

# Constraining PDFs through exclusive processes at the LHCb

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## Work presented appears in

*'Towards a determination of the low  $x$  gluon via exclusive  $J/\psi$  production'*

CAF, S.P.Jones, A.D.Martin, M.G.Ryskin, T.Teubner, 1907.06471

*'How to include exclusive  $J/\psi$  production data in global PDF analyses'*

CAF, S.P.Jones, A.D.Martin, M.G.Ryskin, T.Teubner, 1908.08398

Submitted to EPJC

## Earlier works

*'Exclusive  $J/\psi$  and  $\Upsilon$  production and the low  $x$  gluon'*

S.P.Jones, A.D.Martin, M.G.Ryskin, T.Teubner, 1507.06942

*'The exclusive  $J/\psi$  process at the LHC tamed to probe the low  $x$  gluon'*

S.P.Jones, A.D.Martin, M.G.Ryskin, T.Teubner, 1610.02272

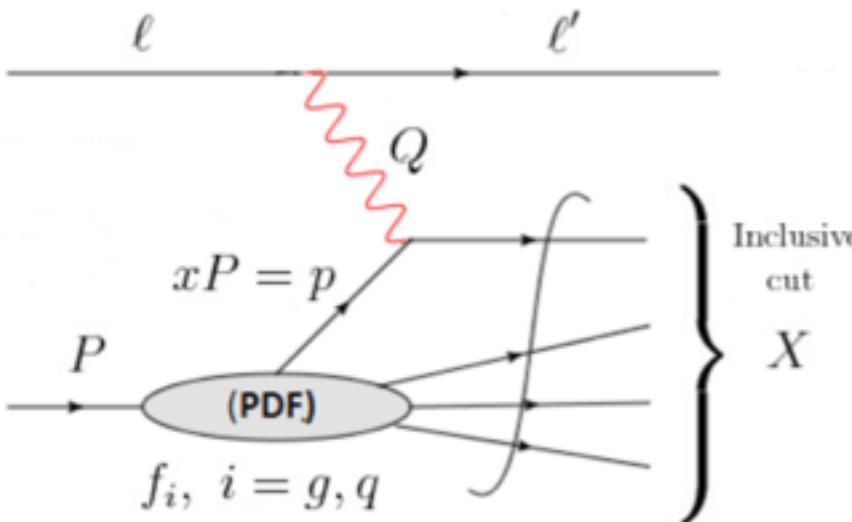
# Introduction

- Inclusive processes do not well constrain small  $x$ /Regge limit domain of PDFs
- Exclusive processes offer sensitive probe of this domain but as of yet not included in global analyses PDF determination - why?
  1. Off forward kinematics imply susceptibility to *GPD* over conventional PDFs
  2. Reliability and stability of theoretical predictions

- As higher CM energies are realised at LHC, pushed towards small  $x$  domain,  $W \sim 1/x$

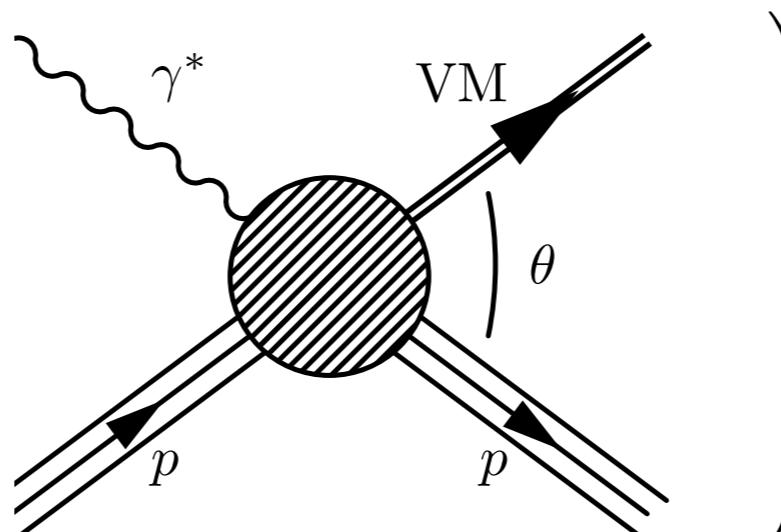
$$\left. \frac{d\sigma}{dt}(\gamma^* p) \right|_{t=0} = \frac{\Gamma_{ee}^{J/\psi} M_{J/\psi}^3 \pi^3}{48\alpha_{\text{em}}} \left[ \frac{\alpha_s(\bar{Q}^2)}{\bar{Q}^4} R_g x g(x, \bar{Q}^2) \right]^2 \left( 1 + \frac{Q^2}{M_{J/\psi}^2} \right)$$

