Where we stand on B-decay Discrepancies – Global Fits –

Diego Guadagnoli CNRS, Annecy

Minimal TH considerations

(before any fit)



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 $dC_9^{(\mu)}$ alone

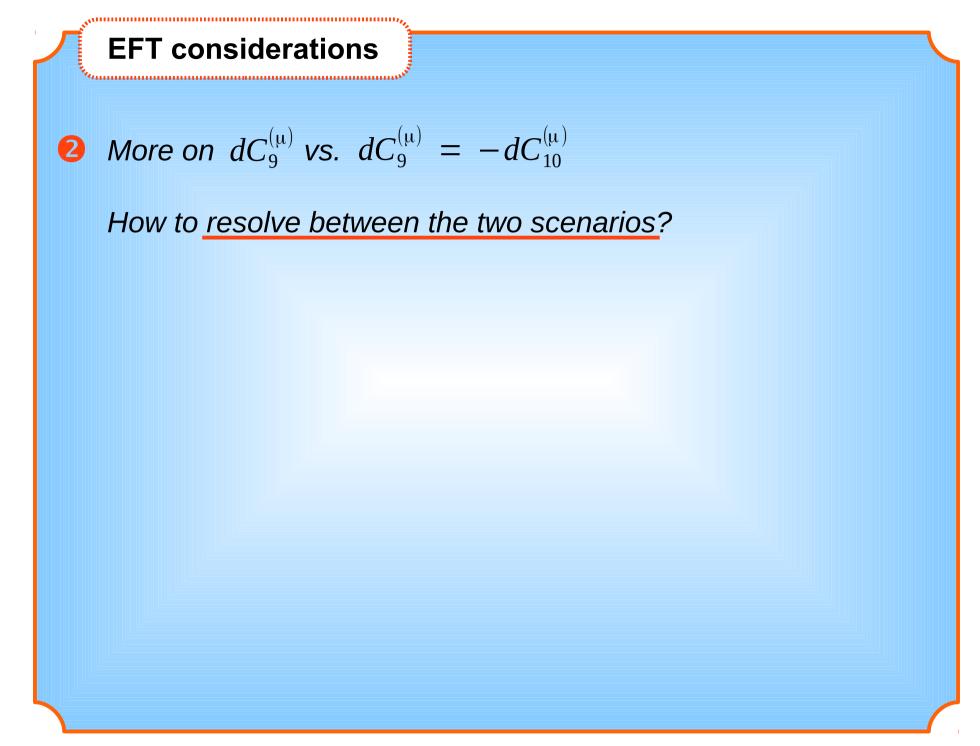
$$dC_9^{(\mu)} = -dC_{10}^{(\mu)}$$

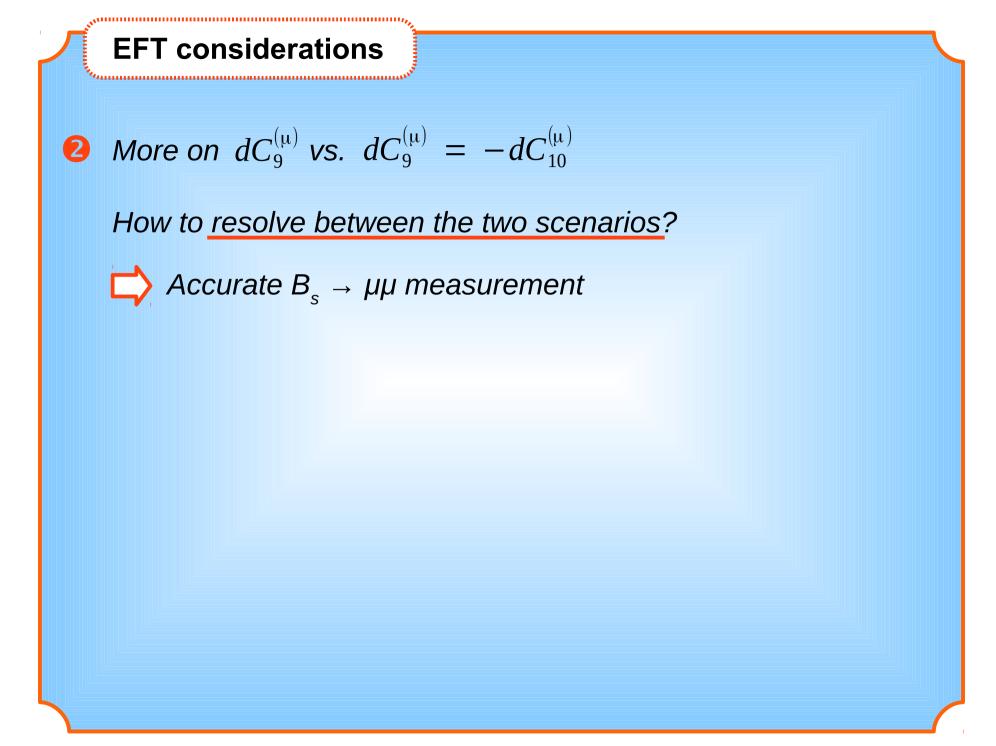
[Hiller-Schmaltz, 2014]

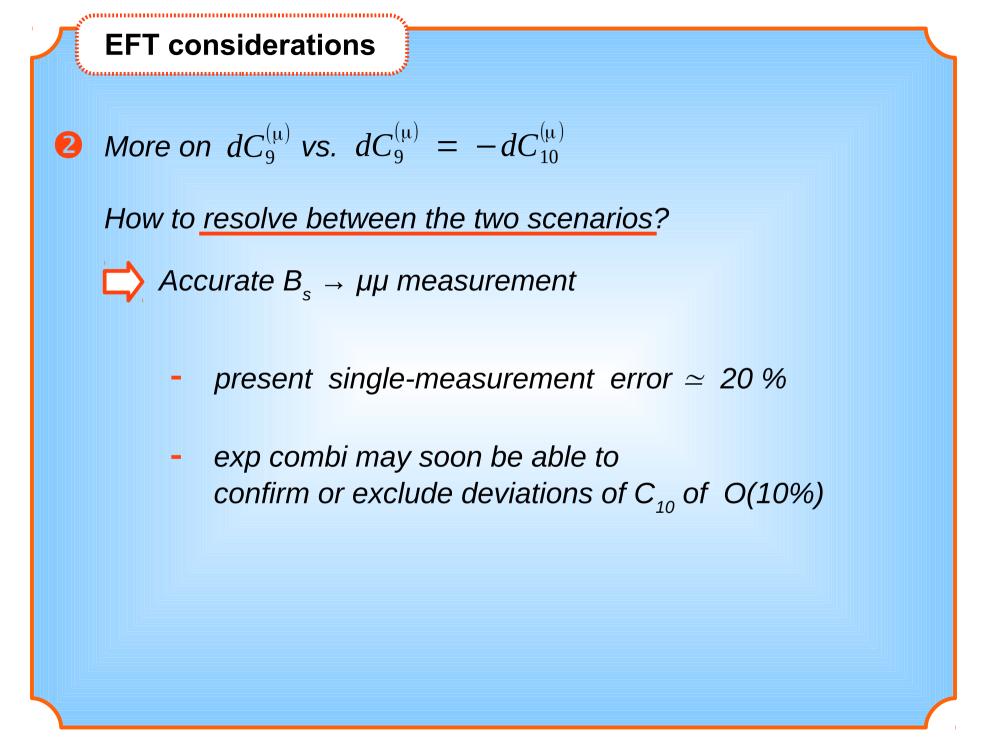
corresponding op. combination can be written in terms of SU(2)_L invariants [Alonso, Grinstein, M.Camalich, 2014]

