

Origin of Hydrodynamic Tunneling

By

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- **During the past two decades, significant progress has been made in the technology of particle accelerators.**
- **This has provided intense particle beams of highly energetic ions from protons up to uranium.**
- **Due to the high intensity, if the beam interacts with matter, substantial amount of specific energy is deposited in the target that leads to high pressures and phase changes. The high pressure generates shocks which lead to hydrodynamic effects.**
- **Hydrodynamic tunneling is one of the consequences of the above effects and will be explained in this talk.**

Specific power deposition:

$$P = \frac{E_s}{\tau}, \quad \text{TW/g}$$

$$E_s = \frac{\frac{1}{\rho} \frac{dE}{dx} N}{\pi r_b^2}, \quad \text{kJ/g}$$

$\frac{1}{\rho} \frac{dE}{dx}$: specific energy loss due to a single ion

N : total number of particles in the beam

r_b : beam radius

To maximize **P**, an optimum set of parameters should be determined.

Beam-Matter Heating

Uranium Beam

- Particle Energy: 1 GeV/u
- Bunch Intensity: 10^{12} ions/bunch

Single Bunch

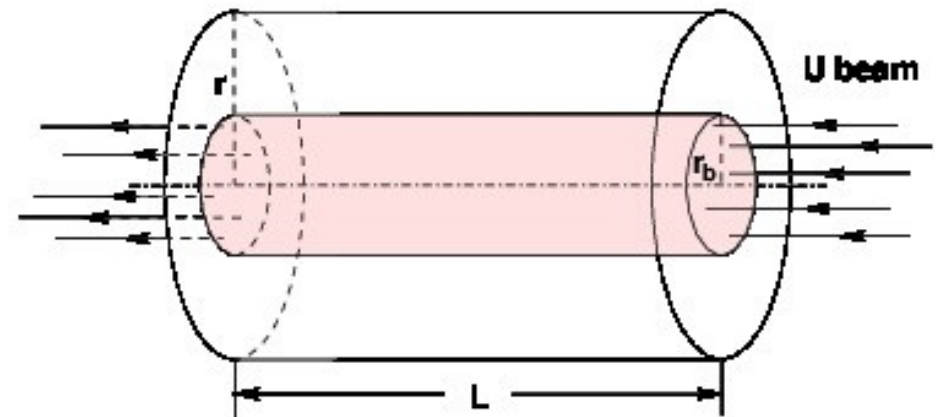
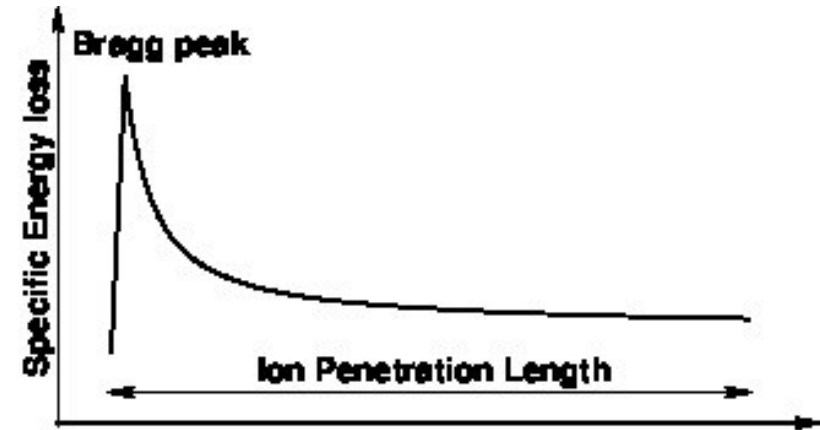
- Length: 50 ns and 1000 ns

Five Bunches

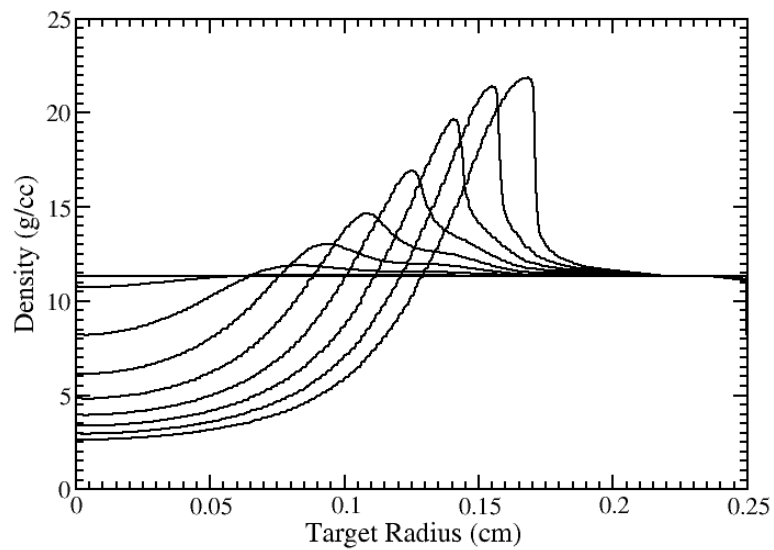
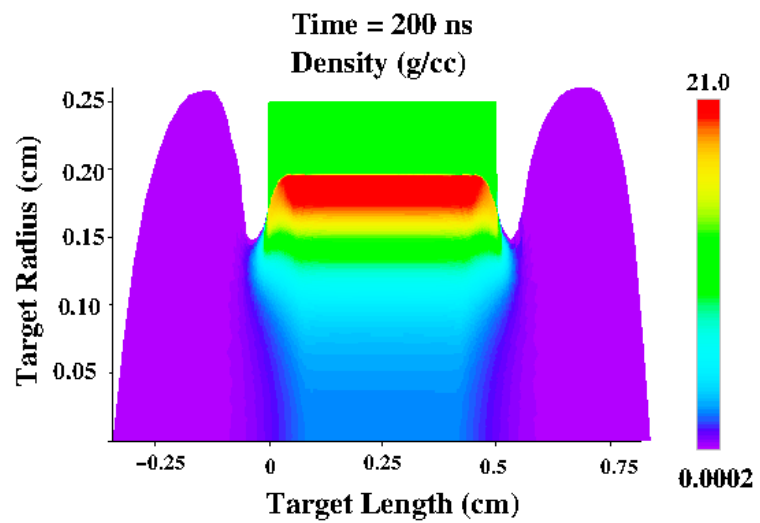
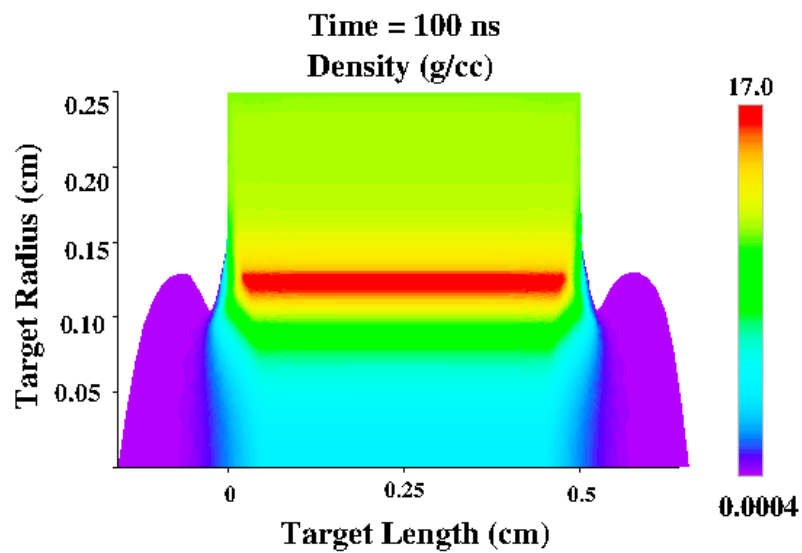
- Bunch length : 140 ns
- Bunch Separation: 140 ns

