Interpreting jet measurements in heavy-ion collisions and latest update of Constituent Subtraction method

Martin Spousta[1]
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Martin Rybář[2], Brian Cole[2]
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[1]Charles University, [2]Columbia University, [3]Mainz University,
[4]University of Chicago

Heavy Ion Jet Substructure Workshop

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Main CS developer / maintainer

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Heavy Ion Jet Substructure Workshop

Part I:
Latest update on Constituent Subtraction

CS: basics

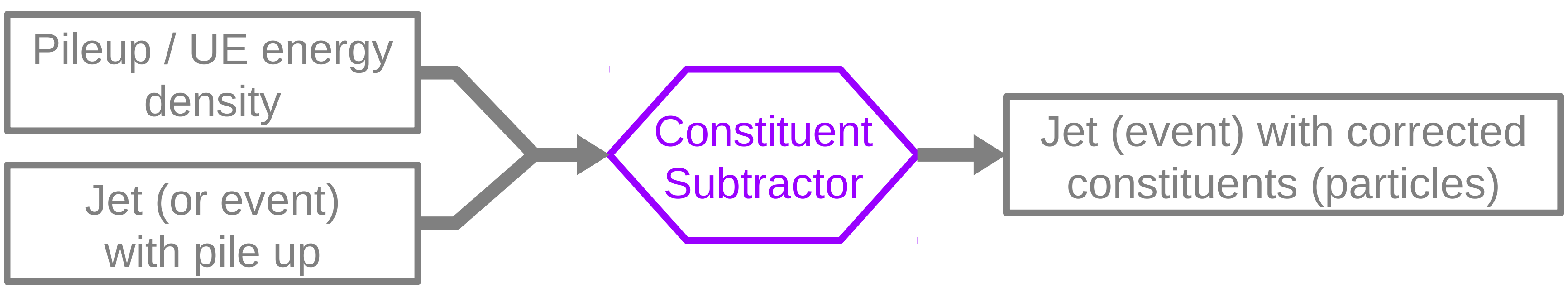
- Two main ingredients of constituent subtraction:
  - background (pileup) p_T density

$$\rho = \mathrm{median}_{\mathrm{patches}} \left\{ \frac{p_{\mathrm{T,patch}}}{A_{\mathrm{patch}}} \right\}$$

  - ghosts

$$p^g_\mu = p^g_{\mathrm{T}} \cdot [\cos\phi^g, \sin\phi^g, \sinh y^g, \cosh y^g]$$

- Constituent subtraction provides a rule for associating the background p_T density with a given constituent.



- Implementation: https://fastjet.hepforge.org/trac/browser/contrib/contribs/ConstituentSubtractor/

CS: algorithm

JHEP 1406 (2014) 092

- For each event:
  1. estimate the background p_T density, ρ
  2. add ghosts ($p_{T,g}$ –> 0) among particles (+ possibly run the jet algorithm)
- For each jet in the event / for each event:
  3. for each ghost set $p_{T,g} = \rho\ A_g$
  4. evaluate distance $\Delta R_{i,k}$ between particle *i* and ghost *k* for all particle-ghost pairs

$$\Delta R_{i,k} = p_{\mathrm{T}i}^{\alpha} \cdot \sqrt{\left(y_i - y_k^g\right)^2 + \left(\phi_i - \phi_k^g\right)^2}$$

  5. sort the distances and iteratively change the momenta as follows

$$\text{If } p_{\mathrm{T}i} \geq p_{\mathrm{T}k}^g : \qquad p_{\mathrm{T}i} \longrightarrow p_{\mathrm{T}i} - p_{\mathrm{T}k}^g, \quad p_{\mathrm{T}k}^g \longrightarrow 0; \qquad \Bigg| \qquad \text{otherwise:} \qquad p_{\mathrm{T}i} \longrightarrow 0, \quad p_{\mathrm{T}k}^g \longrightarrow p_{\mathrm{T}k}^g - p_{\mathrm{T}i}.$$

  until no more pairs remain or $\Delta R_{i,k} > \Delta R^{max}$

  6. discard all particles with zero transverse momentum

CS approach: news compared to the original implementation

- Background rescaling[1]
  - possibility to use pseudorapidity- or phi-modulated background subtraction
- Study of the Event-wide Constituent Subtraction[2]
- Implementation and study of Iterative Constituent Subtraction[3]
- Study of optimization of parameters

arXiv: 1905.03470

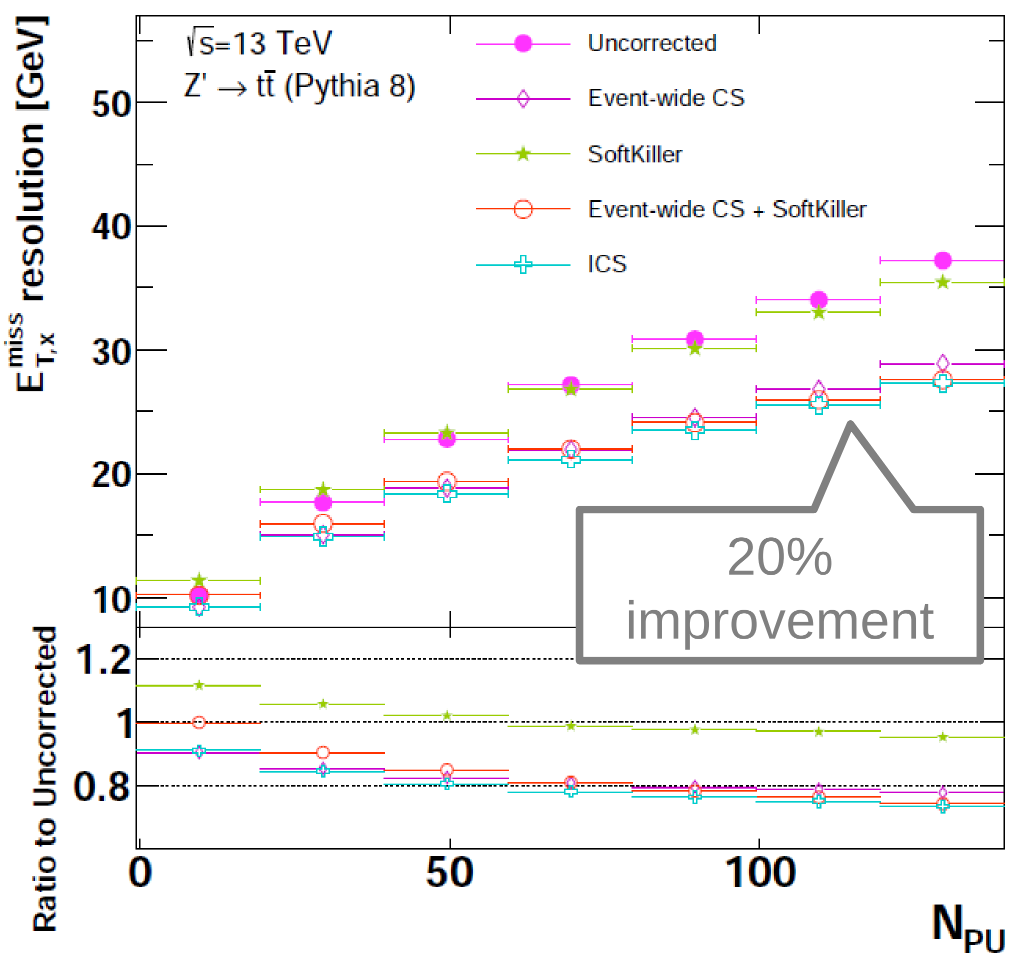
Examples of usage:

[1] example_background_rescaling.cc

[2] example_whole_event.cc

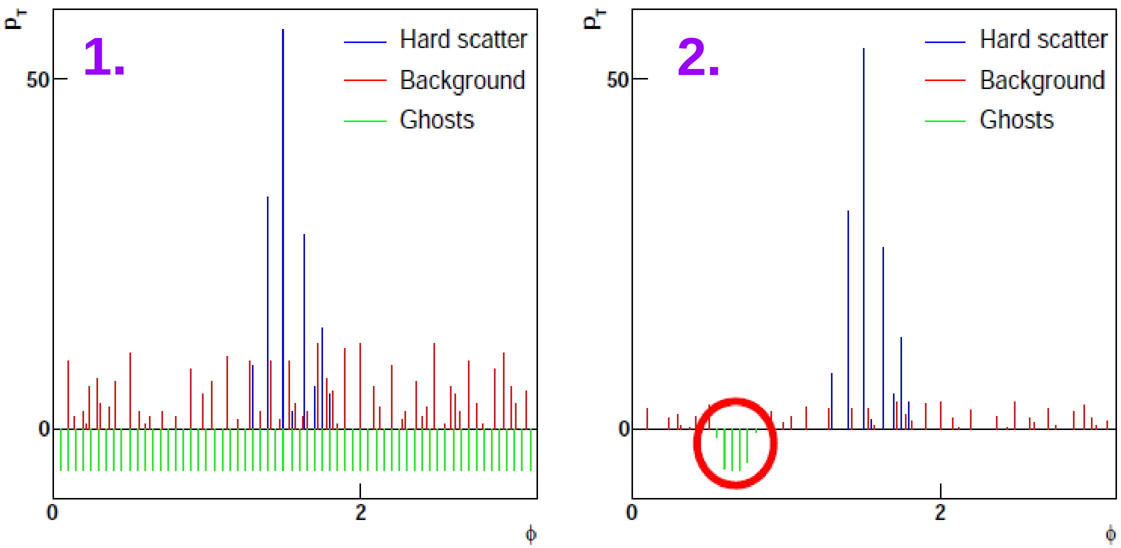
[3] example_whole_event_iterative.cc

Event-wide subtraction



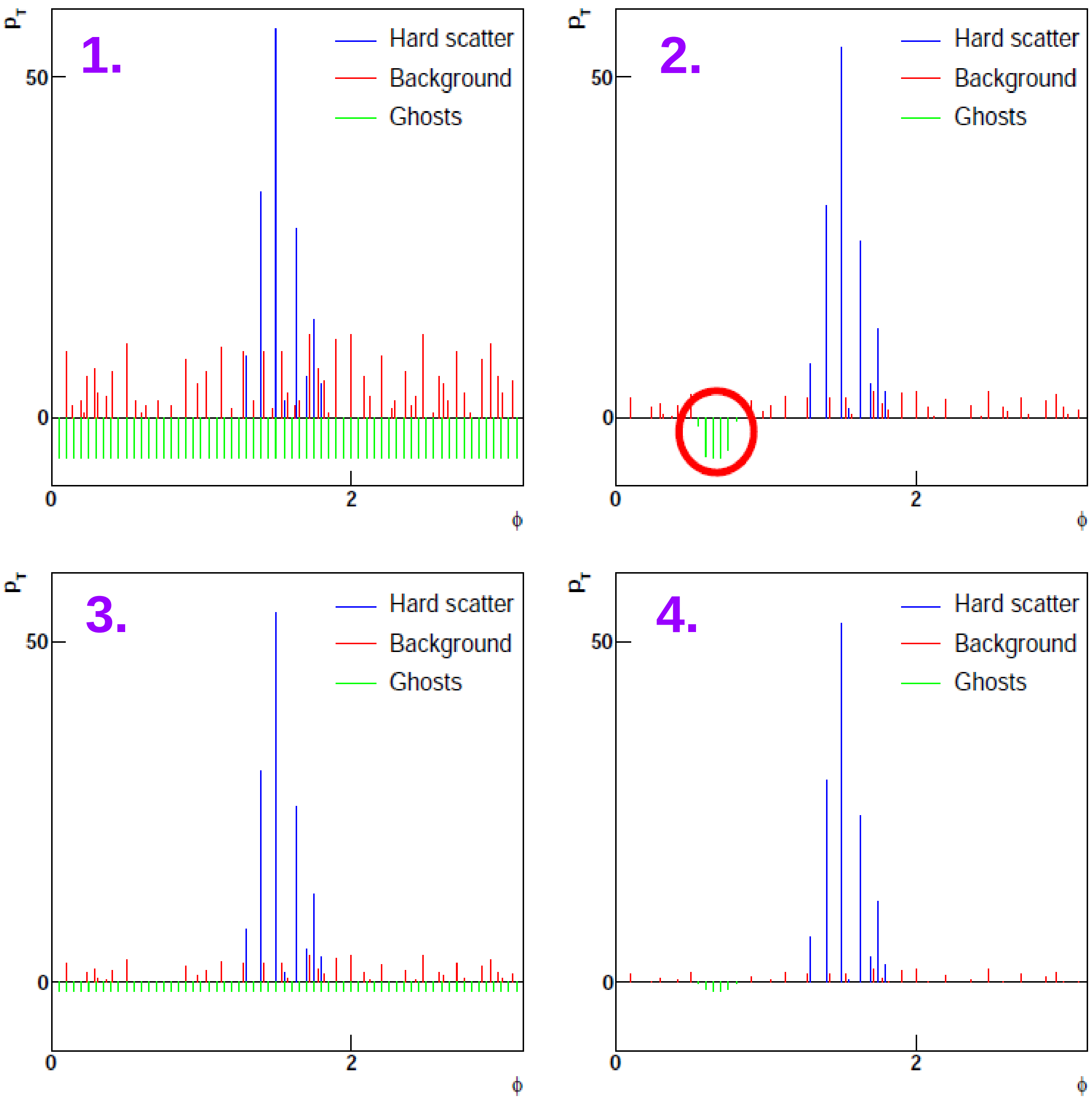
- Idea: **first subtract** and only **then cluster**.
- Jet clustering is **less biased** after whole event correction.
- Improved performance wrt Jet-by-jet CS (not shown here).
- Can be used for **global event observables** such as missing-Et.
- … but ΔR_{max}=infty does not need to be the best choice + a bias may come to a play (next slides) ...

Iterative Constituent Subtraction



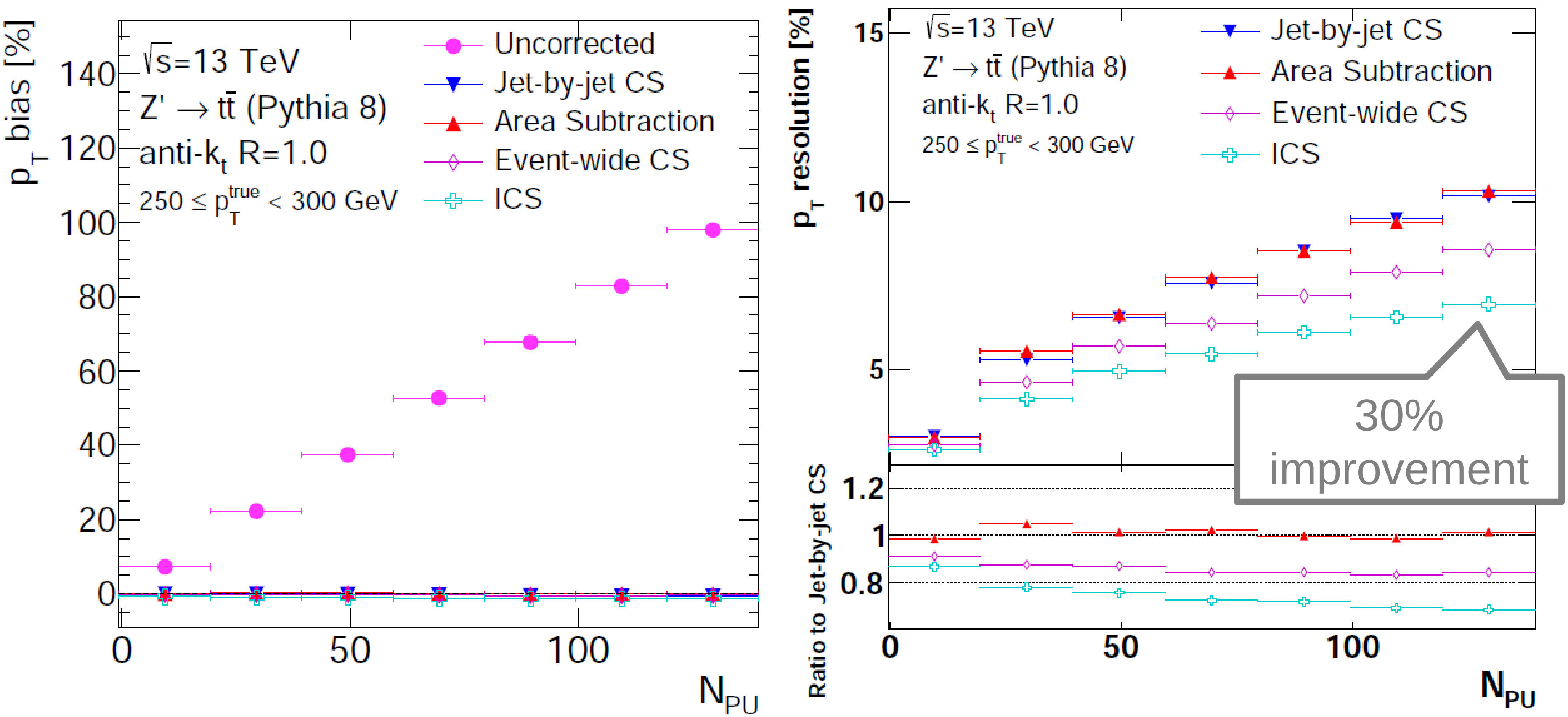
- If ΔR_{max} is finite, spacial fluctuations of background may lead to areas with unsubtracted ghosts

Iterative Constituent Subtraction



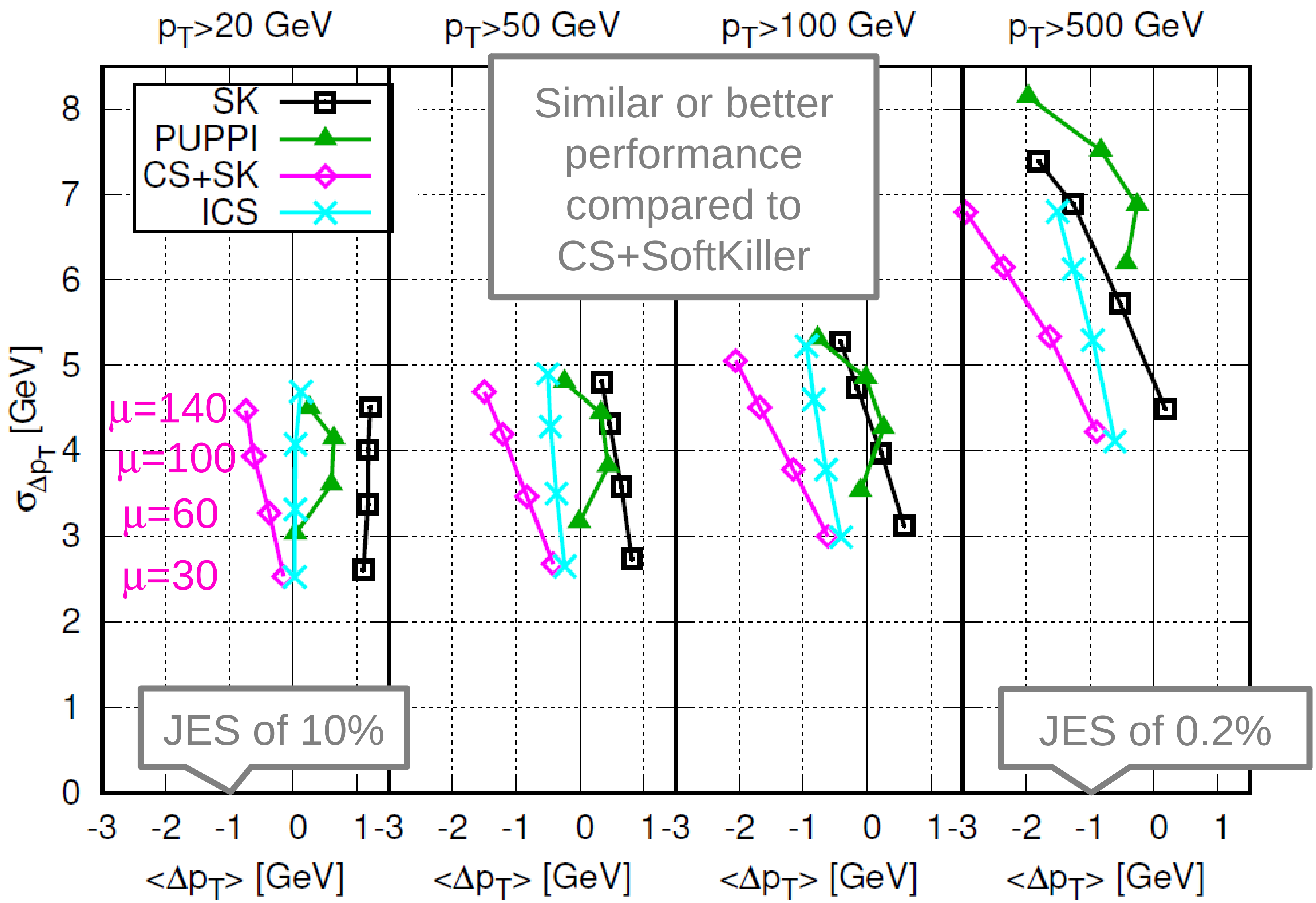
- If ΔR_{max} is finite, spacial fluctuations of background may lead to areas with unsubtracted ghosts
- => perform subtraction iteratively and possibly with different parameters

Example of performance



- Significant improvement for jet kinematics and substructure observables.
- Tested for mass, girth, τ_{21} and τ_{32}.

Comparison to other algorithms

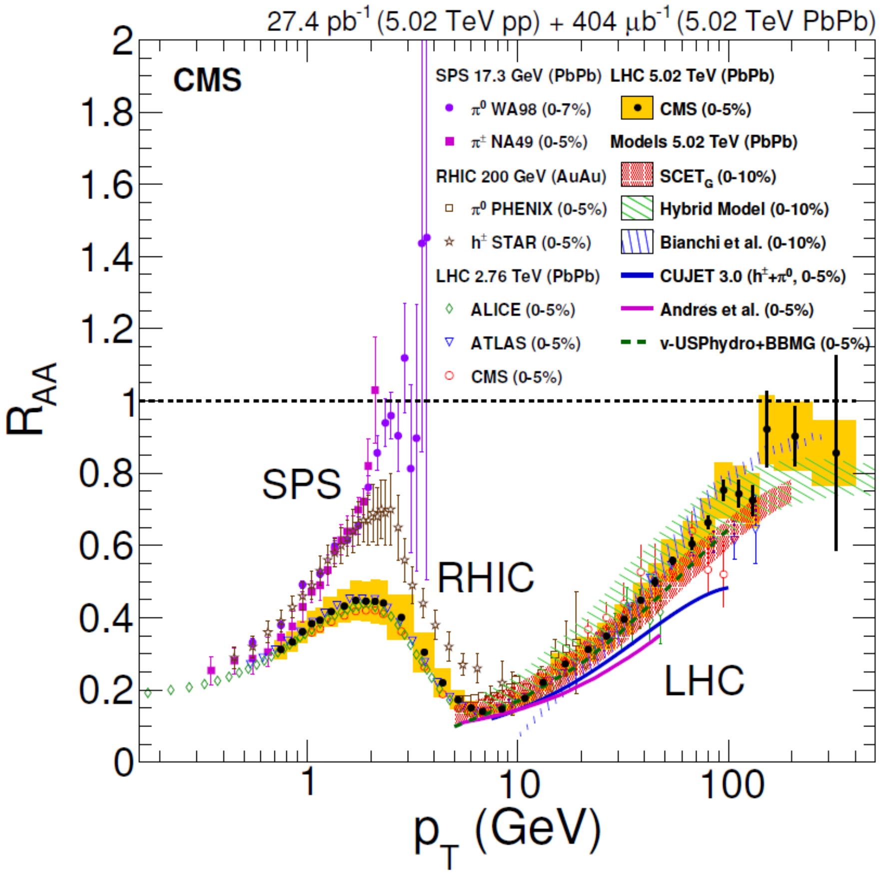


Summary

- CS used / tested by ALICE, CMS, ATLAS, STAR, and sPHENIX.
- Iterative Constituent Subtraction:
  - Can be used for **global event observables** such as missing-Et
  - Provides a significantly **improved performance** with respect to Jet-by-jet CS … **please try**
- Pile-up mitigation:
  - ICS does not use information from charged particles
    => including **charged particles** – next step to explore
- UE subtraction in HI:
  - Can one crossbreed **with other approaches?**
    E.g. arXiv:1904.12815 by Yacine & Alba & Marta
- A guideline for **choice of parameters** provided in Appendix B of arXiv:1905.03470
- Hope to continue providing a useful tool with a good support.

Part II:
Interpreting jet measurements in heavy-ion collisions

Introduction



- Physics of the suppression = **parton** energy loss in **fluctuating** hot nuclear matter
- => Some observables (I_{AA}, R_{AA} of particles, ...) result from a **complicated convolution** of: hard parton spectra, dependence of the loss on the flavor and parton shower shapes, path-length...
- => All observables are convolutions of (non-trivial) initial conditions and (non-trivial) energy loss

Introduction

Two paths:

- Be as realistic as one can:
  - MC generators
  - JETSCAPE Collaboration
  - theory calculations of parton enegy loss

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- Be simple and try to identify what plays a major role for a given observable (e.g. flavor, coherence, path length, fluctuations, …):
  - theory calculations of parton enegy loss
  - parametric modeling of parton energy loss

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Phys.Rev. Lett. 119 (2017) 062302

Phys.Lett B767 (2017) 10

Eur.Phys.J. C76 (2016) no.2, 50

arXiv:1702.01931

arXiv:1906+ε.XXXX

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This talk

Phys.Rev. Lett. 119 (2017) 062302

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Eur.Phys.J. C76 (2016) no.2, 50

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This talk ...

… use a simple model with minimal assumptions on the quenching physics to extract basic properties of the jet quenching

The simplest modeling of parton energy loss

$$\frac{dN}{dp_{\mathrm{T}}^{\mathrm{jet}}} = A\left[f_{q0}\left(\frac{p_{T0}}{p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n_q} + (1-f_{q0})\left(\frac{p_{T0}}{p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n_g}\right]$$

Jet spectra parameterized by a power law

Fraction of jets of a given flavor (i.e. quark or gluon initiated)

$$f_q\left(p_{\mathrm{T}}^{\mathrm{jet}}\right) = \frac{1}{1+\left(\frac{1-f_{q0}}{f_{q0}}\right)\left(\frac{p_{T0}}{p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n_g-n_q}}$$

The simplest parametric modeling of parton energy loss

$$\frac{dn_{\mathrm{Q}}(p_{\mathrm{T}}^{\mathrm{jet}})}{dp_{\mathrm{T}}^{\mathrm{jet}}} = \frac{dn\left(p_{\mathrm{T}}^{\mathrm{jet}} + S(p_{\mathrm{T}}^{\mathrm{jet}})\right)}{dp_{\mathrm{T}}^{\mathrm{jet}}} \times \left(1 + \frac{dS}{dp_{\mathrm{T}}^{\mathrm{jet}}}\right)$$

Yield of quenched jets of a given flavor at given pt

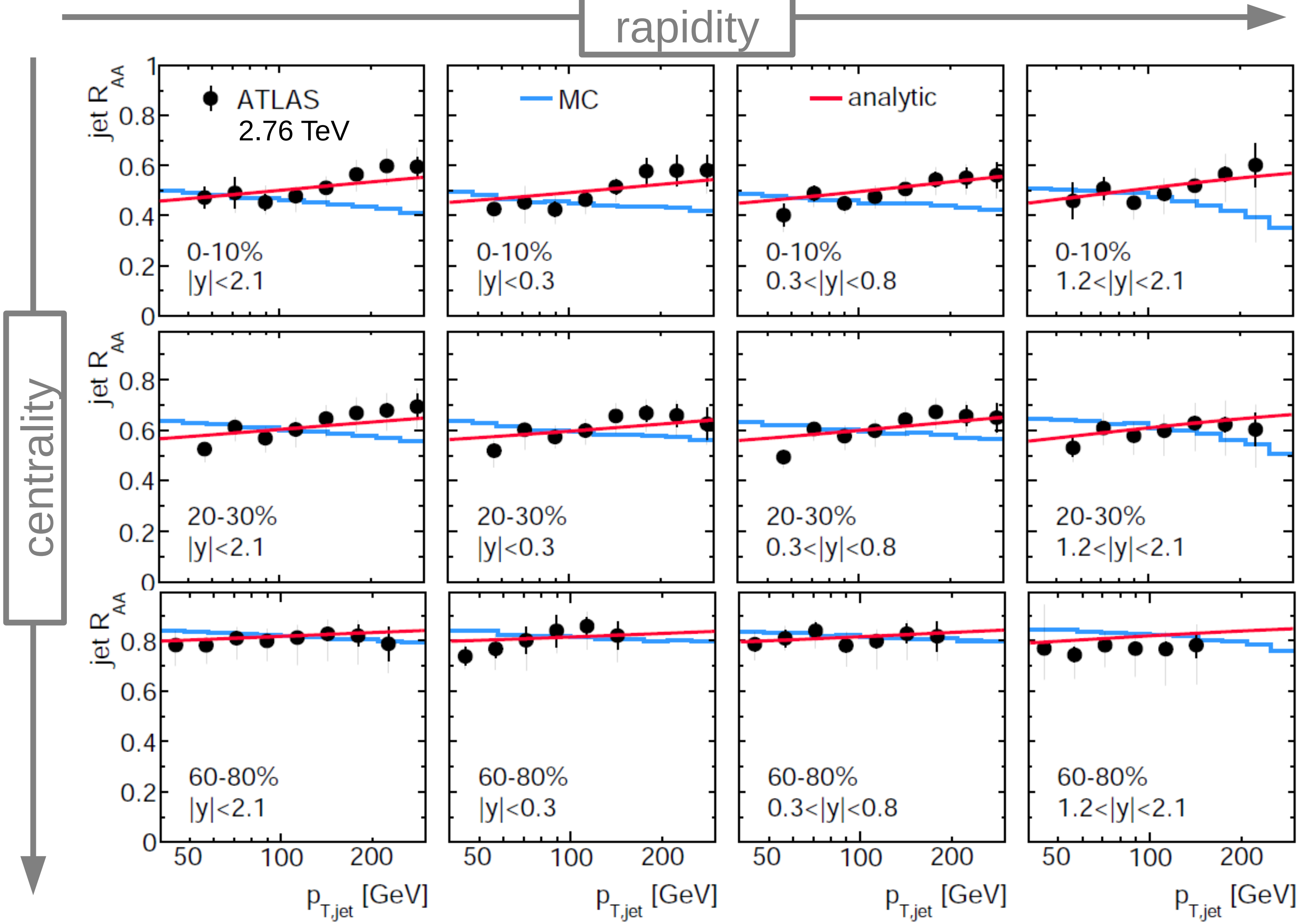
$$\begin{aligned} S_q &\equiv sp_T \\ S_g &= c_{\mathrm{F}} \times S_q \end{aligned}$$