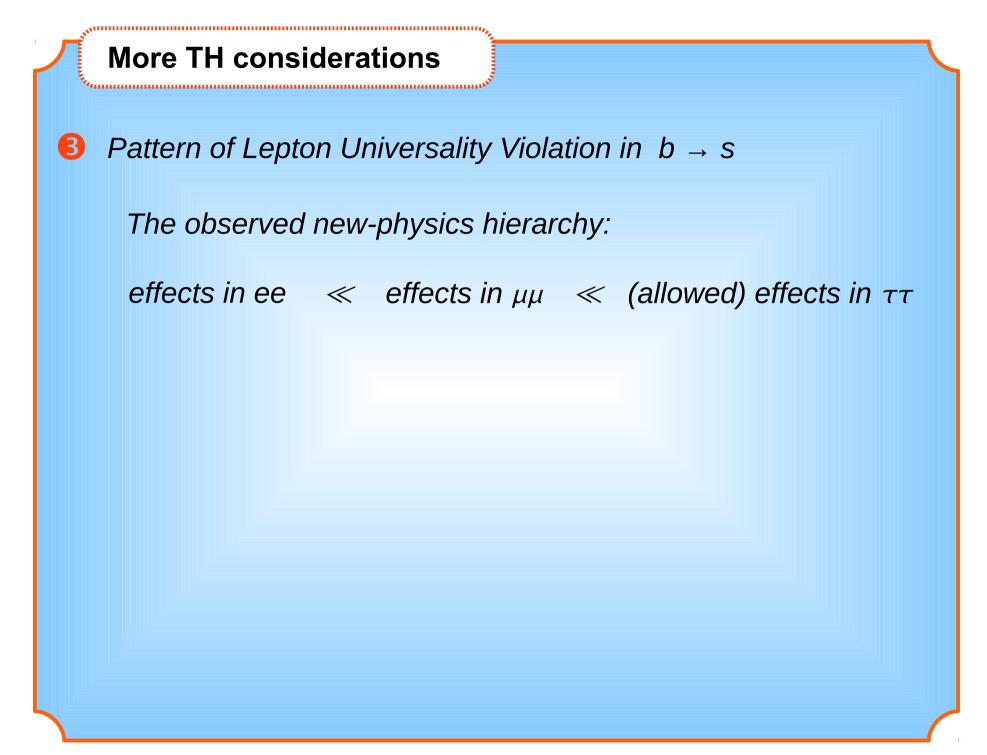
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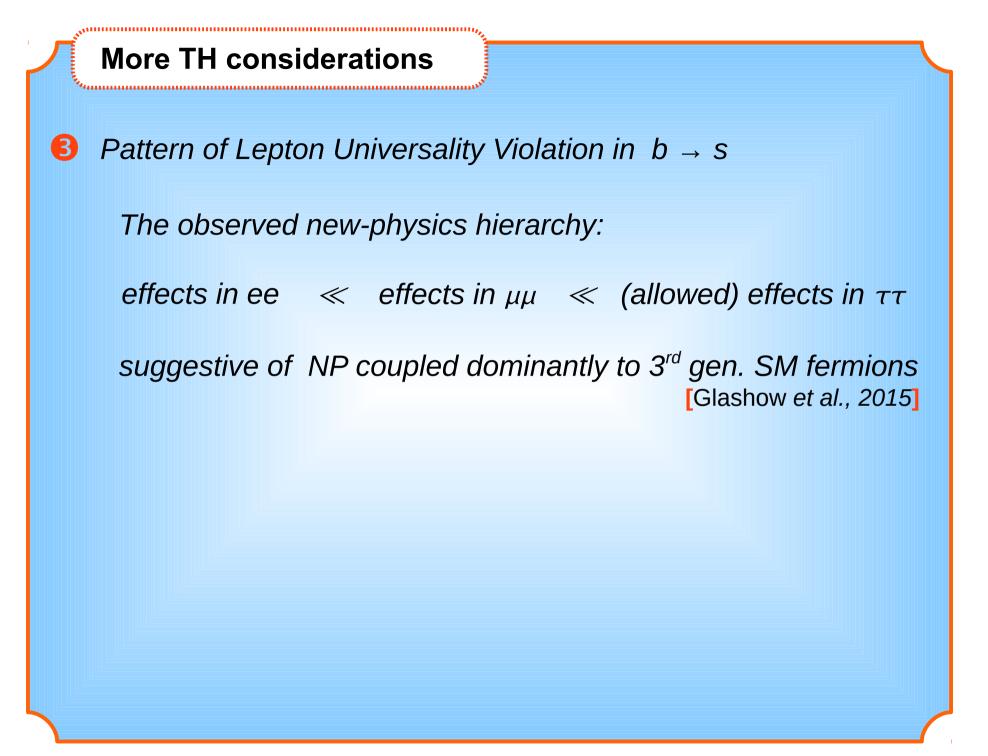
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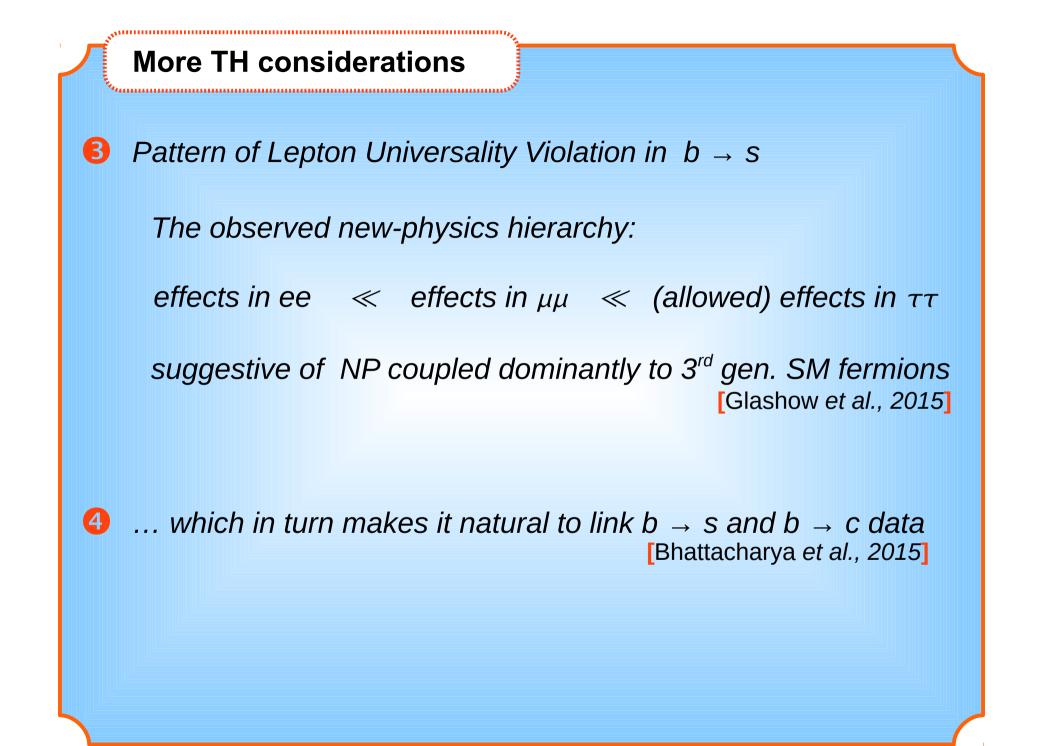


