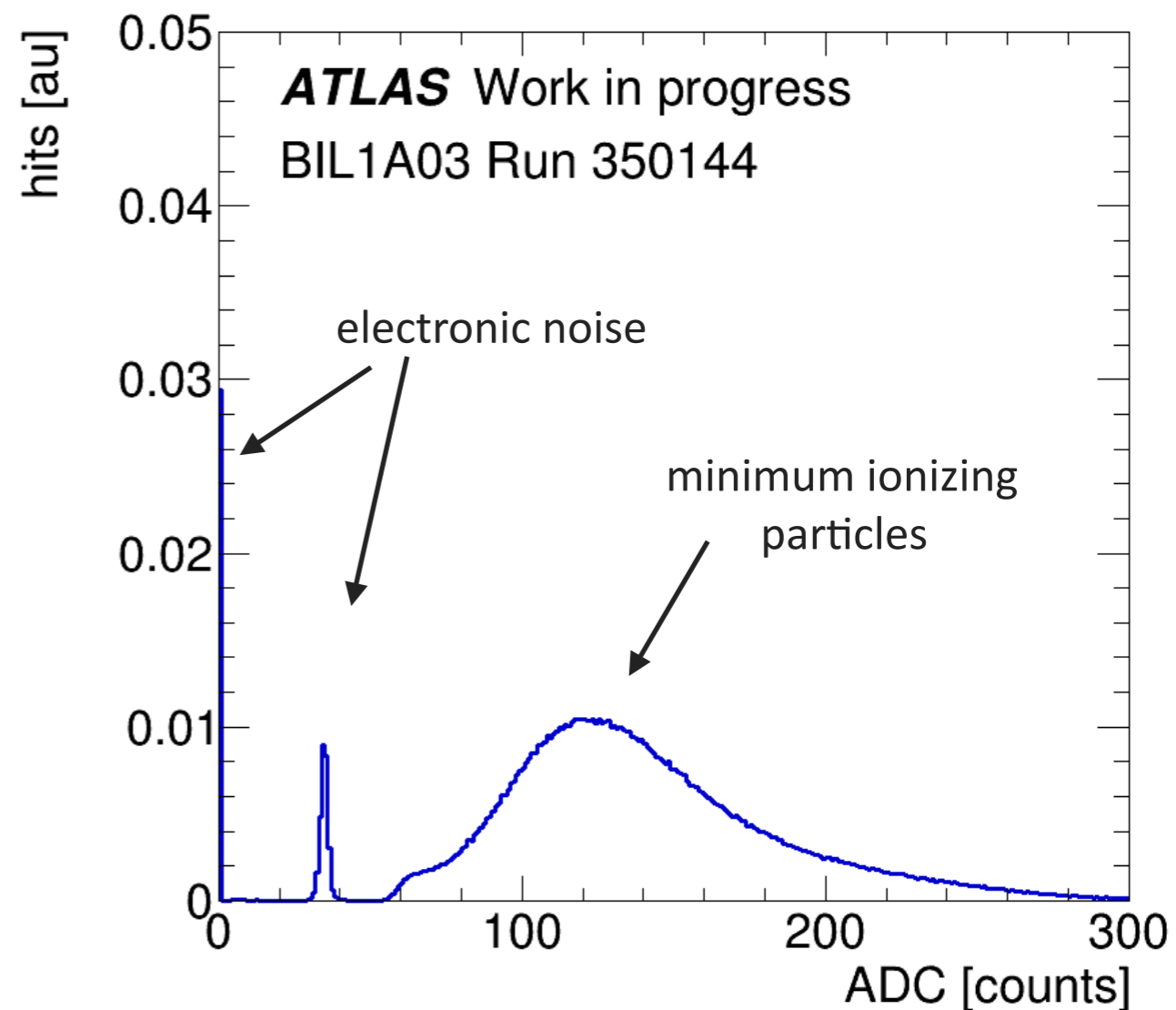
MDTs have a maximum drift time of 750 ns

dead time: if two muons pass through the same tube in <750 ns of each other you MISS the second muon

