

# Jet structure studies in small systems

(arXiv:1805.03101, 1809.10102)

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in collaboration with

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# Motivation - collectivity in small systems

- Surprise at LHC energies:  $v_2$  and long-range correlations in high-multiplicity p+p events;
  - Long range correlations. **ATLAS, Nucl. Phys. A932, 357 (2014)**
  - Azimuthal anisotropy ( $v_n$ , flow).  
**L. Yan, J. Y. Ollitrault, Phys. Rev. Lett. 112, 082301 (2014)**
  - Enhanced heavy quark production depending on the multiplicity. **ALICE Collaboration, JHEP 1608, 078 (2016)**
- QGP is not necessary for collectivity.
  - Vacuum-QCD effects can produce such behaviour in the soft-hard regime.
  - Multi-parton interactions (MPI) can qualitatively explain enhanced heavy flavour production.

# Jet modification in small systems

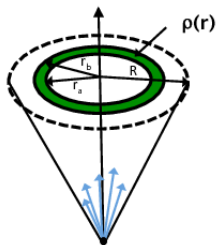
- Jet structure may be modified even in small systems without large volume of QGP.
- **We are looking for non-trivial jet shape dependence on event multiplicity.**
- We are simulating p+p collisions at  $\sqrt{s} = 7$  TeV.
- Monte Carlo event generator: PYTHIA 8.2 with different PDF sets: NNPDF2.3lo for Monash and Monash\*, CTEQ6L1 for 4C.
- Jet reconstruction: Fastjet software package with anti- $k_t$  algorithm.
- Full jet reconstruction with  $R = 0.7$ .  
( $R = \sqrt{\Delta\phi^2 + \Delta\eta^2}$ )

# Different tunes and settings

- Tunes: Monash, Monash\*, 4C.
- Multi parton interactions (MPI).
- Colour reconnection (CR): In PYTHIA this is an in-built mechanism that allows interactions between partons originating in MPI and initial/final state radiations, by minimizing color string lengths.
  - 0: MPI-based scheme,
  - 1: QCD-based string length minimisation scheme,
  - 2: gluon-move scheme.
  - off: we don't use it.

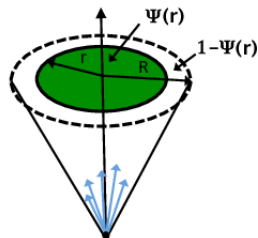
# Differential and integral jet shapes

Differential jet shape:



$$\rho(r) = \frac{1}{\delta r} \frac{\sum_{r_a < r_i < r_b} p_t^{(i)}}{p_t^{\text{jet}}}.$$

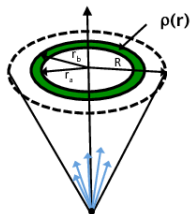
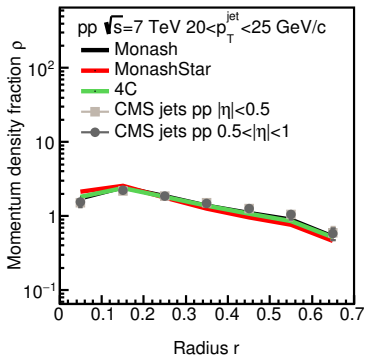
Integral jet shape:



$$\psi(r) = \frac{\sum_{r_i < r} p_t^{(i)}}{p_t^{\text{jet}}}.$$

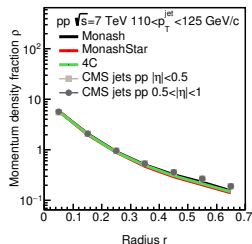
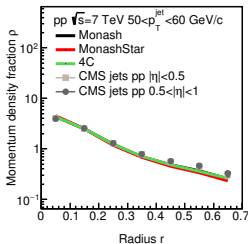
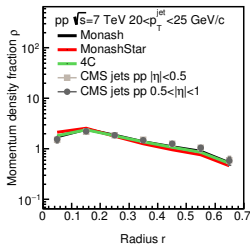
$$\psi(R) = \int_0^R \rho(r') dr' = 1.$$

# A reality check: Comparison with CMS for $\rho(r)$



- The three different tunes reproduce the 7 TeV  $|\eta| < 1$  pp CMS data within uncertainty.
- We investigated different  $p_T^{jet}$  windows between  $15 \text{ GeV}/c < p_T^{jet} < 400 \text{ GeV}/c$ .

