



# A-formulation method for full 3D FEM computation of the superconductor magnetization

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# Motivation

#### Magnetic cloak:

Straight tapes

# ReBCO **BSCCO** FM shell Ferrite in epoxy

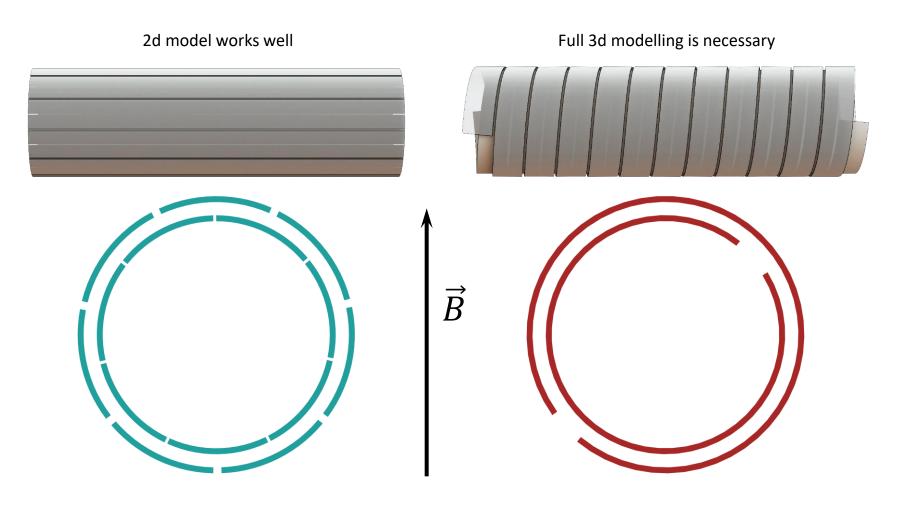
**SC** inserts

2

Helicoidally wound tapes

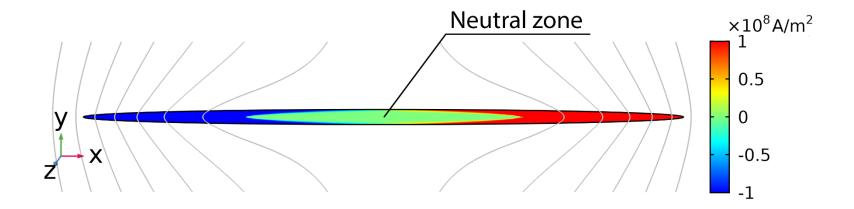
## Motivation

#### How to build a numerical model?



Numerical modelling. From 2d to 3d.

2D: the neutral zone is beneficial for the correct solution<sup>1</sup>



Maxwell equation says

In the cross-section of superconducting wire

$$\vec{E} = -\frac{\partial \vec{A}}{\partial t} - \nabla \varphi \tag{1}$$

$$\vec{E} = -\frac{\partial \vec{A}}{\partial t} \tag{2}$$

Power law definition

Bean model

$$\vec{E} = E_c \left(\frac{\vec{J}}{J_{c0}}\right)^n \tag{3}$$

$$\vec{E} = E_c \left(\frac{\vec{J}}{J_{c0}}\right)^n$$
(3) 
$$\vec{J} = \begin{cases} +J_{c0}, & for & \partial A/\partial t < 0 \\ -J_{c0}, & for & \partial A/\partial t > 0 \\ 0, & for & \partial A/\partial t = 0 \end{cases}$$
(4)

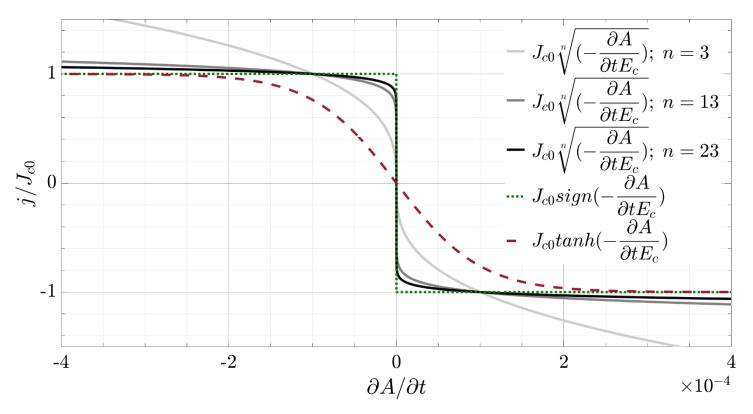
$$\vec{J} = J_{c0} \sqrt[n]{-\frac{\partial \vec{A}}{\partial t E_c}}$$
 (5) 
$$\vec{J} = J_{c0} sign\left(-\frac{\partial \vec{A}}{\partial t E_c}\right)$$
 (6)