Target

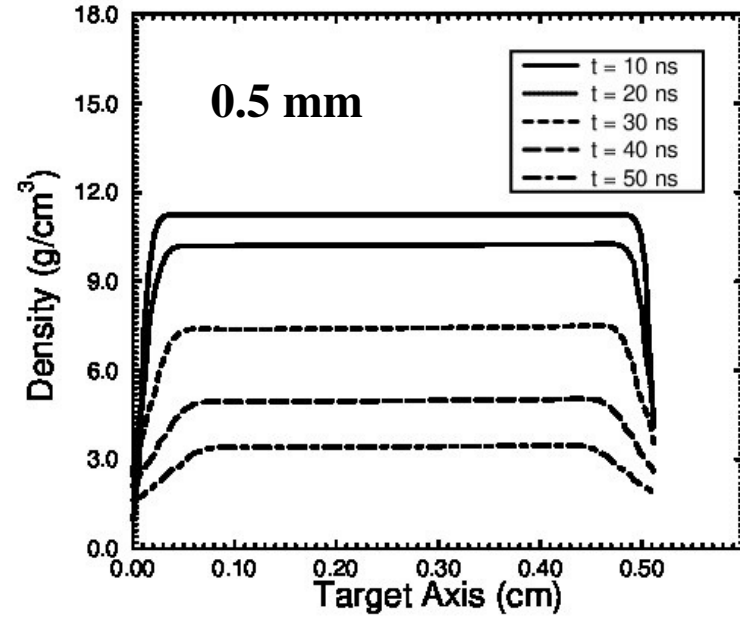
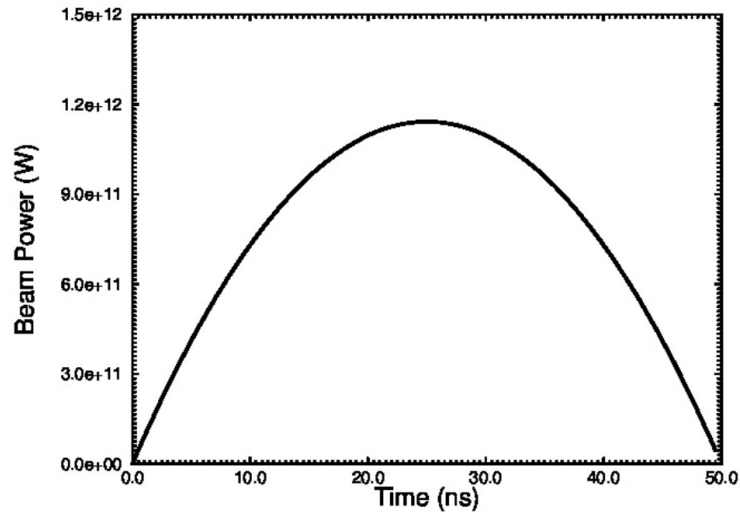
- Lead Cylinder
- Length = 5 mm
- Radius = 2.5 mm



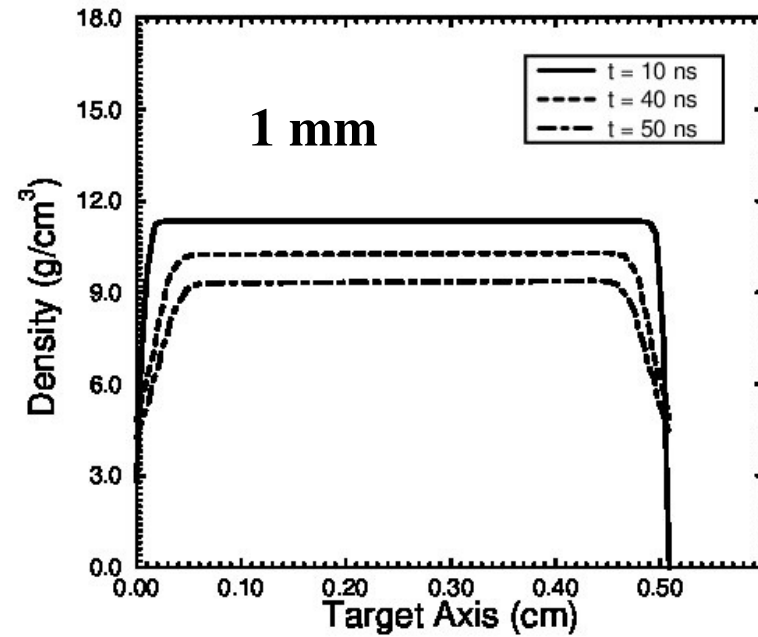
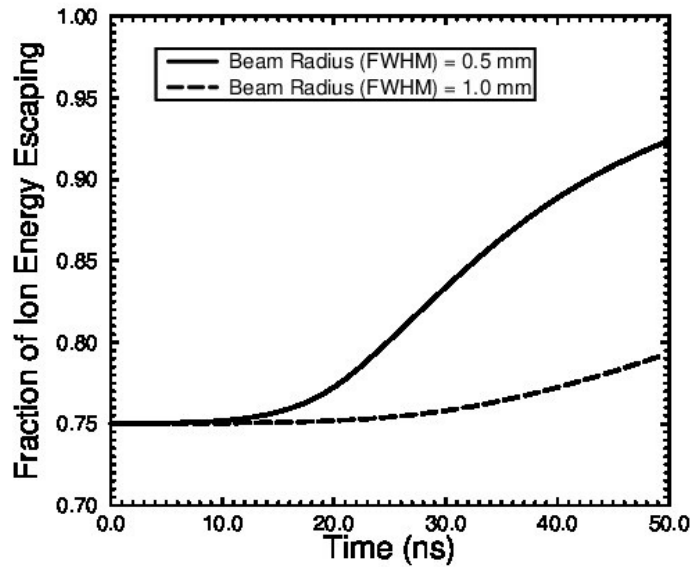
Simulations using 2D
Hydrodynamic code
BIG2 [*N.A. Tahir et al.,
PRE 63 (2001) 036407*]



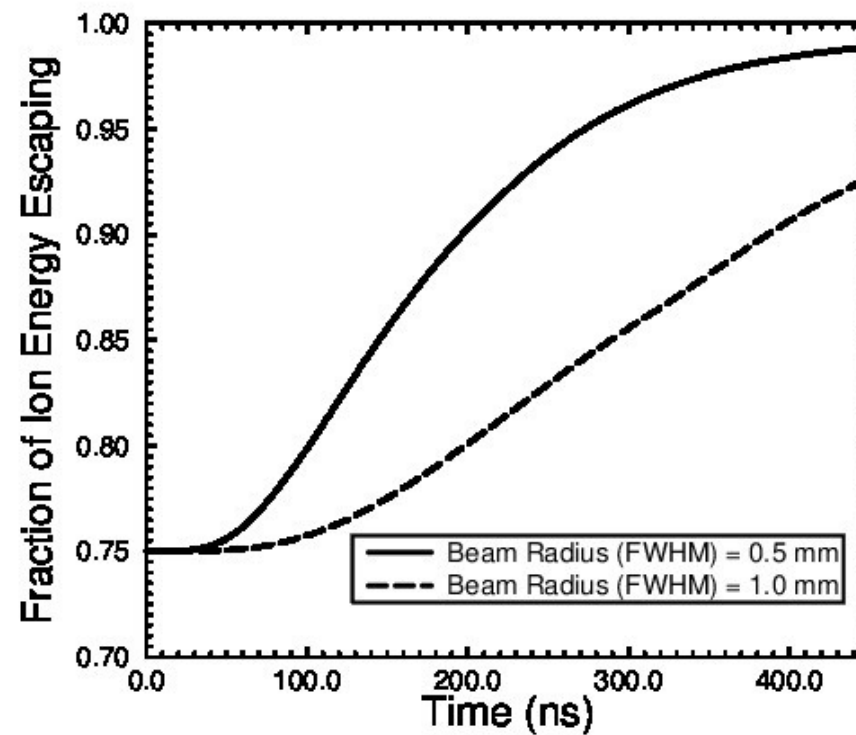
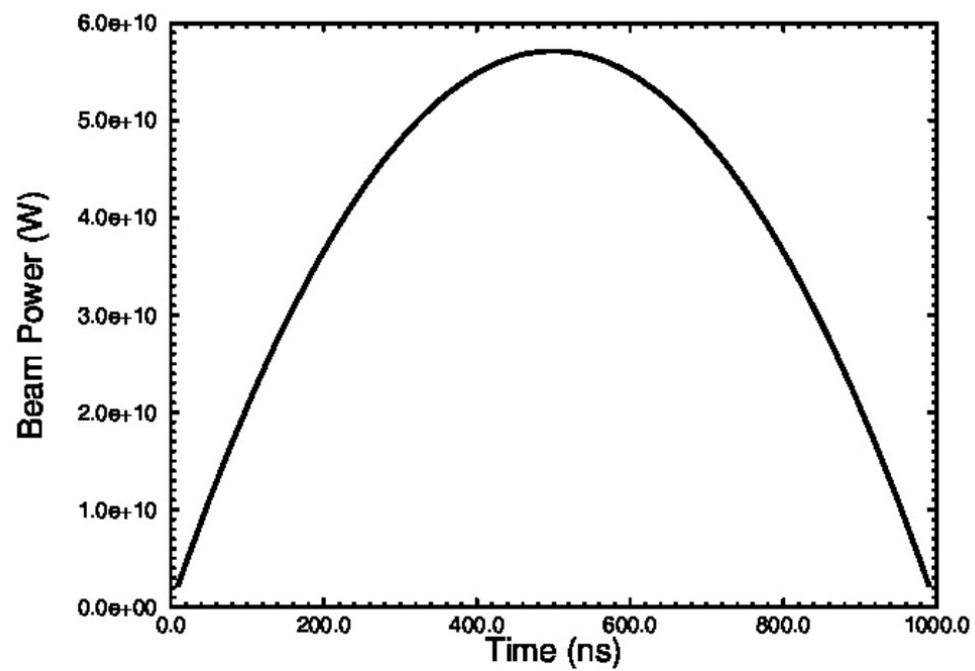
Single Bunch with a Length = 50 ns