TDC = time of leading edge
~arrival time

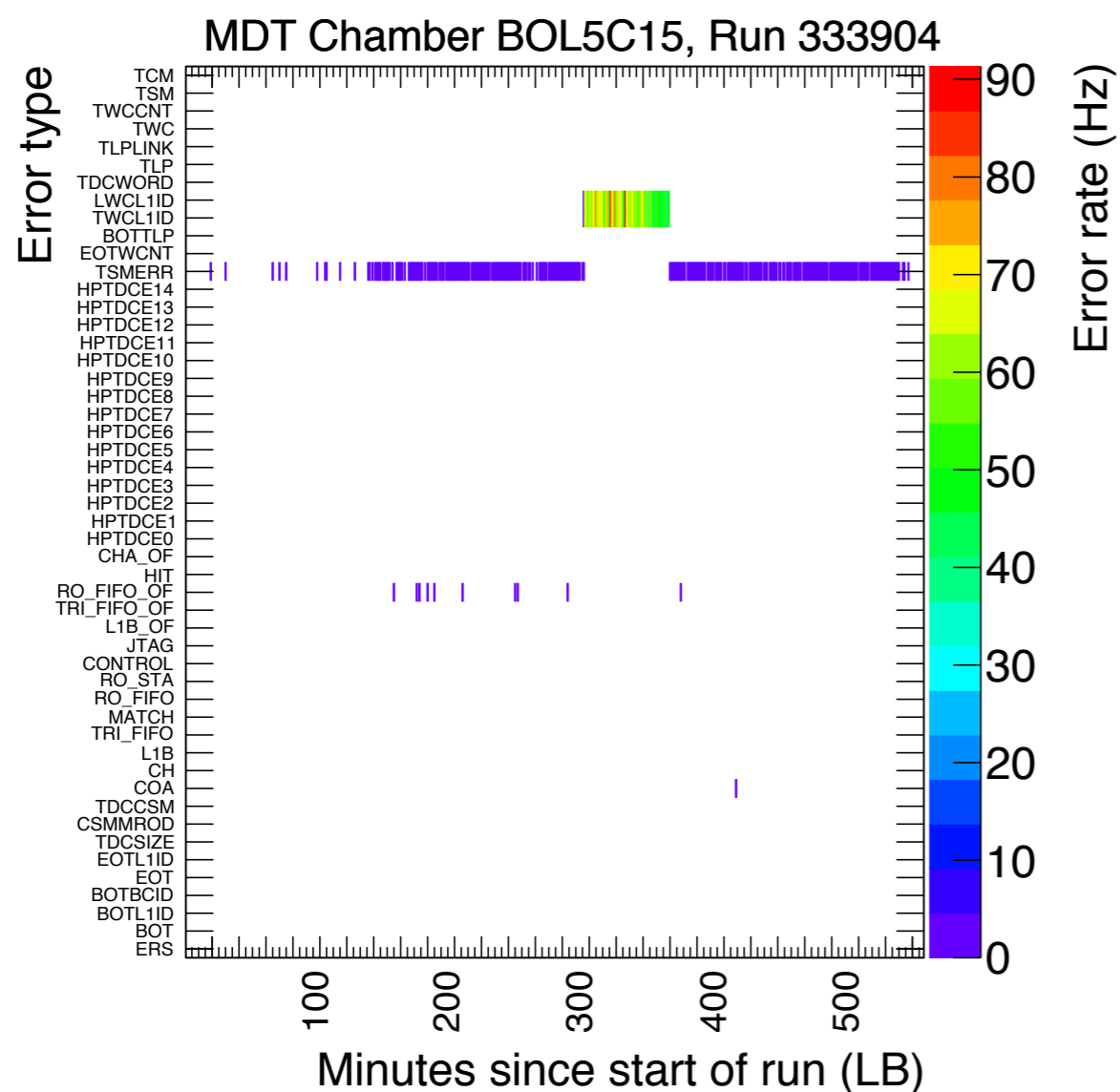


ADC = leading edge - trailing edge
~deposited charge



from calibration stream: hits on track and inside L1 ROIs

also receive electronic error information
useful for determining when to intervene during a run



eg. TWCL1ID, LWCL1ID

MDT hits are present, but timing offset
with respect to ATLAS

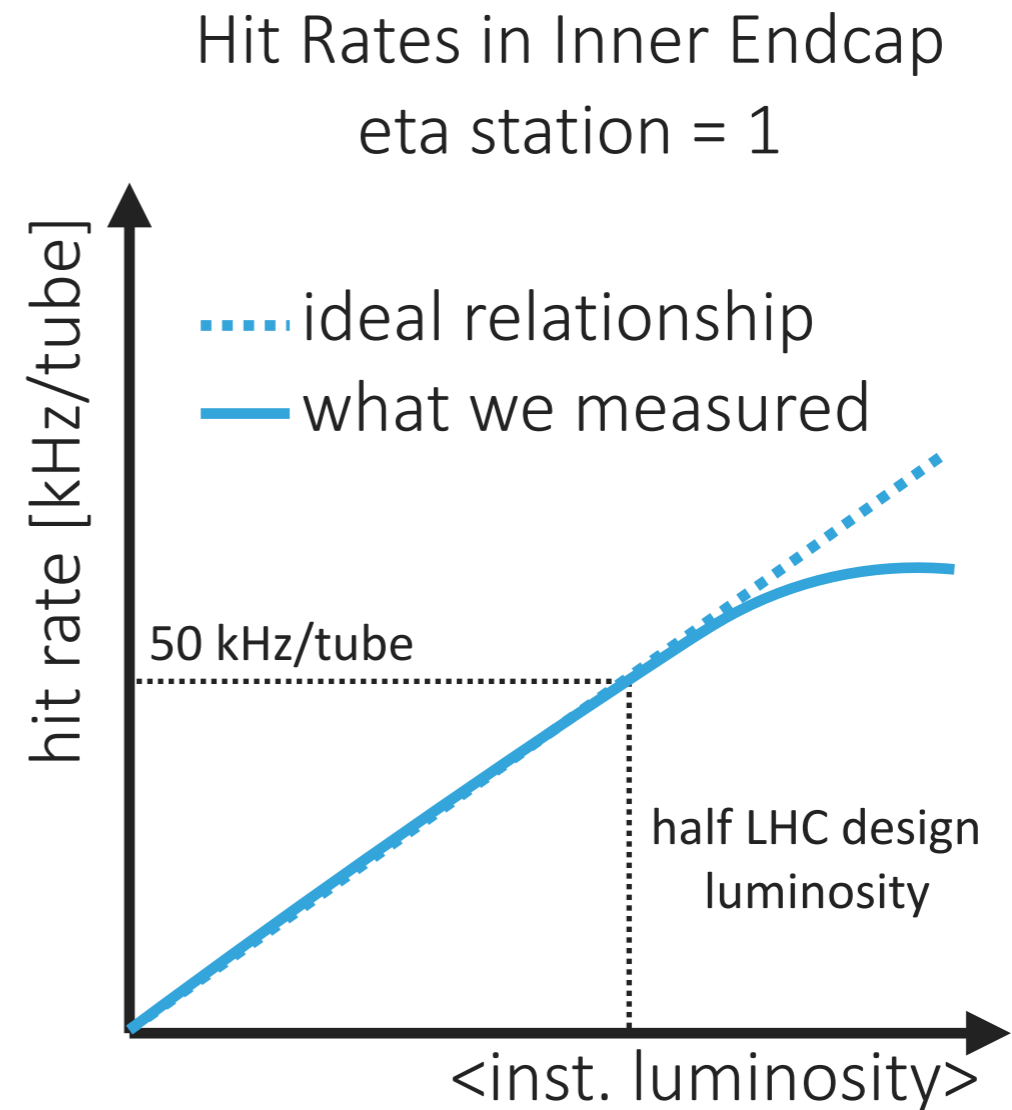
hits either aren't reconstructed as part
of muon or we lose muon altogether

usually fixed by chamber re-initialization

MDT readout saturation in the Inner Endcap

data loss began at hit rates much lower
than expected from dead time alone
~50 kHz/tube versus ~300 kHz/tube

The problem: the electronics,
not the detector



*not a real plot!

