

# Recent Directions in Neutrino Theory and Phenomenology

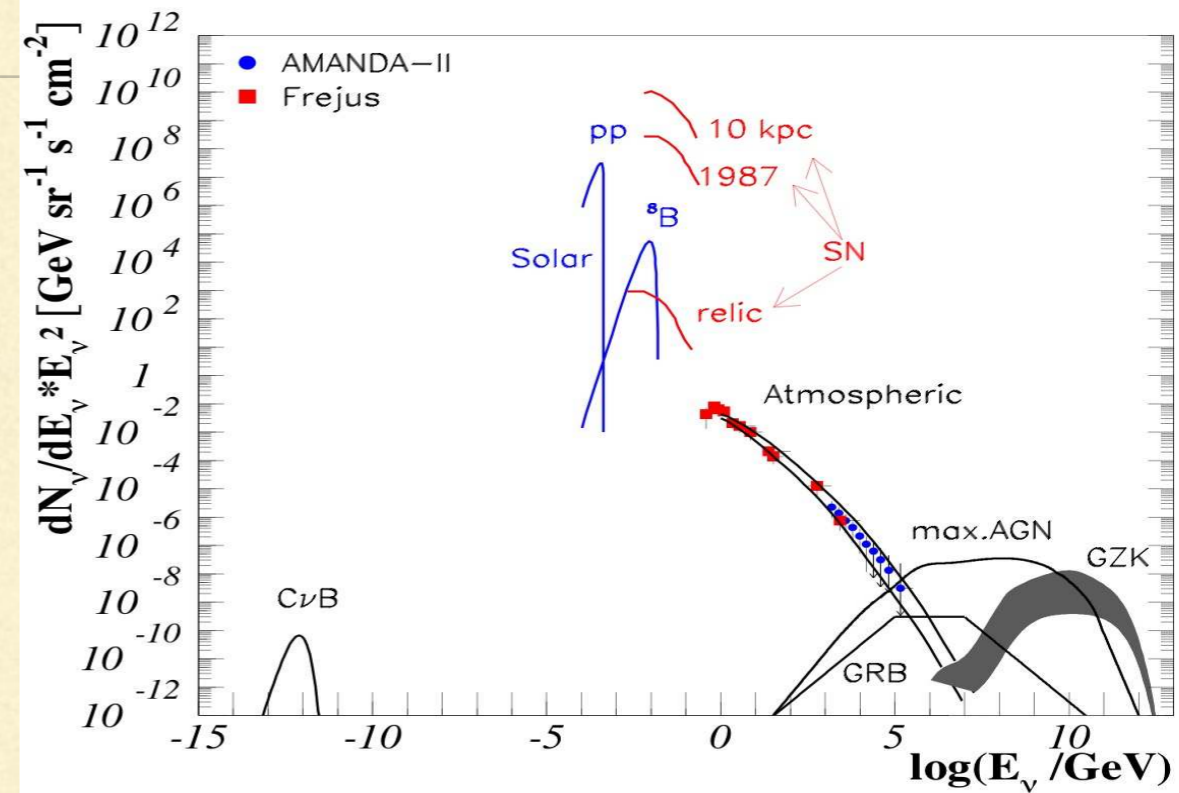
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Neutrino Physics is an extensive field, with experiments spanning a very wide range in energy and sources. Similarly, theoretical and phenomenological areas of study span a wide range



Impossible to review comprehensively all active theoretical areas

Focus (for most part) on topics where there is recent theoretical interest and new ideas motivated by recent results of experiments.

Given the number of young people in the audience, also guided by the question:  
What are good questions to work on now?

Choose 4 topics of current theoretical/phenomenological activity, 2 in low energy neutrino physics and two in high energy neutrino physics.

Choose signals/ issues which cannot be easily explained away or dismissed.



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▶ The MiniBooNE excess and related theoretical developments

▶ The Dirac vs Majorana nature of neutrinos

▶ The IceCube events and Dark matter

▶ The ANITA observations of the highest energy events

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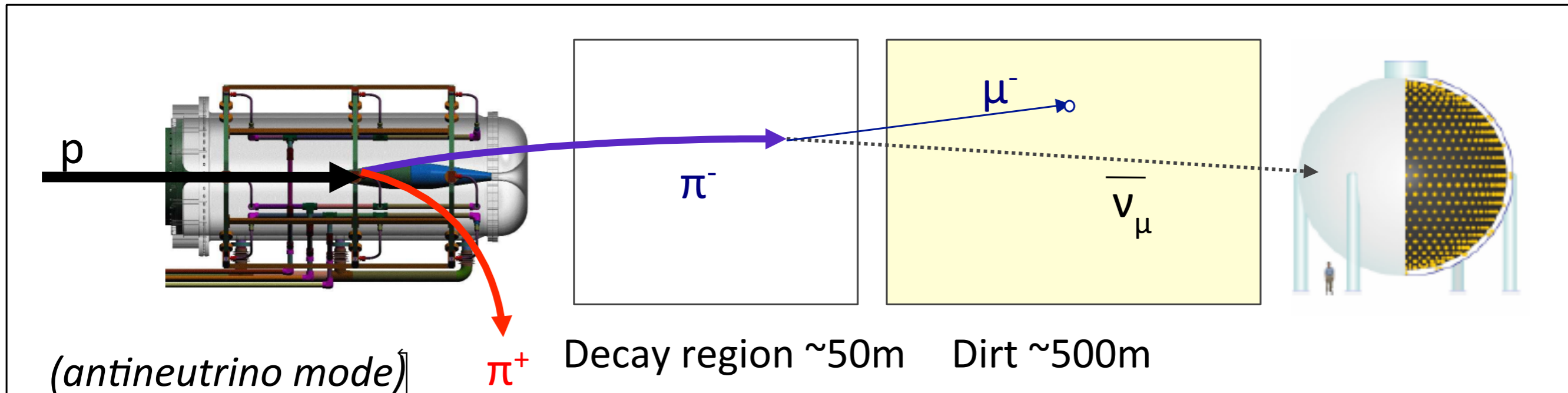
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▶ The MiniBooNE excess and related theoretical developments

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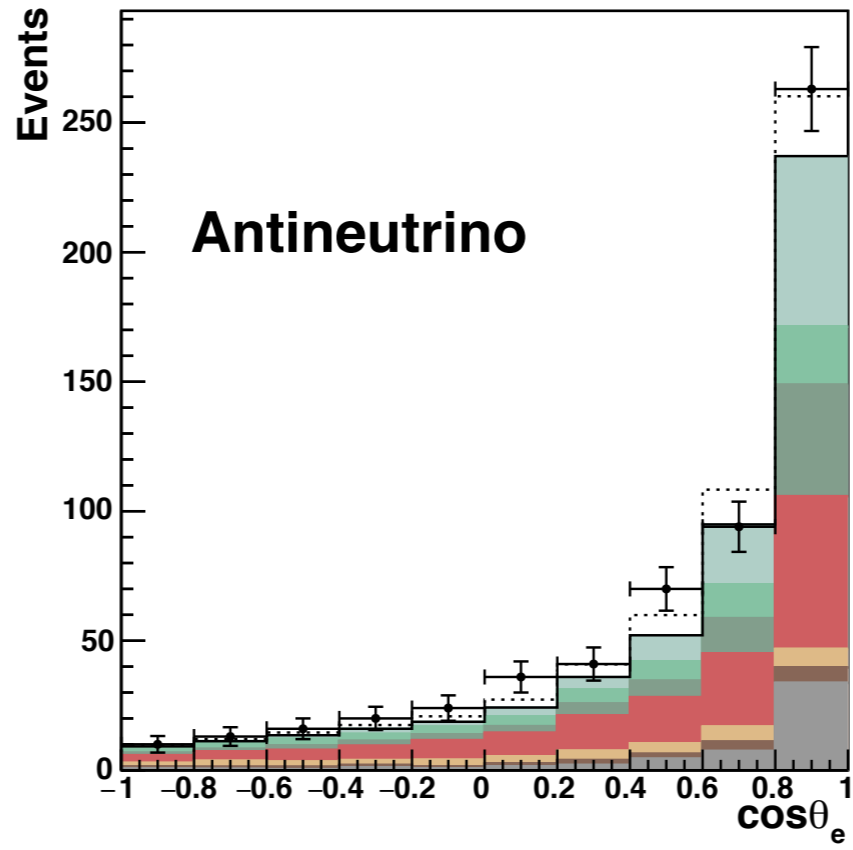
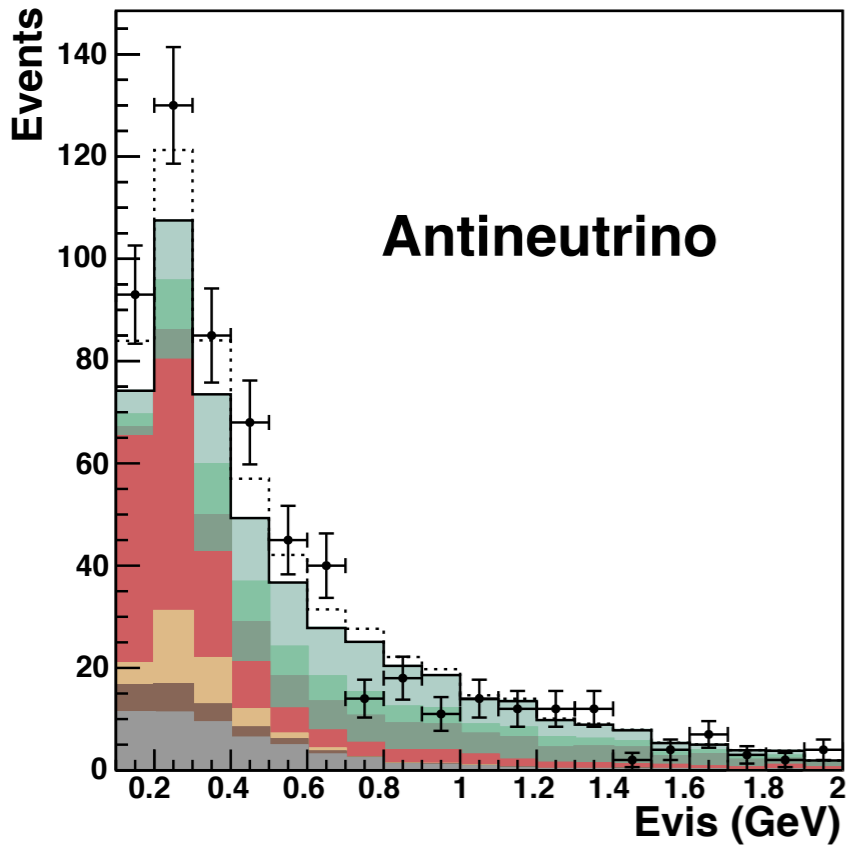
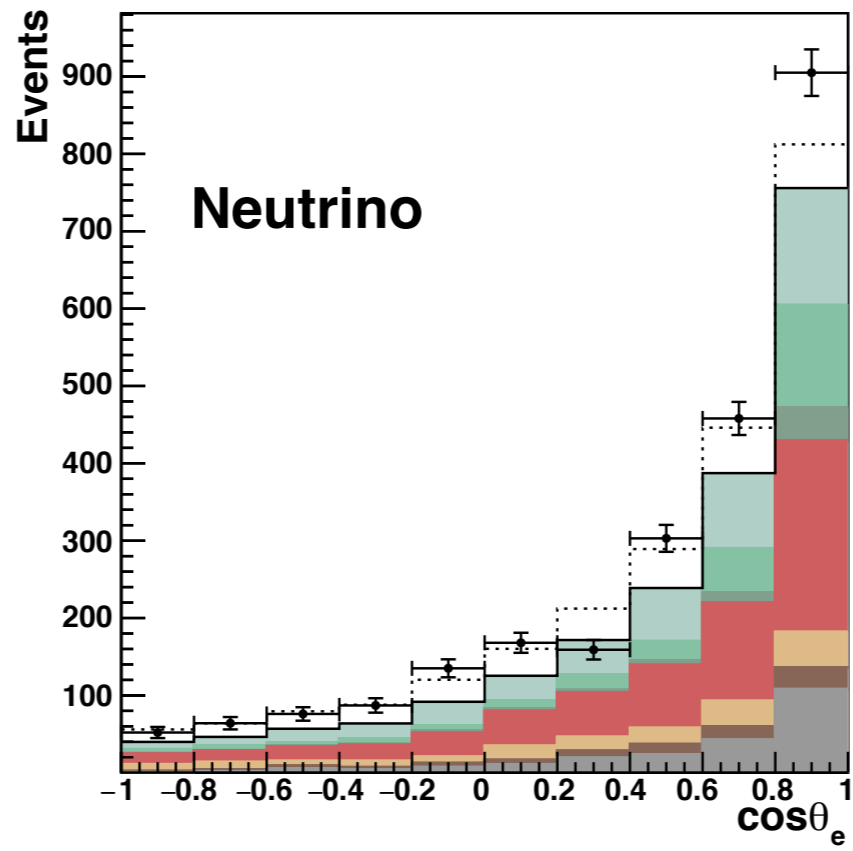
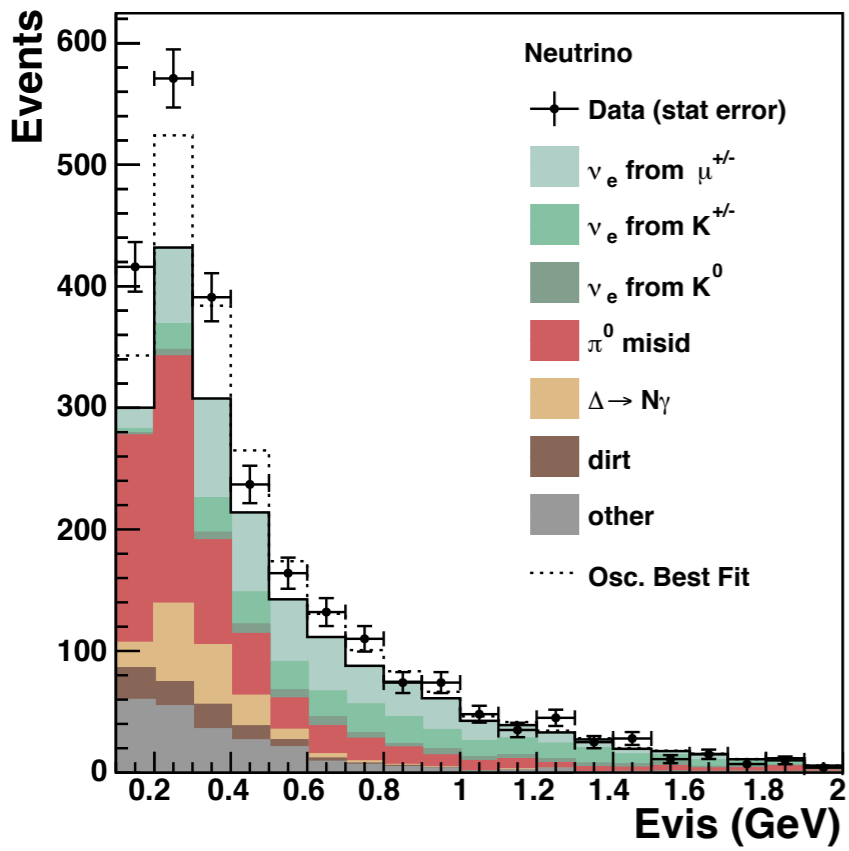
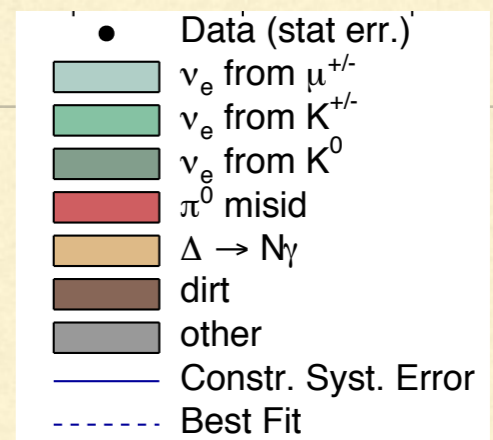
# MiniBooNE Experiment