$$\vec{J} = J_{c0} \sqrt[n]{-\frac{\partial \vec{A}}{\partial t E_c}}$$

$$\vec{J} = J_{c0} sign\left(-\frac{\partial \vec{A}}{\partial t E_c}\right)$$

The alternative function proposed by A. Campbell<sup>2</sup>

$$\vec{J} = J_{c0} tanh \left( -\frac{\partial \vec{A}}{\partial t E_c} \right) \qquad (7)$$



$$\vec{J} = J_{c0} tanh\left(\frac{\vec{E}}{E_c}\right) \tag{7}$$

In 3D problem is assumed that the relation (7) is valid for each component of the vectors  $\vec{j}$  and  $\vec{E}$   $(\vec{A})$  separately.

Correct solution requires also collinearity between  $\vec{j}$  and  $\vec{E}$   $(\vec{A})$ 

$$E_{x}: E_{y}: E_{z} = j_{x}: j_{y}: j_{z}$$
 (8)

And, if model assumes the isotropic  $J_c$ 

$$\sqrt{j_x^2 + j_y^2 + j_z^2} \le J_{c0} \tag{9}$$

After combining eq. (7), (8) and (9) the final expression for current density is:

$$\vec{J} = \frac{J_{c0}}{|E|} \left( |E_x| tanh\left(\frac{E_x}{E_c}\right) \hat{\imath} + |E_y| tanh\left(\frac{E_y}{E_c}\right) \hat{\jmath} + |E_z| tanh\left(\frac{E_z}{E_c}\right) \hat{k} \right)$$
 (10)

Model verification

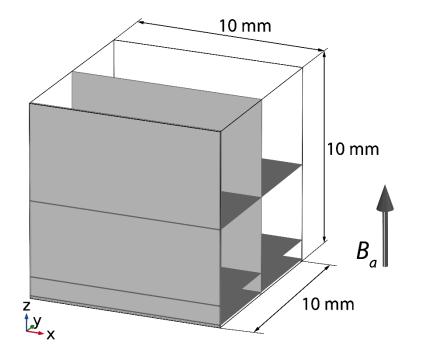
#### www.htsmodelling.com

Cubic bulk as benchmark for 3D modelling of superconductors under slowly varying magnetic fields

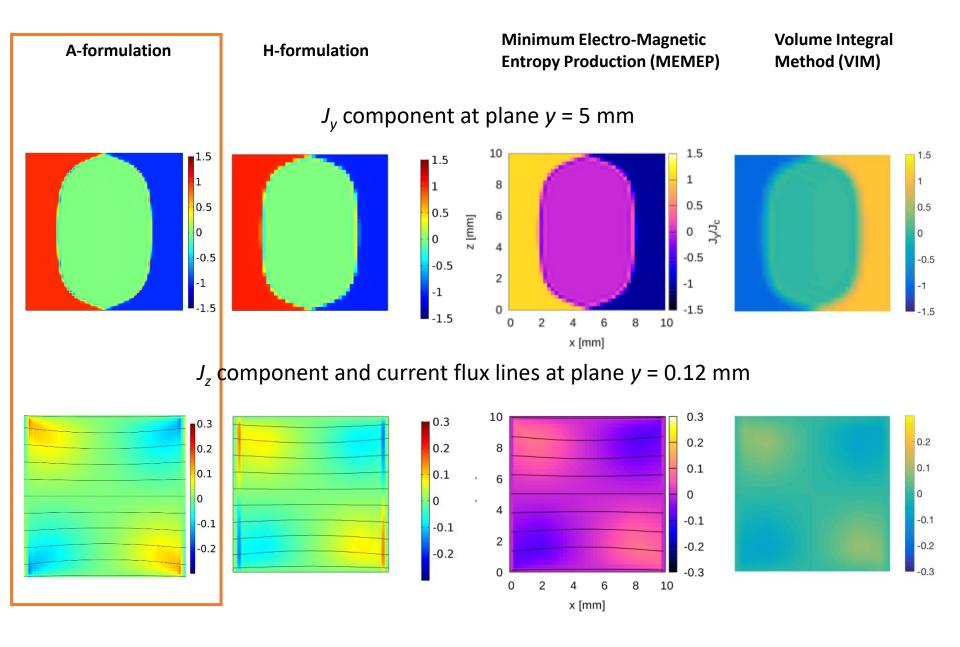
M. Kapolka, E. Pardo, V. M. R. Zermeño, F. Grilli, A. Morandi, P. L. Ribani

