

The BIRD model

Insights and Future Developments

Author E. Spada

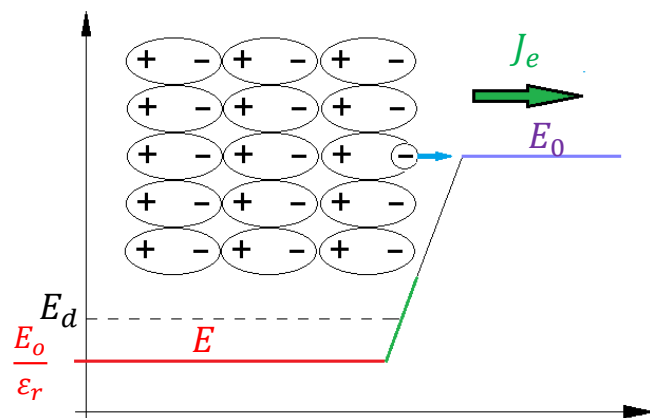
Co-authors: A. De Lorenzi, L. Lotto, N. Pilan, M. Zuin

- **SEMI-CLASSICAL** Approach to the BIRD Model
 - The BIRD model
 - The Electron Cold Emission from a Dielectric Layer
- **EXPERIMENTAL** Evidence
 - Old Data (Furuta et al. 2005)
 - New Data HVSGTF (Padova 2019)
- **THEORETICAL** development
 - Infinite well + Triangular barrier
 - Calculation of λ (*probability/time*)
- **SUMMARY**

Breakdown Induced by Rupture of Dielectric layer

The semi-classical Approach (2019)

IEEE TRANSACTIONS ON PLASMA SCIENCE, VOL. 47, NO. 5, MAY 2019



$$J_e = env$$

$$E > E_d$$

$$J_e = k(E_0 - E)(E - E_d)$$

$$J_e \approx k(E_0 - E)E \exp\left(-\frac{E_d}{E}\right)$$

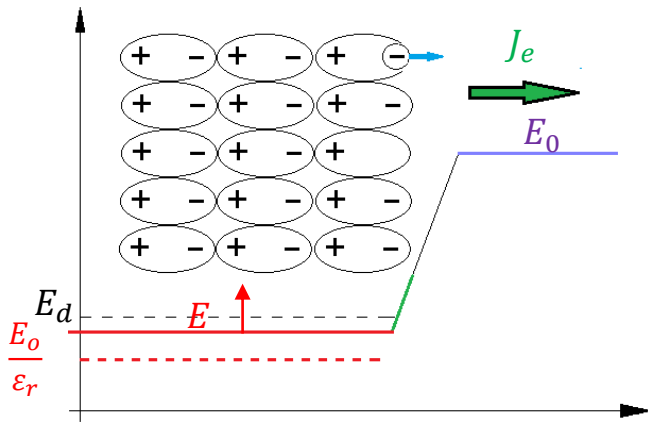
$$J_e(t = 0) \approx k_1 E_0^2 \exp\left(-\frac{k_2}{E_0}\right)$$

- The same law as F.N.
- Different constants k_1 and k_2
- Higher CURRENT DENSITY
- Lower β needed

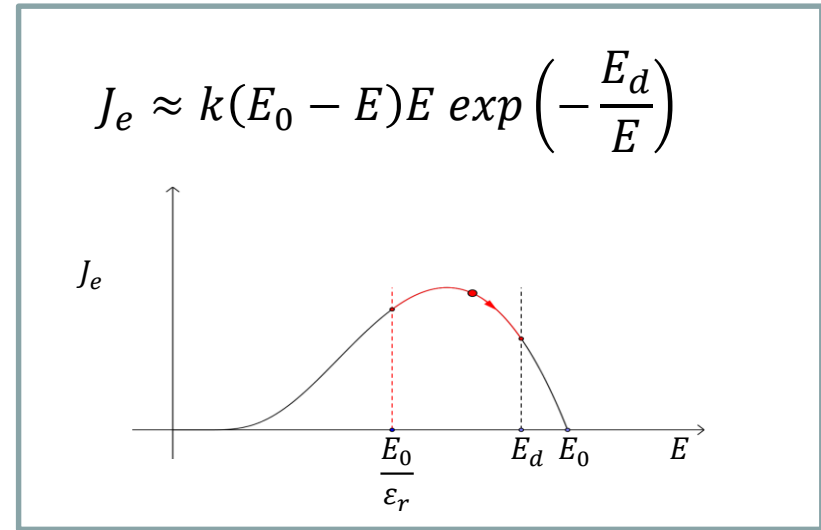
$$k_1 = \frac{\epsilon_r - 1}{\epsilon_r^2} k \quad k_2 = \epsilon_r E_d$$

$$k = \frac{e\epsilon_0 b}{\hbar} \approx 3.8 \cdot 10^{-11} \frac{A}{V^2}$$

The semi-classical Approach (2019)



...at constant VOLTAGE...

