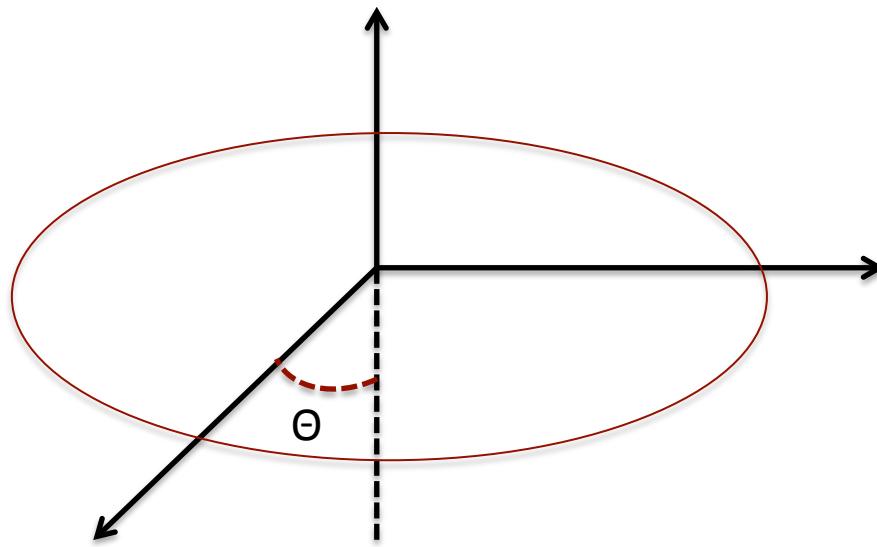
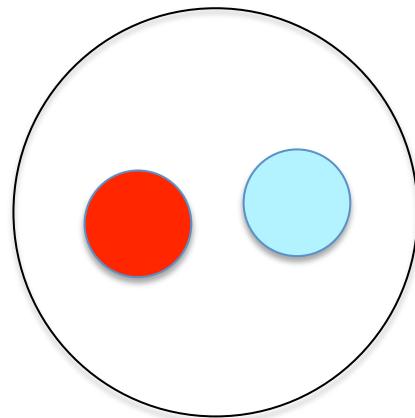


Dark Matter in Composite Models



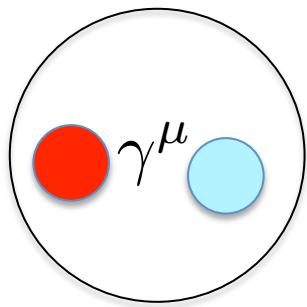
Mads Toudal Frandsen

Motivation

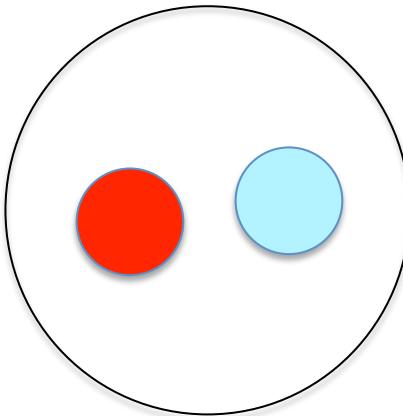


Motivation

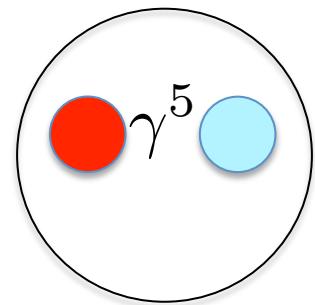
Resonances



Higgs

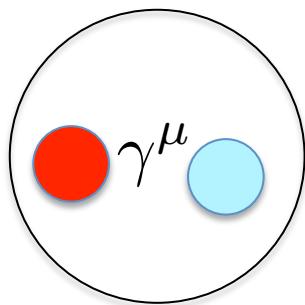


DM

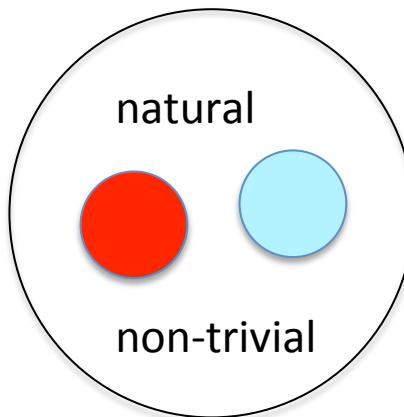


Motivation

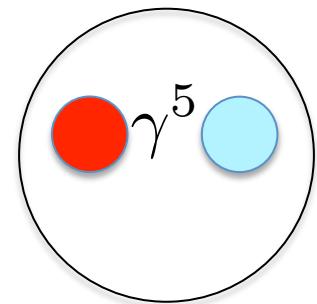
Resonances



Higgs



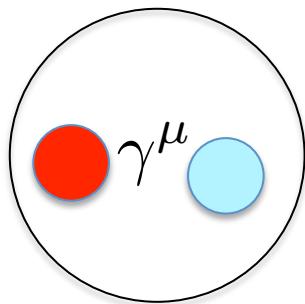
DM



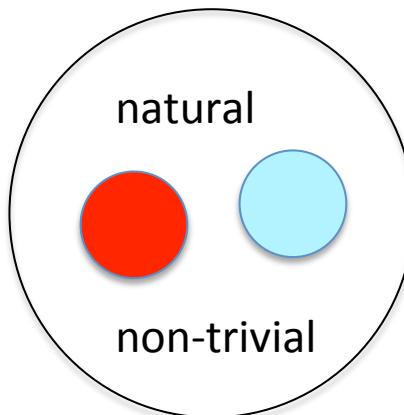
(Weinberg; Susskind & Dimopoulos;
Eichten & Lane; 't Hooft; Kaplan & Georgi;
Nussinov)

Motivation

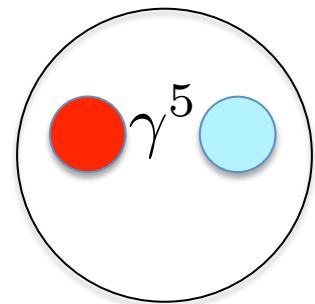
Resonances



Higgs



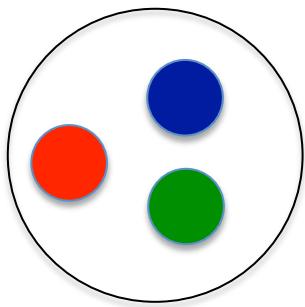
DM



	SU(2) _{TC}	SU(2) _W	U(1) _Y
(U_L, D_L)	□	□	0
\tilde{U}_L	□	1	-1/2
\tilde{D}_L	□	1	+1/2

Motivation

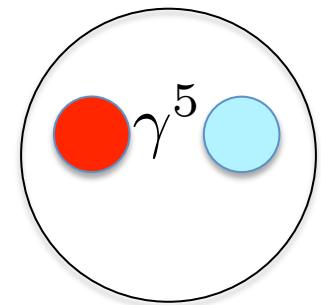
Baryons



Baryon and DM asymmetries

Sphalerons

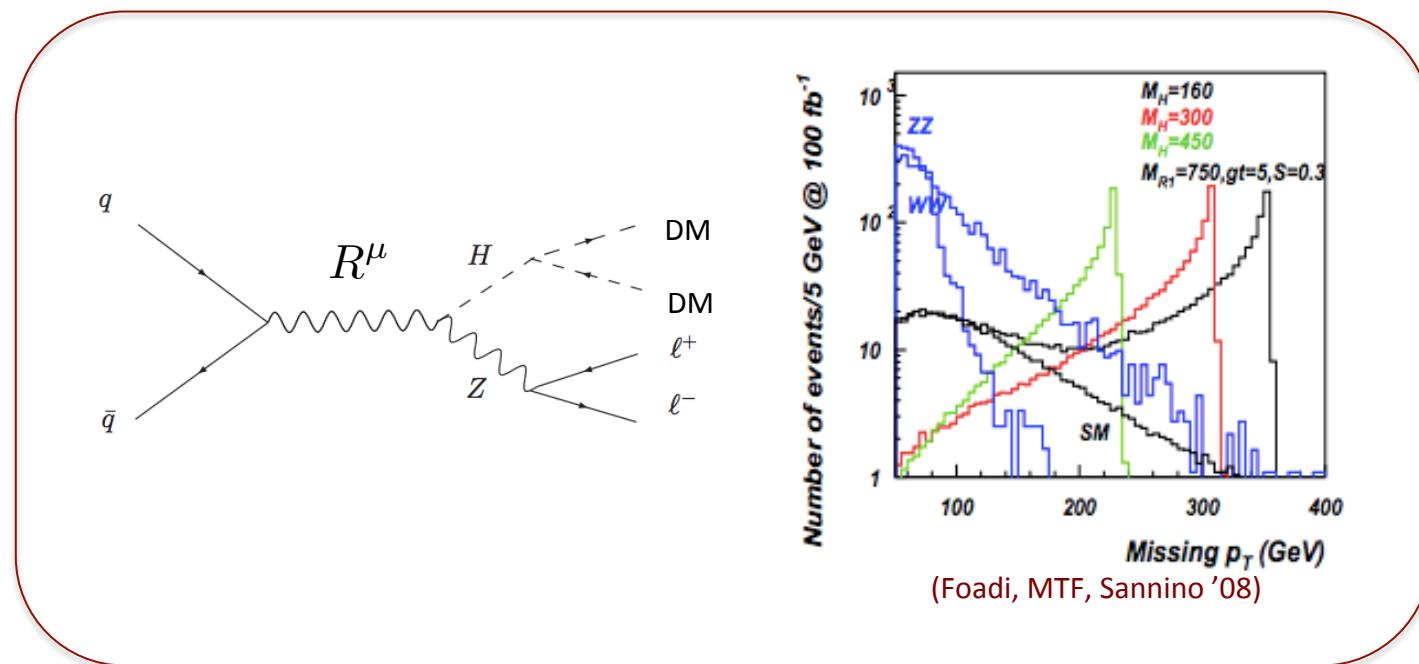
DM



(Barr, Chivukula & Farhi)

Composite DM Probes

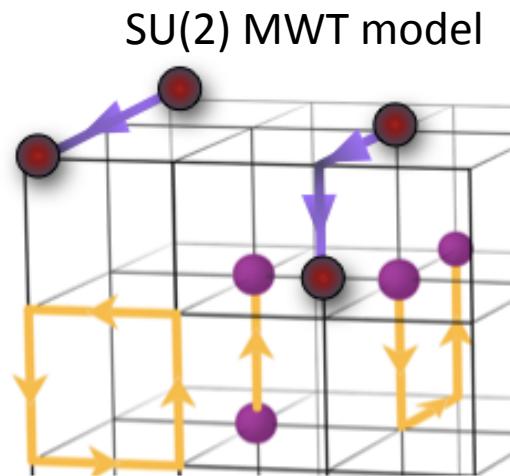
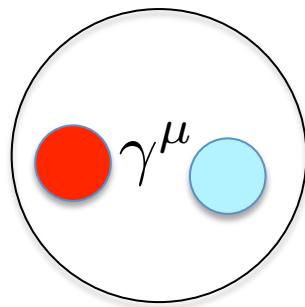
- Resonances in e.g LHC missing E_T searches (U. Haisch talk ZPW19)



Composite DM Probes

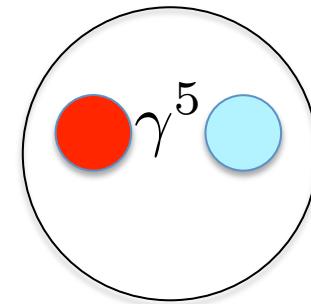
- lattice predictions, low mass DM (F. Petricca talk ZPW19)

$m_R = 2.5 \text{ (0.5) TeV}$,



(Arthur, Drach, Hansen, Hietanen,
Lewis, Pica & Sannino '14)

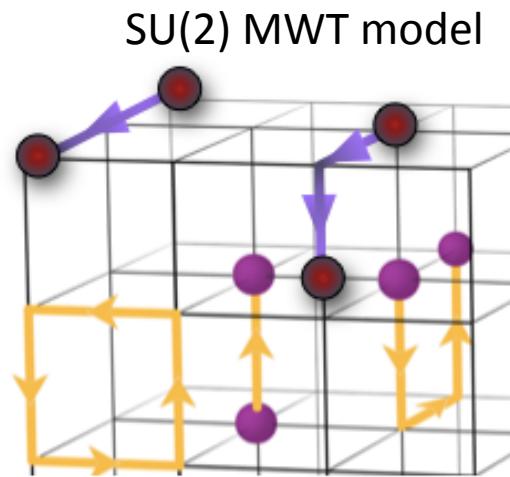
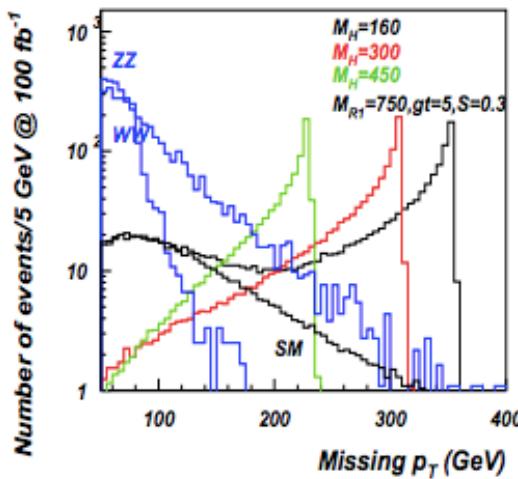
DM pNGB,



Composite DM Probes

- lattice predictions, low mass DM (F. Petricca talk ZPW19)

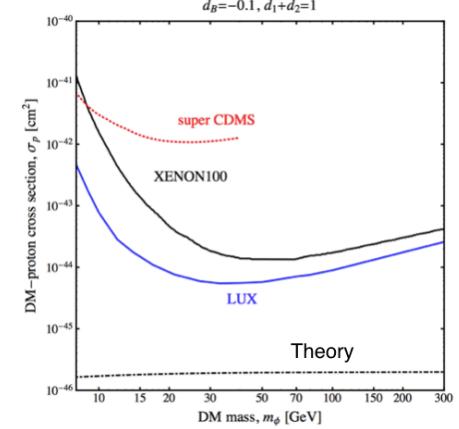
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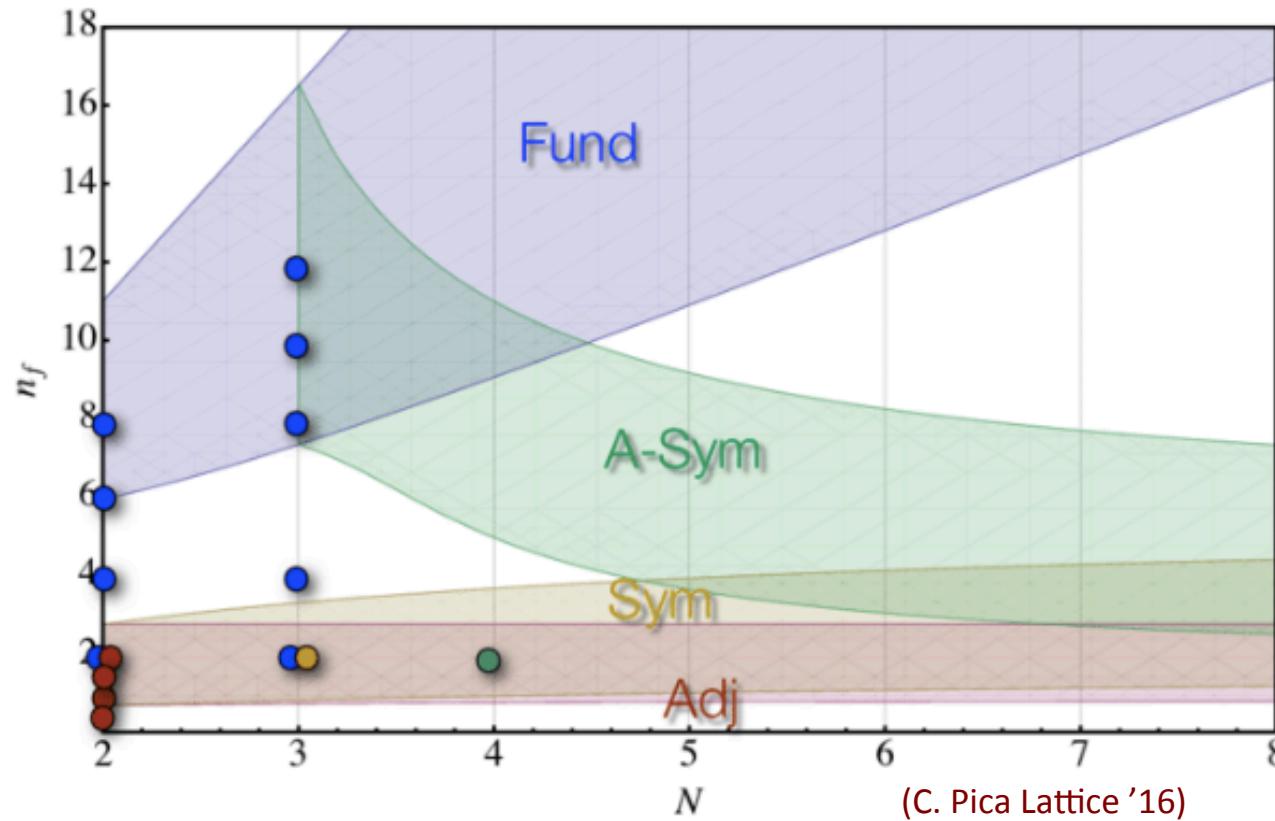
DM pNGB, $\Lambda = m_R$

$$ie \frac{d_B}{\Lambda^2} \Pi_{UD}^* \partial_\mu \Pi_{UD} \partial_\nu F^{\mu\nu}$$
$$d_B = (m_U - m_D)/(m_U + m_D)$$



