

First consideration of RF scheme to mitigate longitudinal instabilities in Super Proton-Proton Collider

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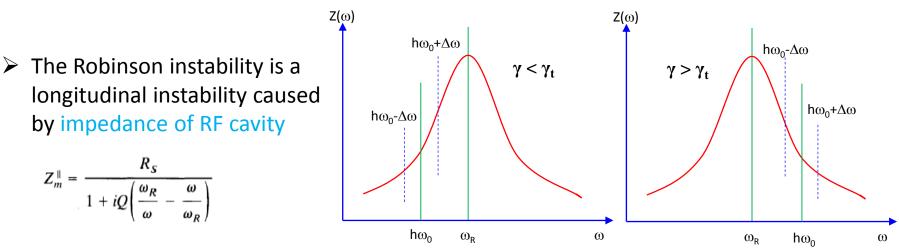
Outline

- A simple review of the main longitudinal instabilities
- RF scheme to mitigate longitudinal instability in SPS/LHC
- Introduction to SPPC
- RF scheme to mitigate longitudinal instability in SPPC
- Summary and outlook



A simple review of the main longitudinal instabilities

Single bunch effects: Robinson Instability [1]



In essence:

✓ the Robinson instability comes from the fact that the revolution frequency of an off-momentum beam is slightly different from the on-momentum revolution frequency, depending on both the slip factor and the energy deviation.

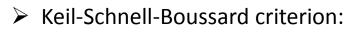
$$\frac{\Delta\omega}{\omega_0} = -\eta \frac{\Delta p}{p_0}$$

- Stable condition:
 - ✓ If η >0, above transition, Robinson damping → $h\omega_0 > \omega_R$
 - ✓ If η <0, below transition, Robinson damping → $h\omega_0 < \omega_R$



A simple review of the main longitudinal instabilities

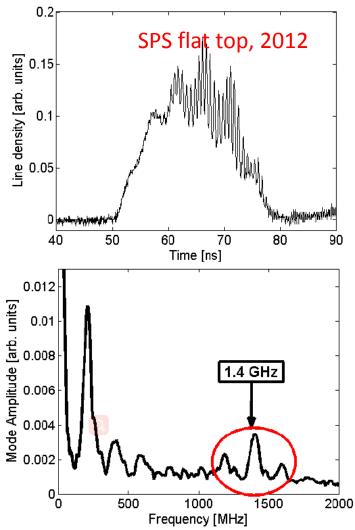
Single bunch effects: Longitudinal microwave instability [2]



$$\frac{Z_{II}}{n} \le \frac{2\pi\beta^2 E\sigma_{\delta}^2 |\eta| F}{e\hat{I}}$$

- grows usually fast;
- the bunch will split into many micro bunches and the momentum spread will increase.
- In proton accelerators the microwave instability is observed as a fast increase of the bunch length and thus of the longitudinal emittance ε_l .
- $f_r \tau >> 1$, microwave signals (100MHz~3GHz)
- Usually, longitudinal microwave instability will not be worsened for larger hadron machines.[3]

 $\left|\frac{Z_{\prime\prime}}{n}\right| \propto \frac{\sigma_{\delta}^2 \sigma_{\tau}}{N_b}$



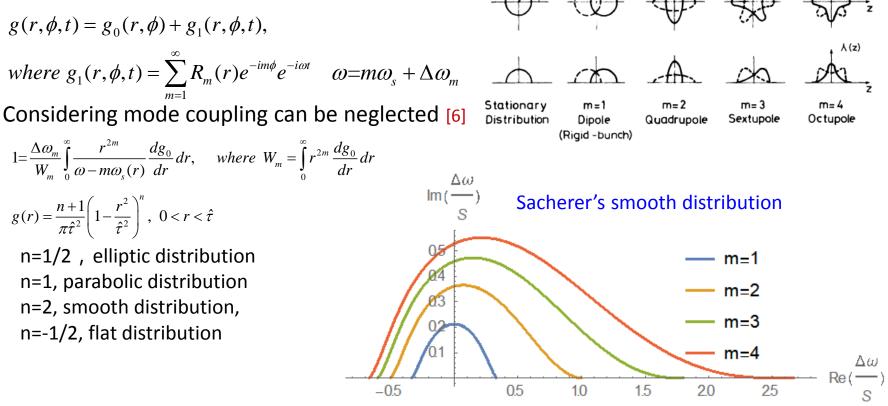


Loss of Landau damping

Landau damping:

• Landau damping of the instability can come from the spread in the synchrotron frequency due to the nonlinearity of the synchrotron force in the bunch.

Sacherer formalism [4,5]



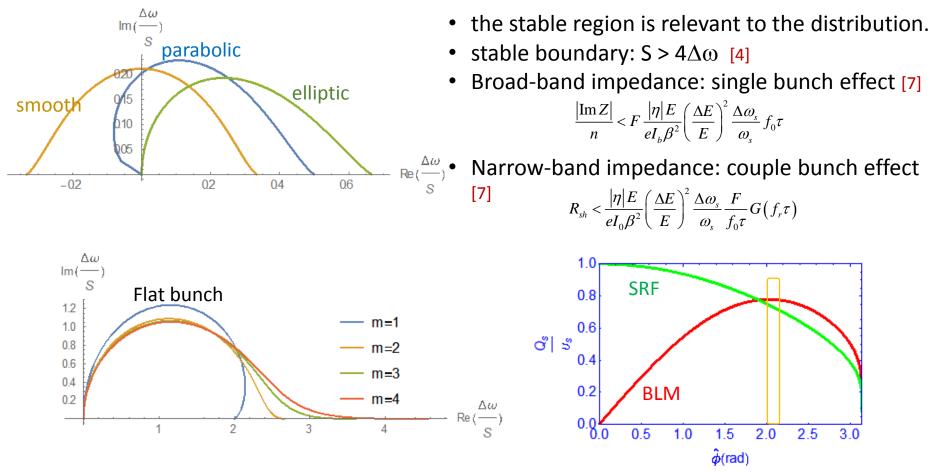
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[6] Ibon Santiago González, Loss of Longitudinal Landau Damping in the LHC Injectors,



A simple review of the main longitudinal instabilities

Loss of Landau damping



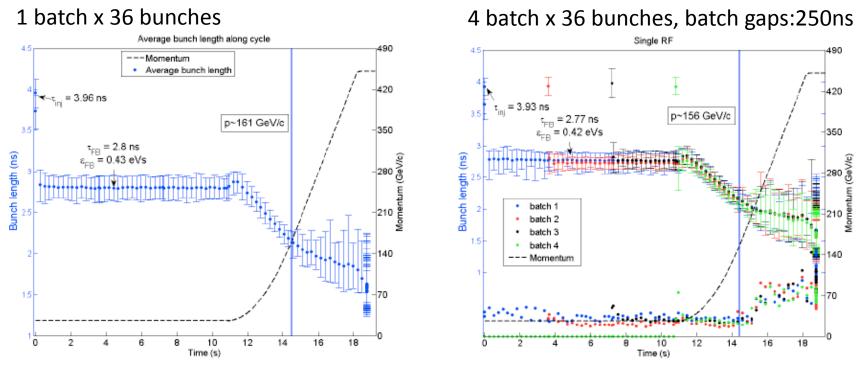
Iandau damping is lost when the derivative of synchrotron frequency is zero outside the bunch center. [8]

[8] E. Shaposhinikova, Bunched beam transfer matrices in single and double RF systems, CERN-SL-94-19



Longitudinal instabilities in the SPS [9,10]