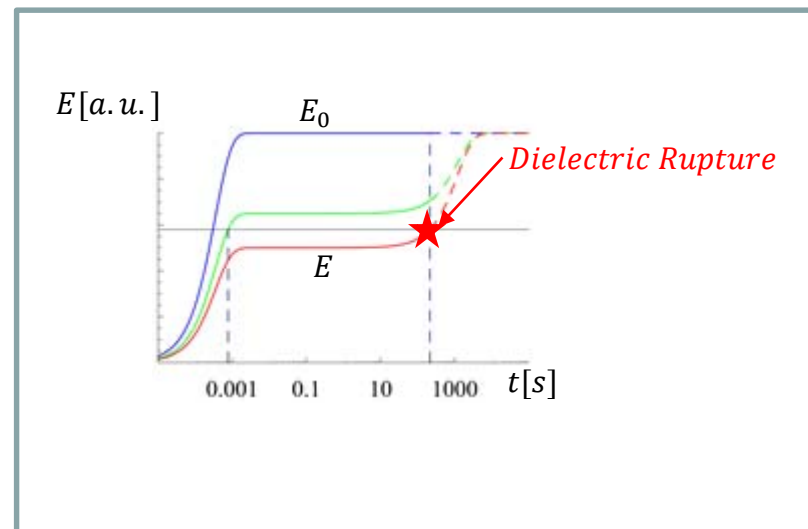
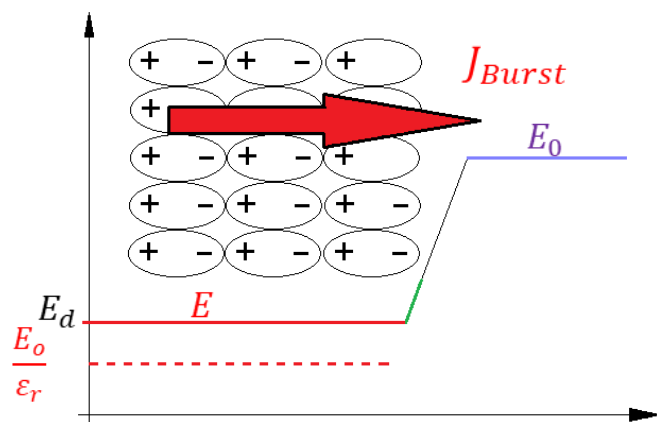


- The Electric Current initially increases and then decreases
- Eventually it reaches the dielectric strength E_d

Breakdown Induced by Rupture of Dielectric layer

The semi-classical Approach (2019)

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When the internal Field E reaches the Dielectric Strength E_d we have a

- Local rupture of the dielectric layer
- Burst of Current

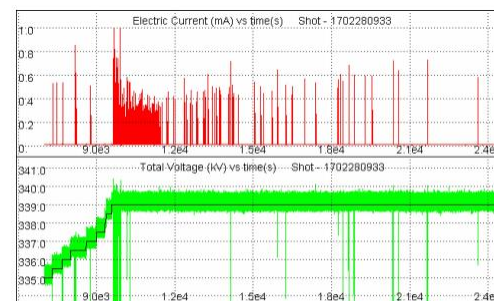
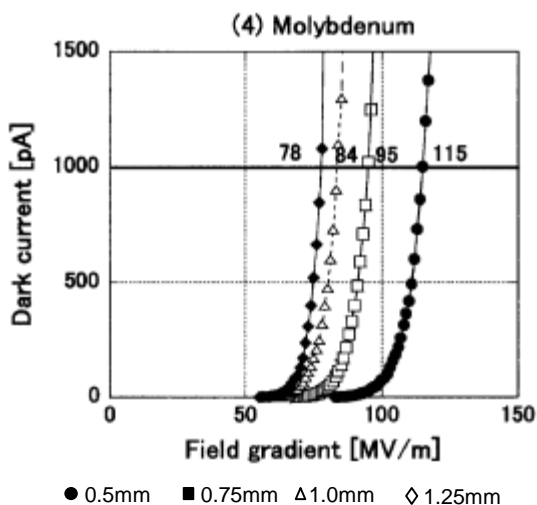


Fig. 5. Experimental current bursts (red line), reference (black line), and measured (green line) voltage in HVPTF during conditioning. A constant voltage $V = 339$ kV has been applied at time $t = 9588$ s.

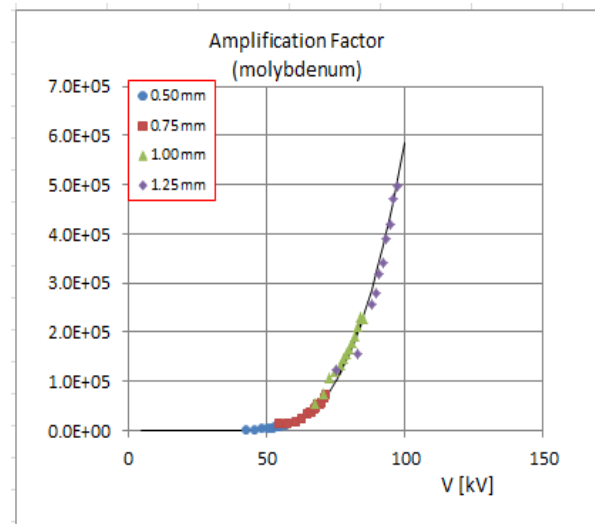
FURUTA et al. data (2005)

