Study of phase transition in two flavour quark matter at finite volume

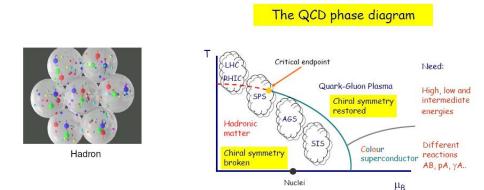
Tamal K Mukherjee

Adamas University

Collaborator: Anirban Lahiri, Rajarshi Ray

- Context of the work
- Model Set up
- Result
- Summary.

QCD Phase Diagram



- Presence of Critical End Point (CEP)!
- CEP is important: indirect evidence of 1st order phase transition and crossover.

(a)

Where to look for the transition



Relativistic Heavy Ion Collider [RHIC]



Large Hadron Collider [LHC]

(日) (周) (日) (日)

Where to look for the transition



- Only one control parameter CM energy \sqrt{S}_{NN}
- Important: fluctuations of the conserved quantities
- Fluctuations can be used to connect the experimental observation to theoretical predictions

Theory side

- From theory side, thermodynamics quantity of interest: susceptibility
- nth order cumulant of the fluctuations of the baryon number is related to the nth order baryon number susceptibility:

$$[B^n] = V T^3 T^{n-4} \chi_B^{(n)}$$

where V is the volume of the observed part of the fireball

- Can't employ above relation directly: V is hard to determine from experiment!!!
- Ratios are useful: can be used to compare between experimental results with that of theory

$$S\sigma = \frac{T \ \chi_B^{(3)}}{\chi_B^{(2)}}, \ \kappa \sigma^2 = \frac{T^2 \ \chi_B^{(4)}}{\chi_B^{(2)}}, \ \frac{\kappa \ \sigma}{S} = \frac{T \ \chi_B^{(4)}}{\chi_B^{(3)}}$$

 Applicability of thermodynamics → See: S. Gupta et. al., Science, 332, 2011; R.V. Gavai, S. Gupta, Phys. Lett. B, 696, 2011

Model: MRE approximation

- The finite-size effects are taken into account by using the Multiple Reflection Expansion (MRE). (for details please see: R. Balian and C. Bloch, Ann. Phys. (N.Y.) **60**, 401 (1970))
- MRE approximation: the finite-size effects are included by modifying the density of states:

$$p_{MRE}(k,\kappa,R) = \frac{k^2}{2\pi^2} \left[1 + \frac{6\pi^2}{kR} f_s\left(\frac{k}{\kappa}\right) + \frac{12\pi^2}{(kR)^2} f_C\left(\frac{k}{\kappa}\right) \right]$$

where R is the radius of the sphere and the functions:

$$f_s(x) = -\frac{1}{8\pi} \left(1 - \frac{2}{\pi} \arctan x\right), \ f_C(x) = \frac{1}{12\pi^2} \left[1 - \frac{3}{2}x\left(\frac{\pi}{2} - \arctan x\right)\right]$$

the subscript "s" and "C" correspond surface and curvature contribution respectively.

• Note: curvature contribution has not been derived within the MRE framework. The form of the f_C is a proposal from J. Madsen, Phys. Rev. D **50**, 3328 (1994).

Model: PNJL model

Lagrangian:

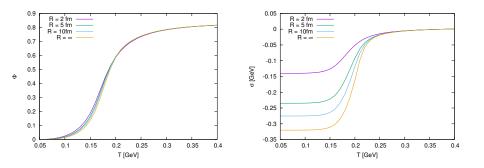
$$\mathcal{L}_{PNJL} = \bar{q}(i\gamma^{\mu}D_{\mu} - m_0 + \mu\gamma^0)q + \frac{G}{2}[(\bar{q}q)^2 + (\bar{q}i\gamma^5\vec{\tau}q)^2] - V(\bar{\Phi},\Phi)$$

Where $D_{\mu} = \partial_{\mu} - ig \mathcal{G}_{\mu}$ and $\mathcal{G}_{\mu} = \delta_{\mu 0} \mathcal{G}_{0}$

- Introducing auxiliary field variables σ and $\vec{\pi}$ an \mathcal{L}_{eff} is obtained, with the replacement $\exp[-\mathcal{G}_0/T] \rightarrow \Phi$.
- The mean fields $<\sigma>= G < \bar{q}q >$ and $<\vec{\pi}>= 0$ for $\mu_I < m_{\pi}$.
- Thermodynamic properties studied with Φ(T) and σ from the thermodynamic potential Ω[Φ, Φ, σ, T, μ₀, μ_I].

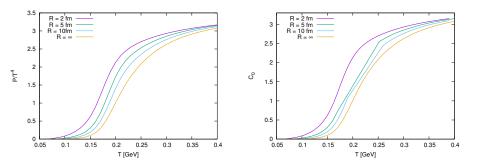
A (10) A (2) A

Order Parameters



- Order Parameter for Chiral Transition is affected by finite size!
- At a critical size of the system value of Chiral Condensate will always be zero!

Pressure



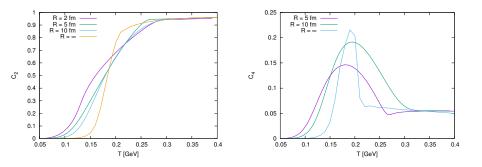
• Approaches ideal gas value at higher temperature.

•
$$\frac{P(T,\mu_q)}{T^4} = \sum_{n=0}^{\infty} C_n(T) \left(\frac{\mu_q}{T}\right)^n$$
, where $C_n(T) = \frac{1}{n!} \frac{\partial^n \left(\frac{P}{T^4}\right)}{\partial \left(\frac{\mu_q}{T}\right)^n}$

3

(日) (同) (三) (三)

Fluctuation



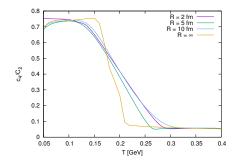
•
$$\frac{P(T,\mu_q)}{T^4} = \sum_{n=0}^{\infty} C_n(T) \left(\frac{\mu_q}{T}\right)^n$$
, where $C_n(T) = \frac{1}{n!} \frac{\partial^n \left(\frac{P}{T^4}\right)}{\partial \left(\frac{\mu_q}{T}\right)^n}$

- Nature of the transition does not change with the finite size
- Peak height of C_4 decreases with the decreasing value of the radius.

< (T) >

• Results need to be improved around the transition temperature.

Fluctuation



•
$$\frac{P(T,\mu_q)}{T^4} = \sum_{n=0}^{\infty} C_n(T) \left(\frac{\mu_q}{T}\right)^n$$
, where $C_n(T) = \frac{1}{n!} \frac{\partial^n \left(\frac{P}{T^4}\right)}{\partial \left(\frac{\mu_q}{T}\right)^n}$

- Distinct values in two different phases!
- Effects of volume cancels out (more or less)!

.∋...>

< 一型

Summary

- We propose to study the transition in the 2-flavour quark matter by using Multiple Reflection Expansion along with PNJL model.
- Qualitatively the model reproduces the known expectations.
- Fluctuations near the transition region needs to be addressed more carefully.
- Ongoing work: to study the fluctuations as the system approaches the critical point from crossover and first-order phase transition sides respectively.

- 4 回 ト - 4 回 ト

THANK YOU!

э