

Lorentz angle in the ATLAS pixel detector: Effects of operating parameters and radiation damage

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Workshop on Advanced Silicon Radiation Detectors
February 25, 2019

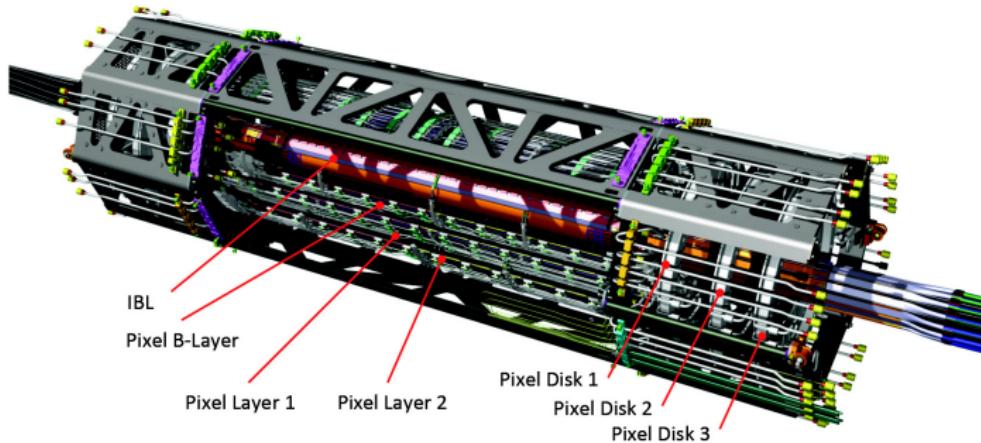


Outline of the talk

- The ATLAS pixel detector and radiation damage digitizer.
- Definition and properties of the Lorentz angle.
- Measurement strategy and fit definition.
- Measurements with cosmic muons in Run I. Temperature dependence.
- Radiation damage after Run I: Results with cosmic muons.
- Electron and hole mobility models.
 - Low-field parameterisations of $\mu(T)$.
 - High-field extrapolations. The extended Thomas model.
- Radiation damage on Run II: IBL Lorentz angle in pp collisions.
- Simulation: Chiocchia model and ATLAS radiation damage digitizer.
- Summary and conclusions.

The ATLAS pixel detector

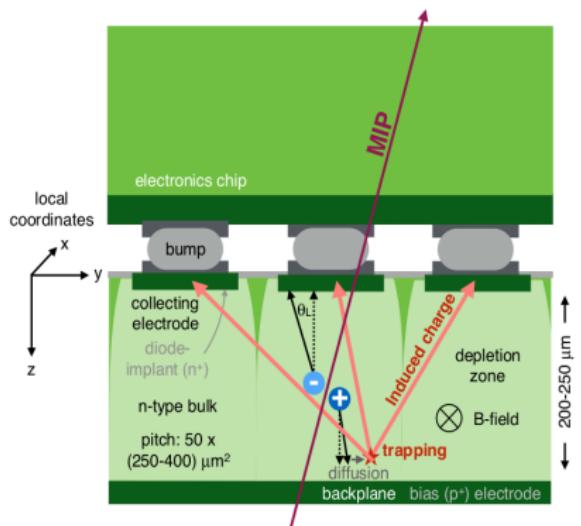
The ATLAS pixel detector is composed of four layers and three endcap disks



| Layer | $\langle r \rangle$ [mm] | Staves | Modules | Channels | Disk | $\langle z \rangle$ [mm] | Sectors | Modules | Channels |
|-------|--------------------------|--------|---------|--------------------|------|--------------------------|---------|---------|-------------------|
| 0 | 33.3 | 14 | 224 | 6.02×10^6 | 0 | 495 | 8 | 48 | 2.2×10^6 |
| 1 | 50.5 | 22 | 286 | 13.2×10^6 | 1 | 580 | 8 | 48 | 2.2×10^6 |
| 2 | 88.5 | 38 | 494 | 22.8×10^6 | 2 | 650 | 8 | 48 | 2.2×10^6 |
| 3 | 122.5 | 52 | 676 | 31.2×10^6 | | | | | |

The ATLAS pixel detector digitizer

Signal digitization is done using the ATHENA framework interfaced with G4.

