



Latest results from electron beam simulations for the HEL

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Outline

- Target-setting for HEL parameters
- Stability analysis
 1. Critical current
 2. Diocotron instability consideration
- Particle tracking and its analysis
- Conclusion

Target-setting for HEL parameters

The ideal hollow electron beam doesn't induce any electric or magnetic fields inside, but generates strong nonlinear fields outside. The transverse kick applied to hollow particles can be expressed as:

$$\theta = \frac{2I_{er}L(1 \pm \beta_e\beta_p)}{r\beta_e\beta_p c^2 (B\rho)_p} \left(\frac{1}{4\pi\epsilon_0} \right)$$

The whole problem, therefore, is to choose all parameters to maximize the efficiency of collimation and at the same time try to avoid different undesirable things which leading to any fields appearing inside hollow.

Stability analysis



HEL Current restrictions

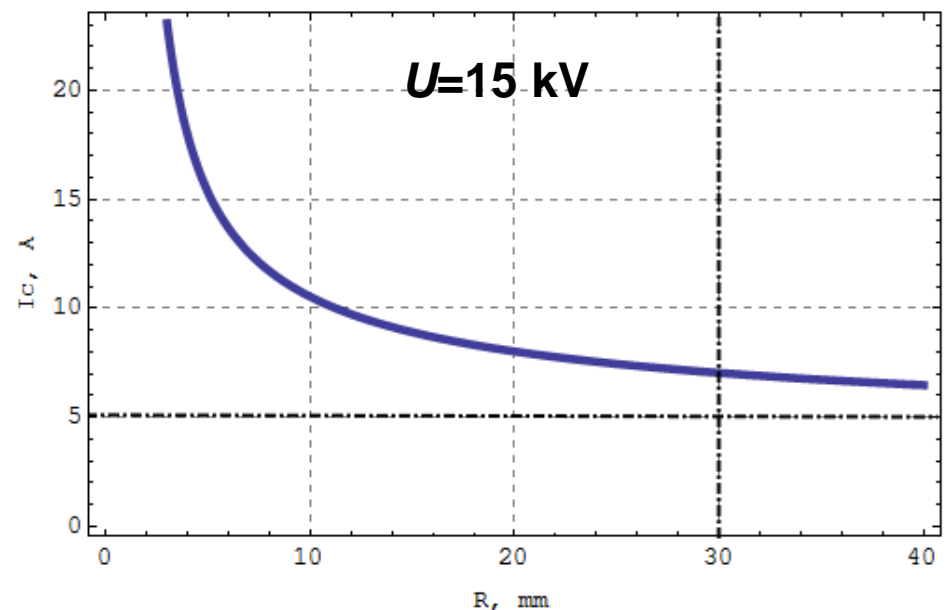
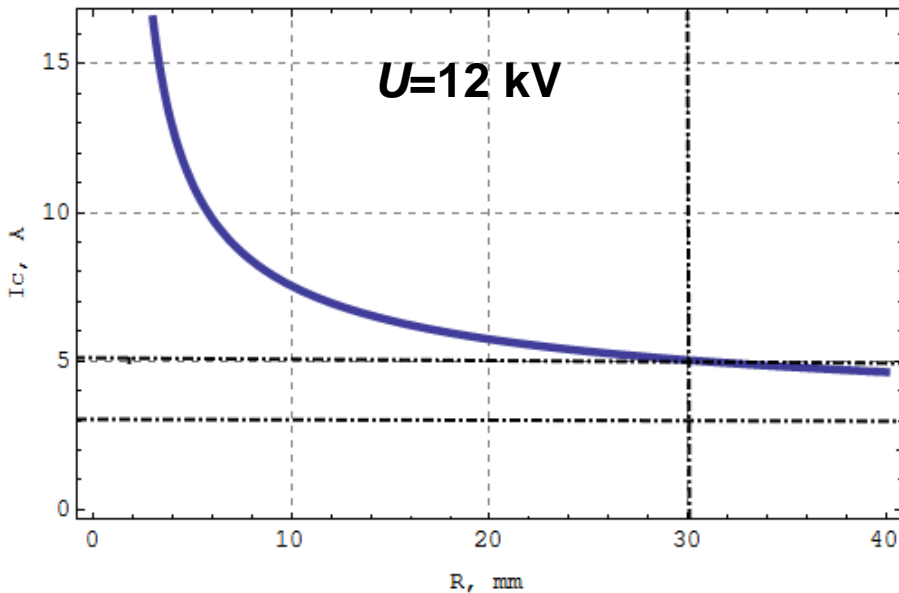
First restriction - gun design and cathode type/size (not a big problem)

Second restriction - Maximum current that can be transported through the vacuum chamber with radius R . In case current exceeds critical value, part of the beam is reflected, i.e. beam potential sags to zero.

$$I_{max} = \frac{4 \cdot \pi \cdot \epsilon_0 \cdot \sqrt{8}}{\sqrt{27}} \cdot \sqrt{\frac{e}{m}} \cdot \frac{U^{3/2}}{1 + 2 \ln\left(\frac{R}{r}\right)}$$

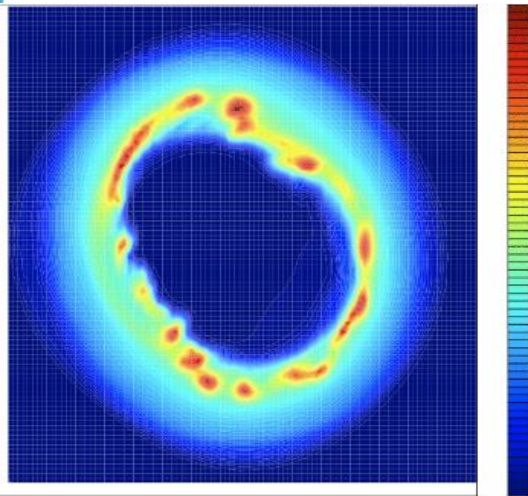
Pipe radius $R = 30$ mm

Outer beam radius $r = 1.8$ mm



Diocotron instability consideration

On Fermilab test stand the destruction of hollow beam into specific microstructure was observed (very similar to diocotron instability):

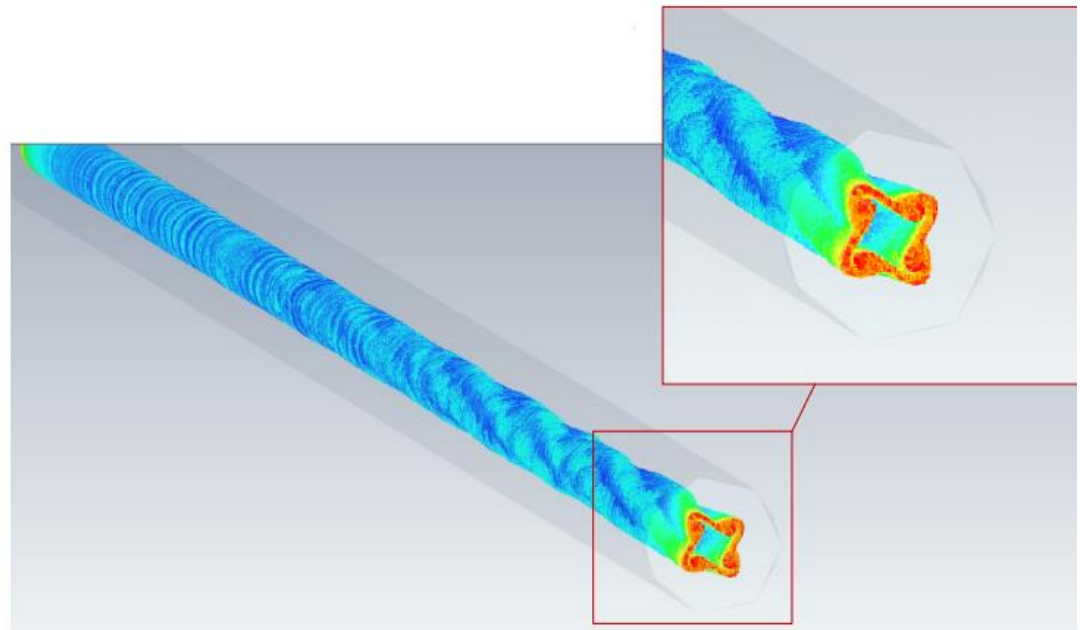


Origin of the diocotron instability: Different angular velocities for different radii provide relative motion of layers (**See (*)**). Small initial asymmetry may lead to the significant density equilibrium violation and cluster origin.