Nuclear Instruments and Methods in Physics Research A 538 (2005) 33–44



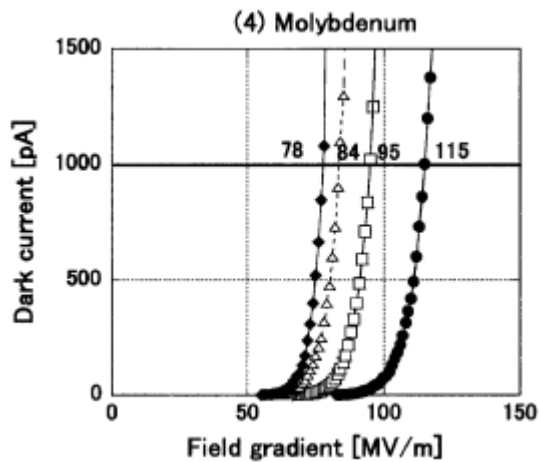
$$\Rightarrow MoO_3 \left\{ \begin{aligned} k_1 &\approx \frac{\epsilon_r - 1}{\epsilon_r^2} k = 0.91 \cdot 10^{-11} \frac{A}{V^2} \\ k_2 &\approx \epsilon_r E_d = 800 \text{ kV/mm} \end{aligned} \right.$$

$$\frac{I_{Furuta}}{J_e}$$



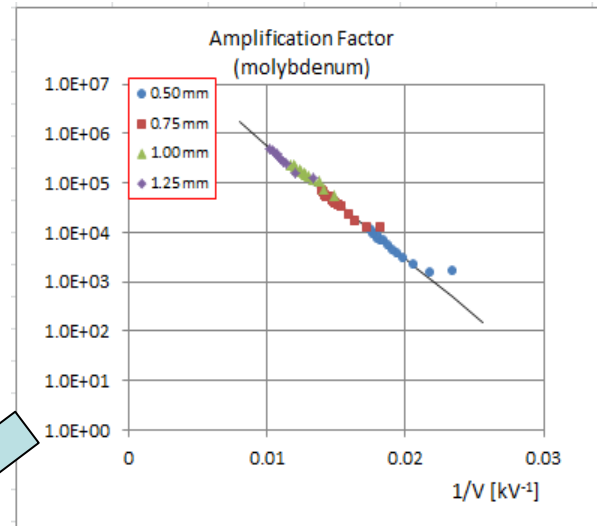
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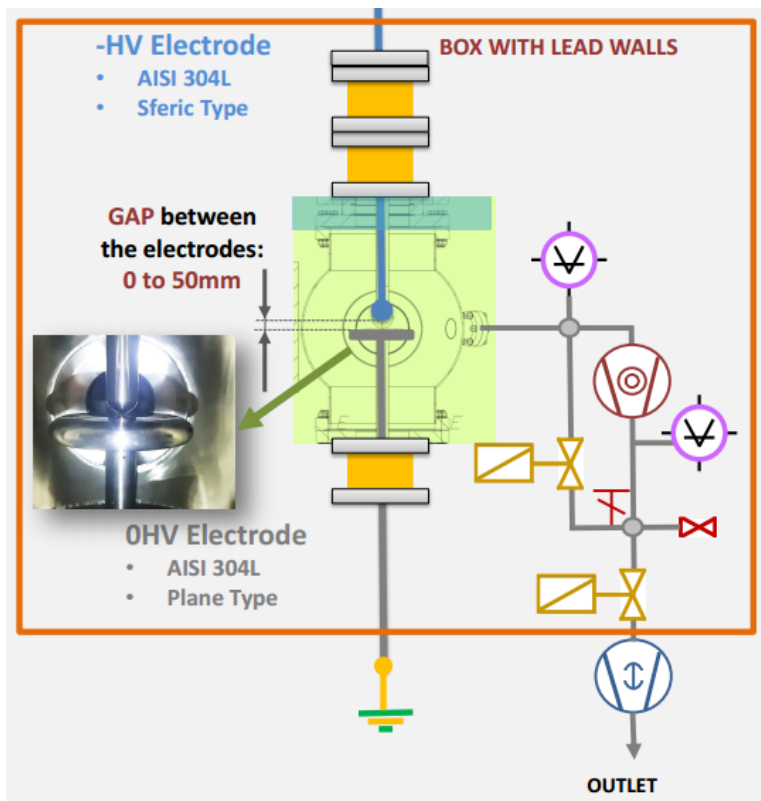
$$\frac{I_{Furuta}}{J_e}$$



$$A(V) = \frac{I_{Furuta}}{J_e} \sim \exp\left(-\frac{k_3}{V}\right)$$

HVSGTF (2019)

