Reducing Uncertainties in $\bar{B} \rightarrow X_s \gamma$ Decay

AYESH GUNAWARDANA

Wayne State University

May , 2019

Based on A.G and Gil Paz "Reducing Uncertainties in $\bar{B} \rightarrow X_s \gamma$ " (To appear)

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $\bar{B} \rightarrow X_s \gamma$ Decay

Introduction

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $\bar{B} \to X_s \gamma$ Decay

• $\bar{B} \rightarrow X_s \gamma$ decay is an important **New Physics** probe

- $\bar{B} \rightarrow X_s \gamma$ decay is an important **New Physics** probe
 - It is suppressed at tree level in SM
 - Can receive contributions from SM extensions.



Figure: $b \rightarrow s\gamma$ flavor changing neural current (FCNC) in SM

AYESH GUNAWARDANA

- $\bar{B} \rightarrow X_s \gamma$ decay is an important **New Physics** probe
 - It is suppressed at tree level in SM
 - Can receive contributions from SM extensions.



Figure: $b \rightarrow s\gamma$ flavor changing neural current (FCNC) in SM

- SM extensions modify the $C_{7\gamma}$ Wilson coefficient

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $\bar{B} \rightarrow X_s \gamma$ Decay

- $\bar{B} \rightarrow X_s \gamma$ decay is an important **New Physics** probe
 - It is suppressed at tree level in SM
 - Can receive contributions from SM extensions.



Figure: $b \rightarrow s\gamma$ flavor changing neural current (FCNC) in SM

- SM extensions modify the $C_{7\gamma}$ Wilson coefficient
- CP violation in $ar{B}
 ightarrow X_s \gamma$ can be enhanced by new physics

AYESH GUNAWARDANA

Wayne State University

Photon production

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $\bar{B} \to X_{\rm s} \gamma$ Decay

Photon production

• Photon can be produced directly:

$$Q_{7\gamma}=rac{-e}{8\pi^2}m_bar{s}\sigma_{\mu
u}F^{\mu
u}(1+\gamma_5)b$$

Photon production

• Photon can be produced directly:

$$Q_{7\gamma} = \frac{-e}{8\pi^2} m_b \bar{s} \sigma_{\mu\nu} F^{\mu\nu} (1+\gamma_5) b$$

• Also, gluon or quark pair can convert to photon

$$Q_{8g}=\frac{-e}{8\pi^2}m_b\bar{s}\sigma_{\mu\nu}G^{\mu\nu}(1+\gamma_5)b$$

$$Q_1^q = (\bar{q}b)_{V-A}(\bar{s}q)_{V-A}$$

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $\bar{B} \rightarrow X_s \gamma$ Decay

• The effective Lagrangian to describe $ar{B} o X_{s} \gamma$

• The effective Lagrangian to describe $ar{B} o X_{s} \gamma$

$$\begin{split} \mathcal{H}_{\text{eff}} &= \; \frac{G_F}{\sqrt{2}} \sum_{q=u,c} \lambda_q \\ & (C_1 Q_1^q + C_2 Q_2^q + \sum_i C_i Q_i + C_{7\gamma} Q_{7\gamma} + C_{8g} Q_{8g}) + \text{h.c.} \end{split}$$

• The effective Lagrangian to describe $ar{B} o X_{s} \gamma$

$$\begin{aligned} \mathcal{H}_{\text{eff}} &= \; \frac{G_F}{\sqrt{2}} \sum_{q=u,c} \lambda_q \\ & \left(C_1 Q_1^q + C_2 Q_2^q + \sum_i C_i Q_i + C_{7\gamma} Q_{7\gamma} + C_{8g} Q_{8g} \right) + \text{h.c.} \end{aligned}$$

-
$$\lambda_q = V_{qb}^* V_{qs}$$

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $ar{B}
ightarrow X_s \gamma$ Decay

• The effective Lagrangian to describe $ar{B}
ightarrow X_{s} \gamma$

$$\begin{split} \mathcal{H}_{\text{eff}} &= \; \frac{G_F}{\sqrt{2}} \sum_{q=u,c} \lambda_q \\ & \left(C_1 Q_1^q + C_2 Q_2^q + \sum_i C_i Q_i + C_{7\gamma} Q_{7\gamma} + C_{8g} Q_{8g} \right) + \text{h.c.} \end{split}$$

- $\lambda_q = V^*_{qb} V_{qs}$

- At leading power: $Q_{7\gamma} Q_{7\gamma}$ contribution only
- At higher orders: $Q_i Q_j$ contributions
- Most important operators are $Q_{7\gamma}$, Q_{8g} and Q_1^q .