Fractional energy loss

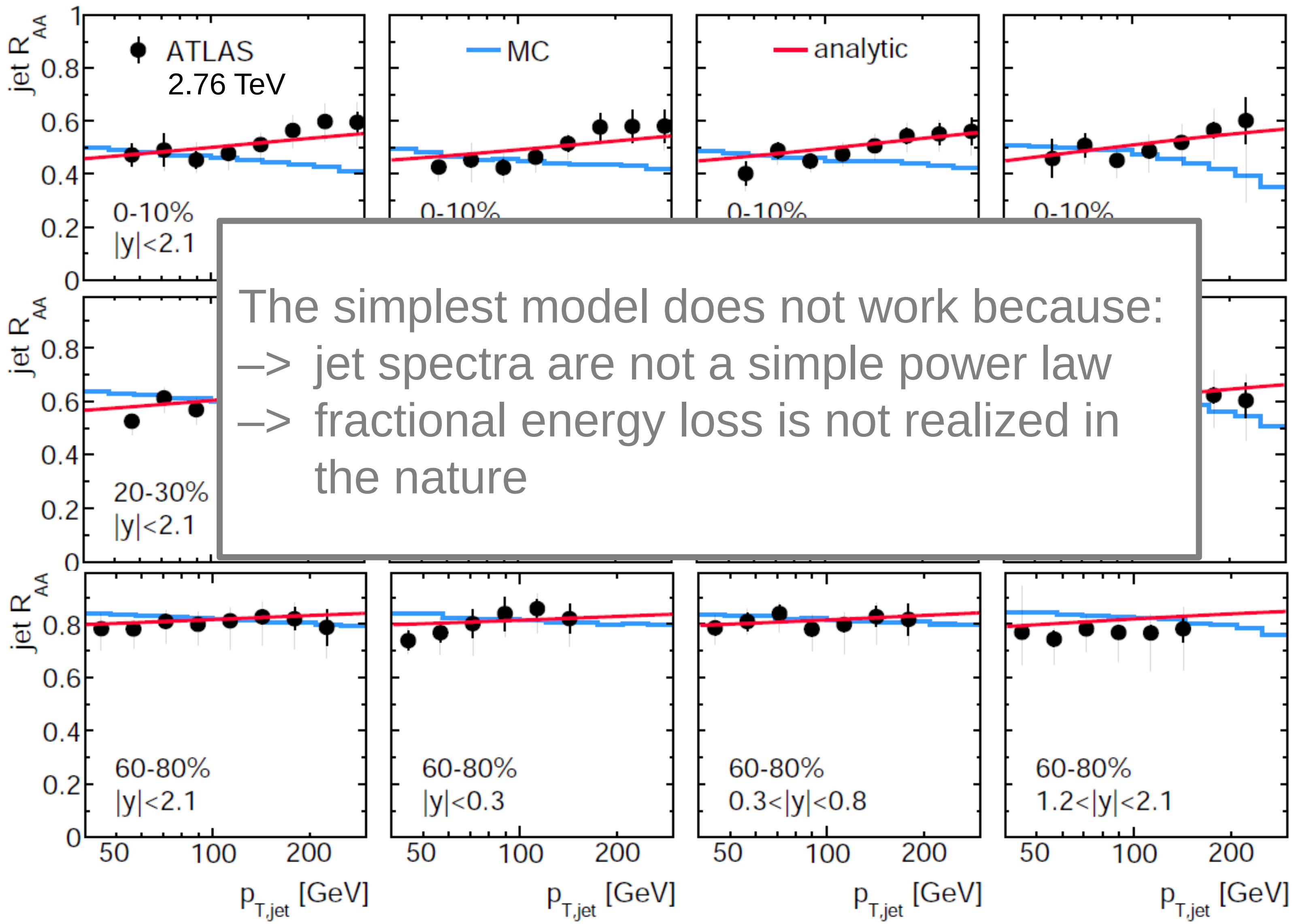
R_{AA} in the approximation of fractional energy loss

$$R_{\mathrm{AA}} = f_q \left(\frac{1}{1 + S_q/p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n_q} \times \left(1 + \frac{dS_q}{dp_T}\right) + (1 - f_q)\left(\frac{1}{1 + S_g/p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n_g} \times \left(1 + \frac{dS_g}{dp_T}\right)$$

Jet R_{AA} in the simplest model



Jet R_{AA} in the simplest model

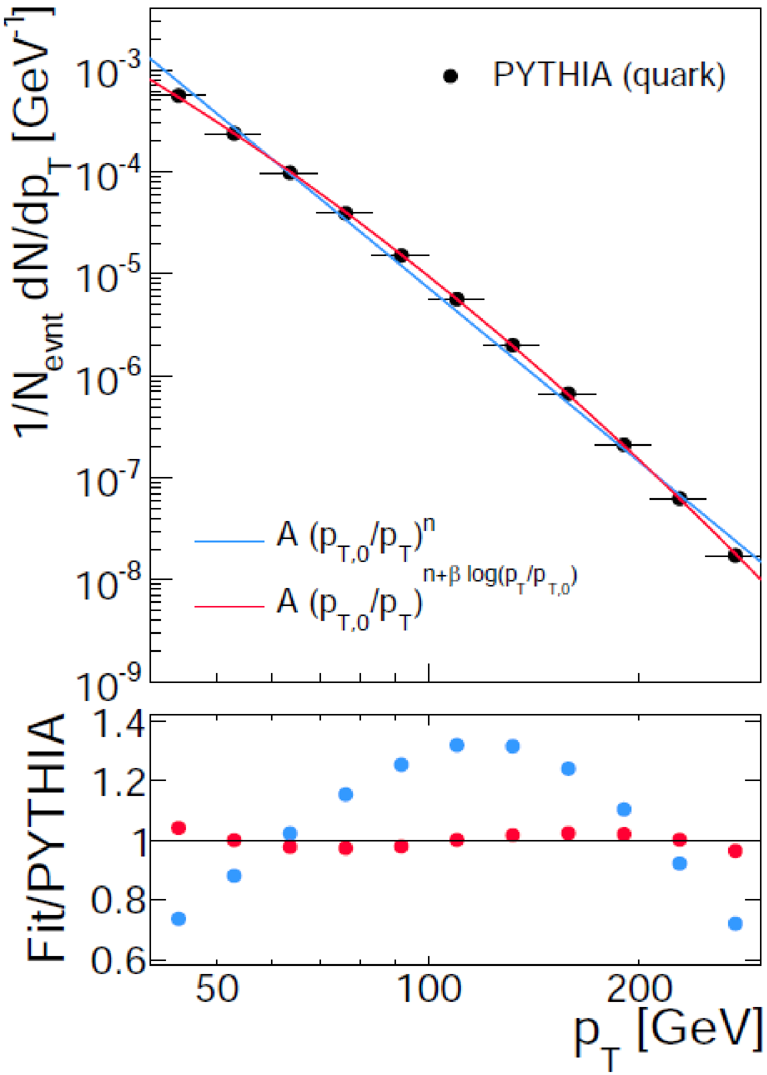


The simplest model does not work because:

–> jet spectra are not a simple power law

–> fractional energy loss is not realized in the nature

Realistic parametric model



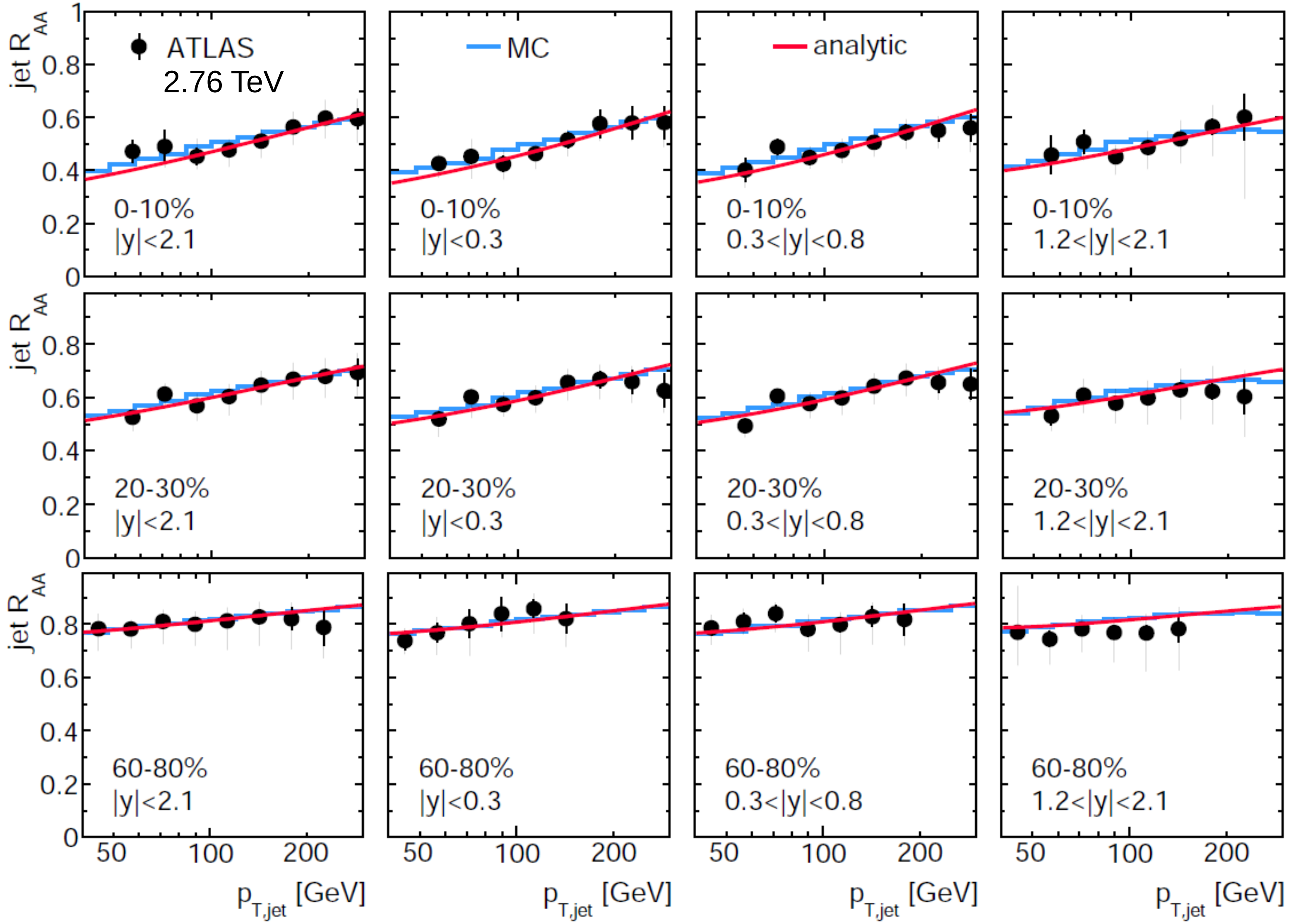
$$\frac{dn}{dp_{\mathrm{T}}^{\mathrm{jet}}} = A\left(\frac{p_{T\,0}}{p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n+\beta\log\left(p_T^{jet}/p_{T\,0}\right)}$$

Realistic parameterization of input jet spectra

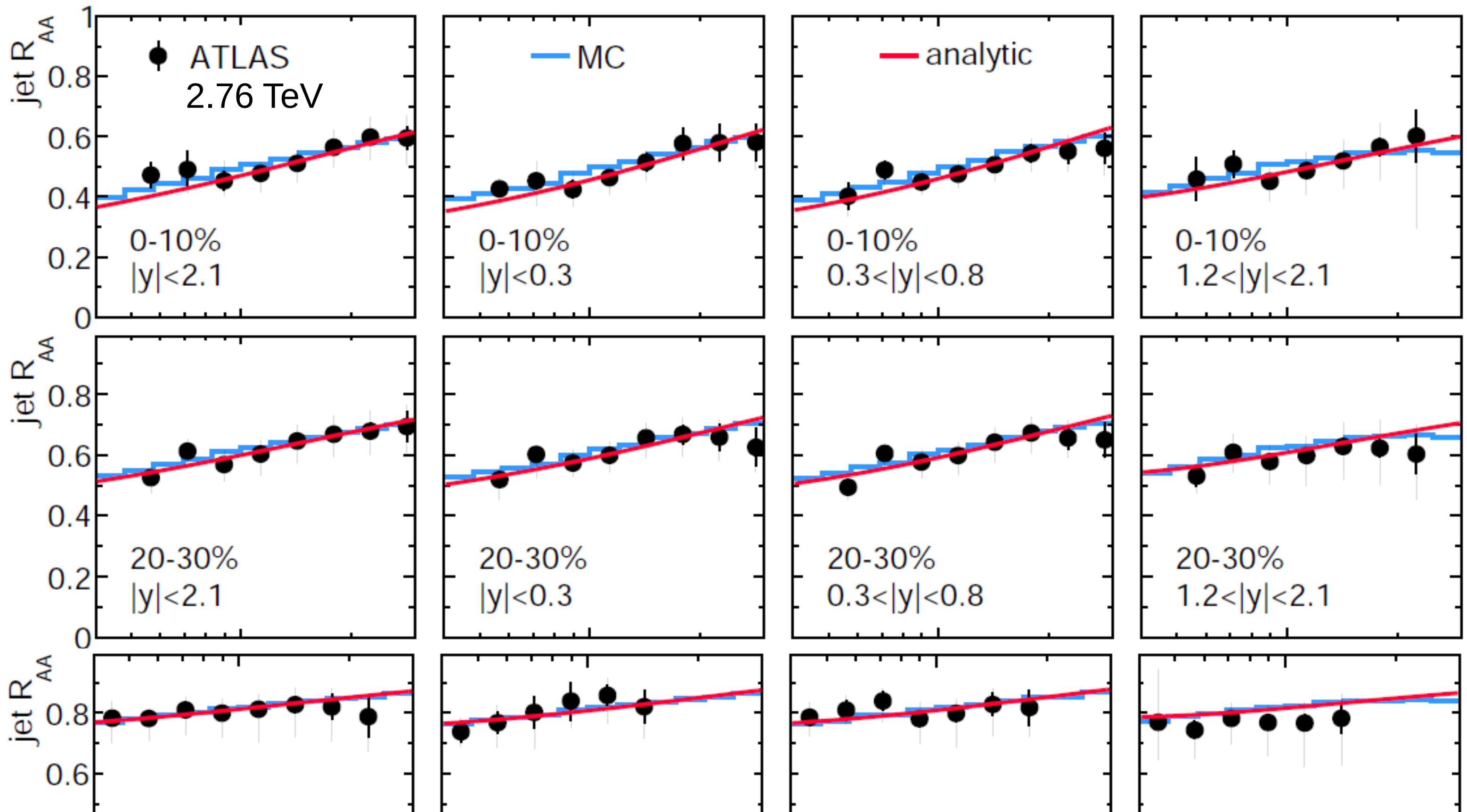
General modeling of jet energy loss

$$S = s'\left(\frac{p_T^{\mathrm{jet}}}{p_{T\,0}}\right)^{\alpha}$$

Jet R_{AA} in realistic model



Jet R_{AA} in realistic model



–> Slow evolution with p_T and no rapidity dependence of jet R_{AA} can be interpreted as a result of different energy loss of quark and gluon initiated jets

See also talk by Xin-Nian for similar conclusions within LBT

Quantifying the parton energy loss

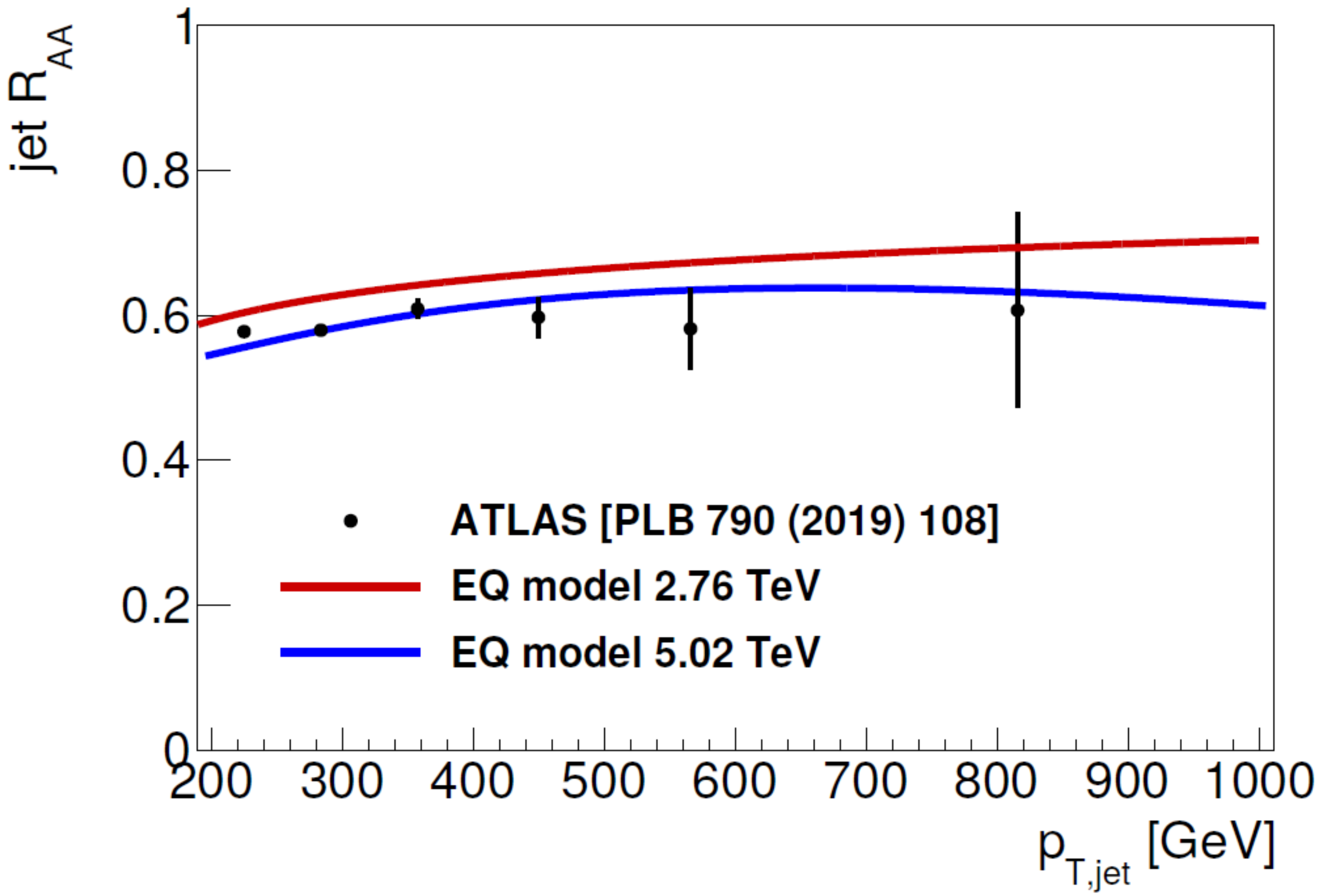
$$S_q = s' \left(\frac{p_T^{\mathrm{jet}}}{p_{T,0}} \right)^{\alpha} \qquad S_g = c_{\mathrm{F}} \times S_q$$

- Use rapidity differential jet R_{AA} measurement to perform a **multidimensional fit** and extract α, s' and c_F simultaneously (Input: **NLO** spectra – POWHEG+PYTHIA8 + 3 variations of PDFs)
- Result:

$s' = x \cdot N_{\mathrm{part}} + y$	$x = (12.3 \pm 1.4) \cdot 10^{-3}$ GeV, $y = 1.5 \pm 0.2$ GeV
α	0.52 ± 0.02
c_{F}	1.78 ± 0.12

- *Average* jet quenching encapsulated **in 4 parameters**.
- Value of **c_F** seems consistent with the value **in the vacuum** (Useful discussion on c_F also in arXiv:1812.06019)

5.02 TeV versus 2.76 TeV

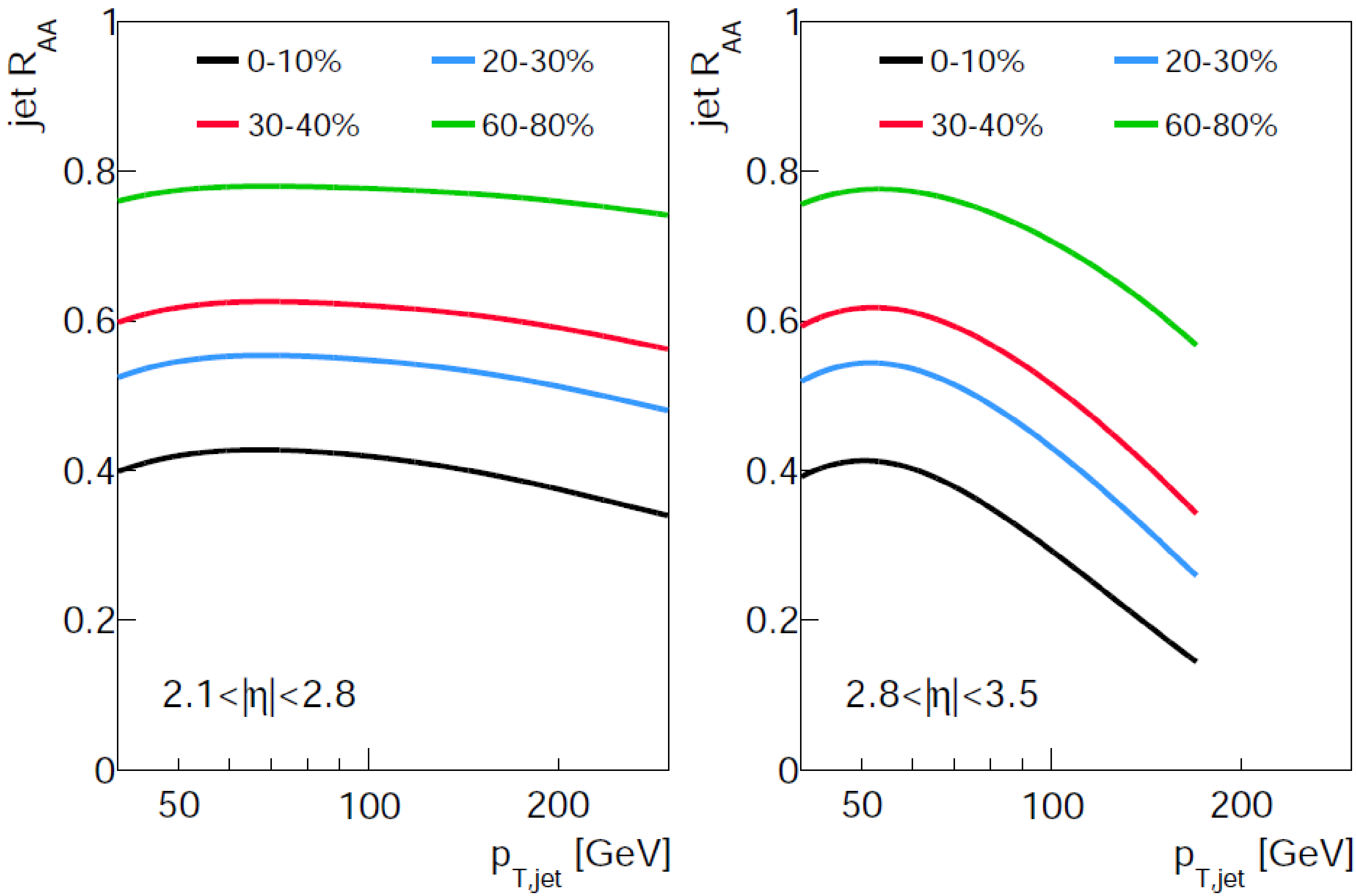


- Same jet R_{AA} … but that does not imply same energy loss.
- Spectra shape and flavor admixture are different
  => energy loss must be different.
- About 10% larger energy loss at 5.02 TeV compared to 2.76 TeV.

Note

- … jet R_{AA} … as a result of different energy loss of quark- and gluon-initiated jets
- Alternative: **shower shape** – wide jets lose more than narrow.
- How to **distinguish**?
- Do as **many comparisons** with data as possible (in the kinematic region insensitive to in-cone radiation / recoil effects). Here:
  - Rapidity dependence of the R_{AA},
  - Behavior of the R_{AA} in the forward region,
  - Jet fragmentation,
  - Jet shapes.
- More info in the **backup**: charged particle R_{AA}, b-jet R_{AA}, z_g, high-p_T charmonia, … and more to come
- Do as many comparisons as possible … and **look for a failure** (by seeing a failure of the model one can learn new stuff)

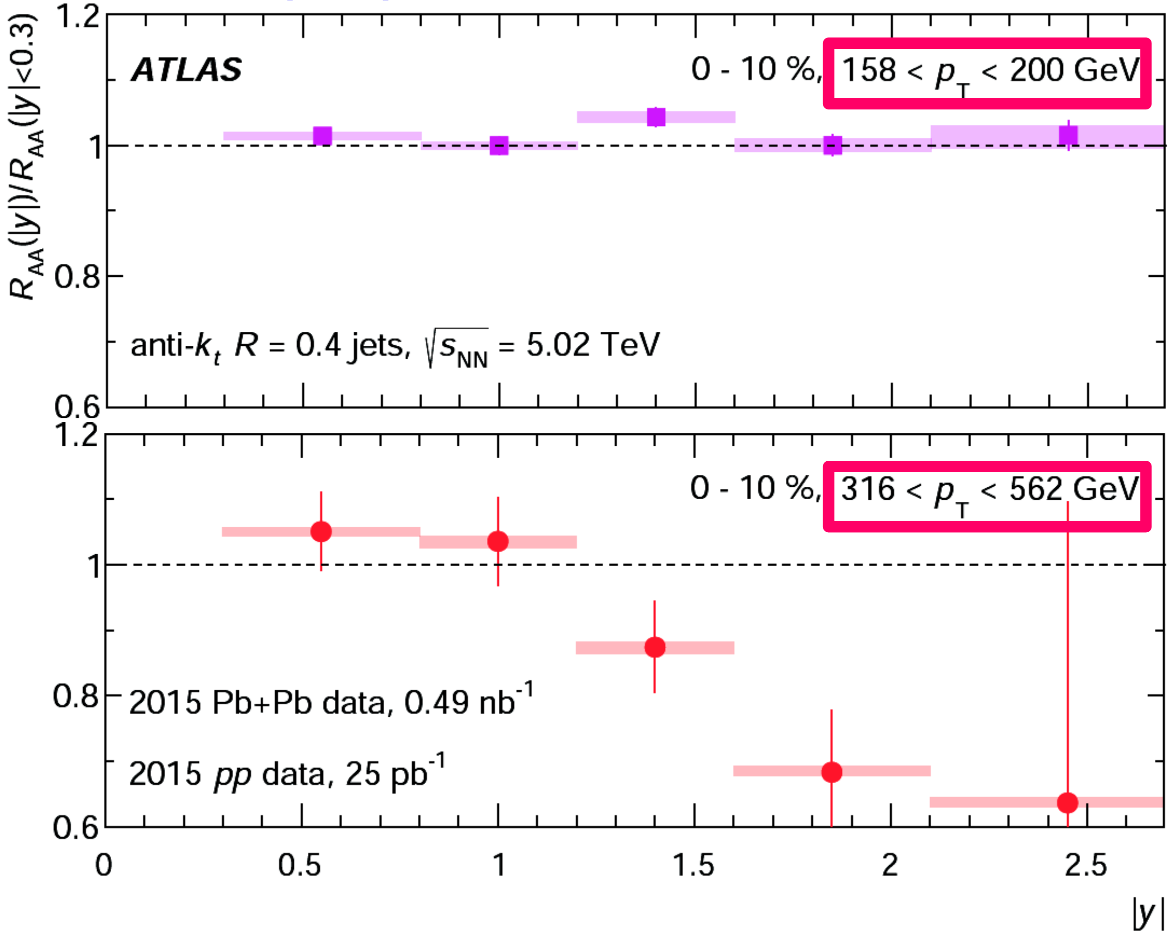
Predictions for the forward region



–> The jet R_{AA} **should decrease** in the forward region

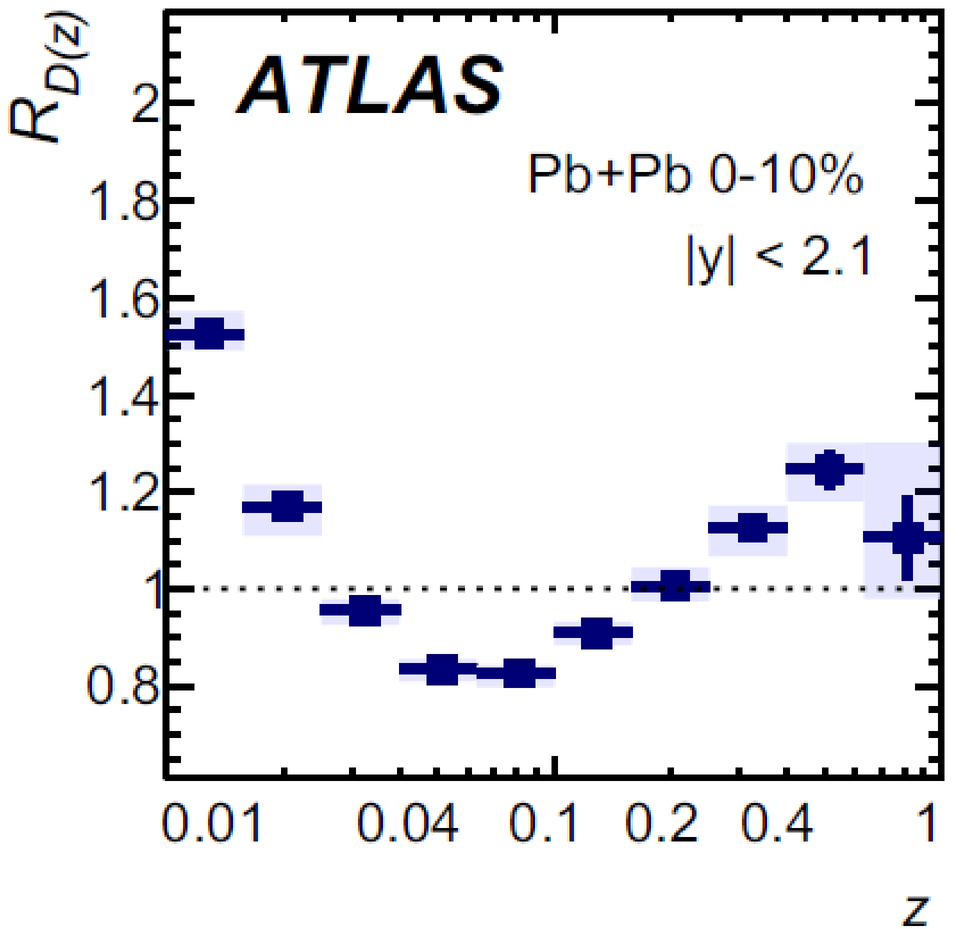
Measurement in the forward region

PLB 790 (2019) 108



–> The jet R_{AA} **does decrease** in the forward region

Modification of longitudinal structure of jet (fragmentation function)



PRC 98 (2018) 024908
EPJC 77 (2017) 379
PRC 90 (2014) 024908
…

Modifications of fragmentation functions

- How is the parton shower modified by the QCD medium?
- Basic picture …

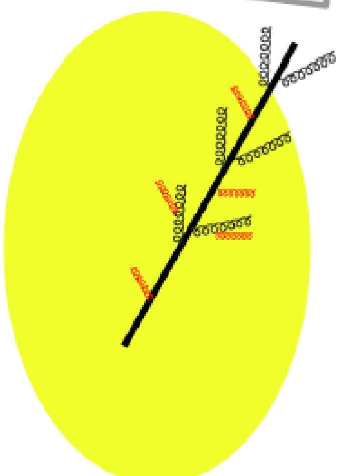
Black = vacuum component of PS
Red = medium induced radiation

c) = a && b

For some configurations medium resolves parton shower

! d) = a || b || c + more

a)



Medium resolves parton shower

e.g.: Phys. Lett. B345 (1995), 277
Nucl. Phys. B582 (2000), 409
Phys. Rev. Lett. 85 (2000), 5535
Phys. Rev. D50 (1994), 1951
JHEP 12 (2001), 009

b)

Emission is coherent

e.g.: Phys. Rev. Lett. 106 (2011)
Phys. Lett. B725 (2013), 357
Phys. Rev. Lett. 111 (2013), 052001

in-cone radiation + jet excites medium => "recoiling" particles from the medium

e.g.: Phys. Rev. C80 (2009) 054913
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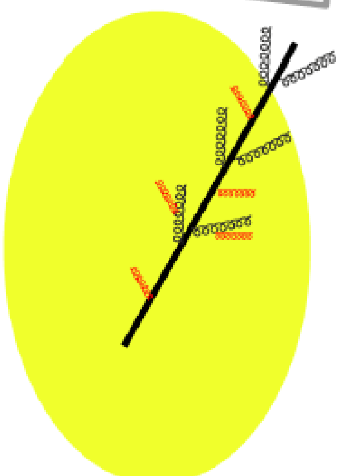