Composite DM Probes

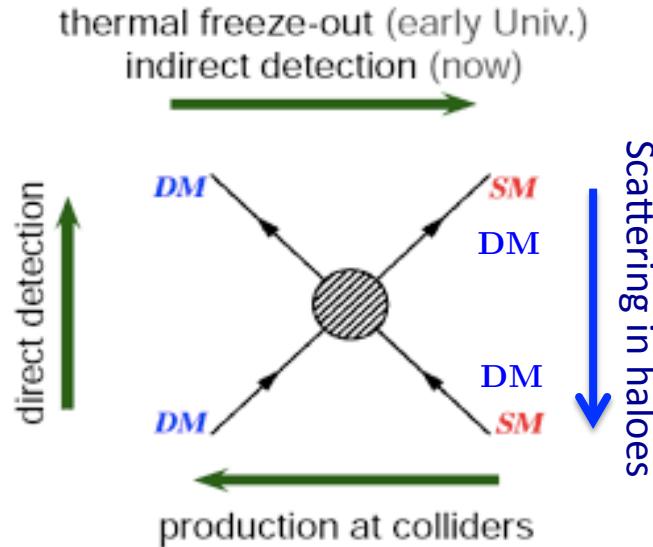
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Composite DM Probes

- Large (in some cases) self-interactions in clusters and galaxies

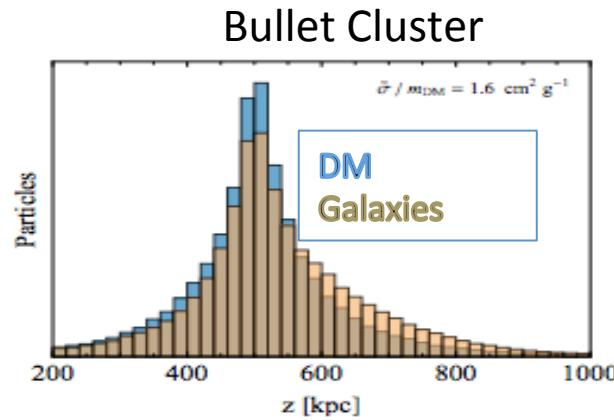
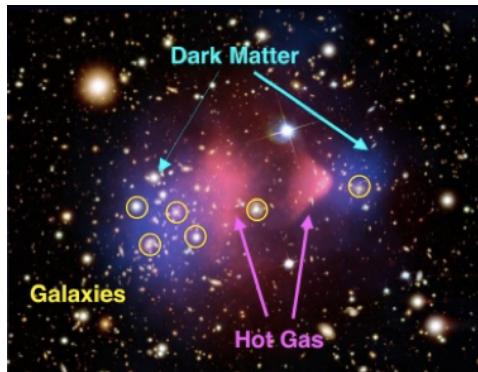
(J. Read talk ZPW19)



Composite DM Probes

- Large (in some cases) self-interactions in clusters and galaxies

(J. Read talk ZPW19)

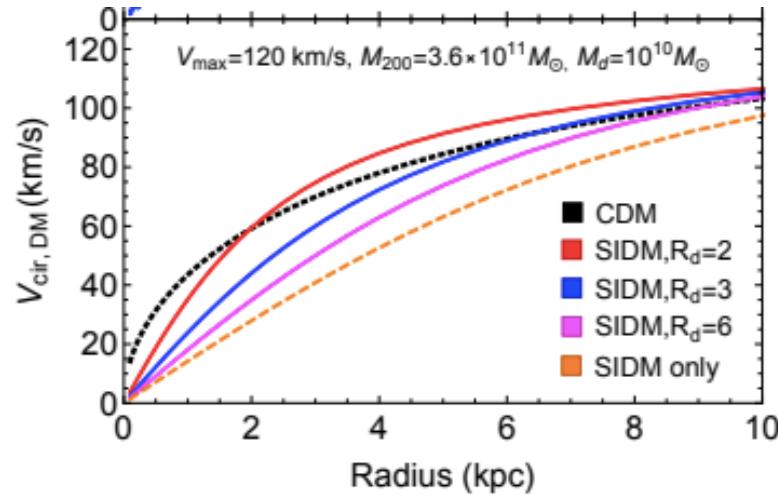


(Kahlhoefer, Schmidt-Hoberg, M.T. F & Sarkar '13)

Composite DM Probes

- Large (in some cases) self-interactions in clusters and galaxies

(J. Read talk ZPW19)

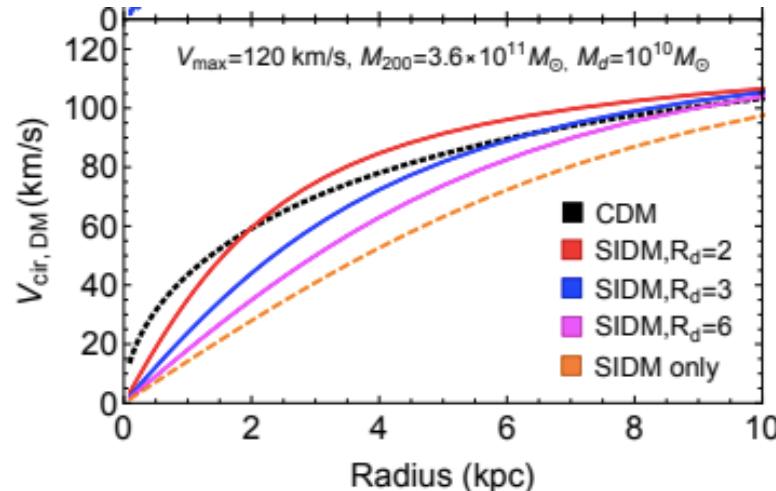
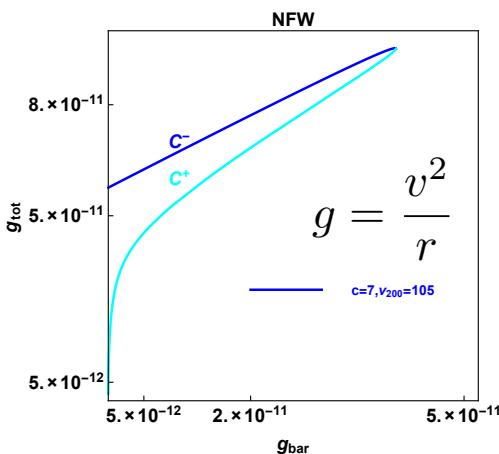


(Kamada, Kaplinghat Pace & Yu '16)

Composite DM Probes

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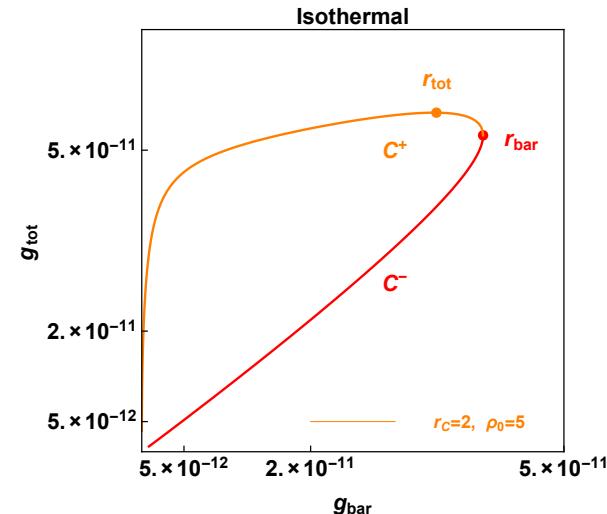
(J. Read talk ZPW19)



(MTF & Petersen '18)

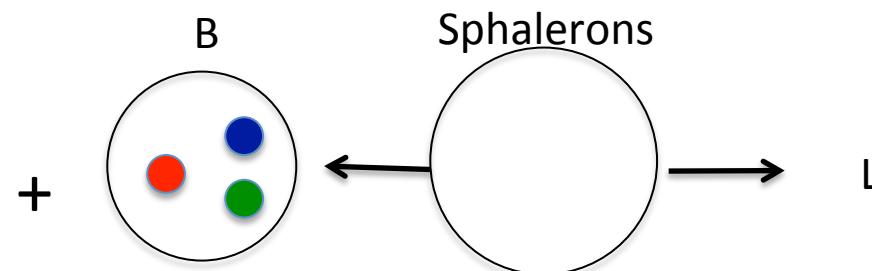
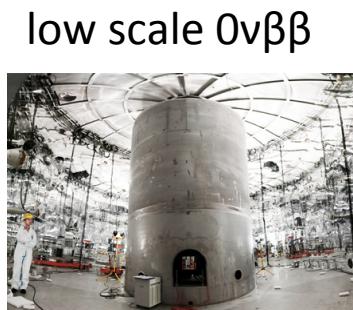
(Kamada, Kaplinghat Pace & Yu '16)

(MTF & Petersen '18)



Composite DM Probes

- Non-standard signatures from asymmetry:
decaying DM, no annihilation signals, possible low scale $0\nu\beta\beta$

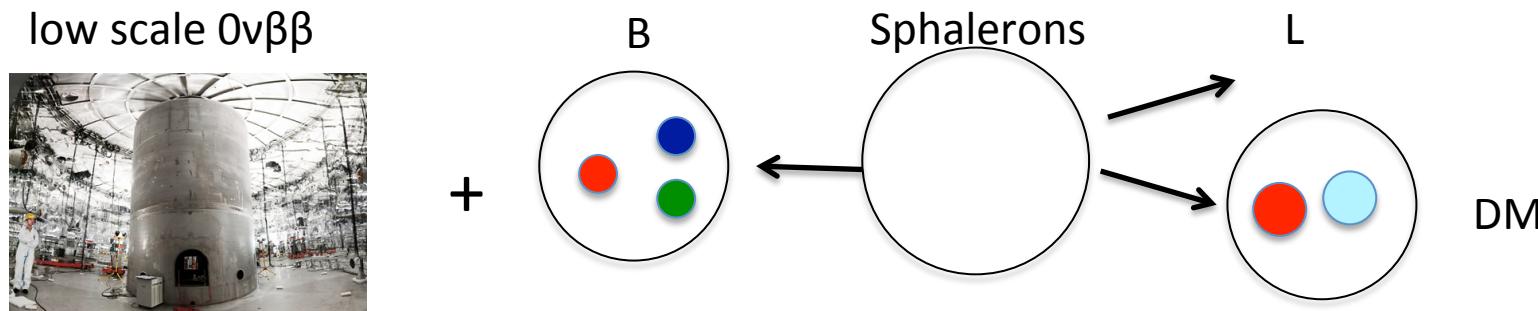


$$\mu_\ell + \mu_H = 0 . \quad \& \quad 3(3\mu_q + \mu_\ell) = 0 \quad \longrightarrow \quad B=0$$

(Harvey & Turner)

Composite DM Brobes

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decaying DM, no annihilation signals, possible low scale $0\nu\beta\beta$



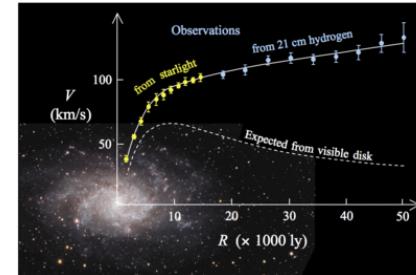
$$\mu_\ell + \mu_H = 0 . \quad \& \quad 3(3\mu_q + \mu_\ell) + n_F\mu_F = 0, \longrightarrow B \neq 0, DM \neq 0$$

(MTF, Hagedorn, Huang,
Molinaro & Päs '18)

The missing mass problem

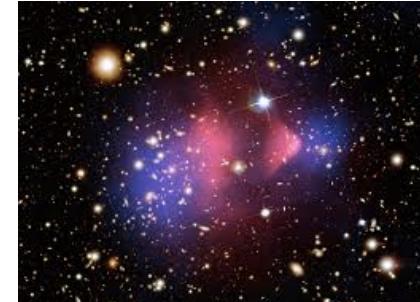
Galactic scales

(Freeman '70, Bosma '78, Rubin et al '78)



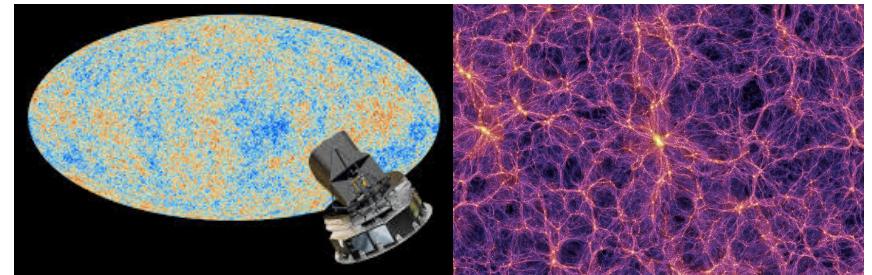
Cluster Scales

(Zwicky '33, Clowe et al '06)



Cosmological scales

(Davis et al '82, Peebles '82)



Universality vs Diversity

NFW scaling of radial DM density profiles from dwarf galaxies to galaxy clusters

DM only cosmological N-body simulations. No baryons

(Navarro, Frenk & White '95)

MOND/RAR: Total acceleration in circular motion correlates with that from baryons

At galactic scales

(Tully & Fisher '77; McGaugh '11)

Cusp/Core

Flat or cored DM profiles in dwarf and Low Surface Brightness galaxies

(Moore '94; Flores & Primack '94)

Missing Satellites & Too-big-to-fail

Too many satellites that are too dense and massive are predicted

(Moore, Quinn, Governato, Stadel & Lake '99; Klypin, Kravtsov, Valenzuela & Prada '99; Boylan-Kolchin, Bullock, Kaplinghat '11)

Universality vs Diversity

Baryonic physics

Discrepancies arise from comparing to DM-only simulations

(J. Read talk ZPW19)

Adiabatic contraction

Supernova feedback

AGN

DM (self-) interactions

Small scales are high DM density and DM interaction rates

Velocity dependence

Long range, short range

Gravity

Discrepancies arise assuming Newtonian gravity

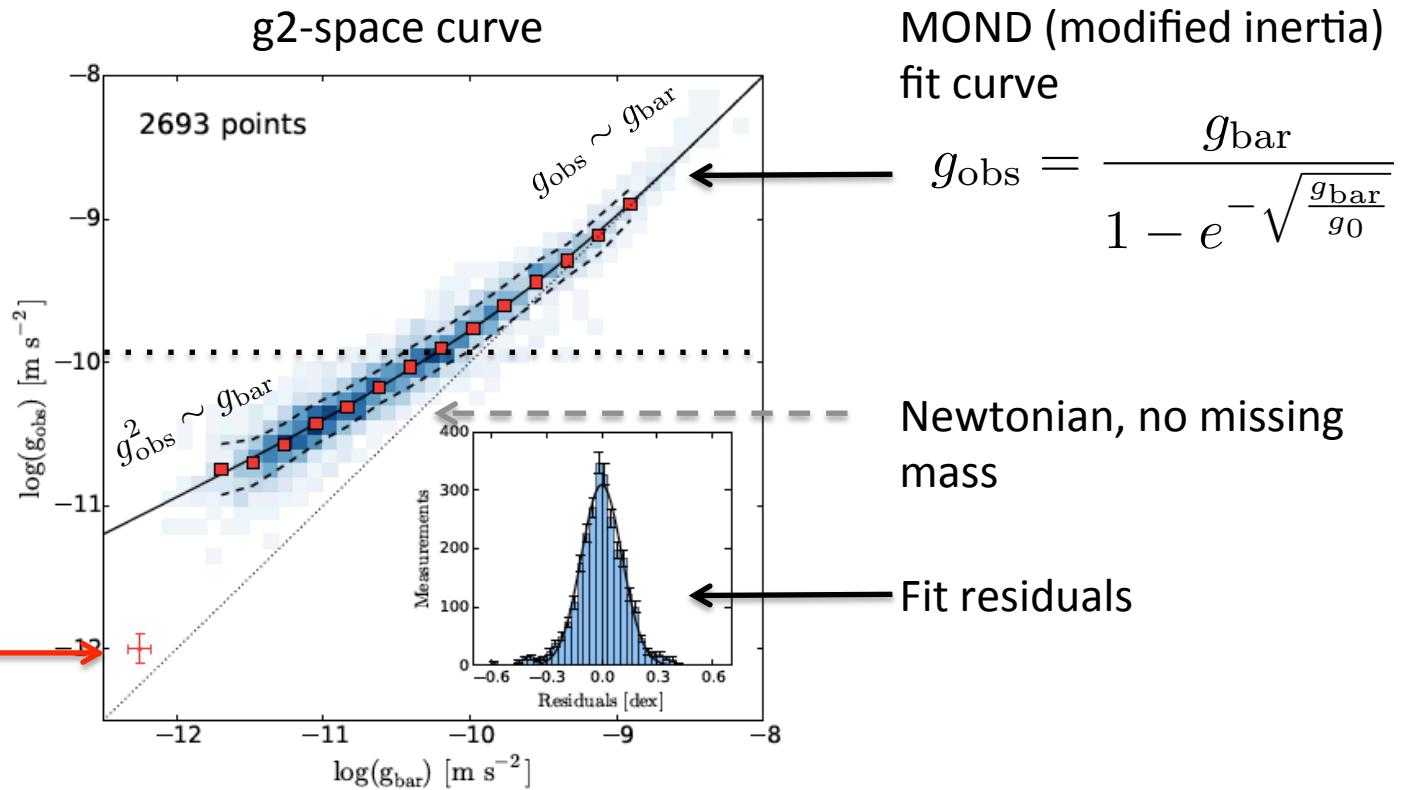
With or without DM

Modified Newtonian Dynamics

Modified Gravity

SPARC Radial Acceleration Relation

(McGaugh, Lelli & Schombert '16)



MOND

$$\mu\left(\frac{g_M}{g_0}\right)g_M = g_{\text{bar}}$$

MOND modified inertia

(Milgrom '83, '93)

Inverse interpolation function ν gives *exact* MOND modified inertia acceleration

$$g_M = \nu\left(\frac{g_{\text{bar}}}{g_0}\right)g_{\text{bar}}$$

Or *simplest* approximation for MOND modified gravity

Example inverse interpolation function: (Famaey and McGaugh '11)