$$\omega_e(r) = \frac{\omega_p}{\omega_c} \left(1 - \frac{r_{out}^2}{r^2}\right)$$

$$\frac{\partial \omega_e(r)}{\partial r} \neq 0 \quad (*)$$

$$r_{in} < r < r_{out}$$



Davidson stability criteria*

$$\left\{ -l \left[1 - \left(\frac{r_1}{r_2} \right)^2 \right] + 2 - \left[\left(\frac{r_1}{a} \right)^{2l} + \left(\frac{r_2}{a} \right)^{2l} \right] \right\}^2 \geq 4 \left(\frac{r_1}{r_2} \right)^{2l} \left[1 - \left(\frac{r_2}{a} \right)^{2l} \right] \quad l = 1, 2, 3, \dots$$

r_1 – beam inner radius, r_2 – beam outer radius, a – radius of the vacuum chamber

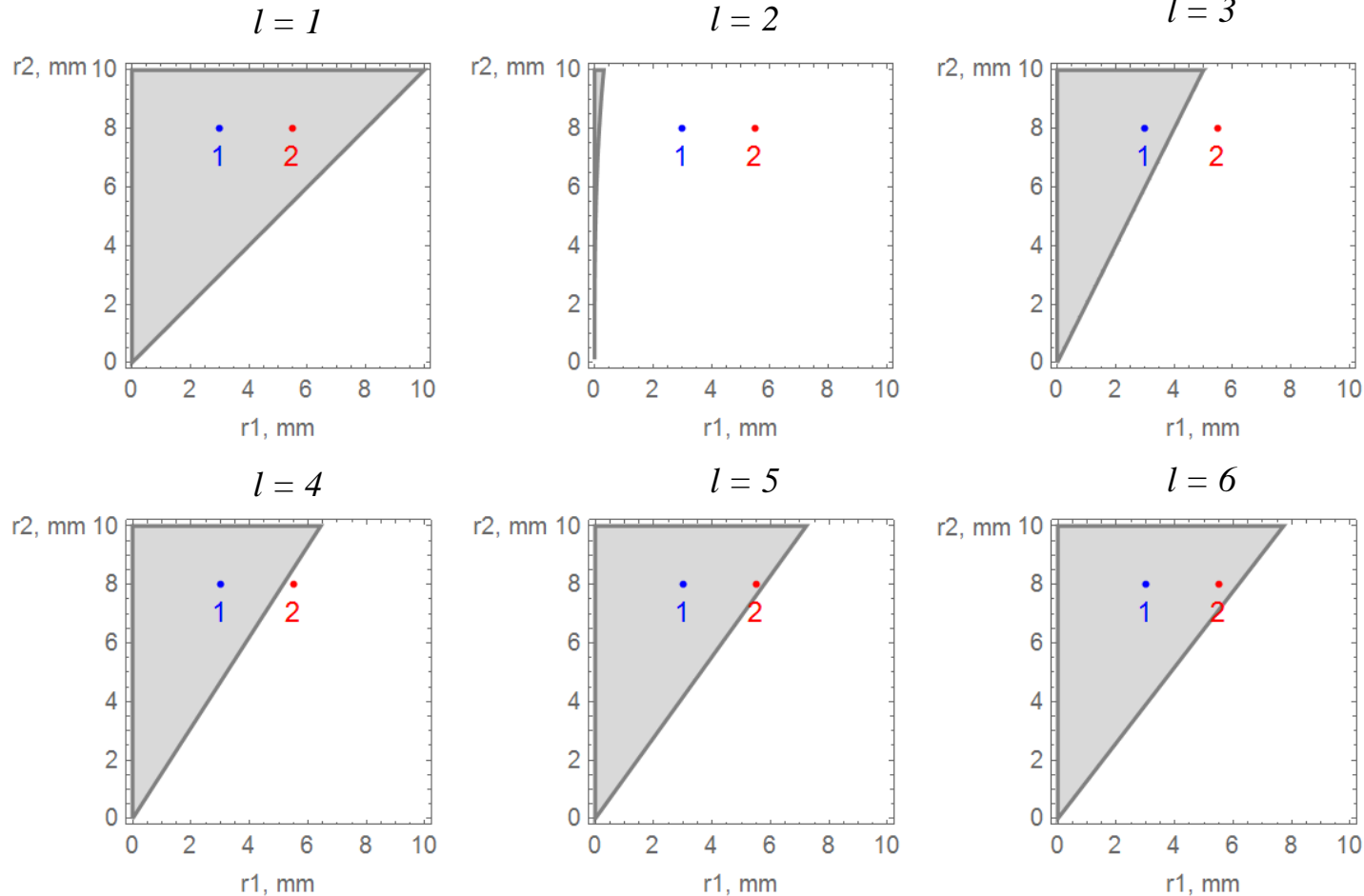
In case the beam is unstable, beam current and external magnetic field influence rate of the instability growth T :

$$T \sim \frac{I_b}{B_z} M$$

I_b is the beam current, B_z is external longitudinal magnetic factor M is the geometry factor (depends on mode number l , r_1 , r_2 , a)

* R. C. Davidson, "Physics of Non-neutral Plasmas"

Stability diagram



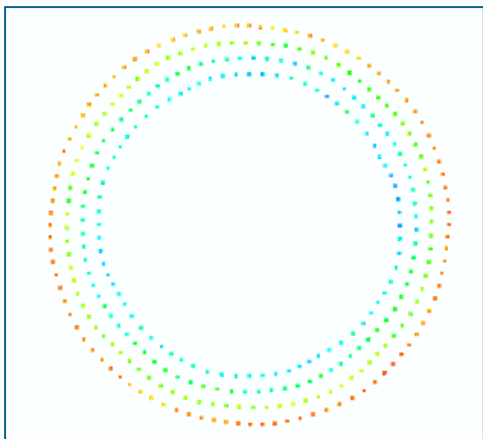
“1” corresponds to the beam with radii $r_1 = 3$ mm, $r_2 = 8$ mm, “2” corresponds to the beam with radii $r_1 = 5.5$ mm, $r_2 = 8$ mm.