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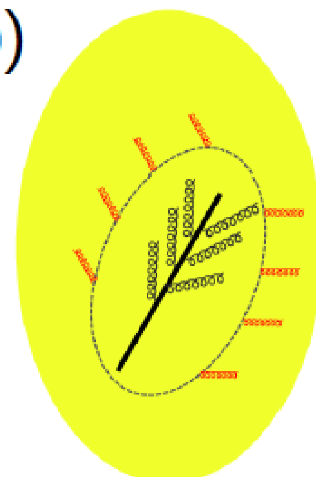

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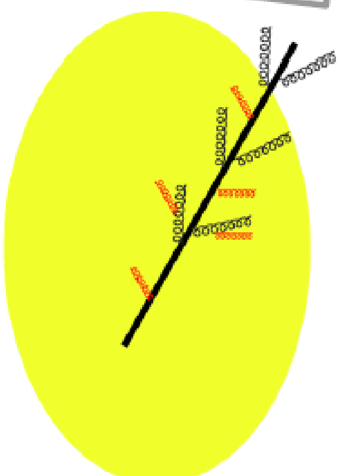

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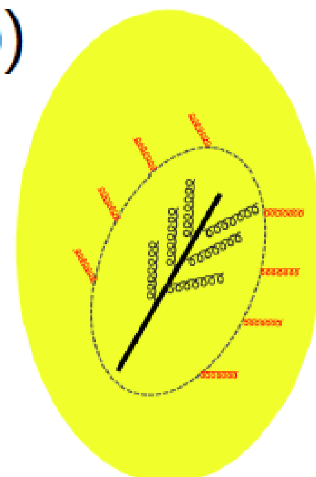

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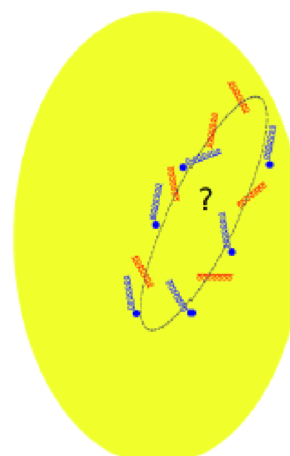

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See e.g. talk by Yacine, Mon

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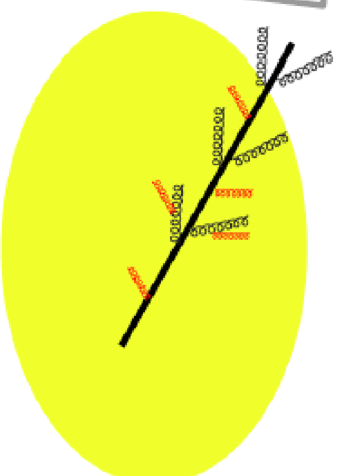

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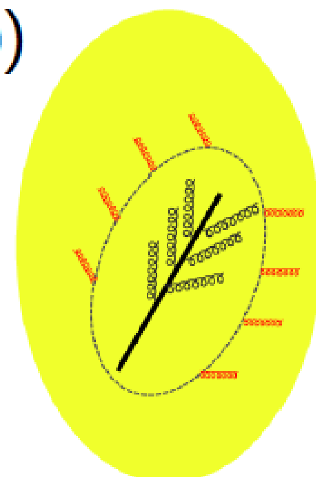

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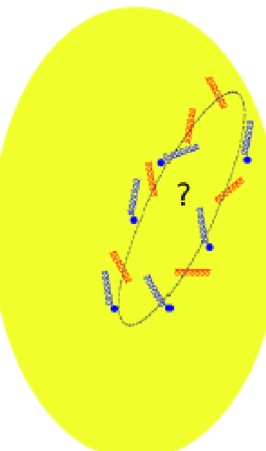

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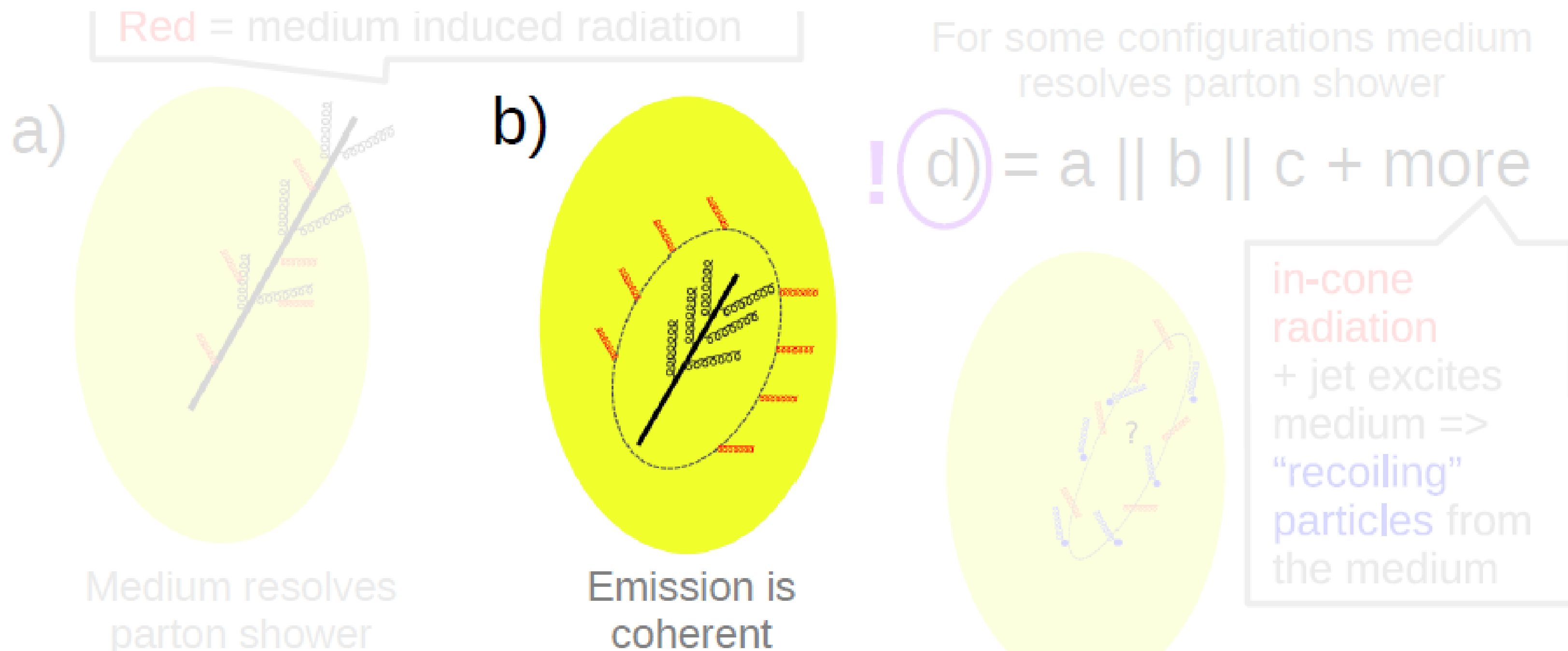
Modifications of fragmentation functions

–> Subtract the energy from the jet / initial parton and then let it fragment as in the vacuum



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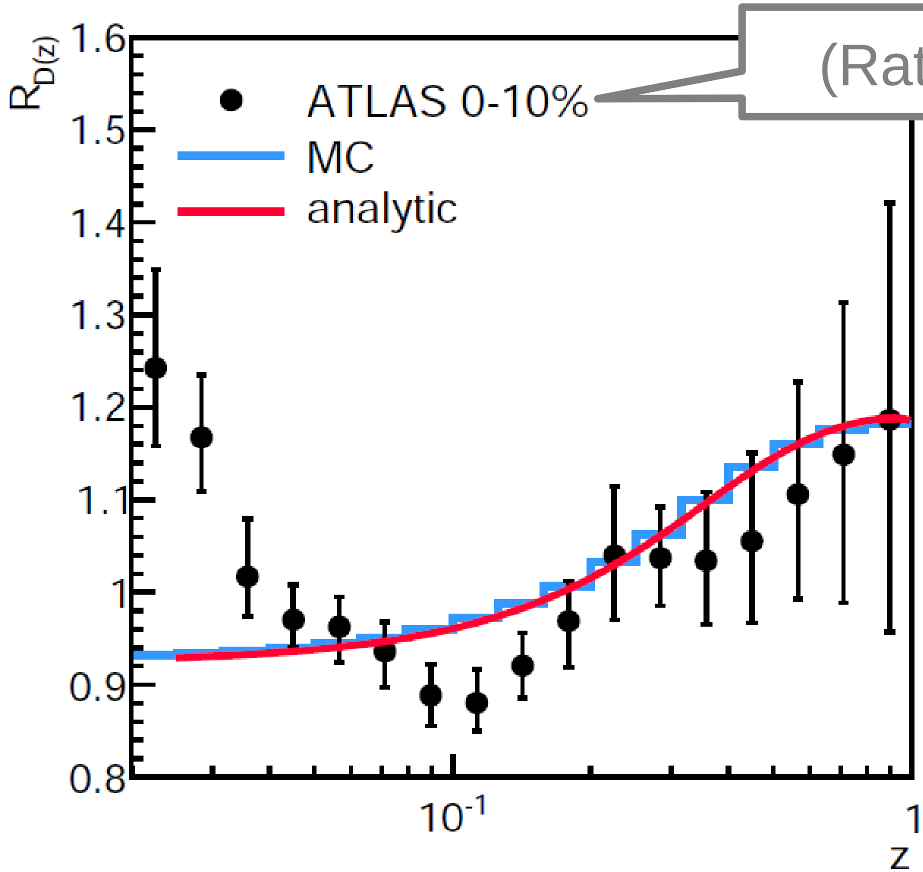
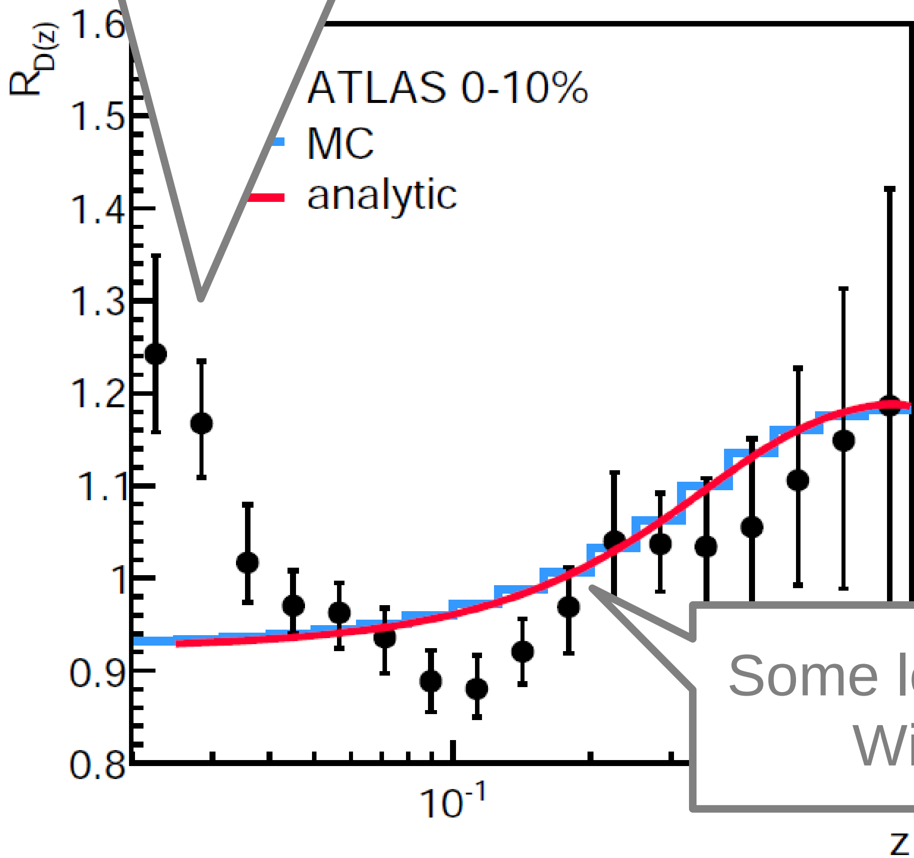


(Ratio of fragmentation functions)

Modifications of fragmentation functions

Excess of low-*z* not due to flavor effects (due to in-cone radiation or recoil effects)

...he jet / initial parton and
...e vacuum



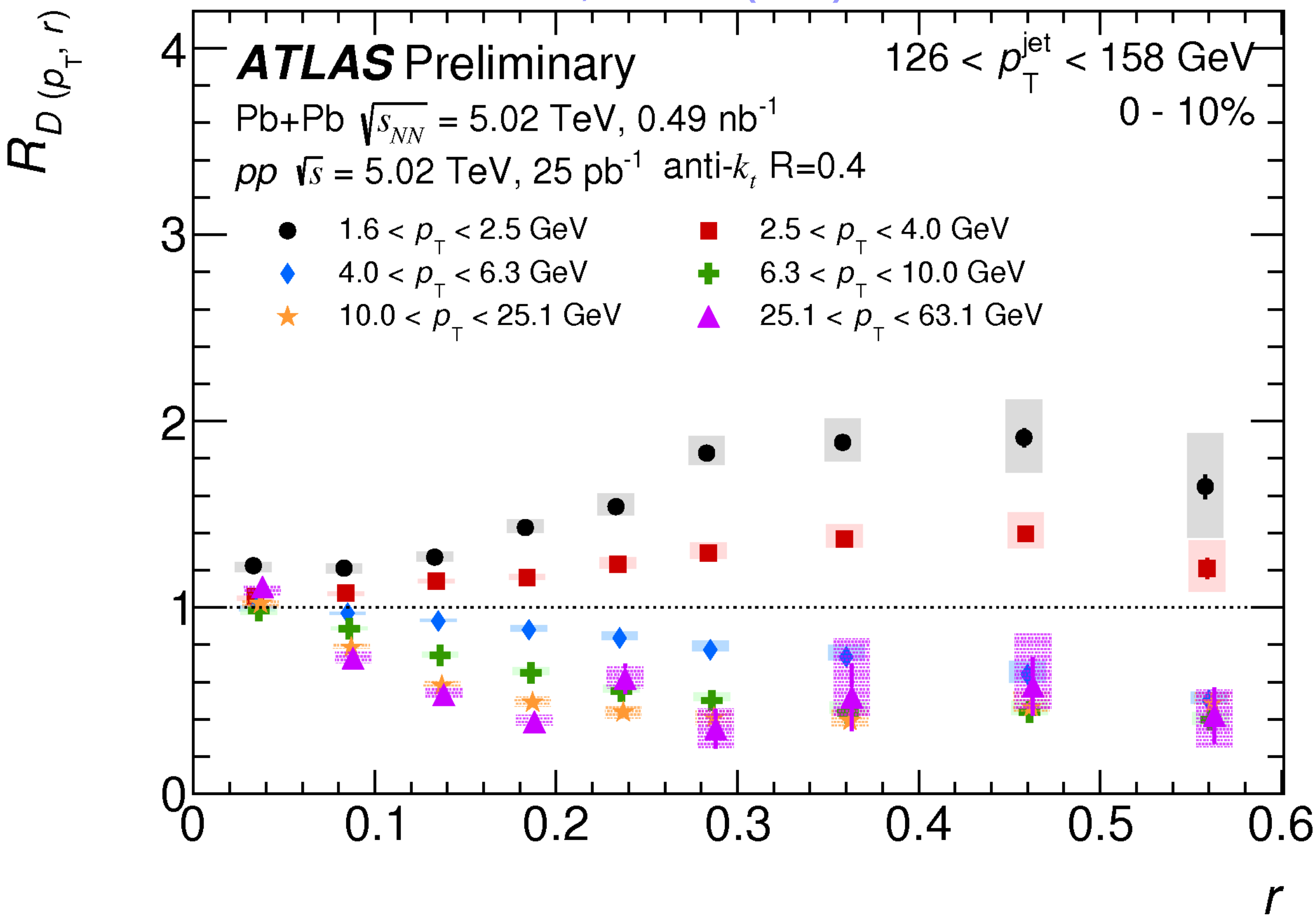
Some level of disagreement? Will get back to it ...

–> Structure seen at intermediate and high-z is due to the difference in quenching of quark and gluon initiated jets

–> Speaks in favor of presence of color coherence effects in the data

Transverse structure of jet (jet shape)

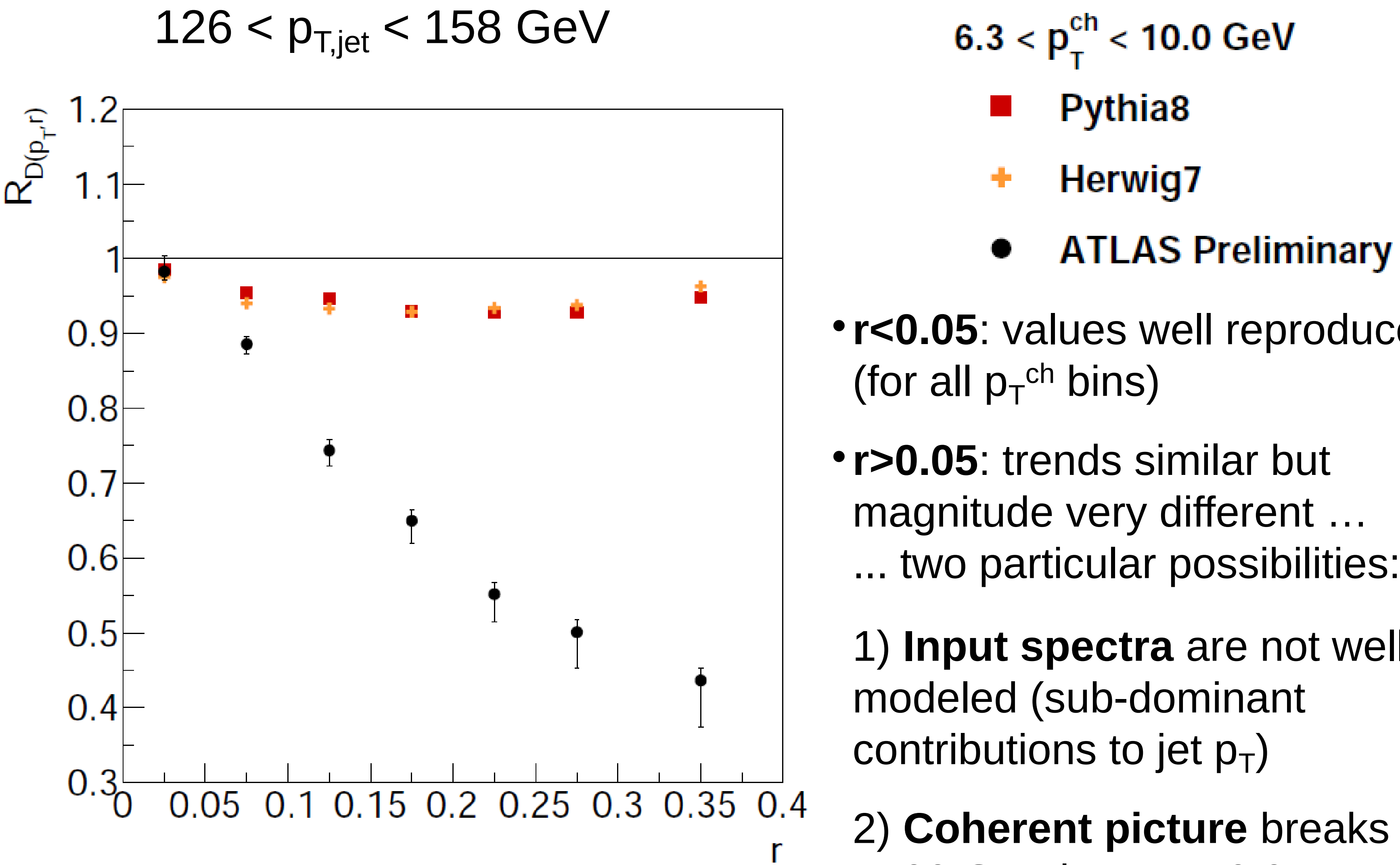
ATLAS-CONF-2018-010, JHEP 05 (2018) 006



Modification of the jet shape

–> Subtract the energy from the jet / initial parton and then let it fragment as in the vacuum

Modification of the jet shape

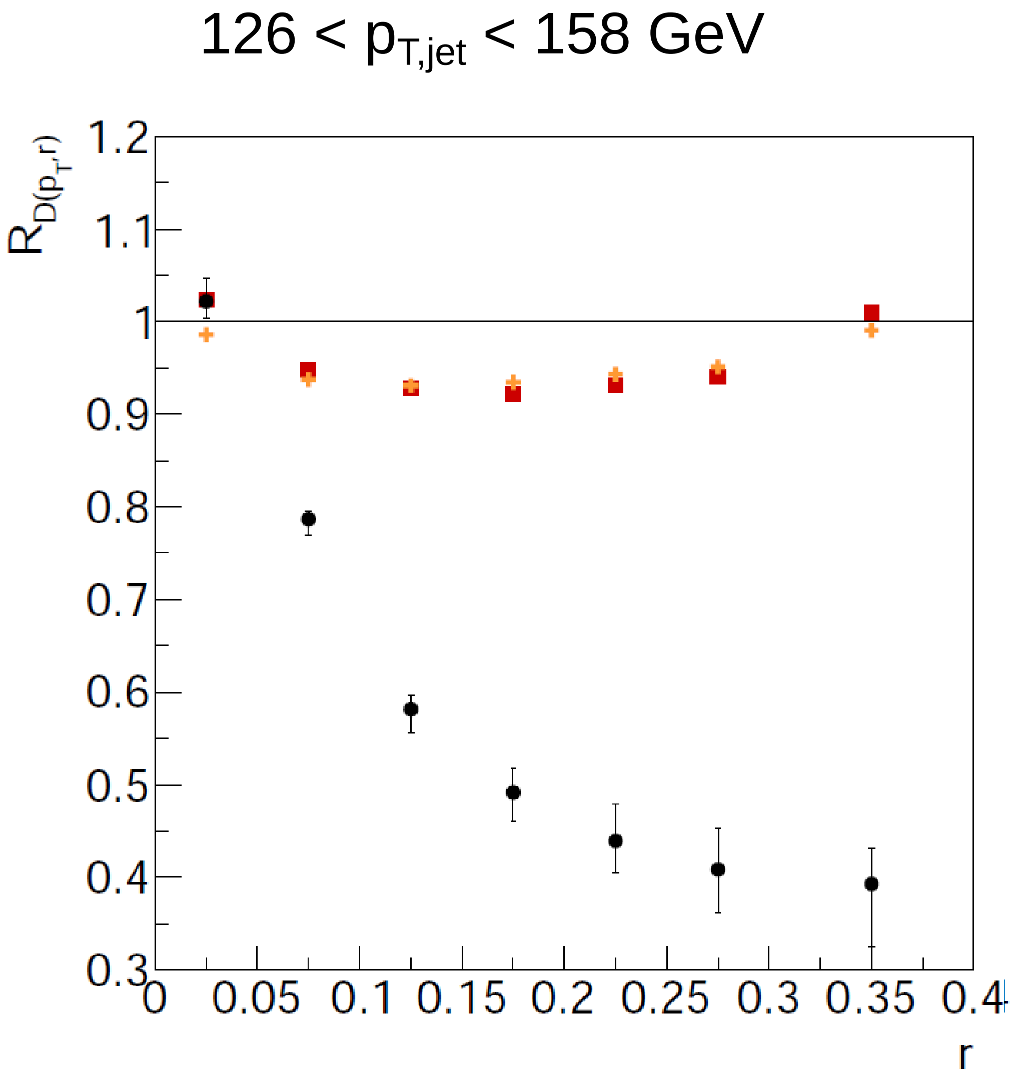


- **r<0.05**: values well reproduced (for all p_T^{ch} bins)
- **r>0.05**: trends similar but magnitude very different … … two particular possibilities:

1) **Input spectra** are not well modeled (sub-dominant contributions to jet p_T)

2) **Coherent picture** breaks for ~100 GeV jets at r~0.05

Modification of the jet shape



- **r<0.05**: values well reproduced (for all p_T^{ch} bins)
- **r>0.05**: trends similar but magnitude very different … ... two particular possibilities:

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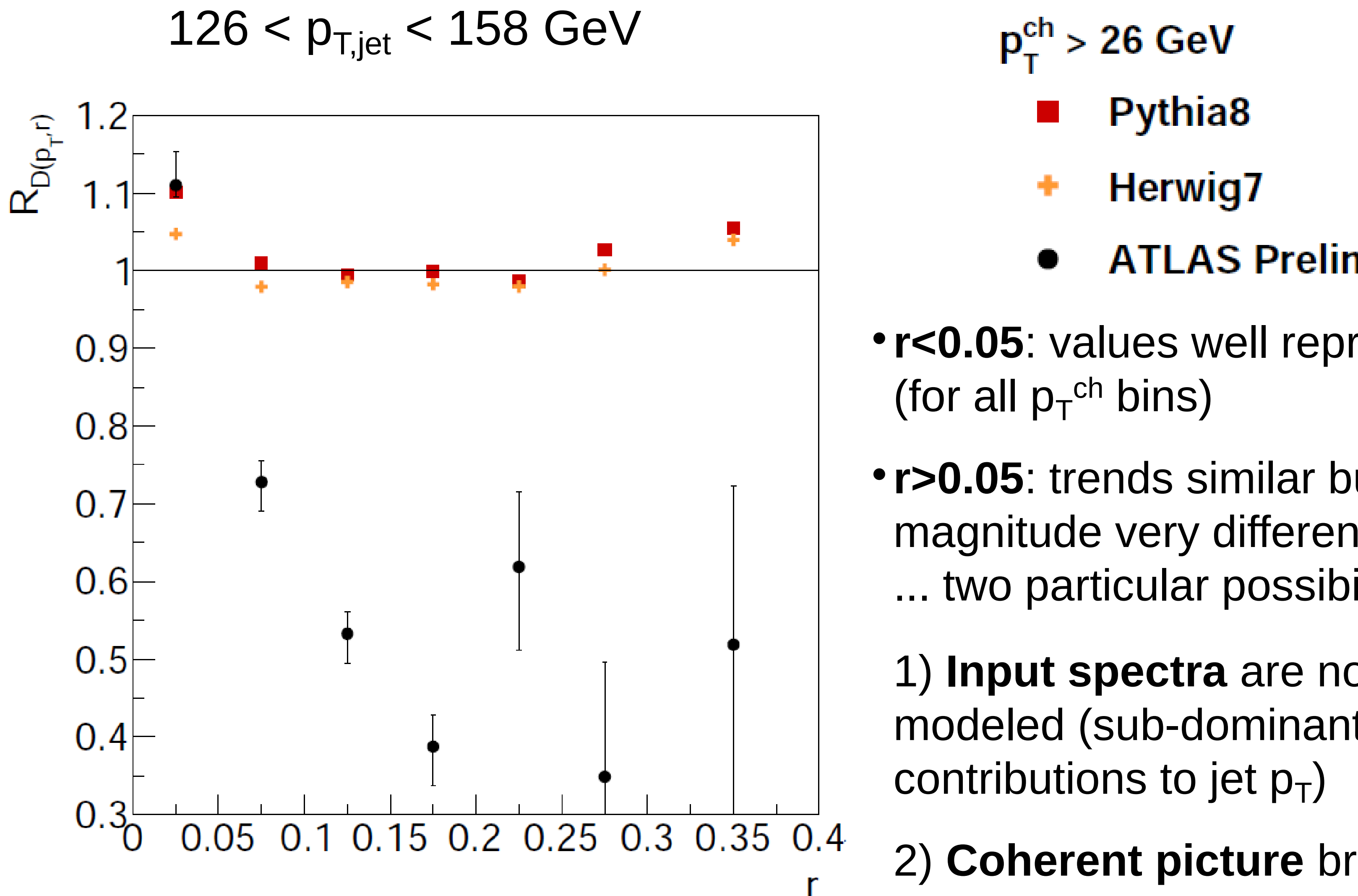

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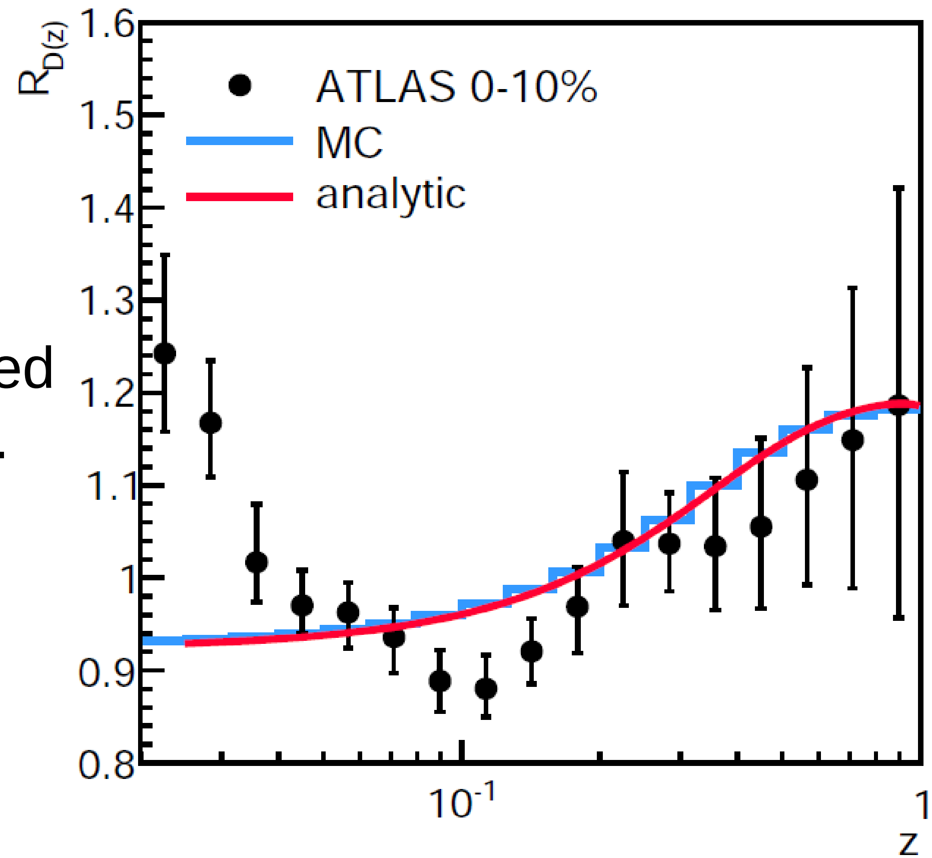
Modifications of fragmentation functions – a detail

Excess of low-z not due to flavor effects (due to in-cone radiation or recoil effects)

- These low-z hadrons contribute to the measured jet energy. Parameter s' contains this soft part.
- Soft part contributes to the measured fragmentation via denominator of z.

$$p_{\mathrm{T,jet}}^{\mathrm{measured}} = p_{\mathrm{T,jet}}^{\mathrm{quenched}} + p_{\mathrm{T}}^{\mathrm{soft}}$$