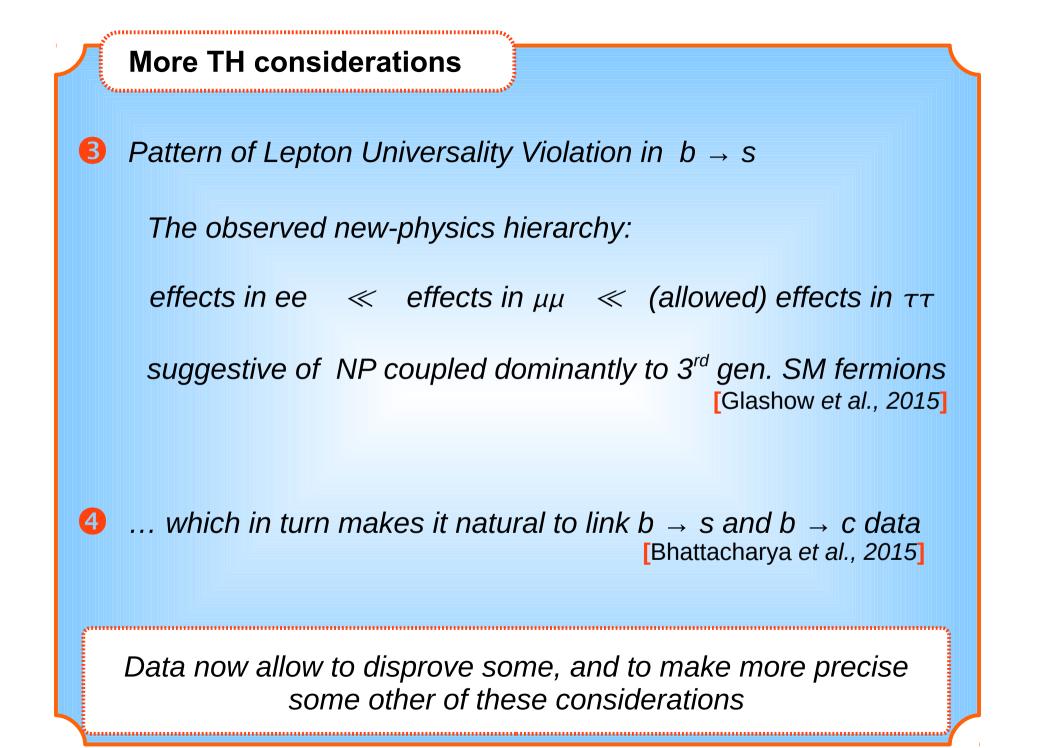












Weak-Effective-Theory

Global fits

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[Aebischer et al., 2019]

Coeff.	best fit	1σ	2σ	pull
$C_9^{bs\mu\mu}$	-0.97	[-1.12, -0.81]	[-1.27, -0.65]	5.9σ
$C_9^{\prime bs\mu\mu}$	+0.14	[-0.03, +0.32]	[-0.20, +0.51]	0.8σ
$C_{10}^{bs\mu\mu}$	+0.75	[+0.62, +0.89]	[+0.48, +1.03]	5.7σ
$C_{10}^{\prime bs\mu\mu}$	-0.24	[-0.36, -0.12]	[-0.49, +0.00]	2.0σ
$C_9^{bs\mu\mu}=C_{10}^{bs\mu\mu}$	+0.20	[+0.06, +0.36]	[-0.09, +0.52]	1.4σ
$C_9^{bs\mu\mu}=-C_{10}^{bs\mu\mu}$	-0.53	[-0.61, -0.45]	[-0.69, -0.37]	6.6σ

Similar fits (w/ non-identical conclusions) performed in:

[Algueró et al., 1903.09578] [Ciuchini et al., 1903.09632]

[Datta et al., 1903.10086] [Kowalska et al., 1903.10932]

[Arbey et al., 1904.08399]

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- Two scenarios stand out: C_9 alone or $C_9 = -C_{10}$ (µµ-channel only)
- $C_9 = -C_{10}$ now better than C_9 alone
- C_{10} alone also ok, but $B \rightarrow K^* \mu \mu$ unresolved

 C_{9} vs. $C_{9} = -C_{10}$

What makes $C_9 = -C_{10}$ more significant than C_9 alone?

