



Instituto de
Física
Teórica
UAM-CSIC

Non-WIMP dark matter

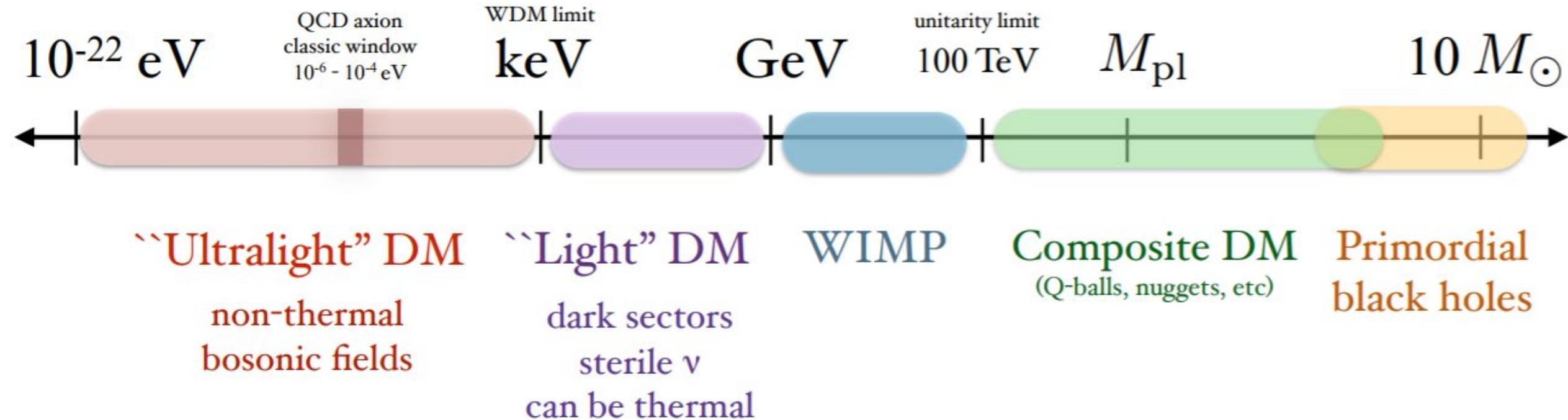
Zhen Liu
U of Maryland

IFT2019
future collider program
Jun. 12



Mass scale of dark matter

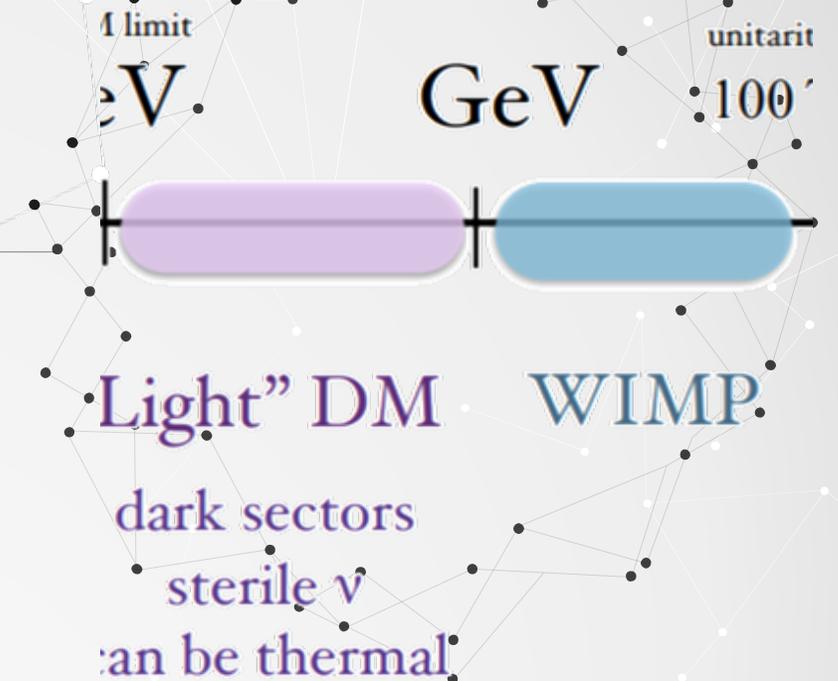
(not to scale)



Light Dark Matter

Moving beyond WIMPs, the broad vicinity of the weak scale is still an excellent place to focus on:

- An important scale!
- Familiar stable matter resides here!
- Thermal DM works well here!



Q: What's so great about equilibrium?

A: Generic and easy to achieve

Compare interaction rate
to Hubble expansion

$$\mathcal{L}_{\text{eff}} = \frac{g^2}{\Lambda^2} (\bar{\chi} \gamma^\mu \chi) (\bar{f} \gamma_\mu f)$$

$$H \sim n \sigma v \quad \Longrightarrow \quad \frac{T^2}{m_{Pl}} \sim \frac{g^2 T^5}{\Lambda^4} \Big|_{T=m_\chi}$$

Equilibrium is reached in the early universe if

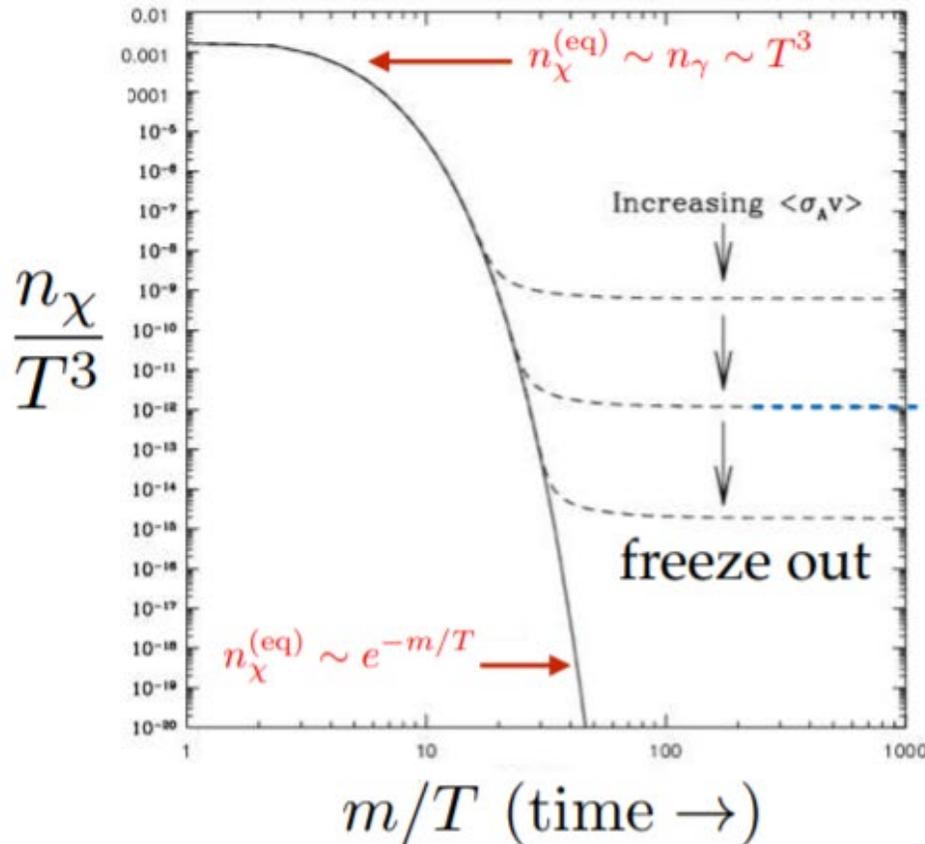
$$g \gtrsim 10^{-8} \left(\frac{\Lambda}{10 \text{ GeV}} \right)^2 \left(\frac{\text{GeV}}{m_\chi} \right)^{3/2}$$

Nearly all testable models feature equilibrium at early times

Q: What's so great about equilibrium?

A: Minimum annihilation rate

$$n_{\chi}^{(\text{eq})} = \int \frac{d^3 p}{(2\pi)^3} \frac{g_i}{e^{E/T} \pm 1} \propto \begin{cases} T^3 & (T \gg m) \\ e^{-m/T} & (T \ll m) \end{cases}$$



Observed density requires

$$\sigma v > 2 \times 10^{-26} \text{ cm}^3/\text{s}$$

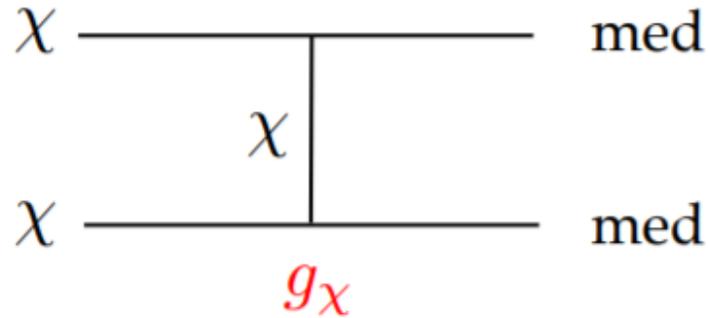
However, minimum target
in all equilibrium scenarios

... asymmetric DM
... coannihilating DM

Who's Heavier: DM or Mediator?

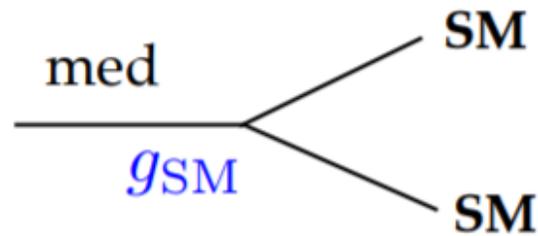
Secluded Annihilation

$$m_\chi > m_{\text{med}}$$



No clear experimental target

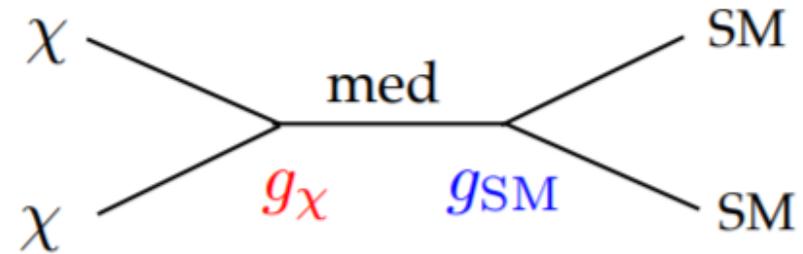
Abundance set by g_χ



Mediator decays **visibly**
Motivates hidden force searches

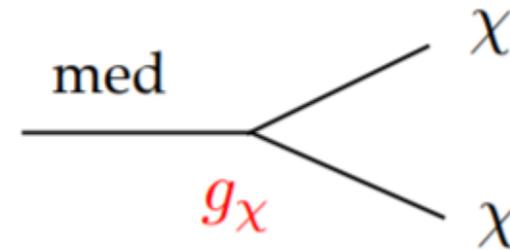
Direct Annihilation

$$m_\chi < m_{\text{med}}$$



Predictive thermal targets

Abundance depends on g_{SM}



Mediator decays **invisibly***
Motivates missing energy probes

Neutrality and Renormalizability require “portal” interactions

$\epsilon \phi H^\dagger H \longrightarrow$ **Scalar ϕ mixes with Higgs after EWSB**
Couples to SM masses $\epsilon \phi \frac{m_f}{v} \bar{f} f$

$\epsilon F'_{\mu\nu} F^{\mu\nu} \longrightarrow$ **Dark photon A' mixes with SM photon**
Couples to EM current $\epsilon A'_\mu J_{\text{EM}}^\mu$

$\epsilon V_\mu J_{\text{SM}}^\mu \longrightarrow$ **Vector V directly couples to DM & SM**
Couples to **different** current J_{SM}^μ

Anomaly free options $B - L$, $L_i - L_j$, $B - 3L_i$

Vector models all similar, but also couple to neutrinos

Kinetic mixing portal

1 limit

eV

GeV

unitarit

100'

"Light" DM

WIMP

dark sectors

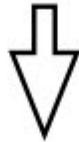
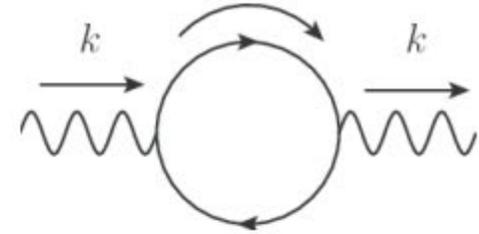
sterile ν

can be thermal

Millicharged particles (motivation)

$\mathcal{L} =$ a particle with charge ϵ .

heavy particles charged under both.



$$\epsilon F_{\mu\nu} F'_{\mu\nu}$$

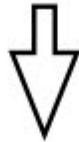
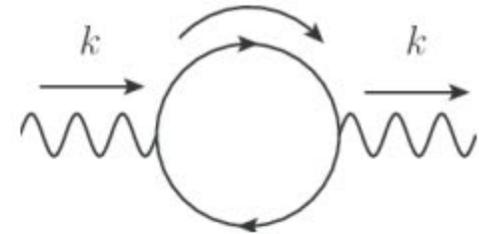
Our $U(1)_{EM}$
+ our matter

another massless $U(1)'$
+ matter'

Millicharged particles (motivation)

$\mathcal{L} =$ a particle with charge ϵ .

heavy particles charged under both.



$$\epsilon F_{\mu\nu} F'_{\mu\nu}$$

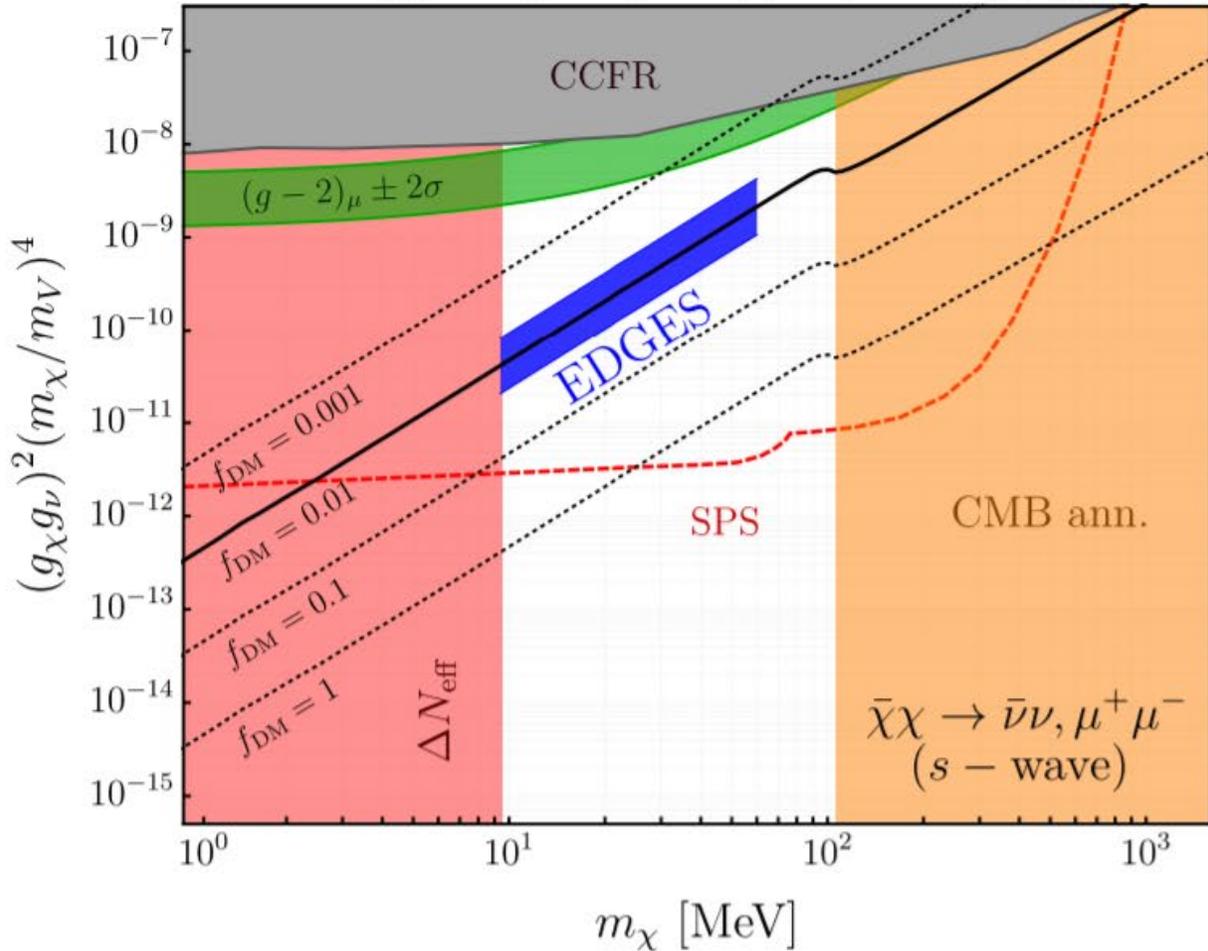
Our $U(1)_{EM}$
+ our matter

After we diagonalized everything, matter ' picks up a charge of ϵ under our EM.

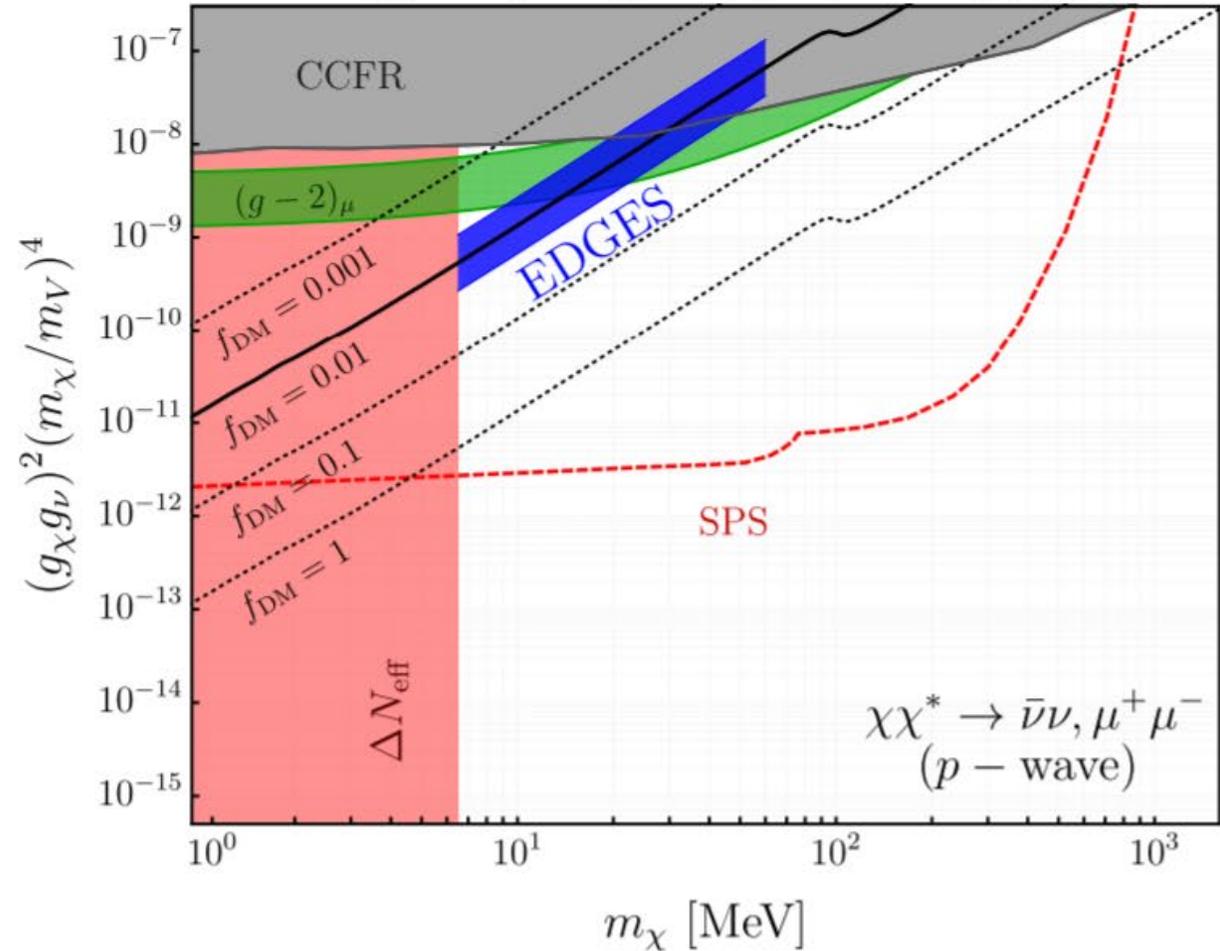
Millicharged particles (motivation)

$\mathcal{L} =$ a particle with charge ε .

