

Pulsed Muon Beam Experiments

Yoshitaka Kuno
Department of Physics,
Osaka University, Japan

August 28th 2019
NuFACT2019,
Daegu, Korea

Introduction



Muon particle physics experiments and Beam time structure

Muon particle physics experiments and Beam time structure

Pulsed muon beam

for the experiments of

- delayed measurements ($\mu \rightarrow e$ conversion)
- with storage ring and acceleration (muon g-2)
- with pulsed instruments (laser)

proton synchrotron
(J-PARC, FNAL, RAL)

Muon particle physics experiments and Beam time structure

Pulsed muon beam

for the experiments of

- delayed measurements ($\mu \rightarrow e$ conversion)
- with storage ring and acceleration (muon g-2)
- with pulsed instruments (laser)

proton synchrotron
(J-PARC, FNAL, RAL)

DC muon beam

for the experiments of

- coincidence measurements with low instantaneous intensity ($\mu \rightarrow e\gamma$, $\mu \rightarrow eee$)

proton cyclotron
(PSI, TRIUMF, MuSIC)

Introduction: World Map



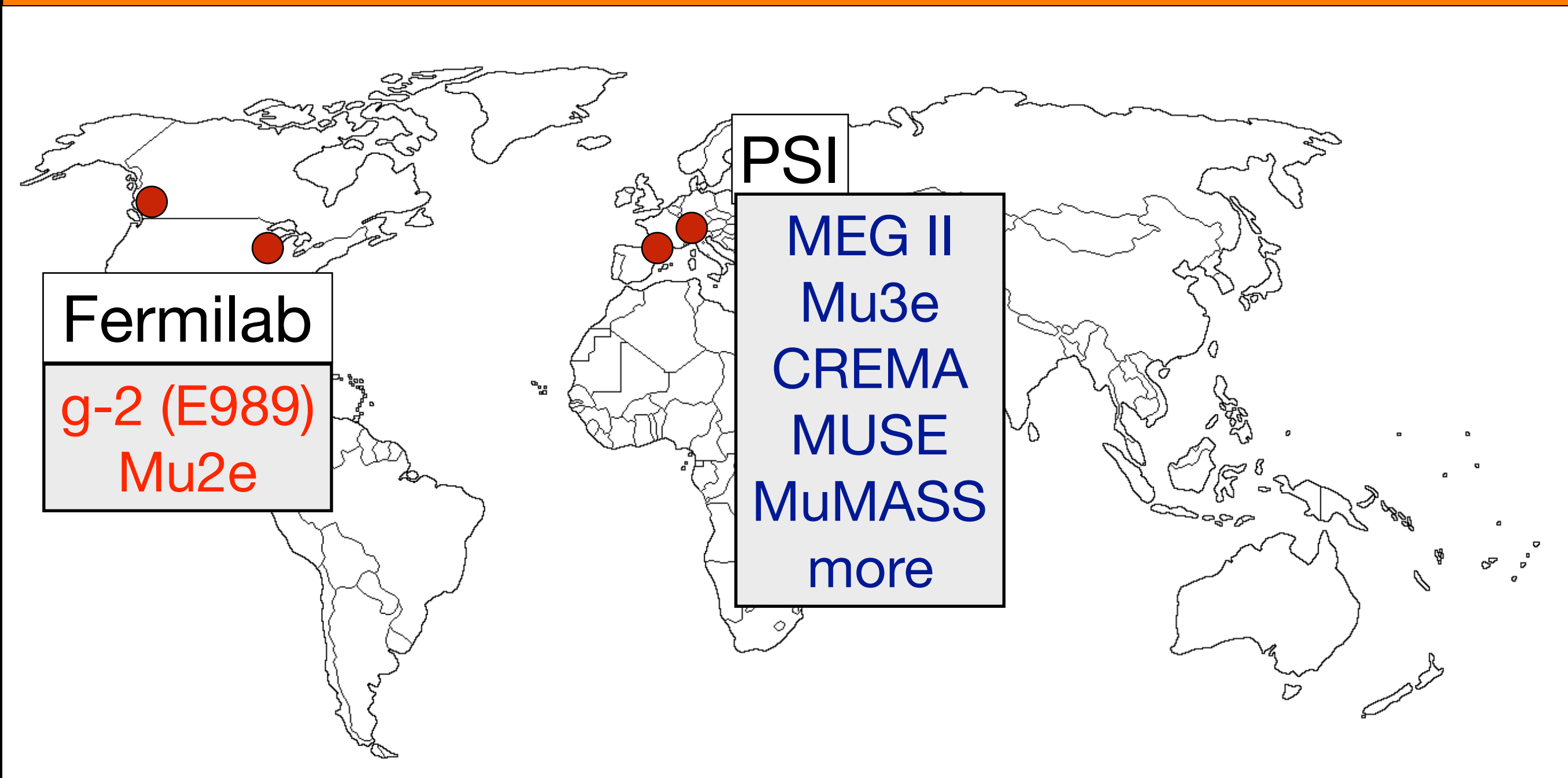
DC muon beam experiments
pulsed muon beam experiments

Introduction: World Map



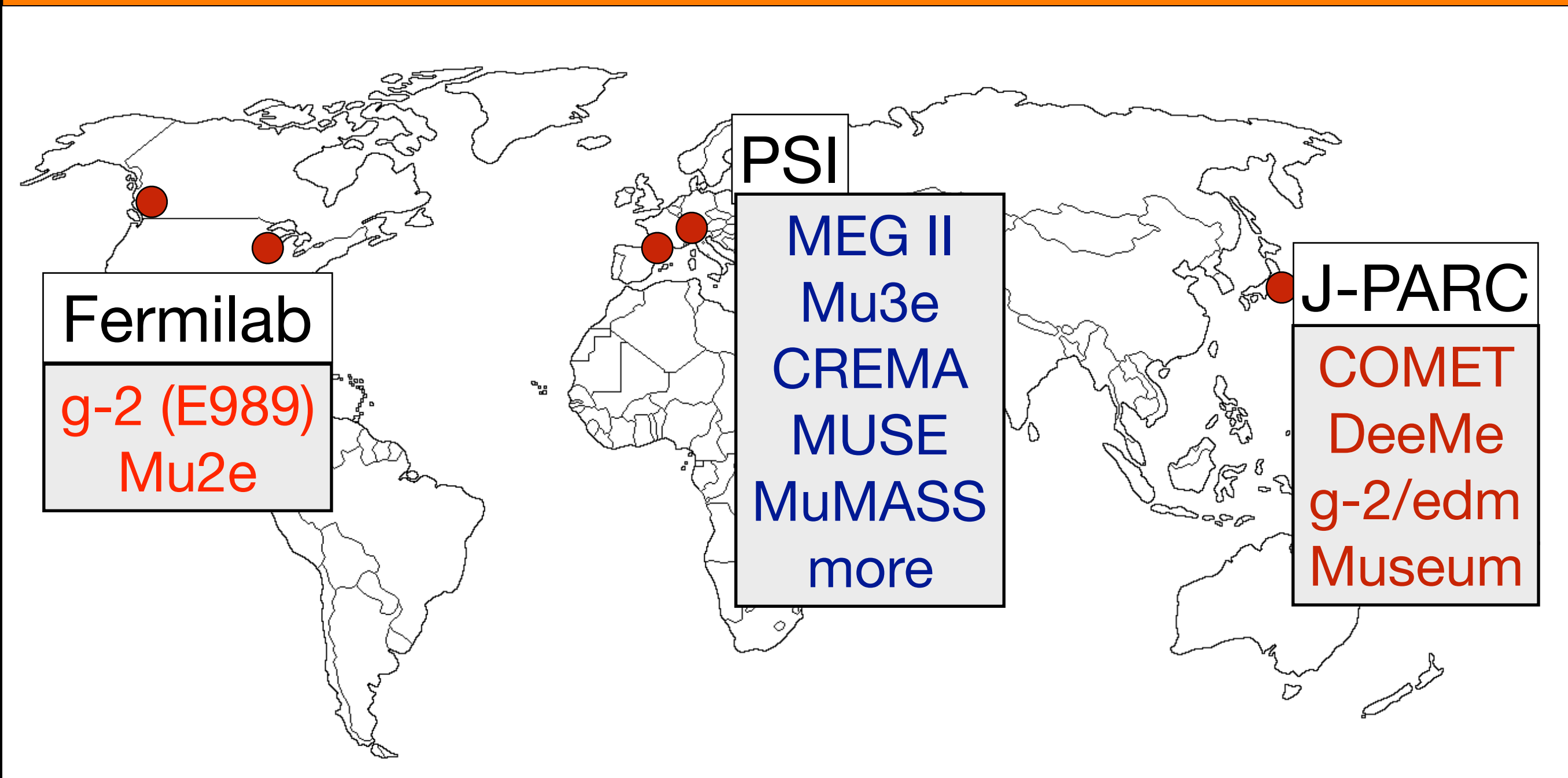
DC muon beam experiments
pulsed muon beam experiments

Introduction: World Map



DC muon beam experiments
pulsed muon beam experiments

Introduction: World Map



DC muon beam experiments
pulsed muon beam experiments

- Charged Lepton Flavour Violation (CLFV) with Muons
 - Physics Motivation
 - $\mu^- \rightarrow e^-$ conversion
 - $\mu^+ \rightarrow e^+ \gamma$, $\mu^+ \rightarrow e^+ e^+ e^-$ (just brief)
- Muon $g-2$
- Summary

- Charged Lepton Flavour Violation (CLFV) with Muons
 - Physics Motivation
 - $\mu^- \rightarrow e^-$ conversion
 - $\mu^+ \rightarrow e^+ \gamma$, $\mu^+ \rightarrow e^+ e^+ e^-$ (just brief)
- Muon $g-2$
- Summary

Details are given in the WG4 sessions.

Charged Lepton
flavour Violation
(CLFV) with Muons

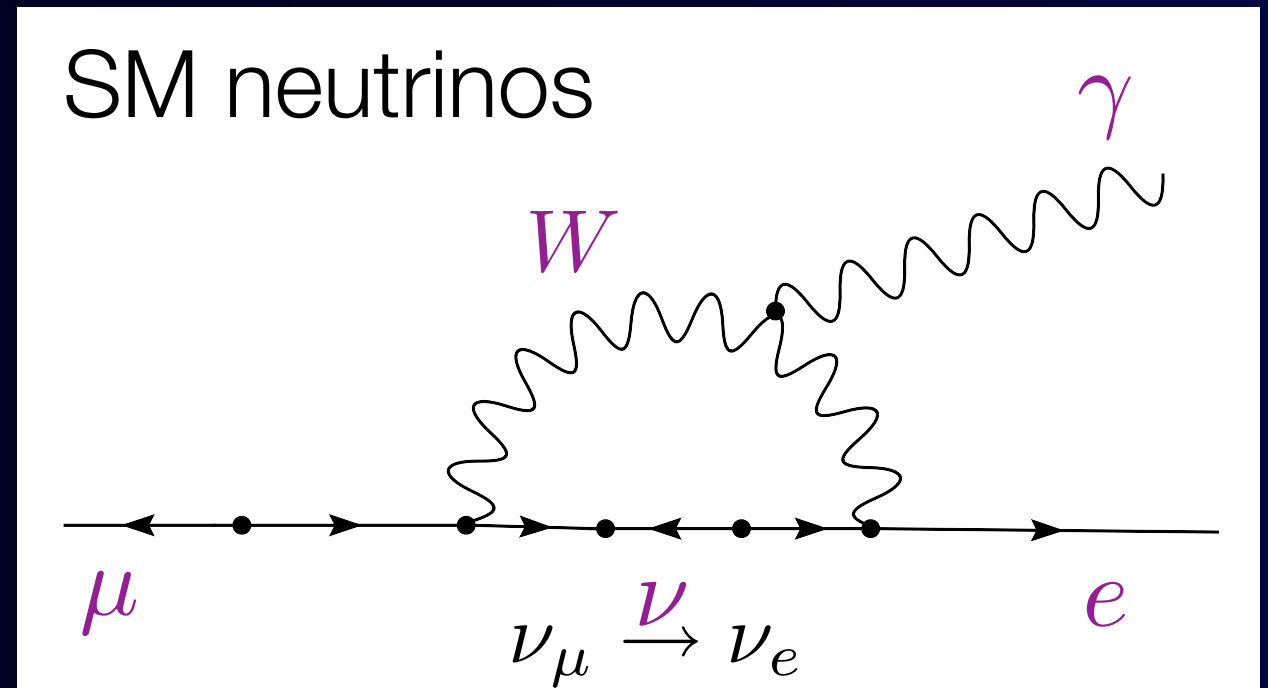


Lepton Mixing in the SM to CLFV

Neutral lepton flavour violation has been observed.
Lepton mixing in the SM has been established.

Lepton Mixing in the SM to CLFV

Neutral lepton flavour violation has been observed. Lepton mixing in the SM has been established.

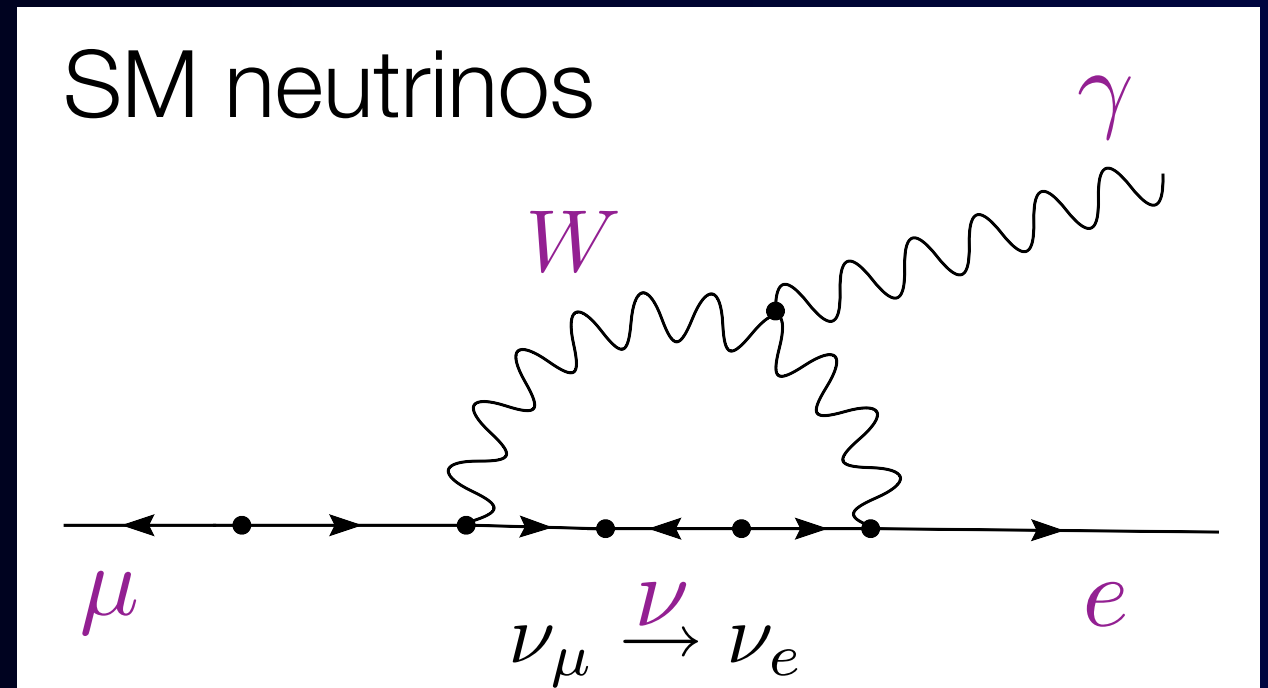


S.T. Petcov, Sov.J. Nucl. Phys. 25 (1977) 340

$$B(\mu \rightarrow e \gamma) = \frac{3\alpha}{32\pi} \left| \sum_l (V_{MNS})_{\mu l}^* (V_{MNS})_{el} \frac{m_{\nu_l}^2}{M_W^2} \right|^2$$

Lepton Mixing in the SM to CLFV

Neutral lepton flavour violation has been observed. Lepton mixing in the SM has been established.



S.T. Petcov, Sov.J. Nucl. Phys. 25 (1977) 340

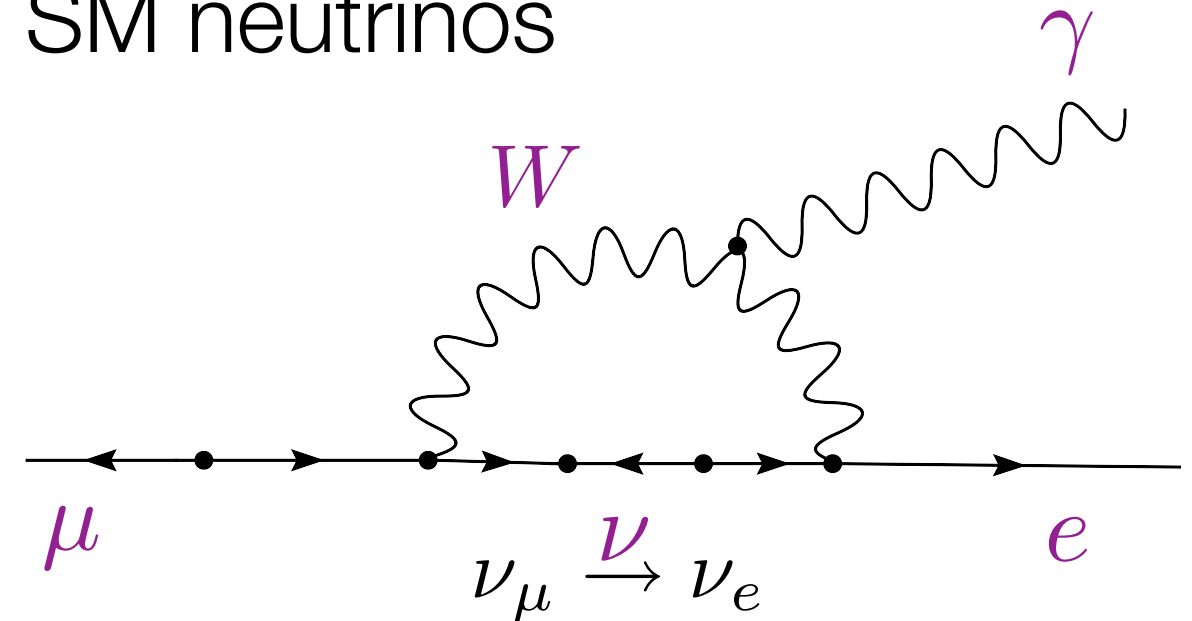
$$B(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_l (V_{MNS})_{\mu l}^* (V_{MNS})_{el} \frac{m_{\nu_l}^2}{M_W^2} \right|^2$$

$$\text{BR} \sim \mathcal{O}(10^{-54})$$

Lepton Mixing in the SM to CLFV

Neutral lepton flavour violation has been observed. Lepton mixing in the SM has been established.

SM neutrinos



S.T. Petcov, Sov.J. Nucl. Phys. 25 (1977) 340

$$B(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_l (V_{MNS})_{\mu l}^* (V_{MNS})_{el} \frac{m_{\nu_l}^2}{M_W^2} \right|^2$$

BR ~ O(10⁻⁵⁴)

Current upper limits on \mathcal{B}_i



Large window for BSM search without SM backgrounds

New Physics Energy Scale of CLFV Search

Effective Field Theory (EFT) Approach

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \sum_{d>4} \frac{C^{(d)}}{\Lambda^{d-4}}$$

Λ is the energy scale of new physics
 $C^{(d)}$ is the coupling constant.

New Physics Energy Scale of CLFV Search

Effective Field Theory (EFT) Approach

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \sum_{d>4} \frac{C^{(d)}}{\Lambda^{d-4}}$$

Λ is the energy scale of new physics
 $C^{(d)}$ is the coupling constant.

from $BR(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13}$

$$\frac{C^6}{\Lambda^2} \mathcal{O}^6 \rightarrow \frac{C^6}{\Lambda^2} \bar{e}_L \sigma^{\rho\nu} \mu_R \Phi F_{\rho\nu}$$



$$\Lambda \sim \mathcal{O}(10^4) \text{ TeV}$$

c.f. $\Delta m_K, \epsilon'$

	$ C_a [\Lambda = 1 \text{ TeV}]$	$\Lambda \text{ (TeV)} [C_a = 1]$	CLFV Process
$C_{e\gamma}^{\mu e}$	2.1×10^{-10}	<u>6.8×10^4</u>	$\mu \rightarrow e\gamma$
$C_{le}^{\mu\mu\mu e, e\mu\mu\mu}$	1.8×10^{-4}	75	$\mu \rightarrow e\gamma$ [1-loop]
$C_{le}^{\mu\tau\tau e, e\tau\tau\mu}$	1.0×10^{-5}	312	$\mu \rightarrow e\gamma$ [1-loop]
$C_{e\gamma}^{\mu e}$	4.0×10^{-9}	<u>1.6×10^4</u>	$\mu \rightarrow eee$
$C_{ll,ee}^{\mu eee}$	2.3×10^{-5}	207	$\mu \rightarrow eee$
$C_{le}^{\mu eee, ee\mu e}$	3.3×10^{-5}	174	$\mu \rightarrow eee$
$C_{e\gamma}^{\mu e}$	5.2×10^{-9}	<u>1.4×10^4</u>	$\mu^- \text{Au} \rightarrow e^- \text{Au}$
$C_{lq,ld,ed}^{e\mu}$	1.8×10^{-6}	745	$\mu^- \text{Au} \rightarrow e^- \text{Au}$
$C_{eq}^{e\mu}$	9.2×10^{-7}	1.0×10^3	$\mu^- \text{Au} \rightarrow e^- \text{Au}$
$C_{lu,eu}^{e\mu}$	2.0×10^{-6}	707	$\mu^- \text{Au} \rightarrow e^- \text{Au}$

F. Feruglio, P. Paradisi and A. Pattori, Eur. Phys. J. C 75 (2015) no.12, 579

G. M. Pruna and A. Signer, JHEP 1410 (2014) 014

New Physics Energy Scale of CLFV Search Future



Future planned CLFV experiments (with muons) expecting improvements by an additional factor of $>10,000$ or more (will be described later) would probe

New Physics Energy Scale of CLFV Search Future



Future planned CLFV experiments (with muons) expecting improvements by an additional factor of $>10,000$ or more (will be described later) would probe

$$\Lambda \sim \mathcal{O}(10^5) \text{ TeV}$$

$$R \propto \frac{1}{\Lambda^4}$$

CLFV would explore scales way beyond the energies that our present and future colliders can directly reach.

It is crucial in establishing where is the next fundamental scale above the electroweak symmetry breaking.

“Golden” $\mu \rightarrow e$ CLFV Transition Processes



$$\mu^+ \rightarrow e^+ \gamma$$

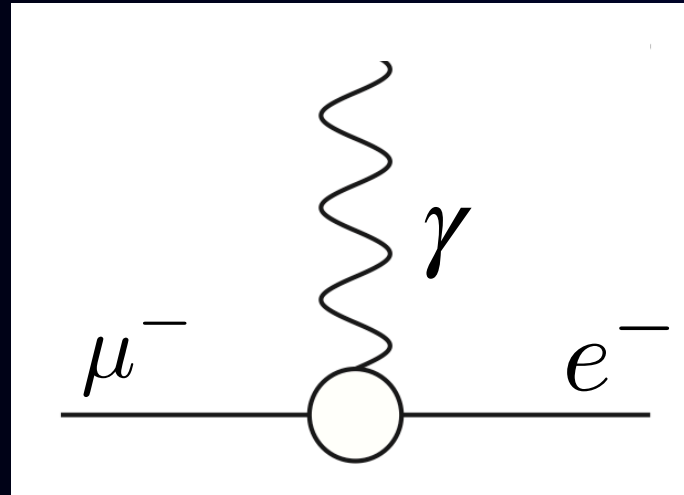
$$\mu^+ \rightarrow e^+ e^+ e^-$$

$$\mu^- N \rightarrow e^- N$$

“Golden” $\mu \rightarrow e$ CLFV Transition Processes

dipole interaction

$$\mu^+ \rightarrow e^+ \gamma$$



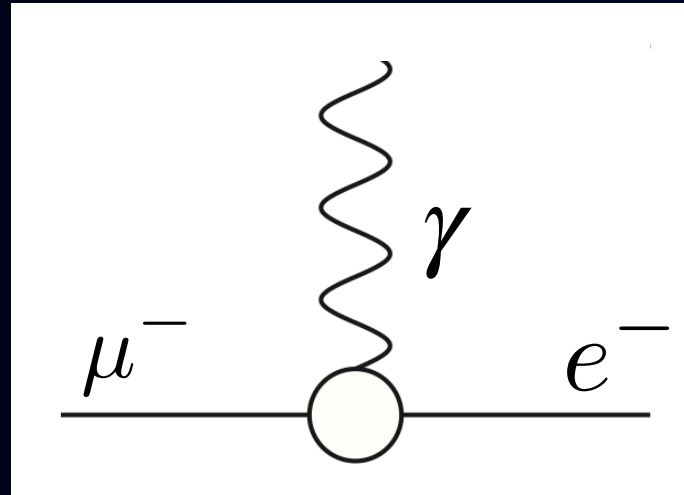
$$\mu^+ \rightarrow e^+ e^+ e^-$$

$$\mu^- N \rightarrow e^- N$$

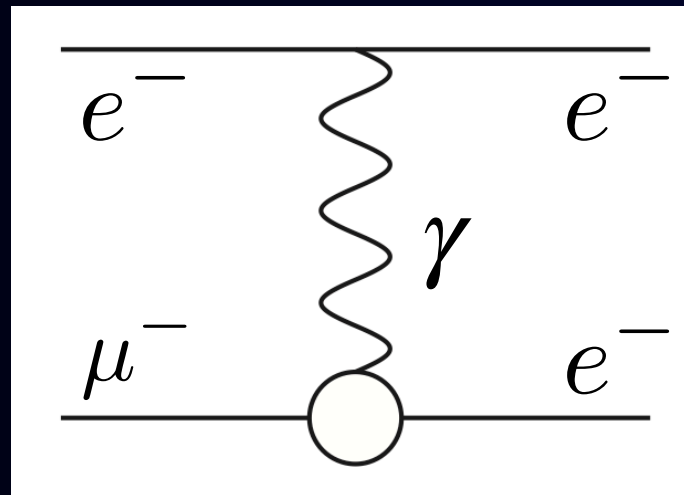
“Golden” $\mu \rightarrow e$ CLFV Transition Processes

dipole interaction

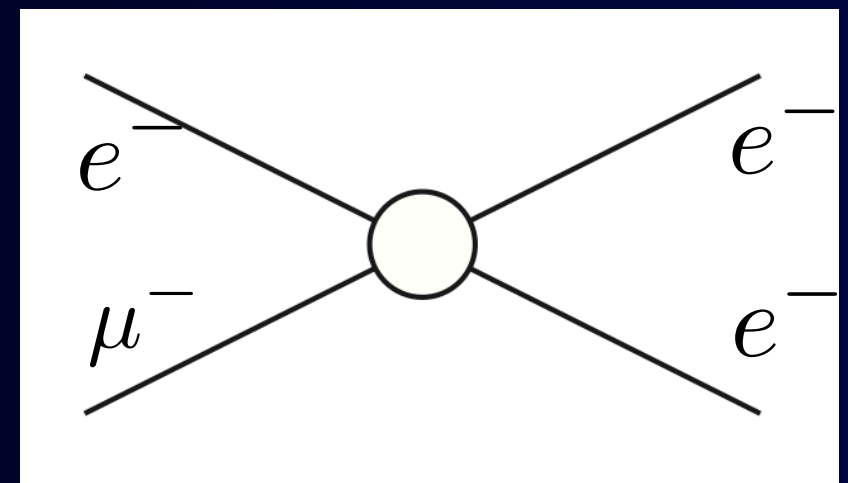
$$\mu^+ \rightarrow e^+ \gamma$$



$$\mu^+ \rightarrow e^+ e^+ e^-$$



contact interaction

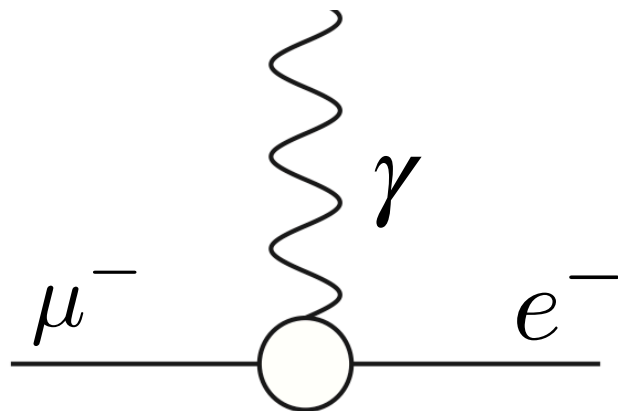


$$\mu^- N \rightarrow e^- N$$

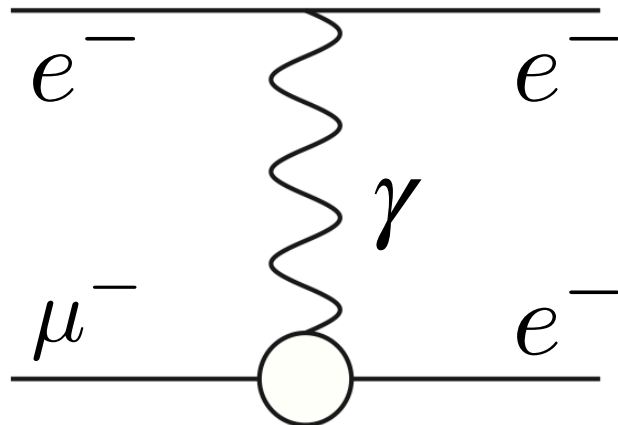
“Golden” $\mu \rightarrow e$ CLFV Transition Processes