**Inclusive** - included in  
global parton analyses

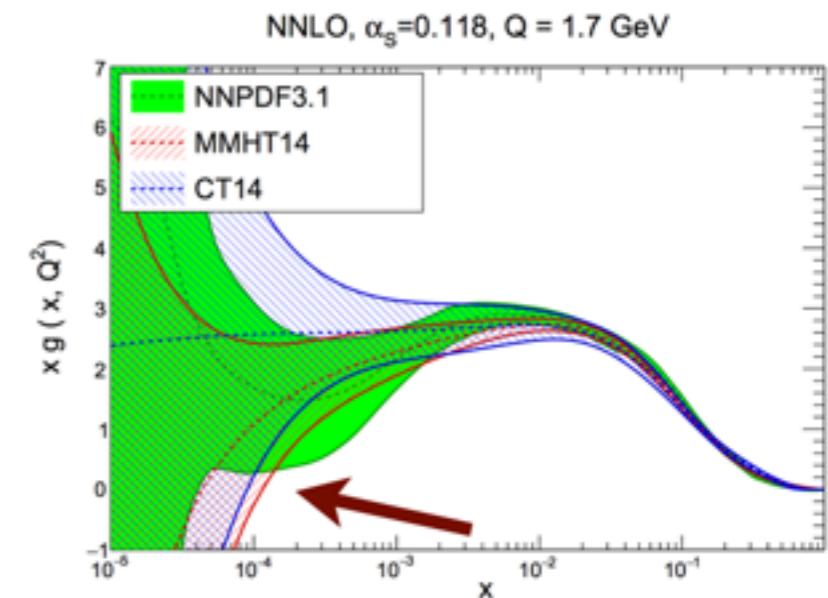


e.g DIS

**Exclusive** - can we use  
the data?

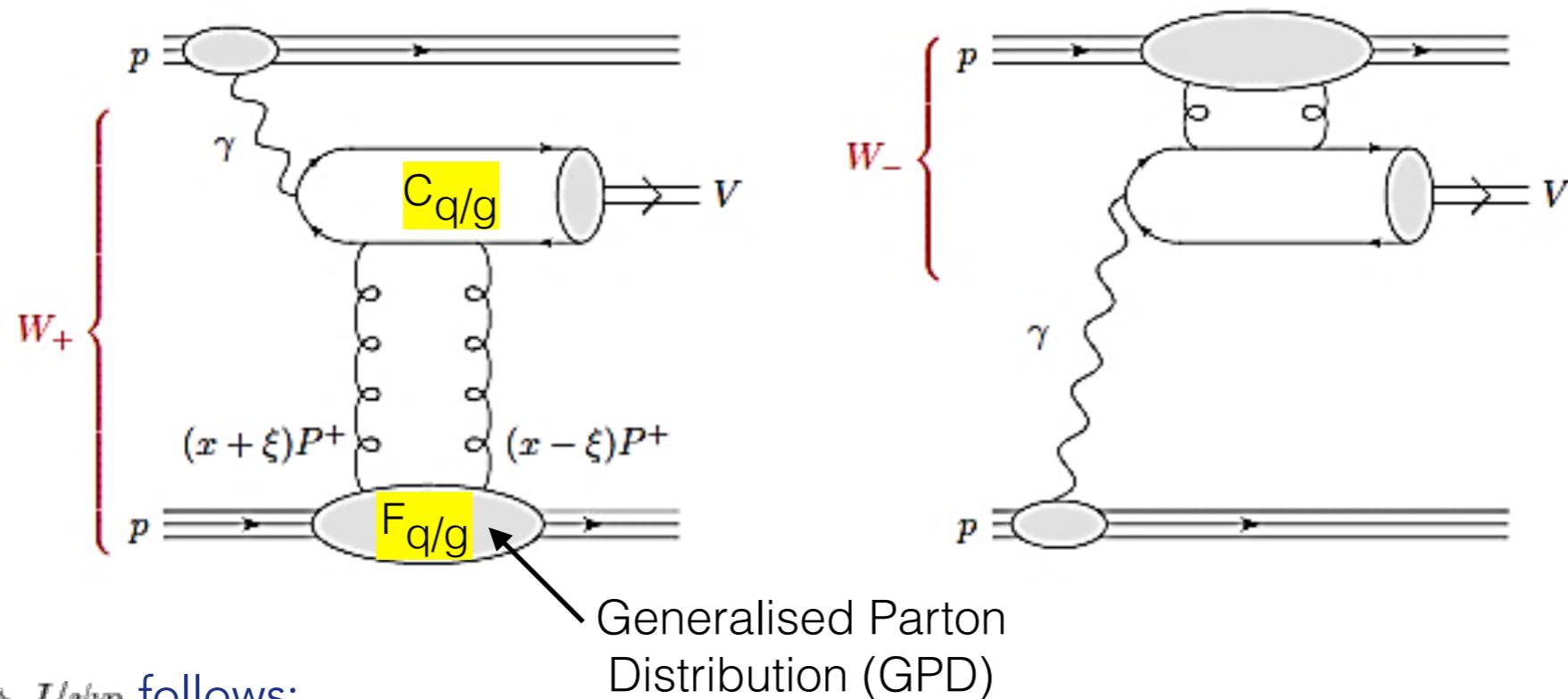


Ryskin 1993



[nnpdf.mi.infn.it](http://nnpdf.mi.infn.it)

# General Set up and assumptions



Setup for  $\gamma p \rightarrow J/\psi p$  follows:

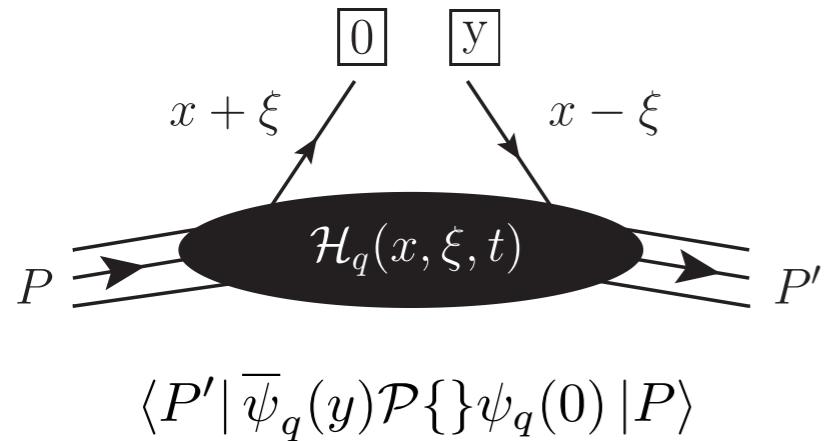
Ivanov, Schäfer, Szymanowski, Krasnikov, 04

- Assume a factorisation  $F_{q/g} \otimes C_{q/g} \otimes \phi_{Q\bar{Q}}^V$
- Leading zeroth order term in rel. velocity (NRQCD)
- Colour singlet exchange between hard and soft sectors

$$A \propto \int_{-1}^1 dx \left[ C_g(x, \xi) F_g(x, \xi) + \sum_{q=u,d,s} C_q(x, \xi) F_q(x, \xi) \right]$$

# GPDs and the Shuvaev transform

GPDs generalise PDFs: outgoing/incoming partons carry different momentum fractions Müller 94; Radyushkin 97; Ji 97



**Shuvaev:** Relates GPDs to PDFs at small  $x$  under physically motivated assumptions c.f analyticity

Shuvaev 99 Martin et al. 09

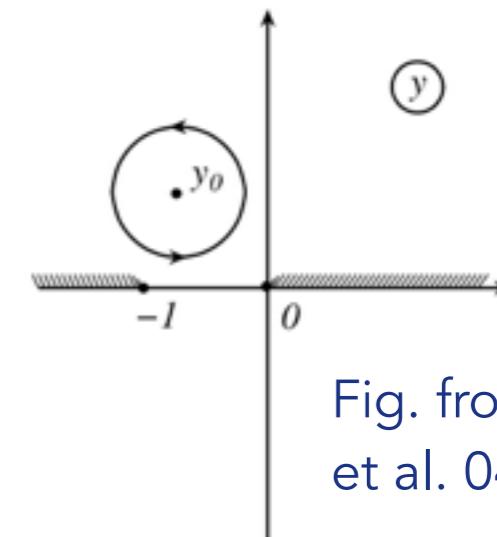


Fig. from Ivanov et al. 04

Idea: Conformal moments of GPDs = Mellin moments of PDFs

(up to corrections of order  $x^{1/2}$ )