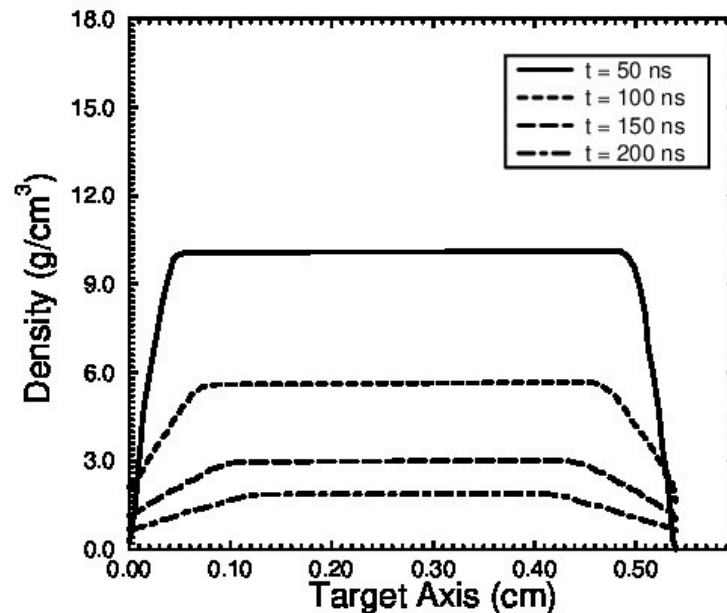
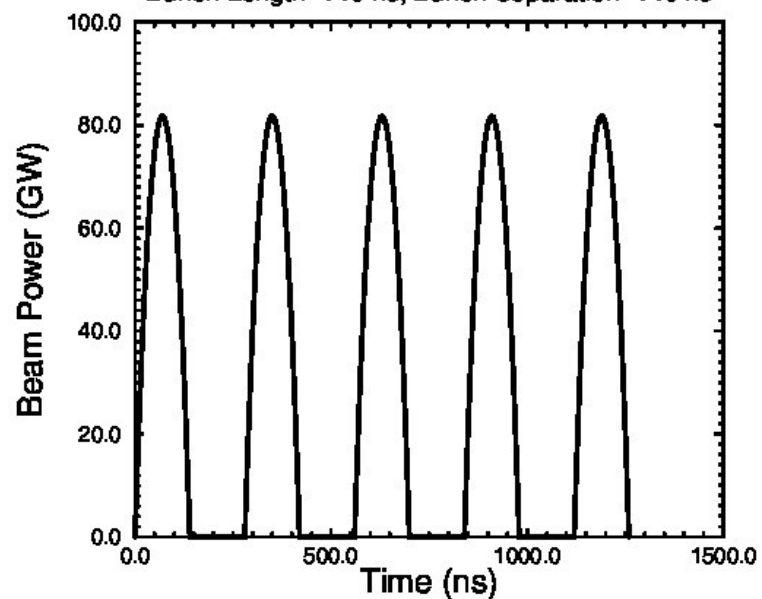
Single Parabolic Bunch, Length = 50 ns



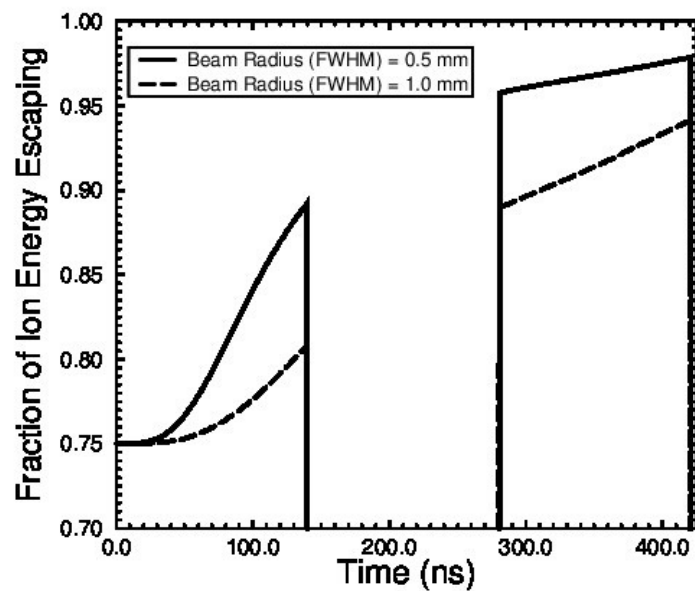
Single Bunch with a Length = 1000 ns



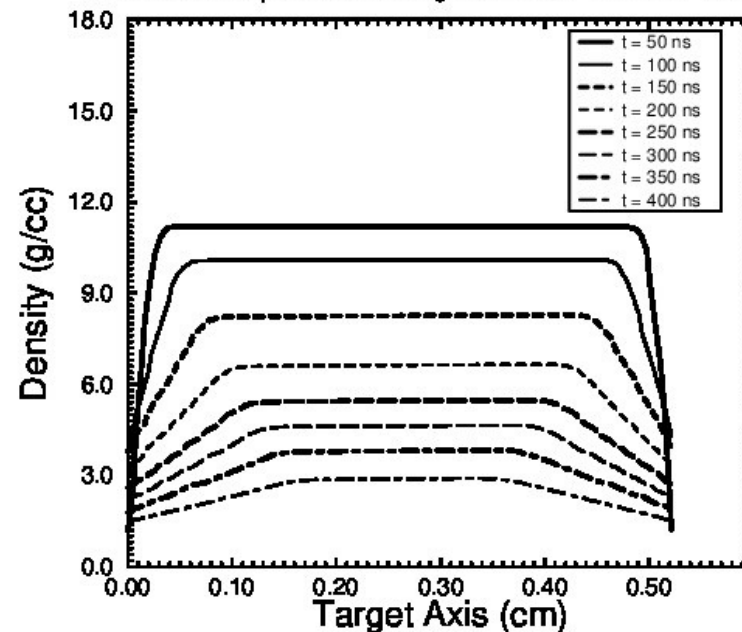
Five Identical Bunches in a Chain, 2×10^{11} Particles per Bunch
 Bunch Length=140 ns, Bunch Separation=140 ns



Five Identical Parabolic Bunches



Gaussian Deposition Along Radius, FWHM = 1.0 mm



Role of Hydrodynamic Tunneling in Case of SPS, LHC and FCC

<u>Parameters</u>	SPS	<u>LHC</u>	<u>FCC</u>
Proton Energy	450 GeV	7 TeV	50 TeV
Bunch Intensity	1.5x10¹¹	1.15x10¹¹	10¹¹
Bunches / Beam	<u>288</u>	<u>2808</u>	<u>10600</u>
Bunch Length	0.5	0.5 ns	0.5 ns
Bunch Separation	25/50 ns	25 ns	25 ns
Beam Duration	7.2 μs	89 μs	265 μs
Focal Spot σ	0.2 mm	0.2 mm	0.2 mm
Energy / Bunch	15.57 kJ	128.8 kJ	800 kJ
Energy / beam	3.8 MJ	362 MJ	8.5 GJ
Tunnel	6.9 km	28 km	100 km

- **An accidental release of even a small fraction of the beam energy can cause considerable damage to the equipment. This may happen by deflecting the beam with a wrong angle (towards collimators, magnets and other equipment).**
- **A very serious scenario could be the loss of the entire beam at a single point.**
- **Chances of happening of such an accident are remote.**
- **The consequences of such an accident needs to be estimated. The risk must remain acceptable. Depending on the consequences the probability for such failure has to be adequate to keep the risk acceptable.**
- **For this purpose we carried out numerical simulations of the full impact of the LHC beam and the FCC beam with solid targets.**

Beam-Target Coupling Parameters

[Solid Density]

	Energy (TeV)	E_p (GeV/g/p) [FLUKA]	I_{bunch}	E_s (kJ/g) [BIG2]	T(K)
SPS	0.44	3.6	1.50×10^{11}	0.08	515
LHC	7	134	1.15×10^{11}	2.4	5019
FCC	40	920	10^{11}	14.4	27440

In case of a long bunched beam, energy deposited by a certain number of bunches [few tens in case of LHC and a few in case of FCC] launches an outgoing radial shock wave that depletes the density along and around the axis.