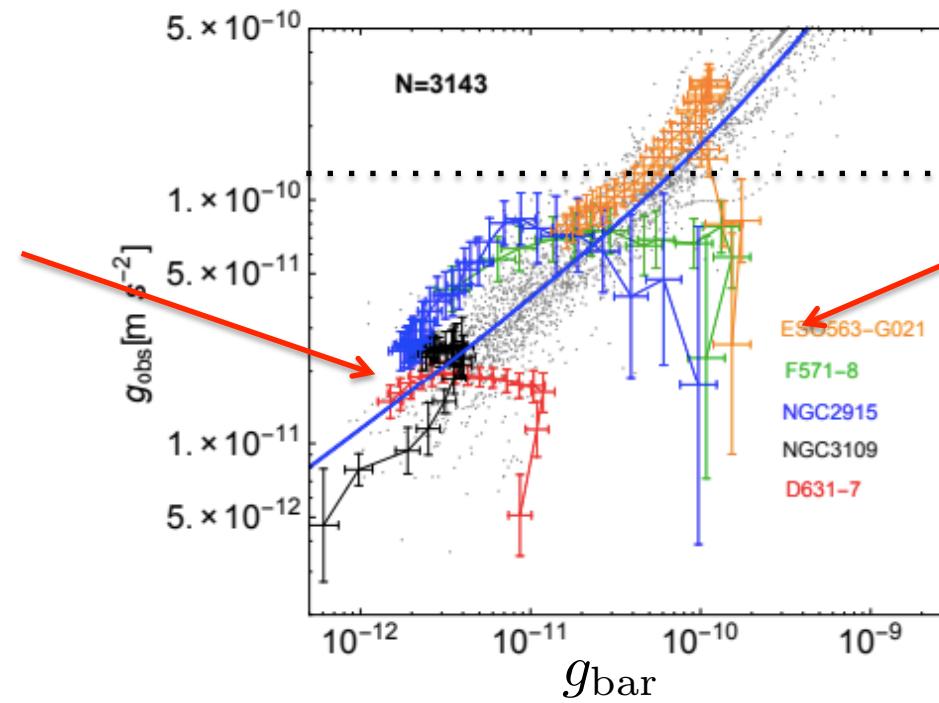
$$g_M = \frac{g_{\text{bar}}}{1 - e^{-\sqrt{\frac{g_{\text{bar}}}{g_0}}}}$$

$$\nu\left(\frac{g_{\text{bar}}}{g_0}\right) = \frac{1}{1 - e^{-\sqrt{\frac{g_{\text{bar}}}{g_0}}}}$$

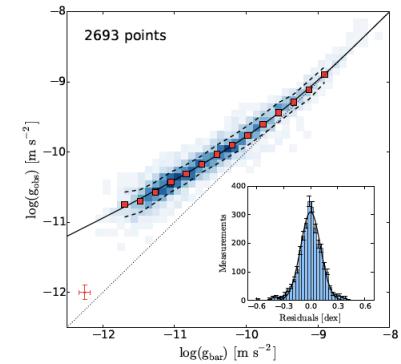
SPARC Individual galaxies highlighted

(J. Petersen and MTF '17
MTF and J. Petersen '18)

Large radii



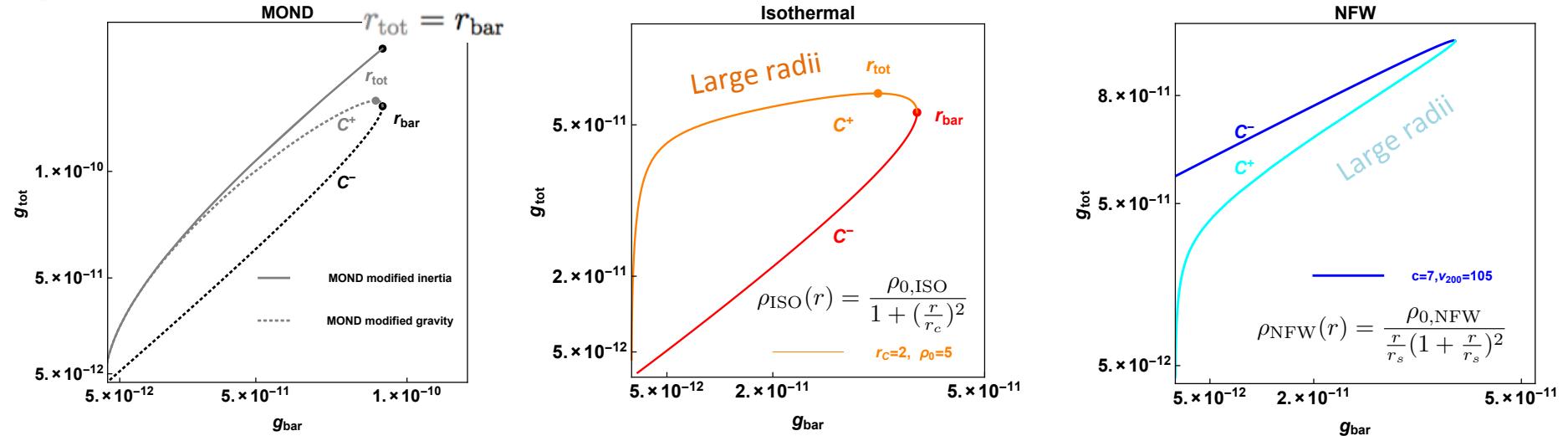
Small radii



Grey dots are all data points from SPARC in g2-space

Individual galaxies don't follow Radial Acceleration Relation at smaller radii

MOND and DM geometry in g2-space



Geometric Classification

Mond Modified Inertia

MOND Modified Gravity

DM Pseudo-Isothermal

DM Navarro-Frenk-White

Models	Reference radii	Curve segments	Curve Area ^a
MOND-MI	$r_{\text{tot}} = r_{\bar{r}}$	$\mathcal{C}^+ = \mathcal{C}^-$	$\mathcal{A}(\mathcal{C}) = 0$
MOND-MG	$r_{\text{tot}} > r_{\bar{r}}$	$\mathcal{C}^+ > \mathcal{C}^-$	$\mathcal{A}(\mathcal{C}) > 0$
DM-ISO	$r_{\text{tot}} > r_{\bar{r}}$	$\mathcal{C}^+ > \mathcal{C}^-$	$\mathcal{A}(\mathcal{C}) > 0$
DM-NFW	$r_{\text{tot}} < r_{\bar{r}}$	$\mathcal{C}^+ < \mathcal{C}^-$	Curves open

MOND models in \hat{g} -space

$$\hat{g}(r)_{\text{bar,tot}} = \frac{g(r)_{\text{bar,tot}}}{g(r_{\text{bar}})_{\text{bar,tot}}}$$

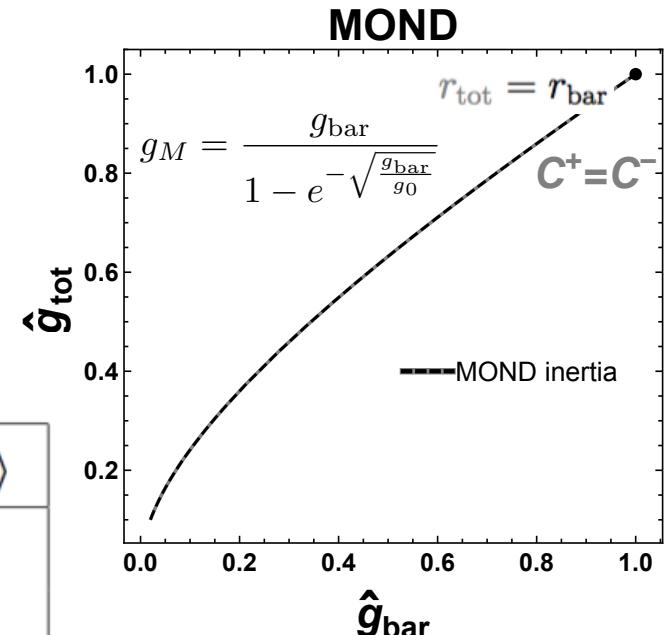
(MTF and J. Petersen '18)

MOND modified Inertia consequences:

$$r_{\text{tot}} = r_{\text{bar}} \text{ and } C^+ = C^- \text{ so } \hat{g}(r_{\text{tot}})_{\text{bar,tot}} = 1$$

SPARC data for $\hat{g}(r_{\text{tot}})_{\text{bar,tot}}$

Data selection	Points	$\langle \hat{g}_{\text{obs}} \pm \delta \hat{g}_{\text{obs}} \rangle$	$\langle \hat{g}_{\text{bar}} \pm \delta \hat{g}_{\text{bar}} \rangle$
$r_j = r_{\text{obs}}$	152	1.39 ± 0.12	0.83 ± 0.01
$r_j = r_{\text{obs}}, \frac{\delta v_{\text{obs}}}{v_{\text{obs}}} < 0.1$	146	1.12 ± 0.02	0.91 ± 0.01
$r_{j,\text{bar}} \in \Delta r_{\text{obs,bar}}$	~ 400	1.23 ± 0.04	0.89 ± 0.01



(For data points $g_{\text{tot}} = g_{\text{obs}}$, $r_{\text{tot}} = r_{\text{obs}}$)

MOND models in \hat{g}^2 -space

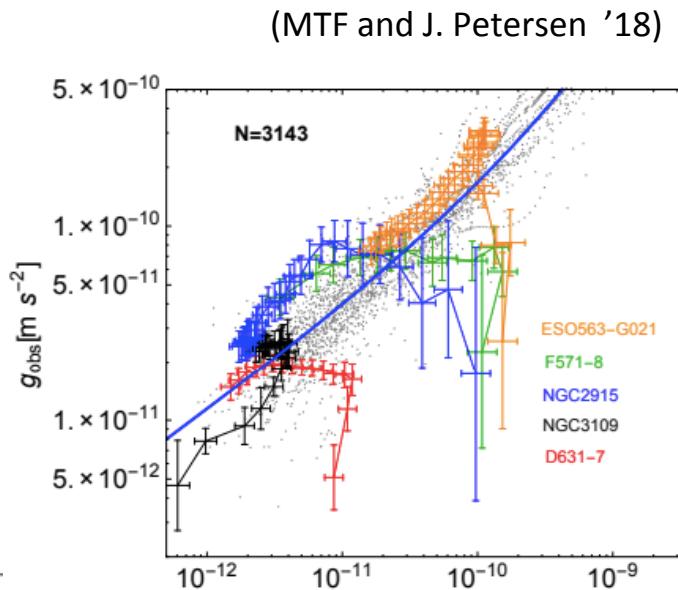
$$\hat{g}(r)_{\text{bar,tot}} = \frac{g(r)_{\text{bar,tot}}}{g(r_{\text{bar}})_{\text{bar,tot}}}$$

MOND modified Inertia consequences:

$$r_{\text{tot}} = r_{\text{bar}} \text{ and } C^+ = C^- \text{ so } \hat{g}(r_{\text{tot}})_{\text{bar,tot}} = 1$$

SPARC data for $\hat{g}(r_{\text{tot}})_{\text{bar,tot}}$

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(For data points $g_{\text{tot}} = g_{\text{obs}}$, $r_{\text{tot}} = r_{\text{obs}}$)

Particle DM

- Assumption: (All) the missing mass is particle DM
- Strategies:
 - **Generalize the SM baryonic relic (neutron)**
 - Generalize the SM neutrino relic
 - Something different

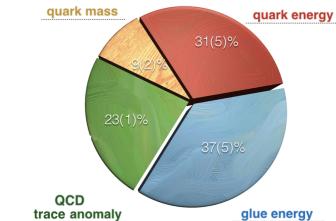
Baryonic relic density

- Proton stability (longevity) due to a **U(1) symmetry**

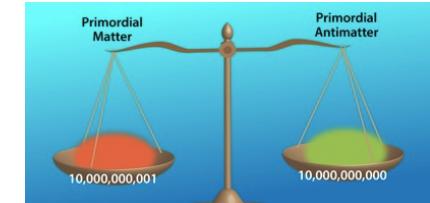
2 flavor massless QCD

$$SU(2)_L \times SU(2)_R \times U(1)_B \\ \rightarrow SU(2)_V \times U(1)_B$$

- Proton mass from **strong dynamics** (and Higgs)



- Proton relic density from some **asymmetry**



- **Neutron** lightest table baryon for zero current quark masses.
Self-interactions from strong dynamics

New Composite Dynamics

- 4d Gauge-Yukawa model with fermions and strong interactions
- May also (partially) break EW symmetry $\langle Q^I Q^J \rangle \sim f^3 E_Q^{IJ}$

$$\mathcal{L} = \mathcal{L}_{\text{SD}} + \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{SD+SM}}$$

CD breaks EW:
e.g TC, CH

$$\mathcal{O}_{\text{CD}} \sim QQ$$

$$\mathcal{L} = \mathcal{L}_{\text{SD}} + \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{SD+SM}}$$

CD induces EW breaking:
e.g. PCH,
 $\mathcal{O}_{\text{CD+SM}} \sim QHQ$
 $\mathcal{O}_{\text{CD+SM}} \sim QQH^\dagger H$

$$\mathcal{L} = \mathcal{L}_{\text{SD}} + \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{SD+SM}}$$

SM breaks EW:
e.g. SIDM
(Spergel & Steinhardt)

$$\mathcal{O}_{\text{SM}} \sim H^\dagger H$$

New Composite Dynamics

- 4d Gauge-Yukawa model with fermions and strong interactions
- May also (partially) break EW symmetry $\langle Q^I Q^J \rangle \sim f^3 E_Q^{IJ}$

$$\mathcal{L} = \mathcal{L}_{CD} + \mathcal{L}_{CD+SM} + \mathcal{L}_{SM}$$

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SM breaks EW:
e.g. SIDM
(Spergel & Steinhardt)

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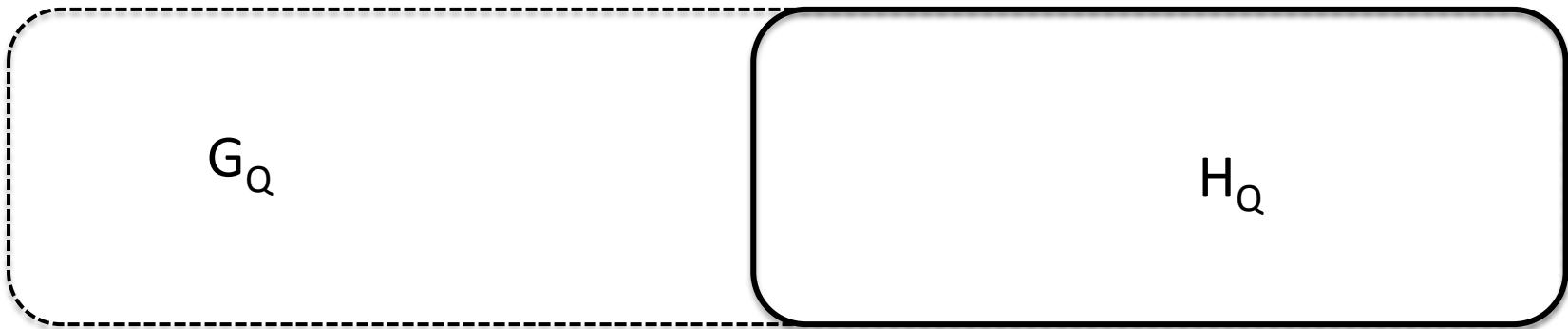
New Composite Dynamics

$$\text{UV: } \mathcal{L}_{UV} = \bar{Q}\gamma^\mu D_\mu Q + \mathcal{L}_{SM-Higgs} + \delta\mathcal{L}$$

$$\text{IR: } \langle Q^I Q^J \rangle \sim f^3 E_Q^{IJ}$$

$$G_Q \supset SU(2) \times SU(2) \times U(1) \quad \rightarrow \quad H_Q \supset SU(2) \times U(1)$$