Fermion χ , Gauged $L_\mu - L_\tau$, $m_V = 3m_\chi$, $g_\chi = 1$

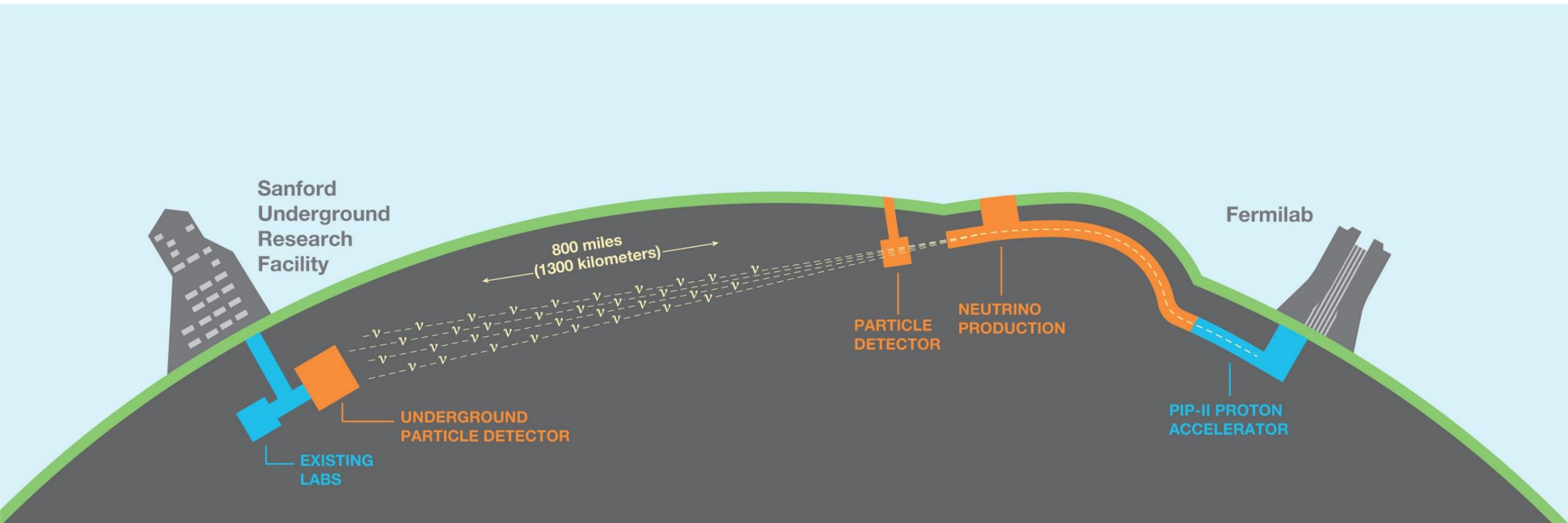


Scalar χ , Gauged $L_\mu - L_\tau$, $m_V = 3m_\chi$, $g_\chi = 1$



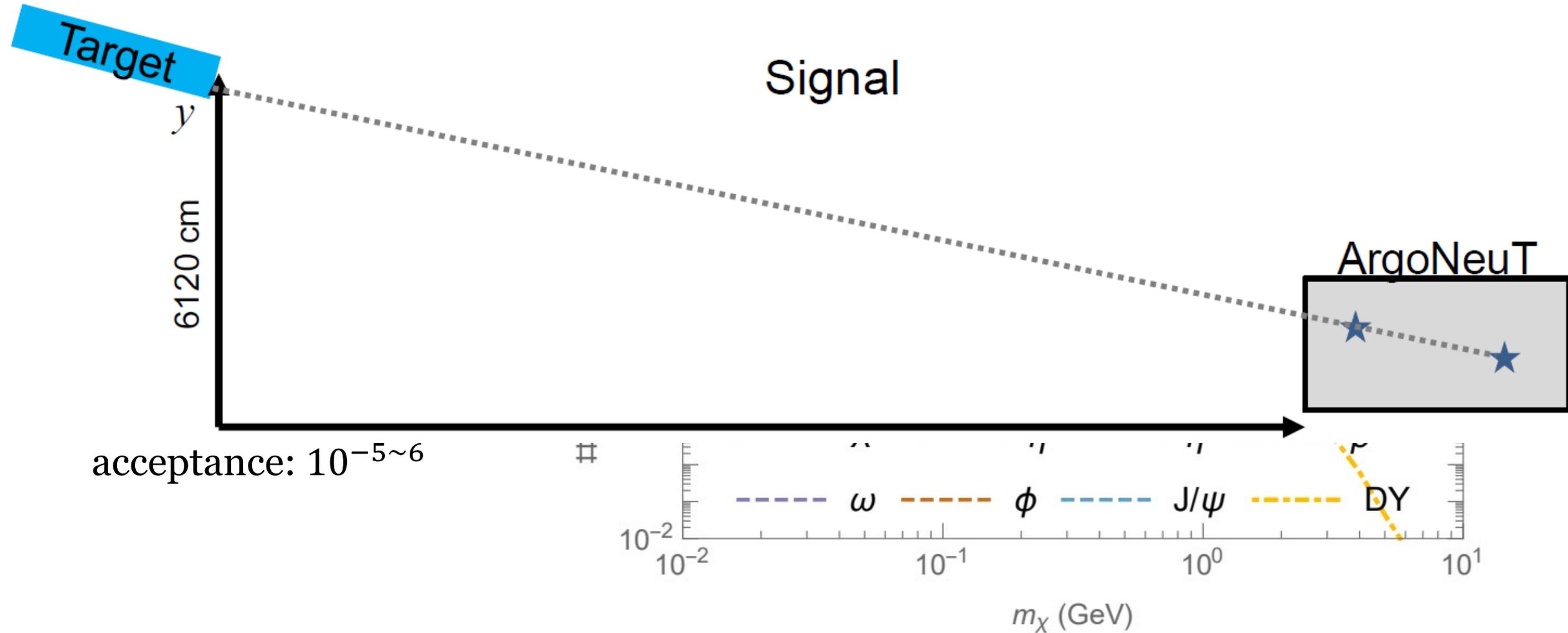
Neutrino experiments → High quality beamdump experiments

E.g., NuMI beam: good source for Millicharged particles



Neutrino experiments \rightarrow High quality beamdump experiments

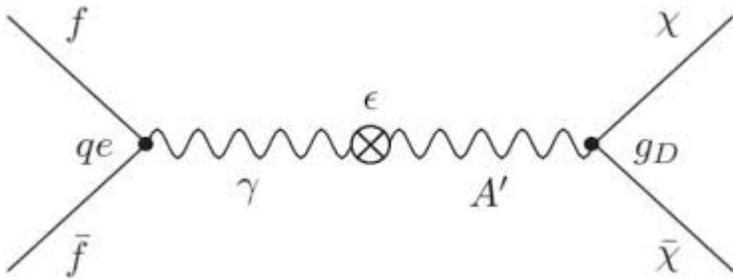
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Neutrino experiments \rightarrow High quality beamdump experiments

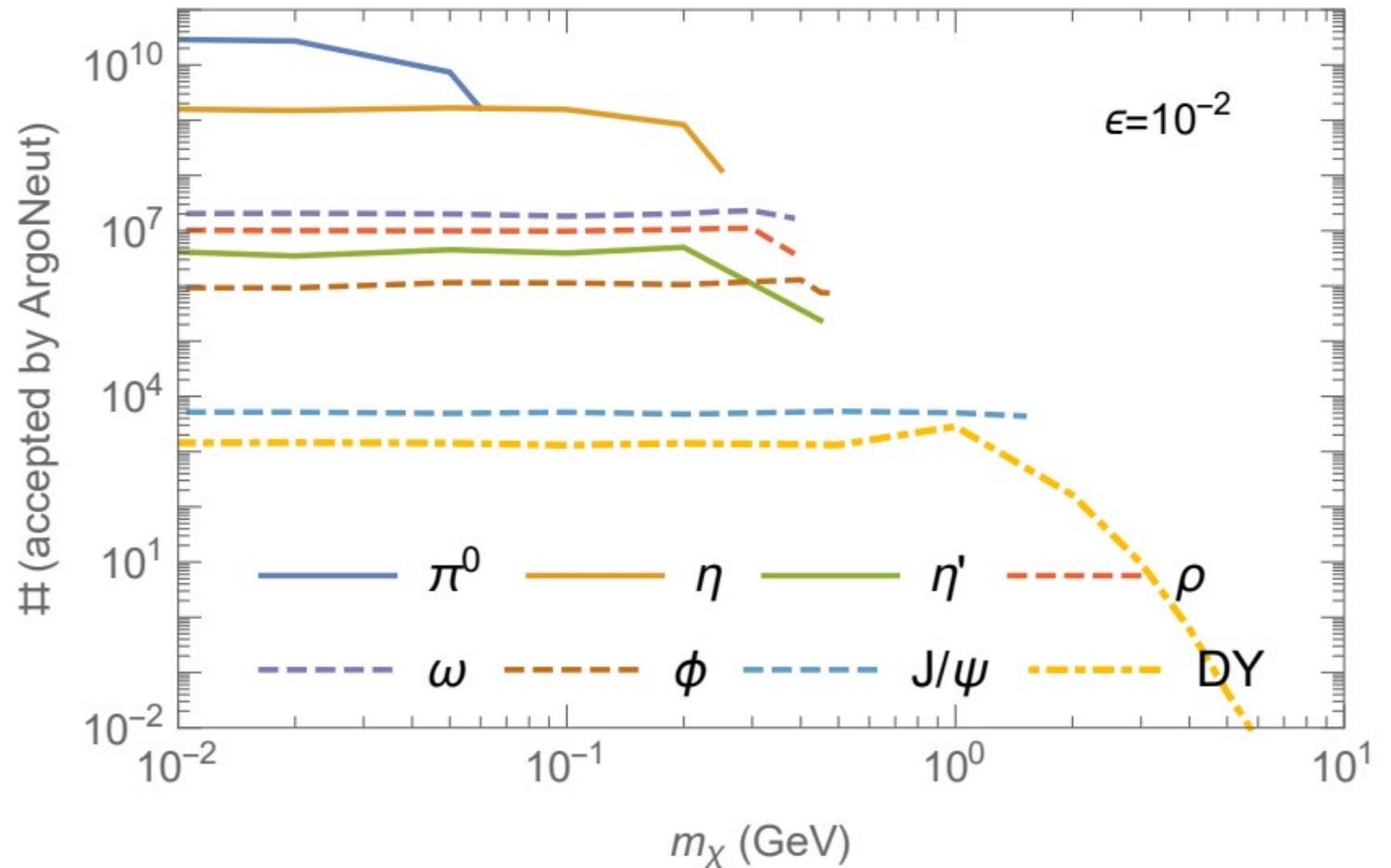
E.g., NuMI beam: good source for Millicharged particles

High beam energy
High POT



Typical geometric acceptance: $10^{-5\sim 6}$

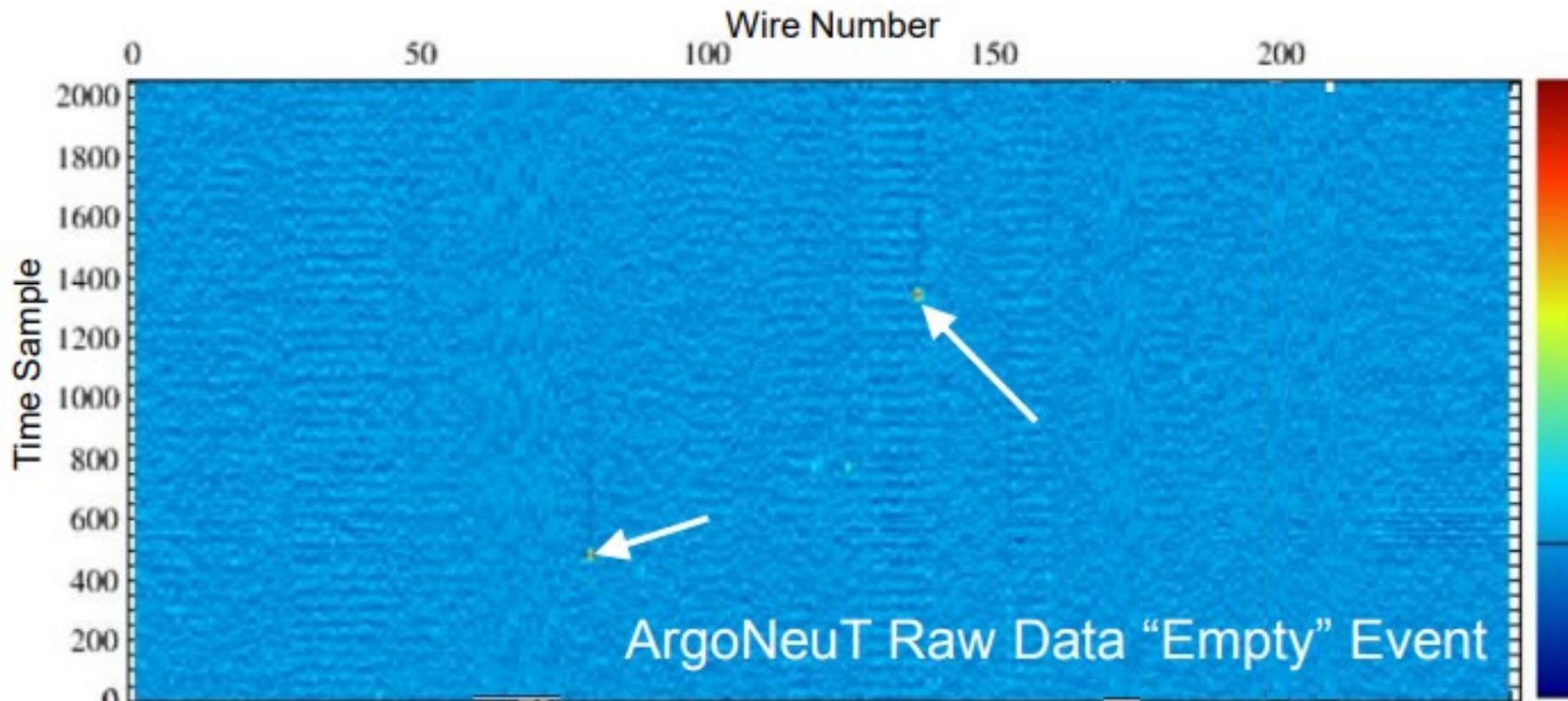
Milli-charged particle production 10^{20} POT 120 GeV



ArgoNeuT detector: low threshold & high resolution

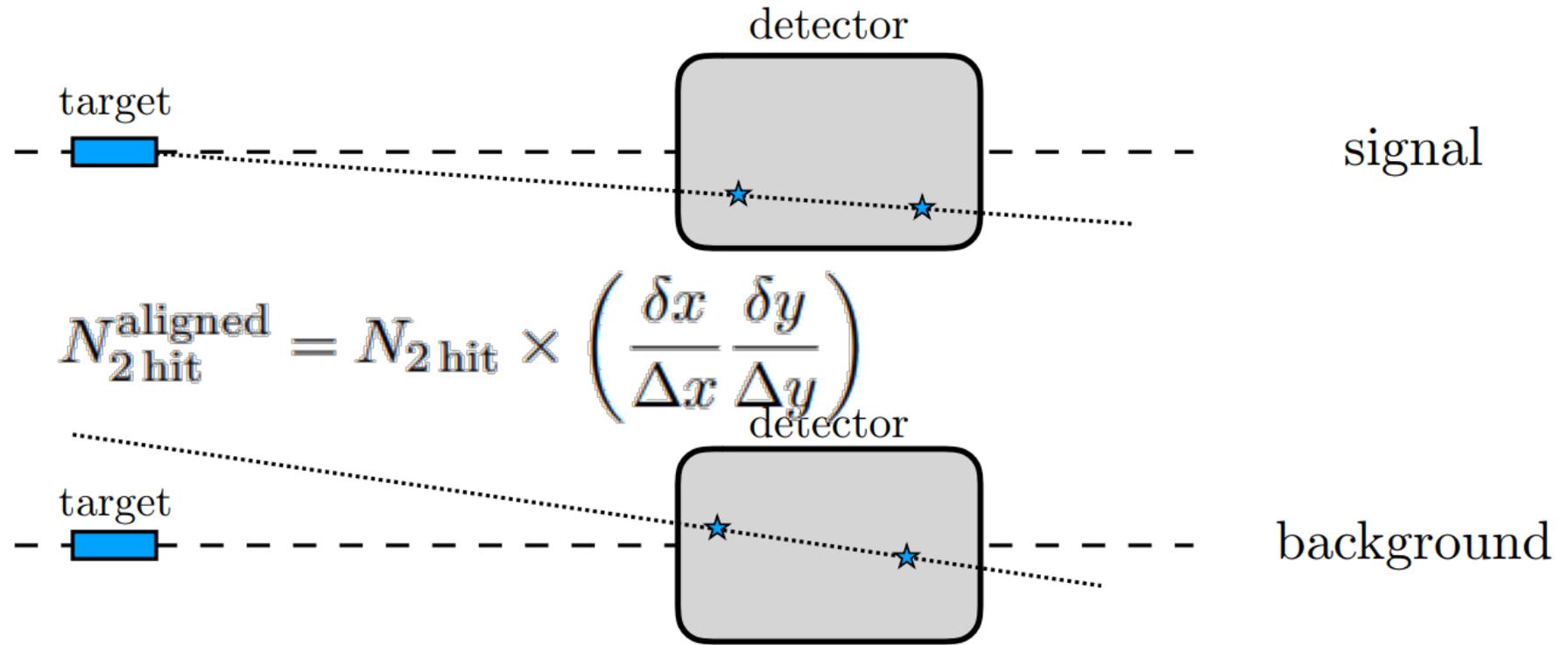
Look for energy depositions (“clusters”) above threshold (~300 keV).
Obtain 3D positions of these depositions.

$$\delta y \times \delta x \times \delta z = 5.6 \text{ mm} \times 0.3 \text{ mm} \times 3.2 \text{ mm}.$$