dipole interaction

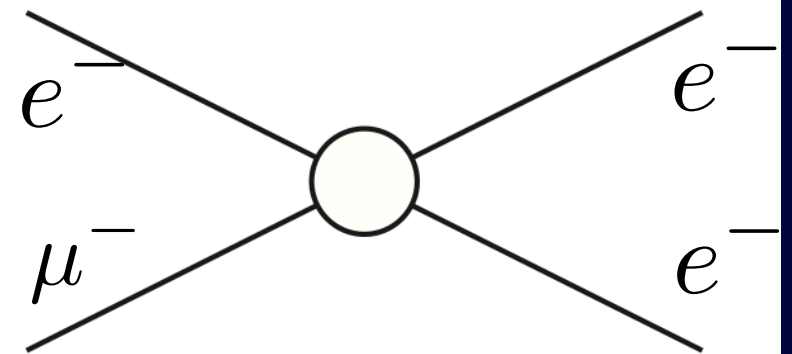
$$\mu^+ \rightarrow e^+ \gamma$$



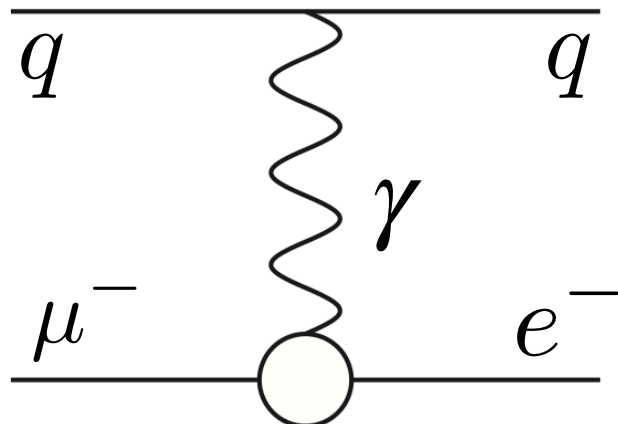
$$\mu^+ \rightarrow e^+ e^+ e^-$$



contact interaction



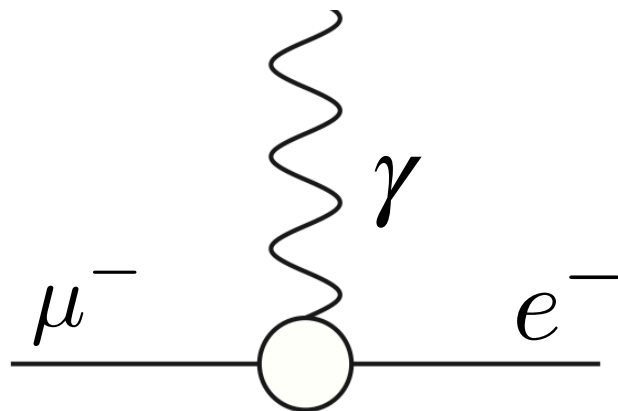
$$\mu^- N \rightarrow e^- N$$



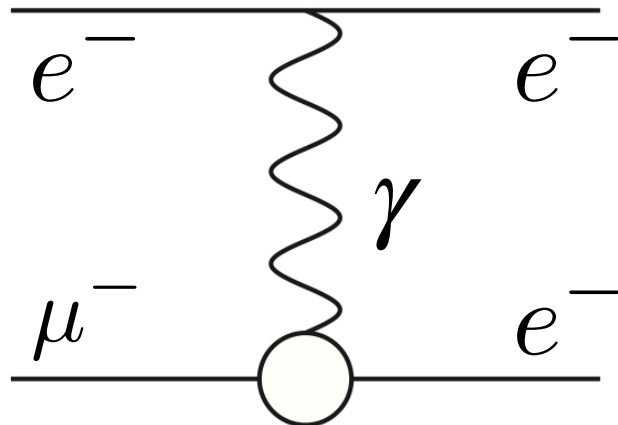
“Golden” $\mu \rightarrow e$ CLFV Transition Processes

dipole interaction

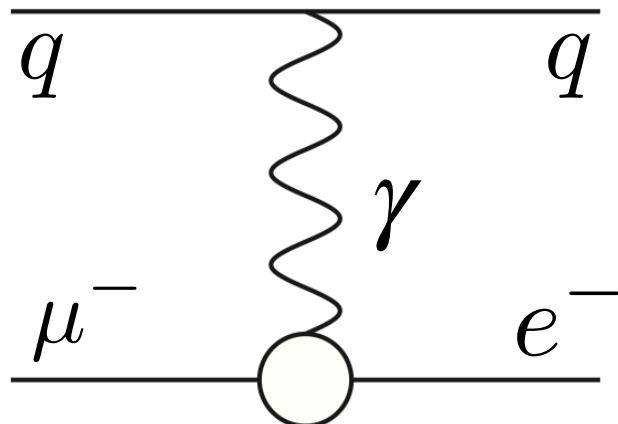
$$\mu^+ \rightarrow e^+ \gamma$$



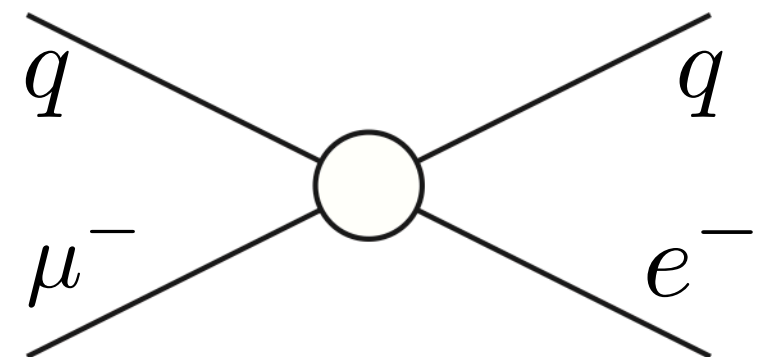
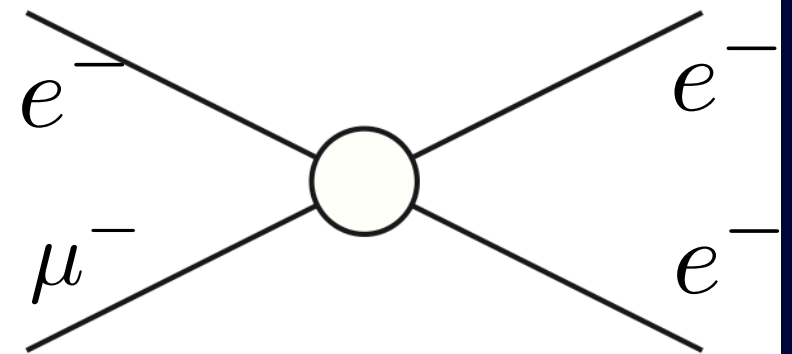
$$\mu^+ \rightarrow e^+ e^+ e^-$$



$$\mu^- N \rightarrow e^- N$$



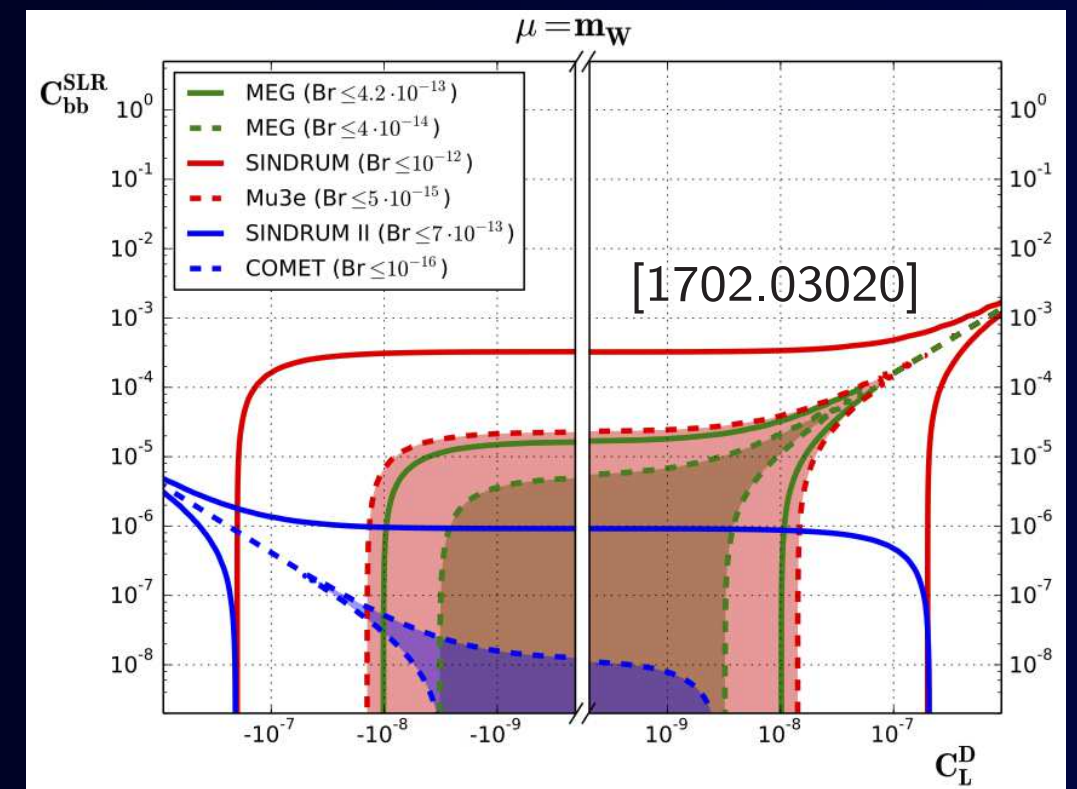
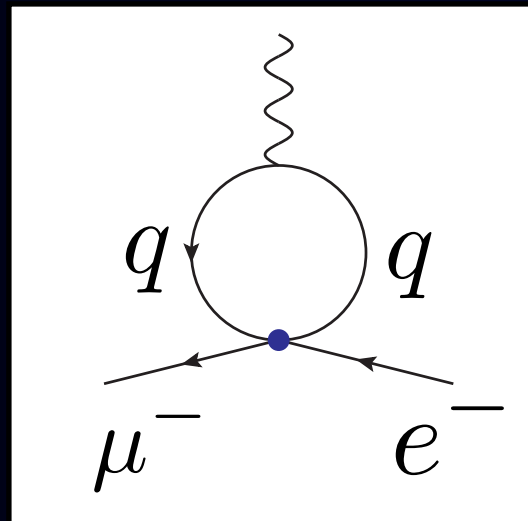
contact interaction



Operator Mixing via RGE

EFT at high physics scale

The operators are mixed in RGE at the experiment scale



All processes are equally important (complementary).

A. Crivellin, S. Davidson, G.M. Pruna and A. Signer, arXiv:1611.03409

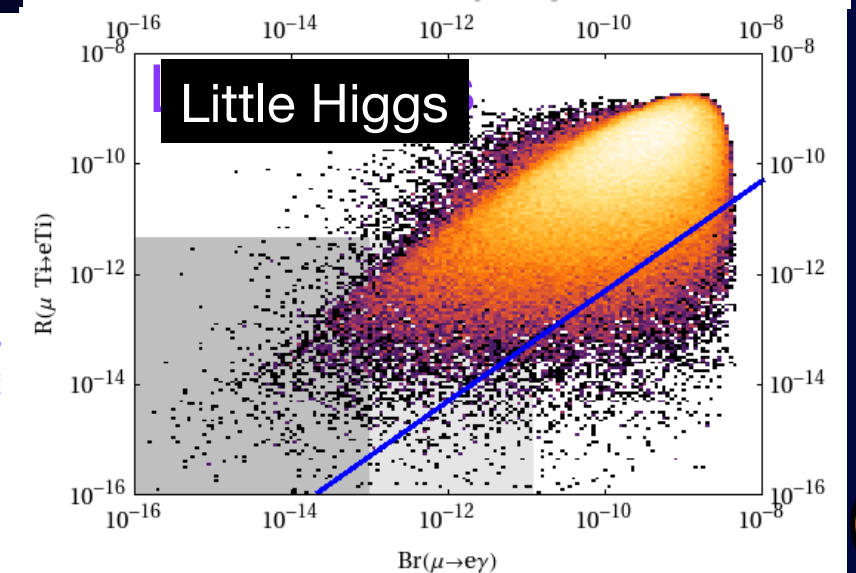
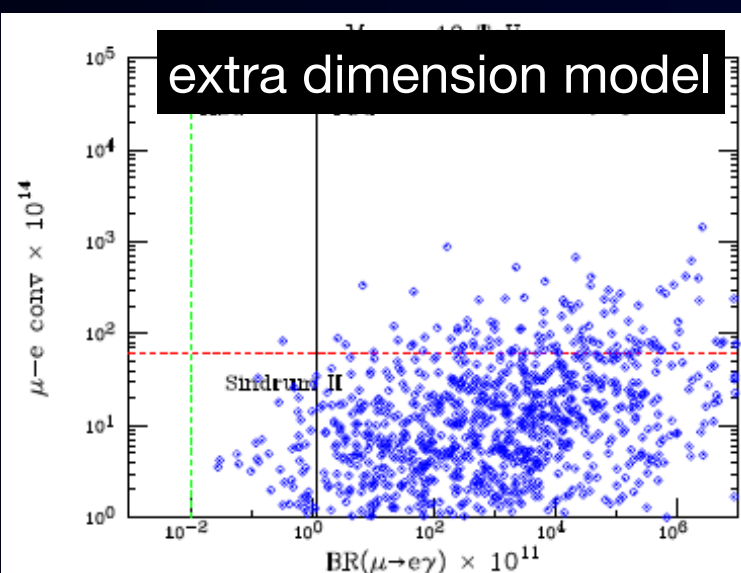
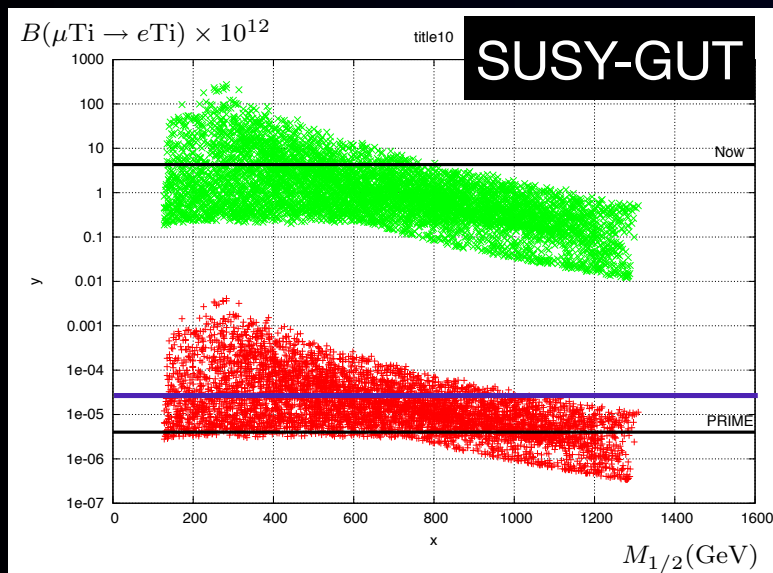
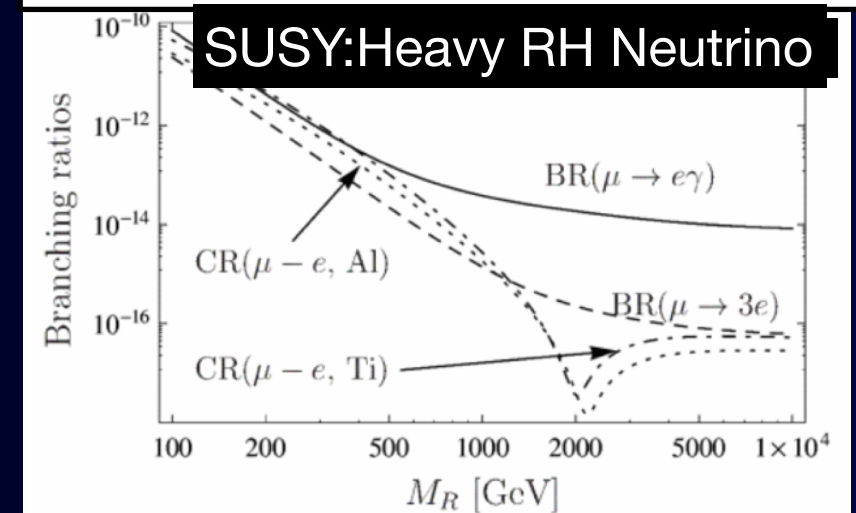
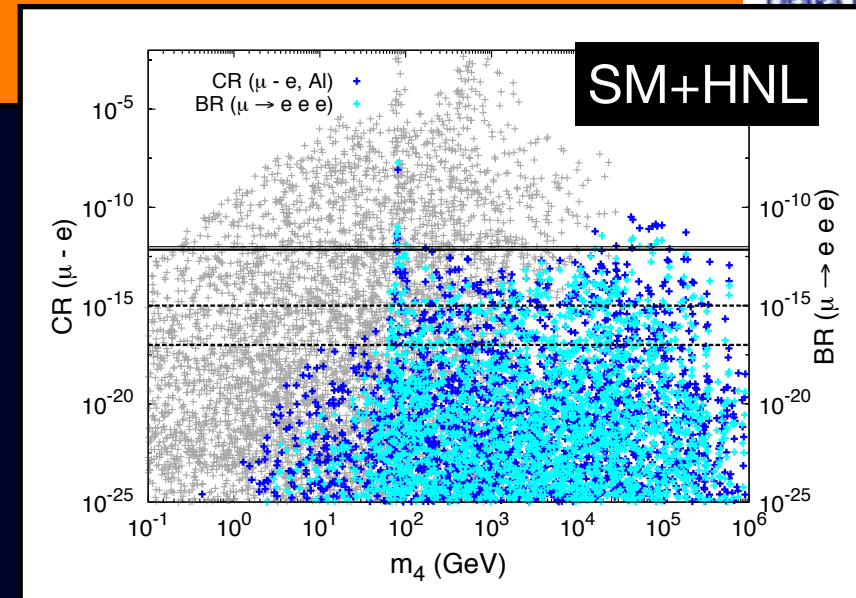
A. Crivellin, S. Davidson, G.M. Pruna and A. Signer, JHEP 117 (2017) no.5

S. Davidson, Eur. Phys. J. C76 (2016) 370

Model dependent CLFV



- SM + NHL (neutral heavy lepton)
- large extra dimensions
- extended Higgs sector
- additional vector boson (Z')
- leptoquark
- SUSY-GUT and SUSY seesaw
- R-parity violating SUSY
- low-energy seesaw
- etc. etc.

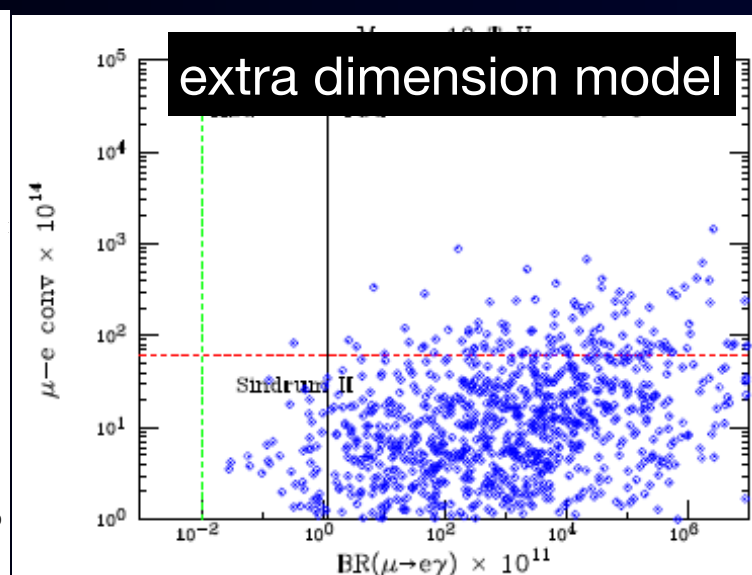
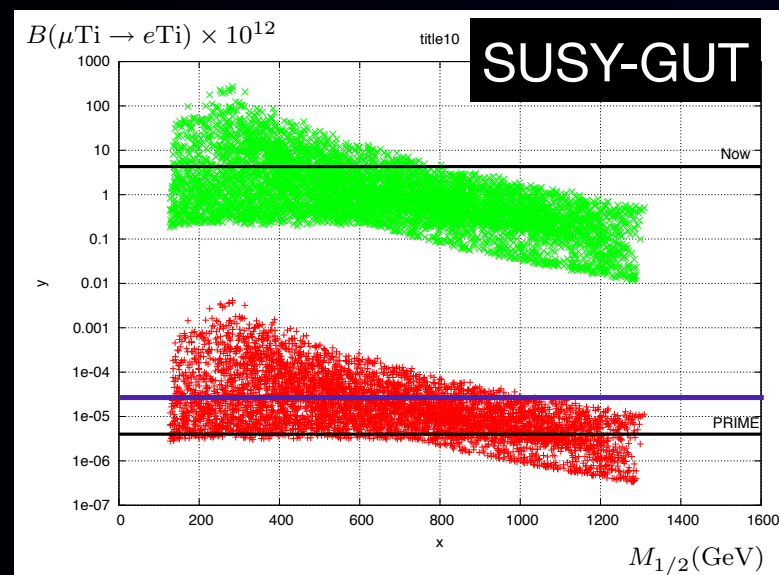
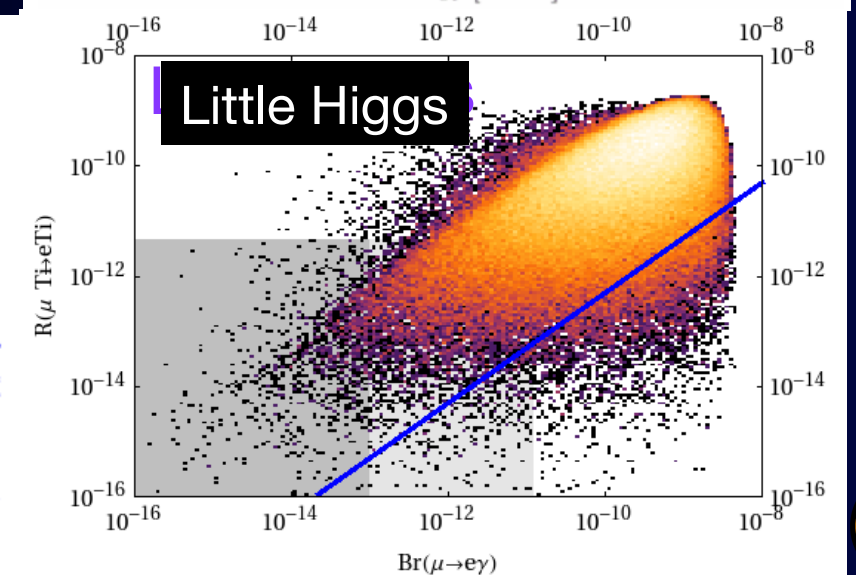
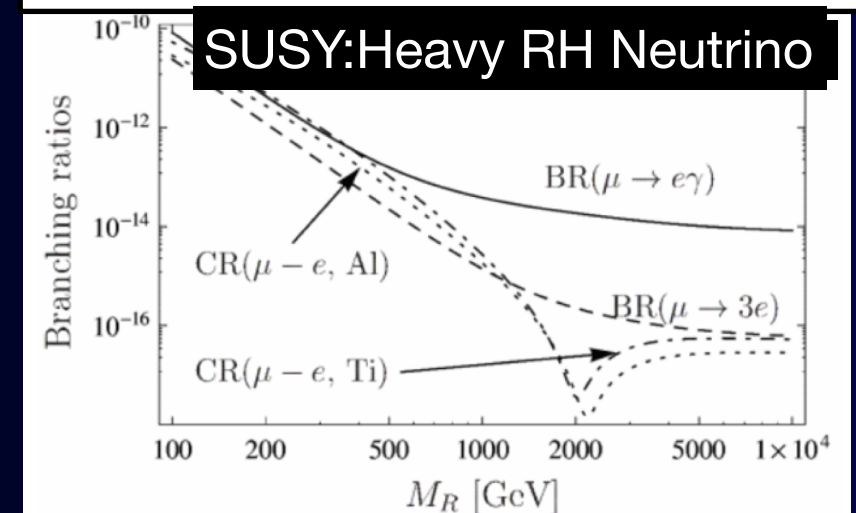
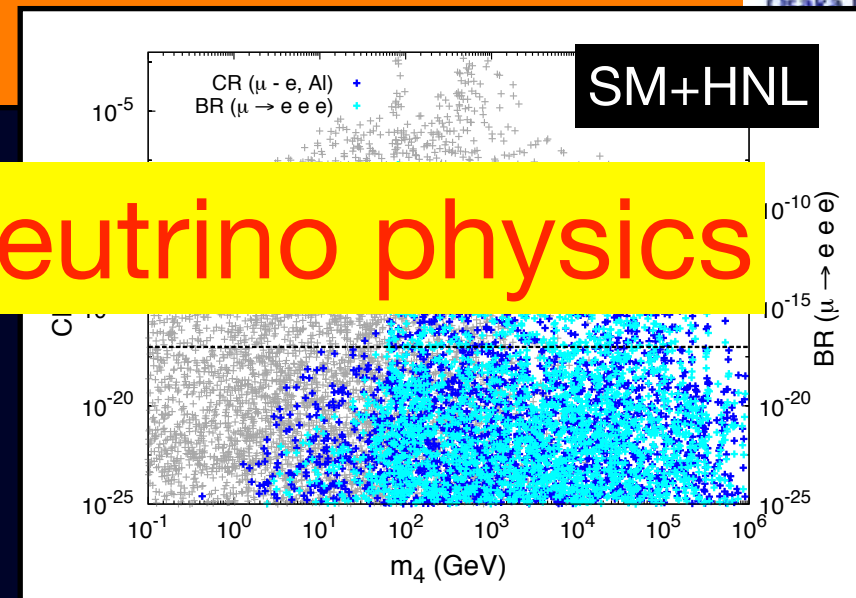


Model dependent CLFV



- SM + NHL (neutral heavy lepton)
- large extra dimensions
- extended Higgs sector
- additional vector boson (Z')
- leptoquark
- SUSY-GUT and SUSY seesaw
- R-parity violating SUSY
- low-energy seesaw
- etc. etc.

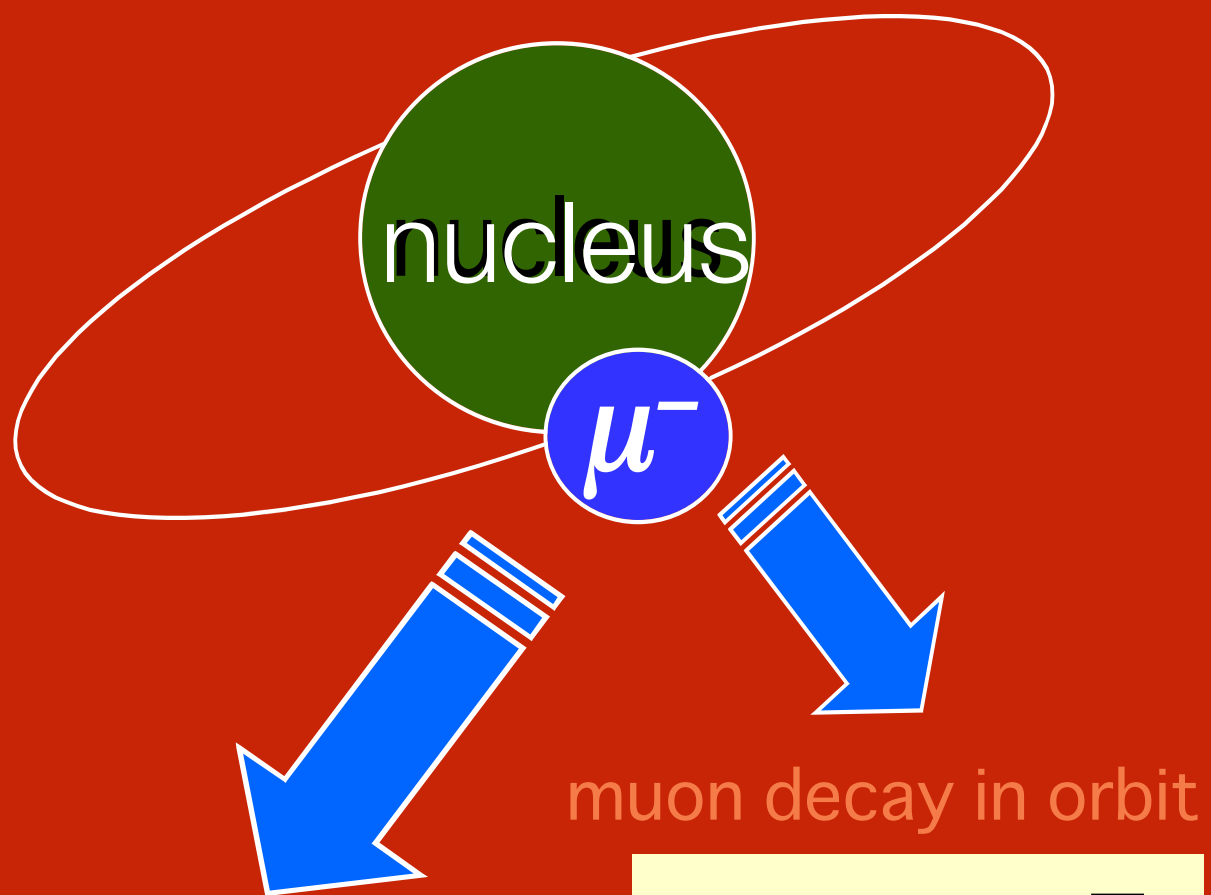
neutrino physics



$\mu \rightarrow e$ conversion
in
a muonic atom

What is $\mu \rightarrow e$ Conversion ?

1s state in a muonic atom



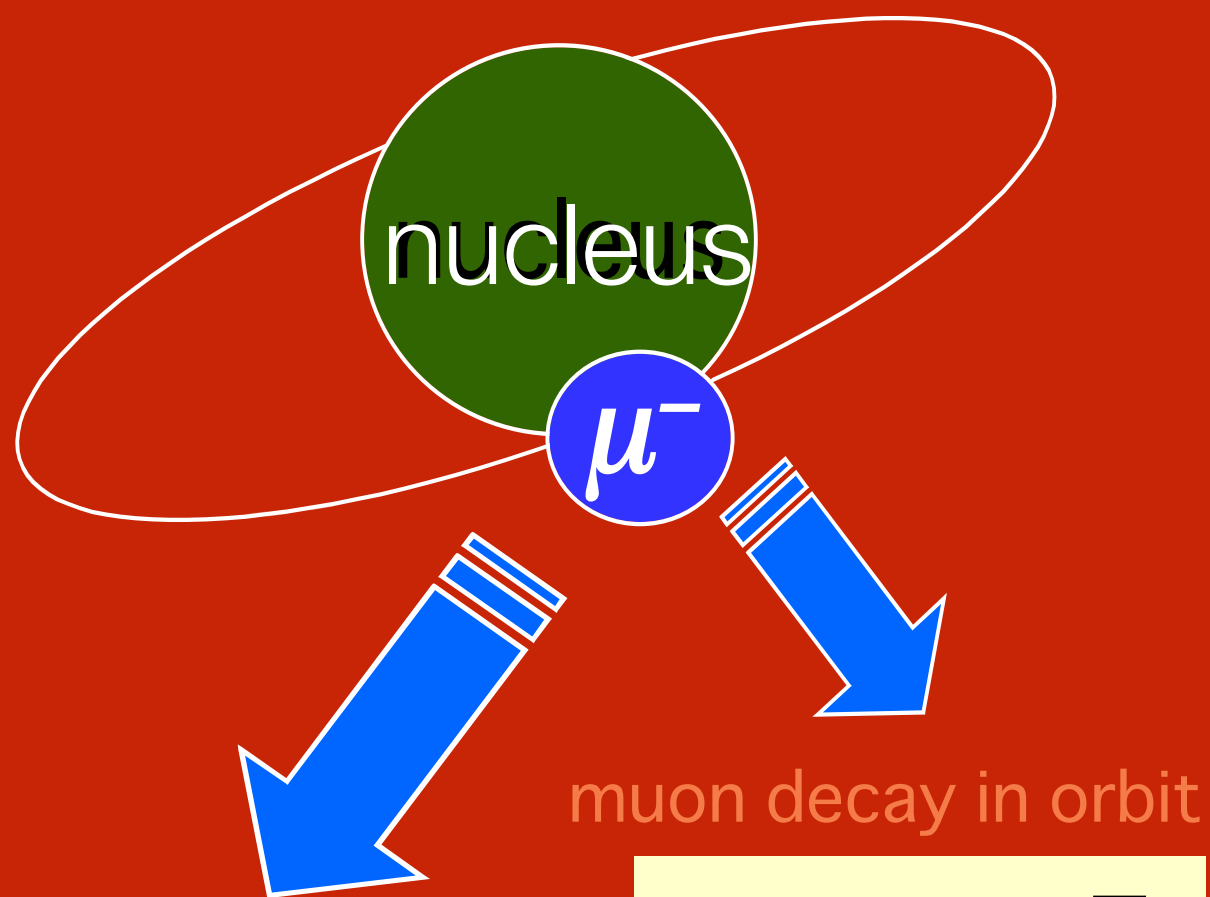
$$\mu^- \rightarrow e^- \nu \bar{\nu}$$

nuclear muon capture

$$\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1)$$

What is $\mu \rightarrow e$ Conversion ?