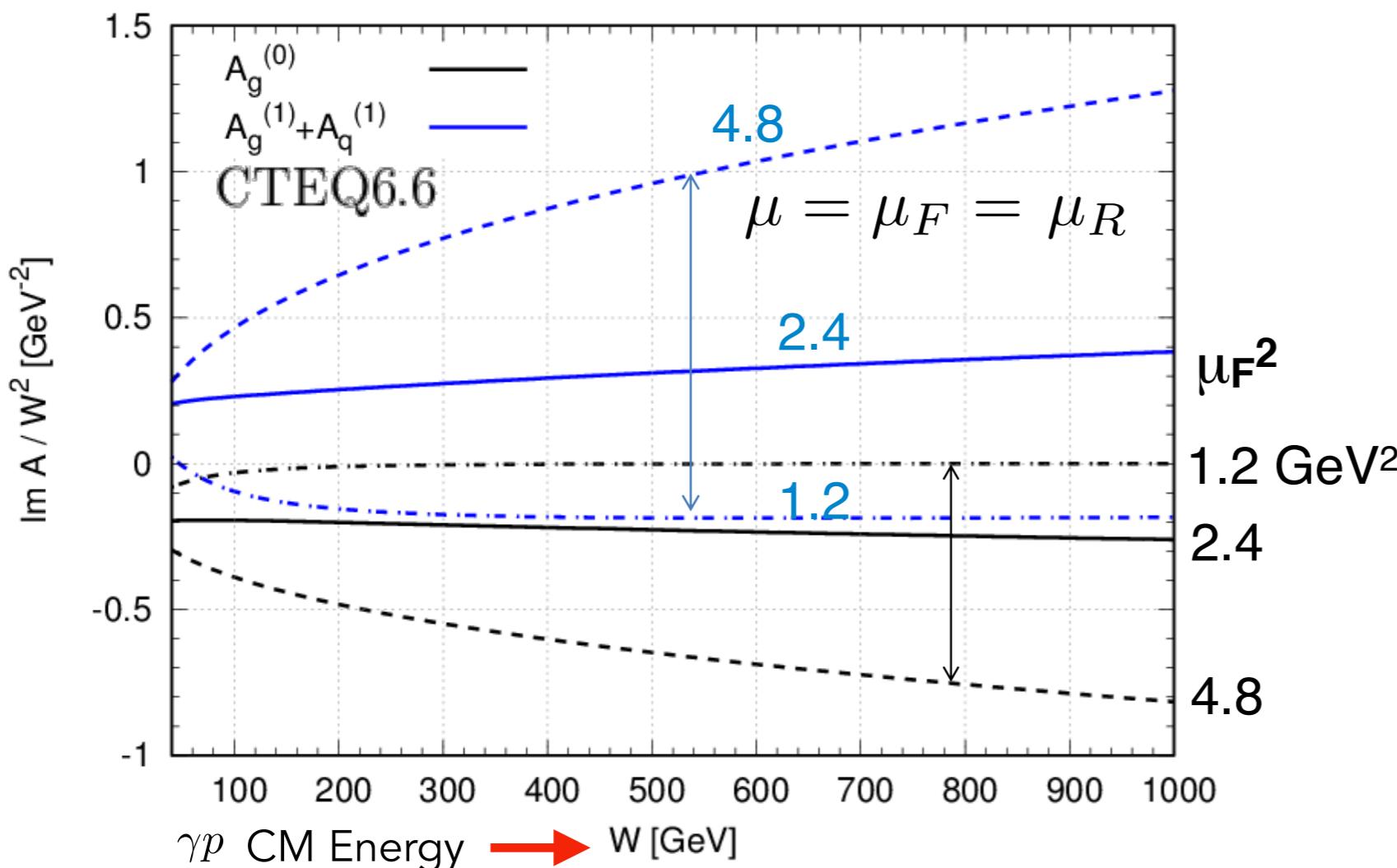
- Construct GPD grids in multidimensional parameter space  $x, \xi/x, q^2$  with forward PDFs from LHAPDF
- Costly computationally due to slowly converging double integral transform
- Regge theory considerations => Shuvaev transform valid in space like (DGLAP) region only. In time like (ERBL) region imaginary part of coefficient is zero

# Stability of prediction I

NLO in MSbar scheme

D. Ivanov, B.Pire, L.Szymanowski, J.Wagner, 1411.3750  
S.P.Jones, PhD thesis, Liverpool (2014)

- A. **Bad perturbative convergence**  $|NLO_{\text{correctn.}}| > |LO|$  and
- B. **Strong dependence on scale  $\mu_F$**  **opp. sign**



**Disclaimer:** Plots generated using existing global partons.  
Here, CTEQ6.6

Can do better...

# Stability of prediction II

## 'Scale Fixing'

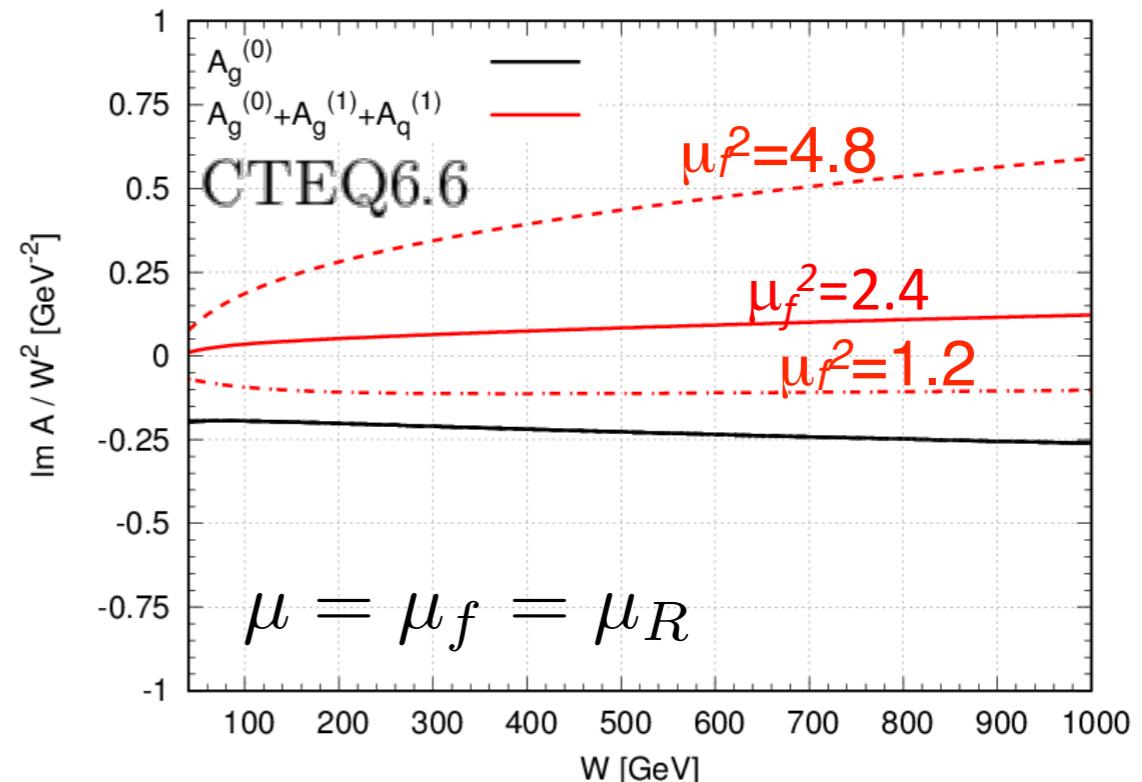
`Optimal' factorisation scale  $\mu_F = m$   
eliminates large logs at NLO