Signature

Dominated by low recoil energy scattering

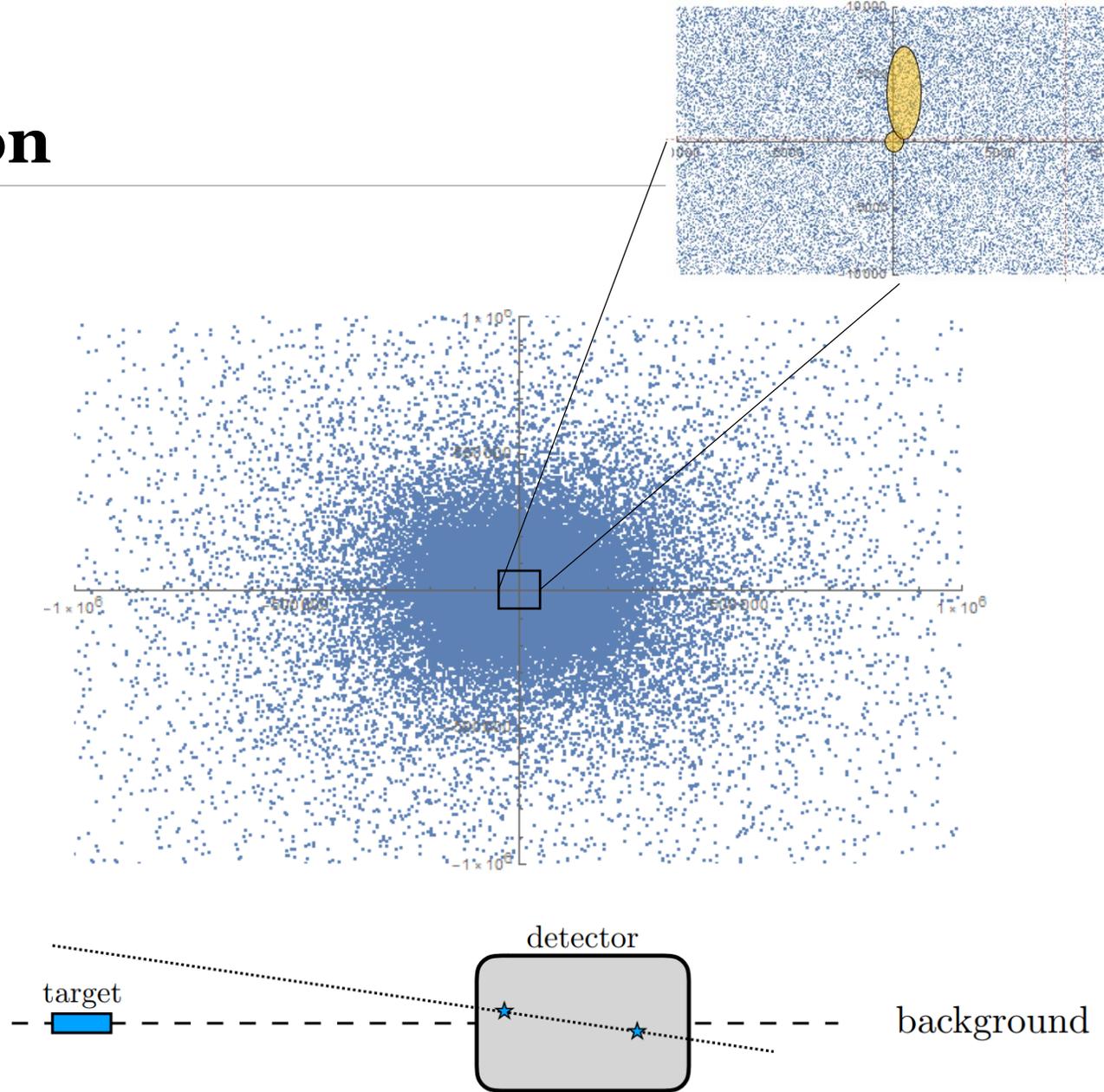
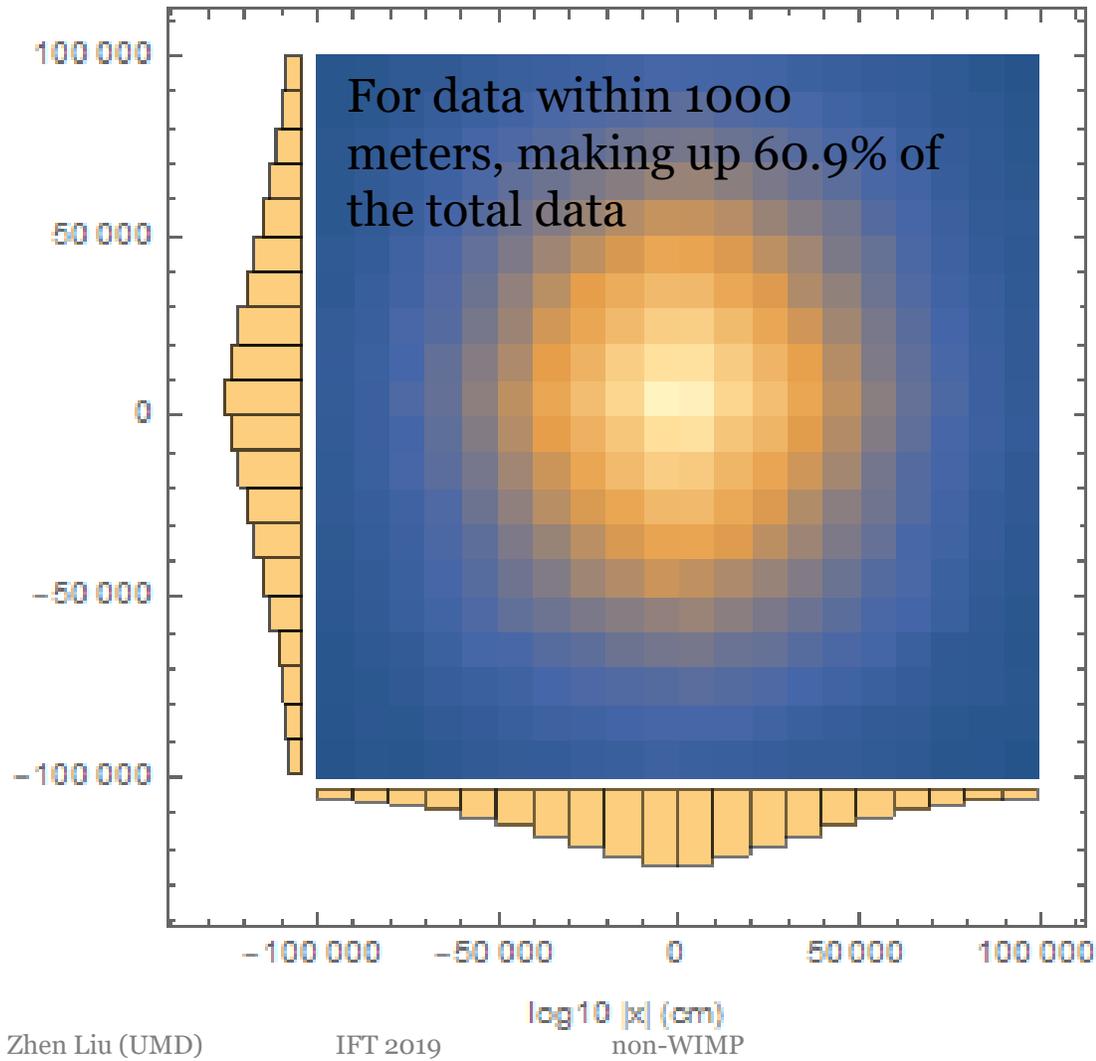


Signal scattering probability and mean free path

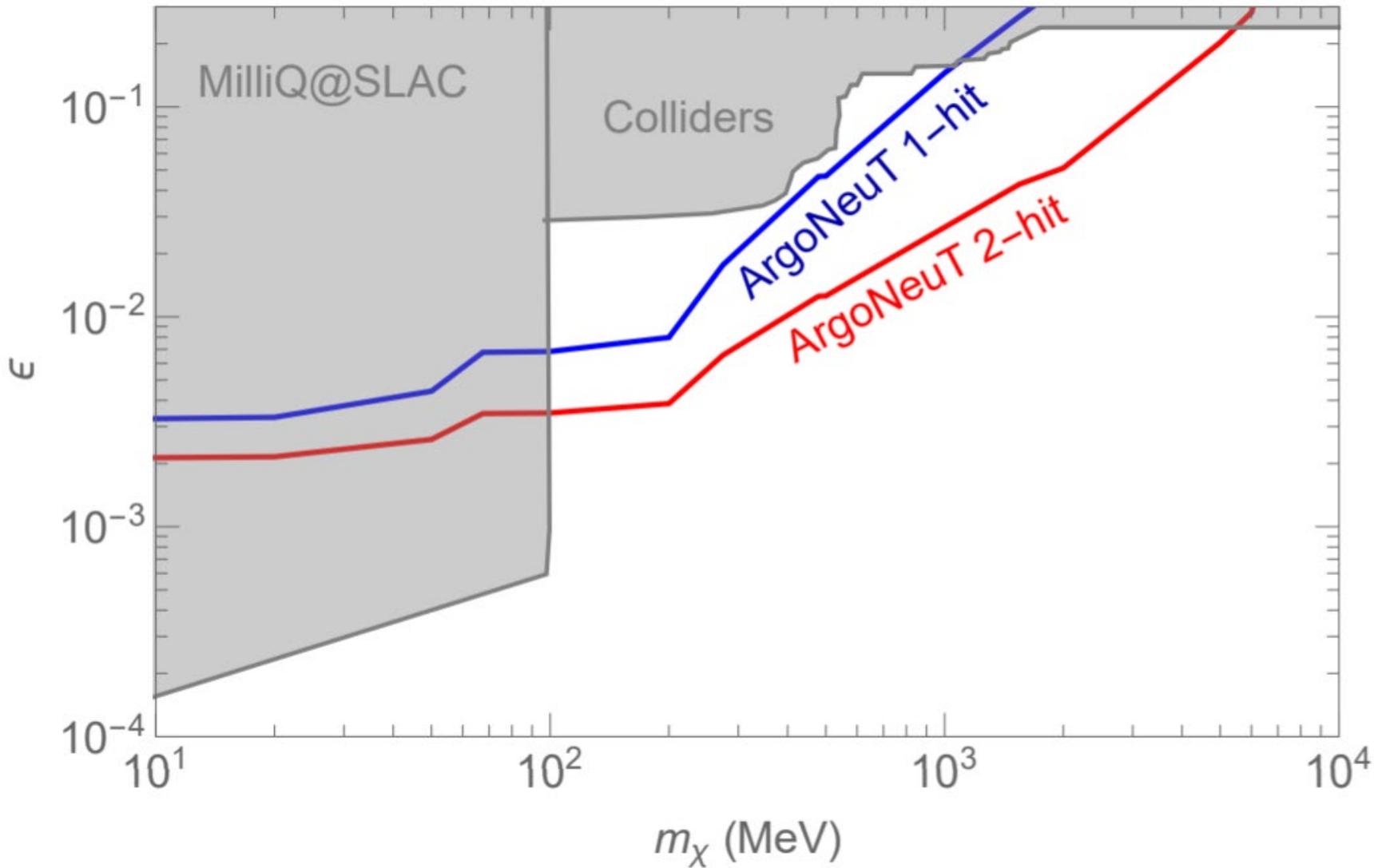
$$\left. \frac{d\sigma}{dE_r} \right|_{E_\chi \gg m_\chi, m_e, E_r} \simeq \frac{2\pi\alpha^2\epsilon^2}{E_r^2 m_e}$$

$$\lambda(E_r^{\text{min}}) \simeq \left(\frac{10^{-2}}{\epsilon} \right)^2 \left(\frac{E_r^{\text{min}}}{1 \text{ MeV}} \right) 1 \text{ km}$$

BKG target screen distribution



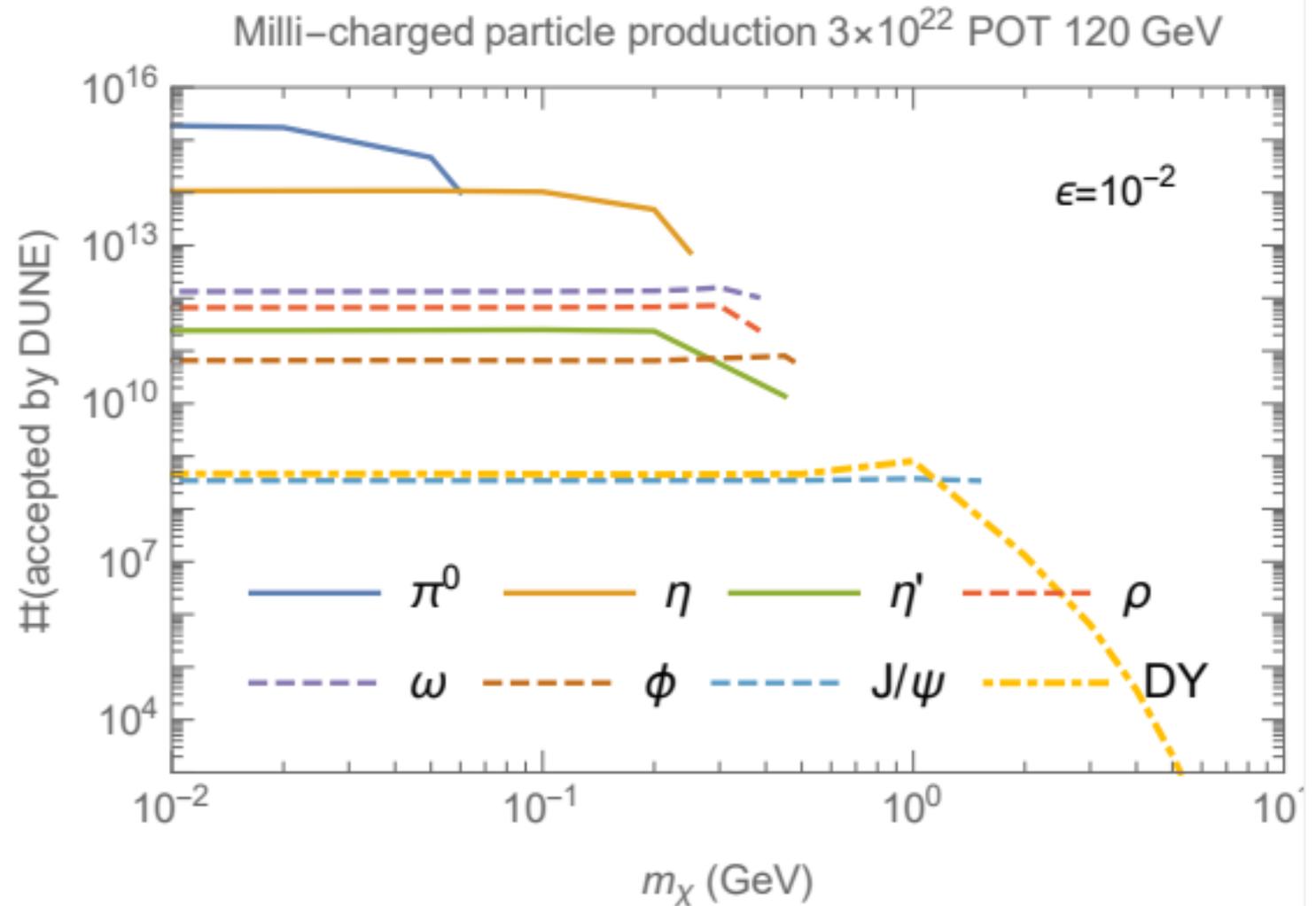
Expected reach



How about DUNE ND?

300 times more POT
240 times larger
detector
Factor of two closer
to the target

Should be very
promising!



Background scale non-trivially for DUNE ND

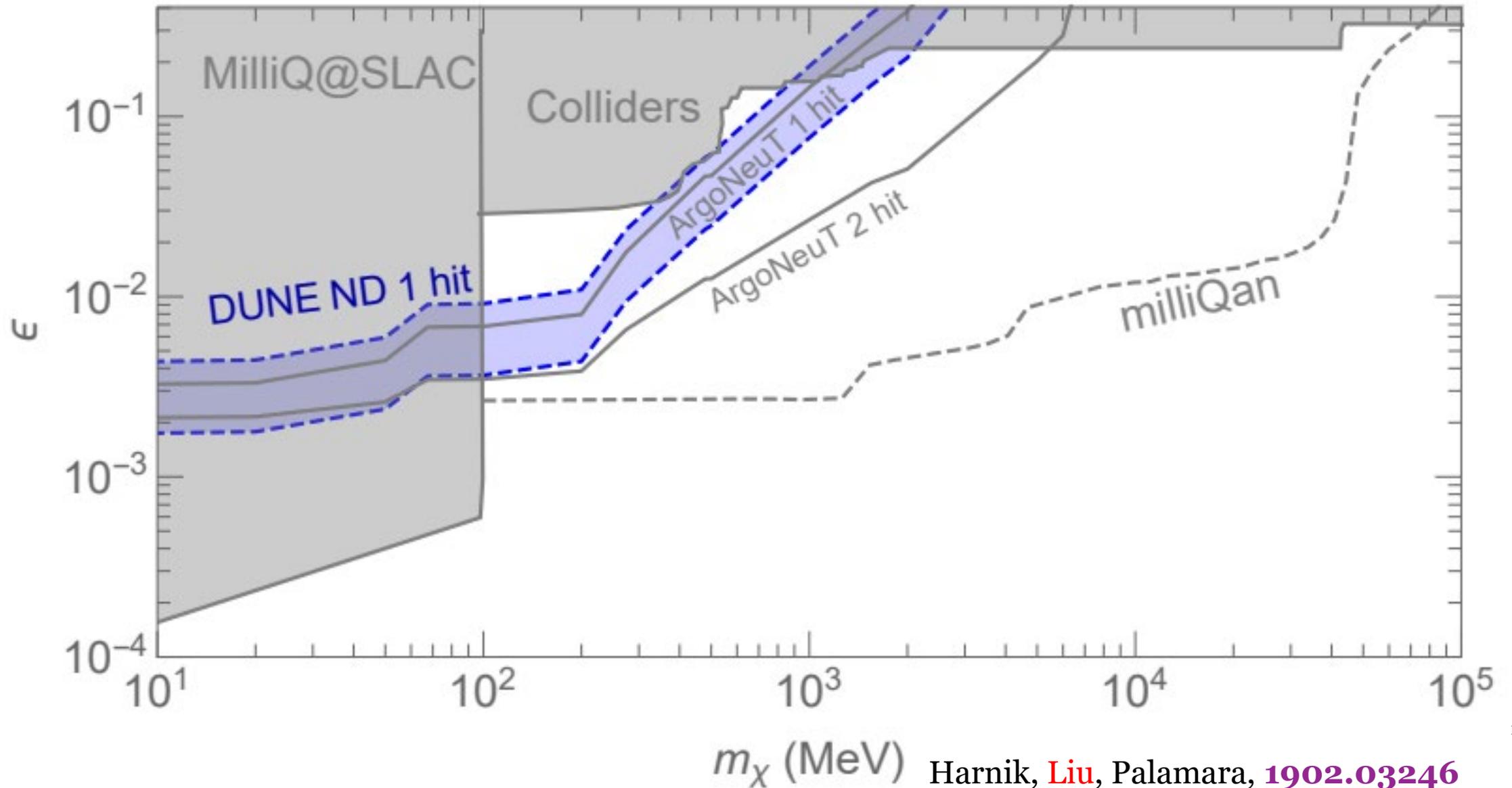
	Bkg Scaling	Bkg reduction	# frames with			# Background events				
			≥ 0 hit	≥ 1 hits	≥ 2 hits	Singlets	Doublets	Aligned doublets	Triplets	Aligned triplets
ArgoNeuT	Reference	Systematic	3.3×10^6	3.9×10^5	2.4×10^4	4.2×10^5	2.7×10^4	0.24	1.1×10^3	9.1×10^{-8}
DUNE ND	Volume	Systematic	1×10^8			4.5×10^9	1.0×10^{11}	1.4×10^4	1.6×10^{12}	0.030
		Statistic						$\sqrt{1.4 \times 10^4}$		0.030
		Timed	1×10^{10}	3.6×10^9	7.6×10^8	1.0×10^9	$\sqrt{1.4 \times 10^2}$	1.6×10^8	3.0×10^{-6}	
	Vol. \times Int.	Systematic	1×10^8			1.8×10^{11}	1.6×10^{14}	2.2×10^7	9.3×10^{16}	1.8×10^3
		Statistic						$\sqrt{2.2 \times 10^7}$		$\sqrt{1.8 \times 10^3}$
		Timed	1×10^{10}			1.6×10^{12}	$\sqrt{2.2 \times 10^5}$	9.3×10^{12}	0.18	

Average occupation number per frame:

ArgoNeuT: 0.13

DUNE ND: 45–1800

DUNE ND 1-hit projections



Background scale non-trivially for DUNE ND

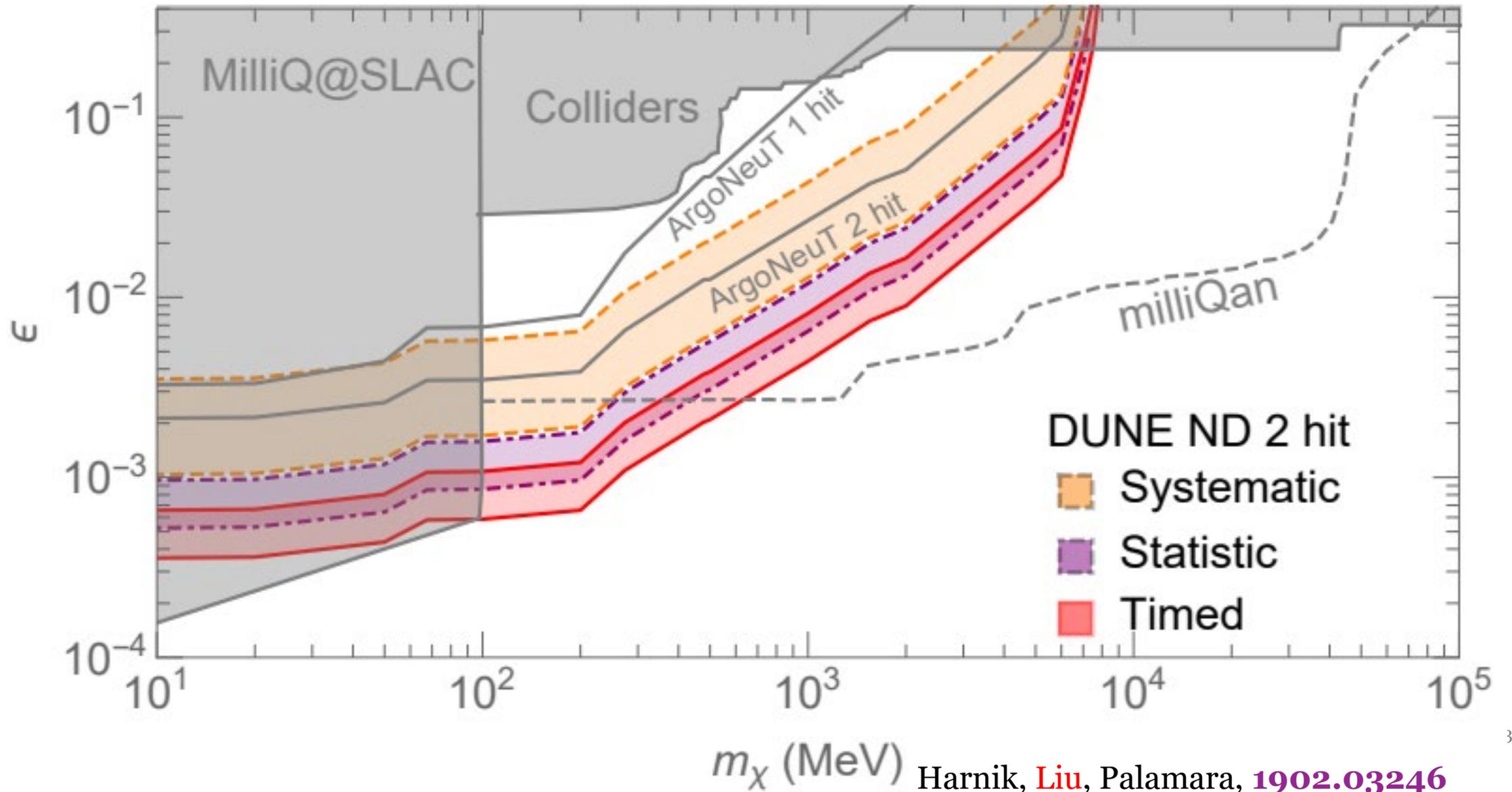
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Average occupation number per frame:

ArgoNeuT: 0.13

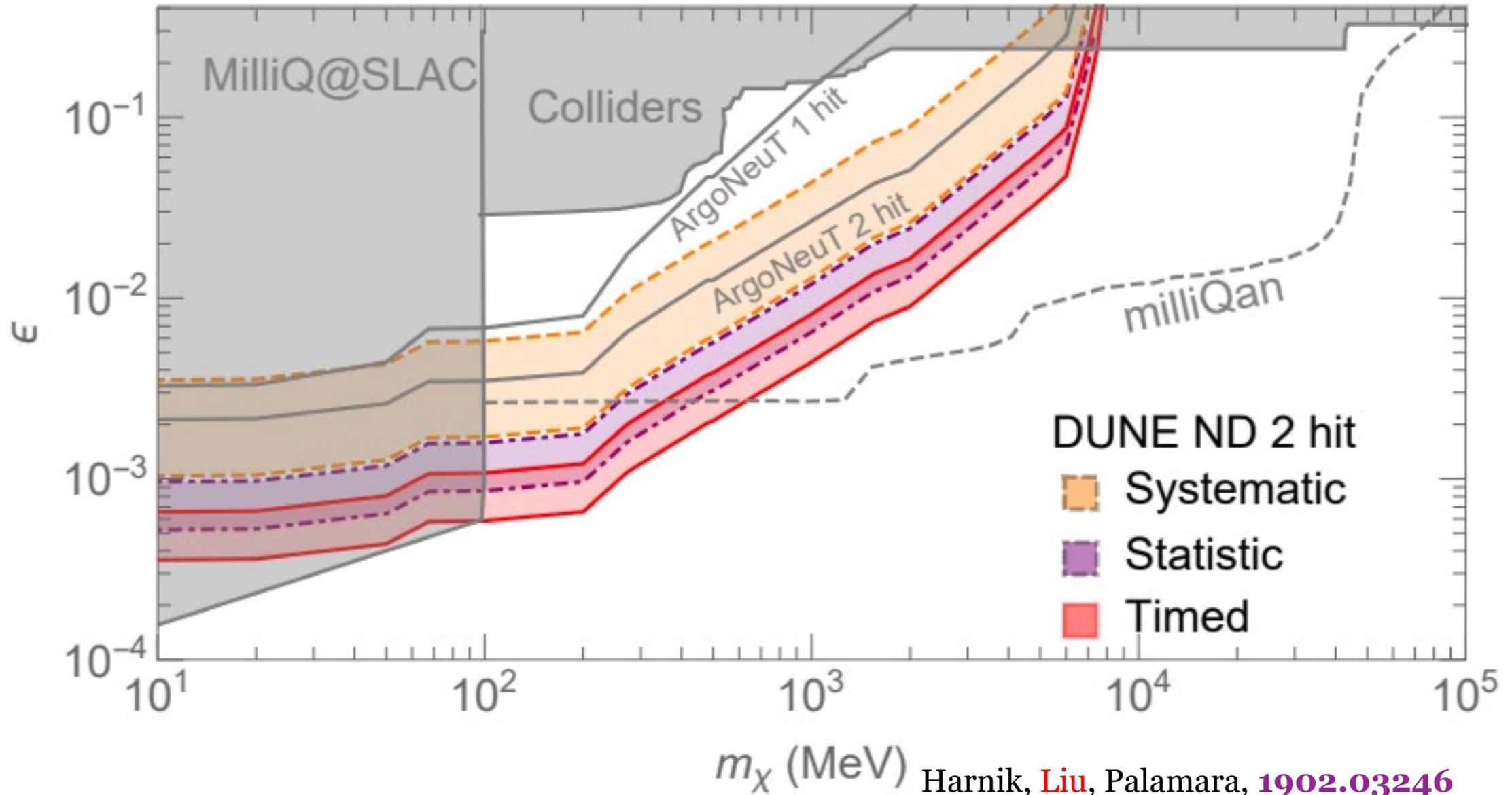
DUNE ND: 45–1800

DUNE ND double hit coverage

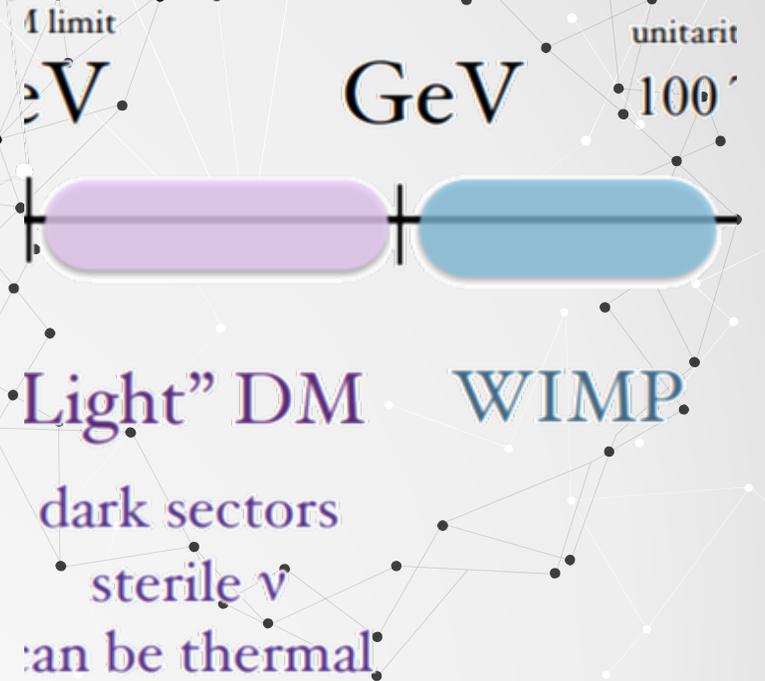


Active research field, see new works
ArgoNeuT result coming, I. Lepetic

DUNE ND double hit coverage



Heavy Neutral Lepton Portal

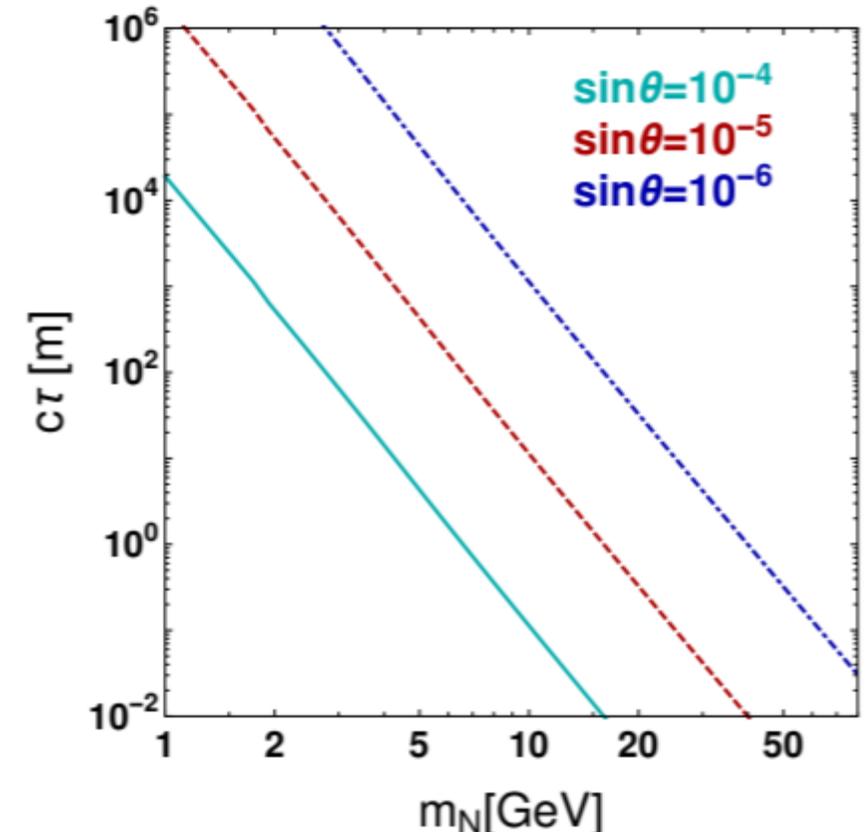


HNL a unique Physics case & LLP benchmark

**HNL are a well-motivated prototype LLP
they have to be studied as thoroughly as possible!**

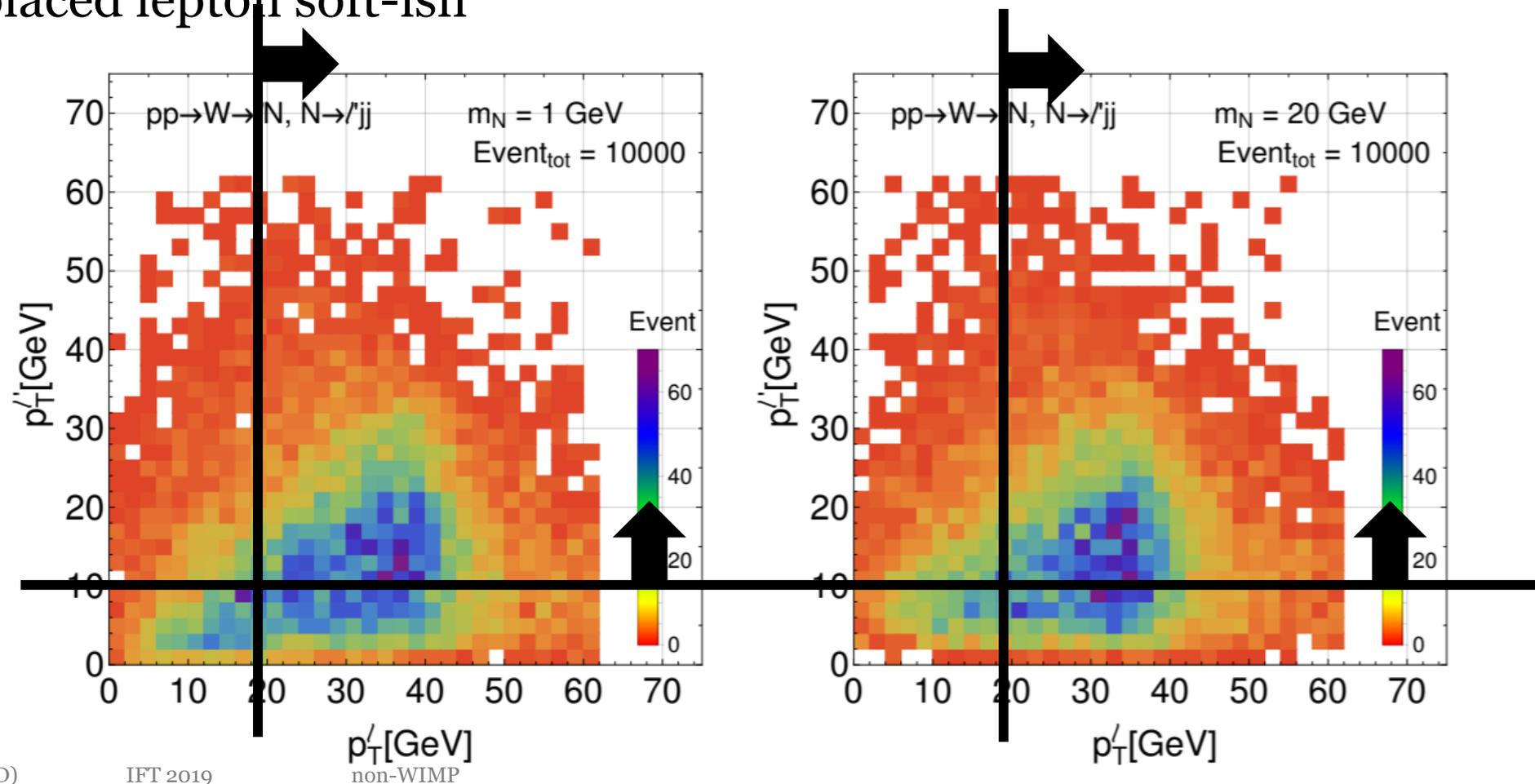
- Singly produced LLP
- Low mass LLP
- Prompt lepton trigger highly efficient
- Boosted low mass LLPs “fails”
“traditional” hard & slow displaced vertex searches

$$\Delta\mathcal{L}_\nu = -\lambda_\nu \bar{L}\tilde{H}N - \frac{m_N}{2}\bar{N}^c N + h.c.$$



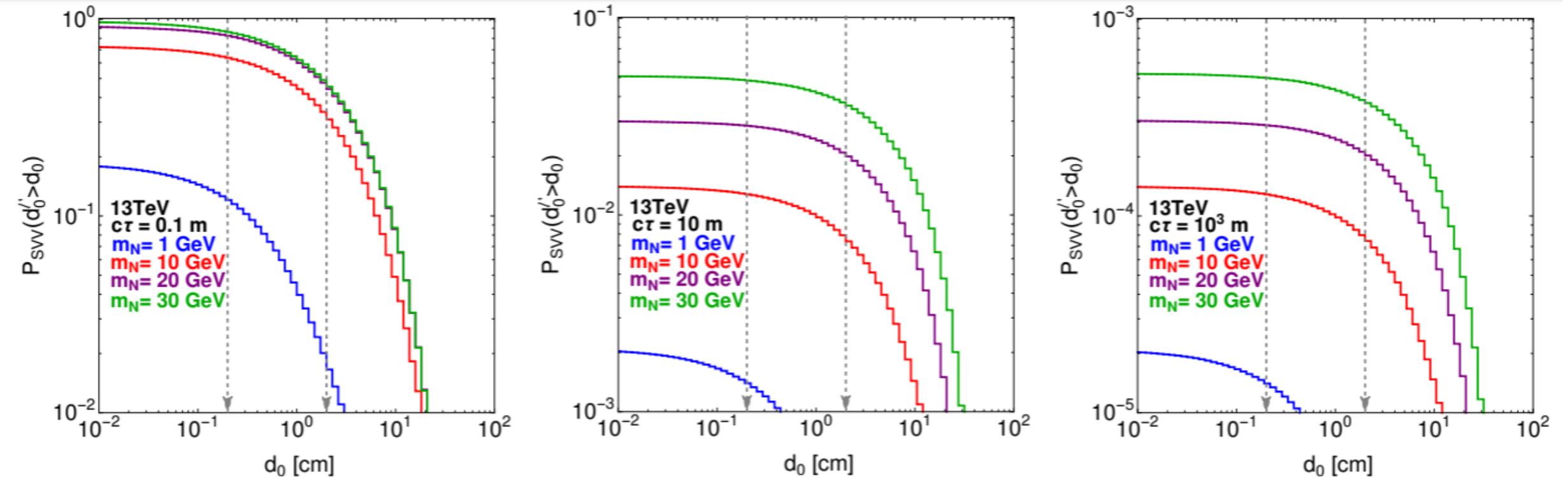
The lepton behaviors

- Prompt lepton hard-ish
- Displaced lepton soft-ish



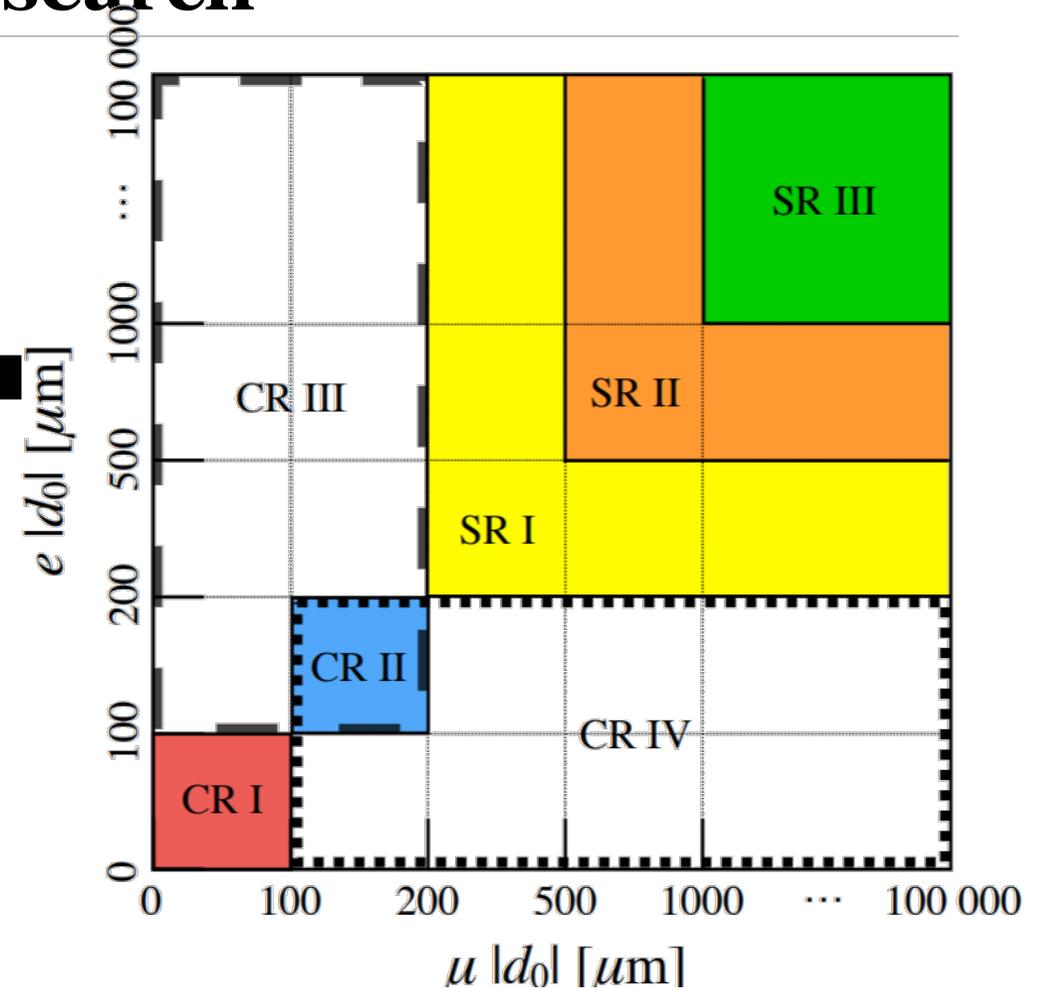
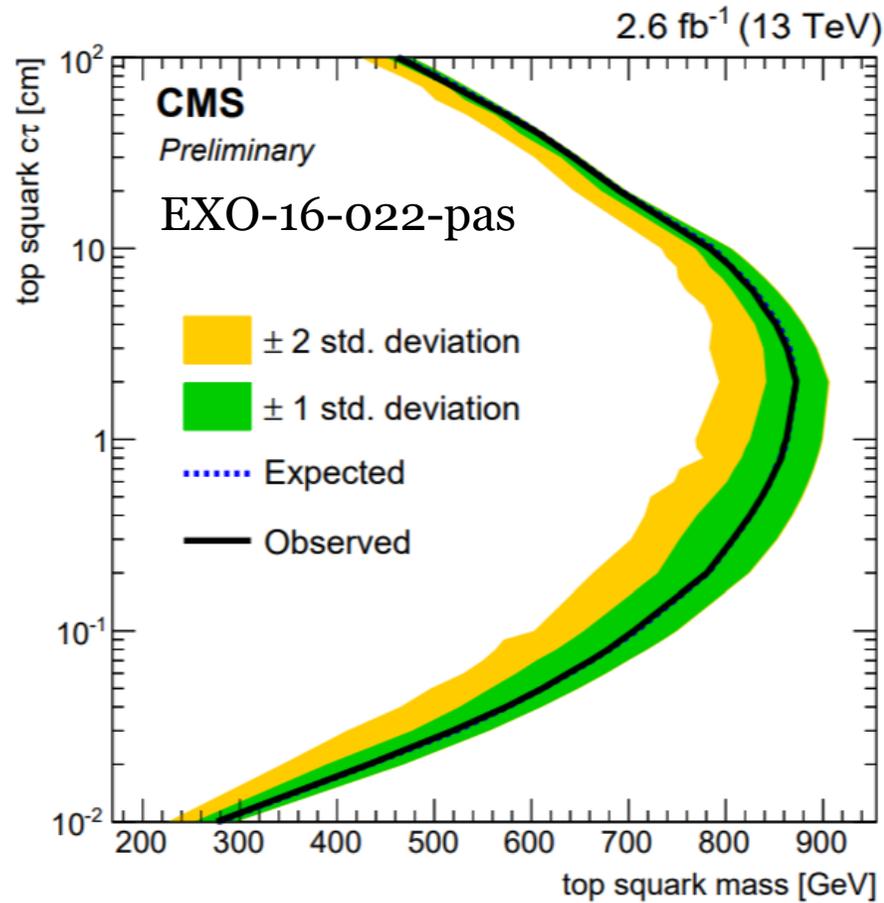
The lepton behaviors