Protons that are delivered in subsequent bunches together with their hadronic shower penetrate deeper into the target. This is called hydrodynamic tunneling of the ultra-relativistic protons.

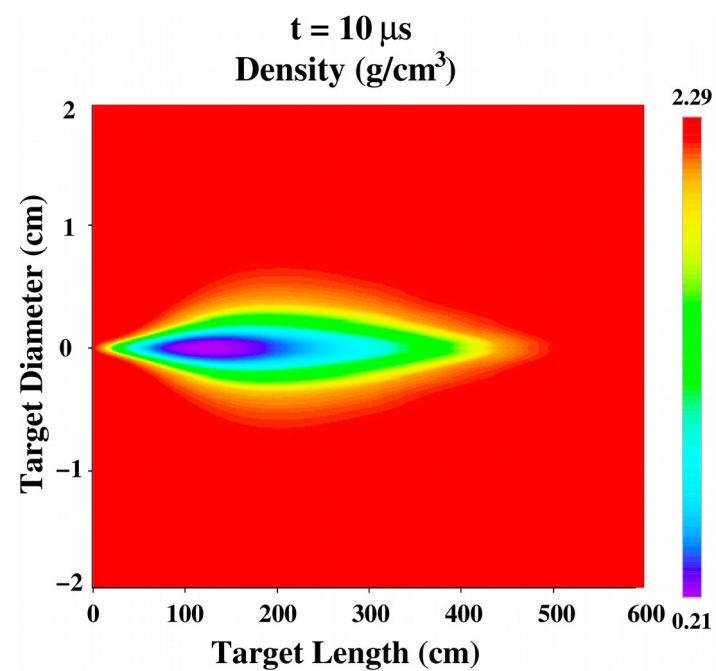
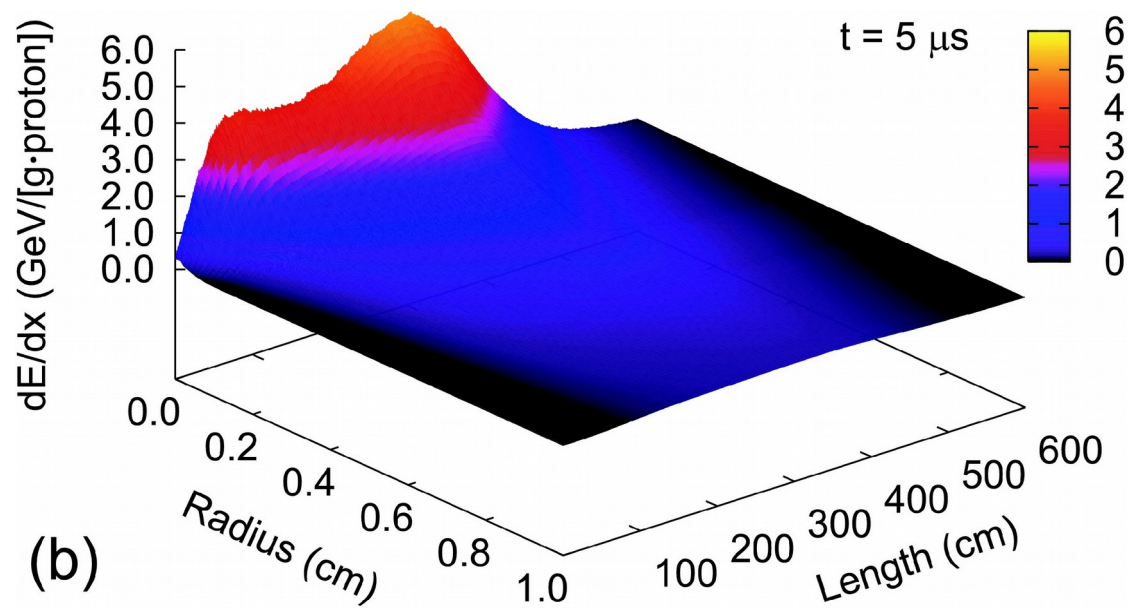
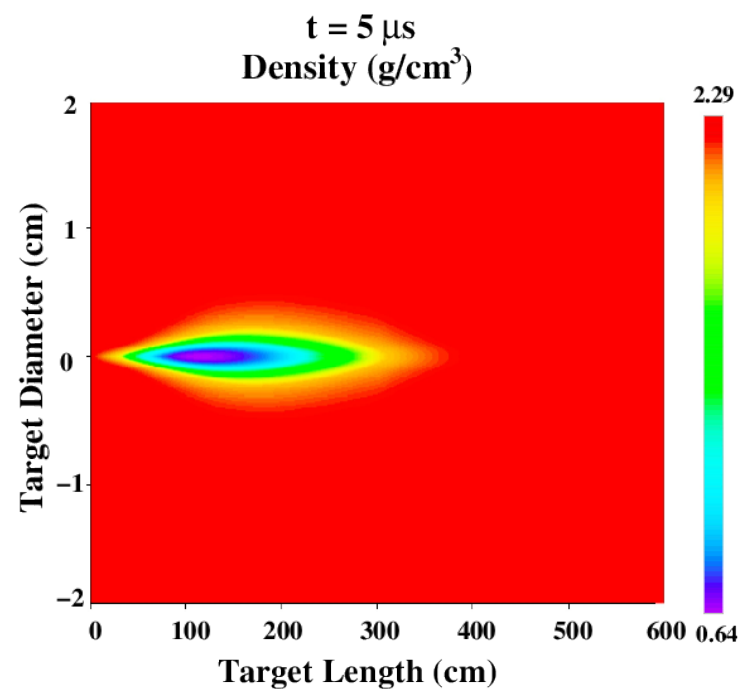
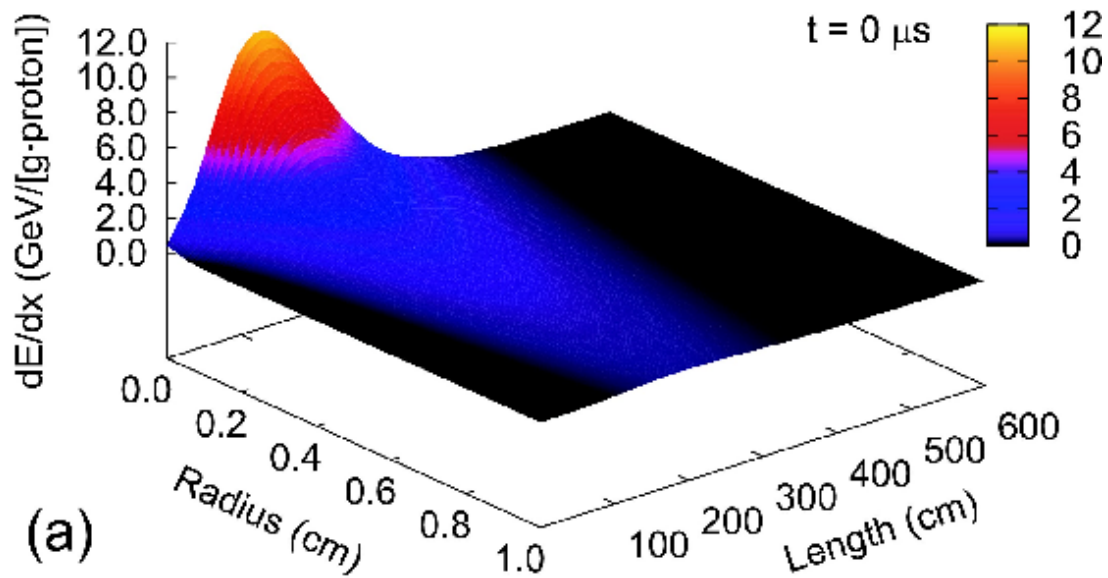
Continuation of this process leads to very substantial range lengthening of the protons and their shower.