High Voltage Short Gap Test Facility

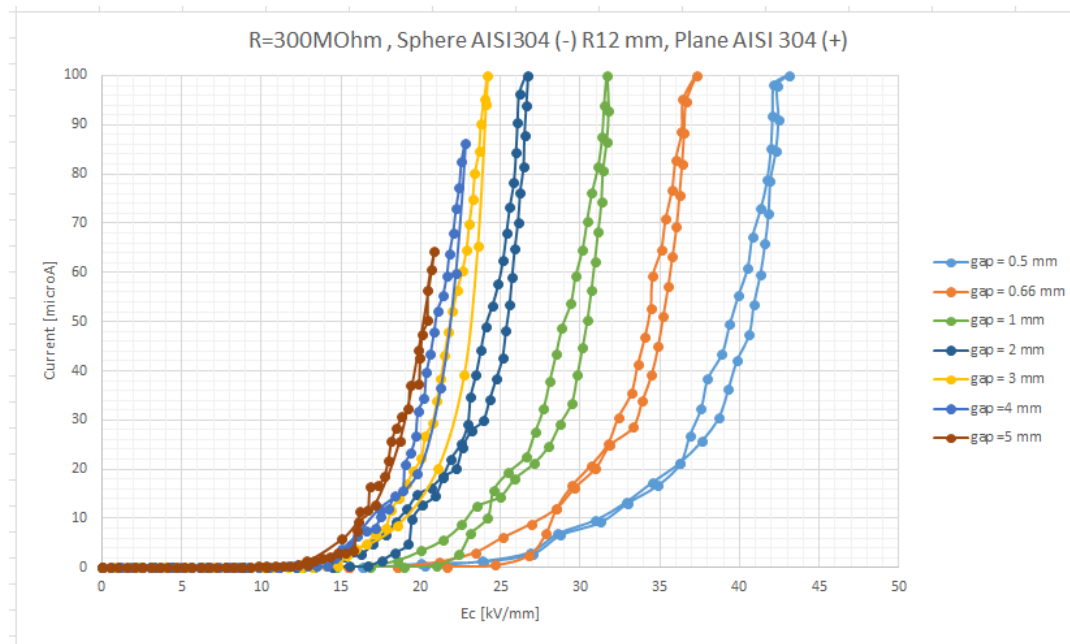


A new Facility has been recently developed in Padova, to study dark-current at short gap.

HVSGTF is designed to carry out experiments up to 100 kVdc with metal and non metal electrodes

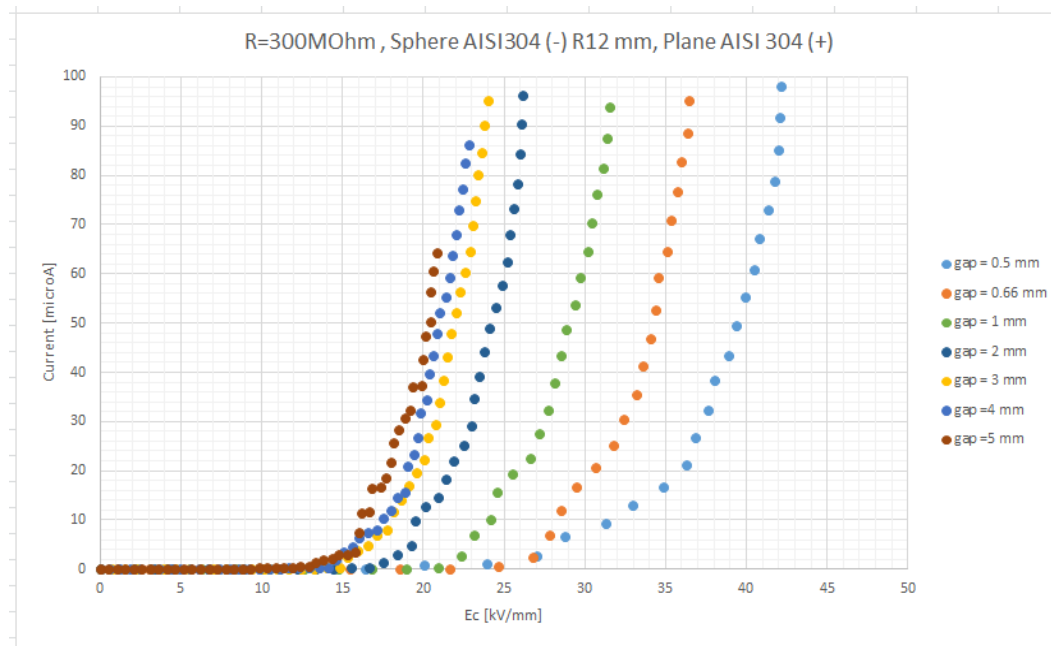
HVSGTF (2019)

First Result

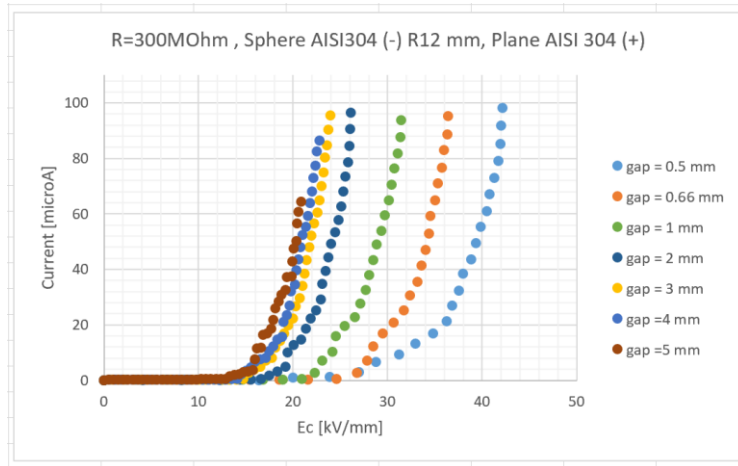


HVSGTF (2019)