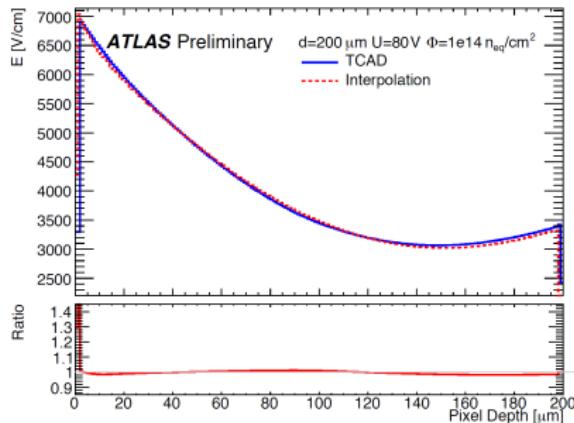
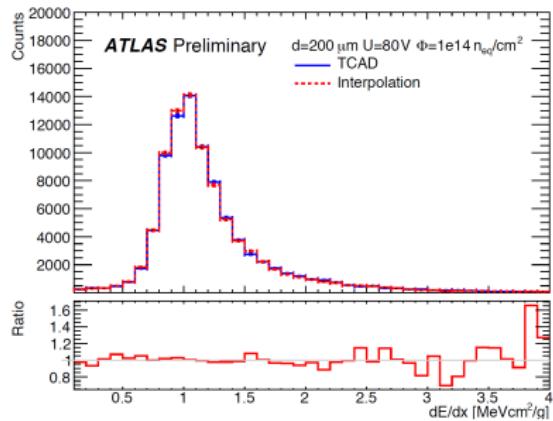


- Simulation takes energy and position of MIP as inputs.
- Ionization energy converted into electrons and holes ($\sim 3.65 \text{ eV/pair}$)
- Charges suffer thermal diffusion and drifts in E , B fields.
- Charge trapping is randomly set, by comparing trapping time (Φ -dependent) with drift time.

After digitization, the collected charge in each module is read out.

Electric field interpolation

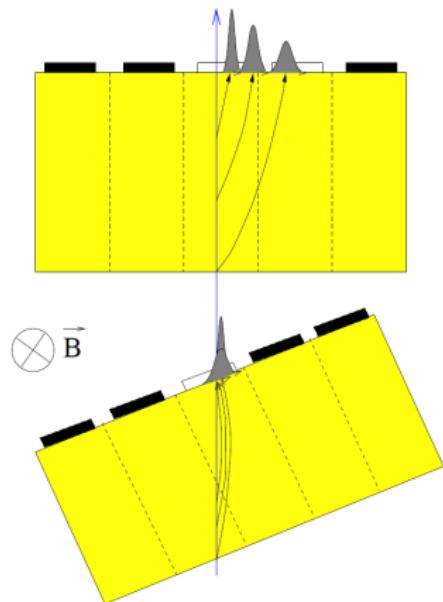
- Electric field is calculated by interpolation from benchmark simulations.
- Simulations use TCAD maps for a given voltage and fluence.
- Irradiated sensors with non-linear profiles simulated using Chiochia model.
- Interpolation is done using cubic splines in both dimensions (V, Φ)
- Closure of the interpolation and TCAD simulation is shown below



Very good closure observed between the simulated and interpolated results.

Lorentz angle: Definition and properties

Lorentz Angle: angle of deviation of drifting electrons in the Si modules under the effect of the solenoidal magnetic field



- Cluster transverse size is a function of the incidence angle φ .
- With $\vec{B} \neq 0$, the cluster size has a minimum for $\varphi = \theta_L$.
- Knowledge of θ_L can be used to improve spatial resolution.
- Important for understanding mobility models in Monte Carlo.
- Indicator of radiation damage.

