

**XXIII DAE-BRNS High Energy Physics Symposium 2018**

Probing new signature using Jet substructure at  
the LHC

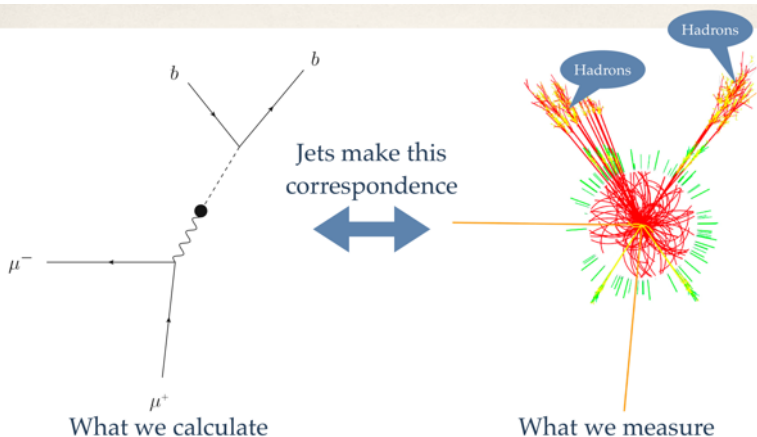
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Thalapillil**

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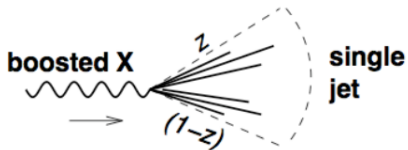
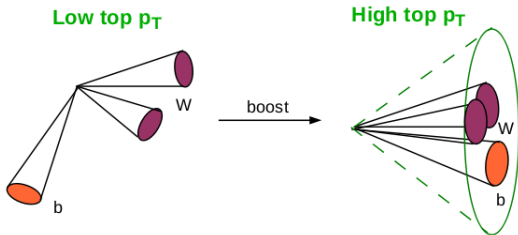
Dec 11, 2018

# What we calculate and What we see at the LHC?



$h \rightarrow b\bar{b}$  57% , Z and W boson decay almost 70% to hadronic channel.

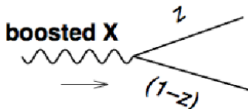
# Boosted topology



$$R \gtrsim \frac{m}{p_t} \frac{1}{\sqrt{z(1-z)}}$$

## How to look for Jet Substructure ....

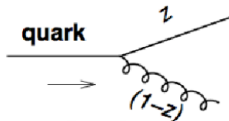
A possible approach for reducing the QCD background is to identify the two prongs of the heavy particle decay, and put a cut on their momentum fraction.



**Signal:**

$$P(z) \sim 1$$

Will split mainly  
**symmetrically**



**Background:**

$$P(z) \sim \frac{1+z^2}{1-z}$$

$$P(z) \sim \frac{1+(1-z)^2}{z}$$

Will split mainly  
**asymmetrically**

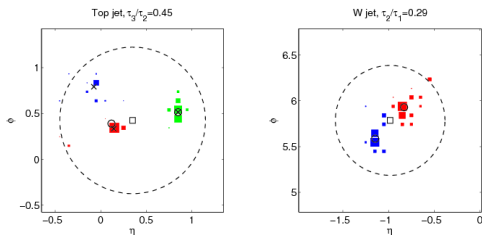
Selection of the symmetric splitting and identify two legs.