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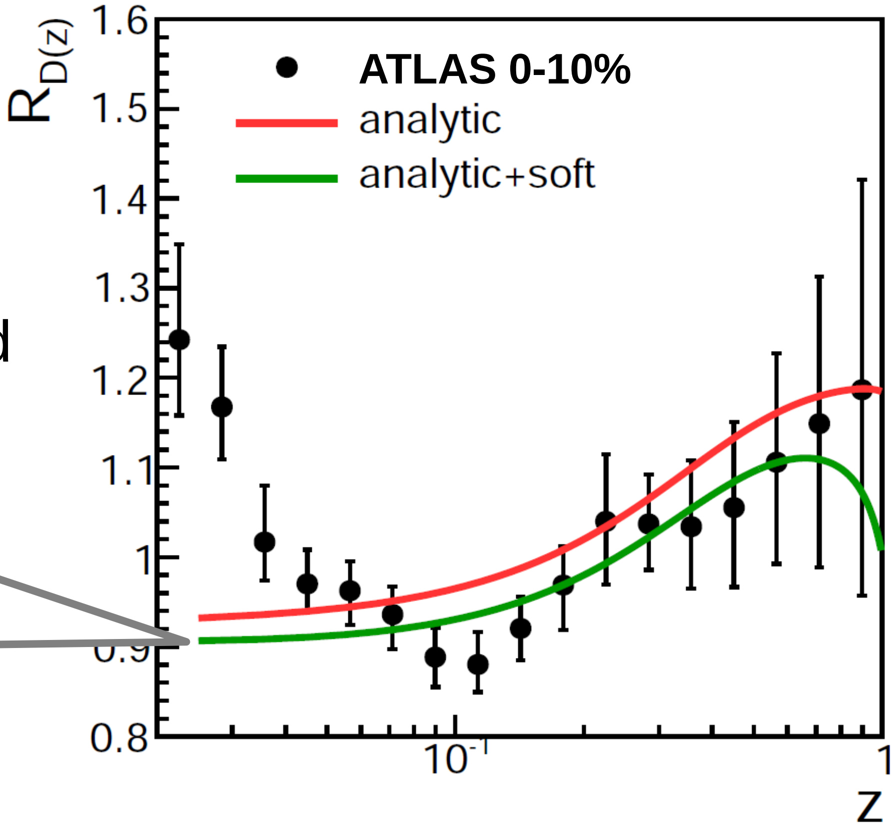
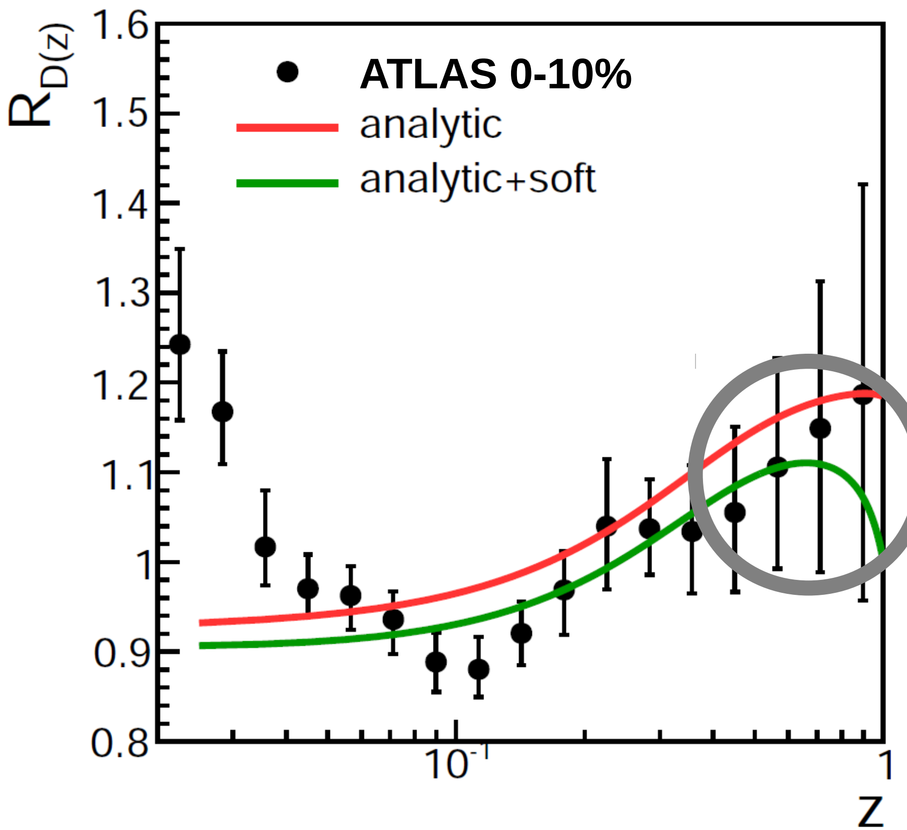


Contribution of soft hadrons to the jet energy can be estimated from the measurement at low-z => fragmentation distributions w/ correct soft contribution

Modifications of fragmentation functions – a detail

–> Prediction: detailed measurement of fragmentation at the highest-*z* (or lowest-ξ) should exhibit a depletion

$$p_{\mathrm{T,jet}}^{\mathrm{measured}} = p_{\mathrm{T,jet}}^{\mathrm{quenched}} + p_{\mathrm{T}}^{\mathrm{soft}}$$



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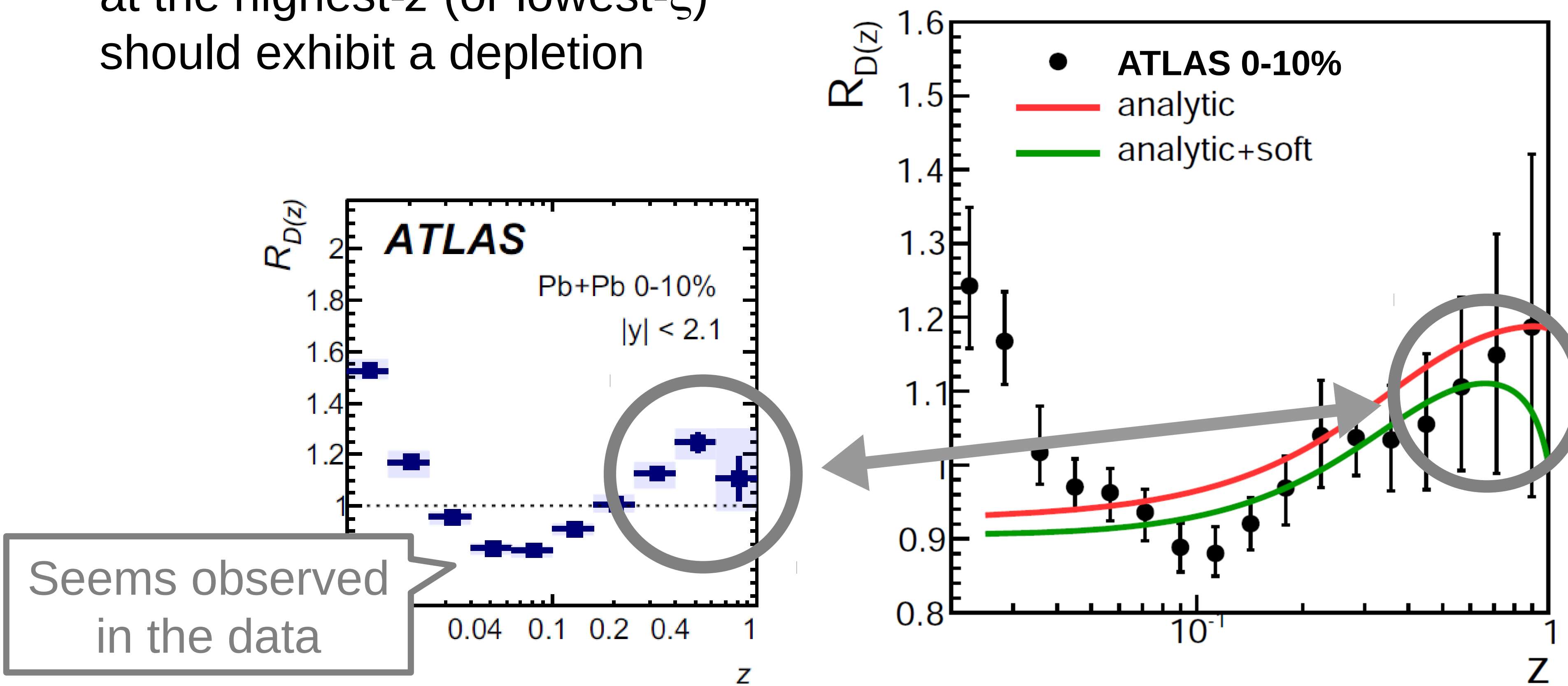
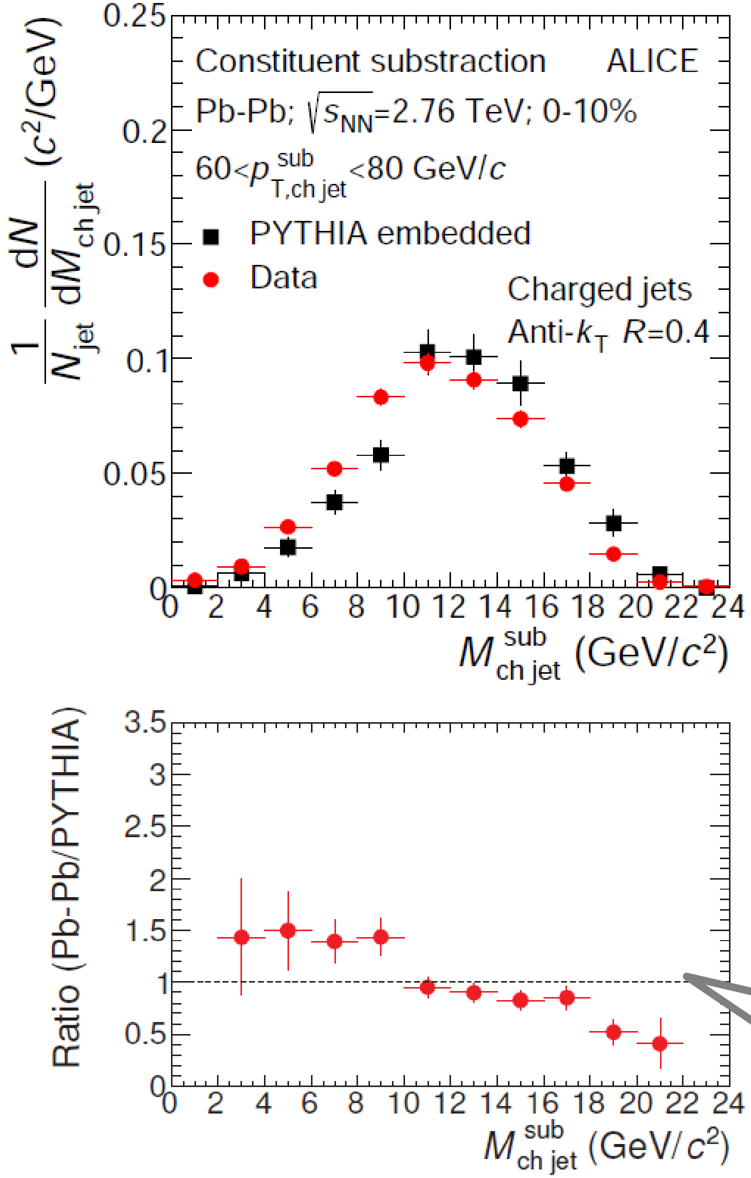
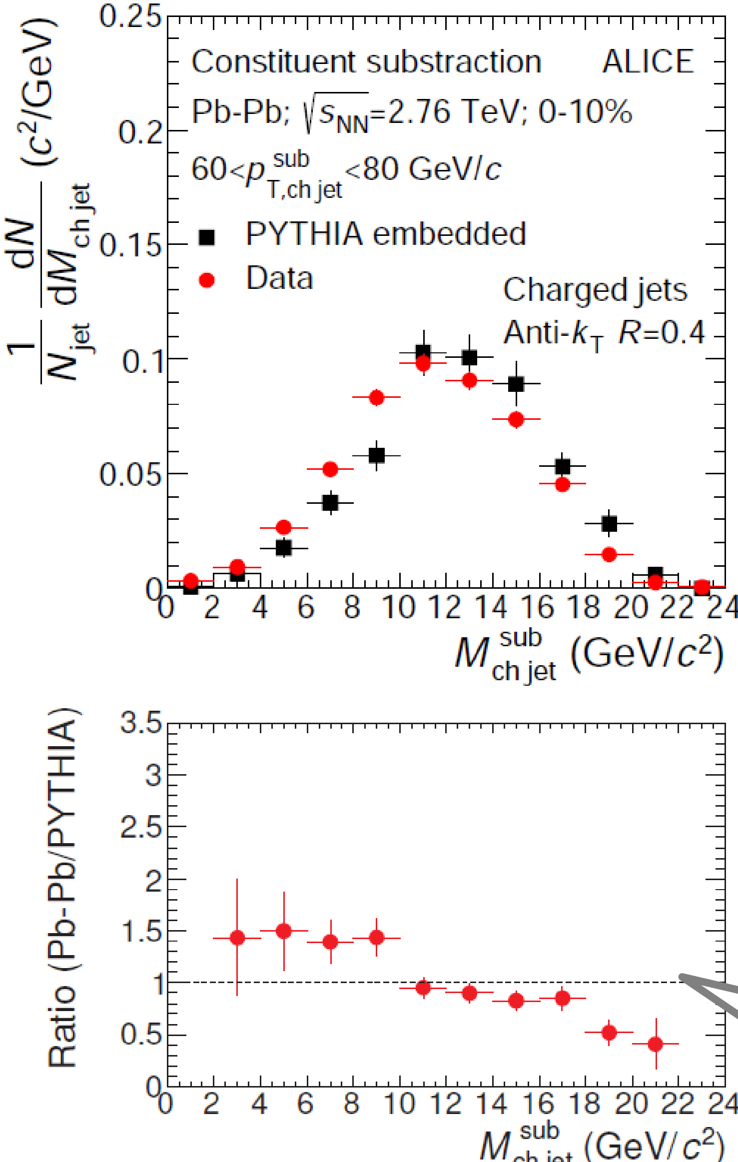


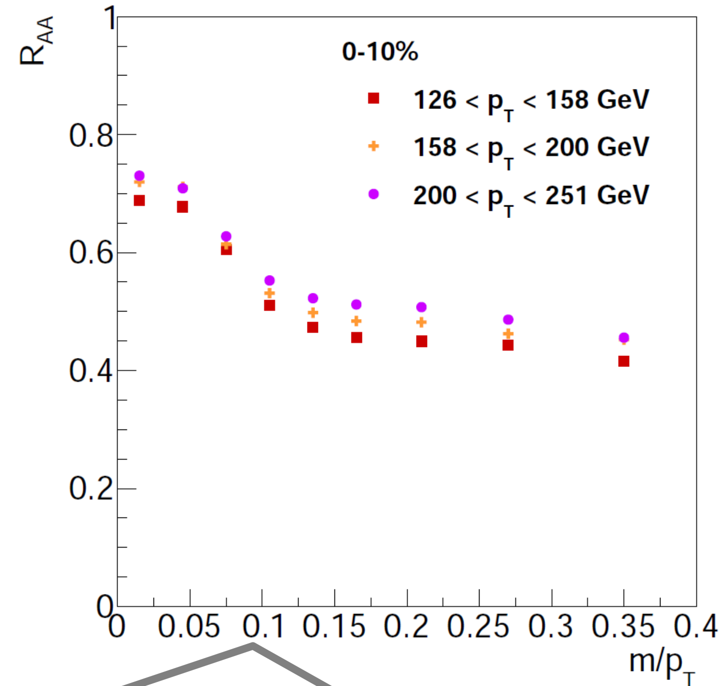
Example of other observables: non-groomed jet mass



A hint of possible shift to lower jet mass values seen in the data

Example of other observables: non-groomed jet mass





… but rather complicated observable: significant flavor dependence + dependence on recoil at low-p_T

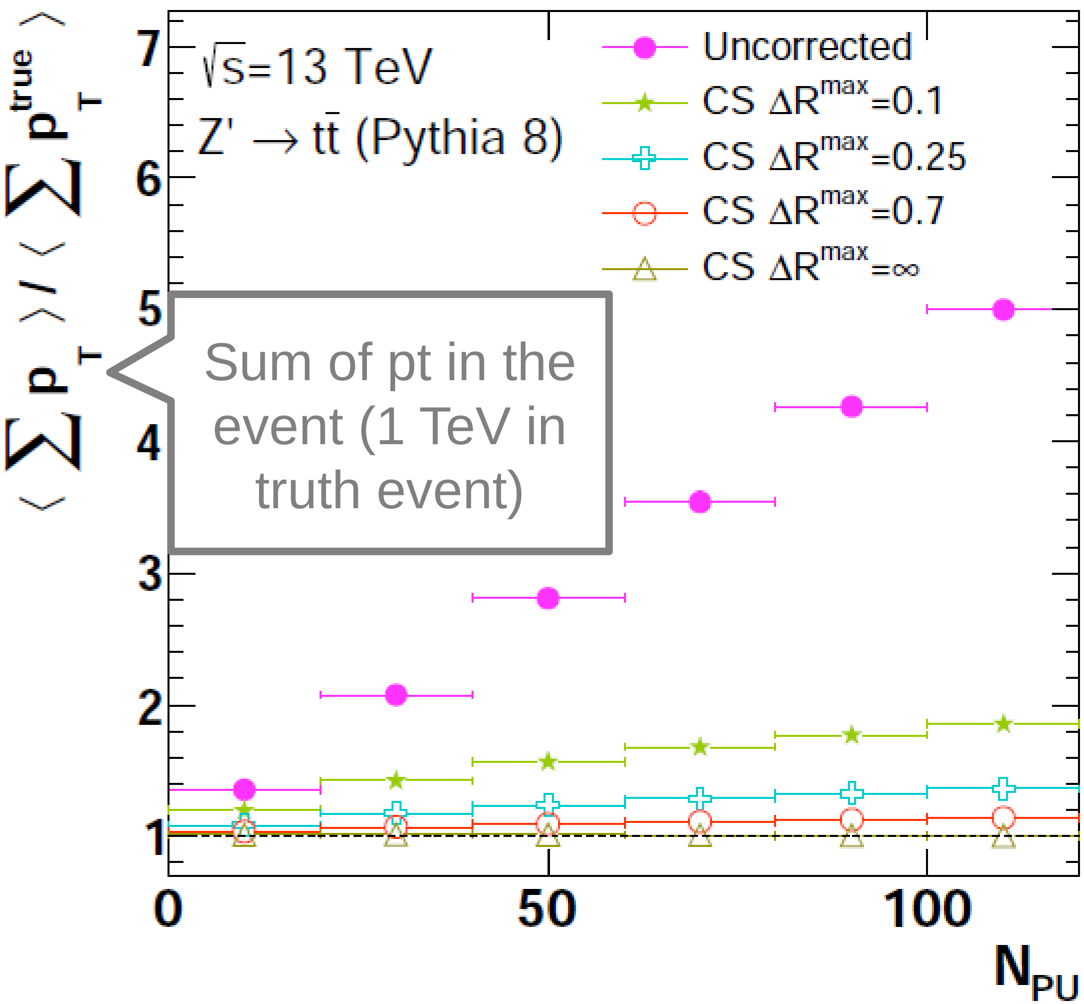
A hint of possible shift to lower jet mass values seen in the data

Summary

- **Flavor** dependence of the jet quenching seems to drive quite a lot of what we see in the data.
- **Coherence** effects seem to be important, but for jets with pt~100 GeV they seem to **break at r~0.05**.
- **Recoil** (or in-cone radiation) can modify kinematic regions where one would not expect that (e.g. high-z fragmentation).
- **Precision** is really needed:
  - precision **data** are needed **to understand details** (recoil via high-z fragmentation, jet shapes at low r; flavor via V-jets).
  - precision **MC** is needed to have the **reference under the control**.
- Perhaps one has already sufficient information at hand **with longitudinal & transverse structure** of the jet, esp. if they are done more differentially and fully corrected to particle level. Perhaps new substructure observables are not critically need.
- (More info in the backup.)

Slides with more information

Event-wide subtraction



- Idea: **first subtract** and only **then cluster**.
- Jet clustering is **less biased** after whole event correction.
- Improved performance wrt Jet-by-jet CS (not shown here).
- Can be used for **global event observables** such as missing-Et.

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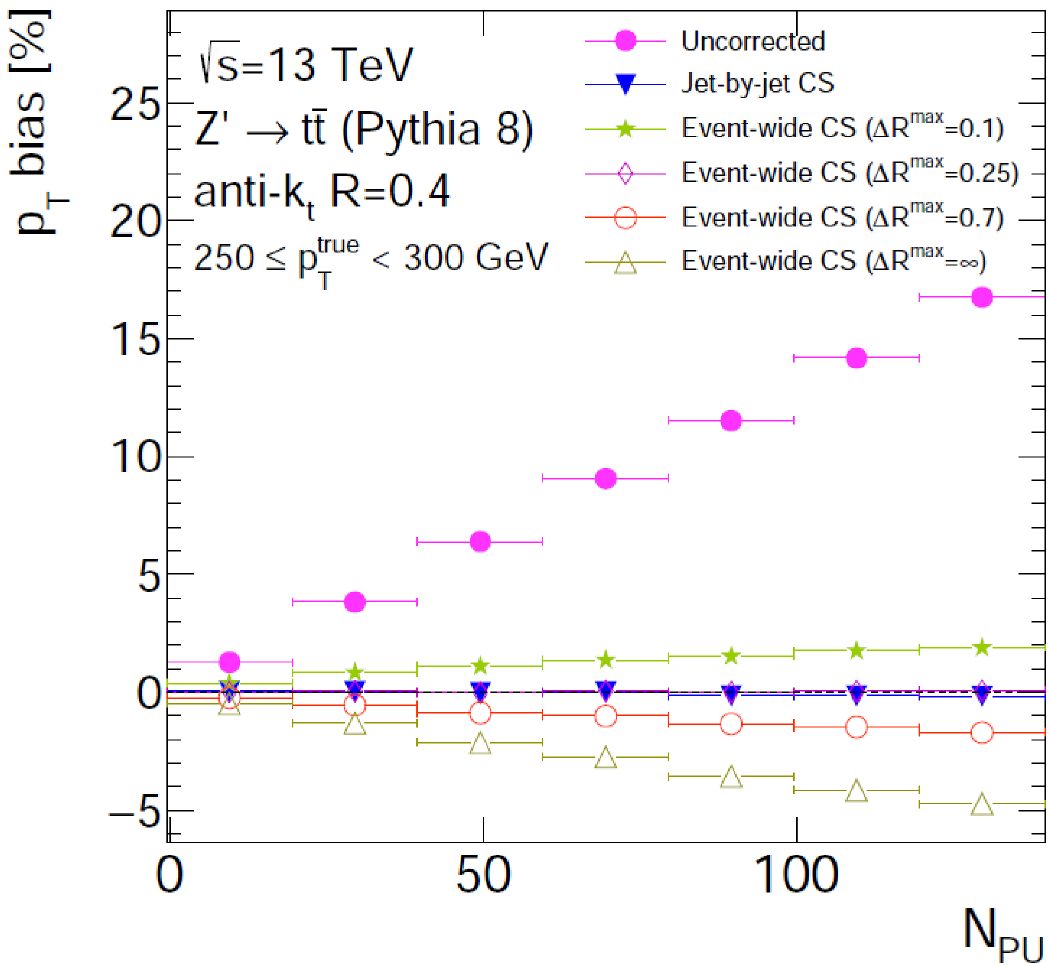

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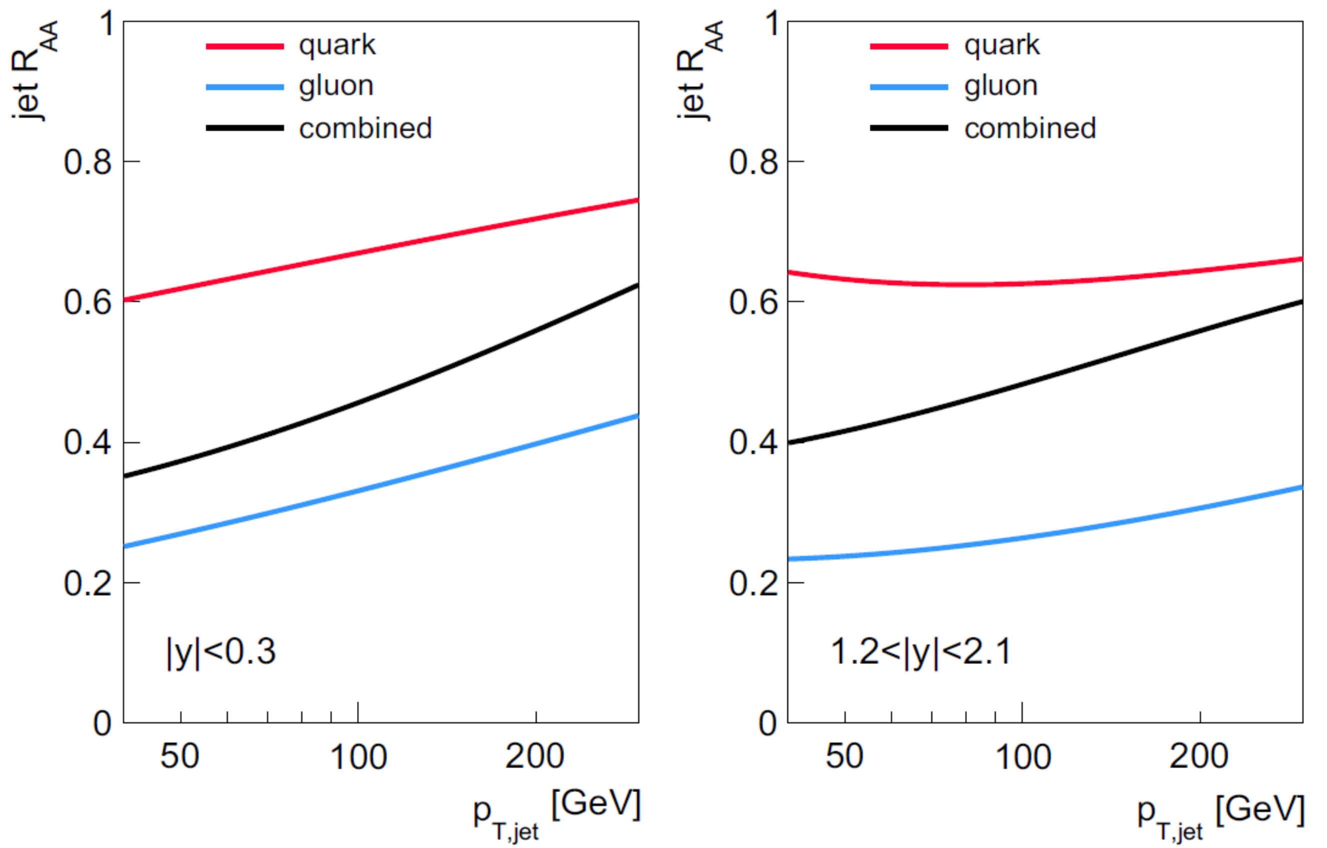
- Jet clustering is **less biased** after whole event correction.
- Improved performance wrt Jet-by-jet CS (not shown here).
- Can be used for **global event observables** such as missing-Et.
- … but a bias may come into the play ...

Jet R_{AA} in realistic model

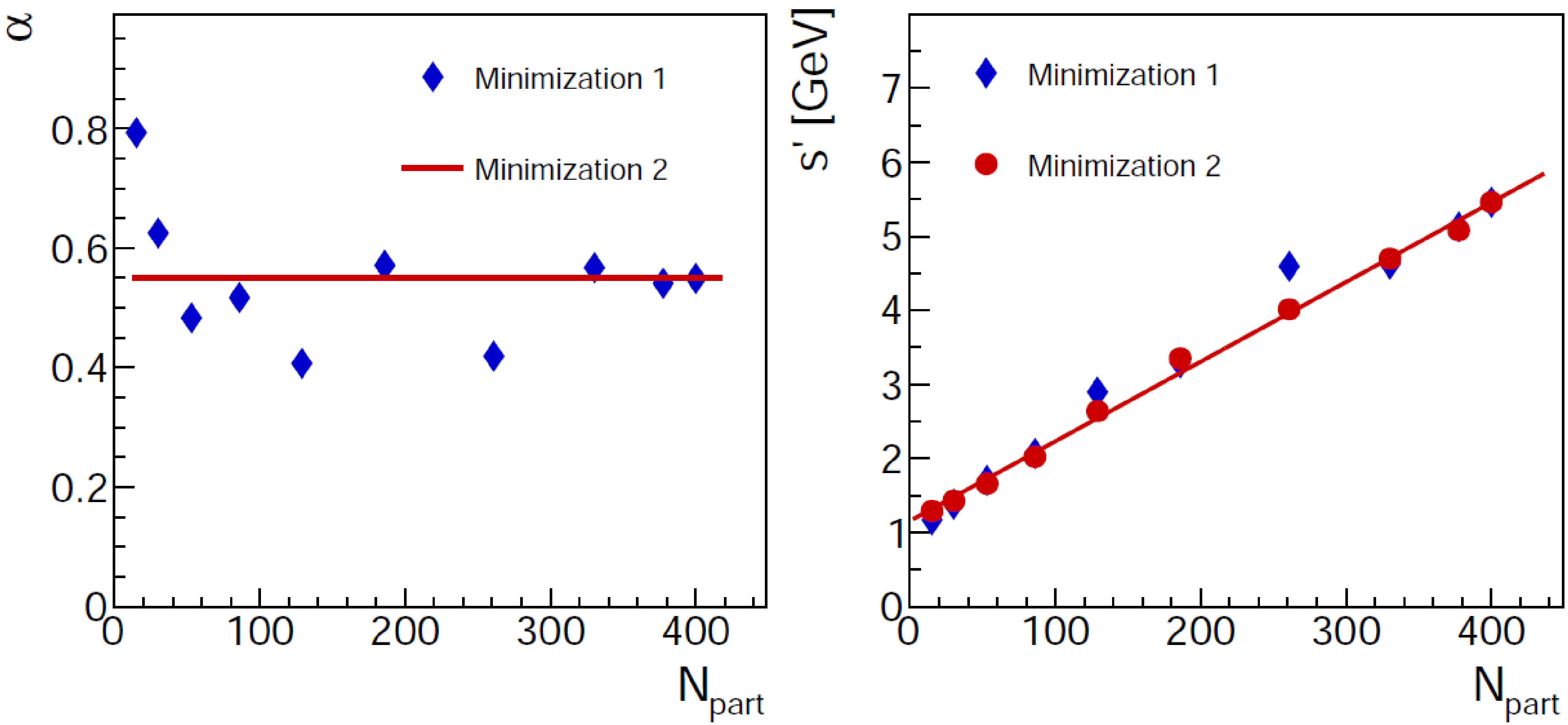


–> Flatness and no rapidity dependence of jet R_{AA} can be interpreted to be a result of different energy loss of quark and gluon initiated jets

Quantifying the parton energy loss (I.)

$$S = s' \left(\frac{p_T^{\text{jet}}}{p_{T0}} \right)^{\alpha}$$

Energy loss parameterized = encapsulated into two free parameters



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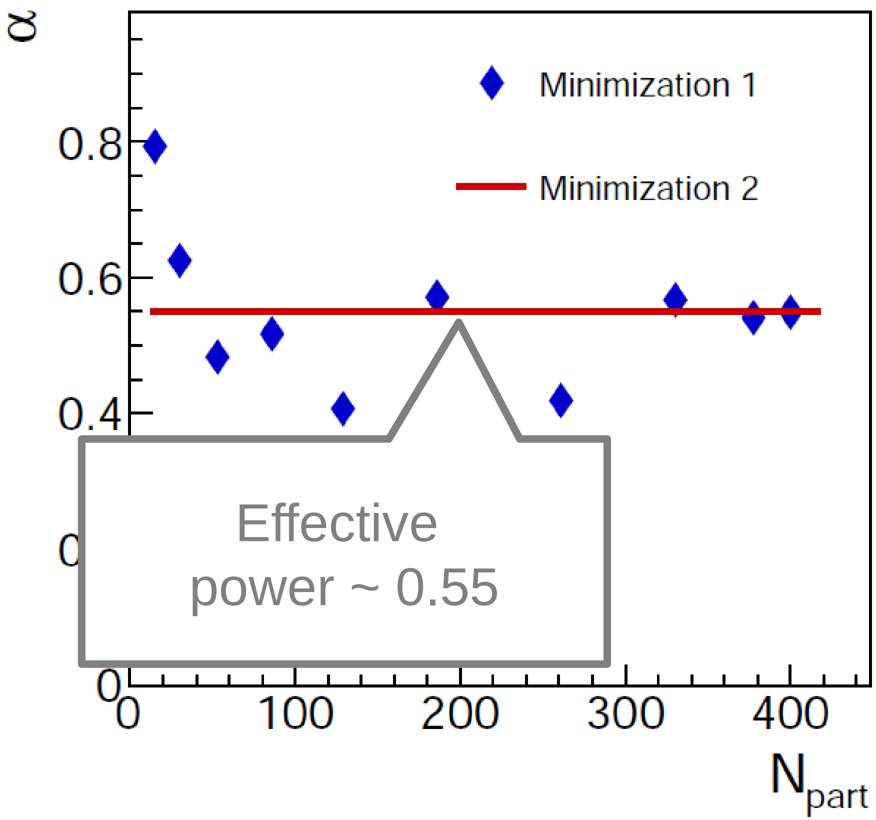
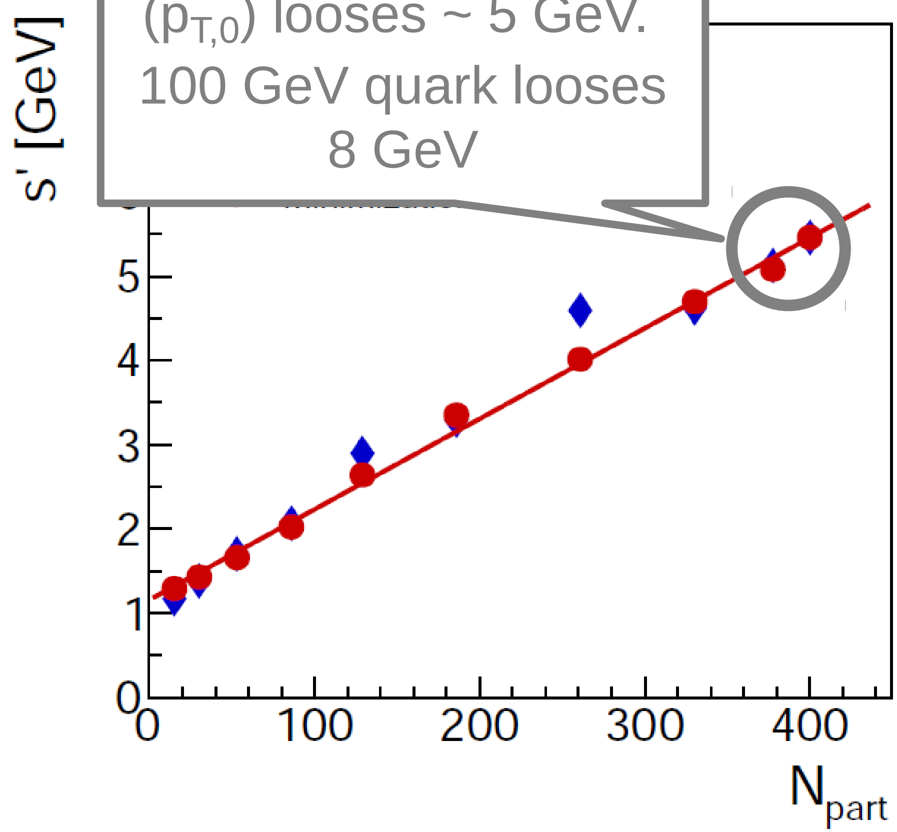


Effective power ~ 0.55



Quark with p_T=40 GeV ($p_{T,0}$) looses ~ 5 GeV. 100 GeV quark looses 8 GeV

Quantifying the parton energy loss (I.)