- Similar L/E as LSND for  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  &  $\nu_\mu \rightarrow \nu_e$  oscillations
  - MiniBooNE  $\sim 500\text{m}/\sim 500\text{MeV}$
  - LSND  $\sim 30\text{m}/\sim 30\text{MeV}$
- Horn focused neutrino beam ( $p+\text{Be}$ )
  - Horn polarity  $\rightarrow$  neutrino or anti-neutrino mode
- 800t mineral oil Cherenkov detector



# The MiniBooNE excess.....



Excesses in neutrino and antineutrino mode qualitatively consistent

MiniBooNE confirms LSND excess at  $4.7\sigma$ , and their combined significance is  $6.0\sigma$

MiniBooNE unable to tell if excess is electrons or photons.

MicroBooNE will help resolve this question. SBN program will resolve whether excess due to sterile neutrino oscillations.

# The MiniBooNE excess.....

Possible

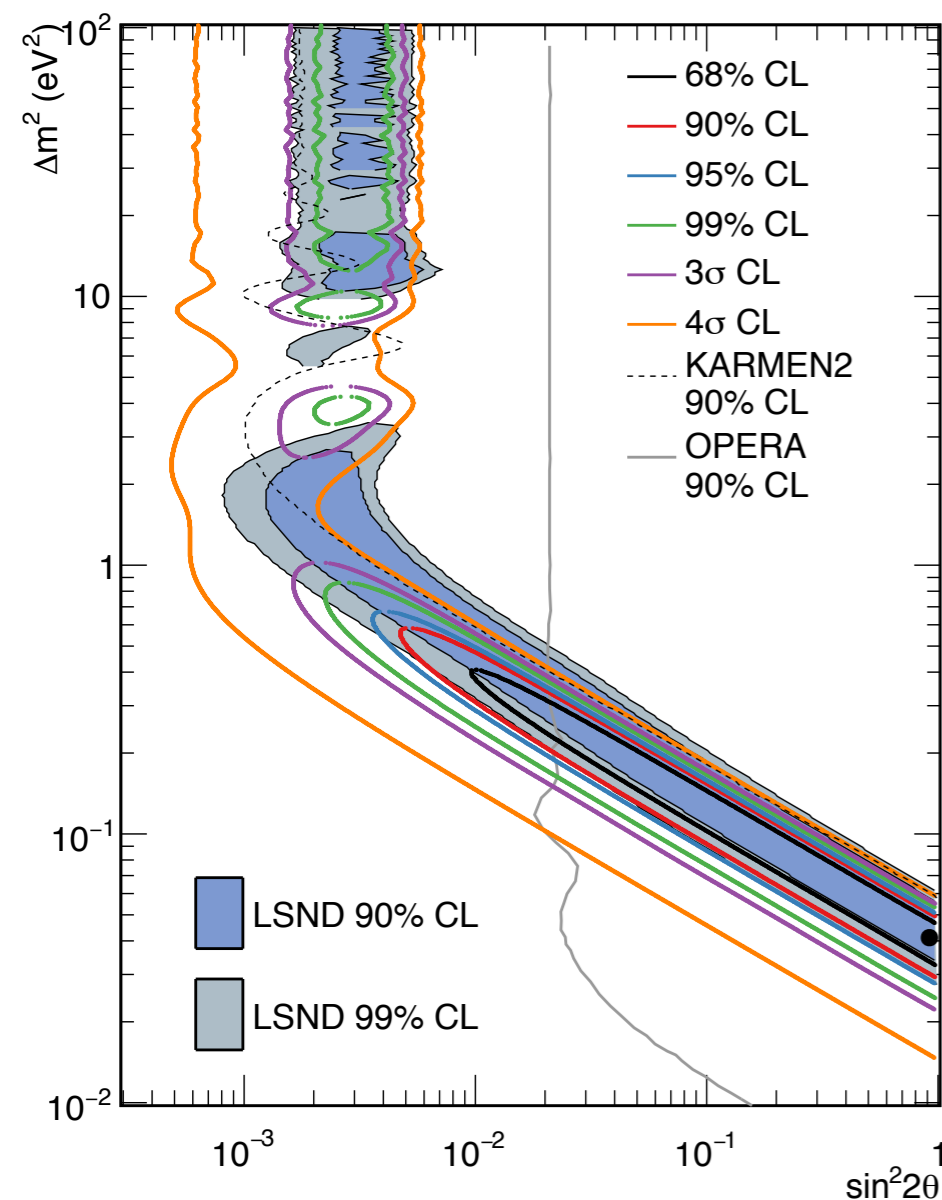
explanation: 3+1  
sterile neutrino

Problem: Best fit  
ruled out by other  
experiments.

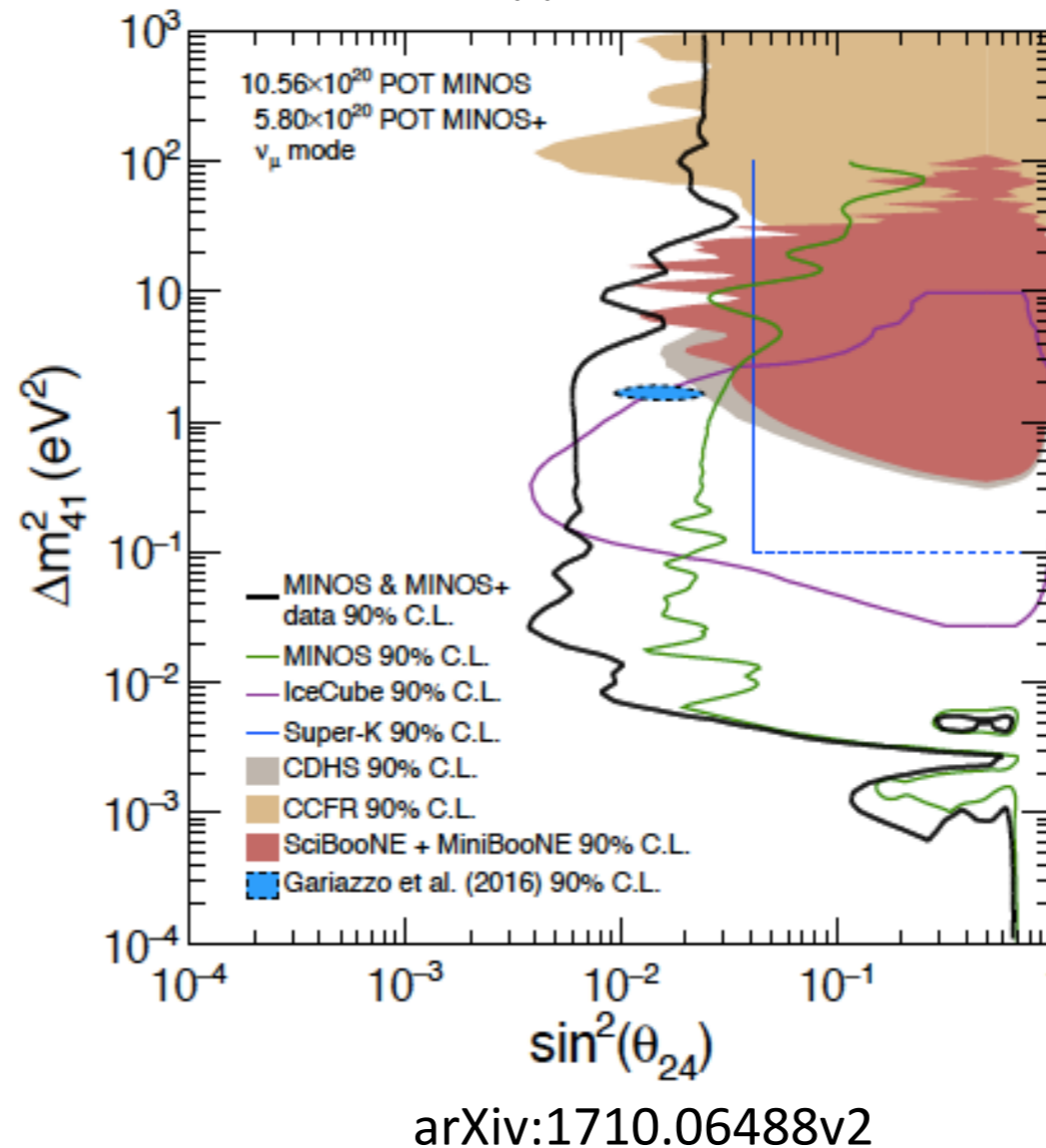
Problem: Solution  
implies  $\nu_\mu$   
disappearance in  
addition to  $\nu_e$   
appearance, but  
this is ruled out by  
IceCube and  
MINOS/MINOS+

Problems not  
mitigated by going  
to 3+n sterile  
neutrinos

Appearance



Disappearance



MiniBooNE backgrounds have been well-checked and measured for the most part.

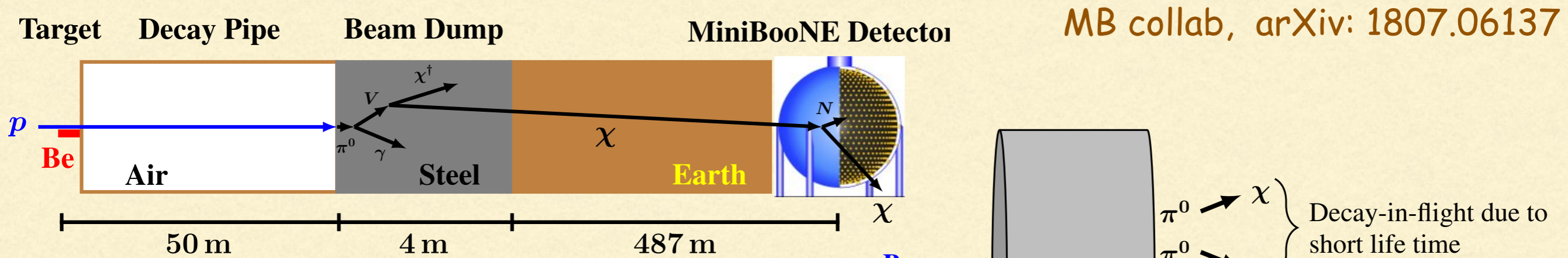
All of this may be pointing to a non-oscillation new physics explanation



# The MiniBooNE excess.....

Before considering an example of new physics, it is important to note that this is a very constrained situation, and any new physics explanation must satisfy many conditions.

As part of a dark matter search, MiniBooNE did an off target run, where beam hits the dump.



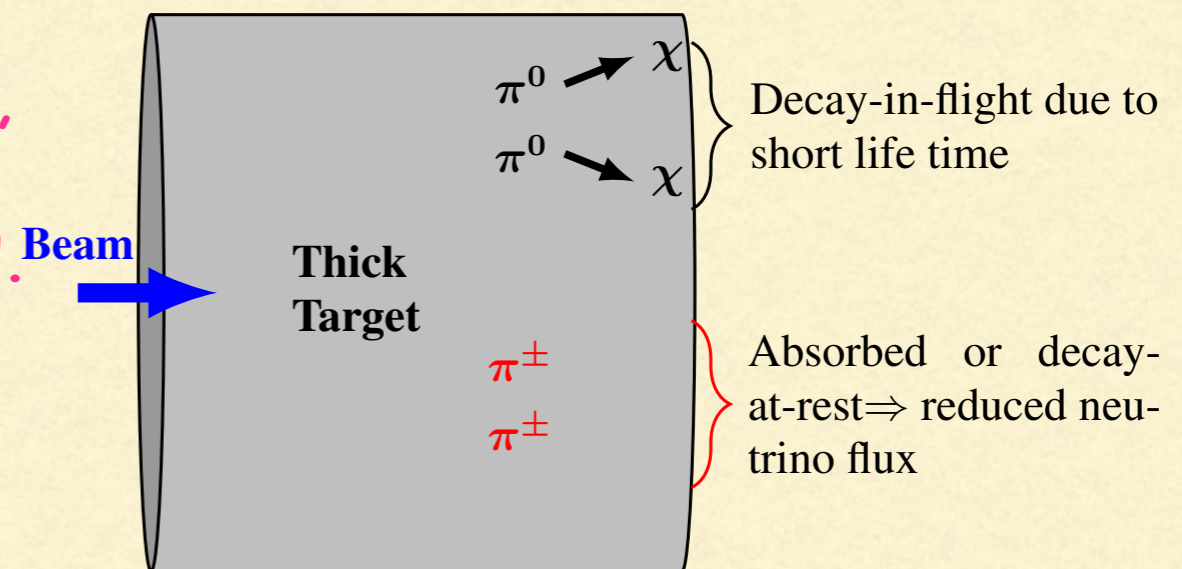
MB collab, arXiv: 1807.06137

Importantly, the excess disappeared when this was done.

Thus, excess cannot be due to new particle (e.g. DM,  $\chi$ ) produced in dump/target and scattering off electrons, via a portal.