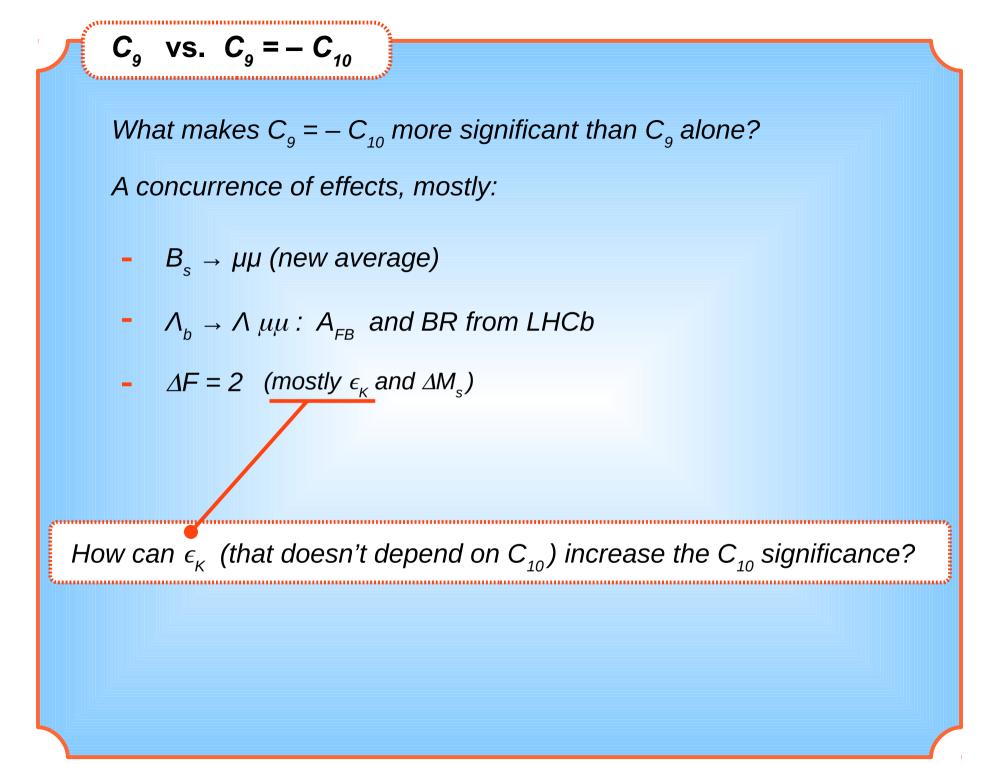
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A concurrence of effects, mostly:

- $B_s \rightarrow \mu\mu$ (new average)
- $\Lambda_b \rightarrow \Lambda \mu \mu$: A_{FB} and BR from LHCb

-
$$\Delta F = 2$$
 (mostly ϵ_{κ} and ΔM_{s})



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Now consider the corresponding terms in the χ^2 function

(with $D_i \equiv O_i^{exp} - O_i^{th}$)

$$\chi^{2} = \frac{D_{1}^{2}(C_{10})}{\sigma_{1}^{2}} + \frac{D_{2}^{2}}{\sigma_{2}^{2}} + \rho_{12} \frac{D_{1}(C_{10})D_{2}}{\sigma_{1}\sigma_{2}}$$

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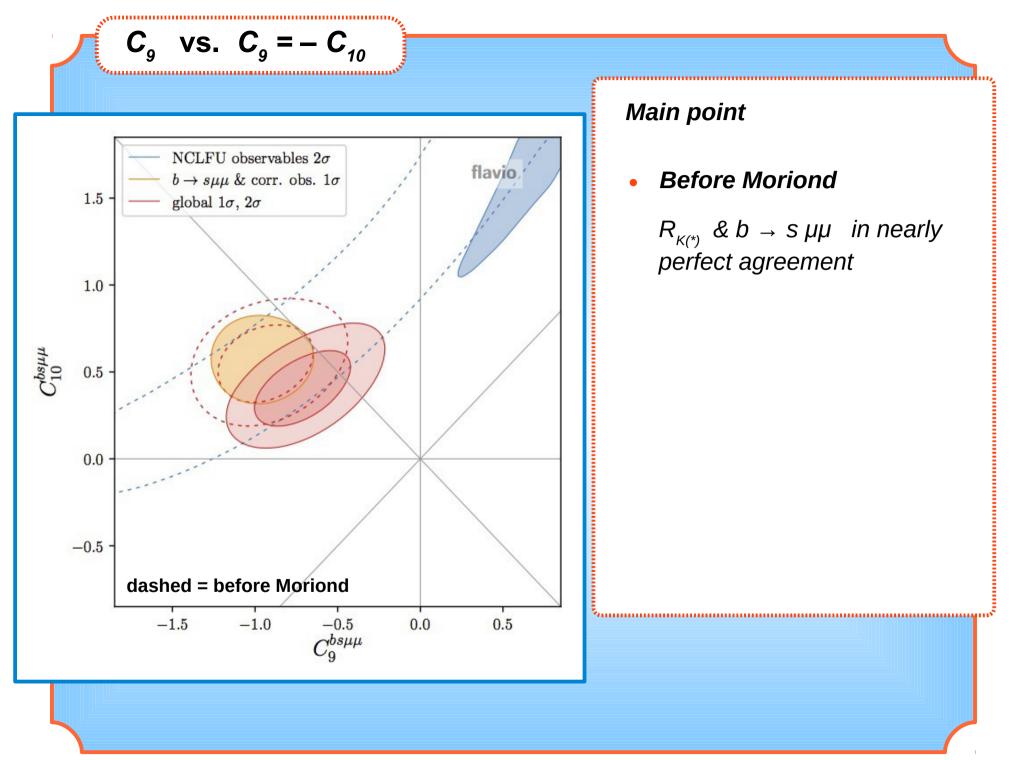
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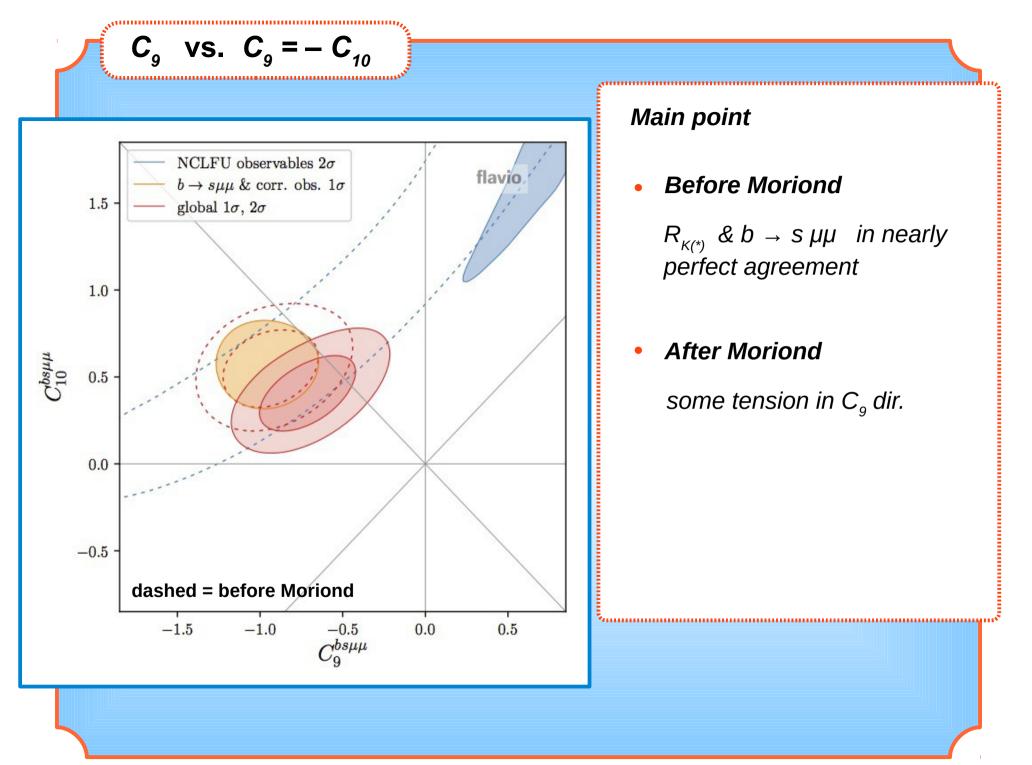
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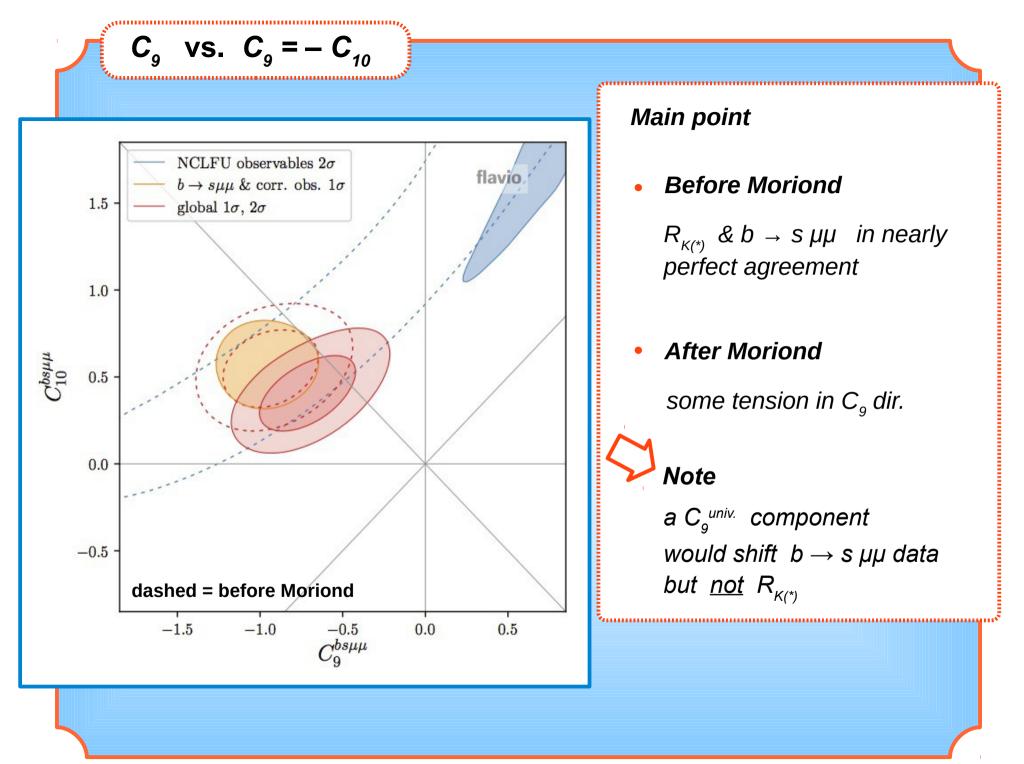
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- The 3^{rd} term depends on C_{10} through D_1