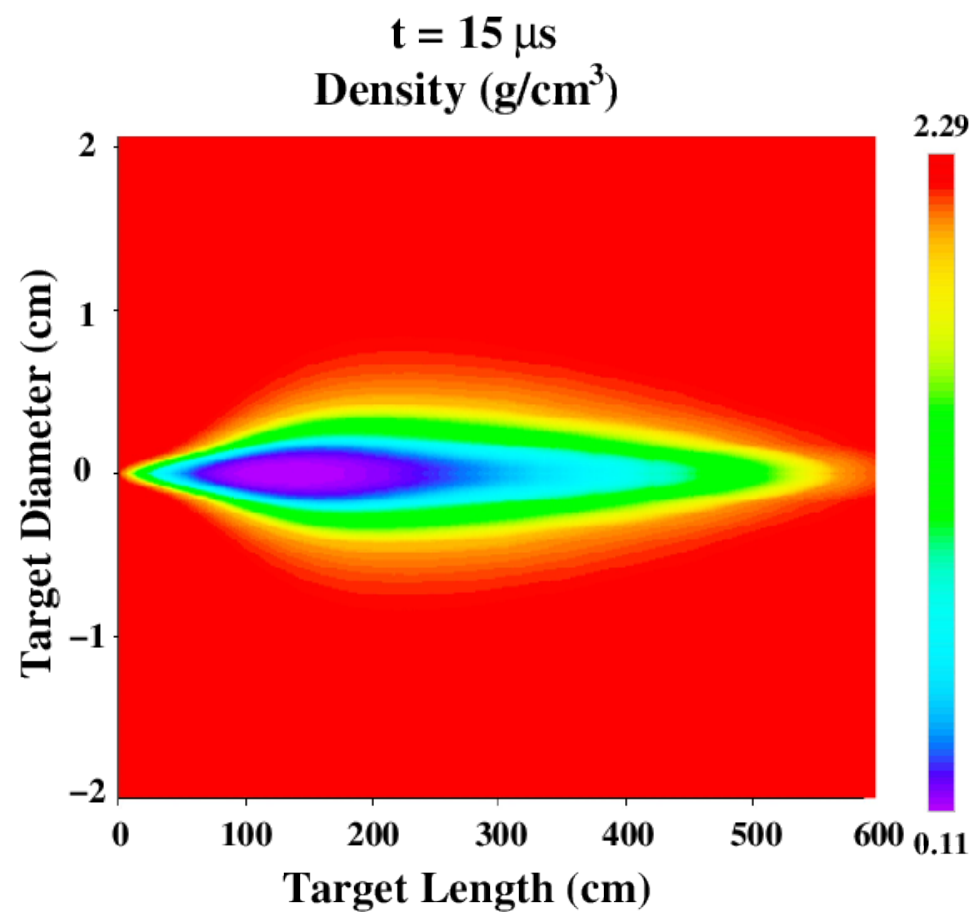
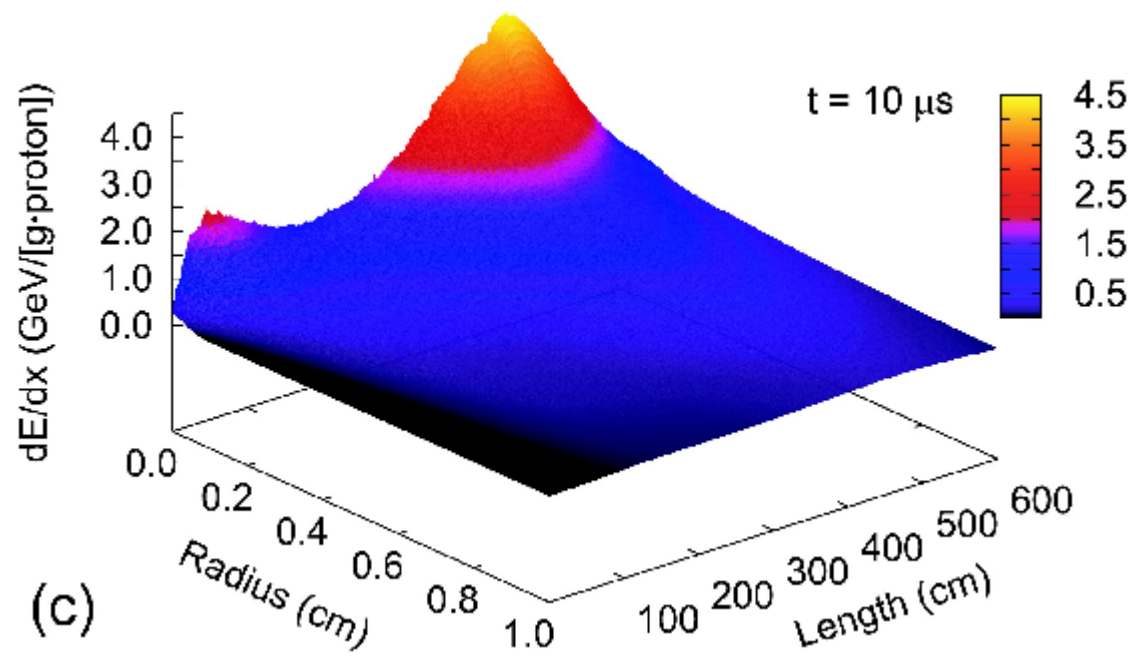
Simulations Strategy

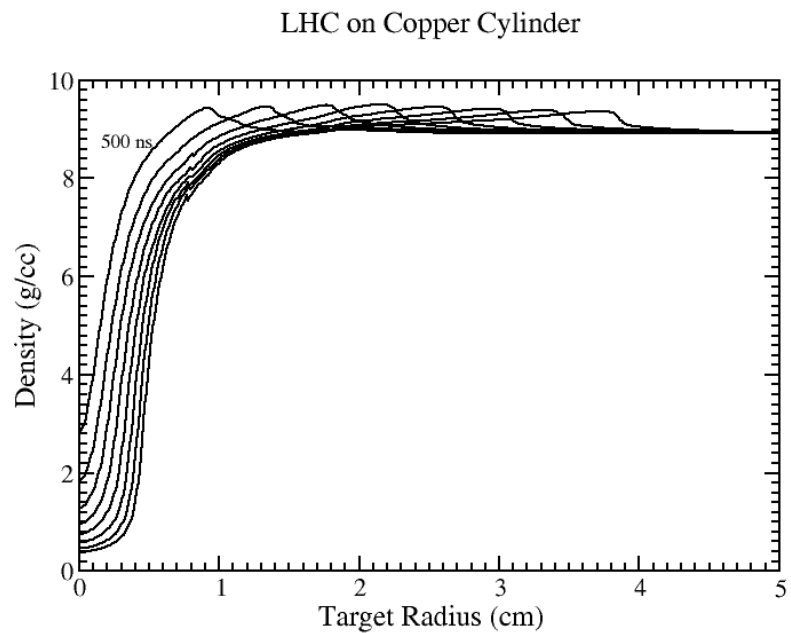
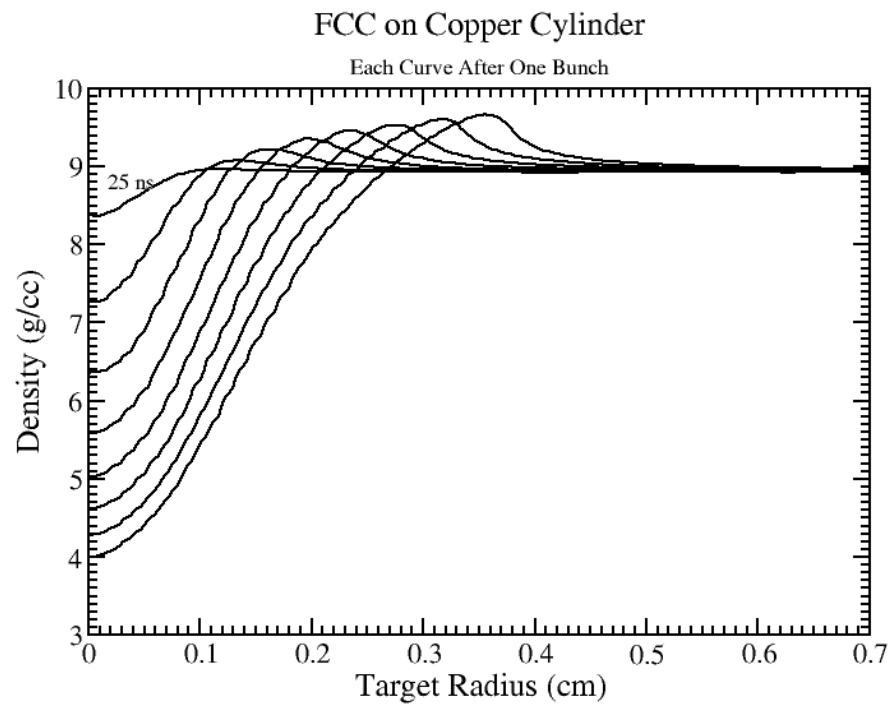
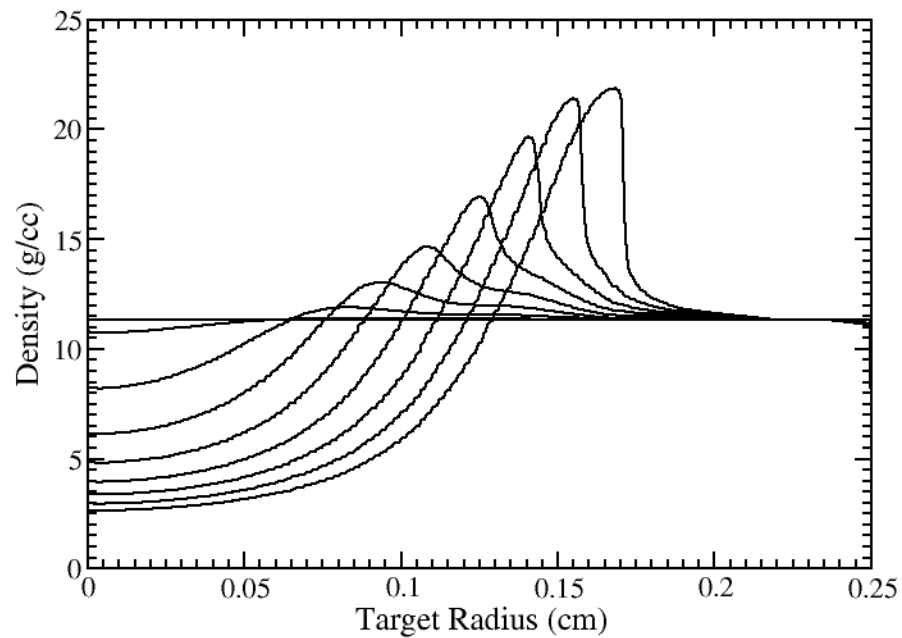
The simulations are carried out using the energy deposition code FLUKA and a 2D hydrodynamic code BIG2, iteratively.