Required decay within $R=0.5\text{m}$
to have good tracks



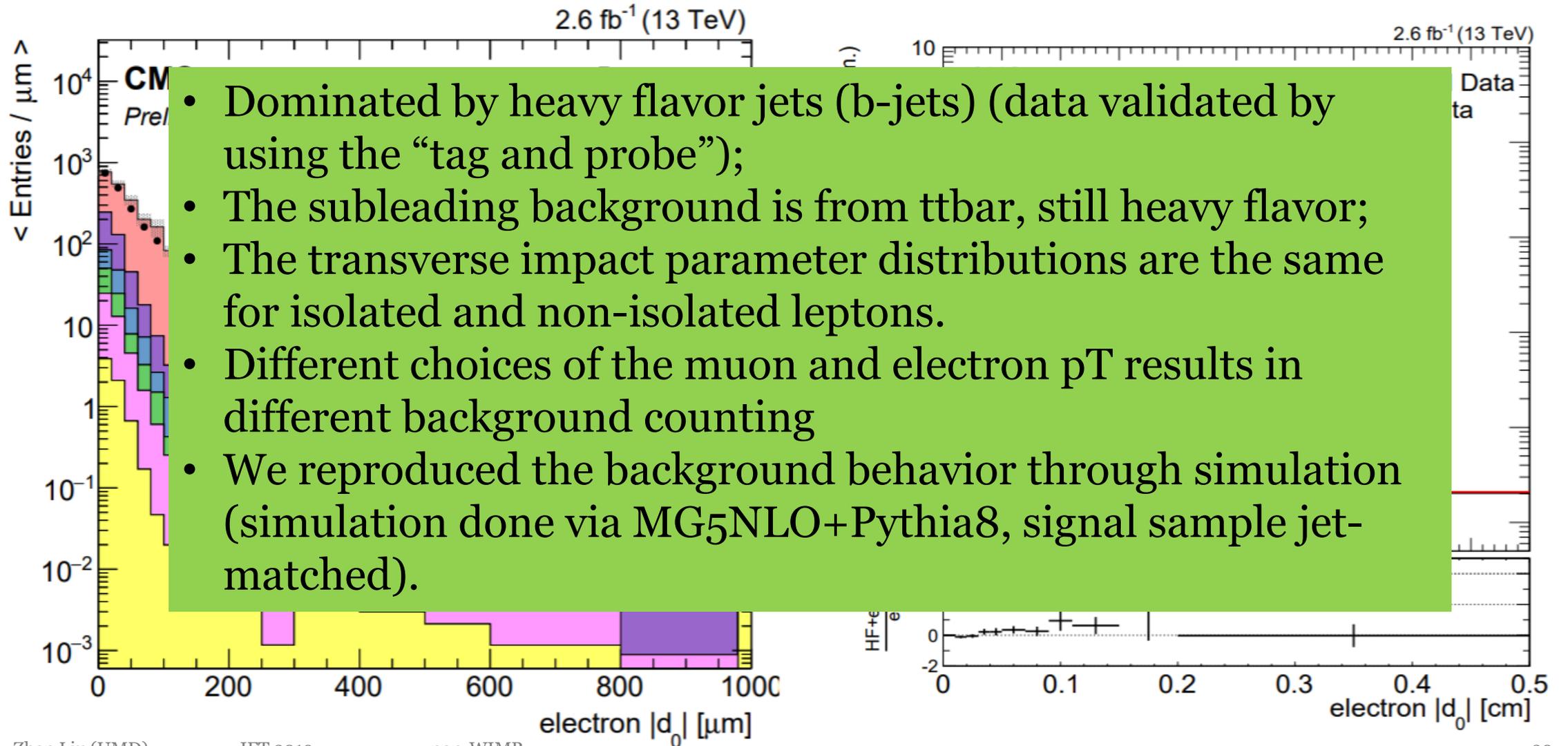
- Large d_0 cut, smaller signal efficiency;
- For short lifetime, $>10 \text{ GeV}$ sterile neutrinos behave similarly;
- For long lifetime, heavier sterile neutrinos are slower and hence higher decay probability within the tracker;
- For $m_N=1 \text{ GeV}$, decay product too collimated, suffering low d_0 ;

Valuable knowledge from a SUSY search



Displaced Electron Region (CR III)	Validation of HF Estimation	$ d_0 _e > 100 \mu\text{m}$ $ d_0 _\mu < 200 \mu\text{m}$
Displaced Muon Region (CR IV)	Validation of HF Estimation	$ d_0 _\mu > 100 \mu\text{m}$ $ d_0 _e < 200 \mu\text{m}$

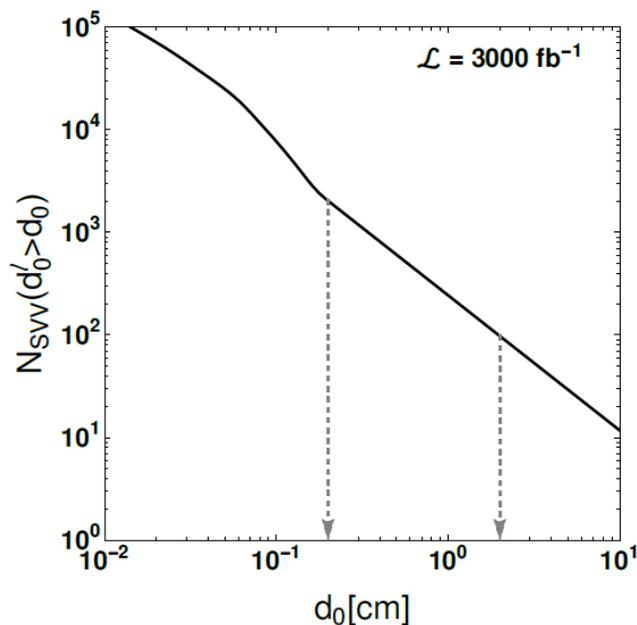
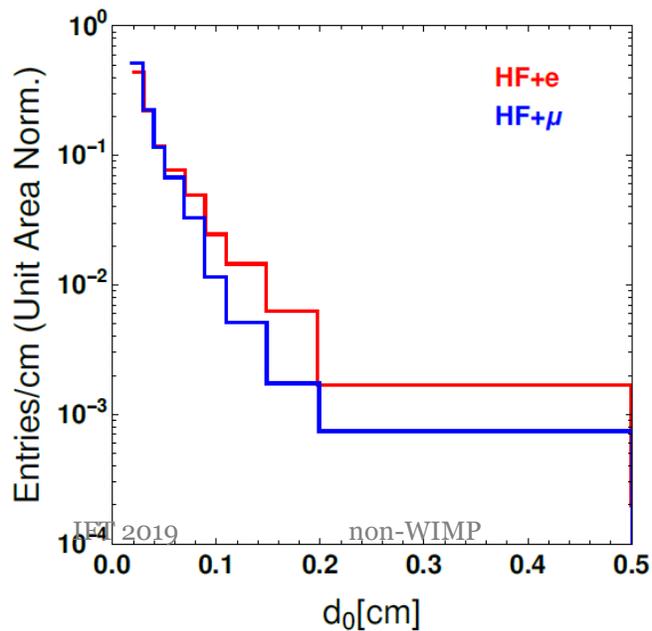
Valuable knowledge from a SUSY search



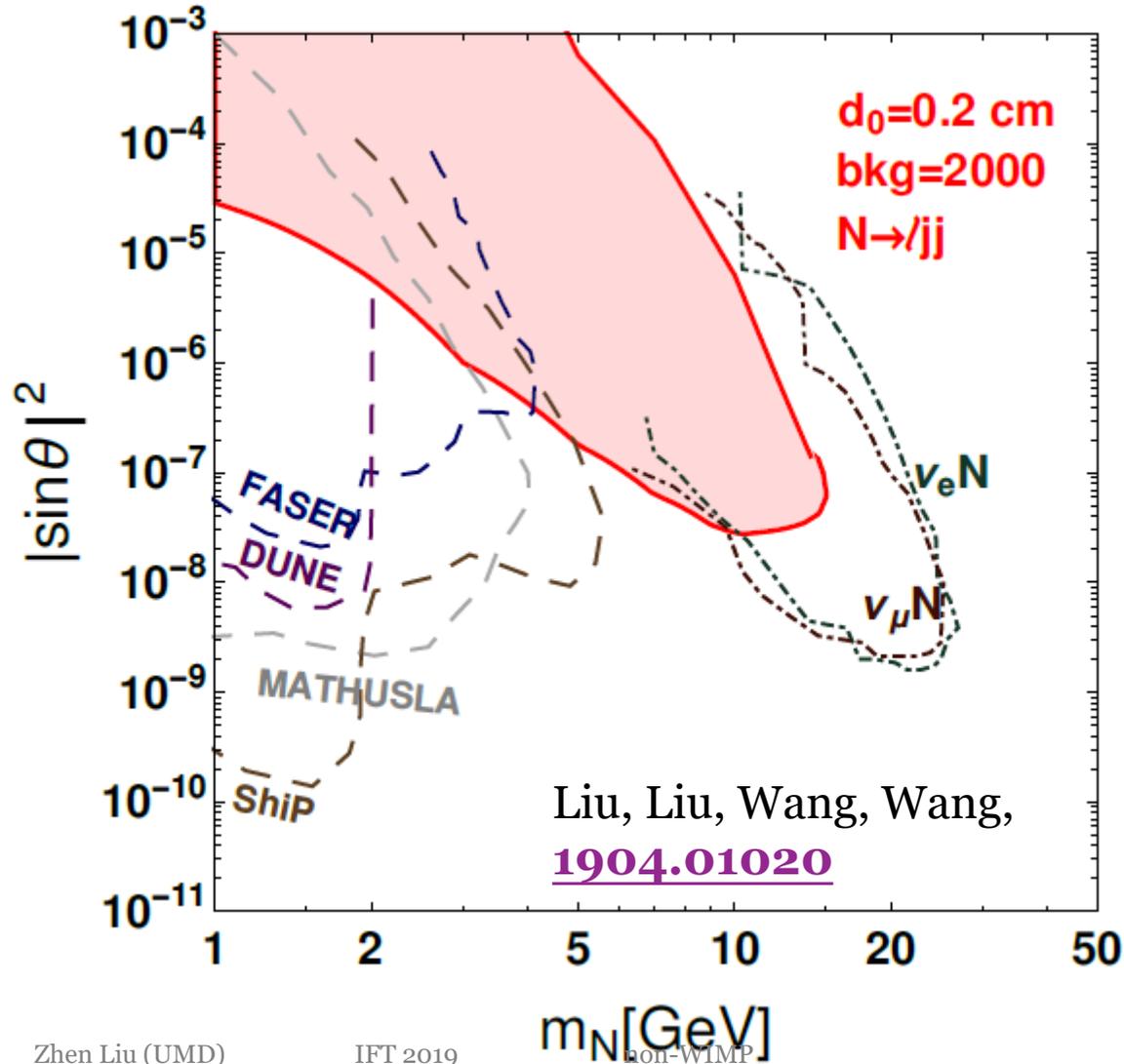
Our search and projected sensitivity

Efficiency	σ^{ncut} (pb)	$N_b^{30} = 0$	$N_j^{20} < 2$	$N_j^{50} = 0$	$H_T^{\text{vis}} < 100$ GeV	$p_T^{\ell_1} > 19$ GeV	$p_T^{\ell_2} > 10.5$ GeV	ϵ_{opt}
$t\bar{t} \rightarrow b\bar{b} + \ell + X$	136	0.25	0.08	0.62	0.43	0.055	0.42	1.2×10^{-4}
$W + b\bar{b}, W \rightarrow \ell\nu$	3.8	0.40	0.60	0.76	0.40	0.27	0.29	5.7×10^{-3}

$$N_{\text{bkg}} = \frac{\sigma_{\text{HF}+e}^{\text{CMS}} + \sigma_{\text{HF}+\mu}^{\text{CMS}}}{\sigma_{b\bar{b}(e)}^{\text{icut}} + \sigma_{b\bar{b}(\mu)}^{\text{icut}}} \left(\sigma_{W+b\bar{b}, W \rightarrow \ell\nu}^{\text{ncut}} \times \epsilon_{\text{opt}}^{W+b\bar{b}} + \sigma_{t\bar{t} \rightarrow b\bar{b} + \ell + X}^{\text{ncut}} \times \epsilon_{\text{opt}}^{t\bar{t}} \right) \times \mathcal{L}_{\text{HL-LHC}}$$



Results



- Non-zero background at HL-LHC: $\sim 2\text{K}$
- Interesting expansion of the LHC coverage to lower masses ($< 5 \text{ GeV}$) by taking the heavy flavor background directly;
- An example of “serious” phenomenological new search studies on LLPs can be done at the LHC;

Axion-like particle portal

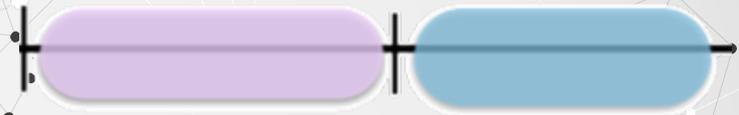
1 limit

eV

GeV

unitarity

100'



Light" DM

WIMP

dark sectors

sterile ν

can be thermal

Strong CP puzzle can lead to long-lived axions

$$L \supset \frac{\alpha_s}{4\pi} \theta \tilde{G}G + y_u \bar{Q}_L \tilde{H} u_R + y_d \bar{Q}_L H d_R$$

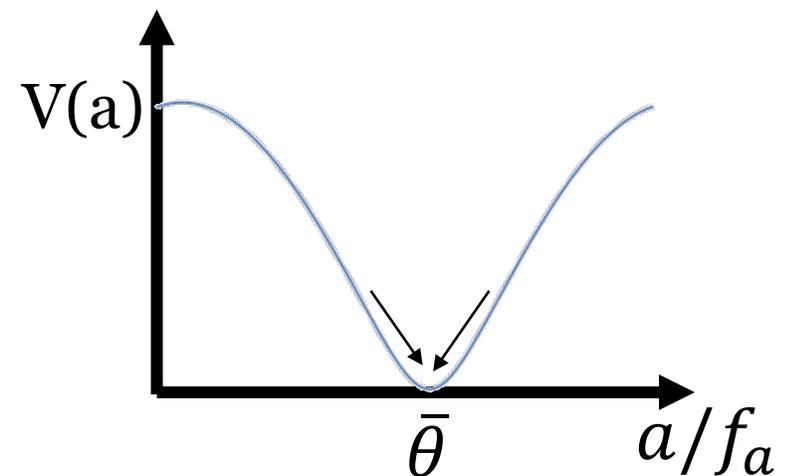
$$\bar{\theta} \equiv \theta + \text{ArgDet}[Y_u Y_d] \leq 10^{-10}$$

While $\text{ArgDet}[Y_u Y_d]$ anticipated around $\delta_{CKM} \sim O(1)$

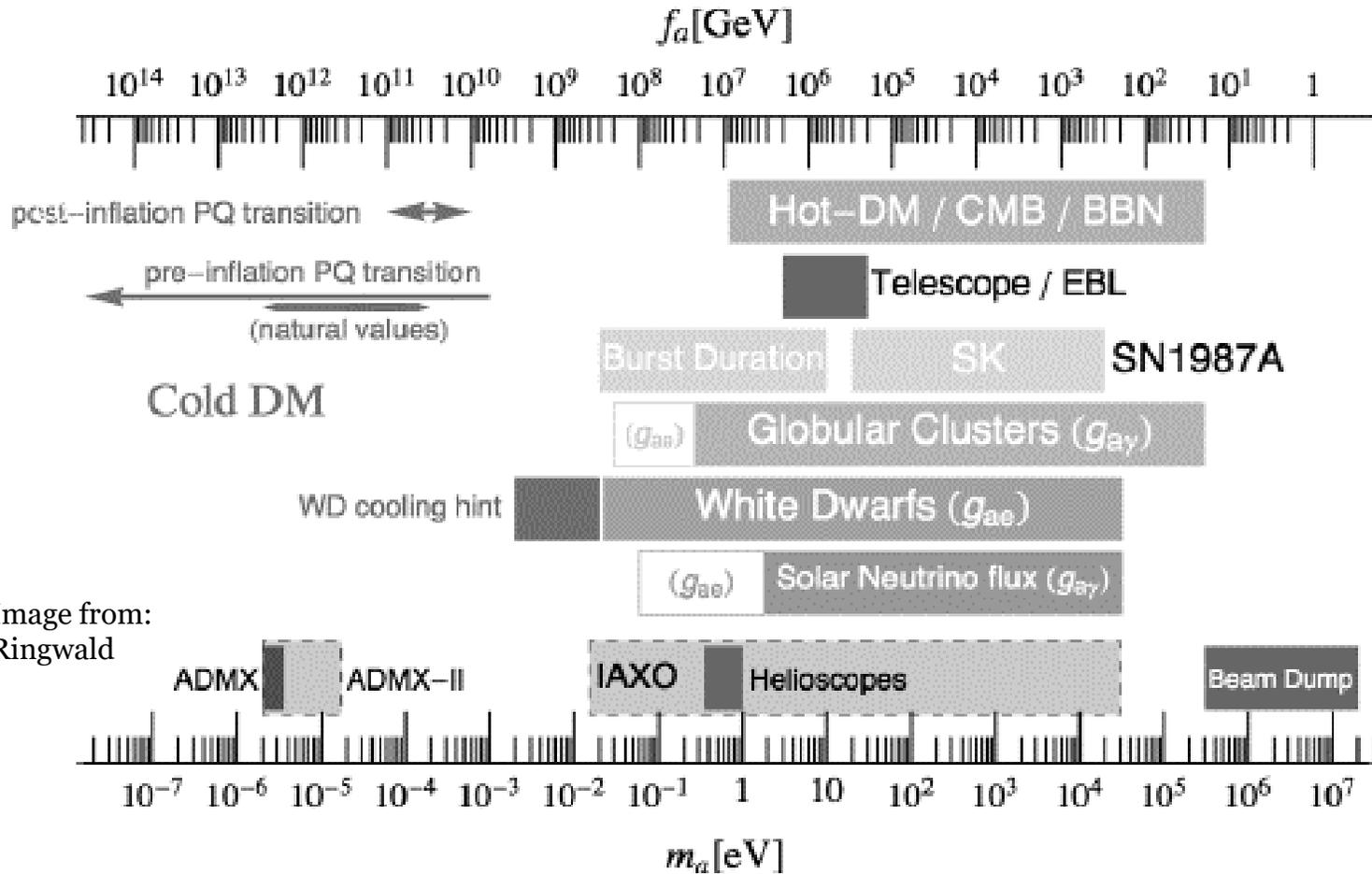
Strong CP puzzle of QCD

Dynamical solution:
QCD Axion a as a pseudo
Nambu-Goldstone boson

$$\frac{\alpha_s}{4\pi} \left(\theta - \frac{a}{f_a} \right) \tilde{G}G$$



Strong CP puzzle can lead to long-lived axions



Dynamical solution:
 QCD Axion a as a pseudo
 Nambu-Goldstone boson

$$\frac{\alpha_s}{4\pi} \left(\theta - \frac{a}{f_a} \right) \tilde{G}G + \frac{\alpha}{4\pi} \frac{a}{f_a} \tilde{F}F$$