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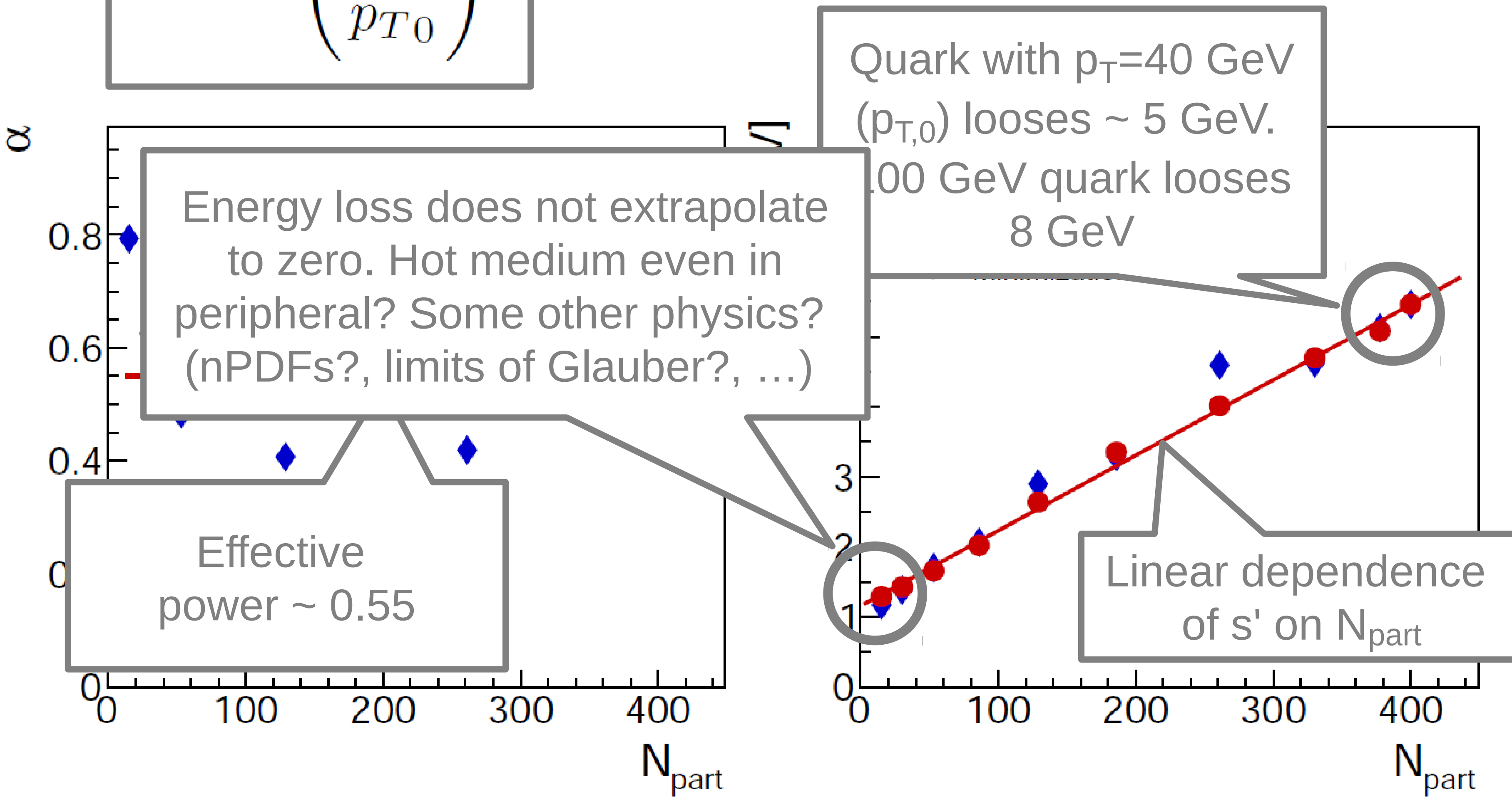
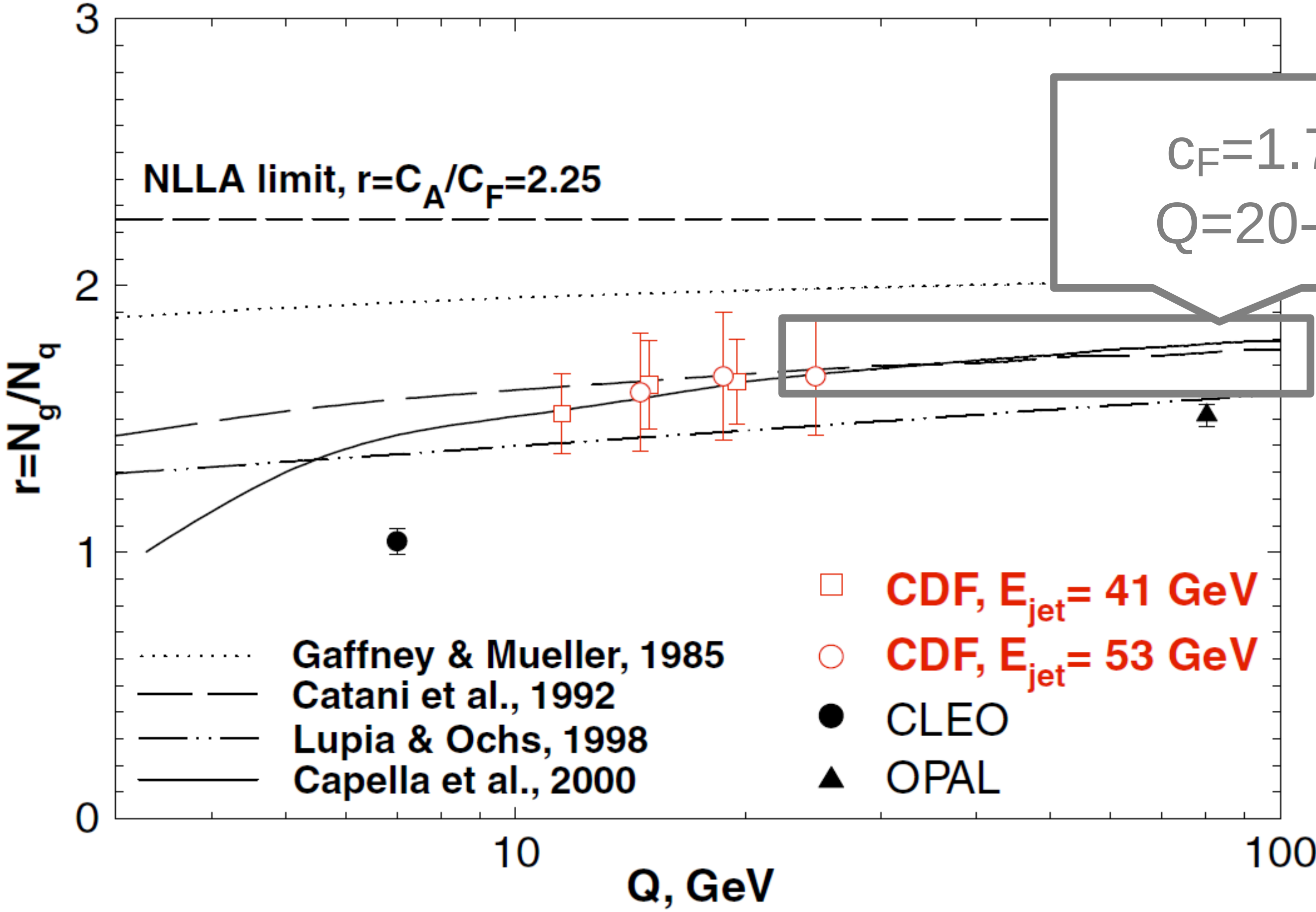


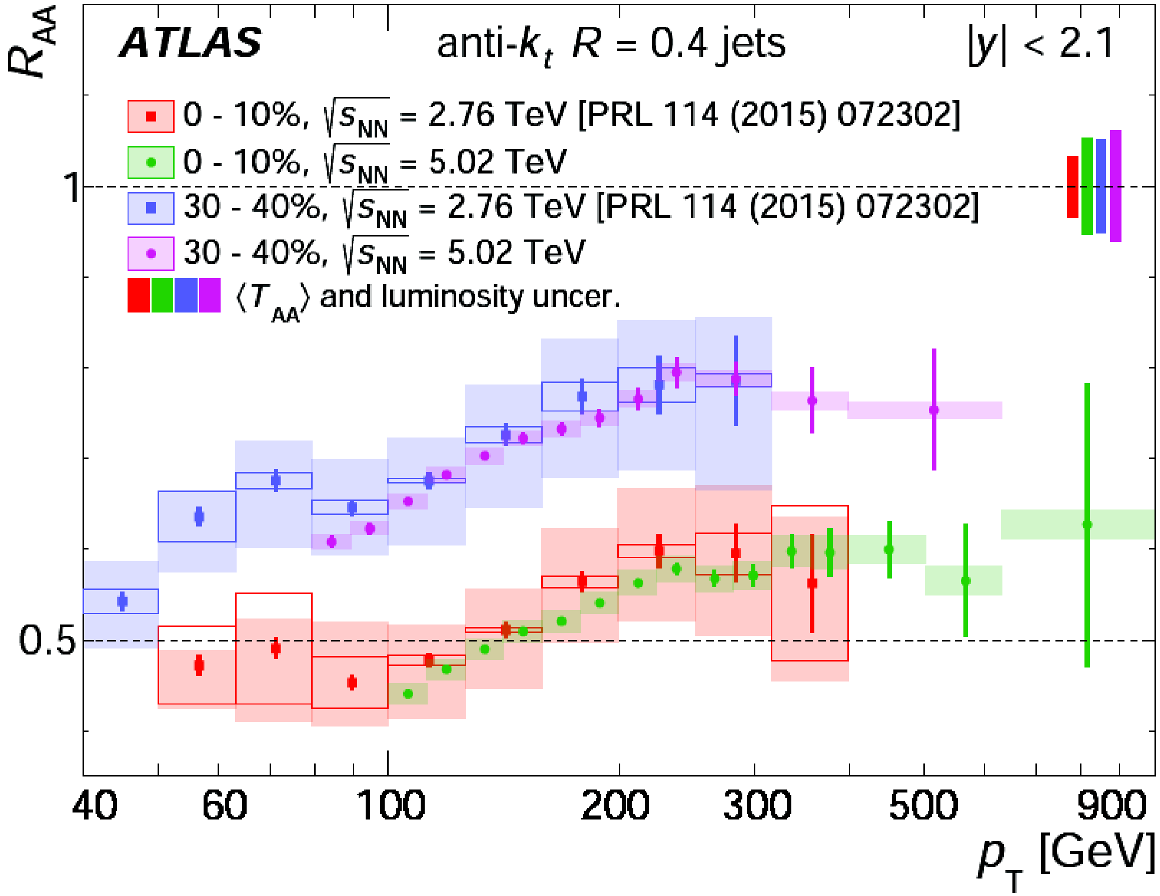
Quantifying the parton energy loss: c_F



c_F=1.7-1.8 for Q=20-100 GeV

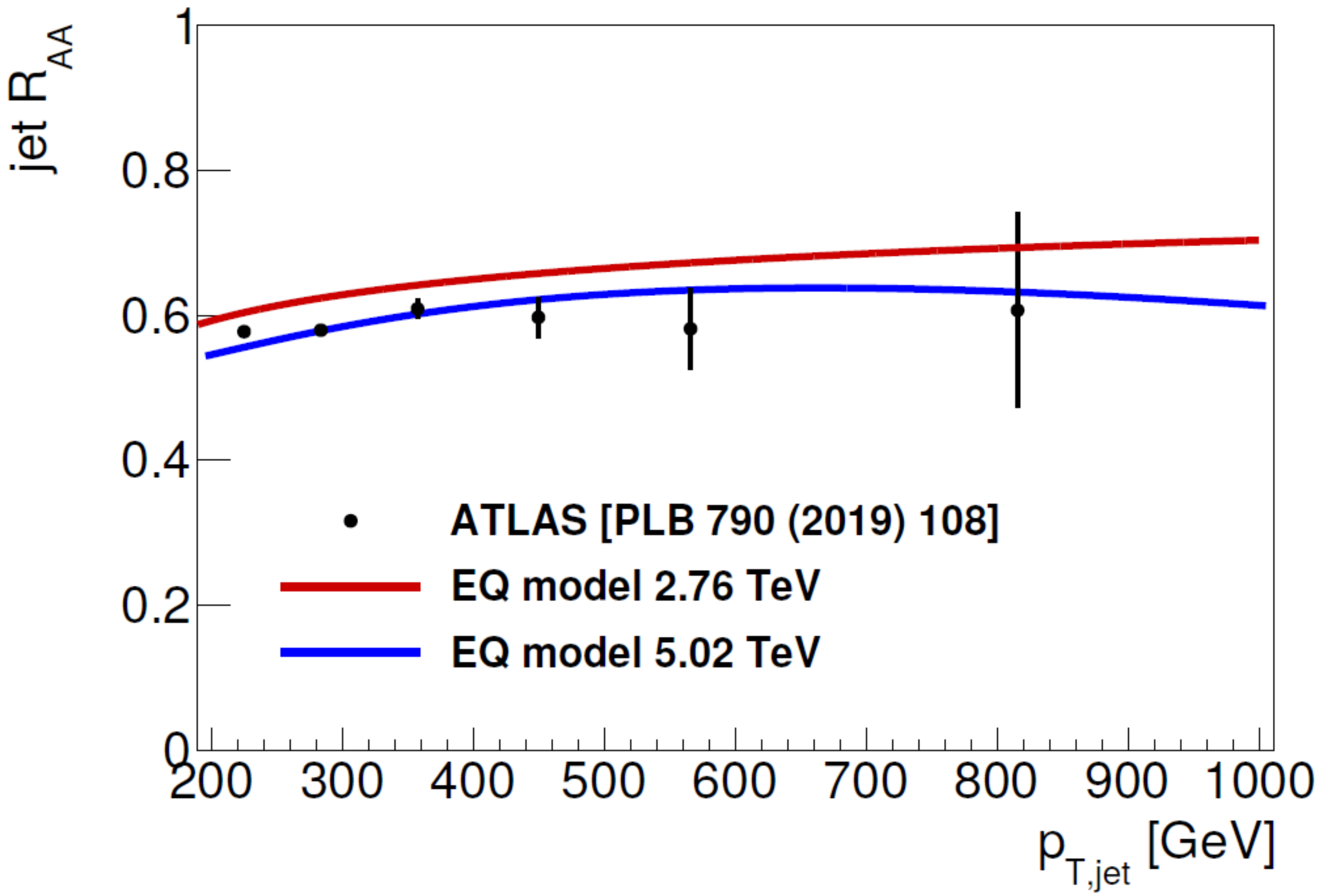
- Vacuum value of c_F measured and calculated in pQCD (MLLA)

5.02 TeV versus 2.76 TeV



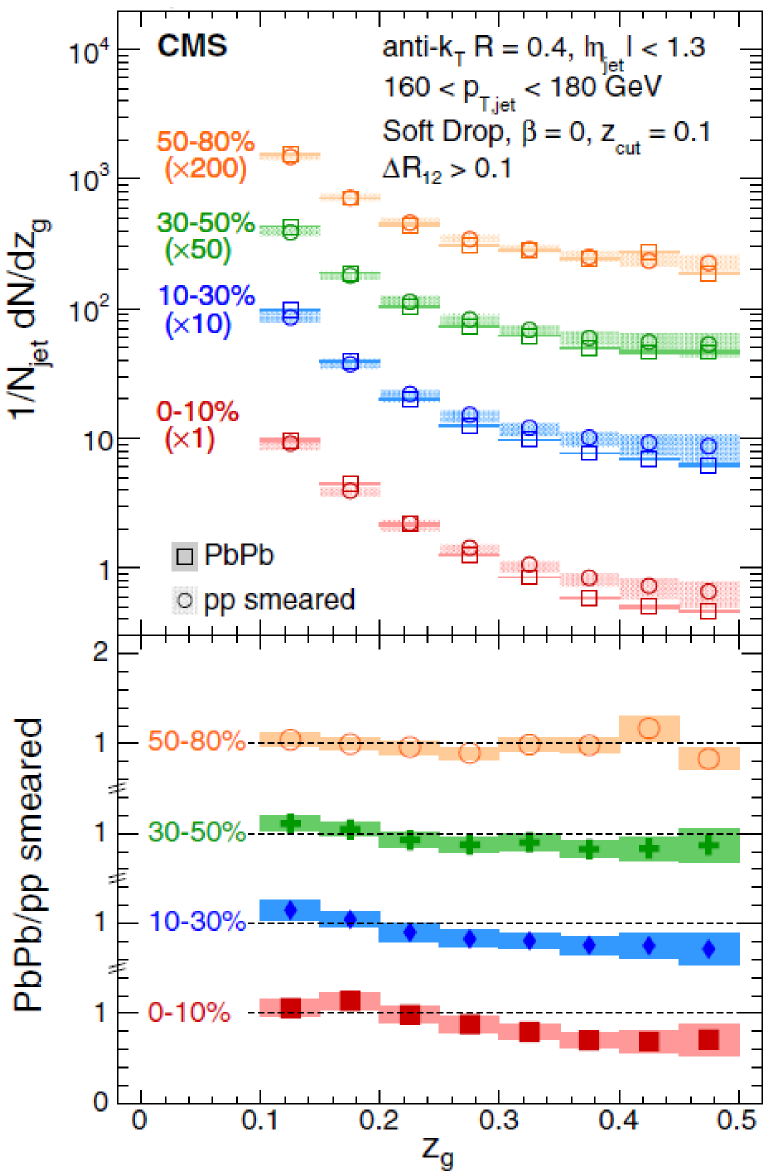
- Same jet R_{AA} … but that does not imply same energy loss.
- Spectra shape and flavor admixture are different => energy loss must be different.

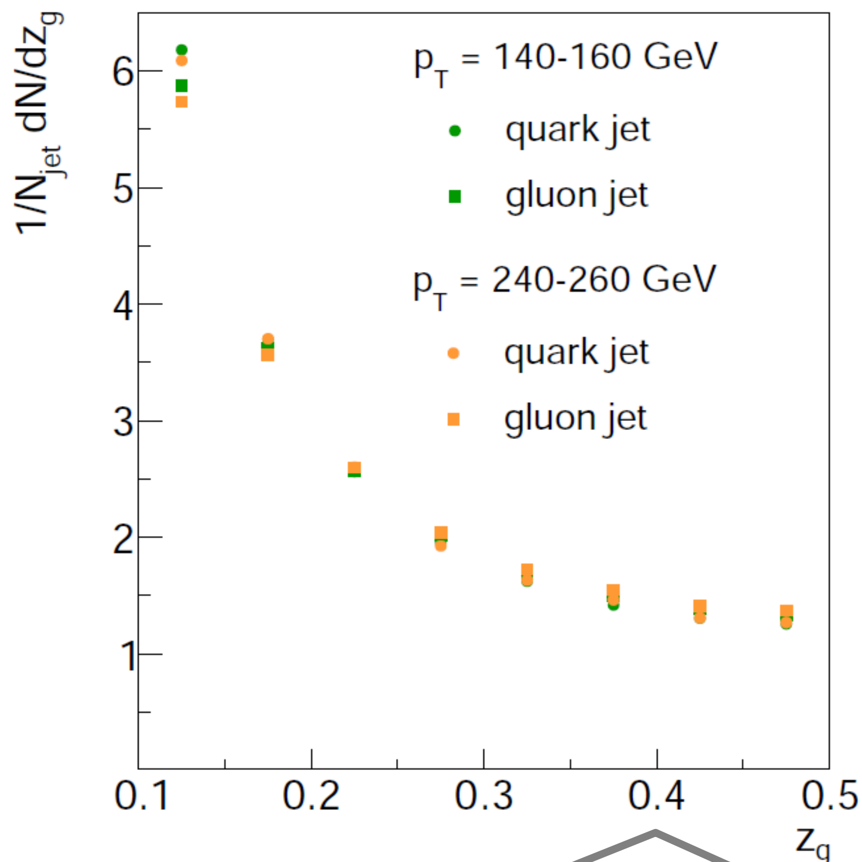
5.02 TeV versus 2.76 TeV



- Same jet R_{AA} … but that does not imply same energy loss.
- Spectra shape and flavor admixture are different
  => energy loss must be different.
- About 10% larger energy loss at 5.02 TeV compared to 2.76 TeV.

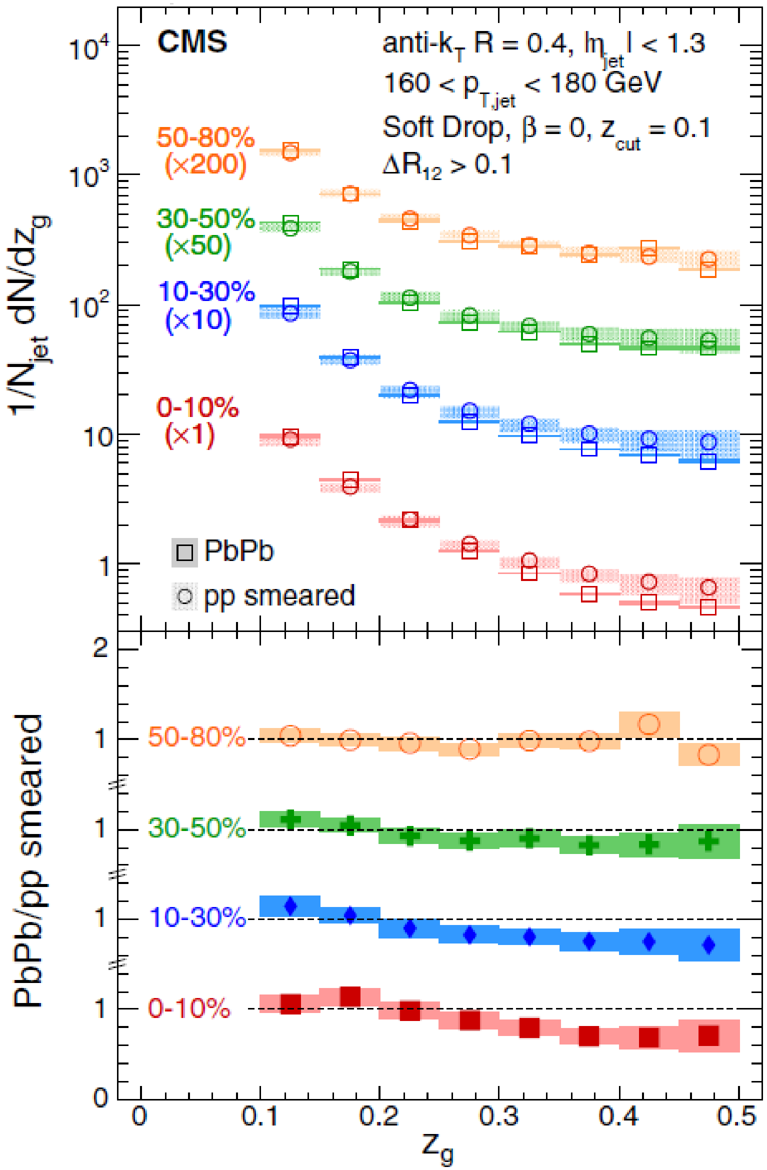
Groomed z_g … checking the impact of jet flavor

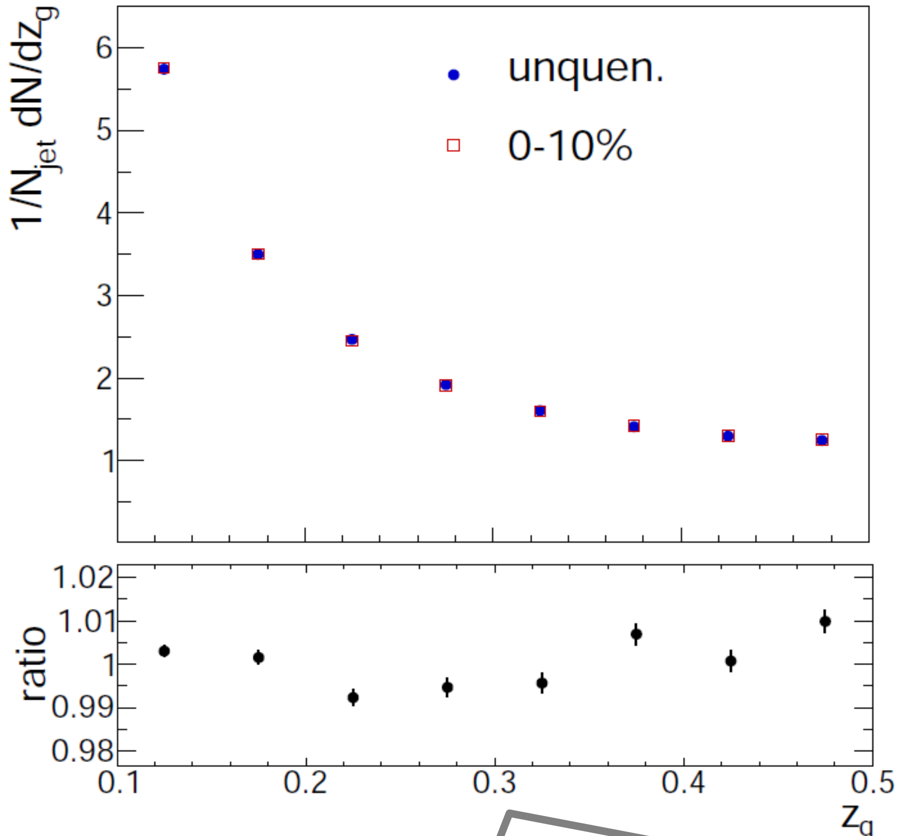




Using PYTHIA:
z_g does not depend much on the flavor or jet p_T

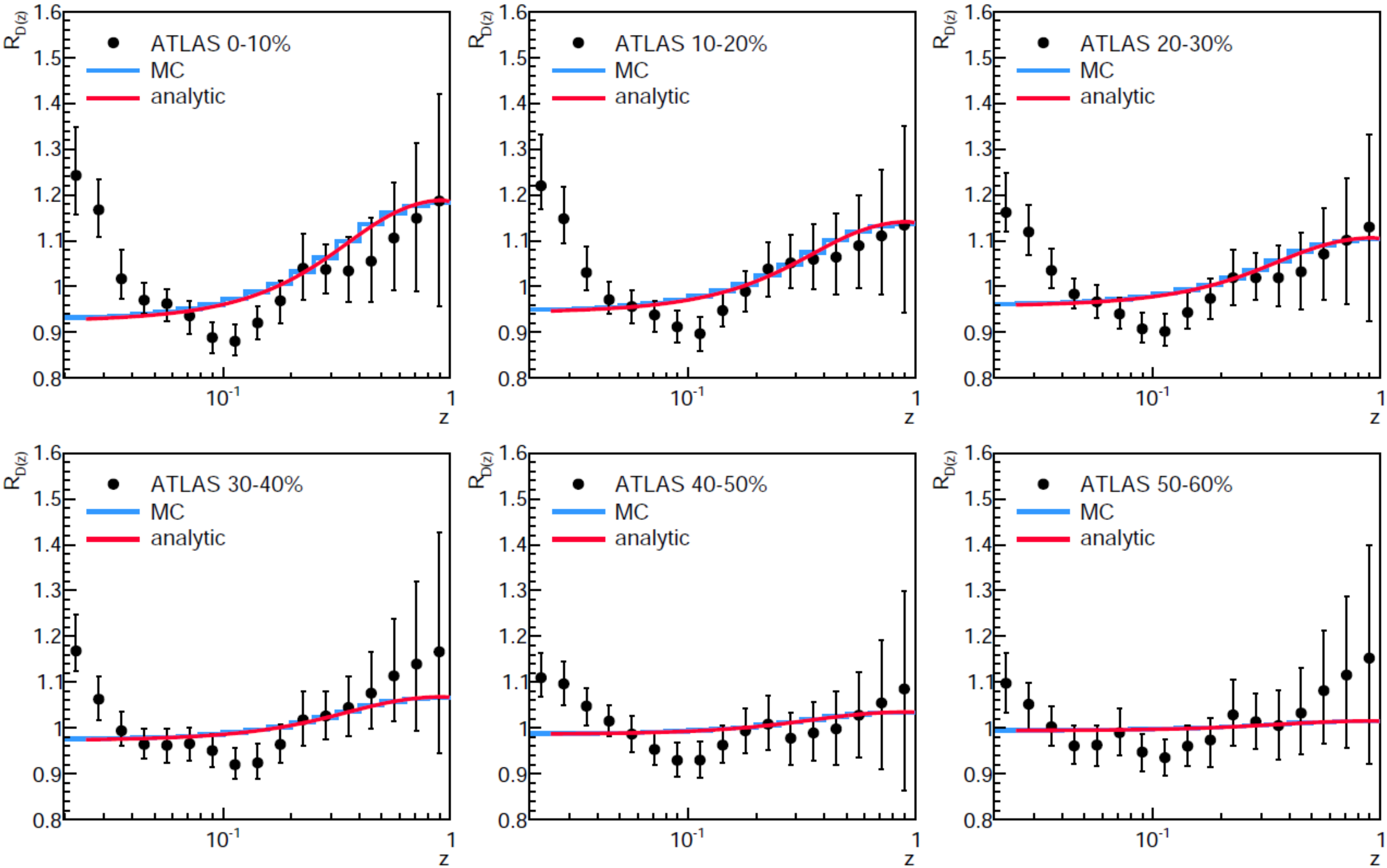
Groomed z_g … checking the impact of jet flavor