# Identifying the 2 or 3 prong nature of the Fat-jet

## N-subjettiness

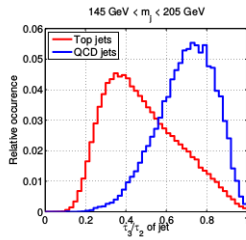
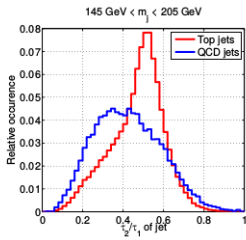
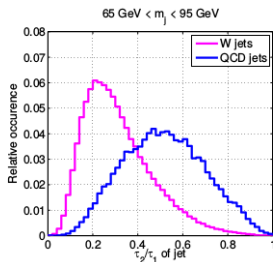
$$\tau_N^\beta = \frac{1}{\mathcal{N}_0} \sum_i p_{i,T} \min \left\{ \Delta R_{i1}^\beta, \Delta R_{i2}^\beta, \dots, \Delta R_{iN}^\beta \right\}.$$

$i$  runs over the constituent particles in a given jet.  
Where,  $\mathcal{N}_0 = \sum_i p_{i,T} R_0$  is the normalizing factor,  
 $R_0$  is the jet radius parameter.



# N-subjettiness Ratio

- Quantify how original jet seem to be composed of N daughter subjets
- $\tau_N \sim 0 \Rightarrow$  Original jet consist of N or fewer sub-jets
- $\tau_N \gg 0 \Rightarrow$  Large fraction of energy diluted from candidate N sub-jets
- A good discriminant the ratios of adjacent N-subjettiness values



## Inverse SeeSaw

The relevant part of the Lagrangian can be written as

$$\mathcal{L} \supset -Y_D^{\alpha\beta} \bar{\ell}_L^\alpha H N_R^\beta - M^{\alpha\beta} \bar{S}_L^\alpha N_R^\beta - \frac{1}{2} \mu_{\alpha\beta} \bar{S}_L^\alpha S_L^{\beta c} + \text{H.c.}$$

After the electroweak symmetry breaking we obtain the neutrino mass matrix as

$$M_\nu = \begin{pmatrix} 0 & M_D & 0 \\ M_D^T & 0 & M^T \\ 0 & M & \mu \end{pmatrix}.$$

we obtain the [inverse seesaw](#) formula for the light neutrinos as

$$M_\nu \simeq M_D M^{-1} \mu M^{-1T} M_D^T.$$

# Inverse SeeSaw : Two different cases

## Single flavor (SF)

- One heavy Dirac pair is at the electroweak scale, while other heavy pairs are assumed to be beyond the kinematic reach of the LHC.
- The lightest heavy Dirac neutrino mass eigenstates dominantly couple to a single lepton flavor *Assume  $\mu$*



# Inverse SeeSaw : Two different cases

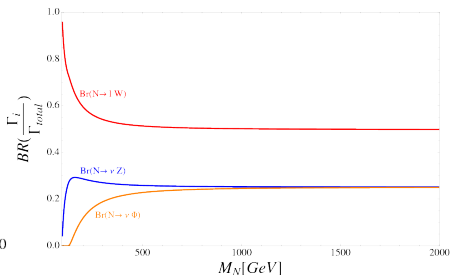
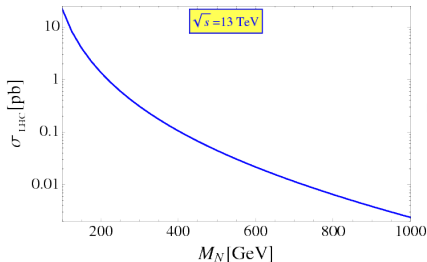
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## Flavor diagonal (FD)

- Two of the heavy Dirac neutrino pairs are degenerate with a equal mass  $m_N$ .
- Assuming that one pair dominantly couples to electrons, and the other one to muons with equal strength.

