

# Extended Gauge Mediation in the NMSSM with Displaced LHC Signals

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Based on:

MB, N. Desai, C. Hugonie, R. Ziegler, arXiv:1810.05618

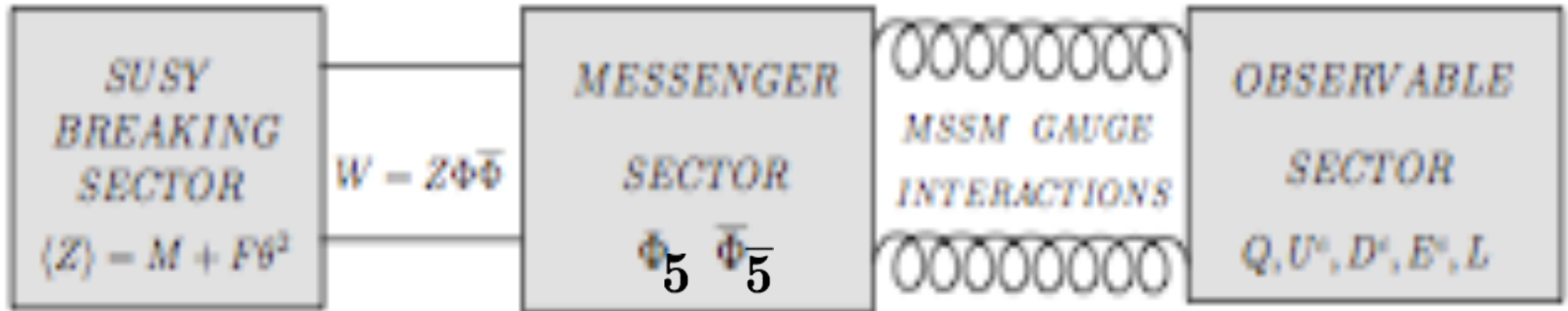
B. Allanach, MB, G. Cottin, N. Desai, C. Hugonie, R. Ziegler, arXiv:1606.03099

B. Allanach, MB, C. Hugonie, R. Ziegler, arXiv:1502.05836



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# Minimal GM with 125 GeV Higgs



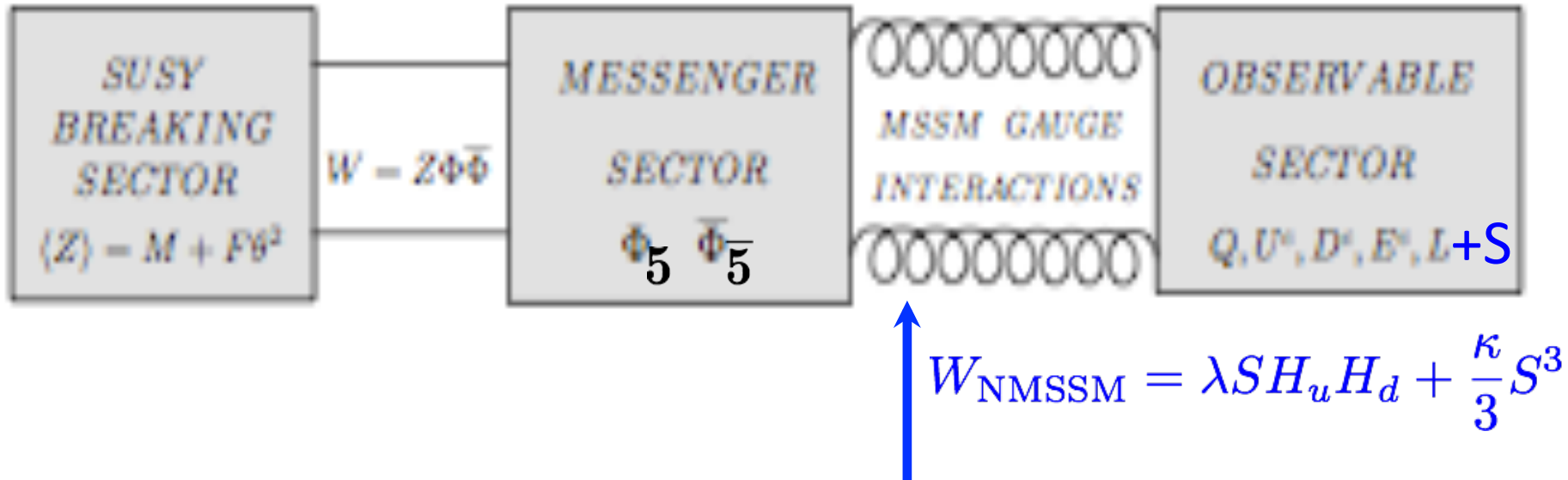
Very predictive framework that solves SUSY flavor problem

But with other **problems**:

- **A-terms suppressed**: very unnatural sparticle spectrum with no chance to be probed at the LHC
- $\mu$  and  $B_\mu$  typically generated at same loop order, therefore  $B_\mu$  too large for correct EWSB ( **$\mu$ - $B_\mu$  problem**)

Need to go to non-minimal models of gauge mediation  $\rightarrow$  NMSSM

# Next-to-minimal gauge mediation



- New contributions to the Higgs mass
- $\mu$  and  $B_\mu$  generated dynamically

However: naive combination NMSSM + minimal GM does not provide correct EWSB (because singlet soft mass too small)

# The minimal model of gauge mediation in NMSSM:

Delgado, Giudice, Slavich '07

Main ingredient:



direct coupling of the singlet to (two pairs) of messengers

$$W_{\text{DGS}} = S \left( \xi_D \bar{\Phi}_1^D \Phi_2^D + \xi_T \bar{\Phi}_1^T \Phi_2^T \right) \quad \xi_D(M_{\text{GUT}}) = \xi_T(M_{\text{GUT}}) \equiv \xi$$

generates NMSSM A-terms and singlet masses (correct EWSB)

very predictive model with only 4 parameters:

$\lambda, \xi, \tilde{m}, M$

soft SUSY breaking scale  messenger scale 

Note:  $\tan \beta$  is NOT a free parameter

# The minimal model of gauge mediation in NMSSM

Delgado, Giudice, Slavich '07

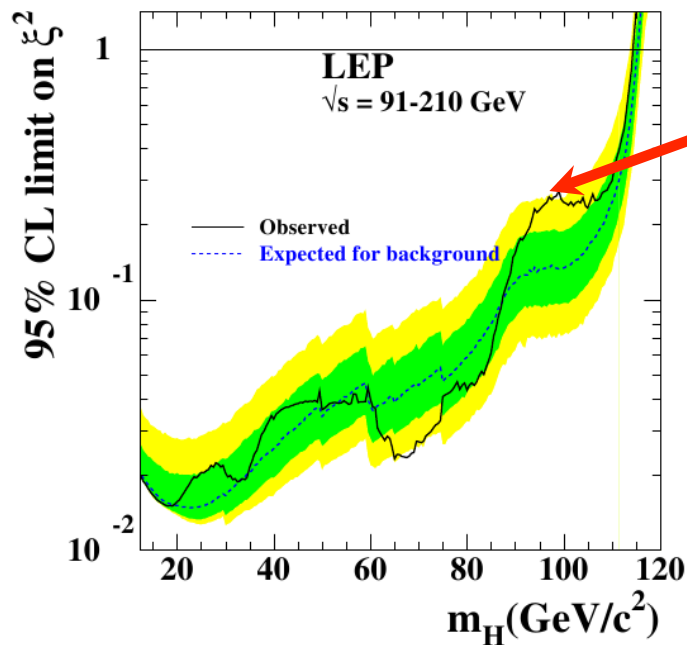
**Conclusion of the original DGS paper:** sparticles are heavier than in Minimal Gauge Mediation in MSSM due to the Higgs mass constraint

Caveat:

The above assumes **singlet heavier than the Higgs**

# Enhanced Higgs mass from mixing with a light singlet

$$\begin{pmatrix} m_h^2 & m_{hs}^2 \\ m_{hs}^2 & m_s^2 \end{pmatrix} \xrightarrow{m_h^2 > m_s^2} m_{h_2}^2 \approx m_h^2 + \frac{m_{hs}^4}{m_h^2}$$



can be large close to the LEP excess

# How large the mixing contribution to the Higgs mass can be?

$$M^2 = \begin{pmatrix} M_{hh}^2 & M_{hs}^2 \\ M_{hs}^2 & M_{ss}^2 \end{pmatrix}$$

MSSM Higgs mass

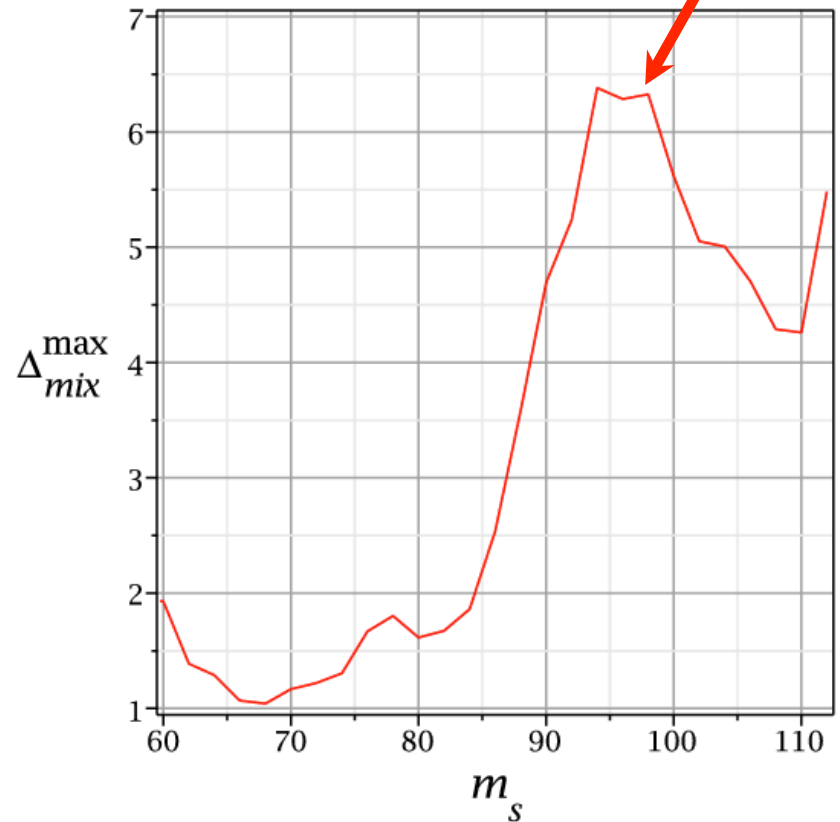
$$m_h \equiv M_{hh} + \Delta_{\text{mix}}$$

$$\Delta_{\text{mix}} \approx \frac{\bar{g}_s^2}{2} \left( m_h - \frac{m_s^2}{m_h} \right)$$

5-6 GeV correction to the Higgs mass is possible when  $m_s$  close to the LEP excess

MB, Olechowski, Pokorski '13

the LEP excess



# Light singlet scenario in DGS model

Allanach, MB, Hugonie, Ziegler '15

Maximizing the contribution to the Higgs mass from mixing essentially fixes most model parameters

$$\left. \begin{array}{l} \cos \theta \approx 0.9 \longrightarrow \lambda \sim 10^{-2} \\ m_{h_1} \approx 90 - 100 \text{ GeV} \longrightarrow \xi \sim 10^{-2} \end{array} \right\} \longrightarrow \begin{array}{l} 100 \text{ GeV Singlino NLSP} \\ 20\text{-}40 \text{ GeV pseudoscalar} \end{array}$$
$$m_{h_2} \approx 125 \text{ GeV} \longrightarrow \tilde{m} \sim 1 \text{ TeV}$$