1s state in a muonic atom



$$\mu^- \rightarrow e^- \nu \bar{\nu}$$

nuclear muon capture

$$\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1)$$

Neutrino-less muon nuclear capture

$$\mu^- + (A, Z) \rightarrow e^- + (A, Z)$$

coherent process

$$\propto Z^5$$

(for the case that the final nucleus is the ground state.)

Event Signature :

a single mono-energetic electron of 105 MeV

$$\text{CR}(\mu^- \text{N} \rightarrow e^- \text{N}) \equiv \frac{\Gamma(\mu^- \text{N} \rightarrow e^- \text{N})}{\Gamma(\mu^- \text{N} \rightarrow \text{all})}$$

Backgrounds for μ -e conversion

intrinsic physics
backgrounds

Muon decay in orbit (DIO)
Radiative muon capture (RMC)
neutrons from muon nuclear capture
Protons from muon nuclear capture

Backgrounds for μ -e conversion

intrinsic physics
backgrounds

Muon decay in orbit (DIO)
Radiative muon capture (RMC)
neutrons from muon nuclear capture
Protons from muon nuclear capture

beam-related
backgrounds

Radiative pion capture (RPC)
Beam electrons
Muon decay in flights
Neutron background
Antiproton induced background

Backgrounds for μ -e conversion

intrinsic physics backgrounds

Muon decay in orbit (DIO)
Radiative muon capture (RMC)
neutrons from muon nuclear capture
Protons from muon nuclear capture

beam-related backgrounds

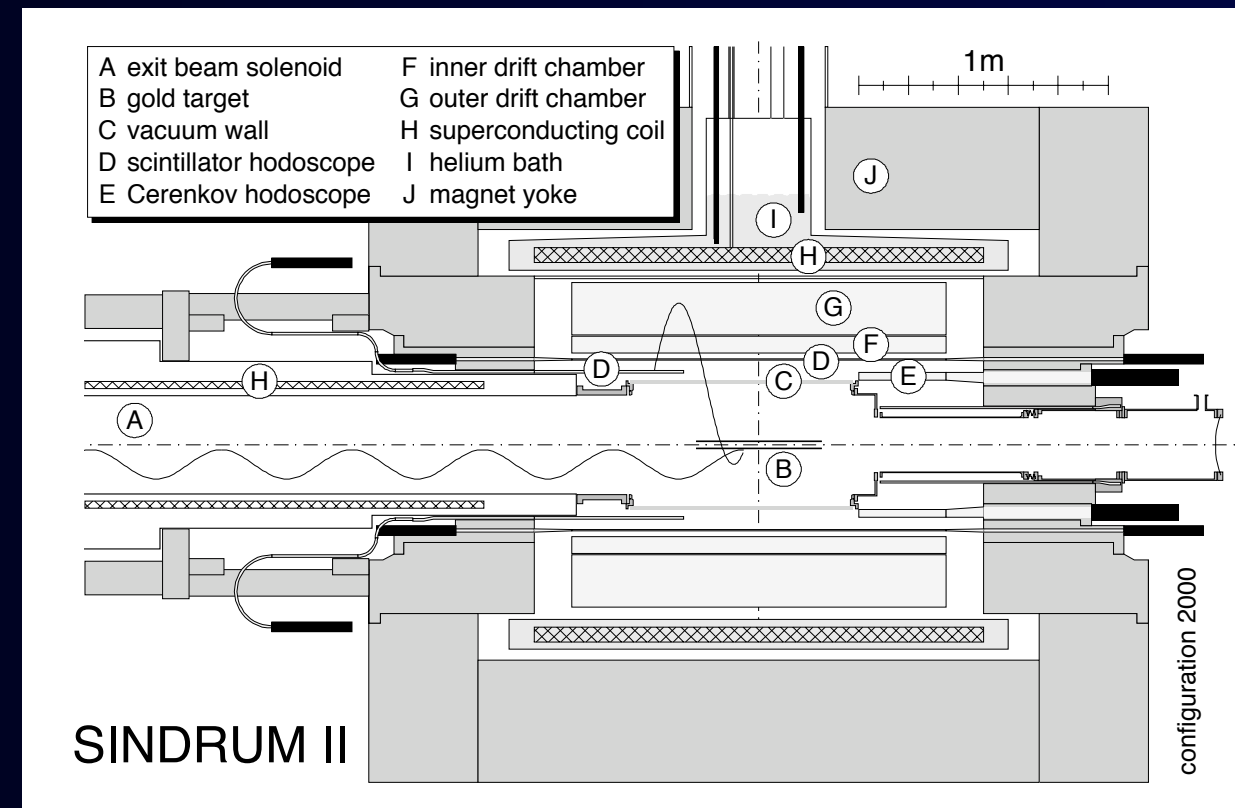
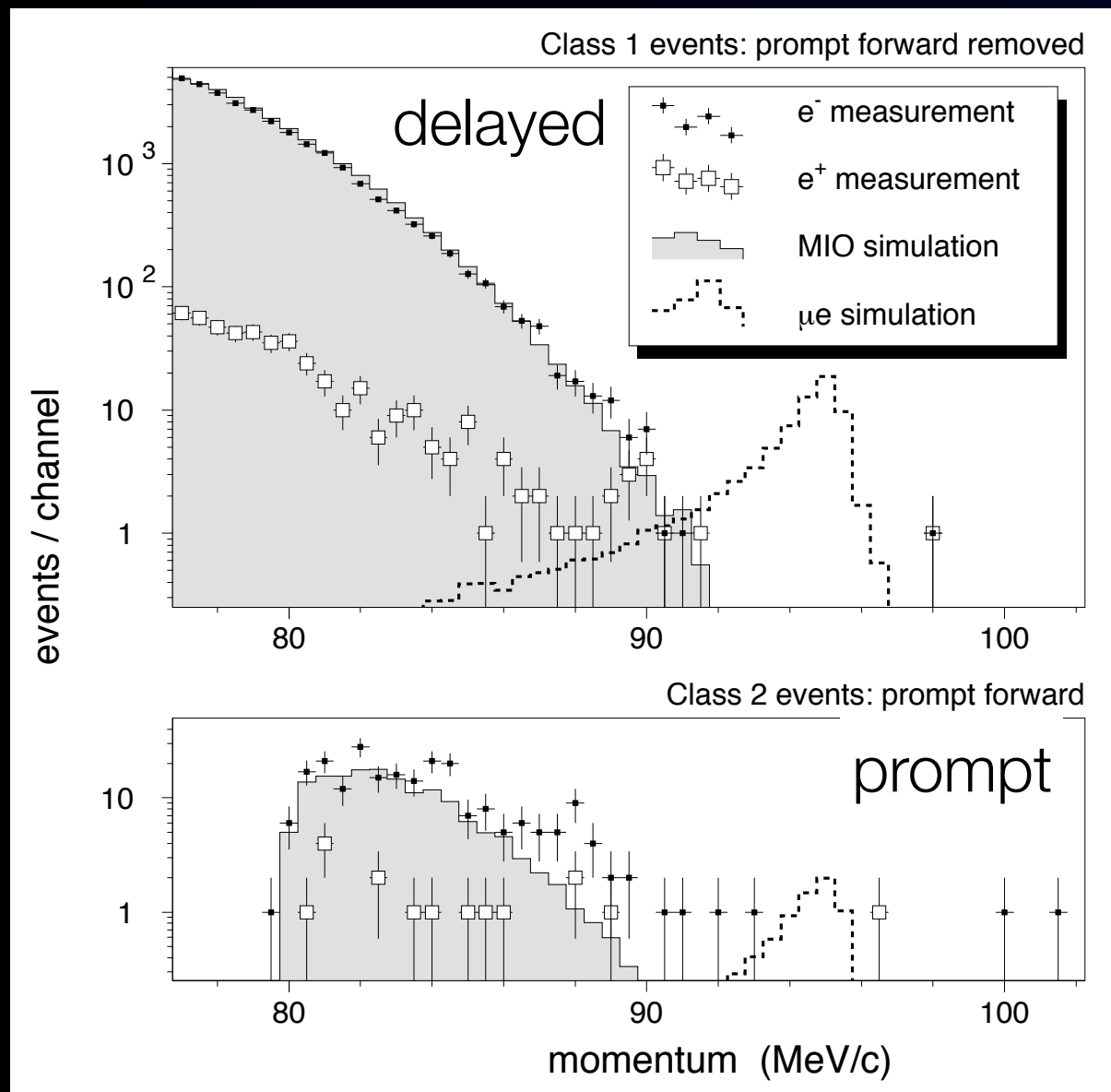
Radiative pion capture (RPC)
Beam electrons
Muon decay in flights
Neutron background
Antiproton induced background

cosmic-ray and other backgrounds

Cosmic-ray induced background
False tracking

Current Limits on $\mu \rightarrow e$ Conversion

SINDRUM-II (PSI)



	Z	S	CR limit
sulfur	16	0	7×10^{-11}
titanium	22	0,5/2,7/2	4.3×10^{-12}
copper	39	3/2	1.6×10^{-8}
gold	79	0,5/2	7×10^{-13}
lead	82	0 (1/2)	4.6×10^{-11}

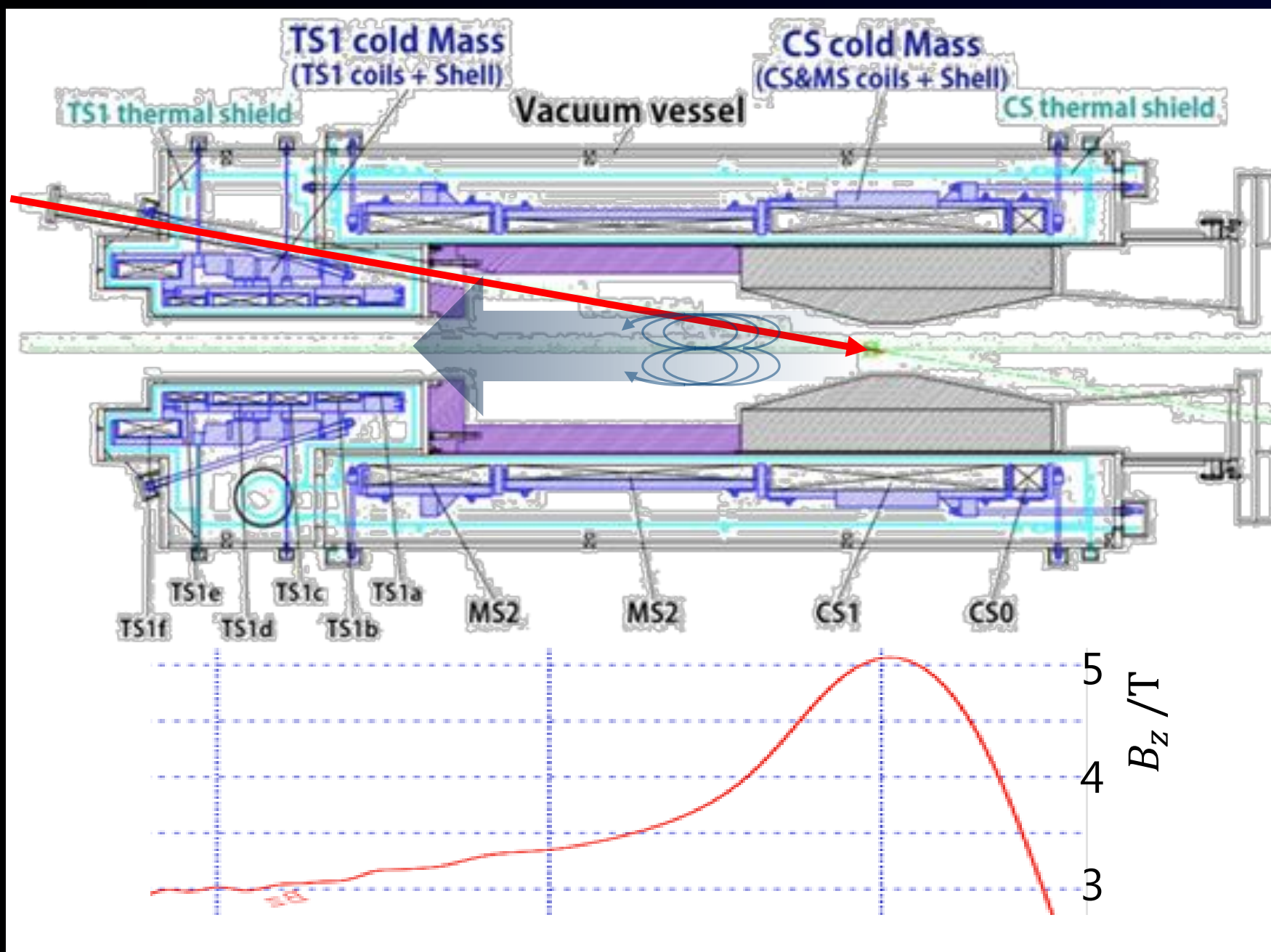
$$B(\mu^- + Au \rightarrow e^- + Au) < 7 \times 10^{-13}$$

In order to make a new-generation experiment to
search for μ -e conversion ...

$$B(\mu N \rightarrow e N) \leq 10^{-16}$$

Highly Intense Muon Source

pion capture in superconducting solenoids



- proton target in a solenoidal field (~ 5 T)
- a long proton target (1.5~2 interaction length) of heavy material)

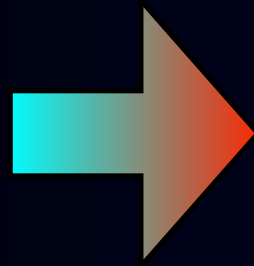
$2 \times 10^{11} \mu^-/\text{s}$
for 56 kW protons

muon yield
>1000 increase

Improvements for Background Rejection



Beam-related
backgrounds

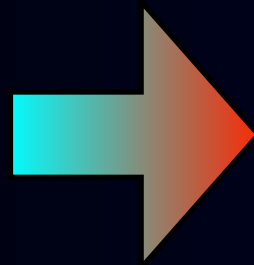


Beam pulsing with
separation of 1 μ sec

measured
between beam
pulses

proton extinction = #protons between pulses/#protons in a pulse $< 10^{-10}$

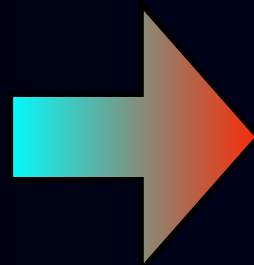
Muon DIO
background



low-mass trackers in
vacuum & thin target

improve
electron energy
resolution

Muon DIF
background

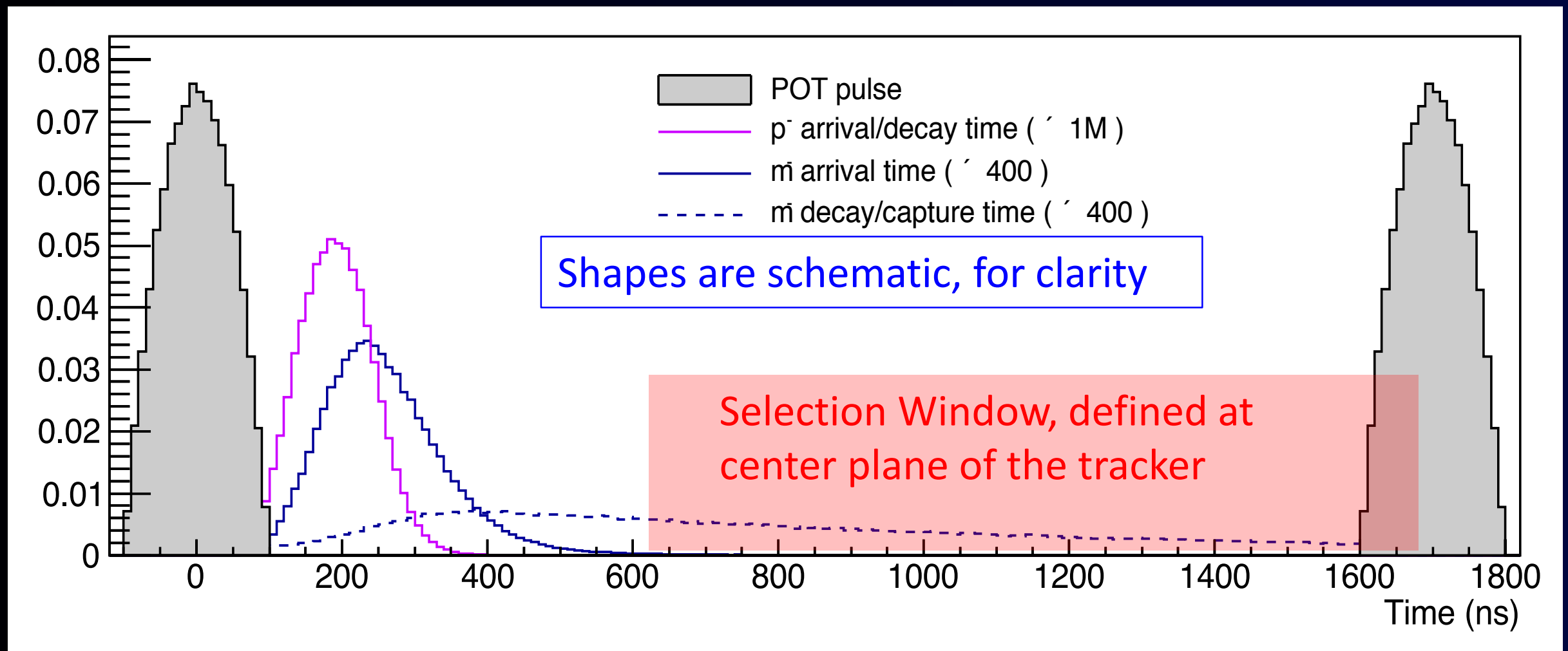


curved solenoids for
momentum selection

eliminate
energetic muons
(>75 MeV/c)

based on the MELC proposal at Moscow Meson Factory

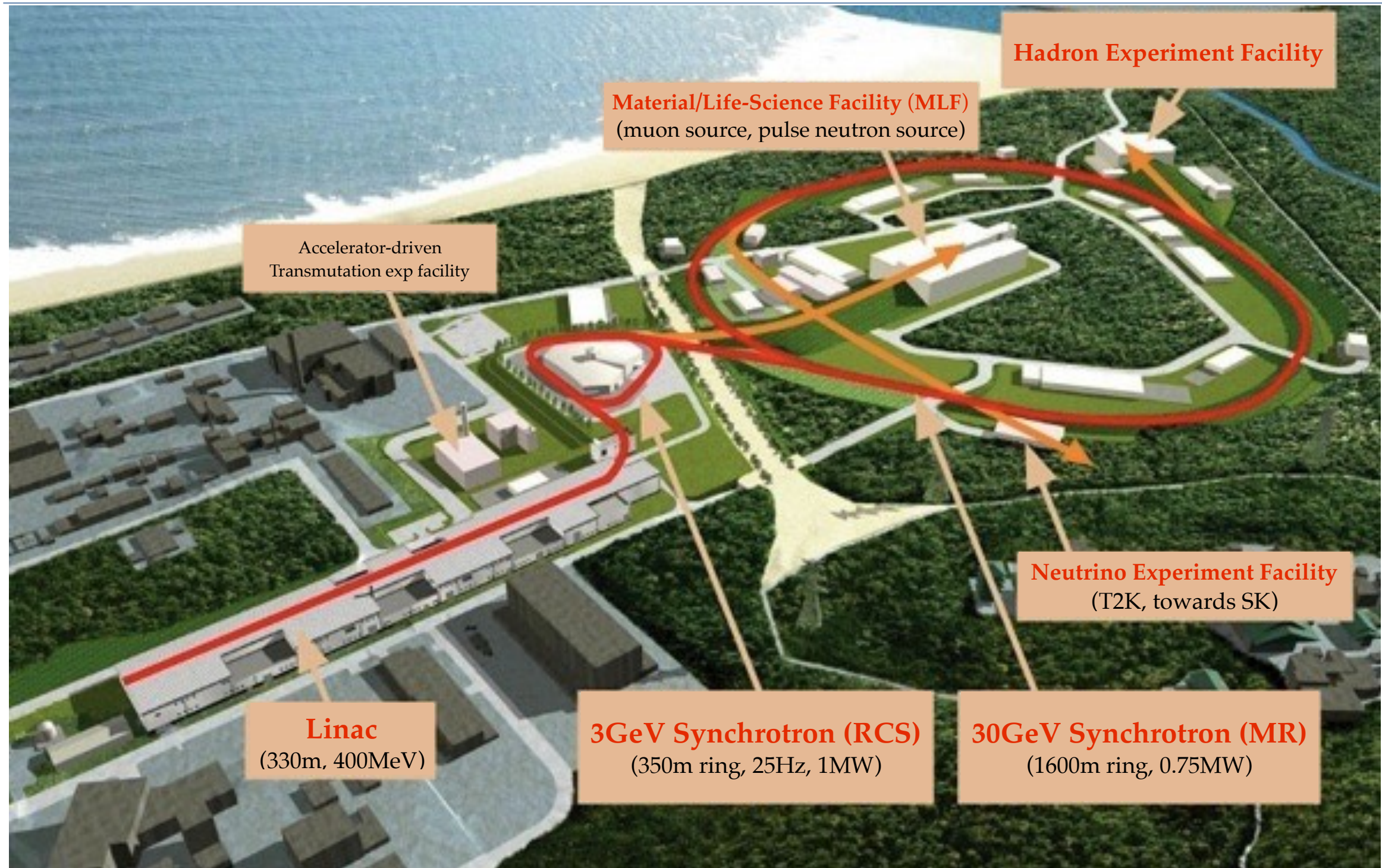
Pulsed Muon Beam for μ -e conversion



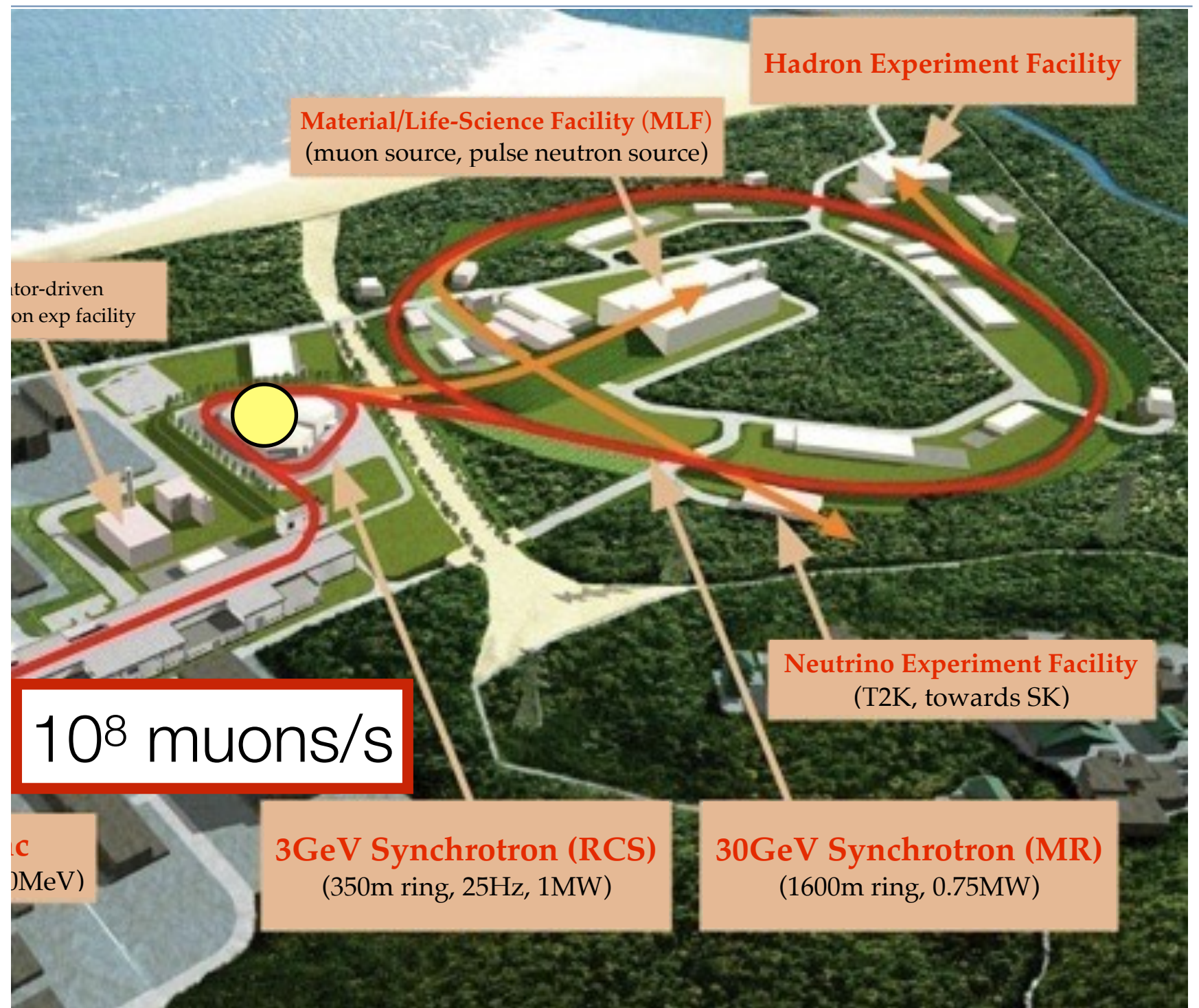
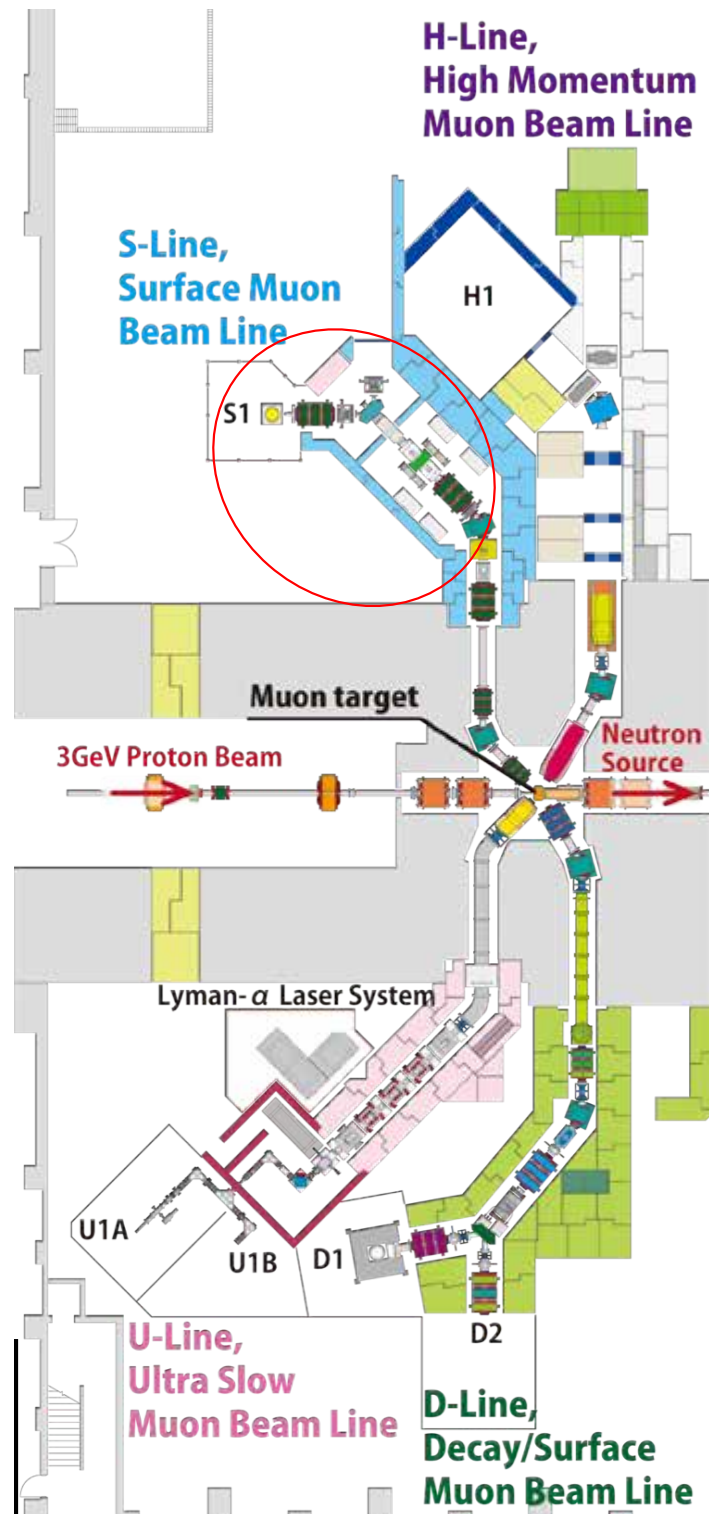
Selection windows turns on late, significantly reducing prompt backgrounds.