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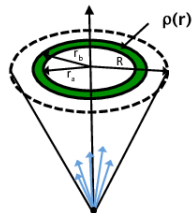
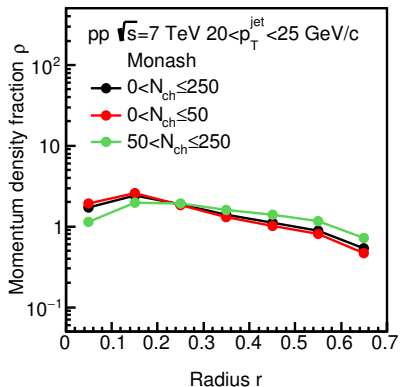


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# New study: Selection of jets by multiplicity!

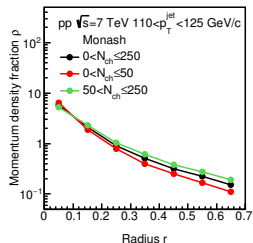
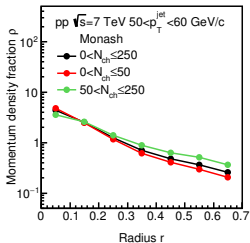
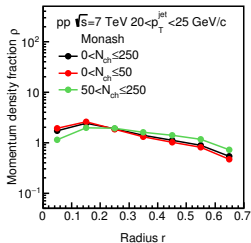


# $\rho(r)$ distribution for different multiplicities



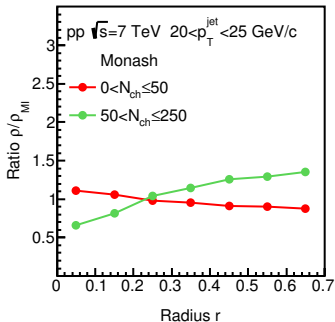
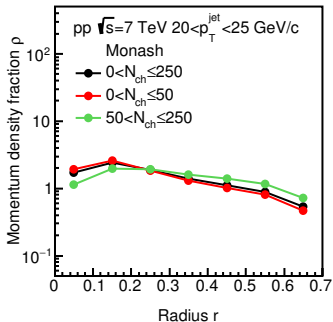
- We see a multiplicity dependence in the jet shapes,
- but it is the trivial multiplicity dependence we expected.
- **Lower** multiplicity jets are more collimated.

# $\rho(r)$ distribution for different multiplicities



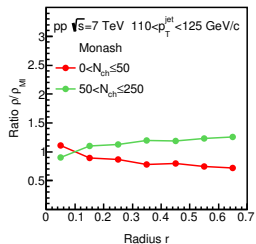
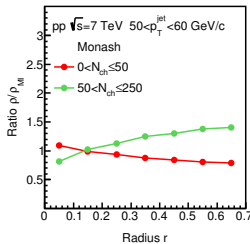
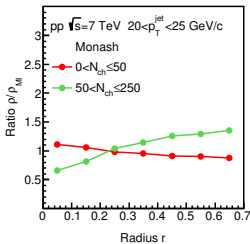
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- **Lower** multiplicity jets are more collimated.

# Ratio with Multiplicity Integrated (MI) $\rho(r)$



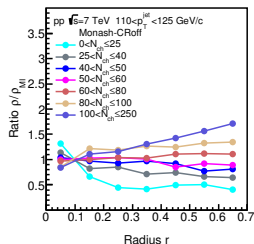
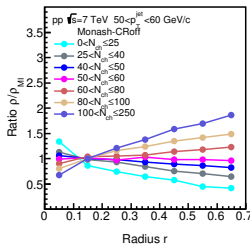
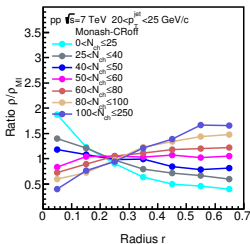
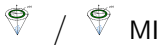
- We divide the curves in the **left** figure with the multiplicity integrated (MI) curves, so
- we gain the **right** figure.

# Ratio with Multiplicity Integrated (MI) $\rho(r)$



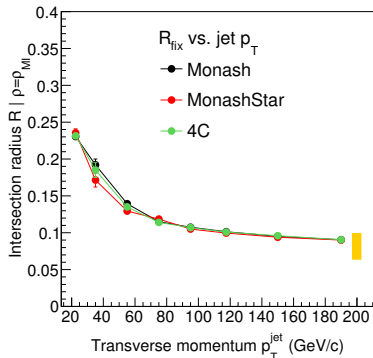
- Two curves:  $\rho_{low-Nch}/\rho_{MI}$  and  $\rho_{high-Nch}/\rho_{MI}$ .
- The low- and high-multiplicity curves intersect each other at unity.
- The intersection point depends on the  $p_T^{jet}$  (three examples above).