AYESH GUNAWARDANA

Decay rate

• World average for experimental value:

 $\mathcal{B}\left(B
ightarrow X_{s}\gamma
ight)\left(E_{\gamma}>1.6~{
m GeV}
ight)=\left(3.32\pm0.15
ight) imes10^{-4}$

[Y. Amhis et. al. EPJC 77, 895 (2017)]

NNLO prediction

$$\Gamma\left(\overline{B} \to X_q \gamma\right) = \underbrace{\Gamma\left(b \to X_q^p \gamma\right)}_{\text{Perturbatively calculable}} + \underbrace{\delta\Gamma_{\text{nonp}}}_{\mathcal{O}(\frac{\Lambda_{\text{QCD}}}{m_b})}$$

• SM prediction (2015) [Misiak et. al. PRL 114, 221801 (2015)]

$$\mathcal{B}^{
m SM}_{s\gamma}=(3.36\pm0.23) imes10^{-4}$$

for $E_{\gamma} > 1.6 \text{ GeV}$

AYESH GUNAWARDANA

Decay rate

• World average for experimental value:

 $\mathcal{B}\left(B
ightarrow X_{s}\gamma
ight)\left(E_{\gamma}>1.6~{
m GeV}
ight)=\left(3.32\pm0.15
ight) imes10^{-4}$

[Y. Amhis et. al. EPJC 77, 895 (2017)]

NNLO prediction

$$\Gamma\left(\overline{B} \to X_q \gamma\right) = \underbrace{\Gamma\left(b \to X_q^p \gamma\right)}_{\text{Perturbatively calculable}} + \underbrace{\delta\Gamma_{\text{nonp}}}_{\mathcal{O}(\frac{\Lambda_{\text{QCD}}}{m_b})}$$

• SM prediction (2015) [Misiak et. al. PRL 114, 221801 (2015)]

$${\cal B}_{s\gamma}^{
m SM} = (3.36\pm 0.23) imes 10^{-4}$$

for $E_{\gamma} > 1.6 \text{ GeV}$

- $\delta\Gamma_{\rm nonp}\equiv$ Non-perturbative contribution
 - The largest contribution to the error 5% from $\mathcal{O}(\frac{\Lambda_{QCD}}{m_b})$

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $\bar{B} \rightarrow X_s \gamma$ Decay

Order $1/m_b$ power corrections to $\Gamma(\bar{B} \to X_s \gamma)$

Non-perturbative effects arise from Resolved Photon Contributions

 $\Delta\Gamma\sim$ $\underbrace{\bar{J}}_{\text{Perturbatively calculable}} \otimes \underbrace{h}_{\text{Non perturbative}}$

Order $1/m_b$ power corrections to $\Gamma(\bar{B} \to X_s \gamma)$

Non-perturbative effects arise from Resolved Photon Contributions



- 2010 estimates for non-perturbative contribution
 - From $Q_1^c Q_{7\gamma} \in [-1.7, +4.0]\%$
 - From $Q_{8g}-Q_{8g}\in [-0.3,+1.9]\%$
 - From $\mathit{Q}_{7\gamma} \mathit{Q}_{8g} \in [-4.4, +5.6]\%$

- 2010 estimates for non-perturbative contribution
 - From $Q_1^c Q_{7\gamma} \in [-1.7, +4.0]\%$
 - From $Q_{8g} Q_{8g} \in [-0.3, +1.9]\%$
 - From $\mathit{Q}_{7\gamma} \mathit{Q}_{8g} \in [-4.4, +5.6]\%$
- The contribution from $Q_{7\gamma}-Q_{8g}$
 - Obtained on experiment with 95% confidence level range

[M. Benzke, S. J. Lee, M. Neubert and G. Paz JHEP 1008, 099(2010)]

- 2010 estimates for non-perturbative contribution
 - From $Q_1^c Q_{7\gamma} \in [-1.7, +4.0]\%$
 - From $Q_{8g} Q_{8g} \in [-0.3, +1.9]\%$
 - From $\mathit{Q}_{7\gamma} \mathit{Q}_{8g} \in [-4.4, +5.6]\%$
- The contribution from $Q_{7\gamma}-Q_{8g}$
 - Obtained on experiment with 95% confidence level range

[M. Benzke, S. J. Lee, M. Neubert and G. Paz JHEP 1008, 099(2010)]

- New Belle result for $Q_{7\gamma}-Q_{8g}$ contribution $\sim 2\%$
- [S. Watanuki et. al. PRD 99, 032012(2019)]

- 2010 estimates for non-perturbative contribution
 - From $Q_1^c Q_{7\gamma} \in [-1.7, +4.0]\%$
 - From $Q_{8g} Q_{8g} \in [-0.3, +1.9]\%$
 - From $\mathit{Q}_{7\gamma} \mathit{Q}_{8g} \in [-4.4, +5.6]\%$
- The contribution from $Q_{7\gamma}-Q_{8g}$
 - Obtained on experiment with 95% confidence level range

[M. Benzke, S. J. Lee, M. Neubert and G. Paz JHEP 1008, 099(2010)]

- New Belle result for $Q_{7\gamma}-Q_{8g}$ contribution $\sim 2\%$
- [S. Watanuki et. al. PRD 99, 032012(2019)]
- Now $Q_1^c Q_{7\gamma}$ is the largest contribution to the error! Can we reduce it?