Gray region corresponds to the beam stable state (up to the stability criterion)

Checking of the stability criterion

Unstable up to the criterion

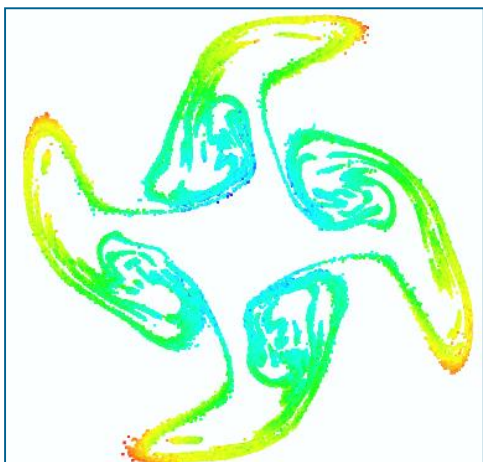
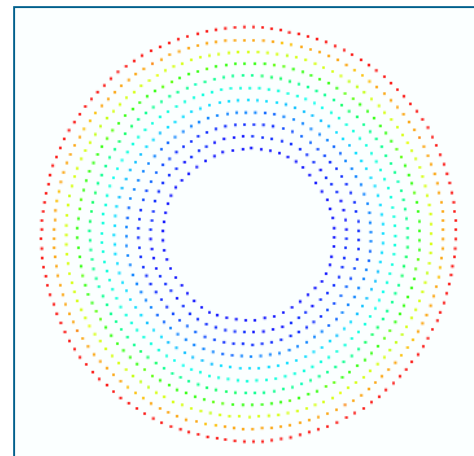
$$r_1 = 5.5 \text{ mm}, r_2 = 8 \text{ mm}$$



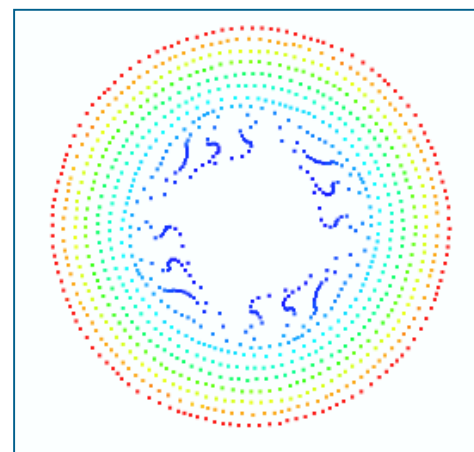
$$\begin{aligned} I &= 10 \text{ A} \\ U &= 15 \text{ keV} \\ B &= 0.2 \text{ T} \end{aligned}$$

Stable up to the criterion

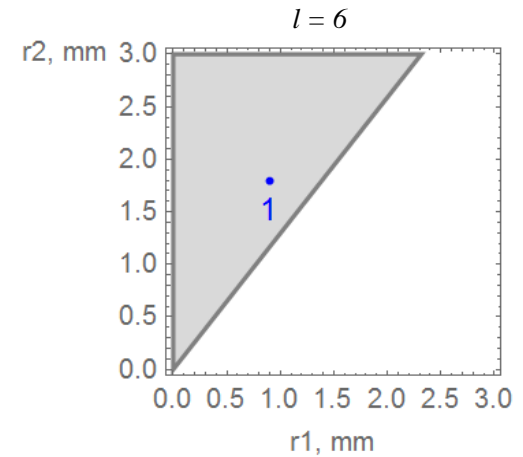
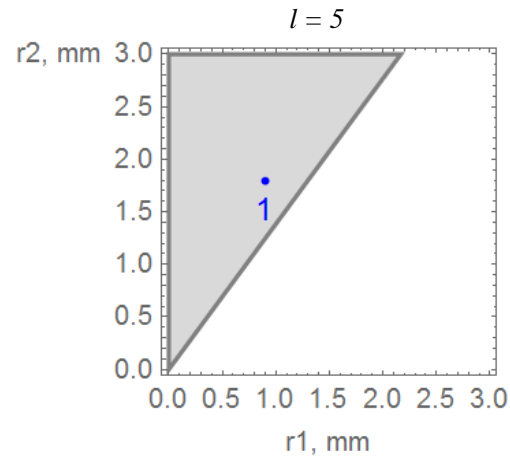
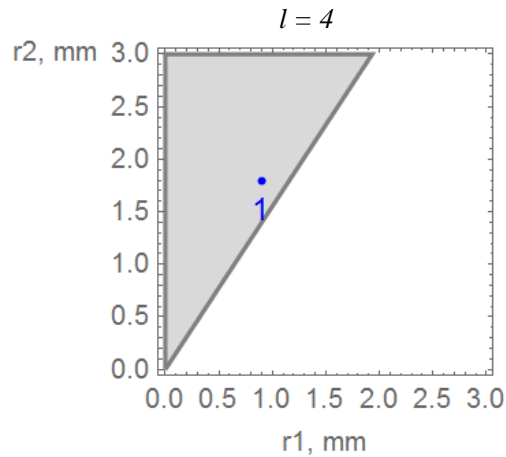
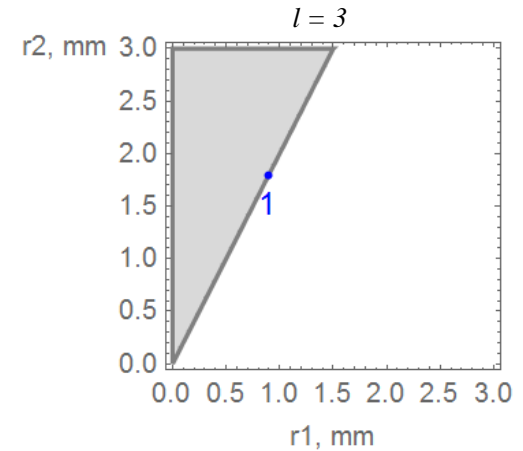
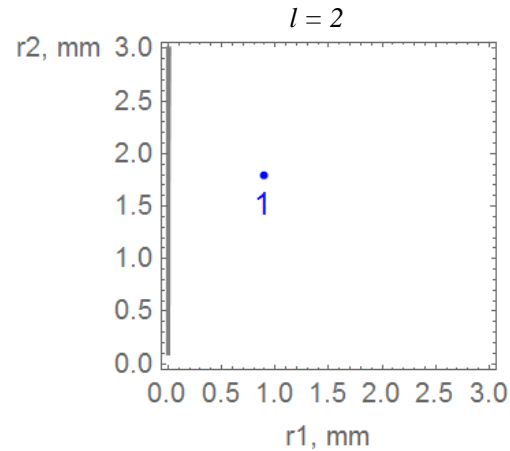
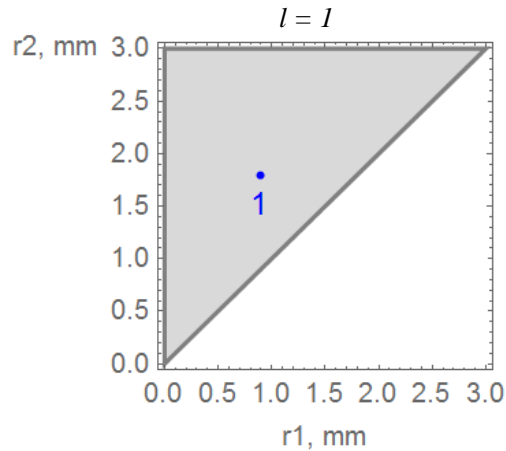
$$r_1 = 3 \text{ mm}, r_2 = 8 \text{ mm}$$