Same procedure as for modeling fragmentation functions => no modification seen => measured modification not due to a flavor

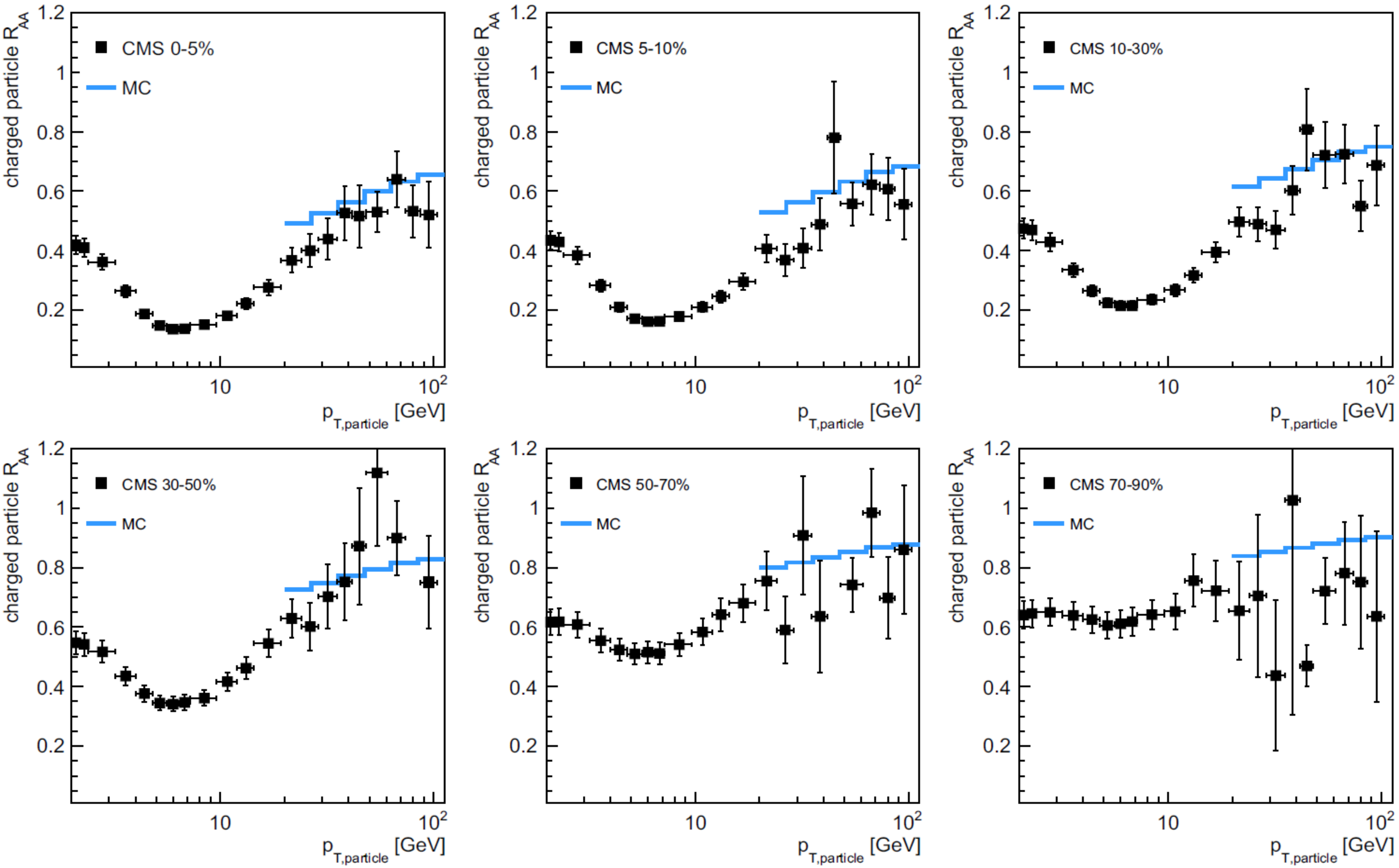
Modifications of fragmentation functions



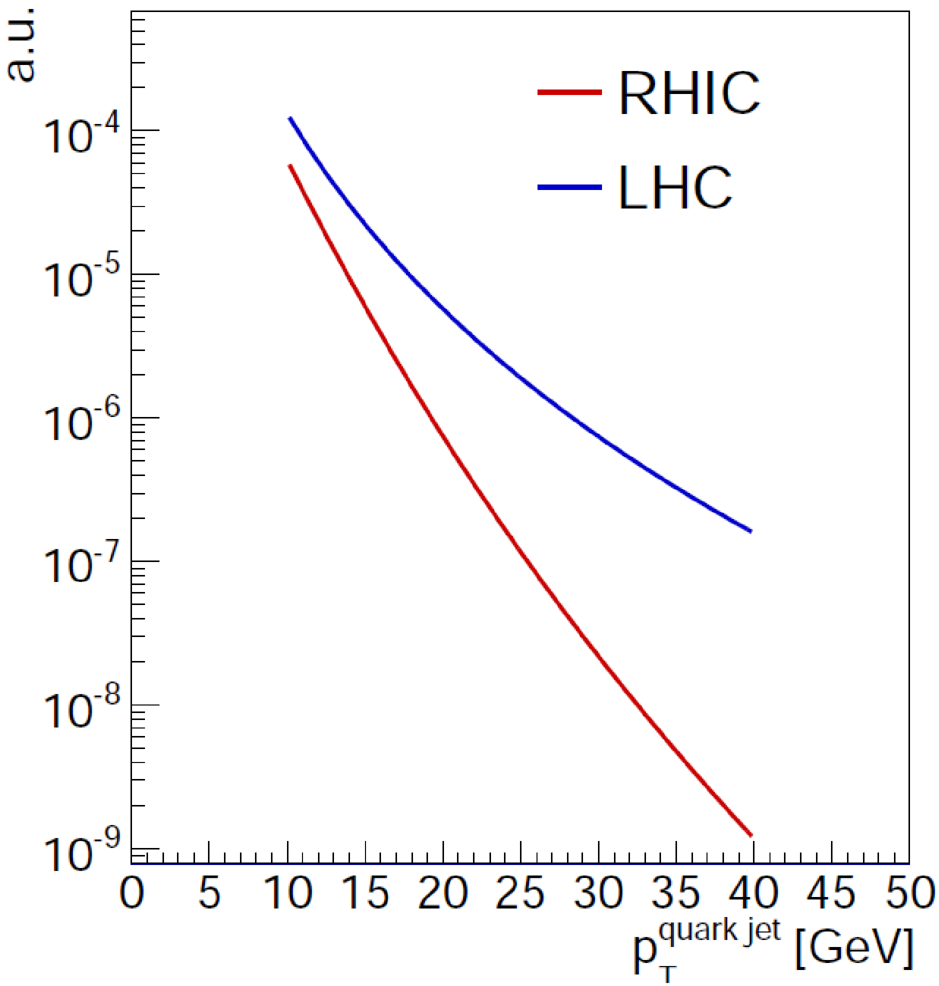
From jet internal structure to charged particle R_{AA}

Each particle of a given p_T must be in a jet of the same or higher p_T
=> Charged particle R_{AA} (at high-p_T) = convolution of flavor dependent jet suppression and fragmentation functions

From jet internal structure to charged particle R_{AA}



Jets at RHIC versus LHC

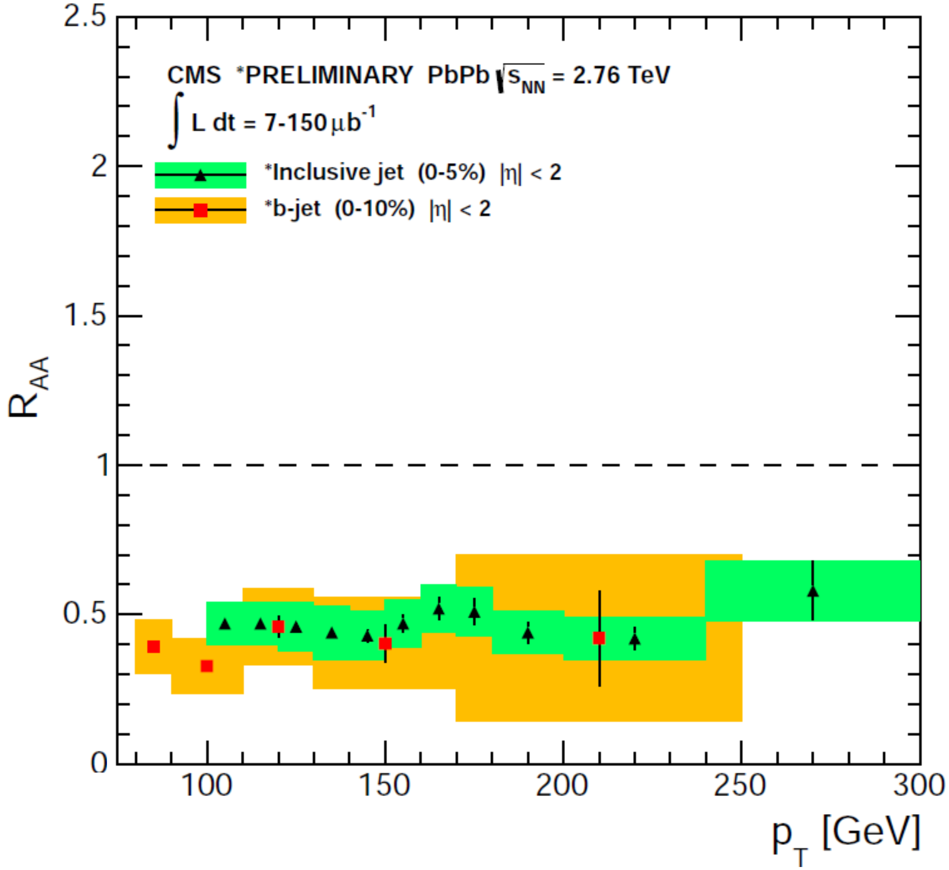


- Jets very different between LHC and RHIC
- Jet spectra for a given flavor more steep at RHIC
- Flavor composition also different

–> Will impact charged particle R_{AA}

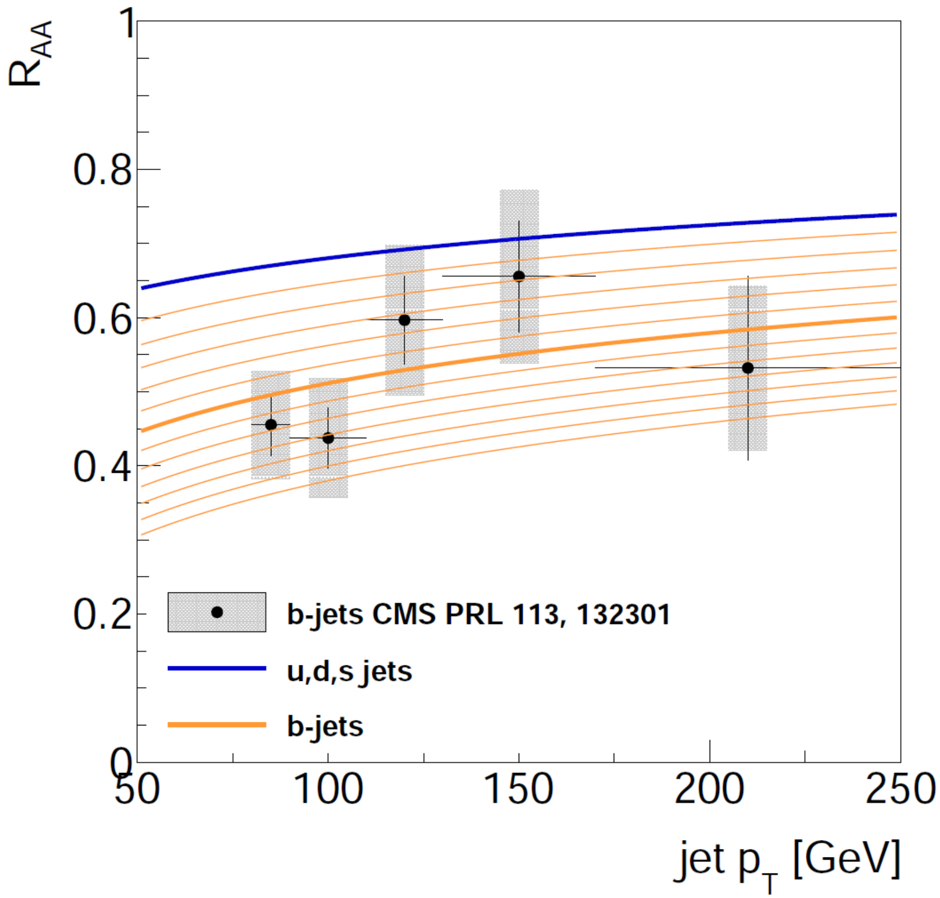
–> Apply the effective quenching factors extracted at the LHC to RHIC jets

b-jets suppression



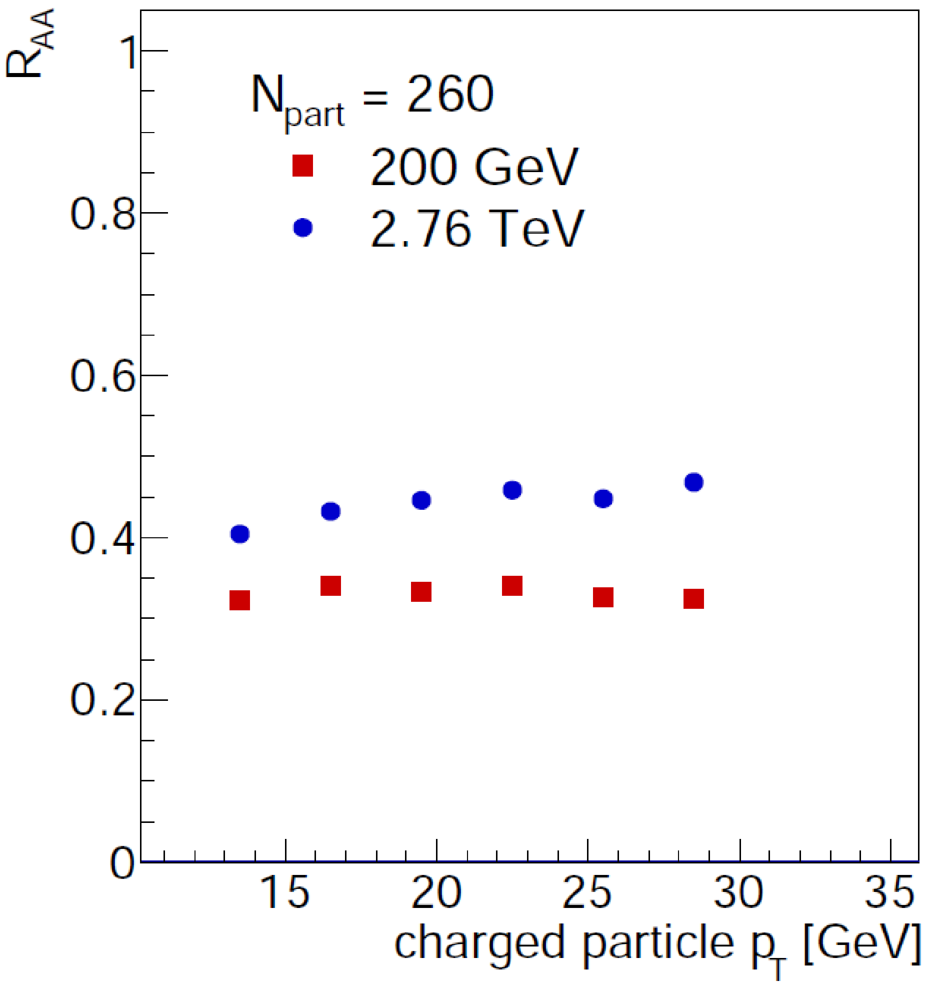
- b-jet R_{AA} … comparable with inclusive jet R_{AA} … but again, spectral shapes are different
- Moreover, just one flavor => direct comparison misleading

b-jets suppression



- b-jet R_{AA} … comparable with inclusive jet R_{AA} … but again, spectral shapes are different
- Moreover, just one flavor => direct comparison misleading
- Use the model + b-jet cross-section measurement to quantify the difference between inclusive jets and b-jets.
- Results of minimization wrt to (statistically limited) data + including role of gluon splitting: ***b*-jets are suppressed by 1.5±0.4 more** than light quark jets.

Charged particle R_{AA}
RHIC vs LHC



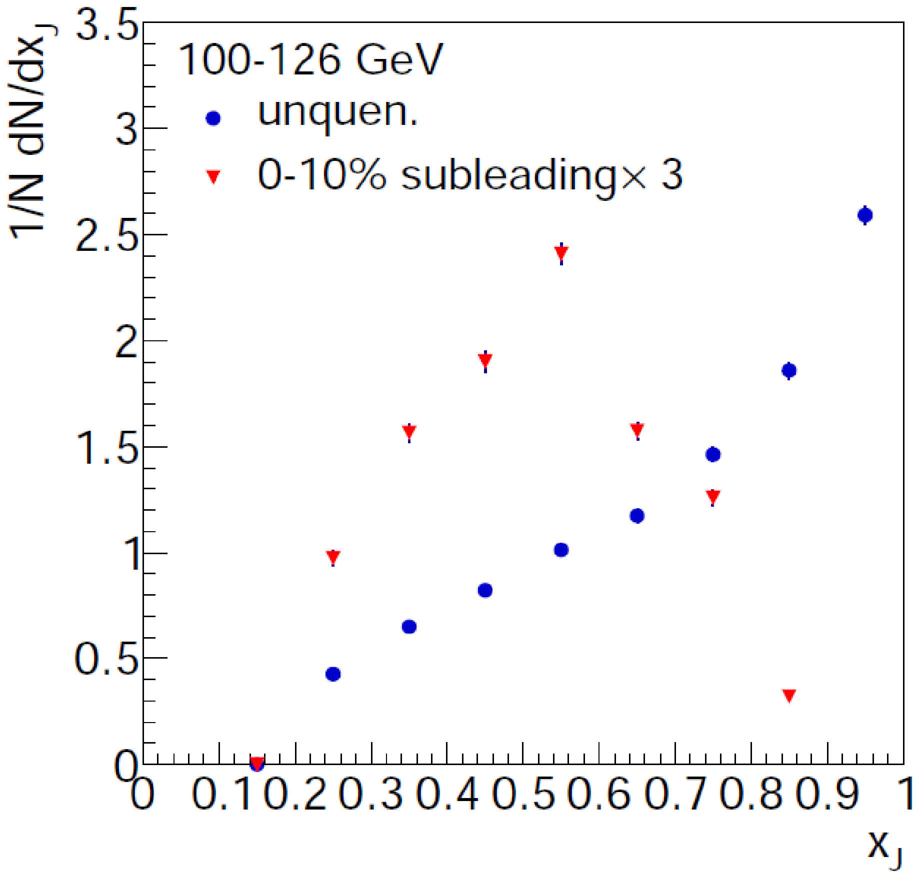
- Underlying jet spectra very different between RHIC and LHC
- Effective quenching factors from LHC applied to RHIC parton/jet spectra
- **Same quenching leads to smaller R_{AA} in the case of RHIC**

=> Initial parton spectra and flavor composition are very important for the extraction of the size of jet quenching

Dijet asymmetry

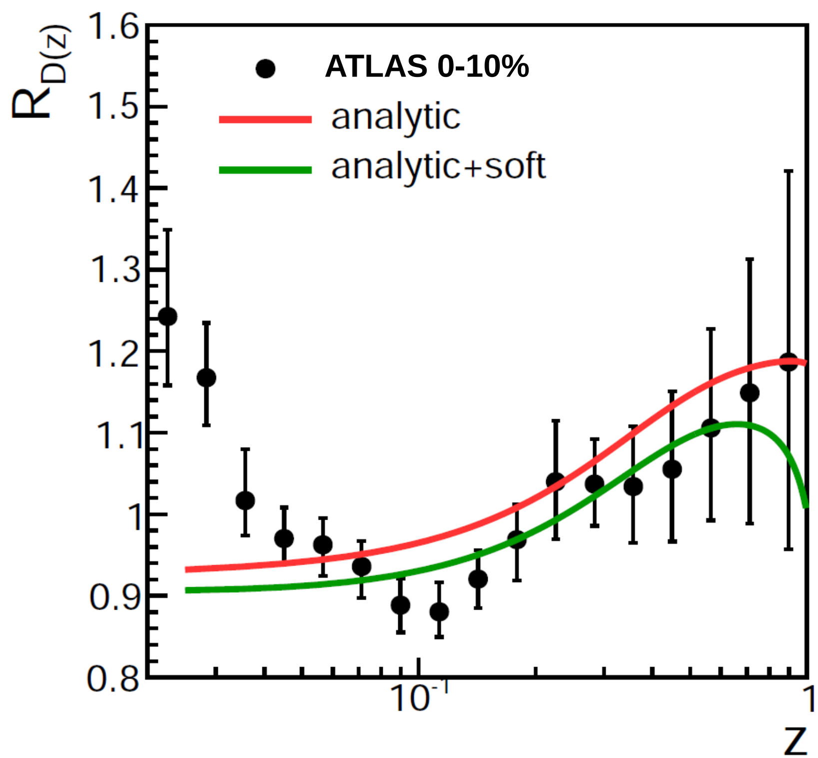
–> The subleading jet is quenched very differently then the leading jet –> quantify

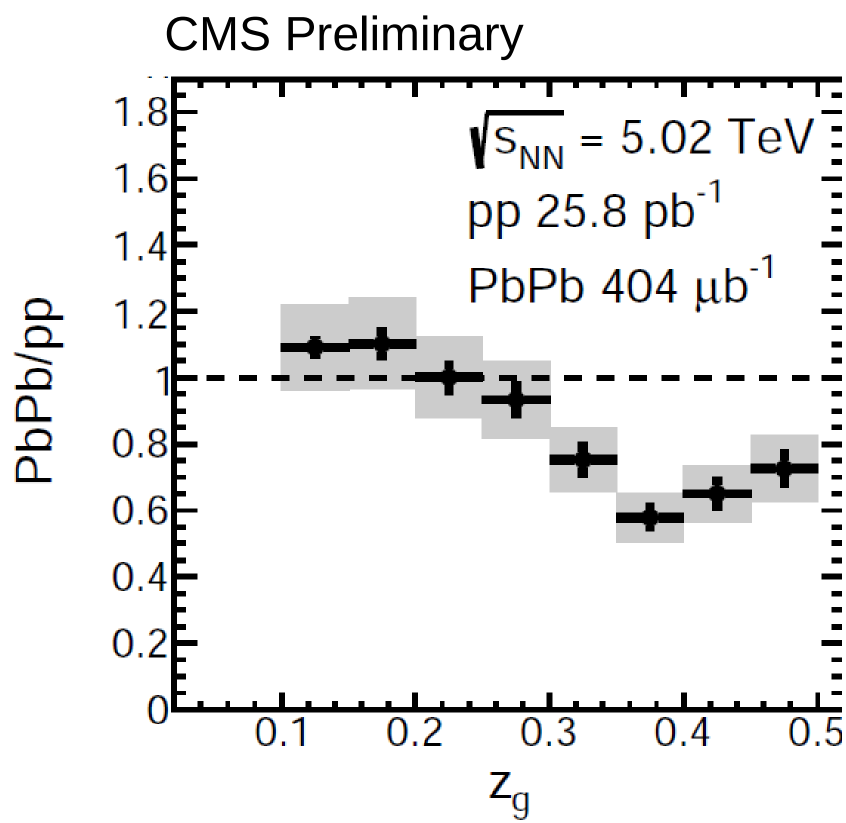
–> The subleading jet in the maximum of the x_J is suppressed by a factor of ~3 larger than the leading jet



What about other objects?

Data tell us that the medium largely sees a jet as one object
=> what about other objects with a structure that are suppressed?





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J/Ψ & $\Psi(2S)$

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J/Ψ & $\Psi(2S)$

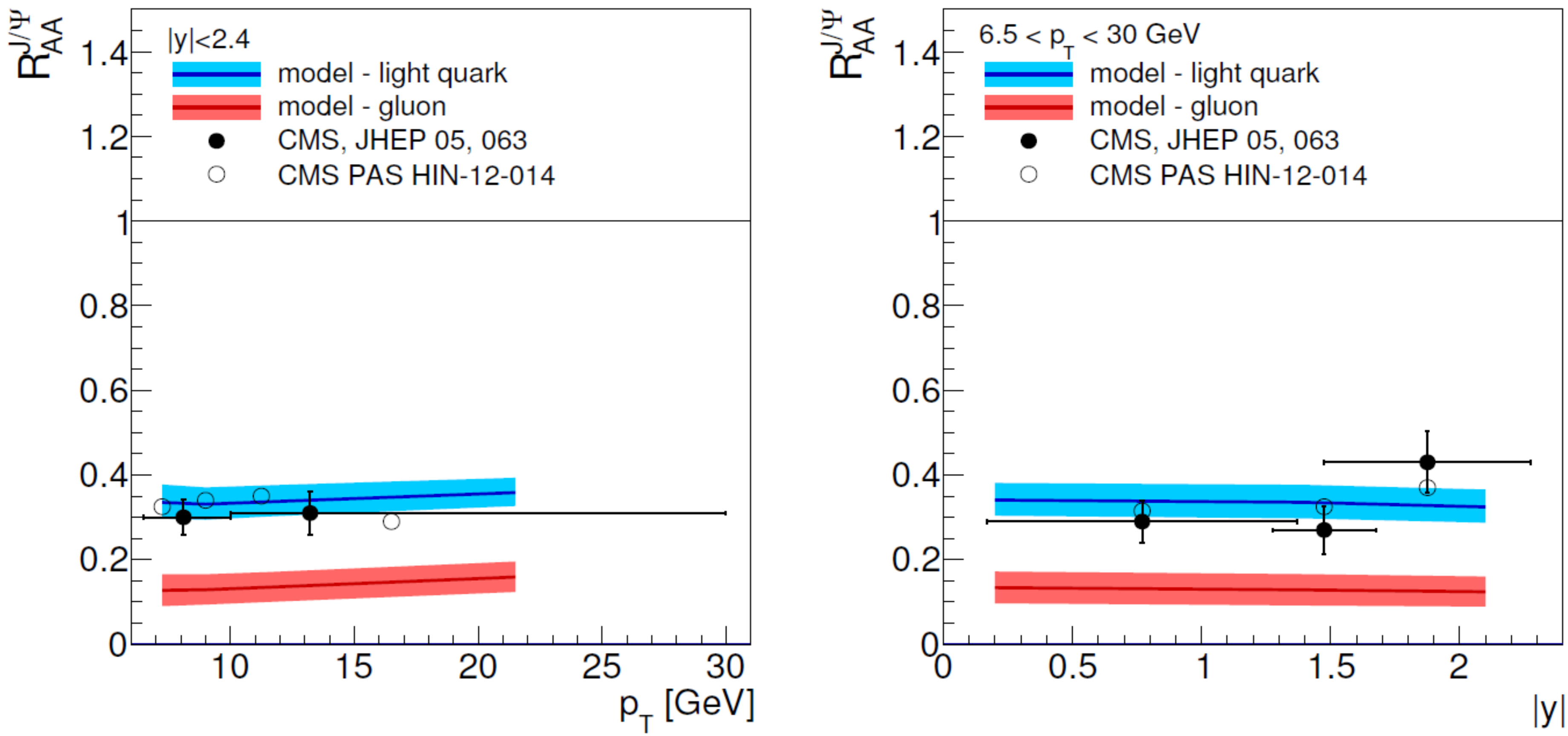
… check the differences between the suppression of jets and charmonia at high-p_T (at the LHC at mid-rapidity)

Input:

– Measured pp spectra of charmonia (cannot rely on out of the box PYTHIA or other generator)

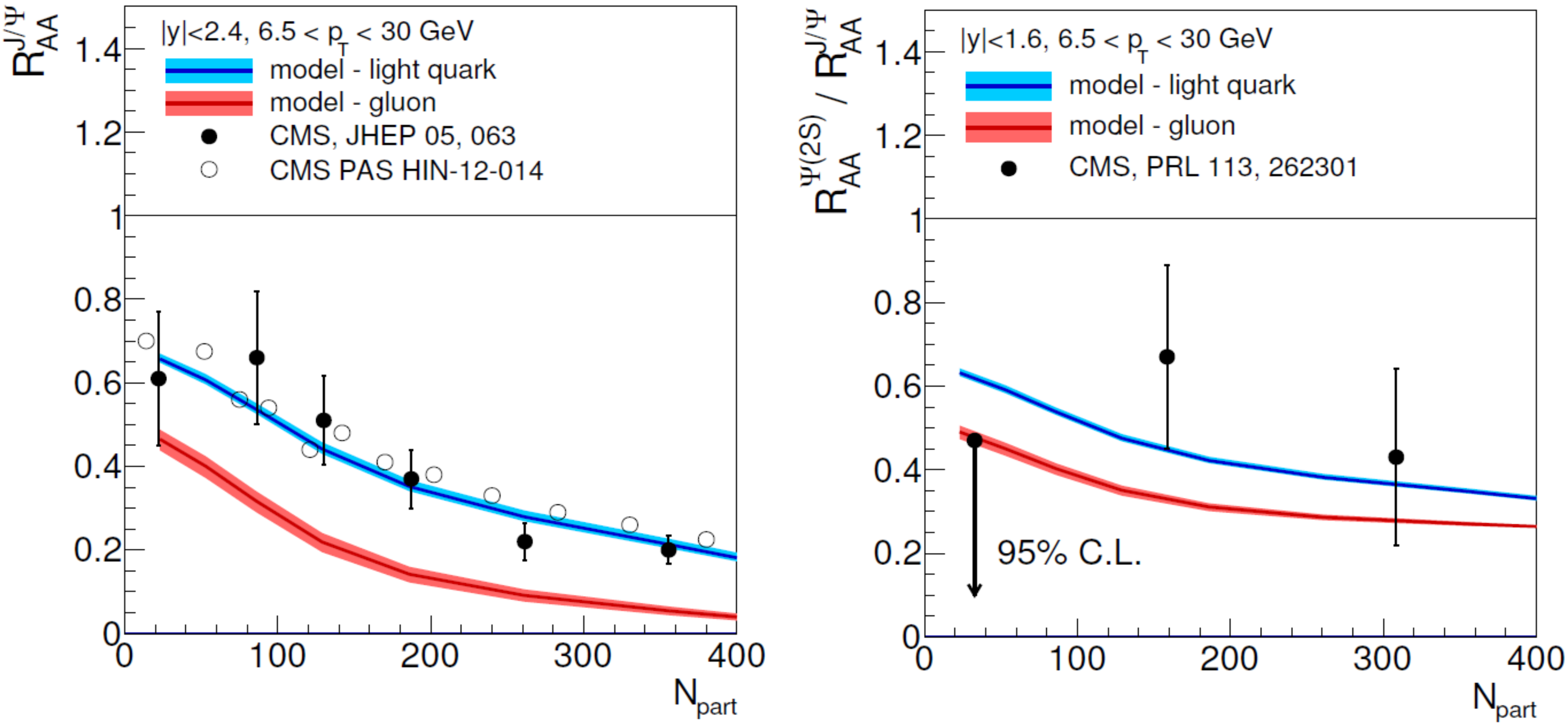
– Energy loss extracted from jets

Charmonia



$R_{AA}^{J/\Psi}$
|y|<2.4
model - light quark
model - gluon
CMS, JHEP 05, 063
CMS PAS HIN-12-014
p_T [GeV]
6.5 < p_T < 30 GeV
|y|

Charmonia



… suppression of both charmonia at p_T>6.5 GeV is similar to the suppression of light quark jets

Summary I.