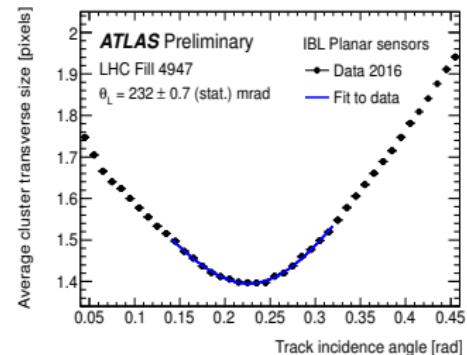
Measurement strategy

- Tracks with hits in the pixel detector are selected for the analysis.
- Cluster size is presented as a function of the incidence angle in $R\phi$ plane.
- The distributions are fitted to the following function

$$W(\varphi) = \left(a |\tan \varphi - \tan \theta_L| + \frac{b}{\sqrt{\cos \varphi}} \right) \otimes G(\varphi),$$

where \otimes is the convolution. The free parameters are a , θ_L , b and σ :

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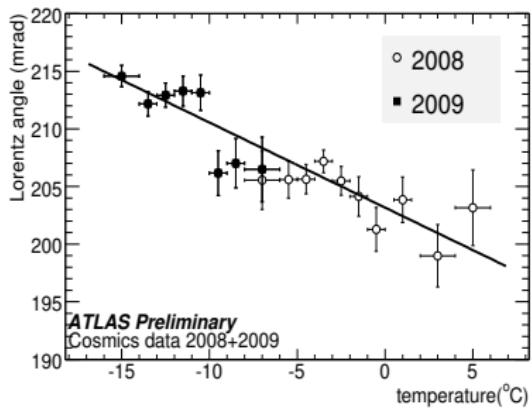
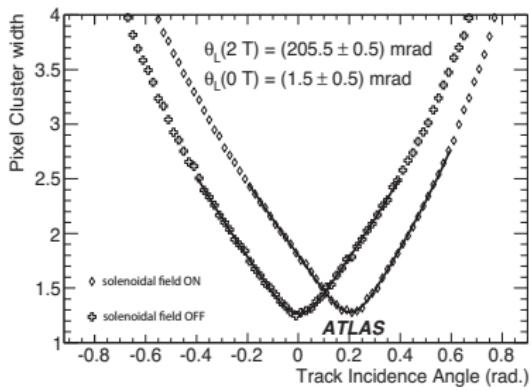


- $a \equiv$ Related to the depletion depth.
- $\theta_L \equiv$ The Lorentz angle.
- $b \equiv$ The minimum cluster size.
- $\sigma \equiv$ The resolution on the incidence angle.

First measurements with cosmic muons (2008-2009)

Lorentz angle from cosmic muons [arXiv:1004.5293]

- Fixed conditions: $V = -150$ V, $T = -3$ °C.
- Temperature scan keeping $V = -150$ V.



Measured dependence with T : $\frac{\partial \theta_L}{\partial T} = (-0.74 \pm 0.03) \text{ mrad/K}$

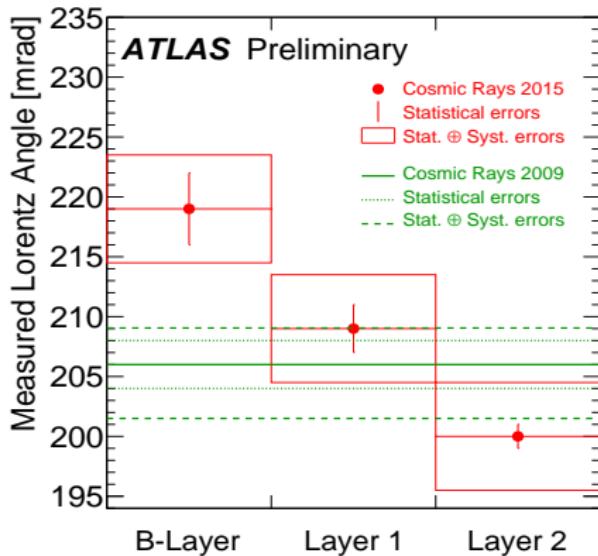
Predicted dependence with T : $\frac{\partial \theta_L}{\partial T} = -0.74 \text{ mrad/K}$

[C. Jacoboni *et al.*: Solid State Electronics 20, 77 (1977)]

Lorentz angle in cosmic muons (2015)

Lorentz angle in cosmic data taken in two different periods ([IDTR-2017-002](#))