- First, the FLUKA code is run to calculate the energy deposition distribution considering solid target density.
- Second, this energy deposition data is used as energy input to BIG2 and thermodynamic and the hydrodynamic response of the material is simulated.
- The BIG2 code is stopped when the density along the target axis is reduced by 15 % due to the outgoing radial shock wave.
- The modified target density distribution is then used in FLUKA to calculate new energy deposition table that is then used in BIG2. The process is continued till the end.

Iteration step is determined by the beam parameters







FLUKA: is a fully integrated particle physics and multi-purpose Monte Carlo simulation package, capable of simulating all components of the particle cascades in matter from as low as a few MeV/u up to 10000 TeV/u.

More details about the applied models and their performance, can be found in

[*Fasso A et al 2005 FLUKA: a multi-particle transport code CERN-2005-10 , INFN/TC-05/11, SLAC-R-773*]

BIG2 Computer Code

- **BIG2 is a 2D hydrodynamic computer code based on a Godunov type numerical scheme.**
- **It can work in Eulerian as well as Lagrangian algorithm.**
- **It can handle simple and complicated geometries and can deal with single as well as multi-layered targets.**
- **It uses sophisticated, very versatile and stable mesh.**
- **It includes thermal conduction.**
- **It includes ion energy deposition [uses **SRIM** for heavy ion energy deposition and **FLUKA** for SPS, LHC and FCC protons].**

- It treats elastic plastic effects using Prandtl-Reuss model, which means Hooks law complemented with von Mises yield criterion:

$$M = [3(\bar{\sigma}_{xx}^2 + 2\bar{\sigma}_{xy}^2 + \bar{\sigma}_{yy}^2)/2Y_0^2]$$

If $M < 1$, elastic regime, Otherwise the material is plastified.

$\bar{\sigma}$ represents components of the stress tensor Y_0 is yield strength.

- Different phases of the target material are treated using a semi-empirical EOS model [*I.V. Lomonosov LPB 25 (2007)567*]. This package includes treatment for solid, melting, liquid, 2-phase liquid-gas and gaseous states. In addition to that, it includes data of all the experimental measurements made for a given material. *Effects of first ionization level are included by combining it with a plasma chemical model.*

Hydrodynamic Simulations

1. Lagrangian algorithm

Coordinate system moves
with the fluid (velocity “v”)

Advantages

Well defined material interfaces

Free boundary conditions
applicable

High degree of compression with
high resolution can be studied

Disadvantages

Strong distortions of the
computational mesh

2. Eulerian algorithm

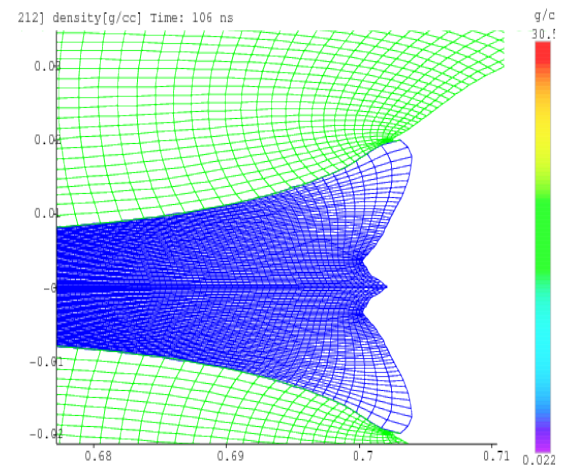
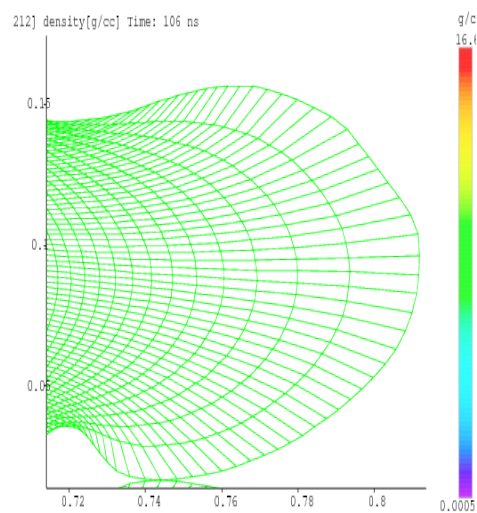
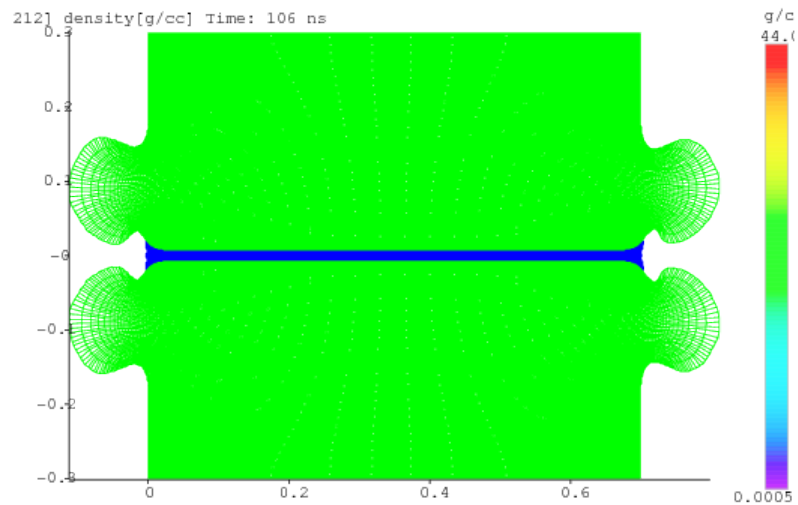
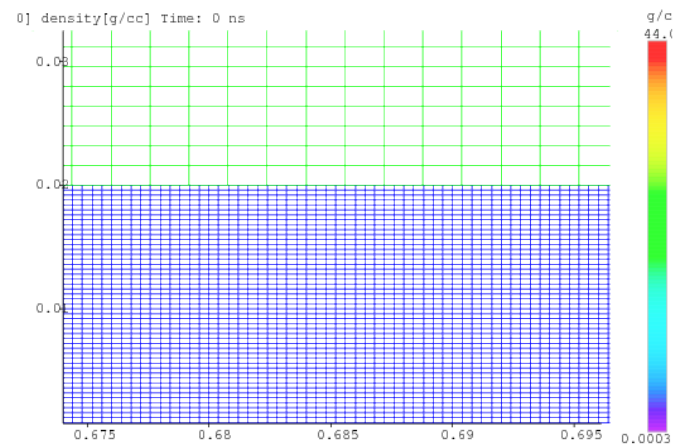
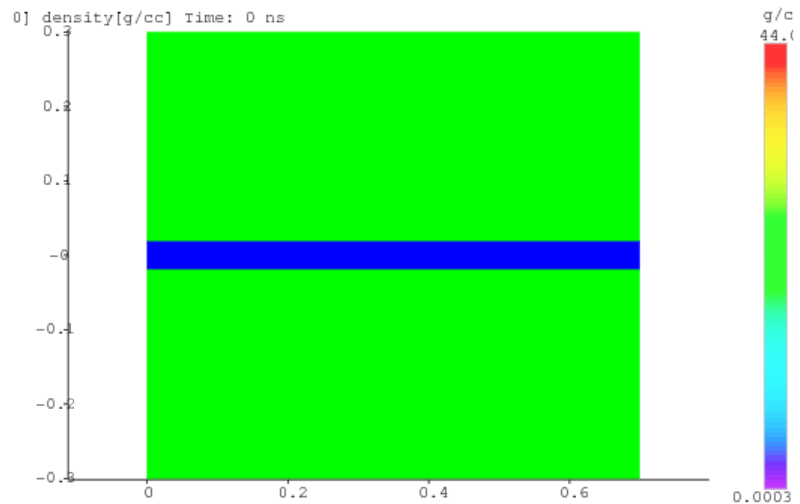
Coordinate system is
fixed (velocity “zero”)