Only the messenger scale remains free and determines collider phenomenology

The 125 GeV Higgs mass obtained for stops and gluinos below 2 TeV



# LHC Phenomenology

new feature is Singlino NLSP & Gravitino LSP

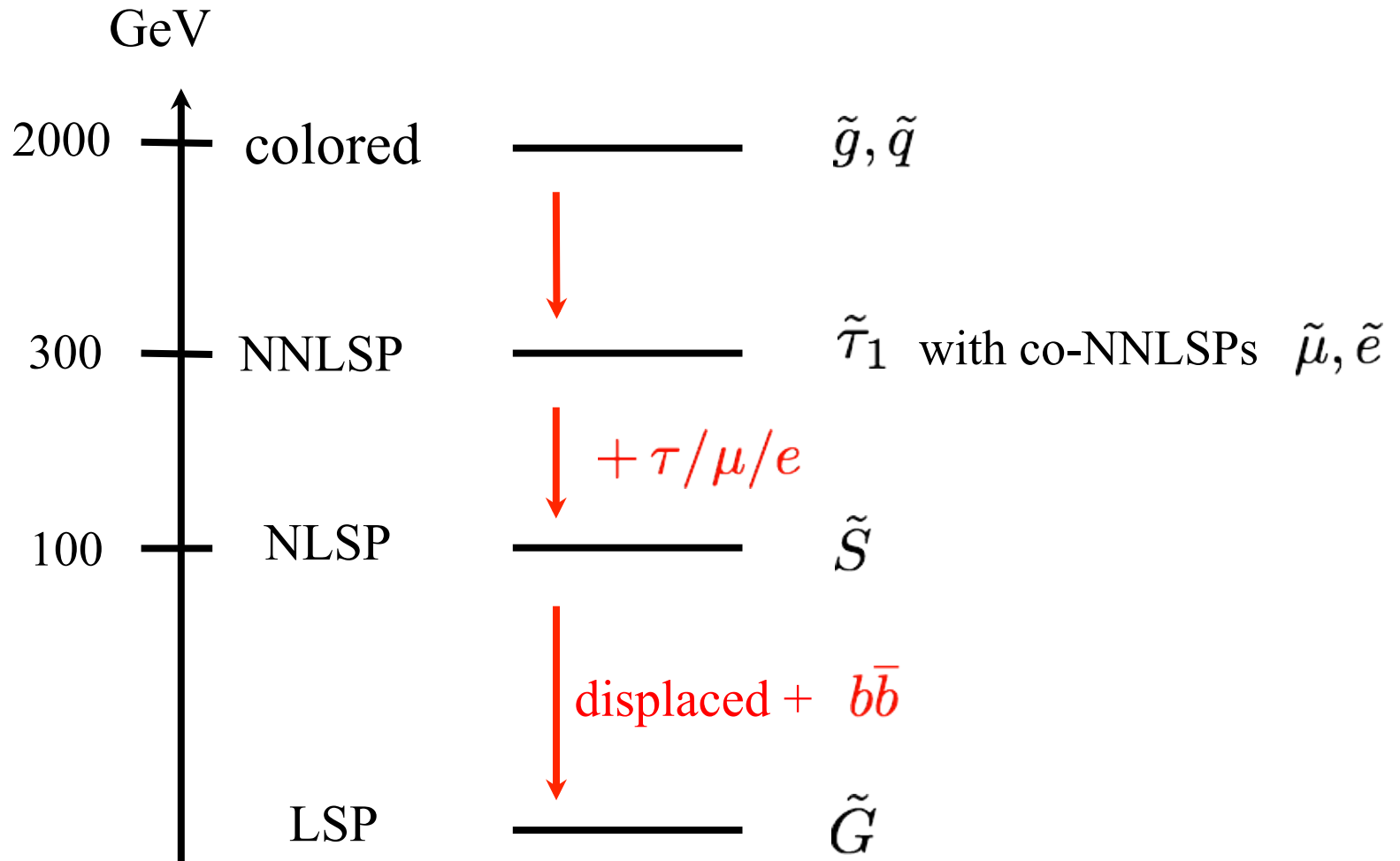
$$\tilde{N}_1 \rightarrow a_1 \tilde{G} \rightarrow b\bar{b}\tilde{G}$$

Messenger scale determines>NNLSP  
(bino or stau) and singlino decay length

$$c\tau_{\tilde{N}_1} \approx 2.5 \text{ cm} \left( \frac{100 \text{ GeV}}{M_{\tilde{N}_1}} \right)^5 \left( \frac{M}{10^6 \text{ GeV}} \right)^2 \left( \frac{\tilde{m}}{\text{TeV}} \right)^2$$

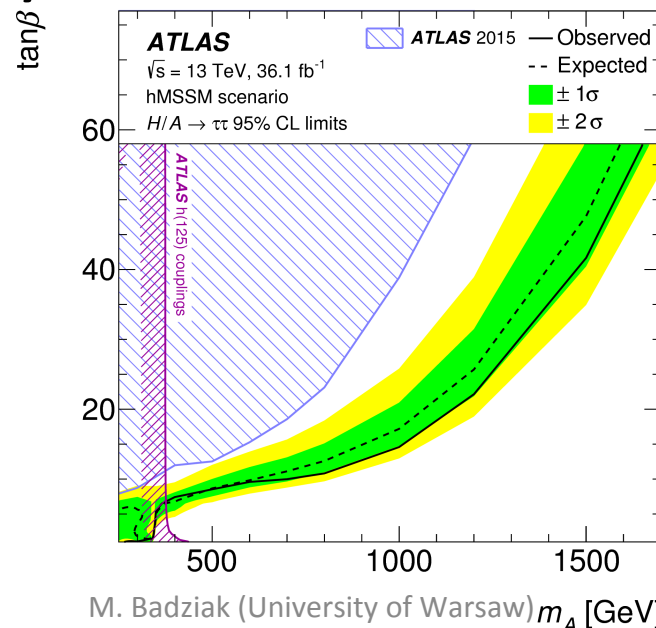
Singlino and Gravitino essentially decoupled:  
all SUSY decay chains to LSP proceed through  
NNLSP and NLSP

# Displaced vertices from long-lived singlino for $M < 10^7$ GeV



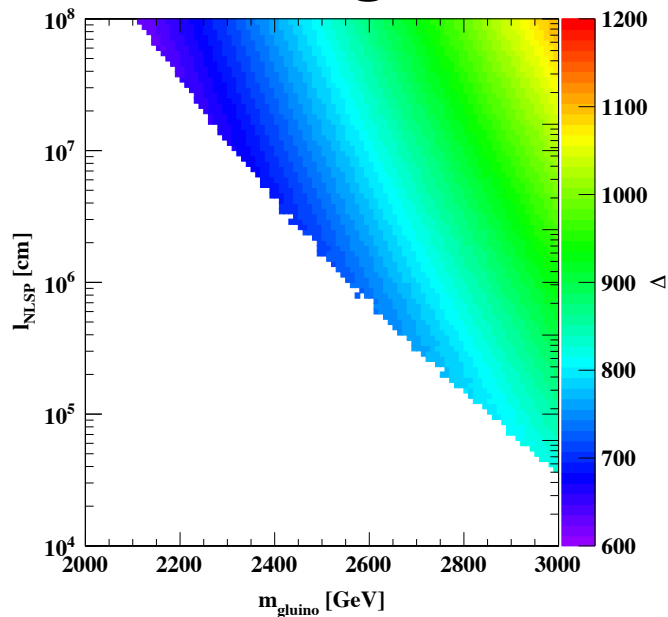
# Heavy Higgs searches hurt DGS model

- The DGS model in a light singlet scenario predicts large  $\tan \beta \gtrsim 30 \div 40$
- The MSSM-like Higgses are strongly constrained by the LHC searches