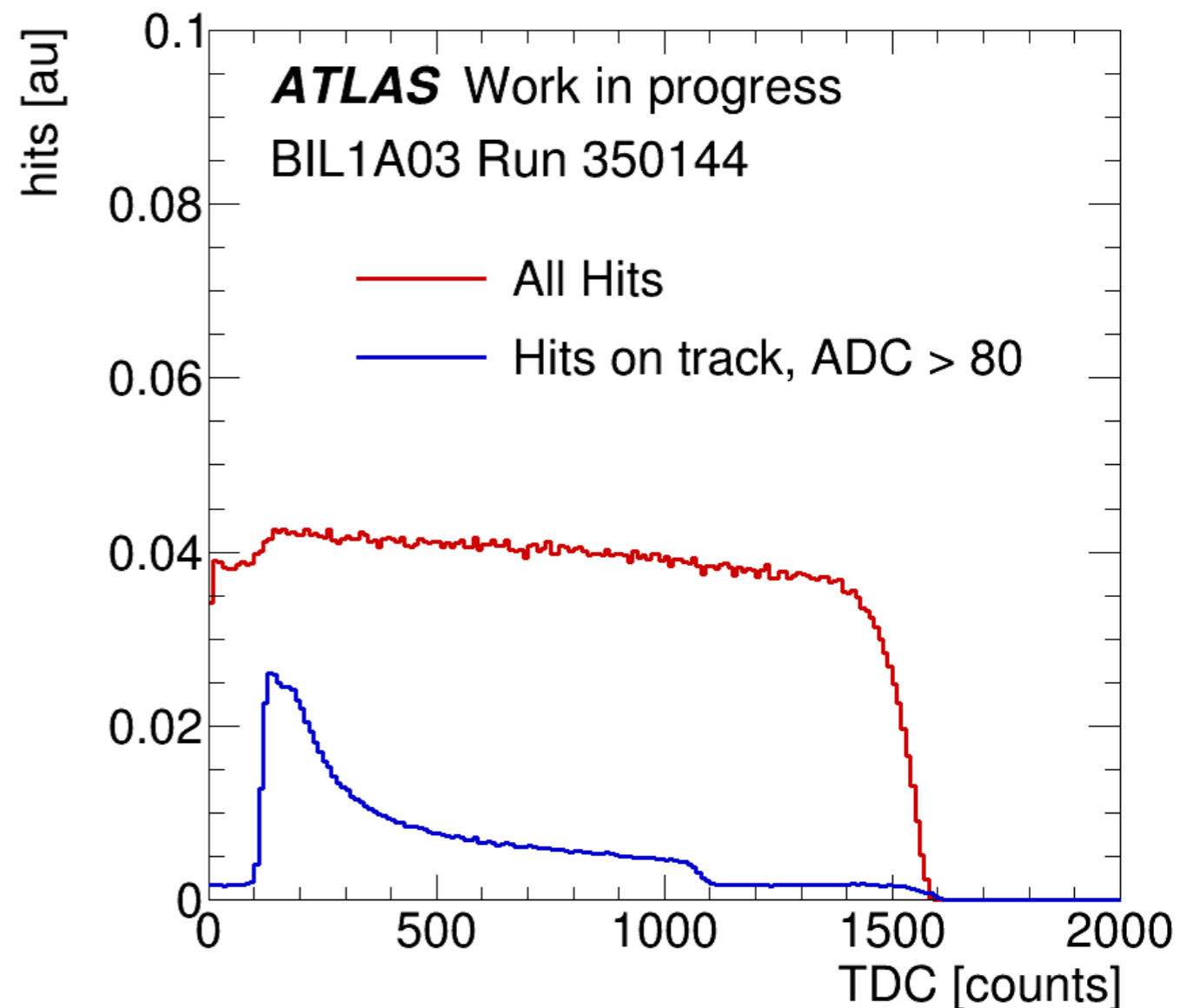
Solution 1:

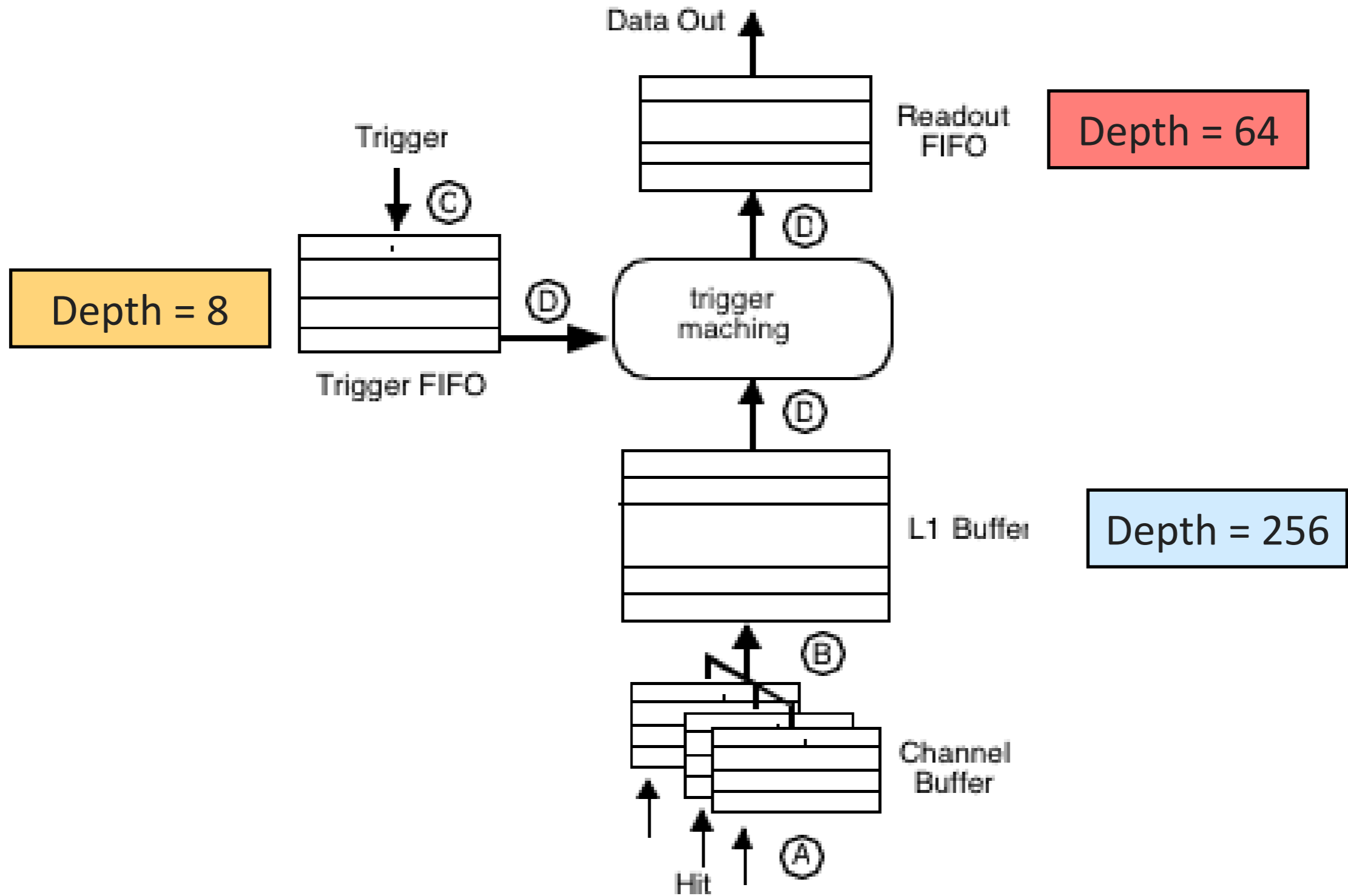
Reduce MDT readout window
from 2500 ns to 1300 ns