SM custodial symmetry DM symmetry



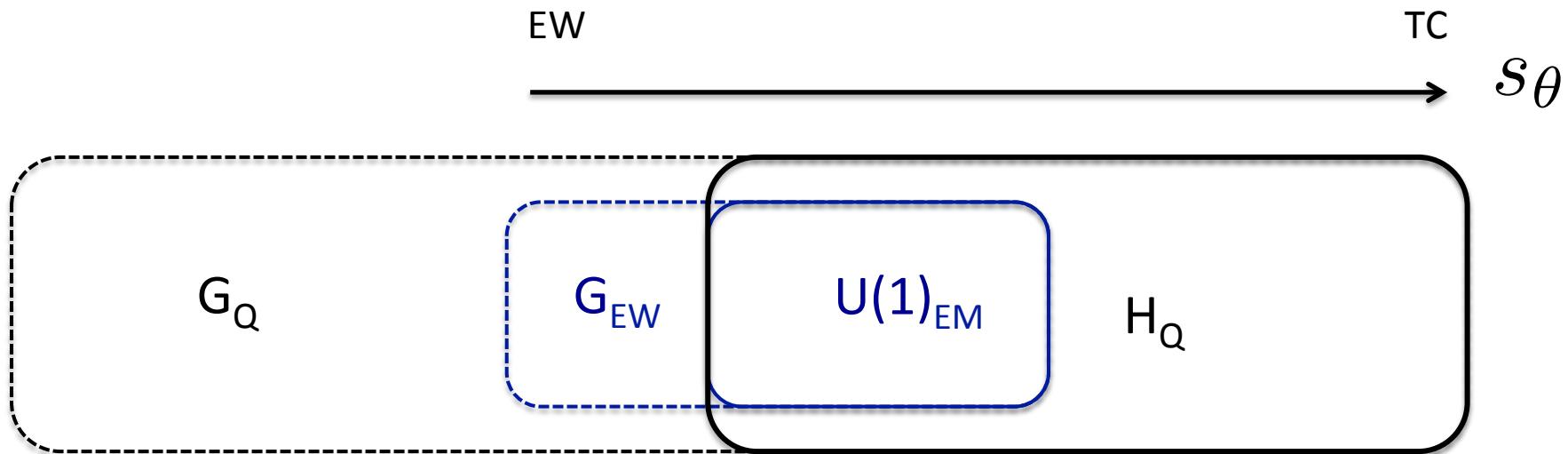
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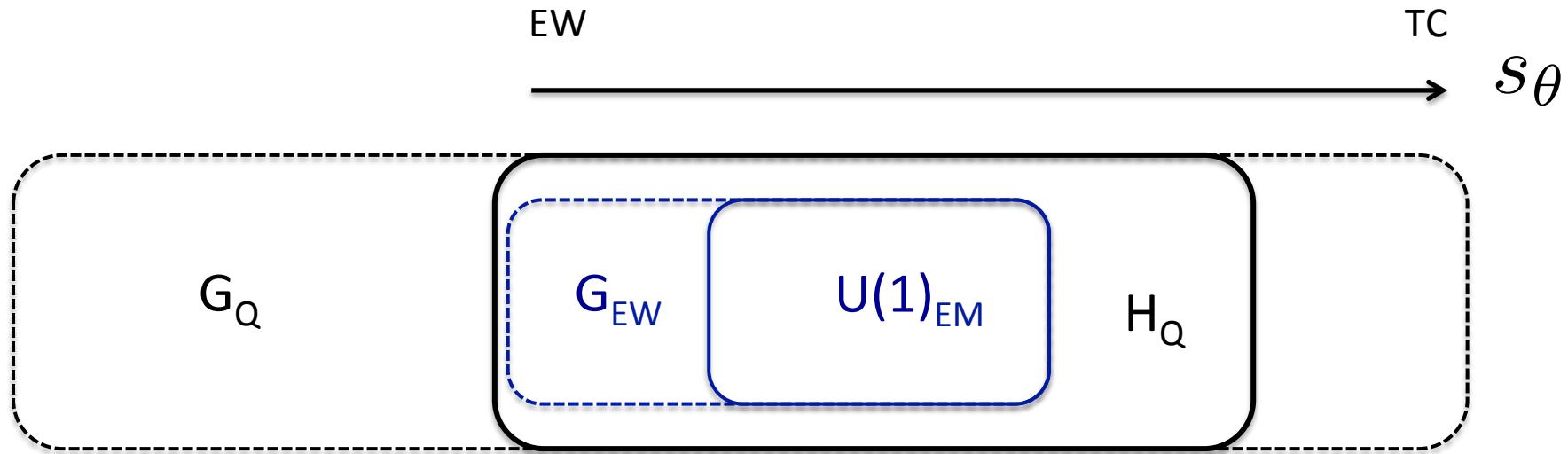
Framework: Composite Dynamics

$$\text{UV: } \mathcal{L}_{UV} = \bar{Q}\gamma^\mu D_\mu Q + \mathcal{L}_{SM-Higgs} + \delta\mathcal{L}$$

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SM custodial symmetry DM symmetry



Framework: Composite Dynamics

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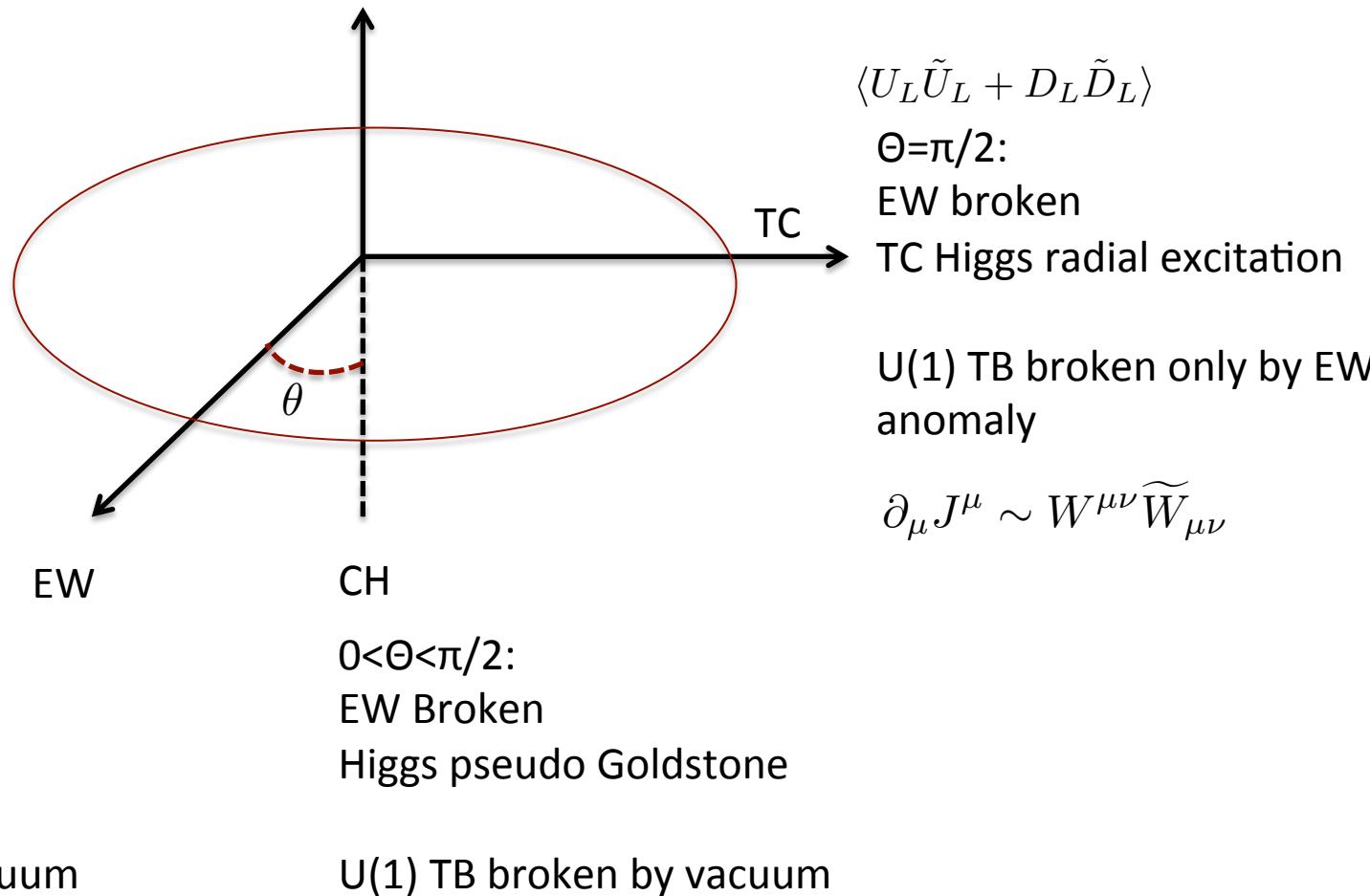
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SM custodial symmetry DM symmetry

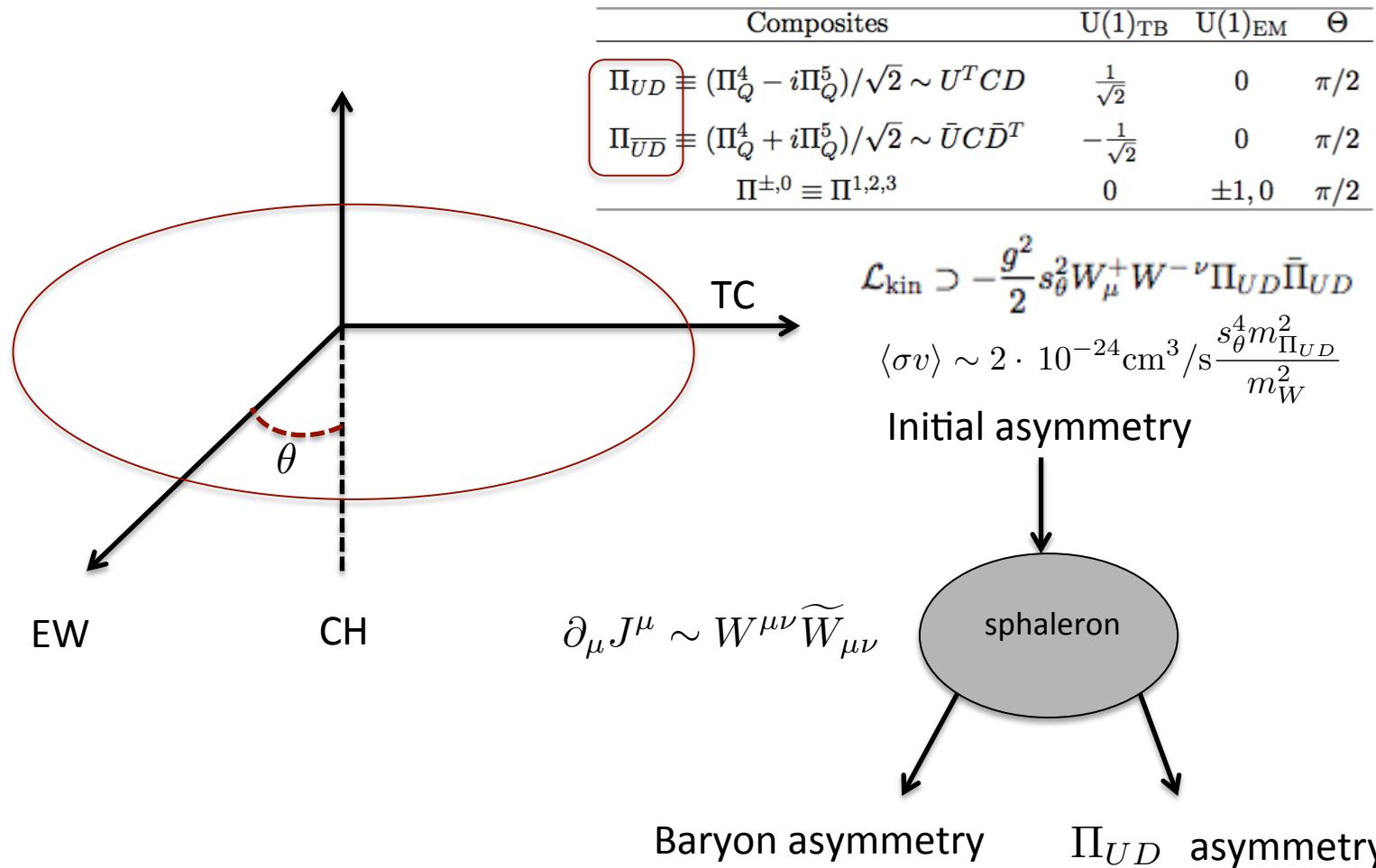
Example: $G_Q = SU(4)$ \rightarrow $H_Q = Sp(4)$

	$SU(2)_{TC}$	$SU(2)_W$	$U(1)_Y$	$Q = \begin{pmatrix} U_L \\ D_L \\ \tilde{U}_L \\ \tilde{D}_L \end{pmatrix}$
(U_L, D_L)	□	□	0	
\tilde{U}_L	□	1	-1/2	
\tilde{D}_L	□	1	+1/2	

Framework: Composite Dynamics



Framework: Composite Dynamics



Technicolor limit

- Technibaryon stable due to (anomalous) **$U(1)$ TB symmetry**
- mass from new **TC strong dynamics** (and Higgs)
- relic density from **sphaleron asymmetry transfer?**
- Self-interactions from TC strong dynamics
- **Technineutron** (pNGB) of some kind DM candidate

Technicolor limit

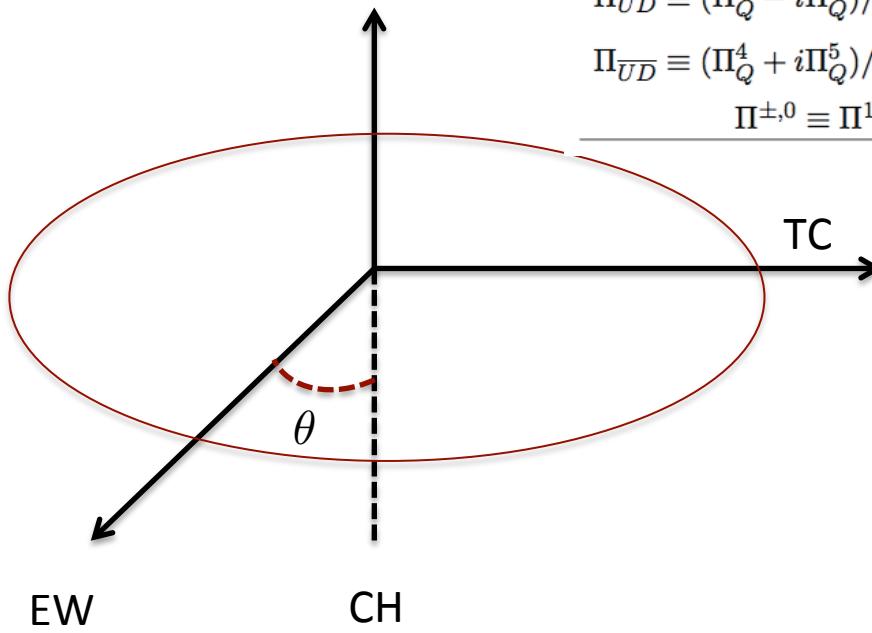
- Technibaryon stability due to (anomalous) **$U(1)$ TB symmetry**
- mass from new **TC strong dynamics** (and Higgs)
- relic density from **sphaleron asymmetry transfer?**
- Self-interactions from TC strong dynamics
- **Technineutron** of some kind DM candidate

Extremely compelling, apparently wrong limit

Observed Higgs properties not obviously consistent with a TC Higgs

Framework: Composite Dynamics

Composites	$U(1)_{TB}$	$U(1)_{EM}$	Θ
$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T CD$	$\frac{1}{\sqrt{2}}$	0	$\pi/2$
$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U} C \bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	$\pi/2$
$\Pi^{\pm,0} \equiv \Pi^{1,2,3}$	0	$\pm 1, 0$	$\pi/2$



Composites	$U(1)_{TB}$	$U(1)_{EM}$	Θ
$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	-	0	0
$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T CD$	-	0	0
$\Pi^{\pm,0} \equiv \Pi^{1,2,3}$	0	$\pm 1, 0$	0

No underlying stabilizing symmetry
for the CH range of parameters

(Galloway, Evans, Luty & Tacchi '10; Ferretti & Karateev '13;
Cacciapaglia & Sannino'14; Alanne, Buarque Franzosi & MTF '17)

Inert extensions

(Luty & Okui '04; Dietrich, Sannino & Tuominen '06;
Ryttov & Sannino '08; Luty '09)

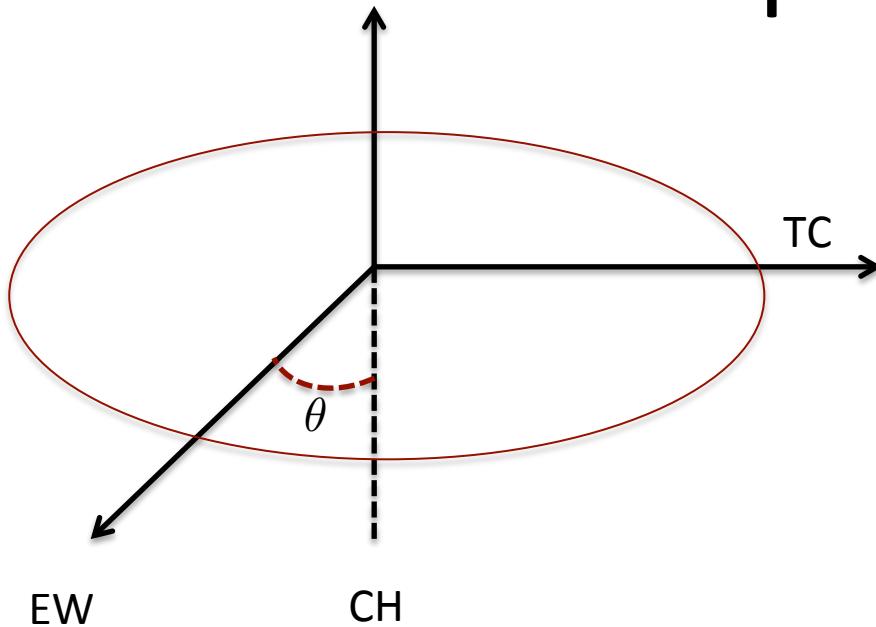
	SU(2) _{TC}	SU(2) _W	U(1) _Y
(U_L, D_L)	□	□	0
\tilde{U}_L	□	1	-1/2
\tilde{D}_L	□	1	+1/2
λ_L	Adj	1	0
$\tilde{\lambda}_L$	Adj	1	0

$$Q = \begin{pmatrix} U_L \\ D_L \\ \tilde{U}_L \\ \tilde{D}_L \end{pmatrix} \quad \text{SU}(4)$$
$$\Lambda = \begin{pmatrix} \lambda_L \\ \tilde{\lambda}_L \end{pmatrix} \quad \text{SU}(2)$$

(M.T.F, Sarkar Schmidt-Hoberg '11; Ryttov & Sannino '08;
Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

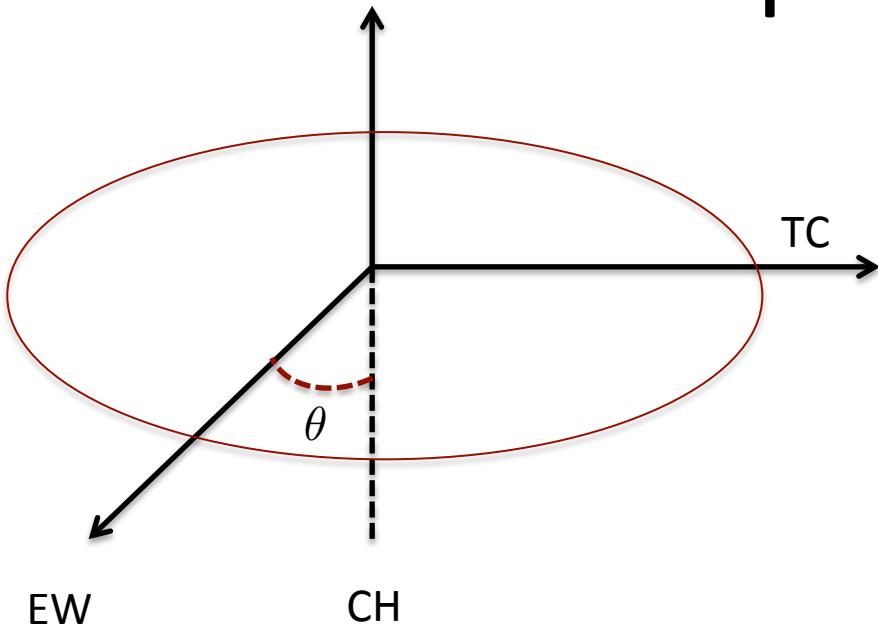
Global symmetry breaking pattern: $\text{SU}(4) \times \text{SU}(2) \times \text{U}(1) \rightarrow \text{Sp}(4) \times \text{U}(1)_{\text{DM}} \times \text{Z}_2$