S.P.Jones, A.D.Martin, M.G.Ryskin, T.Teubner, 1507.06942

Resummation of  $(\alpha_s \ln(1/\xi) \ln(\mu_F/m)^n)$   
terms into LO PDF, leaving  
remnant NLO coefficient  
and residual,  $\mu_f$ , scale  
dependence

$$A(\mu_f) = C^{\text{LO}} \times \text{GPD}(\mu_F) + C^{\text{NLO}}(\mu_F) \times \text{GPD}(\mu_f)$$

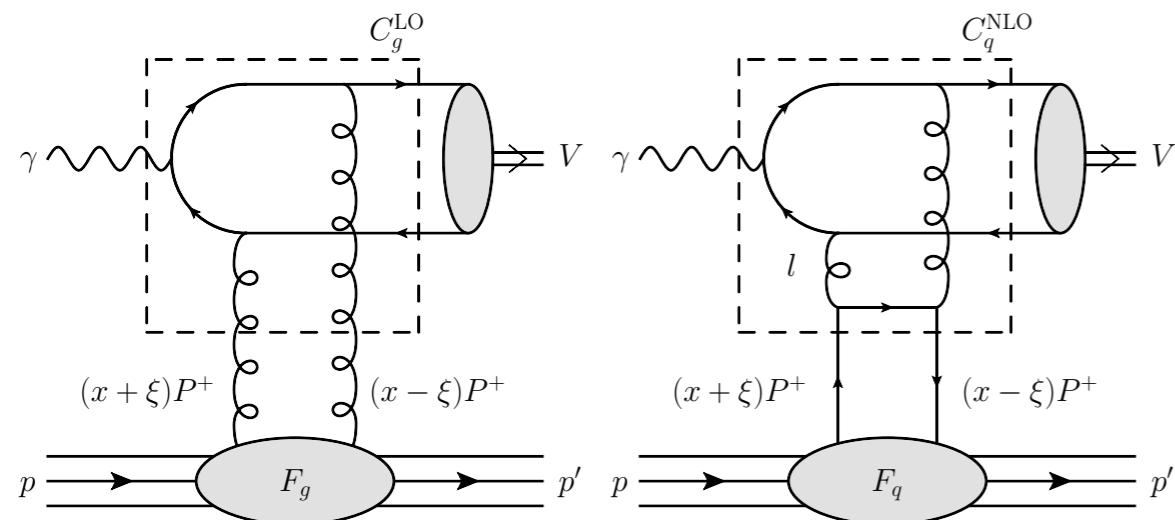
Fix:  $\mu_F^2 = 2.4 \text{ GeV}^2$



Look for another sizeable correction that can reduce variations  
further -> implementation of a '**Q0**' cut

# Stability of prediction III

‘Q0’ cut S.P.Jones, A.D.Martin, M.G.Ryskin, T.Teubner, 1610.02272



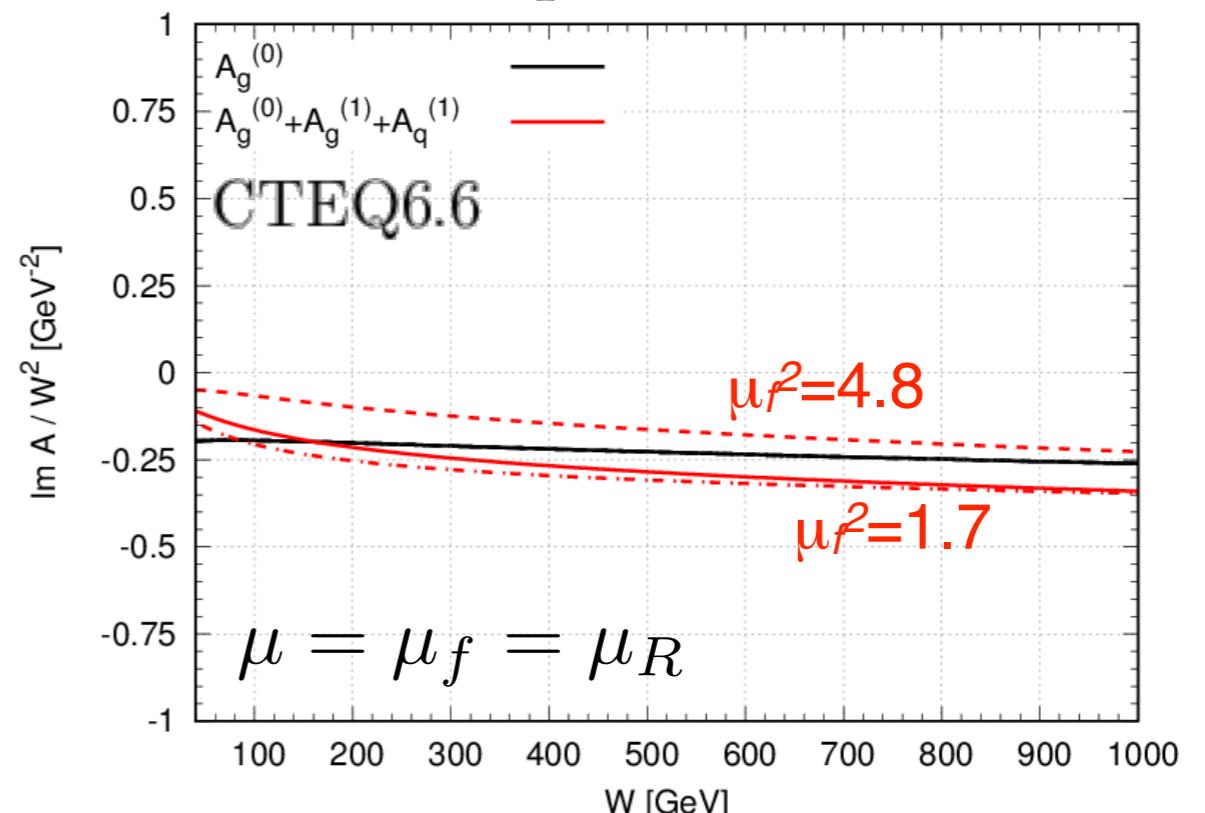
Typically power suppressed, but sizeable here

$$\mathcal{O}(Q_0^2/M_{J/\psi}^2)$$

→  
How do these predictions map onto the data at HERA and LHCb?