Single RF, LHC beam with 50 ns spacing, 1.6x10¹¹ p/b

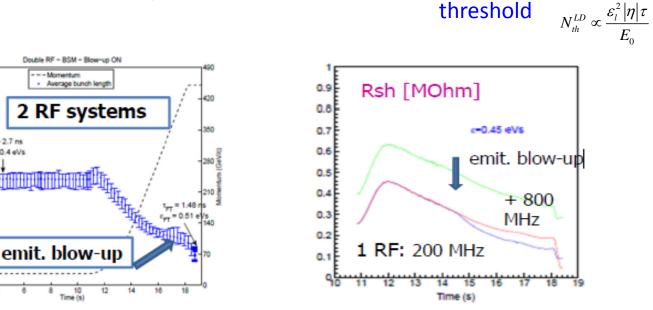


- Longitudinal multi-bunch instability: very low threshold
- The instability starts almost at the same energy.



Cures to stabilize beam in the SPS [9, 10]

- Active damping: RF feedback, feed-forward, longitudinal damper;
- 4th Harmonic RF system (800MHz) in bunch-shortening mode (BSM);
- Controlled emittance blow-up;
- > Optics with lower transition energy, $N_{th} \sim |\eta|$
- Impedance reduction;

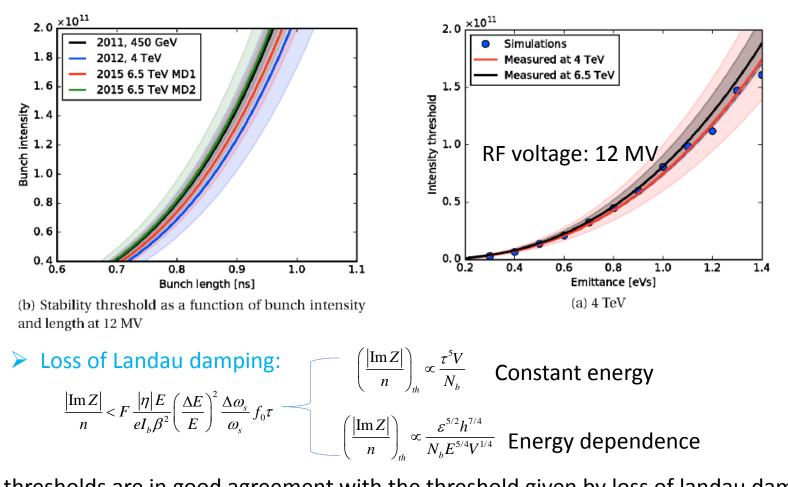


Loss of landau damping

[10] T. ARGYROPOULOS, Longitudinal beam instabilities in a double RF system, CERN-Thesis-2015



Longitudinal instabilities in the LHC [11]



The thresholds are in good agreement with the threshold given by loss of landau damping.

[11] Juan Federico ESTEBAN MÜLLER, Longitudinal intensity effects in the CERN Large Hadron Collider, CERN-Thesis-2016



RF scheme to mitigate Longitudinal instabilities in the LHC

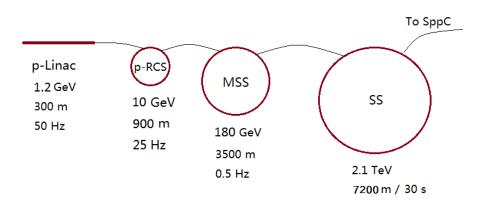
- > 400MHz (main) + 200MHz(capture) [7]
- At injection into the LHC :
 - 200 MHz RF system only for capture.
 - Four cavities, 0.75 MV each, which can be pushed up to 1 MV.
 - For capture the operational total voltage at 200 MHz is 3 MV;
- After capture:
 - 400 MHz RF system adiabatically is increased up to 8 MV
 - 200 MHz RF system is decreased to zero;
 - Acceleration only with 400 MHz.
- On the flat top:
 - the emittance is 2.5 eVs.
 - f_{rf} =400MHz, $V_{rf(max)}$ =16MV to produce ~1 ns long bunches.
- ➤ Cure for :
- The main RF system provides some damping of injection errors and natural Landau damping;
- Controlled emittance blowup by excitation with band-limited noise