Excess cannot be due to photons, for instance, which are produced via decay of new particle, since background of entering photons measured. Also,  $\pi^0$  to  $2\gamma$  background well measured.

Thus, production of neutrinos, present in target case and absent/reduced in off target, must be playing a role in the "true" explanation





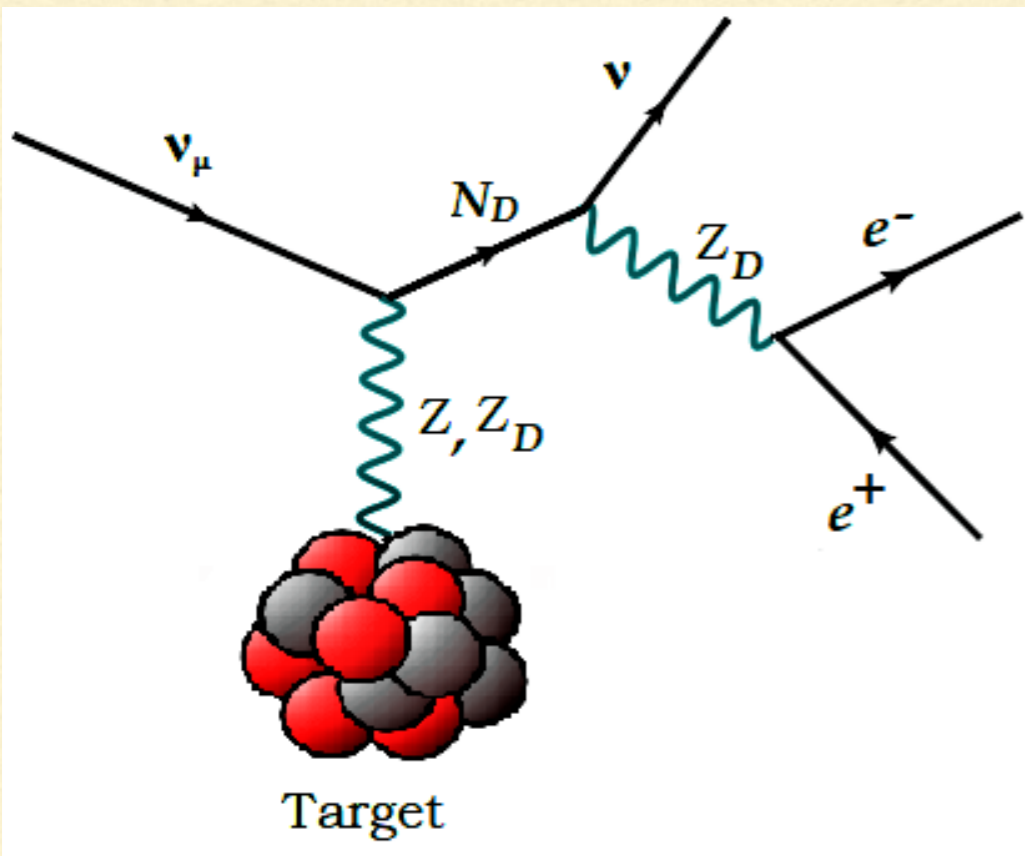
# The MiniBooNE excess.....

In addition, any new physics explanation must reproduce both the energy and angular distribution seen in the data.

Example of possible new physics explanation :

Bertuzzo et al arXiv: 1807.09877;

Ballett et al arXiv: 1808.02915;



Introduce a dark sector composed by a new vector boson,  $Z_D$ , coupling directly solely to a dark neutrino,  $\nu_D$

$$\nu_\alpha = \sum_{i=1}^3 U_{\alpha i} \nu_i + U_{\alpha 4} N_D, \quad \alpha = e, \mu, \tau, D,$$

$$\mathcal{L}_D \supset \frac{m_{Z_D}^2}{2} Z_{D\mu} Z_D^\mu + g_D Z_D^\mu \bar{\nu}_D \gamma_\mu \nu_D + e\epsilon Z_D^\mu J_\mu^{\text{em}} + \frac{g}{c_W} \epsilon' Z_D^\mu J_\mu^Z,$$

kinetic mixing
mass mixing

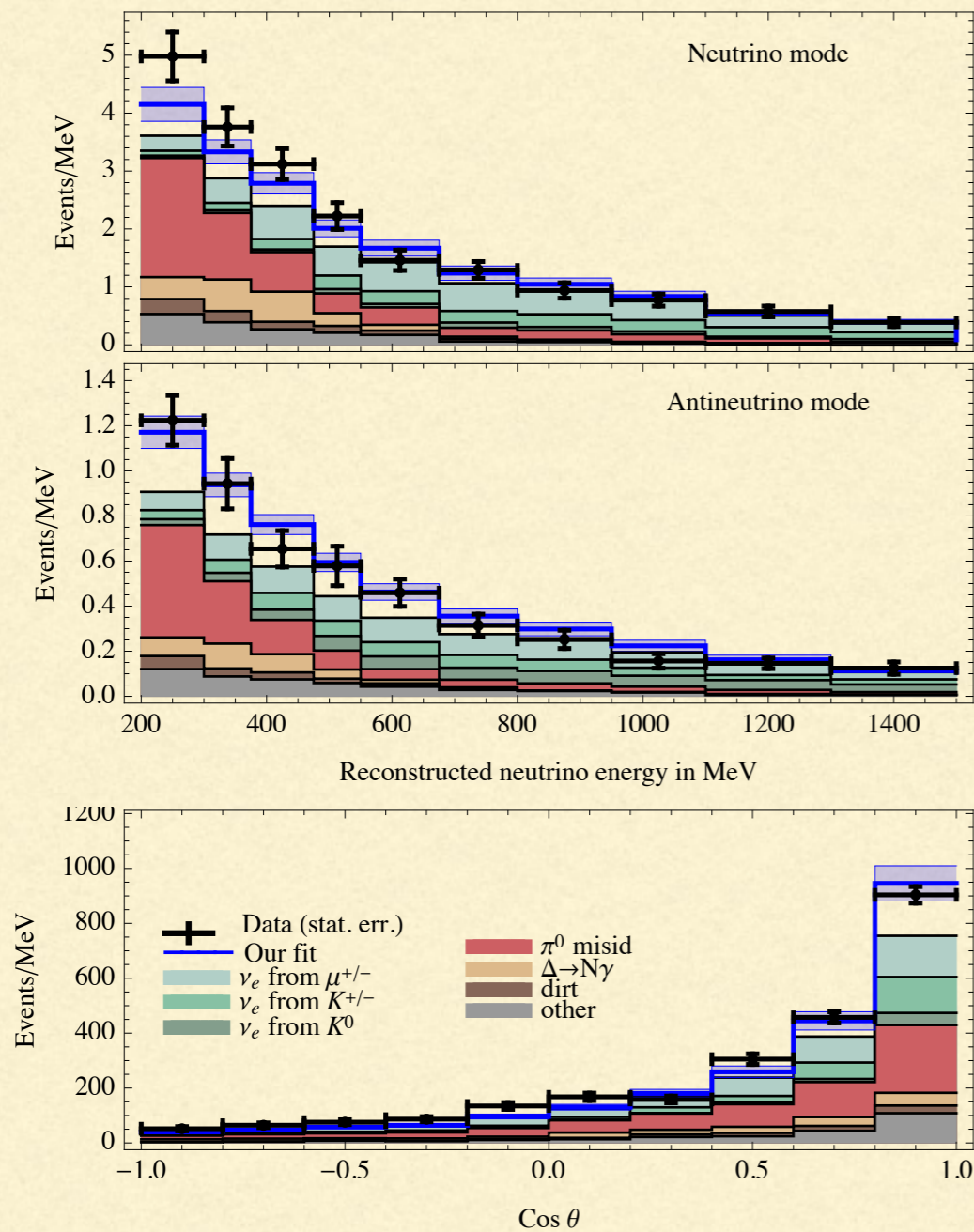
$$m_{N_D} > m_{Z_D}; \quad m_{Z_D} < 2m_\mu$$

Kinetic mixing: from  $B^{\mu\nu}X_{\mu\nu}$  term, which is gauge invariant

Kinematics is such that many of the  $e^+e^-$  pairs will be collimated, and Mini-BooNE would interpret  $Z_D \rightarrow e^+e^-$  decays as electron-like events



Good agreement with observed energy and angular distribution :



Questions open for exploration:

Is this the right solution? How does it fit into the larger framework of BSM physics?

What are the ways in which it can be tested?

Would signals of this have been already seen in existing detectors? What will planned/ upcoming neutrino detectors see if this is true?



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▶ The Dirac vs Majorana nature of neutrinos

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# The Dirac vs Majorana nature of neutrinos.....Dirac and Majorana mass terms

This remains one of the most important unanswered questions in neutrino physics.

Majorana neutrinos offer, in principle, a window to physics at very high scales, and thus an opportunity to better understand what lies beyond the Standard Model (SM).

For Dirac neutrinos, if we add RH neutrinos to the SM, the Yukawa term

$$L_Y = -y H^0 \bar{\nu}_R \nu_L + \text{h.c.}$$

$$\nu = \nu_L + \nu_R$$

leads, after electroweak symmetry breaking, to the mass term

$$L_D = -m_D \bar{\nu}_R \nu_L + \text{h.c.}, = -m_D \bar{\nu} \nu,$$

only minimal extension  
of SM necessary

For Majorana neutrinos, however, 2 types of terms are possible, both connected in different ways to possible BSM/high scale physics

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# The Dirac vs Majorana nature of neutrinos.....Dirac and Majorana mass terms

$$\mathbf{L}_R = -m_R/2 (\bar{\nu}_R)^c \nu_R + \text{h.c.},$$

$$\nu = \nu_R + (\nu_R)^c.$$

breaks total lepton number

Since no other conserved quantum number carried by  $\nu_R$ ,  $m_R$  can be very large in principle.

Even if there are no RH neutrinos, physics at high scales ( $\Lambda$ ) can induce a LH majorana mass term via the effective operator, via an interaction with the SM Higgs field,

$$(\bar{\nu}_L)^c \mathbf{H}^0 \mathbf{H}^0 \nu_L / \Lambda$$

Or, if there is a weak isospin (BSM) Higgs triplet,  $\Delta$ , it can induce a similar term via

$\Delta^0 (\bar{\nu}_L)^c \nu_L$ , where  $\Delta^0$  is the neutral member of the triplet which acquires a vev

In either case, the mass term has the form

$$\mathbf{L}_L = -1/2 m_L (\bar{\nu}_L)^c \nu_L + \text{h.c.},$$

LH Majorana mass term

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# The Dirac vs Majorana nature of neutrinos.....Distinguishing between Dirac

This is very hard to do.....Why?

and Majorana neutrinos experimentally

Chirality and helicity

For any fermionic field,  $\Psi_L \simeq \Psi_- + m/E \Psi_+$  and  $\Psi_R \simeq \Psi_+ + m/E \Psi_-$ .

Thus, for a relativistic fermion, chirality and helicity are almost identical.

Consider  $\pi^+ \rightarrow \mu^+ + \nu_\mu$  and  $\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$  in the Dirac case,

LH and essentially -ve helicity

RH and essentially +ve helicity

$$\mathcal{L}_{CC} \propto \bar{\mu} \gamma^\lambda \frac{(1 - \gamma_5)}{2} \nu_\mu J_\lambda + \bar{\nu}_\mu \gamma^\lambda \frac{(1 - \gamma_5)}{2} \mu J_\lambda^\dagger ,$$

In the case when neutrinos are Dirac, only one term contributes to each of the two decays, since neutrinos and antineutrinos are distinct particles



# The Dirac vs Majorana nature of neutrinos.....Distinguishing between Dirac and Majorana neutrinos experimentally

In the majorana case, neutrinos and antineutrinos are one and the same particle, hence both terms can contribute. However, contribution of the second term is severely helicity suppressed, by a factor  $m/E$ .

Thus, since in almost all circumstances, neutrinos are ultra-relativistic, whether neutrinos are majorana or Dirac makes no practical difference in an experiment.

However, if the neutrino is non-relativistic, from  $\Psi_L \approx \Psi_- + m/E \Psi_+$  and  $\Psi_R \approx \Psi_+ + m/E \Psi_-$ .

$$\mathcal{L}_{CC} \propto \bar{e} \gamma^\lambda \frac{(1 - \gamma_5)}{2} \nu_e J_\lambda + \bar{\nu}_e \gamma^\lambda \frac{(1 - \gamma_5)}{2} e J_\lambda^\dagger .$$

Now, in the Majorana case, the contribution from the "second" term is no longer small, since each chirality is a mix of both helicities.

In the decay of a heavy, sterile neutrino, this leads to different energy and angular distributions for the daughters in the Dirac and Majorana cases.



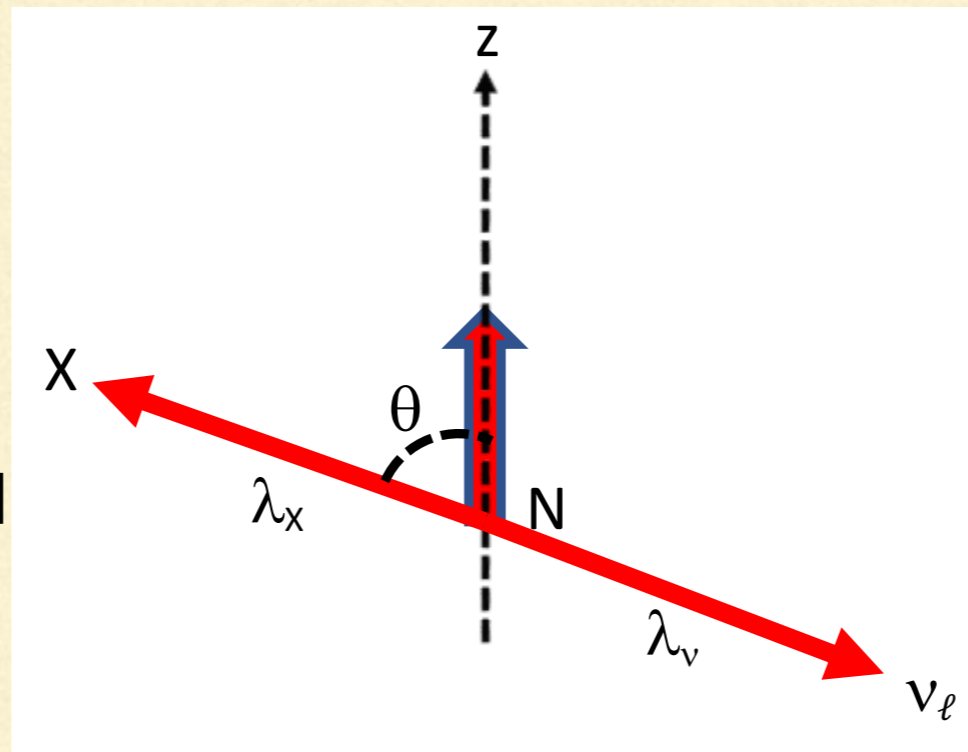
# The Dirac vs Majorana nature of neutrinos.....