no loss of efficiency for **REAL MUONS**
but reduces **EVENT DATA RATE**
by >40%

Solution 2:

change electronics configuration,
flush readout FIFOs when nearly full

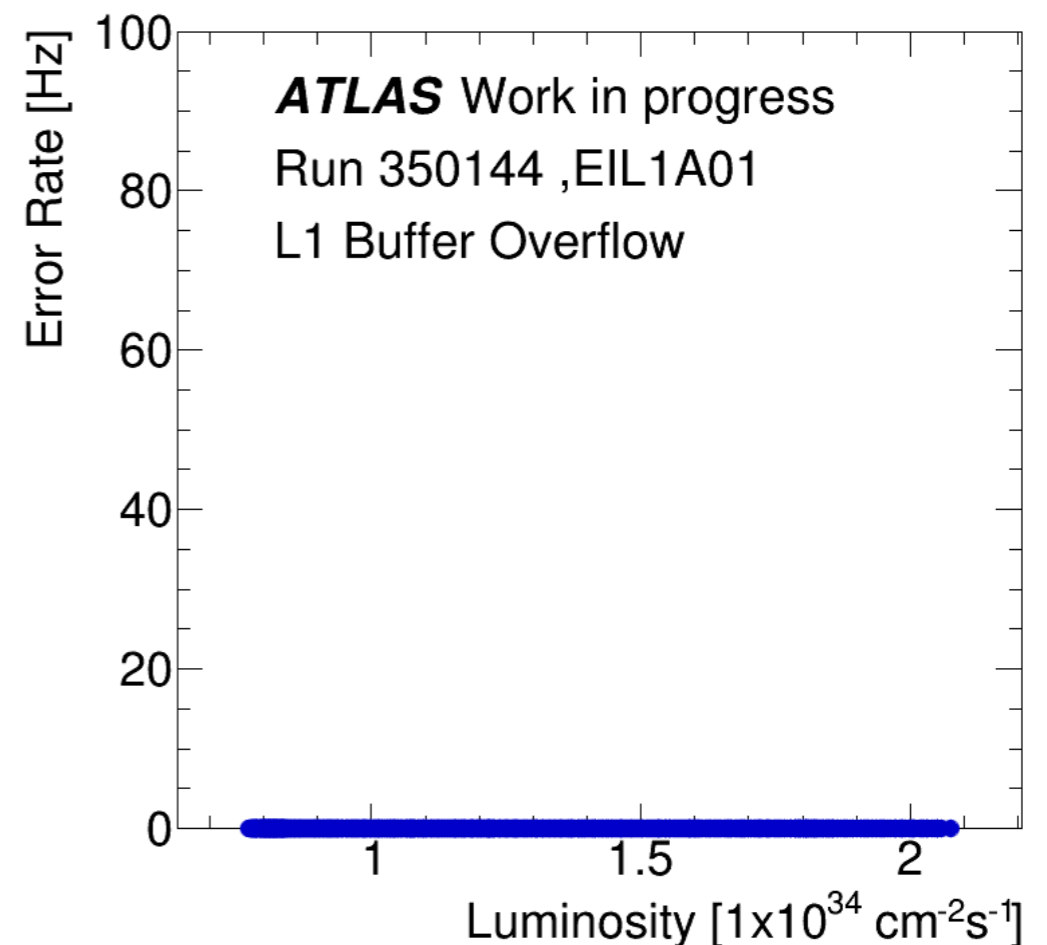
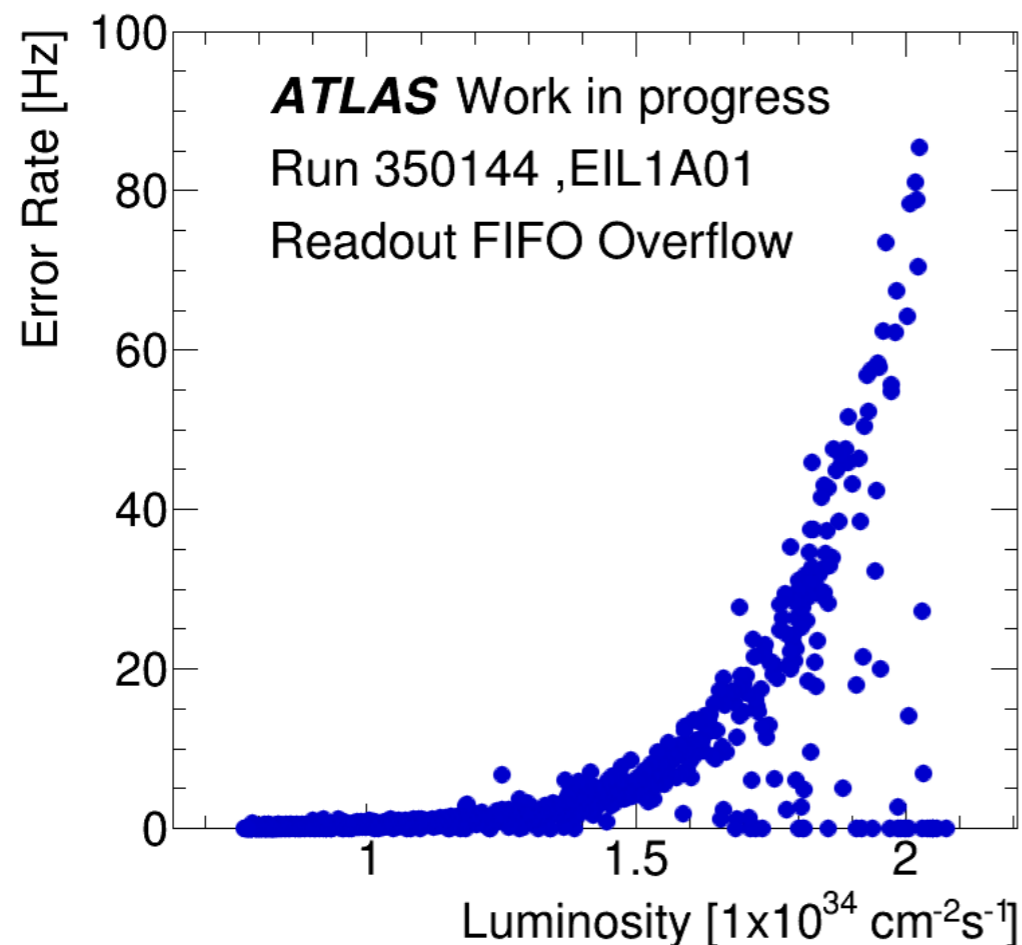