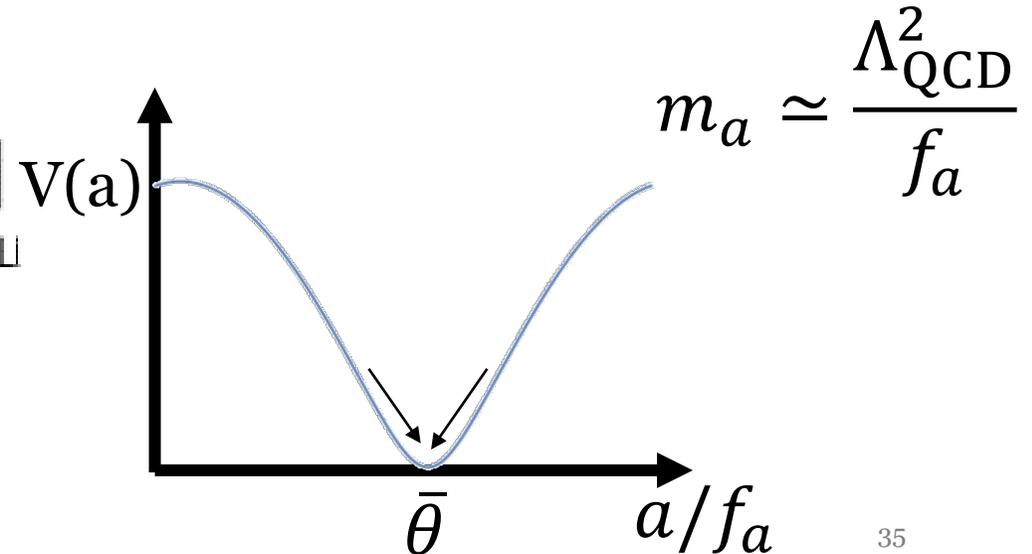
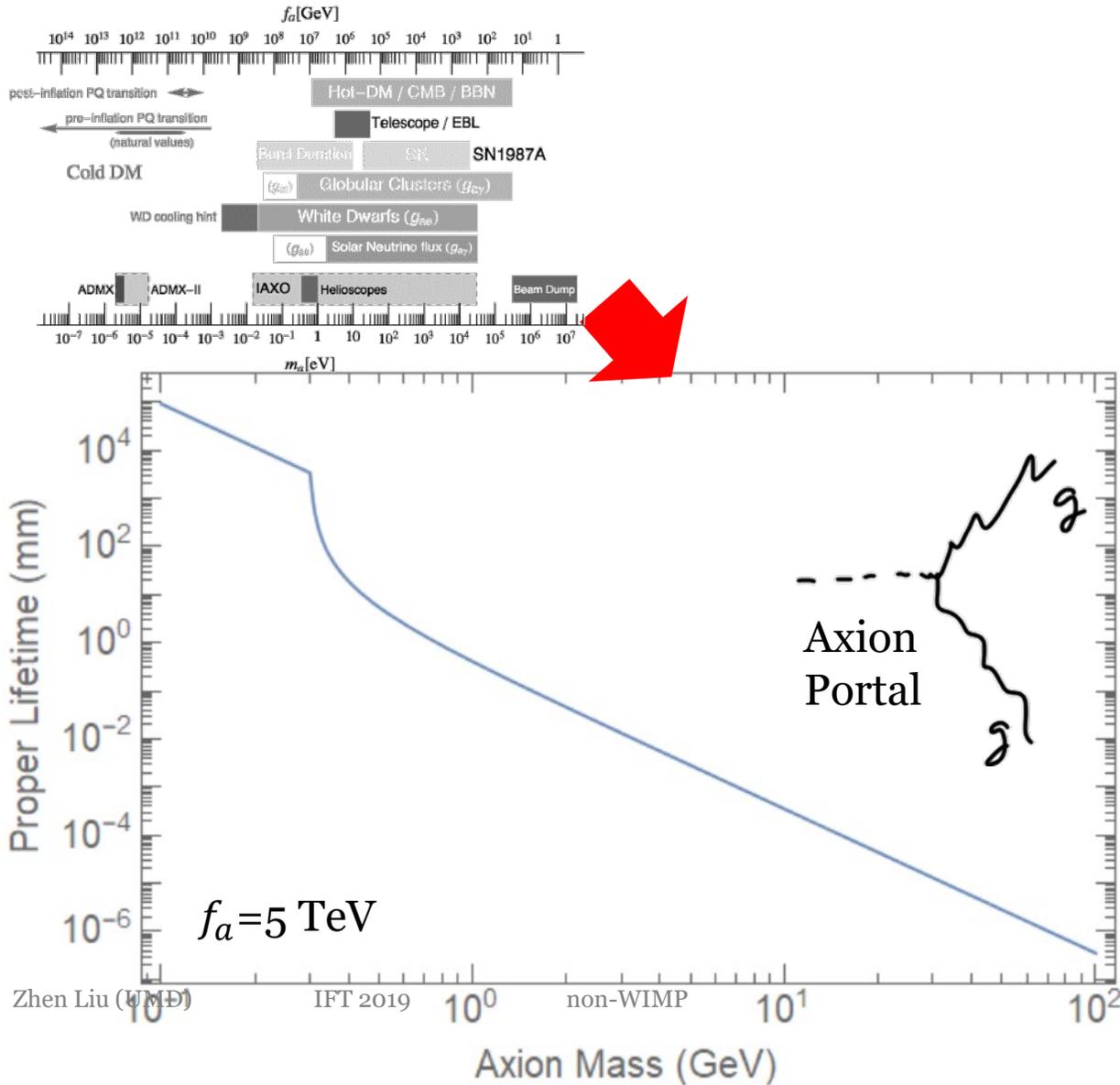


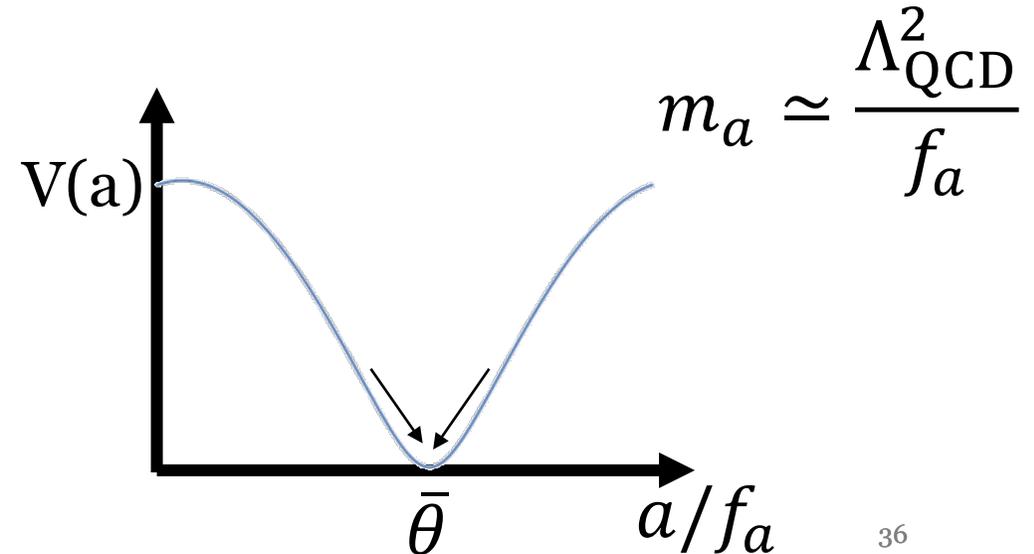
Image from:
 Ringwald

Strong CP puzzle can lead to long-lived axions !



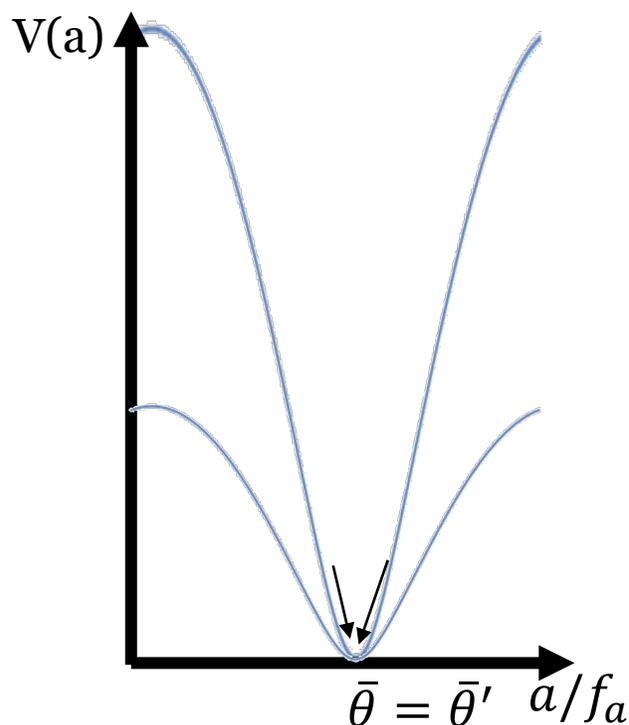
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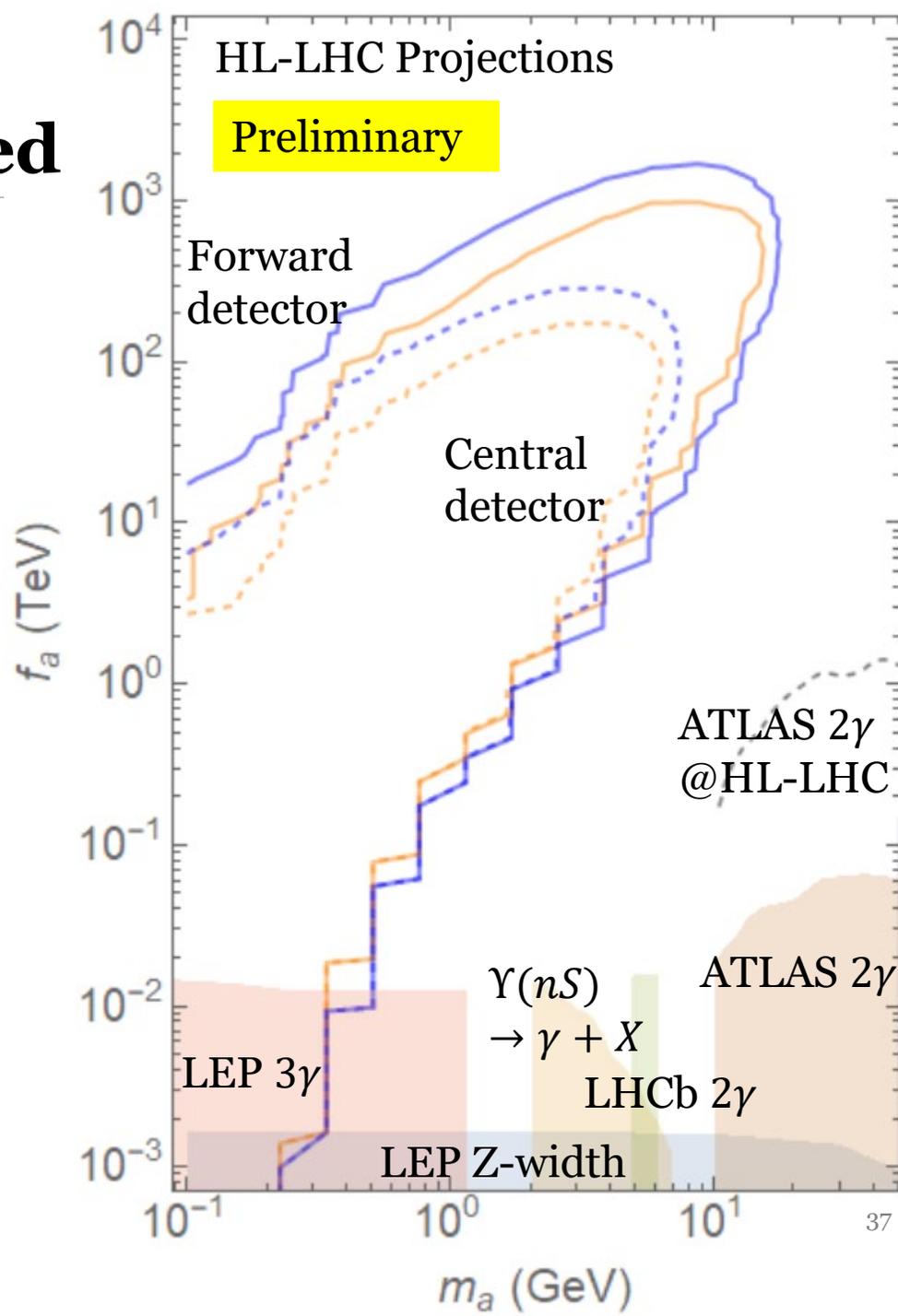


A high-quality axion can be probed

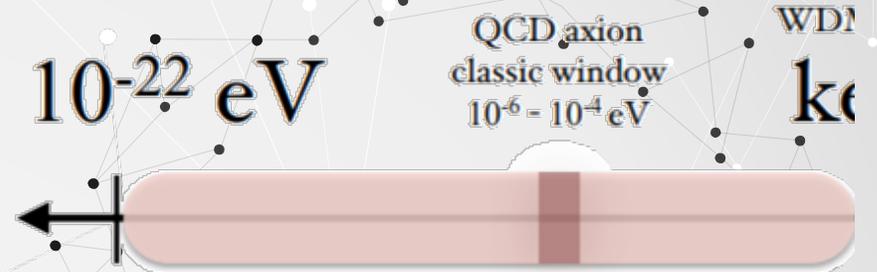
- Gluon factory \rightarrow Rate high
- LLPs \rightarrow background low
- Novel trigger and selection proposed focusing on displaced tracks



$$\frac{\alpha_s}{4\pi} \left(\theta - \frac{a}{f_a} \right) \tilde{G} G$$



Ultralight DM: dark photon & ALP



“Ultralight” DM

non-thermal
bosonic fields

Well motivated theories:

- Dark Photons
- Axions and ALPs

Well motivated searches:

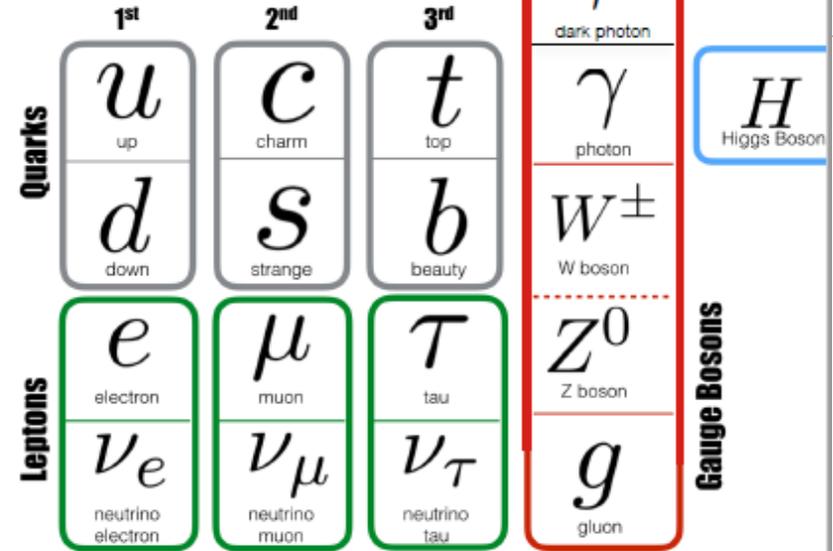
- Light mediators
- Dark matter

Dark Photons

- Imagine another photon, with a different mass.
- Common in top-down frameworks.
- Any* heavy particle that is charged both photons will generate mixing.



Nature has already ordered extra copies of fermions. Why not gauge bosons?



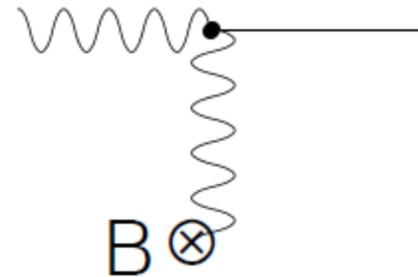
$$\mathcal{L} = -\frac{1}{4} (F_{\mu\nu}F^{\mu\nu} + F'_{\mu\nu}F'^{\mu\nu} - 2\epsilon F_{\mu\nu}F'^{\mu\nu}) + \dots \supset \epsilon (\vec{E} \cdot \vec{E}' + \vec{B} \cdot \vec{B}')$$

An oscillating EM field is a source of dark photons, and vice versa. (reminiscent of neutrino oscillations)

Axion-like particles

- Imagine an approximate symmetry broken at a high scale f .
→ a pseudo-Goldstone Boson \approx an axion-like particle.
- Common in top-down constructions, the axion is invoked to solve the strong CP problem.
- Loops of heavy charged particles can generate interaction:

$$\mathcal{L} = \frac{\alpha}{f} a F_{\mu\nu} \tilde{F}^{\mu\nu} = g_{a\gamma\gamma} a \vec{E} \cdot \vec{B}$$



Axions and photons mix in a magnetic field.

An oscillating $\vec{E} \cdot \vec{B}$ is a source of dark photons.

Longer Range Interactions and Wave-like Dark Matter

- Both axion-like particles and dark photons are well motivated as mediators of long range interactions that can be searched for.

$\mathcal{L} \supset$ dark photons? axions?

- Both axion-like particles and dark photons are dark matter candidates.
- In the Wave-like DM category. Oscillating at $\omega = m_{\text{DM}}$.



\supset dark photons? axions?

Searches with SRF Cavities

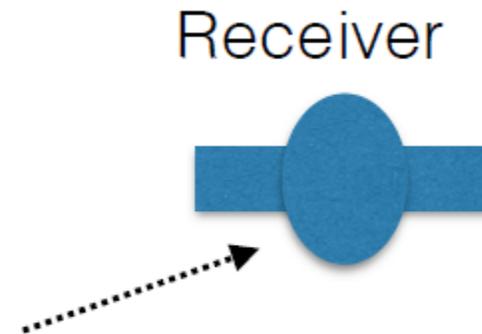
- Fermilab's SRF Cavities are world's highest quality photon resonators, with Q as high as 10^{11} :
 - Large fields when excited \rightarrow can source dark fields.
 - Resonant response \rightarrow can amplify a feeble signal.

Light Shining through wall:



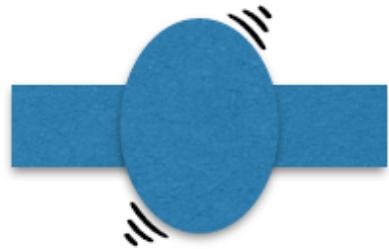
a search for a mediator.

A dark matter search:



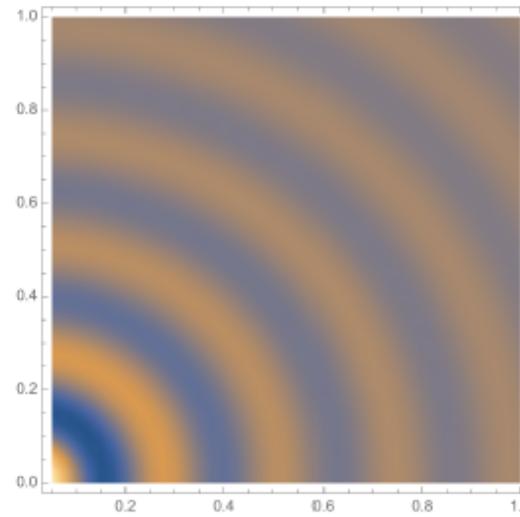
the DM filled Universe
is the emitter

Dark Photon Search

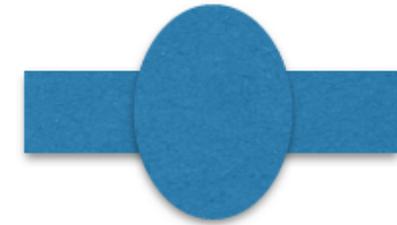


Emitter Cavity

Frequency of 1.3 GHz,
excited to ~ 35 MV/m.
Thats $\sim 10^{25}$ Photons!



a dark photon
field is radiated
at 1.3 GHz.