# A characteristic jet size measure?



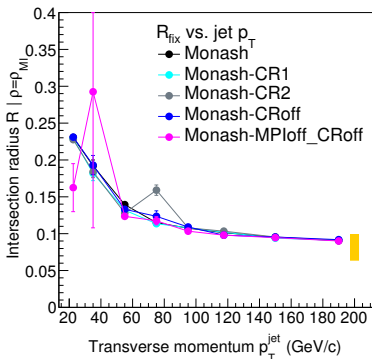
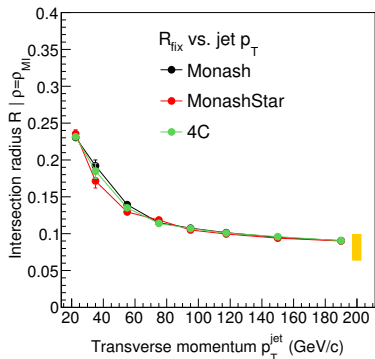
- Pythia describes the multiplicity distributions well.
- The intersection does not depend on our multiplicity bin choice, but it depends on the  $p_T^{\text{jet}}$ .
- Our finding: a non-trivial scaling behaviour.

# The $p_T^{jet}$ dependence of $R_{fix}$



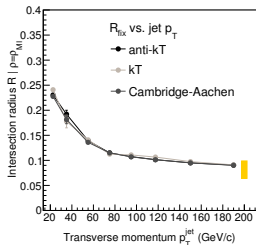
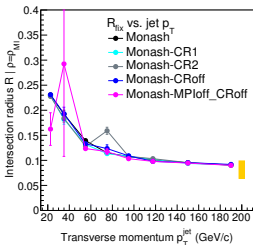
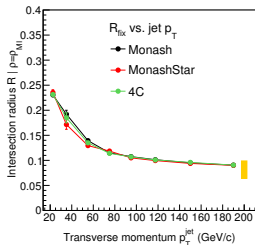
- The  $R_{fix}$  depends on the  $p_T^{jet}$ .
- The shape of the curve can be qualitatively explained by a Lorentz boost.

# The $p_t^{\text{jet}}$ dependence of $R_{\text{fix}}$



- Good agreement between tunes (left) and settings (right).

# The $p_t^{jet}$ dependence of $R_{fix}$

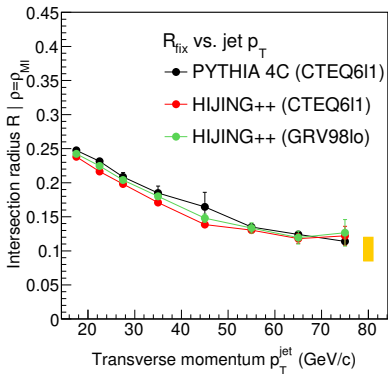
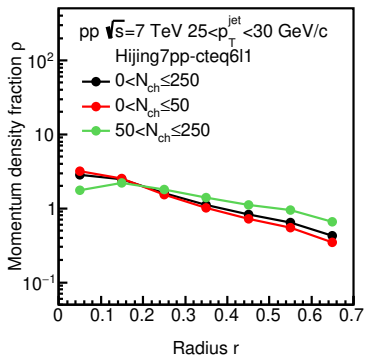


- Good agreement between tunes (left) and settings (right).
- The  $R_{fix}$  does not depend on the three jet reconstruction algorithms (right).
- $R_{fix}$  is a **characteristic jet size** at a given  $p_t^{jet}$ .



**Let's try out a different MC generator!**

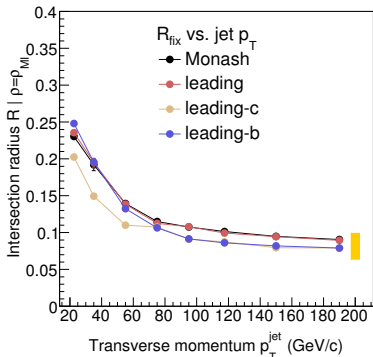
# Jet structure with Hijing++



- We had a look at the new Hijing++ code.
- The  $R_{\text{fix}}$  phenomenon is visible.