Subtract DGLAP contribution  
NLO ( $|q^2| < Q_0^2$ )  
from known NLO MSbar coefficient function to avoid a double count with input GPD at Q0.

Fix:  $\mu_F^2 = 2.4 \text{ GeV}^2$



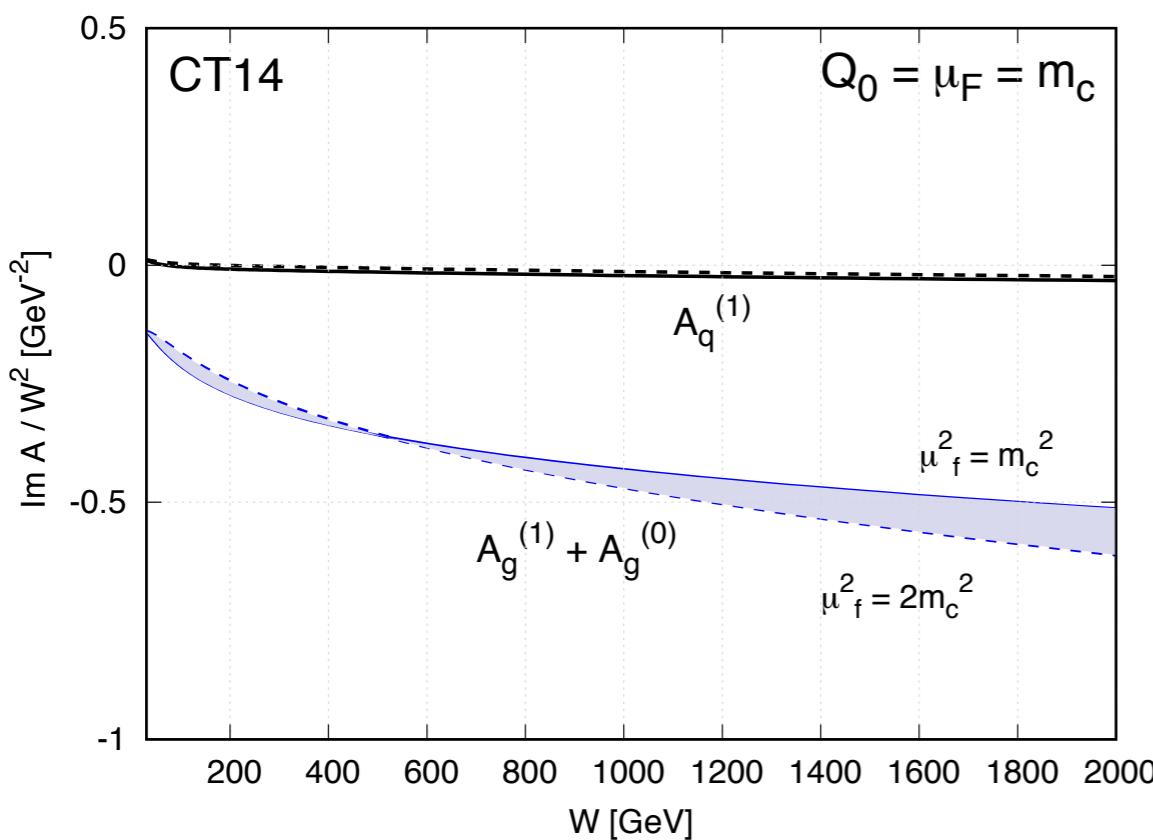
# Philosophy

CAF, S.P.Jones, A.D.Martin, M.G.Ryskin, T.Teubner, 1907.06471 & 1908.08398

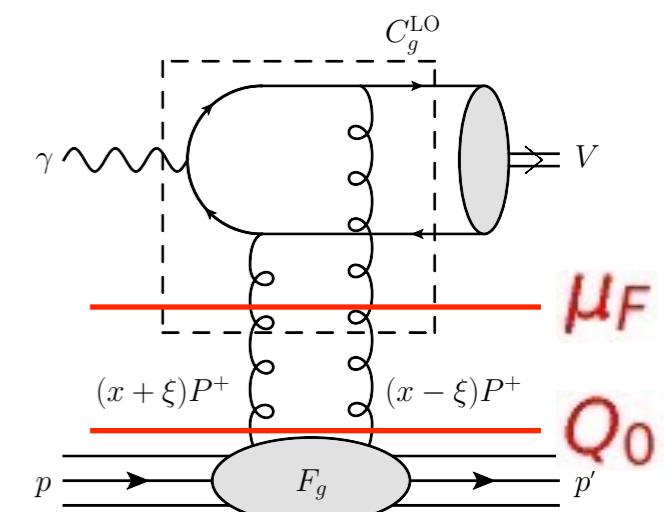
- Philosophy:
- I) Achieve cross section stability in small parameter space
  - II) Any deficiency in LHCb regime attributed to small  $x$  global fitter behaviour

Choices:

- $Q_0^2$  IR transition parameter
- $\mu_F = m_c$  to resum large logarithm at high energy
- $\mu_f = \mu_R$  in accordance with BLM prescription