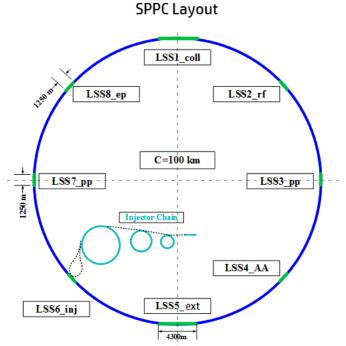
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Layout and main parameters

main parameters	Value
Circumference (km)	100
Dipole field (T)	12
Injection energy (TeV)	2.1
Collision energy (TeV)	37.5
Transition gamma γ_{tr}	99.21
Bunch intensity	1.5 x10 ¹¹
Number of bunches	10080
Bunch spacing (ns)	25
Bunch length during physics (m)	0.0755
RF frequency (MHz)	400

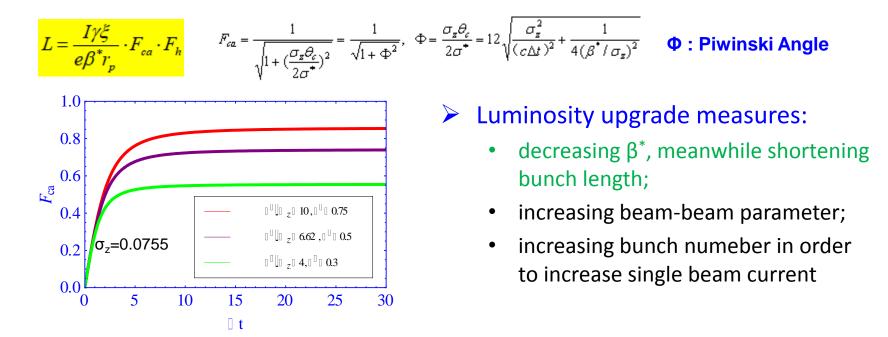




- LSS1:collimation system (comnbine β and momentum)
- LSS5: extraction system
- LSS2: RF system
- LSS6: Injection system
- LSS3/LSS7: pp collision
- LSS4: AA collision
- LSS8: ep collision



Luminosity



The requirements for SPPC longitudinal dynamics

- Two constraints:
 - Intra-beam Scattering
 - Beam instability limits

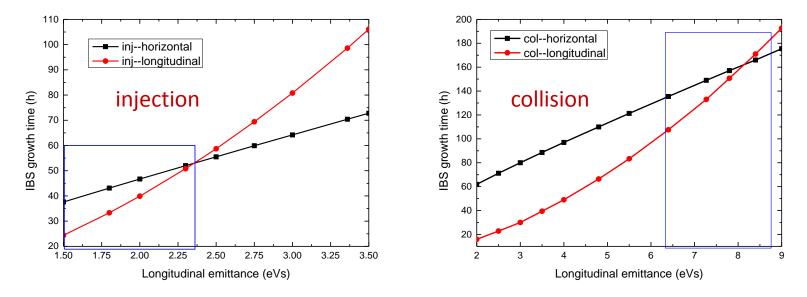
➤The goals:

- achieve this goal of bunch length 7.55cm.
- The variable range of bunch length for luminosity upgrade



Intra-beam Scattering (IBS)