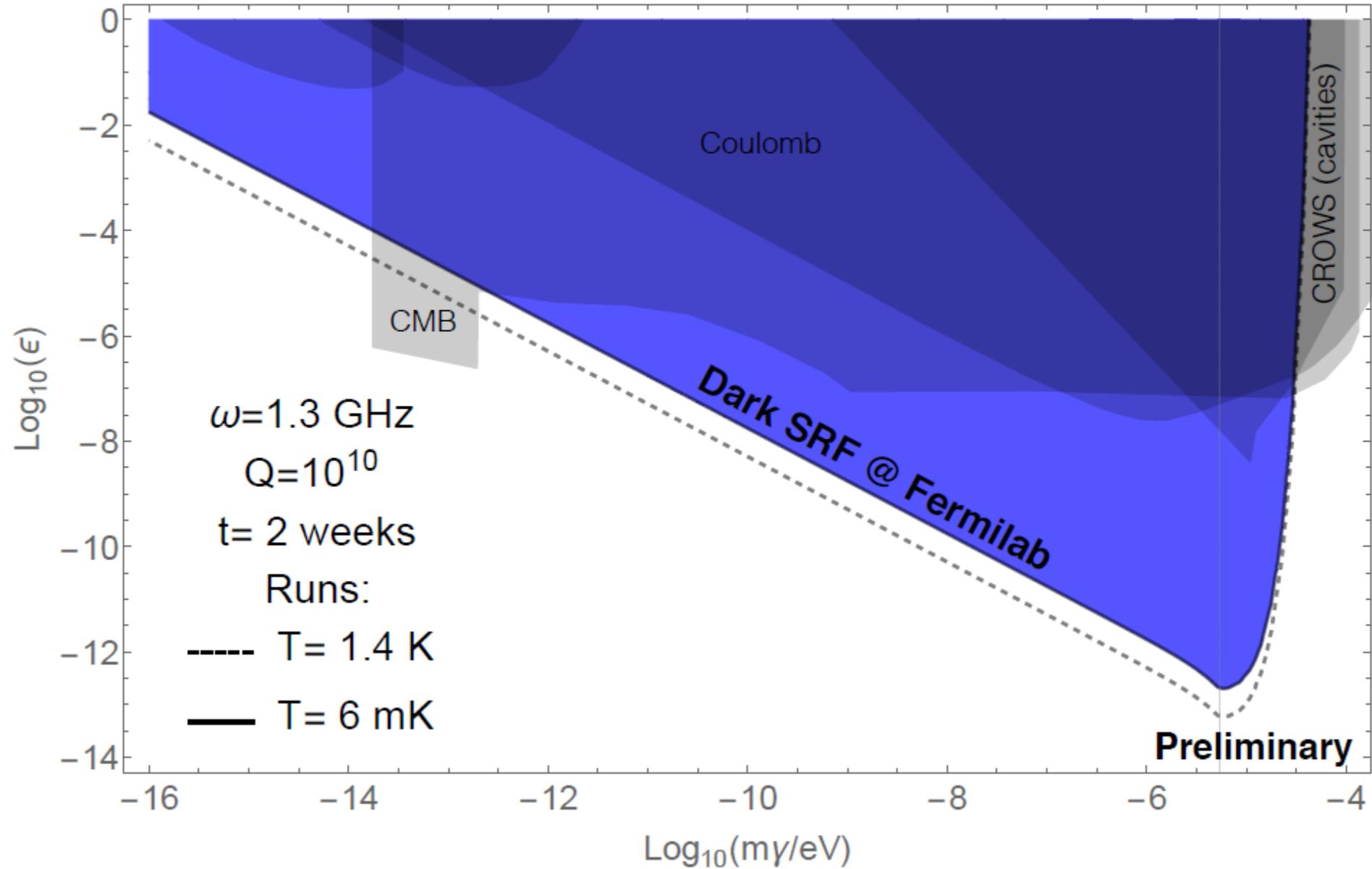


Receiver Cavity

Tuned to 1.3 GHz.
Responds to dark field.
Contains only thermal
noise ($T=1.4$ K).

For correct cavity positioning $P_{\text{rec}} \sim G^2 \epsilon^4 \left(\frac{m_{\gamma'}}{\omega} \right)^4 Q_{\text{rec}} Q_{\text{em}} P_{\text{em}}$

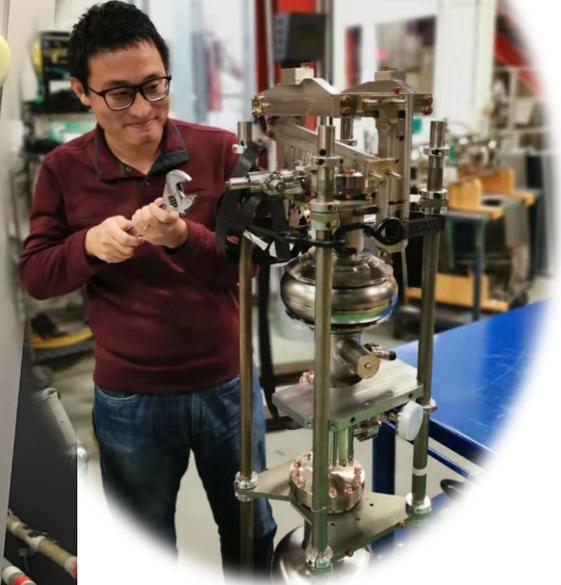
A Dark Photon Search



“Run 0” results summary

Everything works!

- ✓ Design
- ✓ Tuner operation
- ✓ Microwave scheme for matching the frequencies
- ✓ Actual data – first acquisition



Dielectric haloscopes

Image from R. Lasenby

DM can Bragg-convert in medium, producing photons:

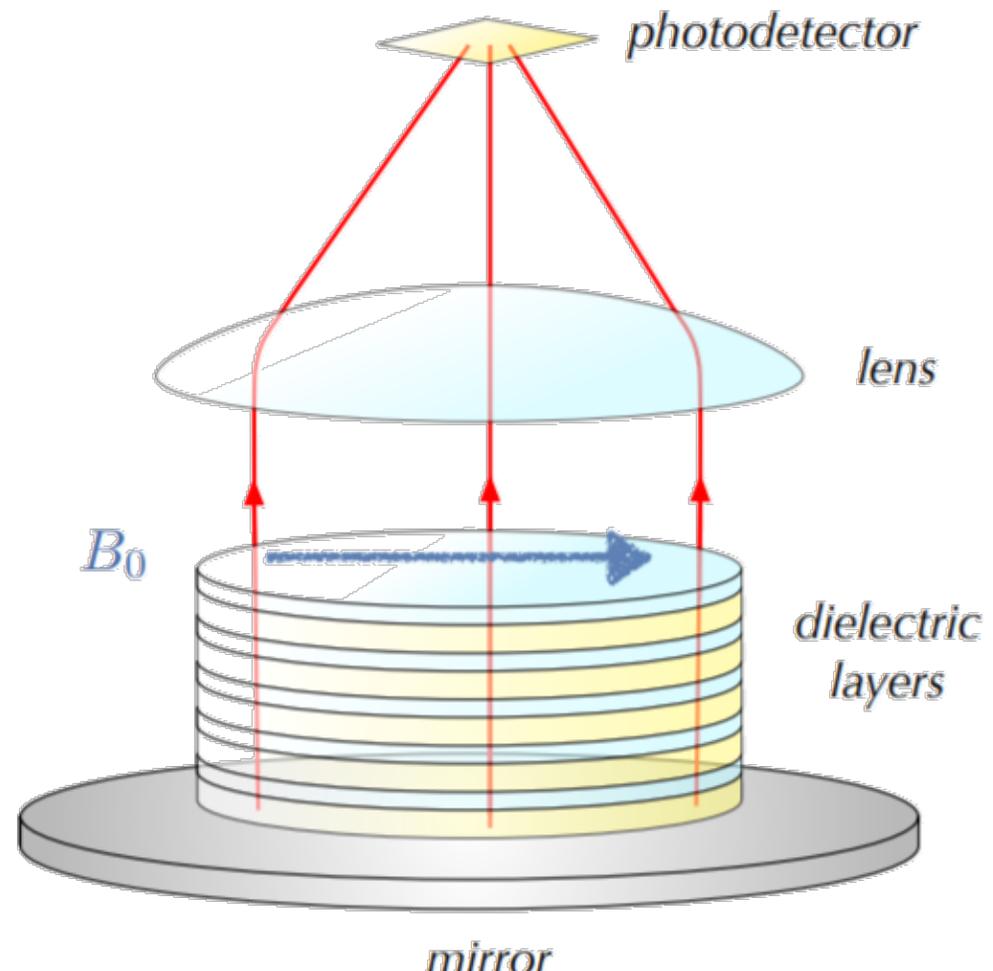
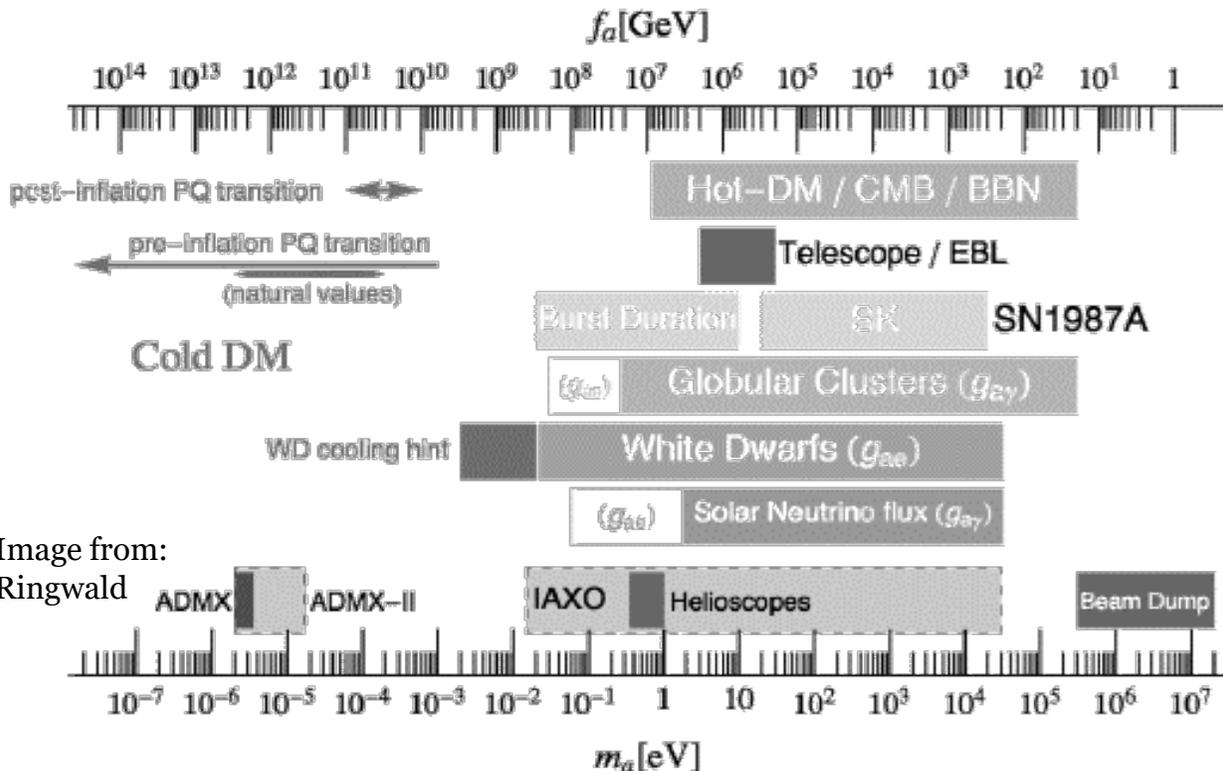
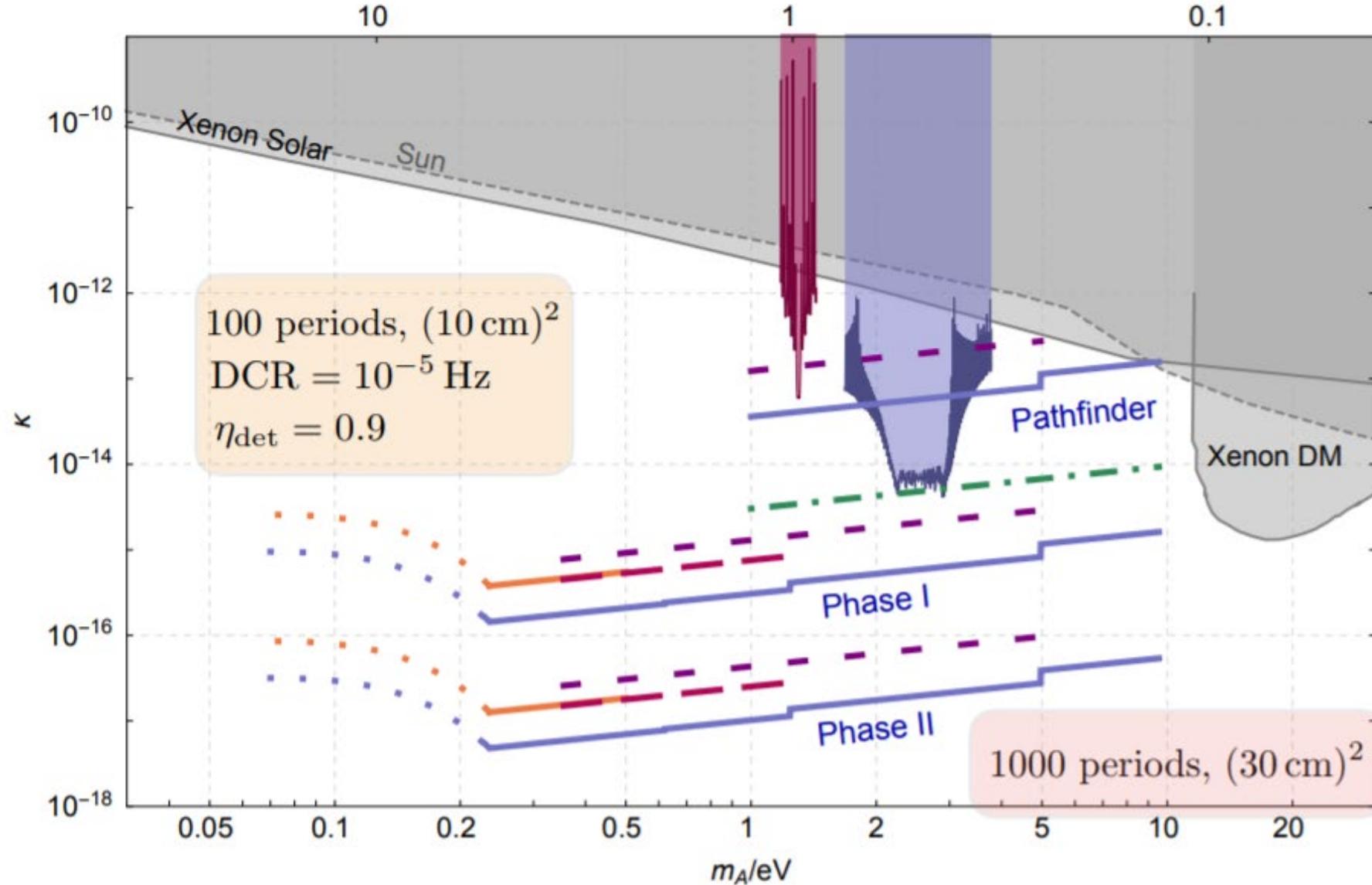


Image from:
Ringwald

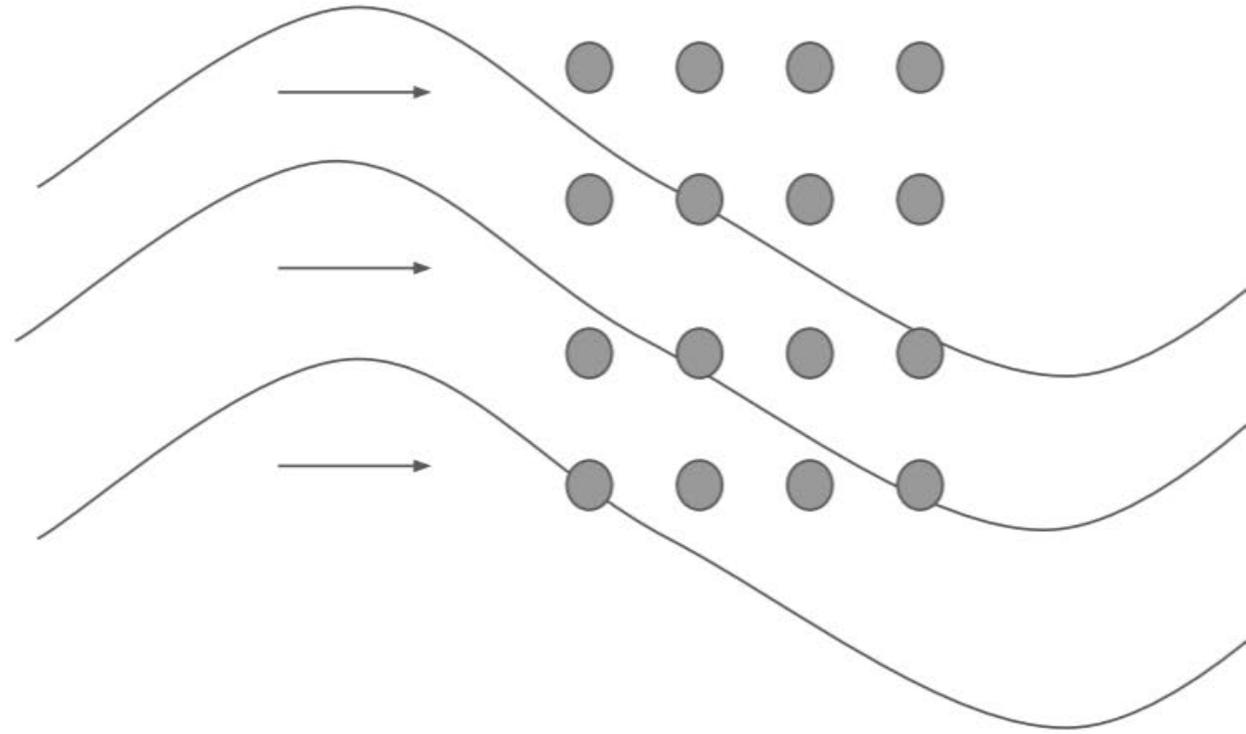
Future sensitivity

Phys.Rev. D98 (2018) no.3, 035006 (1803.11455) - M. Baryakhtar, J. Huang, RL $t_{\text{int}} = 10^6 \text{ sec}$



Ultralight DM

Suppose DM consists entirely of a single, very light field.

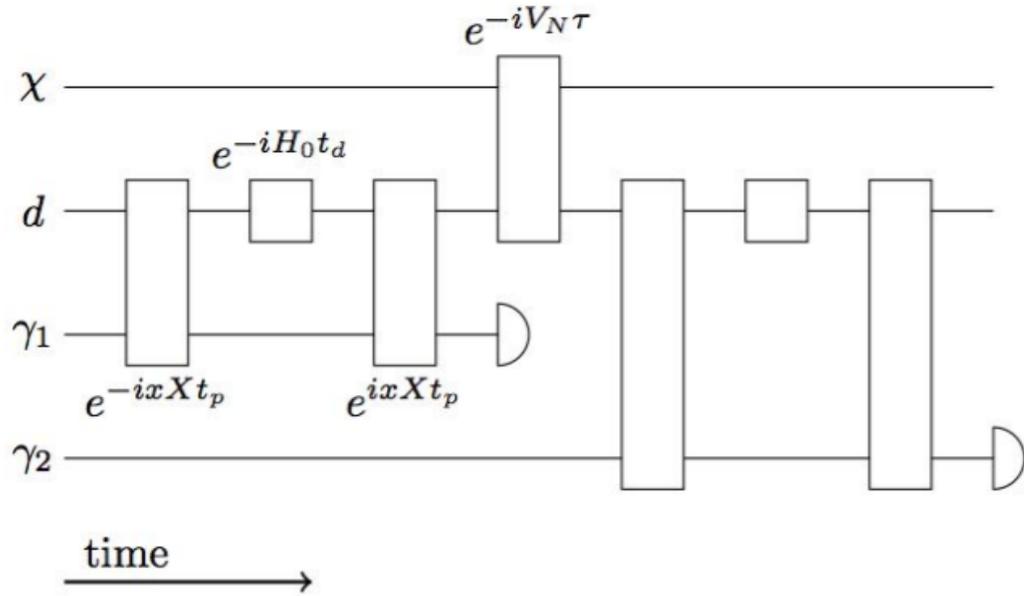


Locally, this will essentially look like a coherent wave acting across the whole array.