... after 3 m



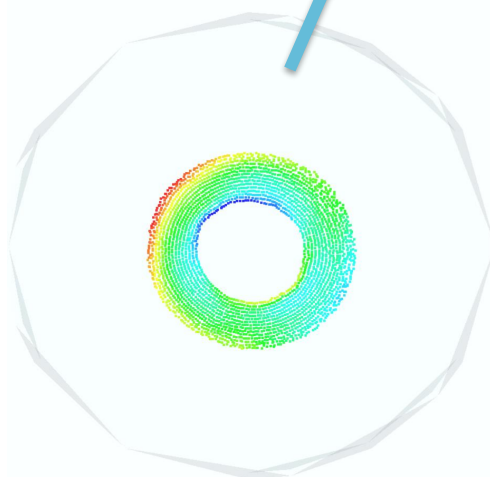
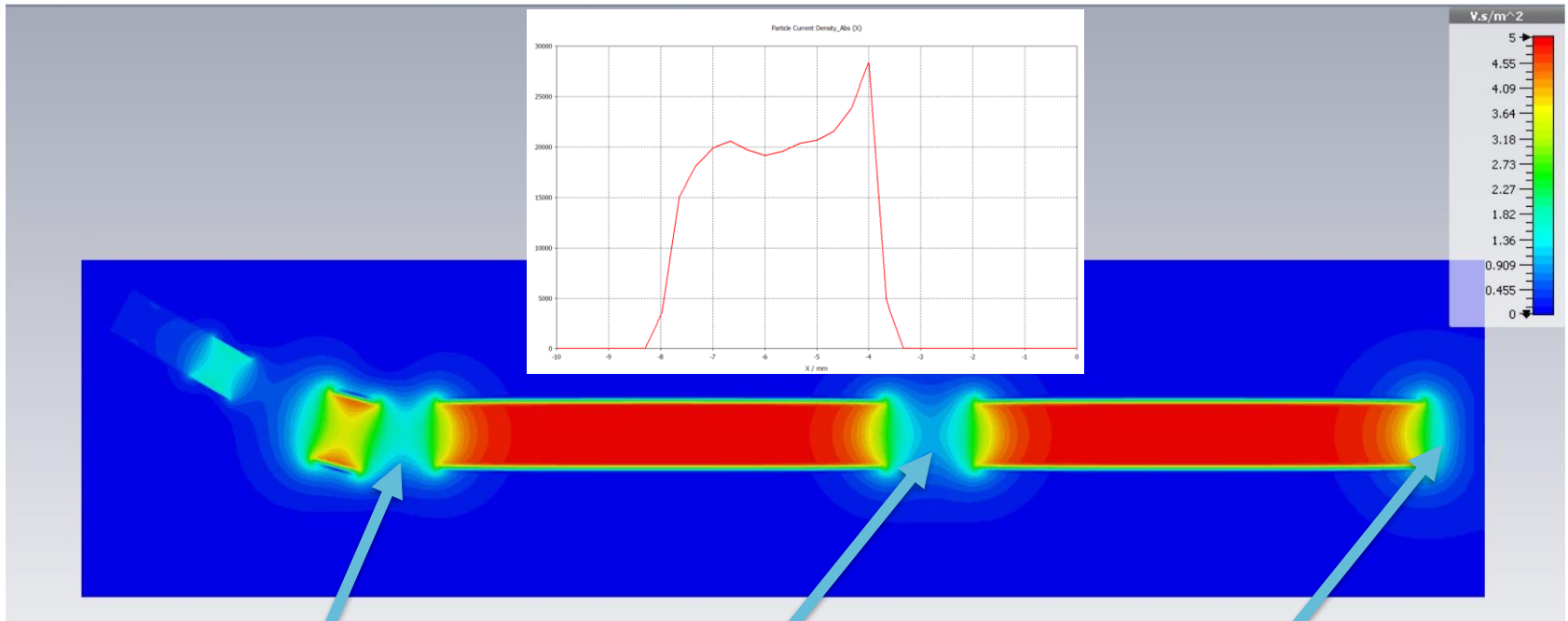
Stability diagram (HEL parameters)



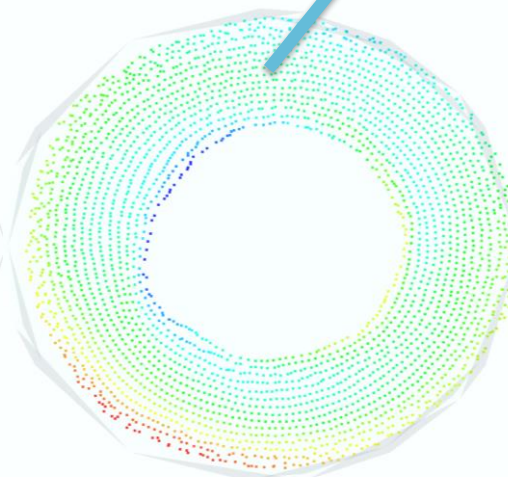
"1" corresponds to the beam with radii $r_1 = 0.9$ mm, $r_2 = 1.8$ mm
Gray region corresponds to the beam stable state (up to the stability criterion)
(Only the region above the 45° line has meaning because $r_2 > r_1$)

Particle tracking

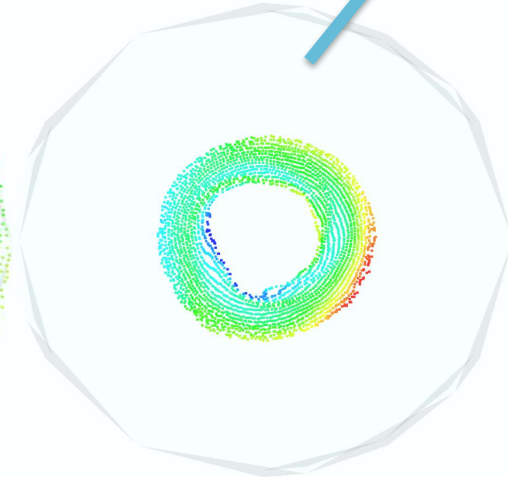




At the beginning of the interaction space



At the middle of the interaction space



At the end of the interaction space

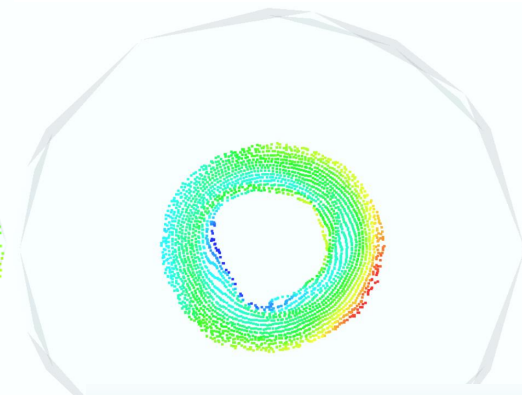
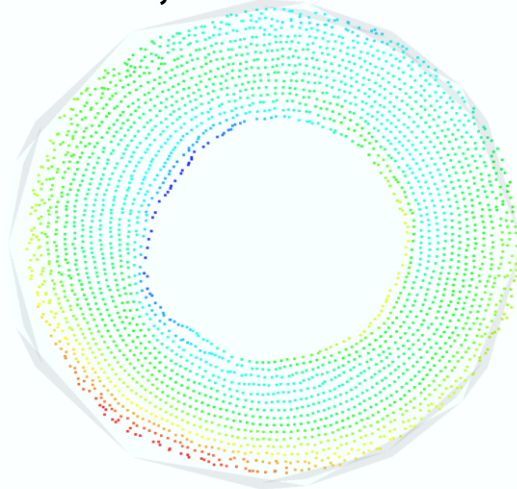
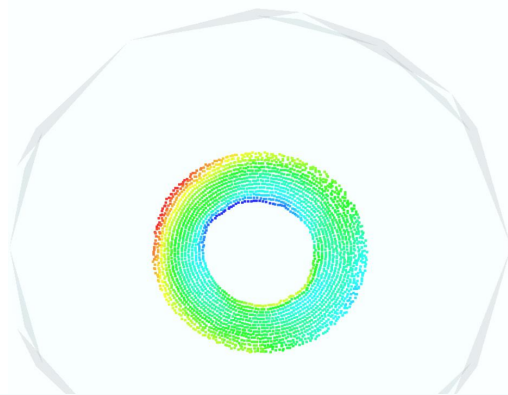
Main sol. with TRK and PIC solvers