Requirements ;
Proton extinction factor during pulses $< 10^{-10}$

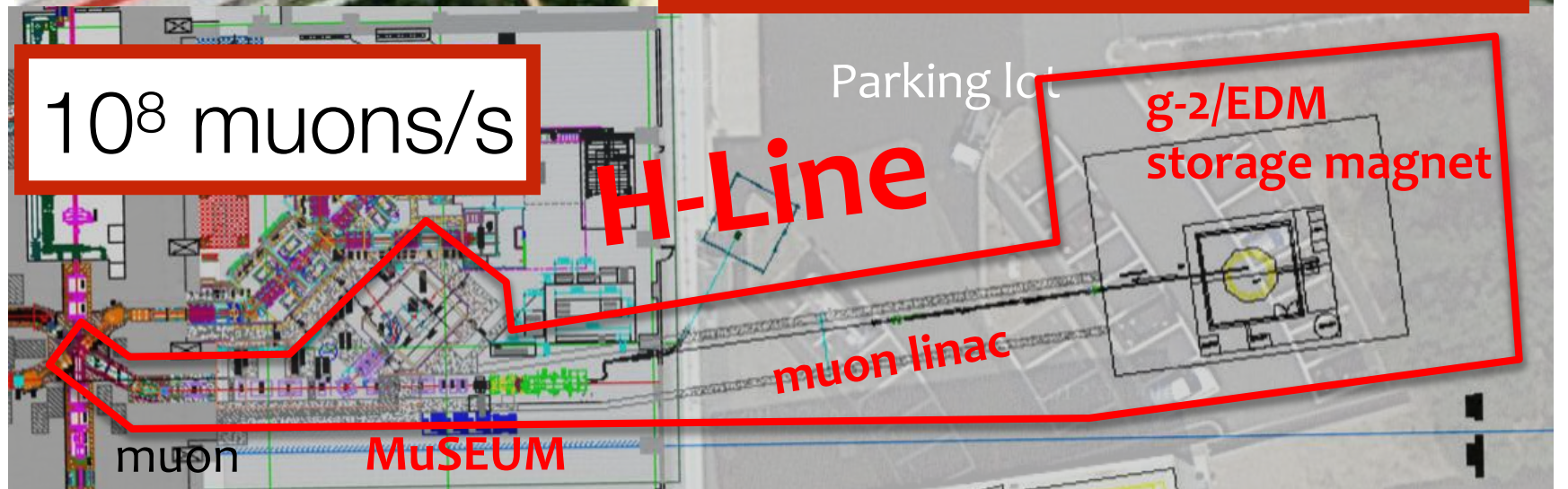
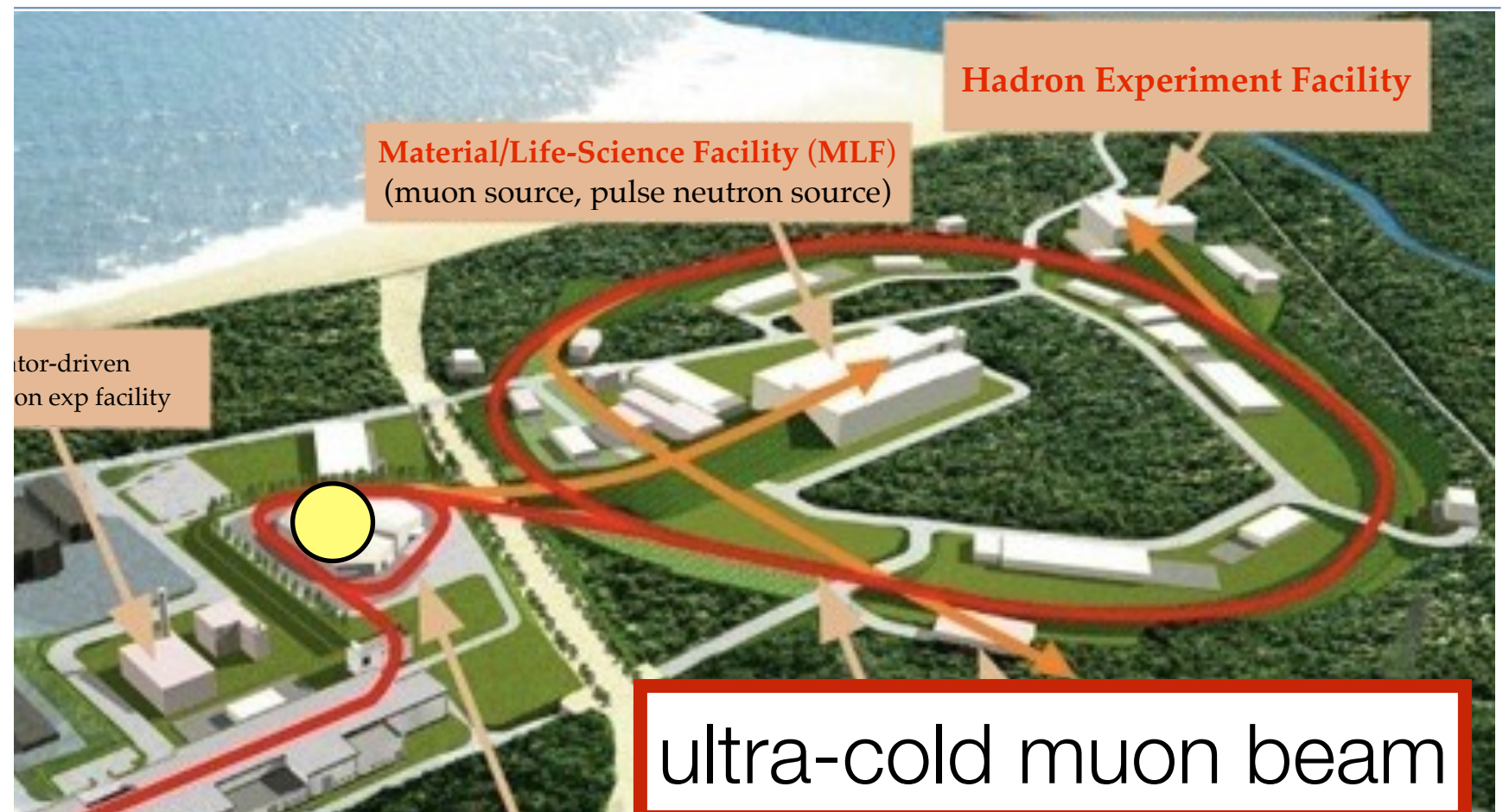
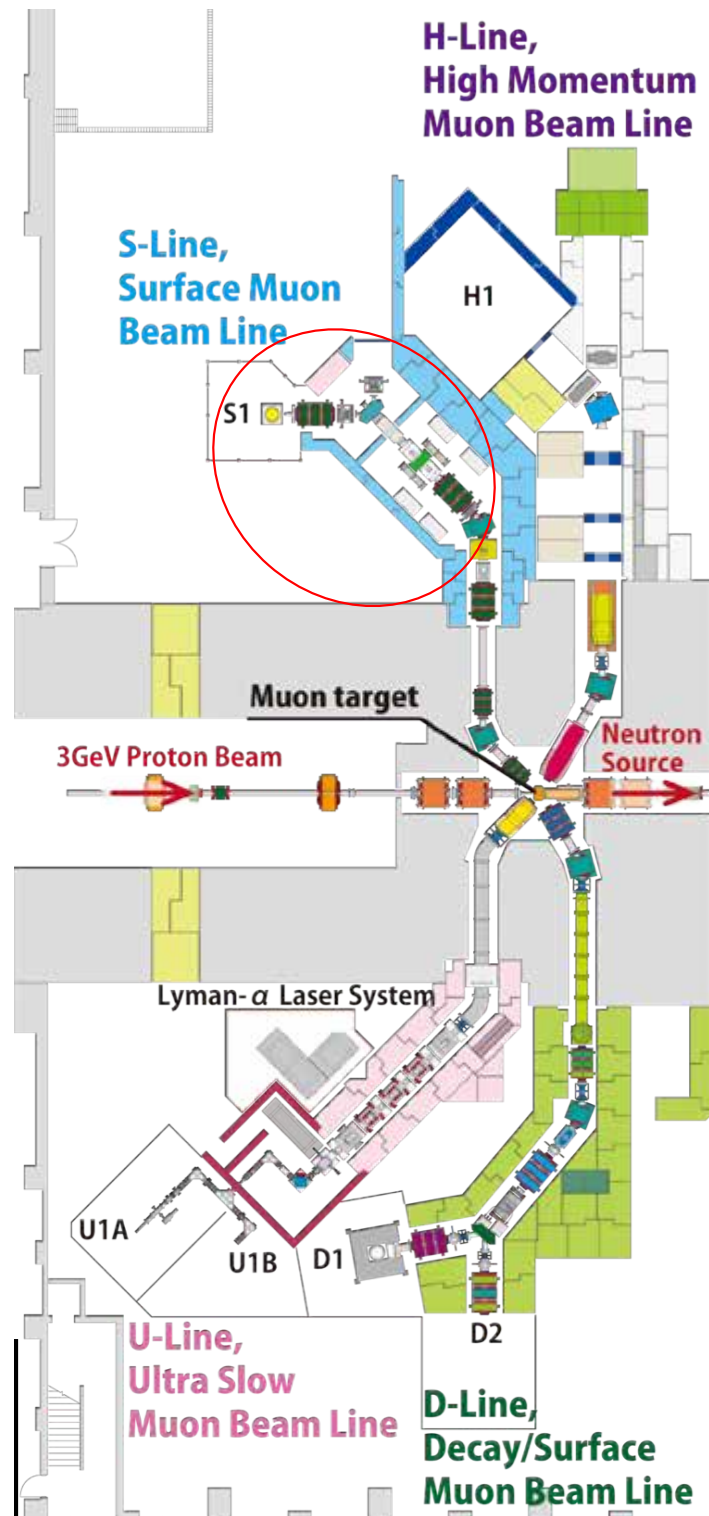
J-PARC (MUSE@MLF)



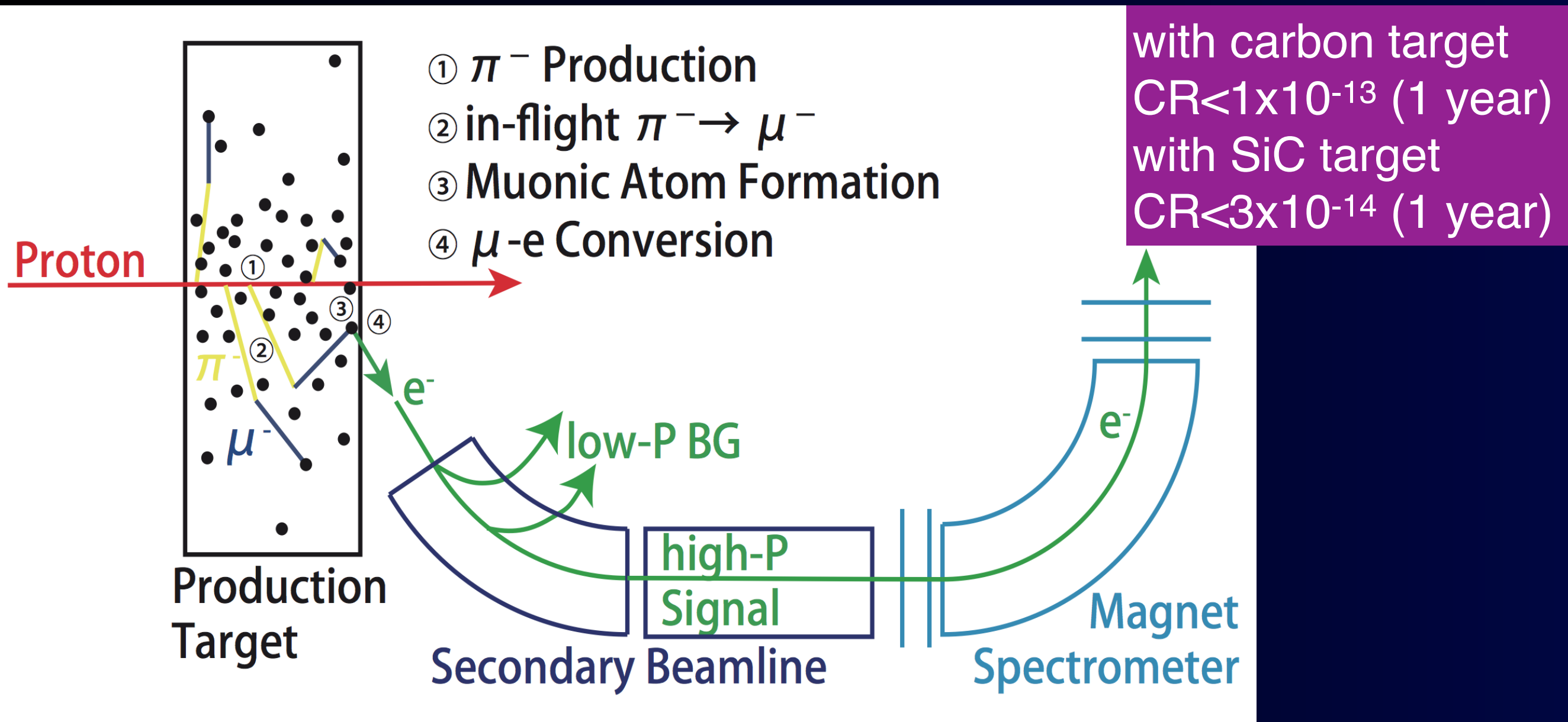
J-PARC (MUSE@MLF)



J-PARC (MUSE@MLF)



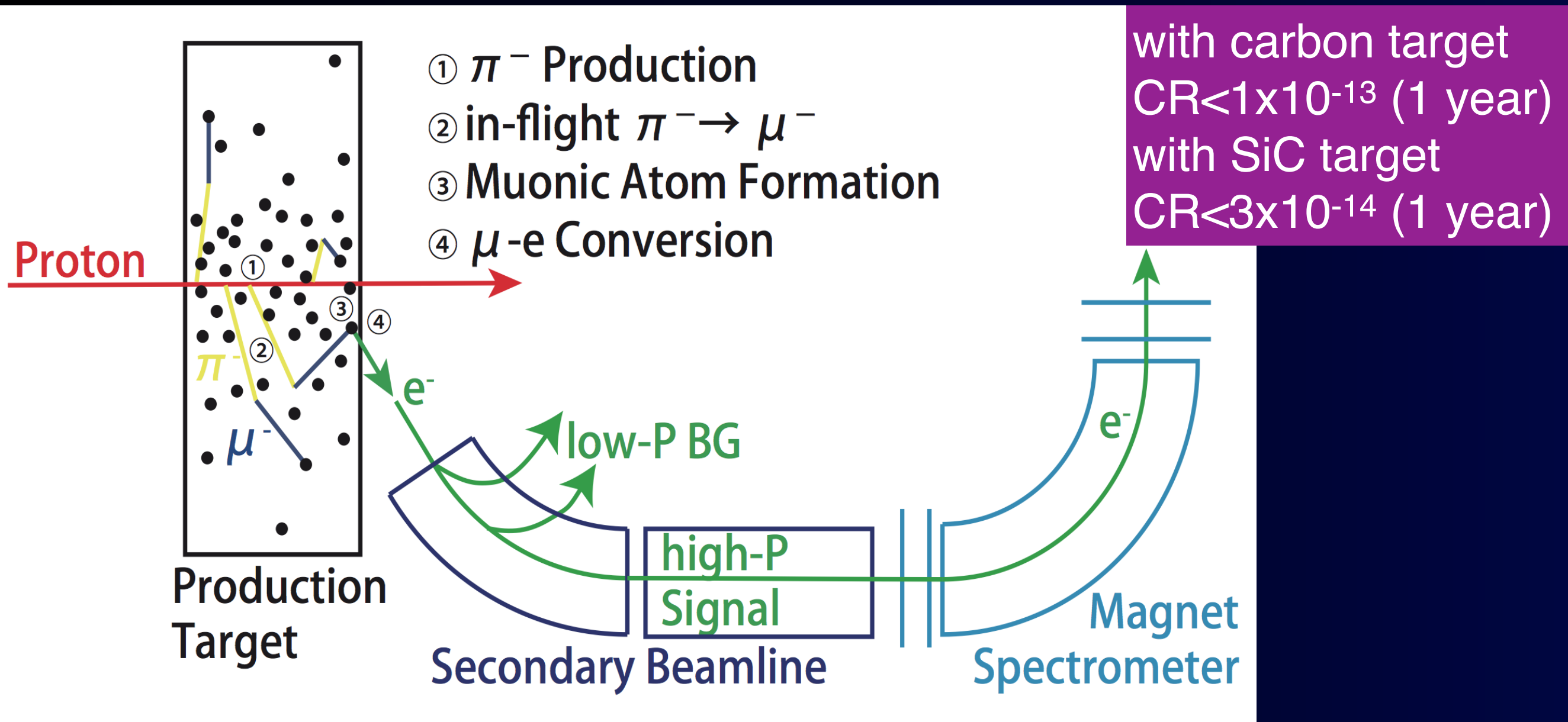
DeeMe at J-PARC MLF



3 GeV J-PARC RCS protons

- beamline and spectrometer to select 100 MeV e^-
- 4 MWPCs with $\Delta p = 0.5$ MeV

DeeMe at J-PARC MLF

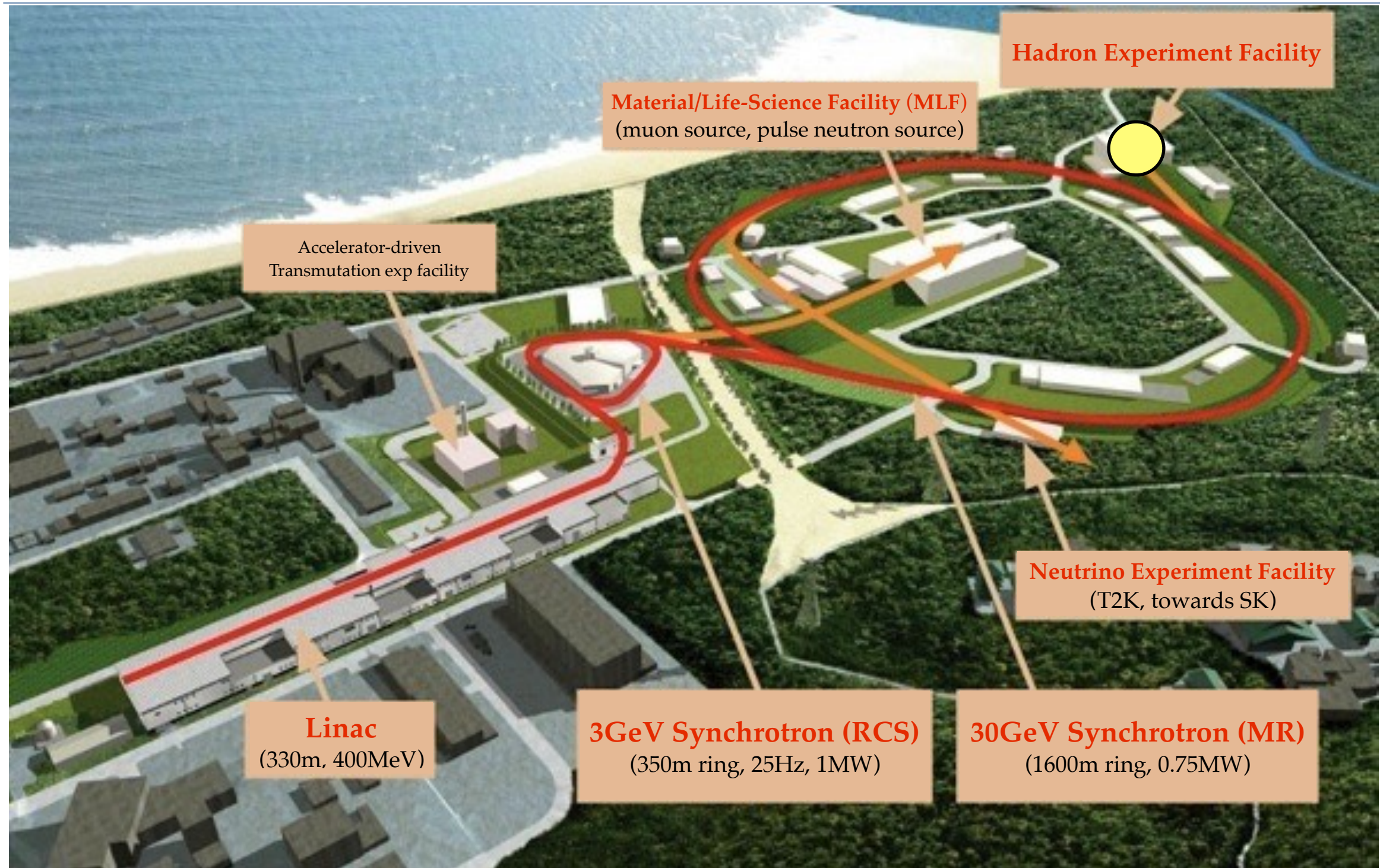


3 GeV J-PARC RCS protons

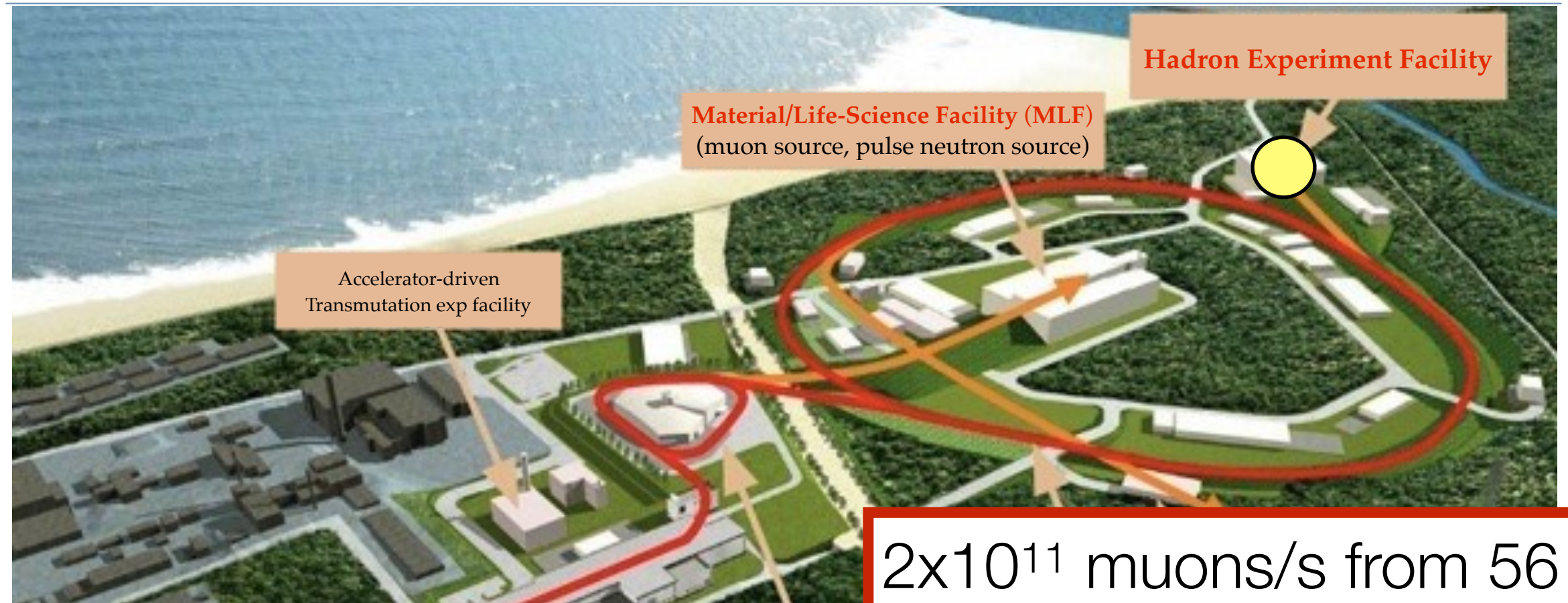
- beamline and spectrometer to select 100 MeV e^-
- 4 MWPCs with $\Delta p = 0.5$ MeV

N. Teshima on Thursday

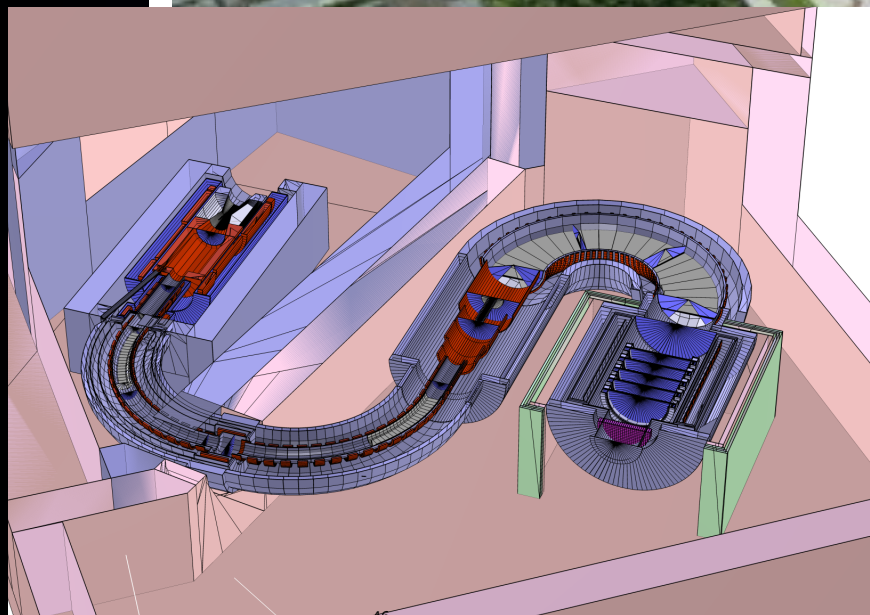
J-PARC (COMET@Main Ring)



J-PARC (COMET@Main Ring)



2×10^{11} muons/s from 56 kW



3GeV Synchrotron
(350m ring, 250 MHz)



COMET = COherent Muon to Electron Transition



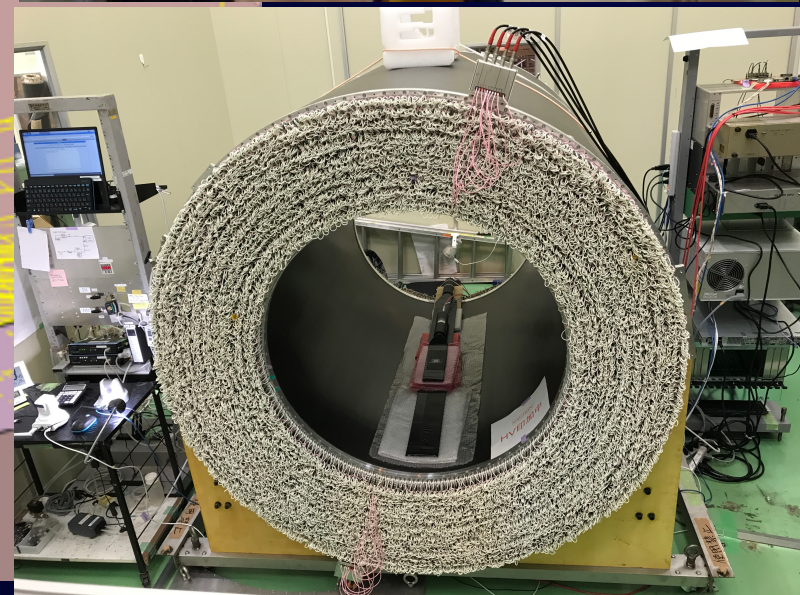
COMET Phase-I : J-PARC E21

Phase-I

proton beam power = 3.2 kW

Single event sensitivity : 2×10^{-15}
a factor of 100 improvement
Running time: 0.4 years (1.2×10^7 s)

proton
beam

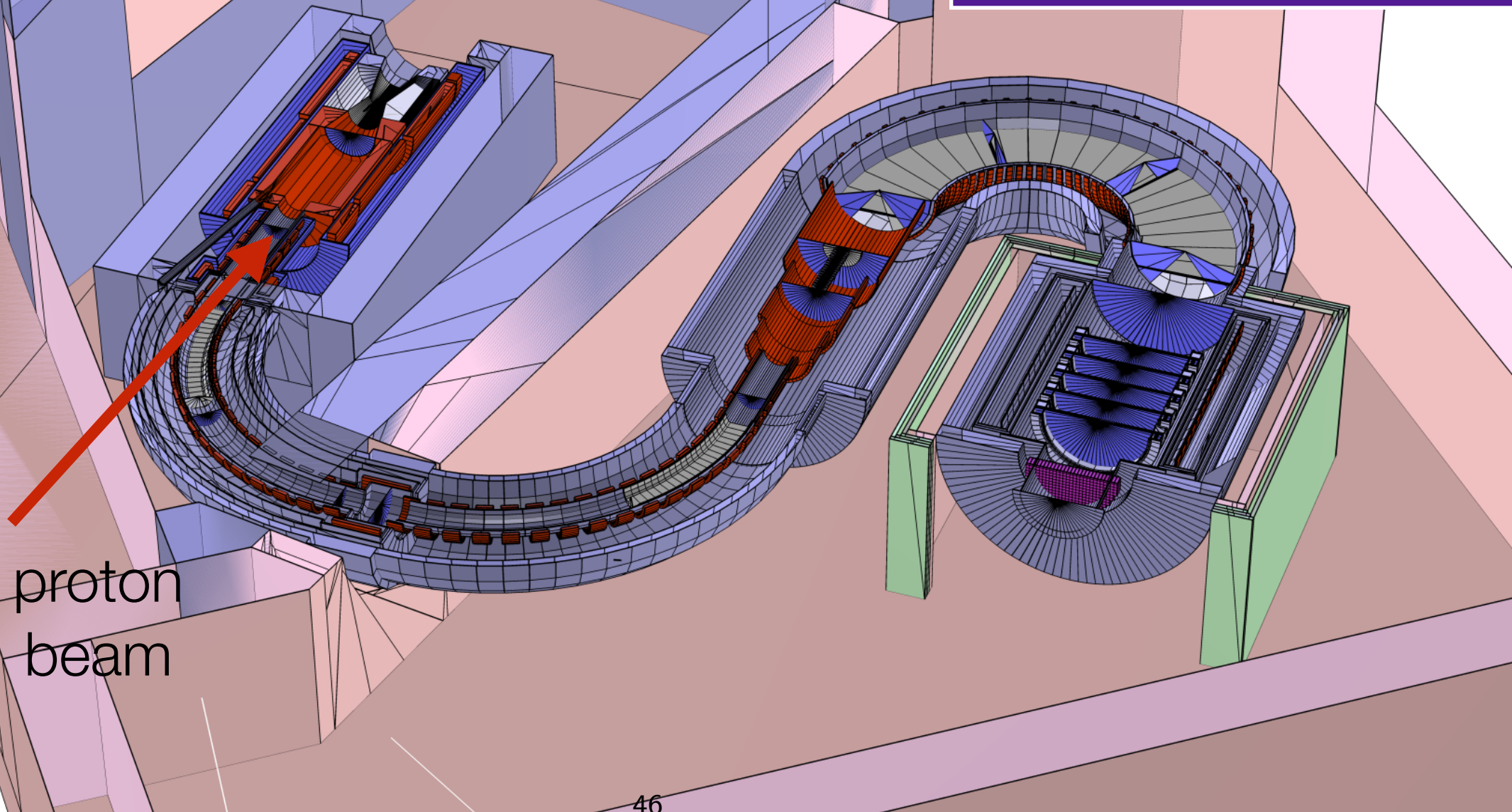


COMET Phase-II : J-PARC E21

Phase-II

proton beam power = 56 kW

Single event sensitivity : 2.6×10^{-17}
 a factor of 10,000 improvement
 Running time: 1 years (2×10^7 sec)