**Is the jet structure from Heavy  
Flavour any different?**

# $R_{fix}$ with Heavy Flavour

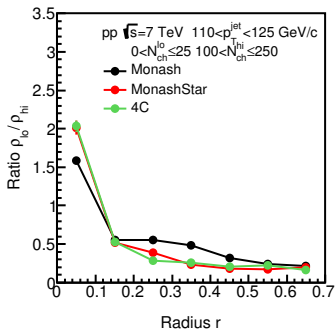


- Selection of leading jets does not make a difference.
- For smaller  $p_T^{jet}$  the charm leading jets appear narrower.
- For higher  $p_T^{jet}$  both charm and bottom jets are narrower.

**Back to the question:**

**Does the jet shape depend on the multiplicity?**

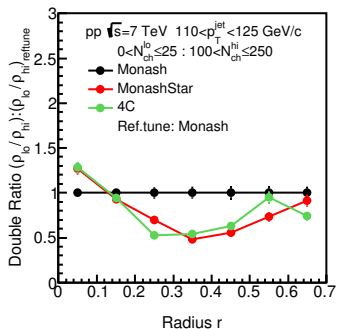
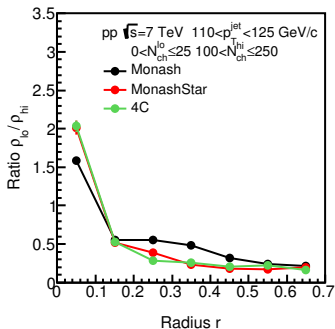
# Applying a double ratio for $\rho(r)$



We want to apply a double ratio to cancel out multiplicity bias:

$$\frac{(\rho_{\text{low}}/\rho_{\text{high}})}{(\rho_{\text{low}}/\rho_{\text{high}})_{\text{ref.tune}}}$$

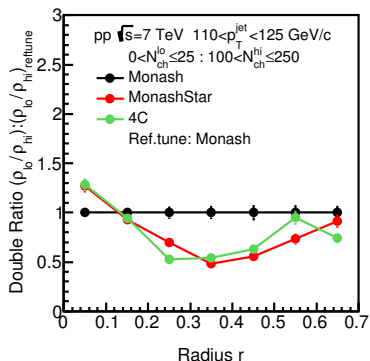
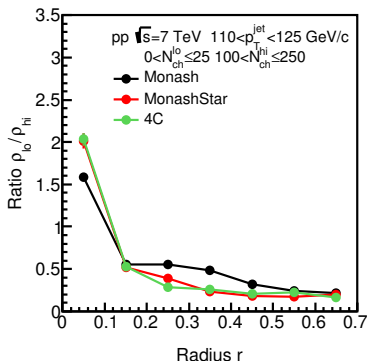
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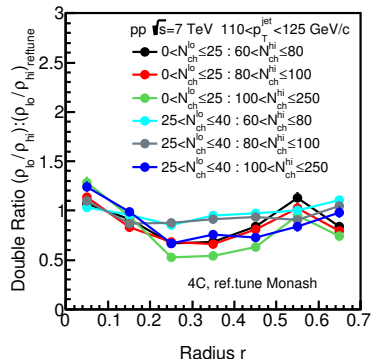
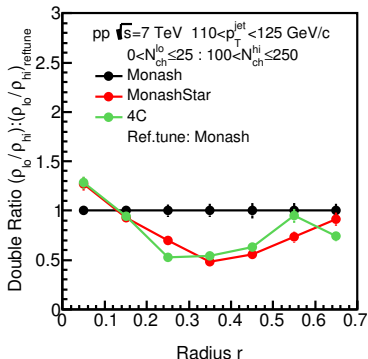
# Applying a double ratio for $\rho(r)$



- Cancelling out trivial multiplicity bias:  $\frac{(\rho_{\text{low}}/\rho_{\text{high}})}{(\rho_{\text{low}}/\rho_{\text{high}})_{\text{ref.tune}}}$ .
- We find a significant effect at given  $p_T^{\text{jet}}$  windows.
- Non-trivial dependence on  $p_T^{\text{jet}}$ , origin of the effect needs further investigation.



# Double ratio for different multiplicity bins



- Same calculations for several different multiplicity bins.
- The effect is larger as the separation in multiplicity is larger.
- Statistically independent samples  $\rightarrow$  not fluctuations.