First Result



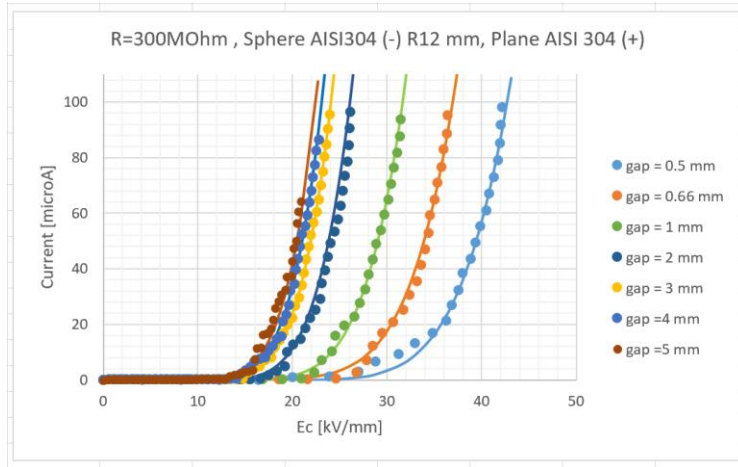
HVSGTF (2019)



$$I = S J_e(E_0) A(V)$$

$$S k_1 \beta^2 E_0^2 \exp\left(-\frac{k_2}{\beta E_0}\right) \cdot \left[1 + \xi \exp\left(-\frac{k_3}{V}\right)\right]$$

HVSGTF (2019)



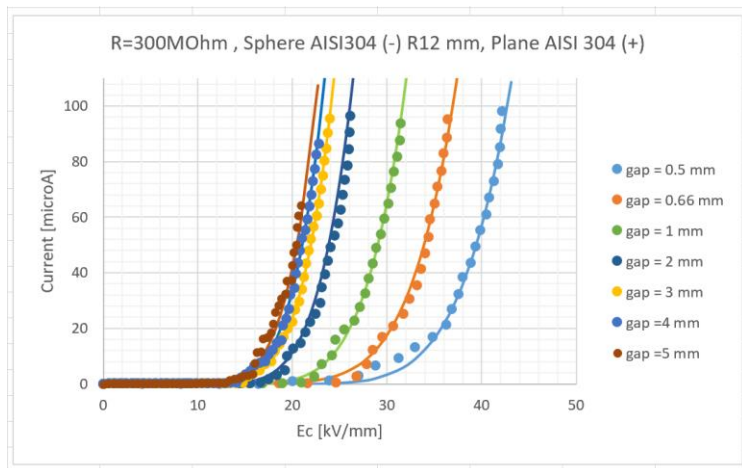
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Const	Value	
$S k_1 \beta^2$	5.20	$\cdot 10^{-20} \frac{Am^2}{V^2}$
$\frac{k_2}{\beta}$	82.7	kV/mm
ξ	400	
k_3	87.9	kV

Bisquare Method – Matlab Tool

HVSGTF (2019)