Framework: Composite Dynamics



Composites	$U(1)_{\text{TB}}$	$U(1)_{\Lambda}$	Z_2
$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T C D$	$\frac{1}{\sqrt{2}}$	0	0
$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U} C \bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	0
$\Phi \sim \Lambda^T C \Lambda$	0	1	0
$\bar{\Phi} \sim \bar{\Lambda} C \bar{\Lambda}^T$	0	-1	0
$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	0	0	0

Framework: Composite Dynamics



Composites	$U(1)_{TB}$	$U(1)_\Lambda$	Z_2
$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T C D$	$\frac{1}{\sqrt{2}}$	0	0
$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U} C \bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	0
$\Phi \sim \Lambda^T C \Lambda$	0	1	0
$\bar{\Phi} \sim \bar{\Lambda} C \bar{\Lambda}^T$	0	-1	0
$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	0	0	0

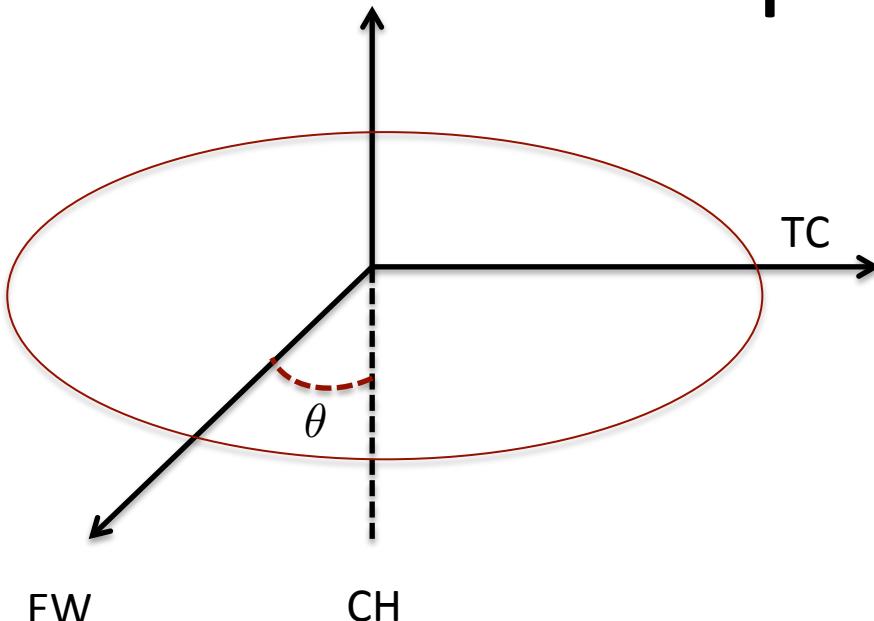
$$\mathcal{L}_{\text{kin}} = \frac{f^2}{8} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] + \frac{f_\Lambda^2}{8} \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda],$$

$$+ \frac{c_1}{4\pi} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda]$$

$$\langle \sigma v \rangle \simeq 2 \cdot 10^{-25} \text{cm}^3/\text{s} \frac{c_1^2 m_\Phi^4}{f_\Lambda^4} \frac{s_\theta^4 m_\Phi^2}{m_W^2}$$

'WIMP' from strong interactions

Framework: Composite Dynamics



Composites	$U(1)_{\text{TB}}$	$U(1)_{\Lambda}$	Z_2
$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T C D$	$\frac{1}{\sqrt{2}}$	0	0
$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U} C \bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	0
$\Phi \sim \Lambda^T C \Lambda$	0	1	0
$\bar{\Phi} \sim \bar{\Lambda} C \bar{\Lambda}^T$	0	-1	0
$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	0	0	0

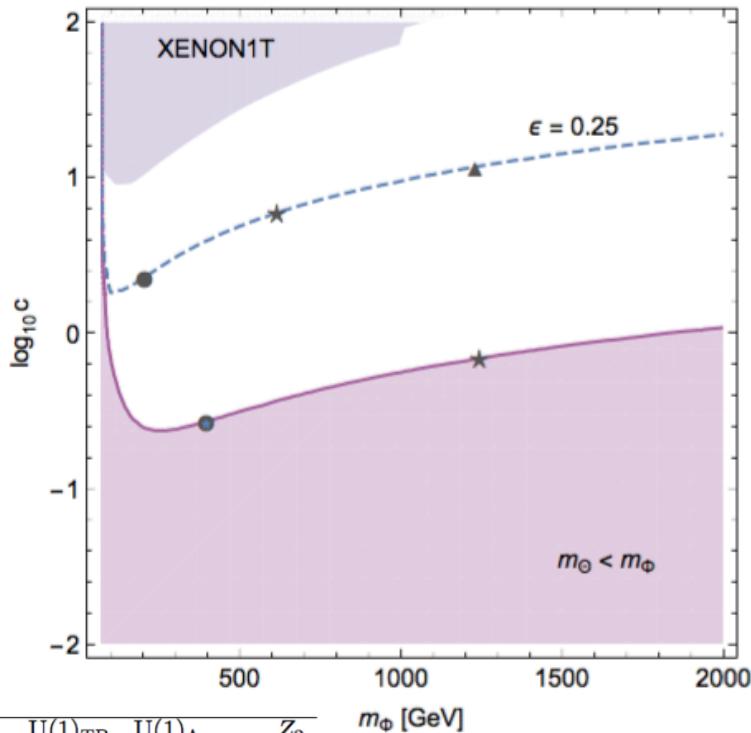
$$\mathcal{L}_{\text{kin}} = \frac{f^2}{8} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] + \frac{f_\Lambda^2}{8} \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda],$$

$$+ \frac{c_1}{4\pi} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda]$$

$$\langle \sigma v \rangle \simeq 2 \cdot 10^{-25} \text{cm}^3/\text{s} \frac{c_1^2 m_\Phi^4}{f_\Lambda^4} \frac{s_\theta^4 m_\Phi^2}{m_W^2}$$

'WIMP' from strong interactions in TC and CH limits

Constraints



$$\epsilon = m_\Phi / f$$

Increasing $f=f_\Lambda$

$$m_\Phi^2 = -16\pi c_\Lambda f_\Lambda m$$

$$m_\eta^2 = m_h^2 / s_\theta^2$$

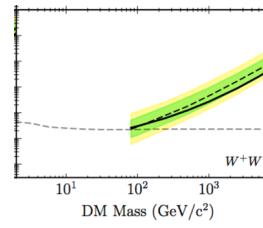
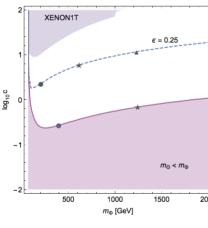
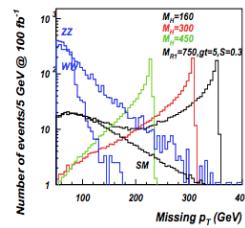
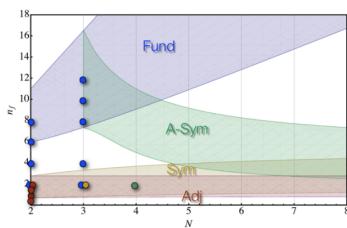
$$m_\Theta^2 = \frac{8}{9} m_\eta^2 c_\theta^2 + \frac{1}{9} m_\Phi^2$$

pNGB spectrum
From vacuum angle

Composites	$U(1)_{TB}$	$U(1)_\Lambda$	Z_2
$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	—	0	0
$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T C D$	—	0	0
$\Phi \sim \Lambda^T C \Lambda$	—	1	0
$\bar{\Phi} \sim \bar{\Lambda} C \bar{\Lambda}^T$	—	-1	0
$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	—	0	0

Summary

- Composite Dynamics compelling framework for DM and EW
U(1) stabilizing symmetries, dynamical symmetry breaking, naturalness, non-triviality, predictability (lattice)
- Vast space of models from 4d gauge-fermion-Yukawa theories
CH limit yields pNGB Higgs with properties tunably close to the SM Higgs.
Correlated with DM turning WIMPy (in studied model)
- Lattice and diverse experiments test underlying models



Model example

(Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

Composites	$U(1)_{TB}$	$U(1)_\Lambda$	Z_2
$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	–	0	0
$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T CD$	–	0	0
$\boxed{\Phi} \sim \Lambda^T C \Lambda$	–	1	0
$\bar{\Phi} \sim \bar{\Lambda} C \bar{\Lambda}^T$	–	–1	0
$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	–	0	0

EW breaking sector Inert strong sector

$$\mathcal{L}_{\text{kin}} = \frac{f^2}{8} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] + \frac{f_\Lambda^2}{8} \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda], \quad + \quad \frac{c_1}{4\pi} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda]$$

$$\langle \sigma v \rangle \simeq 2 \cdot 10^{-25} \text{cm}^3/\text{s} \frac{c_1^2 m_\Phi^4}{f_\Lambda^4} \frac{s_\theta^4 m_\Phi^2}{m_W^2}$$

annihilation cross-section 'WIMP' from strong interactions

Model example

(Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

	$U(1)_{TB}$	$U(1)_\Lambda$	Z_2	
DM candidate	$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	–	0	0
	$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T CD$	–	0	0
	$\Phi \equiv (\Pi_\Lambda^1 - i\Pi_\Lambda^2)/\sqrt{2} \sim \Lambda^T C\Lambda$	–	1	0
	$\bar{\Phi} \equiv (\Pi_\Lambda^1 + i\Pi_\Lambda^2)/\sqrt{2} \sim \bar{\Lambda}C\bar{\Lambda}^T$	–	–1	0
	$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	–	0	0
ADM candidate (Sannino & Ryttov '08)	$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T CD$	$\frac{1}{\sqrt{2}}$	0	0
	$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U}C\bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	0
	Φ	0	1	0
DM candidate (MTF, Sarkar & Schmidt-Hoberg '11)	$\bar{\Phi}$	0	–1	0
	Θ	0	0	0

Effective Lagrangian

$$\mathcal{L}_{\text{kin}} = \frac{f^2}{8} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] + \frac{f_\Lambda^2}{8} \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda],$$

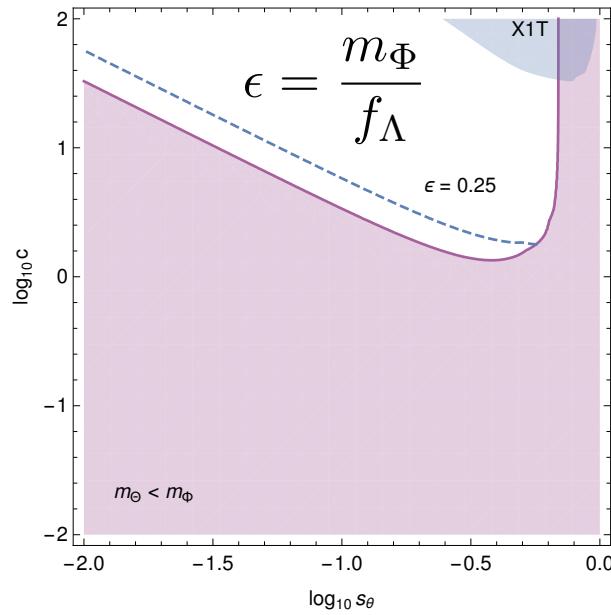
$$+ \frac{c_1}{4\pi} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda] + \dots$$

annihilation cross-section

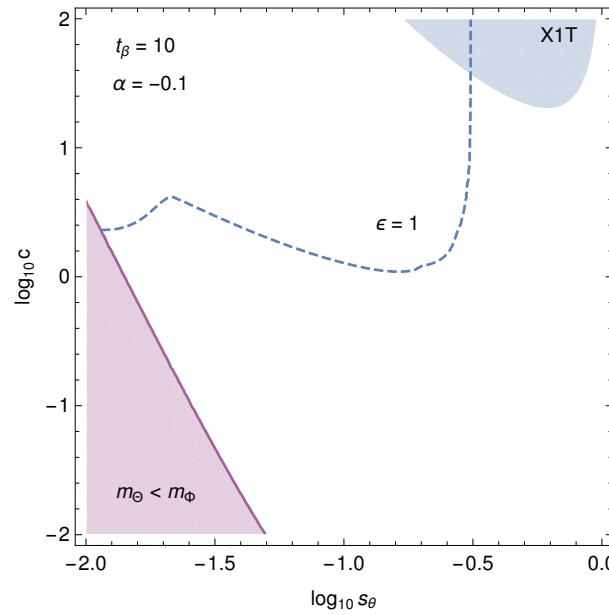
$$\langle \sigma v \rangle \simeq 2 \cdot 10^{-25} \text{cm}^3/\text{s} \frac{c_1^2 m_\Phi^4}{f_\Lambda^4} \frac{s_\theta^4 m_\Phi^2}{m_W^2}$$

Thermal relic density & constraints

Fermion partial compositeness



Partially composite Higgs



Framework: Composite Dynamics

	$SU(2)_{TC}$	$SU(2)_W$	$U(1)_Y$
(U_L, D_L)	□	□	0
\tilde{U}_L	□	1	-1/2
\tilde{D}_L	□	1	+1/2

$$Q = \begin{pmatrix} U_L \\ D_L \\ \tilde{U}_L \\ \tilde{D}_L \end{pmatrix}$$

$U(1) : Q \rightarrow e^{-\alpha} Q$

Framework 2: Composite Higgs

- pNGB Higgs properties tunably close to the SM Higgs

Framework: Composite Dynamics

- 4d Gauge-Yukawa model with strongly interacting fermions
- Technicolor and Composite Higgs models simply different vacuum alignment limits
- Vacuum alignment determined by interactions external to the strong dynamics

Framework 2: Composite Higgs

- Technibaryon stability (longevity) due to a U(1) symmetry ?
 - Technibaryon mass from new strong dynamics (and Higgs)
 - Technibaryon relic density origin?
 - Self-interactions from new strong dynamics
-
- pNGB Higgs properties tunably close to the SM Higgs

Baryonic relic density

- Proton stability (longevity) due to a U(1) symmetry
- Proton mass from strong dynamics (and Higgs)
- Proton relic density from some asymmetry
- Proton self-interactions from strong dynamics

Composite Higgs

- Technibaryon stability (longevity) due to a U(1) symmetry ?
 - Technibaryon mass from new strong dynamics (and Higgs)
 - Technibaryon relic density origin?
 - Self-interactions from new strong dynamics
-
- pNGB Higgs properties tunably close to the SM Higgs

Model example

(Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

	$SU(2)_{TC}$	$SU(2)_W$	$U(1)_Y$	
$SU(4)$	(U_L, D_L)	□	□	0
	\tilde{U}_L	□	1	-1/2
	\tilde{D}_L	□	1	+1/2
$SU(2)$	λ_L	Adj	1	0
	$\tilde{\lambda}_L$	Adj	1	0

(TC: Appelquist, Da Silva, Sannino & Duan '99, '00
CH: Galloway, Evans, Luty & Tacchi '10;
Ferretti & Karateev '13;
Cacciapaglia & Sannino '14)
(TC: Sannino & Ryttov '08; MTF, Sarkar & Schmidt-Hoberg '11)

Global symmetry breaking pattern: $SU(4) \times SU(2) \times U(1) \rightarrow Sp(4) \times U(1)_{DM} \times Z_2$
 Technicolor limit:

$$\langle QQ \rangle \sim f^3 E_Q, \langle \Lambda \Lambda \rangle \sim f_\Lambda^3 E_\Lambda$$

Model example

(Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

	U(1) _{TB}	U(1) _{Λ}	Z_2
$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	–	0	0
$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T C D$	–	0	0
$\Phi \equiv (\Pi_\Lambda^1 - i\Pi_\Lambda^2)/\sqrt{2} \sim \Lambda^T C \Lambda$	–	1	0
$\bar{\Phi} \equiv (\Pi_\Lambda^1 + i\Pi_\Lambda^2)/\sqrt{2} \sim \bar{\Lambda} C \bar{\Lambda}^T$	–	-1	0
$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	–	0	0
<hr/>			
$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T C D$	$\frac{1}{\sqrt{2}}$	0	0
$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U} C \bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	0
Φ	0	1	0
$\bar{\Phi}$	0	-1	0
Θ	0	0	0

EW unbroken $s_\theta=0$

TC limit $c_\theta=0$

Vacuum alignment in Q-sector $\langle QQ \rangle \sim f^3 E_Q$, $\langle \Lambda \Lambda \rangle \sim f_\Lambda^3 E_\Lambda$ $E_Q = c_\theta E_Q^- + s_\theta E_Q^B$.

$$E_Q^\pm = \begin{pmatrix} i\sigma_2 & 0 \\ 0 & \pm i\sigma_2 \end{pmatrix}, \quad E_Q^B = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad E_\Lambda = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Model example

(Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

	$U(1)_{TB}$	$U(1)_\Lambda$	Z_2
DM candidate	$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	–	0
	$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T CD$	–	0
	$\Phi \equiv (\Pi_\Lambda^1 - i\Pi_\Lambda^2)/\sqrt{2} \sim \Lambda^T C\Lambda$	–	1
	$\bar{\Phi} \equiv (\Pi_\Lambda^1 + i\Pi_\Lambda^2)/\sqrt{2} \sim \bar{\Lambda}C\bar{\Lambda}^T$	–	–1
	$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	–	0
ADM candidate (Sannino & Ryttov '08)	$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T CD$	$\frac{1}{\sqrt{2}}$	0
	$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U}C\bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0
DM candidate (MTF, Sarkar & Schmidt-Hoberg '11)	Φ	0	1
	$\bar{\Phi}$	0	–1
	Θ	0	0