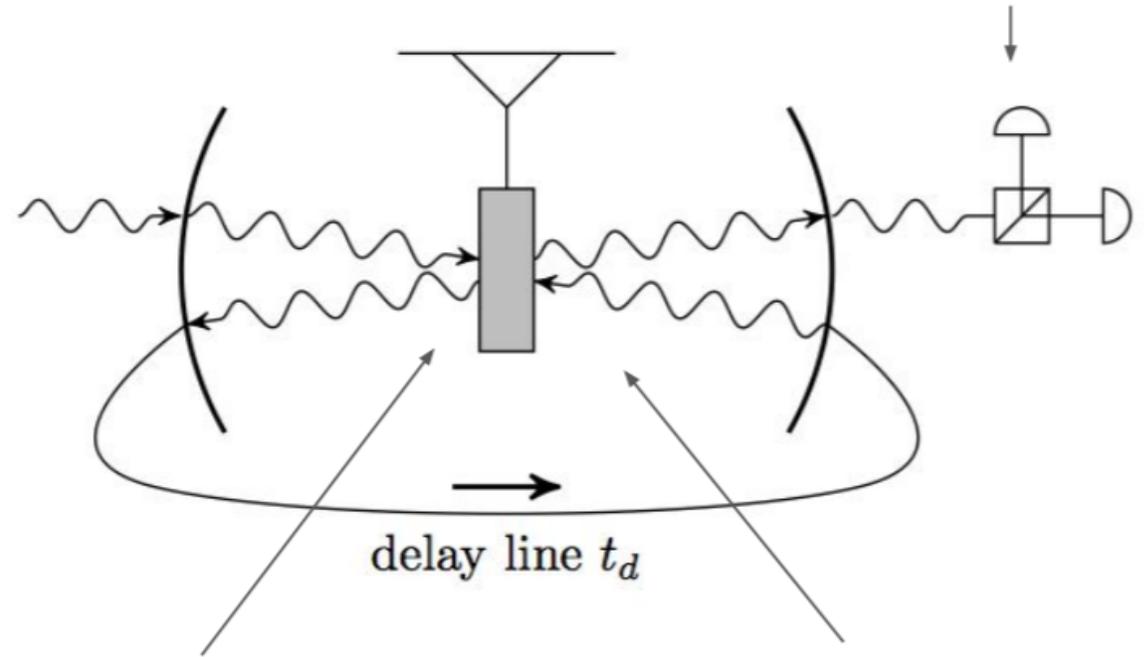
$$\mathcal{L}_{int} = g_{B-L} A \bar{n} n$$

$$F = g_{B-L} N_n F_0 \sin(\omega_s t)$$

Back-action evasion measurement of velocity



readout light phase via interferometer

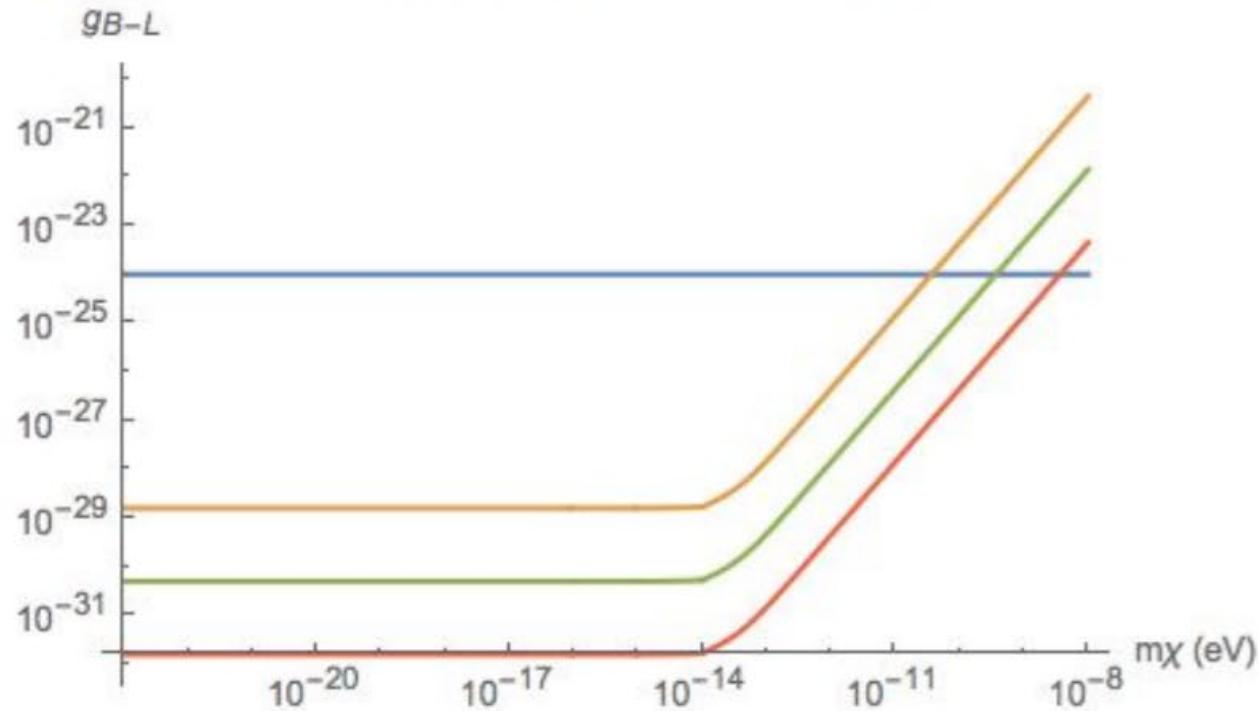


light phase $\sim x(t_1)$
impart $+p$ to mirror

light phase $\sim -x(t_2)$
impart $-p$ to mirror

→ Light phase $\sim x(t_1) - x(t_2) \sim v$, momentum transfer ~ 0 .

Achievable g_{B-L} detection with N sensors
(gas-limited SQL, incoherent readout)



WIP w/ A. Hook, Z. Liu, J. Taylor (UMD),
Y. Zhao (Utah) & discussions w/ D. Moore
(Yale)

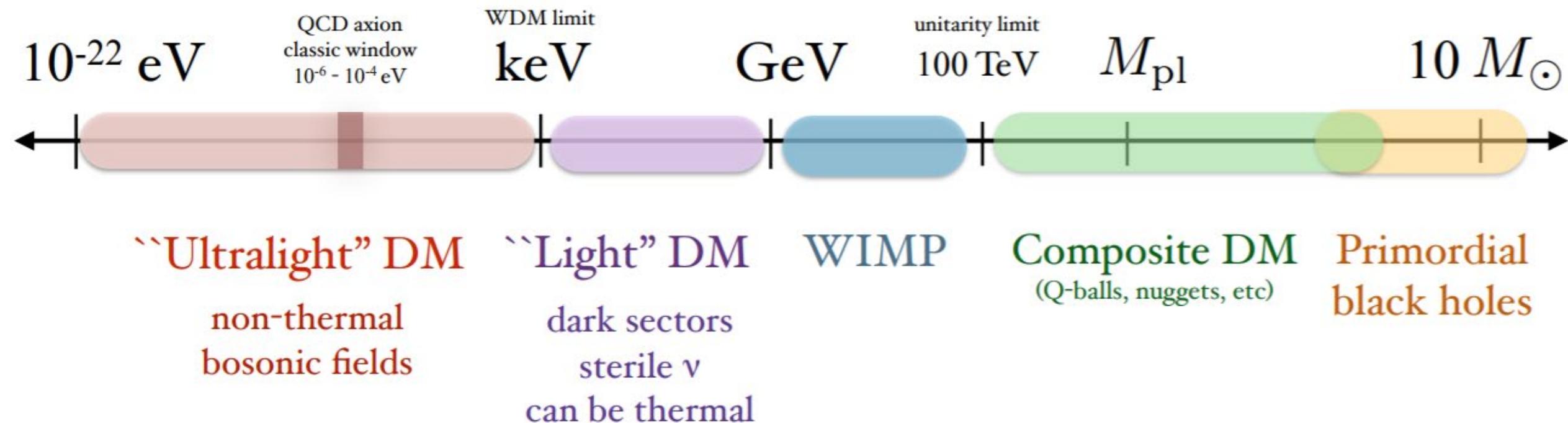
- Eotvos-Wash
- N=1
- N=10³
- N=10⁶

$\text{SNR} \sim 1/\sqrt{N_{\text{total}}}$ \rightarrow huge win via scaling (signal = coherent oscillation of entire lattice)

Here plotting mg-scale detectors operating “just” at SQL.

Backaction-evasion improves high-frequency limits.

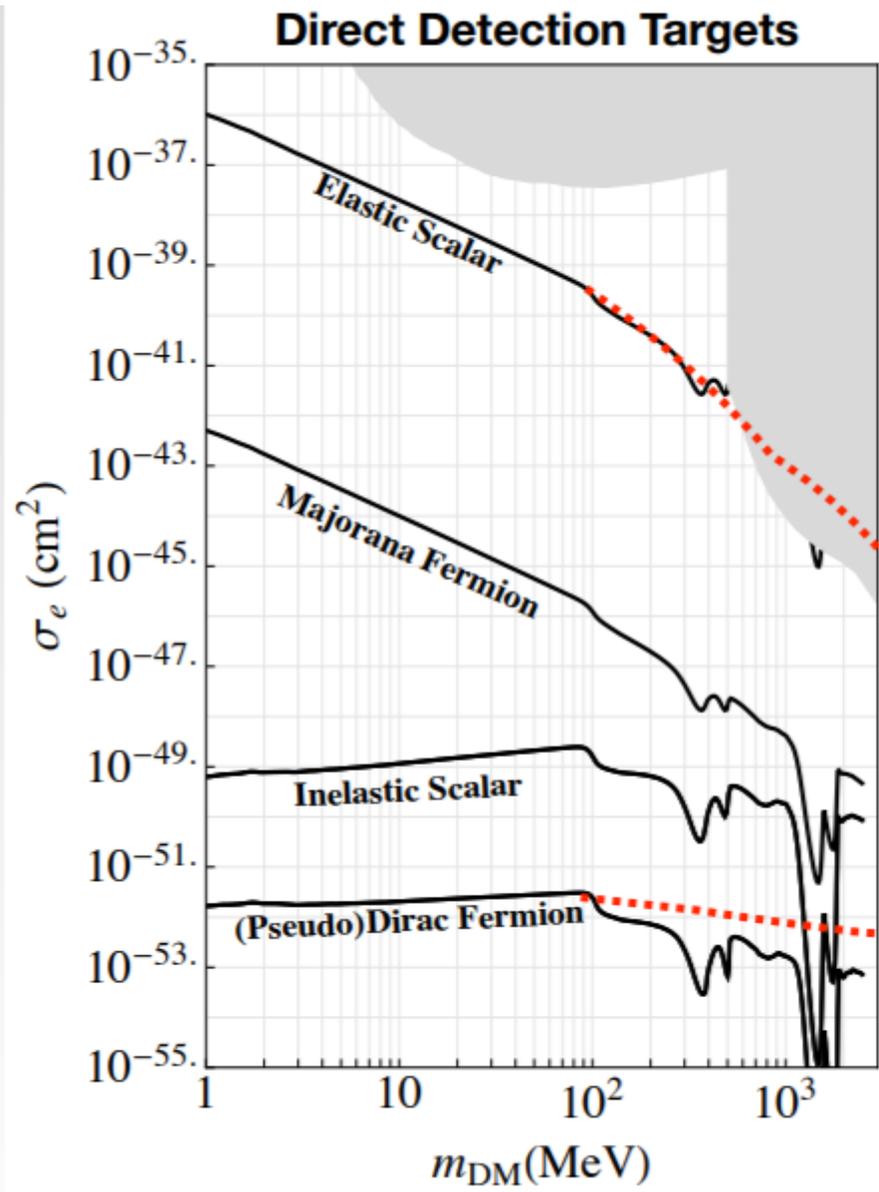
Summary



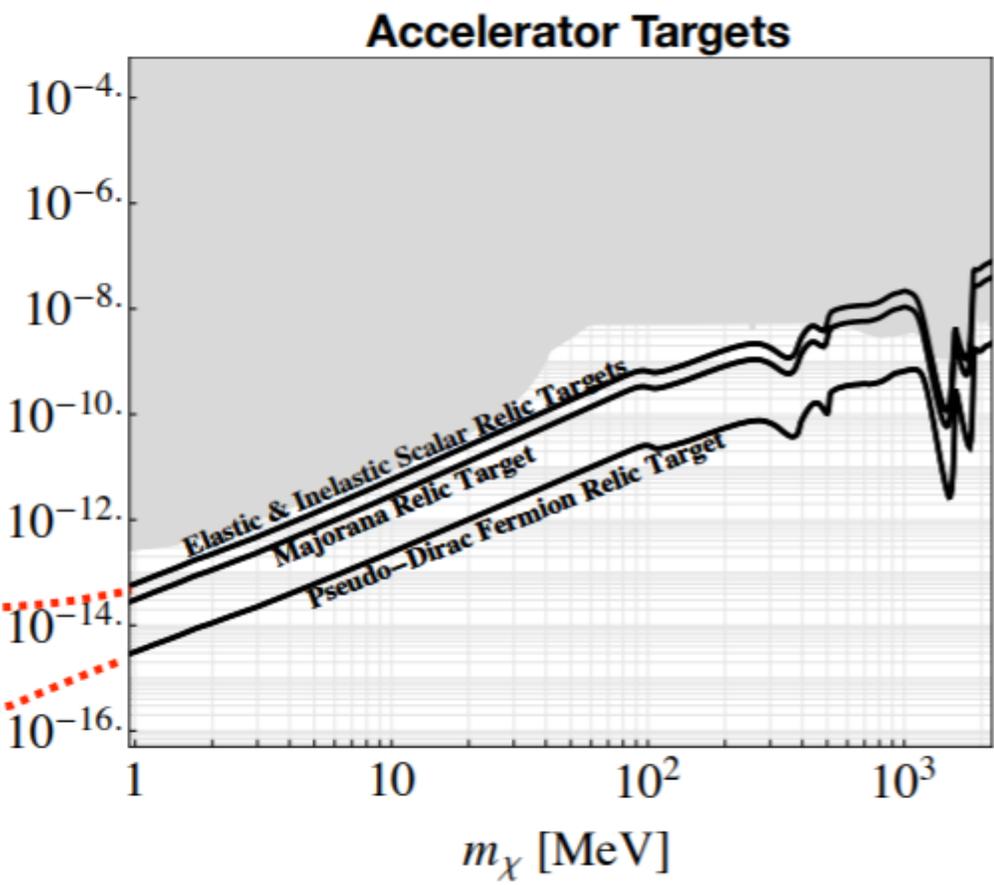
Lots of opportunities!

Backup

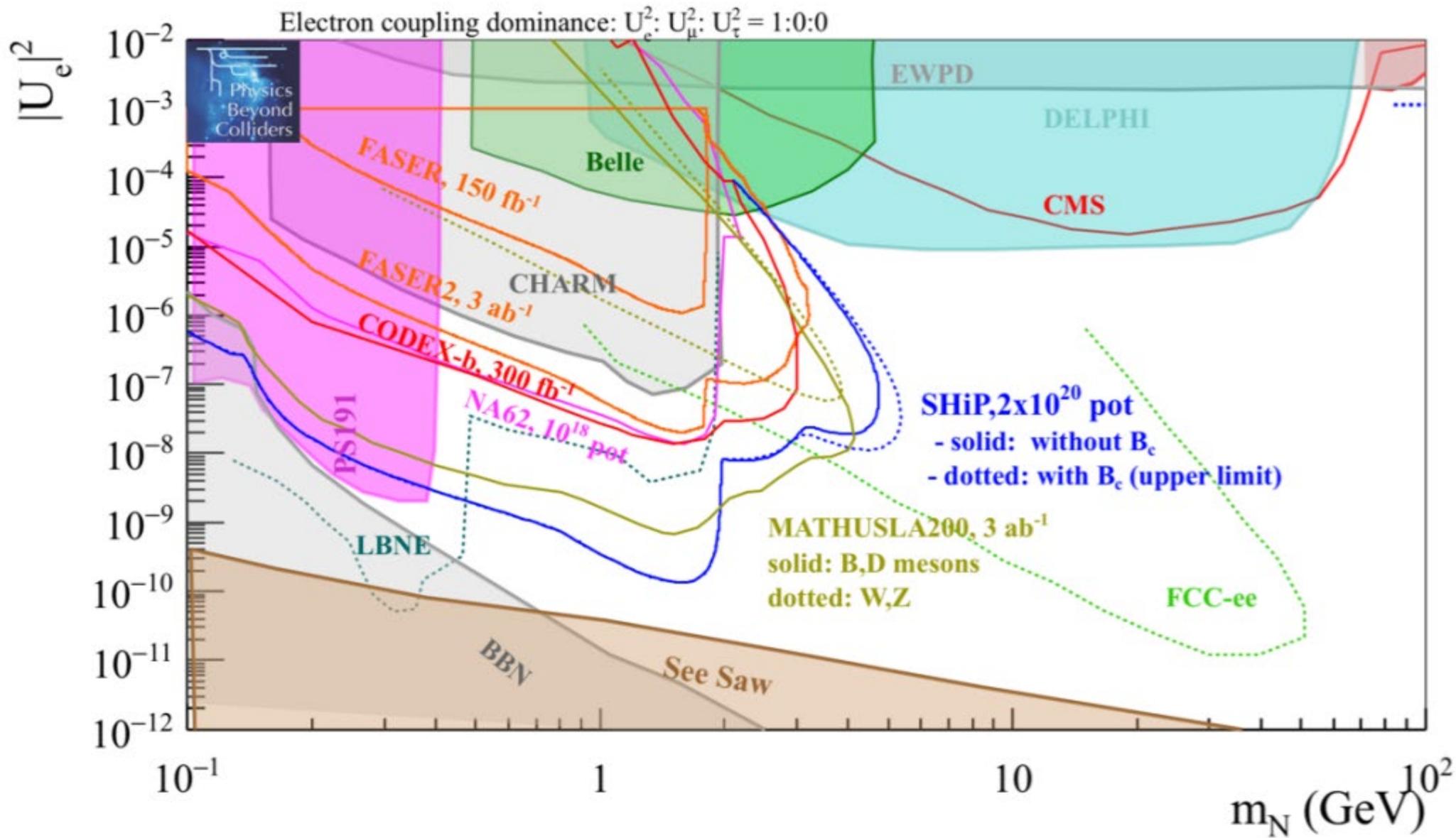
Courtesy of G.Krnjaic

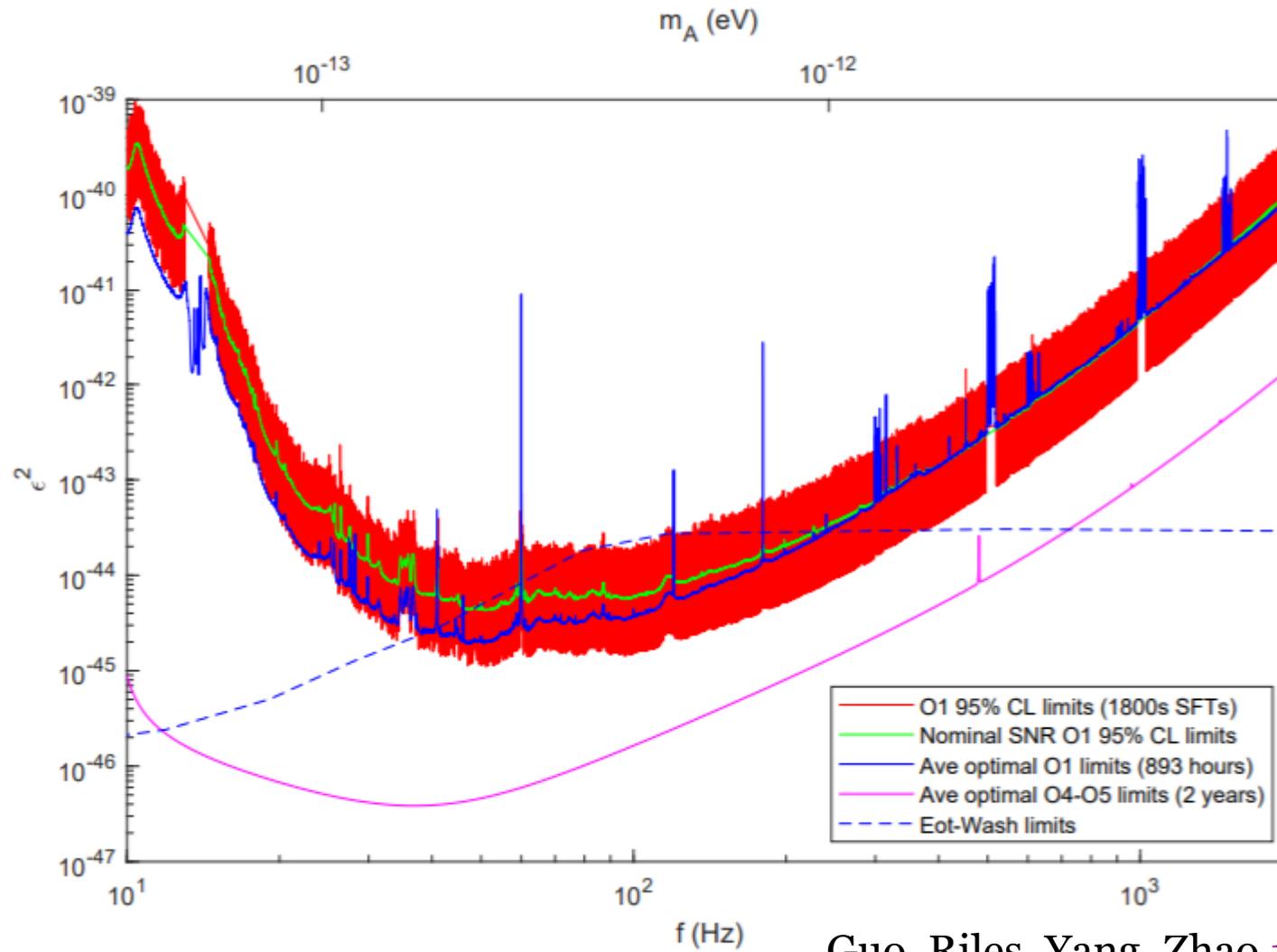


$$y = \epsilon^2 \alpha_D (m_\chi / m_{A'})^4$$



non-relativistic cross sections can be loop- or velocity- suppressed





Guo, Riles, Yang, Zhao [1905.04316](#)