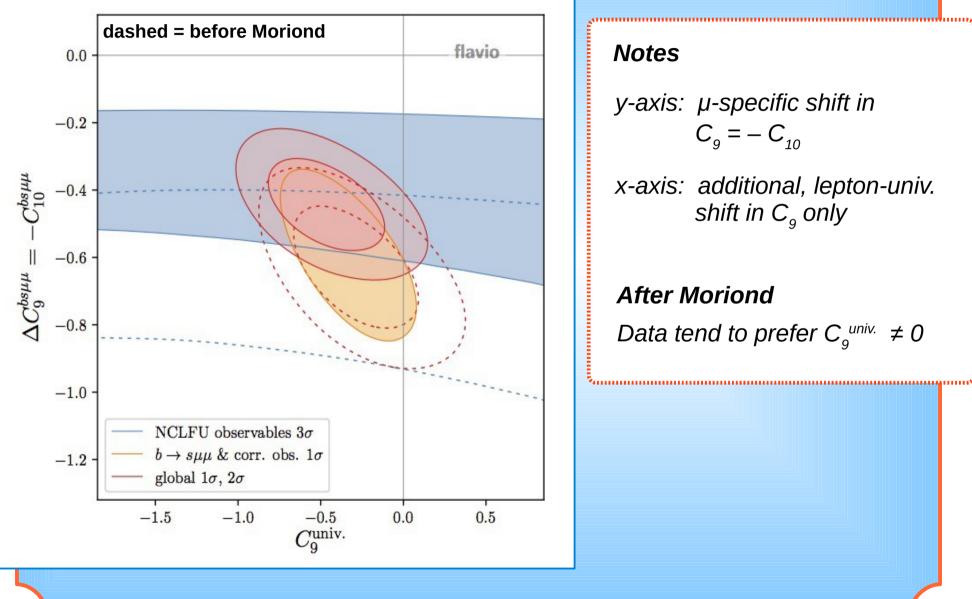
- This term will influence the C_{10} value, depending on sign($D_2 \rho_{12}$)







Univ. vs. non-univ. Wilson coeffs.



y-axis: µ-specific shift in $C_{9} = -C_{10}$ x-axis: additional, lepton-univ. shift in C_{q} only After Moriond

Data tend to prefer $C_9^{\text{univ.}} \neq 0$

Both $C_9 = -C_{10}$ and $C_9^{univ.}$ well justified above the EW scale. The SMEFT picture

Assume BSM d.o.f. to occur at a scale $\Lambda \gg M_{_{EW}}$ Dynamics below Λ described by ops.

