# B-anomalies: status and implications

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## IFT-Workshop

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## Outline & Questions

- 1. Diagnosis of anomalies: Where we stand?
- 2. A comparative study of Pre and Post Moriond

-Are now all the global significances smaller?-Are new emerging hypothesis?-Brief Comparison with other analysis.

- 3. Lepton Flavour Universal (LFU) New Physics -Two kinds of New Physics? Maybe two scales?
- 4. Linking charge, neutral and LFU New Physics.
- 5. Solutions proposed to the anomalies
- 6. What's next?  $Q_5$
- 7. Conclusions

## Diagnosis of anomalies in $b \to s \ell \ell$

## Model independent approach to $b ightarrow s\ell\ell$

$$\begin{aligned} \mathcal{H}_{\text{eff}} &= -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \mathcal{C}_i \mathcal{O}_i \\ \mathcal{O}_7 &= \frac{e}{16\pi^2} m_b (\bar{s}\sigma_{\mu\nu} P_R b) F^{\mu\nu}, \\ \mathcal{O}_{7'} &= \frac{e}{16\pi^2} m_b (\bar{s}\sigma_{\mu\nu} P_L b) F^{\mu\nu}, \\ \mathcal{O}_{9\ell} &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \ell), \\ \mathcal{O}_{9\ell'} &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_R b) (\bar{\ell}\gamma^\mu \ell), \\ \mathcal{O}_{10\ell} &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell), \\ \mathcal{O}_{10\ell'} &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_R b) (\bar{\ell}\gamma^\mu \gamma_5 \ell), \\ \mathcal{A} t \text{ the } \mu_b = 4.8 \text{ GeV scale:} \end{aligned}$$

$$\begin{aligned} \mathcal{C}_{7}^{\text{SM}} &= -0.29, \ \mathcal{C}_{9}^{\text{SM}} &= 4.1, \ \mathcal{C}_{10}^{\text{SM}} &= -4.3 \end{aligned}$$

## The starting point: Angular distribution

4-body angular distribution  $\bar{\mathbf{B}}_{\mathbf{d}} \rightarrow \bar{\mathbf{K}}^{*0} (\rightarrow \mathbf{K}^{-} \pi^{+}) \mathbf{l}^{+} \mathbf{l}^{-}$  with three angles, invariant mass of lepton-pair  $q^{2}$ .



 $\theta_{\ell}$ : Angle of emission between  $\bar{K}^{*0}$ and  $\mu^{-}$  in di-lepton rest frame.  $\theta_{\mathbf{K}}$ : Angle of emission between  $\bar{K}^{*0}$ and  $K^{-}$  in di-meson rest frame.  $\phi$ : Angle between the two planes.

q<sup>2</sup>: dilepton invariant mass square.

$$\frac{d^4\Gamma(\bar{B}_d)}{dq^2 \, d\cos\theta_\ell \, d\cos\theta_K \, d\phi} = \frac{9}{32\pi} \sum_i J_i(q^2) f_i(\theta_\ell, \theta_K, \phi)$$

$$J_i(q^2) \text{ function of transversity (helicity) amplitudes of K^*: } A^{L,R}_{\perp,\parallel,0} \text{ but also } A_t, A_S$$

 $A_{\perp,\parallel,0}^{L,R} = C_i$  (short) × Hadronic quantities (long)

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 $J_i(q^2)$  function of transversity (helicity) amplitudes of K\*:  $A_{\perp,\parallel,0}^{L,R}$  but also  $A_t, A_S$  $A_{\perp,\parallel,0}^{L,R} = C_i$  (short) × Hadronic quantities (long)

$$\frac{1}{\Gamma_{full}'} \frac{d^4 \Gamma}{dq^2 d\cos \theta_K d\cos \theta_l d\phi} = \frac{9}{32\pi} \left[ \frac{3}{4} \mathbf{F_T} \sin^2 \theta_K + \mathbf{F_L} \cos^2 \theta_K + (\frac{1}{4} \mathbf{F_T} \sin^2 \theta_K - \mathbf{F_L} \cos^2 \theta_K) \cos 2\theta_l + \sqrt{\mathbf{F_T} \mathbf{F_L}} \left( \frac{1}{2} \mathbf{P'_4} \sin 2\theta_K \sin 2\theta_l \cos \phi + \mathbf{P'_5} \sin 2\theta_K \sin \theta_l \cos \phi \right) + 2 \mathbf{P_2} \mathbf{F_T} \sin^2 \theta_K \cos \theta_l + \frac{1}{2} \mathbf{P_1} \mathbf{F_T} \sin^2 \theta_K \sin^2 \theta_l \cos 2\phi_l + \sqrt{\mathbf{F_T} \mathbf{F_L}} \left( \mathbf{P'_6} \sin 2\theta_K \sin \theta_l \sin \phi - \frac{1}{2} \mathbf{P'_8} \sin 2\theta_K \sin 2\theta_l \sin \phi \right) - \mathbf{P_3} \mathbf{F_T} \sin^2 \theta_K \sin^2 \theta_l \sin 2\phi_l \left( 1 - \mathbf{F_S} \right) + \frac{1}{\Gamma_{full}'} \mathbf{W_S}$$

#### [SDG,JM,JV,1207.2753]



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Theoretical framework: QCDF/SCET+**robust large-recoil symmetries** +breaking (pert+non-pert)  $\hookrightarrow$  independent of LCSR details

$$\mathcal{T}_{a} = \xi_{a} \left( C_{a}^{(0)} + \frac{\alpha_{s} C_{F}}{4\pi} C_{a}^{(1)} \right) + \frac{\pi^{2}}{N_{c}} \frac{f_{B} f_{K^{*},a}}{M_{B}} \Sigma_{a} \sum_{\pm} \int \frac{d\omega}{\omega} \Phi_{B,\pm}(\omega) \int_{0}^{1} du \Phi_{K^{*},a}(u) T_{a,\pm}(u,\omega). \quad a = \bot, \|$$

 $\xi_a$  (soft FF).  $C_i = 1 + O(\alpha_s)$  hard-vertex renormalization and  $T_i$  hard-scattering kernels computed in  $\alpha_s$ -expansion.  $\Phi_i$  light-cone wave functions. Two types of non-factorizable contributions:

• Hard spectator scattering  $(T_a)$ : matrix elements of 4-quark op. and the chromomagnetic  $O_8$  operator



#### Perturbative and non-perturbative charm

**Problem:** Charm-loop yields a (most likely)  $q^2$  – and process-dependent contribution with  $O_{7,9}$  structures that may (in a local analysis of data) mimic New Physics.