Distinguishing between Dirac and Majorana neutrinos experimentally. Consider, in the parent's rest frame, the decay  $N \rightarrow \nu_i + X$  of a heavy neutrino  $N$  that is fully polarized by its production mechanism, with its spin pointing along  $+z$ .  $X$  is a self-conjugate boson and  $\nu_i$  is a SM neutrino

$X$  and emerges at an angle  $\theta$  with respect to the  $+z$  direction (with  $\nu_i$  emerging oppositely), with helicities  $\lambda_X$ , and  $\lambda_\nu$ , respectively. With  $\lambda \equiv \lambda_X - \lambda_\nu$ , rotational invariance dictates that the

$\Gamma \equiv$  Total decay rate

$$\alpha = (\Gamma_{\lambda=+1/2} - \Gamma_{\lambda=-1/2})/\Gamma \in [-1, +1]$$



Example:  $N \rightarrow \nu_i \pi^0$ .  $\nu_i$  is always -ve helicity, so  $\lambda = \lambda_X - \lambda_\nu = 1/2$ ; and  $\alpha = 1$ . Thus this decay will have angular distribution of the form  $(1 + \cos\theta)$

$$\frac{d\Gamma(N \rightarrow \nu + X)}{d(\cos \theta)} = \frac{\Gamma_{\lambda=+1/2}}{2} (1 + \cos \theta) + \frac{\Gamma_{\lambda=-1/2}}{2} (1 - \cos \theta)$$

$$= \frac{\Gamma_0}{2} (1 + \alpha \cos \theta); \quad -1 \leq \alpha \leq +1$$

$$\Gamma_0 = \Gamma_{\lambda=+1/2} + \Gamma_{\lambda=-1/2} > 0$$



From CPT and rotational invariance, it can be shown that for Majorana neutrinos,  $\alpha=0$ , whereas for Dirac neutrinos

$$\frac{d\Gamma(N \rightarrow \nu + X)}{d(\cos \theta)} = \frac{\Gamma_0}{2} (1 + \alpha \cos \theta)$$

$\mu$  and  $d$  are the magnetic and electric transition dipole moments

X	$\gamma$	$\pi^0$	$\rho^0$	$Z^0$	$H^0$
$\alpha$	$\frac{2\Im m(\mu d^*)}{ \mu ^2 +  d ^2}$	1	$\frac{m_N^2 - 2m_\rho^2}{m_N^2 + 2m_\rho^2}$	$\frac{m_N^2 - 2m_Z^2}{m_N^2 + 2m_Z^2}$	1

Thus, if mass of N is known, the angular distribution can in principle determine if SM neutrinos are Majorana or Dirac, because even if one neutrino that mixes with the others is Majorana, all of them are majorana.

This requires experiments to look for heavy sterile neutrinos, and possibly measure their decays.



# The Dirac vs Majorana nature of neutrinos.....Detection and Challenges

There are several experiments planning to or looking for heavy sterile neutrinos, some of them are MicroBooNE, SHiP, DUNE, NA48/2, and NA62

What are some of the possible challenges in such a program to determine the Dirac or Majorana nature of neutrinos?

▶ Such a neutrino should exist! (Models?) (More work and investigation needed!) ↓

▶ They must be produced in a sufficient number (e.g. say by meson decays) and be massive enough (few hundred MeV or more) to decay quickly in the detector and give a statistically significant number of events.

▶ The sample must be polarized. If produced in a weak decay, this is the case.

▶ If there is a charged lepton in decay final state, its charge identification needs to be made to get supplementary information on lepton number violation. This is often not possible in many neutrino experiments, e.g, Super-K, Hyper-K, NovA, DUNE, etc

▶ In order to do an angular distribution analysis in the rest frame, the momentum of  $N$  in the lab frame must be accurately reconstructed. This can be difficult if there is a neutrino in the final state, whose momentum cannot be directly measured.

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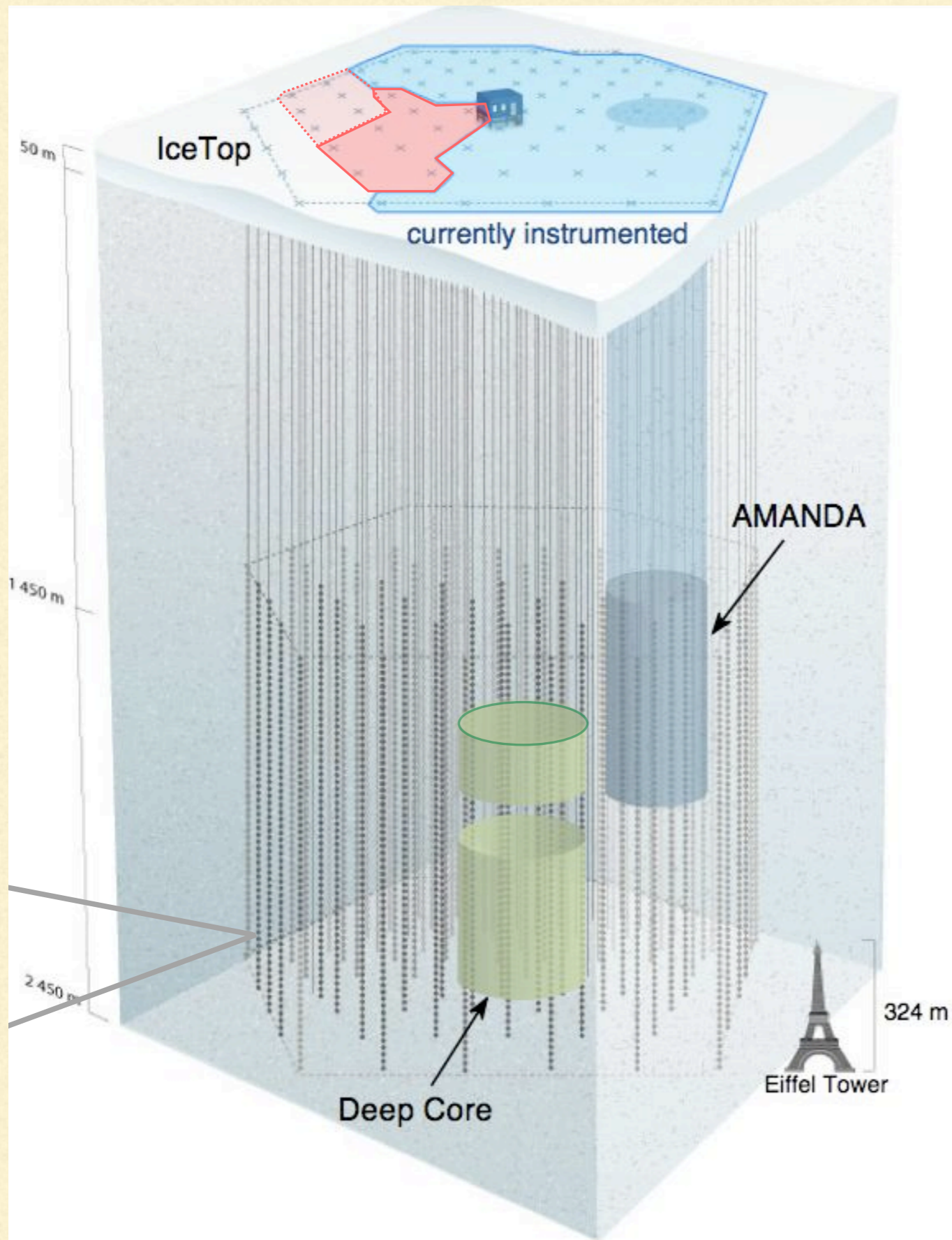
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► The IceCube events and Dark matter

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# The IceCube Detector



86 strings, 60 OM/string

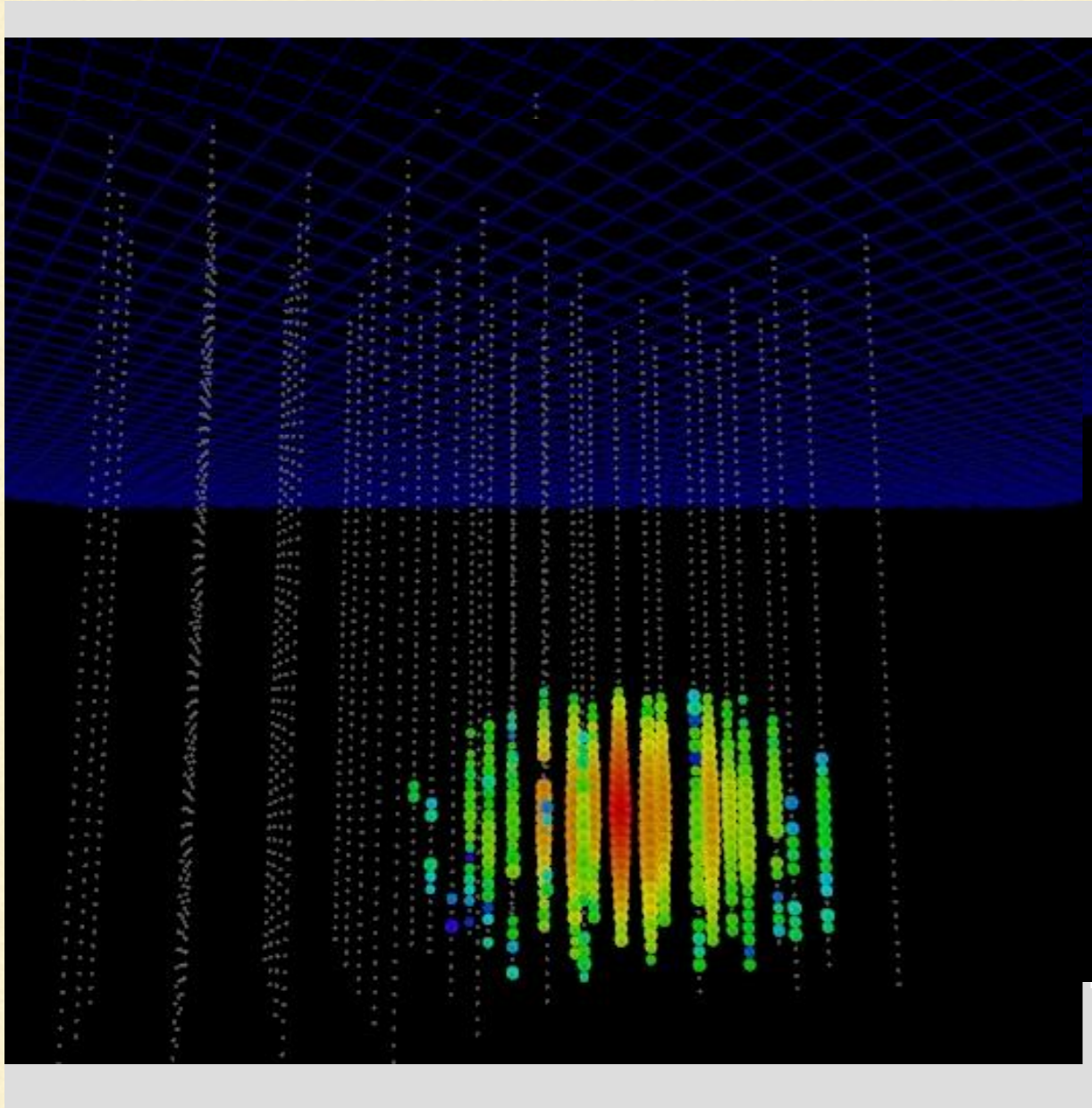
17 m distance between 2 OM on same string

125 m distance between 2 consecutive strings

1 km<sup>3</sup> instrumented volume



# Signals in Icecube..... Showers/

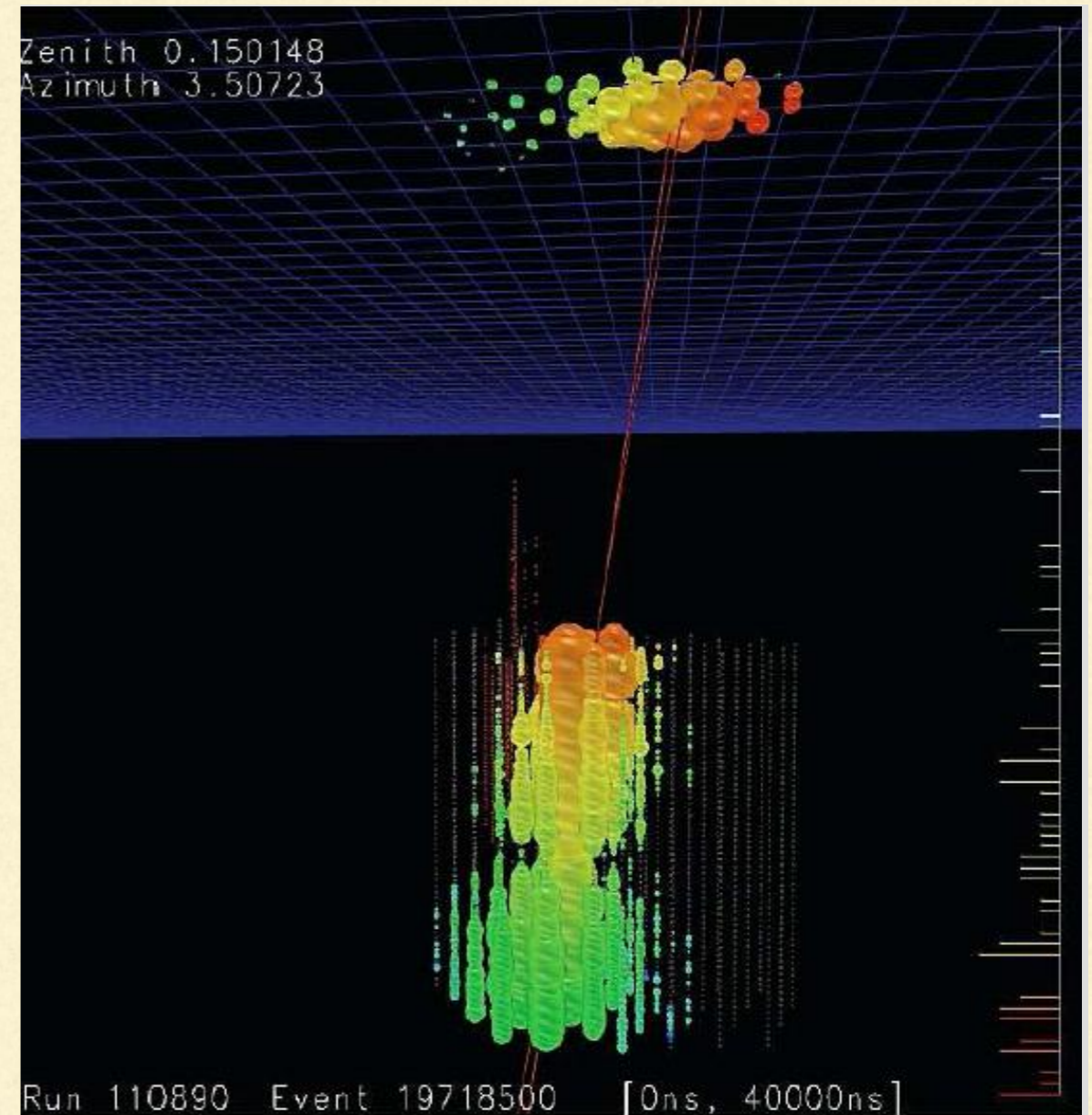


Shower/Cascade

All NC, most CC  $\nu_\tau$  all CC  $\nu_e$

15 % resolution on the deposited energy

$10^\circ$  angular resolution (above 100 TeV)