1000 mm

2240 mm

3500 mm

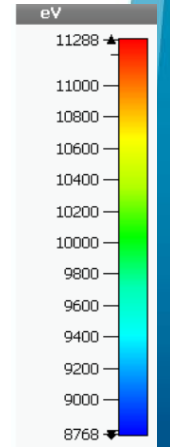
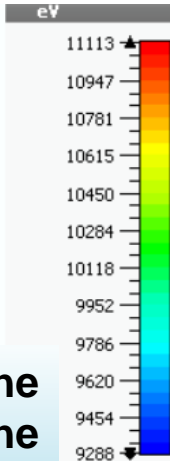
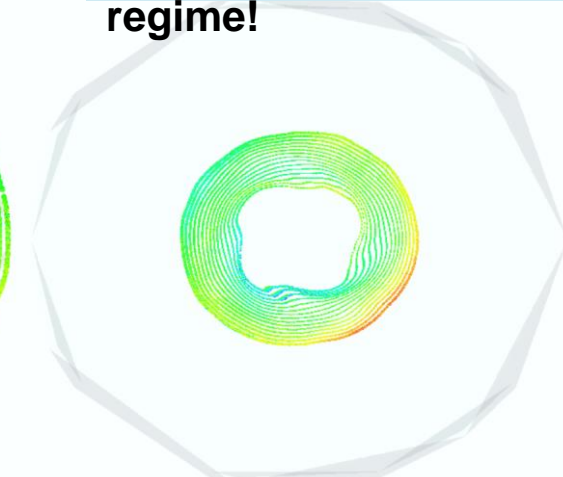
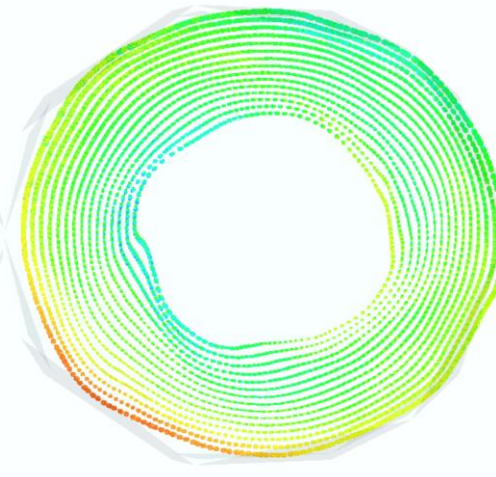
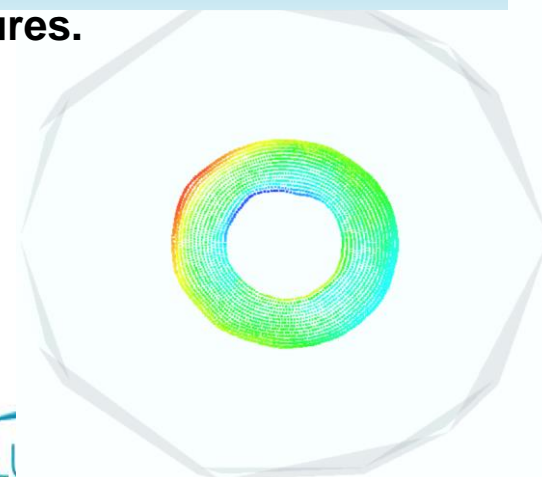
TRK, 0.5x0.5x0.5 mm



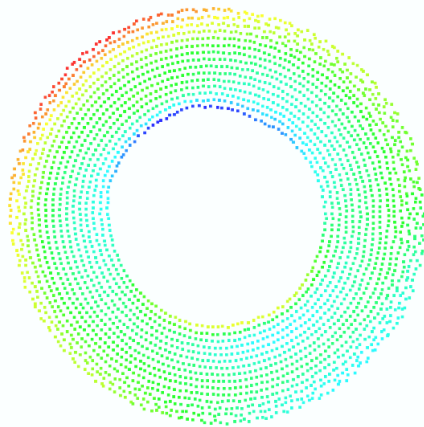
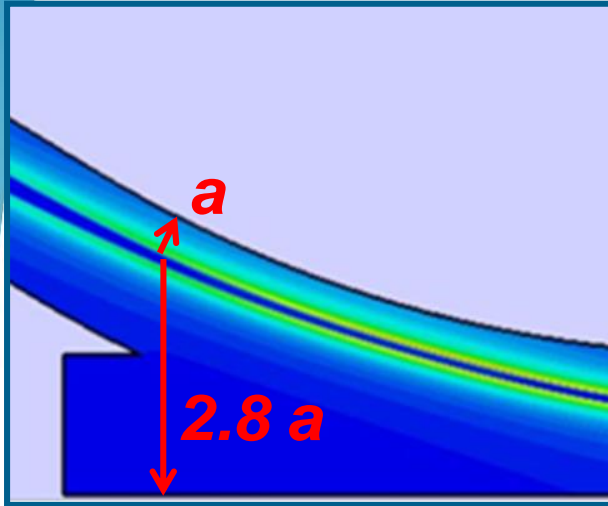
Particles with this maximum potential has to be rotated on 190 deg, what one can see on the figures.

PIC, 0.75x0.75x0.75 mm

The difference of the potentials results in the beam shape distortion. But this is not unstable regime!



Potential changing due to the asymmetry of the vacuum chamber



Beam is not centered between walls of the vacuum tube



Walls of vacuum chamber influence the beam = beam is influenced by external electric field

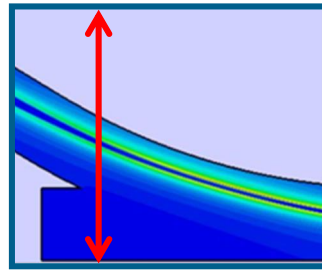
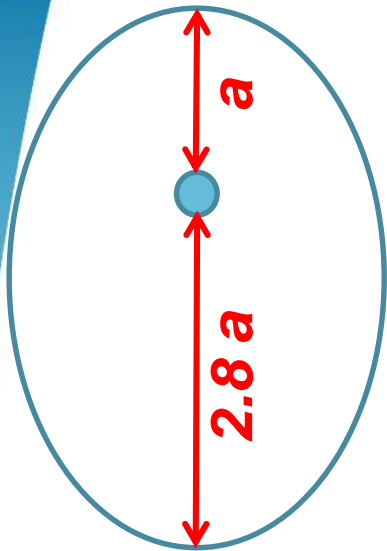


Beam tries to compensate influence of external electric field, but it is “frozen” in high magnetic field of 5 T so that particle density cannot change