# Heavy Higgs searches hurt DGS model

- Correlation between MSSM-like Higgs and sparticle masses results in a strong lower limit on gluino mass



Glauino mass is pushed above 3 TeV for singlino lifetimes leading to displaced vertices

# Extended gauge mediation in NMSSM

MB, Desai, Hugonie, Ziegler '18

For direct Higgs-messenger couplings e.g.

$$\Delta W = \lambda_{ij} Q_i U_j \Phi_{H_u}$$

- Sparticle masses may not be correlated with MSSM-like Higgs mass
- The Higgs mass may be additionally enhanced via stop mixing due to contribution to A-terms

# The U model

One pair of messengers:

5 free parameters (1 more than DGS):

$$W_U = \lambda_t Q_3 U_3 \Phi_u + \lambda_{S_d} S \Phi_u H_d$$

$\lambda, \tilde{m}, M, \lambda_t, \lambda_{S_d}$

➤ Large  $\lambda_t$  makes the MSSM-like Higgs heavier:

$$m_{a_2}^2 \approx -m_{H_u}^2 \supset 9y_t^2 \lambda_t^2 \tilde{m}^2$$

➤ allows for stop-mixing correction to the Higgs mass

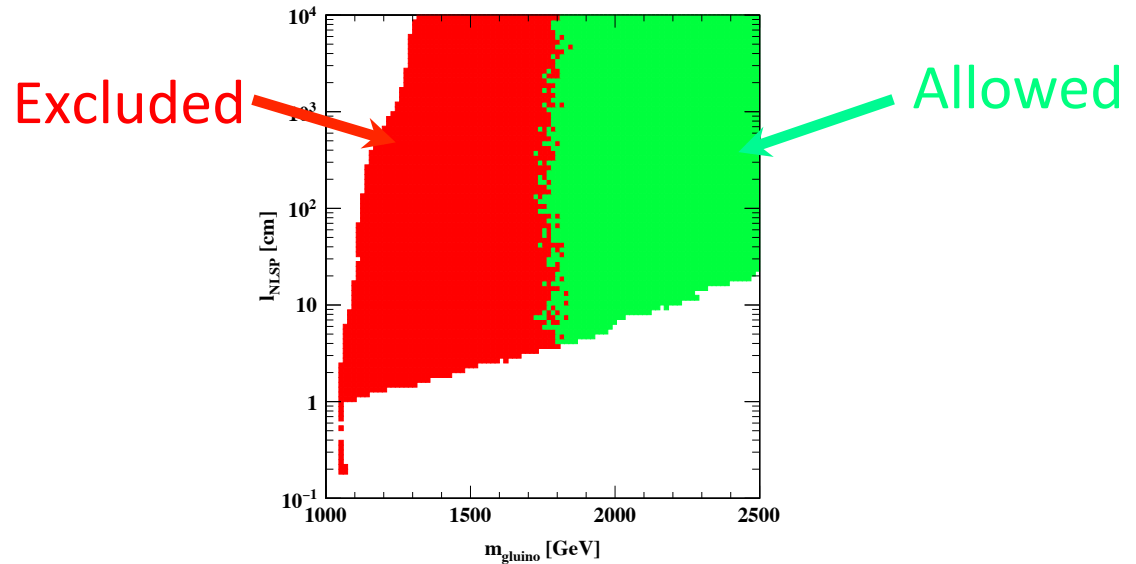
$$A_t \supset -3\lambda_t^2 \tilde{m}$$

➤ suppresses  $\tan \beta$  (e.g. due to RG effect of  $A_t$  on  $A_\lambda$ )

**Heavy Higgs searches no longer constrain SUSY spectrum!**

# The U model

Very light gluinos allowed by the (light and heavy) Higgs data



- Gluino mass down to 1.7 TeV allowed by SUSY searches at the LHC
- singlino decay length as small as  $O(1 \text{ cm})$

# The DGSU model

Combines U model with DGS model (2 pairs of messengers):

$$W_{\text{DGSU}} = S \left( \xi_D \Phi_u^{(1)} \Phi_d^{(2)} + \xi_T \Phi_T^{(1)} \Phi_{\overline{T}}^{(2)} \right) + \lambda_t Q_3 U_3 \Phi_u^{(2)} + \lambda_{S_d} S \Phi_u^{(2)} H_d$$

$$\xi_D(M_{\text{GUT}}) = \xi_T(M_{\text{GUT}}) \equiv \xi$$

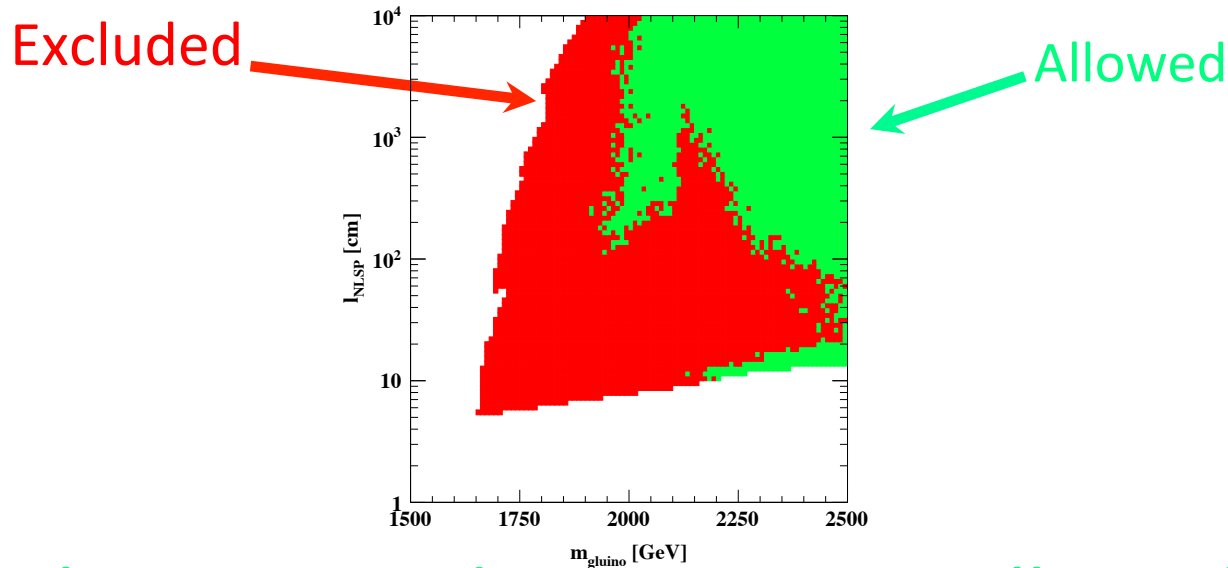
$$\lambda_{S_d}(M) y_t(M) = \lambda_t(M) \lambda(M)$$

Higgs-messenger mixing condition

- Very predictive (5 free parameters)
- The same singlino/singlet sector leading to displaced pheno as in the DGS model but with much lighter colored sparticles!