- Flavor dependence of the jet quenching drives quite a lot of what we see in the data.
- Coherence effects are seen in the data.
- Recoil (or in-cone radiation) modifies also the measured high-z fragmentation (same holds for all other observables).
- *b*-jets are quenched by 1.5±0.4 more than light quark jets.
- Dijet asymmetry: peaking can emerge in the situation when quark jet suppression is more non-linear in path-length than the gluon jet suppression.
- Same magnitude of jet / charged particle R_{AA} between 2.76 TeV and 5 TeV implies slightly larger suppression at 5 TeV (natural $\sqrt{s_{NN}}$-dependence of the energy loss).

Summary II.

- Average jet quenching can be quantified from the data as follows:

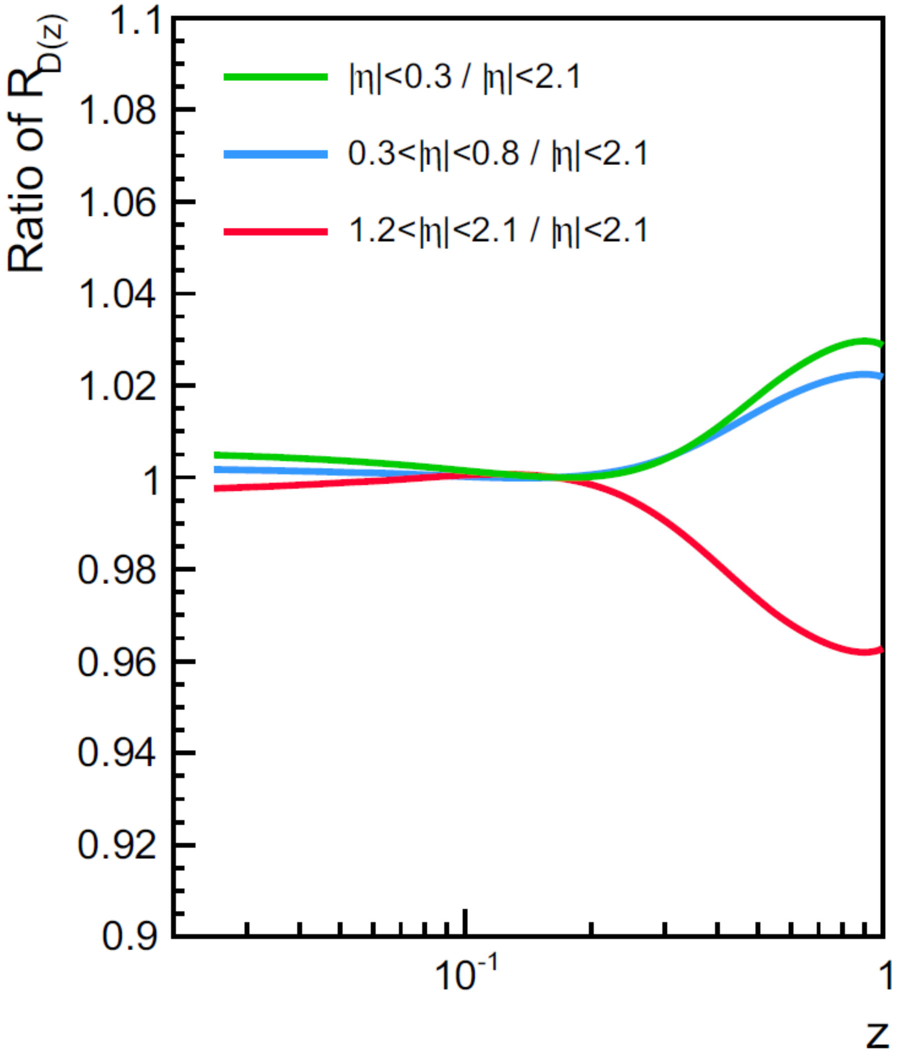
$s = x \cdot N_{\text{part}} + y$	$x = 12.3 \pm 1.4$ GeV, $y = 1.5 \pm 0.2$ GeV
α	0.52 ± 0.02
c_{F}	1.78 ± 0.12

$$S_q = s' \left(\frac{p_T^{\text{jet}}}{p_{T,0}} \right)^{\alpha}$$

- Color factor, extracted for the first time in HI, seems not to be modified by the medium ($c_F = 1.78 \pm 0.12$).
- Parton energy loss does not extrapolate to 0 for $N_{part} \rightarrow 0$.
- Suppression of charmonia at p_T>6.5 GeV at midrapidity behaves like the suppression of light quark jets.

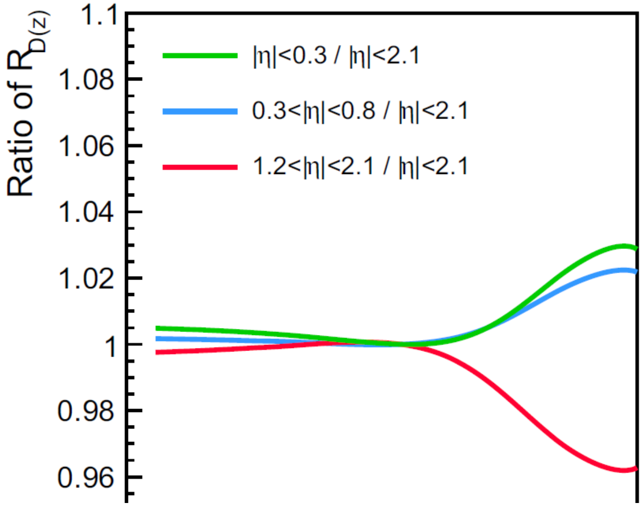
Unused slides, technical details

Modifications of fragmentation functions – prediction



… central rapidity – higher yields at high-z (but not by much)

Modifications of fragmentation functions – prediction



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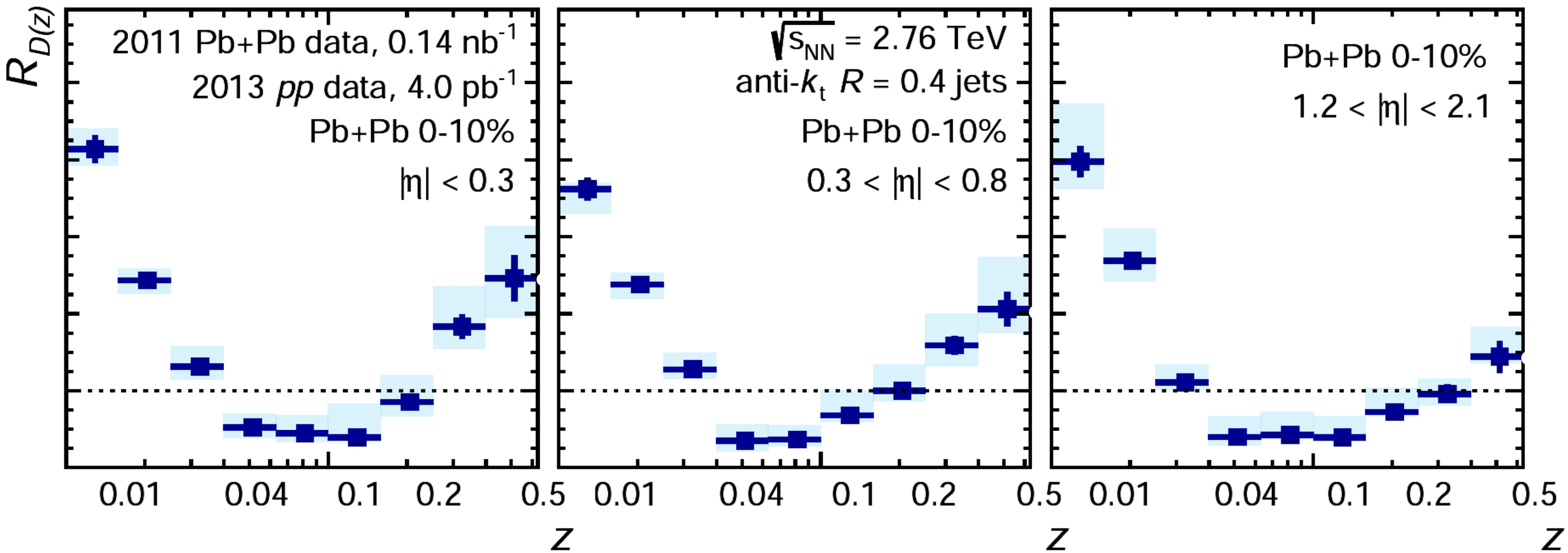

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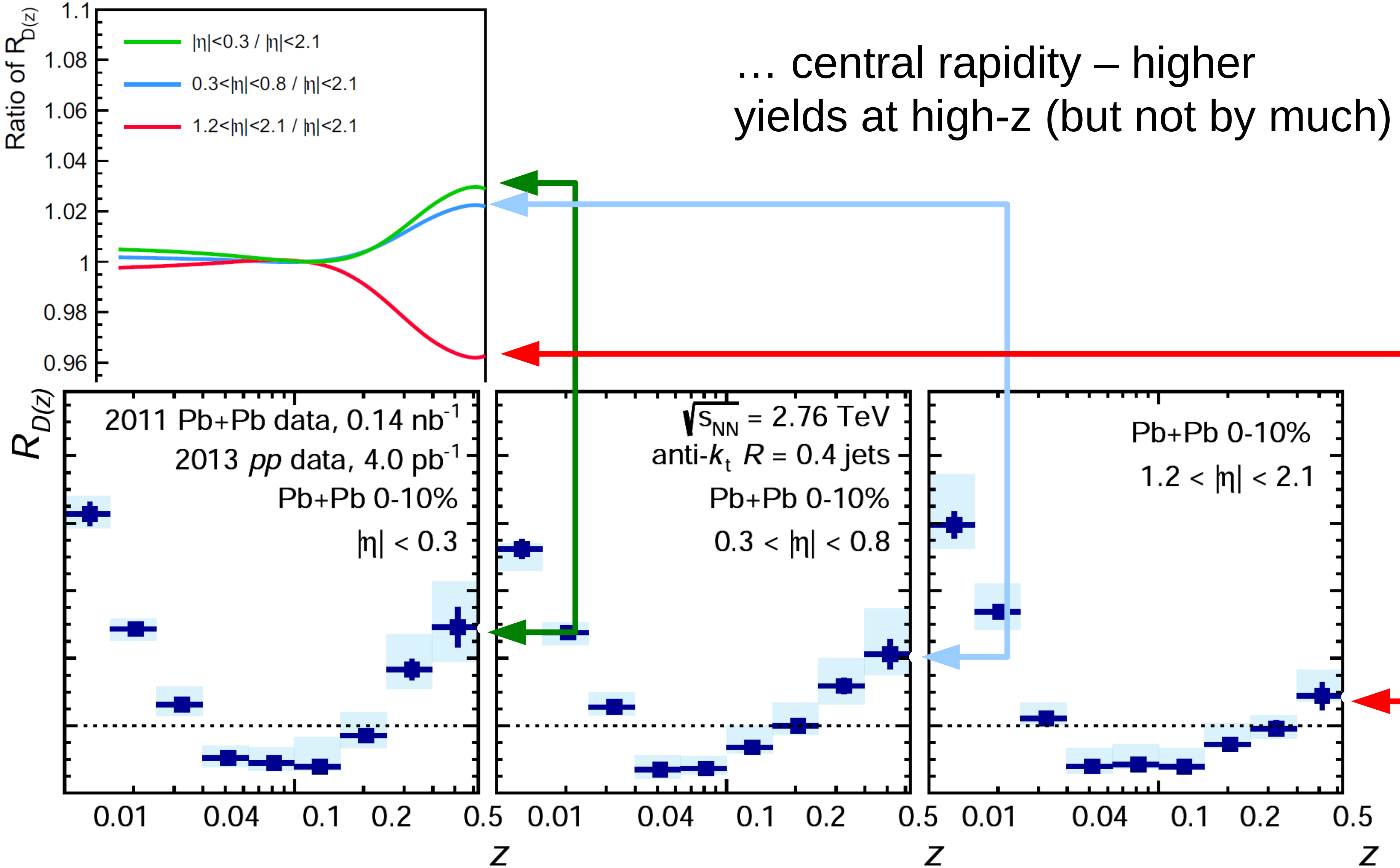
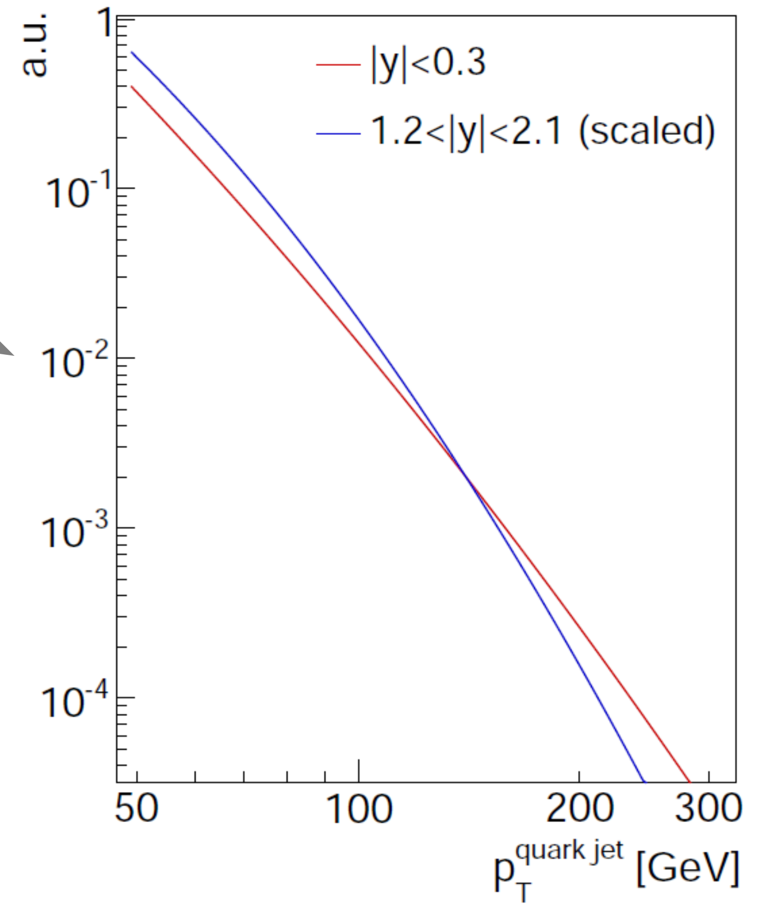


Start: Two basic questions

- Why do have the jet and charge particle R_{AA} almost **no rapidity dependence** given quite different input parton spectra and flavor composition at different rapidities?
- What is responsible for the **enhancement** (= not suppression) at high *z* seen in the fragmentation?



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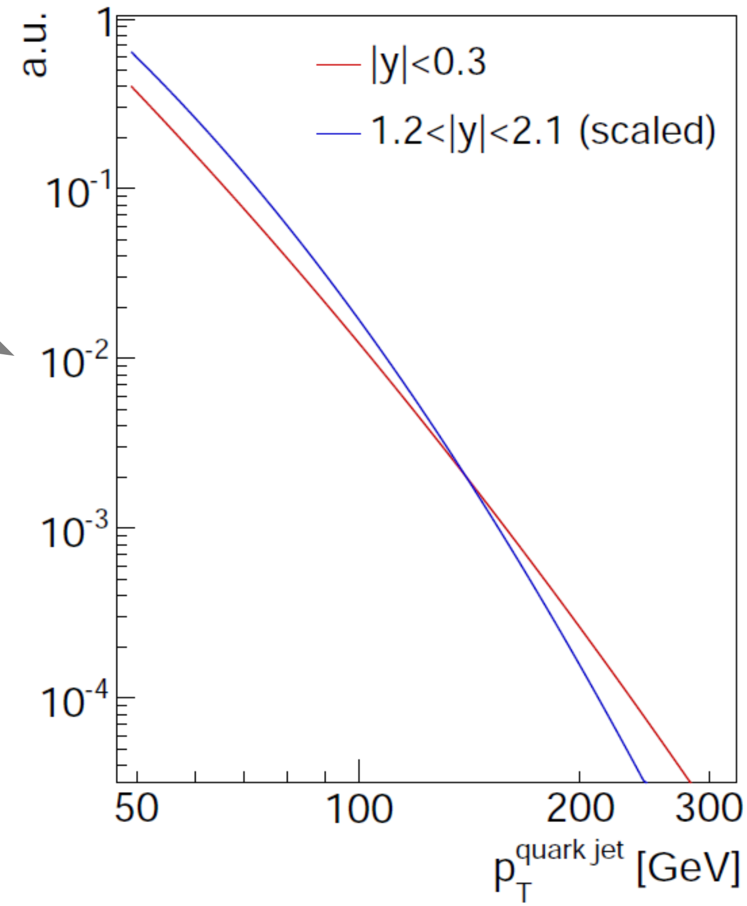
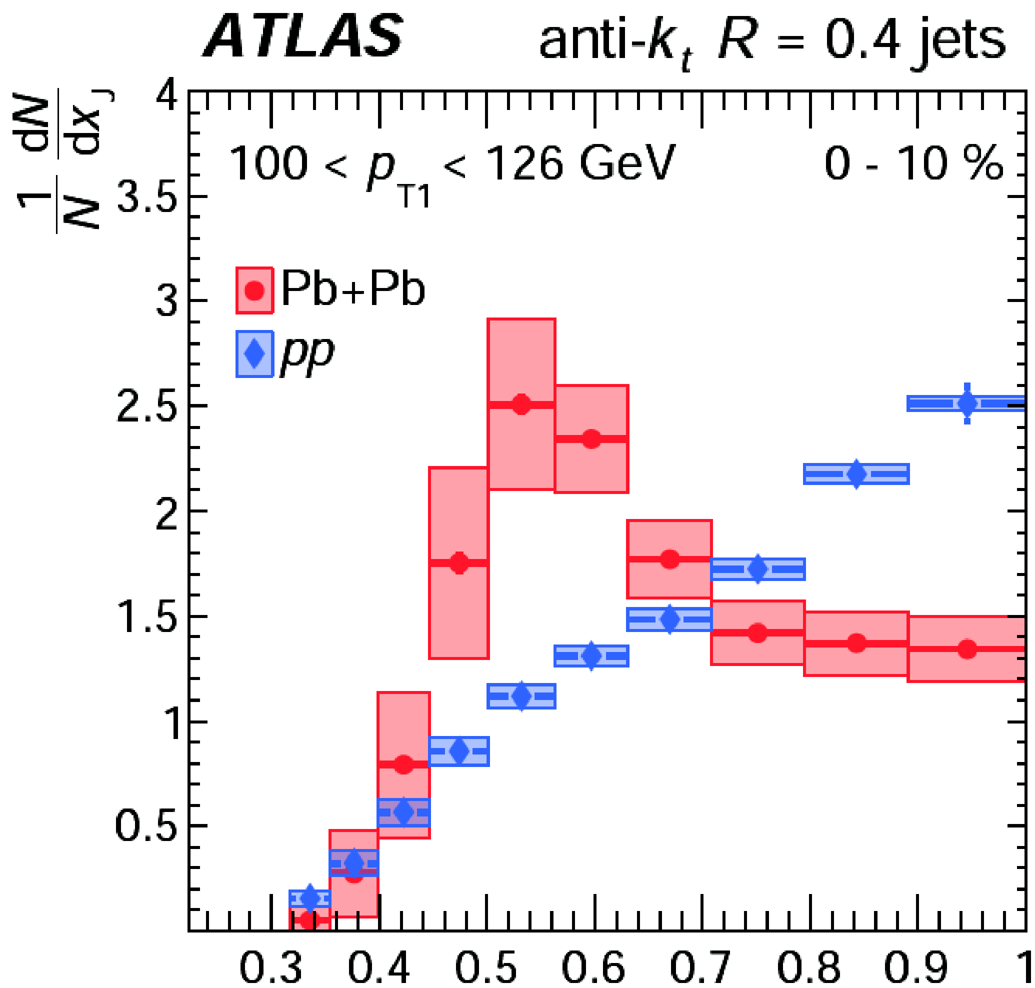


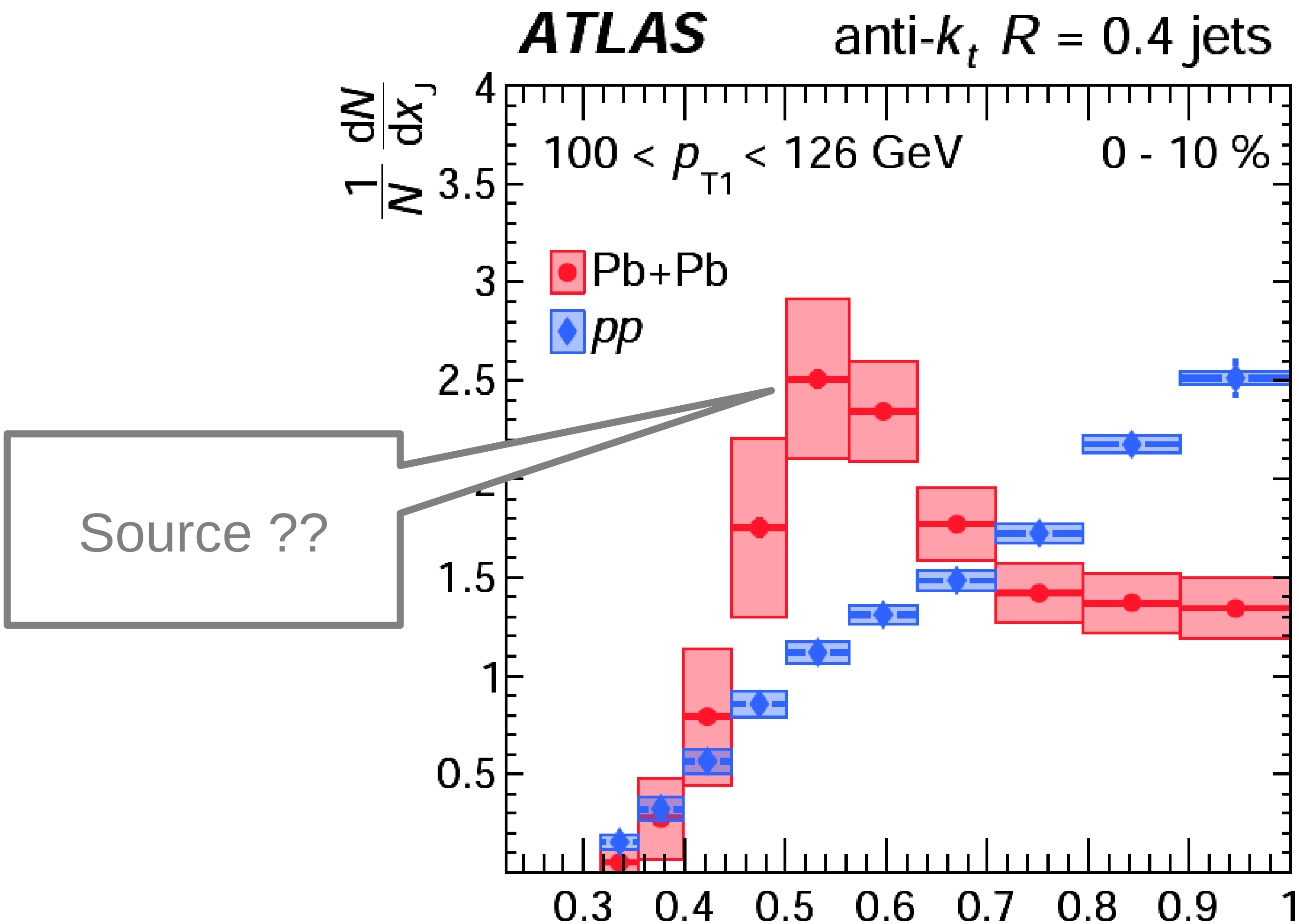
–> Use a simple model with minimal assumptions on the quenching physics to extract basic properties of the jet quenching

Dijet asymmetry



$$x_J = \frac{p_{T,\mathrm{subleading}}}{p_{T,\mathrm{leading}}}$$

Dijet asymmetry



$$x_J = \frac{p_{T,\mathrm{subleading}}}{p_{T,\mathrm{leading}}}$$

Dijet asymmetry

- Test the role of path-length dependence

$$S(p_{T,\mathrm{ini}}, l) =$$

$$= \frac{c_{\mathrm{F}} s}{\langle l \rangle} \left(\frac{p_{T,\mathrm{ini}}}{p_{T,0}} \right)^{\alpha} f(l)$$

$$l^k$$

$$k = 0.5, 1, 2, 3$$

$$l = \int \mathrm{d}\tau \, \tau \rho(\vec{r})$$

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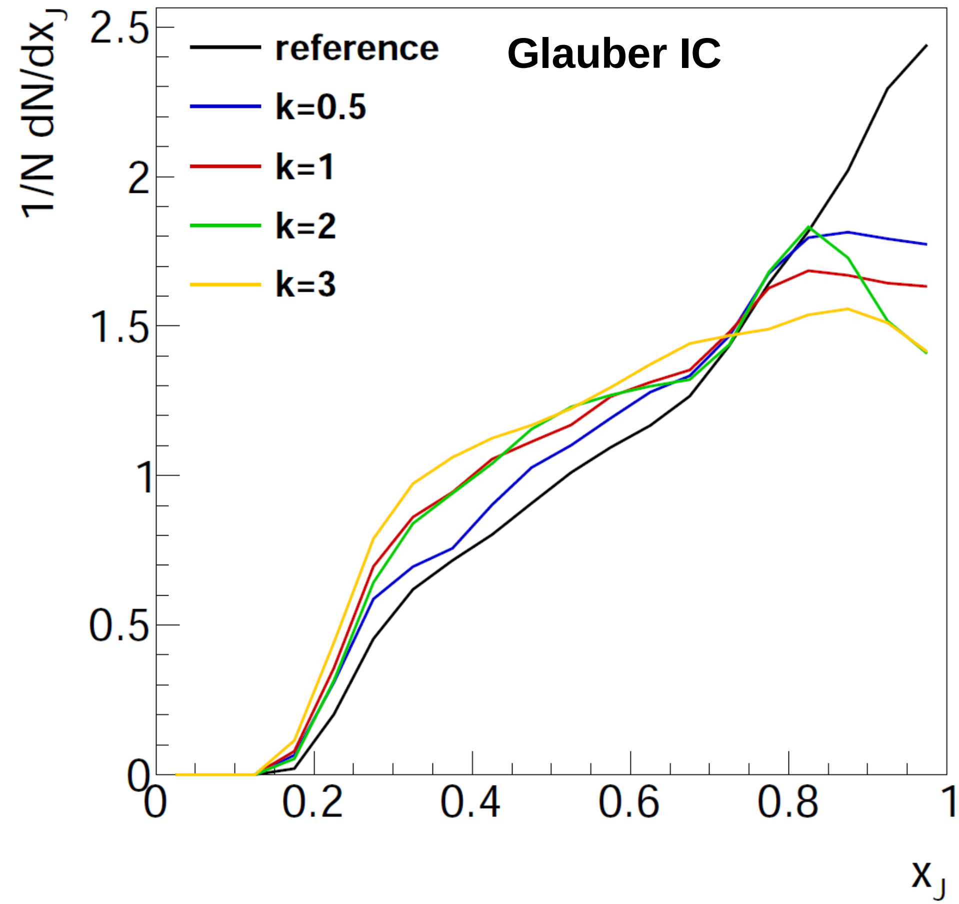

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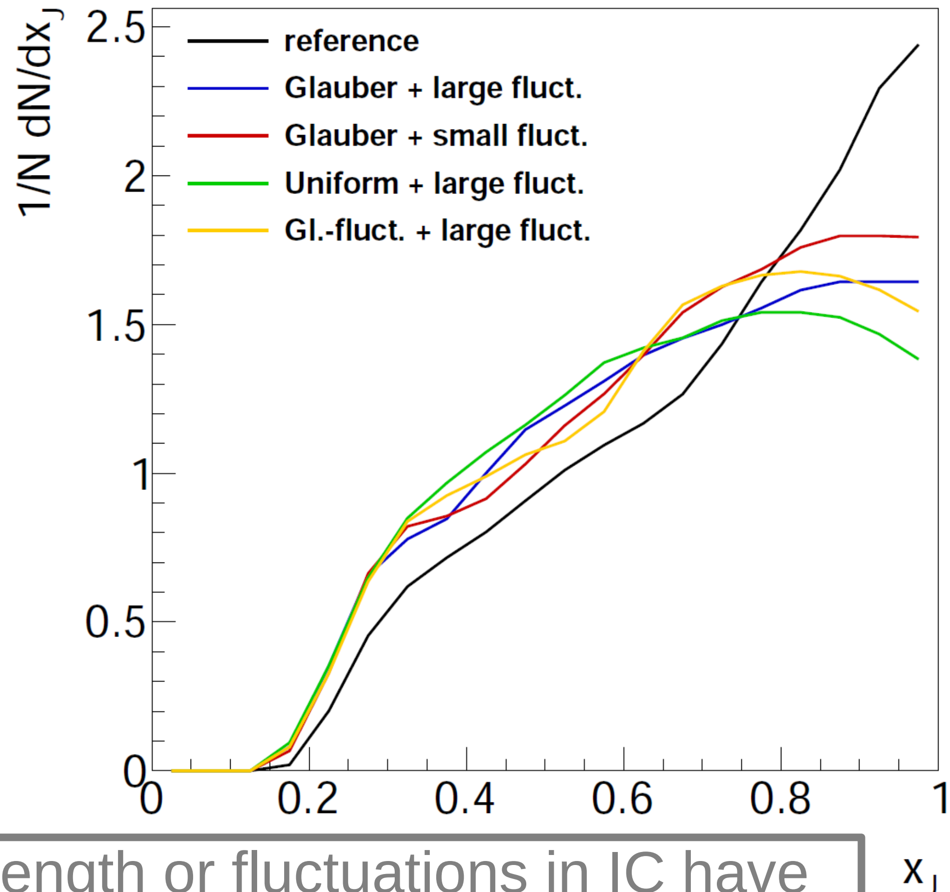
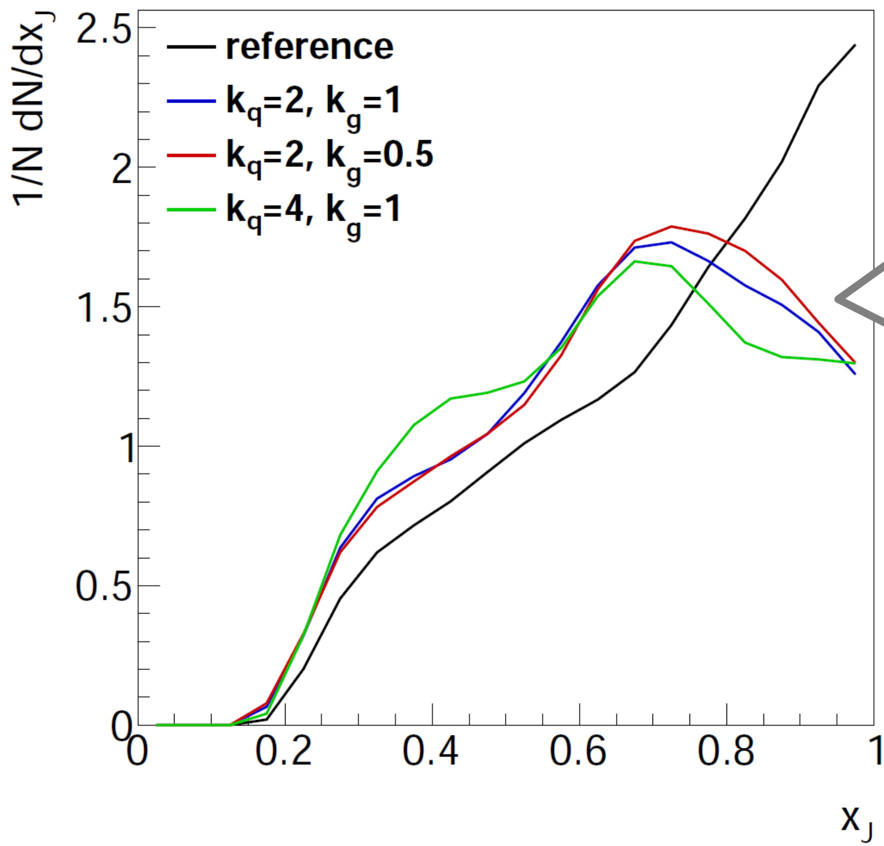