Beam potential becomes asymmetric to compensate external electric field

Potential asymmetry at the outer radius

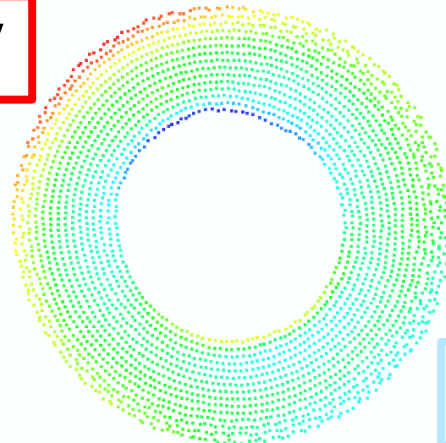


$$\Delta U_b \approx \frac{I}{2\pi v_b \epsilon_0} \ln \frac{r_b}{a} - \frac{I}{2\pi v_b \epsilon_0} \ln \frac{r_b}{2a + \sqrt{3}/2a} = \frac{I}{2\pi v_b \epsilon_0} \ln \frac{2a + \sqrt{3}/2a}{a}$$

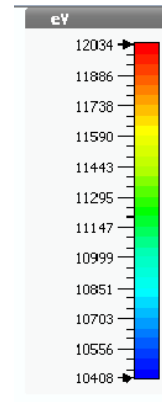
If the beam was solid with parameters $U_0 = 15$ keV, $v_b = 0.2$ c
 $I = 5$ A, $a = 30$ mm, $r_b = 1.8$ mm, potential difference is

$$\Delta U_b \approx 1.3 \text{ kV}$$

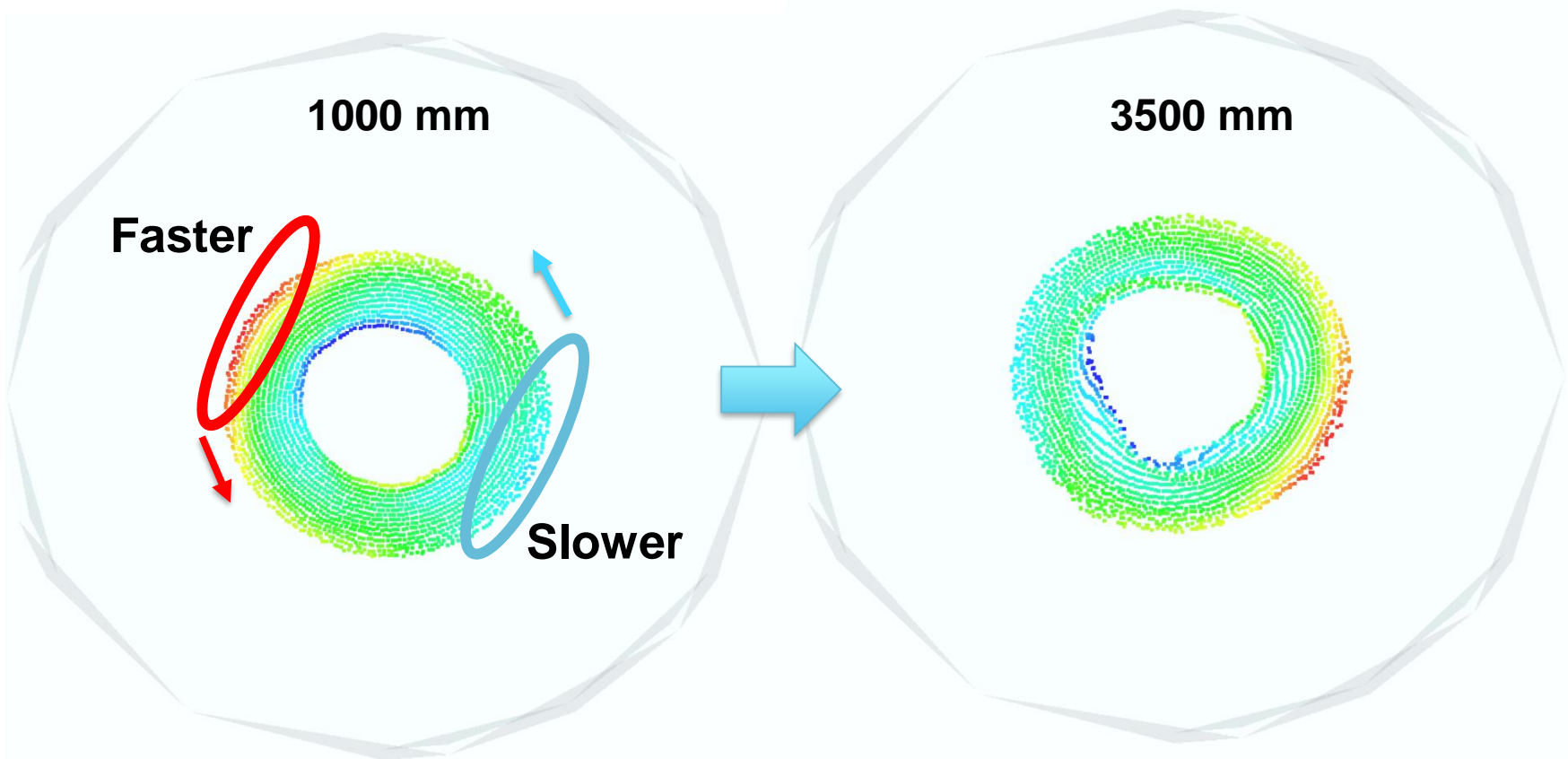
12 kV



10.8 kV



Potential asymmetry and beam shape perturbation



Particles starting from different azimuthal angles have different rotation velocity