SMEFT picture

- constructed with SM fields only
- and invariant under the full SM group

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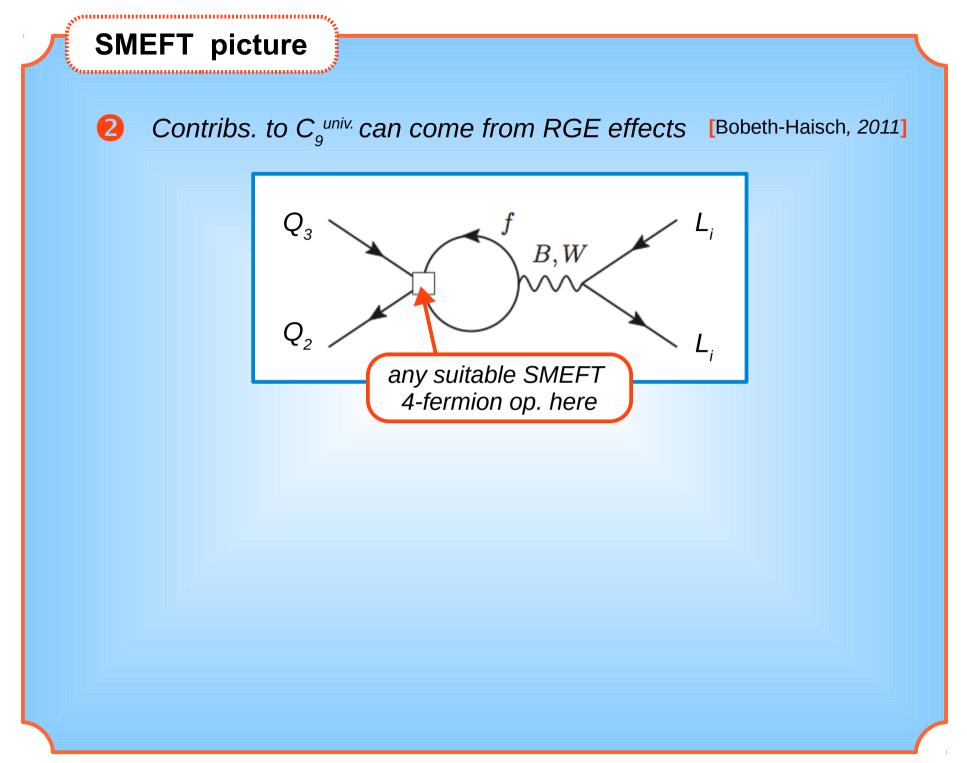
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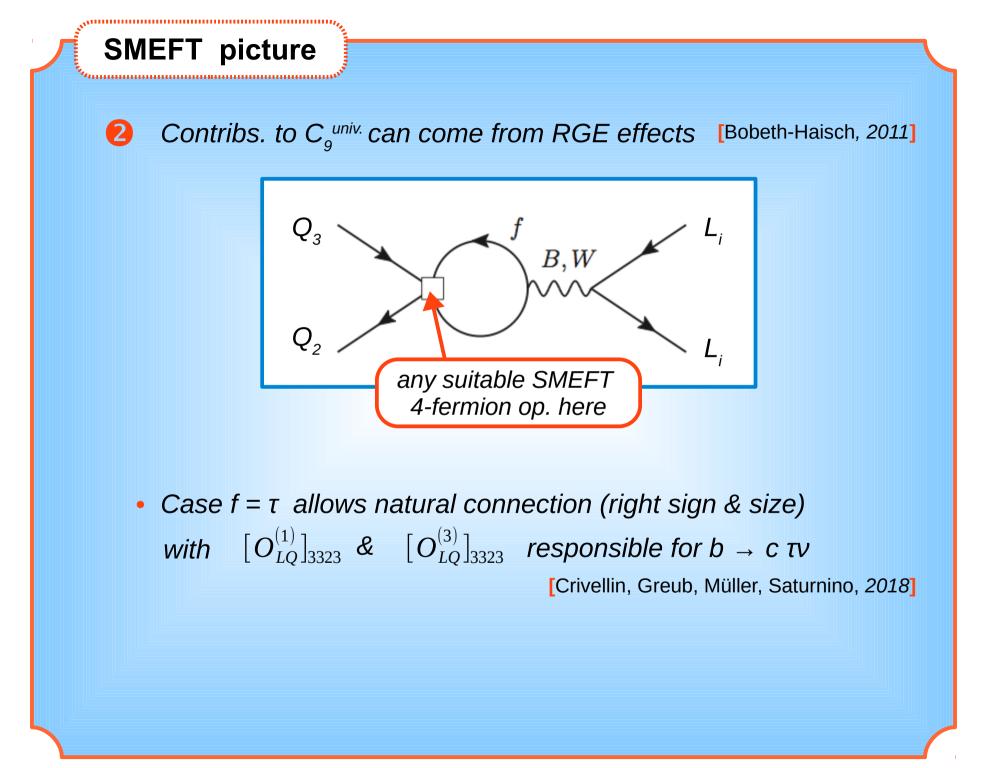
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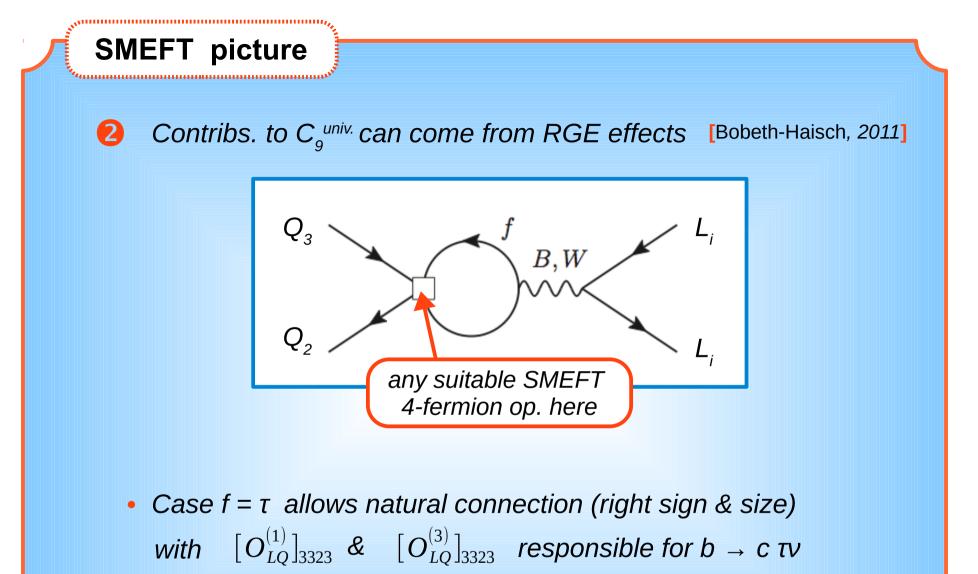
Contributions to muonic $C_9 = -C_{10}$ may come from SMEFT ops. directly matching onto $O_{9,10}$

$$[O_{LQ}^{(1)}]_{2223} = \overline{L}_2 \gamma^{\lambda} L_2 \cdot \overline{Q}_2 \gamma_{\lambda} Q_3$$

$$[O_{LQ}^{(3)}]_{2223} = \overline{L}_2 \gamma^{\lambda} \sigma^a L_2 \cdot \overline{Q}_2 \gamma_{\lambda} \sigma^a Q_3$$