Path-length or fluctuations in IC have no major impact (similar conclusions in EPJC 76 (2016) no.5, 288)

Dijet asymmetry

- Test the role of flavor

$$S(p_{T,\mathrm{ini}}, l) = \frac{c_\mathrm{F}\, s}{\langle l^{k(c_\mathrm{F})}\rangle} \left(\frac{p_{T,\mathrm{ini}}}{p_{T,0}}\right)^{\alpha} f(l^{k(c_\mathrm{F})})$$

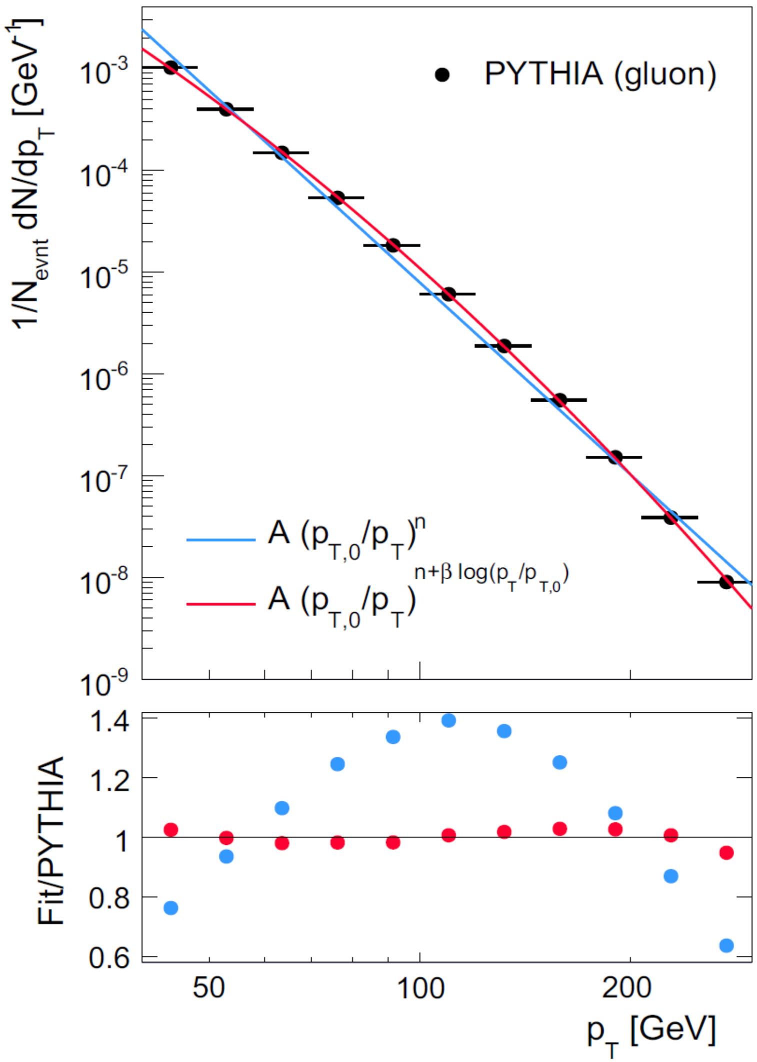


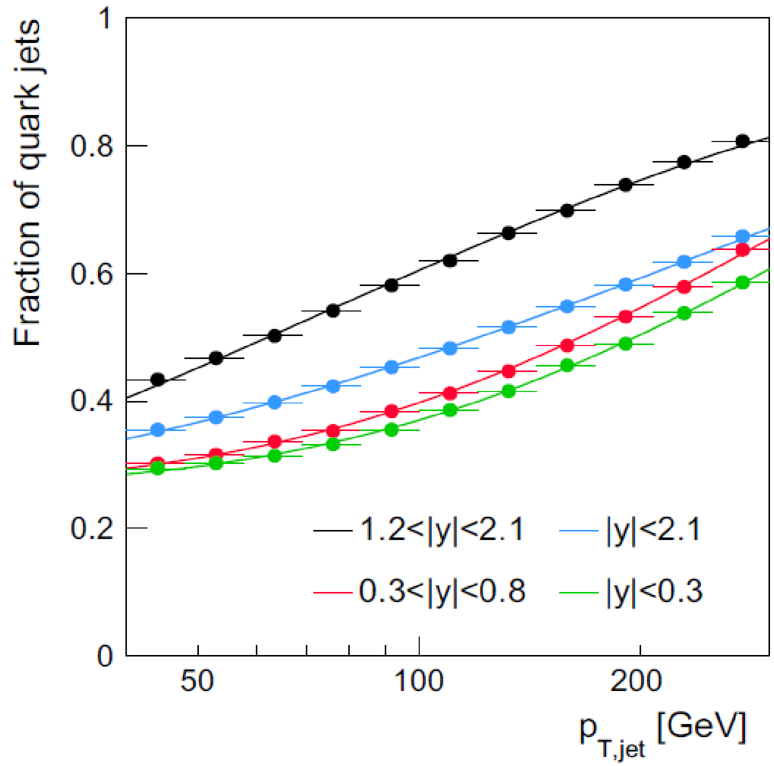
Peaking in the configurations when the loss of quark jets is more non-linear than the loss of gluon jets

… contra-intuitive

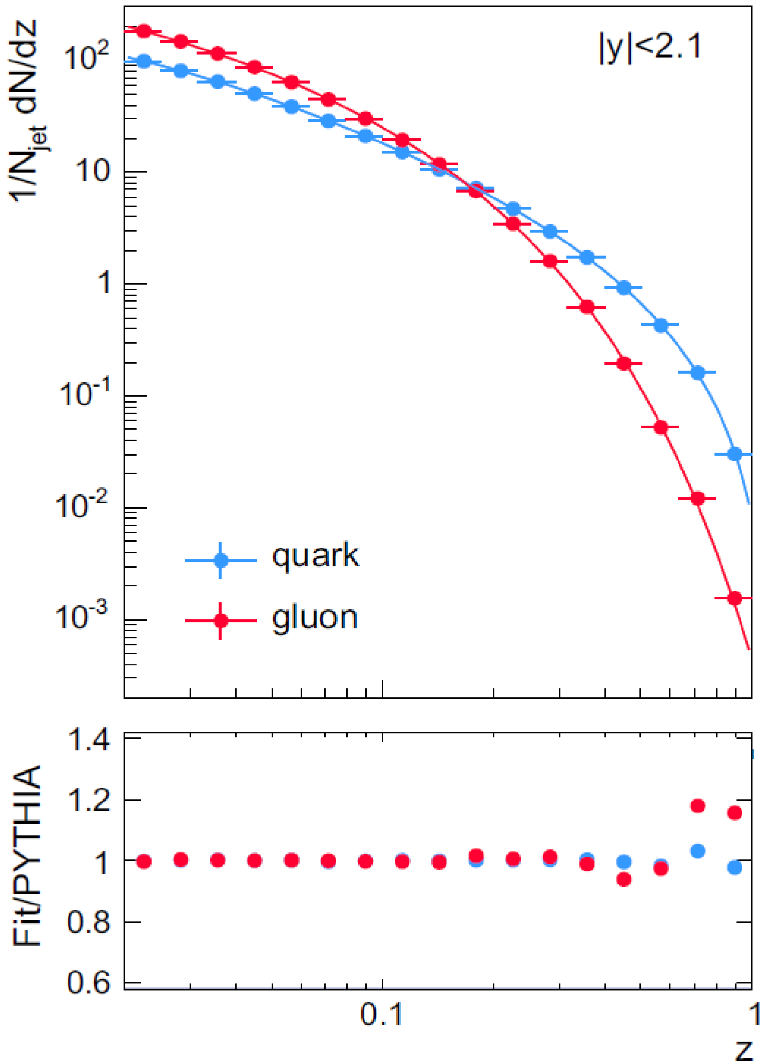
Flavor fractions and fit parameters

Fit type	Parameter	$\|y\| < 2.1$	$\|y\| < 0.3$	$0.3 < \|y\| < 0.8$	$1.2 < \|y\| < 2.1$
All	f_{q0}	0.34	0.28	0.29	0.40
Power law	n_q	5.66	5.37	5.40	6.15
	n_g	6.25	5.97	6.09	6.92
Extended power law	n_q	4.19	4.34	4.27	3.75
	β_q	0.71	0.49	0.54	1.2
	n_g	4.69	4.55	4.57	4.60
	β_g	0.80	0.71	0.76	1.2





D(z) parameterization



$$D(z) = a \cdot \frac{(1+dz)^b}{(1+ez)^c} \cdot \exp(-fz)$$

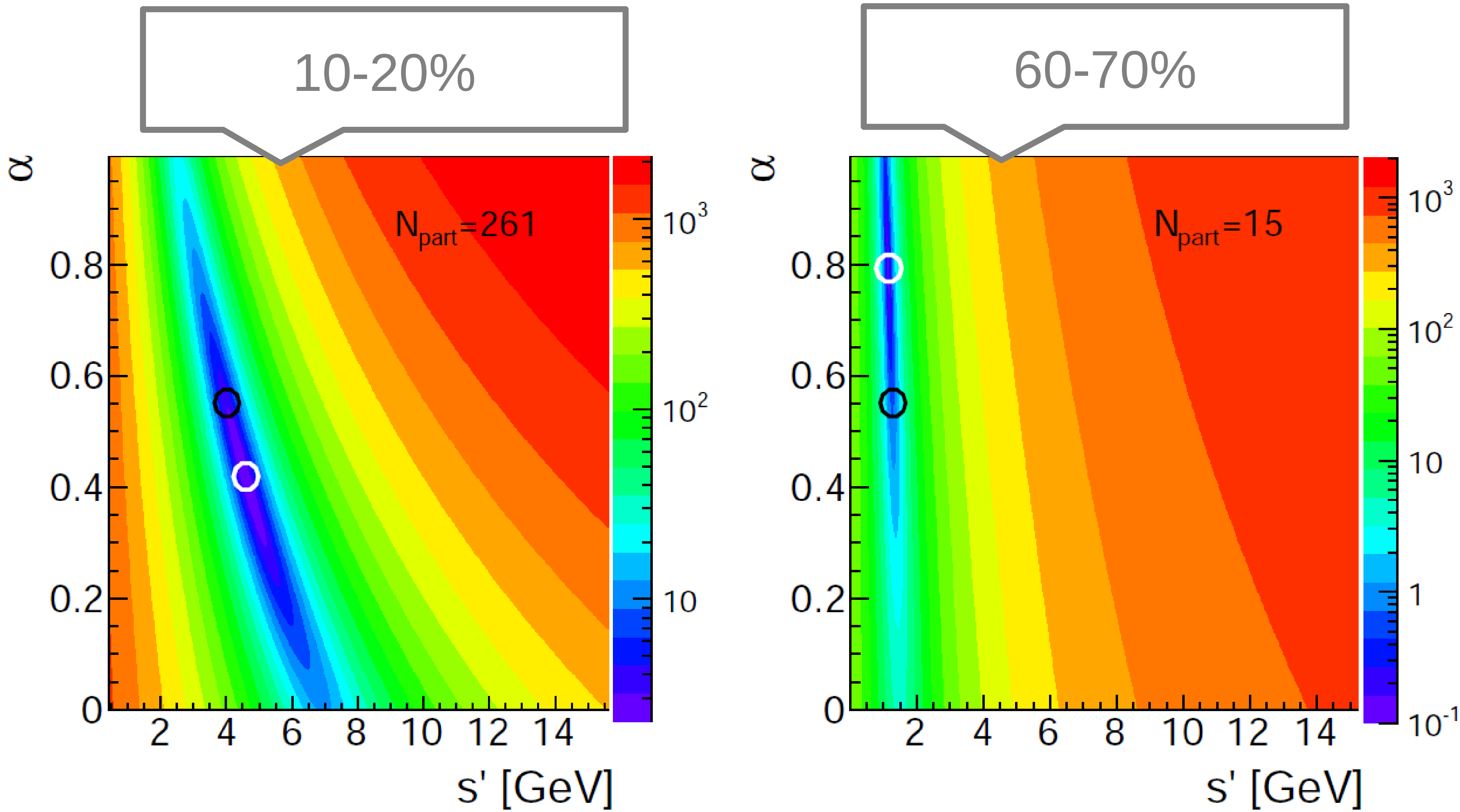
	a	b	c	d	e	f
Quark	318	2.51	1.44	−0.85	52.4	0
Gluon	574	1.87	2.32	9.09	32.0	10.3

R_{AA} – full analytic expression

$$
\begin{aligned}
R_{\mathrm{AA}} = f_q & \left(\frac{1}{1+S_q/p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n_q+\beta_q \log((p_{\mathrm{T}}^{\mathrm{jet}}+S_q)/p_{\mathrm{T}0})} \\
& \times \left(\frac{p_{\mathrm{T}0}}{p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{\beta_q \log(1+S_q/p_{\mathrm{T}}^{\mathrm{jet}})} \left(1+\frac{dS_q}{dp_{\mathrm{T}}^{\mathrm{jet}}}\right) \\
& + (1-f_q)\left(\frac{1}{1+S_g/p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n_g \beta_g \log((p_{\mathrm{T}}^{\mathrm{jet}}+S_g)/p_{\mathrm{T}0})} \\
& \times \left(\frac{p_{\mathrm{T}0}}{p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{\beta_g \log(1+S_g/p_{\mathrm{T}}^{\mathrm{jet}})} \left(1+\frac{dS_g}{dp_{\mathrm{T}}^{\mathrm{jet}}}\right),
\end{aligned}
$$

$$
f_q\left(p_{\mathrm{T}}^{\mathrm{jet}}\right) = \frac{1}{1+\left(\frac{1-f_{q0}}{f_{q0}}\right)\left(\frac{p_{\mathrm{T}0}}{p_{\mathrm{T}}^{\mathrm{jet}}}\right)^{n_g-n_q+(\beta_g-\beta_q)\log\left(p_{\mathrm{T}}^{\mathrm{jet}}/p_{\mathrm{T}0}\right)}}.
$$

Minimization in (I.)



Modifications of fragmentation functions – a detail

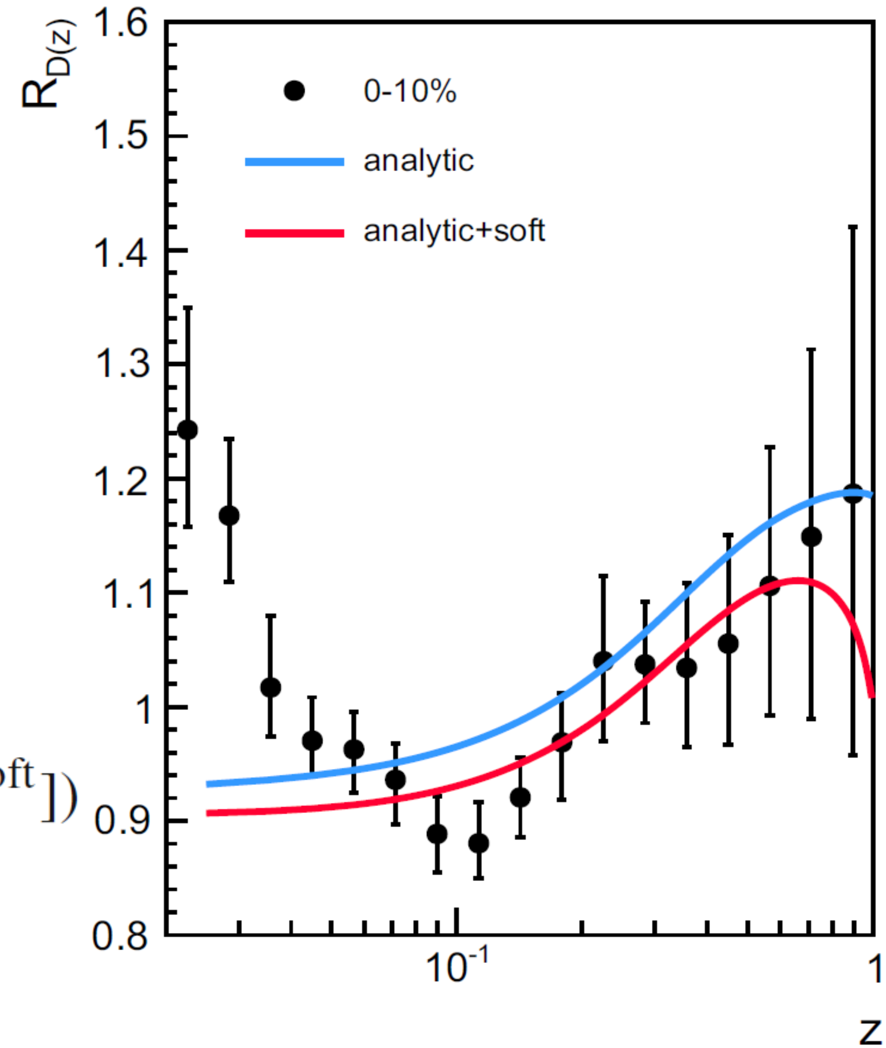
How is the soft excess estimated:

Measured
(at least partially)

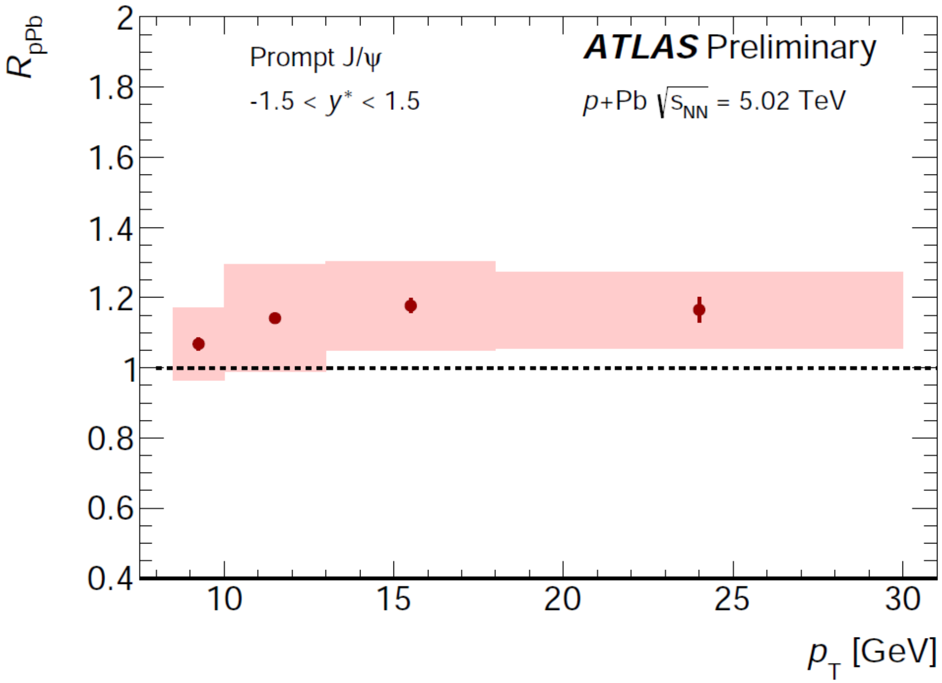
$$\Phi_{\rm inc}^{\rm soft} = f_q^{\rm int}\Phi_q^{\rm soft} + (1 - f_q^{\rm int})\Phi_{\rm g}^{\rm soft}$$

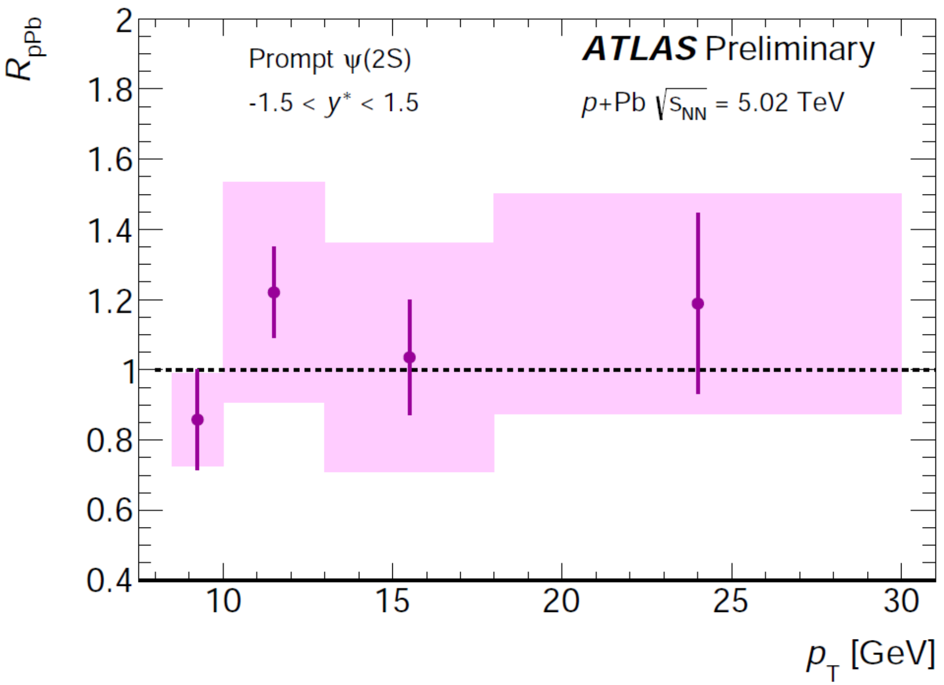
$$\Phi_{\rm g}^{\rm soft} = c_F \Phi_q^{\rm soft}$$

$$D^{\rm meas}(z) = f_q^{\rm int} D_q(z[1 + \Phi_q^{\rm soft}]) + (1 - f_q^{\rm int}) D_g(z[1 + \Phi_{\rm g}^{\rm soft}])$$

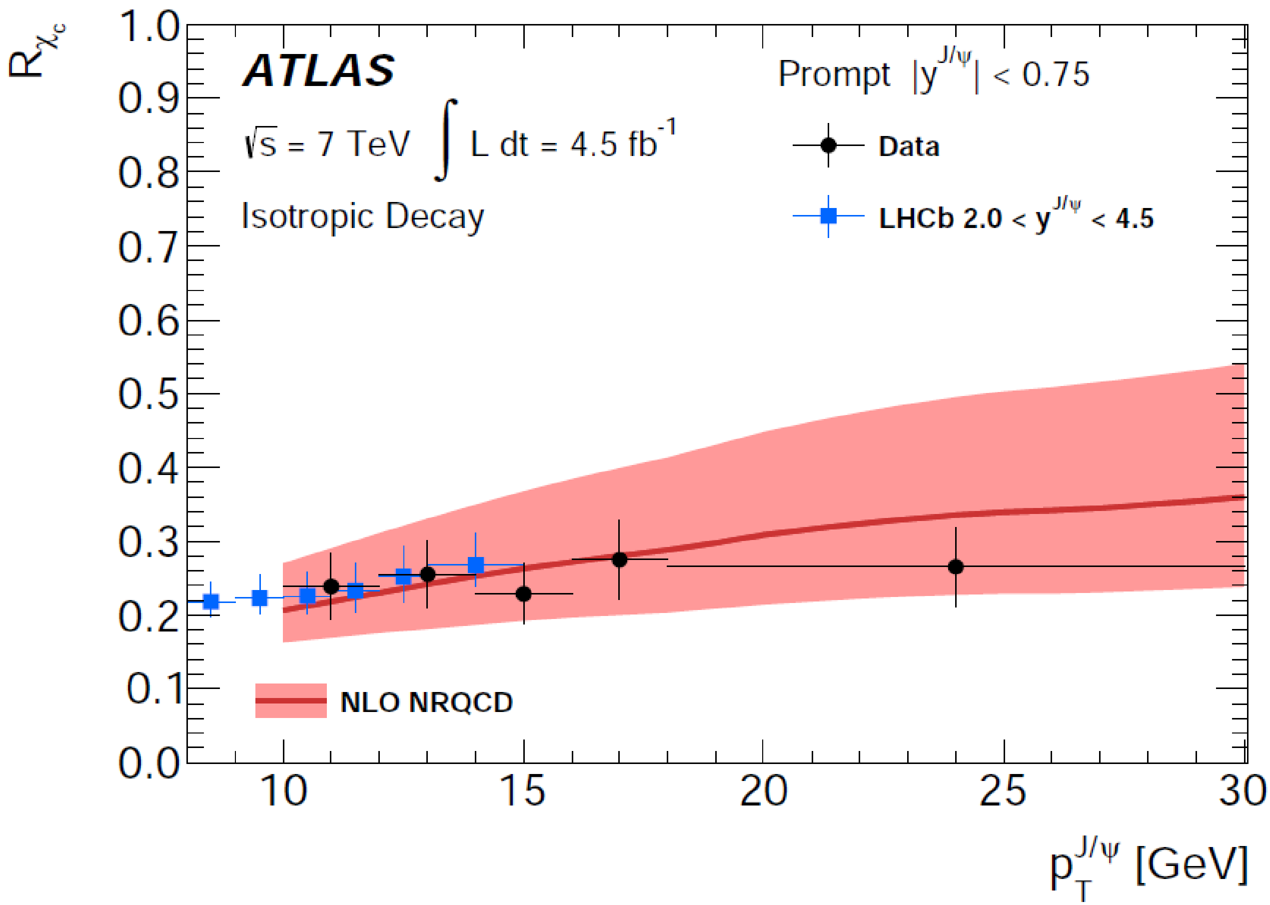


Charmonia in p+Pb



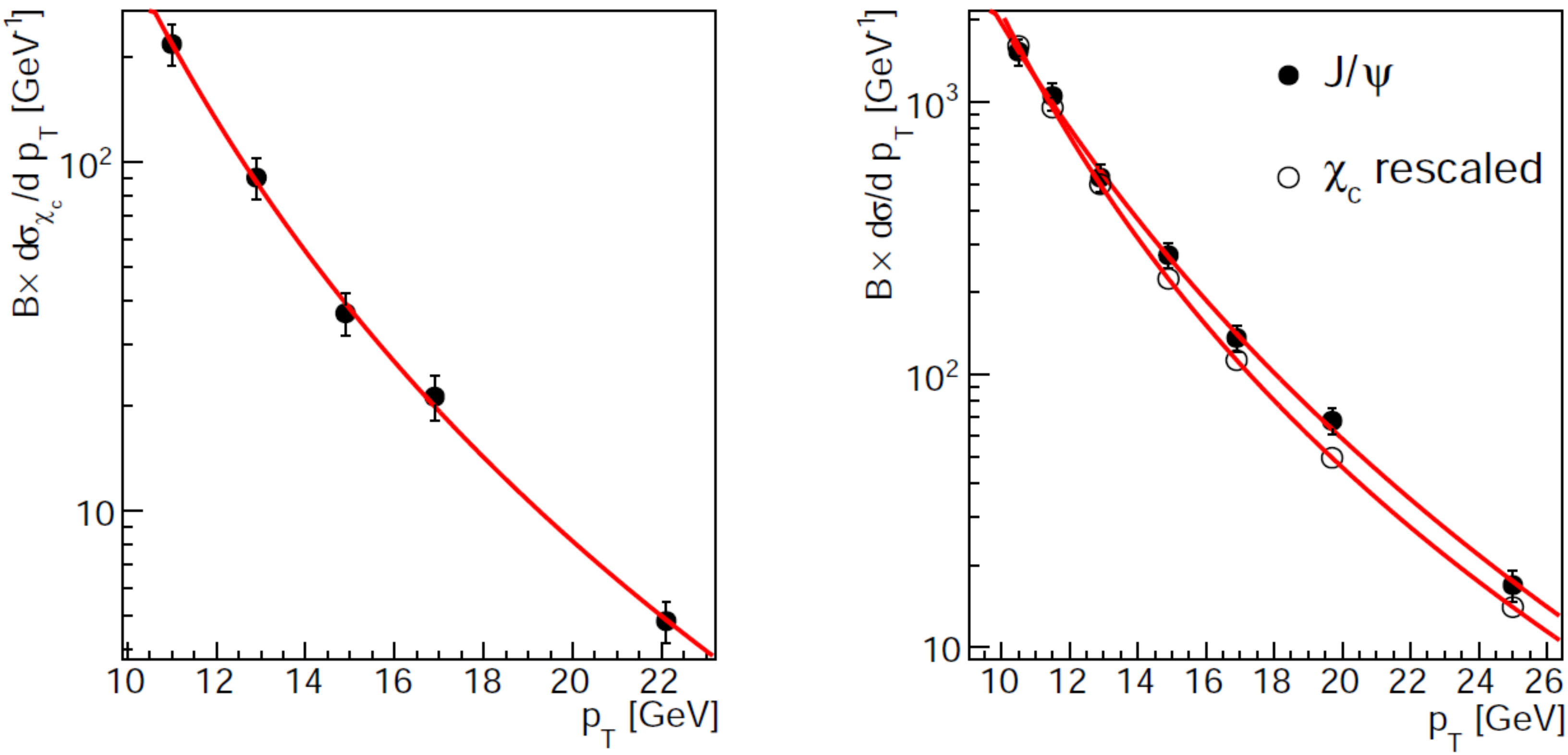


Feed down



ATLAS, JHEP 07 (2014) 154

Feed down



ATLAS, JHEP 07 (2014) 154

Introduction

Two paths:

- Be as realistic as one can:
  - MC generators
    - JEWEL (jewel.hepforge.org)
    - QPythia (igfae.usc.es/qatmc)
    - Martini (PRC 80 (2009) 054913)
  - JETSCAPE Collaboration
  - theory …tions of parton enegy loss

The Jet Energy-loss Tomography with a Statistically and Computationally Advanced Program Envelope (goal provide modular software which includes: modeling of initial state + dynamical evolution of QGP + jet energy loss + advanced statistical tools; http://jetscape.wayne.edu/)

:

  - parametric modeling of parton energy loss

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