Track event (muons)

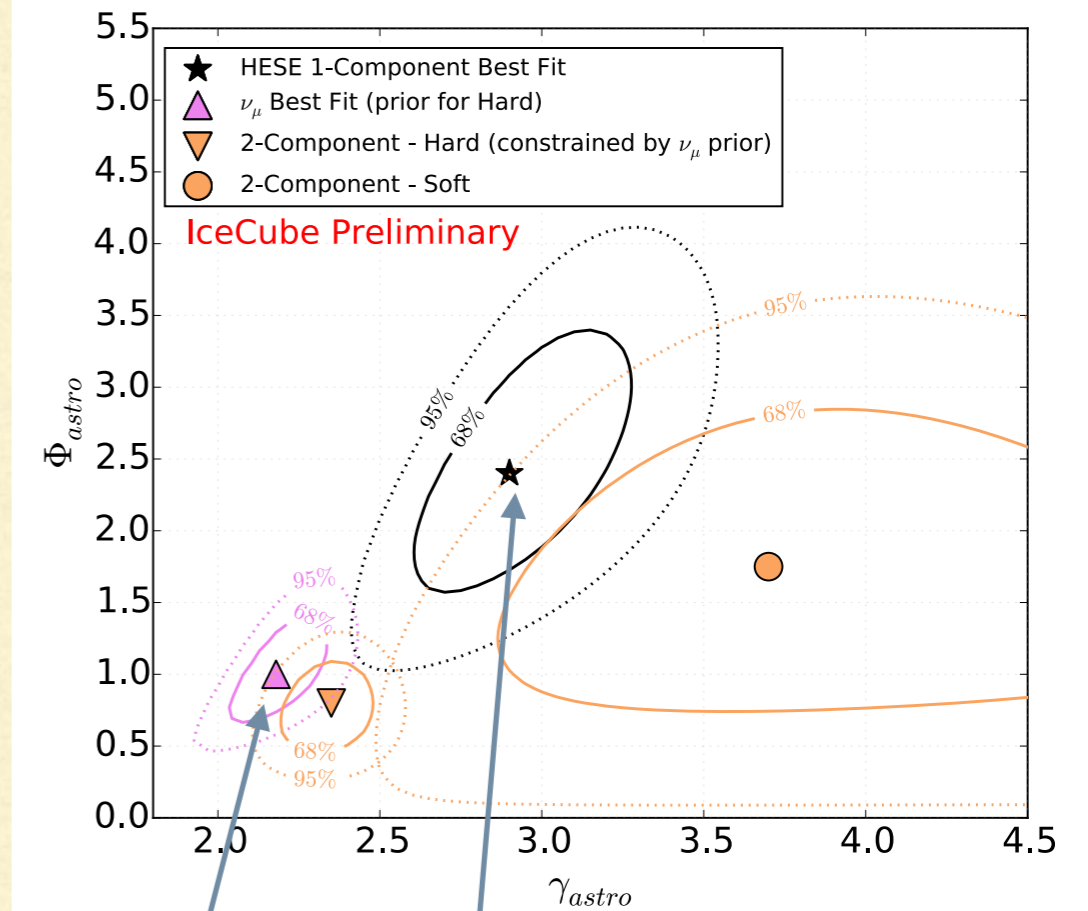
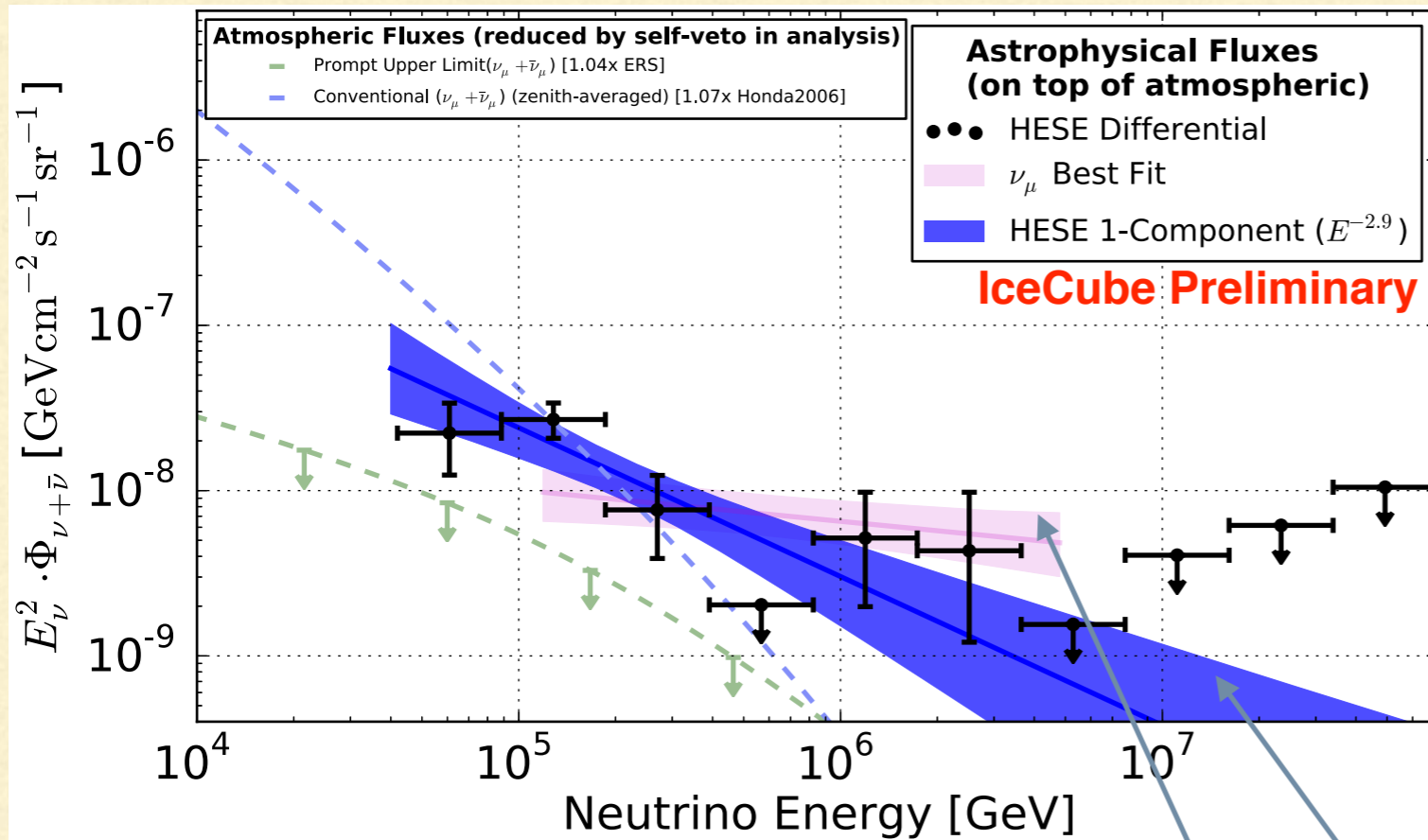
Charged current  $\nu_\mu$

Factor  $\sim 2$  energy resolution

$< 1^\circ$  angular resolution



# Questions/Issues: Power-law behavior of observed neutrino fluxes....



It is widely believed that UHE neutrinos are produced in charged pion decays produced in pp and or p $\gamma$  interactions in the source. Such neutrinos are expected to follow a  $E^{-2}$  spectrum

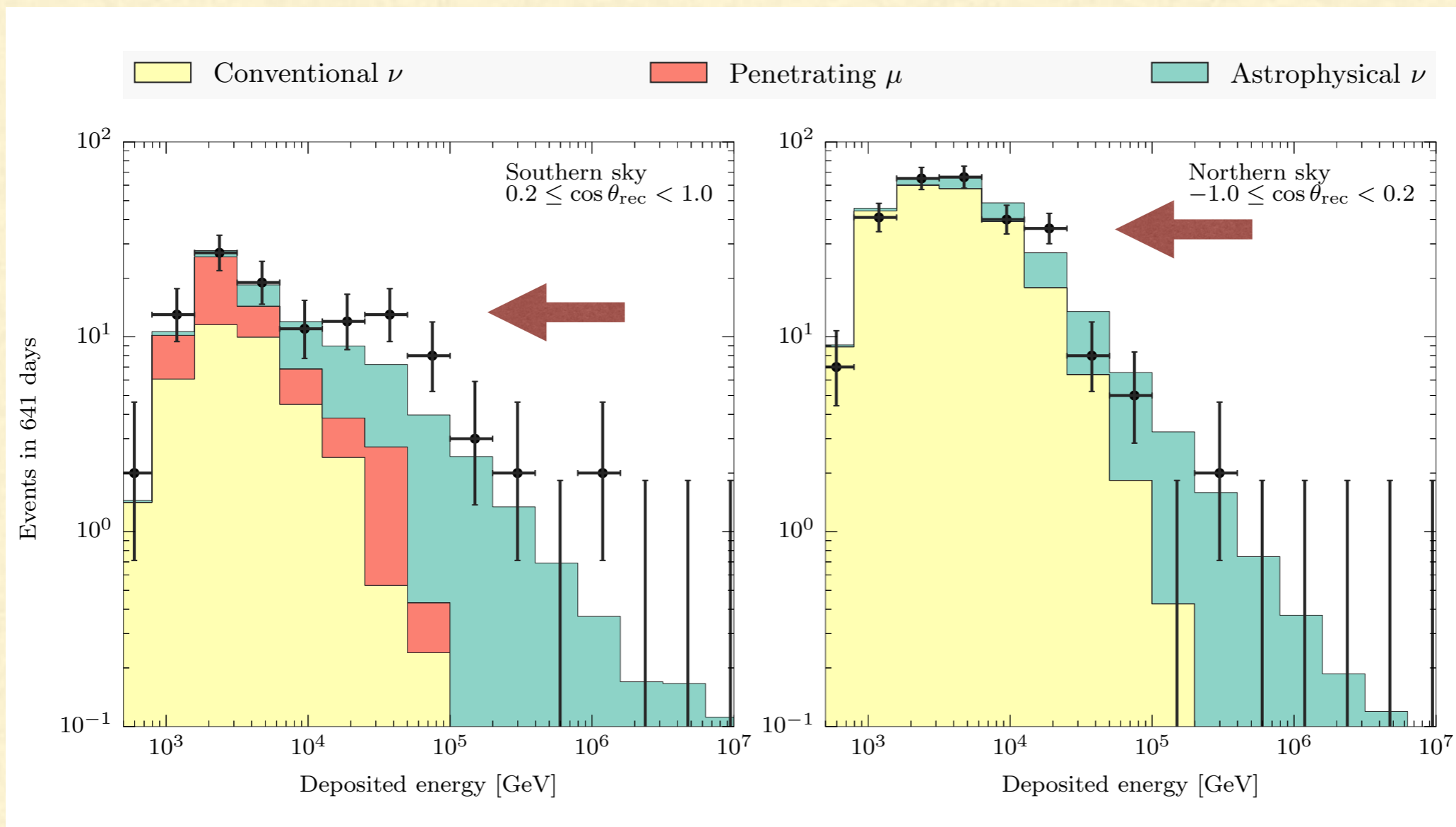
However....

Power-law behavior (index) of 8 yr up-going muon data and HESE data significantly different.



# Questions/Issues: Excess in 30-100 GeV region.....

At lower energies, in the range of 50 - 200 TeV, there appears to be an excess, with a bump-like feature (compared to a simple power-law spectrum), which is prominently present in events from the southern hemisphere, but also visible in events from the northern hemisphere. The maximum local significance of this excess is about  $2.3\sigma$ .