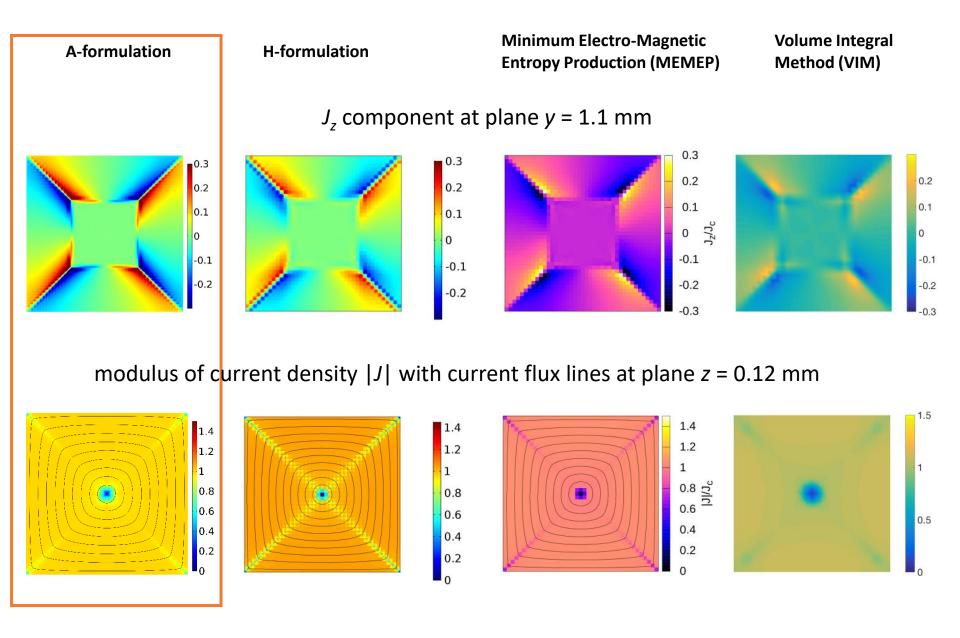
$$J_{c0} = 10^8 \text{A/m}^2$$

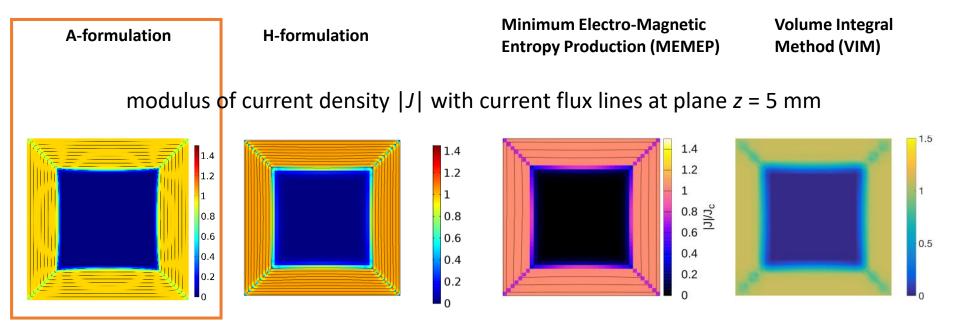
$$B_a = 200 \text{ mT} / 50 \text{ Hz}$$



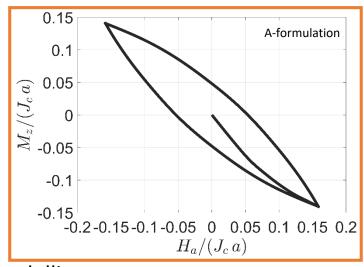
- cube was placed into the sphere with 60 mm of diameter (free tetrahedral mesh);
- cubic regular mesh 216 000 domain elements (60x60x60) for the superconductor;
- linear mesh discretization gave 1 357 645 DOF;
- **64** hours (1¼ T) on PC on Intel® Core™ i7-7740X CPU @ 4.3 GHz (8 cores) and 64GB RAM;
- Comsol Multiphysics© version 5.4 running under Windows 10 64-bit.