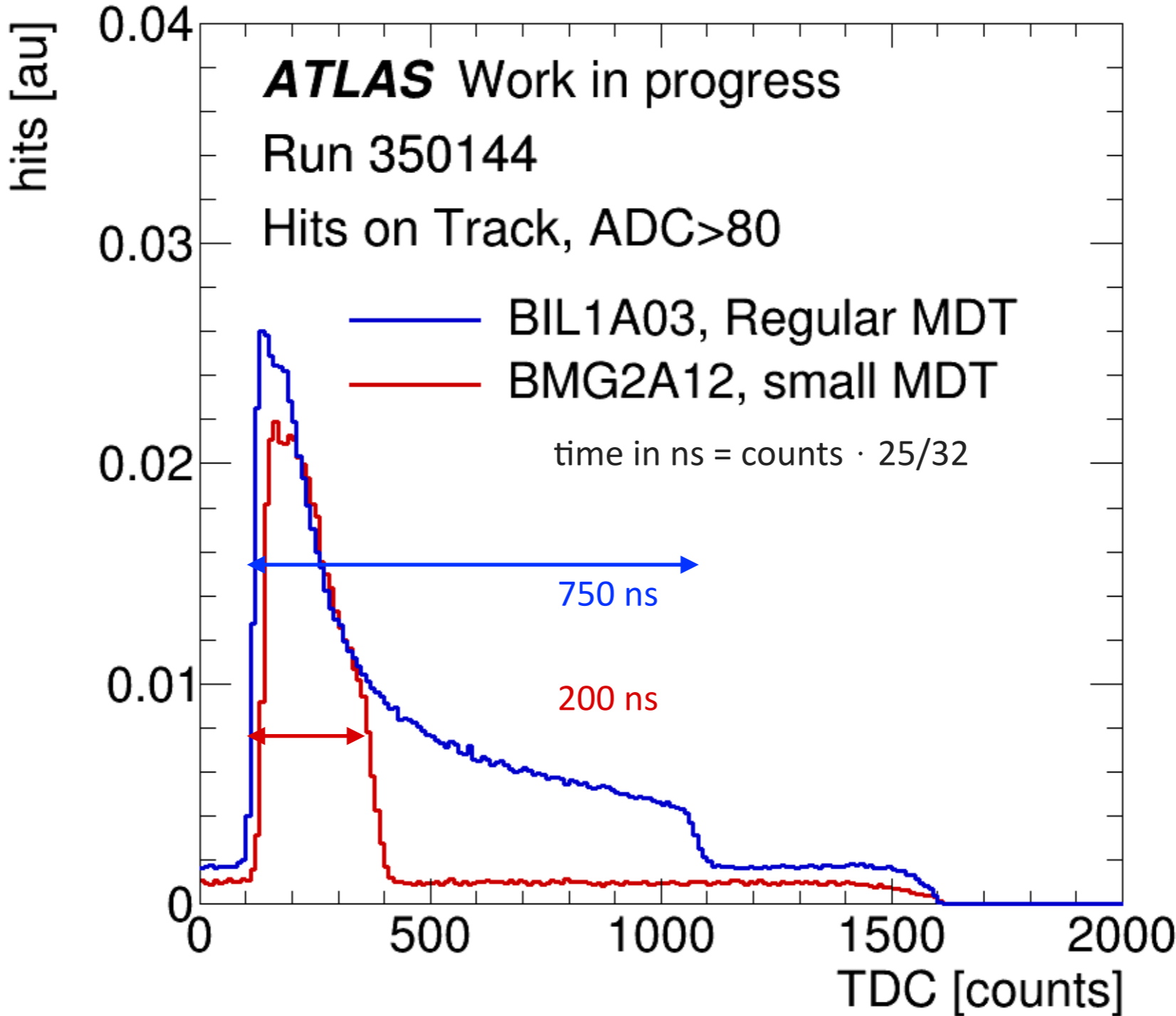


- Run 1 Readout FIFO would overflow, backing up hit buffers takes long time to clear buffers
- Run 2: clear readout FIFOs when nearly full, prevent hit buffers from ever overflowing

Assumptions

1. readout buffer overflow is RARE
2. clearing buffers has negligible impact on real muons





Accounting for MDT deadtime in hit rate measurement

observed hit rate = delivered hit rate · single hit efficiency

$$\text{delivered hit rate} = A + \mathcal{L} \cdot B$$

$$\text{single hit efficiency} = 1 - \text{delivered hit rate} \cdot C$$

gives

$$\text{observed hit rate} = (1-AC)A + (1-2AC)B\mathcal{L} - C B^2 \mathcal{L}^2$$