 $C_{9i}^{\text{eff}}(q^2) = \mathbf{C}_{9SMpert} + C_9^{NP} + \mathbf{s_i} \delta \mathbf{C}_{9i}^{\mathbf{c}\bar{\mathbf{c}}\mathbf{L}\mathbf{D}}(\mathbf{q^2}). \qquad \mathbf{i} = \bot, \|, \mathbf{0}$ 

Perturbative:  $C_{9 \text{ SMpert}} = C_9^{\text{SM}} + Y(q^2)$ with  $Y(q^2)$  stemming from one-loop matrix elements of 4-quark operators  $O_{1-6}$ . ... $\mathcal{O}(\alpha_s)$  corrections to  $C_{7,9}^{\text{eff}}$  of  $Y(q^2)$  included via  $C_{\perp,\parallel}^{1 \text{ (nf)}}$  but only  $O_{1,2}$  (previous slide)

#### Non-perturbative: $\delta C_{9i}^{c\bar{c}LD}(q^2)$

More difficult to make progress here:

- 1 Use LCSR to estimate long-distance contribution with soft-gluon exchange.  $\Rightarrow$
- 2 Or use fits to the same data you want to explain [Ciuchini, Silvestrini et al.]  $\Rightarrow$





### A bright future: LHCb ultimate precision expected in RUNII

Projections from LHCb for  $P'_5$  in Phase-II Upgrade.





A large number of small bins open the window in  $P'_5$  for another observable: zero of  $P'_5$ .

At LO:

$$q_0^2 = -rac{m_b m_B^2 \mathcal{C}_7^{ ext{eff}}}{m_b \mathcal{C}_7^{ ext{eff}} + m_B \mathcal{C}_9^{ ext{eff}}(q_0^2)}$$

zero not sensitive to  $C_{10}$  (at LO).

#### At NLO:

• Large shift of zero of  $P_5'$  from  $q_0^{2SM} \simeq 2 \text{ GeV}^2$ to  $q_0^{C_9^{\rm NP} = -1.76} \simeq 3.8 \text{ GeV}^2$ .

## Diff. Branching Ratios: Lepton Flavour Dependent



## $B_s \rightarrow \phi \mu \mu vs B \rightarrow K^* \mu \mu$ : Lepton Flavour Dependent



with corrected BSZ FF

**Not yet significant**: FF at low-q<sup>2</sup> for  $B_s \rightarrow \phi$  (BSZ) larger than  $B \rightarrow K^*$ , while data is reversed. Ok at high-q<sup>2</sup>. **BSZ problem or statistical fluctuation**?

Our prediction for  $B \rightarrow K^*$  with KMPW has larger errors so **no problem in our case**.

More data will clarify it....

## $R_K$ : Lepton Flavour Universality Violation

**FCNC**, **test of universality** of lepton coupling, potential high sensitivity to NP contributions.

First possible signal of LFUV ... after LHCb update

$$R_K^{[1.1,6]} = \frac{\mathcal{B}(B \to K\mu^+\mu^-)}{\mathcal{B}(B \to Ke^+e^-)} = 0.846^{+0.060}_{-0.054} + 0.016_{-0.014}$$



Simple structure of BR:  $f_{+,0,T} \rightarrow f_+$ 

dominates while the other two suppressed by lepton mass or C<sub>7</sub>. => **Good observable in presence NP** => tensions cannot be explained by FF or charm. Electromagnetic small. [Isidori et al.]

#### Does a more SM-like central value imply a reduction in significance?



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## $R_{K^*}$ : Lepton Flavour Universality Violation



## Updated global analysis of $\,b \to s\ell\ell\,$





... hopefully now the race for the right pattern

include additional interesting horses than just the old guys:  $C_9$  and  $C_9=-C_{10}$ 

178 observables from (LHCb, Belle, ATLAS and CMS, no CP-violating obs)

•  $B \to K^* \mu \mu$  ( $P_{1,2}, P'_{4,5,6,8}, F_L$  in 5 large-recoil bins + 1 low-recoil bin)+available electronic obs.

...latest update  $Br(B \rightarrow K^* \mu \mu)$  in small bins.

...LHCb results on  $R_{K^*}$ 

- $B_s \rightarrow \phi \mu \mu$  ( $P_1, P'_{4,6}, F_L$  in 3 large-recoil bins + 1 low-recoil bin)
- $B^+ \to K^+ \mu \mu$ ,  $B^0 \to K^0 \ell \ell$  (BR) ( $\ell = e, \mu$ ) (new average  $R_K = 0.846^{+0.060+0.016}_{-0.054-0.014}$ )
- $B \to X_s \gamma, B \to X_s \mu \mu, B_s \to \mu \mu$  (BR).
- Radiative decays:  $B^0 \to K^{*0}\gamma$  ( $A_I$  and  $S_{K^*\gamma}$ ),  $B^+ \to K^{*+}\gamma$ ,  $B_s \to \phi\gamma$
- ► Belle measurements for the isospin-averaged but lepton-flavour dependent  $(Q_{4,5} = P_{4,5}'^{\mu} P_{4,5}'^{e})$ : [3rd test of LFUV]

$$P_i^{\prime \ell} = \sigma_+ P_i^{\prime \ell}(B^+) + (1 - \sigma_+) P_i^{\prime \ell}(\bar{B}^0) \qquad \sigma_+ = 0.5 \pm 0.5$$

similar treatment of new Belle isospin-averaged result on  $R_{K^*}$  (3-bins)

▶ ATLAS measurement of whole basis of  $P_i$  and CMS measurements of  $P_1$  and  $P'_5$ .

► ATLAS update of  $B_s \rightarrow \mu\mu$  (averaged with LHCb & CMS) and latest  $f_{Bs}$  lattice update.

