

$H \rightarrow b\bar{b}$  @ N<sup>3</sup>LO

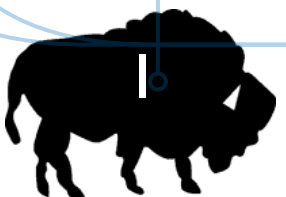
Ciaran Williams (Buffalo)

1904.08960 (with R. Mondini and M. Schiavi)

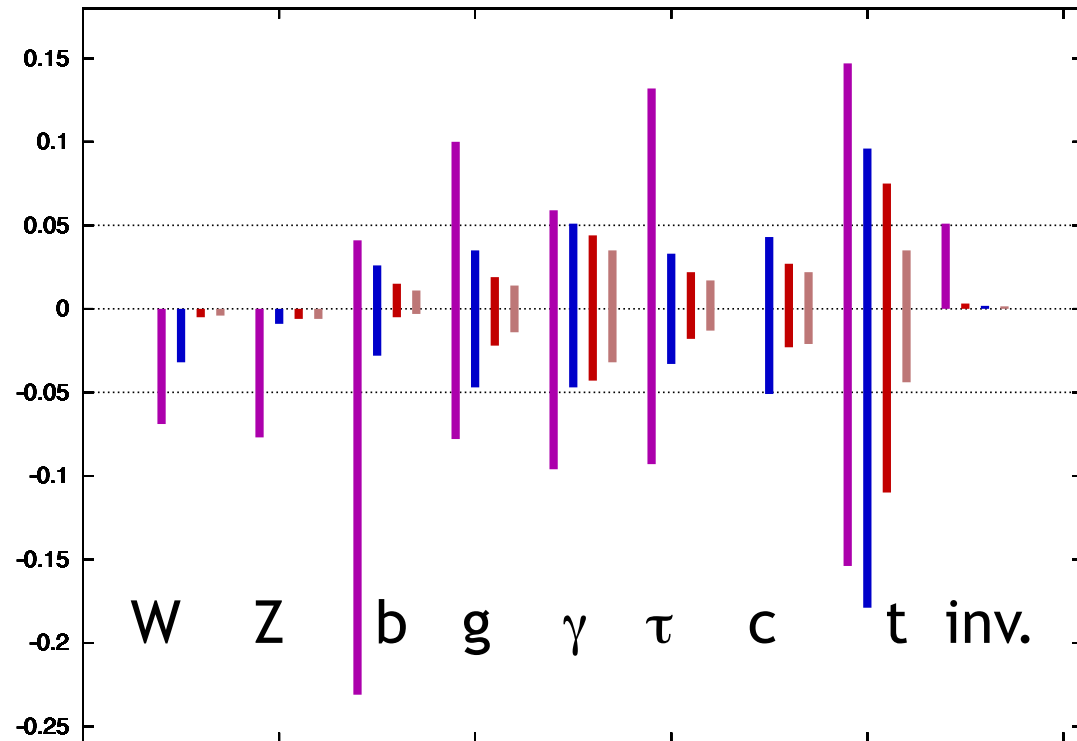
1904.08961 (with R. Mondini)



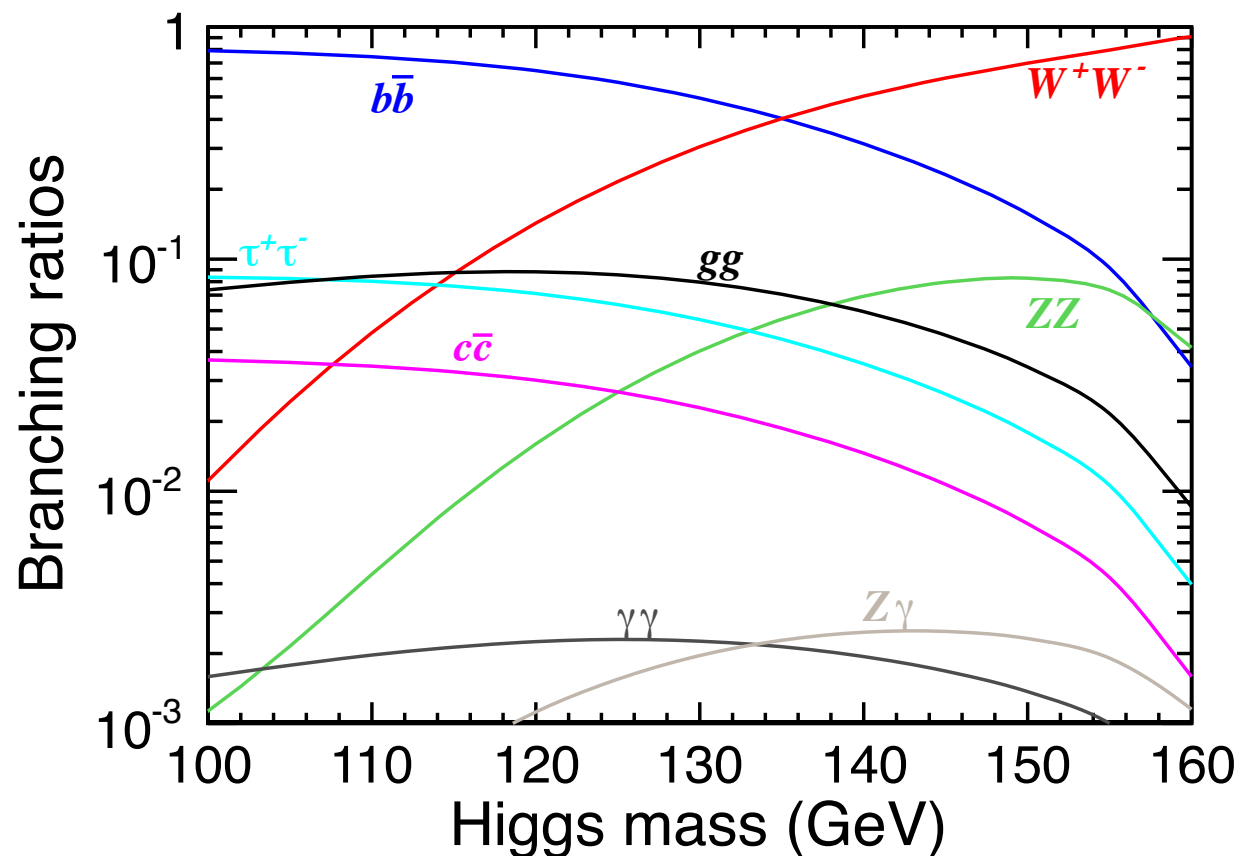
We're grateful for funding from  
the National Science Foundation



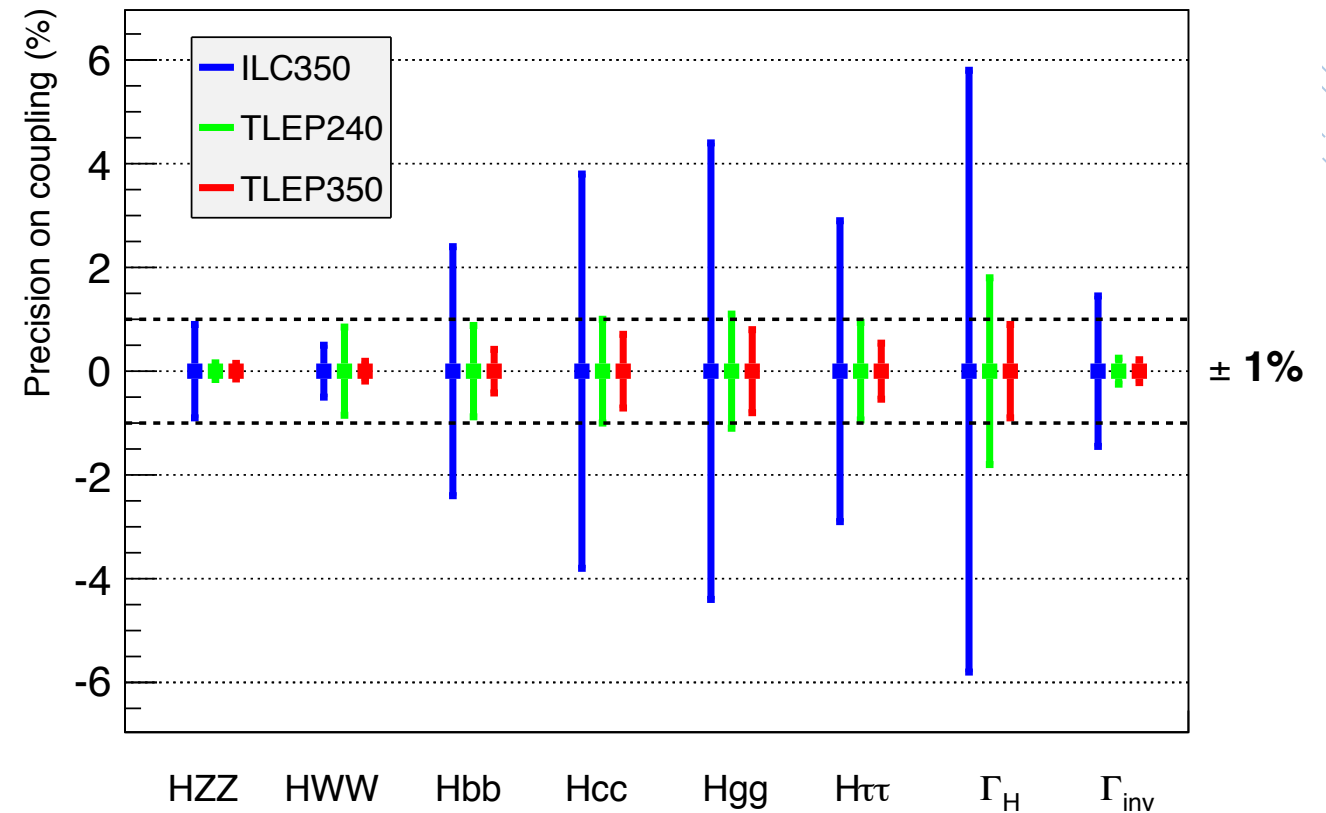
$g(hAA)/g(hAA)|_{SM} - 1$  LHC/ILC1/ILC/ILCTeV



ILC 1308.6176



TLEP 1308.6176

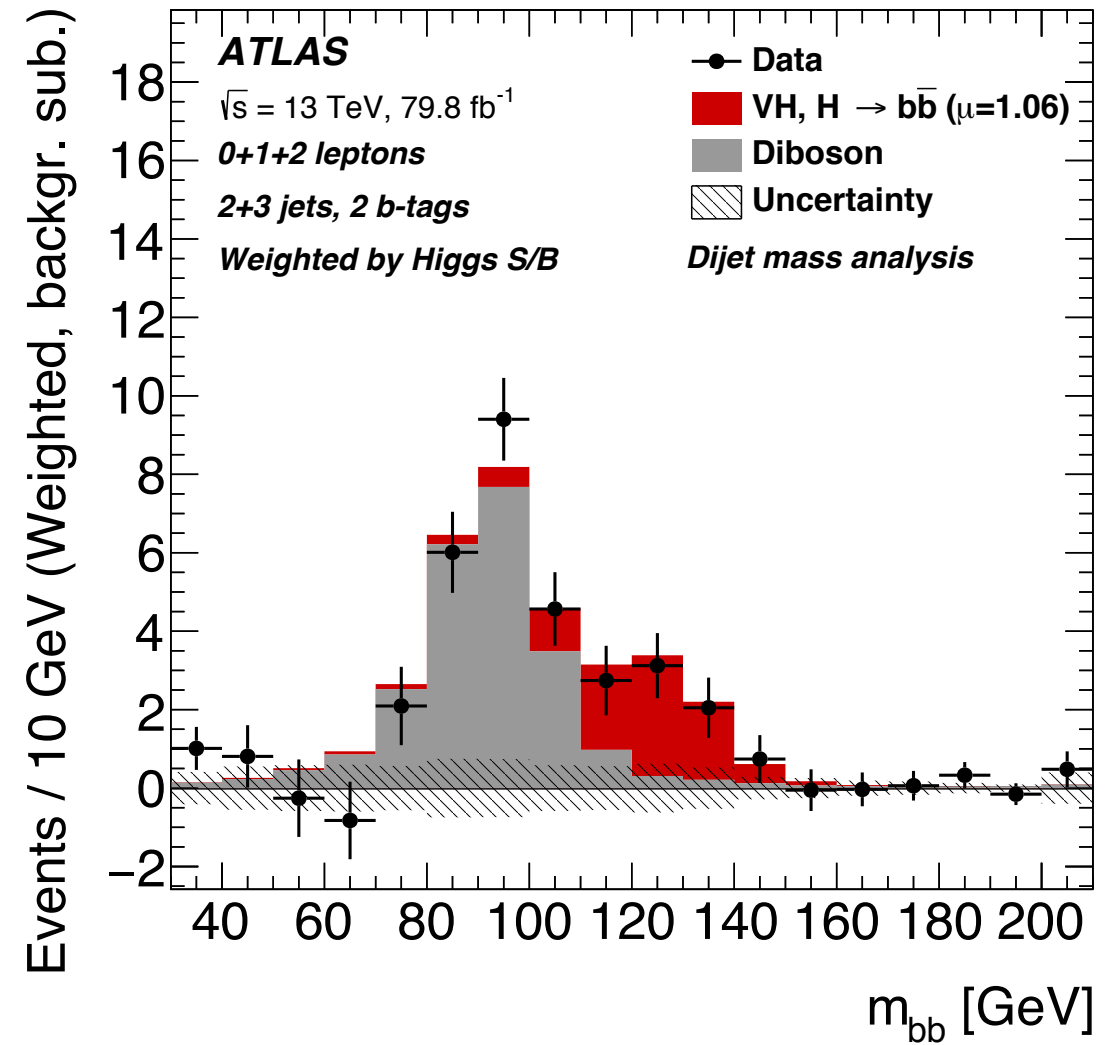
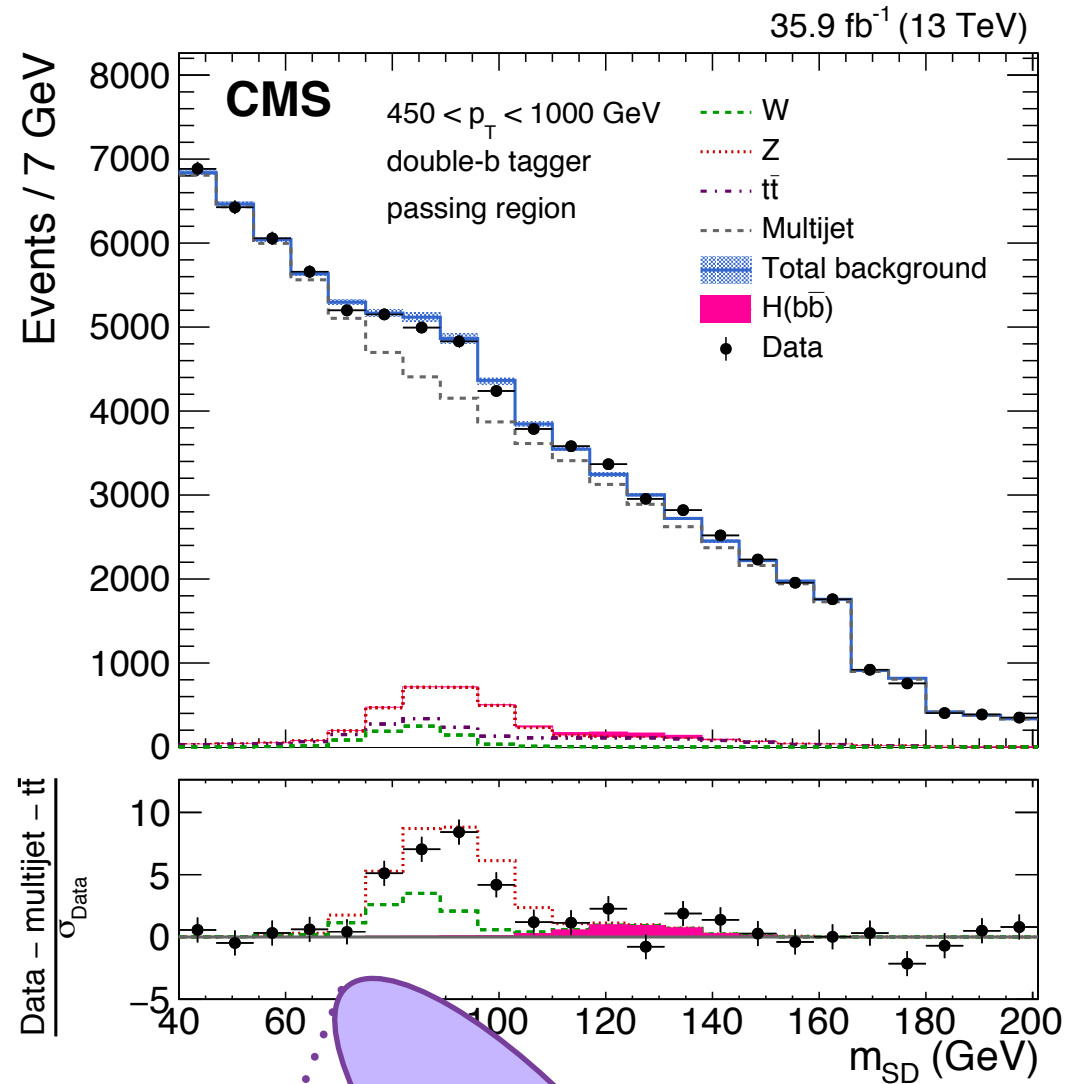