140 ns dead time

5.6 mm η pitch, 13 mm ϕ pitch

1 strip = 5 mm * 1 m

Loss of hits with a rate of

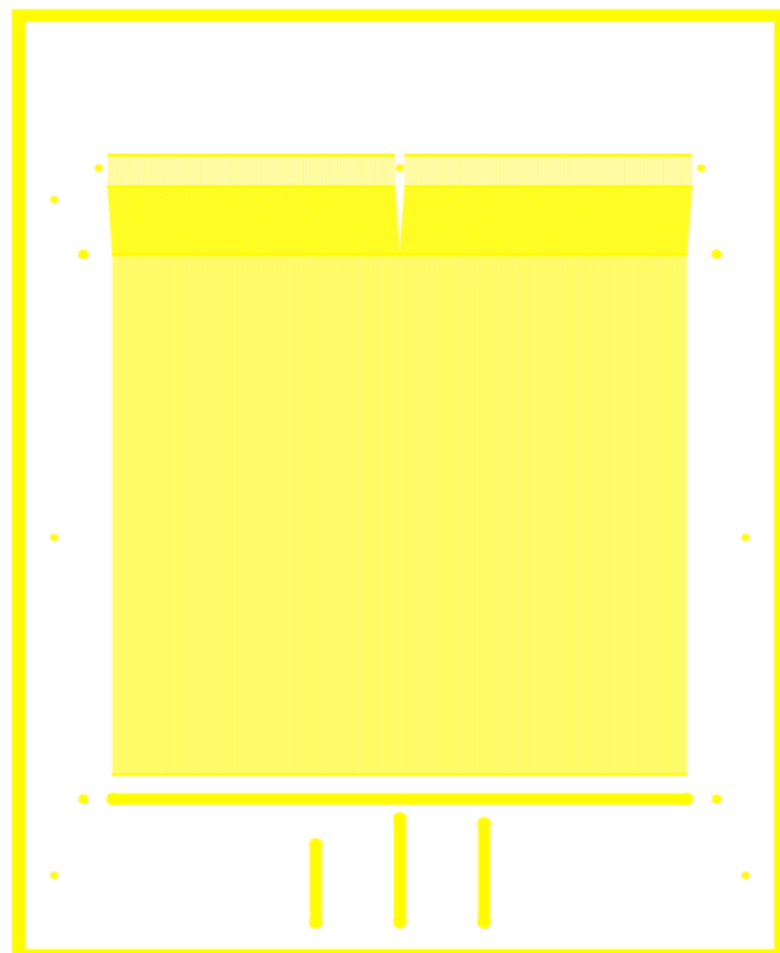
$\sim 1 / 140$ ns

~ 7 MHz / strip

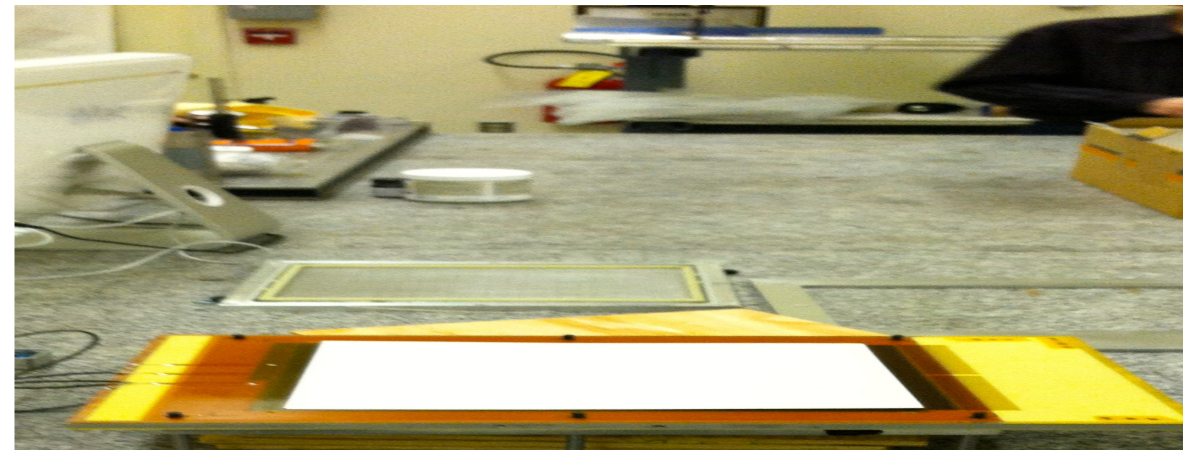
~ 140 kHz/cm²

Readout Strip Design

PCBs designed at Harvard



Sent to CERN for manufacture
resistive strips, pillar, and mesh deposition
final board layout below