- Cosmics 2009: Inclusively in all layers.
- Cosmics 2015: Separated in three layers (L1 and L2 compatible with 2009)



- Same temperature and voltage conditions: $T = -10 \text{ }^{\circ}\text{C}$, $V = -150 \text{ V}$.
- After Run-I, B -layer received more radiation than L1, L2 \Rightarrow Larger θ_L

Modelling the Lorentz angle

- Equivalent to modelling the mobility μ ($\tan \theta_L \sim \mu |\vec{B}|$).
- 10% discrepancy found in Run I: [ATL-INDET-PUB-2018-001](#)
- Mobility $\mu(T, \vec{E})$ defines the drift velocity: $\vec{v}_d = \mu \vec{E}$.
- Parameterisations of electron and hole low and high-field mobility:
 - Low-field mobility parameterised by a power law:

$$\mu(T) = a T_n^{-b}; \quad T_n = \frac{T}{300 \text{ K}}$$

- High-field mobility extrapolated using [Thomas parameterisation](#)

$$\mu(T, E) = \mu_0(T) \left[1 + \left(\frac{\mu_0(T) E}{\nu_s(T)} \right)^\beta \right]^{-\frac{1}{\beta}}$$

Electron and hole mobility: Low-field models

For low electric field, three main models are studied

- Jacoboni-Canali:

$$\mu_{e,0}(T) = 1533.7 \text{ cm}^2/(\text{V}\cdot\text{s}) T_n^{-2.42}$$

$$\mu_{h,0}(T) = 463.9 \text{ cm}^2/(\text{V}\cdot\text{s}) T_n^{-2.20}$$

- Canali:

$$\mu_{e,0}(T) = 1437.7 \text{ cm}^2/(\text{V}\cdot\text{s}) T_n^{-2.42}$$

$$\mu_{h,0}(T) = 463.9 \text{ cm}^2/(\text{V}\cdot\text{s}) T_n^{-2.20}$$

- Hamburg-Thomas

$$\mu_{e,0}(T) = 1440(15) \text{ cm}^2/(\text{V}\cdot\text{s}) T_n^{-2.260(7)}$$

$$\mu_{h,0}(T) = 474(10) \text{ cm}^2/(\text{V}\cdot\text{s}) T_n^{-2.619(7)}$$

Electron and hole mobility: High-field models

High-field mobility is extrapolated using

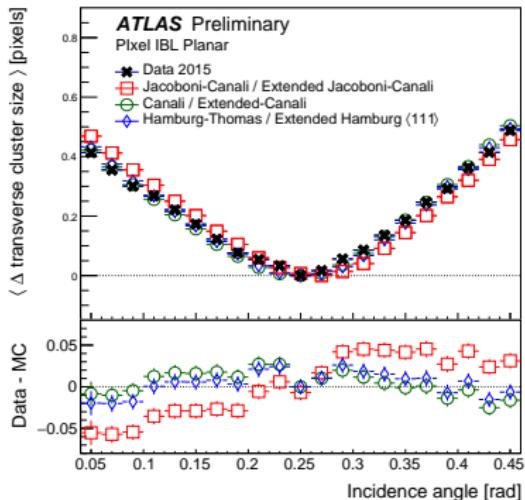
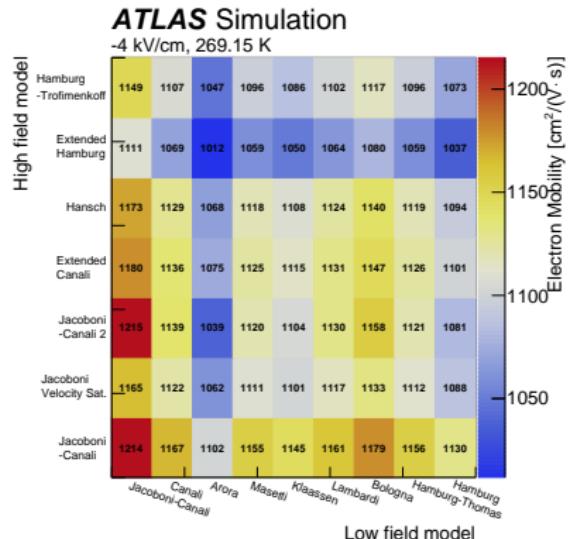
$$\mu(T, E) = \mu_0(T) \left[1 + \left(\frac{\mu_0(T)E}{\nu_s(T)} \right)^\beta \right]^{-\frac{1}{\beta}}$$