$H \rightarrow b\bar{b}$  is the largest BR for the Higgs.

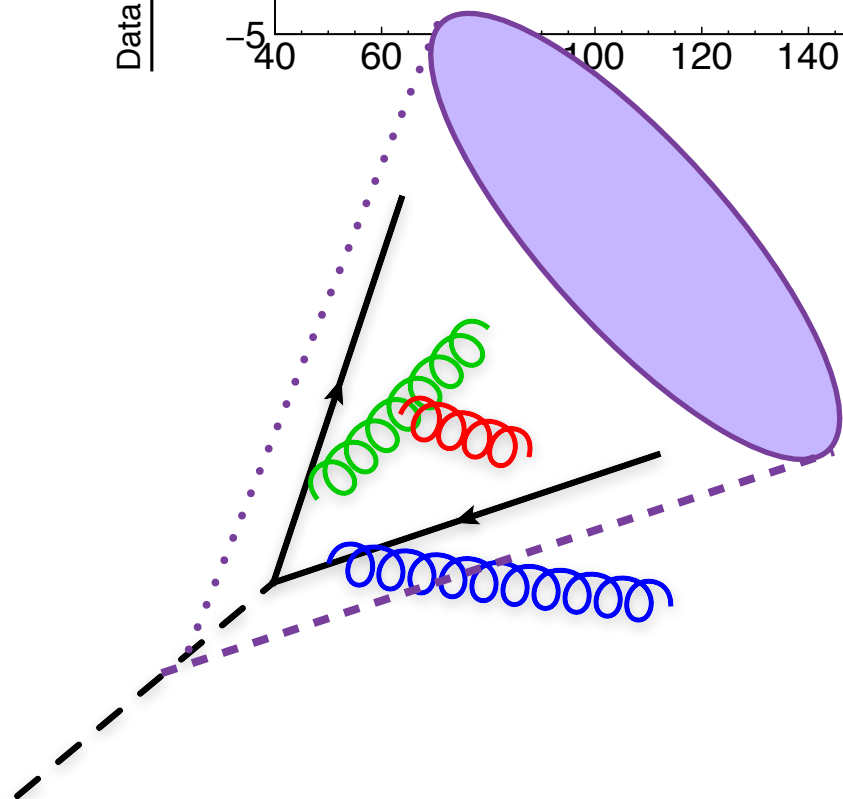
It affects every (on-shell) Higgs measurement through the width.

Will be measured, to sub percentage at a future Lepton Collider





Traditional and boosted searches at the LHC also need precise QCD predictions.



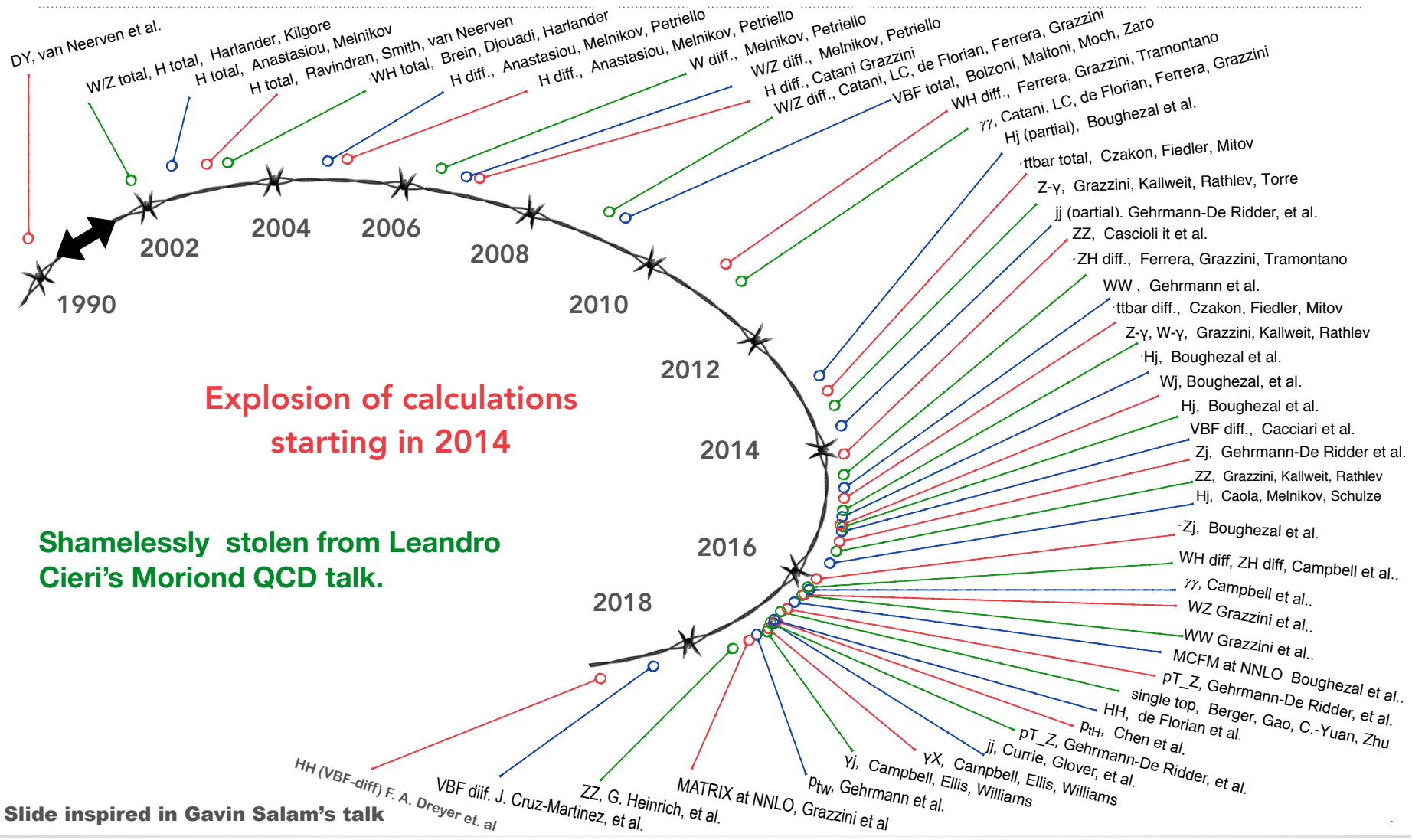




Its good to have an indoor project



# NNLO HADRON-COLLIDER CALCULATIONS VS. TIME



# Progress in QCD

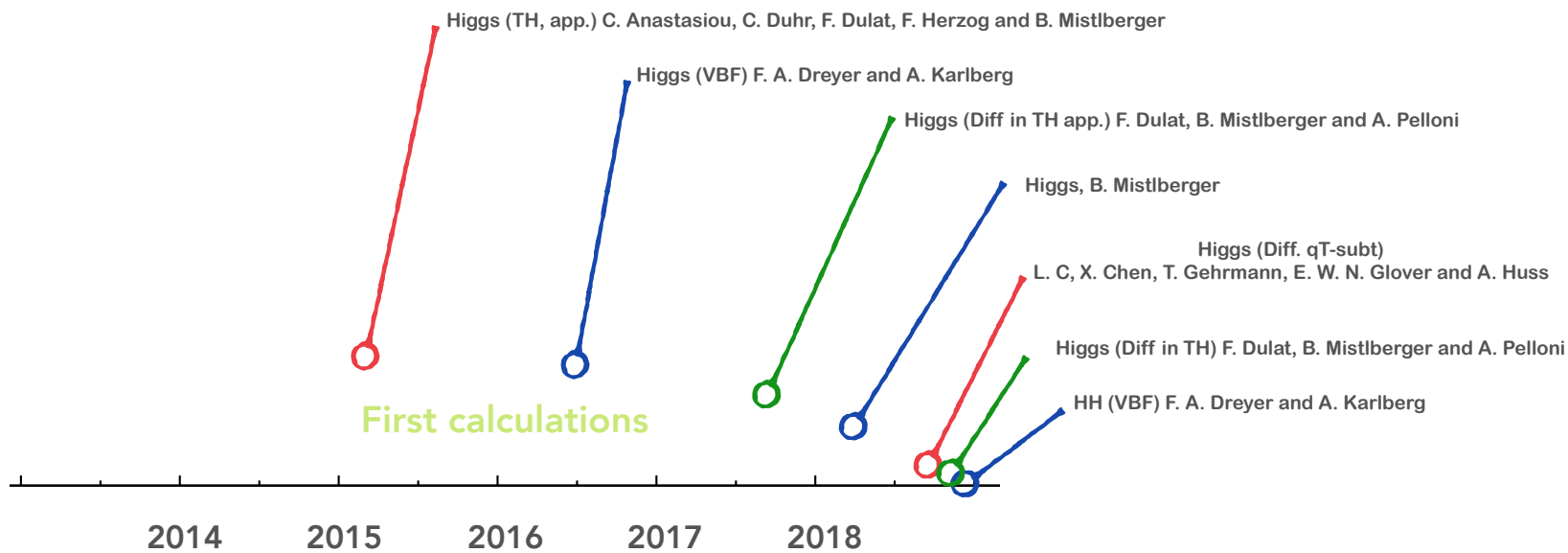
NNLO 2->2 for LHC now standard.





**N<sup>3</sup>LO HADRON-COLLIDER CALCULATIONS VS. TIME**

Shamelessly stolen from Leandro Cieri's Moriond QCD talk.



N3LO is still something we are learning to do!

Handful of different calculations for cross sections.

Higgs production known differentially.



