Finding Heavy Scalar bosons at the LHeC

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Motivation

- ► Additional CP even scalars come with many theory extensions.
- Can be connected to neutrino mass physics, Dark Matter, vacuum stability ...
- \blacktriangleright There are hints in the LHC data for a \sim 300 GeV scalar.
- This is difficult to detect and almost impossible to study at the LHC.

The Model

- SM extended with a complex neutral scalar boson S
- ► *S* is a singlet under the SM gauge group
- Most general renormalizable scalar potential:

 $m_1^2 H^{\dagger} H + m_2^2 S^{\dagger} S + \lambda_1 (H^{\dagger} H)^2 + \lambda_2 (S^{\dagger} S)^2 + \lambda_3 (H^{\dagger} H) (S^{\dagger} S)$

- The scalar fields H, S develop vevs v, x.
- \blacktriangleright Mass matrix mixes the fields, mixing angle α
- The neutral mass eigenstates:

$$\left(\begin{array}{c}h_1\\h_2\end{array}\right) = \left(\begin{array}{c}\cos\alpha & -\sin\alpha\\\sin\alpha & \cos\alpha\end{array}\right) \left(\begin{array}{c}H\\S\end{array}\right)$$

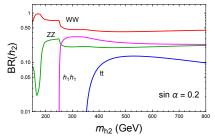
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2 / 13

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Physical scalars



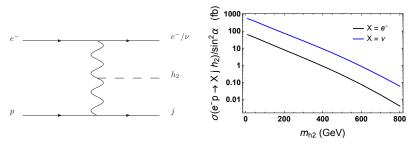
- In the limit of $\alpha \rightarrow 0$ the SM Higgs boson is recovered.
- We define h_1 being "SM-like" Higgs, meaning:
 - $m_{h_1} = 125 \text{ GeV}$
 - interaction strength = $SM \times \cos^2 \alpha$
 - Limit from current LHC data: α < 0.3.
- h₂ is a heavy Higgs-like scalar. We posit:
 - $m_{h_2} > m_{h_1}$
 - Interaction strength = (SM Higgs with $m_h = m_{h_2}$) × sin² α .
 - No other decay modes $(h_2
 ightarrow 2 \, h_1$ possible above threshold)

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Scalar production at the LHeC



- Production modes identical to SM Higgs boson.
- CC cross section is larger than NC one.

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Considered signal modes

We focus on the leading decay modes, $h_2 \rightarrow VV$. The three signal modes are:

1. $\mu_{\ell \ell}^{Z} := h_{2} \rightarrow ZZ \rightarrow 4\ell$ 2. $\mu_{\ell q}^{Z} := h_{2} \rightarrow ZZ \rightarrow 2\ell 2q$ 3. $\mu_{\ell q}^{W} := h_{2} \rightarrow WW \rightarrow \nu \ell 2j$

Quote from the article:

... other interesting decay channels ... extremely useful to characterise ... the extra scalar, ... discriminate among different models, are the di-higgs and di-top decay modes [and the hadronic decays of the vector bosons].

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Considered backgrounds

Nr.	final state	$\sigma_{ m LHeC}$ [fb]
1	e ⁻ jWW	23.0
2	$e^{-}jZW^{+}$	4.16
3	e ⁻ jZZ	0.1
4	ujWW	10.4
5	$ u$ jZW $^-$	8.0
6	ν jZZ	2.4

Table: The SM background processes considered in this analysis. The samples have been produced with the following cuts: $P_T(j) > 20$ GeV, $P_T(l) > 2$ GeV and $|\eta(j/l)| < 6$.

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Sidenote on simulation

Cuts:

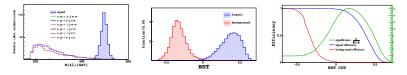
- ► All: |η(j)| < 4.5</p>
- ▶ DIS jet: P_T(j) > 10 GeV
- Electrons and muons: $P_T(I) > 2 \text{ GeV}$

Simulation:

- Madgraph, Pythia6, Delphes
- We patched Pythia as prescribed by Uta
- We used a modified delphes card, including:
 - a muon module
 - adjusted eta acceptance
 - hadron tracking precision as in CDR

 \Rightarrow This needs to be made available publicly to support 'external' BSM studies.

Signal channel 1: 4 leptons

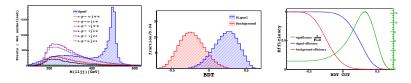


- ► $\ell_{\alpha}^{+}\ell_{\alpha}^{-}\ell_{\beta}^{+}\ell_{\beta}^{-}$ with flavours $\alpha, \beta \in \{e, \mu\}$.
- τ not considered; unsure about τ reconstruction@LHeC.
- ▶ Preselection collects same flavor $\ell^+\ell^-$ pairs with $M_{2\ell}$ closest to m_Z .
- ▶ Main backgrounds are vjZZ and e⁻jZZ.
- MVA BDT with 42 kinematical observables.
- Most important observables are $M_{4\ell}$ and $M_{2\ell}$.

Analysis

- We created in total 42 observables from the visible final states.
- Made use of a Boosted Decision Tree and performed a multivariate analysis (with TMVA).
- Signal significances for each of the signal channels separately.
- Channels were combined with a statistical tool (Combined Higgs limits).

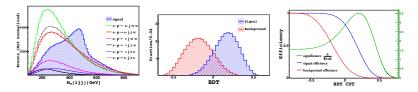
Signal channel 2: semileptonic without MET



- Two leptons and two jets, each with M_{2f} close to m_Z .
- ▶ Main backgrounds are vjZZ and e⁻jZZ, contributions from processes with at least one W boson
- Analysis strategy similar to the one for 4 leptons.
- Most important observables (BDT ranking): total invariant mass, η(ℓ), angular sep. between ℓ and Z.

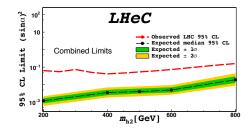
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Signal channel 2: semileptonic with MET



- One lepton, two jets with inv. mass $\sim m_W, m_Z$.
- Difficult to reconstruct because of escaping neutrino.
- Preselection collects among all jets the pair with invariant mass closest to m_Z.
- ▶ We reconstruct from lepton, MET, and *W* candidate the transverse mass of the second vectorboson.
- Most important observable is the total visible transverse mass.

Combination of signal modes and comparison to LHC



- Including a systematic uncertainty of 2% on the background.
- ► LHC results from 35.9/fb at 13 TeV, including 4ℓ, 2ℓ2j and 2ℓ2ν.
- Extrapolations for HL-LHC, without new backgrounds, comparable to LHeC, better for larger masses.
- LHC analysis seems to struggle for masses $\sim 100 \div 300$ GeV.

Conclusions

- Present searches at the LHC are compatible with additional heavy scalar particles that mix on the percent level with the SM Higgs boson.
- Multivariate techniques: h₂ with mass 200÷ 800 GeV and mixings > sin² α ∼ 10⁻³.
- Promising discovery prospects of heavy scalars at the LHeC
- Complementary to the searches at the LHC, especially wrt semileptonic final states.
- Many other interesting channels exist that could improve this sensitivity further.
- \Rightarrow We have a new working and efficient pipeline for BSM@ep studies at the reconstructed level!

4 Leptons

 $MET, P_T(j), \eta(j), P_T(l), \eta(l)$ $P_T(Z)_{ll}, \eta(Z)_{ll}, \Delta R(Z)_{ll}$ $P_T(h)_{llll}, \eta(h)_{llll}, \Delta R(h)_{llll}$ $P_T(h+j), \Delta R(h,j)_{llll}$

2l 2j

$$\begin{split} MET, P_{T}(j), \eta(j), P_{T}(l), \eta(l) \\ P_{T}(j)_{beam}, \eta(j)_{beam}, P_{T}(j_{1}), \eta(j_{1}), P_{T}(j_{2}), \eta(j_{2}) \\ P_{T}(l_{1}), \eta(l_{1}), P_{T}(l_{2}), \eta(l_{2}), P_{T}(Z)_{jj}, \eta(Z)_{jj}, M(Z)_{jj} \\ P_{T}(Z)_{ll}, \eta(Z)_{ll}, M(Z)_{ll}, P_{T}(h)_{jjll}, \eta(h)_{jjll}, M(h)_{jjll} \\ \Delta R(h)_{jjll}, \Delta R(Z_{ll}, Z_{jj}), \Delta R(h, j_{beam}) \\ \bigvee 2j \end{split}$$

 $MET, P_{T}(j)_{beam}, \eta(j)_{beam}, P_{T}(l), \eta(l)$ $P_{T}(j1)_{W}, \eta(j1)_{W}, P_{T}(j_{2})_{W}, \eta(j_{2})_{W}$ $P_{T}(W)_{jj}, \eta(W)_{jj}, M(W)_{jj},$ $P_{T}(W+l), \eta(W+l), M(W+l),$ MT(W+l),

With I = electron or muons

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