Multiple small-angle Coulomb scattering of charged particle beams



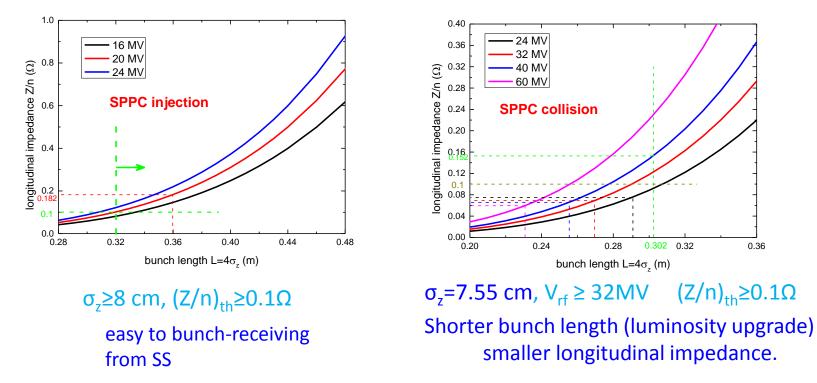
- At injection, longitudinal emittance between 1.5 eVs and 2.5 eVs may be a good choice; (injection time 840s)
- In collision, longitudinal emittance between 6 eVs and 9 eVs may be a good choice; (time during physics 14.2h)
- Using controlled emittance blowup, IBS maybe not a serious matter.



Loss of Landau damping at 400 MHz

Loss of Landau damping:

 $\frac{\left|\operatorname{Im} Z_{jj}\right|}{n} \leq \frac{3\pi^2}{64} \frac{h^3 V_{rf}}{I_b} \left(\frac{L}{2\pi R}\right)^5$



- Injection: $\sigma_z = 9$ cm, 0.182 Ω @20 MV; Collision: $\sigma_z = 7.55$ cm, 0.152 Ω @40 MV
- Longitudinal reactive impedance less than 0.15 Ω needs to be studied! [LHC~0.1 Ω]



RF schemes to mitigate longitudinal instability in SPPC

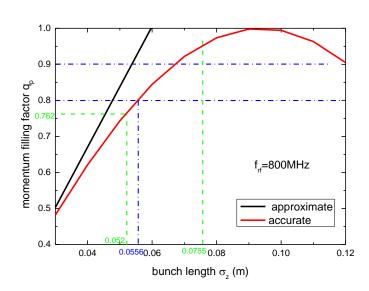
• In order to enhance Landau damping, a large spread in synchrotron frequency inside the bunch is required.



use a higher harmonic cavity (800MHz RF cavity needs to be studied)
 Dual harmonic RF system
 Controlled emittance blow-up

Controlled emittance blow-up

1) 800 MHz RF system



- 7.55 cm @ 800 MHz, q_p > 0.9, i.e. significant beam loss. So the rms. bunch length will have to be reduced if using 800 MHz RF cavity.
- $q_p \le 0.8 --> \sigma_z \le 5.56 \text{ cm};$
- the goal bunch length is set as $\sigma_z = 5.2$ cm, then $q_p = 76.2\%$;
- Correspondingly: ε_s =6.4eVs, V_{RF}=52MV, σ_{δ} =0.79x10⁻⁴

The corresponding impedance threshold

$$\frac{\left|\frac{\operatorname{Im} Z_{\prime\prime}}{n}\right|_{800}}{\left|\frac{\operatorname{Im} Z_{\prime\prime}}{n}\right|_{400}} = \frac{\left(h^{3} V_{rf} L^{5}\right)_{800}}{\left(h^{3} V_{rf} L^{5}\right)_{400}} \approx 1.6 \Rightarrow \left|\frac{\operatorname{Im} Z_{\prime\prime}}{n}\right|_{800} = 0.243\Omega$$



② Dual harmonic RF system (400MHz+800MHz)

- Applied RF voltage is: $V(\phi) = V_0(t) \left[\sin \phi + k \sin(n\phi + \phi_2) \right]$
- in $(\phi, \Delta E/\omega_0)$ phase space, the Hamiltonian is

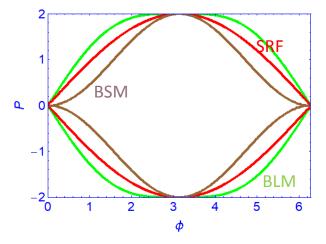
$$H=\frac{h\alpha_0^2\eta}{2\,\beta^2 E}W^2+\frac{eV_0}{2\pi}U(\phi)$$

