# Where we stand on B-decay Discrepancies <br> - Global Fits - 

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# Minimal TH considerations 

(before any fit)

## EFT considerations

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well-suited to UV-complete models

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Accurate $B_{s} \rightarrow \mu \mu$ measurement

- present single-measurement error $\simeq 20 \%$
- exp combi may soon be able to confirm or exclude deviations of $C_{10}$ of $O(10 \%)$


## More TH considerations

(3) Pattern of Lepton Universality Violation in $b \rightarrow s$

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Data now allow to disprove some, and to make more precise some other of these considerations

## Weak-Effective-Theory

## Global fits

## 1-Wilson-coeff. picture

[Aebischer et al., 2019]

| Coeff. | best fit | $1 \sigma$ | $2 \sigma$ | pull |
| :---: | :---: | :---: | :---: | :---: |
| $C_{9}^{b s \mu \mu}$ | -0.97 | $[-1.12,-0.81]$ | $[-1.27,-0.65]$ | $5.9 \sigma$ |
| $C_{9}^{\prime b s \mu \mu}$ | +0.14 | $[-0.03,+0.32]$ | $[-0.20,+0.51]$ | $0.8 \sigma$ |
| $C_{10}^{b s \mu \mu}$ | +0.75 | $[+0.62,+0.89]$ | $[+0.48,+1.03]$ | $5.7 \sigma$ |
| $C_{10}^{\prime b s \mu \mu}$ | -0.24 | $[-0.36,-0.12]$ | $[-0.49,+0.00]$ | $2.0 \sigma$ |
| $C_{9}^{b s \mu \mu}=C_{10}^{b s \mu \mu}$ | +0.20 | $[+0.06,+0.36]$ | $[-0.09,+0.52]$ | $1.4 \sigma$ |
| $C_{9}^{b s \mu \mu}=-C_{10}^{b s \mu \mu}$ | -0.53 | $[-0.61,-0.45]$ | $[-0.69,-0.37]$ | $6.6 \sigma$ |

Similar fits (w/ non-identical conclusions) performed in:
[Algueró et al., 1903.09578] [Ciuchini et al., 1903.09632]
[Datta et al., 1903.10086] [Kowalska et al., 1903.10932]
[Arbey et al., 1904.08399]

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- Two scenarios stand out: $C_{9}$ alone or $C_{9}=-C_{10} \quad(\mu \mu$-channel only)
- $C_{9}=-C_{10}$ now better than $C_{9}$ alone
- $C_{10}$ alone also ok, but $B \rightarrow K^{*} \mu \mu$ unresolved

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A concurrence of effects, mostly:

- $B_{s} \rightarrow \mu \mu$ (new average)
- $\Lambda_{b} \rightarrow \Lambda \mu \mu: A_{F B}$ and BR from LHCb
- $\Delta F=2$ (mostly $\epsilon_{\kappa}$ and $\Delta M_{s}$ )


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- $\Delta F=2$ (mostly $\epsilon_{k}$ and $\Delta M_{s}$ )

How can $\epsilon_{\kappa}$ (that doesn't depend on $C_{10}$ ) increase the $C_{10}$ significance?

## Role of correlated quantities

Consider the following two observables

- $O_{1} \equiv B R\left(B_{s} \rightarrow \mu \mu\right)$
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Now consider the corresponding terms in the $\chi^{2}$ function (with $D_{i} \equiv O_{i}^{\text {exp }}-O_{i}^{\text {th }}$ )

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\chi^{2}=\frac{D_{1}^{2}\left(C_{10}\right)}{\sigma_{1}^{2}}+\frac{D_{2}^{2}}{\sigma_{2}^{2}}+\rho_{12} \frac{D_{1}\left(C_{10}\right) D_{2}}{\sigma_{1} \sigma_{2}}
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- The $3^{\text {rd }}$ term depends on $C_{10}$ through $D_{1}$
- This term will influence the $C_{10}$ value, depending on $\operatorname{sign}\left(D_{2} \rho_{12}\right)$


## $C_{9}$ vs. $C_{9}=-C_{10}$



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- Before Moriond
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## Note

$a C_{g}{ }^{\text {uni. }}$ component would shift $b \rightarrow s \mu \mu$ data but not $R_{\left.K_{( }\right)}$


## Notes

$y$-axis: $\mu$-specific shift in $C_{9}=-C_{10}$
x-axis: additional, lepton-univ. shift in $\mathrm{C}_{9}$ only

After Moriond
Data tend to prefer $C_{g}{ }^{\text {univ. }} \neq 0$

## Both $C_{9}=-C_{10}$ and $C_{9}$ univ.

well justified above the EW scale.

## The SMEFT picture

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$\square$
Dynamics below $\wedge$ described by ops.

- constructed with SM fields only
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Then, two clear possibilities to generate the above pattern:
(1) Contributions to muonic $C_{9}=-C_{10}$ may come from SMEFT ops. directly matching onto $O_{9,10}$

$$
\begin{aligned}
& {\left[O_{L Q}^{(1)}\right]_{2223}=\bar{L}_{2} \gamma^{\lambda} L_{2} \cdot \bar{Q}_{2} \gamma_{\lambda} Q_{3}} \\
& {\left[O_{L Q}^{(3)}\right]_{2223}=\bar{L}_{2} \gamma^{\lambda} \sigma^{a} L_{2} \cdot \bar{Q}_{2} \gamma_{\lambda} \sigma^{a} Q_{3}}
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- Case $f=\tau$ allows natural connection (right sign \& size) with $\left[O_{L Q}^{(1)}\right]_{3323} \&\left[O_{L Q}^{(3)}\right]_{3323}$ responsible for $b \rightarrow c \tau v$
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[Crivellin, Greub, Müller, Saturnino, 2018]
- Caveat: need $\left[C_{L Q}^{(1)}\right]_{3323} \simeq\left[C_{L Q}^{(3)}\right]_{3323}$
to avoid $B \rightarrow K\left(^{*}\right) v v$ constraint [Buras-Girrbach-Niehoff-Straub, 2014]
$\left[C_{L Q}^{(1)}\right]_{3323}=\left[C_{L Q}^{(3)}\right]_{3323} \quad$ vs. $\quad\left[C_{L Q}^{(1)}\right]_{2223}=\left[C_{L Q}^{(3)}\right]_{2223}$


Before Moriond (dashed)
$R_{\left.K^{*}\right)}$ (blue) and $b \rightarrow s \mu \mu$ (orange)
were in perfect agreement
So
in a region close to 0 in the $x$-axis
$\Rightarrow R_{D\left(^{*}\right)}$ not explained


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$R_{\left.K_{( }\right)}$and $b \rightarrow s \mu \mu$ intersect
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## Conclusions

Post-Moriond updates imply a nicely coherent TH picture, with

- $R_{K} \& R_{K^{*}}<$ by O(20\%) than SM
- $R_{D} \& R_{D^{*}}>$ by $O(10 \%)$ (not more) than $S M$
- $B R\left(B_{s} \rightarrow \mu \mu\right)<$ by $O(20 \%)$ than $S M$


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We'll know soon (?) whether this is all just a happy coincidence

with CMS PAS BPH-16-004 (Aug. 2019 update)