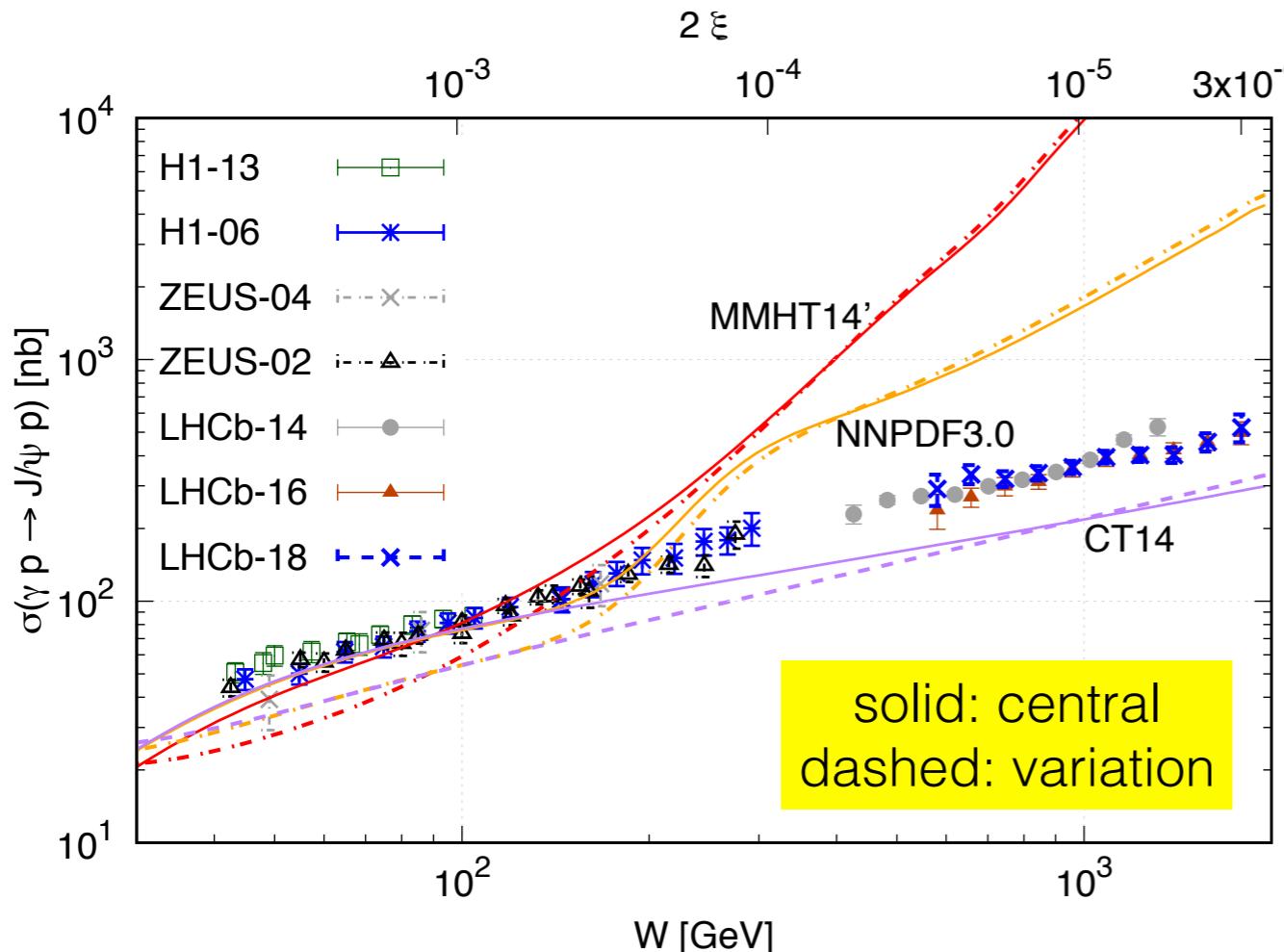


- Exclusive J/psi probe of gluon density, quark contribution effectively absorbed by  $Q_0$  subtraction



# Cross section stability

- Regge based arguments => imaginary part of amplitude dominant
- Nonetheless, may restore real part through dispersion relation - numerically evaluate perturbation  $\frac{\text{Re}\mathcal{M}}{\text{Im}\mathcal{M}} \sim \frac{\pi}{2} \lambda = \frac{\pi}{2} \frac{\partial \ln \text{Im}\mathcal{M}/W^2}{\partial \ln W^2}$  with  $\mathcal{M} \sim x^{-\lambda}$
- Achieved good cross section stability at  $Q_0 = \mu_F = m_c$  and with variations w.r.t  $\mu_f^2, \mu_R^2 \in [m_c^2, 2m_c^2]$