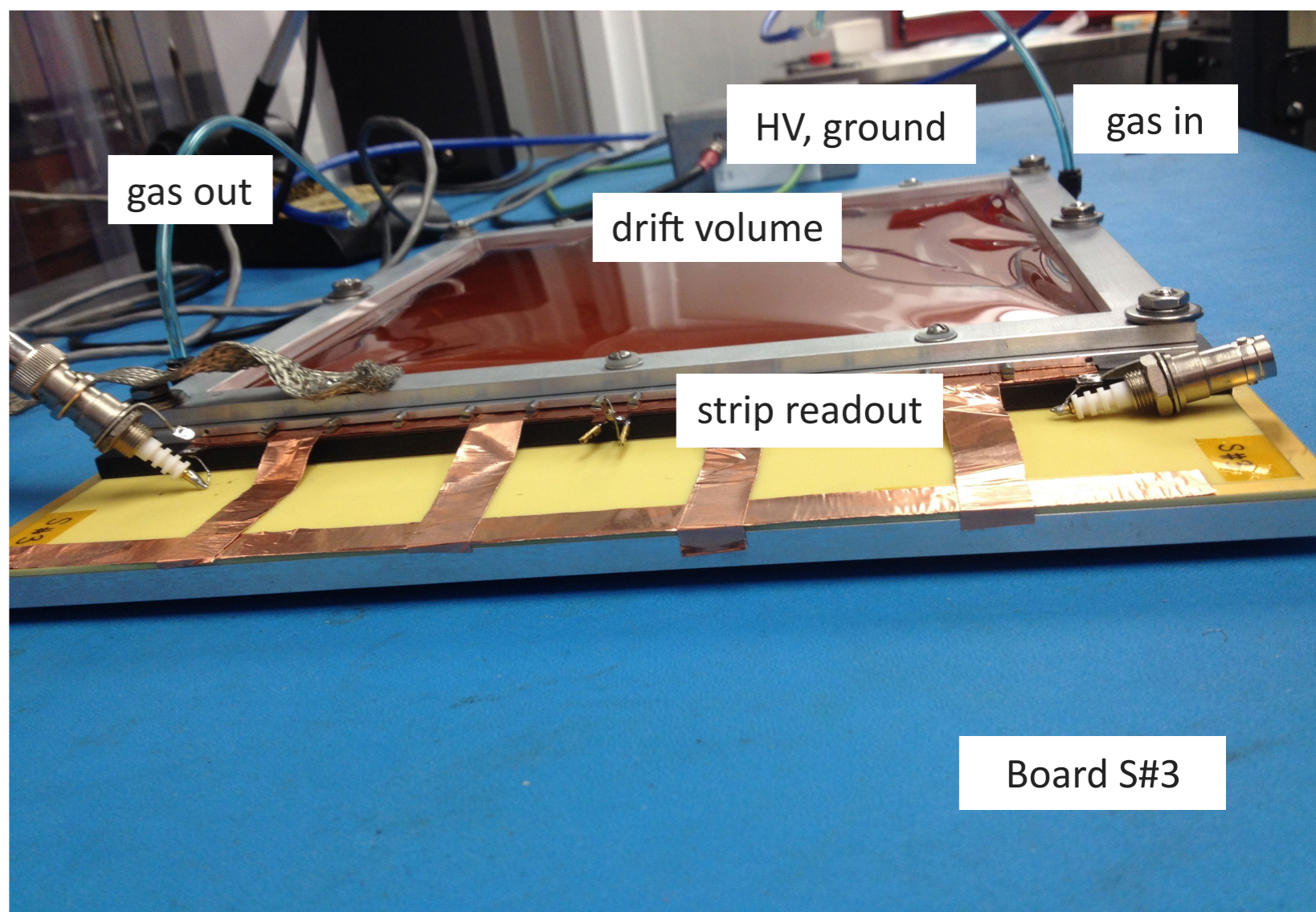
Setup for testing individual boards

Each Board tested individually before assembly

Breakdown voltage in air and in Ar:CO₂

Detector Gain

Shorted Zebra connector to readout all strips at once w/ a single charge amplifier