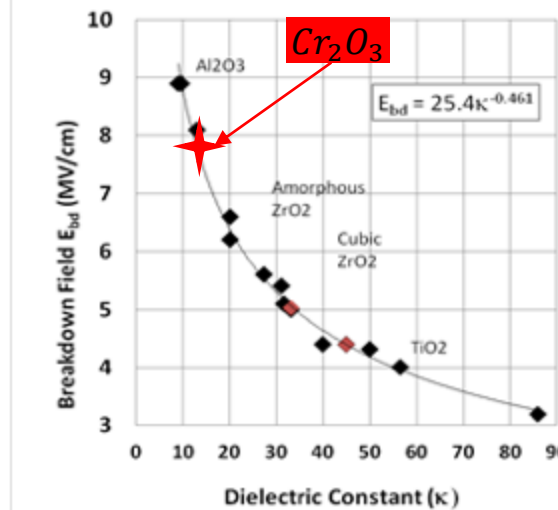


Cr₂O₃

$$\frac{k_2}{\beta} = 82.7 \frac{kV}{mm} = \frac{\epsilon_r E_d}{\beta}$$

$$\epsilon_r = 12.5 \quad E_d = 7.9 \frac{kV}{mm}$$

<https://intermolecular.com/materials/metal-oxides/>

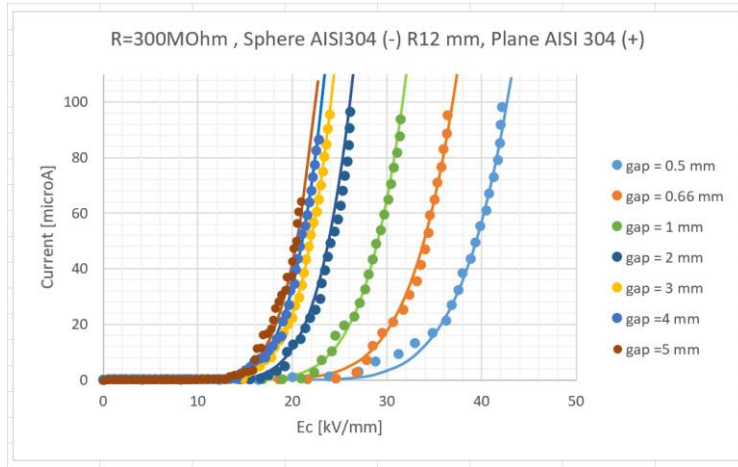


Const	Value	
$Sk_1\beta^2$	5.20	$\cdot 10^{-20} \frac{Am^2}{V^2}$
$\frac{k_2}{\beta}$	82.7	kV/mm
ξ	400	
k_3	87.9	kV

→ $\beta = 1.2$

Bisquare Method – Matlab Tool

HVSGTF (2019)



Cr_2O_3

$$Sk_1\beta^2 = 5.20 \cdot 10^{-20} \frac{Am^2}{V^2}$$

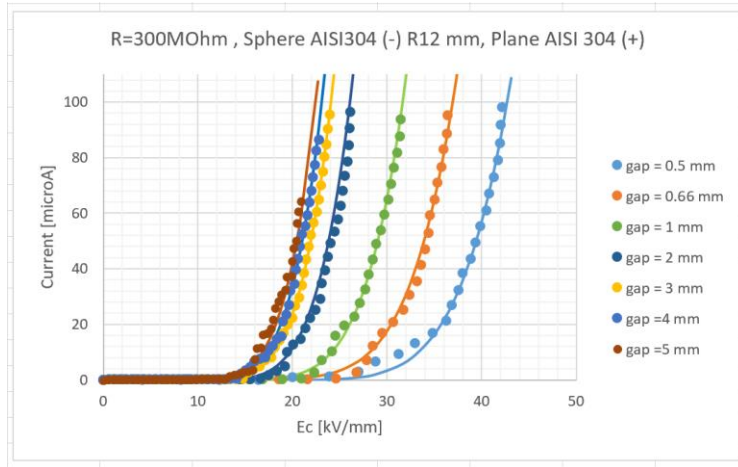
$$k_1 = \frac{\epsilon_r - 1}{\epsilon_r^2} k \sim 2.8 \cdot 10^{-12} \frac{A}{m^2}$$

Const	Value	
$Sk_1\beta^2$	5.20	$\cdot 10^{-20} \frac{Am^2}{V^2}$
$\frac{k_2}{\beta}$	82.7	kV/mm
ξ	400	
k_3	87.9	kV

→ $S = 1.3 \cdot 10^{-2} mm^2$

Bisquare Method – Matlab Tool

HVSGTF (2019)



Cr_2O_3

$\xi = 400$

?

$k_3 = 87.9 \text{ kV}$

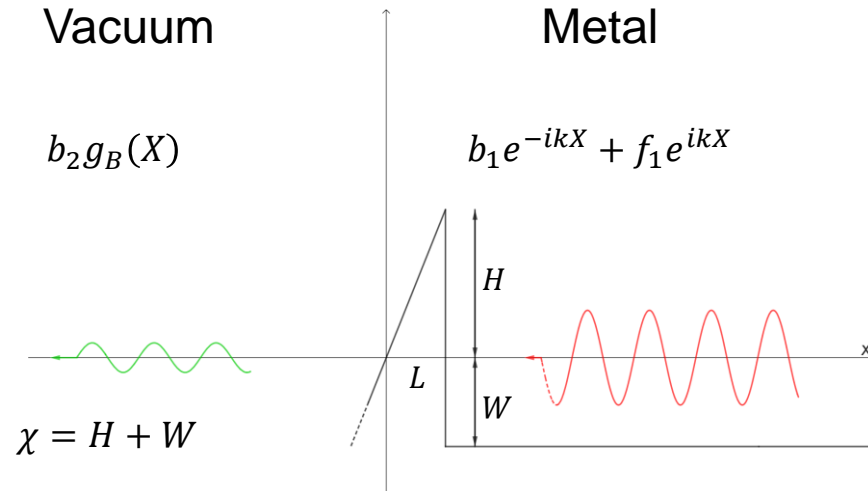
Const	Value	
$Sk_1\beta^2$	5.20	$\cdot 10^{-20} \frac{Am^2}{V^2}$
$\frac{k_2}{\beta}$	82.7	kV/mm
ξ	400	
k_3	87.9	kV

....model for A(V) is needed...