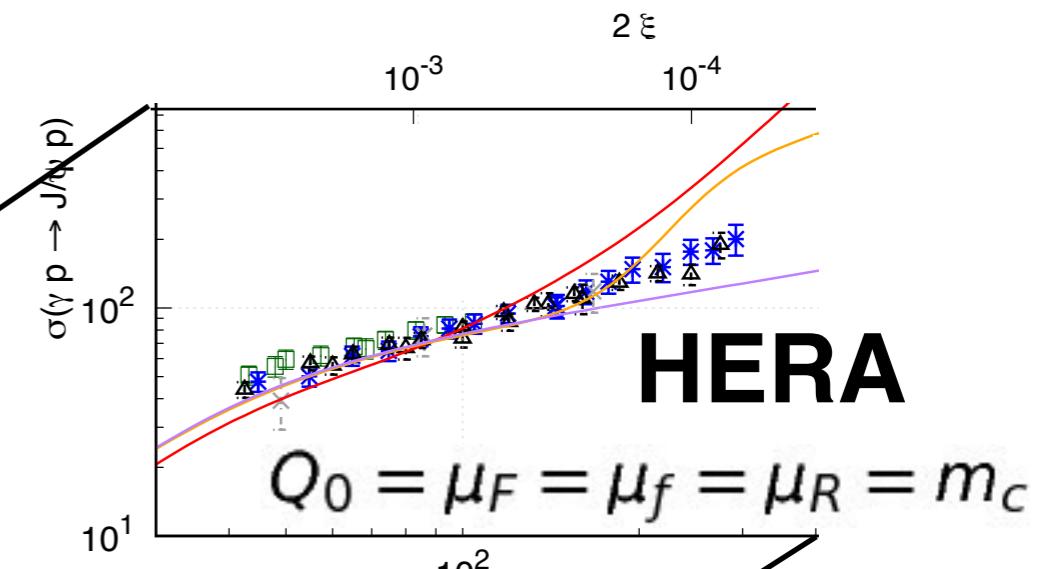
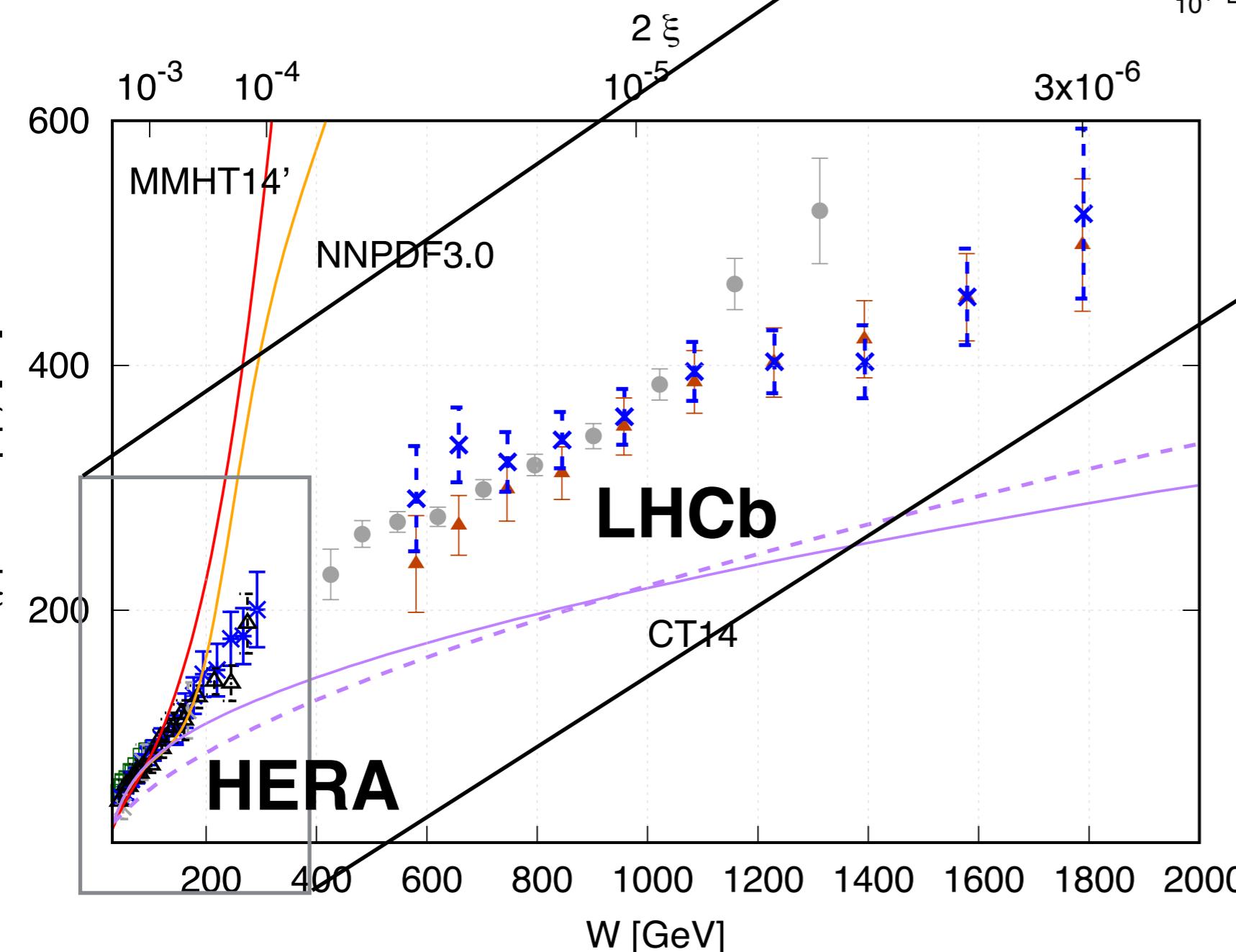
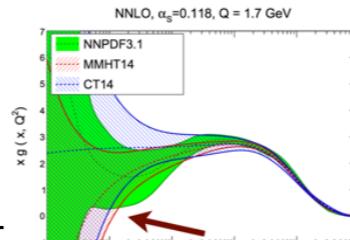
**Repeat Disclaimer:**  
Convoluting with existing global partons.  
Here,  
MMHT14,  
NNPDF3.0 &  
CT14

Plot demonstrates good scale stability of our NLO predictions in LHCb regime

Predictions at optimal scale (solid) agree better with HERA data

# Towards the bigger picture

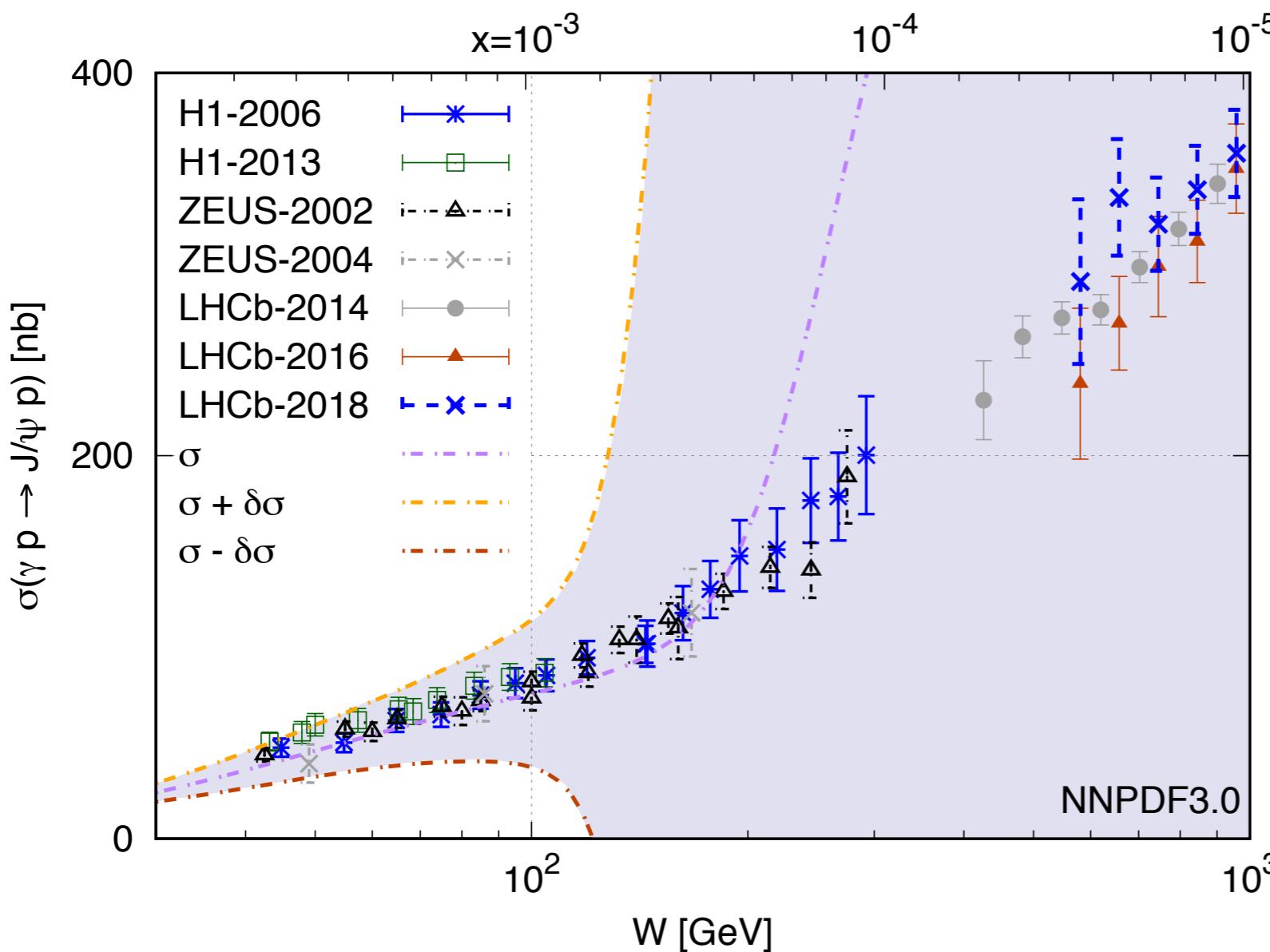
- NNPDF3.0 and MMHT central values overshoot HERA data towards  $x \sim 10^{-4}$ , but is covered by 1sigma error band (see next)



- But wildly different behaviours in LHCb regime! Important message....