The expressions of ν_s and β for electrons and holes are given below

| Extended Model | Parameter | Electrons | Holes |
|-----------------|----------------|--|---|
| Jacoboni-Canali | ν_s (cm/s) | $1.07 \times 10^7 \times T_n^{-0.87}$ | $8.34 \times 10^6 \times T_n^{-0.52}$ |
| | β | $1.109 \times T_n^{0.66}$ | $1.213 \times T_n^{0.17}$ |
| Canali | ν_s (cm/s) | $1.00 \times 10^7 \times T_n^{-0.87}$ | $8.34 \times 10^6 \times T_n^{-0.52}$ |
| | β | $1.109 \times T_n^{0.66}$ | $1.213 \times T_n^{0.17}$ |
| Hamburg | ν_s (cm/s) | $1.054(38) \times 10^7 \times T_n^{-0.602(3)}$ | $9.40(27) \times 10^6 \times T_n^{-0.226(2)}$ |
| | β | $0.992(4) \times T_n^{0.572(3)}$ | $1.181 \times T_n^{0.644(3)}$ |

Mobility models confronted to IBL data

- Mobility at 80 V and -4 °C predicted by many different models (left)
- Description of the data made by the three models under discussion (right)

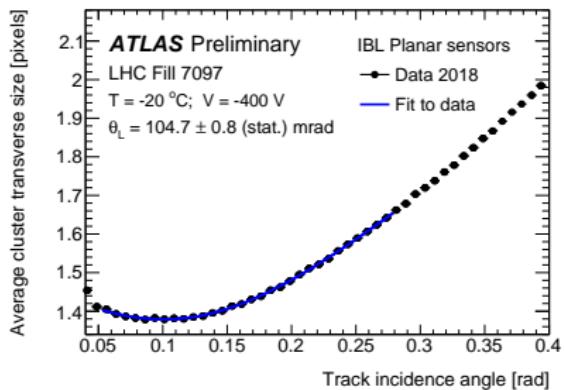
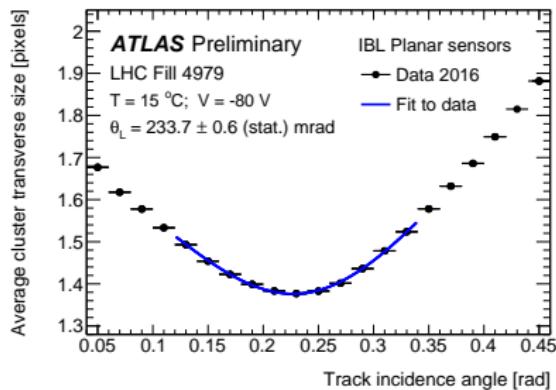


Canali / Extended Canali model is used for Run-II simulation and beyond.

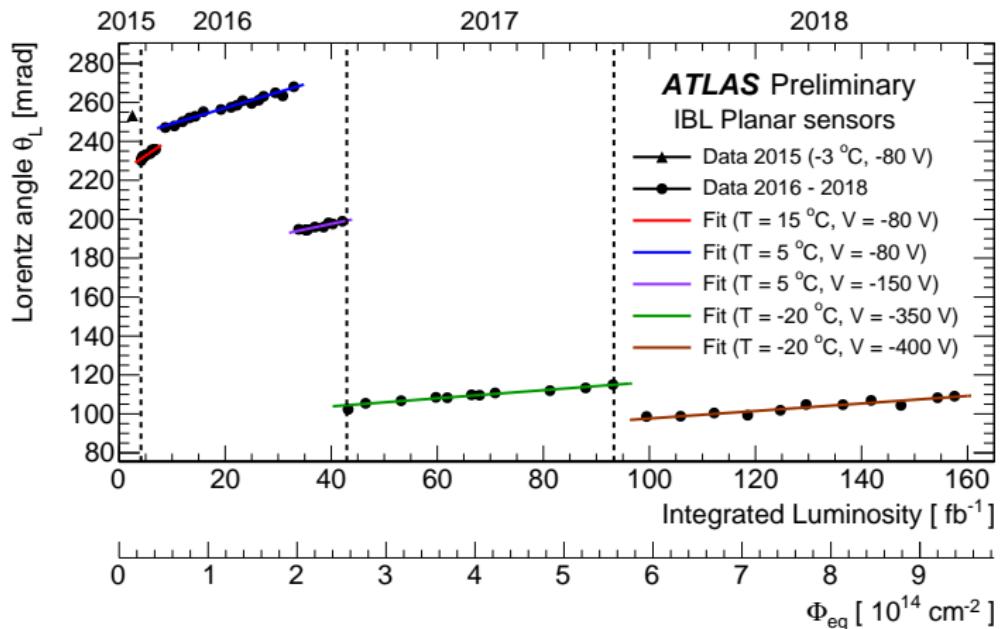
Measurement of the IBL Lorentz angle over Run II