## Implications of the new updates on $R_{\rm K},\,R_{\rm K^*},\,B_{\rm S}{\rightarrow}\mu\mu$

 $\text{Pull}_{\text{SM}}: \chi^2_{\text{SM}}(C_i=0)-\chi^2 \text{min}(C_i^{\text{HIP}}) \text{ considering } N_{\text{dof}}$ 



- Hierarchy remains invariant except  $C_{9\mu} = -C_{9'\mu}$  scenario  $(R_K \approx 1)$ 
  - Scenario  $C_{9\mu}$  preferred in "All" fit Scenario  $C_{9\mu}$  = - $C_{10\mu}$  preferred in "LFUV" fit.
- Best fit points for All and LFUV fits in scen.  $C_{9\mu}$  in nice agreement
- Scenario  $C_{10\mu}$  stays at a significance of  $\approx 4\sigma$  for All and LFUV fits.

## Implications of the new updates on $R_{\rm K},\,R_{\rm K^*},\,B_{\rm s}{\rightarrow}\mu\mu$

#### Interesting surprises in 2D updates...

2017	All			LFUV		
2D Hyp.	Best fit	$\operatorname{Pull}_{\mathrm{SM}}$	p-value	Best fit	$\text{Pull}_{\text{SM}}$	p-value
$(\mathcal{C}_{9\mu}^{\mathrm{NP}},\mathcal{C}_{10\mu}^{\mathrm{NP}})$	(-1.01,0.29)	5.7	72	(-1.30,0.36)	3.7	75
$(\mathcal{C}_{9\mu}^{\mathrm{NP}},\mathcal{C}_{7}')$	(-1.13, 0.01)	5.5	69	(-1.85,-0.04)	3.6	66
$(\mathcal{C}_{9\mu}^{\mathrm{NP}},\mathcal{C}_{9'\mu})$	(-1.15, 0.41)	5.6	71	(-1.99, 0.93)	3.7	72
$\left \left(\mathcal{C}_{9\mu}^{\mathrm{NP}},\mathcal{C}_{10'\mu} ight) ight $	(-1.22, -0.22)	5.7	72	(-2.22,-0.41)	3.9	85
$(\mathcal{C}_{9\mu}^{\mathrm{NP}},\mathcal{C}_{9e}^{\mathrm{NP}})$	(-1.00, 0.42)	5.5	68	(-1.36, 0.46)	3.5	65
Hyp. 1	(-1.16, 0.38)	5.7	73	(-1.68, 0.60)	3.8	78
Hyp. 2	(-1.15, 0.01)	5.0	57	(-2.16, 0.41)	3.0	37
Hyp. 3	(-0.67, -0.10)	5.0	57	(0.61, 2.48)	3.7	73
Hyp. 4	(-0.70, 0.28)	5.0	57	$\left  (-0.74, 0.43) \right $	3.7	72
2019		All			LFUV	
2D Hyp.	Best fit	$\mathrm{Pull}_{\mathrm{SM}}$	p-value	Best fit	Pull <sub>SM</sub>	p-value
$(\mathcal{C}_{9\mu}^{\mathrm{NP}},\mathcal{C}_{10\mu}^{\mathrm{NP}})$	(-0.95, 0.20)	5.7	69.5 %	(-0.30,0.52)	3.6	74.5%
$(\mathcal{C}_{9\mu}^{\mathrm{NP}},\mathcal{C}_{7}')$	(-1.03, 0.02)	5.6	68.2%	(-1.03,-0.04)	3.1	53.7%
$(\mathcal{C}_{9\mu}^{\mathrm{NP}},\mathcal{C}_{9'\mu})$	(-1.13, 0.54)	5.9	74.5%	(-1.88,1.14)	3.6	75.7%
$(\mathcal{C}_{9\mu}^{\rm NP},\mathcal{C}_{10'\mu})$	(-1.17, -0.34)	6.1	78.1%	(-2.07,-0.63)	4.0	92.8%
$(\mathcal{C}_{9\mu}^{ ext{NP}},\mathcal{C}_{9e}^{ ext{NP}})$	(-1.04,-0.11)	5.5	65.3%	(-0.76.9.25)	3.1	50.8%
Hyp. 1	(-1.09, 0.28)	6.0	75.8%	(-1.69,0.32)	3.6	77.1%
Hyp. 2	(-1.00,0.09)	5.4	63.9%	(-2.00,0.26)	3.3	61.2%
Hyp. 3	(-0.50, 0.08)	5.1	55.8%	(-0.43,-0.09)	3.6	74.5%
Hyp. 4	(-0.52,0.11)	5.2	58.7%	(-0.50,0.15)	3.7	81.9%
Hyp. 5	(-1.17,0.24)	6.1	78.2%	$\ $ (-2.20,0.52)	4.1	93.8%

 Increase in significance in scenarios with RHC

- R<sub>K</sub> more SM-like better described if C<sub>9'µ</sub>>0 and C<sub>10'µ</sub><0</li>
- A  $R_q \otimes L_\ell$  structure for primed operators prefers a V over a  $L_\ell$  structure for leptons.
- Hyp.1 is SM-like for  $B_s \rightarrow \mu \mu$  but perfect for  $R_K!$

Hyp. 1:  $(\mathcal{C}_{9\mu}^{NP} = -\mathcal{C}_{9'\mu}, \mathcal{C}_{10\mu}^{NP} = \mathcal{C}_{10'\mu}),$ Hyp. 2:  $(\mathcal{C}_{9\mu}^{NP} = -\mathcal{C}_{9'\mu}, \mathcal{C}_{10\mu}^{NP} = -\mathcal{C}_{10'\mu}),$ Hyp. 3:  $(\mathcal{C}_{9\mu}^{NP} = -\mathcal{C}_{10\mu}^{NP}, \mathcal{C}_{9'\mu} = \mathcal{C}_{0'\mu}),$ Hyp. 4:  $(\mathcal{C}_{9\mu}^{NP} = -\mathcal{C}_{10\mu}^{NP}, \mathcal{C}_{9'\mu} = -\mathcal{C}_{10'\mu})$ Hyp. 5:  $(\mathcal{C}_{9\mu}^{NP}, \mathcal{C}_{9'\mu} = -\mathcal{C}_{10'\mu}).$ 