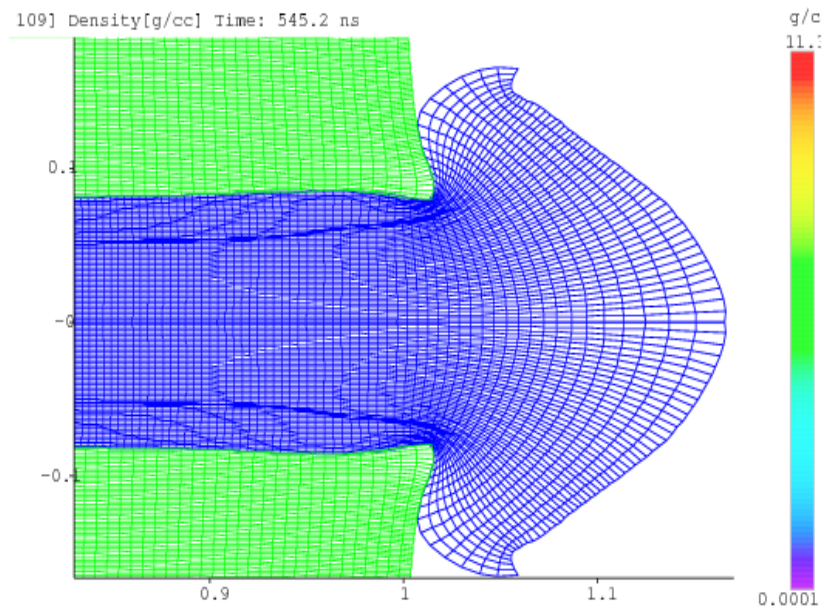
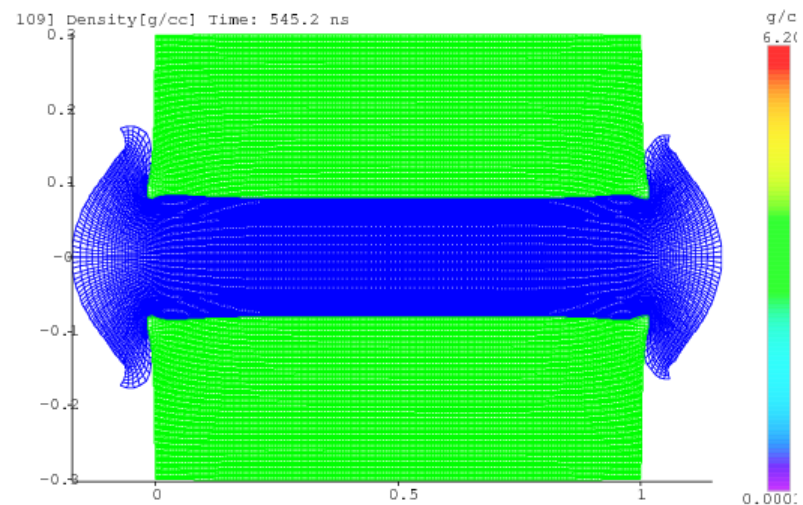
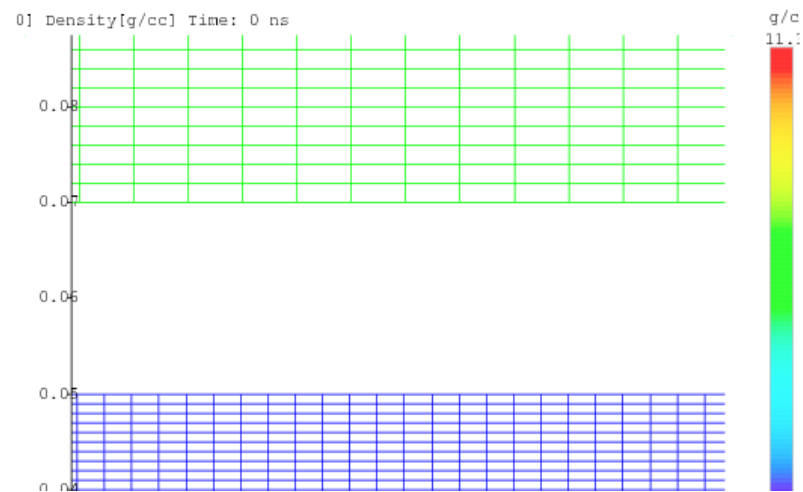
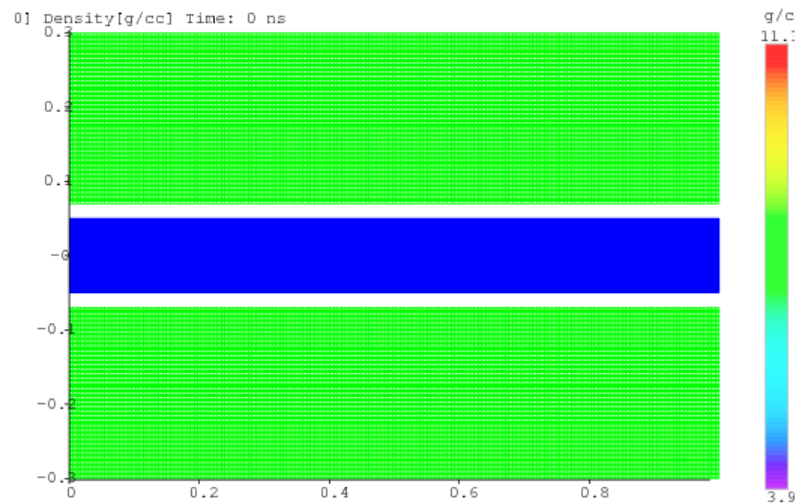
Advantages