Coefficient of N3LO correction to the width is

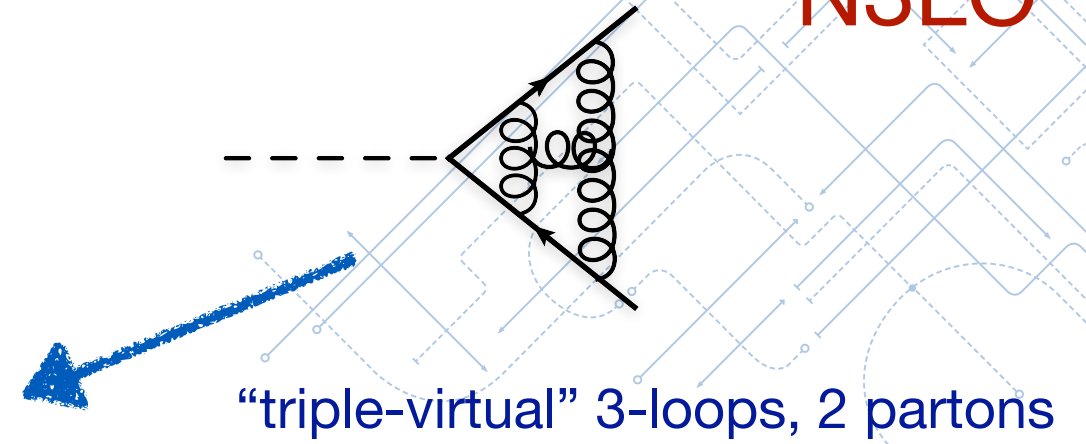
$$\begin{aligned}
 \frac{d \Delta \Gamma_{H \rightarrow b\bar{b}}^{\text{N3LO}}}{d \mathcal{O}_m} &= \int d\Gamma_{H \rightarrow b\bar{b}}^{VVV} F_2^m(\Phi_2) d\Phi_2 \\
 &+ \int d\Gamma_{H \rightarrow b\bar{b}}^{RVV} F_3^m(\Phi_3) d\Phi_3 \\
 &+ \int d\Gamma_{H \rightarrow b\bar{b}}^{RRV} F_4^m(\Phi_4) d\Phi_4 \\
 &+ \int d\Gamma_{H \rightarrow b\bar{b}}^{RRR} F_5^m(\Phi_5) d\Phi_5
 \end{aligned}$$





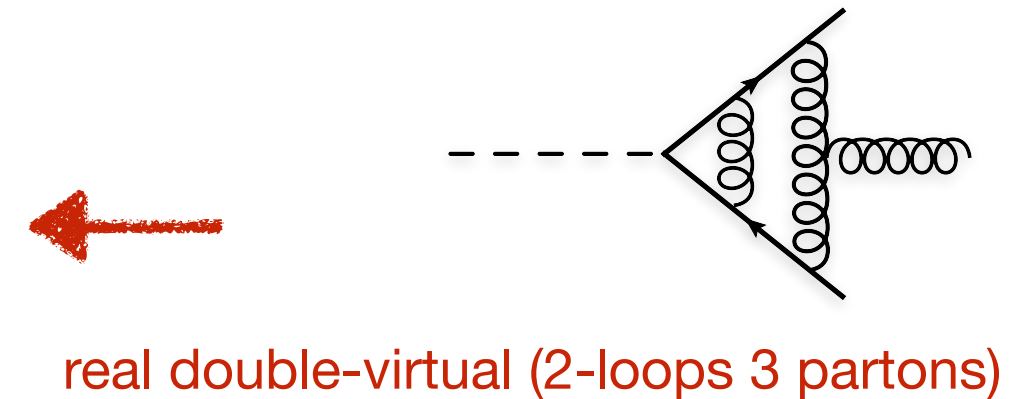
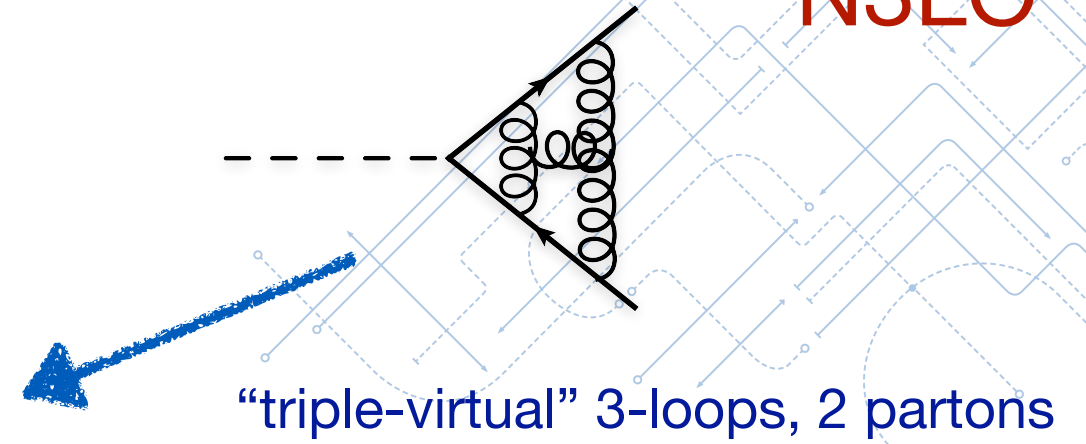
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 \end{aligned}$$



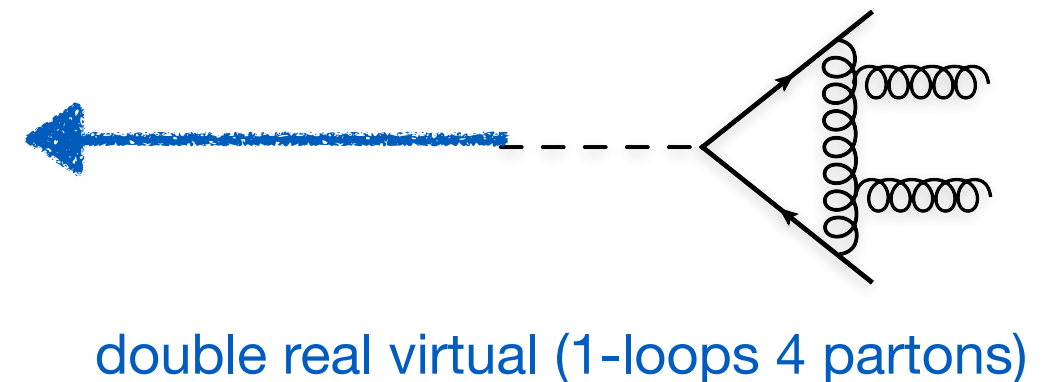
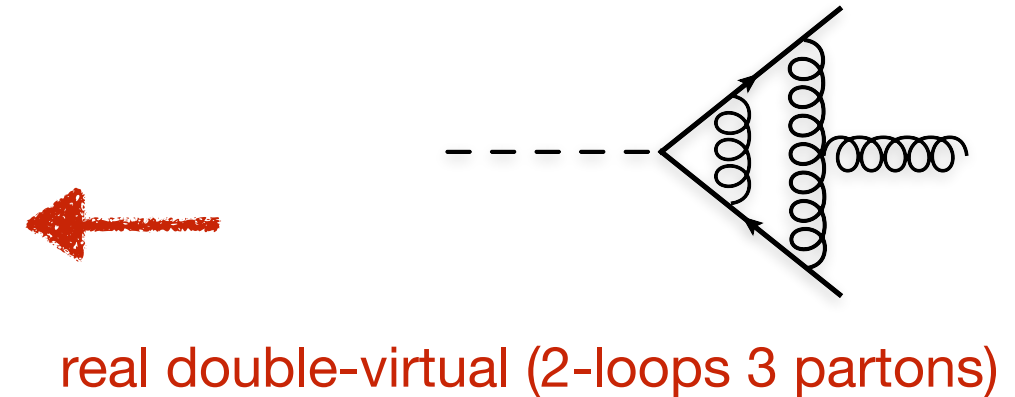
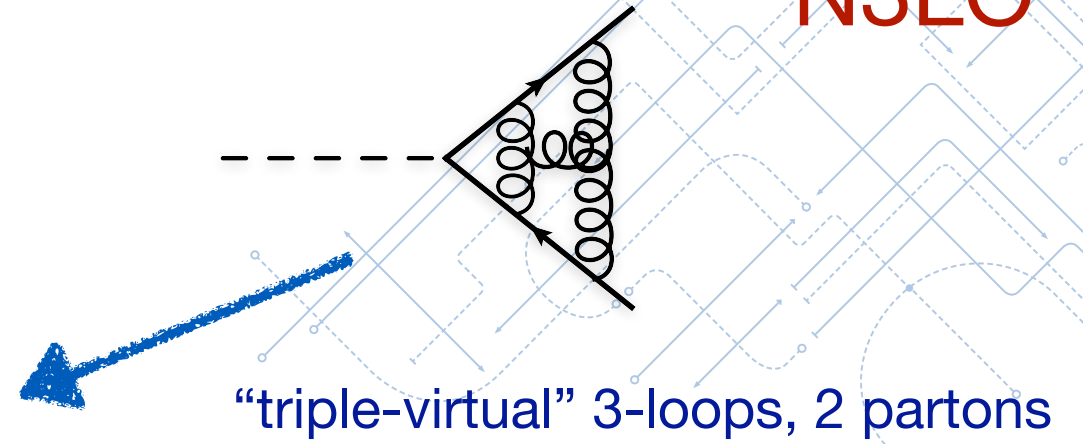
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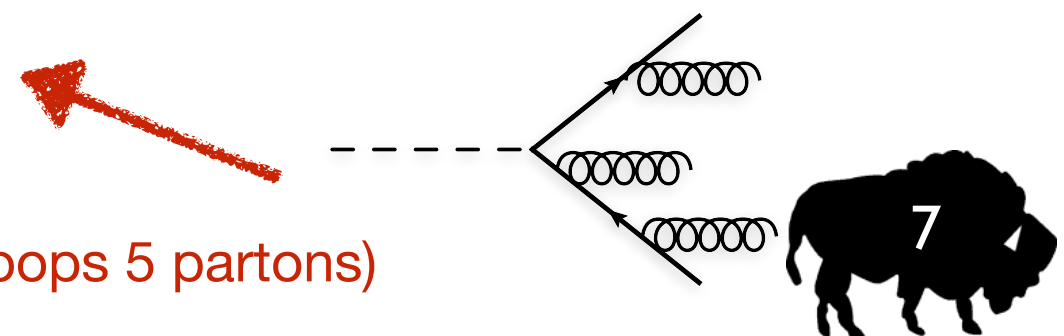
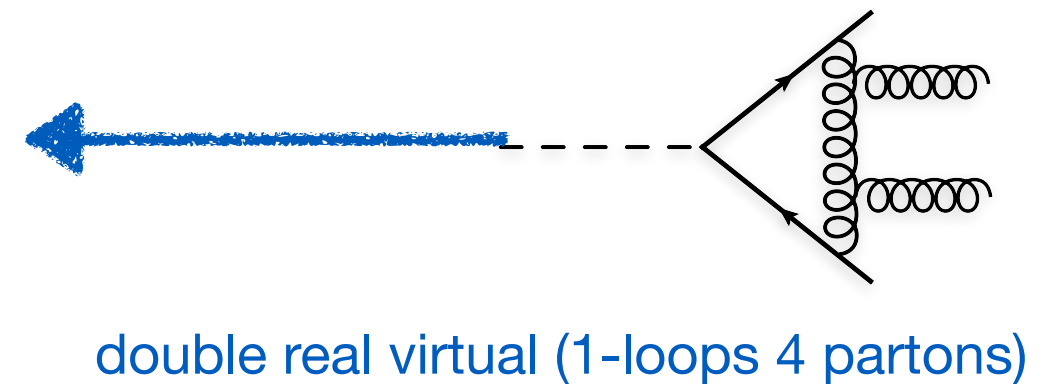
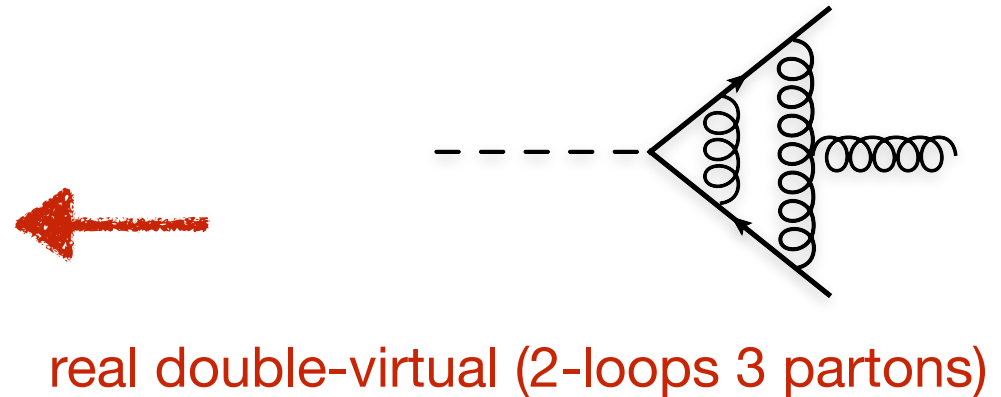
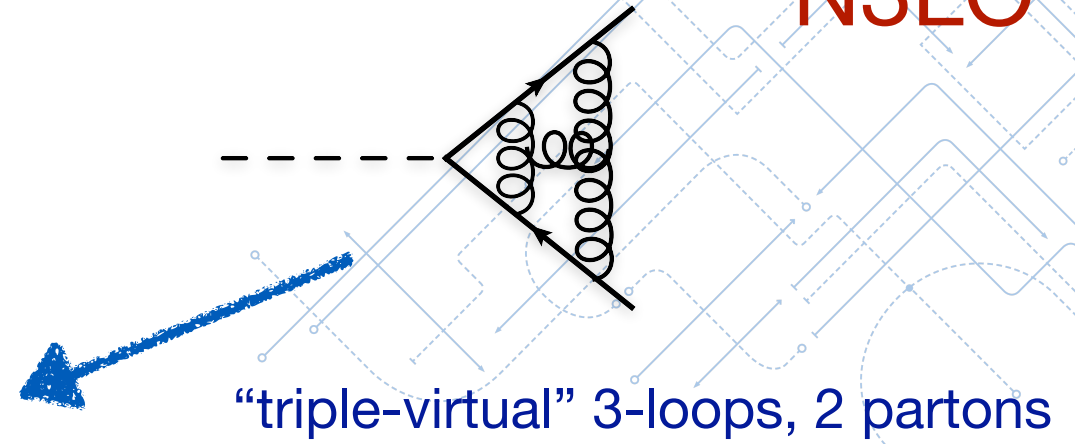
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$$+ \int d\Gamma_{H \rightarrow b\bar{b}}^{RVV} F_3^m(\Phi_3) d\Phi_3$$

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$$+ \int d\Gamma_{H \rightarrow b\bar{b}}^{RRR} F_5^m(\Phi_5) d\Phi_5$$



tripe real -virtual (0-loops 5 partons)

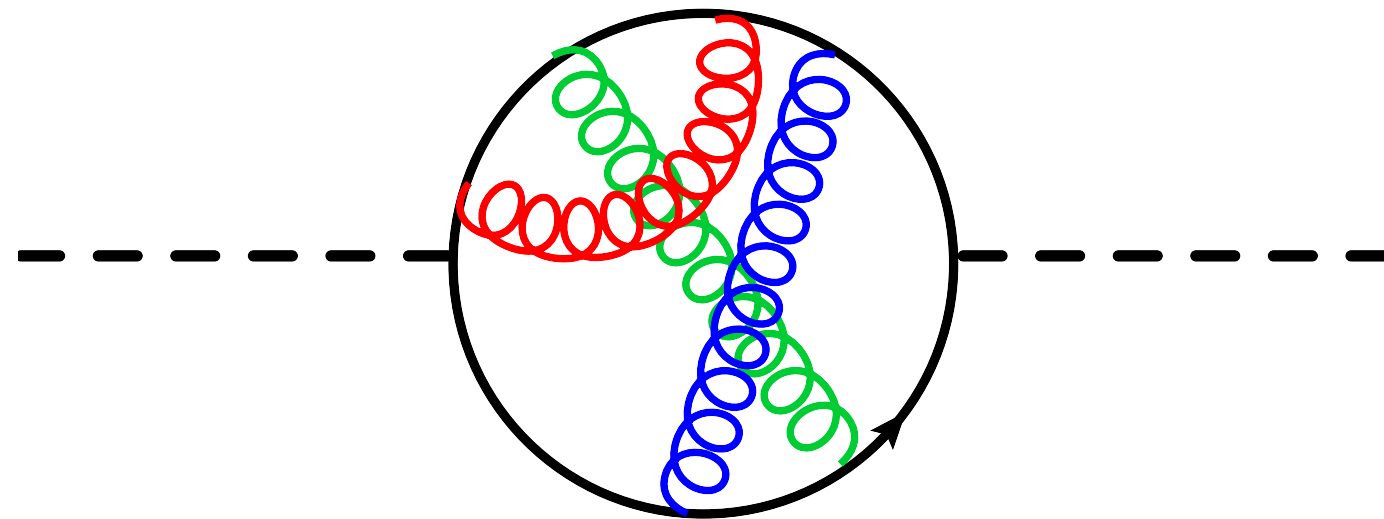
$F_i^m(\Phi_i)$  defines the measurement function and tell us how to make m-jets out of i partons.

$$\begin{aligned}
 \frac{d \Delta\Gamma_{H \rightarrow b\bar{b}}^{\text{N3LO}}}{d\mathcal{O}_m} &= \int d\Gamma_{H \rightarrow b\bar{b}}^{VVV} F_2^m(\Phi_2) d\Phi_2 \\
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 &+ \int d\Gamma_{H \rightarrow b\bar{b}}^{RRR} F_5^m(\Phi_5) d\Phi_5
 \end{aligned}$$



The measurement function adds nearly all of the complexity. If I'm only interested in the total decay rate I can remove it and obtain the inclusive (partial) width.

Can obtain the inclusive partial width coefficient by using the optical theorem applied to the 4-loop two-point correlator



$$\begin{aligned}
 \Gamma_{H \rightarrow b\bar{b}}^{(3)} = & s_3 + L (2s_2\beta_0 + s_1\beta_1 + 2s_2\gamma_m^0 + 2s_1\gamma_m^1 + 2\gamma_m^2) \\
 & + L^2 (s_1\beta_0^2 + 3s_1\beta_0\gamma_m^0 + \beta_1\gamma_m^0 + 2s_1(\gamma_m^0)^2 + 2\beta_0\gamma_m^1 + 4\gamma_m^0\gamma_m^1) \\
 & + L^3 \left( \frac{2}{3}\beta_0^2\gamma_m^0 + 2\beta_0(\gamma_m^0)^2 + \frac{4}{3}(\gamma_m^0)^3 \right).
 \end{aligned}$$

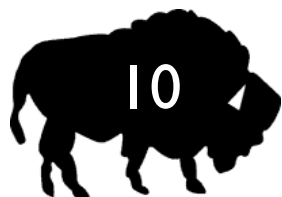
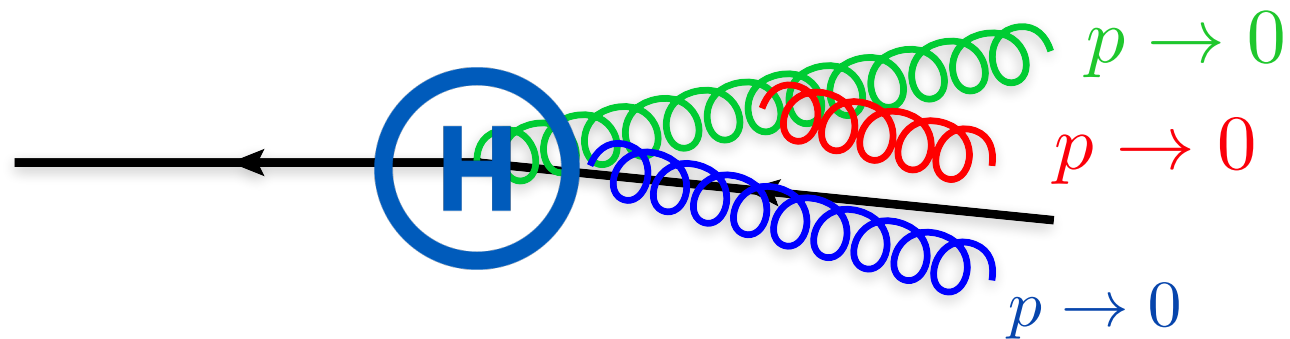
Chetyrkin [hep-ph/9608318](https://arxiv.org/abs/hep-ph/9608318)





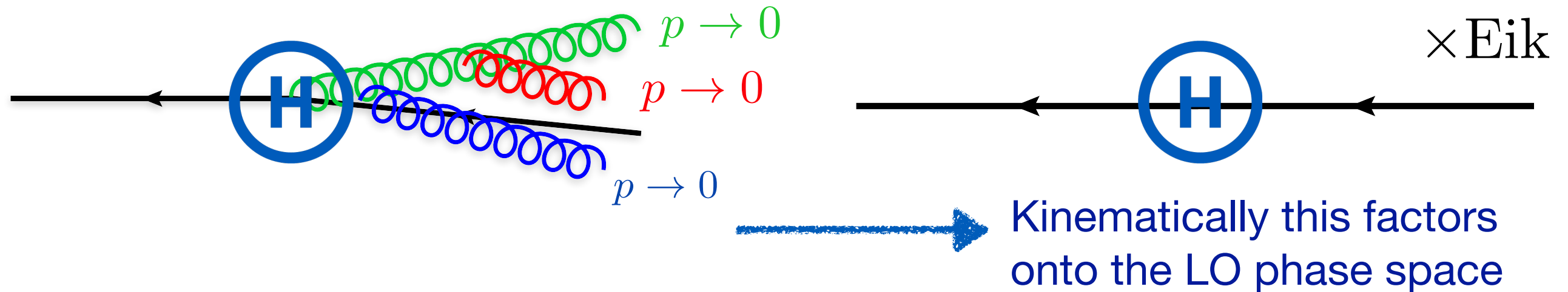
The measurement function makes life difficult since it exposes us to a multitude of Infrared singularities, which exist in the individual parton phase spaces, but cancel upon combination into a suitably inclusive observable.

For example consider the triple-soft limit in which all of the gluons in this diagram have vanishing momentum.



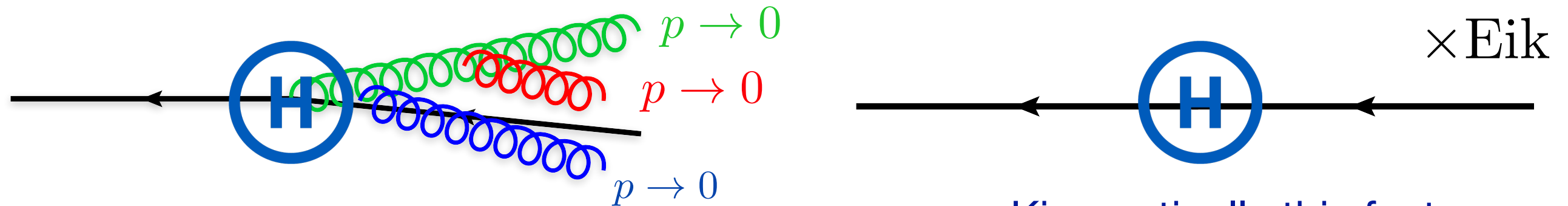
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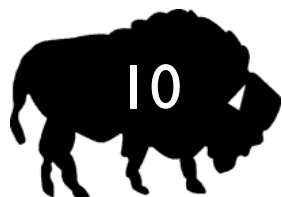
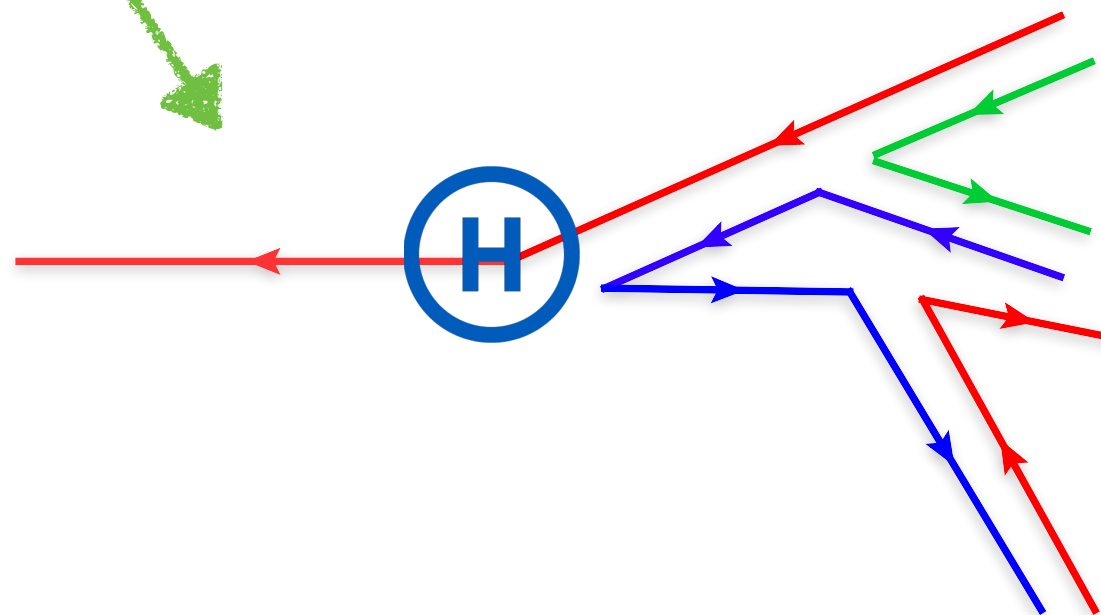
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Kinematically this factors onto the LO phase space

But the color is still present!





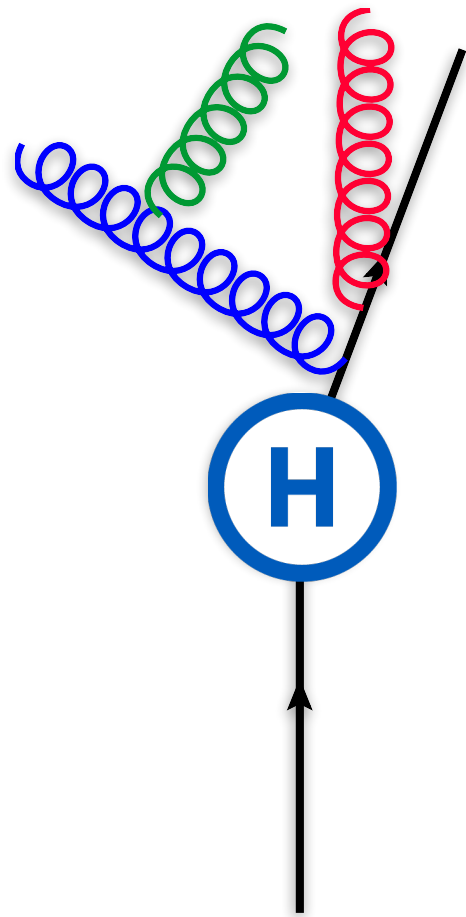
There is no one set technique for dealing with IR issues at either NLO, NNLO or N3LO.

We're going to use Projection to Born Method (Cacciari, Dreyer, Karlberg, Salam, Zanderighi 1506.02660)



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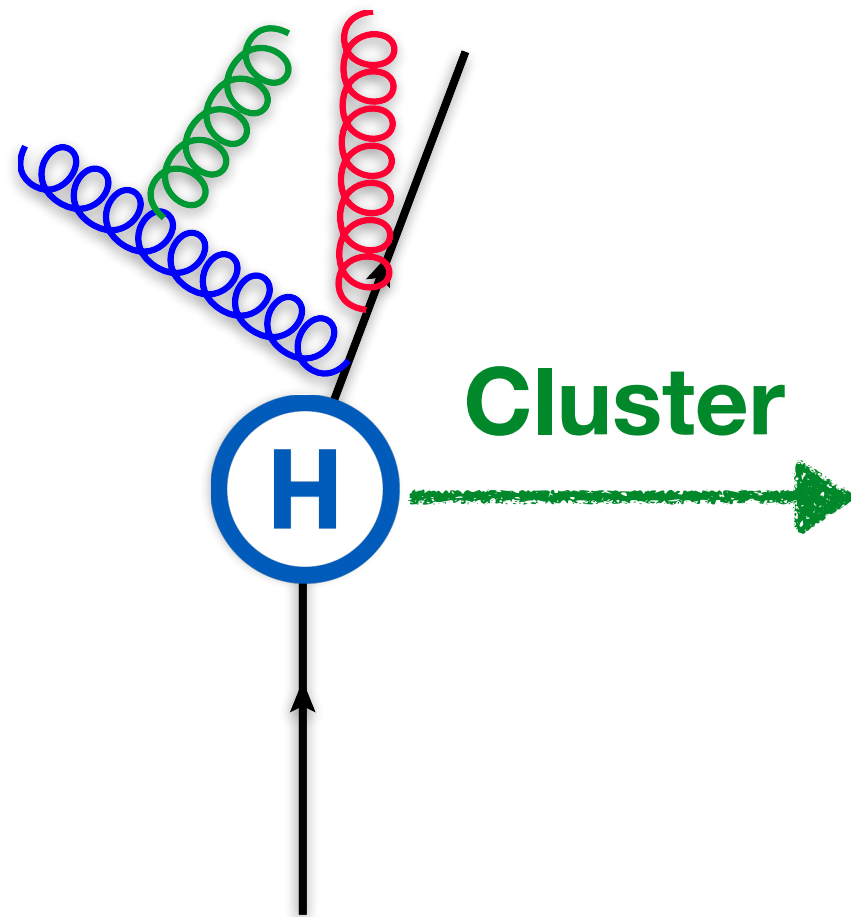


Generated event with  $|\mathcal{M}^2|$



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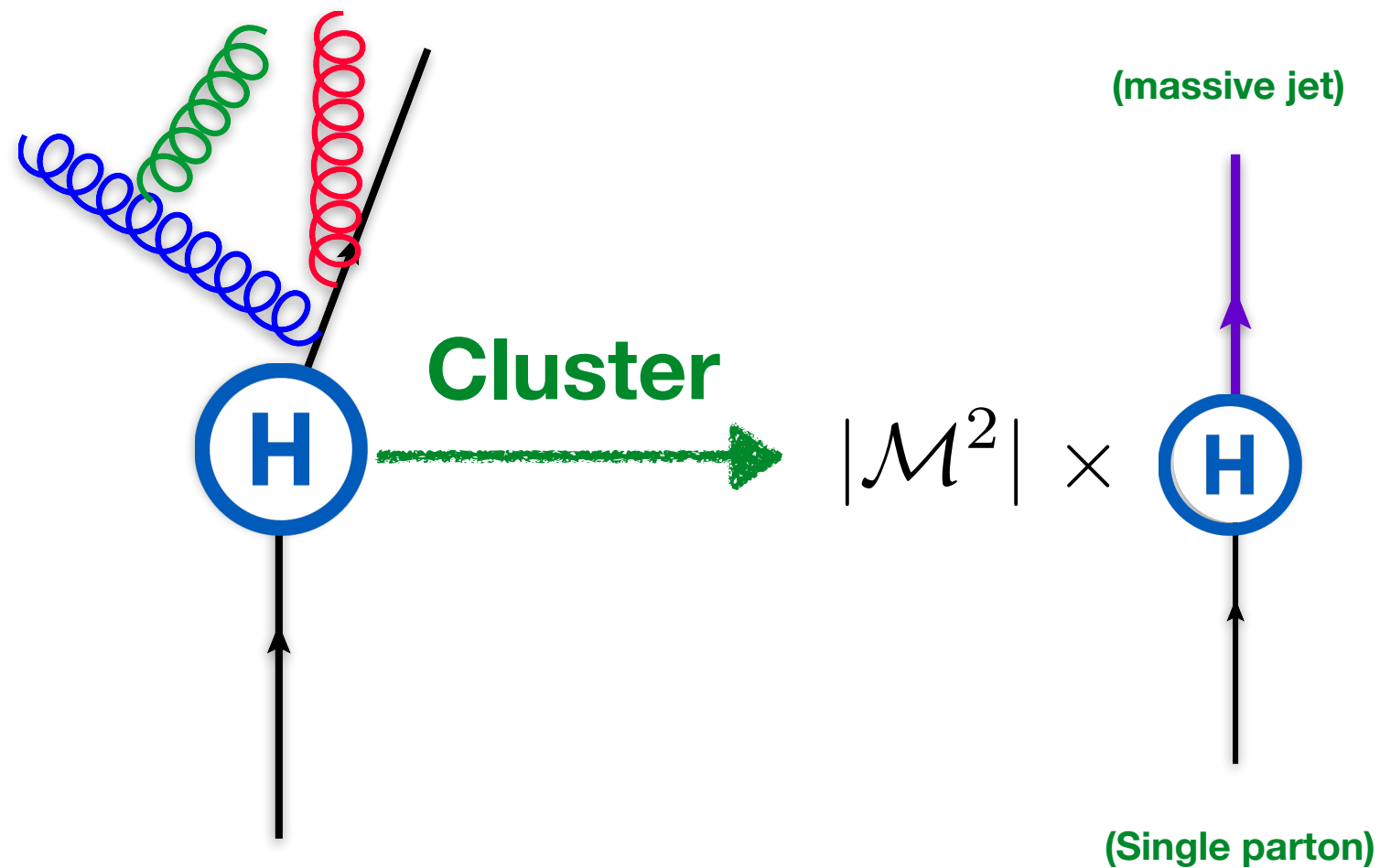


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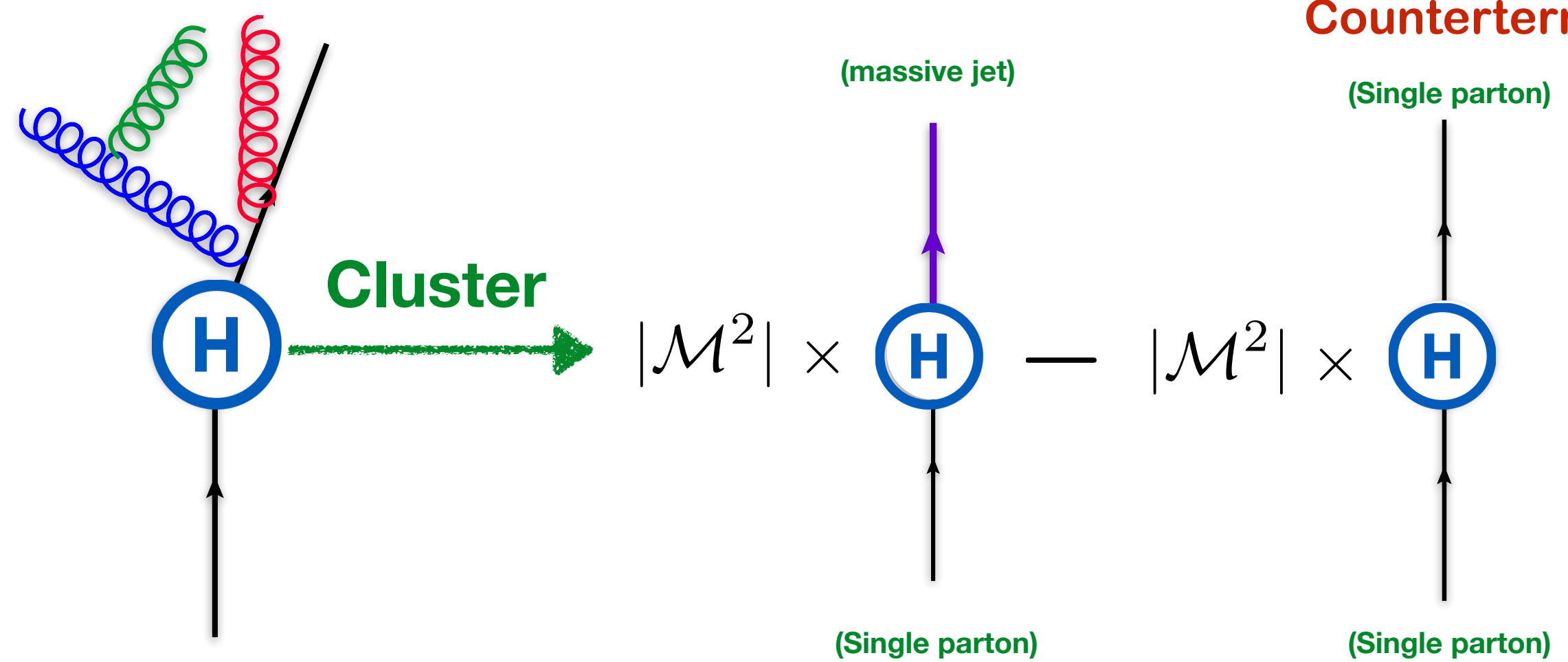
$F_2^5(\Phi_5)$



There is no one set technique for dealing with IR issues at either NLO, NNLO or N3LO.

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Subtract Counterterm



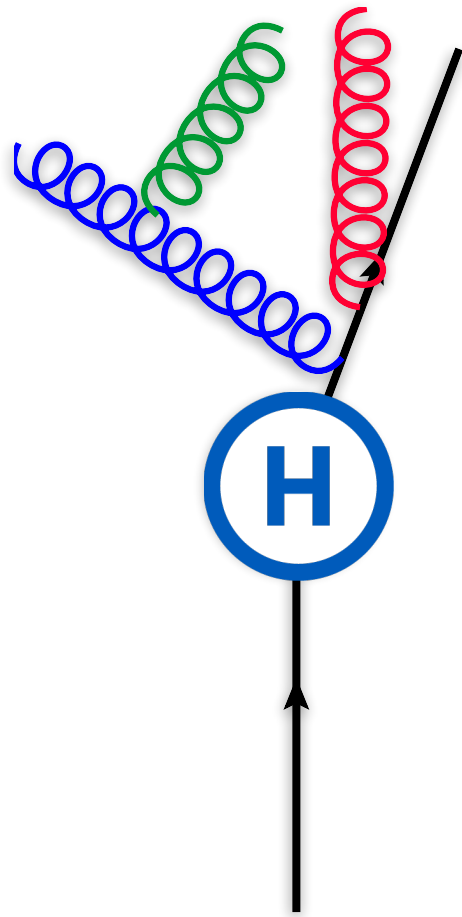
Generated event with  $|\mathcal{M}^2|$

$F_2^5(\Phi_5)$

$F_2^2(\Phi_B)$



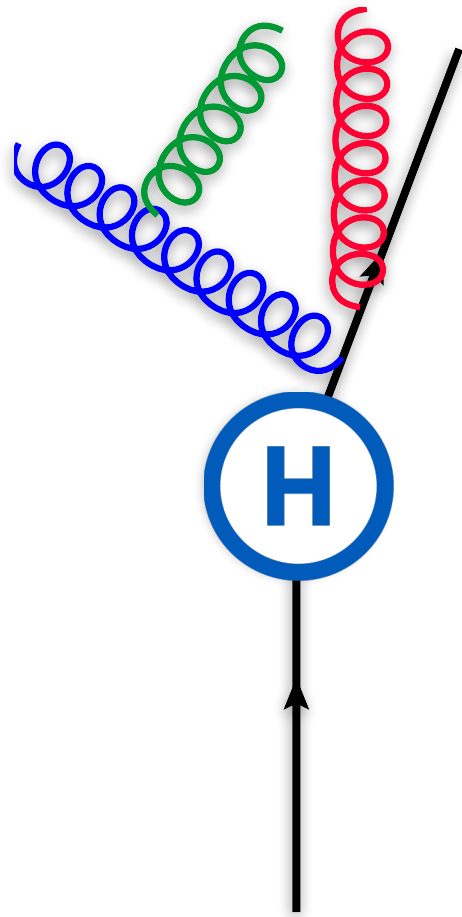
$$|\mathcal{M}^2| \times \left( F_2^5(\Phi_5) - F_2^2(\Phi_B) \right)$$



Cancels exactly when the full phase space **matches** the projected born one.



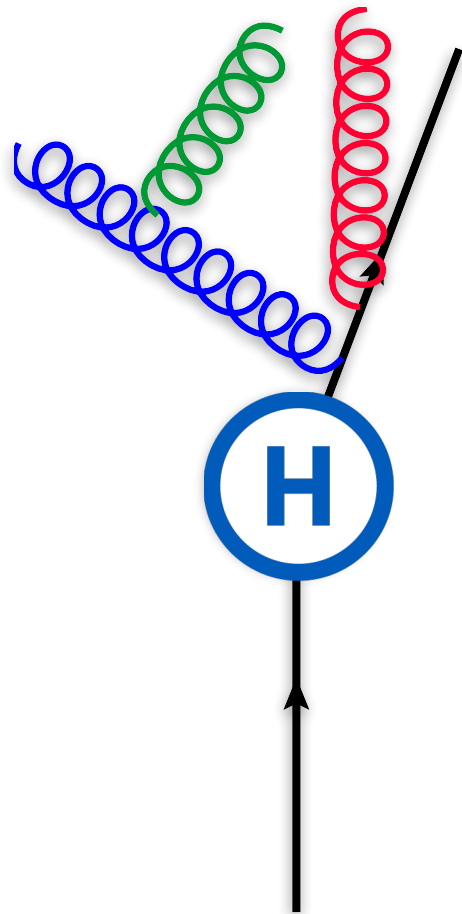
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Cancels exactly when the full phase space **matches** the projected born one.

This is exactly the triple unresolved region.

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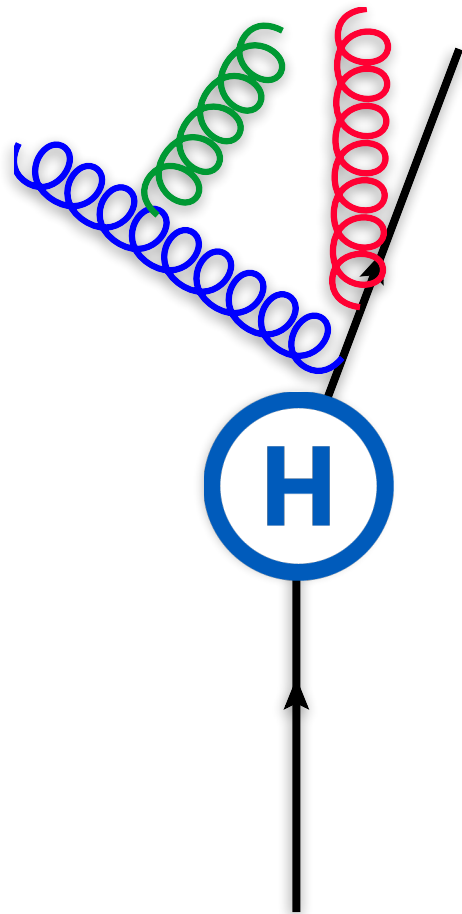


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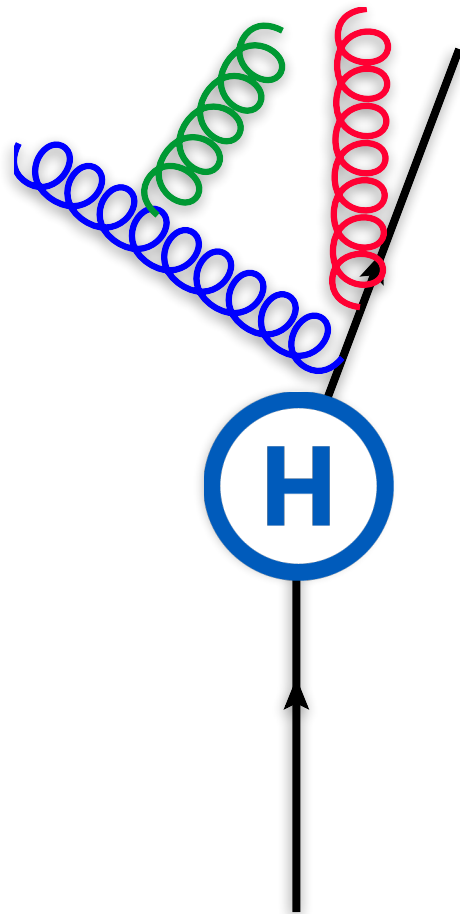
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This is an **NNLO** calculation.

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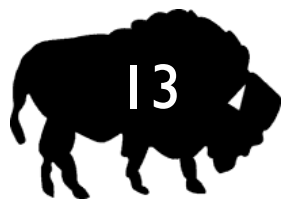
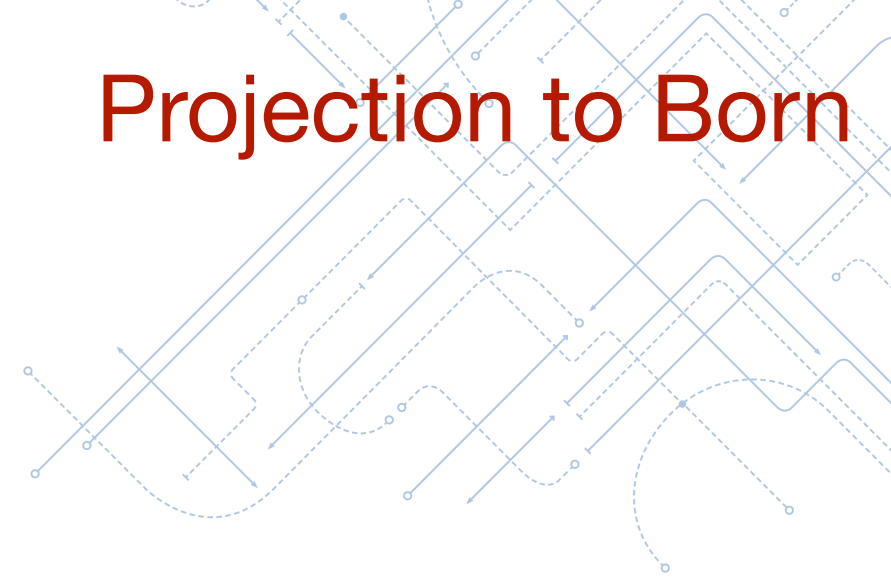
This is an **NNLO** calculation.

So ingredient 1 for a projection to Born Method is an NNLO calculation of  $H \rightarrow 3j$  (**Mondini's** talk)



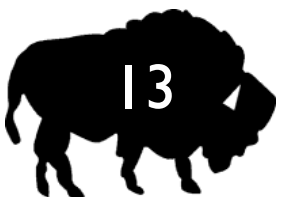
We just arbitrarily subtracted a counterterm.

Projection to Born



We just arbitrarily subtracted a counterterm.

On its own this is not cool, since it is not part of the SM, just a regulator.



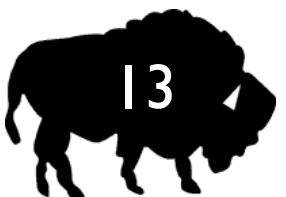


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We had better reintroduce it.

However now we're going to explicitly integrate out the phase space



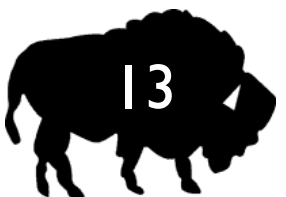
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If we integrate over all the projected Born phase space, we'll just recover the inclusive width (multiplying a LO phase space factor)



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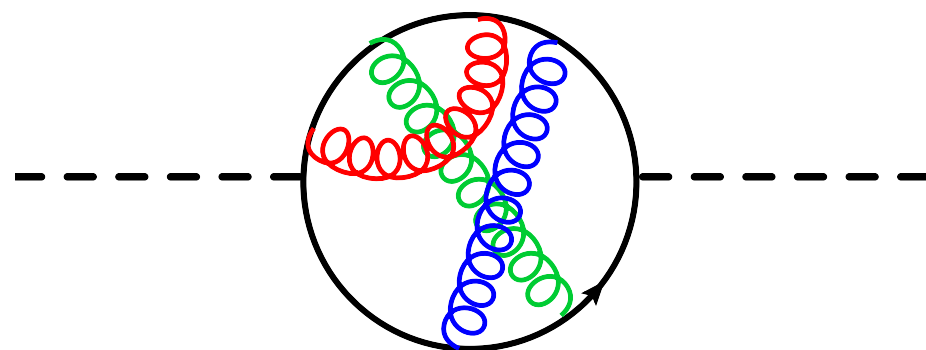
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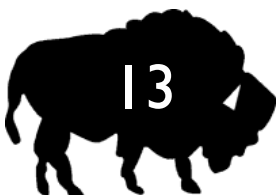
However now we're going to explicitly integrate out the phase space

If we integrate over all the projected Born phase space, we'll just recover the inclusive width (multiplying a LO phase space factor)

So ingredient 2 for a projection to Born Method is the inclusive partial width (cross section) as a function of the LO kinematics.

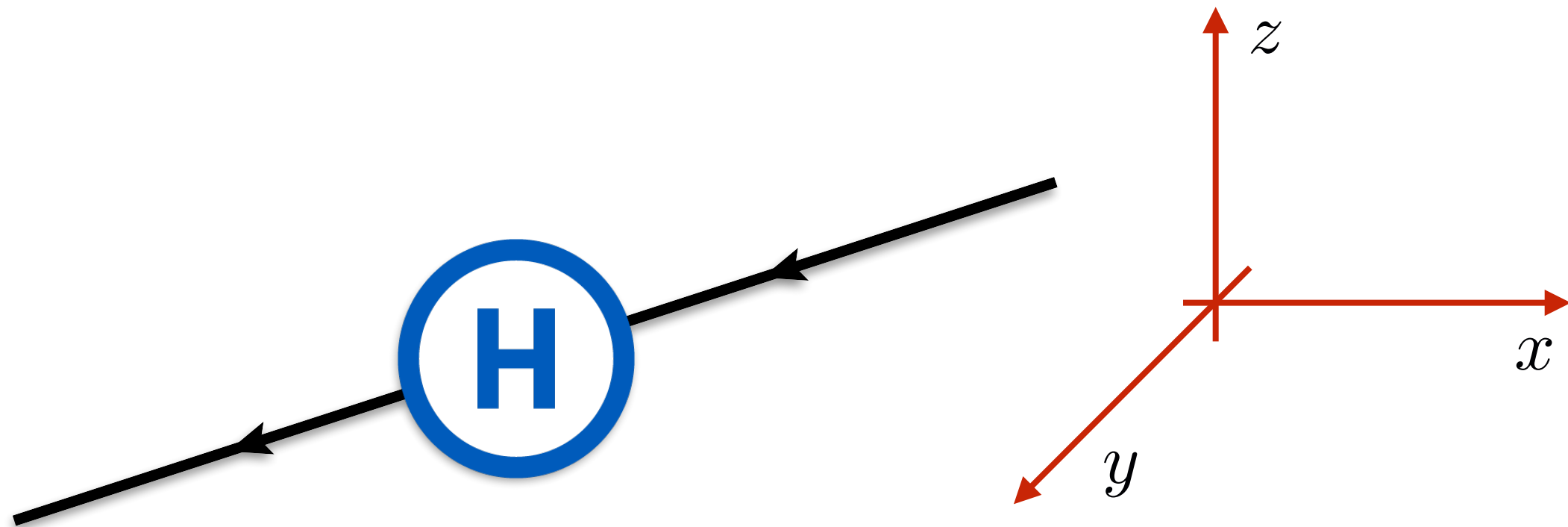


Chetyrkin [hep-ph/9608318](https://arxiv.org/abs/hep-ph/9608318)



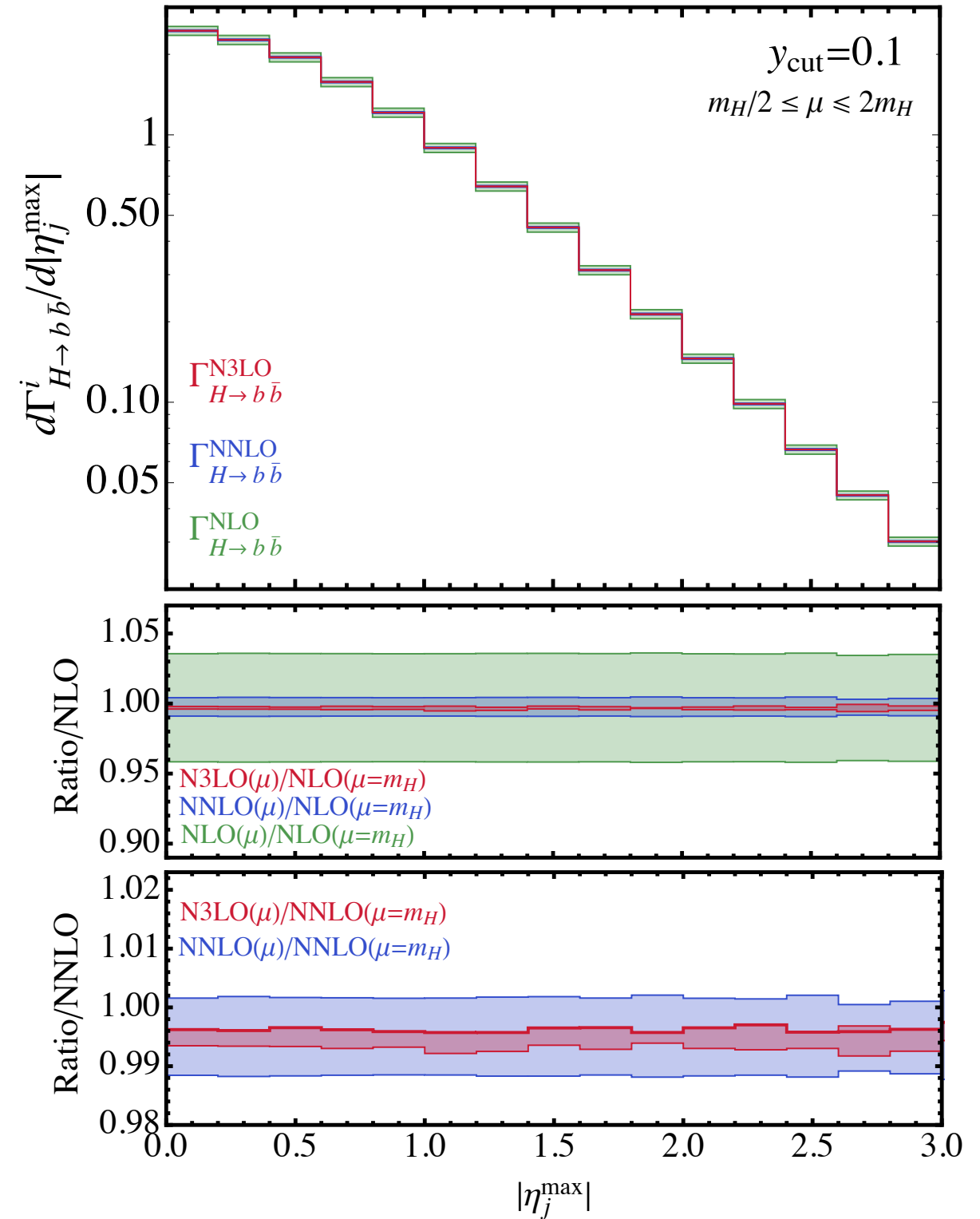
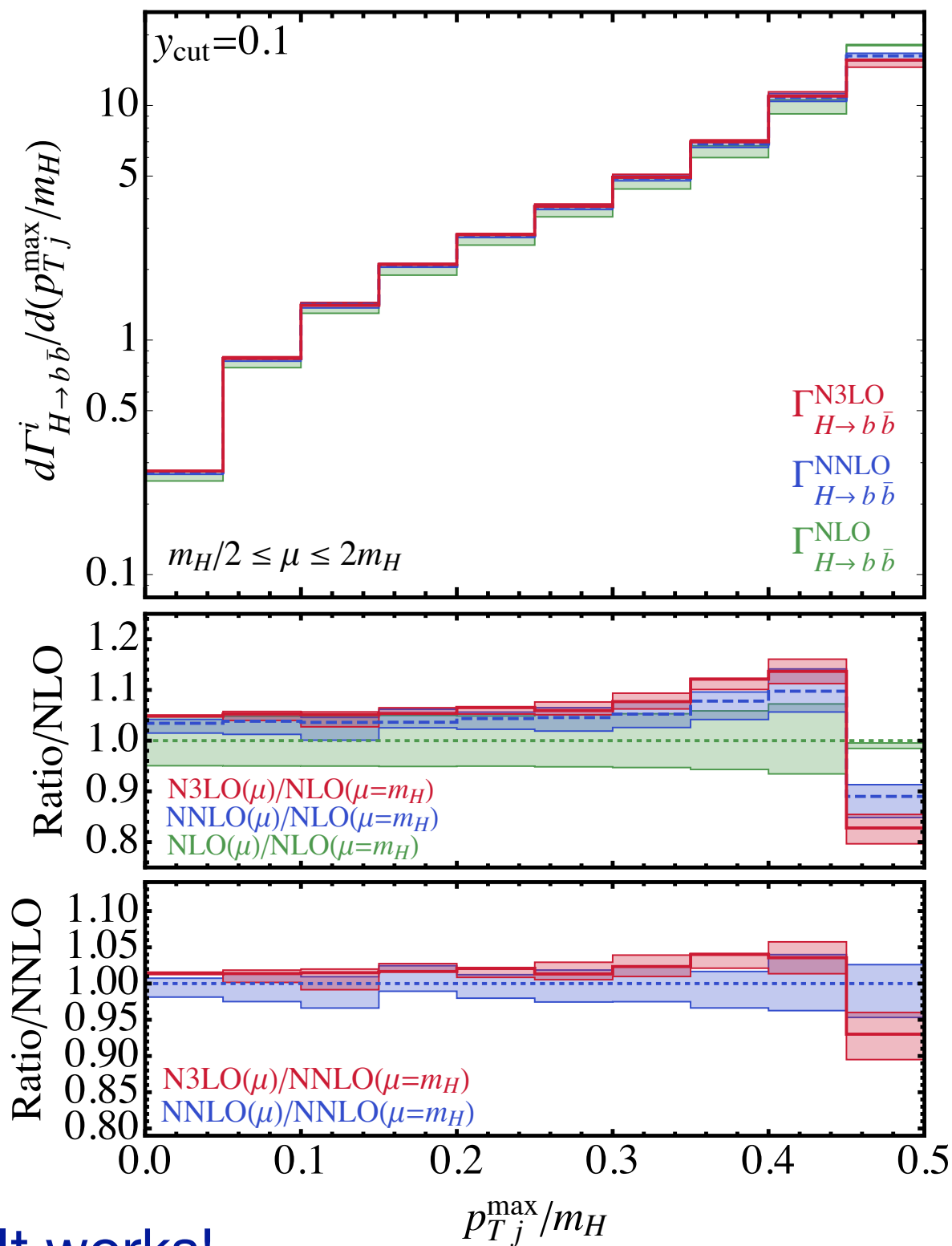
$H \Rightarrow bb$  at LO is super trivial, two massless partons back to back.

We want some interesting observables at LO to test the IR cancellations in our code, so we introduce a fictitious collision axis, and measure relative to that.



Now we can define  $p_t$  and rapidity like at the LHC.

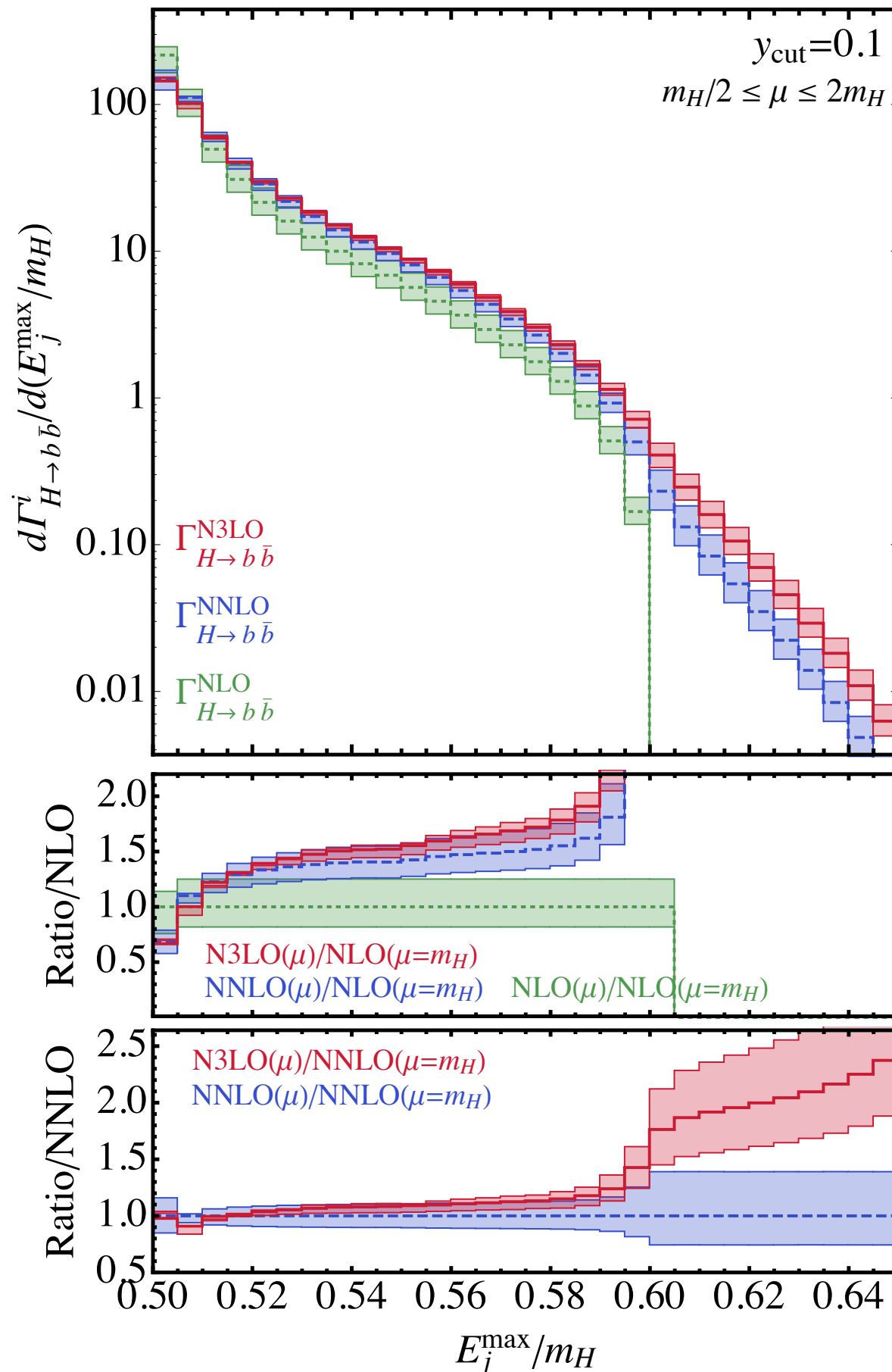
We cluster with the Durham jet algorithm  
(ask your academic grandparents)



It works!

The size of corrections are observable dependent, scale variation is tiny





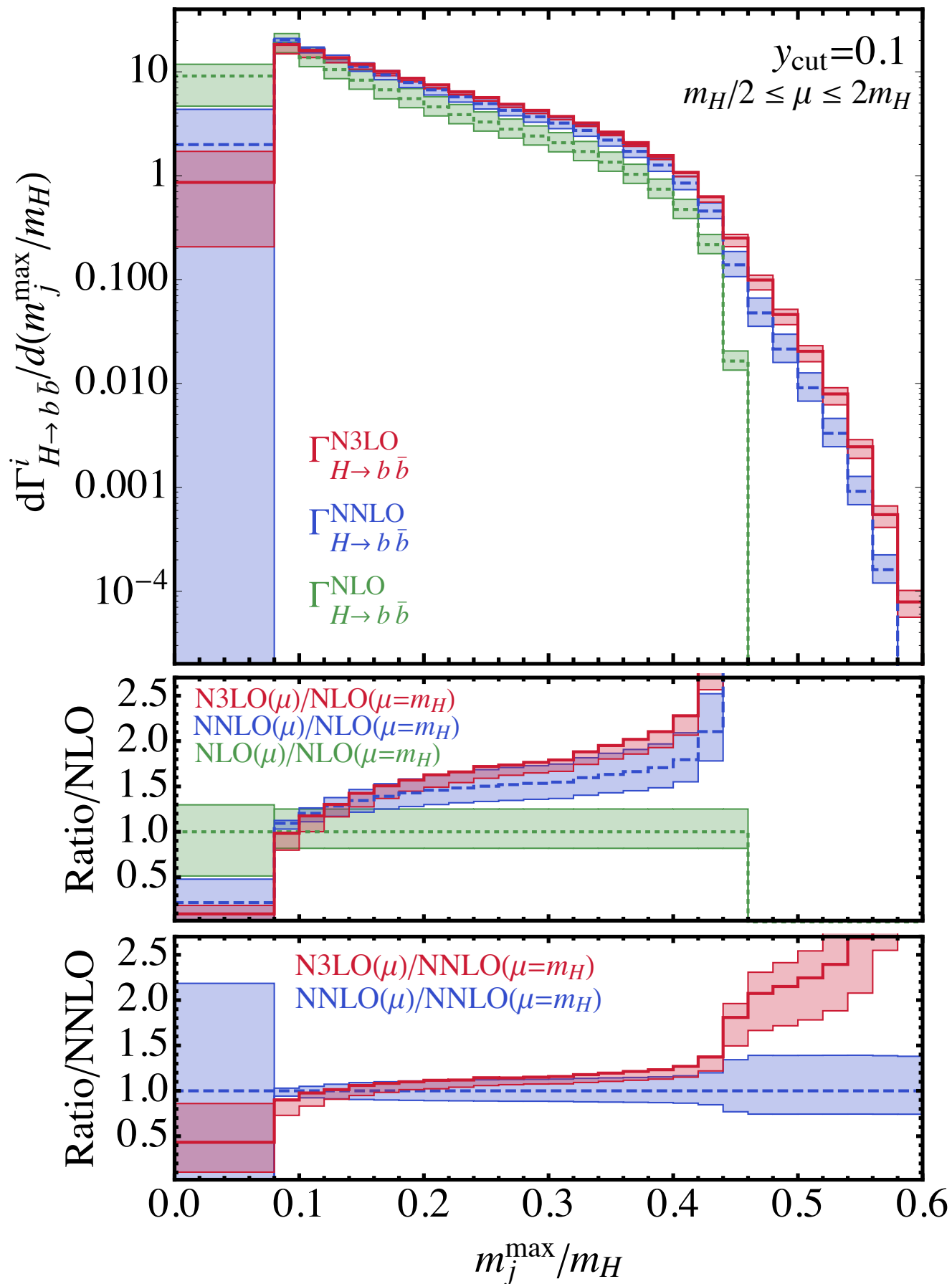
Corrections to “real” observables are a little different.

Three regions :

- 1)  $\delta(\mathcal{O})$  all phases spaces contribute, small uncertainty
- 2) Bulk, 3 parton+ phase spaces contribute, NNLO style scale variation
- 3) Tail 4 parton+ phase spaces contribute, looks like NLO scale variation







Corrections to “real” observables are a little different.

Three regions :

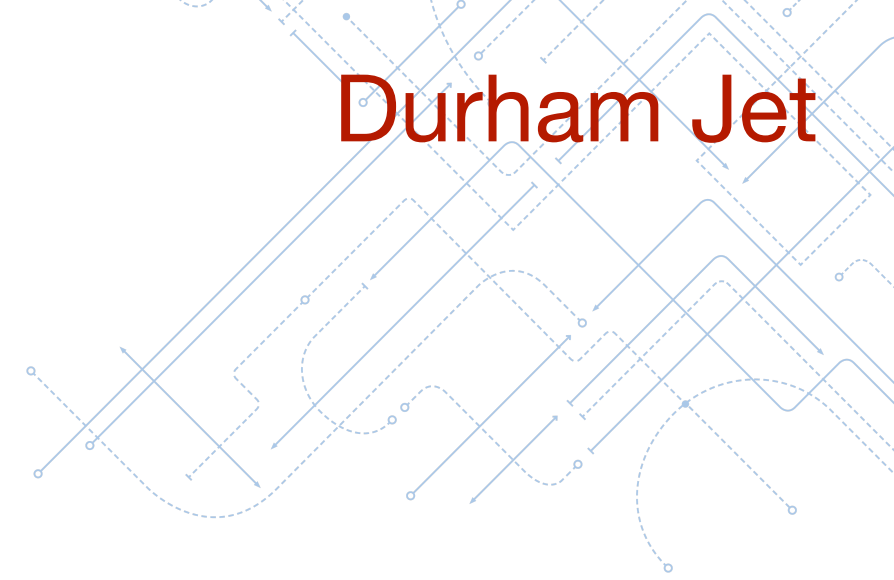
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- 3) Tail 4 parton+ phase spaces contribute, looks like NLO scale variation



- We computed N3LO corrections to  $H \rightarrow bb$
- We used the projection to Born method + N-jettiness slicing to deal with the IR singular structure
- Our calculation is fully differential and could be deployed out of the rest frame for LHC/FCC applications.
- Differential effects at N3LO are larger than small inclusive ones.



# BACKUP



$$y_{ij} = \frac{2 \min(E_i^2, E_j^2)(1 - \cos \theta_{ij})}{Q^2},$$



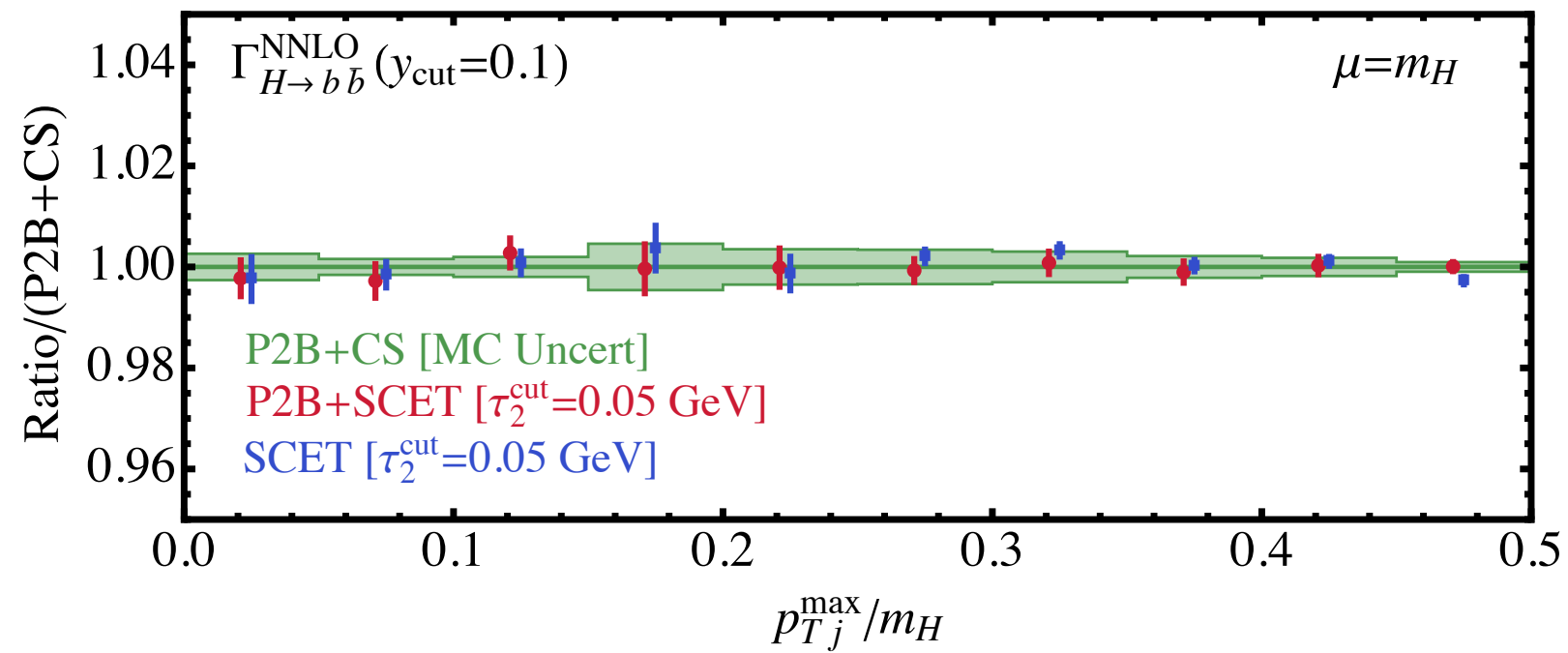
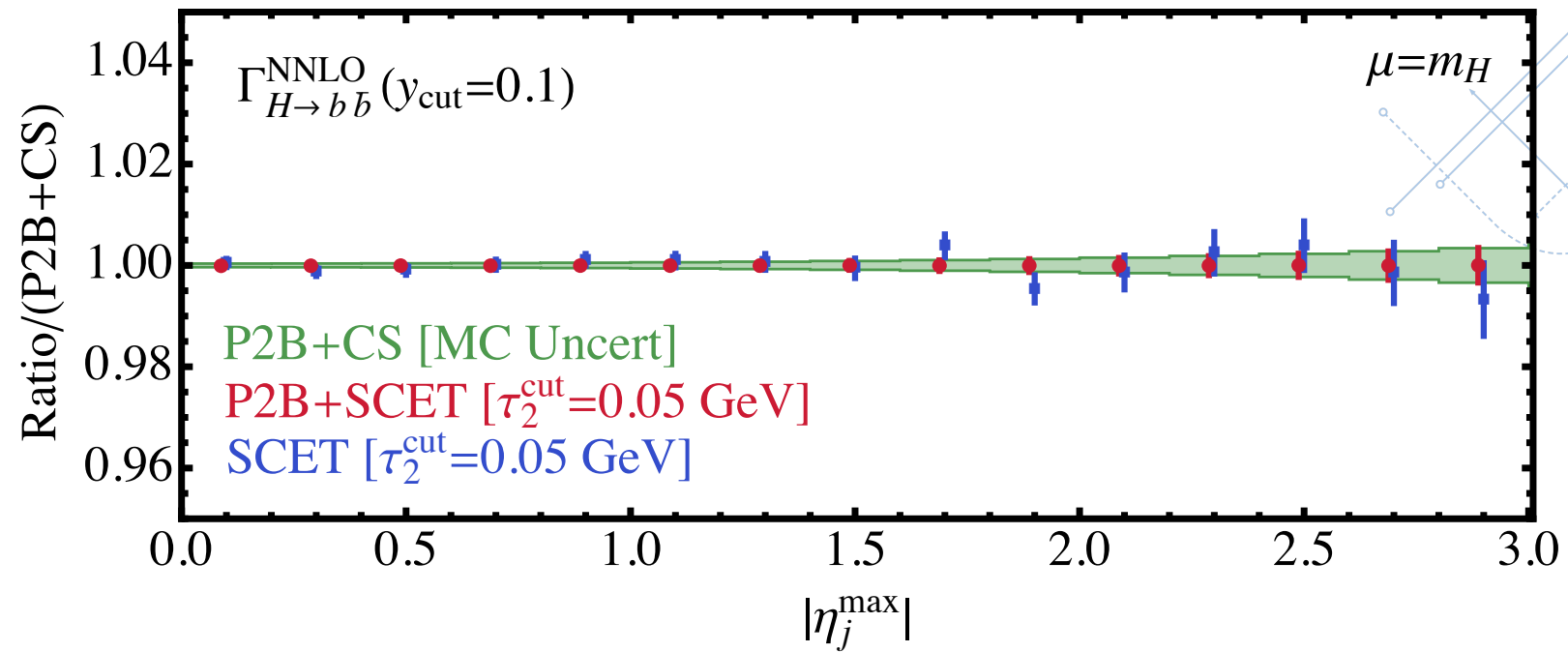
**Inclusive**

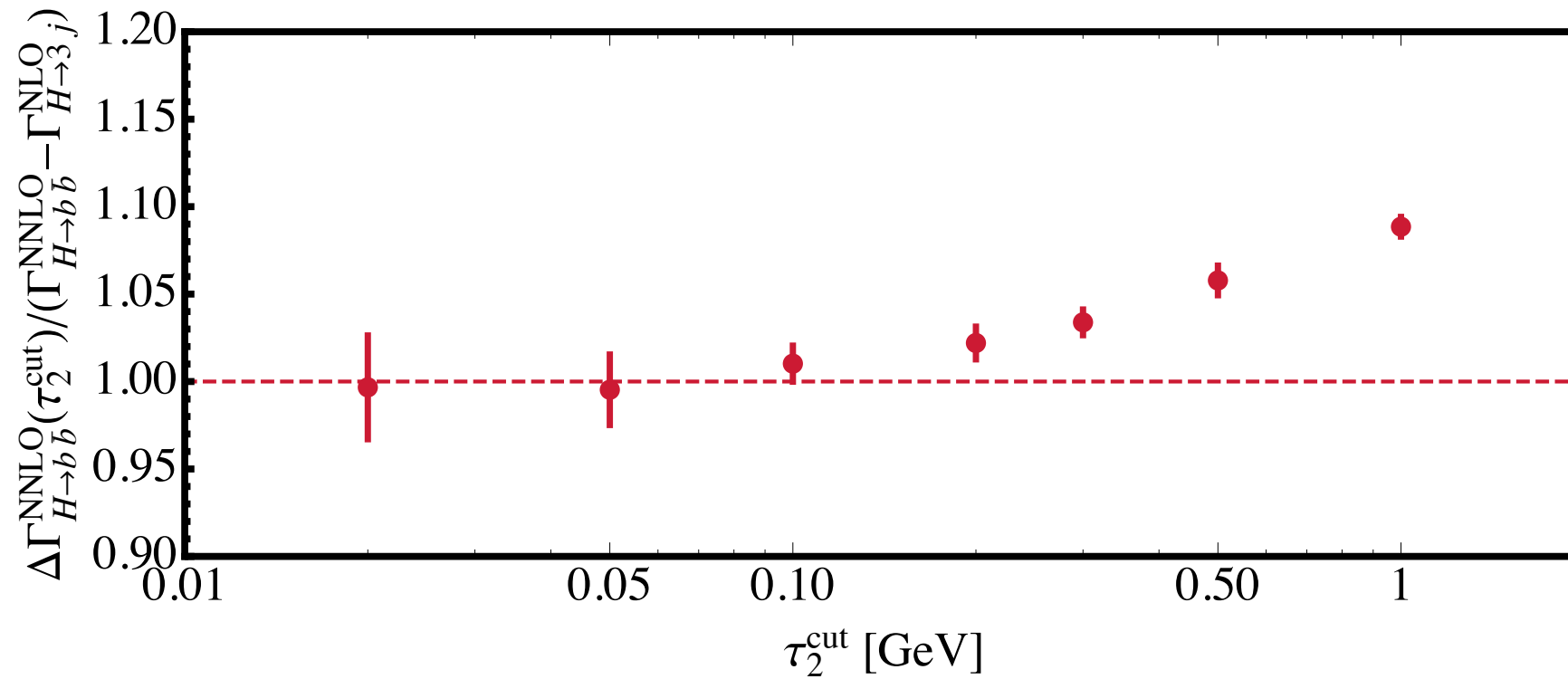
$$\begin{aligned} \Delta\Gamma_{H\rightarrow b\bar{b}}^{\text{N3LO}} &= \left(\frac{\alpha_s}{\pi}\right)^3 \int 8\pi \Gamma_{H\rightarrow b\bar{b}}^{\text{LO}} \Gamma_{H\rightarrow b\bar{b}}^{(3)} d\Phi_2 \\ &= \int \Delta\hat{\Gamma}_{H\rightarrow b\bar{b}}^{\text{N3LO}} d\Phi_2. \end{aligned}$$

**Differential**

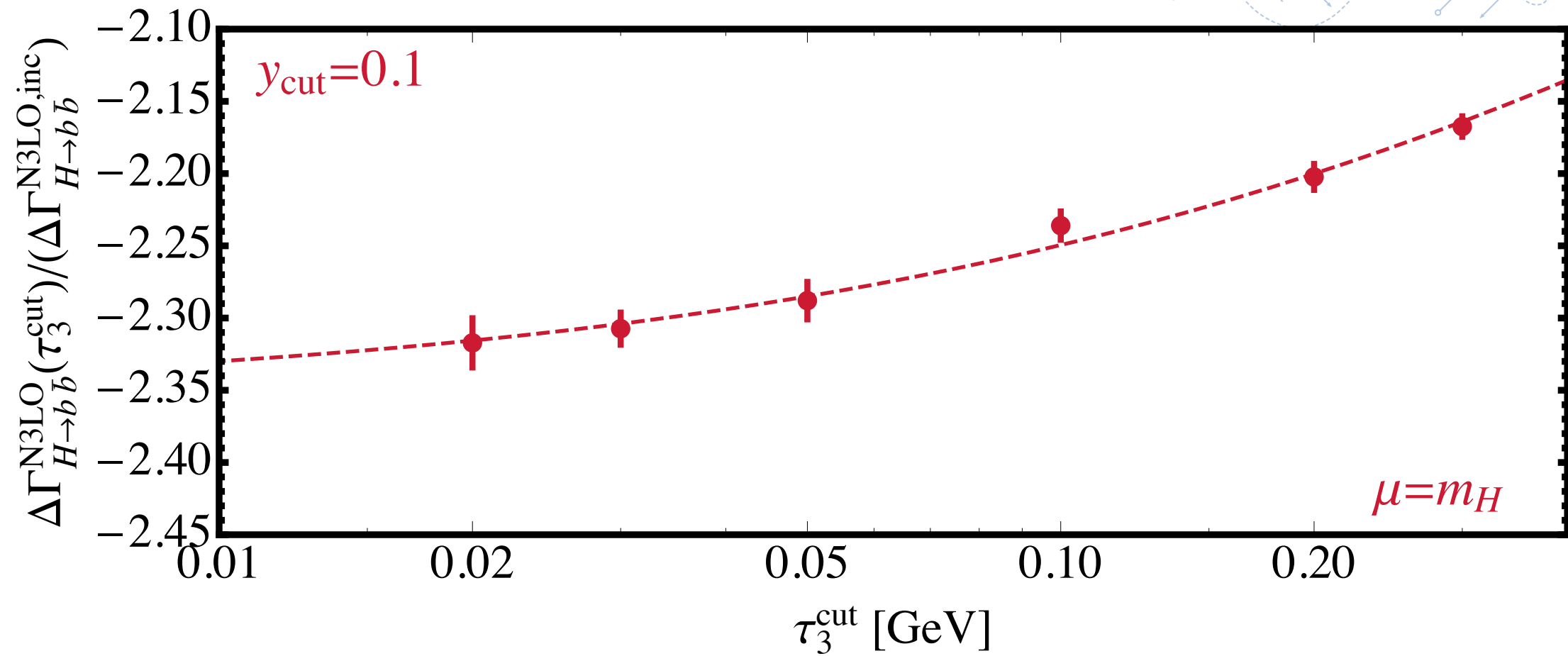
$$\begin{aligned} \frac{d\Delta\Gamma_{H\rightarrow b\bar{b}}^{\text{N3LO}}}{d\mathcal{O}_m} &= \int \Delta\hat{\Gamma}_{H\rightarrow b\bar{b}}^{\text{N3LO}} F_2^m(\Phi_B) d\Phi_B \\ &+ \int d\Gamma_{H\rightarrow b\bar{b}}^{\text{RVV}} [F_3^m(\Phi_3) - F_2^m(\Phi_B)] d\Phi_3 \\ &+ \int d\Gamma_{H\rightarrow b\bar{b}}^{\text{RRV}} [F_4^m(\Phi_4) - F_2^m(\Phi_B)] d\Phi_4 \\ &+ \int d\Gamma_{H\rightarrow b\bar{b}}^{\text{RRR}} [F_5^m(\Phi_5) - F_2^m(\Phi_B)] d\Phi_5. \end{aligned}$$







$$\tau_N = \sum_{j=1,n} \min_{i=1,2,N} \left\{ \frac{2q_i \cdot p_j}{Q_i} \right\}$$



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