- Measurements of θ_L performed over Run II as a function of the luminosity.
- Different periods with different temperature and voltage conditions .

| Year | Bias voltage [V] | Temperature [$^{\circ}\text{C}$] | Luminosity [fb^{-1}] |
|------|------------------|------------------------------------|---------------------------------|
| 2015 | -80 | -3 | 4.1 |
| 2016 | -80 | +15 | 3.0 |
| | -80 | +5 | 26.2 |
| | -150 | +5 | 9.6 |
| 2017 | -350 | -20 | 50.4 |
| 2018 | -400 | -20 | 64.2 |



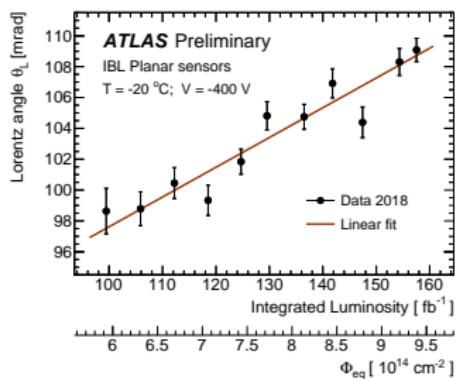
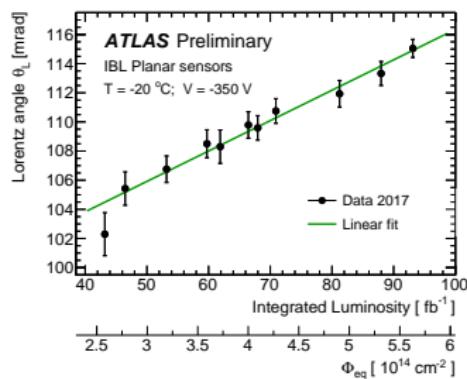
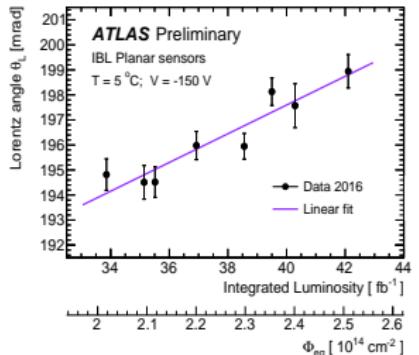
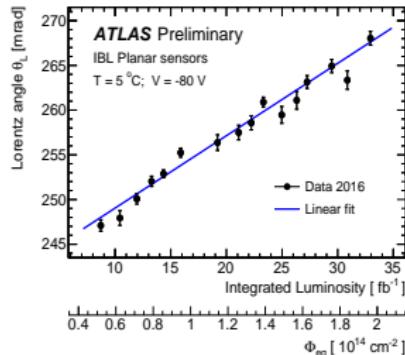
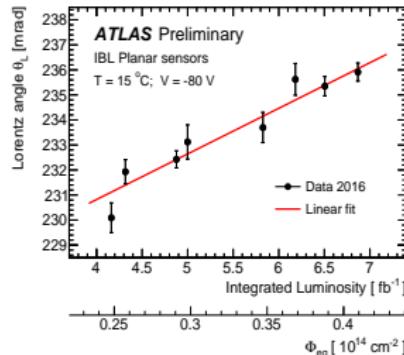
Lorentz angle evolution over Run II