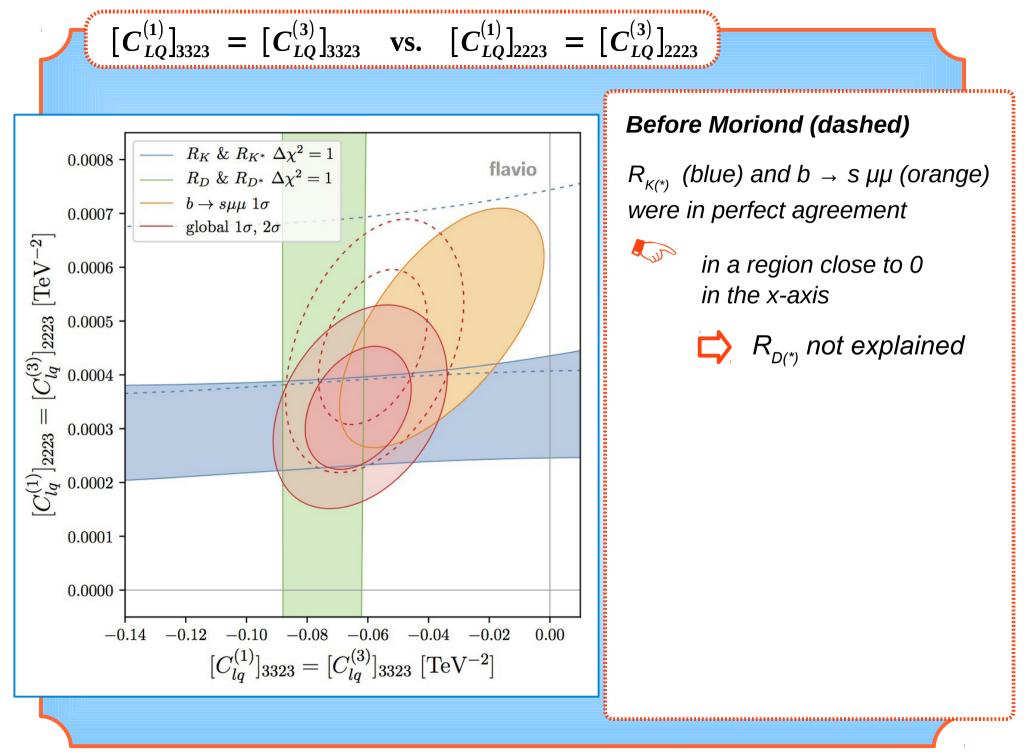


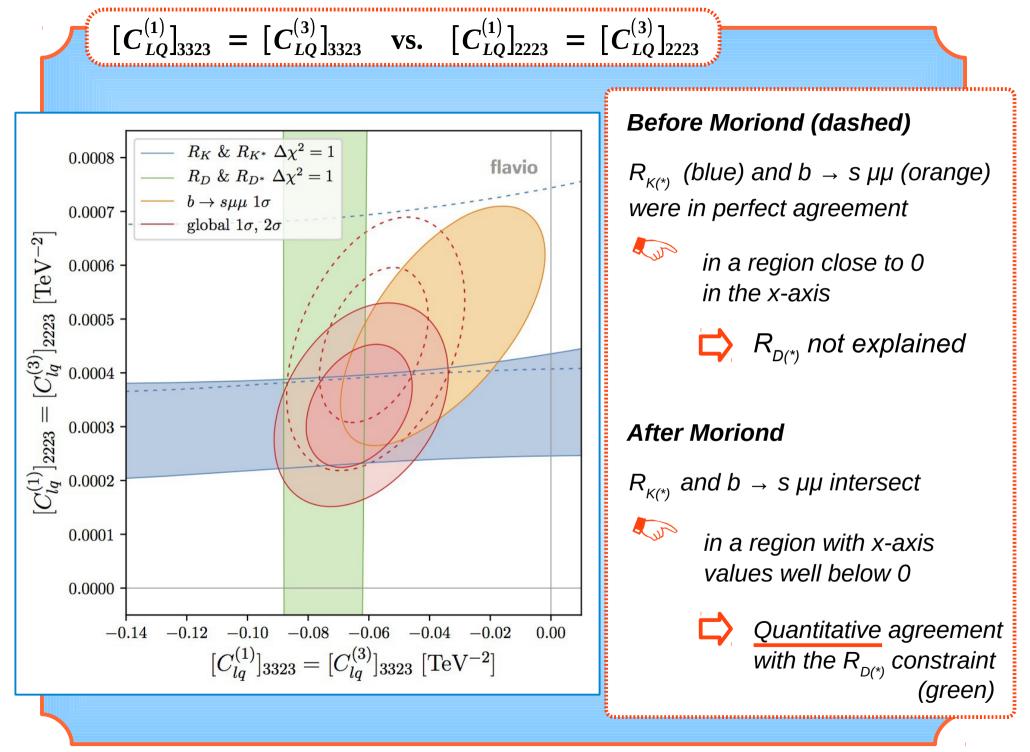


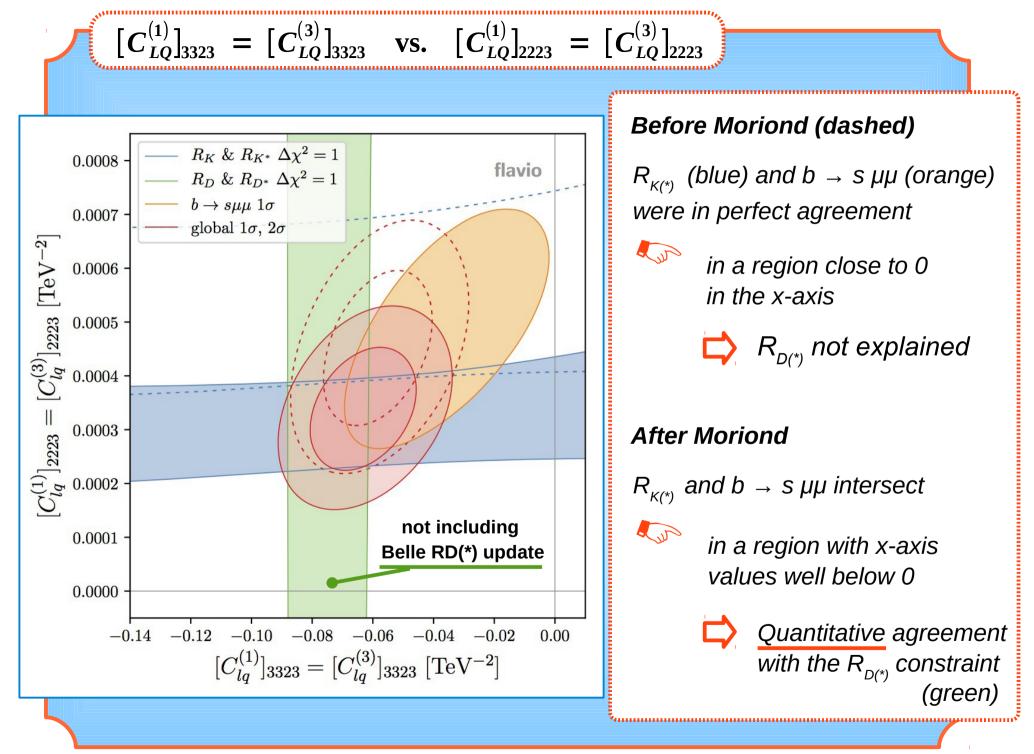


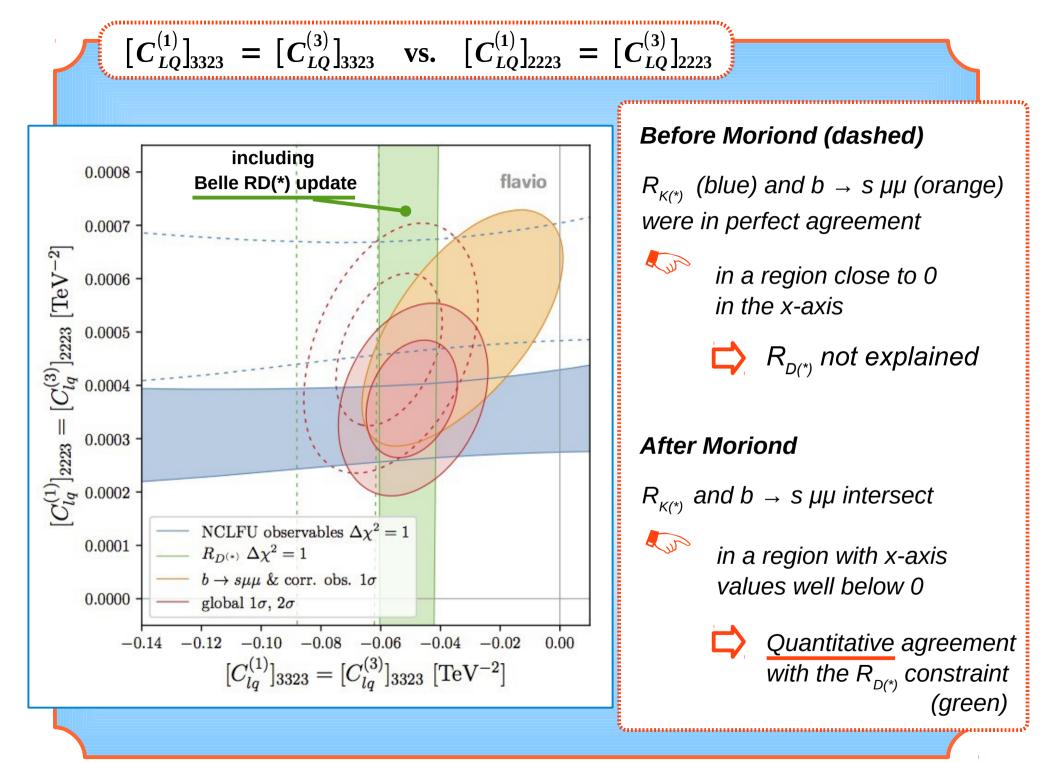
[Crivellin, Greub, Müller, Saturnino, 2018]

• Caveat: need $[C_{LQ}^{(1)}]_{3323} \simeq [C_{LQ}^{(3)}]_{3323}$ to avoid $B \rightarrow K(*) vv$ constraint [Buras-Girrbach-Niehoff-Straub, 2014]









Conclusions

Post-Moriond updates imply a nicely coherent TH picture, with

- $R_{\kappa} \& R_{\kappa^*} < by O(20\%)$ than SM