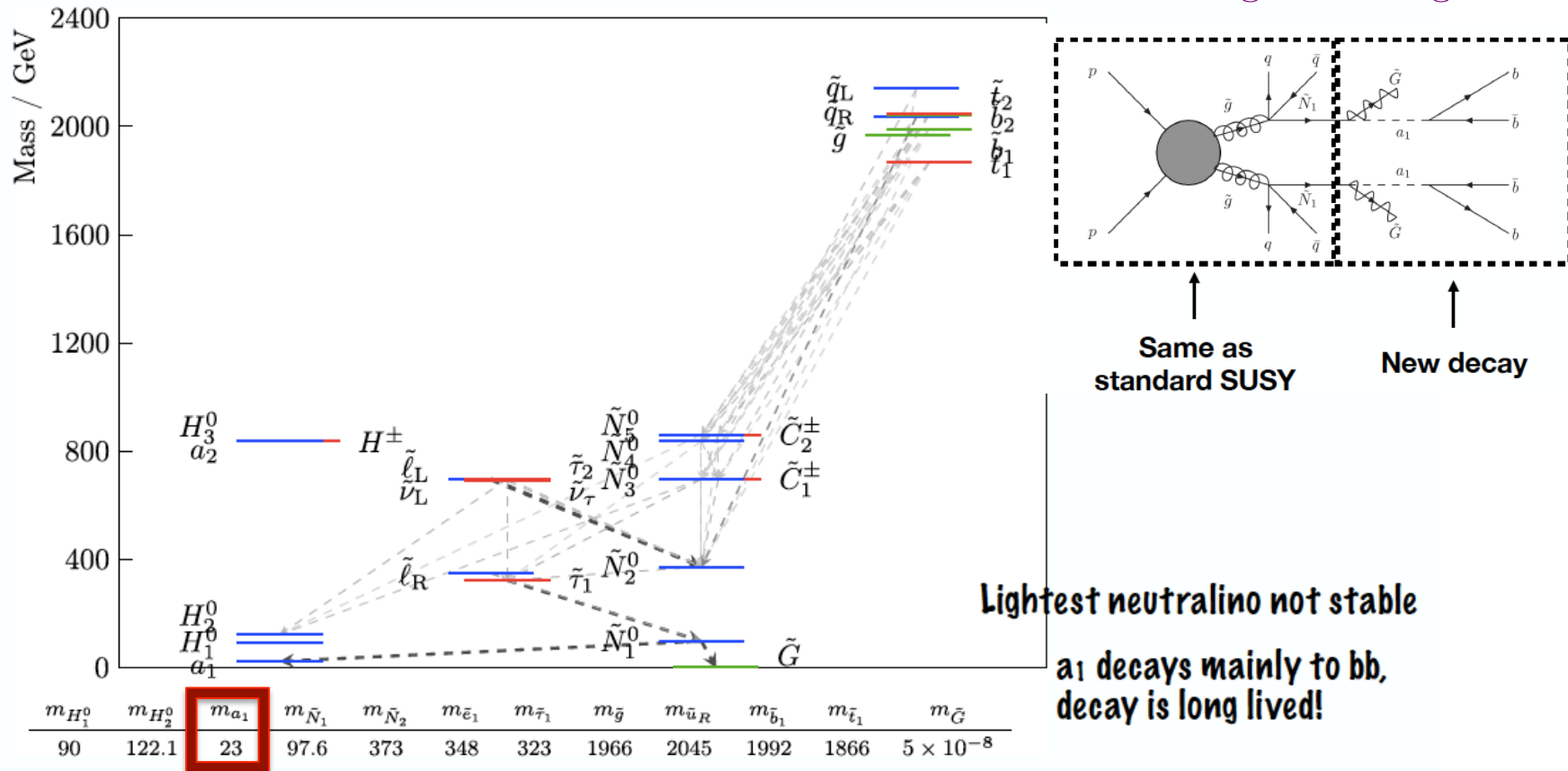
# The DGSU model



- Gluino mass down to 2 TeV allowed by SUSY searches at the LHC (the limit is stronger than in the U model due to lighter squarks)
- singlino decay length as small as  $O(1 \text{ m})$

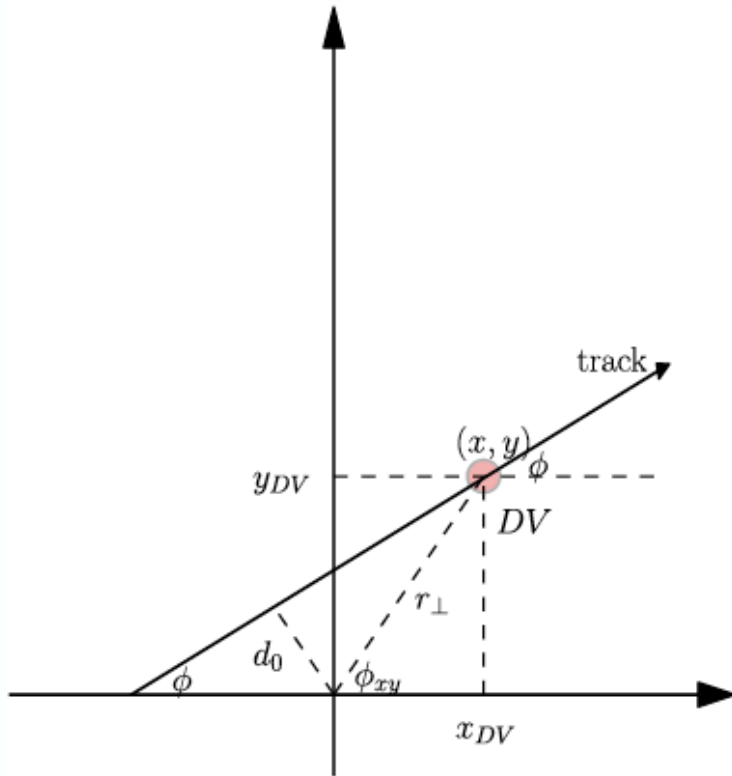
# Looking for the displaced signature at the LHC: benchmark spectrum