### How can we test the presence of RHC $(C_9' \text{ and } C_{10}')$ ?

#### An accurate measurement:

Observable  $P_1$  in two bins



## Implications of the new updates on RK, RK\*, $B_s \rightarrow \mu \mu$



#### Let's check how the 6D fit has evolved:

2017	$\mathcal{C}_7^{\mathrm{NP}}$	$\mathcal{C}_{9\mu}^{ ext{NP}}$	$\mathcal{C}^{\mathrm{NP}}_{10\mu}$	$\mathcal{C}_{7'}$	$\mathcal{C}_{9'\mu}$	${\cal C}_{10'\mu}$	
Best fit	+0.03	-1.12	+0.31	+0.03	+0.38	+0.02	
$1 \sigma$	[-0.01, +0.05]	[-1.34, -0.88]	[+0.10, +0.57]	[+0.00, +0.06]	[-0.17, +1.04]	[-0.28, +0.36]	
$2 \sigma$	[-0.03, +0.07]	[-1.54, -0.63]	[-0.08, +0.84]	[-0.02, +0.08]	[-0.59, +1.58]	[-0.54, +0.68]	
2019	$\mathcal{C}_7^{\mathrm{NP}}$	$\mathcal{C}_{9\mu}^{ ext{NP}}$	$\mathcal{C}^{\mathrm{NP}}_{10\mu}$	$\mathcal{C}_{7'}$	$\mathcal{C}_{9'\mu}$	$\mathcal{C}_{10'\mu}$	
Best fit	+0.02	-1.13	+0.21	+0.02	+0.39	-0.12	
1 σ	[-0.01, +0.05]	[-1.28, -0.91]	[+0.04, +0.42]	[+0.00, +0.04	[-0.09, +0.9]	[-0.40, +0.17]	
$2 \sigma$	[-0.03, +0.06]	[-1.48, -0.71]	[-0.12, +0.61]	[-0.02, +0.06]	$[6] \mid [-0.56, +1.1]$	$[4] \mid [-0.57, +0.34]$	

#### • Again **same picture**,

-except change in sign of bfp of  $C_{10'\mu}$ -except significance 5.0 $\sigma$   $\rightarrow$  **5.3** $\sigma$ 

## Implications of the new updates on $R_{\rm K},\,R_{\rm K^*},\,Bs{\rightarrow}\mu\mu$



It is then natural to expect some impact in the significance of LFUV+LFU scenarios

Flavour observables are sensitive to higher scales than direct searches at colliders

#### ... if NP affects flavour it is not surprising that we detect it first.

What is the scale of NP for  $b \to s\ell\ell$ ? Reescaling the Hamiltonian by  $H_{eff}^{NP} = \sum \frac{\mathcal{O}_i}{\Lambda_i^2}$ 

• Tree-level induced (semi-leptonic) with  $\mathcal{O}(1)$  couplings ( $\times \sqrt{g_{bs} g_{\mu\mu}}$ ):

$$\Lambda_{i}^{\text{Tree}} = \frac{4\pi v}{s_{w}g} \frac{1}{\sqrt{2|V_{tb}V_{ts}^{*}|}} \frac{1}{|C_{i}^{\text{NP}}|^{1/2}} \sim \frac{\mathbf{35\text{TeV}}}{|C_{i}^{\text{NP}}|^{1/2}}$$

• Loop level-induced (semi-leptonic) with  $\mathcal{O}(1)$  couplings:

$$\Lambda_i^{\text{Loop}} \sim \frac{35 \text{TeV}}{4\pi |C_i^{\text{NP}}|^{1/2}} = \frac{2.8 \text{TeV}}{|C_i^{\text{NP}}|^{1/2}}$$

• MFV with CKM-SM, suppression  $\sqrt{|V_{tb}V_{ts}^*|} \sim 1/5$ : Tree level:  $\frac{7 \text{ TeV}}{|C_i^{\text{NP}}|^{1/2}}$  and Loop:  $\frac{0.6 \text{ TeV}}{|C_i^{\text{NP}}|^{1/2}}$ Solution  $C_9^{\text{NP}} \sim -1.1$  (scale is ~ numerator) or  $C_9^{\text{NP}} = -C_{10}^{\text{NP}} \sim -0.6$  (30 % higher scale). Similar exercise for  $b \to c\tau\nu$  taking a 15% enhancement over SM:

$$\Lambda^{\rm NP} \sim 1/(\sqrt{2}G_F |V_{cb}| 0.15)^{1/2} \sim 3.2 \,{\rm TeV}$$

### Are we overlooking Lepton Flavour Universal NP?

[Algueró, Capdevila, SDG, Masjuan, JM, PRD'19]

## Hypothesis: Lepton Flavour Universality

We traded the usual controversy:

[Algueró, Capdevila, SDG, Masjuan, JM, PRD'19]

#### Is this New Physics or long-distance charm?

by a more constructive question:

Are we observing two kinds of New Physics?

$$\mathcal{C}_{i\ell}^{NP} = \mathcal{C}_{i\ell}^{V} + \mathcal{C}_{i}^{U}$$
 with  $i = 9, 10$   $\ell = e, \mu$ 

 $C_{ie}^{V} = 0$  Lepton Flavour **Universal** NP Lepton Flavour Universal **Violating** NP

....extended to primed operators in [Addendum: 1903.09578v3]

#### **Motivation:**

We have LFUV and LFD observables, then it is natural to split:

$$\mathcal{C}_{i\ell}^V \qquad \qquad \mathcal{C}_{i\ell}^V + \mathcal{C}_i^U$$
  
New mechanism to fulfill B<sub>s</sub>  $\rightarrow$   $\mu\mu$ 

## Is this the same as adding NP in electrons?

Many previous works already included NP in electrons: Mahmoudi et al. (large and low recoil data) London et al. (large and low recoil data) Ciuchini et al. (only large recoil data)

#### Which is the difference with our proposal?

All previous analyses explored directions within 2D planes in coordinates  $(C_{9\mu}, C_{10\mu})$  and  $(C_{9e}, C_{10e})$ 