Error budgets: errors due to parameter variations in global fits >> experimental uncertainty and scale variations in the theoretical result

..... exclusive data now in a position to readily improve global analyses



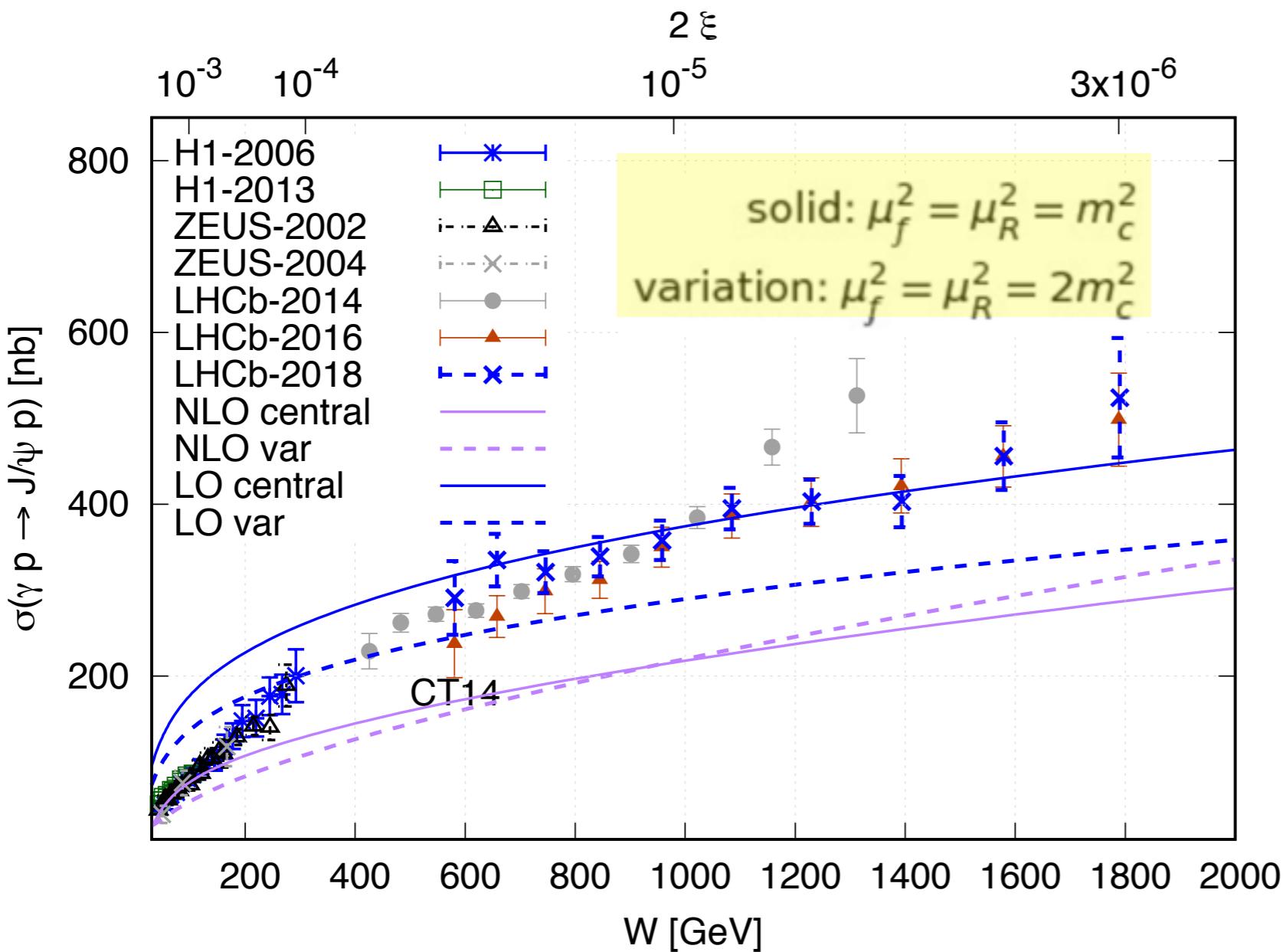
*Exclusive* LHCb data  
will constrain small  $x$   
growth whilst *exclusive*  
HERA data will improve  
determination of partons  
in regime with data  
constraints already from  
diffractive DIS HERA  
data

# Summary

- Naive MSbar NLO coll. fact. result unreliable and unstable
- Systematic taming via ‘Q0’ cut and resummation of large logarithmic contributions collectively reduce wild scale variations
- Mapped predictions to cross section level with good stability observed and central values in agreement of data within 1sigma error bands
- MMHT14’ and NNPDF3.0 largely overshooting data in LHCb regime
- Upshot: Exclusive data constrains the growth of these partons, certain small  $x$  behaviours must be rejected to permit reconciliation with our data constraint.
- In a position to finally use this exclusive data in a global fitter framework

**Thank you**

# Backups



Achieve better stability in going from LO  $\rightarrow$  NLO

NLO correction in better agreement with HERA

Reasonable description of LHCb regime at LO interpreted as coincidental: indeed curve is flatter with wrong  $W$  growth

# Shuvaev Transform

**Full Transform:**

$$\mathcal{H}_q(x, \xi) = \int_{-1}^1 dx' \left[ \frac{2}{\pi} \text{Im} \int_0^1 \frac{ds}{y(s) \sqrt{1 - y(s)x'}} \right] \frac{d}{dx'} \left( \frac{q(x')}{|x'|} \right),$$

$$\mathcal{H}_g(x, \xi) = \int_{-1}^1 dx' \left[ \frac{2}{\pi} \text{Im} \int_0^1 \frac{ds(x + \xi(1 - 2s))}{y(s) \sqrt{1 - y(s)x'}} \right] \frac{d}{dx'} \left( \frac{g(x')}{|x'|} \right),$$

$$y(s) = \frac{4s(1 - s)}{x + \xi(1 - 2s)}.$$