Allanach, MB, Cottin, Desai, Hugonie, Ziegler '16



# The displaced vertex search

ATLAS Coll. Phys. Rev. D (2015)



DV jets

4 or 5 or 6 jets with  $|\eta| < 2.8$  and  $p_T > 90, 65, 55$  GeV, each

DV reconstruction

DV made from tracks with  $p_T > 1$  GeV,  $|\eta| < 2.5$  and  $|d_0| > 2$  mm, satisfying a tracking efficiency

Vertices within 1 mm are merged

DV fiducial

DV within  $4 \text{ mm} < r_{DV} < 300 \text{ mm}$  and  $|z_{DV}| < 300 \text{ mm}$

DV material

No DV in regions near beampipe or within pixel layers:

Discard tracks with  $r_{DV}/\text{mm} \in \{[25, 38], [45, 60], [85, 95], [120, 130]\}$

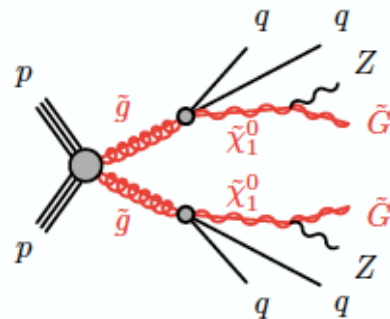
$N_{\text{trk}}$

DV track multiplicity  $\geq 5$

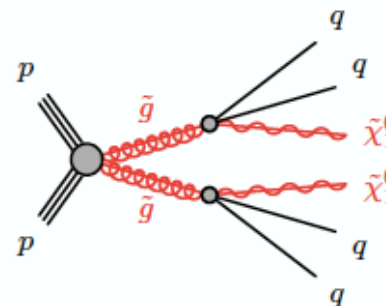
$m_{DV}$

DV mass  $> 10$  GeV

**Gauge Mediated benchmarks**

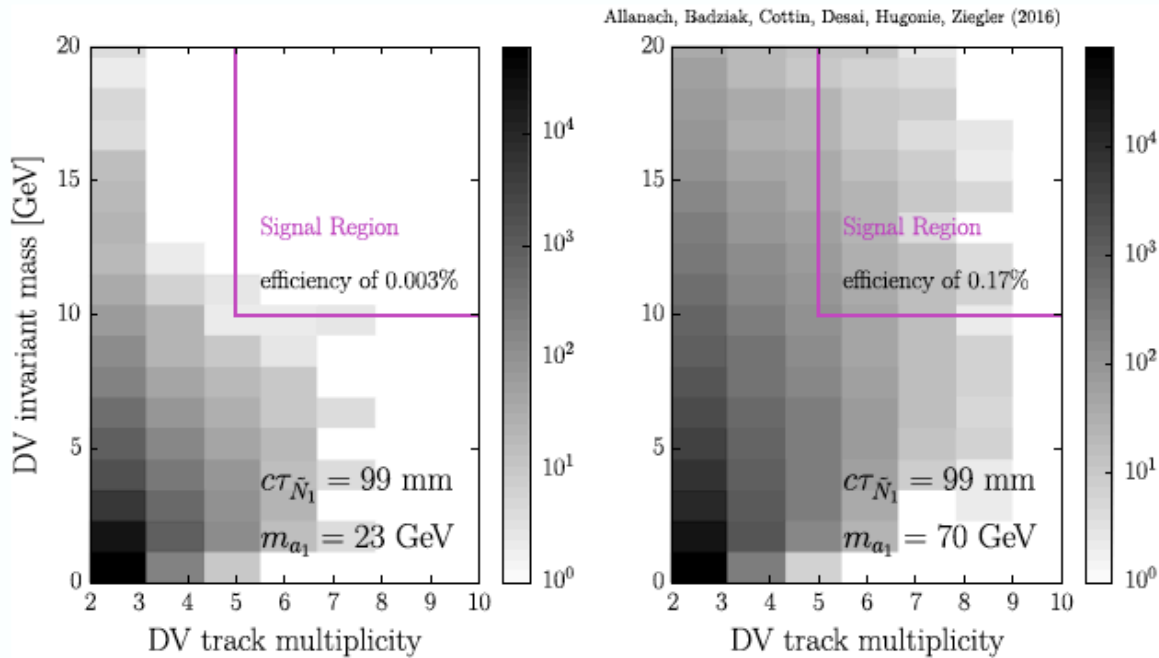


**RPV benchmark**



slide by N. Desai

# Dependence on DV mass and $N_{\text{trk}}$

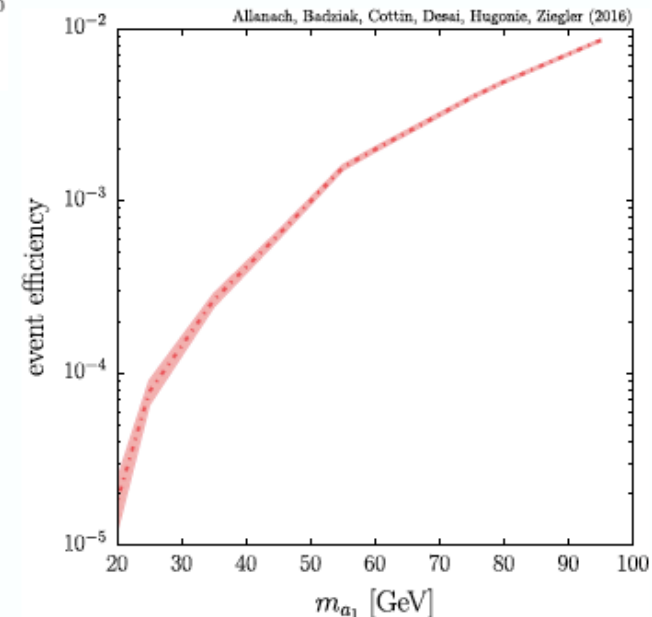


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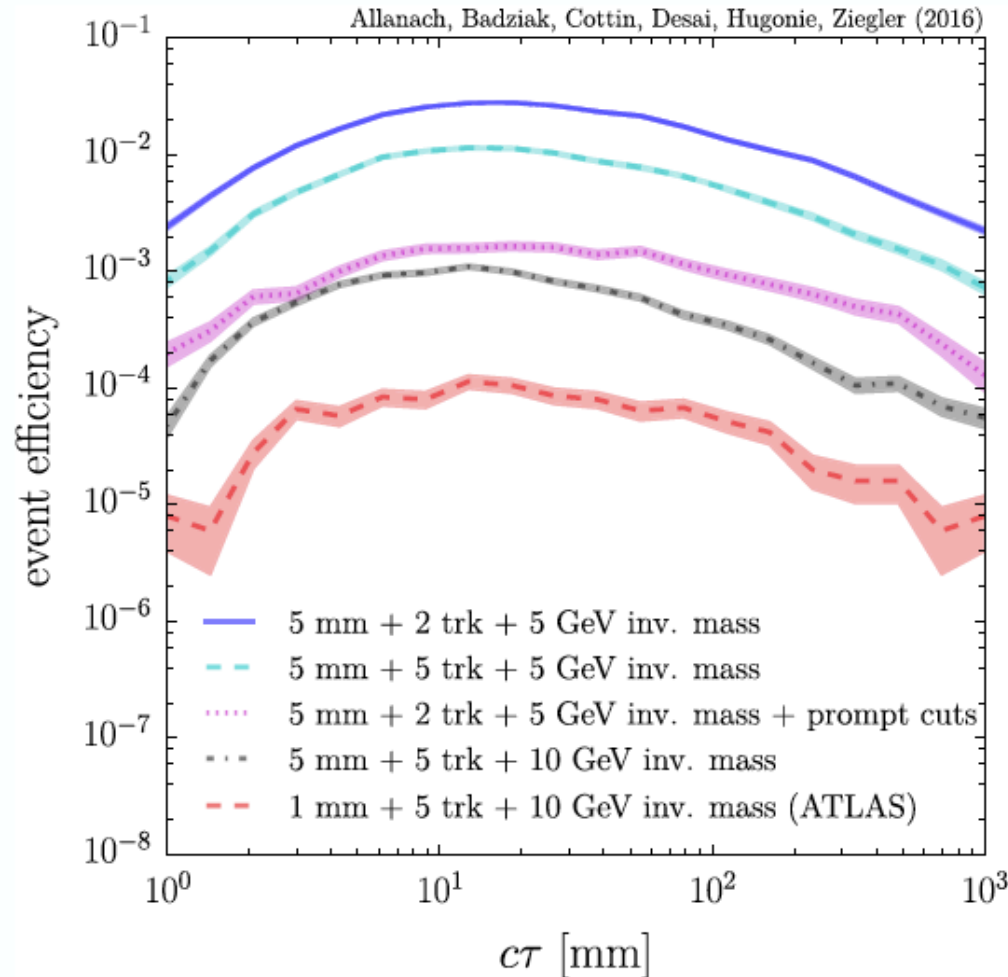
Our model fails to satisfy the DV cuts because of long lived B-mesons in the final state.