#### instead the plane in coordinates $(C_9^v, C_{10}^v)$ in presence for instance of $C_9^v$ LFU can translate in a tilted plane in $(C_{9\mu}, C_{10\mu}, C_{9e})$ coordinates



## LFU updates 2019

1809.08447		Best-fit point	1 <i>σ</i>	Pull <sub>SM</sub>	p-value
Sc. 5	$\mathcal{C}_{9\mu}^{V} \ \mathcal{C}_{10\mu}^{V} \ \mathcal{C}_{9}^{U} = \mathcal{C}_{10}^{U}$	-0.16 +1.00 -0.87	$\begin{matrix} [-0.94, +0.46] \\ [+0.18, +1.59] \\ [-1.43, -0.14] \end{matrix}$	5.8	78%
Sc. 6	$egin{aligned} \mathcal{C}^{\mathrm{V}}_{9\mu} &= -\mathcal{C}^{\mathrm{V}}_{10\mu} \ \mathcal{C}^{\mathrm{U}}_{9} &= \mathcal{C}^{\mathrm{U}}_{10} \end{aligned}$	-0.64 -0.44	[-0.77, -0.51] [-0.58, -0.29]	6.0	79%
Sc. 7	$\mathcal{C}^{\mathrm{V}}_{9\mu}\ \mathcal{C}^{\mathrm{U}}_{9}$	-1.57 +0.56	[-2.14, -1.06] [+0.01, +1.15]	5.7	72%
Sc. 8	$egin{array}{c} \mathcal{C}^{\mathrm{V}}_{9\mu} = -\mathcal{C}^{\mathrm{V}}_{10\mu} \ \mathcal{C}^{\mathrm{U}}_{9} \end{array} \end{array}$	-0.42 -0.67	[-0.57, -0.27] [-0.90, -0.42]	5.8	74%
	2019	Best-fit point	1 σ		p-value
Sc. 5	$\mathcal{C}^{\mathrm{V}}_{9\mu}\ \mathcal{C}^{\mathrm{V}}_{10\mu}\ \mathcal{C}^{\mathrm{U}}_{9}=\mathcal{C}^{\mathrm{U}}_{10}$	-0.34 +0.69 -0.50	$ \begin{bmatrix} -0.93, +0.19 \\ [+0.21, +1.12] \\ [-0.92, +0.02] \end{bmatrix} $	5.5	72%
Sc. 6	$\mathcal{C}^{ m V}_{9\mu} = - \mathcal{C}^{ m V}_{10\mu} \ \mathcal{C}^{ m U}_{9} = \mathcal{C}^{ m U}_{10}$	-0.52 -0.37	[-0.64, -0.41] [-0.52, -0.22]	5.8	71%
Sc. 7	$\mathcal{C}^{V}_{9\mu}\ \mathcal{C}^{U}_{9}$	-0.91 -0.08	[-1.25, -0.58] [-0.46, +0.31]	5.5	65%
Sc. 8	$\mathcal{C}_{9\mu}^{\mathrm{V}}=-\mathcal{C}_{10\mu}^{\mathrm{V}}$ $\mathcal{C}_{9}^{\mathrm{U}}$	-0.33 -0.72	$\begin{bmatrix} -0.45, -0.22 \\ [-0.93, -0.47] \end{bmatrix}$	5.9	74%

#### Changed

Sc. 7: If only V-NP  $(C_9)$  now preference for LFUV-C<sub>9</sub>

$$\mathcal{C}_{9\mu}^V + \mathcal{C}_9^U = -0.98$$

#### Unchanged

Sc. 8: A LFUV left-handed lepton struc.  $(C_9^V=-C_{10}^V)$ **yields a better description** with LFU-NP in C<sub>9</sub>.

#### Still

Sc. 6: A LFUV V-A struc.  $(C_9^{v}=-C_{10}^{v})$  and a LFU V+A struc. provides a good description of data.

#### • LFU-NP is quite dependent on structure of LFUV-NP

## LFU updates 2019

#### [1903.09578]

Scenario		Best-fit point	$1 \sigma$	Pull <sub>SM</sub>	p-value
Sc. 9	$\mathcal{C}_{9\mu}^{\mathrm{V}} = -\mathcal{C}_{10\mu}^{\mathrm{V}}$	-0.63	[-0.79, -0.47]	5.3	73.4%
	$\mathcal{C}_{10}^{0}$	-0.39	[-0.65, -0.13]		
Sc 10	$\mathcal{C}_{9\mu}^{\mathrm{V}}$	-0.99	[-1.17, -0.80]	57	69.7%
00.10	$\mathcal{C}_{10}^{\mathrm{U}}$	+0.29	[0.10, 0.48]	0.1	
Sc. 11	$\mathcal{C}_{9\mu}^{\mathrm{V}}$	-1.07	[-1.25, -0.88]	5.0	73 0 %
50.11	$\mathcal{C}_{10'}^{\mathrm{U}}$	-0.31	[-0.48, -0.13]	0.9	10.970
Sc. 12	$\mathcal{C}_{9'\mu}^{V}$	-0.05	$\left[-0.23, 0.14\right]$	17	131%
50.12	${\cal C}_{10}^{ m U}$	+0.43	$\left[0.22, 0.65\right]$	1.1	10.1 /0
Sc. 13	$\mathcal{C}_{9\mu}^{\mathrm{V}}$	-1.12	[-1.29, -0.94]		
	$\mathcal{C}^{\mathrm{V}}_{9'\mu}$	+0.48	$\left[0.19, 0.85\right]$	56	78 7 %
	$\mathcal{C}_{10}^{U}$	+0.26	[0.01, 0.50]	0.0	10.1 /0
	$\mathcal{C}^{\mathrm{U}}_{10'}$	-0.05	[-0.28, 0.18]		

- Sc. 9 versus Sc.10 preference of  $C_9^V$  versus  $C_9^{V=-}C_{10}^V$  in presence of  $C_{10}^U$ , opposite to the case of  $C_9^U$  (sc.7-8).

Sc. 10 versus Sc.11 shows a slight preference of  $C_{10}$ ,<sup>U</sup> over  $C_{10}$ <sup>U</sup>.