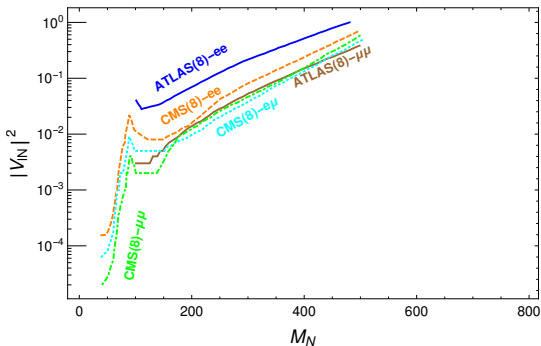
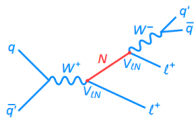
# Heavy Neutrino at LHC



Total production cross-section and branching ratios of heavy Majorana neutrino as a function of its mass at the LHC with  $\sqrt{s} = 13 \text{ TeV}$  and normalised by the  $|V_{\mu N}|^2$ .

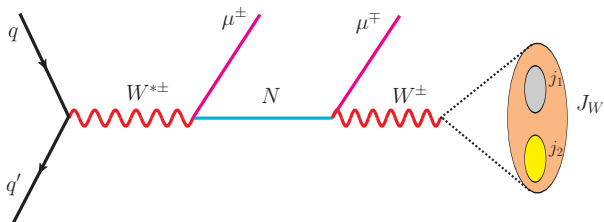
$$BR(N \rightarrow lW) : BR(N \rightarrow \nu Z) : BR(N \rightarrow \nu h) \simeq 2 : 1 : 1$$

# Conventional Experimental searches



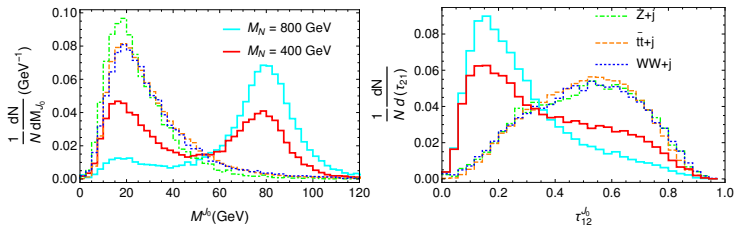
# Fatjet from Heavy Neutrino (OSDL)

$$pp \rightarrow l_1^+ N, N \rightarrow l_2^- W^+, W^+ \rightarrow J_W$$
$$pp \rightarrow l_1^- \bar{N}, \bar{N} \rightarrow l_2^+ W^-, W^- \rightarrow J_W$$



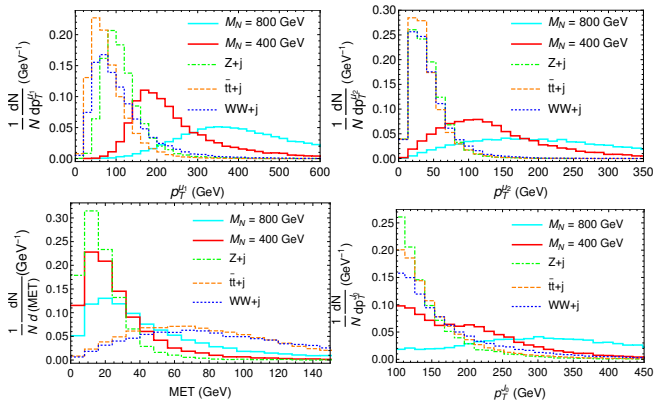
Fatjet signature will be seen if  $W^\pm$  is sufficiently boosted.

# Utilising Jetsubstructure for W-tagging



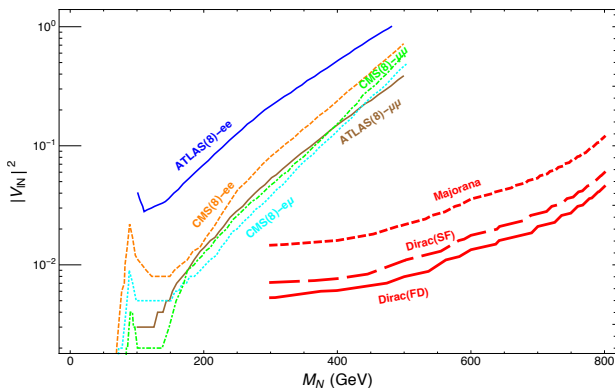
Normalised distribution invariant mass  $M_{J_0}$  (*left*) and N-subjettiness ratio  $\tau_{21}$  (*right*) of the leading fat-jet after the application of the basic selection cuts; including  $p_T^{J_0} > 100 \text{ GeV}$ ..