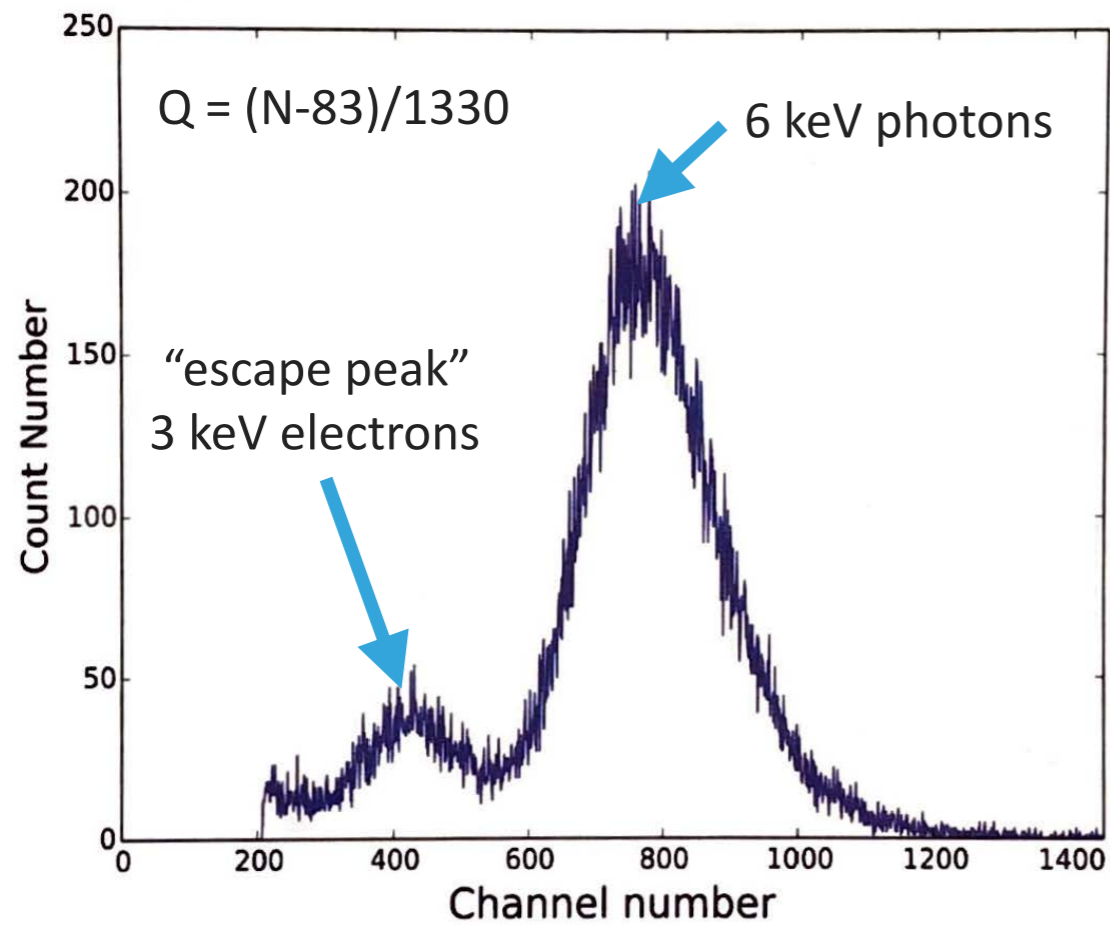


Measuring breakdown voltage in Air & Ar/CO2

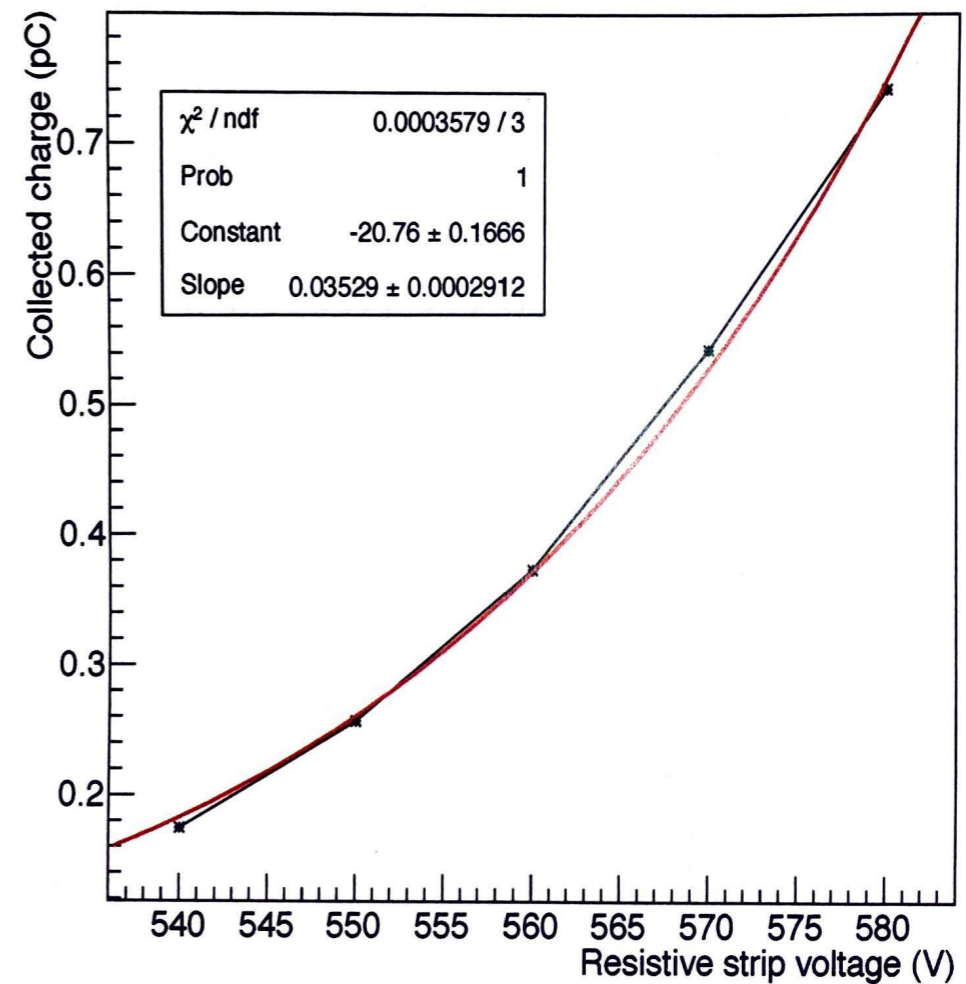
define breakdown if drawing currents more than ~ 10 nA
or 960 V in Air
or 580 V in Ar/CO2

Board	S#1	S#2	S#3	S#4	A#1	A#2	A#3	A#4	A#5
Air	960	950	960	960	880	960	960	960	960
Ar/CO2	580	580	580	580	550	580	580	580	580

Measuring chamber gain v. strip voltage using Fe-55 6 keV photon peak



Collected Charge v. Application Voltage

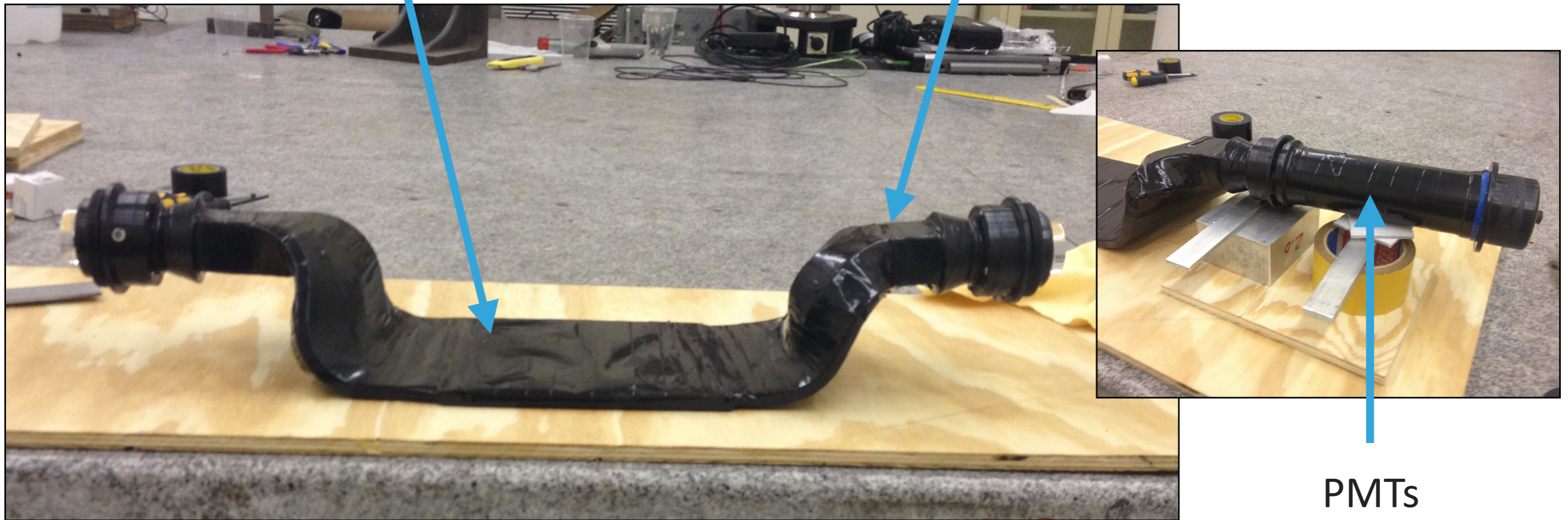


Scintillator

polyvinyl toluene (PVT)
+ “pop-pop” wavelength shifter

plexiglass light guides

glued to scintillator. with epoxy, and PMT
adapter with liquid plexiglass + accelerator



PMTs

recycled from CDF
hadronic calorimeter