- Sc.12 irrelevance of RHC without  $C_9^{v}$ . If  $C_{10}^{v} \rightarrow C_9^{v}$  then  $4\sigma$ 

#### Changed

Sc. 7: If only V-NP  $(C_9)$  now preference for LFUV-C<sub>9</sub>

$$\mathcal{C}_{9\mu}^V + \mathcal{C}_9^U = -0.98$$

#### Unchanged

Sc. 8: A LFUV left-handed lepton struc.  $(C_9^{V}=-C_{10}^{V})$ **yields a better description** with LFU-NP in C<sub>9</sub>.

#### New

Sc.9-13: We extend the universal contribution also to **primed universal coefficients** associated to models.

• Sc.7-10 show LFU-NP is quite dependent on structure of LFUV-NP

### LFU updates 2019



Assuming loop-level  
scale of NP and no MFV  
$$\Lambda_i^L \sim \frac{v}{s_w g} \frac{1}{\sqrt{2|V_{tb}V_{ts}^*|}} \frac{1}{|\mathcal{C}_i^{NP}|^{1/2}}$$
$$Mild preference$$
$$Scenario 6: \qquad \begin{pmatrix} \mathcal{C}_{9\mu}^v = -\mathcal{C}_{10\mu}^v \\ \mathcal{C}_{9}^v = \mathcal{C}_{10}^{U} \end{pmatrix}$$
$$LFUV-NP \quad L_q \otimes L_\ell$$
$$\Lambda_i^{\text{LFUV}} \sim 3.9 \text{ TeV}$$
$$LFU-NP \quad L_q \otimes R_\ell$$
$$\Lambda_i^{\text{LFU}} \sim 4.6 \text{ TeV}$$
$$Scenario 8: \qquad \begin{pmatrix} \mathcal{C}_{9\mu}^v = -\mathcal{C}_{10\mu}^v \\ \mathcal{C}_{9}^v = \mathcal{C}_{10\mu}^v \end{pmatrix}$$
$$LFUV-NP \quad L_q \otimes L_\ell$$
$$\Lambda_i^{\text{LFU}} \sim 4.6 \text{ TeV}$$
$$LFUV-NP \quad L_q \otimes L_\ell$$
$$\Lambda_i^{\text{LFUV}} \sim 4.6 \text{ TeV}$$
$$LFU-NP \quad L_q \otimes V_\ell$$
$$\Lambda_i^{\text{LFUV}} \sim 3.3 \text{ TeV}$$

- If we are in presence of two types and scales of NP, their hierarchy depend on the LFU

## **Results from other analysis**

[Aebischer, Altmannshofer, Guadagnoli, Reboud, Stangl, Straub]

Similar results in general terms **but** 1D differences. Why?

Coeff.	best fit	$1\sigma$	$2\sigma$	$\operatorname{pull}$
$C_9^{bs\mu\mu}$	-0.95	[-1.10, -0.79]	[-1.26, -0.63]	5.8σ
$C_9^{\prime b s \mu \mu}$	+0.09	[-0.07, +0.24]	[-0.23, +0.39]	$0.5\sigma$
$C^{bs\mu\mu}_{10}$	+0.73	[+0.59, +0.87]	[+0.46, +1.01]	$5.6\sigma$
$C_{10}^{\prime bs\mu\mu}$	-0.19	[-0.30, -0.07]	[-0.41, +0.04]	$1.6\sigma$
$C_9^{bs\mu\mu} = C_{10}^{bs\mu\mu}$	+0.20	[+0.05, +0.35]	[-0.09, +0.51]	$1.4\sigma$
$C_9^{bs\mu\mu}=-C_{10}^{bs\mu\mu}$	-0.53	[-0.62, -0.45]	[-0.70, -0.36]	6.5 <i>σ</i>

 Difference in observable sets: BR(b→ sℓℓ) (B, B<sub>s</sub>, h<sub>b</sub> (BR, P<sub>i</sub>), R<sub>K(\*)</sub>, b→ sγ favours mildly C<sub>9µ</sub> = -C<sub>10µ</sub>
 But latest Belle updates on P<sub>5</sub>' and Q<sub>5</sub> are missing
 Extra assumption: no NP in ΔF=2 observables

=> constraints inputs for  $B_s \rightarrow \mu \mu (f_{B_s}, V_{tb} V_{ts}^* \dots)$ 

**Different** question: Is there NP in b $\rightarrow$ Sll assuming no NP in  $\Delta F=2$ 



## $P_5$ under different scenarios



## **Results from other analysis**

[Arbey, Hurth, Mahmoudi, Martinez Santos, Neshatpour]

**Obs**:  $b \to s\ell\ell (B, B_s) (BR, S_i), R_{K(*)}, b \to s\gamma$ not included yet latest Belle's results on P<sub>5</sub>'. FF: light-meson LCSR+lattice

Left-handed hypothesis considered. ... similar 1D and 2D results

Confirm our hierarchy of 1D scenarios

[Alok, Dighe, Gangal, Kumar]

 $\chi^2_{\rm min}$ Pull<sub>SM</sub> b.f. value 99.2  $\delta C_9$  $-1.01 \pm 0.20$  $4.2\sigma$  $\delta C_9^{\mu}$  $5.3\sigma$  $-0.93 \pm 0.17$ 89.4 $\delta C_9^e$  $3.2\sigma$  $0.78 \pm 0.26$ 106.6 $\delta C_{10}$  $0.25 \pm 0.23$ 115.7 $1.1\sigma$  $\delta C_{10}^{\mu}$  $0.53\pm0.17$ 105.8 $3.3\sigma$  $\delta C_{10}^e$  $-0.73 \pm 0.23$ 105.2 $3.4\sigma$  $\delta C^{\mu}_{\rm LL}$  $-0.41\pm0.10$ 96.6 $4.5\sigma$  $\delta C^e_{\mathrm{LL}}$  $0.40 \pm 0.13$  $3.3\sigma$ 105.8

All observables ( $\chi^2_{\rm SM} = 117.03$ )

 $\delta \mathcal{C}_{LL}^{\ell} = \delta \mathcal{C}_{9}^{\ell} = -\delta \mathcal{C}_{10}^{\ell}$ 

122 **Obs**:  $BR(b \rightarrow s\ell\ell)$   $(B, B_s)$ ,  $P'_5 R_{K(*)}$  FF: light-meson LCSR+lattice Flavio based analysis: slight decrease of SM pull for  $(C_{9\mu}, C_{10\mu})$ , at the same level as  $(C_{9\mu}, C_{9'\mu})$  and  $(C_{9\mu}, C_{10'\mu})$  ...very similar results to ours [Ciuchini et al.]