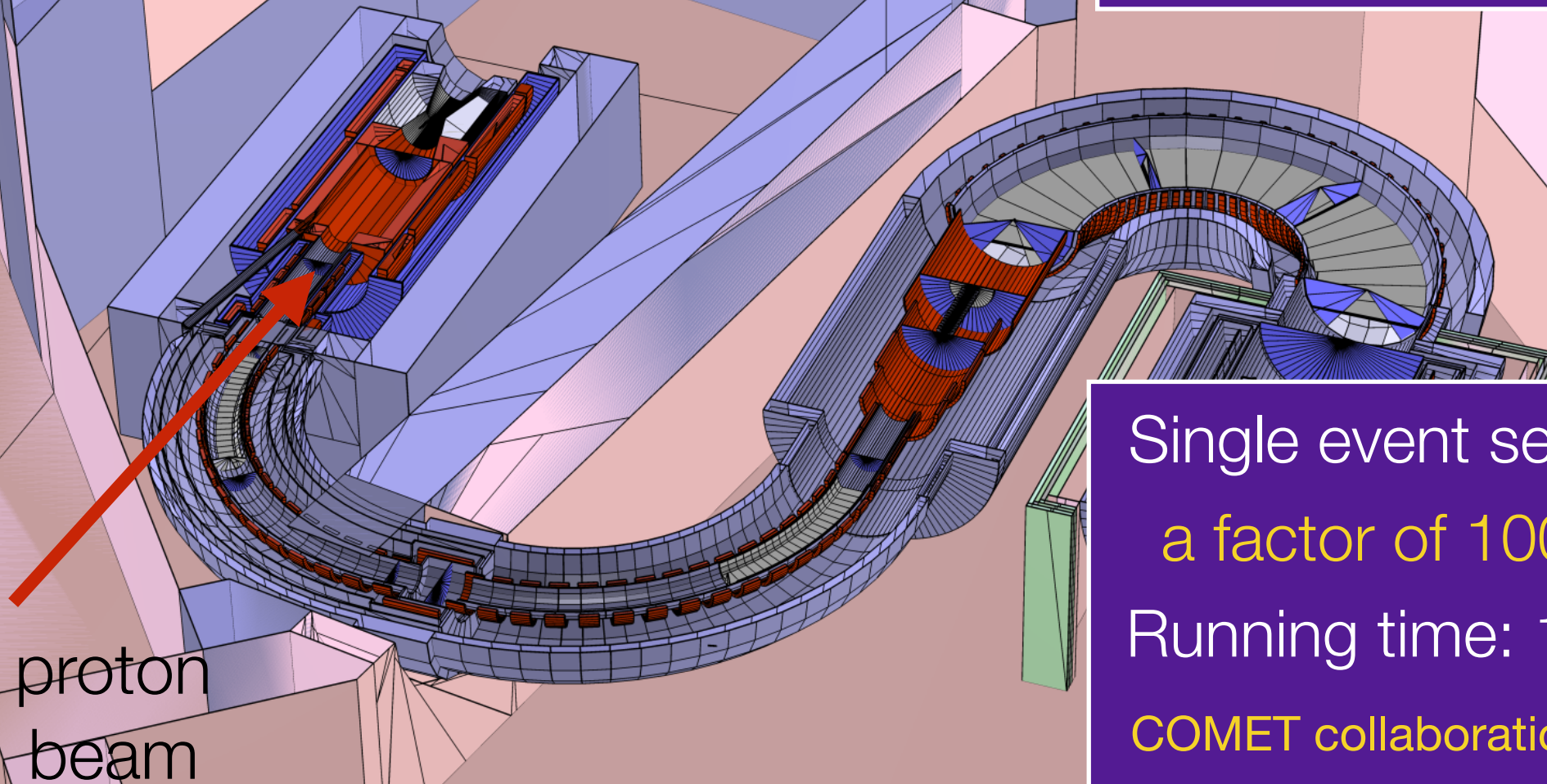


COMET Phase-II : J-PARC E21

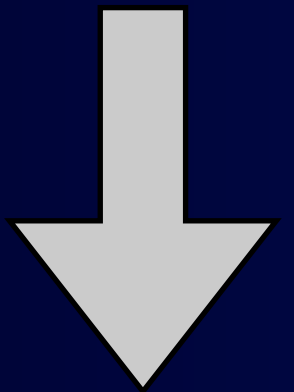
Phase-II

proton beam power = 56 kW

Single event sensitivity : 2.6×10^{-17}
 a factor of 10,000 improvement
 Running time: 1 years (2×10^7 sec)



Single event sensitivity : $O(10^{-18})$
 a factor of 100,000 improvement
 Running time: 1 years (2×10^7 sec)
 COMET collaboration, arXiv:1812.07824, 2018

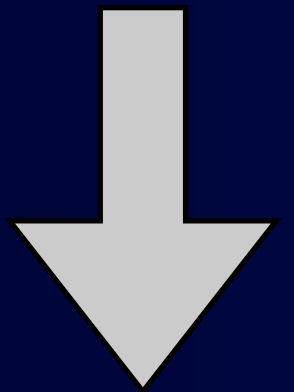
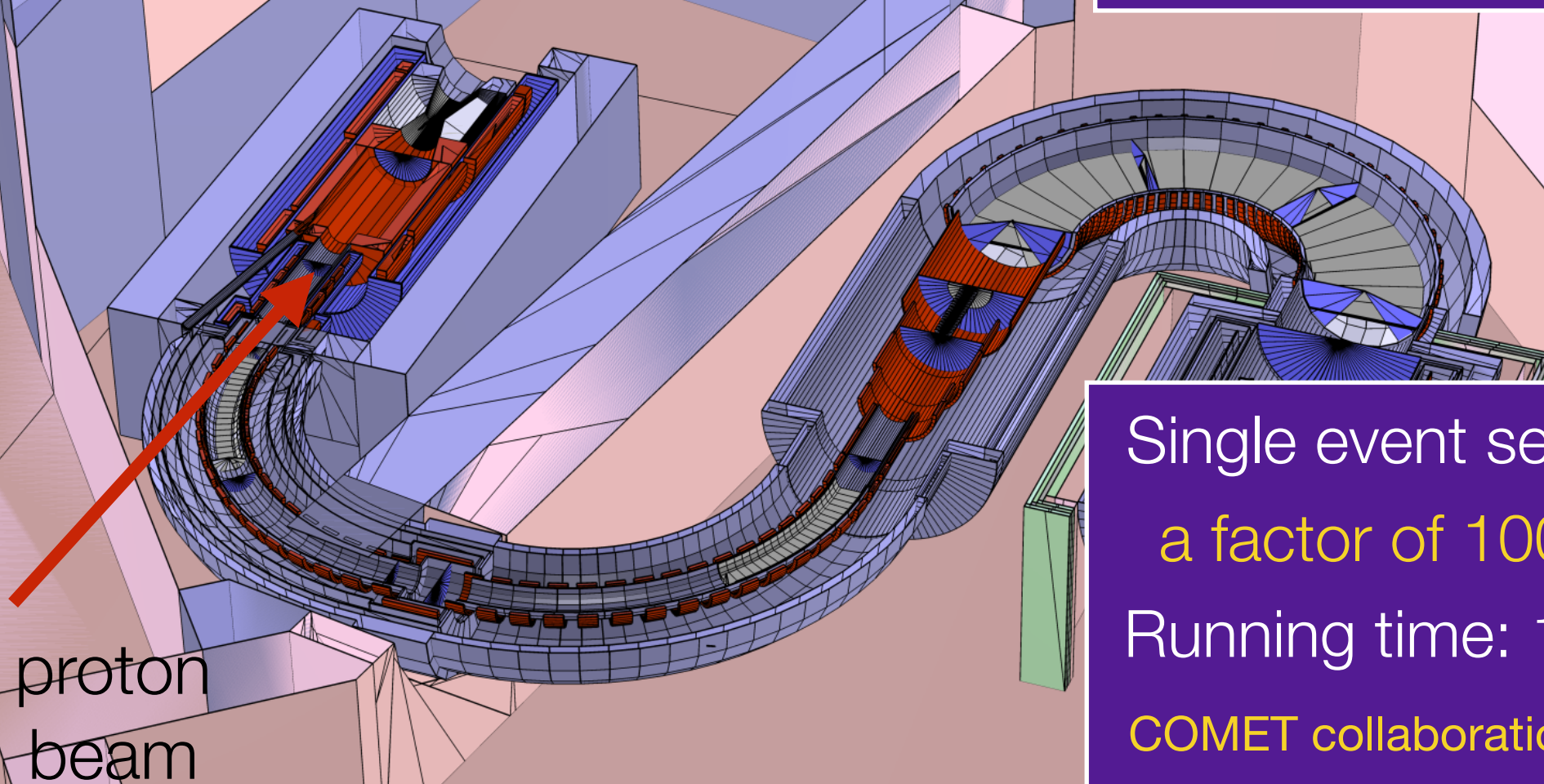


COMET Phase-II : J-PARC E21

Phase-II

proton beam power = 56 kW

Single event sensitivity : 2.6×10^{-17}
a factor of 10,000 improvement
Running time: 1 years (2×10^7 sec)



Single event sensitivity : $O(10^{-18})$
a factor of 100,000 improvement
Running time: 1 years (2×10^7 sec)
COMET collaboration, arXiv:1812.07824, 2018

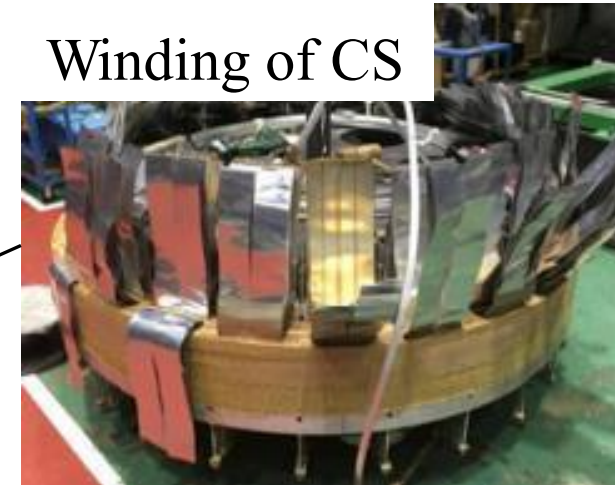
T. Xing on Thursday

COMET Phase-I : Preparation

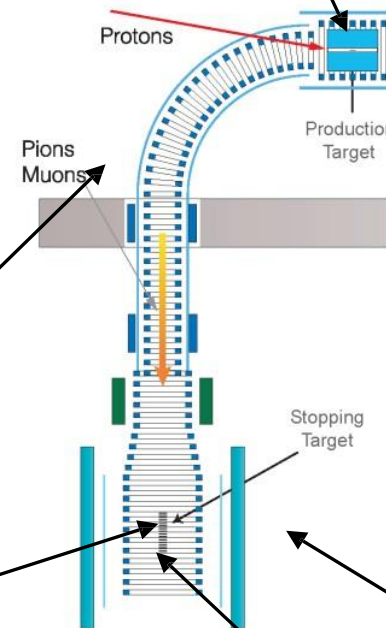
B line in Jan 2018
To be finished in 2019



Winding of CS



Transportation solenoid
Finished in 2015



Inside the experimental hall



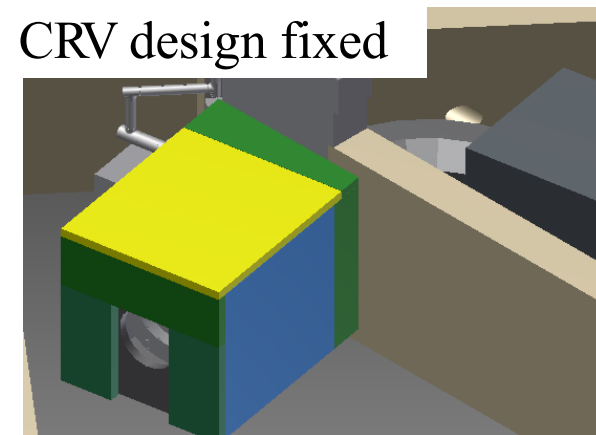
Prototype for stopping target



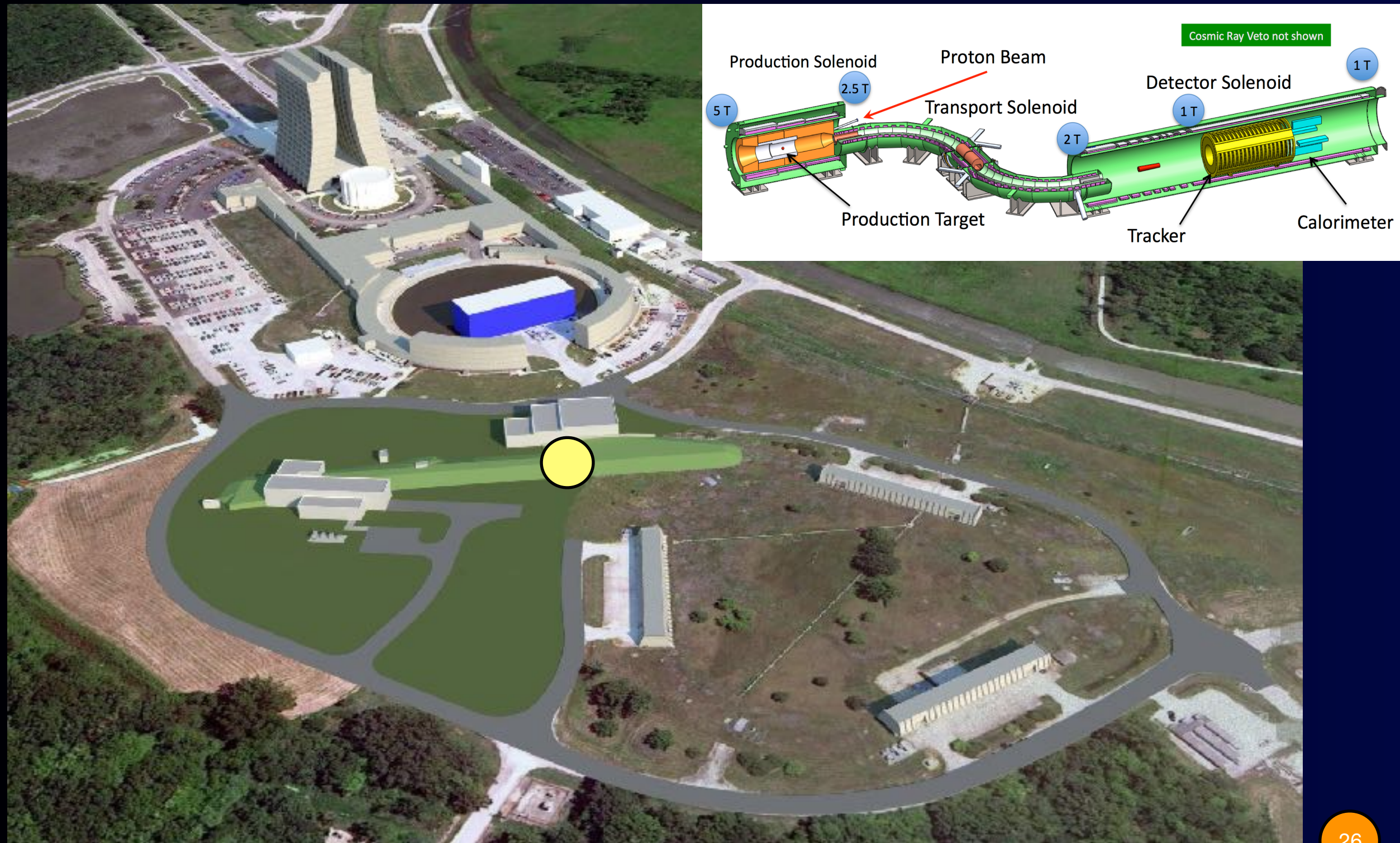
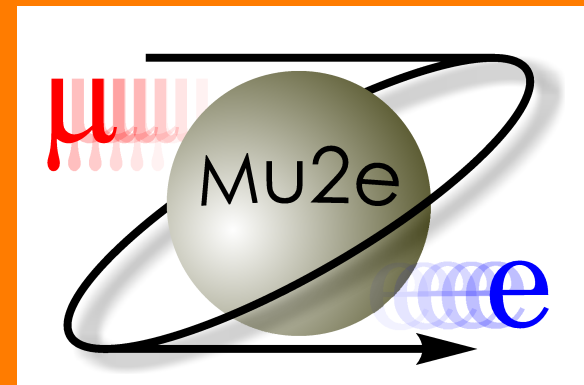
Detector solenoid



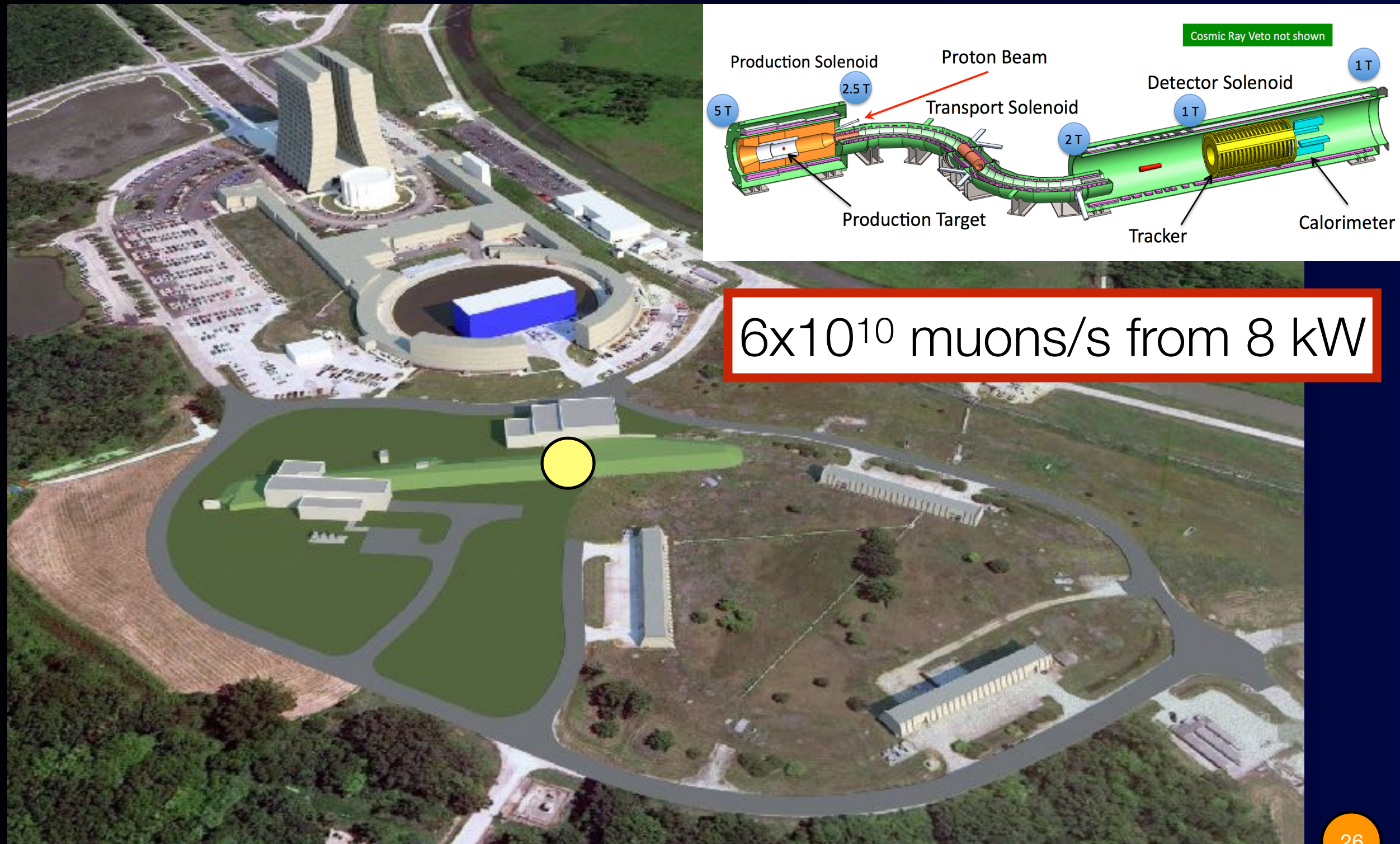
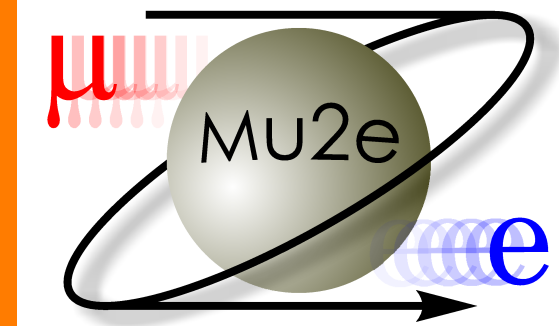
CRV design fixed



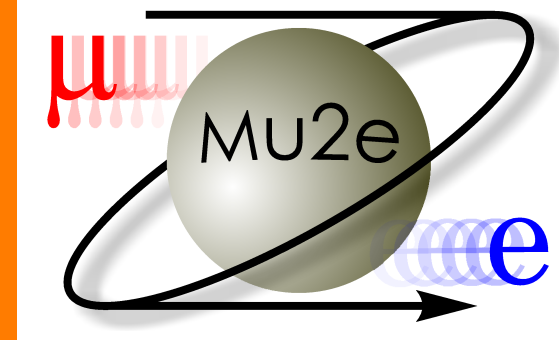
Fermilab Muon Campus



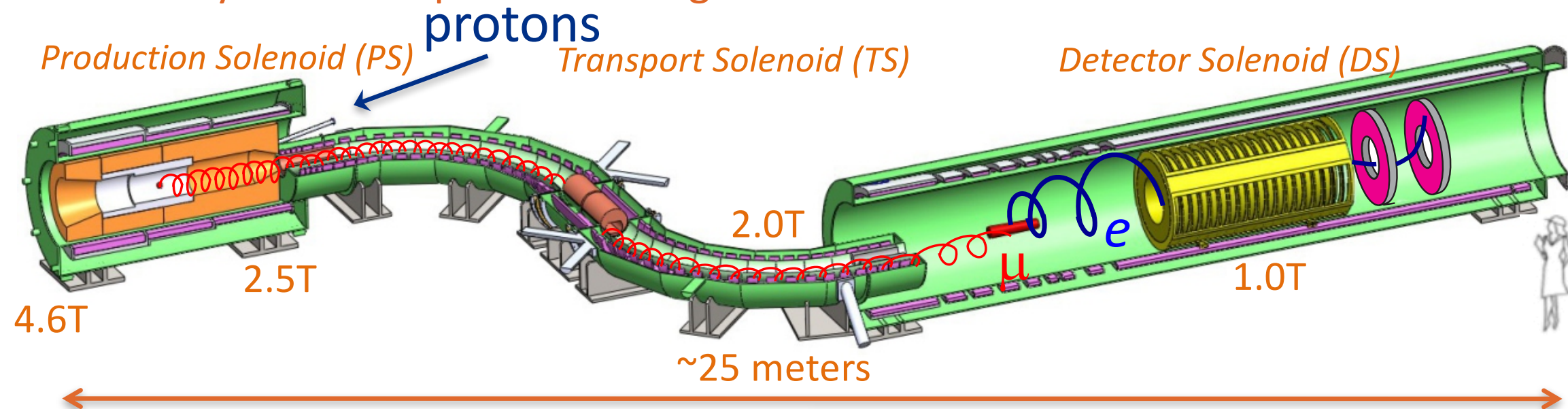
Fermilab Muon Campus



Mu2e at Fermilab



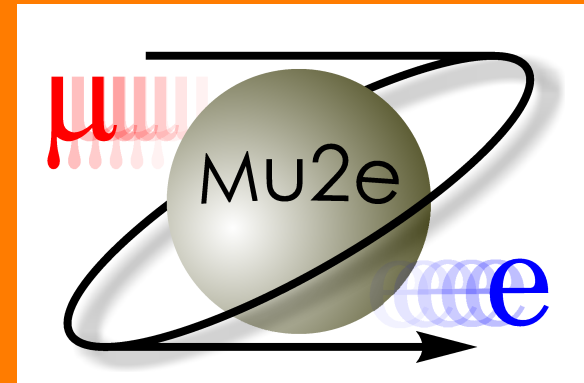
A System of superconducting solenoids and an intense muon beam



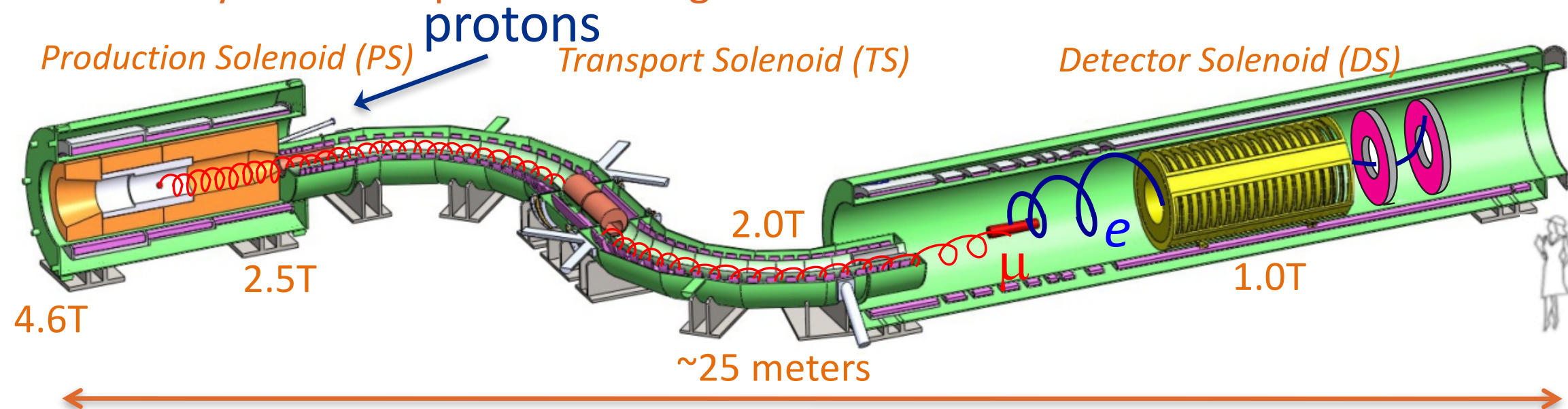
- A search for Charged Lepton Flavor Violation: $\mu N \rightarrow e N$
 - Expected sensitivity of 6×10^{-17} @ 90% CL, x10,000 better than SINDRUM-II
 - Probes effective new physics mass scales up to 10^4 TeV/ c^2
 - *Discovery* sensitivity to broad swath of NP parameter space

- Experiment scope includes
 - Proton Beam line
 - Solenoid systems
 - Detector elements
(tracker, calorimeter, cosmic veto, DAQ, beam monitoring)
 - Experimental hall
 - Commissioning begins in 2022

Mu2e at Fermilab



A System of superconducting solenoids and an intense muon beam



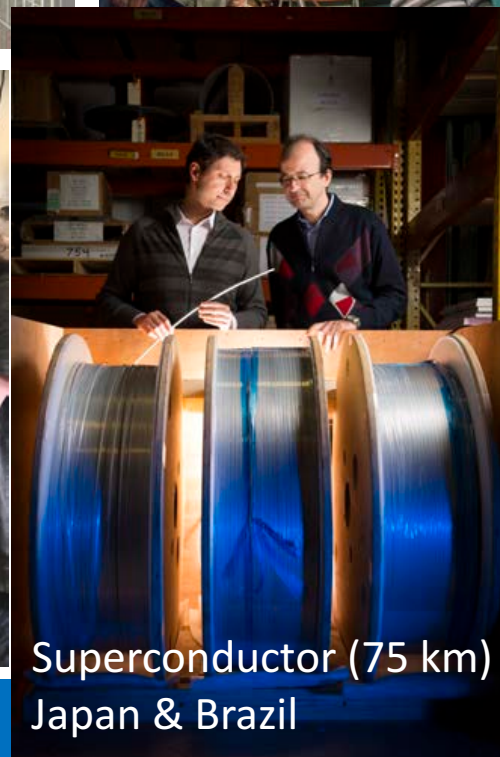
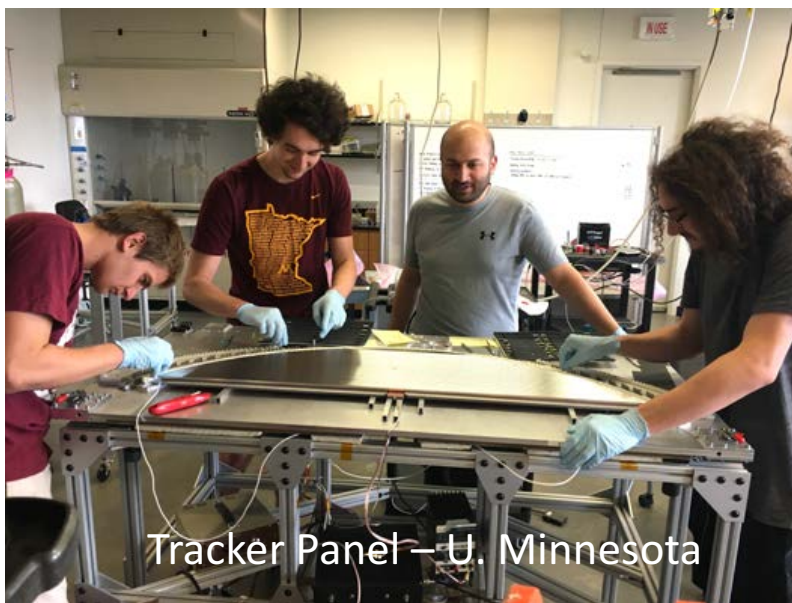
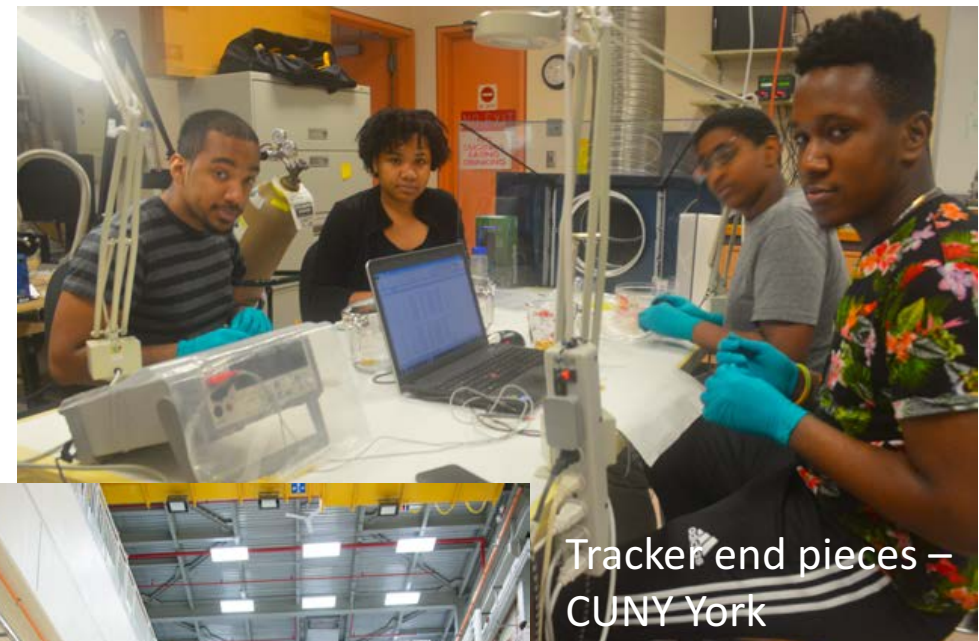
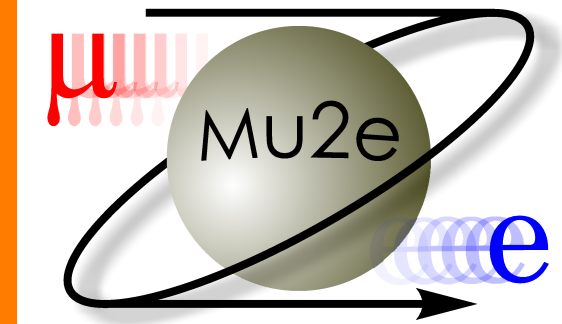
- A search for Charged Lepton Flavor Violation: $\mu N \rightarrow e N$
 - Expected sensitivity of 6×10^{-17} @ 90% CL, x10,000 better than SINDRUM-II
 - Probes effective new physics mass scales up to $10^4 \text{ TeV}/c^2$
 - *Discovery* sensitivity to broad swath of NP parameter space

- Experiment scope includes
 - Proton Beam line
 - Solenoid systems
 - Detector elements (tracker, calorimeter, cosmic veto, DAQ, beam monitoring)
 - Experimental hall

Commissioning begins in 2022

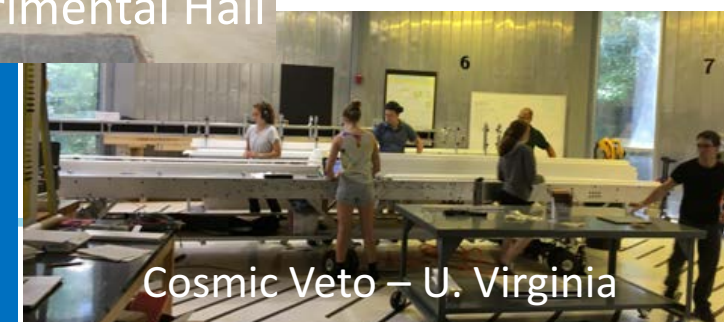
R. Bonvetre on Thursday

Mu2e at Fermilab

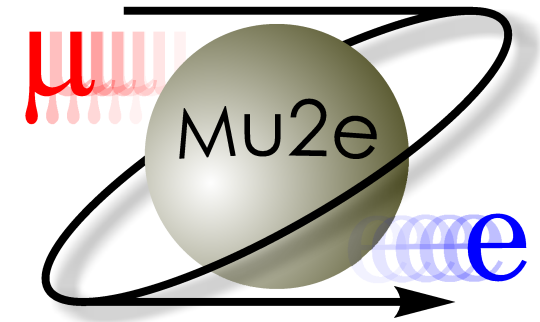


Completed:

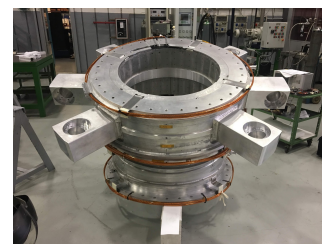
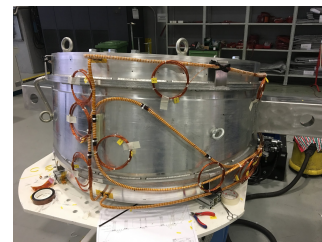
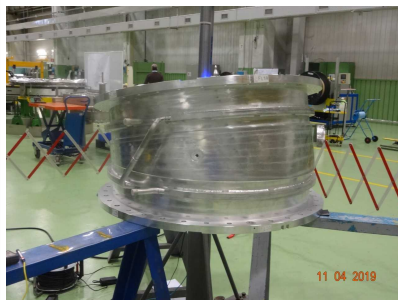
- 75 km of superconductor
- 30 km of extruded scintillator
- 25 km of aluminized mylar straws



Mu2e at Fermilab



Transport Solenoid:
Testing & assembly
in progress!