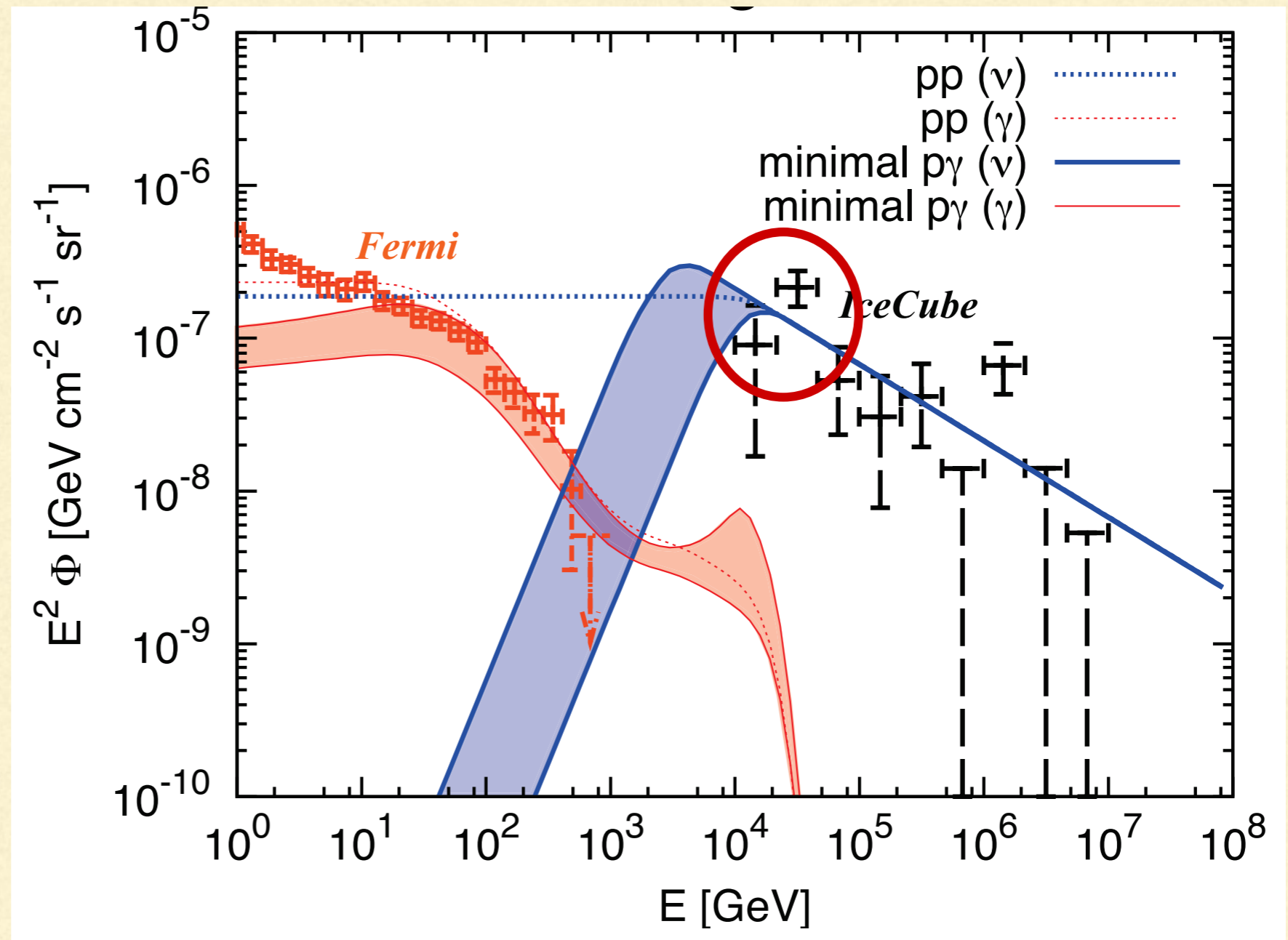




# Fermi Gamma-ray data in tension with IC neutrino data in >30 TeV range.....

For any source, the same processes that produce charged pions which decay to give you the UHE neutrino flux also produce neutral pions which decay to HE photons.

This leads to a natural co-relation between the  $\nu$  and the  $\gamma$  fluxes.



For both pp and p $\gamma$  sources, the observed neutrino flux in IC in the 30-200 TeV region exhibits strong tension with Fermi gamma ray (IRGB) data in GeV region.

This implies either "dark" or opaque sources, or new physics.





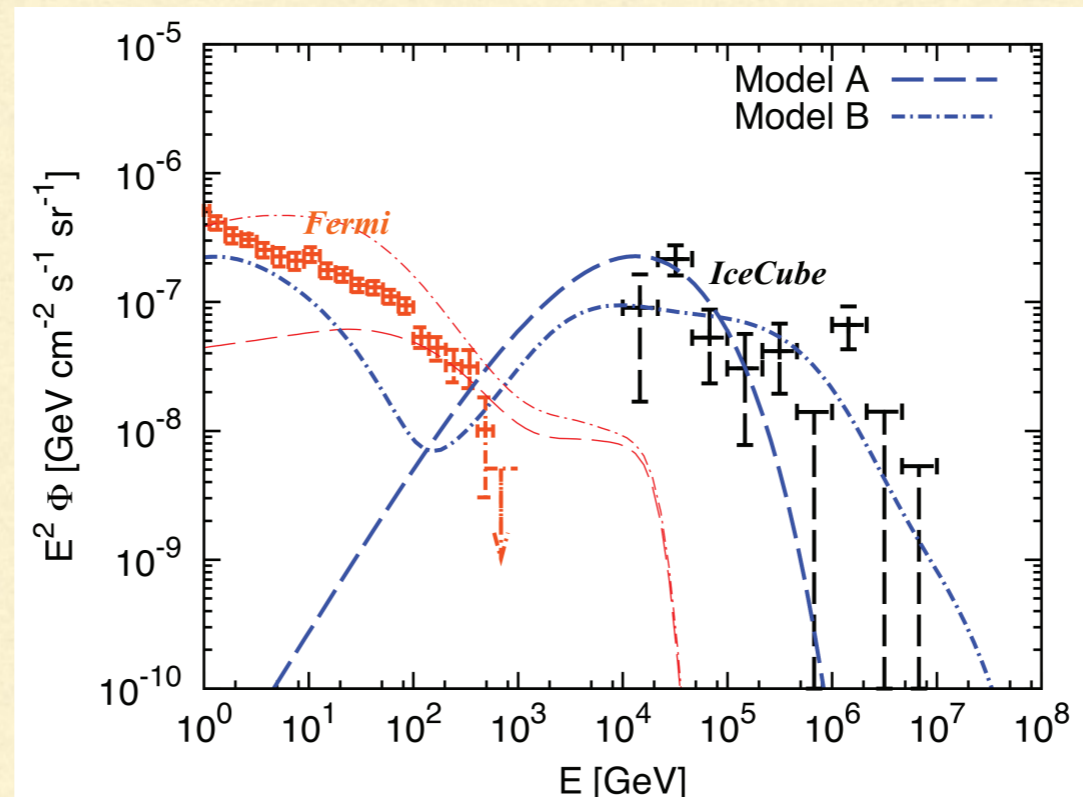
# The "hidden source solution" to the IC signals.....

$\gamma$  rays above TeV energies initiate electromagnetic cascades in the extragalactic background light (EBL) and cosmic microwave background (CMB) as they propagate over cosmic distances. As a result, high-energy  $\gamma$  rays are regenerated at sub-TeV energies, and should have been seen by Fermi.

Thus, assume and study sources are such that two-photon annihilation, inverse-Compton scattering, and synchrotron radiation processes in them can prevent direct  $\gamma$ -ray escape — "dark/hidden sources"

Possible with  $p\gamma$ , but strong tension in case of  $pp$  sources persists.

Conclude that dark  $p\gamma$  sources could alleviate this tension, examples of such sources are models of choked gamma-ray burst (GRB) jets and active galactic nuclei (AGN) cores which are opaque to GeV-TeV  $\gamma$  rays.





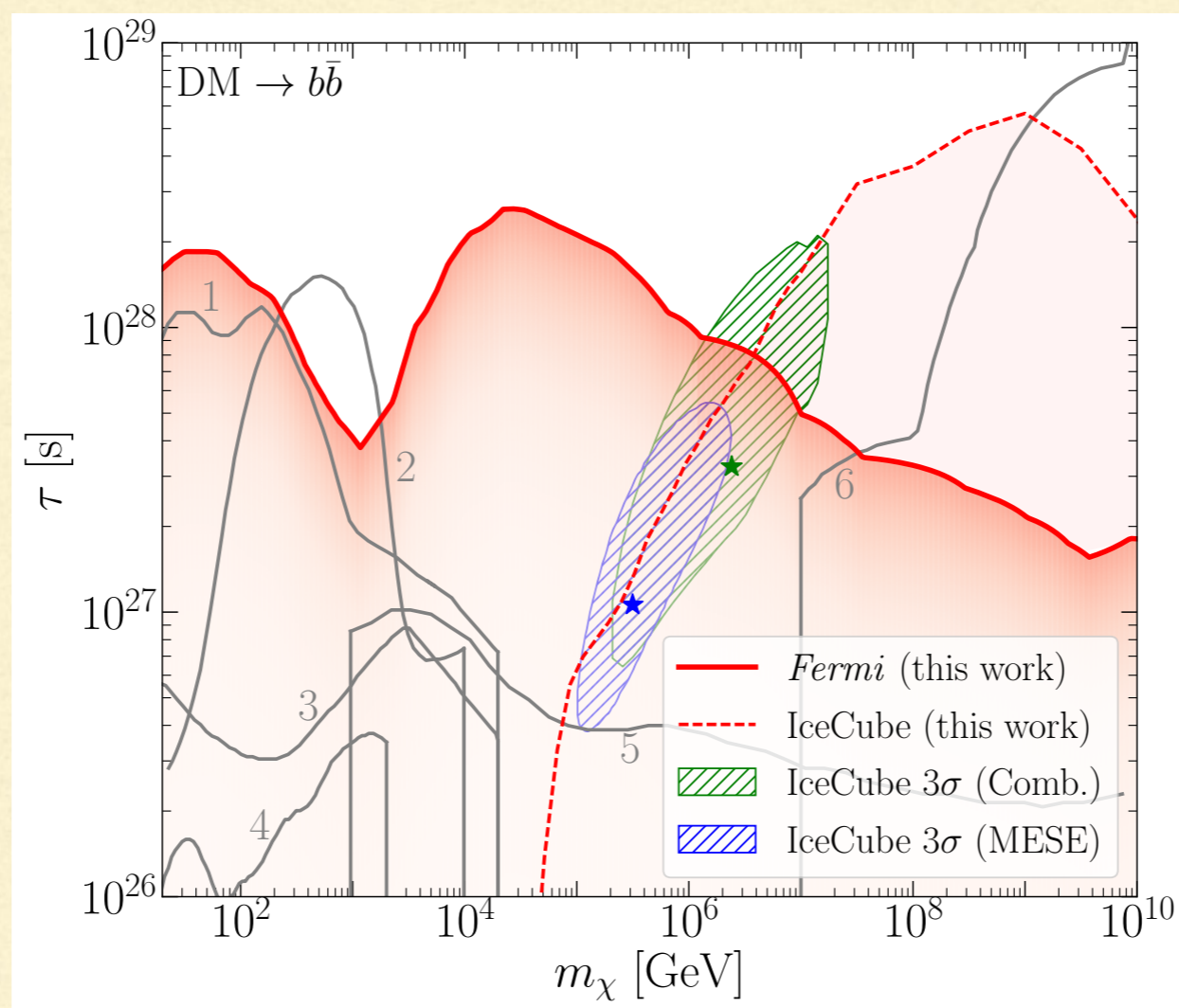
# Power-law incompatibilities and DM.....

The incompatibilities a) between expected  $E^{-2}$  flux and observed spectrum b) between through going muons and HESE spectra, along with proximity of flux to WB bound have led to the speculation that IC sees more than one flux.

Secondly, the second component may not be astrophysical, but due to decay of DM to SM particles leading to neutrinos.

The  $\gamma$ -ray constraints from Fermi can also be used to constrain DM mass and lifetime in this scenario.

Cohen et al, 1612.05638



Explanation of MESE events (30-200 TeV) excess via DM  $\rightarrow$  SM particles very strongly constrained by Fermi-LAT.

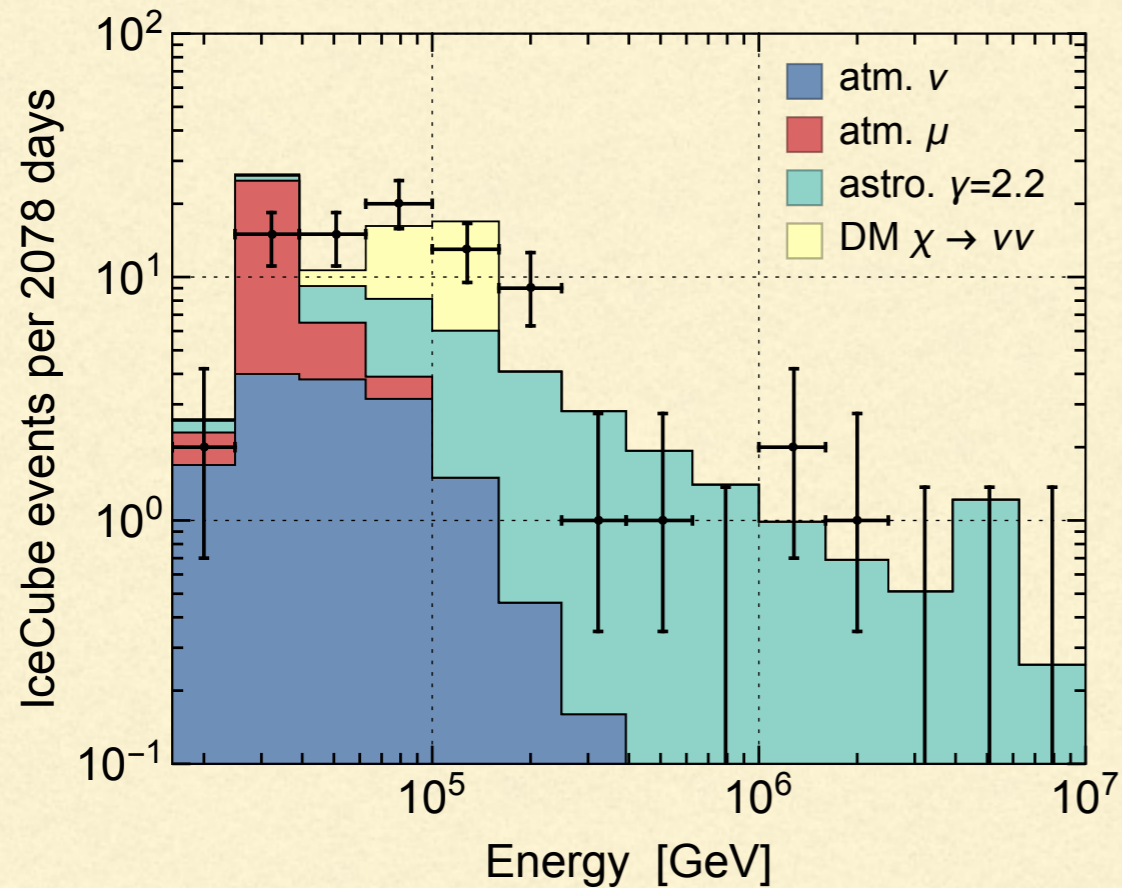


# IceCube Anomalies and DM.....

This implies either "hidden" or opaque sources, or new physics.