Only large-recoil obs. considered, but latest Belle results on  $P_5$ ' incl. Flavio based analysis for FF. Bayesian approach. OK: RHC and not  $C_{10}$ .

### Linking charge and neutral anomalies and LFU NP

### LFUV for charged anomalies $b \rightarrow c \tau v$



SM Semi-tauonic B decays are charged current processes that can probe also New Physics. Experimentally (in analogy to  $R_{K,K^*}$ ) a LFUV ratio:

$$R_{D^{(*)}} = \frac{\mathcal{B}(\bar{B} \to D^{(*)}\tau^-\bar{\nu}_{\tau})}{\mathcal{B}(\bar{B} \to D^{(*)}\ell^-\bar{\nu}_{\ell})}$$

The ratio:

- differs in lepton mass:  $\tau$  versus  $\ell = \mu, e$  mass.
- cancels: form factors,  $V_{cb}$ , experimental systematics

## R(D) and $R(D^*)$ combination



#### [From Julian Garcia Pardiñas, UZH]

New world average for R(D) and R(D\*) at 3.1 $\sigma$  from the SM

## Linking charged and neutral anomalies (step 1)

Let's move to SMEFT ( $\Lambda_{NP} >> m_{t,W,Z}$ )

[Grzadkowski, Iskrzynski, Misiak, Rosiek; Alonso, Grinstein, Camalich]

• **NP contribution to** :  $[\bar{\mathbf{c}}\gamma^{\mu}\mathbf{P}_{\mathbf{L}}\mathbf{b}][\bar{\tau}\gamma_{\mu}\mathbf{P}_{\mathbf{L}}\nu_{\tau}] \longrightarrow R_{J/\psi}/R_{J/\psi}^{\mathrm{SM}} = R_D/R_D^{\mathrm{SM}} = R_{D^*}/R_{D^*}^{\mathrm{SM}}$ 

 $G_F$  rescaling

**BUT** who order that

(at high energy)? Only Two SU(2)<sub>L</sub> invariant operators in SMEFT @ 1st order



## Linking anomalies with LFU NP (step 2)



## Some Solutions to the anomalies

## Solution to anomalies, generation of couplings

Colourless vector  $SU(2)_L$  triplets (W', B') or U(1)' singlet

 $G \equiv SU(3)_c \times SU(2)_L \times U(1)_Y \times G_E$ 



#### **Generating Quark FV Coupling:**

Vector-like quarks: SM-VL couplings

#### $b_L$ $\phi$ Q Z' $s_L$ $\phi^*$

Loop induced: SM FCNC, Z' penguins

 $G_E \equiv SU(2)_L$  could pot. explain anomalies  $(R_K > 0.9$  & conflict with LHC searches)

- $\bar{b}sZ'$  Quark FVC
- $Z'\ell\ell$  LFUV coupling

#### **Generating Couplings to Leptons:**

- Gauged  $U(1)_{\mu-\tau}$  symmetry
- Loop induced with vector-like fermions



## Solution to anomalies: leptoquarks



- Spin 1 (vector)  $SU(2)_L \underline{\text{singlet}}$  or  $\underline{\text{triplet}}$  leptoquarks
- Spin 0 (scalar)  $SU(2)_L$  singlet or triplet leptoquarks
- via loop....

They mainly point in all versions to  $C_9 = -C_{10}$  (left-handed structure like in the SM)

Important constraints:

- $b \rightarrow s \nu \bar{\nu}$  (two scalars LQ can do the job)
- direct bounds (from 0.5-1 TeV)

Colour triplet  
Scalar LQ:  

$$S_1 \sim (3, 1, 1/3)$$
  
 $S_3 \sim (3, 3, 1/3)$   
Vector LQ:  
 $U_1'' \sim (3, 1, 2/3)$   
 $U_3'' \sim (3, 3, 2/3)$ 



## Solution to anomalies: leptoquarks



- Spin 1 (vector)  $SU(2)_L \underline{\text{singlet}}$  or triplet leptoquarks
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Important constraints:

- $b \rightarrow s \nu \bar{\nu}$  (two scalars LQ can do the job)
- direct bounds (from 0.5-1 TeV)

Colour triplet  
Scalar LQ:  

$$S_1 \sim (\overline{3}, 1, 1/3)$$
  
 $S_3 \sim (\overline{3}, 3, 1/3)$   
Vector LQ:  
 $N_1^{\mu} \sim (3, 1, 2/3)$   
 $U_3^{\mu} \sim (3, 3, 2/3)$   
 $b \qquad \nu$   
 $i \phi$   
 $j \phi$   

A very promising candidate is:

Vector leptoquark SU(2) singlet:  $U_1(3,1,2/3)$ Coupled mainly to 3<sup>rd</sup> generation

1. It can explain both charged and neutral anomalies2.  $C_9=-C_{10}$  pattern in  $b \rightarrow s \mu \mu$ 3. No tree level effect for  $b \rightarrow s v v$ 4. No conflict with direct searches

## Good solution, but challenging UV completion

## Possible UV completions

- SU(4)×SU(3)'×SU(2)<sub>L</sub>×U(1)<sub>Y</sub> + Vector-like fermions
   L. Di Luzio, A. Greljo, M. Nardecchia, arXiv:1708.08450
- SU(4)×U(2)<sub>L</sub>×SU(2)<sub>R</sub> + vector-like fermions
   L. Calibbi, AC, T. Li, arXiv:1709.00692
- SU(4) × SU(4) × SU(4)

M. Bordone, C. Cornella, J. Fuentes-Martin, G. Isidori, arXiv:1712.01368

- SU(4) ×U(2)<sub>L</sub>×SU(2)<sub>R</sub> including scalar LQs and light right-handed neutrinos
   J. Heeck, D. Teresi, arXiv:1808.07492
- SU(8) might even explain ε'/ε

S. Matsuzaki, K. Nishiwaki and K. Yamamoto, arXiv:1806.02312

SU(4)×U(2)×SU(2)<sub>R</sub> in RS background
 M. Blanke, AC, arXiv:1801.07256

## Good solution, but challenging UV completion

## Pati-Salam LQ model

## Original PS=SU(4) x SU(2)<sub>L</sub> x SU(2)<sub>R</sub>

It does not work: tight bounds from couplings to light generation:  $K_L \rightarrow \mu e$  and  $K \rightarrow \pi \mu e$ 

... too heavy (Flavour-Blind) to work.