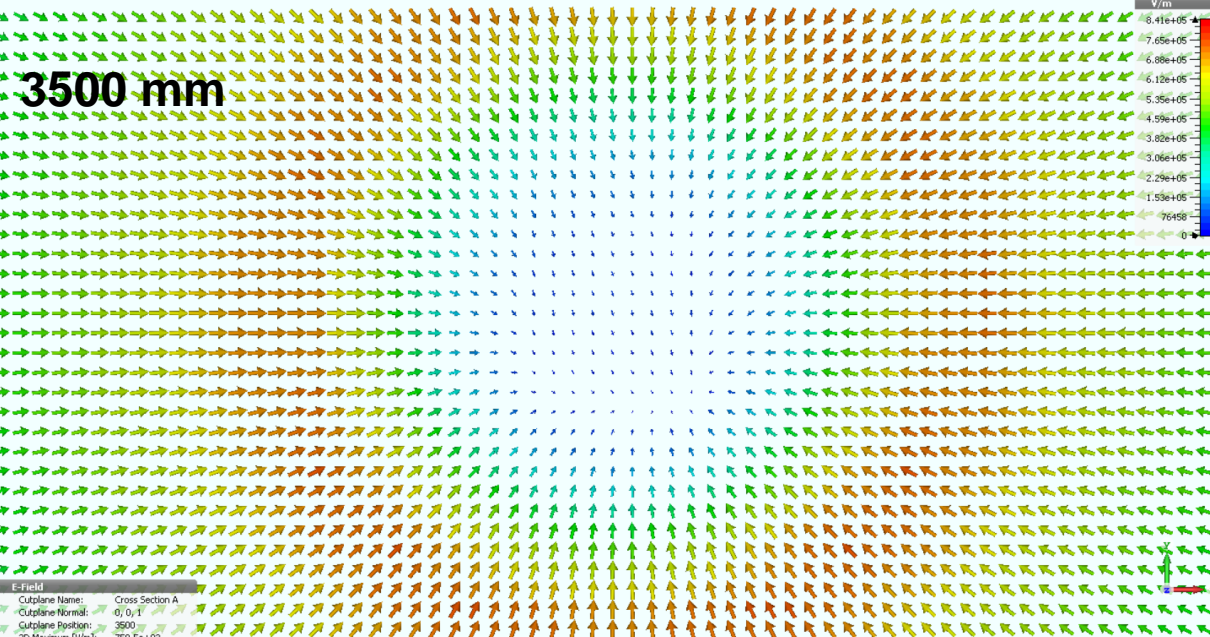
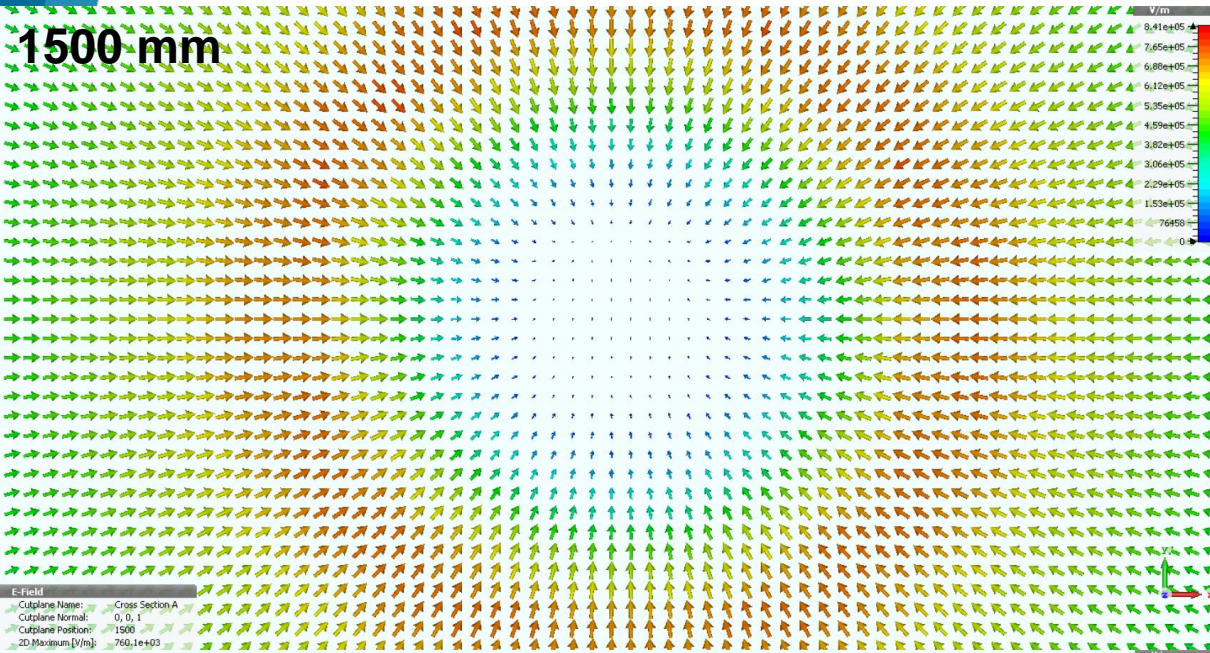


Perturbation of the azimuthally uniform electron density



Non-zero electric field in the hollow increasing with the beam motion through the vacuum chamber

HEL with TRK solver and with grid of 0.5x0.5 mm: fields



According to the simulation the maximum field is about 0.8 MV/m and according to the theory is 0.7 MV/m. The difference can be provided by slightly difference of the beam size and beam current (in the simulation the current is about 5.3 A).

The field irregularity in the middle of the beam is not more than 5% (~4% or about 30 kV/m) at the end of the HEL

Conclusion

Stability:

- The used beam parameters (current is not more than 5 A with voltage of 15 kV) are looked quite reliable. They should provide the stable regime of the HEL
- With exact design parameters $r_1 = 0.9$ mm, $r_2 = 1.8$ mm beam is not influenced by the diocotron instability

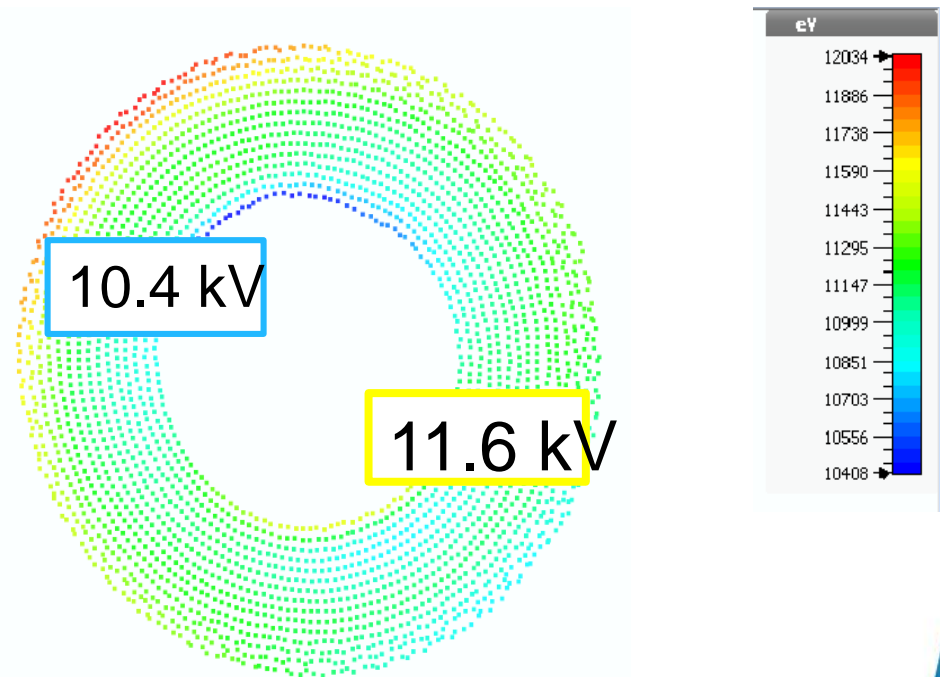
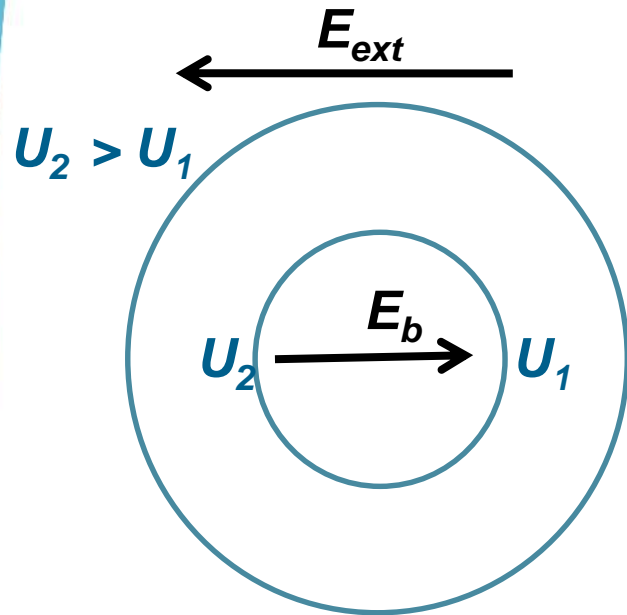
Particle tracking:

- All simulations were performed for the irregular particles distribution with “peak” current density near the inner radius. The influence of this peak on the beam motion is not observed and this peak is saved during the motion.
- The beam simulation was carried out with different PIC and TRK solvers With the same grid they give the same results
- **The amplitude of residual field in the center not more than 5% from the maximum field near the outer beam radius at the end of the HEL.**