Example of new physics: DM which decays only to neutrinos

Chianese et al arXiv: 1808.02486



Extend the Standard Model with a scalar  $SU(2)_L$ -triplet with hyper-charge  $Y = +1$

$$\Delta = \sum_{i=1}^3 \delta_i \tau_i = \begin{pmatrix} \Delta^+ & \sqrt{2}\Delta^{++} \\ \sqrt{2}\Delta^0 & -\Delta^+ \end{pmatrix}$$

Role of Dark Matter is played by the neutral component of this triplet,  $\chi$ , which couples to leptons via

$$\mathcal{L}_\nu = \frac{1}{2} \lambda_{ij} L_i^T C^{-1} i\tau_2 \Delta L_j + \text{h.c.},$$

Impose new global  $U(1)$  and arrange its charges such that it allows the  $\chi$  to decay to only to neutrinos via

$$L_i = \begin{pmatrix} \nu_{iL} \\ \ell_{iL}^- \end{pmatrix}$$

$$\frac{1}{\sqrt{2}} \lambda_{ij} \chi \nu_{iL}^T C^{-1} \nu_{jL} + \text{h.c.},$$

No coupling to quarks due to color conservation



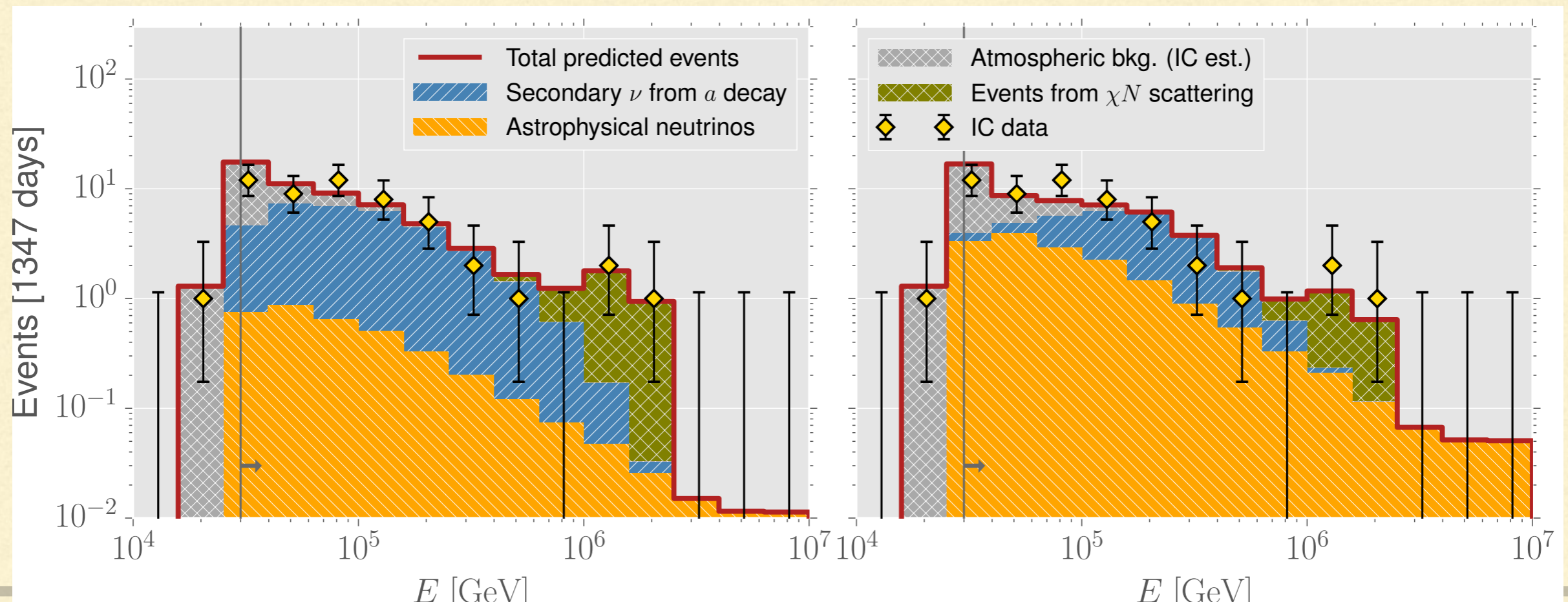
# Another new physics possibility... DM

Bhattacharya, RG, Gupta JCAP  
1503 (2015), 027 (1407.3280)

Study the implications of the premise that any new, relativistic, highly energetic neutral particle that interacts with quarks and gluons would create cascade-like events in the IceCube (IC) detector.

Bhattacharya, RG,  
Gupta, S. Mukhopadhyay  
JCAP 1705 (2017) no.05,  
002 (1612.02834)

Premise: A flux of boosted light dark matter (LDM) particles ( $\chi$ ), which results from the late-time decay of a heavy dark matter (HDM) particle ( $\phi$ ). When  $\chi$  is much lighter than  $\phi$ , its scattering in IC resembles the NC DIS scattering of an energetic neutrino, giving rise to cascade-like events.

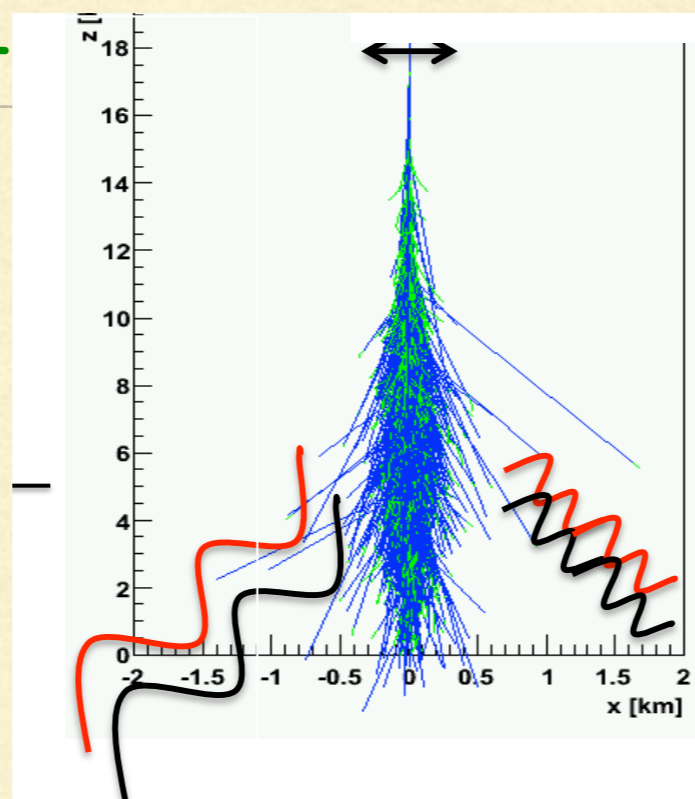




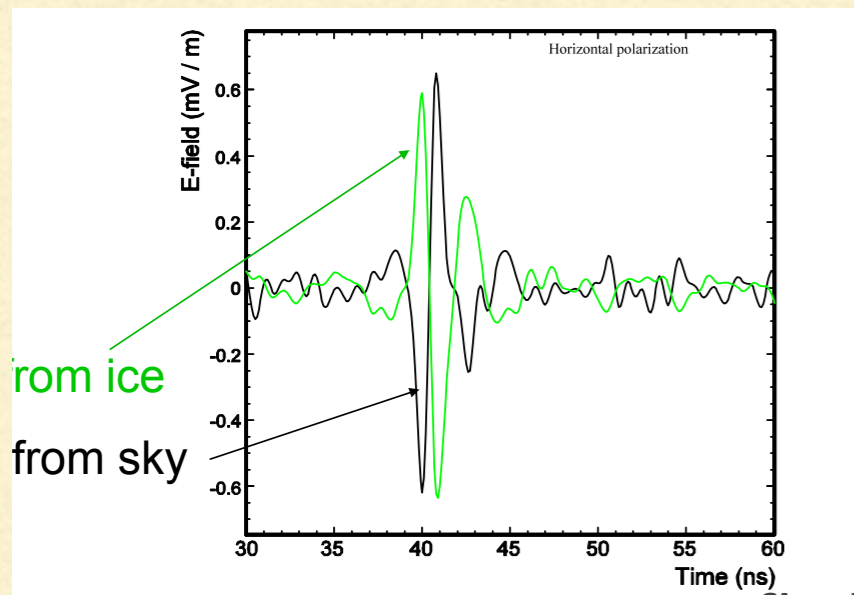
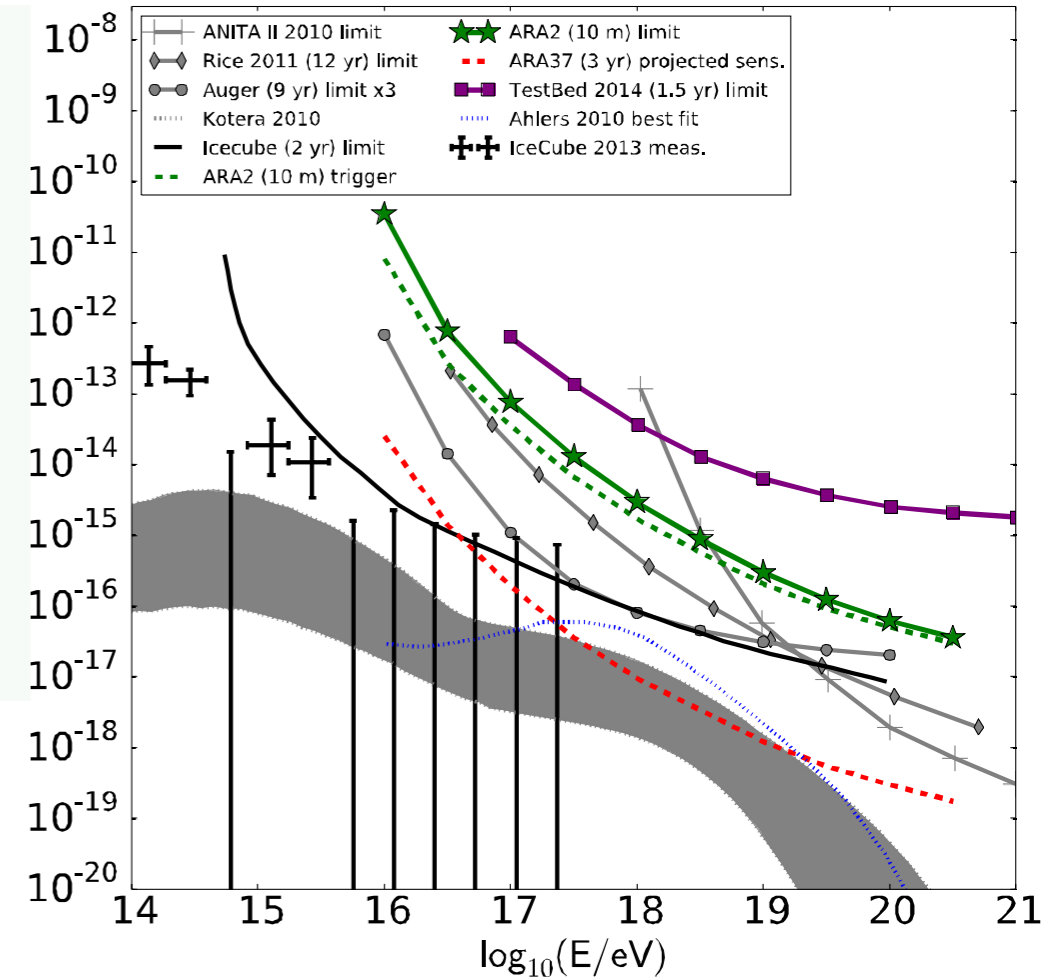
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- ▶ The ANITA observations of the highest energy events
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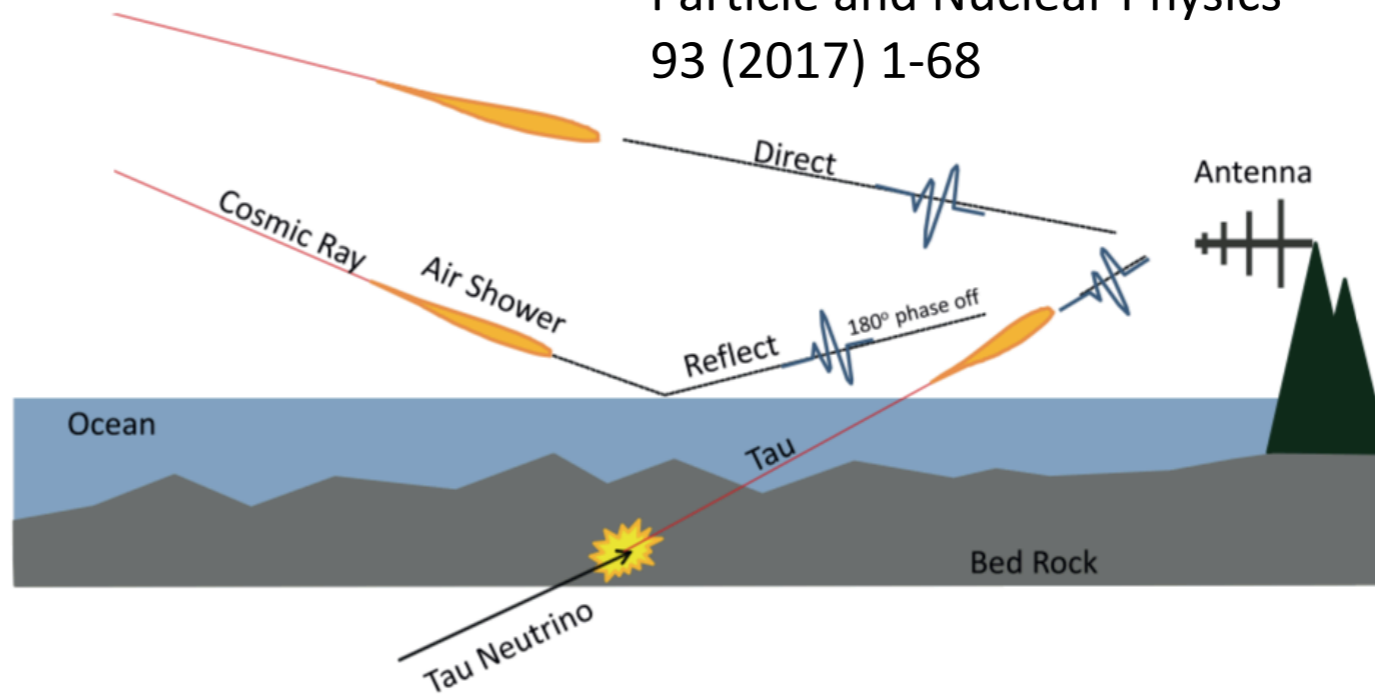
# The ANITA Experiment