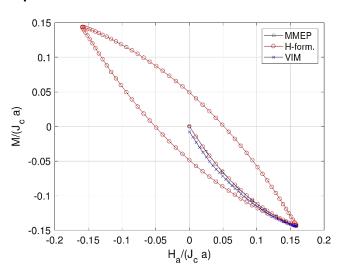




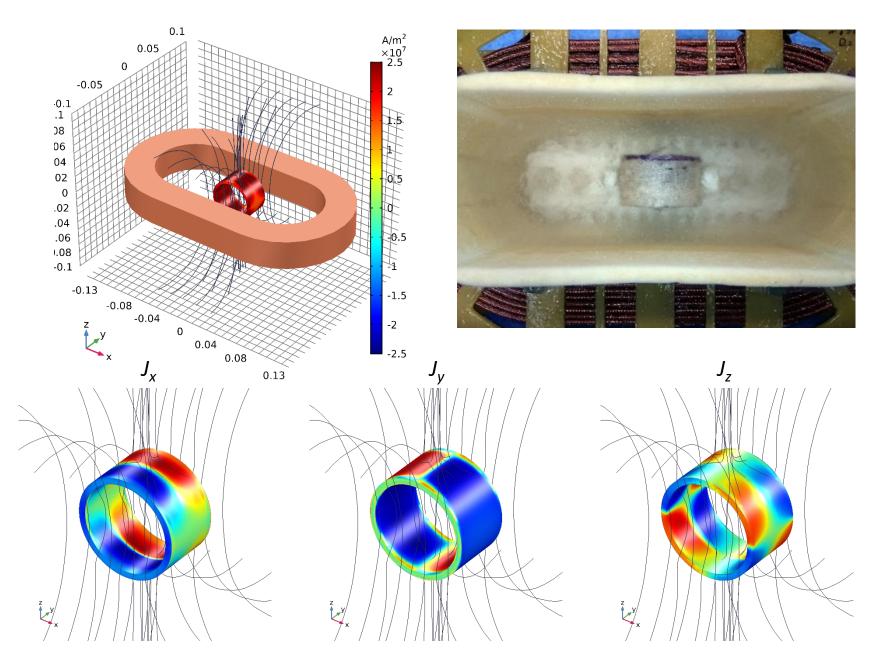


#### magnetization loop of the cube





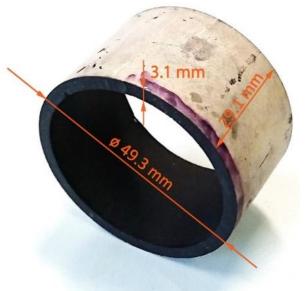
#### Experimental verification, BSCCO ring



#### Experimental verification, BSCCO ring

Melt-cast processed Bi<sub>2</sub>Sr<sub>2</sub>CaCu<sub>2</sub>O<sub>8+δ</sub> (Nexans production for FCL program<sup>3</sup>)

Electric measurements were performed on the small piece, cut from the tube along the main axis:



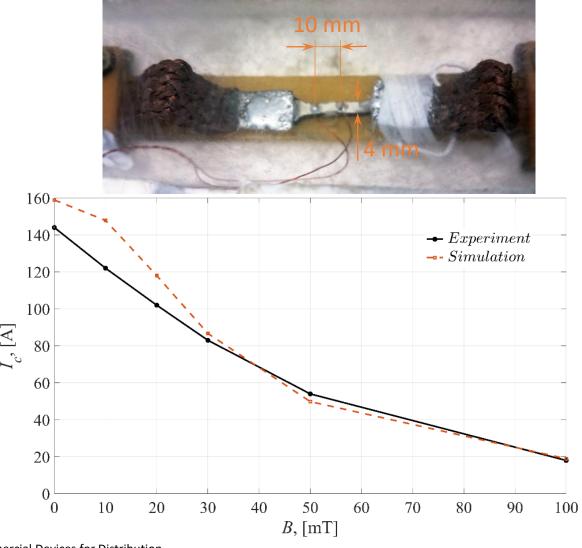
$$J_c(B) = \frac{J_{c0}}{\left(1 + \frac{|B|}{B_c}\right)^{\alpha}}$$
 (11)