Production Solenoid - coil windings
have begun for the first coil!

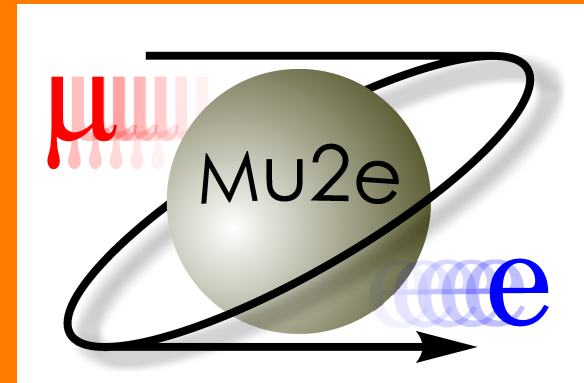
Progress across sub-systems:

- 68% of Calorimeter CsI crystals received & tested
- 37% of Cosmic Veto di-counters built & tested
- 75% of pre-production Tracker panels built & tested
- Delivery Ring installed and working
- M4 Beam line installation in progress



Delivery Ring installation complete
(in operation for Muon g-2 experiment)

Mu2e-II - a next generation $\mu \rightarrow e$ conversion experiment at FNAL



Expression of Interest for Evolution of the Mu2e Experiment[†]

F. Abusalma²³, D. Ambrose²³, A. Artikov⁷, R. Bernstein⁸, G.C. Blazey²⁷, C. Bloise⁹, S. Boi³³, T. Bolton¹⁴, J. Bono⁸, R. Bonventre¹⁶, D. Bowring⁸, D. Brown¹⁶, D. Brown²⁰, K. Byrum¹, M. Campbell²², J.-F. Caron¹², F. Cervelli³⁰, D. Chokheli⁷, K. Ciampa²³, R. Ciolini³⁰, R. Coleman⁸, D. Cronin-Hennessy²³, R. Culbertson⁸, M.A. Cummings²⁵, A. Daniel¹², Y. Davydov⁷, S. Demers³⁵, D. Denisov⁸, S. Denisov¹³, S. Di Falco³⁰, E. Diociaiuti⁹, R. Djilkibaev²⁴, S. Donati³⁰, R. Donghia⁹, G. Drake¹, E.C. Dukes³³, B. Echenard⁵, A. Edmonds¹⁶, R. Ehrlich³³, V. Evdokimov¹³, P. Fabbriatore¹⁰, A. Ferrari¹¹, M. Frank³², A. Gaponenko⁸, C. Gatto²⁶, Z. Giorgio¹⁷, S. Giovannella⁹, V. Giusti³⁰, H. Glass⁸, D. Glenzinski⁸, L. Goodenough¹, C. Group³³, F. Happacher⁹, L. Harkness-Brennan¹⁹, D. Hedin²⁷, K. Heller²³, D. Hitlin⁵, A. Hocker⁸, R. Hooper¹⁸, G. Horton-Smith¹⁴, C. Hu⁵, P.Q. Hung³³, E. Hungerford¹², M. Jenkins³², M. Jones³¹, M. Kargiantoulakis⁸, K. S. Khaw³⁴, B. Kiburg⁸, Y. Kolomensky^{3,16}, J. Kozminski¹⁸, R. Kutschke⁸, M. Lancaster¹⁵, D. Lin⁵, I. Logashenko²⁹, V. Lombardo⁸, A. Luca⁸, G. Lukicov¹⁵, K. Lynch⁶, M. Martini²¹, A. Mazzacane⁸, J. Miller², S. Miscetti⁹, L. Morescalchi³⁰, J. Mott², S. E. Mueller¹¹, P. Murat⁸, V. Nagaslaev⁸, D. Neuffer⁸, Y. Oksuzian³³, D. Pasciuto³⁰, E. Pedreschi³⁰, G. Pezzullo³⁵, A. Pla-Dalmau⁸, B. Pollack²⁸, A. Popov¹³, J. Popp⁶, F. Porter⁵, E. Prebys⁴, V. Pronskikh⁸, D. Pushka⁸, J. Quirk², G. Rakness⁸, R. Ray⁸, M. Ricci²¹, M. Röhrken⁵, V. Rusu⁸, A. Saputi⁹, I. Sarra²¹, M. Schmitt²⁸, F. Spinella³⁰, D. Stratakis⁸, T. Strauss⁸, R. Talaga¹, V. Tereshchenko⁷, N. Tran², R. Tschirhart⁸, Z. Usubov⁷, M. Velasco²⁸, R. Wagner¹, Y. Wang², S. Werkema⁸, J. Whitmore⁸, P. Winter¹, L. Xia¹, L. Zhang⁵, R.-Y. Zhu⁵, V. Zutshi²⁷, R. Zwaska⁸

06 February 2018

Abstract

We propose an evolution of the Mu2e experiment, called Mu2e-II, that would leverage advances in detector technology and utilize the increased proton intensity provided by the Fermilab PIP-II upgrade to improve the sensitivity for neutrinoless muon-to-electron conversion by one order of magnitude beyond the Mu2e experiment, providing the deepest probe of charged lepton flavor violation in the foreseeable future. Mu2e-II will use as much of the Mu2e infrastructure as possible, providing, where required, improvements to the Mu2e apparatus to accommodate the increased beam intensity and cope with the accompanying increase in backgrounds.

Mu2e-II is an upgrade that will:

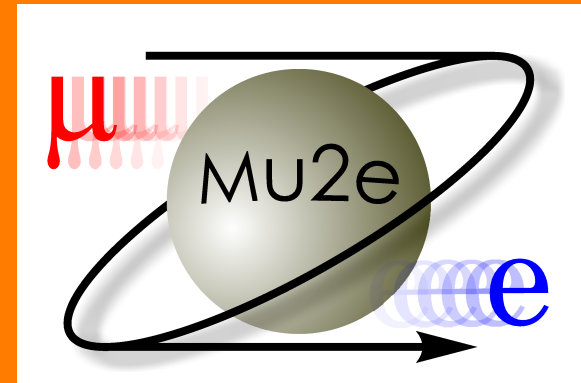
- Use ~100 kW of PIP-II protons @800 MeV
- Achieve an order of magnitude improvement in sensitivity
 - probe $R_{\mu e} \sim 10^{-18}$ level,
 - extend Λ_{NP} reach by x2



- EOI Submitted to Fermilab PAC in 2018
- arXiv:1802.02599, Fermilab-FN-1052
- 130 Signatories, 36 Institutions

PAC: “physics case is compelling” “endorse request for R&D funding”
Status: Pursuing high priority R&D. Data taking ~2030 timescale.

Mu2e-II - a next generation $\mu \rightarrow e$ conversion experiment at FNAL



Expression of Interest for Evolution of the Mu2e Experiment[†]

F. Abusalma²³, D. Ambrose²³, A. Artikov⁷, R. Bernstein⁸, G.C. Blazey²⁷, C. Bloise⁹, S. Boi³³, T. Bolton¹⁴, J. Bono⁸, R. Bonventre¹⁶, D. Bowring⁸, D. Brown¹⁶, D. Brown²⁰, K. Byrum¹, M. Campbell²², J.-F. Caron¹², F. Cervelli³⁰, D. Chokheli⁷, K. Ciampa²³, R. Ciolini³⁰, R. Coleman⁸, D. Cronin-Hennessy²³, R. Culbertson⁸, M.A. Cummings²⁵, A. Daniel¹², Y. Davydov⁷, S. Demers³⁵, D. Denisov⁸, S. Denisov¹³, S. Di Falco³⁰, E. Diociaiuti⁹, R. Djilkibaev²⁴, S. Donati³⁰, R. Donghia⁹, G. Drake¹, E.C. Dukes³³, B. Echenard⁵, A. Edmonds¹⁶, R. Ehrlich³³, V. Evdokimov¹³, P. Fabbriatore¹⁰, A. Ferrari¹¹, M. Frank³², A. Gaponenko⁸, C. Gatto²⁶, Z. Giorgio¹⁷, S. Giovannella⁹, V. Giusti³⁰, H. Glass⁸, D. Glenzinski⁸, L. Goodenough¹, C. Group³³, F. Happacher⁹, L. Harkness-Brennan¹⁹, D. Hedin²⁷, K. Heller²³, D. Hitlin⁵, A. Hocker⁸, R. Hooper¹⁸, G. Horton-Smith¹⁴, C. Hu⁵, P.Q. Hung³³, E. Hungerford¹², M. Jenkins³², M. Jones³¹, M. Kargiantoulakis⁸, K. S. Khaw³⁴, B. Kiburg⁸, Y. Kolomensky^{3,16}, J. Kozminski¹⁸, R. Kutschke⁸, M. Lancaster¹⁵, D. Lin⁵, I. Logashenko²⁹, V. Lombardo⁸, A. Luca⁸, G. Lukicov¹⁵, K. Lynch⁶, M. Martini²¹, A. Mazzacane⁸, J. Miller², S. Miscetti⁹, L. Morescalchi³⁰, J. Mott², S. E. Mueller¹¹, P. Murat⁸, V. Nagaslaev⁸, D. Neuffer⁸, Y. Oksuzian³³, D. Pasciuto³⁰, E. Pedreschi³⁰, G. Pezzullo³⁵, A. Pla-Dalmau⁸, B. Pollack²⁸, A. Popov¹³, J. Popp⁶, F. Porter⁵, E. Prebys⁴, V. Pronskikh⁸, D. Pushka⁸, J. Quirk², G. Rakness⁸, R. Ray⁸, M. Ricci²¹, M. Röhrken⁵, V. Rusu⁸, A. Saputi⁹, I. Sarra²¹, M. Schmitt²⁸, F. Spinella³⁰, D. Stratakis⁸, T. Strauss⁸, R. Talaga¹, V. Tereshchenko⁷, N. Tran², R. Tschirhart⁸, Z. Usubov⁷, M. Velasco²⁸, R. Wagner¹, Y. Wang², S. Werkema⁸, J. Whitmore⁸, P. Winter¹, L. Xia¹, L. Zhang⁵, R.-Y. Zhu⁵, V. Zutshi²⁷, R. Zwaska⁸

06 February 2018

Abstract

We propose an evolution of the Mu2e experiment, called Mu2e-II, that would leverage advances in detector technology and utilize the increased proton intensity provided by the Fermilab PIP-II upgrade to improve the sensitivity for neutrinoless muon-to-electron conversion by one order of magnitude beyond the Mu2e experiment, providing the deepest probe of charged lepton flavor violation in the foreseeable future. Mu2e-II will use as much of the Mu2e infrastructure as possible, providing, where required, improvements to the Mu2e apparatus to accommodate the increased beam intensity and cope with the accompanying increase in backgrounds.

Mu2e-II is an upgrade that will:

- Use ~100 kW of PIP-II protons @800 MeV
- Achieve an order of magnitude improvement in sensitivity
 - probe $R_{\mu e} \sim 10^{-18}$ level,
 - extend Λ_{NP} reach by x2



- EOI Submitted to Fermilab PAC in 2018
- arXiv:1802.02599, Fermilab-FN-1052
- 130 Signatories, 36 Institutions

PAC: “physics case is compelling” “end of the road”
Status: Pursuing high priority R&D. D. I. Oksuzian on Thursday

If found ...

Effective Field Theory for $\mu \rightarrow e$ Conversion

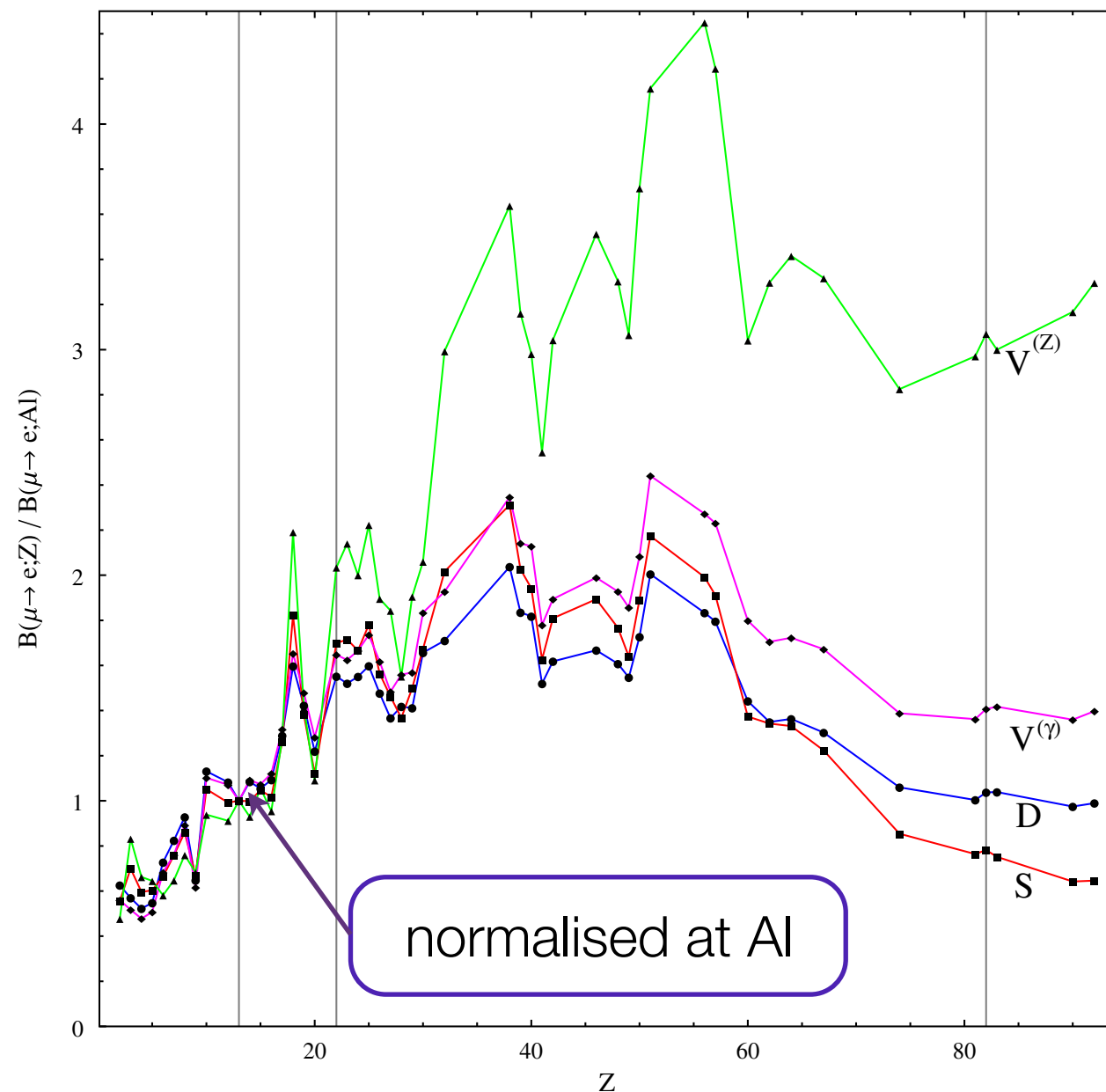


two-lepton and two-nucleon operators and dipole operators

$$\begin{aligned}
 \mathcal{L}_{\mu A \rightarrow e A}(\Lambda_{expt}) = & -\frac{4G_F}{\sqrt{2}} \sum_{N=p,n} \left[m_\mu (C_{DL} \bar{e}_R \sigma^{\alpha\beta} \mu_L F_{\alpha\beta} + C_{DR} \bar{e}_L \sigma^{\alpha\beta} \mu_R F_{\alpha\beta}) \right. \\
 & \text{scalar} \quad + \left(\tilde{C}_{SL}^{(NN)} \bar{e} P_L \mu + \tilde{C}_{SR}^{(NN)} \bar{e} P_R \mu \right) \bar{N} N \\
 & \text{pseudo-scalar} \quad + \left(\tilde{C}_{P,L}^{(NN)} \bar{e} P_L \mu + \tilde{C}_{P,R}^{(NN)} \bar{e} P_R \mu \right) \bar{N} \gamma_5 N \\
 & \text{vector} \quad + \left(\tilde{C}_{VL}^{(NN)} \bar{e} \gamma^\alpha P_L \mu + \tilde{C}_{VR}^{(NN)} \bar{e} \gamma^\alpha P_R \mu \right) \bar{N} \gamma_\alpha N \\
 & \text{axial-vector} \quad + \left(\tilde{C}_{A,L}^{(NN)} \bar{e} \gamma^\alpha P_L \mu + \tilde{C}_{A,R}^{(NN)} \bar{e} \gamma^\alpha P_R \mu \right) \bar{N} \gamma_\alpha \gamma_5 N \\
 & \text{(derivative)} \quad + \left(\tilde{C}_{Der,L}^{(NN)} \bar{e} \gamma^\alpha P_L \mu + \tilde{C}_{Der,R}^{(NN)} \bar{e} \gamma^\alpha P_R \mu \right) i(\bar{N} \overleftrightarrow{\partial}_\alpha \gamma_5 N) \\
 & \text{tensor} \quad + \left(\tilde{C}_{T,L}^{(NN)} \bar{e} \sigma^{\alpha\beta} P_L \mu + \tilde{C}_{T,R}^{(NN)} \bar{e} \sigma^{\alpha\beta} P_R \mu \right) \bar{N} \sigma_{\alpha\beta} N + h.c. \Big] \quad .
 \end{aligned}$$

dipole

Discrimination of the interactions by different targets



vector interaction
(with Z boson)

with Z penguin

vector interaction
(with photon -
charge radius)

left-right models

dipole interaction

SUSY-GUT

scalar interaction

SUSY seesaw

R. Kitano, M. Koike and Y. Okada, Phys.Rev. D66 (2002) 096002; D76 (2007) 059902
V. Cirigliano, R. Kitano, Y. Okada, and P. Tuzon, Phys. Rev. D80 (2009) 013002
S. Davidson, YK, M. Yamanaka, Phys. Lett. B790 (2019) 380-388

Spin Dependent μ -e conversion and Spin Independent μ -e conversion



dipole
interaction

vector
interaction

scalar
interaction

Spin Independent
 μ -e Conversion
(coherent)

Spin Dependent μ -e conversion and Spin Independent μ -e conversion



dipole
interaction

vector
interaction

scalar
interaction

Spin Independent
 μ -e Conversion
(coherent)

Pseudo-
scaler
interaction

axial vector
interaction

tensor
interaction

Spin Dependent
 μ -e Conversion
(incoherent)

compare zero-spin and non-zero-spin nuclear targets

V. Cirigliano, S. Davidson, YK, Phys. Lett. B 771 (2017) 242

S. Davidson, YK, A. Saporta, Eur. Phys. J. C78 (2018) 109

LVN/CLFV

μ^- to e^+ conversion in muonic atom

$$\mu^- + N(A, Z) \rightarrow e^+ + N(A, Z - 2) \quad \text{ground or excited final states.}$$

Lepton number violation (LNV) and
Lepton flavour violation (LFV)

signal signature

$$E_{\mu e^+} = m_\mu - B_\mu - E_{rec} - (M(A, Z - 2) - M(A, Z))$$

backgrounds

- radiative muon nuclear capture (RMC)
 $\mu^- + N(A, Z) \rightarrow N(A, Z - 1) + \nu + \gamma$

$$E_{RMC} = m_\mu - B_\mu - E_{rec} - (M(A, Z - 1) - M(A, Z))$$

μ^- to e^+ conversion in muonic atom

$$\mu^- + N(A, Z) \rightarrow e^+ + N(A, Z - 2) \quad \text{ground or excited final states.}$$

Lepton number violation (LNV) and
Lepton flavour violation (LFV)

signal signature

$$E_{\mu e^+} = m_\mu - B_\mu - E_{rec} - (M(A, Z - 2) - M(A, Z))$$

backgrounds

- radiative muon nuclear capture (RMC)
 $\mu^- + N(A, Z) \rightarrow N(A, Z - 1) + \nu + \gamma$

$$E_{RMC} = m_\mu - B_\mu - E_{rec} - (M(A, Z - 1) - M(A, Z))$$

- **Current limits :**

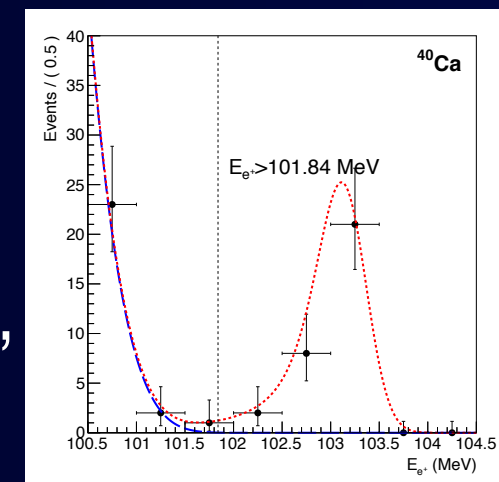
$$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}(\text{gs}) \leq 1.7 \times 10^{-12}$$

$$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}(\text{ex}) \leq 3.6 \times 10^{-11}$$

J. Kaulard et al. (SINDRUM-II)
Phys. Lett. B422 (1998) 334.

- **Future prospects:**

- Mu2e or COMET can improve with proper targets,



T. Geib, A. Merie, K. Zuber, Phys. Lett. B764 (2017) 157-162

B. Yeo, YK, M. Lee and K. Zuber, Phys. Rev. D96 (2017) 075027

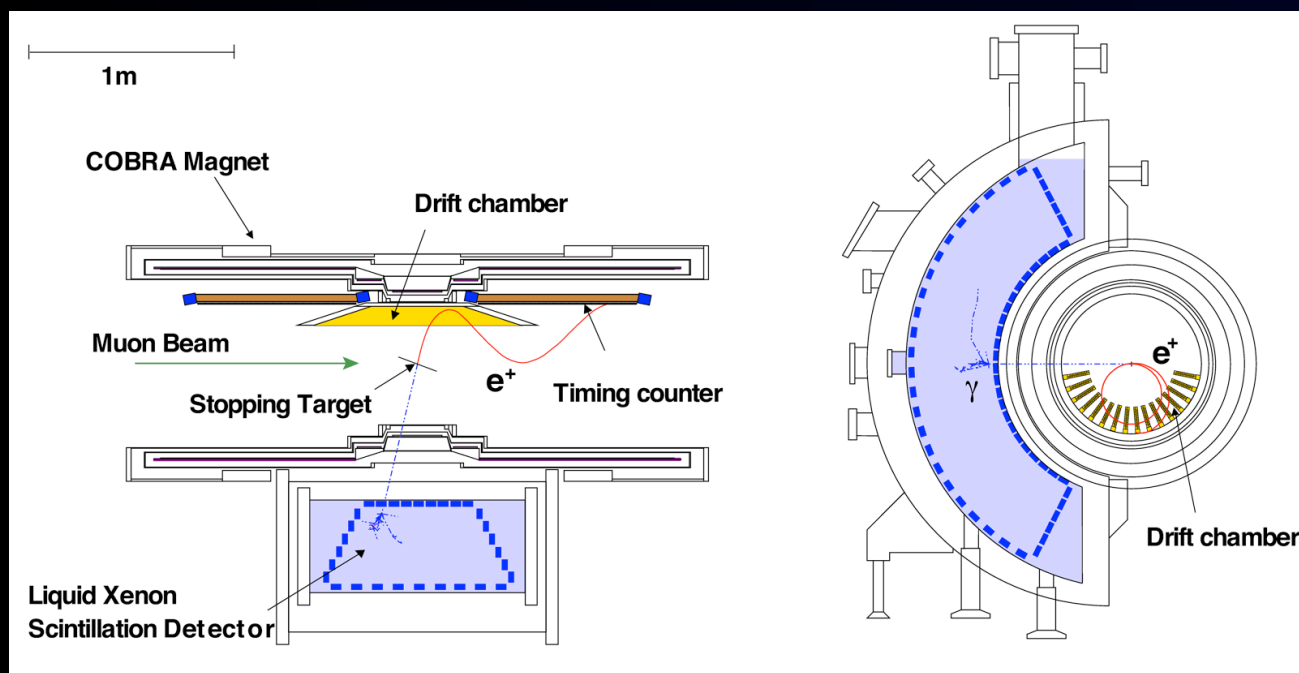
DC muon beam experiments are ...

CLFV Decay of Muons : $\mu^+ \rightarrow e^+ \gamma$

MEG @PSI

(2016)

- drift chamber for positrons
- liquid Xe detector for gammas
- DC muon beam at PSI



$$B(\mu^+ \rightarrow e^+ \gamma) < 4.2 \times 10^{-13}$$

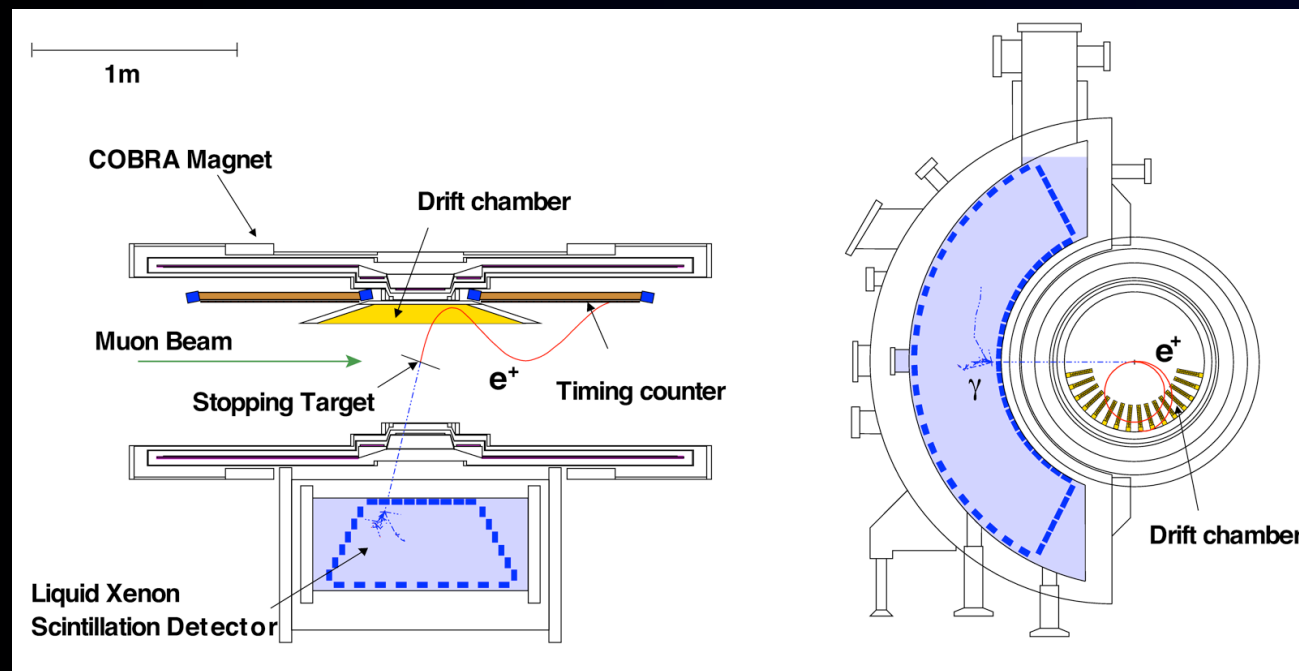
- a factor of 30 improvement

CLFV Decay of Muons : $\mu^+ \rightarrow e^+ \gamma$

MEG @PSI

(2016)

- drift chamber for positrons
- liquid Xe detector for gammas
- DC muon beam at PSI

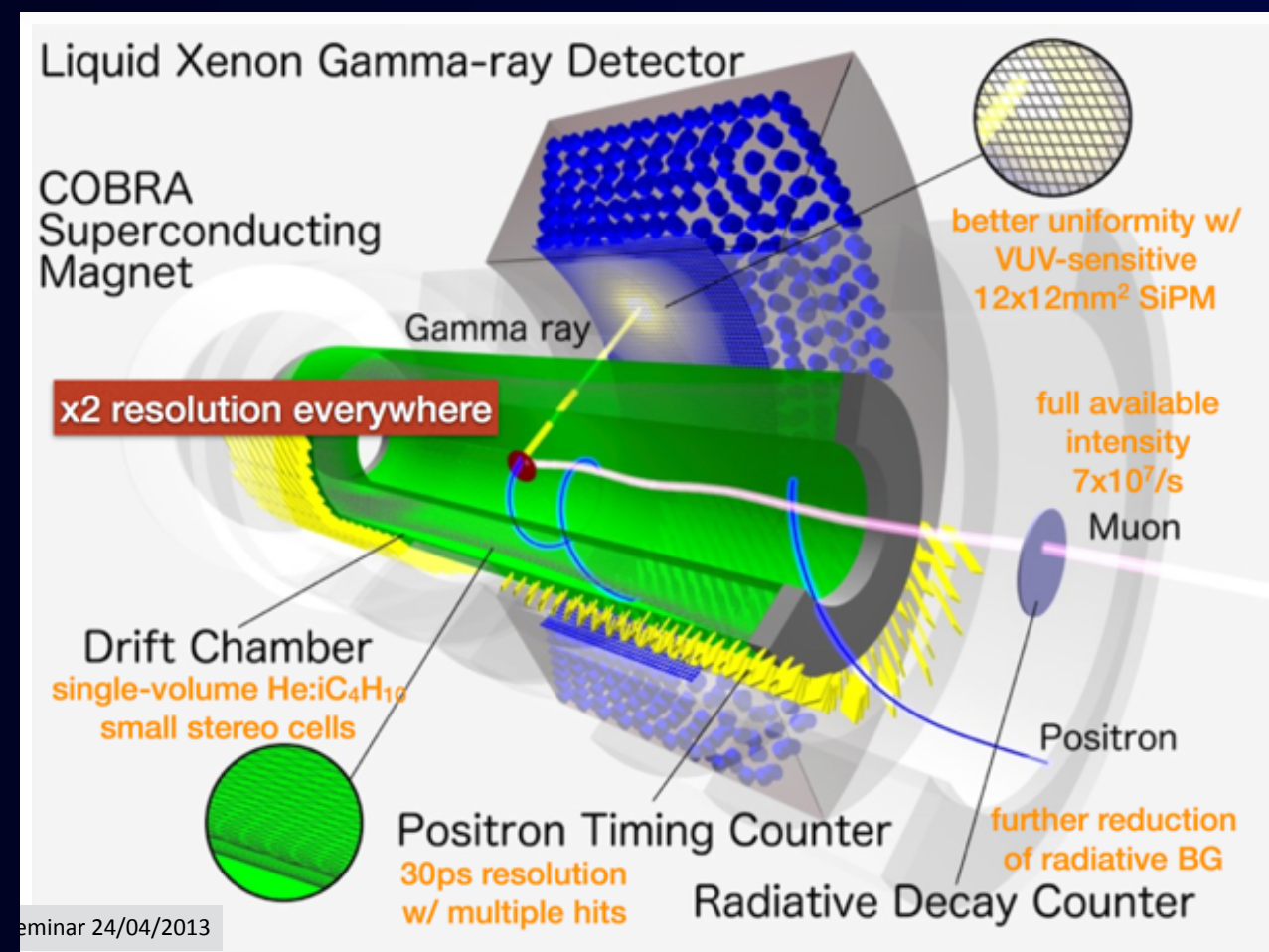


$$B(\mu^+ \rightarrow e^+ \gamma) < 4.2 \times 10^{-13}$$

- a factor of 30 improvement

MEG II

- all detectors upgraded
- full muon beam intensity
- Goal $\sim 6 \times 10^{-14}$ (2020-2023)