Radio pulses produced by UHE showers



## Particle and Nuclear Physics 93 (2017) 1-68



TAROGE-1



# The ANITA Anomalous Events.....

TABLE I. Properties of the ANITA Anomalous Events

Property	AAE 061228	AAE 141220
Flight & Event	ANITA-I #3985267	ANITA-III #15717147
Date & Time (UTC)	2006-12-28 00:33:20	2014-12-20 08:33:22.5
Equatorial coordinates (J2000)	R.A. 282°14064, Dec. +20°33043	R.A. 50°78203, Dec. +38°65498
Energy $\varepsilon_{cr}$	$0.6 \pm 0.4$ EeV	$0.56_{-0.20}^{+0.30}$ EeV
Zenith angle $z'/z$	$117.4 / 116.8 \pm 0.3$	$125.0 / 124.5 \pm 0.3$
Earth chord length $\ell$	$5740 \pm 60$ km	$7210 \pm 55$ km
Mean interaction length for $\varepsilon_\nu = 1$ EeV	290 km	265 km
$p_{SM}(\varepsilon_\tau > 0.1 \text{ EeV})$ for $\varepsilon_\nu = 1 \text{ EeV}$	$4.4 \times 10^{-7}$	$3.2 \times 10^{-8}$
$p_{SM}(z > z_{obs})$ for $\varepsilon_\nu = 1 \text{ EeV}, \varepsilon_\tau > 0.1 \text{ EeV}$	$6.7 \times 10^{-5}$	$3.8 \times 10^{-6}$
$n_\tau(1-10 \text{ PeV}) : n_\tau(10-100 \text{ PeV}) : n_\tau(> 0.1 \text{ EeV})$	34 : 35 : 1	270 : 120 : 1

$\tau$  must be produced close to and inside earth's surface. This implies many interactions for the primary  $\nu_\tau$  given zenith angle, and implies starting energies which are very high. Flux at these energies is very low, and this flux violates bounds by Pierre Auger and IceCube.

**New physics? Source must be inside earth.**

**SM cross sections imply that it is very unlikely that these are tau neutrinos.**

**Example of new physics: Model with heavy  $\nu_R$  neutrino in CPT symmetric universe, mass 480 PeV, which is also DM, decays inside earth to Higgs and SM neutrino. need non-central distribution—assume collision of earth with “dark disk”**



# Summary

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- ▶ The MiniBooNE excess, when combined with previous LSND results, is an intriguing puzzle
  - ▶ Could be due to mundane physics (un-understood background)
    - ▶ But results have been carefully scrutinized over the long term and many important backgrounds measured may also imply that they are signals of new physics. If so, whatever new physics explains one or both is likely to be non-trivial and important.
    - ▶ The question whether neutrinos are Dirac or Majorana in nature remains unanswered. Recent progress offers possibly an additional handle on this if there exists a heavy sterile neutrino, by studying the energy and angular distributions in its decay.
  - ▶ While several experiments are looking for such a neutrino, determining angular and energy distributions in a decay to an active neutrino and a boson is very challenging.
  - ▶ IceCube events show a power-law discrepancy between up going muon and contained HESE/MESE events, and a tension with Fermi-LAT gamma-ray data in GeV region. Hidden sources? or DM? DM decay to SM particles also in tension with Fermi-LAT. DM to DM decay?
  - ▶ ANITA, a balloon experiment, has recorded 2 events which are  $\sim 600$  PeV with non-inversion in polarization of the detected radio signal. While  $\nu_\tau$  can in principle be responsible, angle of approach makes this highly unlikely. If new physics, need fresh ideas.
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*Thank you for your attention*

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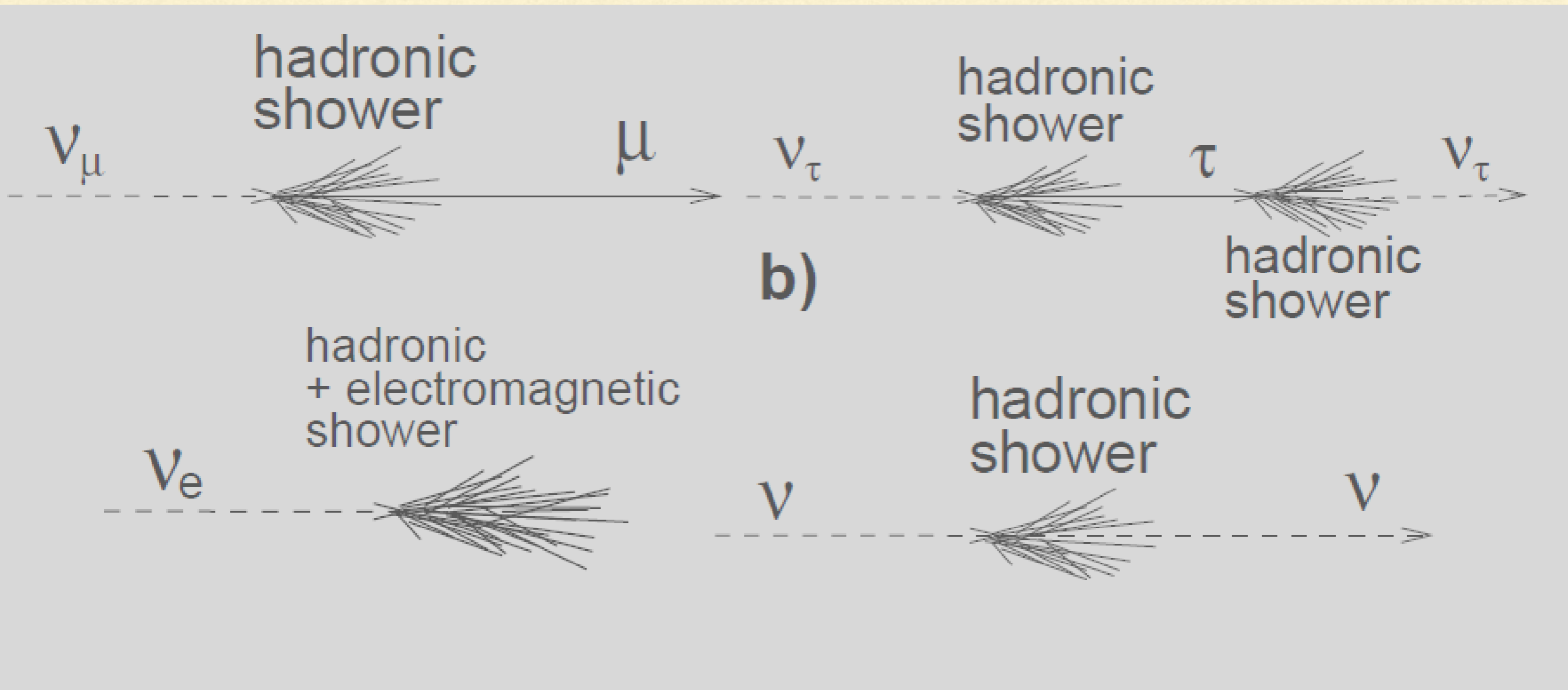
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# Backup Slides

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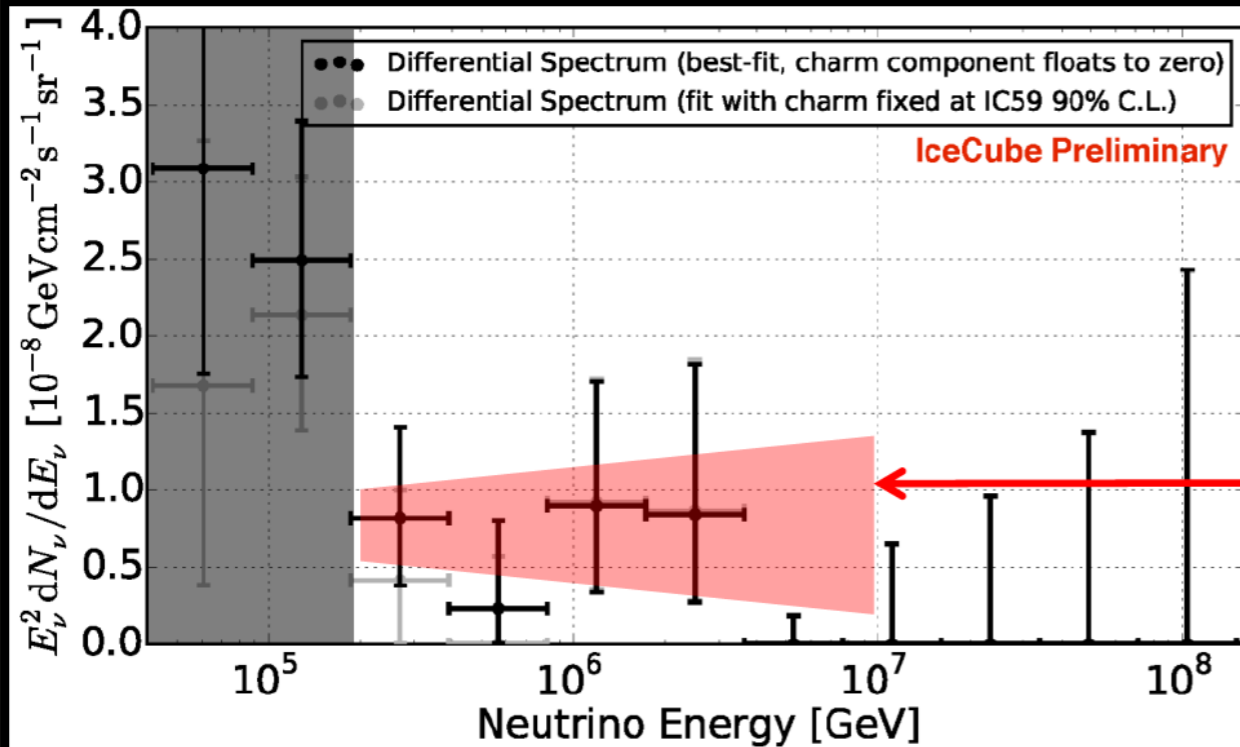
# Neutrino Signals in IceCube.....



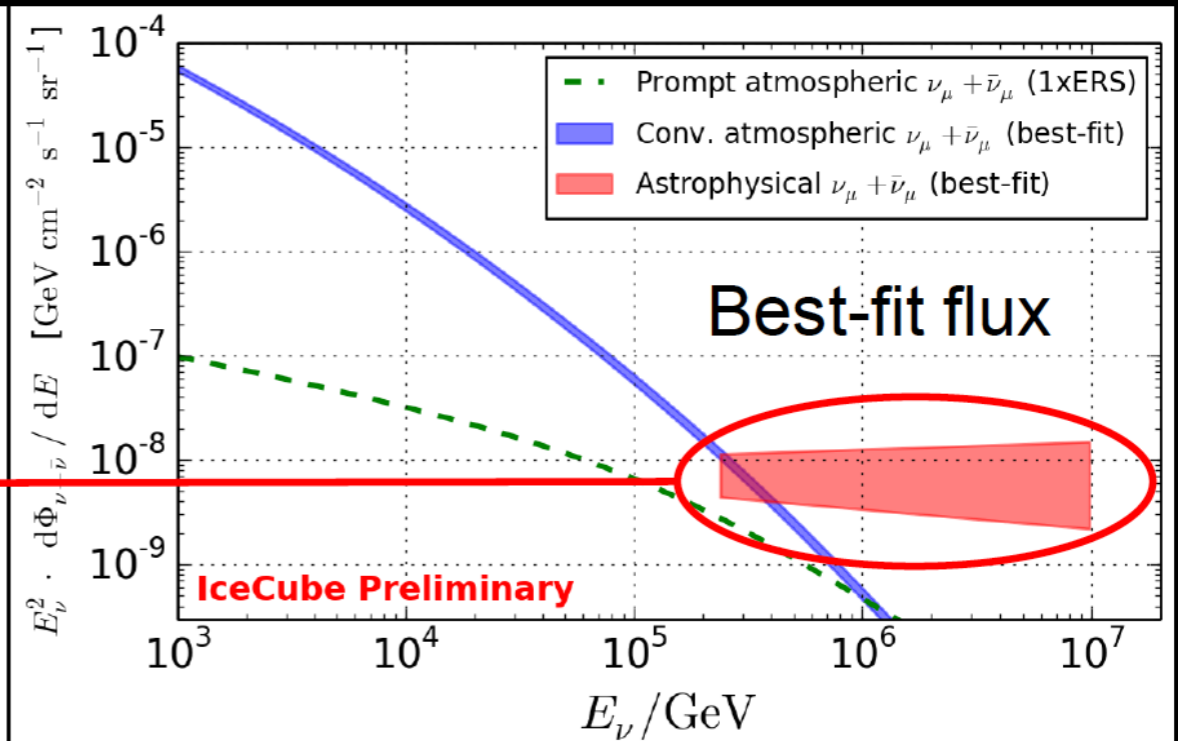


# Questions/Issues: Power-law behavior of observed neutrino fluxes.....

HESE 4 year unfolding  
(→ dominated by shower-like events)



6 year up-going numu analysis



Power-laws of the HESE and thoroughgoing muon fluxes seem consistent with each other only above 100 TeV, and with Fermi shock acceleration.

Difficult, in this way of looking at the data, to understand the 30-100 TeV data (MESE), or use single power-law for all data.



# The "hidden source solution" to the IC signals.....

$\gamma$  rays above TeV energies initiate electromagnetic cascades in the extragalactic background light (EBL) and cosmic microwave background (CMB) as they propagate over cosmic distances. As a result, high-energy  $\gamma$  rays are regenerated at sub-TeV energies, and should have been seen by Fermi.

Thus, assume and study sources are such that two-photon annihilation, inverse-Compton scattering, and synchrotron radiation processes in them can prevent direct  $\gamma$ -ray escape — "dark/hidden sources"

Possible with  $p\gamma$ , but strong tension in case of  $pp$  sources persists.

Conclude that dark  $p\gamma$  sources could alleviate this tension, examples of such sources are models of choked gamma-ray burst (GRB) jets and active galactic nuclei (AGN) cores which are opaque to GeV-TeV  $\gamma$  rays.