- ➔ B-mesons themselves give displaced vertices (not enough tracks  $< 1$ mm)
- ➔ Small momentum transfer

This improves with heavier  $a_1$ .



# Modifying the displaced vertex criteria



- It is possible to significantly improve efficiency by relaxing cuts
  - Not easy to estimate background for these changes
  - Our solution: combine prompt cuts + DV cuts & use prompt background estimate as a conservative upper limit
  - Much better sensitivity possible with better estimate of background
- slide by N. Desai

# Summary

- **Light NMSSM singlet** (which explains the LEP excess) solves the problems of **minimal gauge mediation** in MSSM
  - Direct **singlet-messenger couplings** is crucial for correct EWSB
  - Sparticle masses close to present experimental bounds thanks to **Higgs-messenger couplings**
- Novel signatures from decays of **long-lived singlino NLSP to displaced b-quarks**
  - Displaced vertex signature with hard prompt cuts improves sensitivity of analysis + points to underlying model
  - Much more optimisation of this search possible with a dedicated background estimate

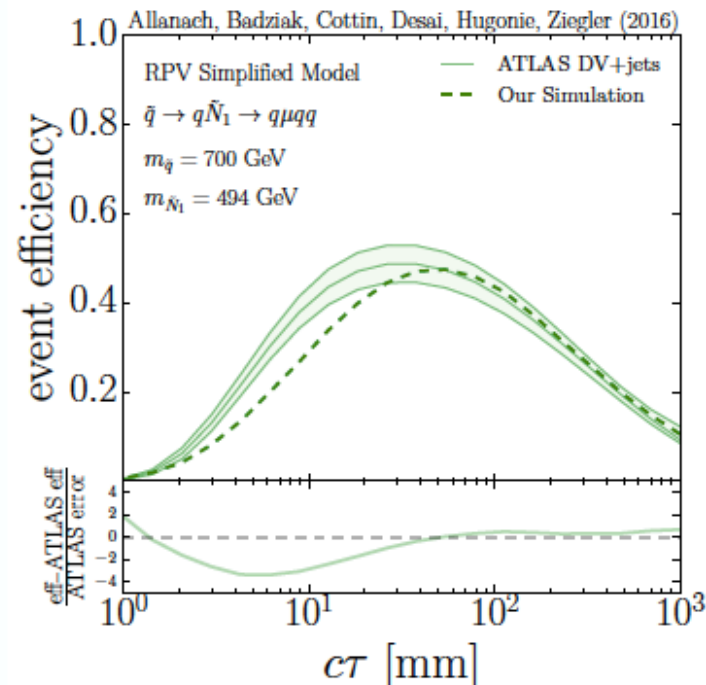
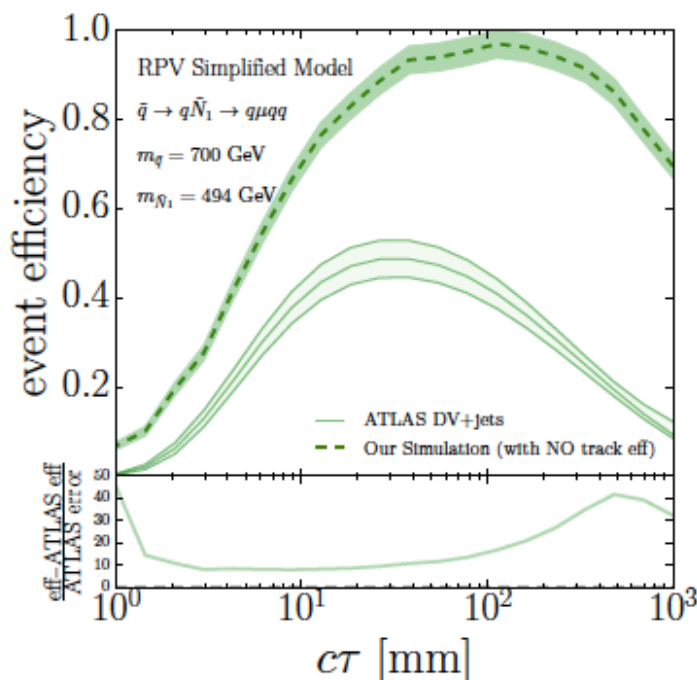
# BACKUP