# Summary

- We gave predictions for several jet structure observables in  $\sqrt{s} = 7$  TeV p+p collisions using PYTHIA 8.
- Multiplicity-dependent experimental jet-structure analyses could differentiate between otherwise well-performing models.
- We suggest  $R_{fix}$  as a multiplicity-independent jet size measure.
- $R_{fix}$  is present in Hijing++, even with different PDF sets.
- Multiplicity-dependent jet structures of Heavy Flavor jets are sensitive probes of flavor-dependent fragmentation.

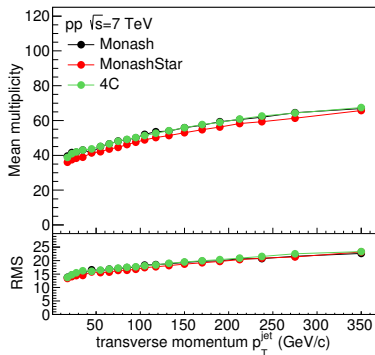
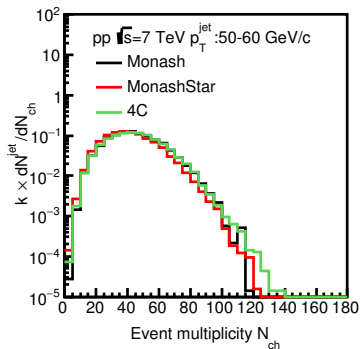
# Thank you for your attention!

This work has been supported by the NKFIH/OTKA K 120660 grant, the János Bolyai scholarship of the Hungarian Academy of Sciences.

The simulations were done with the resources of The Wigner GPU Laboratory and Wigner Datacentre.

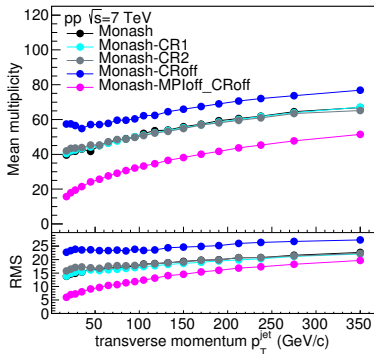
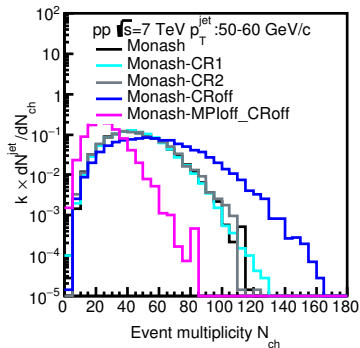
# Backup

# Multiplicity distribution for the tunes



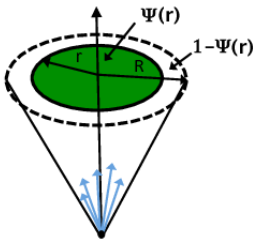
- All three tunes show similar event multiplicity ( $N_{\text{ch}}$ ) distributions (they were calibrated for real data).
- The mean multiplicity increases with the  $p_T^{\text{jet}}$ .

# Multiplicity distribution for the settings



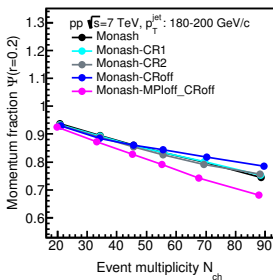
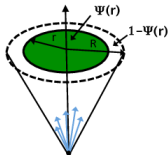
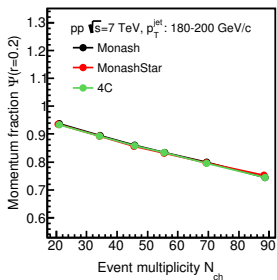
- CR schemes show similar  $N_{\text{ch}}$  distributions.
- The MPI:off and CR:off settings are not physical
  - MPI:off (CR:off)  $\rightarrow$  lower mean multiplicity.
  - CR:off (MPI:on)  $\rightarrow$  higher mean multiplicity.

We also have the integral jet shape  $\psi(r)$



$$\psi(r) = \frac{\sum_{r_i < r} p_t^{(i)}}{p_t^{\text{jet}}}$$

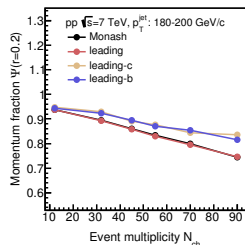
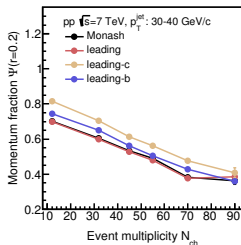
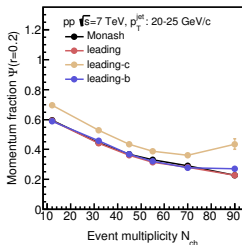
# $\psi(r=0.2)$ dependence on multiplicity