Selected parameters:

$$J_{c0} = 4.8 \times 10^7 \text{ A/m}^2$$

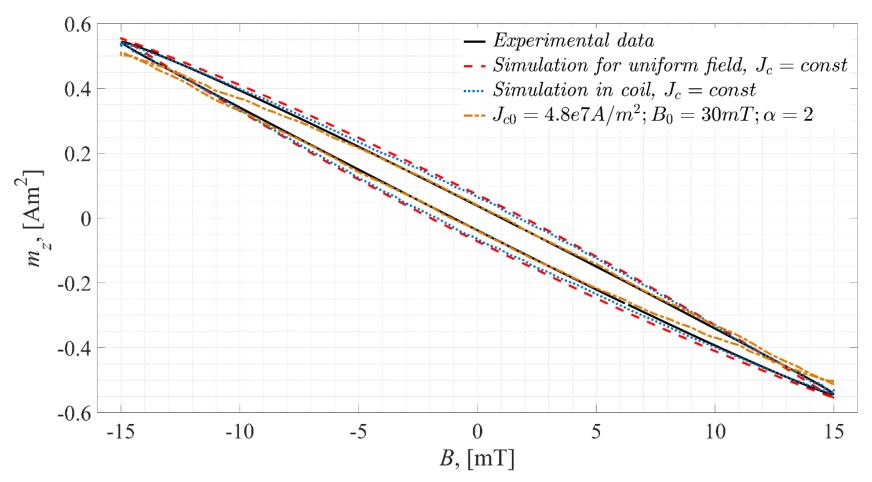
$$B_0 = 30 \text{ mT}$$

$$\alpha = 2$$



<sup>3.</sup> Bock, J. *et al.* HTS Fault Current Limiters-First Commercial Devices for Distribution Level Grids in Europe. *IEEE Trans. Appl. Supercond.* **21**, 1202–1205 (2011).

#### Experimental verification, BSCCO ring

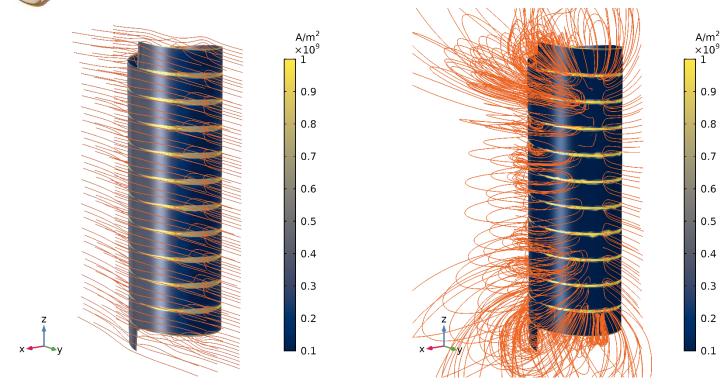


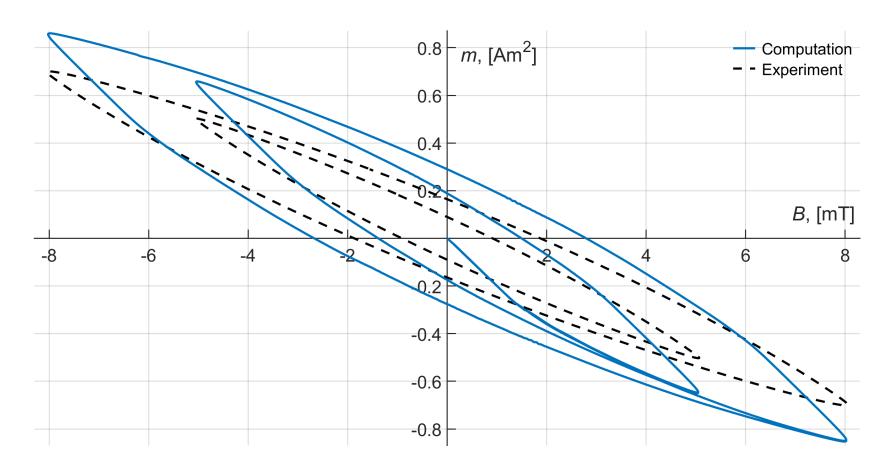
Magnetization loop measured in "calibration-free" system<sup>4</sup>, compared with the computed ones