AYESH GUNAWARDANA

$Q_1^c - Q_{7\gamma}$ contribution

• The contribution to the error from $Q_1^c - Q_{7\gamma}$ is given by

 $\frac{C_1}{C_{7\gamma}}\frac{\Lambda_{17}}{m_b}$

$Q_1^c - Q_{7\gamma}$ contribution

- The contribution to the error from $Q_1^c-Q_{7\gamma}$ is given by

$$\frac{C_1}{C_{7\gamma}}\frac{\Lambda_{17}}{m_b}$$

where

$$\Lambda_{17} = e_c \operatorname{Re} \int_{-\infty}^{\infty} \frac{d\omega_1}{\omega_1} \left[1 - \underbrace{\mathcal{F}\left(\frac{m_c^2 - i\varepsilon}{m_b\omega_1}\right)}_{\text{perturbative}} + \frac{m_b\omega_1}{12m_c^2} \right] \underbrace{h_{17}\left(\omega_1\right)}_{\text{non-perturbative}}$$

$${\it Q}_1^{\it c}-{\it Q}_{7\gamma}$$
 contribution

- The contribution to the error from $Q_1^c-Q_{7\gamma}$ is given by

$$\frac{C_1}{C_{7\gamma}}\frac{\Lambda_{17}}{m_b}$$

where

$$\Lambda_{17} = e_c \operatorname{Re} \int_{-\infty}^{\infty} \frac{d\omega_1}{\omega_1} \left[1 - \underbrace{\mathcal{F}\left(\frac{m_c^2 - i\varepsilon}{m_b\omega_1}\right)}_{\text{perturbative}} + \frac{m_b\omega_1}{12m_c^2} \right] \underbrace{h_{17}\left(\omega_1\right)}_{\text{non-perturbative}}$$

- Need a new model for h_{17} to reduce the error
 - Using new information on moments of $h_{17} \Rightarrow$ New model

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $ar{B}
ightarrow X_s \gamma$ Decay

• h_{17} can be thought of as a gluon PDF of a B meson

- h_{17} can be thought of as a gluon PDF of a B meson
 - Non-local operator matrix element
 - Describe the hadronic effects of the process

- h_{17} can be thought of as a gluon PDF of a B meson
 - Non-local operator matrix element
 - Describe the hadronic effects of the process

$$h_{17}(\omega_1) = \\ = \int \frac{dr}{2\pi} e^{-i\omega_1 r} \frac{\langle \bar{B} | (\bar{h}S_{\bar{n}})(0) \not n(1+\gamma_5) i \gamma^{\perp} \bar{n}_{\beta} (S_{\bar{n}} g G^{\alpha\beta} S_{\bar{n}}) (r\bar{n}) (S_{\bar{n}}^{\dagger} h)(0) | \bar{B} \rangle}{2M_B}$$

- h_{17} can be thought of as a gluon PDF of a B meson
 - Non-local operator matrix element
 - Describe the hadronic effects of the process

$$h_{17}(\omega_1) = \\ = \int \frac{dr}{2\pi} e^{-i\omega_1 r} \frac{\langle \bar{B} | (\bar{h}S_{\bar{n}})(0) \not n(1+\gamma_5) i \gamma^{\perp} \bar{n}_{\beta}(S_{\bar{n}}gG^{\alpha\beta}S_{\bar{n}})(r\bar{n})(S_{\bar{n}}^{\dagger}h)(0) | \bar{B} \rangle}{2M_B}$$

$$S_n(x) = \mathbf{P} \exp\left(ig \int_{-\infty}^0 dun \cdot A_s(x+un)\right)$$

 $n^{\mu}\equiv(1,0,0,1)$ and $\overline{n}^{\mu}\equiv(1,0,0,-1)$

AYESH GUNAWARDANA

-

_

Wayne State University

Reducing Uncertainties in $ar{B}
ightarrow X_s \gamma$ Decay

• k th moment of h_{17} ; Obtained using $\frac{\partial^k}{\partial r^k}e^{-i\omega_1r}$