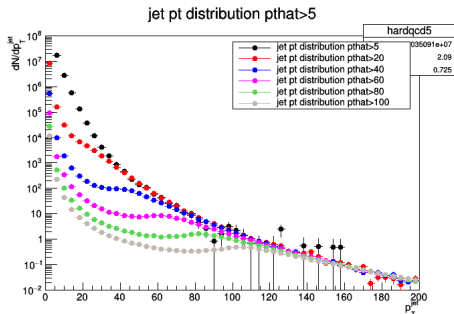
- No significant difference in integral jet structure between the three tunes and between different CR schemes.
- Turning off MPI causes significant differences for higher multiplicities, which can not be explained by bin effects.
- Observation implies MPI influence on jet structure.



# Integral jet structure with Heavy Flavour

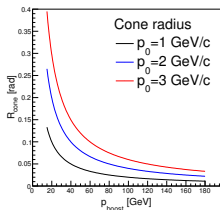


- Charm leading jets always differ (except for very low  $N_{ch}$ ).
- Bottom leading jets also differ for higher multiplicities in the case of high enough  $p_T^{jet}$ .
- For certain  $p_T^{jet}$  (depends on  $r$ ) all curves differ  $\rightarrow$  HF fragmentation is different  $\rightarrow$  model differentiation.

$p_t^{jet}$  distributions for different  $\hat{p}_t$ 


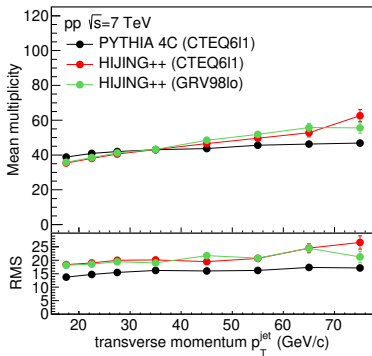
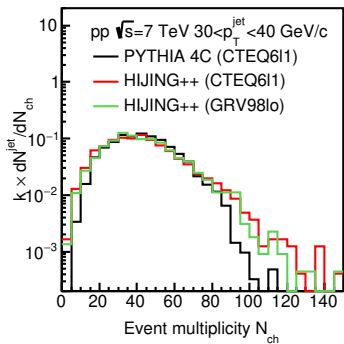
$p_t^{jet}$	$\hat{p}_t$
20 - 25	$5 \leq$
30 - 40	$5 \leq$
50 - 60	$20 \leq$
70 - 80	$20 \leq$
90 - 100	$40 \leq$
110 - 125	$40 \leq$
140 - 160	$80 \leq$
180 - 200	$80 \leq$
225 - 250	$80 \leq$

# A toy model

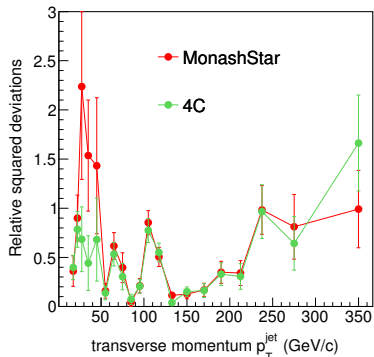
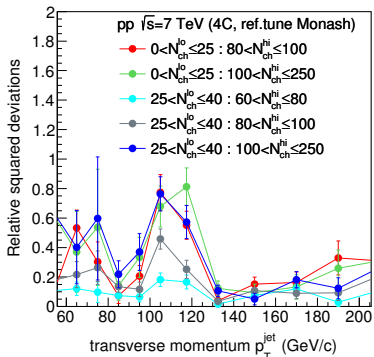


- Jet consisting of particles with equal momentum  $p_0$ .
- We boost with a certain momentum  $p_{\text{boost}}$  towards the jet axis.
- High- $p_T$ : qualitatively similar behaviour of the characteristic jet size with respect to  $p_{\text{boost}}$ .
- Low- $p_T$ : blow-up is not expected in data because jet rec. is limited in R and because of angular cut-off in splitting.

# Multiplicity distribution for Hijing<sub>++</sub>



# Relative squared deviations



- $$RSD = \sqrt{\sum_{0 < r_i < R} (DR(r_i) - 1)^2}$$
- The nature of the  $p_T^{jet}$  dependence is independent of the chosen multiplicity bin.