

Confronting Neutrino Mass Generation Mechanism with MiniBooone Anomaly

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Phenomenology Symposium 2019

University of Pittsburgh

Based on

1. *arXiv: 1807.09877*, *Phys.Rev.Lett.* 121 (2018) no.24, 241801
2. *arXiv: 1808.02500*, *Phys.Lett.B*791 (2019) 210-214

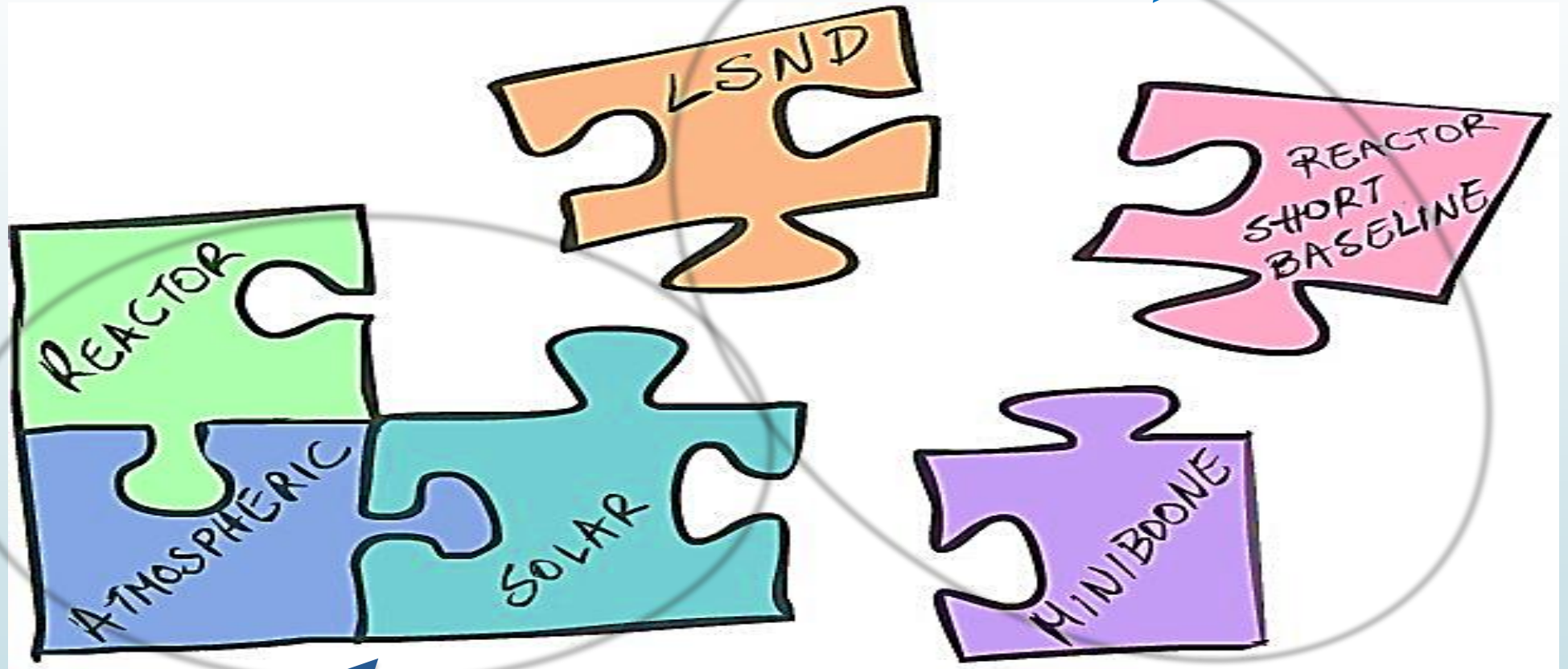
in collaboration with

E. Bertuzzo, Pedro A. N. Machado and R. Zukanovich-Funchal



Three-neutrino oscillation: Not the full picture?

Short Baseline
Anomalies



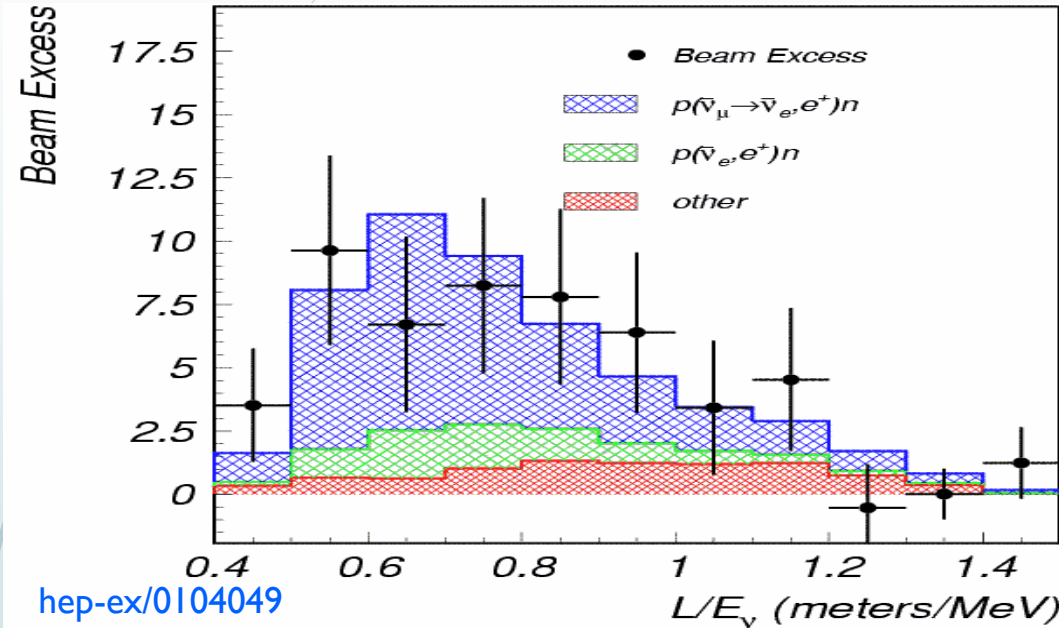
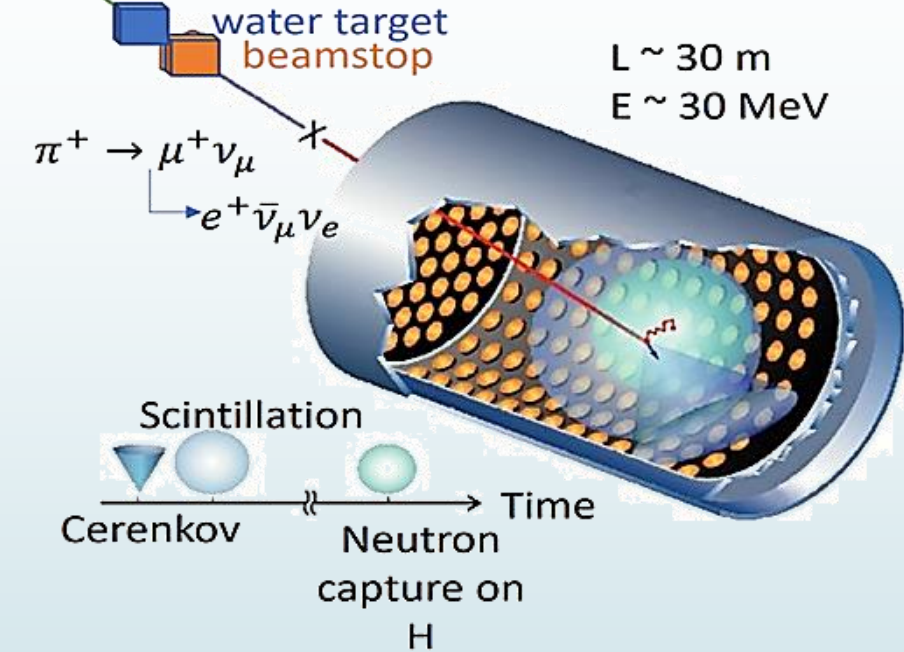
Long and Medium
Baseline



Liquid Scintillator Neutrino Detector (LSND)

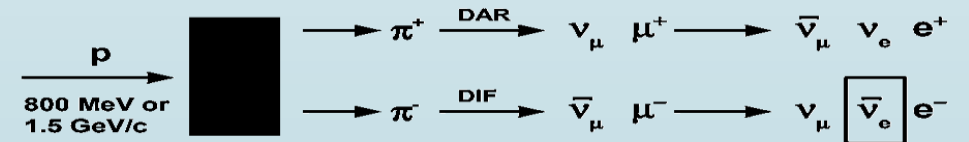
Cherenkov : Scintillation = 1:5

0.8 GeV proton beam

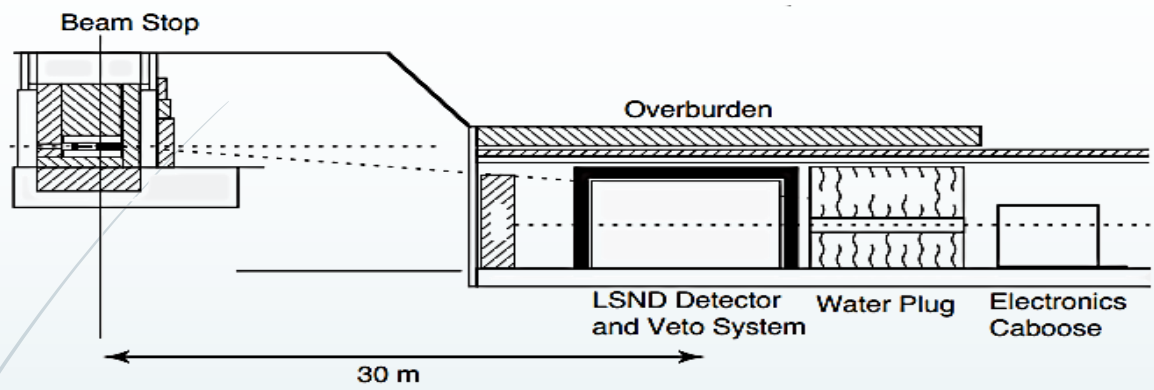


LSND detected more $\bar{\nu}_e$ than expected :
 $87.9 \pm 22.4 \pm 6.0$ events
3.8 σ excess

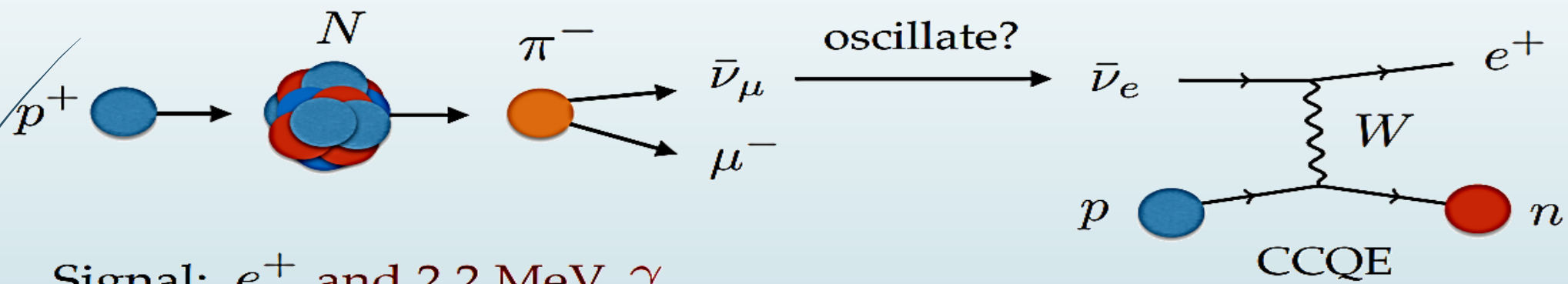
LSND neutrino source



DAR of μ^- competes with $\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z-1)$



10^{22} POT, $E_p = 800$ MeV



Signal: e^+ and 2.2 MeV γ
 Scatter + neutron absorption

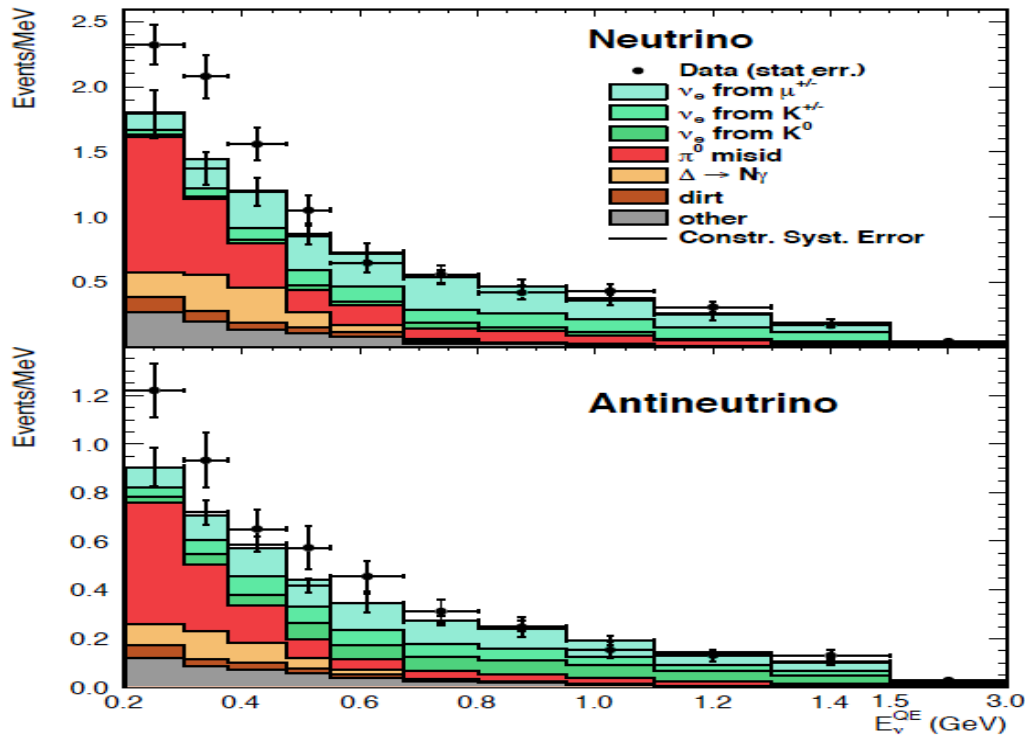
Observed 90 events
 Expectation of 0 events
 3.8σ significance

$$P_{\text{osc}} = \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E_\nu} \right)$$

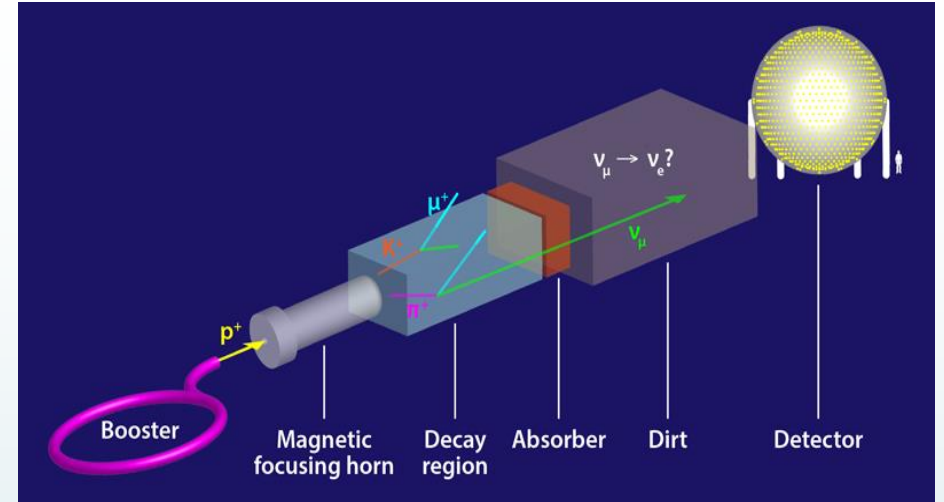
LSND Collaboration hep-ex/0104049



MiniBooNE 1207.4809

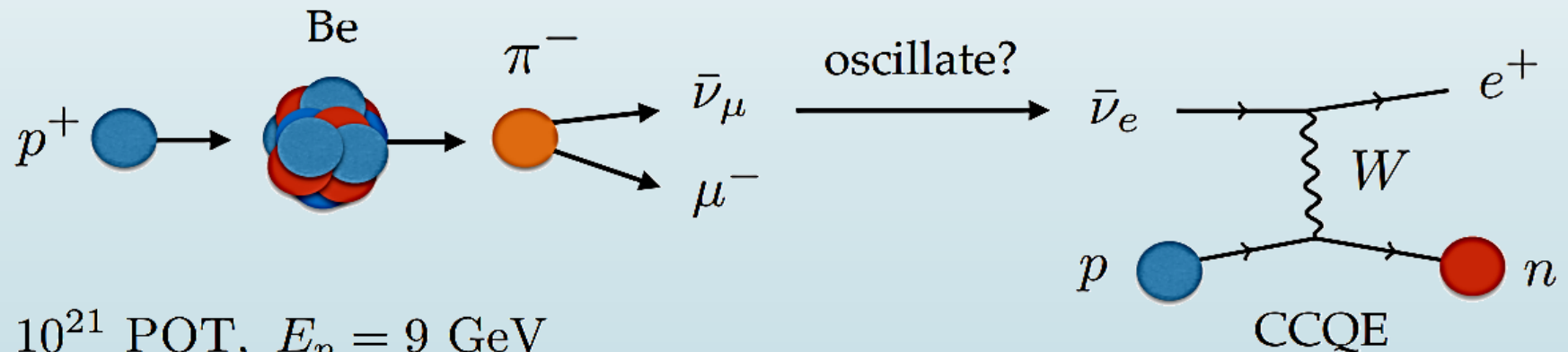
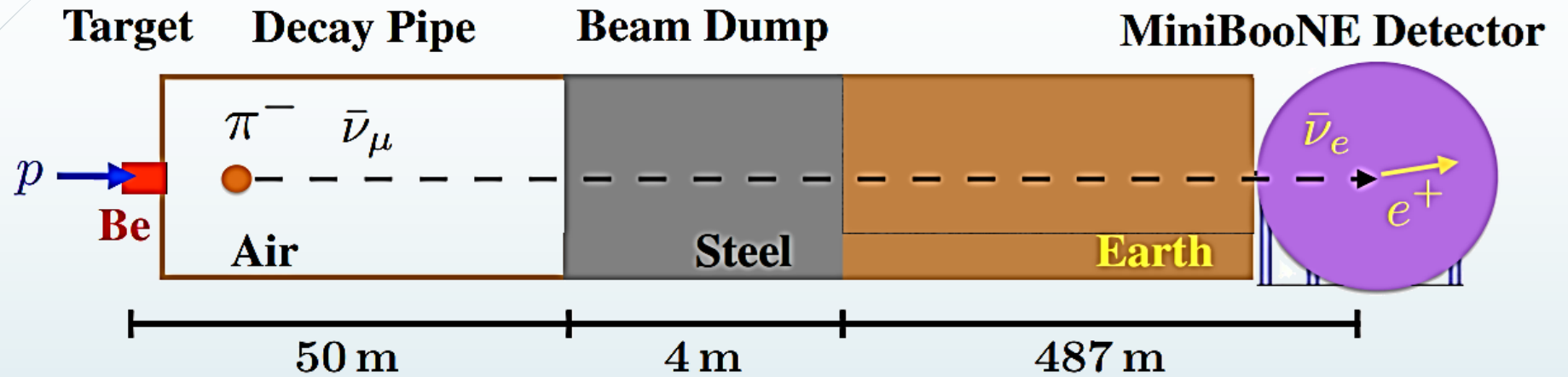
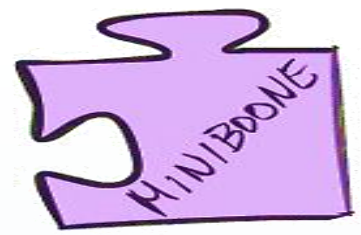


• Neutrino and anti neutrino modes see excesses of ν_e and $\bar{\nu}_e$ (Combined is also 3.8σ excess)



- ❑ To test the LSND indication of anti-electron neutrino oscillations
- ❑ Keep L/E same, change beam, energy, and systematic errors
- ❑ Baseline: $L = 540$ meters, $\sim x 15$ LSND
- ❑ Neutrino Beam Energy: $E \sim x (10-20)$ LSND
- ❑ Different systematics: event signatures and backgrounds different from LSND High statistics: $\sim x 6$ LSND
- ❑ Perform experiment in both neutrino and anti-neutrino modes.