Potential asymmetry at the inner radius

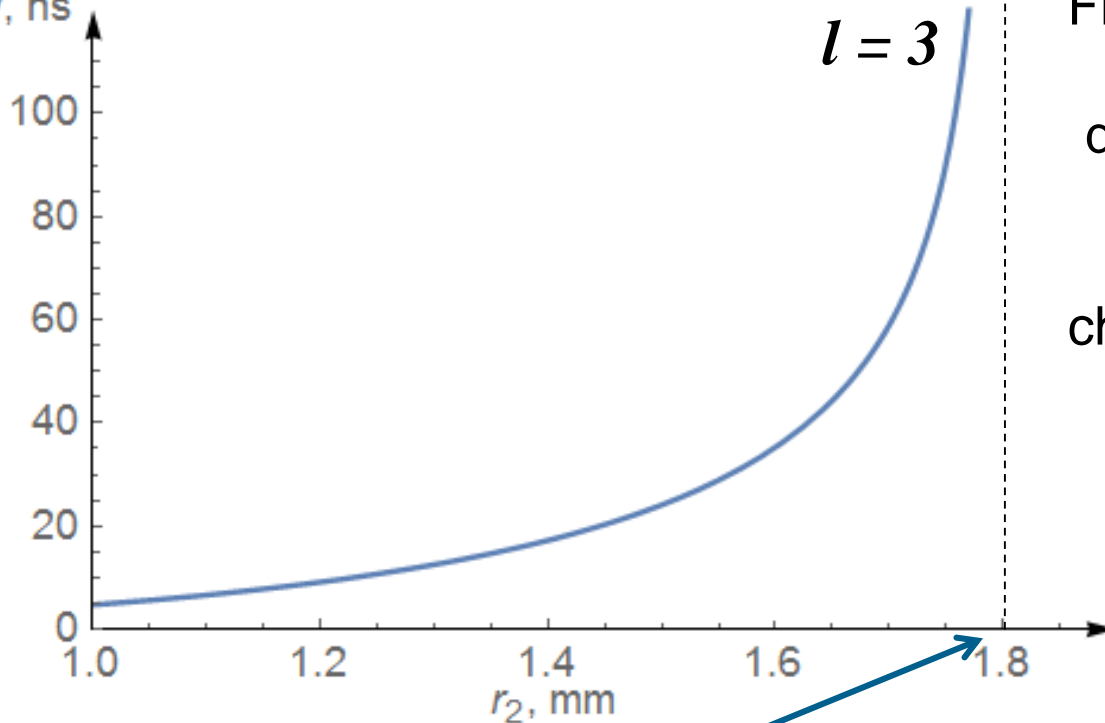
Beam is off-centered = beam is influenced by the equivalent electric field

External electric field does not penetrate into the hollow because beam compensates it therein, changing the potential



Dependence of time of instability growth on beam outer radius

Characteristic time of instability, ns



Fix inner radius $r_1 = 0.9$ mm, consider dependence on beam outer radius (which is more changeable because of beam field)

Design beam radii

$r_1 = 0.9$ mm, $r_2 = 1.8$ mm

correspond to the stable state

Example of deviation of designed parameter

$r_1 = 0.9$ mm, $r_2 = 1.6$ mm

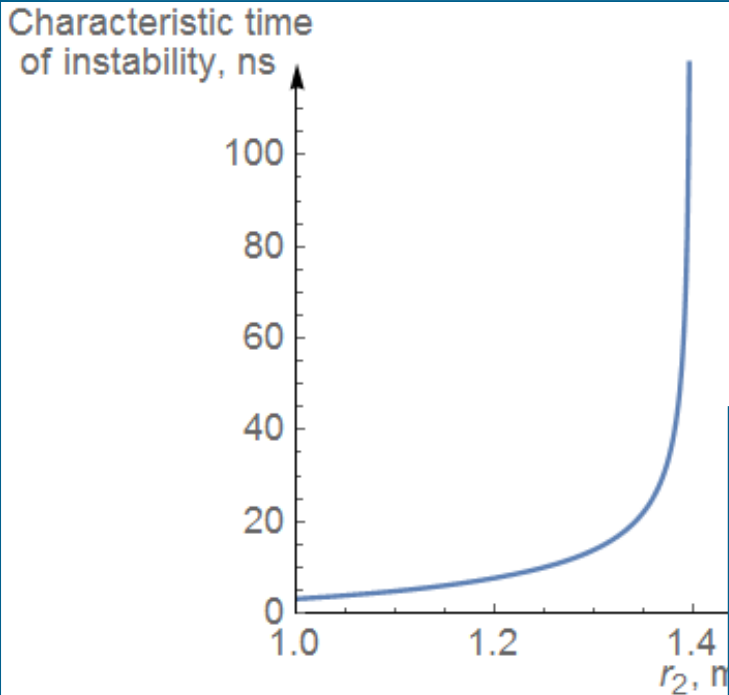
$\tau \approx 40$ ns

Time of flight ≈ 50 ns

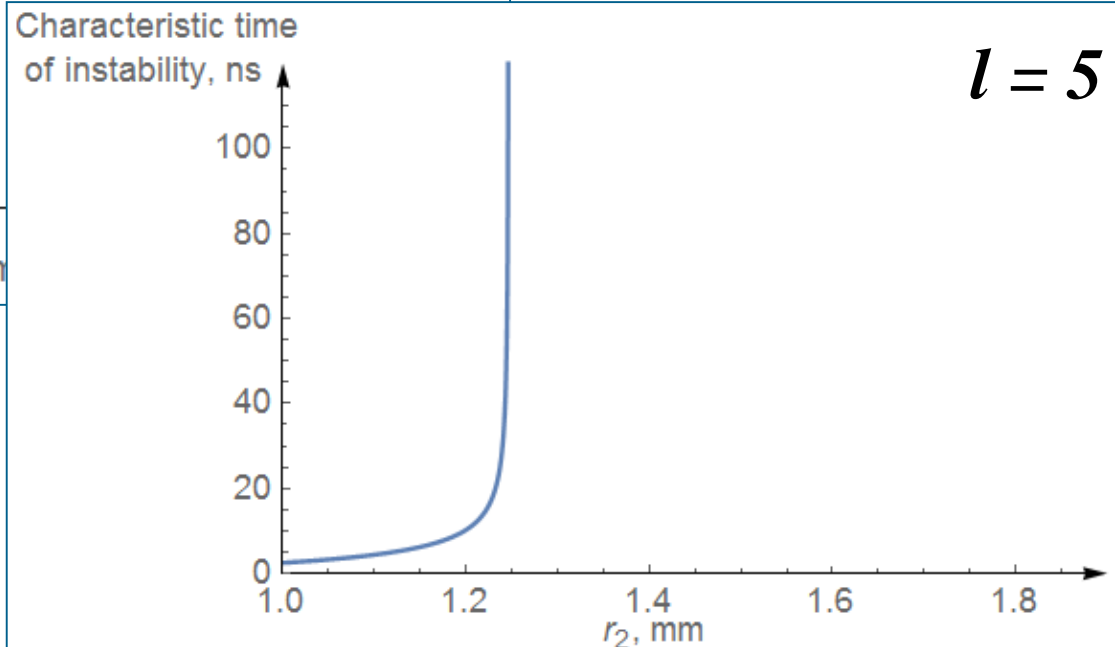
Dependence of time of instability growth vs beam outer radius

Fixed inner
radius $r_1 = 0.9$
mm

$l = 4$



$l = 5$



Modes with $l \geq 6$ are
less dangerous