[M. Bordone et al.]

A recent proposal :  $PS^3 \equiv PS_1 \times PS_2 \times PS_3$ 

 $PS_i = SU(4)_i \times [SU(2)_L]_i \times [SU(2)_R]_i$ 

1. SSB decouple very heavy fields coupled to 1<sup>st</sup> & 2<sup>nd</sup> gen.

2. TeV-scale LQ associated to  $3^{rd}$  gen and LQ coupling to RH SM

3. Higgs of EWSB only on third generation site:

... yukawa hierarchies from hierarchy of breaking vevs

## Near Future next test: $Q_5 = P'_{5\mu} - P'_{5e}$

## What can we learn?

## Q<sub>5</sub> can disentangle relevant scenarios?

 $R_K$  (if no-RHC are included) cannot distinguish among relevant scenarios.



[Alguerò, Capdevila, SDG, Masjuan, JM: 1902.04900]

of  $\mathbf{R}_{\mathbf{K}}$  are still too packed within 0.5  $\sigma$  to disentangle the correct pattern.

## $Q_5$ can disentangle relevant scenarios?

Only Belle has been able to measure  $Q_5$  up to now:  $Q_5[1,6]^{Belle} = 0.656 \pm 0.496$ 

![](_page_49_Figure_2.jpeg)

[S. Wehle et al. PRL118 (2017)]

**Table 2:** Results for the lepton-flavor-universality-violating observables  $Q_4$  and  $Q_5$ . The first uncertainty is statistical and the second systematic.

-		
$q^2$ in GeV <sup>2</sup> / $c^2$	$Q_4$	Q5
[1.00, 6.00]	$0.498 \pm 0.527 \pm 0.166$	$0.656 \pm 0.485 \pm 0.103$
[0.10, 4.00]	$-0.723 \pm 0.676 \pm 0.163$	$-0.097 \pm 0.601 \pm 0.164$
[4.00, 8.00]	$0.448 \pm 0.392 \pm 0.076$	$0.498 \pm 0.410 \pm 0.095$
[14.18, 19.00]	$0.041 \pm 0.565 \pm 0.082$	$0.778 \pm 0.502 \pm 0.065$

## $Q_5$ can disentangle relevant scenarios?

![](_page_50_Figure_1.jpeg)

All scenarios with  $C_{9}^{v}$  are packed as well as those with  $C_{9}^{v} = -C_{10}^{v}$  BUT in two different sets. Also: \* $Q_5$  positive and large would favour scenarios with  $C_{9\mu} < -1$ \* $Q_5 < 0$  or small would favour scenarios with  $C_{9\mu=-}C_{10\mu}$ 

## Conclusions

• After the updates of  $R_K$  (LHCb),  $R_{K^*}$ (Belle) and  $B_s \rightarrow \mu \mu$  we find:

- **no dramatic changes** in the hierarchy of 1D hypothesis:  $C_9$  and  $C_9$ =- $C_9$ ' preferred with All fit [178 obs] significance 5.8 (5.7)  $\sigma$  $C_9$ =- $C_{10}$  preferred with LFUV fit [20 obs] significance 4.0  $\sigma$ 

- 2D **new emerging scenarios including RHC** with C<sub>9</sub>'>0 & C<sub>10</sub>'<0:  $(C_{9\mu}, C'_{9\mu} = -C'_{10\mu})$  (6.1  $\sigma$ )
- LFU-NP structure is **quite dependent** on LFUV-NP structure:  $A C_9^{V=-}C_{10}^{V}$  struct. provides a good description only in presence of  $C_9^{U}$
- We have found a link of charged & neutral anomalies & LFU NP in scn 8.
- While  $R_K$  cannot disentangle scenarios, **a measurement of Q\_5** such that:

-Q<sub>5</sub> **positive and large** would **favour** scenarios with  $C_{9\mu}$ <-1 -Q<sub>5</sub> < 0 or small would **favour** scenarios with  $C_{9\mu=-}C_{10\mu}$ 

.... new data on  $Q_5$ ,  $R_{\phi}$ , updated optimized observables is essential. Belle II inputs are also crucial. BACK-UP

### ... in summary

#### [Courtesy of A. Crivellin]

![](_page_53_Figure_2.jpeg)

### Pati-Salam LQ can explain the flavour anomalies

Different theory approaches to **estimate/predict** "LD charm":

**Long distance non-factorizable** : LCSR by Khodjamirian

+ s<sub>i</sub> const/destr interference.

#### **Empirical model to determine** the impact of resonances :

(amp. analysis+BW) Blake et al. '17

#### LD effects from analyticity:

(fixes q<sup>2</sup> dep. up to pol. & systematic) Bobeth et al.'18

In all theoretical estimates the anomaly remains.

Different theory approaches to **estimate/predict** "LD charm":

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#### **LD effects from analyticity:** (fixes q<sup>2</sup> dep. up to pol. & systematic) Bobeth et al.'18

In all theoretical estimates the anomaly remains.

![](_page_55_Figure_8.jpeg)

Different theory approaches to **estimate/predict** "LD charm":

LOSR by Khodjamirian + s<sub>i</sub> const/destr interference.

Empirical model to determine the impact of resonances :

> (amp. analysis+BW) Blake et al. '17

**LD effects from analyticity:** (fixes q<sup>2</sup> dep. up to pol. & systematic) Bobeth et al.'18

In all theoretical estimates the anomaly remains.

![](_page_56_Figure_7.jpeg)

![](_page_57_Figure_1.jpeg)