Bisquare Method – Matlab Tool

ECE from Metal

«G. Forbes and J. Deane» *Proc. R. Soc. A* (2011) 467, 2927-2947



$$g_B(X) = [Ai(k_A X) - iBi(k_A X)]$$

$$T = \frac{k_A}{\pi k} \left| \frac{b_2}{b_1} \right|^2$$

$$f_1 = \frac{1}{2} \left(g_B(x_L) - i \frac{k_A}{k} g'_B(x_L) \right) e^{-ikL} b_2$$

$$b_1 = \frac{1}{2} \left(g_B(x_L) + i \frac{k_A}{k} g'_B(x_L) \right) e^{ikL} b_2$$

$$x_L = k_A L = \left(\frac{2m}{\hbar^2} eE \right)^{\frac{1}{3}} \frac{H}{eE} = \left(\frac{2m}{\hbar^2} \right)^{\frac{1}{3}} \frac{H}{(eE)^{\frac{2}{3}}}$$

$$x_L^{\frac{3}{2}} = \left(\frac{2m}{\hbar^2} \right)^{\frac{1}{2}} \frac{H^{\frac{3}{2}}}{eE} = \frac{E_d}{E}$$

$$x_L \rightarrow \infty$$



$$E \rightarrow 0$$

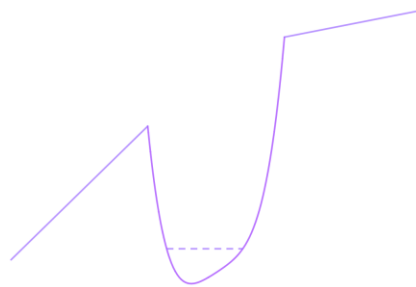
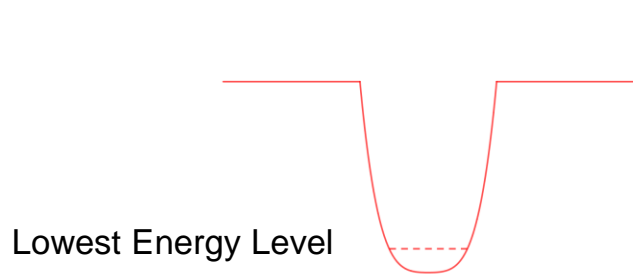
$$T = \frac{4}{\left(\frac{k_A \sqrt{x_L}}{k} \right) + \left(\frac{k_A \sqrt{x_L}}{k} \right)^{-1}} \exp\left(-\frac{4}{3} x_L^{\frac{3}{2}}\right)$$

$$J = \int D(W) T(W) dW$$

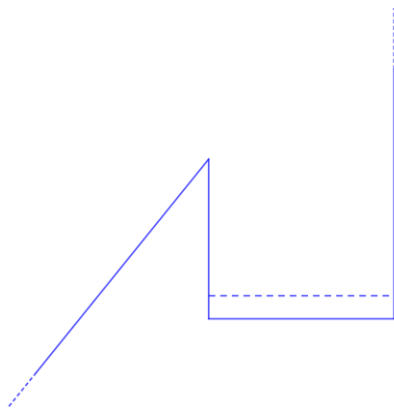
$$J = k_1 E^2 \exp\left(-\frac{k_2}{E}\right)$$

ECE from Dielectrics

Electrons are confined in polarization structures....

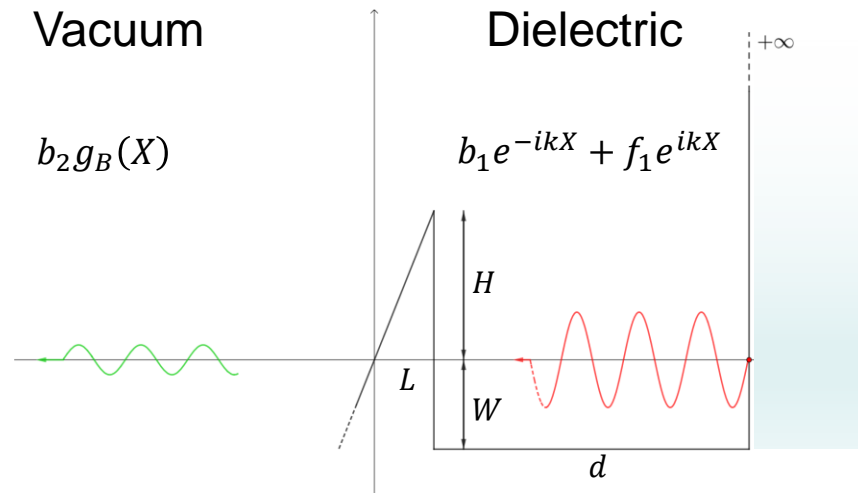


...modified in presence of an Electric Field



The simplest possible model

ECE from Dielectrics



$$g_B(X) = [Ai(k_A X) - iBi(k_A X)]$$

$$T = \frac{k_A}{\pi k} \left| \frac{b_2}{b_1} \right|^2$$

$$f_1 = \frac{1}{2} \left(g_B(x_L) - i \frac{k_A}{k} g'_B(x_L) \right) e^{-ikL} b_2$$

$$b_1 = \frac{1}{2} \left(g_B(x_L) + i \frac{k_A}{k} g'_B(x_L) \right) e^{ikL} b_2$$

$$b_1 e^{-ik(L+d)} + f_1 e^{ik(L+d)} = 0$$



$$\text{tg}(kd) = - \frac{k}{k_A} \frac{g_B(x_L)}{g'_B(x_L)}$$

$$W \rightarrow W + i \frac{\hbar}{2} \lambda$$

$$e^{\frac{i}{\hbar} W} \rightarrow e^{\frac{i}{\hbar} W} \cdot e^{-\frac{\lambda}{2} t}$$