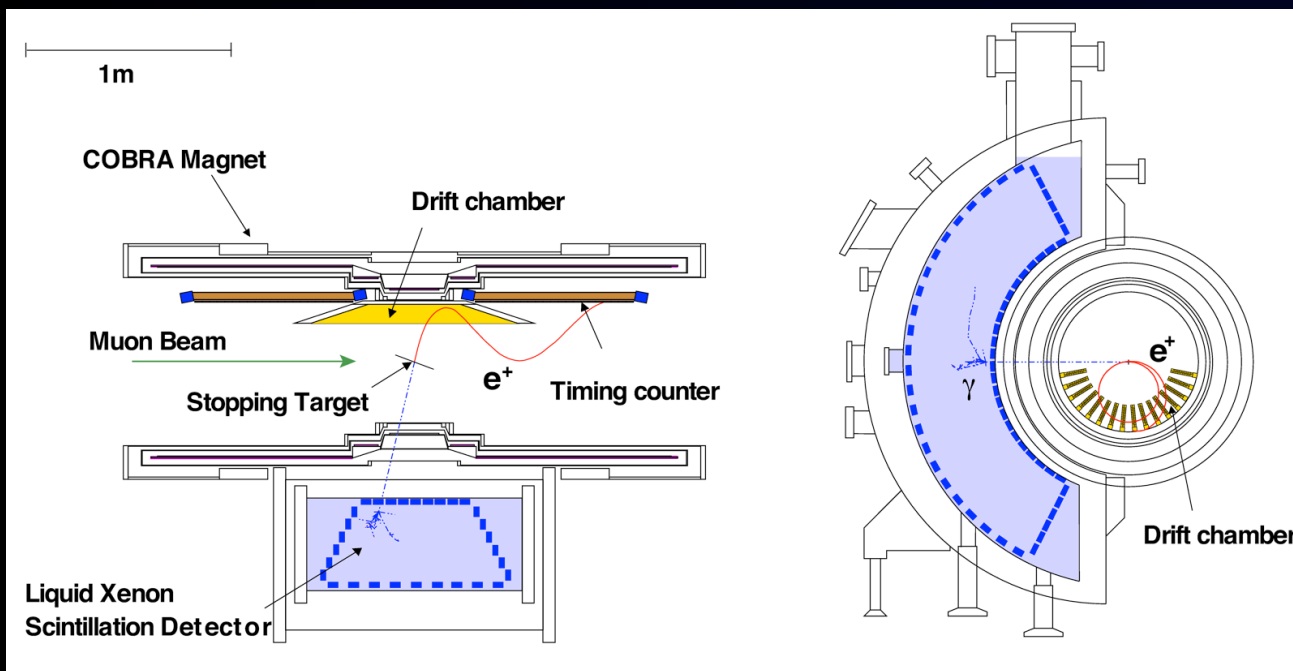


CLFV Decay of Muons : $\mu^+ \rightarrow e^+ \gamma$

MEG @PSI

(2016)

- drift chamber for positrons
- liquid Xe detector for gammas
- DC muon beam at PSI

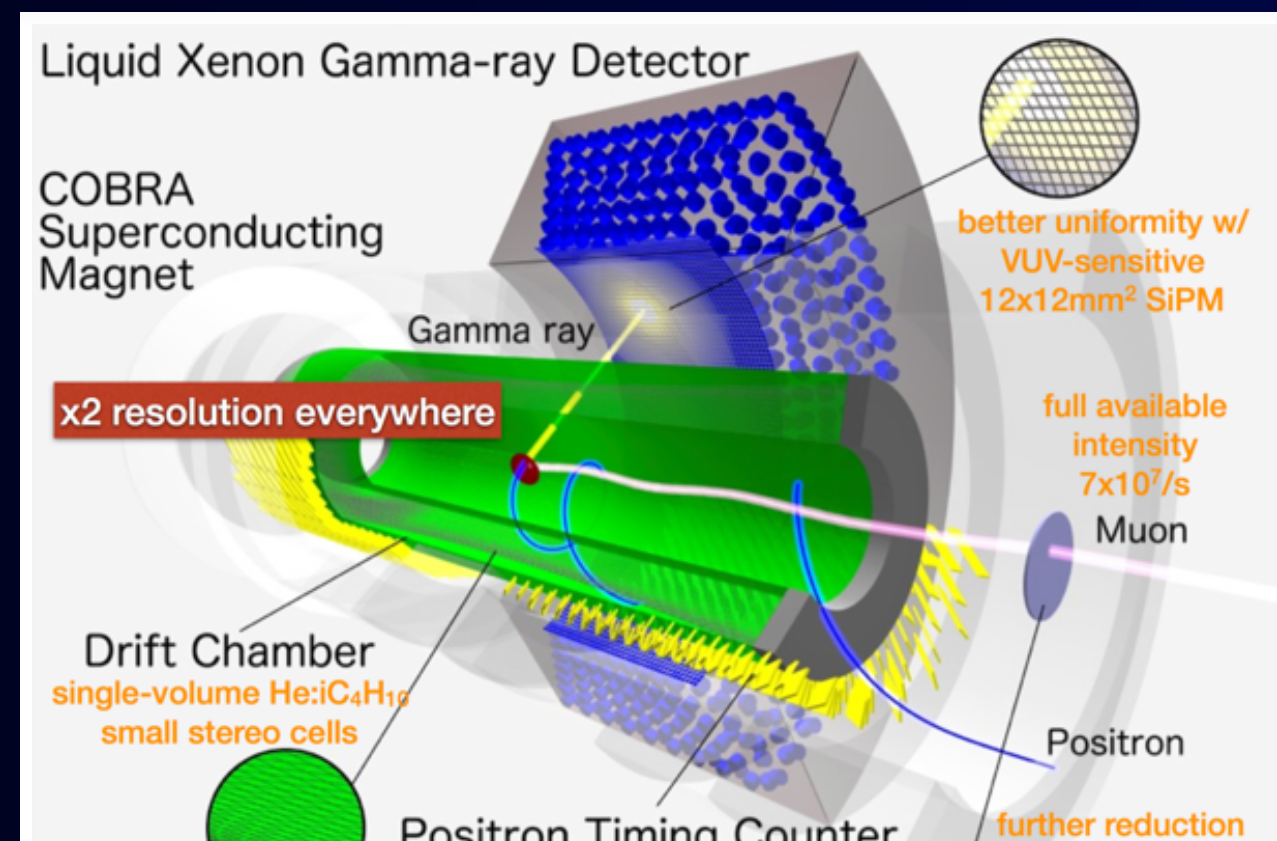


$$B(\mu^+ \rightarrow e^+ \gamma) < 4.2 \times 10^{-13}$$

- a factor of 30 improvement

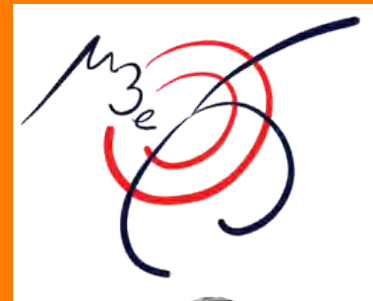
MEG II

- all detectors upgraded
- full muon beam intensity
- Goal $\sim 6 \times 10^{-14}$ (2020-2023)



S. Mihara on Thursday

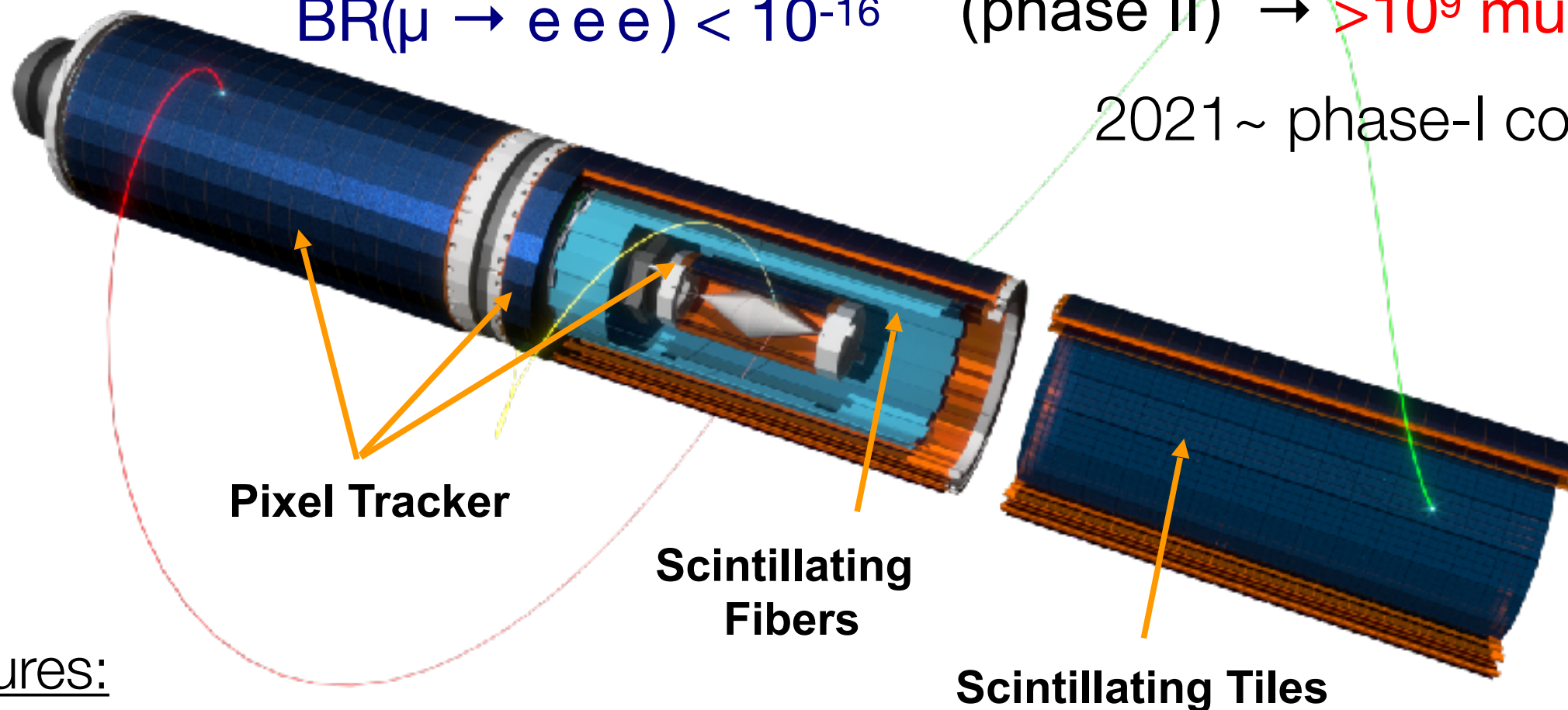
Mu3e at PSI



$BR(\mu \rightarrow e e e) < 2 \cdot 10^{-15}$ (phase I) $\rightarrow 10^8$ muons/s (PiE5)

$BR(\mu \rightarrow e e e) < 10^{-16}$ (phase II) $\rightarrow > 10^9$ muons/s (HiMB)

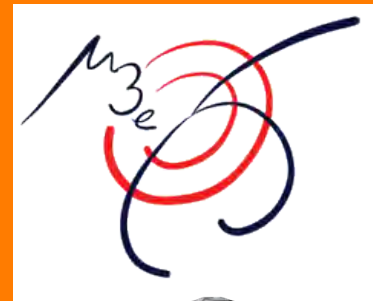
2021 ~ phase-I commissioning



Features:

- surface muons ($p=29$ MeV/c, DC) stopped on target at high rate: $10^8 - 10^9$ /s
- ultra thin **silicon pixel detector** (HV-MAPS) with 1 per mill radiation length / layer
- high precision tracking using recurling tracks in strong magnetic field
- fast timing detectors (**scintillating fibers & tiles**)
- helium gas cooling

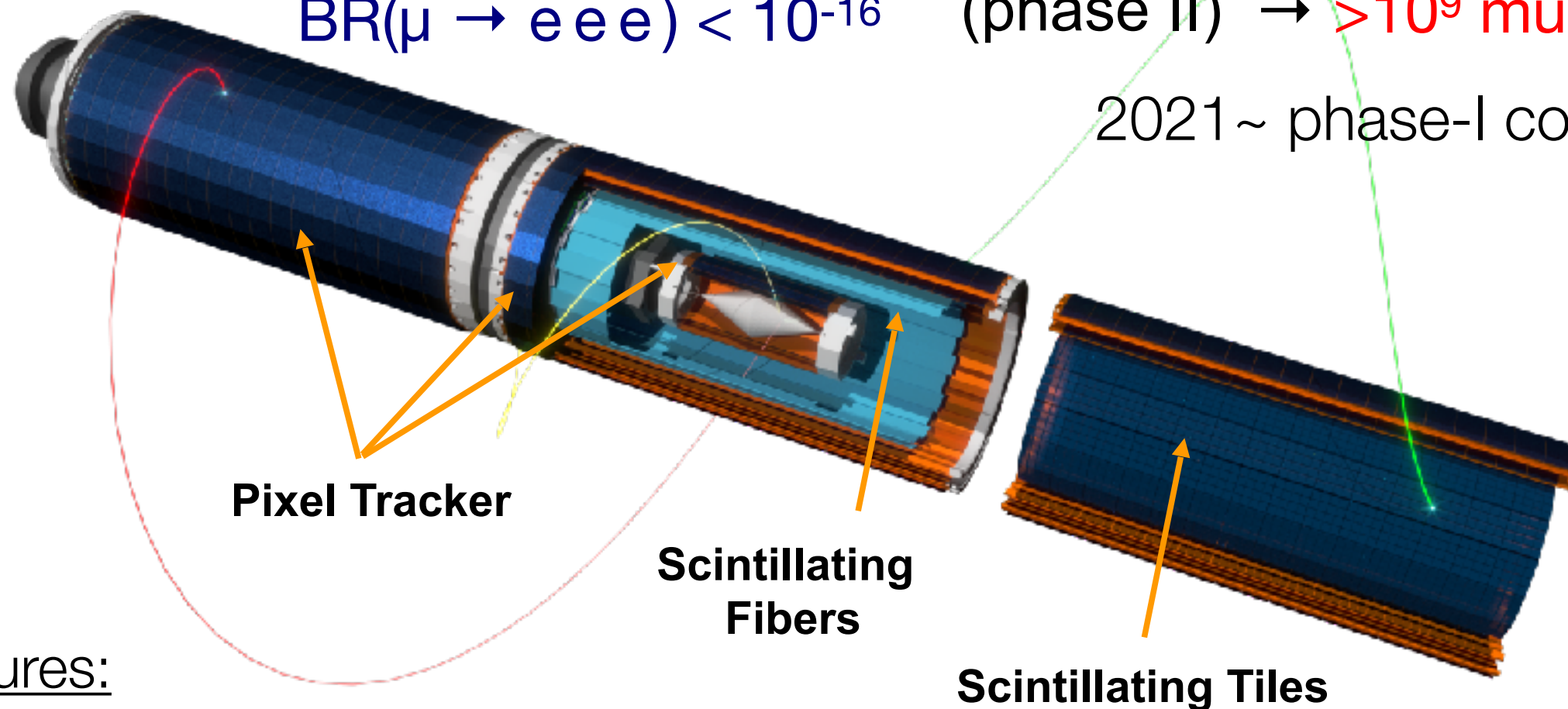
Mu3e at PSI



$BR(\mu \rightarrow e e e) < 2 \cdot 10^{-15}$ (phase I) $\rightarrow 10^8$ muons/s (PiE5)

$BR(\mu \rightarrow e e e) < 10^{-16}$ (phase II) $\rightarrow > 10^9$ muons/s (HiMB)

2021 ~ phase-I commissioning



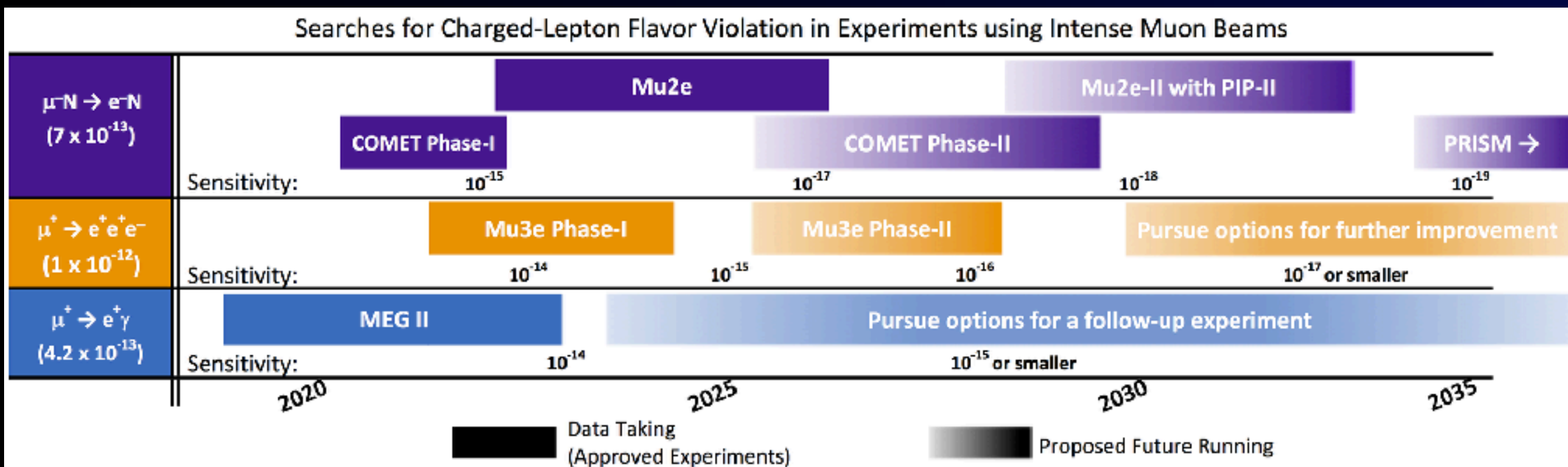
Features:

- surface muons ($p=29$ MeV/c, DC) stopped on target at high rate: $10^8 - 10^9$ /s
- ultra thin **silicon pixel detector** (HV-MAPS) with 1 per mill radiation length / layer
- high precision tracking using recurling tracks in strong magnetic field
- fast timing detectors (**scintillating fibers** & **tiles**)
- helium gas cooling

S. Dittmeier on Thursday

Schedule of “golden” $\mu \rightarrow e$ Transition processes in 2025 and beyond

Timeline submitted to EPPSU 2020



Improvement : $> 10,000$

Muon g-2



Muon g-2 : Spin precession

Under a magnetic field

Spin vector precession

Momentum vector motion

$$\omega_S = \frac{eB}{m_\mu \gamma} \left[1 + \frac{(g-2)}{2} \gamma \right] \quad \omega_C = \frac{eB}{m_\mu \gamma}$$

Muon g-2 : Spin precession

Under a magnetic field

Spin vector precession

Momentum vector motion

$$\omega_S = \frac{eB}{m_\mu \gamma} \left[1 + \frac{(g-2)}{2} \gamma \right] \quad \omega_C = \frac{eB}{m_\mu \gamma}$$

Spin precession with respect to the momentum vector

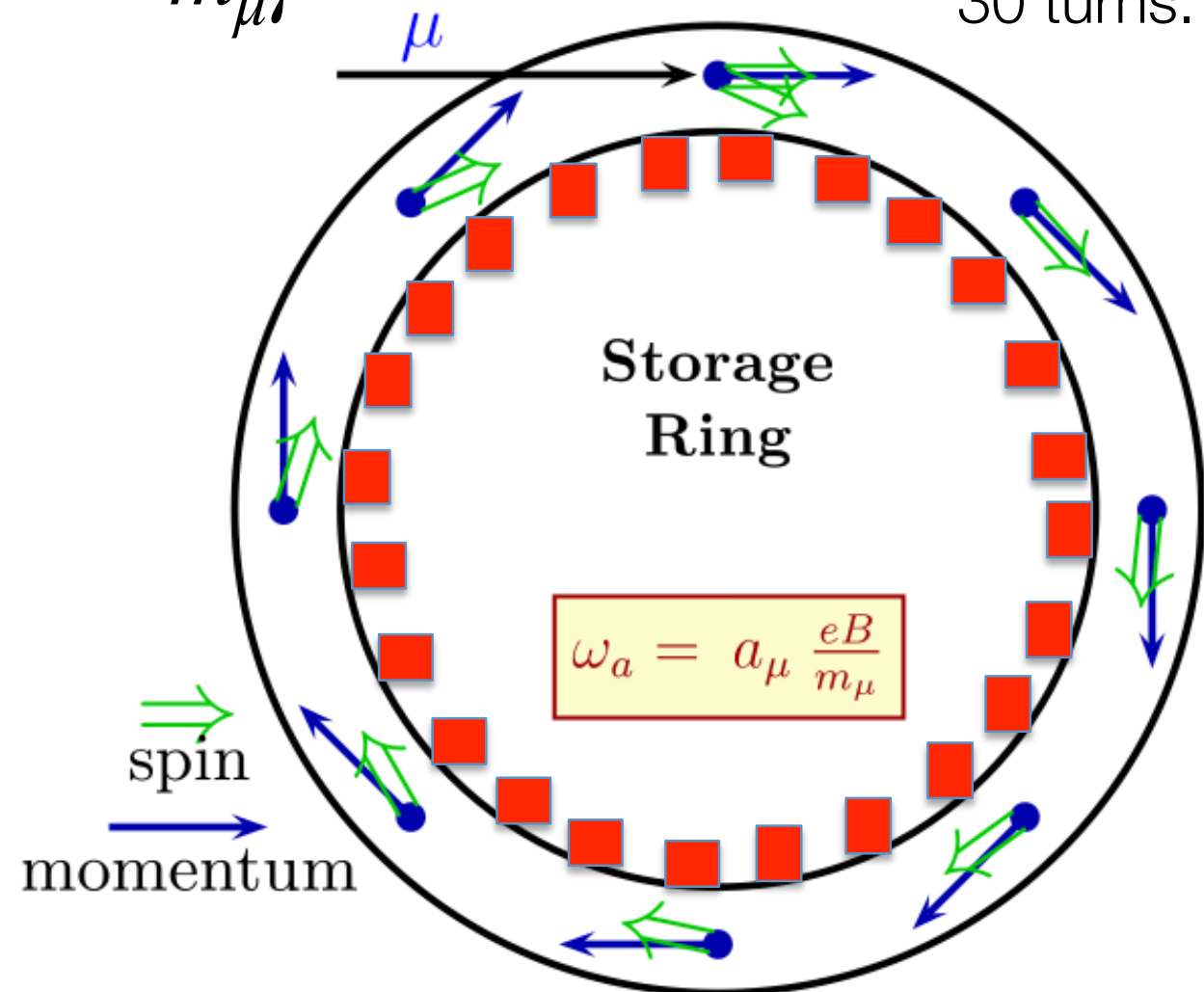
$$\omega_S - \omega_C = \omega_a = \frac{a_\mu eB}{m_\mu}$$

$$a_\mu = \frac{1}{2}(g-2)$$



Francis Farley, Hans Sens, Georges Charpak, Theo Muller, Anton with the 6-meter g-2 magnet



At the BNL experiment, spin precesses around momentum once every 30 turns.



Muon g-2

$$\omega_a$$

$$\omega_a = -\frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

ZERO  ZERO 

FNAL/BNL approach : effect of focussing E-field cancelled by using “magic” 3.09 GeV momenta muons.

J-PARC approach : 300 MeV beam with v. low transverse momenta requiring no E-field to focus.

Muon g-2

$$\omega_a$$

$$\omega_a = -\frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

ZERO ZERO

FNAL/BNL approach : effect of focussing E-field cancelled by using “magic” 3.09 GeV momenta muons.

J-PARC approach : 300 MeV beam with v. low transverse momenta requiring no E-field to focus.

B is measured using NMR in terms of the proton Larmor frequency

$$\omega_p$$

Muon g-2

$$\omega_a$$

$$\omega_a = -\frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

ZERO ZERO

FNAL/BNL approach : effect of focussing E-field cancelled by using “magic” 3.09 GeV momenta muons.

J-PARC approach : 300 MeV beam with v. low transverse momenta requiring no E-field to focus.

$$\omega_p$$

B is measured using NMR in terms of the proton Larmor frequency

$$\frac{g-2}{2} = a_\mu = \frac{\omega_a}{\omega_p} \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

3ppb
22ppb
0.26ppt

Muon g-2 Prediction

$$a_\mu = a_\mu(QED) + a_\mu(had) + a_\mu(weak) + a_\mu(BSM)$$

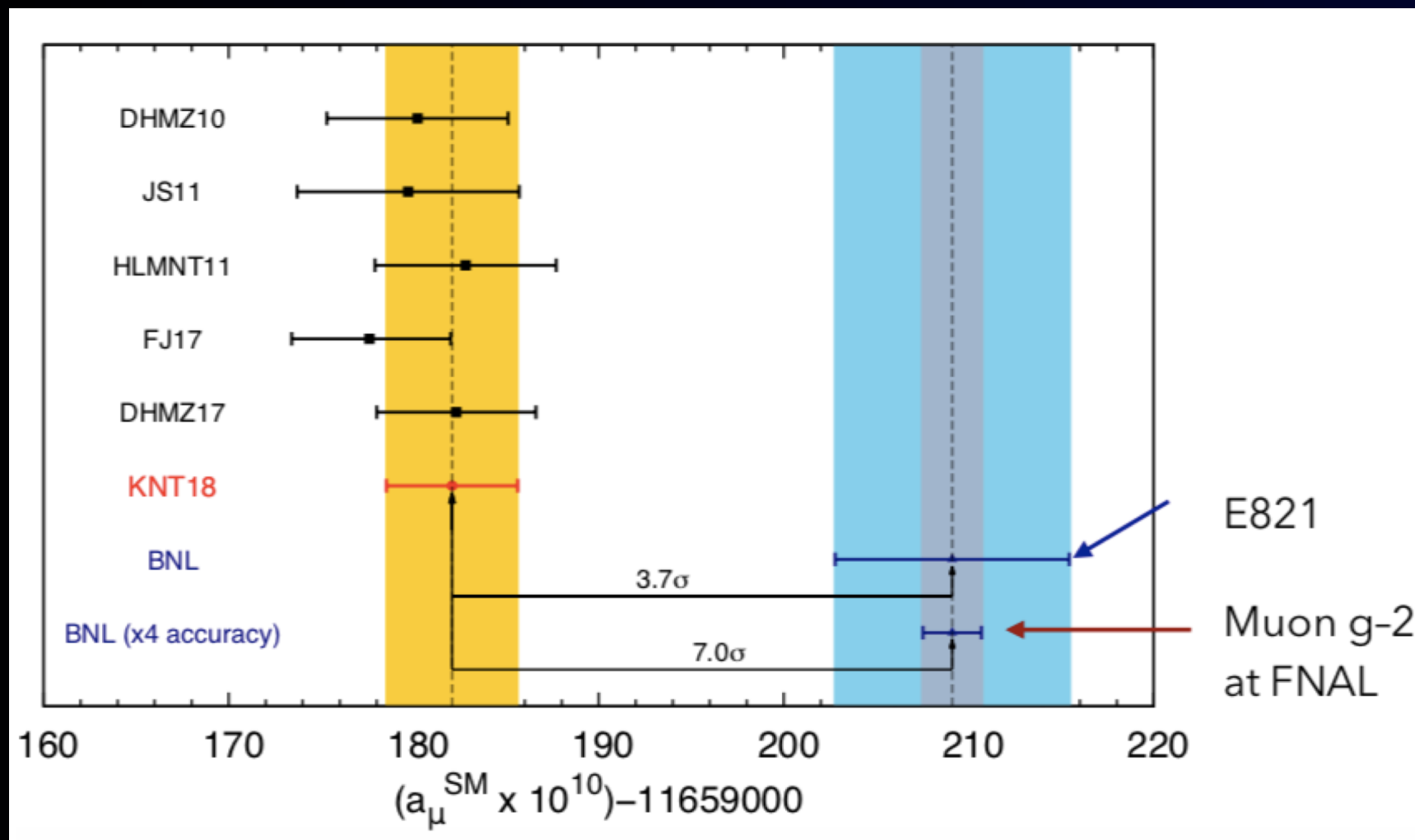
$$a_\mu^{SM} = 116591820.5(35.6) \times 10^{-11}$$