The SF12050AP tape on tubular former

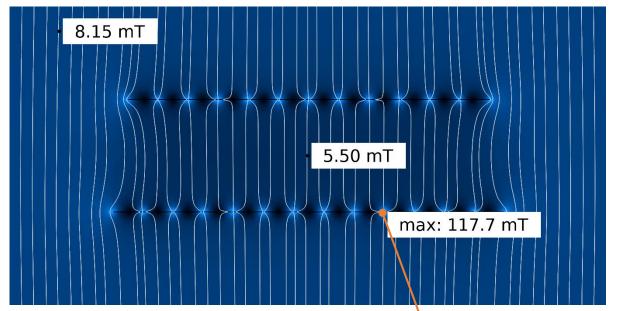
- declared minimal critical current is 387 A
- former diameter 45.1 mm
- former length 145 mm



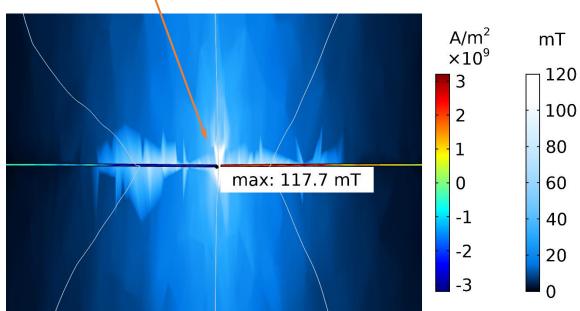


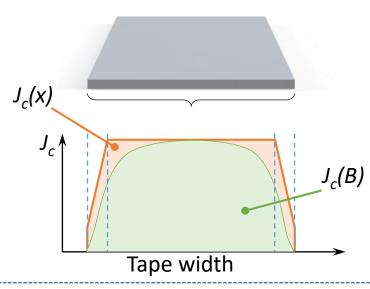
Magnetization loops obtained assuming tape width 12 mm and gap between turns 0.1 mm.

 $J_c$  value corresponds to 387 A of critical current and 10  $\mu$ m layer thickness.

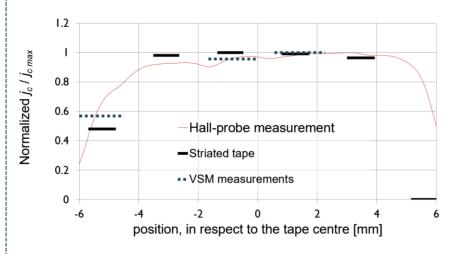


Magnetic flux density distribution in the helix cross-section



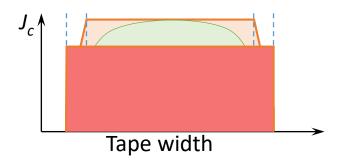


Previously measured  $J_c$  distribution across the tape width, for tape of the same manufacturer<sup>5</sup>

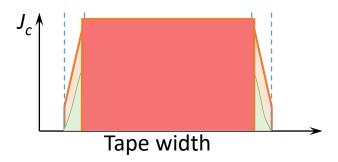


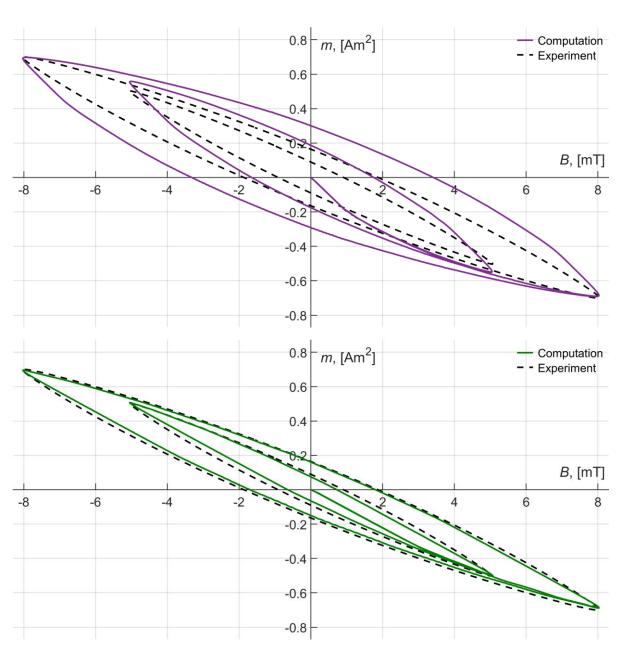
5. Solovyov, M. *et al.* Non-uniformity of coated conductor tapes. *Superconductor Science & Technology* **26**, (2013).