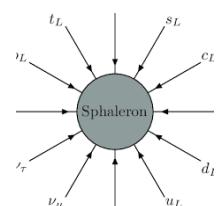
EW unbroken $s_\theta=0$

TC limit $c_\theta=0$

TC limit: $U(1)_{TB}$ current is preserved by condensates but EW anomalous \rightarrow ADM

$$\partial_\mu J_{TB}^\mu = \frac{1}{2\sqrt{2}} \frac{g^2}{32\pi^2} \epsilon_{\mu\nu\rho\sigma} W^{\mu\nu} W^{\rho\sigma}, \quad \text{and} \quad J_{TB}^\mu = \frac{1}{2\sqrt{2}} (\bar{U}\gamma^\mu U + \bar{D}\gamma^\mu D)$$

(TC: Sannino & Ryttov '08; MTF, Sarkar & Schmidt-Hoberg '11)



Model example

(Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

	$U(1)_{TB}$	$U(1)_\Lambda$	Z_2	
DM candidate	$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	–	0	0
	$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T CD$	–	0	0
	$\Phi \equiv (\Pi_\Lambda^1 - i\Pi_\Lambda^2)/\sqrt{2} \sim \Lambda^T C\Lambda$	–	1	0
	$\bar{\Phi} \equiv (\Pi_\Lambda^1 + i\Pi_\Lambda^2)/\sqrt{2} \sim \bar{\Lambda}C\bar{\Lambda}^T$	–	–1	0
	$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	–	0	0
ADM candidate (Sannino & Ryttov '08)	$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T CD$	$\frac{1}{\sqrt{2}}$	0	0
	$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U}C\bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	0
	Φ	0	1	0
DM candidate (MTF, Sarkar & Schmidt-Hoberg '11)	$\bar{\Phi}$	0	–1	0
	Θ	0	0	0

} EW unbroken $s_\theta=0$

} TC limit $c_\theta=0$

Effective Lagrangian

$$\mathcal{L}_{\text{kin}} = \frac{f^2}{8} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] + \frac{f_\Lambda^2}{8} \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda],$$

$$\Sigma_Q = \exp \left[2\sqrt{2}i \left(\frac{\Pi_Q}{f} - \frac{1}{3} \frac{\Theta}{f_\Theta} \mathbb{1}_4 \right) \right] E_Q$$

$$\Sigma_\Lambda = \exp \left[2\sqrt{2}i \left(\frac{\Pi_\Lambda}{f_\Lambda} + \frac{1}{6} \frac{\Theta}{f_\Theta} \mathbb{1}_2 \right) \right] E_\Lambda,$$

Model example

(Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

	$U(1)_{TB}$	$U(1)_\Lambda$	Z_2	
DM candidate	$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	–	0	0
	$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T CD$	–	0	0
	$\Phi \equiv (\Pi_\Lambda^1 - i\Pi_\Lambda^2)/\sqrt{2} \sim \Lambda^T C \Lambda$	–	1	0
	$\bar{\Phi} \equiv (\Pi_\Lambda^1 + i\Pi_\Lambda^2)/\sqrt{2} \sim \bar{\Lambda} C \bar{\Lambda}^T$	–	–1	0
	$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	–	0	0
ADM candidate (Sannino & Ryttov '08)	$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T CD$	$\frac{1}{\sqrt{2}}$	0	0
	$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U}C\bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	0
	Φ	0	1	0
DM candidate (MTF, Sarkar & Schmidt-Hoberg '11)	$\bar{\Phi}$	0	–1	0
	Θ	0	0	0

Effective Lagrangian

$$\mathcal{L}_{\text{kin}} = \frac{f^2}{8} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] + \frac{f_\Lambda^2}{8} \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda],$$

Π_{UD} annihilation cross-section

$$\mathcal{L}_{\text{kin}} \supset -\frac{g^2}{2} s_\theta^2 W_\mu^+ W^{-\nu} \Pi_{UD} \bar{\Pi}_{UD}$$

$$\langle \sigma v \rangle \simeq 2 \cdot 10^{-24} \text{cm}^3/\text{s} \frac{s_\theta^4 m_{\Pi_{UD}}^2}{m_W^2}$$

Model example

(Alanne, Buarque Franzosi, MTF & Rosenlyst '18)

	$U(1)_{TB}$	$U(1)_\Lambda$	Z_2	
DM candidate	$h \equiv \Pi_Q^4 \sim \bar{U}U + \bar{D}D$	–	0	0
	$\eta \equiv \Pi_Q^5 \sim \text{Im } U^T CD$	–	0	0
	$\Phi \equiv (\Pi_\Lambda^1 - i\Pi_\Lambda^2)/\sqrt{2} \sim \Lambda^T C \Lambda$	–	1	0
	$\bar{\Phi} \equiv (\Pi_\Lambda^1 + i\Pi_\Lambda^2)/\sqrt{2} \sim \bar{\Lambda} C \bar{\Lambda}^T$	–	–1	0
	$\Theta \sim i(\bar{U}\gamma^5 U + \bar{D}\gamma^5 D - (1/2)\bar{\Lambda}\gamma^5 \Lambda)$	–	0	0
ADM candidate (Sannino & Ryttov '08)	$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T CD$	$\frac{1}{\sqrt{2}}$	0	0
	$\Pi_{\bar{U}\bar{D}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U}C\bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	0
	Φ	0	1	0
DM candidate (MTF, Sarkar & Schmidt-Hoberg '11)	$\bar{\Phi}$	0	–1	0
	Θ	0	0	0

} EW unbroken $s_\theta=0$

} TC limit $c_\theta=0$

Effective Lagrangian

$$\mathcal{L}_{\text{kin}} = \frac{f^2}{8} \text{Tr}[D_\mu \Sigma_Q^\dagger D^\mu \Sigma_Q] + \frac{f_\Lambda^2}{8} \text{Tr}[\partial_\mu \Sigma_\Lambda^\dagger \partial^\mu \Sigma_\Lambda],$$

Correct size for EW scale mass and $s_\theta = 0.1$

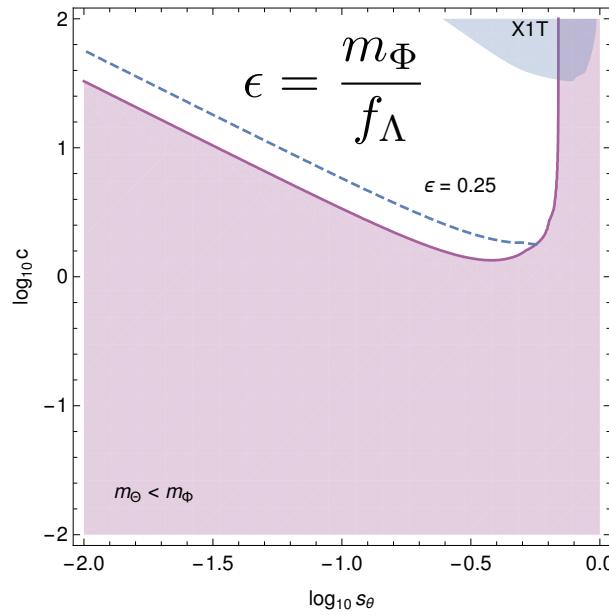
Π_{UD} annihilation cross-section

$$\mathcal{L}_{\text{kin}} \supset -\frac{g^2}{2} s_\theta^2 W_\mu^+ W^{-\nu} \Pi_{UD} \bar{\Pi}_{UD}$$

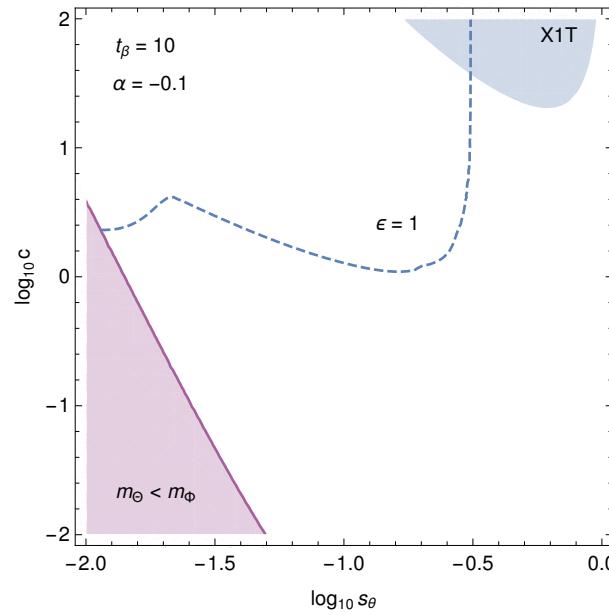
$$\langle \sigma v \rangle \simeq 2 \cdot 10^{-24} \text{cm}^3/\text{s} \frac{s_\theta^4 m_{\Pi_{UD}}^2}{m_W^2}$$

Thermal relic density & constraints

Fermion partial compositeness



Partially composite Higgs



Summary II

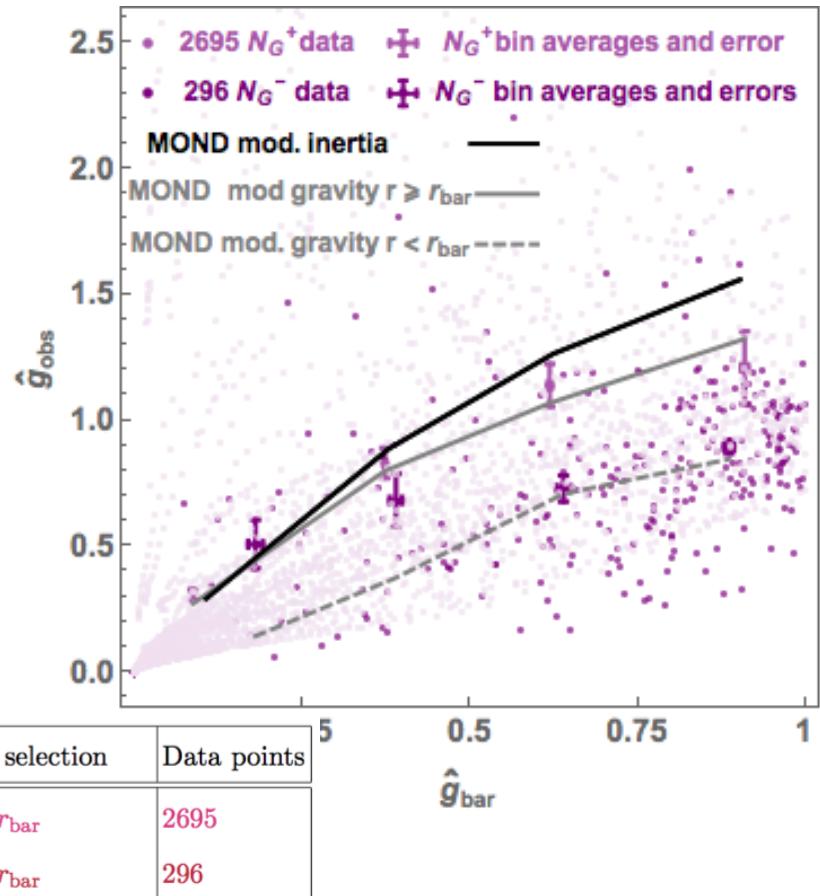
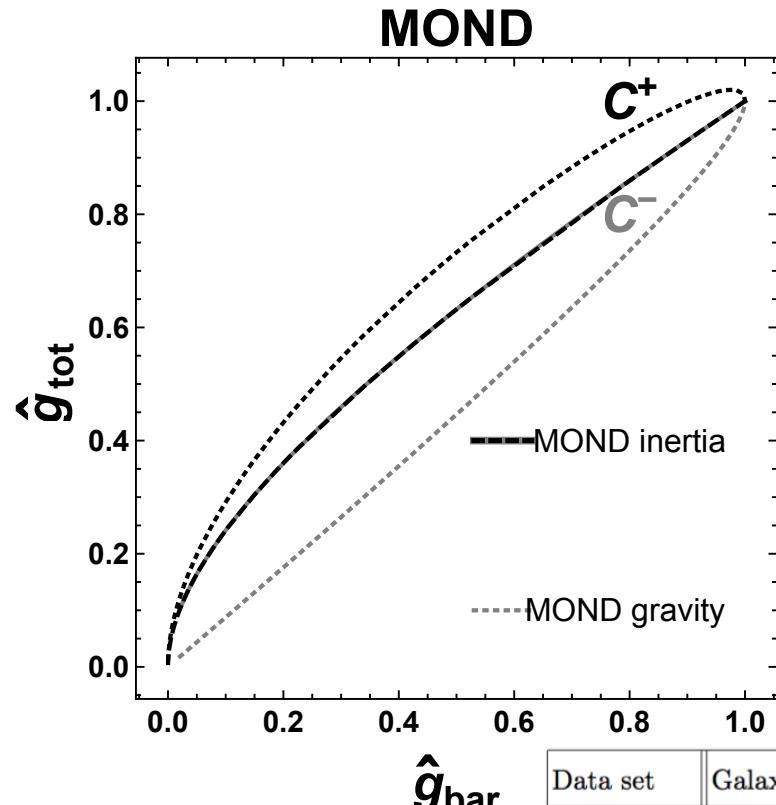
- Composite Dynamics well motivated framework for EWSB & DM
- Nature has been somewhat tough on TC so CH is next best thing
- Simple ADM scenario of TC is not easily transferred to CH
- But U(1) symmetric DM candidates exist

(Ma & Cacciapaglia '15;
Cai, Cacciapaglia, Zhang '18)

SPARC data in \hat{g}_2 -space

(MTF and J. Petersen '18)

All SPARC data after quality criteria organized wrt r_{bar}



Comparison to SPARC data

(MTF and J. Petersen '18)

(NB: For data $g_{\text{tot}} \rightarrow g_{\text{obs}}$)

$g_{\text{obs,bar}}$ in terms of measured quantities:

$$g_{\text{obs}}(r_j) = \frac{v_{\text{obs}}^2(r_j)}{r_j}, \quad g_{\text{bar}}(r_j) = \frac{(v_{\text{gas}}^2(r_j) + \Upsilon_{\text{disk}} v_{\text{disk}}^2(r_j) + \Upsilon_{\text{bul}} v_{\text{bul}}^2(r_j))}{r_j}$$

Comparison to SPARC data

(MTF and J. Petersen '18)

(NB: For data $g_{\text{tot}} \rightarrow g_{\text{obs}}$)

$g_{\text{obs,bar}}$ in terms of measured quantities:

$$g_{\text{obs}}(r_j) = \frac{v_{\text{obs}}^2(r_j)}{r_j}, \quad g_{\text{bar}}(r_j) = \frac{(v_{\text{gas}}^2(r_j) + \Upsilon_{\text{disk}} v_{\text{disk}}^2(r_j) + \Upsilon_{\text{bulge}} v_{\text{bulge}}^2(r_j))}{r_j}$$

g_{obs} uncertainties: random δv_{obs} , Inclination angle δi , Galaxy distance δD

$$\delta g_{\text{obs}}(r_j) = g_{\text{obs}}(r_j) \sqrt{\left[\frac{2\delta v_{\text{obs}}(r_j)}{v_{\text{obs}}(r_j)} \right]^2 + \left[\frac{2\delta i}{\tan(i)} \right]^2 + \left[\frac{\delta D}{D} \right]^2}$$

g_{bar} uncertainties: δv_{gas} , disk and bulge mass to light ratios $\delta \Upsilon_{\text{disk,bulge}}$

$$\delta g_{\text{bar}}(r_j) = \frac{\sqrt{(2v_{\text{gas}}(r_j))^2 \delta v_{\text{gas}}^2 + v_{\text{disk}}^4(r_j) \delta \Upsilon_{\text{disk}}^2 + v_{\text{bulge}}^4(r_j) \delta \Upsilon_{\text{bulge}}^2}}{r_j}$$

Comparison to SPARC data

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$g_{\text{obs,bar}}$ in terms of measured quantities:

$$g_{\text{obs}}(r_j) = \frac{v_{\text{obs}}^2(r_j)}{r_j}, \quad g_{\text{bar}}(r_j) = \frac{(v_{\text{gas}}^2(r_j) + \Upsilon_{\text{disk}} v_{\text{disk}}^2(r_j) + \Upsilon_{\text{bulge}} v_{\text{bulge}}^2(r_j))}{r_j}$$

g_{obs} uncertainties: random δv_{obs} , Inclination angle δi , Galaxy distance δD

$$\delta g_{\text{obs}}(r_j) = g_{\text{obs}}(r_j) \sqrt{\left[\frac{2\delta v_{\text{obs}}(r_j)}{v_{\text{obs}}(r_j)} \right]^2 + \left[\frac{2\delta i}{\tan(i)} \right]^2 + \left[\frac{\delta D}{D} \right]^2}$$

g_{bar} uncertainties: δv_{gas} , disk and bulge mass to light ratios $\delta \Upsilon_{\text{disk,bulge}}$

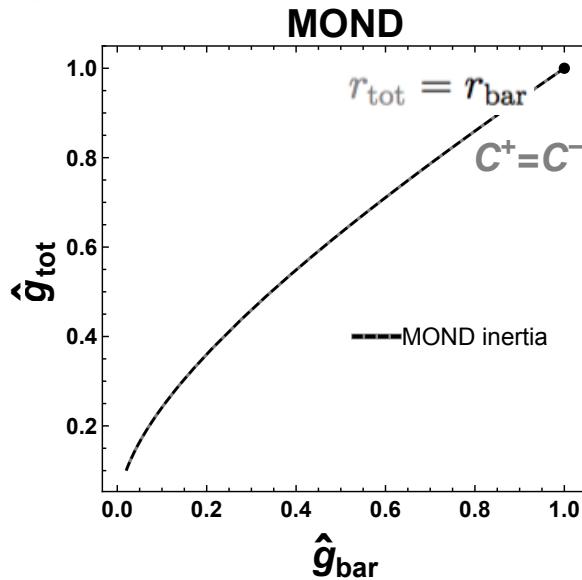
$$\delta g_{\text{bar}}(r_j) = \frac{\sqrt{(2v_{\text{gas}}(r_j))^2 \delta v_{\text{gas}}^2 + v_{\text{disk}}^4(r_j) \delta \Upsilon_{\text{disk}}^2 + v_{\text{bulge}}^4(r_j) \delta \Upsilon_{\text{bulge}}^2}}{r_j}$$

Galaxy specific uncertainties δi , δD , $\delta \Upsilon_{\text{disk,bulge}}$ dominate the RAR scatter

(Li, Lelli, McGaugh & Schombert '18)

Also these will depend on radius

MOND geometry in g2-space



Geometric Classification

MOND Modified Inertia

Models	Reference radii	Curve segments	Curve Area
MOND-MI	$r_{tot} = r_{bar}$	$\mathcal{C}^+ = \mathcal{C}^-$	$\mathcal{A}(\mathcal{C}) = 0$

MOND models in \hat{g} -space

(MTF and J. Petersen '18)

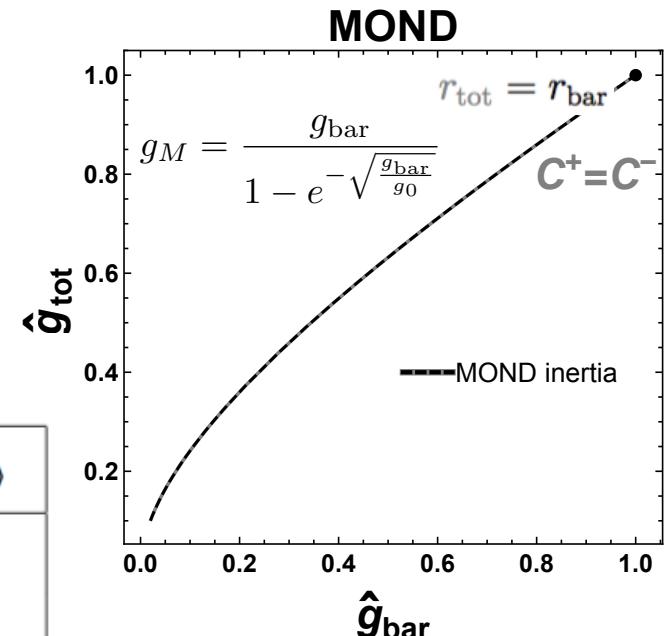
MOND modified Inertia consequences:

$$r_{\text{tot}} = r_{\text{bar}} \text{ and } C^+ = C^- \text{ so } \hat{g}(r_{\text{tot}})_{\text{bar,tot}} = 1$$

Ruled out independent of model $> 5\sigma$

SPARC data analysis for $\hat{g}(r_{\text{tot}})_{\text{bar,tot}}$

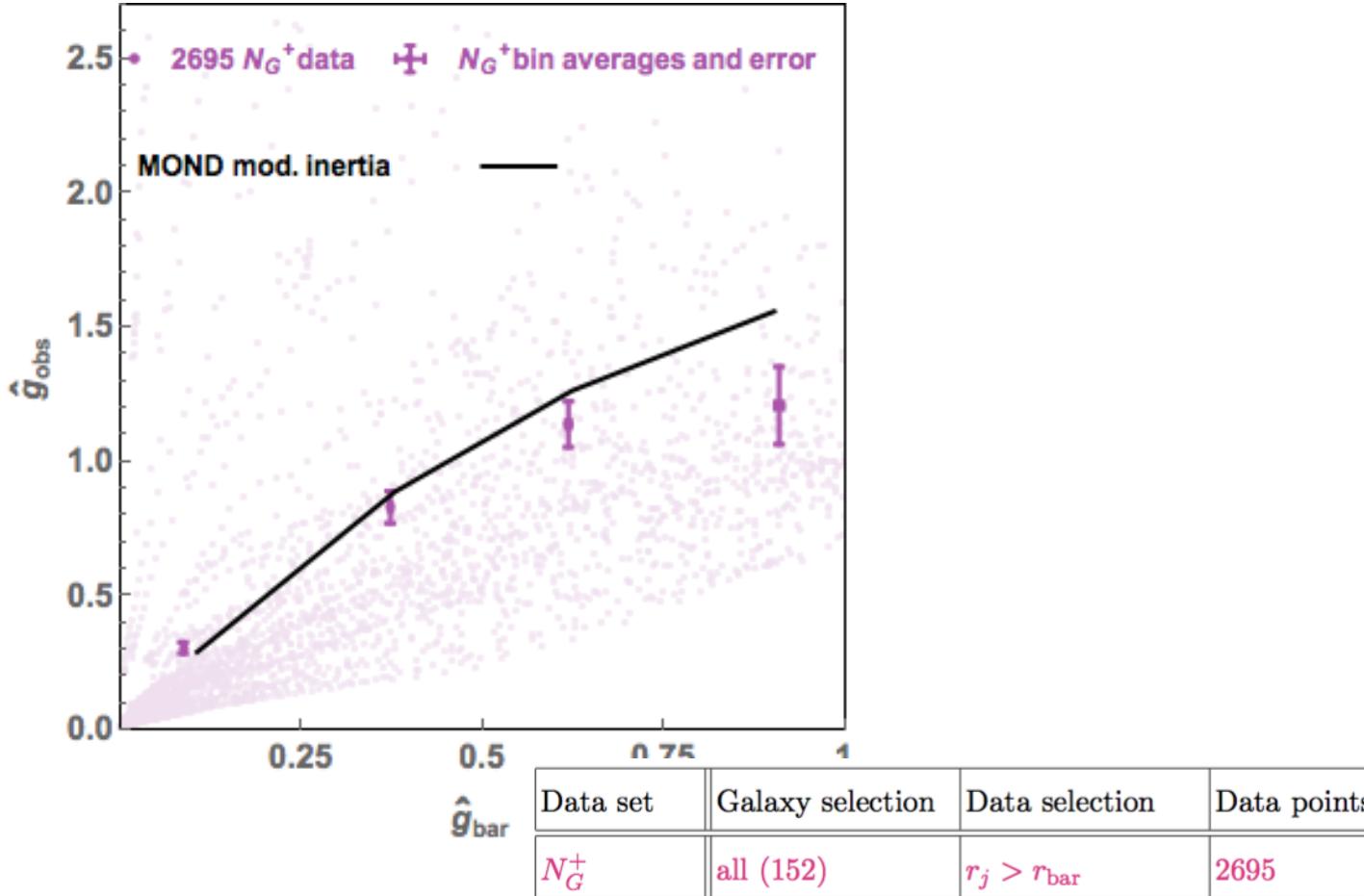
Data selection	points	$\langle \hat{g}_{\text{obs}} \pm \delta \hat{g}_{\text{obs}} \rangle$	$\langle \hat{g}_{\text{bar}} \pm \delta \hat{g}_{\text{bar}} \rangle$
$r_j = r_{\text{obs}}$	152	1.39 ± 0.12	0.83 ± 0.01
$r_j = r_{\text{obs}}, \frac{\delta v_{\text{obs}}}{v_{\text{obs}}} < 0.1$	146	1.12 ± 0.02	0.91 ± 0.01



(NB: For data points $g_{\text{tot}} \rightarrow g_{\text{obs}}$, $r_{\text{tot}} \rightarrow r_{\text{obs}}$)

SPARC data in \hat{g}^2 -space

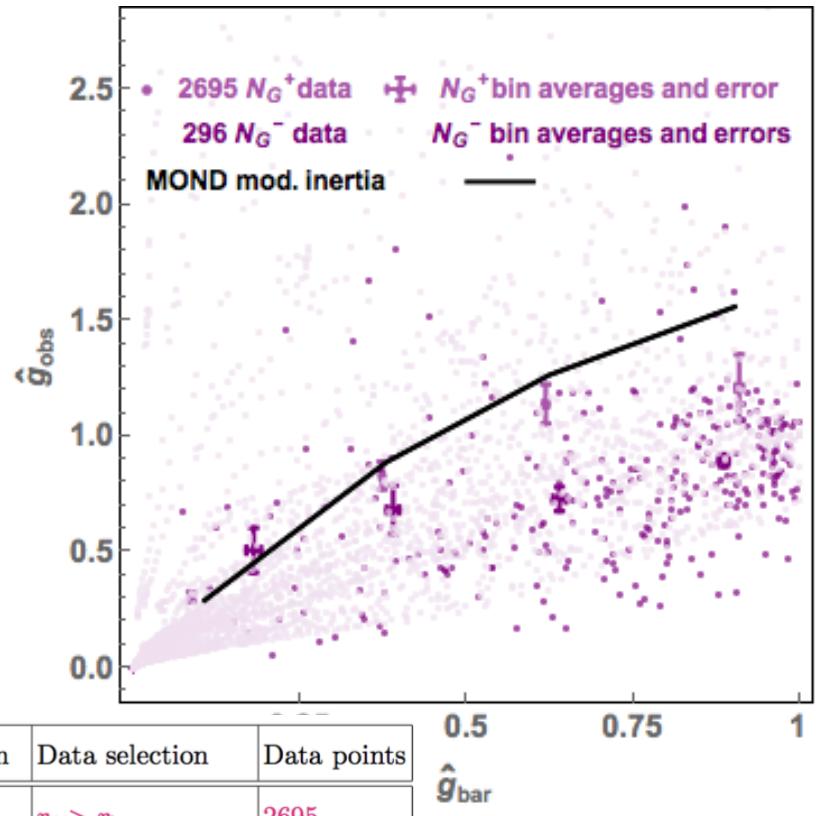
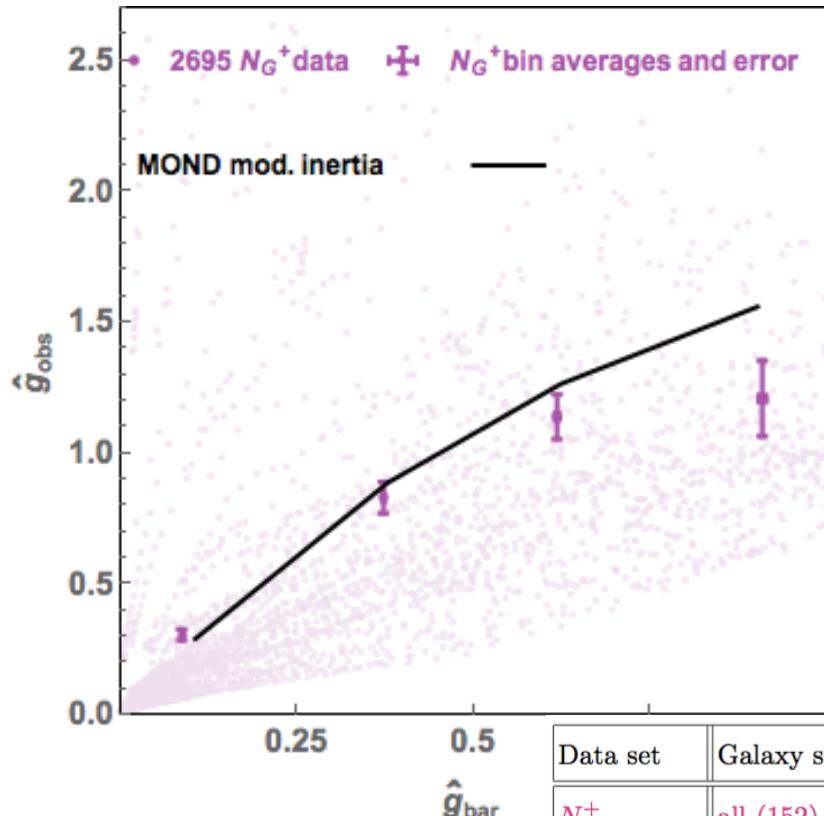
(MTF and J. Petersen '18)



SPARC data in \hat{g}_2 -space

(MTF and J. Petersen '18)

All SPARC data after quality criteria organized wrt r_{bar}



Data set	Galaxy selection	Data selection	Data points
N_G^+	all (152)	$r_j > r_{\text{bar}}$	2695
N_G^-	all (152)	$r_j \leq r_{\text{bar}}$	296

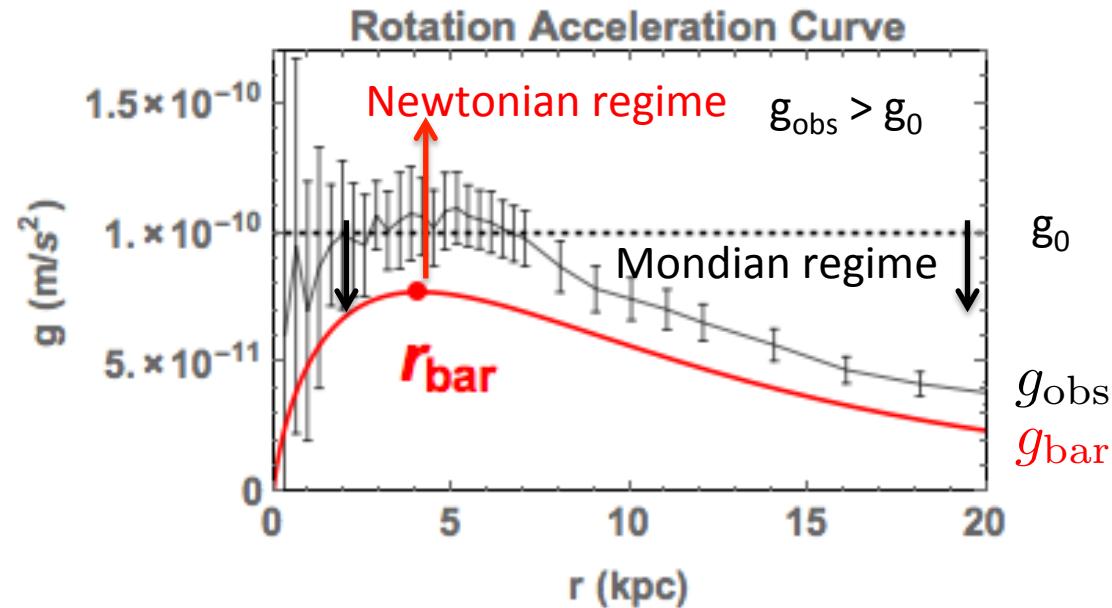
Rotation acceleration curves

Baryonic and total/observed accelerations

MONDian acceleration relation

$$\mu\left(\frac{g_M}{g_0}\right)g_M = g_{\text{bar}}$$

$$\mu(x) = \begin{cases} \mu(x) \simeq 1 & x \gg 1 \\ \mu(x) \simeq x & x \ll 1 \end{cases}$$



Newtonian regime $g_{\text{obs}} > g_0$

$$g_{\text{obs}} \sim g_{\text{bar}}$$

$g_0 \sim 10^{-10} \text{ m/s}^2$ is transition acc. scale

r_{bar} is radii of maximum baryonic acc.

Mondian regime $g_{\text{obs}} < g_0$

$$g_{\text{obs}}^2 \sim g_{\text{bar}}$$

Framework: Composite Dynamics

UV	$\mathcal{L}_{UV} = \bar{Q}\gamma^\mu D_\mu Q + \mathcal{L}_{SM-Higgs} + \delta\mathcal{L}$	$G_Q \in \underbrace{SU(2) \times SU(2)}_{\text{SM Higgs custodial symmetry}} \times \textcolor{red}{U(1)}$
IR	$\langle Q^I Q^J \rangle \sim f^3 E_Q^{IJ}$	$H_Q \in \underbrace{SU(2)}_{\text{}} \times \textcolor{red}{U(1)}$

Example

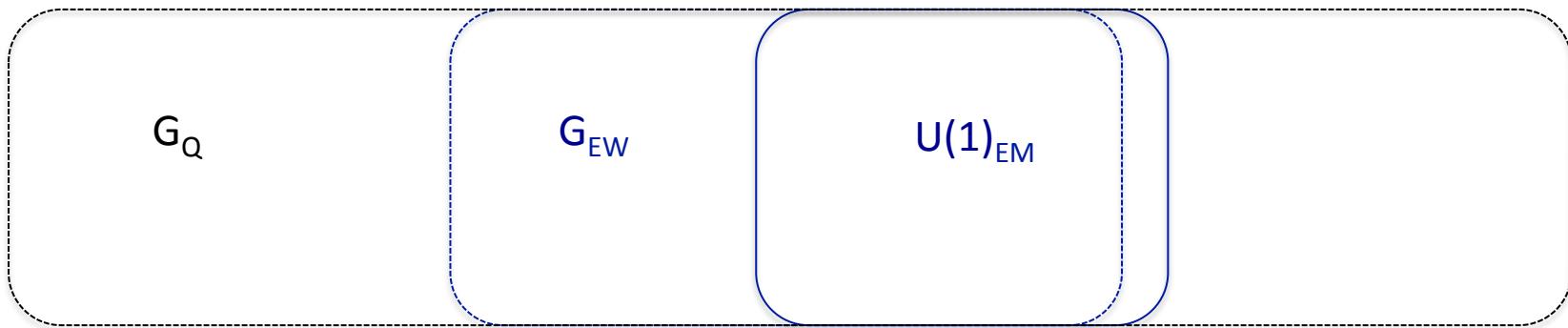
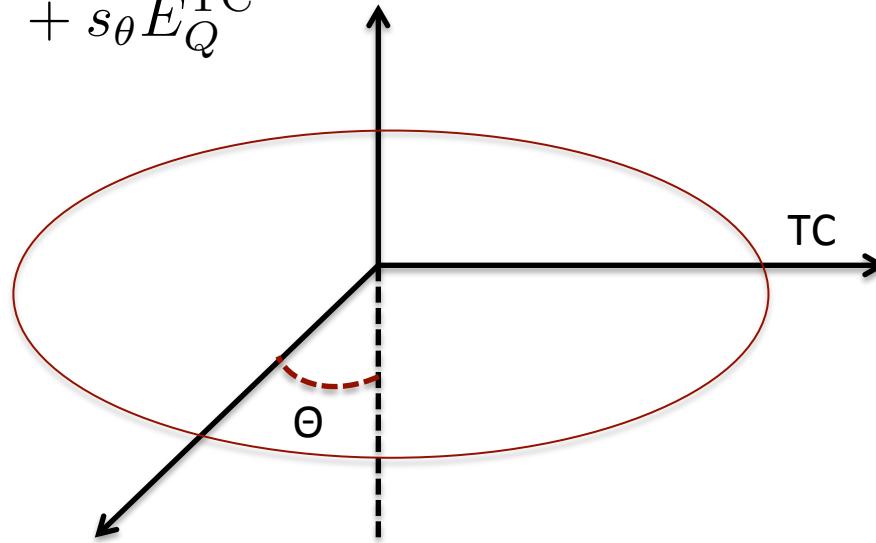
	SU(2) _{TC}	SU(2) _W	U(1) _Y
(U_L, D_L)	□	□	0
\tilde{U}_L	□	1	-1/2
\tilde{D}_L	□	1	+1/2

$$Q = \begin{pmatrix} U_L \\ D_L \\ \tilde{U}_L \\ \tilde{D}_L \end{pmatrix}$$

$$\textcolor{red}{U(1)} : Q \rightarrow e^{-\alpha} Q$$

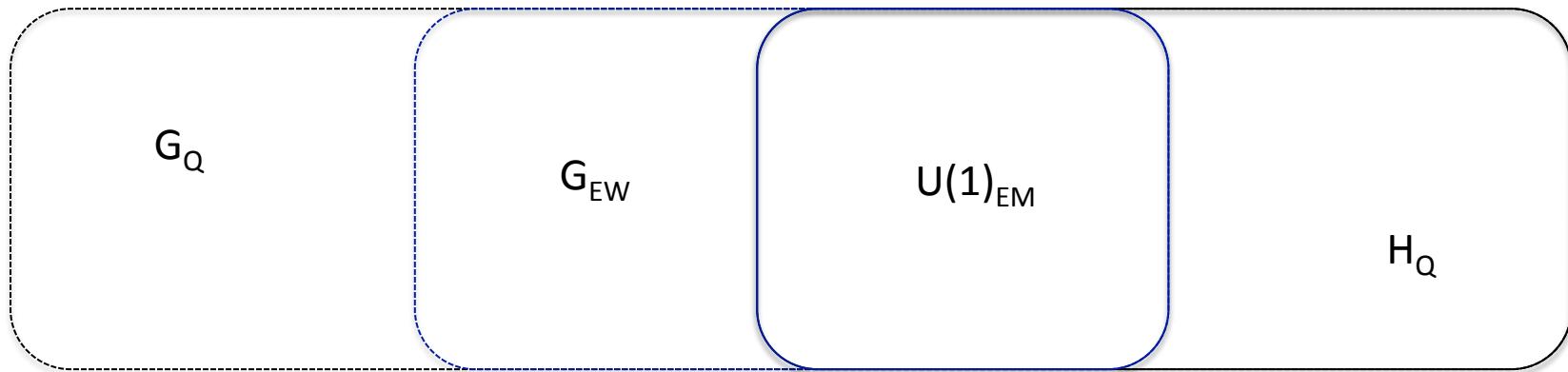
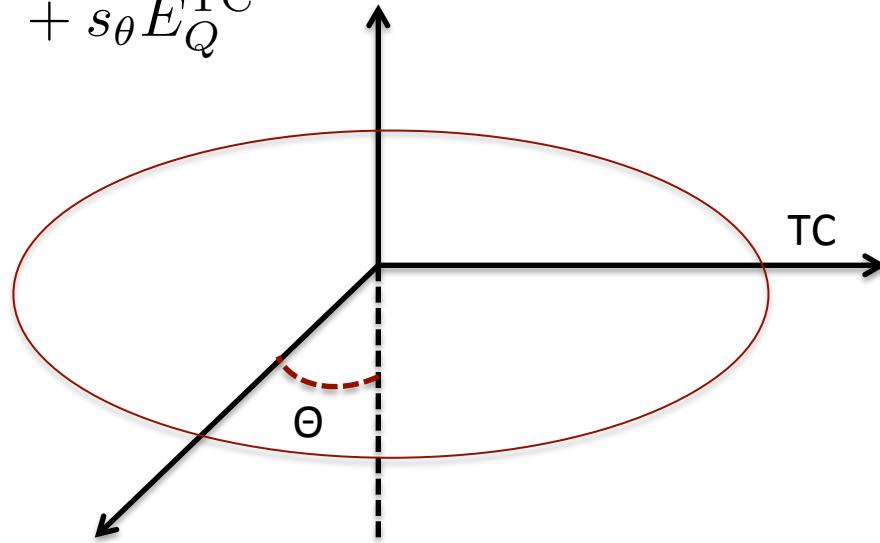
Framework: Composite Dynamics

$$E_Q = c_\theta E_Q^{\text{EW}} + s_\theta E_Q^{\text{TC}}$$



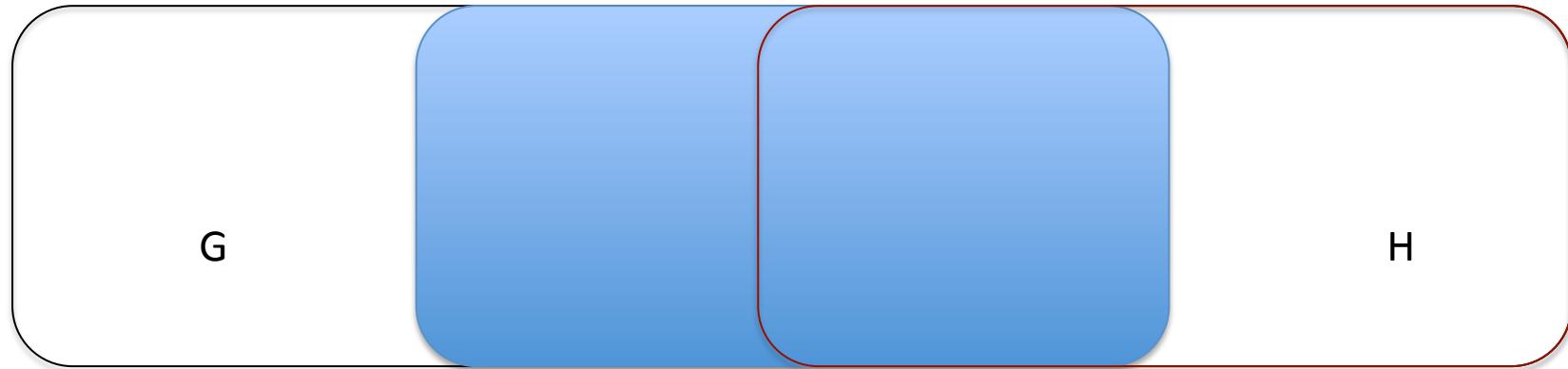
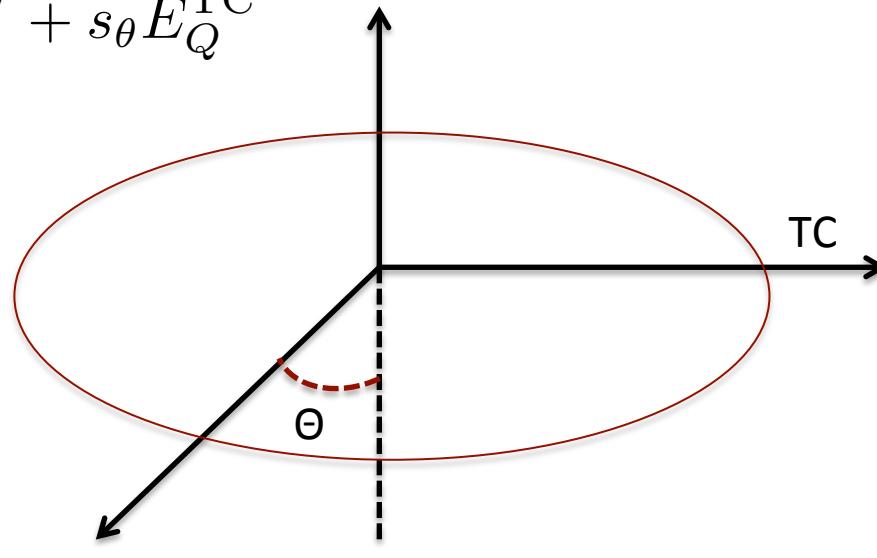
Framework: Composite Dynamics

$$E_Q = c_\theta E_Q^{\text{EW}} + s_\theta E_Q^{\text{TC}}$$



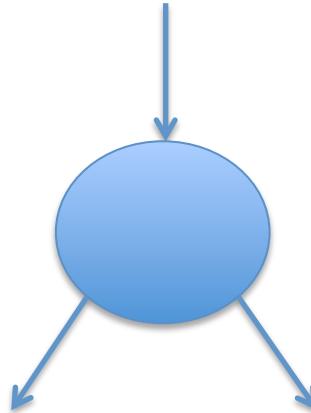
Framework: Composite Dynamics

$$E_Q = c_\theta E_Q^{\text{EW}} + s_\theta E_Q^{\text{TC}}$$



Framework: Composite Dynamics

Initial asymmetry

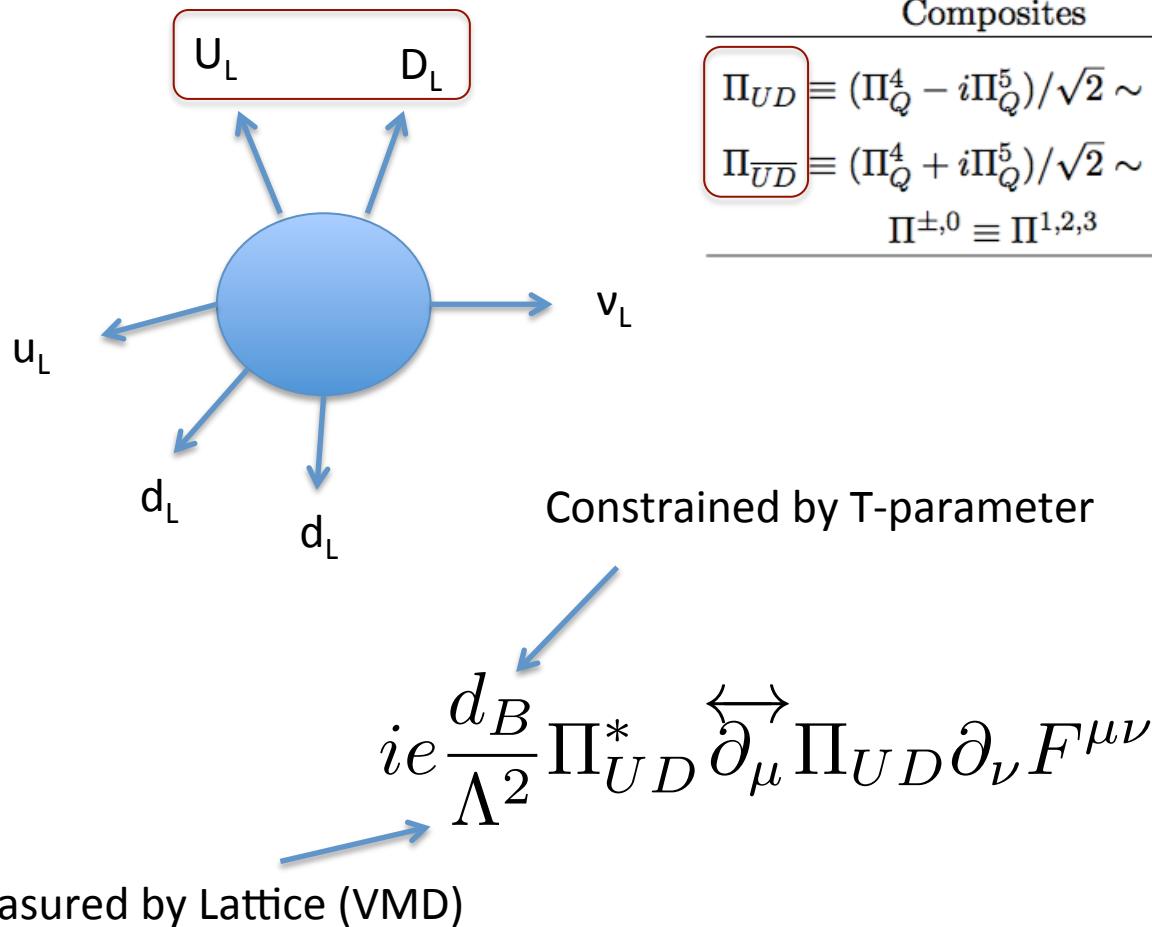


Baryon asymmetry

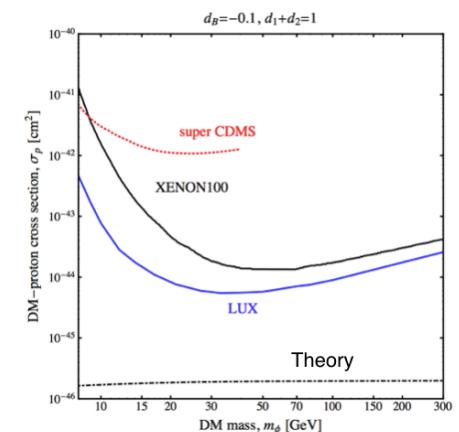
ADM

Composites	$U(1)_{\text{TB}}$	$U(1)_{\text{EM}}$	Θ
$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T CD$	$\frac{1}{\sqrt{2}}$	0	$\pi/2$
$\Pi_{\overline{UD}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U} C \bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	$\pi/2$
$\Pi^{\pm,0} \equiv \Pi^{1,2,3}$	0	$\pm 1, 0$	$\pi/2$

Framework: Composite Dynamics



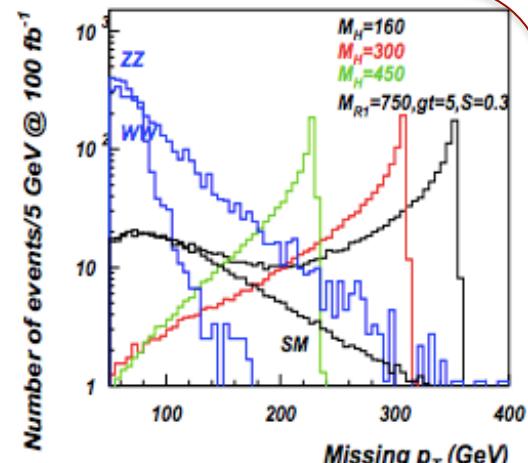
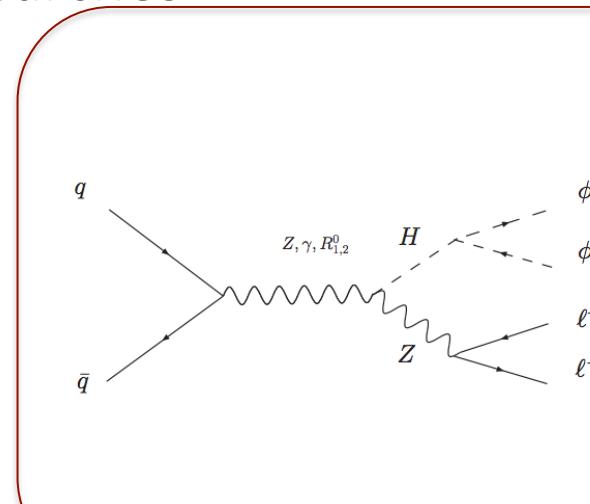
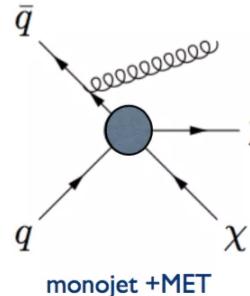
Composites	$U(1)_{\text{TB}}$	$U(1)_{\text{EM}}$	Θ
$\Pi_{UD} \equiv (\Pi_Q^4 - i\Pi_Q^5)/\sqrt{2} \sim U^T C D$	$\frac{1}{\sqrt{2}}$	0	$\pi/2$
$\Pi_{\overline{UD}} \equiv (\Pi_Q^4 + i\Pi_Q^5)/\sqrt{2} \sim \bar{U} C \bar{D}^T$	$-\frac{1}{\sqrt{2}}$	0	$\pi/2$
$\Pi^{\pm,0} \equiv \Pi^{1,2,3}$	0	$\pm 1, 0$	$\pi/2$



Composite DM Possible Features

- Resonances in collider searches ,
e.g missing E_T searches

(U. Haisch talk ZPW19)

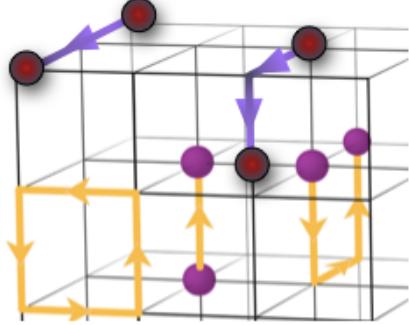


(Foadi, MTF, Sannino '08)

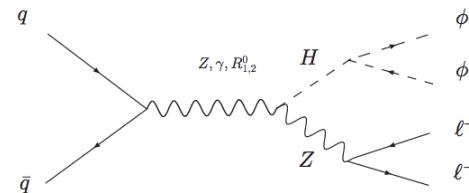
Composite DM Features

- lattice predictions and low mass DM,
e.g. for direct detection, LHC (F. Petricca talk ZPW19)

SU(2) MWT model
 $m_{R1}=2.5$ (0.5) TeV,
 $\Lambda=m_{R1}$, $d_B=(m_U-m_D)/(m_U+m_D)$



(Arthur, Drach, Hansen, Hietanen, Lewis, Pica & Sannino '14)



$$ie \frac{d_B}{\Lambda^2} \Pi_{UD}^* \partial_\mu \Pi_{UD} \partial_\nu F^{\mu\nu}$$