Muon g-2 Prediction

$$a_\mu = a_\mu(QED) + a_\mu(had) + a_\mu(weak) + a_\mu(BSM)$$

$$a_\mu^{SM} = 116591820.5(35.6) \times 10^{-11}$$

$$a_\mu^{exp} = 116595208.9(63) \times 10^{-11} \quad \text{BNL E821}$$



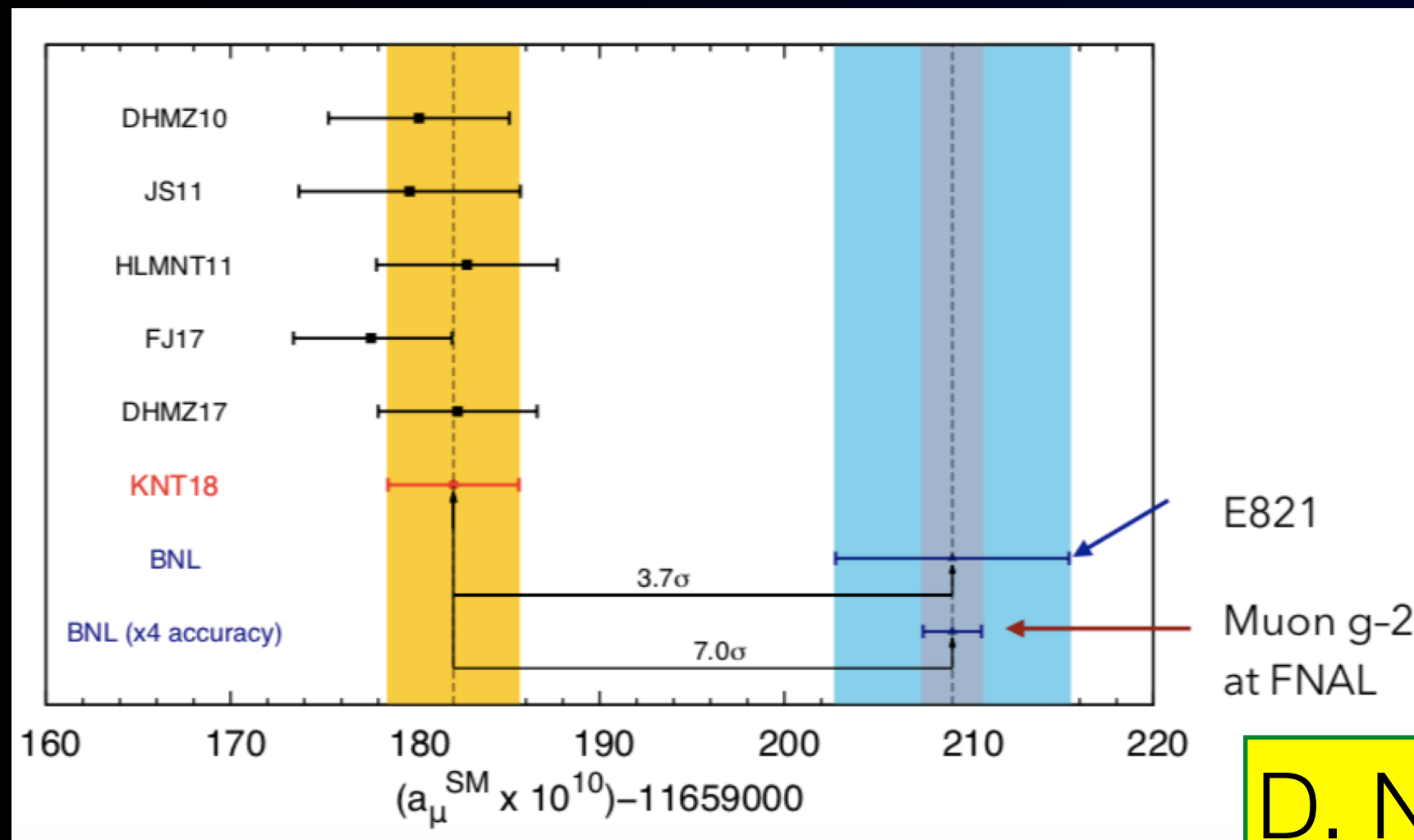
3.7σ deviation
from the SM
expectation

Muon g-2 Prediction

$$a_\mu = a_\mu(QED) + a_\mu(had) + a_\mu(weak) + a_\mu(BSM)$$

$$a_\mu^{SM} = 116591820.5(35.6) \times 10^{-11}$$

$$a_\mu^{exp} = 116595208.9(63) \times 10^{-11} \quad \text{BNL E821}$$



3.7σ deviation
from the SM
expectation

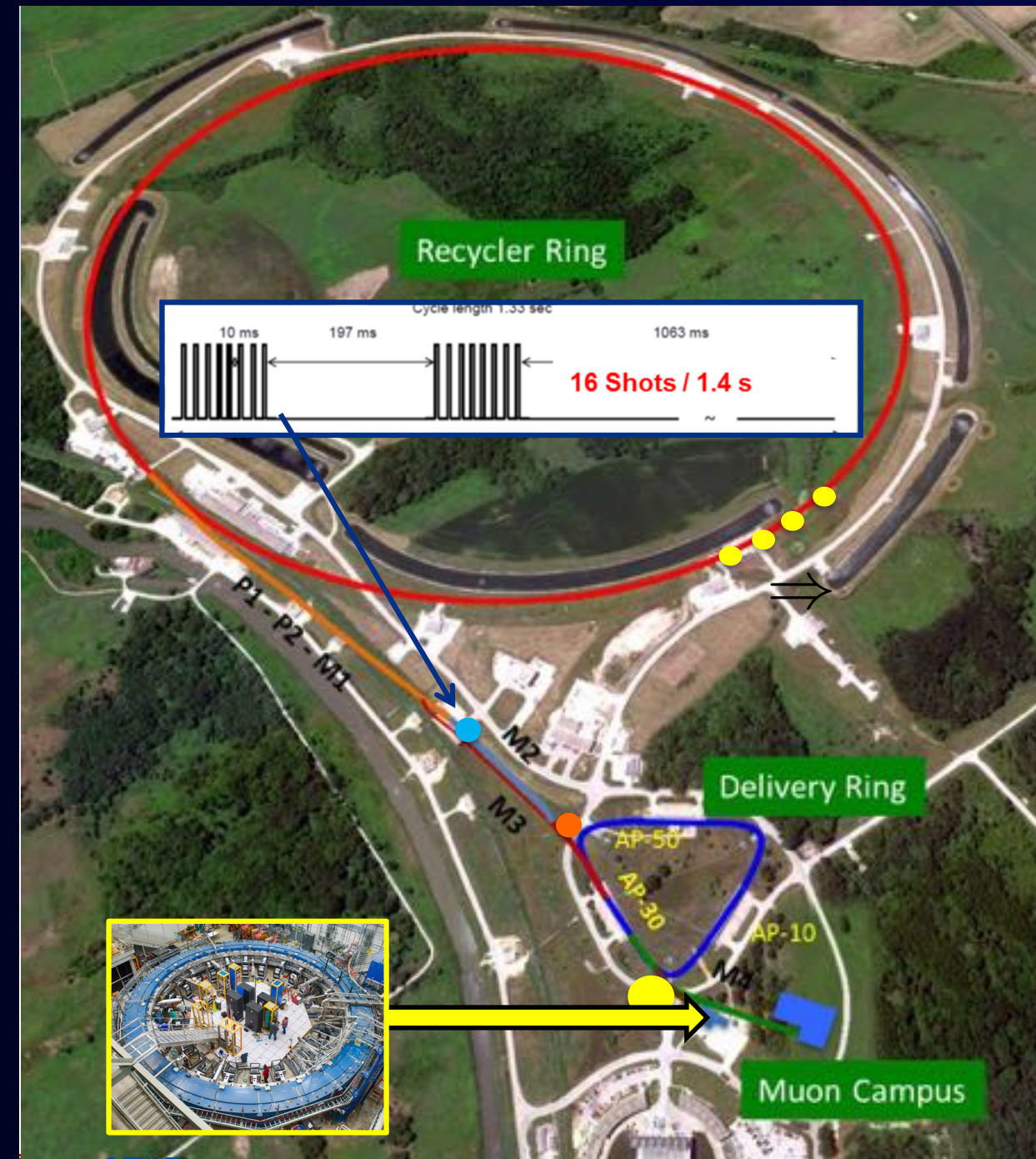
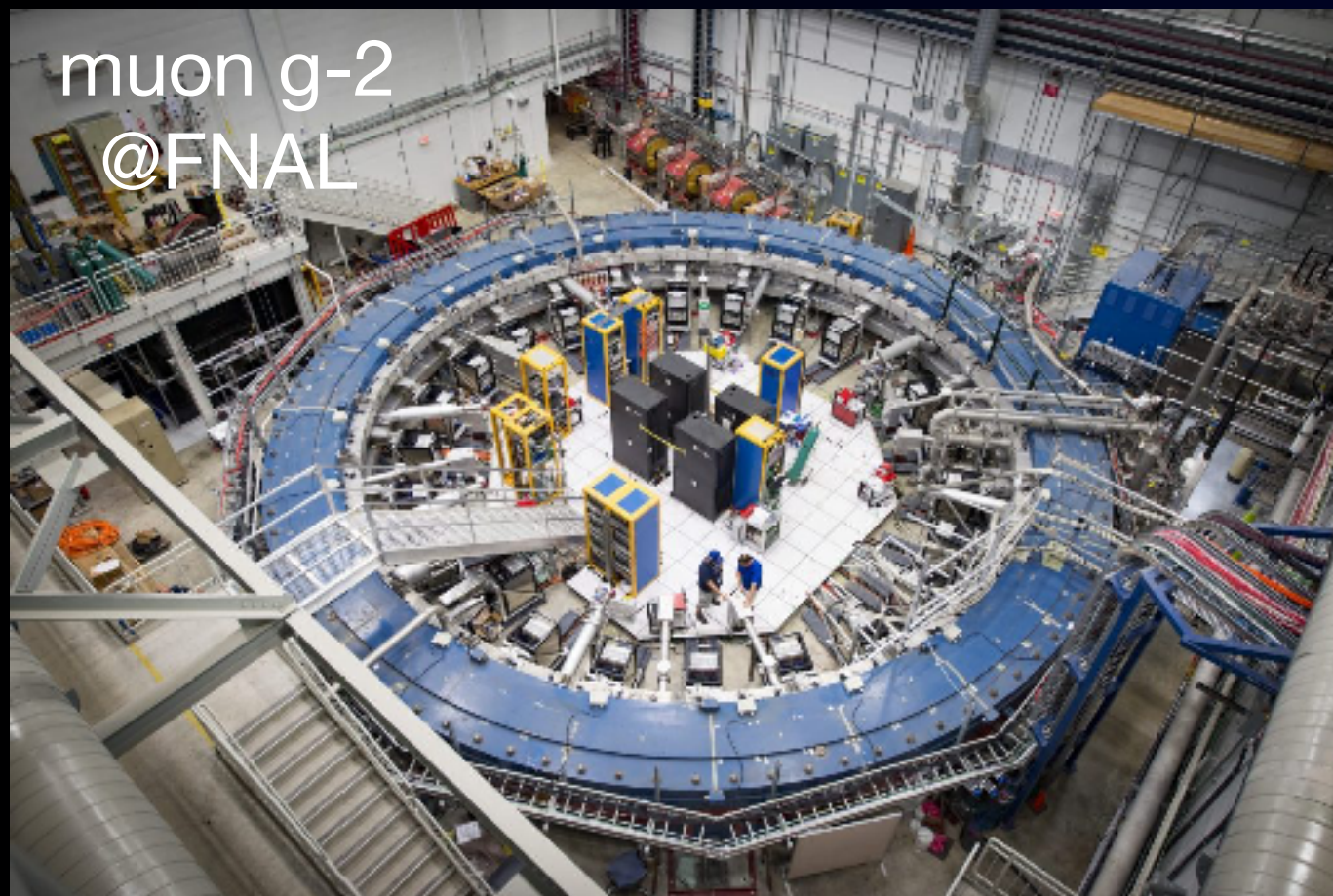
D. Nomura on Tuesday

Muon $g-2$ (E989) at FNAL

- aim at 0.14 ppm (x4 improvement)
- significant improvements over BNL E821
- run 1 data (2018) $\sim 1.4 \times \text{BNL}$
- run 2 data (2019) $\sim 1.8 \times \text{BNL}$
- run 1 result by the end of this year

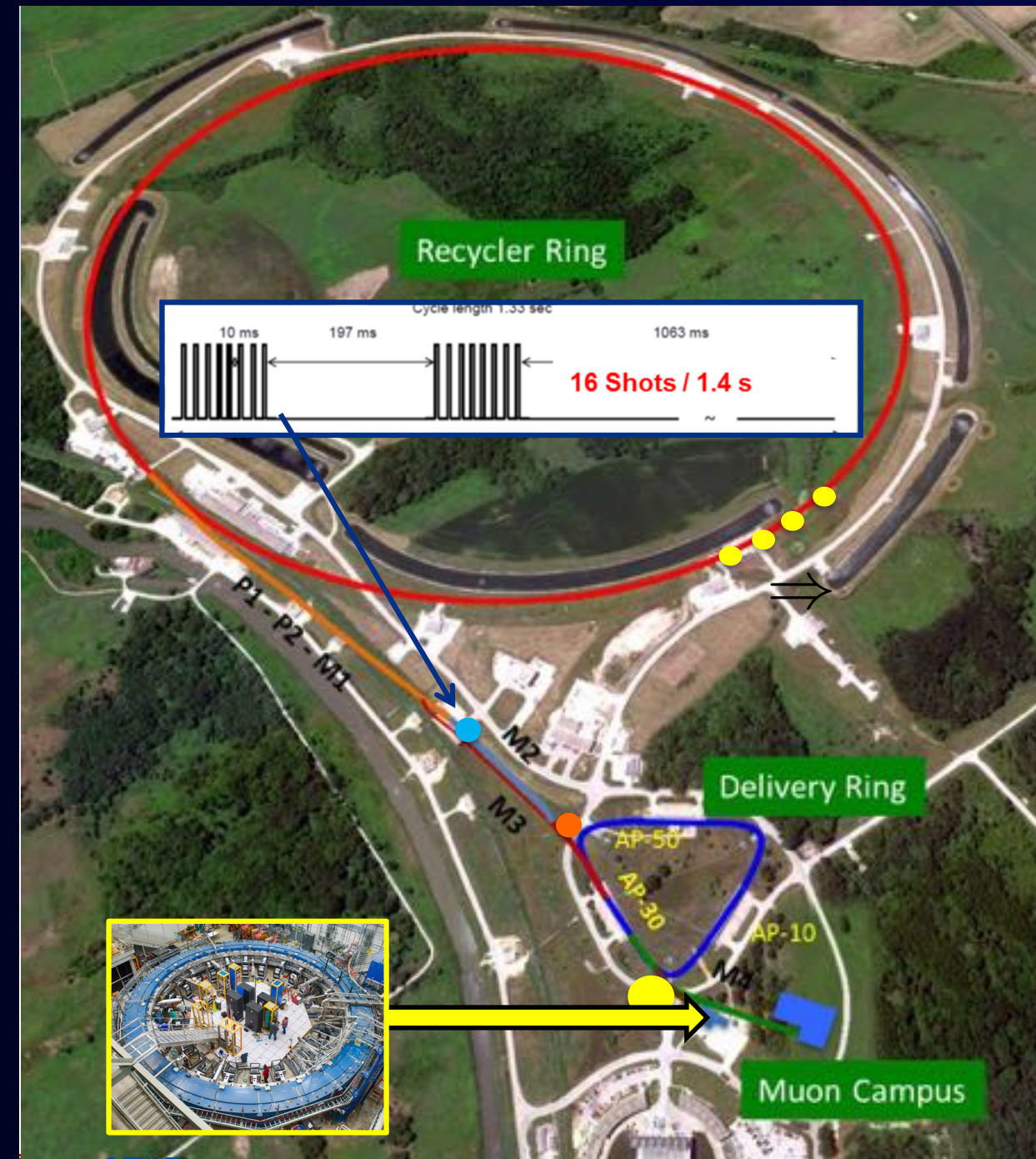
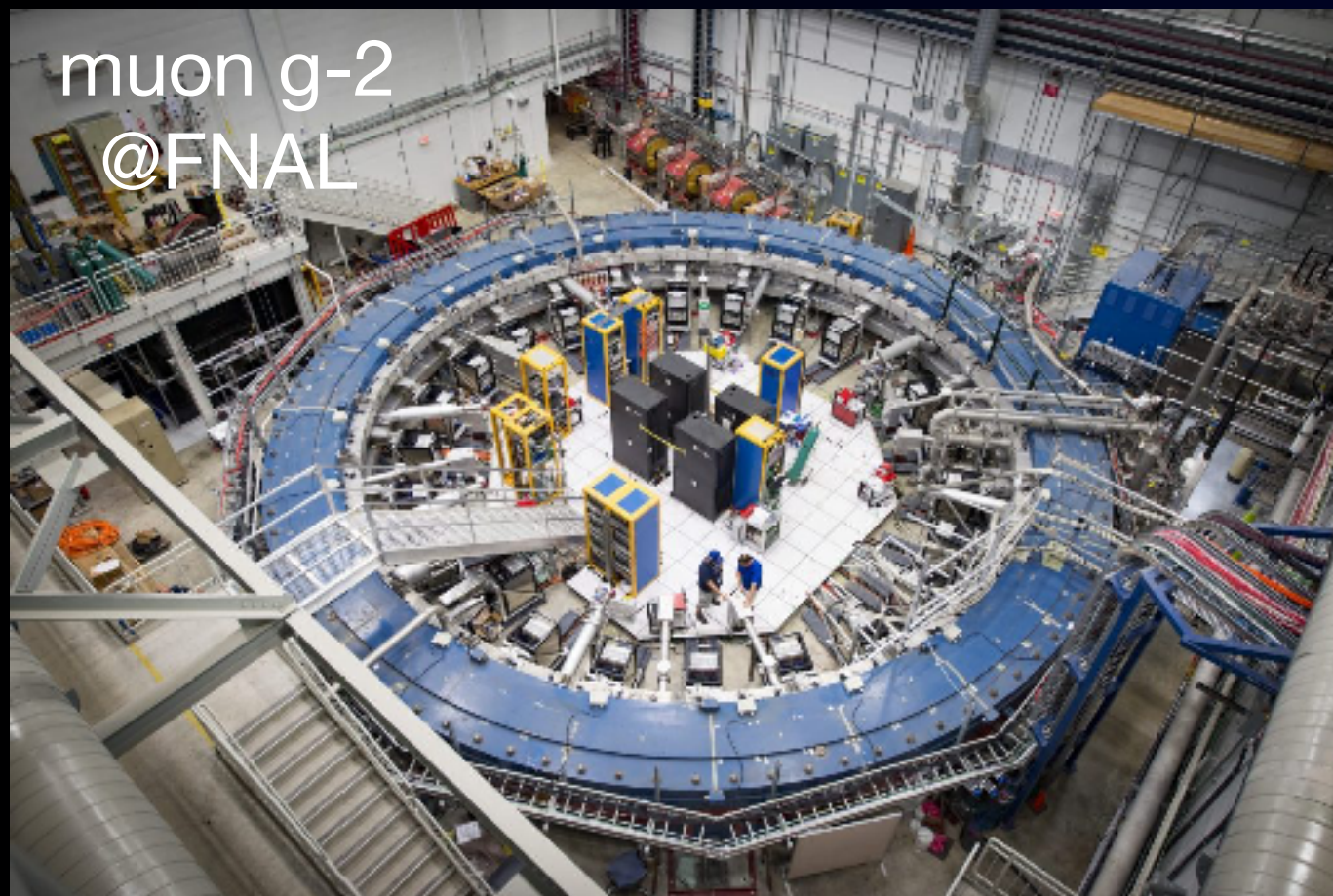
Muon g-2 (E989) at FNAL

- aim at 0.14 ppm (x4 improvement)
- significant improvements over BNL E821
- run 1 data (2018) ~ 1.4xBNL
- run 2 data (2019) ~ 1.8xBNL
- run 1 result by the end of this year



Muon g-2 (E989) at FNAL

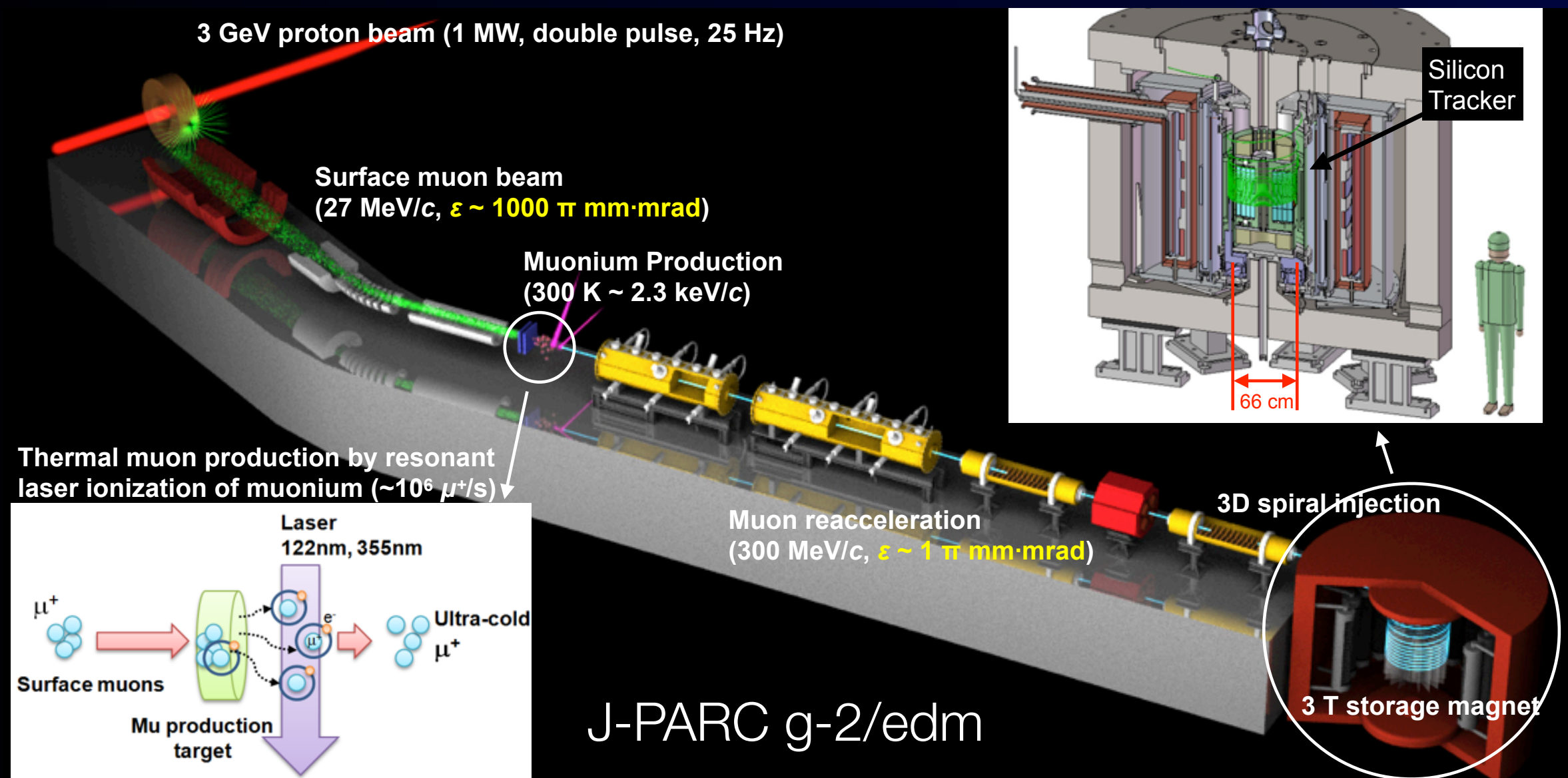
- aim at 0.14 ppm (x4 improvement)
- significant improvements over BNL E821
- run 1 data (2018) ~ 1.4xBNL
- run 2 data (2019) ~ 1.8xBNL
- run 1 result by the end of this year



S. Corrodi on Tuesday

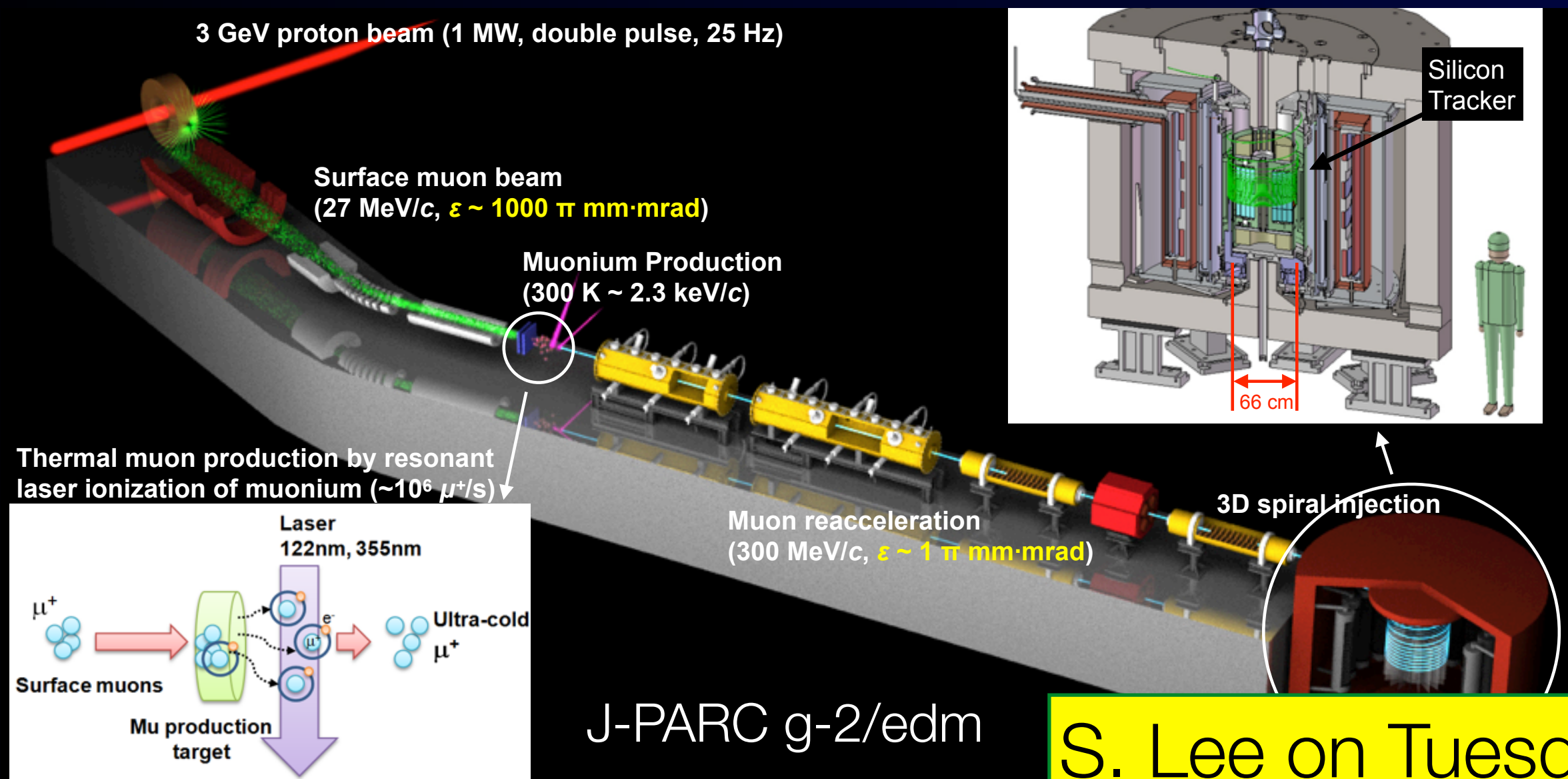
J-PARC g-2/edm

- new technique : slow muons from laser ionization of muonium.
- aim at 0.45 ppm
- different systematics



J-PARC g-2/edm

- new technique : slow muons from laser ionization of muonium.
- aim at 0.45 ppm
- different systematics



Muon $g-2$ Anomaly and Muon CLFV

muon $g-2$ anomaly

flavour conserving component of the BSM dipole operator

muon CLFV ($\mu \rightarrow e\gamma$ etc.)

flavour violating component of the BSM dipole operator

Muon $g-2$ Anomaly and Muon CLFV

muon $g-2$ anomaly

flavour conserving component of the BSM dipole operator

muon CLFV ($\mu \rightarrow e\gamma$ etc.)

flavour violating component of the BSM dipole operator

If the Muon $g-2$ anomaly is confirmed, it will establish the presence of a BSM muon interaction which may induce sizable effects of muon CLFV.

Summary



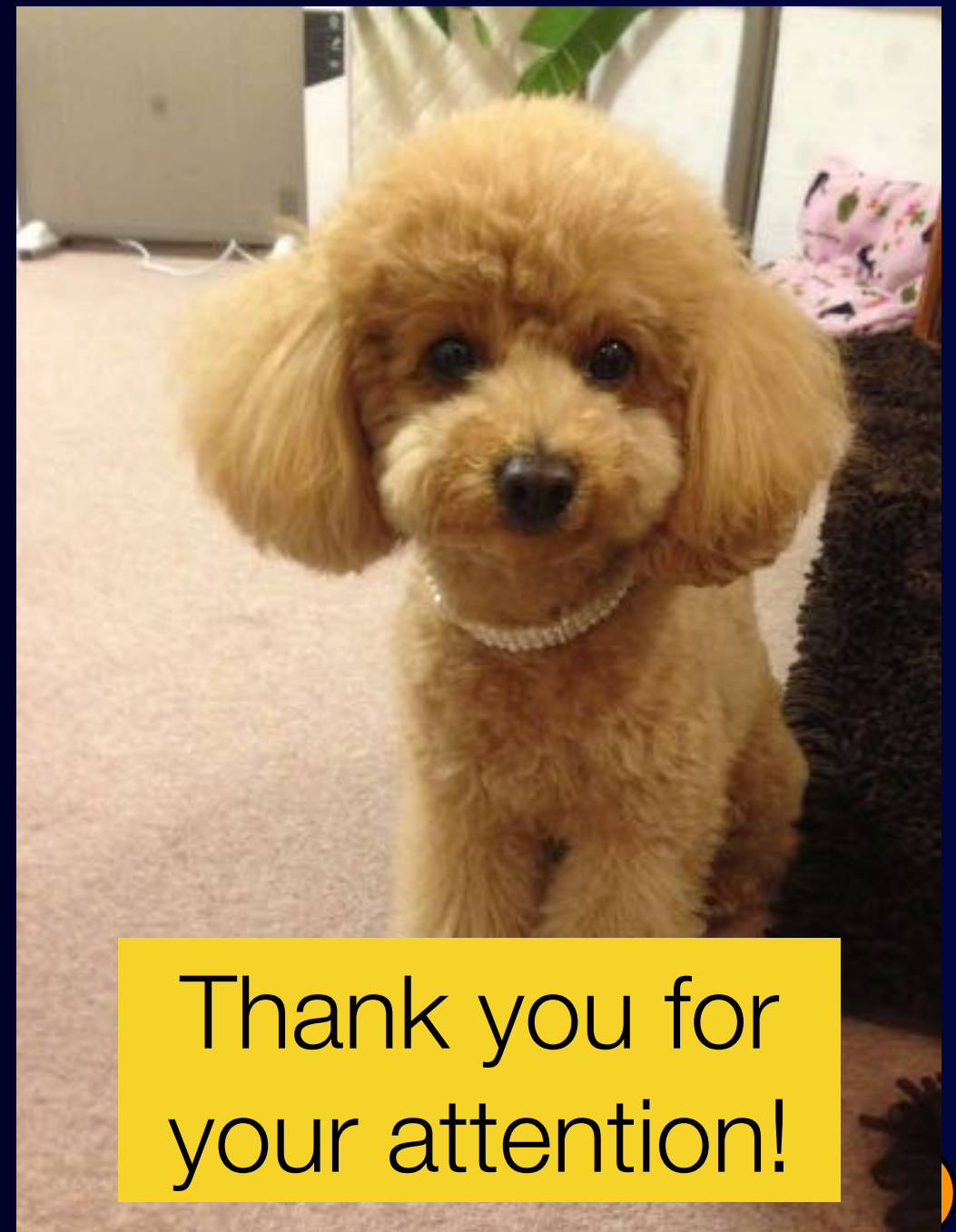
- Physics prospects with muons are rich.
- Study of charged leptons has sensitivity to BSM physics extending and complementing the reach of the LHC with a significant synergy with the neutrino program.
- Many projects will have their first results/start data taking in the next 2-3 years.

Summary



- Physics prospects with muons are rich.
- Study of charged leptons has sensitivity to BSM physics extending and complementing the reach of the LHC with a significant synergy with the neutrino program.
- Many projects will have their first results/start data taking in the next 2-3 years.

my dog, IKU



Thank you for
your attention!