- Linear, increasing trend with $\mathcal{L} \propto \Phi_{eq}$ observed in each period.
- Value of θ_L increases when decreasing temperature.
- Value of θ_L decreases when increasing voltage.
- Slope is milder with increasing voltage (decreasing temperature)

Lorentz angle evolution over Run II

Detailed variation of θ_L for each separate period



Fit results

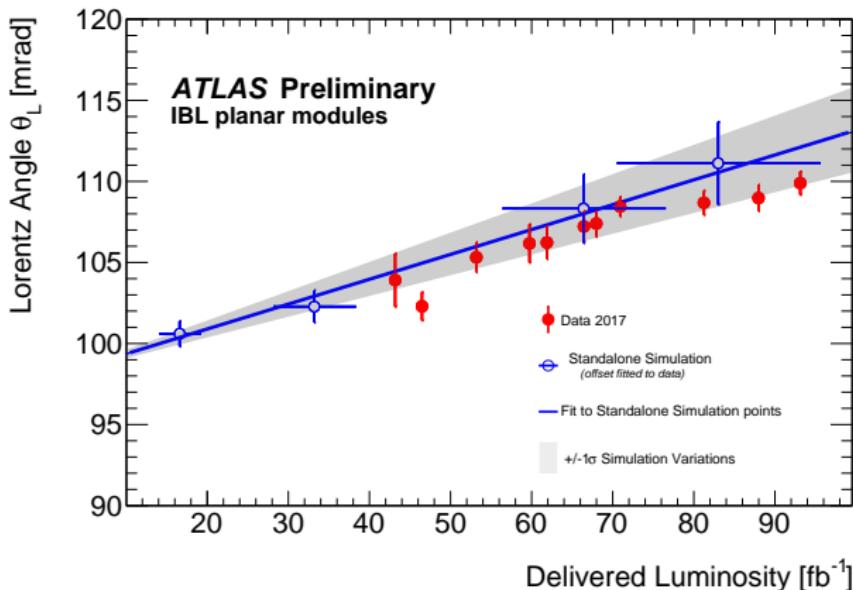
Results of the linear fits to each period are summarised below

| Temperature | Voltage | $\theta_L(\Phi_{\text{eq}} = 0)$ [mrad] | $(\partial \theta_L / \partial \Phi_{\text{eq}})_{T,V}$ [mrad · cm ²] |
|-------------|---------|---|---|
| 15 °C | 80 V | 223.5 ± 1.0 | $(30.6 \pm 3.0) \cdot 10^{-14}$ |
| 5 °C | 80 V | 240.9 ± 0.7 | $(13.6 \pm 0.6) \cdot 10^{-14}$ |
| | 150 V | 174.6 ± 3.6 | $(9.6 \pm 1.6) \cdot 10^{-14}$ |
| -20 °C | 350 V | 95.5 ± 1.3 | $(3.5 \pm 0.3) \cdot 10^{-14}$ |
| | 400 V | 78.3 ± 2.8 | $(3.2 \pm 0.4) \cdot 10^{-14}$ |

- Lorentz angle at zero fluence increases with decreasing temperature.
- Lorentz angle at zero fluence decreases with increasing voltage.
- Radiation damage degradation milder for lower temperatures.
- Radiation damage degradation milder for higher voltages.

Simulation results compared to data

Chiocchia model is used as input to ATLAS radiation damage digitizer.



- Vertical error bars on data are the fit results.
- Slope is fitted to the simulation, while offset is fitted from data.
- Grey uncertainty band: statistical and systematic (simulation parameters).
- Horizontal error bars: luminosity to fluence conversion.

Summary and conclusions

- Measurements and simulations of the Lorentz angle have been presented.
- Hints of radiation damage on *B*-Layer observed in cosmic data after Run I.
- Different mobility models have been investigated.
- Best description by Canali + Extended Canali model, used from now on.
- Measurements of the Lorentz angle as a function of the fluence presented.
- Good description of the effects by the Chiocchia model is observed.