• The potential well is

$$U(\phi) = -\frac{1}{V_0} \int_{\phi_s}^{\phi} \left[V(\phi) - V(\phi_s) \right] d\phi = \left(\cos\phi - \cos\phi_s \right) + \frac{k}{n} \left[\cos(n\phi + \phi_2) - \cos(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + k\sin(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + b\cos(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + b\cos(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + b\cos(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + b\cos(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + b\cos(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + b\cos(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + b\cos(n\phi_s + \phi_2) \right] + (\phi - \phi_s) \left[\sin\phi_s + b\cos(n\phi_s + \phi_2) \right] + (\phi - \phi_$$

• In normalised phase space (ϕ , $P = -h|\eta|/\nu_s \delta$), the Hamiltonian is

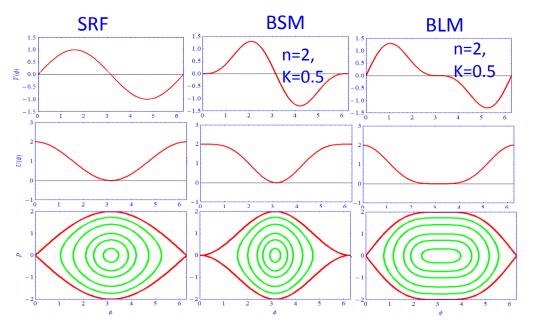
$$H = -\frac{1}{2}\nu_s P^2 + \nu_s U(\phi) \qquad \qquad \nu_s = \sqrt{\frac{heV_0|\eta|}{2\pi\beta^2 E}}$$



- Bunch lengthening mode (BLM): n=2, k=0.5, φ₂=0, for η>0
- > Bunch shortening mode (BSM): n=2, k=0.5, $\phi_2 = \pi$, for $\eta > 0$
- A_{BLM}/A_{single}~1.1478, A_{BSM}/A_{single}~0.7854;
- The bucket height keeps the same.



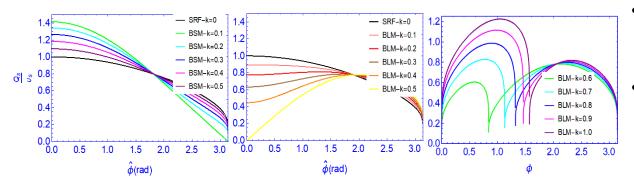
② Dual harmonic RF system (400MHz+800MHz)



The BLM is used much more often, because for the same voltage and harmonic ratios it gives :

- larger bucket area;
- smaller peak line density reduced space charge effects.
- larger synchrotron frequency spread

Synchrotron frequency distribution: n=2

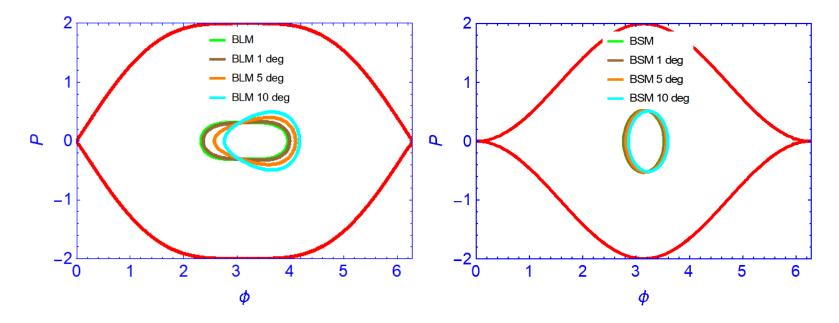


- In BL-mode region ω_s'(φ)=0
 exists for n=2 and~ k>0.2 →
 loss of Landau damping
- <Qs> is smaller for BLM, bad for TMCI(
 Qs).



② Dual harmonic RF system (400MHz+800MHz)

The effect of relative phase shift error



- BLM is extremely sensitive to the relative phase error between two rf system;
- BSM seems more reasonable.