No mesh distortion

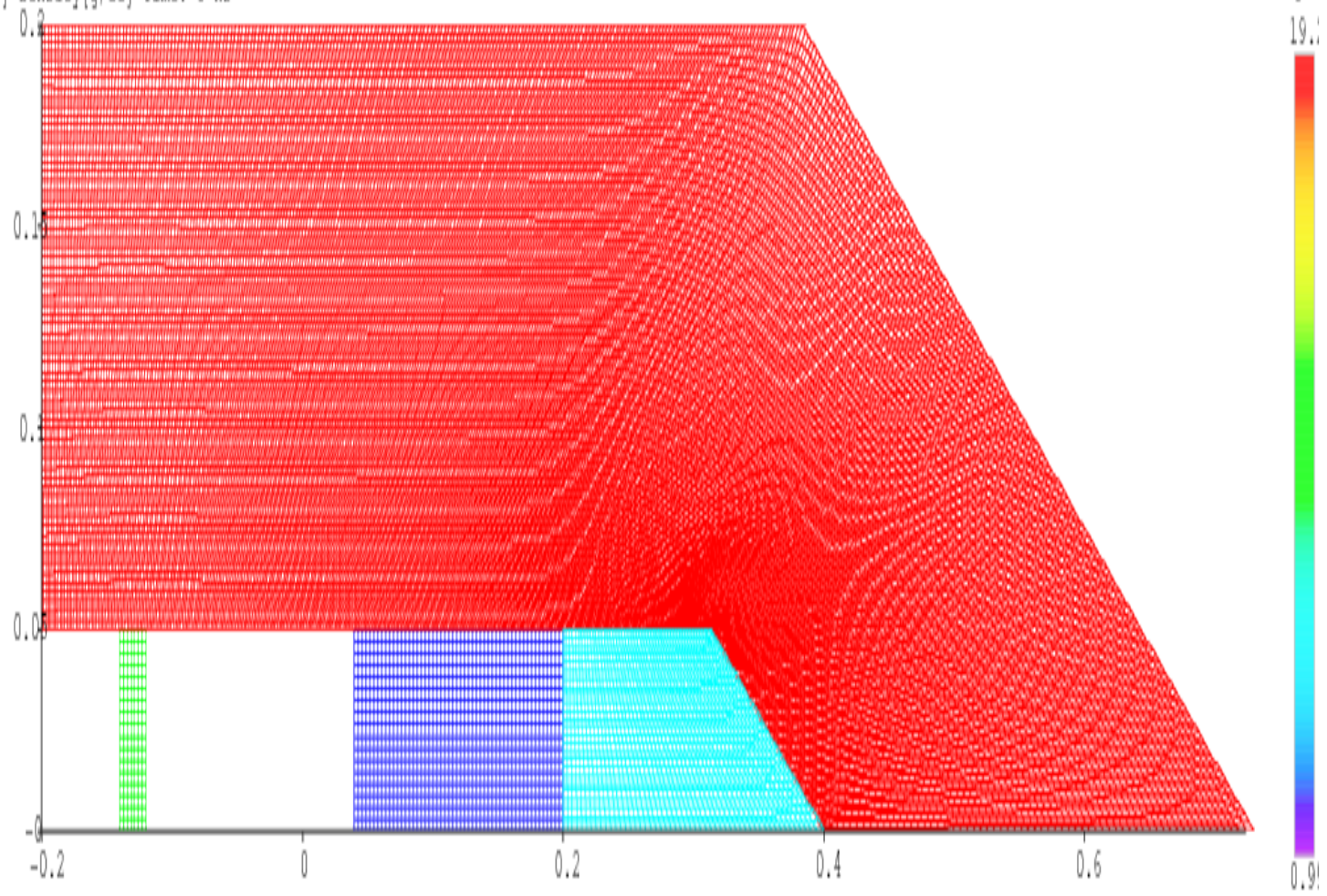
Disadvantages

Material interfaces lose their
sharp definitions, special treatment
at interfaces, complicated, inaccuracies.
Local regions of fine resolution difficult
to achieve





0] Density[g/cc] Time: 0 ns



Historical View of Hydrodynamic Tunneling

Simulations were done in the 80's to study the impact of the **20 TeV Superconducting Super Collider (SSC)** proton beam with a carbon beam dump.

Energy deposition code **MARS** and hydrodynamic code, **MESA** (2D and 3D) were used for this purpose.

Beam duration = 290 μ s

$\delta = 2$ mm

Two fluences considered: 4×10^{17} and 10^{19} protons/s

First case: wave speed: **7 cm/s**, penetration length = **20 m**

Second case: wave speed: **70 cm/s**, penetration length = **200 m**

Beam dump length = **8 m**

1. Impact of LHC 7 TeV proton beam with

**1. Copper cylinder: $L = 5$ m; $r = 5$ cm; $\delta = 0.2$ mm;
Beam duration = 89 μ s.**

**Solid density energy deposition data normalized
with line density in each simulation cell at every
time step.**

Penetration distance = 35 m.

**2. Carbon cylinder: $L = 6$ m; $r = 5$ cm; $\delta = 0.4$ mm;
FLUKA and BIG2 iteratively.**

Penetration distance = 25 m.

2. Design of Experiment at HiRadMat Facility

450 GeV SPS proton beam with Cu cylinder

FLUKA and BIG2 used iteratively

L = 1.5 m; r = 5 cm

288 proton bunches, bunch length = 0.5 mm

Bunch separation 25 ns.

Beam duration 7.2 μ s.

$\bar{\sigma} = 0.1, 0.2$ and 0.5 mm

Penetration distance longer for shorter $\bar{\sigma}$.

It was concluded that 150 cm long Cu cylinder

Should be sufficient to stop the beam.

3. Interpretation of Experiments

440 GeV proton beam

Bunch length = 0.5 ns, Bunch separation = 50 ns,

Beam Duration = 7.2 μ s.

Focal spot $\bar{\sigma}$ = 2 mm, 144 bunches.

$\bar{\sigma}$ = 0.2 mm, 108 bunches and 144 bunches.

15 solid copper cylinders, each 10 cm long having a radius = 4 cm with 1 cm gap in between.

For simplicity we used a single cylinder 150 cm long.

FLUKA and BIG2 used iteratively. Simulation

results agreed with the experimental measurements within a few per cent.

4. Impact of FCC 40 TeV Proton Beam on Copper Target

40 TeV protons

**10600 bunches, bunch length = 0.5 ns,
bunch separation = 25 ns, $\bar{\sigma} = 0.2$ mm.**

Impact on a solid Cu cylinder

Length = 5 m, radius = 2 cm

FLUKA and BIG2 used iteratively.

Penetration distance about 310 m!

For 50 TeV it could be around 350 m.

5. Impact of FCC 50 TeV Proton Beam on Water Target

50 TeV protons

**10600 bunches, bunch length = 0.5 ns,
bunch separation = 25 ns, $\sigma = 0.4$ mm.**

Impact on water cylinder

Length = 15 m, radius = 15 cm

FLUKA and BIG2 used iteratively.

Penetration distance about 1.3 km.

Journal Publications

1. N.A. Tahir, F. Burkart, R. Schmidt et al., NIMB 427 (2018) 70.
2. N.A. Tahir, F. Burkart, R. Schmidt et al., Phys. Plasmas 24 (2017) 072712.
3. N.A. Tahir, F. Burkart, R. Schmidt et al., Contributions to Plasma Phys. 57 (2017) 452.
4. N.A. Tahir, F. Burkart, R. Schmidt et al., PRAB 19 (2016) 081002.
5. N.A. Tahir, F. Burkart, R. Schmidt et al., High Energy Density Phys. 21 (2016) 27.
6. F. Burkart, R. Schmidt, V. Raginal et al., J. Appl. Phys. 118 (2015) 055902.
7. R. Schmidt, J. Blanco Sancho, F. Burkart et al., Phys. Plasmas 21 (2014) 080701.
8. N.A. Tahir, F. Burkart, A. Shutov et al., Phys. Rev. E 90 (2014) 063112.
9. N.A. Tahir, J. Blanco Sancho, R. Schmidt et al., High Energy Density Phys. 21 (2013) 269.

10. N.A. Tahir, L. Blanco Sancho, A. Shutov et al., PRSTAB 15 (2012) 051003.
11. N.A. Tahir, R. Schmidt, A. Shutov et al., Contributions to Plasma Phys. 51 (2011) 299.
12. N.A. Tahir, R. Schmidt, A. Shutov et al., Phys. Rev E 79 (2009) 046410.
13. N.A. Tahir, R. Schmidt, M. Brugger et al., Laser Part. Beams 27 (2009) 475.
14. N.A. Tahir, R. Schmidt, M. Brugger et al., Phys. Plasmas 16 (2009) 082703.
15. N.A. Tahir, R. Schmidt, M. Brugger et al., New J. Phys. 10 (2008) 073028.
16. N.A. Tahir, A. Shutov, I.V. Lomonosov et al., J. Phys. IV France 133 (2006) 1085.
17. N.A. Tahir, B. Goddard, V. Kain et al., J. Appl. Phys. 97 (2005) 083532.
18. N.A. tahir, V. Kain, R. Schmidt et al., Phys. Rev. Lett. 94 (2005) 0135004.