# Benchmarks for U and DGSU models

	P1	P2	P3	P4	P5	P6
$\tilde{m}$ [TeV]	1.7	1.5	1.5	0.87	1.0	1.0
$M$ [GeV]	$2.8 \times 10^6$	$3.1 \times 10^6$	$2.5 \times 10^6$	$5.6 \times 10^6$	$5.1 \times 10^6$	$1.6 \times 10^6$
$\lambda$	$4.6 \times 10^{-3}$	$4.4 \times 10^{-3}$	$1.1 \times 10^{-3}$	$4.9 \times 10^{-3}$	$5.4 \times 10^{-3}$	$2.5 \times 10^{-3}$
$\kappa$	$1.4 \times 10^{-4}$	$1.2 \times 10^{-4}$	$4.3 \times 10^{-5}$	$1.5 \times 10^{-4}$	$2.1 \times 10^{-4}$	$6.5 \times 10^{-5}$
$\lambda_t$	-0.33	-0.48	-0.73	-0.64	-0.38	0.76
$\lambda_{S_d}$	0.022	0.028	0.17	-	-	-
$\xi$	-	-	-	0.012	0.010	0.010
$\tan \beta$	18	11	10	9.1	17	8.7
$m_{\tilde{g}}$ [TeV]	2.0	1.7	1.8	2.0	2.3	2.2
$m_{\tilde{d}_R}$ [TeV]	2.6	2.2	2.3	2.0	2.4	2.3
$m_{\tilde{t}_1}$ [TeV]	2.1	1.7	2.2	1.8	2.0	2.2
$m_{\tilde{N}_1}$ [GeV]	95	96	200	96	97	106
$m_{\tilde{N}_2}$ [GeV]	370	320	330	380	440	430
$m_{a_1}$ [GeV]	26	32	290	26	24	26
$m_{a_2}$ [TeV]	1.7	1.9	2.7	1.6	1.4	2.1
$m_{h_1}$ [GeV]	89	89	110	89	91	101
$m_{\mu L}$ [GeV]	960	830	820	710	840	790
$m_{\mu R}$ [GeV]	480	430	520	390	420	450
$m_{\chi_1^\pm}$ [GeV]	720	620	640	720	840	830
$m_{H_u}^2$ [TeV <sup>2</sup> ]	-2.2	-2.8	-6.8	-2.2	-1.5	-3.9
$m_{H_d}^2$ [TeV <sup>2</sup> ]	0.75	0.62	0.59	0.47	0.58	0.58
$m_S^2$ [GeV <sup>2</sup> ]	$-4.3 \times 10^3$	$-4.3 \times 10^3$	$7.8 \times 10^3$	$-4.3 \times 10^3$	$-4.5 \times 10^3$	$-5.3 \times 10^3$
$A_\lambda$ [GeV]	64	130	180	150	52	210
$A_\kappa$ [GeV]	-4.9	-7.2	-280	-4.8	-3.9	-4.3
$c\tau_{\tilde{N}_1}$ [cm]	200	200	10	200	200	12
$\Delta$	530	680	37000	440	310	720



# Finding the track efficiency



**Tracking efficiency determined by fitting parameters of an empirical function**

$$\varepsilon_{\text{trk}} = 0.5 \times (1 - \exp(-p_T/[4.0 \text{ GeV}]))$$

$$\times \exp(-z/[270 \text{ mm}])$$

$$\times \max(-0.0022 \times r_{\perp}/[1 \text{ mm}] + 0.8, 0)$$

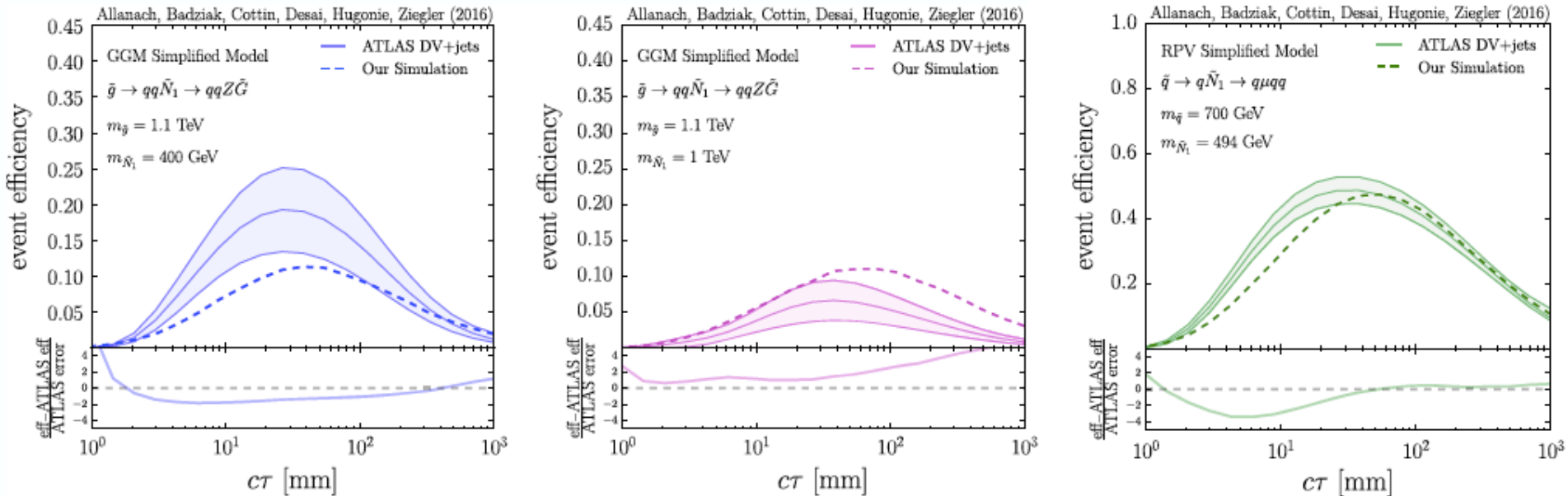
**Remove low  $p_T$**

**Dependence on  $z$  of DV (i.e. truth of decay vertex)**

**Dependence on radial distance of DV**

slide by N. Desai

# Finding the track efficiency



- Fitting any one benchmark gives vary bad fit for others
  - ➔ Not the right parameters? (we tried  $d_0, z_0$  with no improvement)
  - ➔ Hidden dependence on extra variables?
- Three benchmarks used to fit tracking parameters as a compromise

slide by N. Desai