$$|\psi|^2 \rightarrow |\psi|^2 e^{-\lambda t}$$

$$\lambda = \text{decay constant} = \frac{\text{Probability}}{\text{time}}$$

ECE from Dielectrics

for $E \rightarrow 0$ ($x_L \rightarrow \infty$)

$$T = \frac{k_A}{\pi k} \left| \frac{b_2}{b_1} \right|^2$$

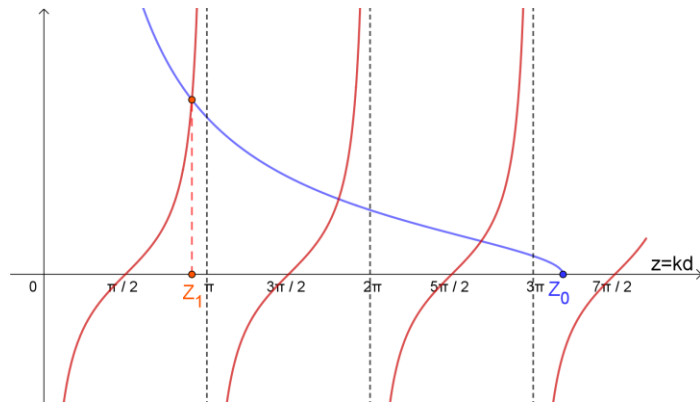


$$T = \frac{4k}{k_A \sqrt{x_L}} \exp\left(-\frac{4}{3} x_L^{\frac{3}{2}}\right)$$

$$\text{tg}(kd) = -\frac{k}{k_A} \frac{g_B(x_L)}{g'_B(x_L)}$$



$$\text{tg}(kd) = -\frac{k}{\sqrt{k_0^2 - k^2}}$$



$$\text{tg}\left(z - \frac{\pi}{2}\right) = \sqrt{\frac{z_0^2}{z^2} - 1}$$

$$z_0 = k_0 d = \frac{\sqrt{2m\chi}}{\hbar} d$$

$$z = kd$$

$$\frac{\pi}{2} < z_1 < \pi$$

$$k_1 \sim \frac{1}{d}$$

$$W_1 \sim \frac{\hbar^2}{md^2}$$

ECE from Dielectrics

How can we calculate λ ?

Gamow theory of alpha decay - 1928

G. Gamow, ZP, **51**, 204

We can now solve the problem for two symmetrical potential barriers [Fig. 3]. We shall seek two solution.

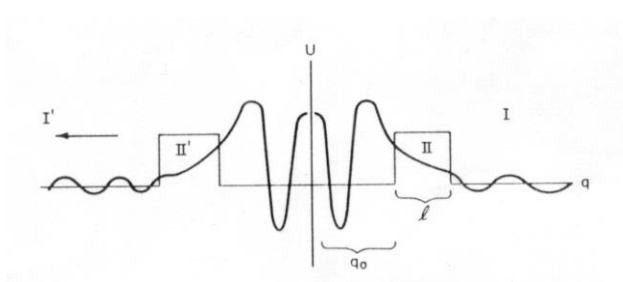


Figure 3:

The conservation principle then reads

$$\frac{\partial}{\partial t} e^{-\lambda t} \int_{-(q_0+l)}^{+(q_0+l)} \Psi_{II,III}^{(q)} \cdot \Psi_{II,III}^{(q)} dq = -2 \frac{A^2 h}{4\pi \cdot i \cdot m} \cdot 2ik \cdot e^{-\lambda t},$$

from which we obtain

$$\lambda = \frac{4hk \sin^2 \theta}{\pi m \left[1 + \left(\frac{k'}{k_0} \right)^2 \right] 2(l + q_0)k} \cdot e^{-\frac{4\pi l \sqrt{2m}}{h} \sqrt{U_0 - E}}, \quad (5)$$

where k is a number of order of magnitude one.

$$\lambda \sim \frac{\hbar k}{m(L + d)} T$$

ECE from Dielectrics

$$\lambda = \frac{\hbar k}{mL} T = \frac{\hbar k}{mL} \cdot \frac{k_A}{\pi k} \left| \frac{b_2}{b_1} \right|^2 = \frac{8W_1}{\hbar \sqrt{x_L^3}} \exp\left(-\frac{4}{3} \sqrt{x_L^3}\right) \propto E \cdot \exp\left(-\frac{4}{3} \frac{E_d}{E}\right)$$

$$\sigma = \varepsilon_0(E_0 - E) \quad J_e = \sigma \lambda$$

$$J_e = \sigma \lambda = k(E_0 - E)E \exp\left(-\frac{4}{3} \frac{E_d}{E}\right)$$

...the same functional shape for J_e has been recovered !

BIRD model

- Explain, at least qualitatively, some features (ECE – Bursts)
- The semi-classical model is consistent with the first experimental results
- An amplification factor $A(V)$ is needed
- A simple quantum model gives the expected shape of ECE

Future investigations require:

- More experimental work (Current evolution – Different electrodes -..)
- The development of a much more solid quantum model
- To understand the physical origin of $A(V)$

THANK YOU FOR THE ATTENTION