③ Controlled emittance blowup

(1) Band limited RF phase noise: [12-14]

- injected into the main RF system through phase Loop;
- the choices of excitation noise spectrum: $0.85\Omega_{s0} \le \Omega \le 1.1\Omega_{s0}$
- Steady state between resonance and diffusion
- Example: LHC, SPS, PSB

(2) adding a phase-modulated high frequency RF to the main RF. [15-16]

• The high-frequency phase-modulated voltage :

 $V_H = \hat{V}_H \sin(h_H \omega_R t + \alpha \sin \omega_M t + \theta_H)$

- RF phase-modulated higher harmonic rf system drive bunch near resonant island and cause the bunch density redistribution.
- Example: PS;

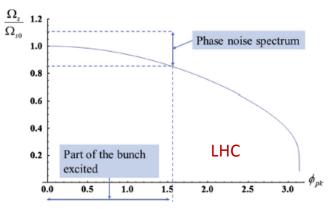


Fig. 2. Ω_s/Ω_{s0} as a function of the maximum phase deviation in radians. Stationary bucket.

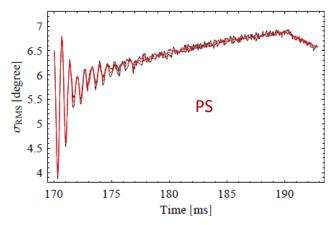


Figure 4: Measured RMS bunch length along the BU with $V_b = 28 \,\text{kV}$ at 200 MHz, $f_{\text{mod}} = 10 \,\text{kHz}$, $\Delta \phi = 3.1 \,\text{rad}$; black lines: individual cycle, red line: average over several cycles.



		Comparison among 400MHz, 800MHz, 400+800MHz	
		 Advantage: h increase → the frequency spread increase → instability improve; 	
	800 MH	• shorter bunch \rightarrow luminosity increased by 7%.	
L		• both the bucket area and bucket height are reduced \rightarrow beam loss	
Advantage: larger bucket area;		larger bucket area;	
	E L	 smaller peak line density reduced space charge effects. larger synchrotron frequency spread. Disadvantage: 	
400 MHz +800		 ω_s'(φ)=0 exists which cause loss of Landau damping average synchrotron tune reduced extremely sensitive to the relative phase error between two rf systemeters 	
MHz	z E S N	 Advantage: the average synchrotron tune and tune spread are increased. Disadvantage: the bucket area is reduced to some extent and momentum filling factor increased, but bunch length is shorter 	



Other questions need to be considered

Synchrotron radiation

- Energy loss per turn U0~1.48MeV,
- Longitudinal damping time: $\tau_{\epsilon} \sim 1.17h$ (physics running 14.2 h)
- The choice of controlled emittance growth
 - Like LHC or SPS, band-limited phase noise?
- Transverse Mode Coupling Instability (TMCI) [17]

$$\left(\beta_{av}Z_{T}\right)_{th}\propto\frac{E}{I_{b}}Q_{s}$$

- No impedance model
 - macroparticle tracking cannot proceed.



[1] Qing Qin, High Energy Circular Accelerator Physics, Chinese edition

[2] T. Argyropoulos, LONGITUDINAL MICROWAVE INSTABILITY IN A MULTI-RF SYSTEM, Proceedings of HB2014, East-Lansing, MI, USA

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 [4] F.J. Sacherer, A LONGITUDINAL STABILITY CRITERIONFOR BUNCHED BEAMS, IEEE Trans. Nuclear Sci. NS 20,3,825 (1973)

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[6] Ibon Santiago González, Loss of Longitudinal Landau Damping in the LHC Injectors,

[7] E. Shaposhinikova, Longitudinal beam parameters during acceleration in LHC, LHC Project Note 242

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[9] E. Shaposhinikova et, 52nd ICFA Advanced Beam Dynamics Workshop HB2012;

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Thanks for your attention!