MINIBOONE'S LOW ENERGY EXCESS



10^{21} POT, $E_p = 9$ GeV

Energy and baseline chosen to test LSND

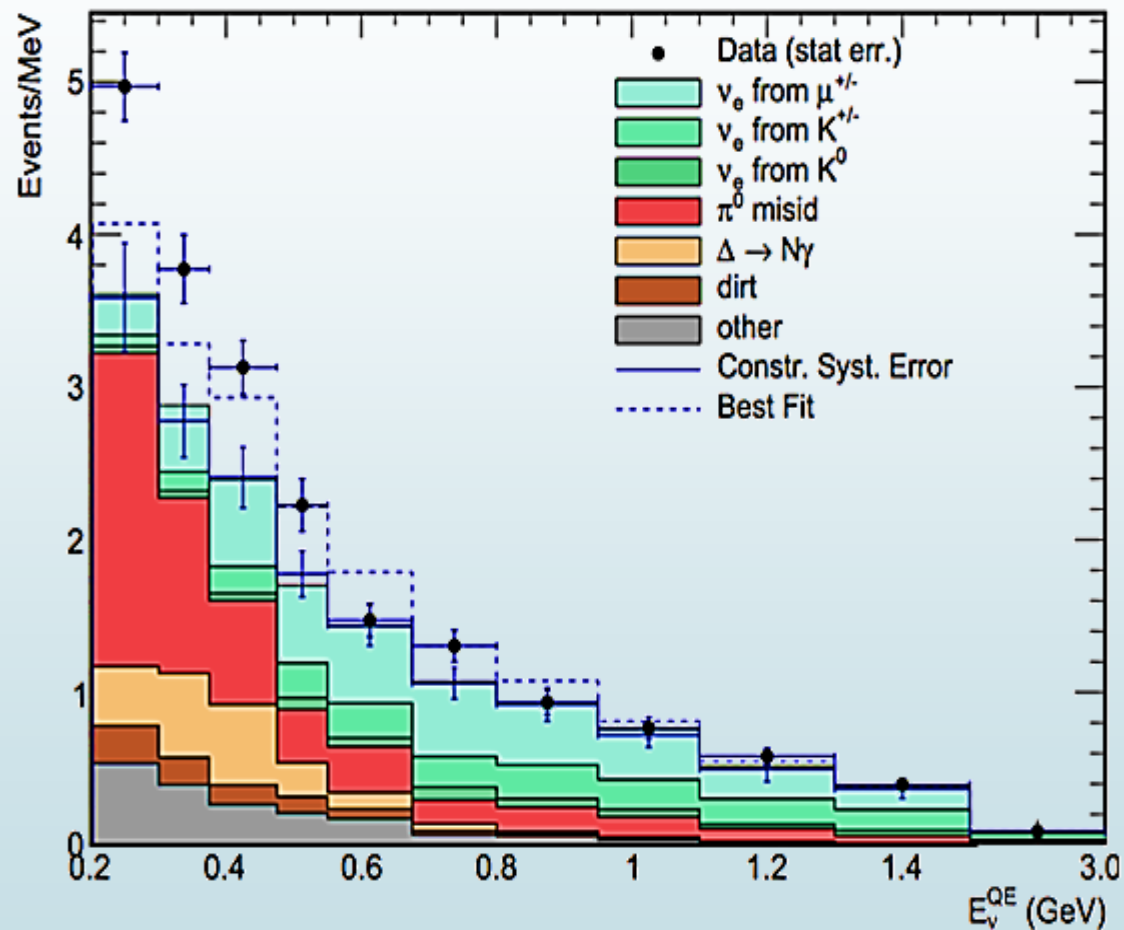
Comparable oscillation probabilities



MiniBooNE's Low Energy Excess

- ❖ Observation of a Significant Excess of Electron-Like Events in the MiniBooNE Short Baseline Neutrino Experiment
- ❖ Double neutrino-mode data in 2016-2017 ($6.46 \times 10^{20} + 6.38 \times 10^{20}$ POT)
- ❖ Event excess: 381.2 ± 85.2 (4.5σ)

MiniBooNE Collaboration [hep-ex/1805.12028](https://arxiv.org/abs/hep-ex/1805.12028)



What is going on???

- What is the nature of the excess?
- Possible detector anomalies or reconstruction problems?
- Incorrect estimation of the background?
- New sources of background?
- **New physics** including/excluding exotic oscillation scenarios?

The origin of such excess is unclear - it could be the presence of new physics, or a large background mismodeling. However, the MiniBooNE result, if due to new physics, would revolutionize the field of particle physics.

What sort of new physics can explain these anomalies?

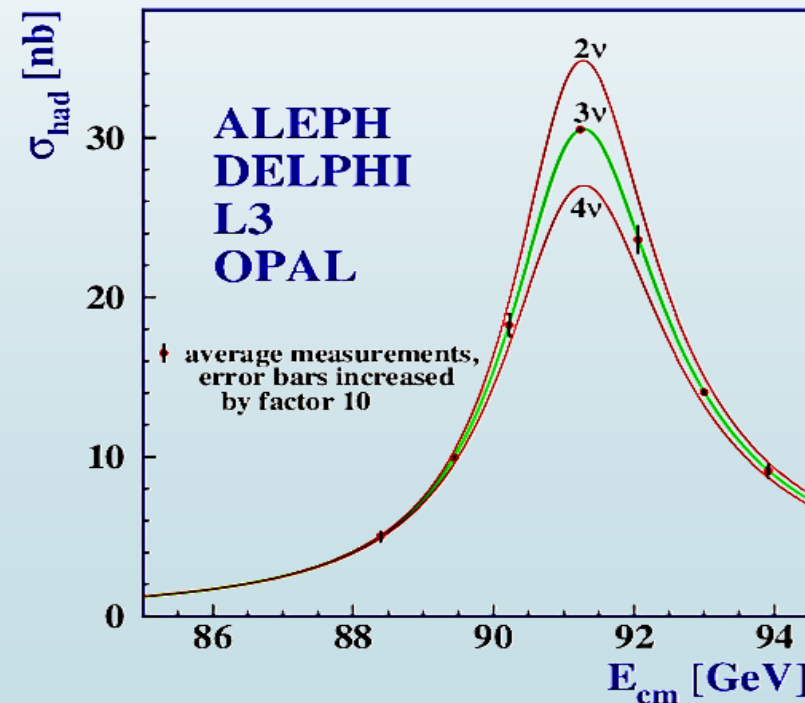


❖ *What about $e\nu$ Sterile Neutrino Interpretation ???*

Beyond three-neutrino oscillations

- We can add a fourth neutrino
- This neutrino must be sterile, which means it is a singlet under all standard model gauge groups
- A fourth active neutrino is excluded by observations of invisible Z-decays

$$e^+e^- \rightarrow Z \rightarrow \sum_{j=e,\mu,\tau} \nu_j$$



❖ What about $e\nu$ Sterile Neutrino Interpretation ???

Effective 3+1 oscillations

We extend the mixing matrix

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \Rightarrow \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

APPEARANCE

$$P_{\alpha\beta}^{\text{SBL}} \approx \sin^2(2\theta_{\alpha\beta}) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

$$\sin^2(2\theta_{\alpha\beta}) = 4|U_{\alpha 4}|^2 |U_{\beta 4}|^2$$

$$\nu_{\mu} \rightarrow \nu_e : \sin^2(2\theta_{\mu e}) = 4|U_{e 4}|^2 |U_{\mu 4}|^2$$

@LSND, Karmen, MiniBoone,
Opera

DISAPPEARANCE

$$P_{\alpha\alpha}^{\text{SBL}} \approx 1 - \sin^2(2\theta_{\alpha\alpha}) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

$$\sin^2(2\theta_{\alpha\alpha}) = 4|U_{\alpha 4}|^2 (1 - |U_{\alpha 4}|^2)$$

$$\nu_e \rightarrow \nu_e : |U_{e 4}|^2 = \sin^2 \theta_{14}$$

@Reactors and Gallium

$$\nu_{\mu} \rightarrow \nu_{\mu} : |U_{\mu 4}|^2 = \sin^2 \theta_{24} \cos^2 \theta_{14}$$

@atmospherics and accelerators

❖ *What about $e\nu$ Sterile Neutrino Interpretation ???*

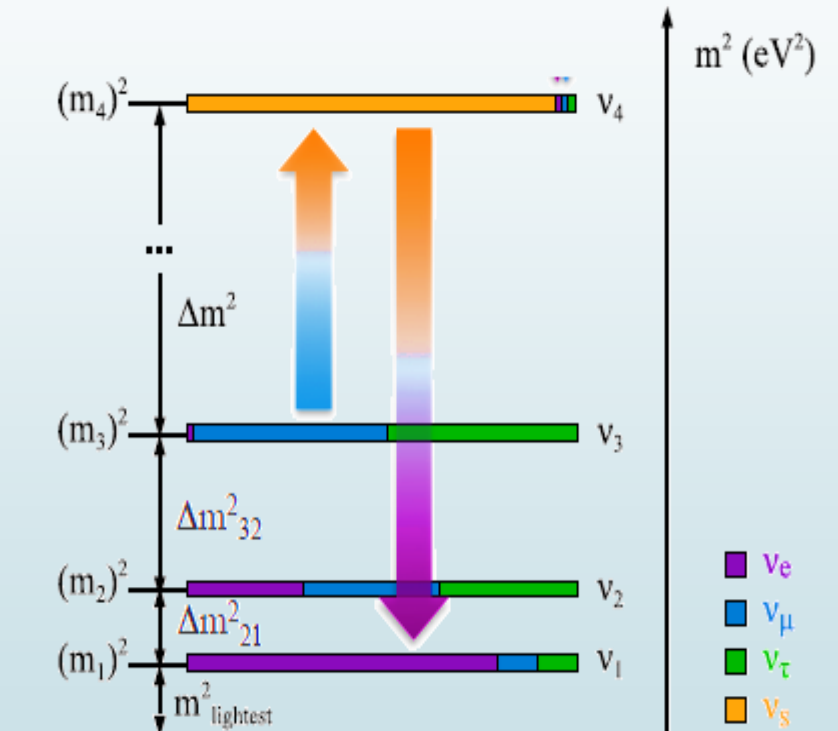
$$P_{\alpha\alpha}^{\text{SBL}} = 1 - 4|U_{\alpha 4}|^2(1 - |U_{\alpha 4}|^2) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

$$P_{\alpha\beta}^{\text{SBL}} = 4|U_{\alpha 4}|^2|U_{\beta 4}|^2 \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right).$$

$$\sin^2 2\theta_{\mu e} = 4 |U_{e4}|^2 |U_{\mu 4}|^2$$

Leads to ν_e
disappearance

Leads to ν_μ
disappearance

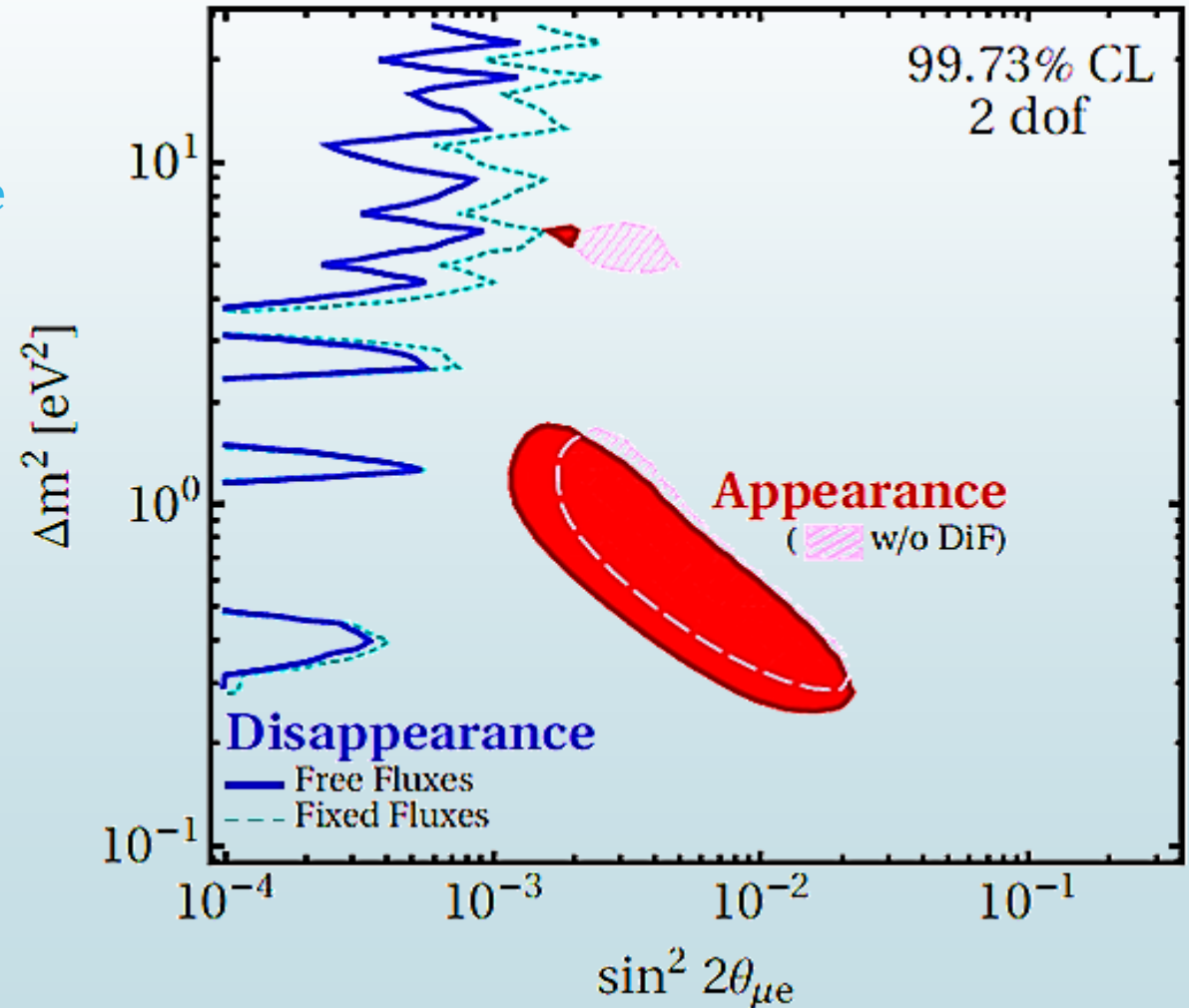


❖ *What about $e\nu$ Sterile Neutrino Interpretation ???*

$$\sin^2 2\theta_{\mu e} = 4 \left| \underbrace{U_{e4} U_{\mu 4}} \right|^2$$

Leads to ν_μ to ν_e disappearance

- 2 variables: $U_{e4}, U_{\mu 4}$
- 3 data sets: ν_e - Disappearance
 ν_μ - Disappearance
 ν_e - Appearance

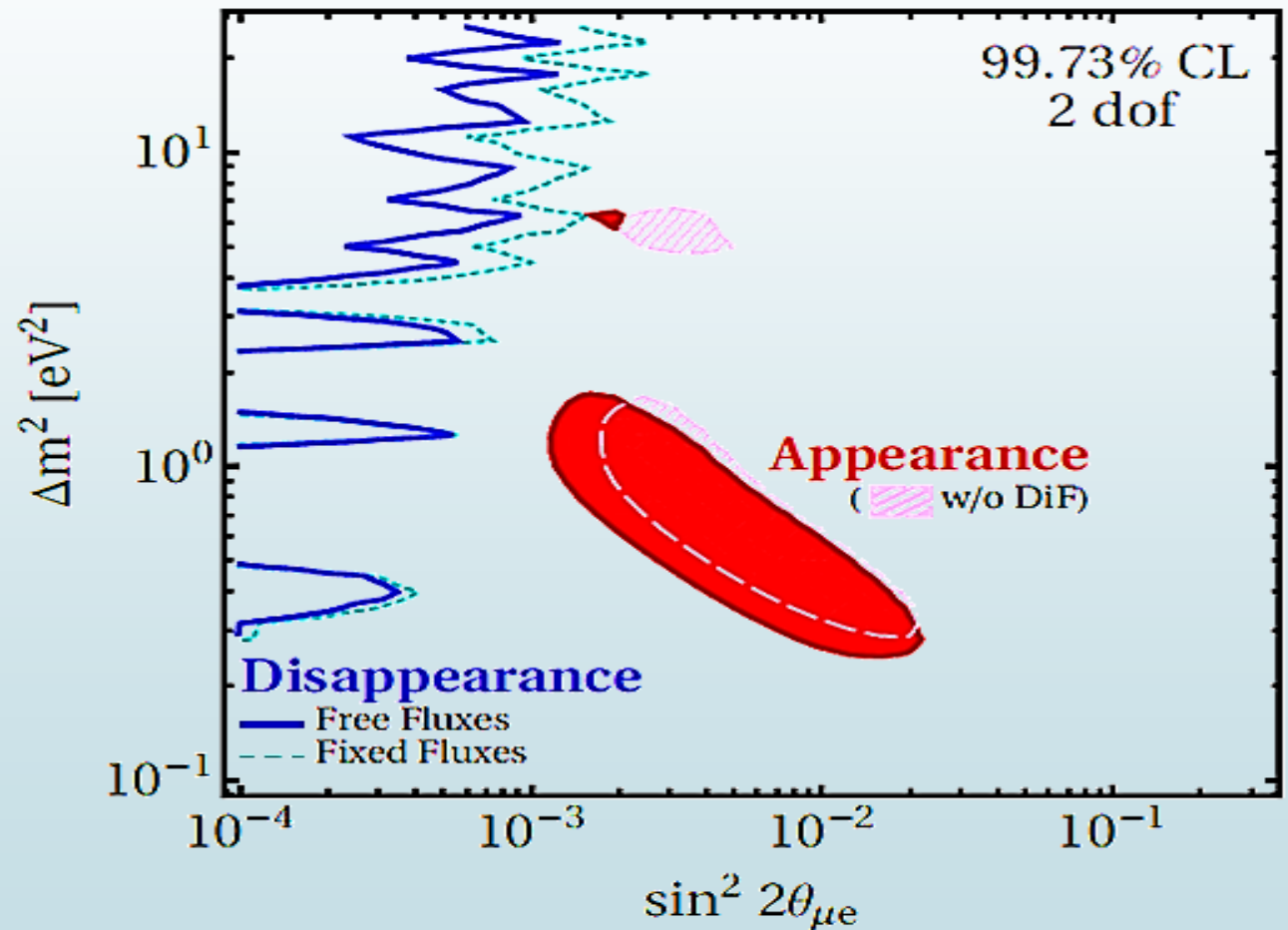


❖ *What about eV Sterile Neutrino Interpretation ???*

$$\sin^2 2\theta_{\mu e} = 4 |U_{e4} U_{\mu 4}|^2$$

4.7 σ tension between
Appearance and
Disappearance data sets
under eV sterile
interpretation

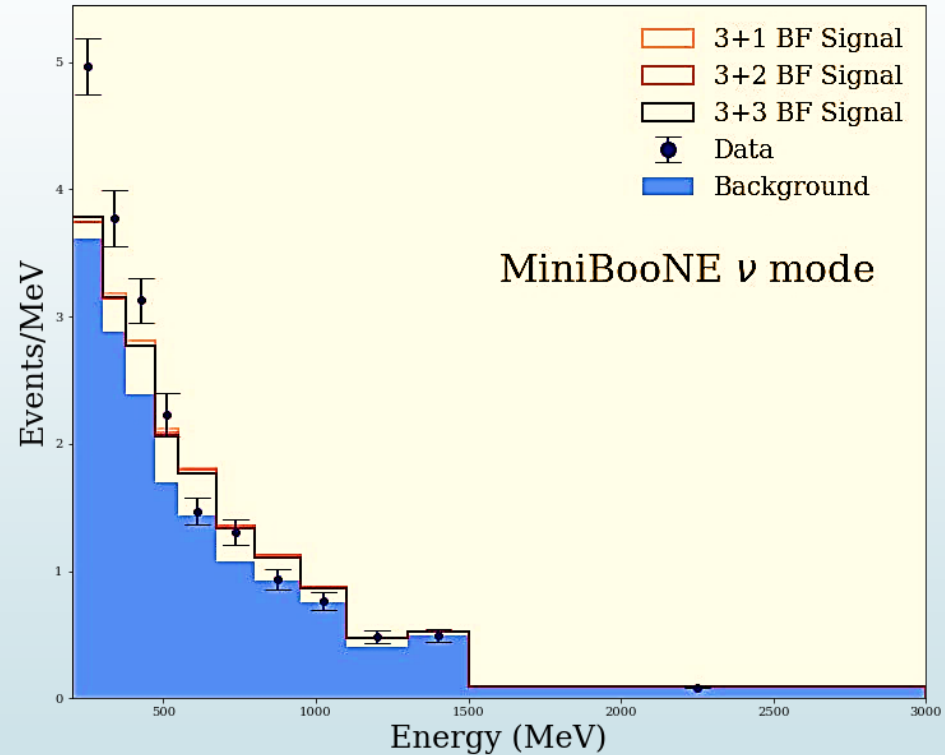
Mona Dentler et al. JHEP 1808 (2018) 010
Collin et al. 1602.00671
Gariazzo et al 1703.00860



➤ 3+N GLOBAL FITS

Shortcoming:
Failure to accommodate
MiniBooNE low-energy excess.

“3+N
STANDARD
STERILE
NEUTRINOS”:
INSUFFICIENT



D. Cianci, et al. (Talk presented at Applied Antineutrino Physics Workshop 2018)

❖ What about eV Sterile Neutrino Interpretation ???

Sterile neutrinos require $\sin^2 2\theta_{\mu e} > 10^{-3}$, $m_4 < \text{few eV}$

Generic early universe thermalization

$$\Gamma > H \implies \sin^2 2\theta_{\mu e} G_F^2 T^5 > \sqrt{g_*} \frac{T^2}{m_{\text{Pl}}} \implies n_4 \sim n_\nu$$

Excluded by BBN/CMB $N_{\text{eff}} = 2.99 \pm 0.17$ Planck 1807.06209

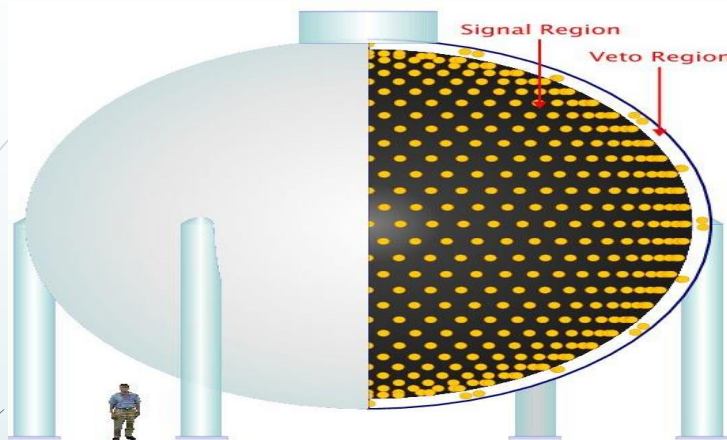
Unless max temperature satisfies $T_{\text{max}} \lesssim 15 \text{ MeV} \left(\frac{10^{-3}}{\sin^2 2\theta_{\mu e}} \right)^{1/3}$



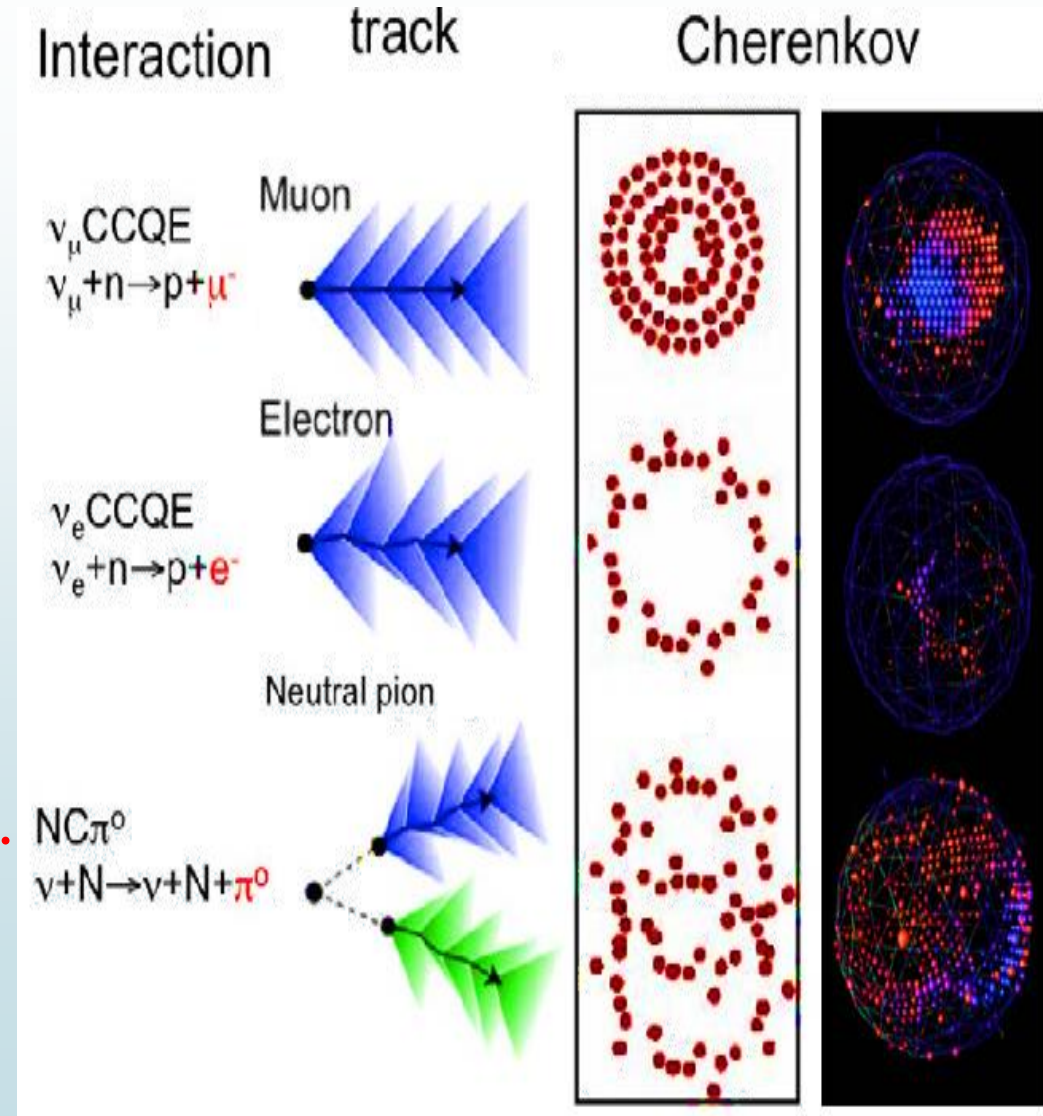
□ *Explanation of MiniBooNE's low energy excess*

- ❖ **Sterile ν at the eV scale present strong tension between data sets**
- ❖ **Cosmological bounds further threaten the eV sterile ν hypothesis**
- ❖ **Is there an explanation that is not ruled out?**
- ❖ **Is there a “real model” for these explanations?**
- ❖ **Can this relate to any of the theoretical problems of the SM?**

❖ *Explanation of MiniBooNE's low energy excess*

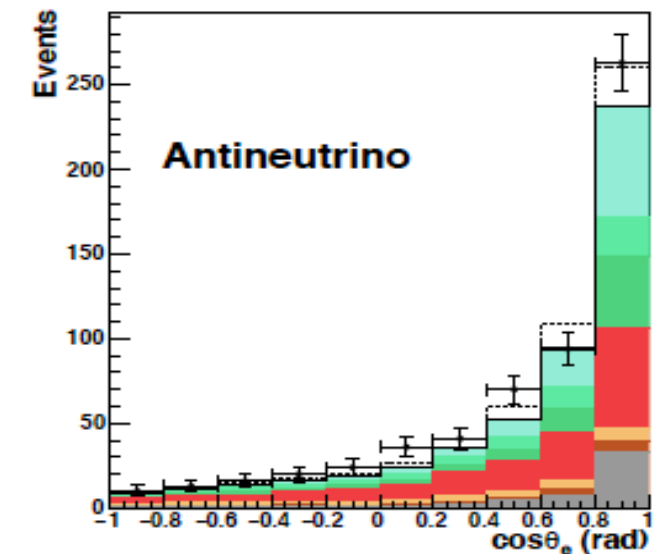
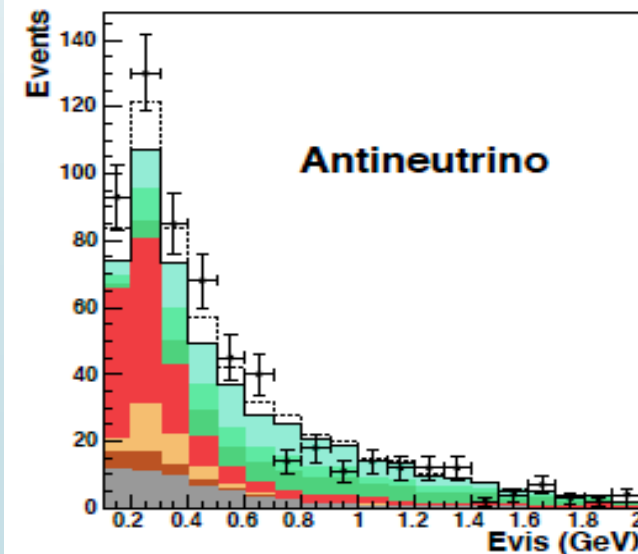
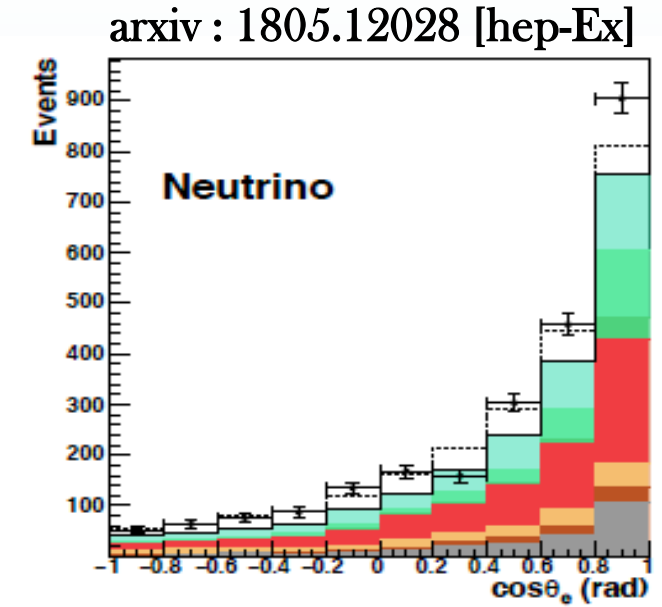
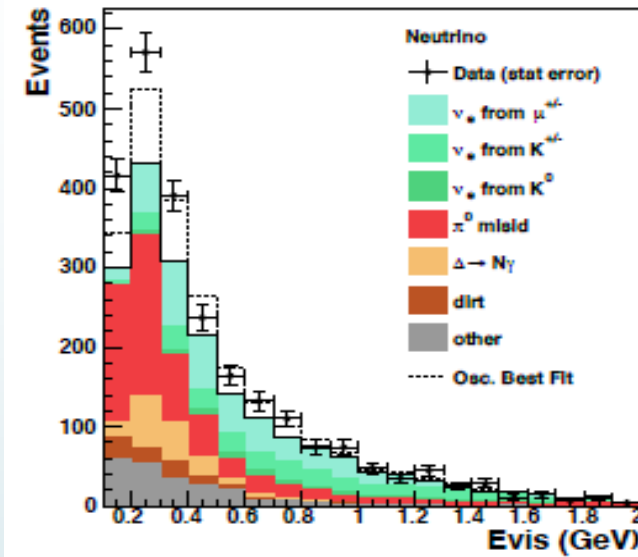


- MiniBooNE is a mineral oil (CH_2) detector that can observe Cherenkov radiation of charged particles.
- Crucially, it could not distinguish electron induced Cherenkov cones from photon induced Cherenkov cones.
- Excess is correlated with beam in power, angle and timing. It is present in positive and negative horn polarities. It is not present in beam dump configuration



❖ *Explanation of MiniBooNE's low energy excess*

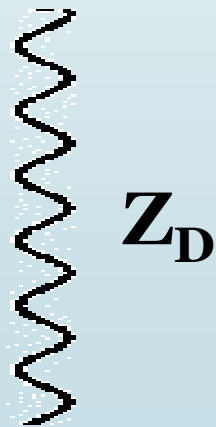
- Angular spectrum is forward, but not that much
- Scattering on electrons would typically lead to $\cos\theta > 0.99$
- Decays of invisible light (< 10 MeV) particles produced in the beam would also lead to forward spectrum



❖ *Explanation of MiniBooNE's low energy excess*

A LIGHT DARK SECTOR - THE IDEA

- There is a dark sector with a novel interaction



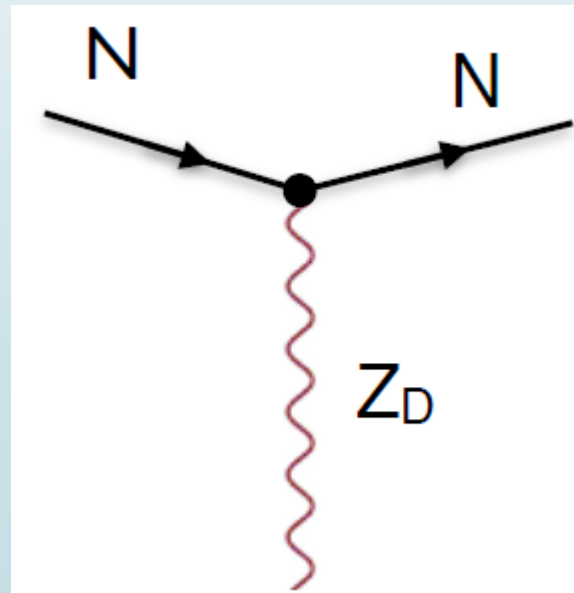
Bertuzzo et al 1807.09877

Bertuzzo et al 1808.02500

❖ *Explanation of MiniBooNE's low energy excess*

A LIGHT DARK SECTOR - THE IDEA

- There is a dark sector with a novel interaction
- Right-handed neutrinos are part of the dark sector and are subject to new interaction



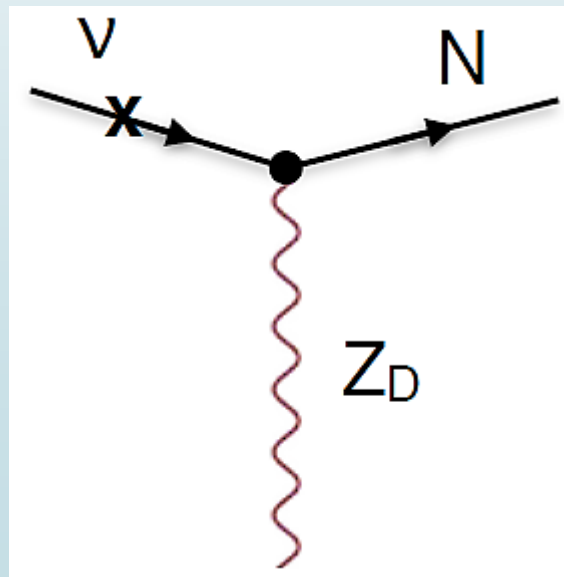
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- There is a dark sector with a novel interaction
- Right-handed neutrinos are part of the dark sector and are subject to new interaction
- Mixing between RH and LH neutrinos leads to interaction in active neutrino sector



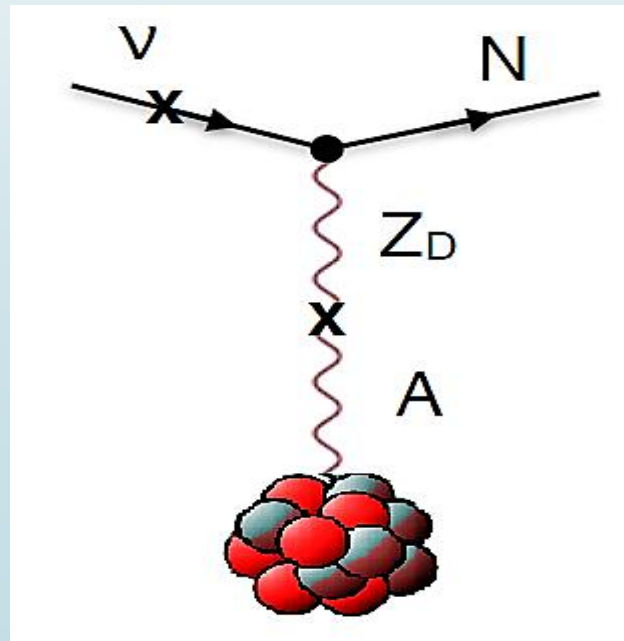
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- Mixing between Z_D and photon leads to interaction with protons



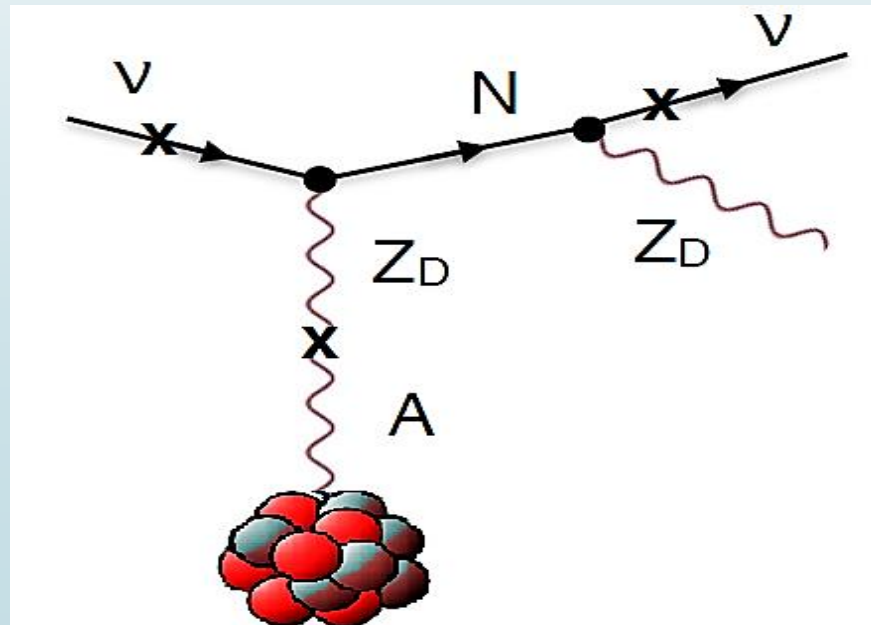
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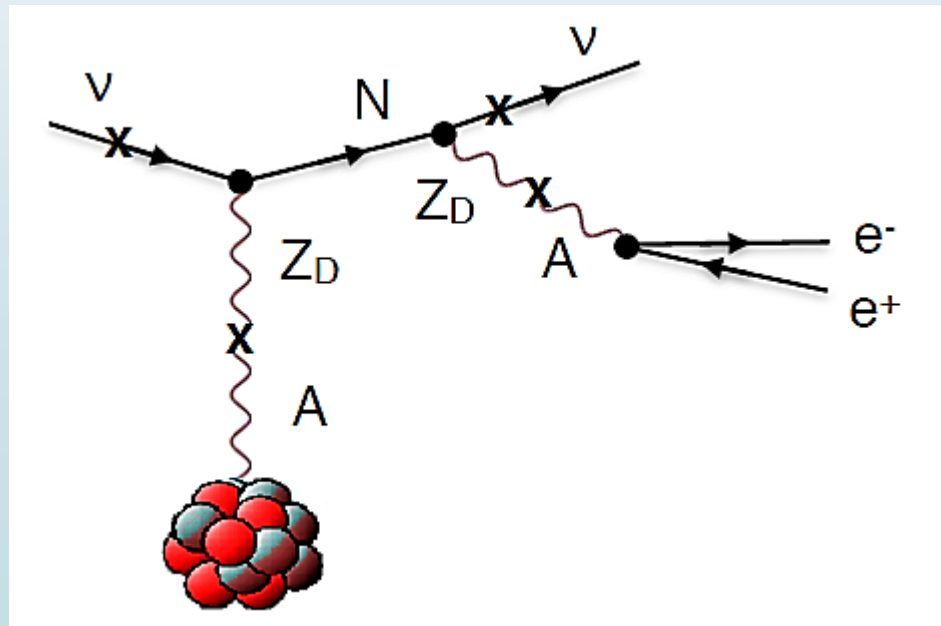
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Bertuzzo et al 1807.09877

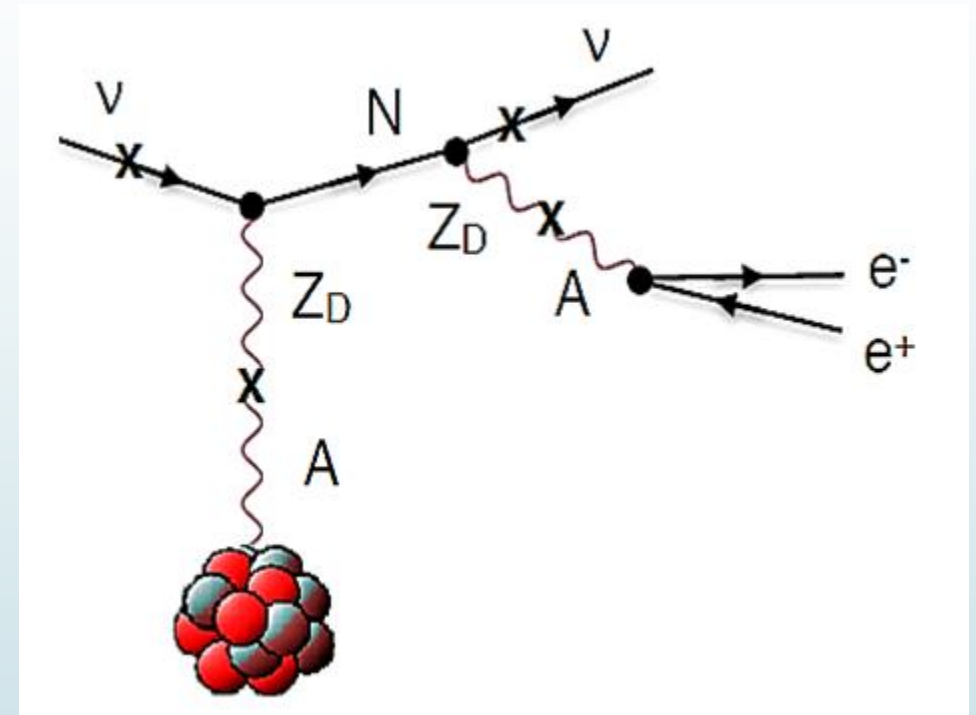
Bertuzzo et al 1808.02500

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- Relevant part of the Lagrangian :

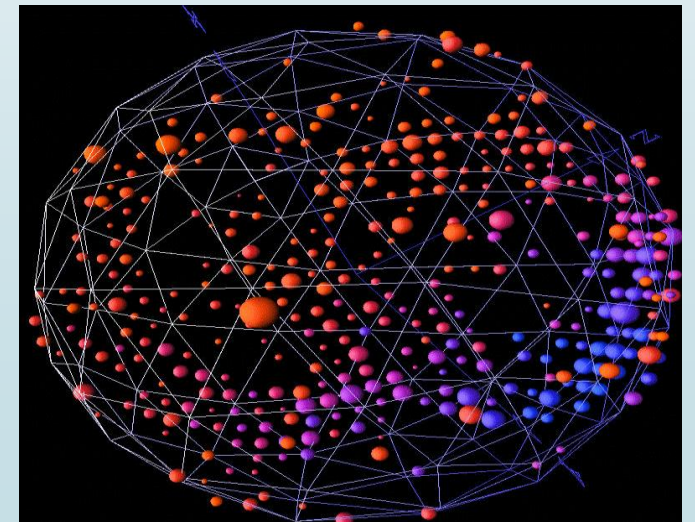
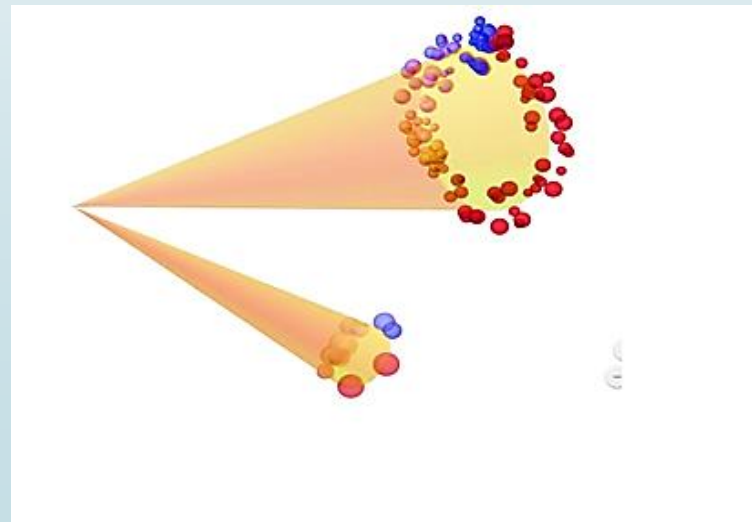
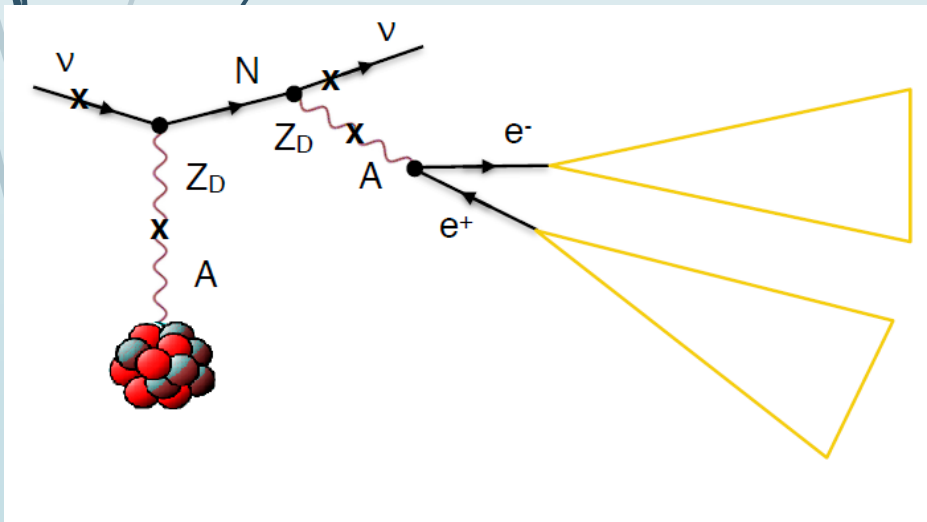
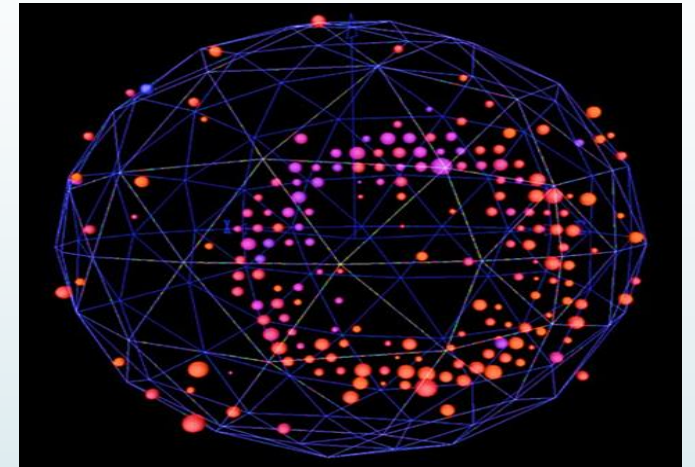
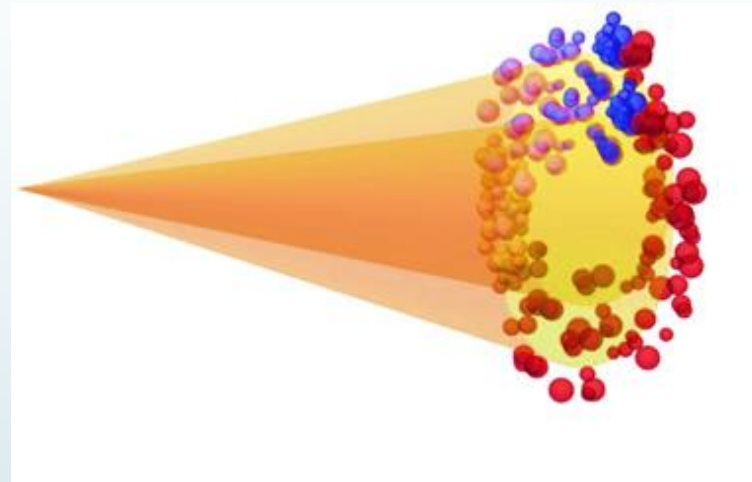
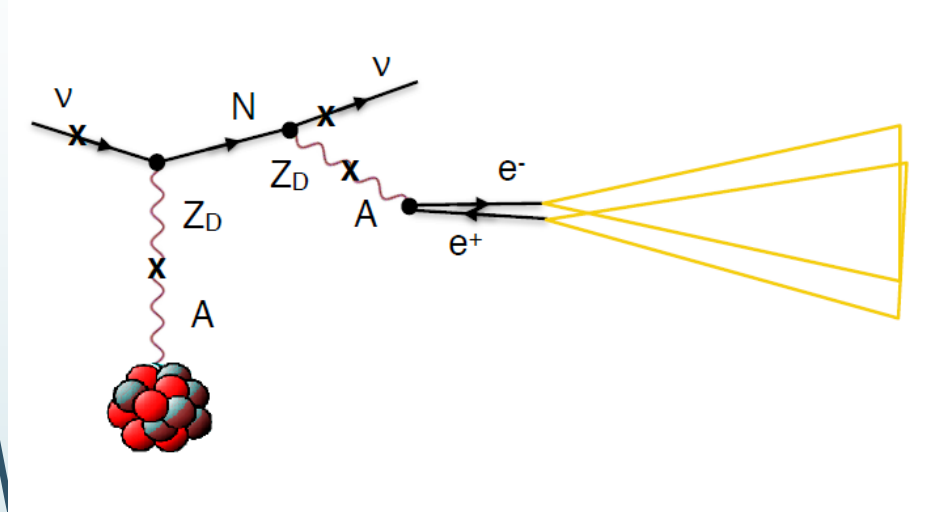
$$\mathcal{L}_D \supset \frac{m_{Z_D}^2}{2} Z_{D\mu} Z_D^\mu + g_D Z_D^\mu J_{D\mu} + e\epsilon Z_D^\mu J_\mu^{\text{em}} + \frac{g}{c_W} \epsilon' Z_D^\mu J_\mu^Z$$



Bertuzzo et al 1807.09877
Bertuzzo et al 1808.02500

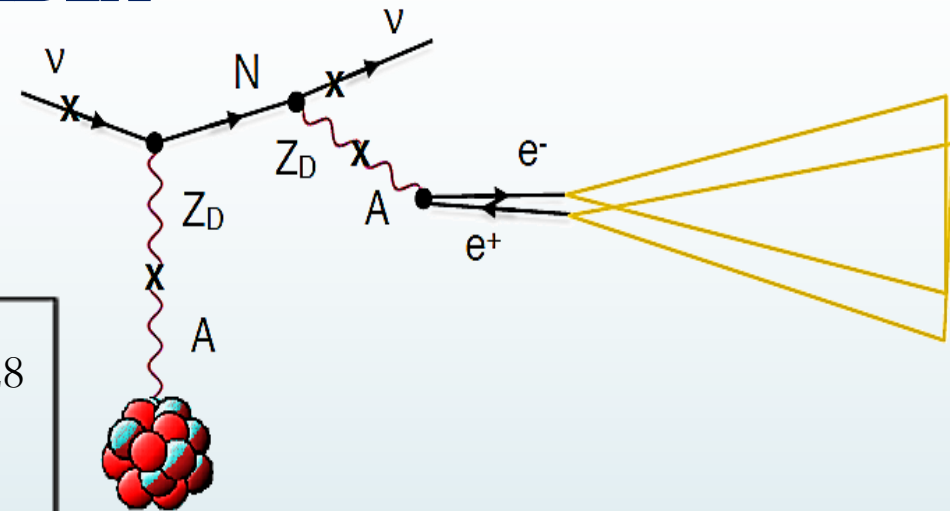
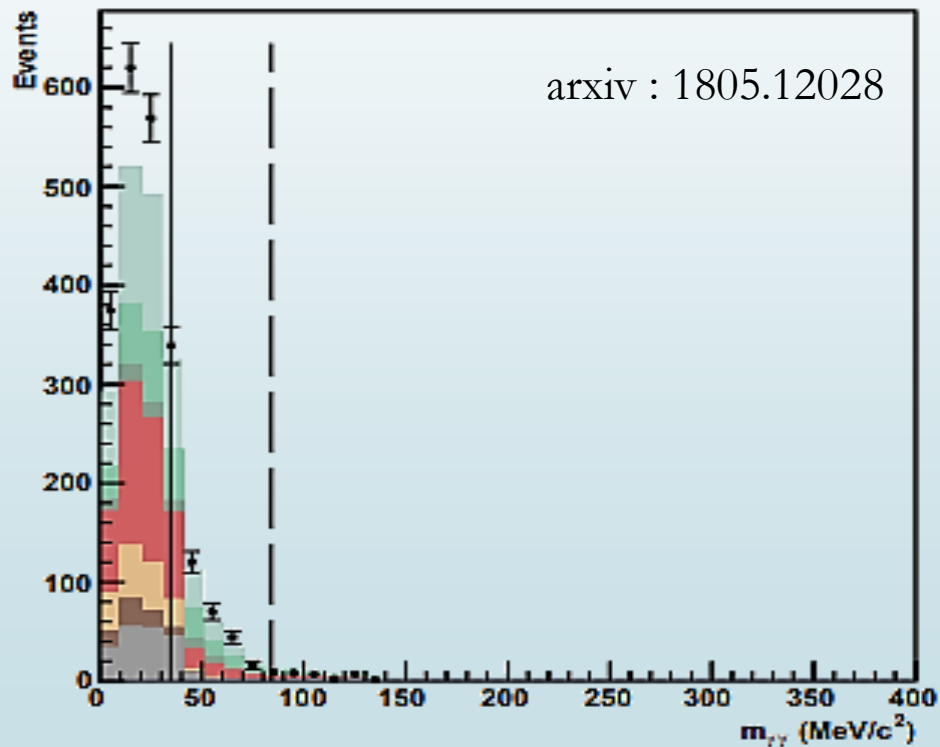
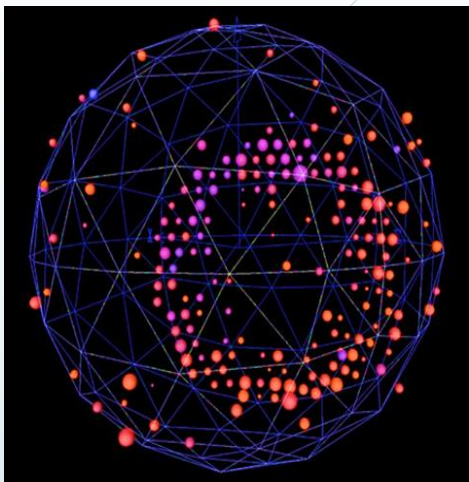
❖ *Explanation of MiniBooNE's low energy excess*

A LIGHT DARK SECTOR - THE IDEA



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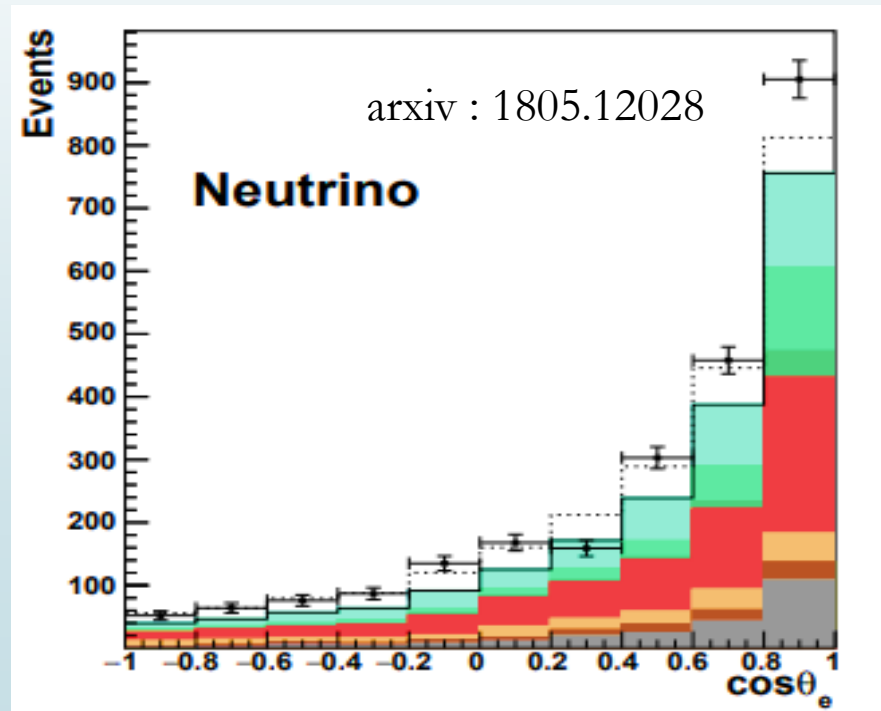
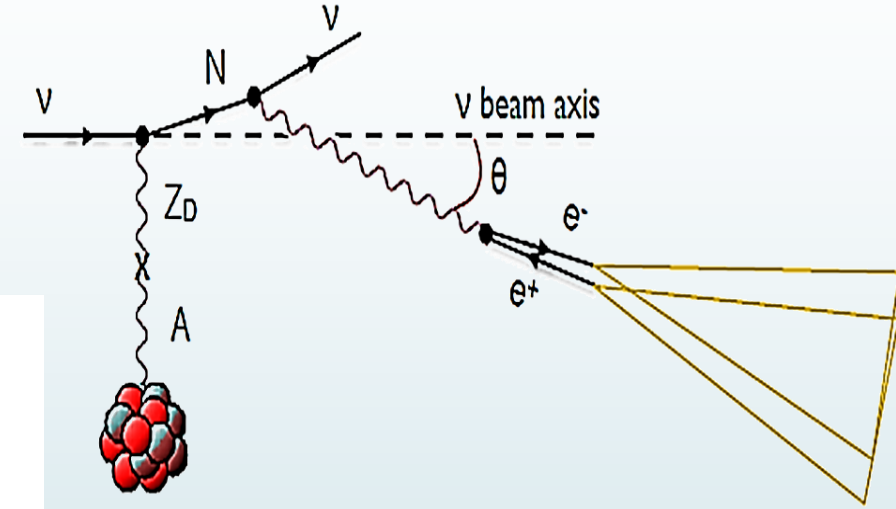
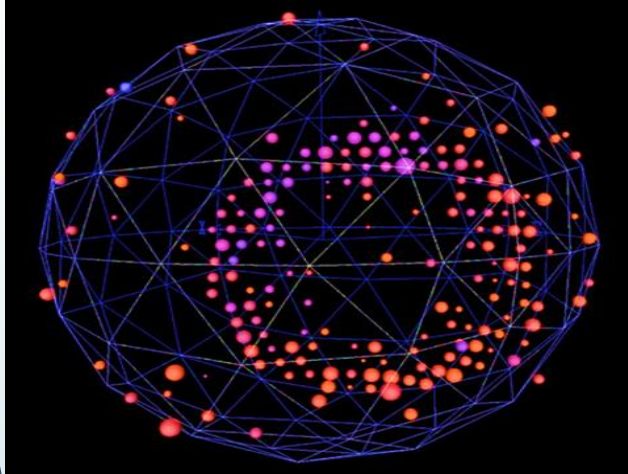
A LIGHT DARK SECTOR - THE IDEA



If e^+e^- pair is collimated ($\cos\theta_{ee} > 0.99$ -ish), it will be classified as e-like

❖ *Explanation of MiniBooNE's low energy excess*

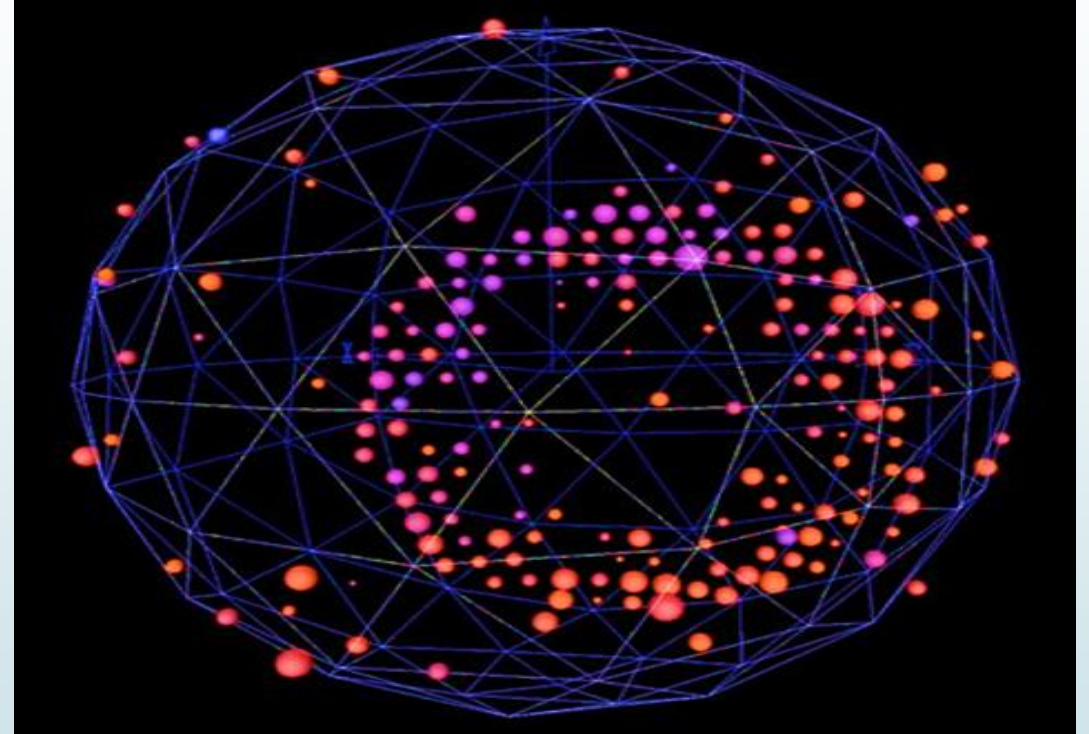
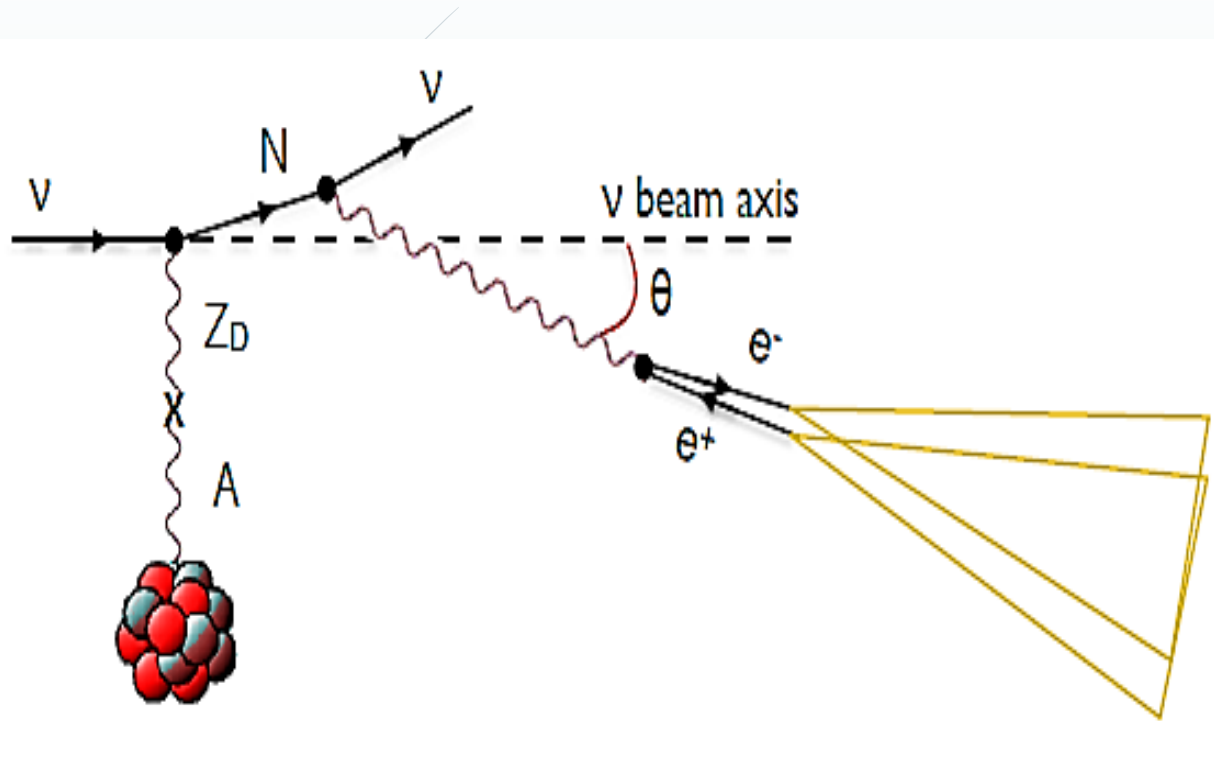
A LIGHT DARK SECTOR - THE IDEA



We have to get this angular spectrum

❖ *Explanation of MiniBooNE's low energy excess*

A LIGHT DARK SECTOR - THE IDEA



- (1) N_D should be heavy (> 100 MeV) so its decay products are not so boosted
- (2) Z_D should be light (< 60 MeV) so that the e^+e^- pair is collimated

❖ *Explanation of MiniBooNE's low energy excess*

Fit to energy spectrum only
(Official MB data release)

Benchmark Points :

$$m_N = 420 \text{ MeV}$$

$$m_{ZD} = 30 \text{ MeV}$$

$$|U_{\mu 4}|^2 = 9 \times 10^{-7}$$

$$\alpha_D = 0.25$$

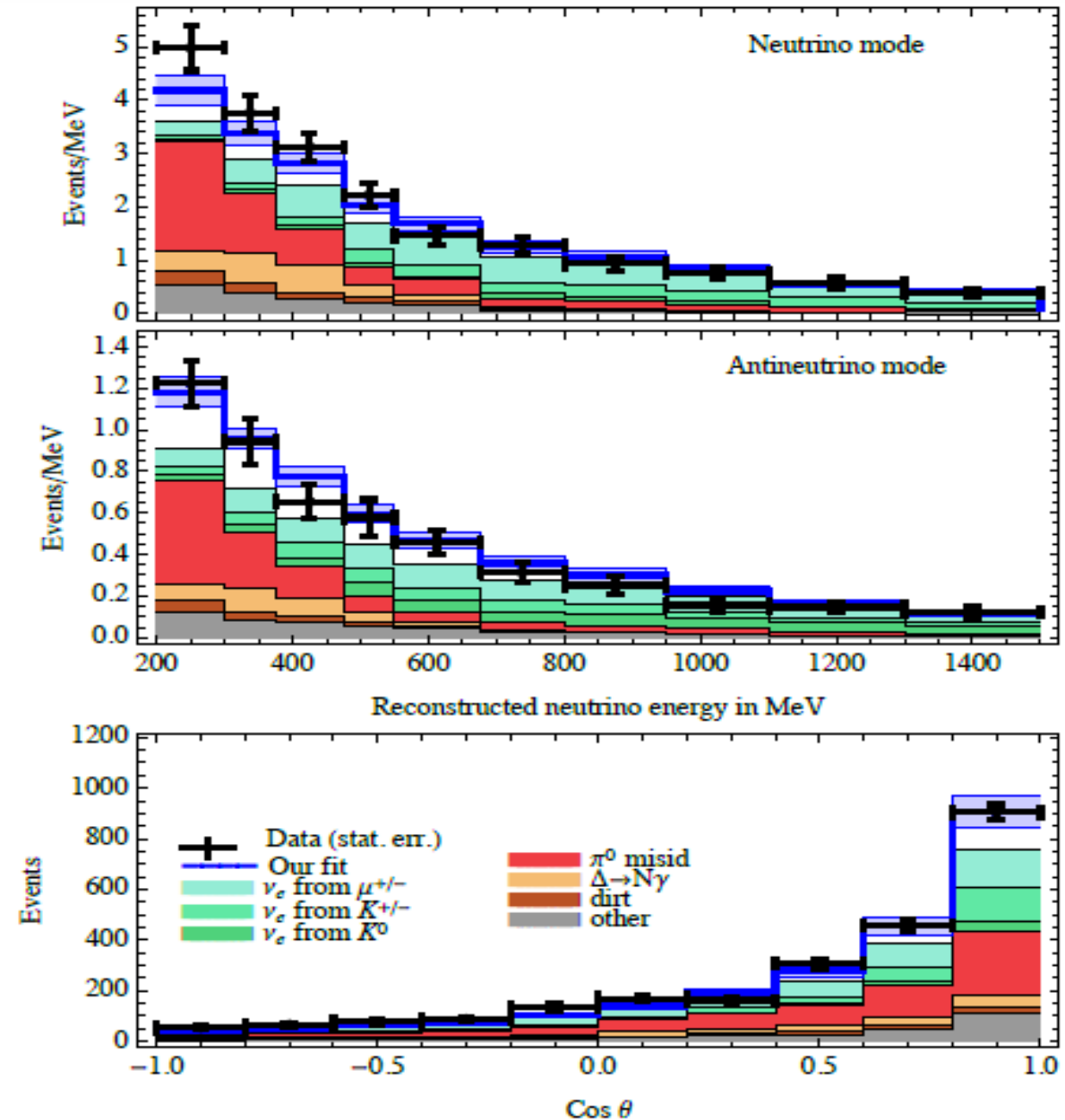
$$\alpha \varepsilon^2 = 2 \times 10^{-10}$$

$$\chi^2/\text{dof} = 33.2/36$$

[Bertuzzo et al 1807.09877](#)

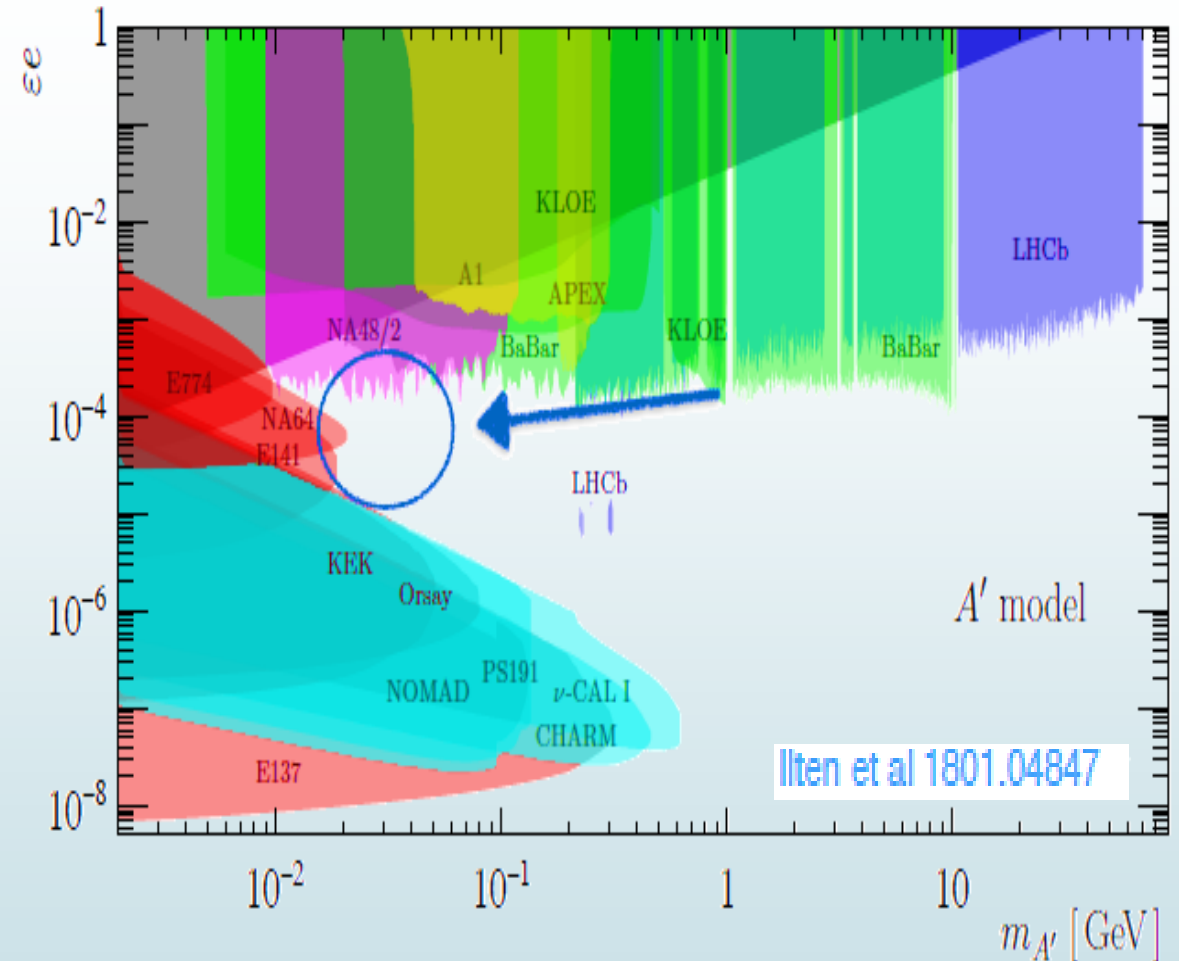
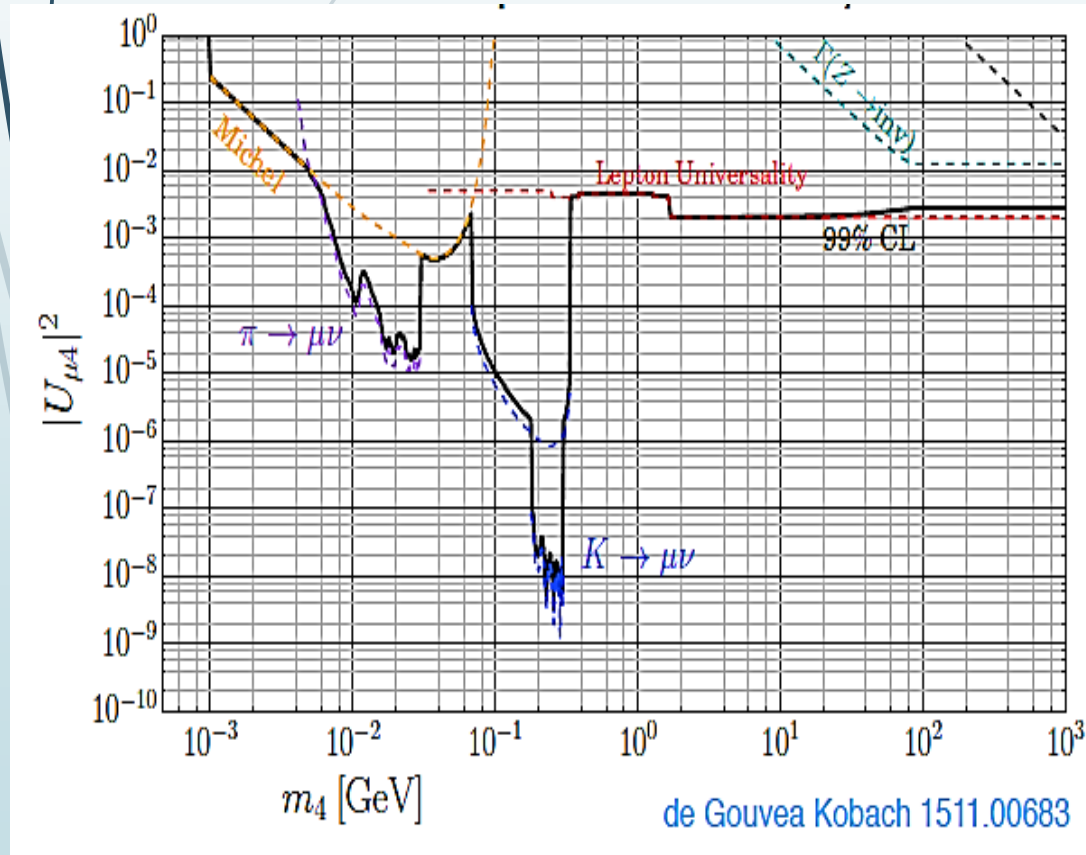
See also [Ballett et al 1808.02915](#)

for different realization of the mechanism



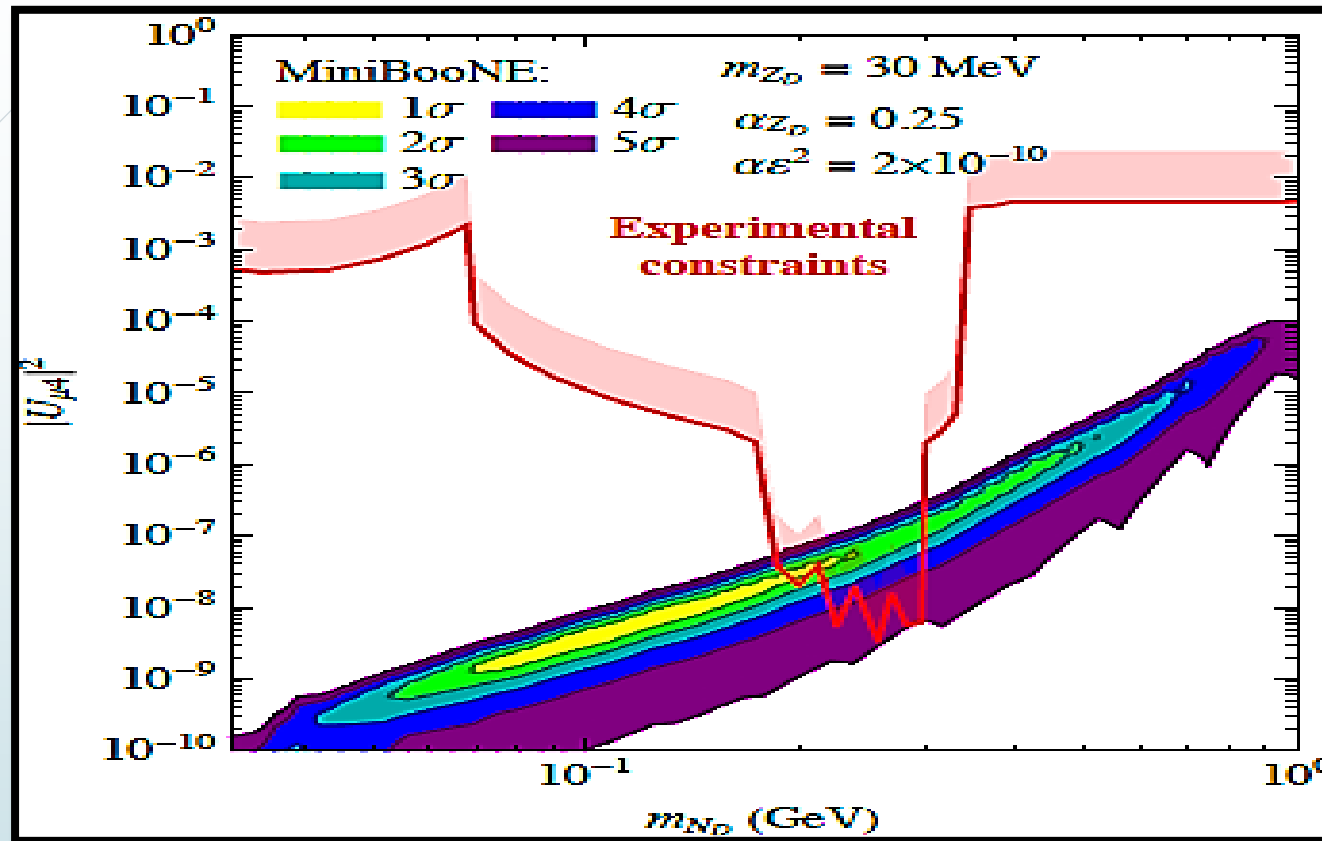
❖ Constraint on Light Dark Sector

Model Independent Constraint on Heavy Sterile Neutrino



- Z_D phenomenology is similar to dark photon case
- LHC constraints are not expected to be stringent below 1 GeV

❖ Explanation of MiniBooNE's low energy excess



Bertuzzo et al 1807.09877

Region of our model in the $|U_{\mu 4}|^2$ versus m_{N_D} plane satisfying MiniBooNE data at 1 σ to 5 σ CL, for the hypothesis $m_{Z_D} = 30 \text{ MeV}$, $\alpha_{Z_D} = 0.25$ and $\alpha \epsilon^2 = 2 \times 10^{-10}$. The region above the red curve is excluded at 99% CL by meson decays, the muon decay Michel spectrum and lepton universality

A dark blue arrow points to the right from the left edge of the slide. Several thin, light blue lines curve upwards from the bottom left corner towards the center of the slide.

Connection to Neutrino Mass Generation Mechanism

Scale of Seesaw Mechanism

- ❖ **Despite numerous searches for neutrino mass models (at TeV scale) at high-energy colliders, no compelling evidence has been found so far.**
- ❖ **Is it really sufficient to search for new physics scale behind neutrino mass generation mechanism at LHC only ?**
- ❖ **The new physics scale behind neutrino mass generation mechanism might be at low scale and which is less sensitive to high energy collider experiments**
- ❖ **It may show up at low energy neutrino experiments at near future.**

Scale of Seesaw Mechanism

❖ Despite numerous searches for neutrino mass models (at TeV scale) at high-energy colliders, no compelling evidence has been found.

❖ Is it really sufficient to search for physics scale behind neutrino mass?

Can neutrino masses come from light physics?
Seesaw mechanism might be at low scale and is less sensitive to high energy collider experiments

❖ It may show up at low energy neutrino experiments at near future.

Neutrino masses from light physics

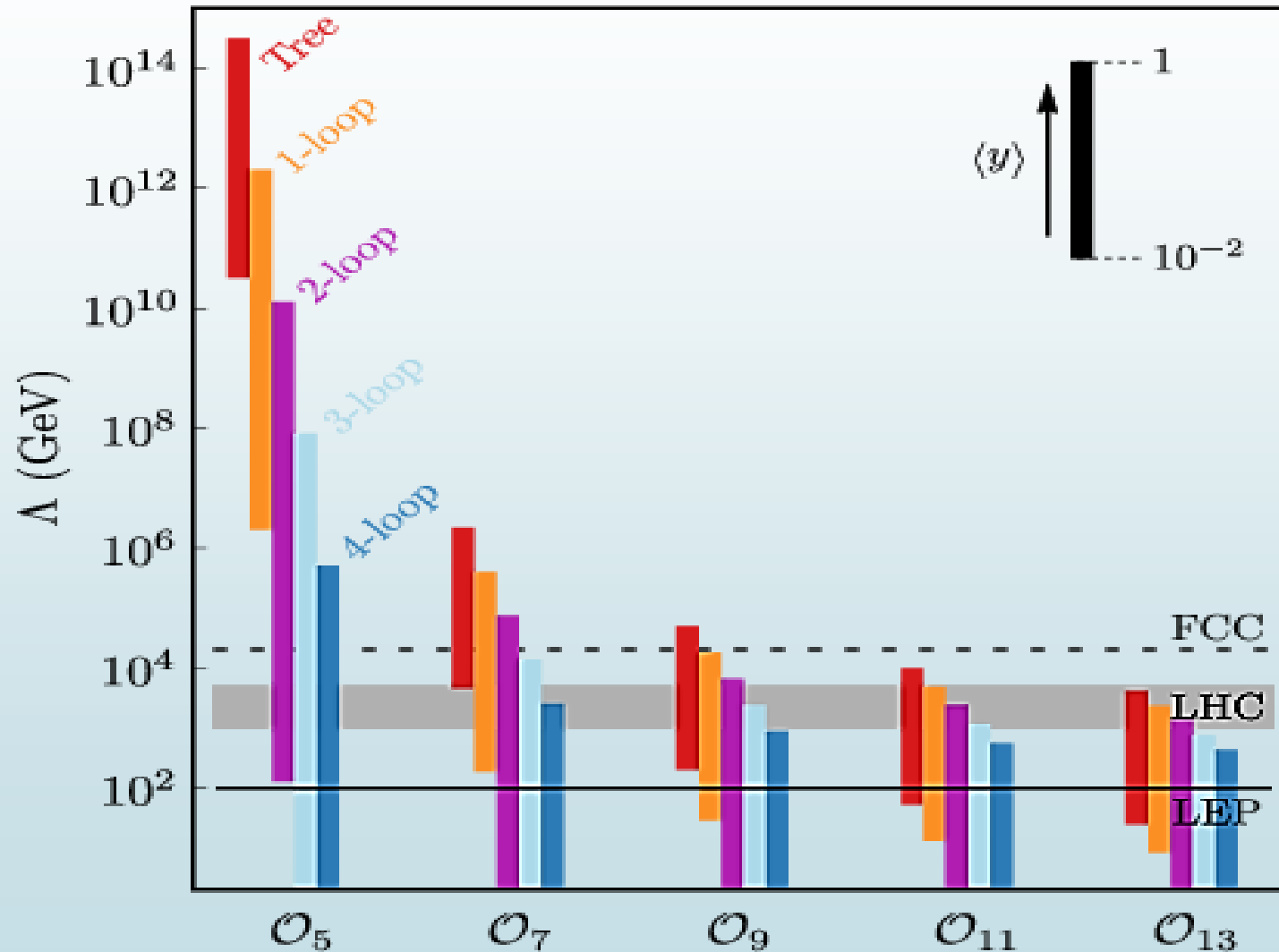
In an effective theory, the Lagrangian should be described as

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda_{\text{NP}}} \mathcal{O}^{d=5} + \frac{1}{\Lambda_{\text{NP}}^2} \mathcal{O}^{d=6} + \frac{1}{\Lambda_{\text{NP}}^3} \mathcal{O}^{d=7} + \dots$$

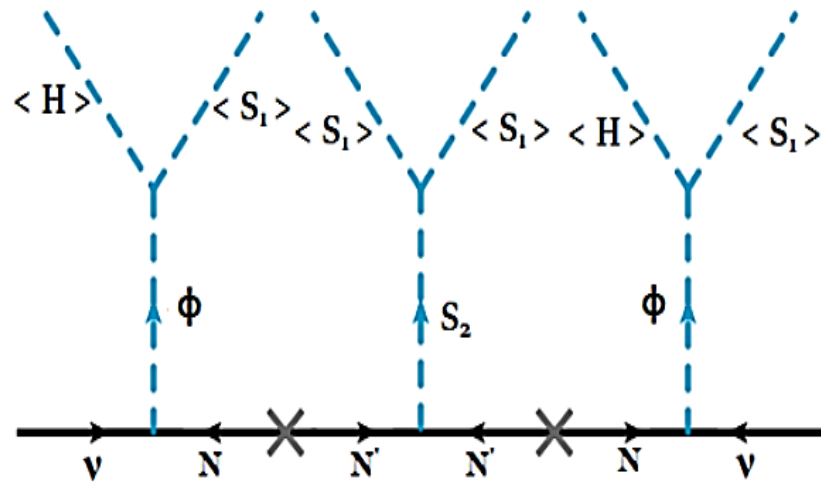
Neutrino masses from a n -loop-induced dim- d operator

$$m_\nu = v \times \left(\frac{1}{16\pi^2} \right)^n \times \left(\frac{v}{\Lambda_{\text{NP}}} \right)^{d-4}$$

Scale of Seesaw Mechanism



Neutrino masses from light physics



$$\mathcal{L}_\nu^{d=9} \sim y_\nu^2 y_N \frac{\mu^2}{M_{H_D}^2} \frac{\mu'}{M_{S'_D}^4} \frac{(\overline{L^c H})(H^T L)}{m^2} (S_1^* S_1)^2$$

Neutrino masses from D=9 operator

All scales involved may be below electroweak

Light Z_D , ν - N mixing, Z_D - ν - N coupling, kinetic mixing unavoidable

Neutrino masses from light physics

Gauge $U(1)_D$: SM has no charge, RH neutrinos N have charge $+1$

Anomaly cancellation: N' with opposite charge should be included

anomaly cancellation is a requirement to have a consistent QFT

Walks and quarks like inverse seesaw

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & m & 0 \\ m & 0 & M \\ 0 & M & \mu \end{pmatrix} \begin{matrix} \nu \\ N \\ N' \end{matrix} \begin{matrix} 0 \\ + \\ - \end{matrix} \implies m_\nu = \mu \frac{m^2}{M^2}$$

m and μ are forbidden by dark symmetry, they need to be generated dynamically

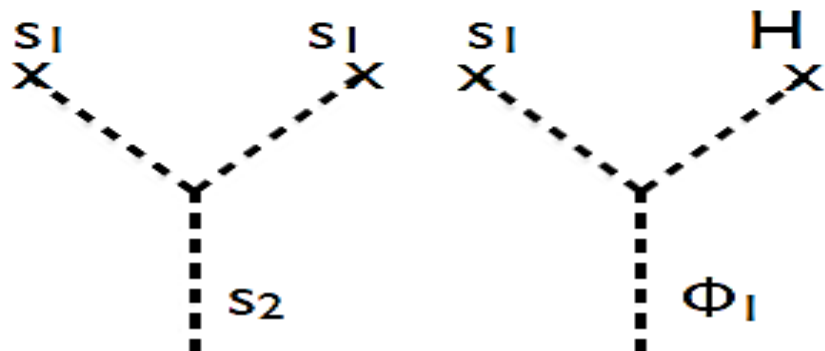
Neutrino masses from light physics

Minimum scalar content

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & y\phi_1 & 0 \\ y\phi_1 & 0 & M \\ 0 & M & y's_2 \end{pmatrix}$$

Φ_1 = doublet with dark charge +1
 s_2 = singlet with dark charge +2

Add s_1 with charge +1 and something special happens:
 Φ_1 and s_2 start with no vevs, s_1 develops a vev like the Higgs



Φ_1 and s_2 vevs are **induced**, like in type II seesaw, and thus can be naturally very small!

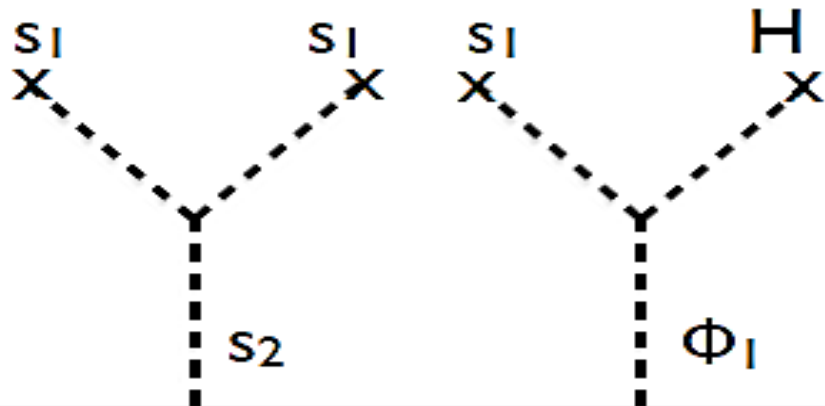
Neutrino masses from light physics

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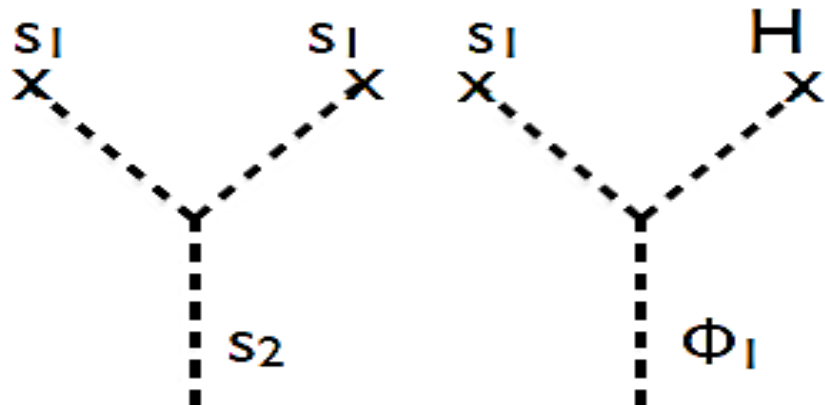
Φ_1 and s_2 vevs are **induced**, like in type II seesaw, and thus can be naturally very small!

Neutrino masses from light physics

Minimum scalar content

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & y\phi_1 & 0 \\ y\phi_1 & 0 & M \\ 0 & M & \underbrace{y's_2}_{\text{red circle}} \end{pmatrix} \quad \mathcal{L}_\nu = -y_\nu \bar{L}\tilde{\phi}N + y_N S_2 \bar{N}N^c + y_{N'} S_2^* \bar{N}'N'^c + m \bar{N}'N^c + \text{h.c.}$$

Add s_1 with charge +1 and something special happens:
 Φ_1 and s_2 start with no vevs, s_1 develops a vev like the Higgs



Φ_1 and s_2 vevs are **induced**, like in type II seesaw, and thus can be naturally very small!

Scale of Seesaw Mechanism

▶ Seesaw I mechanism with TeV scale heavy neutrinos

- Standard Seesaw with small Yukawa couplings

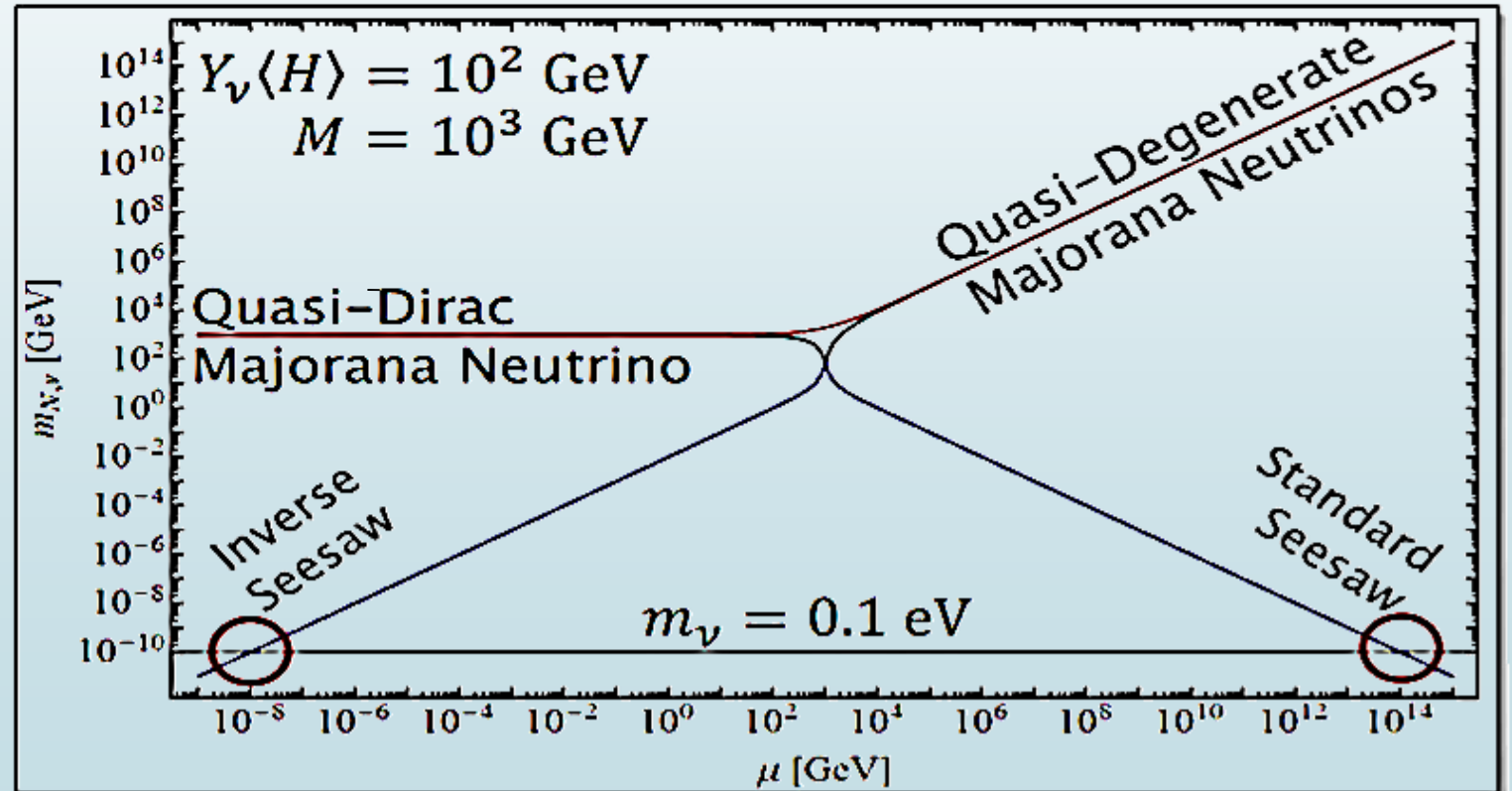
$$Y_\nu \approx 10^{-6} \sqrt{M_N/\text{TeV}}$$

- “Bent” Seesaw I mechanisms (e.g. Inverse Seesaw)

- Decouple Δ_{LNV} from heavy neutrino mass
- Example

$$\begin{pmatrix} 0 & Y_\nu \langle H \rangle & 0 \\ Y_\nu \langle H \rangle & \mu & M \\ 0 & M & \mu \end{pmatrix}$$

- Large Yukawa couplings $\approx 10^{-2}$
- Quasi-Dirac heavy neutrino



❖ *Inverse Seesaw*

$$y\Psi^c H\ell + m_\Psi \Psi\Psi^c + \frac{1}{2}\mu\Psi\Psi$$

Ψ, Ψ^c **Pseudo-Dirac**

$$m_\nu \sim \frac{y^2 v^2}{m_\Psi^2} \mu$$

$$y \sim 0.1 \quad m_\Psi \sim 1\text{TeV} \quad \mu \sim 1\text{keV}$$

[Mohapatra, 86]
[Mohapatra, Valle, 86]



- Why μ is much smaller than TeV scale?

Neutrino masses from light physics

Vacuum Expectation Values

v (GeV)	ω_1 (MeV)	v_ϕ (MeV)	ω_2 (MeV)
246	136	0.176	0.65

Coupling Constants

λ_H	$\lambda_{H\phi} = \lambda'_{H\phi}$	λ_{HS_1}	λ_{HS_2}
0.129	10^{-3}	10^{-3}	-10^{-3}
$\lambda_{\phi S_1}$	$\lambda_{\phi S_2}$	λ_{S_1}	$\lambda_{S_1 S_2}$
10^{-2}	10^{-2}	2	0.01
μ (GeV)	μ' (GeV)	α	g_D
0.15	0.01	10^{-3}	0.22

Bare Masses

m_ϕ (GeV)	m_2 (GeV)
100	5.51

$$\begin{aligned}
 V = & -m_H^2(H^\dagger H) + m_\phi^2(\phi^\dagger \phi) - m_1^2 S_1^* S_1 + m_2^2 S_2^* S_2 \\
 & - \left[\frac{\mu}{2} S_1(\phi^\dagger H) + \frac{\mu'}{2} S_1^2 S_2^* + \frac{\alpha}{2} (H^\dagger \phi) S_1 S_2^* + \text{h.c.} \right] \\
 & + \lambda'_{H\phi} \phi^\dagger H H^\dagger \phi + \sum_{\varphi \in \{H, \phi, S_1, S_2\}} \lambda_\varphi (\varphi^\dagger \varphi)^2 \\
 & + \sum_{\varphi < \varphi'}^{\{H, \phi, S_1, S_2\}} \lambda_{\varphi\varphi'} (\varphi^\dagger \varphi) (\varphi'^\dagger \varphi').
 \end{aligned}$$

$$v_\phi \simeq \frac{1}{8\sqrt{2}} \left(\frac{\alpha \mu' v \omega_1^3}{M_{S_D'}^2 M_{H_D}^2} + 4 \frac{\mu \omega_1 v}{M_{H_D}^2} \right) \quad \omega_2 \simeq \frac{1}{8\sqrt{2}} \left(\frac{\alpha \mu v^2 \omega_1^2}{M_{S_D'}^2 M_{H_D}^2} + 4 \frac{\mu' \omega_1^2}{M_{S_D'}^2} \right)$$

Masses of the Physical Fields

$m_{h_{SM}}$ (GeV)	m_{H_D} (GeV)	m_{S_D} (MeV)	$m_{S_D'}$ (MeV)	$m_{H_{\mathbb{D}}^\pm}$ (GeV)	m_{A_D} (GeV)	m_{a_D} (MeV)	m_{Z_D} (MeV)	m_{N_D} (MeV)
125	100	272	320	100	100	272	30	150

Mixing between the Fields

$\theta_{H\phi}$	θ_{HS_1}	θ_{HS_2}	$\theta_{\phi S_1}$	$\theta_{\phi S_2}$	$\theta_{S_1 S_2}$	$\epsilon\epsilon$	ϵ'	$ U_{\alpha N} ^2$
1.3×10^{-6}	2.1×10^{-6}	10^{-8}	1.2×10^{-3}	8.3×10^{-7}	3.4×10^{-2}	2×10^{-4}	3.6×10^{-14}	$\mathcal{O}(10^{-6})$

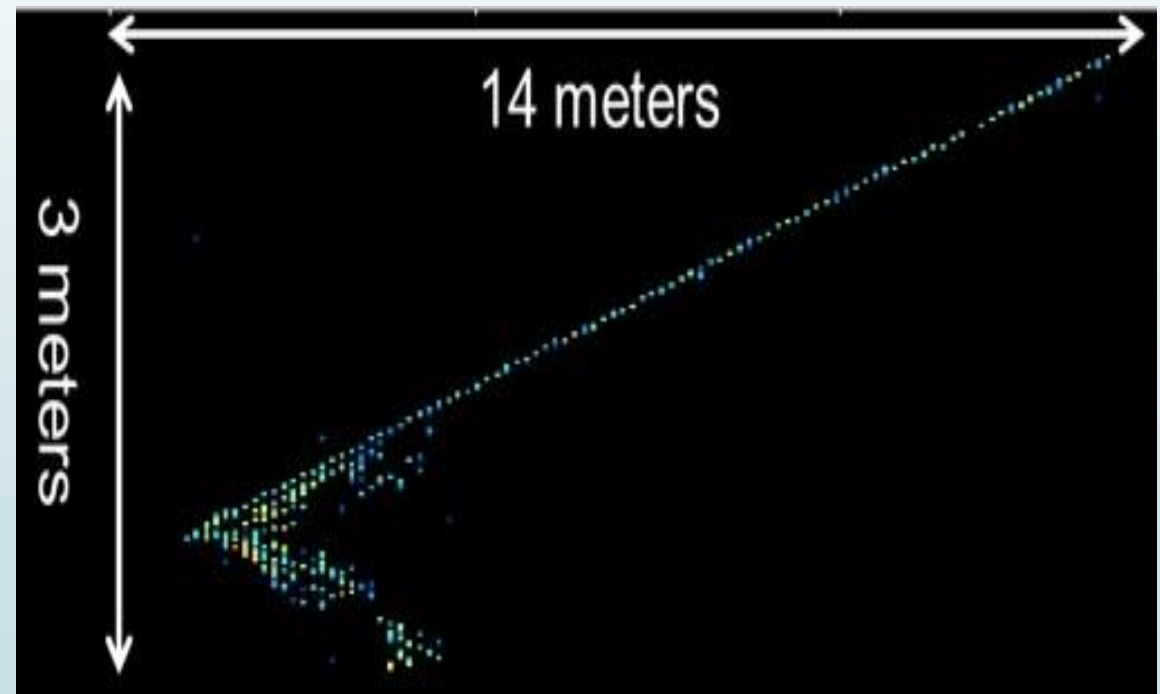
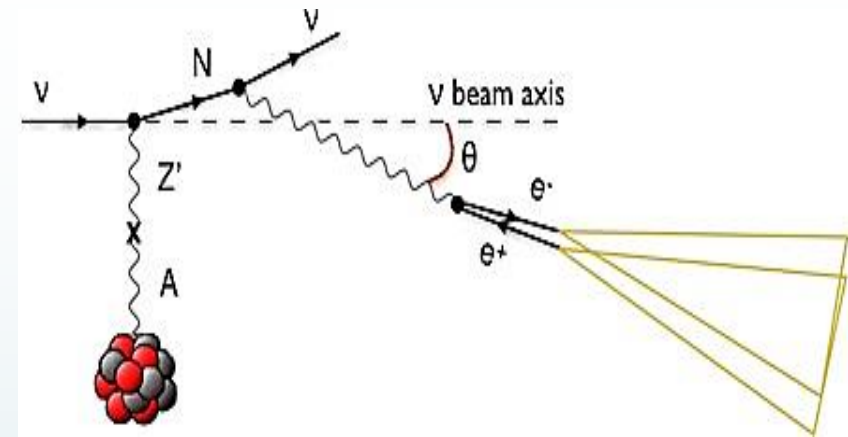
❖ Phenomenology on other neutrino experiment

MiniBooNE's signature: Collimated e^+e^- pair in MINOS+, NOvA, or T2K is likely be tagged as ν_e event

General signature:

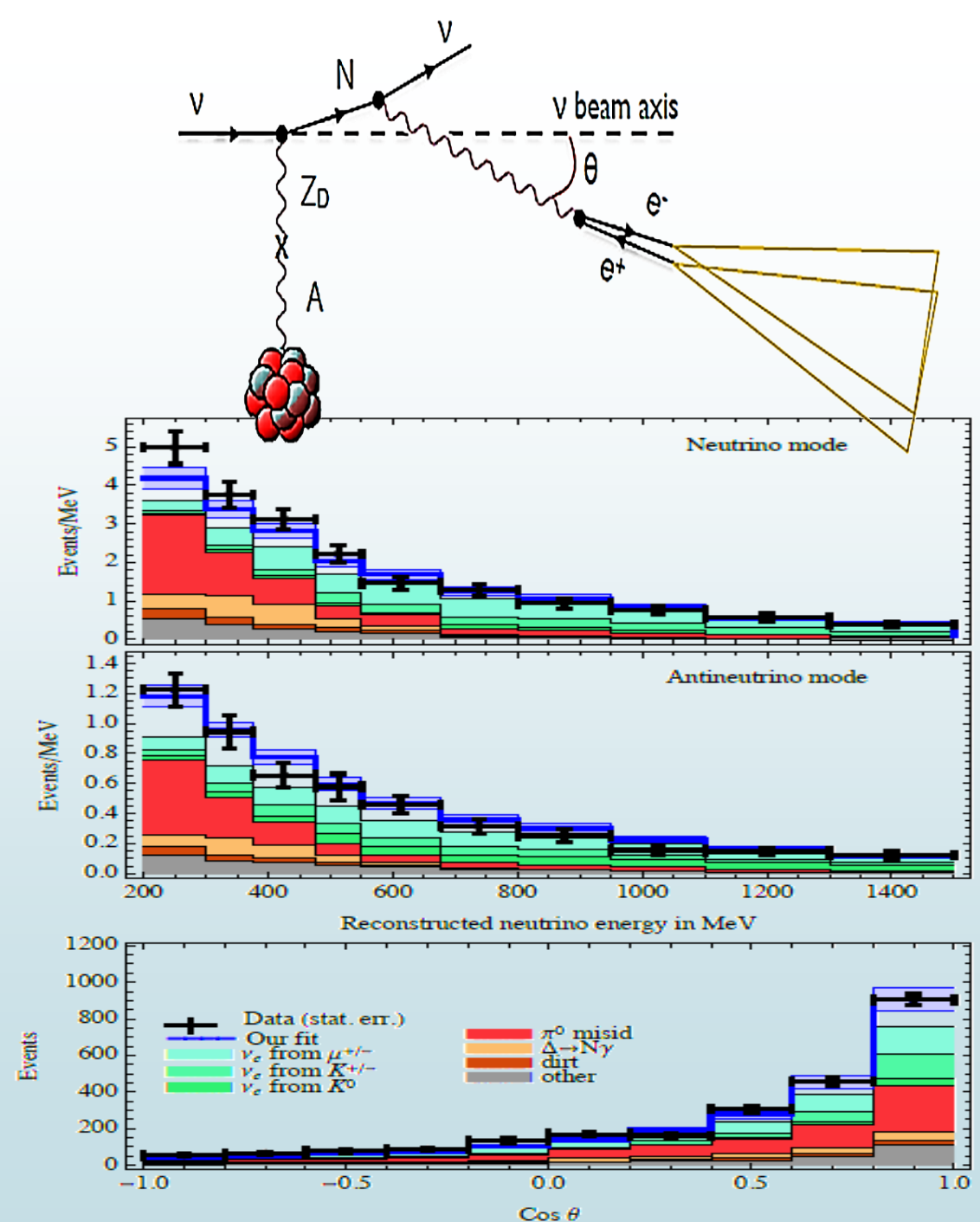
Heavy enough Z_D can decay to $\mu^+\mu^-$ or $\pi^+\pi^-$ pair, much easier signature (MINOS+ is magnetized...)

Lower energy experiments (reactor and solar neutrinos) as well as electron scattering may lack energy to produce N



Conclusions :

- ❖ Novel explanation of MiniBooNE
- ❖ Agreement with all EXP data
- ❖ Novel, simple frameworks
- ❖ Deep connection to neutrino mass generation mechanism
- ❖ A realistic “complete” model below EW scale to explain neutrino mass generation
- ❖ Solves the hierarchy of Inverse Seesaw
- ❖ Rich phenomenology



DO YOU
KNOW MY
ORIGIN ?

PROBABLY
YES !

TOP
SECRET



DO YOU
KNOW MY
ORIGIN ?

PROBABLY
YES !



MiniBooNE

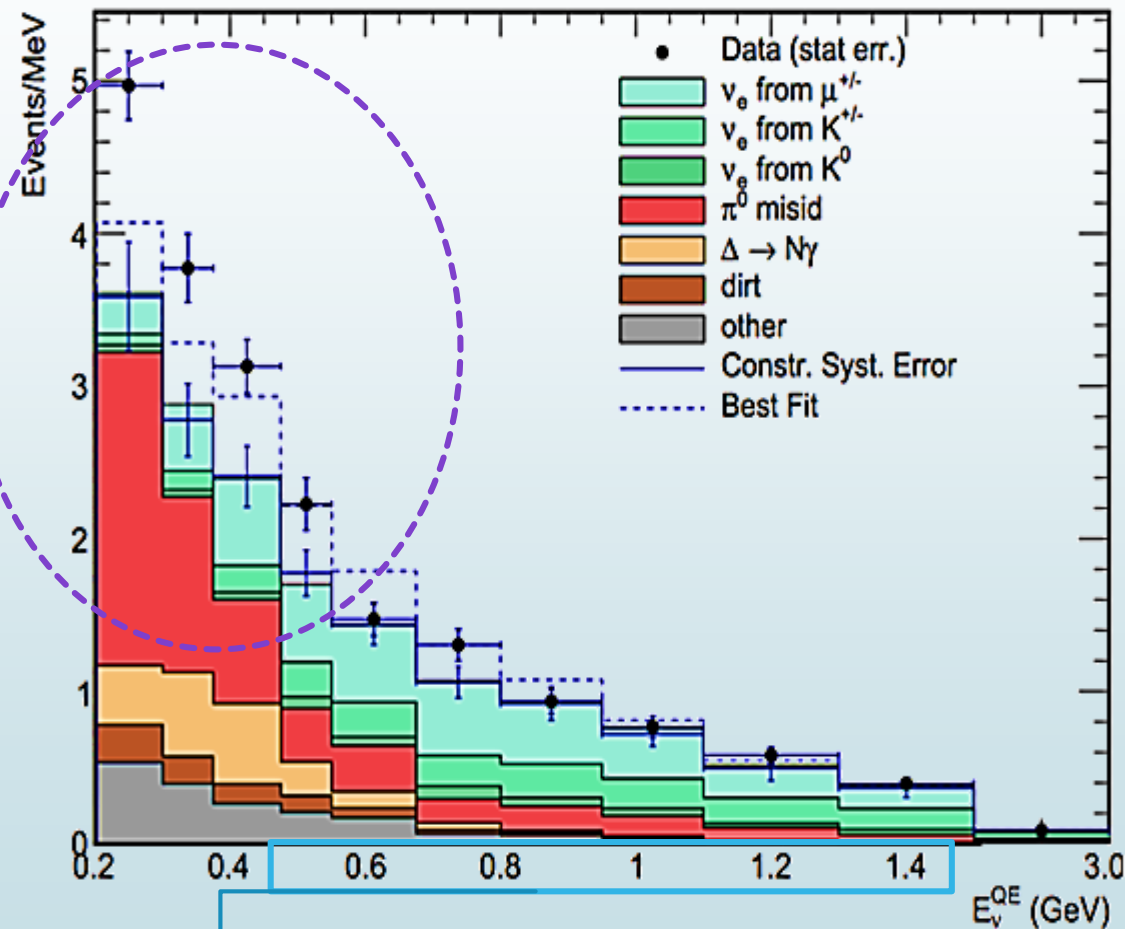




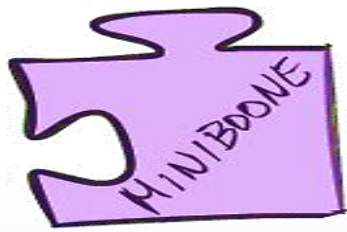
MiniBooNE's Low Energy Excess

MiniBooNE Collaboration [hep-ex/1805.12028](https://arxiv.org/abs/hep-ex/1805.12028)

MiniBooNE
observed a large
excess in the much
lower energy region

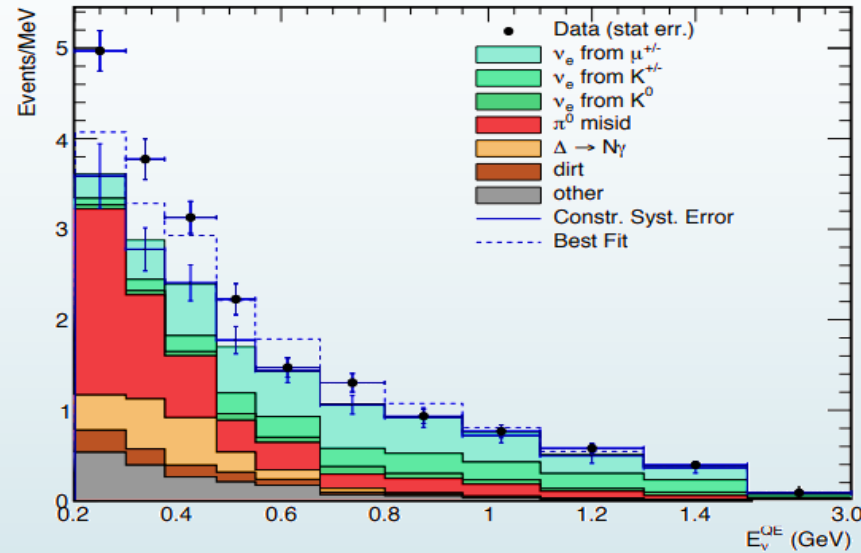


**This is the region that
the LSND oscillation
should have appeared**

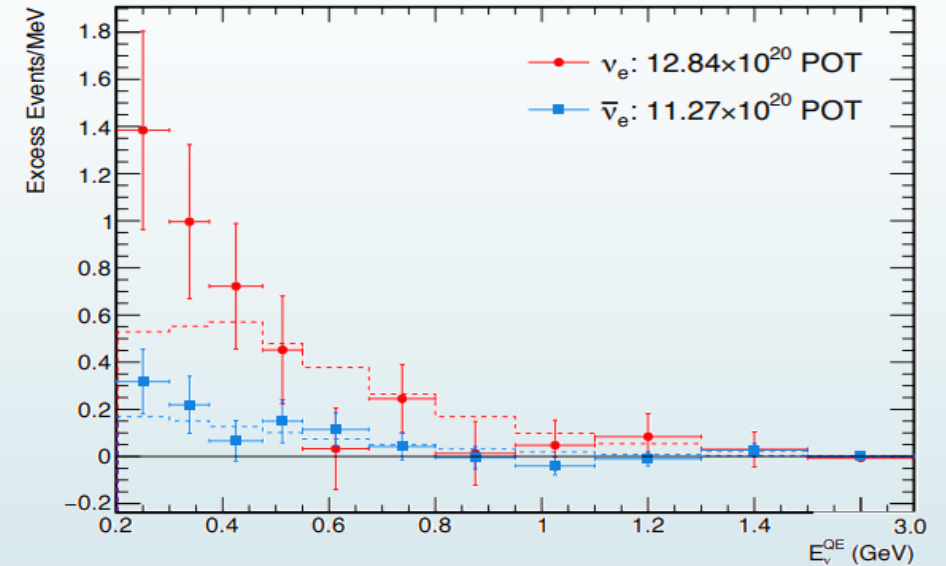


MiniBooNE's Low Energy Excess

Neutrino mode only



Both excesses, BG subtracted

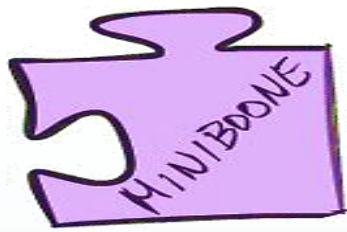


$$E_{\nu}^{(\text{reconst.})} = \frac{2m_n E_e + m_p^2 - m_n^2 - m_e^2}{2(m_n - E_e + \cos \theta_e \sqrt{E_e^2 - m_e^2})}$$

Measure charged lepton energy/angle

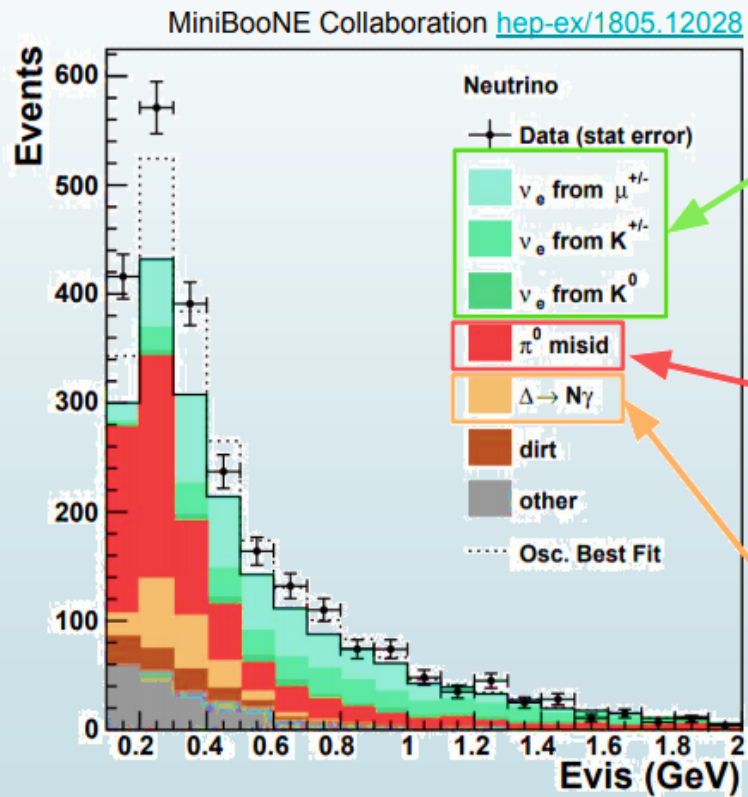
Observed ~ 400 events, PMNS predicts 0

Combined $\nu/\bar{\nu}$ modes : 4.8σ excess



MiniBooNE's Low Energy Excess

Possible Explanations: Motivated by backgrounds



Intrinsic ν_e in the beam? Constrained by measuring ν_μ which come from the same π decay as the μ 's that subsequently produce the ν_e .

π^0 misidentification? In which the second shower was missed or incorrectly reconstructed. MiniBooNE measured the largest sample of NC π^0 events ever collected and used this to constrain the exact rate of π^0 's for the CCQE analysis.

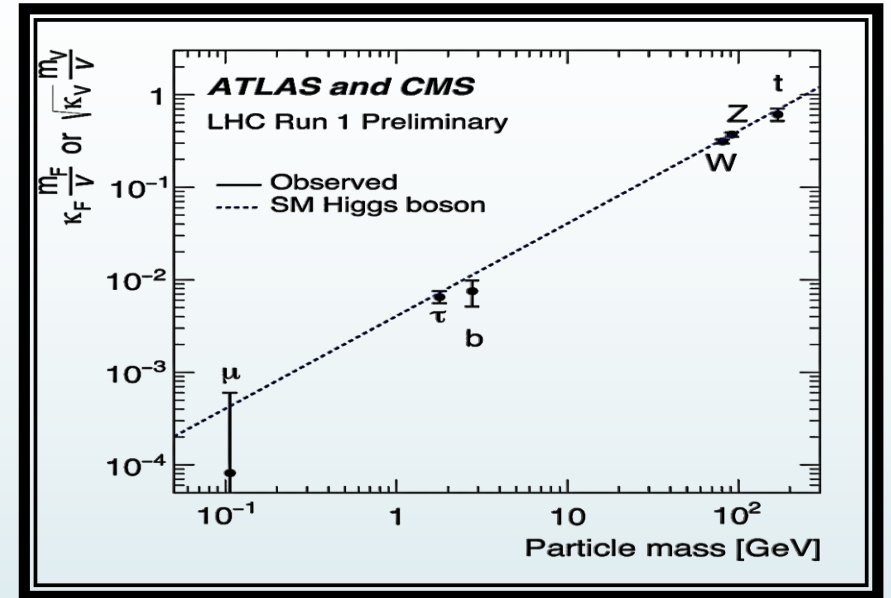
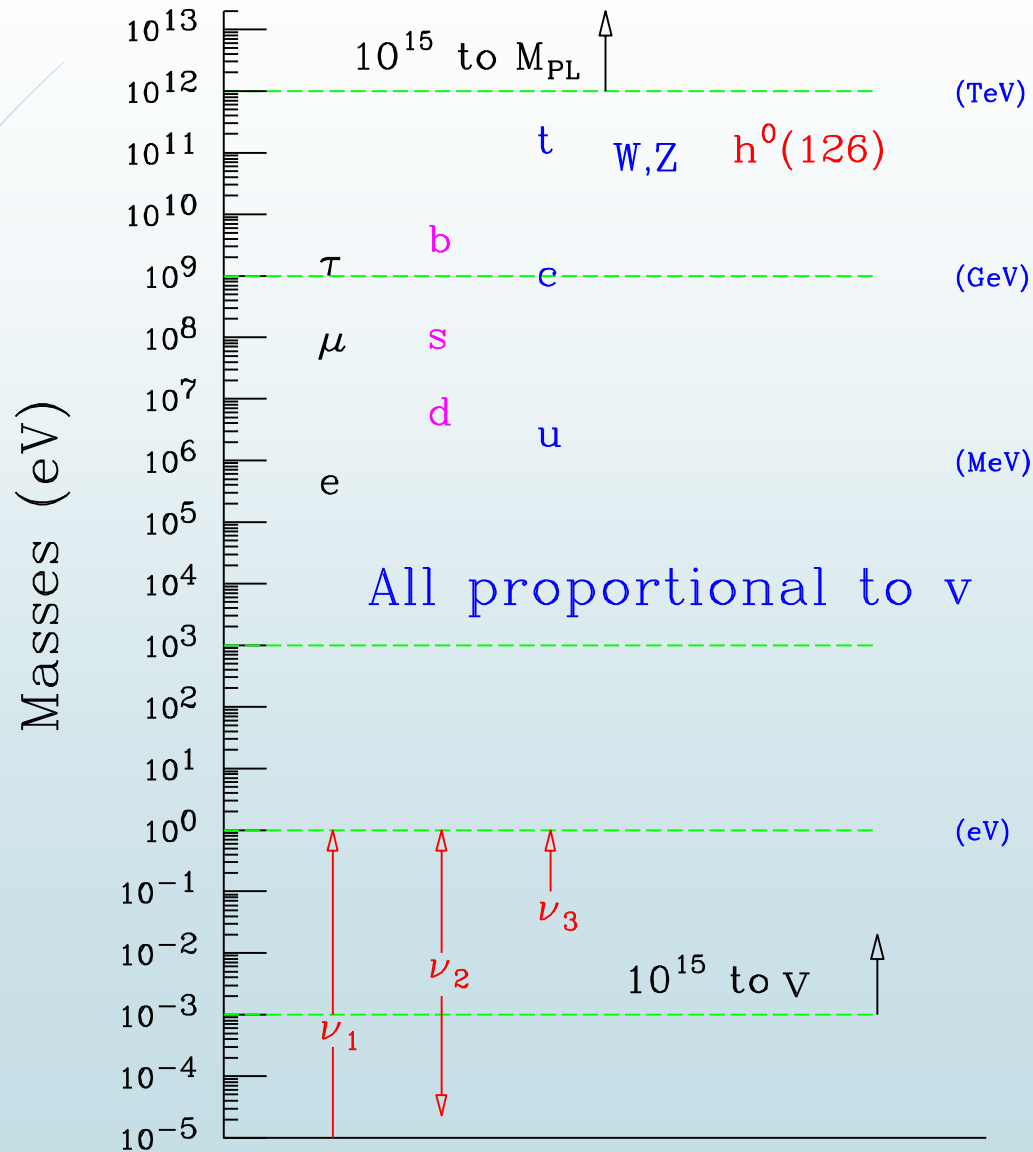
Radiative Δ decay? This has never been observed in the neutrino sector. MiniBooNE bound it using their NC π^0 measurements which agrees well with best theoretical calculations. The biggest channel of interest to MiniBooNE's photon LEE analysis.

❖ *Explanation of MiniBooNE's low energy excess*

A LIGHT DARK SECTOR – THE PRESCRIPTION

- How **low-energy** does the subleading electron have to be in an e^+e^- pair in order for an “**Asymmetric**” pair to look like a single ring? $E_{\text{True}} < 30 \text{ MeV}$
- How **small an opening angle** does the e^+e^- pair have to have before it is “**Overlapping**” sufficiently to look like a single ring? $\theta_{\text{SEP}} < 5^\circ$
- When forcing a two-ring fit to an event, the associated invariant mass should be sufficiently **non- π^0** like: $m_{\gamma\gamma} < 80 \text{ MeV}$

Neutrino Mass \longrightarrow New physics beyond SM:



$m_\nu \sim y_\nu^{eff} v$

$y_\nu^{eff} < 10^{-12}$

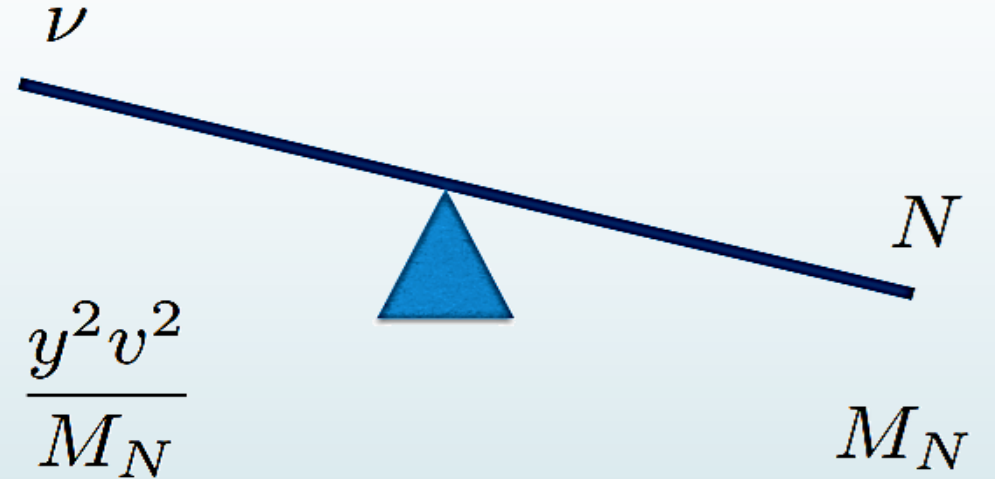
❖ Standard/Type I Seesaw

56

$$yN H \ell + M_N N N$$

$$\frac{y^2 \ell H \ell H}{M_N}$$

$$m_\nu \sim \frac{y^2 v^2}{M_N}$$



$$m_\nu \sim 0.1 \text{eV} \quad y \sim 0.1 \quad M_N \sim 10^{12} \text{GeV}$$

Lepton number is broken at very high scale M_N

❖ Phenomenology on other neutrino experiment

U(1)' models in Future and Current LArTPCs

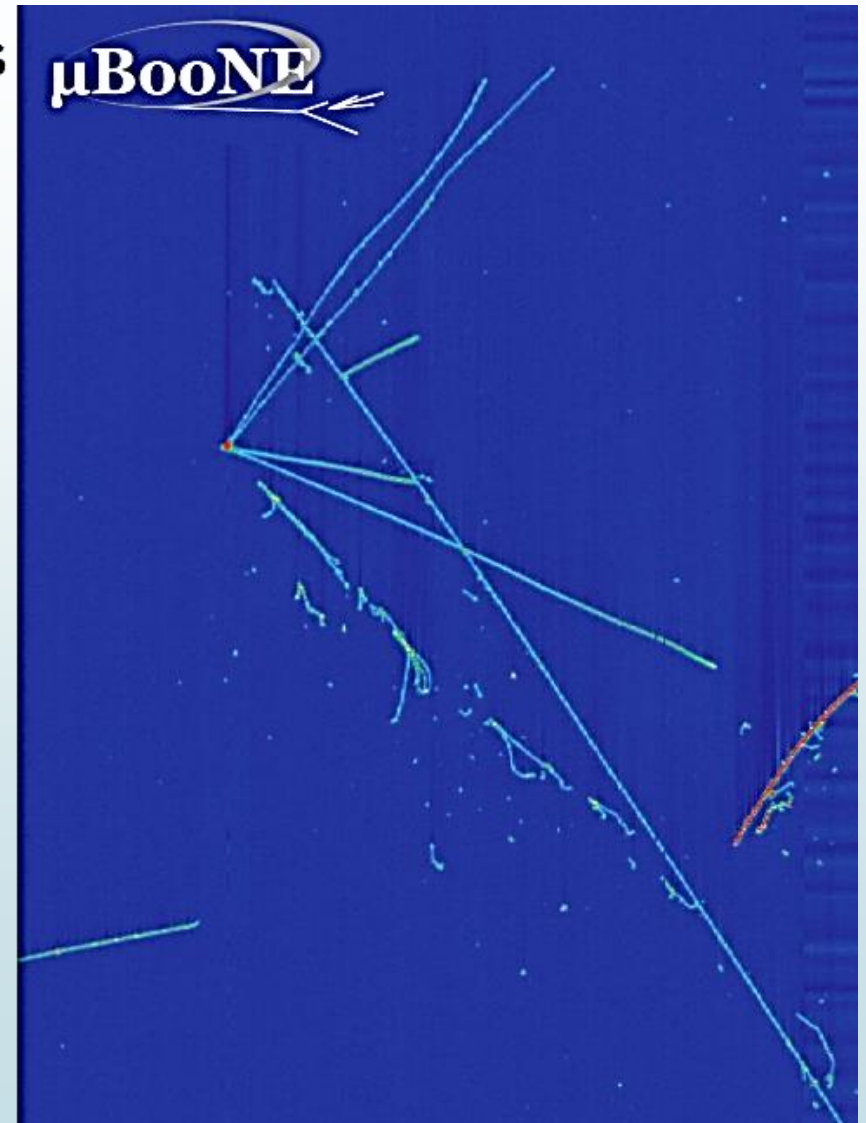
This class of models has has incredibly **rich phenomenology** at LArTPCs such as **MicroBooNE, SBND or the DUNE near detector**:

LArTPCs have the distinct advantage that one can tell photons and electron showers apart via two methods:

- Directly look for the **conversion gap**

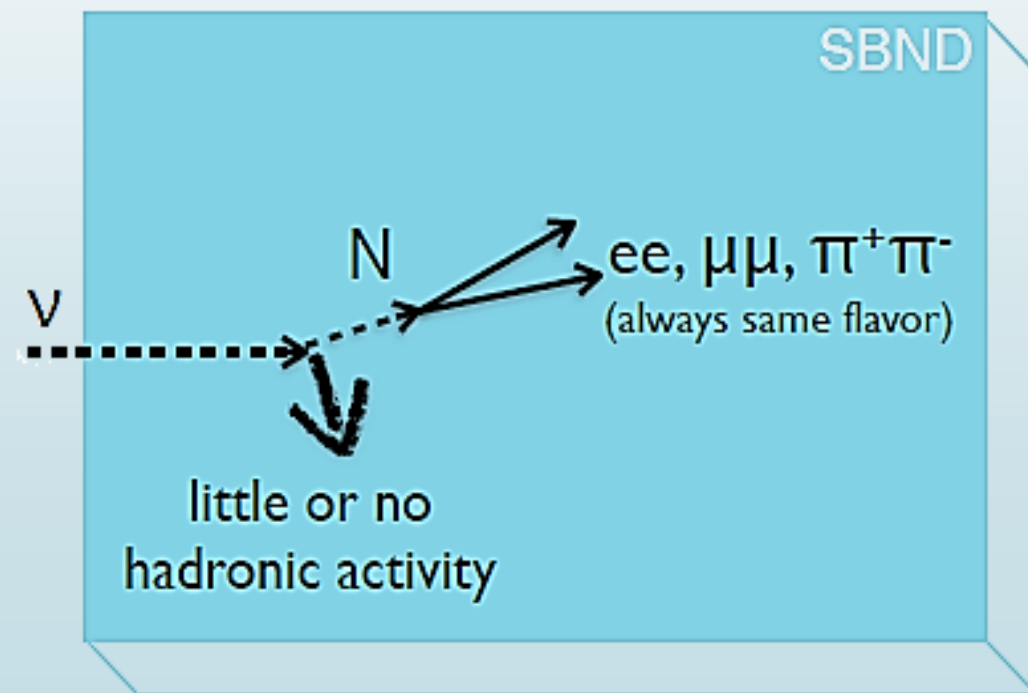
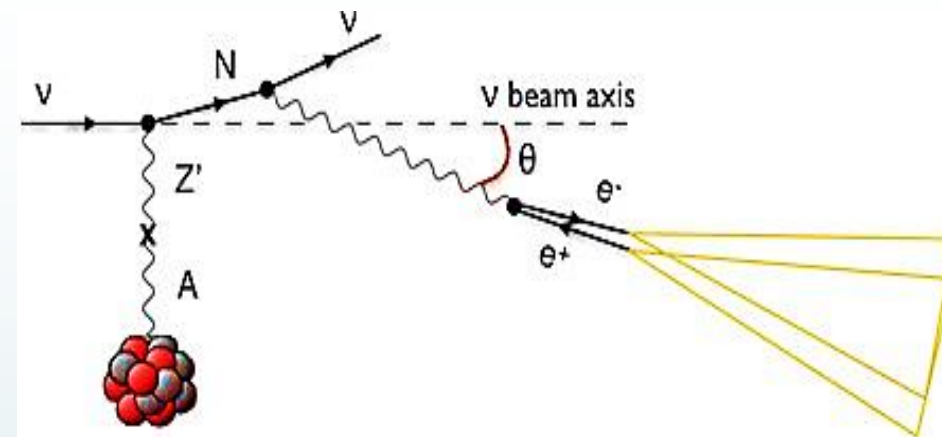


- Use **Calorimetric measurements** to see rate of energy deposition (dE/dx). **Photons** that pair convert to e^+e^- **deposit x2 as much energy**.



❖ What happens at the SBN program?

- ✓ **No baseline dependence**
- ✓ **Almost no hadronic activity to tag interaction vertex**
- ✓ Decays to collimated e^+e^- pairs
- ✓ **More events due to coherence:**
- ✓ ${}^6\text{C}$ vs ${}_{18}\text{Ar}$ ~ 3 times more events for same exposure
- ✓ **Hard to probe !!!**



Severe Constraints on New Physics Explanations of the MiniBooNE Excess

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²*Princeton University, Princeton, NJ*

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angle is weighted by the track energies; by momentum conservation, this sum is simply the original X 4-vector, which must satisfy $\cos \theta_e > 0.9999$ in order for X to enter the MiniBooNE detector, a sphere of fiducial radius 5.75 m located 541 m away from the target. This is highly inconsistent with the $\cos \theta_e$ distribution of the excess (see Fig. 2), which shows significant contributions from $\cos \theta_e < 0.8$. In particular, a model which matches the size of the neutrino mode excess (381.2 events), but predicts all events to have $\cos \theta_e > 0.8$ is incompatible with the observed excess of 150 ± 31 in this bin (in consideration of statistical errors only; systematics and bin-to-bin correlations are not available, noting that the angular resolution is 3-5° for 100-600 MeV electron energies in ν_e CCQE events [28]).

body decay where X decays into a lighter dark- X' and a photon ($X \rightarrow X' + \gamma$). Three- and higher-body decays are also allowed but will be increasingly suppressed; regardless, we consider decays to X' with an arbitrary number of electromagnetic tracks. Since electromagnetic tracks must be well-collimated to contribute to the excess, we will treat this scenario as a quasi-two-body decay where the electromagnetic energy is considered as a vector p_{EM} with $0 \leq p_{EM}^2 \leq (30 \text{ MeV})^2$.

In the X rest frame, the electromagnetic energy is $(m_X^2 - m_{X'}^2)/2m_X$. Electromagnetic energy with a small invariant mass compared to the beam energy, emitted in the X rest frame, will be boosted to very small energies,