Using the "effective value" of critical current:



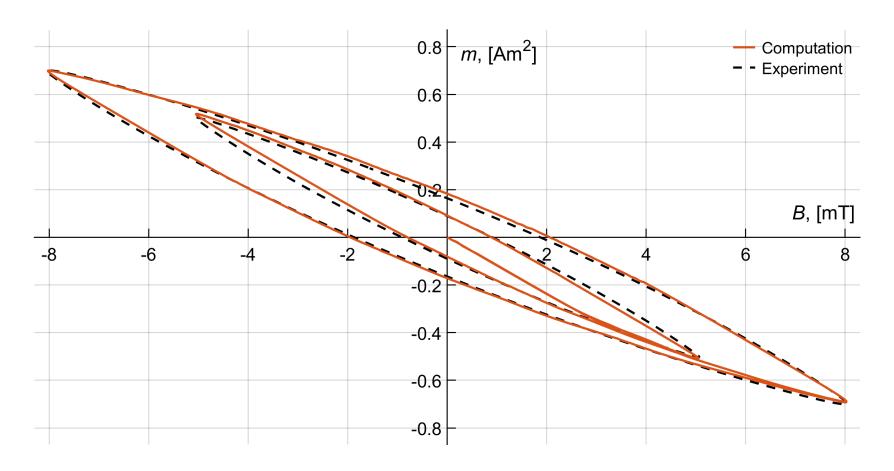
Using the "effective value" of tape width:





Magnetization loops for original geometry and 260 A of critical current

Original critical current (387 A) and tape width reduced by 1.2 mm



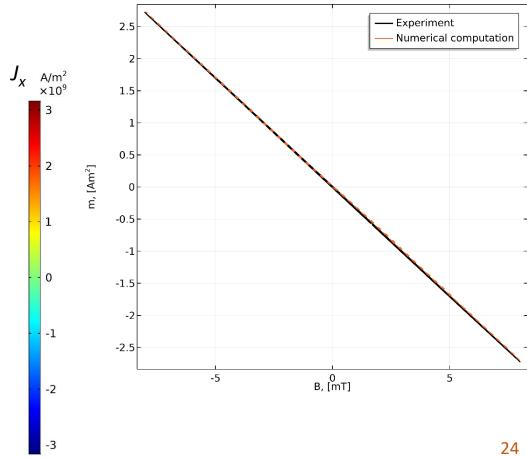
Combination of tape width reduction (by 1 mm) and critical current reduction (to 360 A)

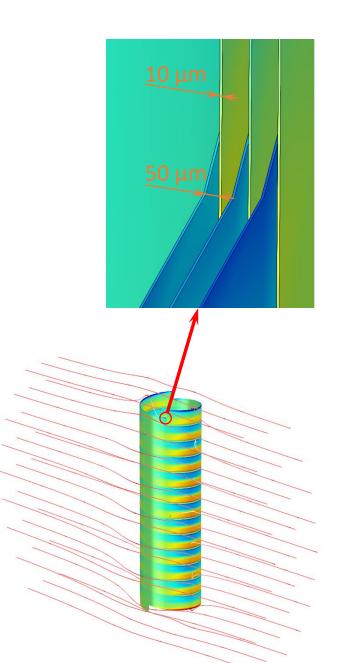
#### Conclusion

- We developed novel numerical procedure for modeling of electromagnetic problems in superconductors in 3D geometry.
- The results were validated by numerical and experimental benchmarks.
- Our method is based on A-formulation that is the native approach for electromagnetic computation used in FEM and usually is implemented better into commercial codes than other formulations.
- Unfortunately, incorporation of detailed material properties rapidly increases the model complexity. As a consequence, the computation time grows, and some time it may cause the solver convergence problem.
- A general guide to be followed in search for minimal but still valid electromagnetic representation of superconducting object in 3D geometry remains and open problem.

#### Recent results







#### Recent results

### Current transport in CORC™ cable:

