

# Field electron emission in an external magnetic field parallel to the surface

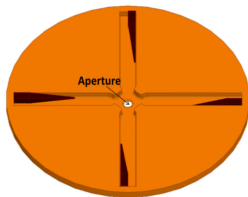
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The emission current "switching-off" by means of  $\vec{B} \perp \vec{E}$  is a possible way of preventing breakdowns.

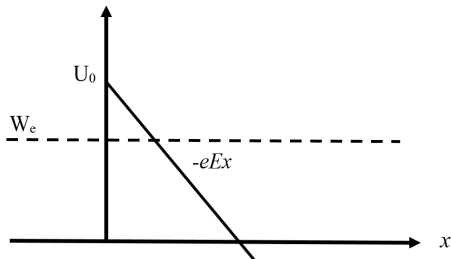


Parameters for the Main Linac RF structures  
[\[The CLIC Project Implementation Plan, Geneva, 2018\]](#)

	3 TeVstage
Number of structures	143,232
Active structure length [mm]	230
Aperture diameter [mm]	6.3-4.7
Accelerating voltage loaded [MV]	100

- The distance which an electron can cover during the elementary act of the acceleration is  $\sim 6 \text{ cm} \gg \text{Aperture}$   $\rightarrow$  the electric field is constant during the electron motion in the aperture.
- The electron Lorentz factor is  $\sim 1 \rightarrow$  it is possible neglects the electron's spin.

- Using the Klein–Gordon equation instead of the Schrödinger equation for wave functions of an electron motion at relativistic speed.
- The Klein–Gordon equation allows consideration of the external uniform magnetic field parallel to the metal surface in the case  $cB < E$ .



$$\left[ \frac{\partial^2}{\partial x^2} + \frac{x^2}{4} - 1 \right] \psi(x) = 0, \quad (1)$$

where  $x$  is a dimensionless coordinate,  
 $D_\nu(x)$  is a parabolic cylinder function.

*Fig. 1. The potential barrier in presence of an electric field.*

This equation is a parabolic cylinder equation. The solution is

$$\psi(x) = D_{-\frac{1}{2}-i} \left( x e^{-\frac{\pi i}{4}} \right).$$

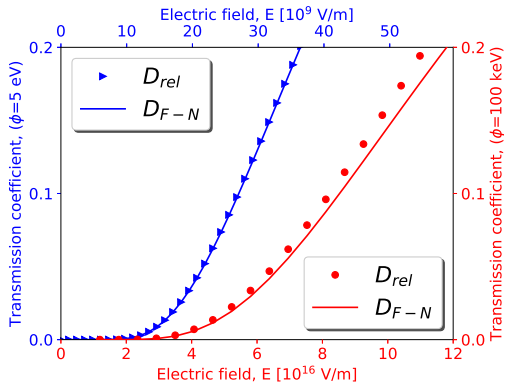


Fig. 2. Fowler-Nordheim  $D_{F-N}$  and relativistic generalized  $D_{rel}$  transmission coefficients for different work function values. A typical values for laboratory conditions  $E \cong 10^9 \frac{\text{V}}{\text{m}}$ ,  $\phi \cong 5\text{eV}$  are in blue color and for the field emission from neutron stars surface  $E \cong 10^{16} \frac{\text{V}}{\text{m}}$ ,  $\phi \cong 100\text{keV}$  are in red color.

$$\begin{aligned}
 D_{rel} = D_{F-N} & \left( 1 + \frac{\sqrt{2}}{5} \frac{(U_0 - W_e)^{\frac{5}{2}}}{\sqrt{mc^2} e E \hbar} + \frac{\sqrt{2}}{48} \frac{(7U_0 - 12W_e) e E \hbar}{U_0 \sqrt{m} (U_0 - W_e)^{\frac{3}{2}}} + \right. \\
 & \left. \frac{37U_0 - 79W_e + 12 \frac{W_e^2}{U_0}}{120 mc^2} + \frac{(49U_0^2 - 216U_0 W_e + 192W_e^2) e^2 E^2 \hbar^2}{1536U_0^2 m (U_0 - W_e)^3} \right) \quad (2)
 \end{aligned}$$

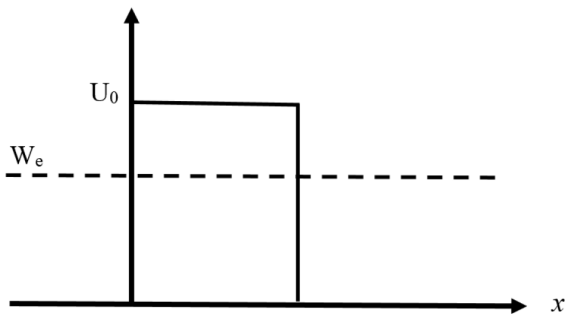


Fig. 3 The potential barrier in absence of an electric field.

The potential barrier width in relativistic case can be written as

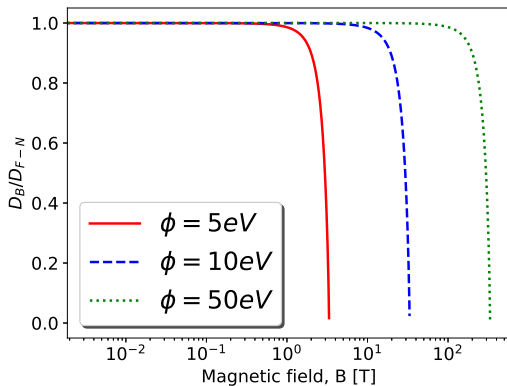
$$h_{rel} = h\sqrt{1 - V^2/c^2}, \quad (3)$$

where  $h$  is the potential barrier width in non-relativistic case and

$$V^2 = \frac{U_0 - W_e}{2m}$$

The transmission coefficient in presence of a magnetic field  $D_B$  in the first nonvanishing approximation can be written as

$$D_B = e^{-\frac{4\sqrt{2}\left(E^2 - 2\left(\frac{U_0 - W_e}{mc^2}\right)c^2 B^2\right)\sqrt{m}(U_0 - W_e)^{\frac{3}{2}}}{3E^3 e h}} \frac{4\sqrt{U_0 - W_e}(E^2 - B^2 c^2)^{\frac{3}{4}} E^{\frac{3}{2}} \sqrt{W_e}}{W_e E^3 + (E^2 - B^2 c^2)^{\frac{3}{2}} (U_0 - W_e)} \quad (4)$$



*Fig. 4. Dependence of the transmission coefficient upon magnetic field for the different work function values  $\phi$ .*

- The Fowler-Nordheim equation for field emission current density has been generalized to the relativistic case.
- The effect of Lorentz contraction of a potential barrier at the metal-vacuum interface was found. This effect results in increasing of the transmission coefficient by 0.015% for field emission in laboratory conditions and by 15% for emission from neutron stars surface.
- An expression for the transmission coefficient was found when the condition  $cB < E$  is satisfied. The transmission coefficient decreases by 0.1%, 5% and 50% when  $cB$  is equal to  $0.1 E$ ,  $0.5 E$  and  $0.9 E$  respectively.
- The magnetic induction is needed for "switching-off" an emission current during the elementary act of the acceleration is  $B \geq 0.5 T$ .

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Thank you for attention!