# Utilising Jetsubstructure for W-tagging



Normalised distribution as a function of observables after the application of the basic selection cuts; including  $p_T^{J_0} > 100 \text{ GeV}$ .

## Results and Discussion (OSDL)



Exclusion limit in terms of heavy neutrino mass  $M_N$  and  $|V_{\mu N}|^2$  for an integrated luminosity of  $3000 \text{ fb}^{-1}$  at the 13 TeV LHC with other available limits.

# Conclusion

- Jet Substructure technique: one of the most efficient and popular discovery tool in the search of new physics signatures at the LHC.
- As the available center-of-mass energy at the LHC increases, more and more boosted particles, understanding of the physics in this boosted regime is extremely crucial.
- One of the most active field of research in particle physics phenomenology !



## Conclusion

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THANK YOU

# Cut flow for OSDL + a fatjet

Cut	Signal	Background					
		$Z^l + j$	$t\bar{t} + j$	$W^l W^l + j$	$Z^l W^h + j$	$Z W^l + j$	$Z^l W^l W^h + j$
Pre-selection + $\mu^+ \mu^- + 1J$	621.69 [100%]	$5.1 \times 10^7$ [100%]	$5.6 \times 10^6$ [100%]	$2.2 \times 10^5$ [100%]	$5.7 \times 10^5$ [100%]	$2.9 \times 10^4$ [100%]	120.8 [100%]
$p_T(t_1) > 100$ GeV + $p_T(t_2) > 60$ GeV	532.1 [85.60%]	$8.9 \times 10^6$ [16.88%]	$6 \times 10^5$ [10.68%]	$3.6 \times 10^4$ [16.86%]	$10.2 \times 10^4$ [18.27%]	5556 [19.48%]	33.64 [27.85%]
$M_{\mu^+ \mu^-} > 200$ GeV	446.4 [71.82%]	583.7 [0.0010%]	$4.2 \times 10^5$ [7.37%]	$2.2 \times 10^4$ [9.87%]	9.4 [0.0016%]	1471.2 [5.15%]	10.42 [8.62%]
b-veto	414.6 [66.71%]	530.6 [9.9 $\times 10^{-4}$ %]	$14.8 \times 10^4$ [2.67%]	$2.0 \times 10^4$ [9.52%]	8.54 [0.0014%]	1384.5 [4.85%]	9.15 [7.58%]
MET < 60	369.7 [59.49%]	371.4 [6.9 $\times 10^{-4}$ %]	$5.4 \times 10^4$ [0.982%]	8104 [3.70%]	6.0 [10.3 $\times 10^{-4}$ %]	751.41 [2.63%]	4.04 [3.35%]
$p_T^J > 150$ GeV	240.3 [38.65%]	265.32 [4.9 $\times 10^{-4}$ %]	19606 [0.351%]	4032 [1.842%]	4.27 [7.3 $\times 10^{-4}$ %]	419.9 [1.47%]	2.75 [2.28%]
$M^{J0} > 50$ GeV	171.2 [27.54%]	42.4 [7 $\times 10^{-5}$ %]	3560 [0.063%]	887.4 [0.40%]	0.68 [1.2 $\times 10^{-4}$ %]	105.7 [0.37%]	1.48 [1.233%]
$\tau_{21} < 0.4$	152.3 [24.50%]	21.18 [3.9 $\times 10^{-5}$ %]	2062 [0.036%]	433.8 [0.198%]	0.34 [5.8 $\times 10^{-5}$ %]	71.03 [0.24%]	1.13 [0.94%]