- $R_D \& R_{D^*} > by O(10\%)$ (not more) than SM
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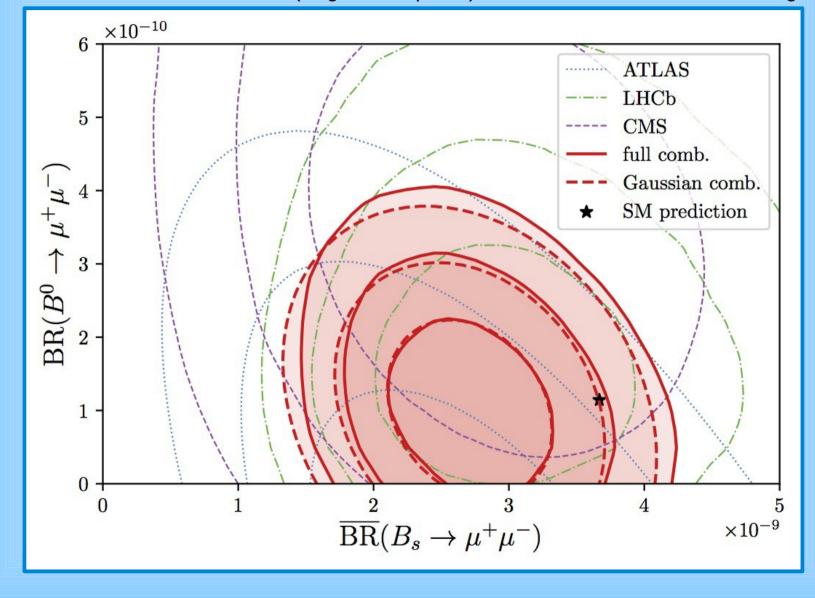
We'll know soon (?) whether this is all just a happy coincidence

$B_s \rightarrow \mu\mu$ average

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w/o CMS PAS BPH-16-004 (Aug. 2019 update)

Credit: Peter Stangl



$B_s \rightarrow \mu\mu$ average

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with CMS PAS BPH-16-004 (Aug. 2019 update)

Credit: Peter Stangl