• k th moment of h_{17} ; Obtained using $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$

$$(-1)^{k}\frac{1}{2M_{B}}\left\langle \overline{B}\left|\left(\bar{h}S_{\bar{n}}\right)\left(0\right)\vec{p}\left(1+\gamma_{5}\right)i\gamma_{\alpha}^{\perp}\bar{n}_{\beta}\left(i\bar{n}\cdot\partial\right)^{k}\left(S_{\bar{n}}^{\dagger}gG_{s}^{\alpha\beta}S_{\bar{n}}\right)\left(r\bar{n}\right)\left(S_{\bar{n}}^{\dagger}h\right)\left(0\right)\left|\bar{B}\right\rangle\right|_{r=0}$$

• k th moment of h_{17} ; Obtained using $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$

$$(-1)^{k}\frac{1}{2M_{B}}\left\langle \overline{B}\left|\left(\bar{h}S_{\bar{n}}\right)\left(0\right)\vec{p}\left(1+\gamma_{5}\right)i\gamma_{\alpha}^{\perp}\bar{n}_{\beta}(i\bar{n}\cdot\partial)^{k}\left(S_{\bar{n}}^{\dagger}gG_{s}^{\alpha\beta}S_{\bar{n}}\right)(r\bar{n})\left(S_{\bar{n}}^{\dagger}h\right)(0)\right|\bar{B}\right\rangle\right|_{r=0}$$

• Using the (new) identity

$$i\bar{n}\cdot\partial\left(S_{\bar{n}}^{\dagger}(x)O(x)S_{\bar{n}}(x)\right)=S_{\bar{n}}^{\dagger}(x)[i\bar{n}\cdot D,O(x)]S_{\bar{n}}(x)$$

• k th moment of h_{17} ; Obtained using $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$

$$(-1)^{k}\frac{1}{2M_{B}}\left\langle \overline{B}\left|\left(\bar{h}S_{\bar{n}}\right)\left(0\right)\vec{p}\left(1+\gamma_{5}\right)i\gamma_{\alpha}^{\perp}\bar{n}_{\beta}(i\bar{n}\cdot\partial)^{k}\left(S_{\bar{n}}^{\dagger}gG_{s}^{\alpha\beta}S_{\bar{n}}\right)(r\bar{n})\left(S_{\bar{n}}^{\dagger}h\right)(0)\right|\bar{B}\right\rangle\right|_{r=0}$$

• Using the (new) identity

$$i\bar{n}\cdot\partial\left(S_{\bar{n}}^{\dagger}(x)O(x)S_{\bar{n}}(x)\right)=S_{\bar{n}}^{\dagger}(x)[i\bar{n}\cdot D,O(x)]S_{\bar{n}}(x)$$

• k th moment of h_{17} ; Obtained using $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$

$$(-1)^{k}\frac{1}{2M_{B}}\left\langle \overline{B}\left|\left(\bar{h}S_{\bar{n}}\right)\left(0\right)\vec{p}\left(1+\gamma_{5}\right)i\gamma_{\alpha}^{\perp}\bar{n}_{\beta}(i\bar{n}\cdot\partial)^{k}\left(S_{\bar{n}}^{\dagger}gG_{s}^{\alpha\beta}S_{\bar{n}}\right)(r\bar{n})\left(S_{\bar{n}}^{\dagger}h\right)(0)\right|\bar{B}\right\rangle\right|_{r=0}$$

• Using the (new) identity

$$i\bar{n}\cdot\partial\left(S_{\bar{n}}^{\dagger}(x)O(x)S_{\bar{n}}(x)\right)=S_{\bar{n}}^{\dagger}(x)[i\bar{n}\cdot D,O(x)]S_{\bar{n}}(x)$$

- Apply this for k derivatives \Rightarrow k commutators of $i\bar{n} \cdot D$

-
$$[iD^{\mu}, iD^{\nu}] = igG^{\mu\nu}$$

AYESH GUNAWARDANA

• k th moment of h_{17} ; Obtained using $\frac{\partial^k}{\partial r^k}e^{-i\omega_1 r}$

$$(-1)^{k}\frac{1}{2M_{B}}\left\langle \overline{B}\left|\left(\bar{h}S_{\bar{n}}\right)\left(0\right)\vec{p}\left(1+\gamma_{5}\right)i\gamma_{\alpha}^{\perp}\bar{n}_{\beta}(i\bar{n}\cdot\partial)^{k}\left(S_{\bar{n}}^{\dagger}gG_{s}^{\alpha\beta}S_{\bar{n}}\right)(r\bar{n})\left(S_{\bar{n}}^{\dagger}h\right)(0)\right|\bar{B}\right\rangle\right|_{r=0}$$

Using the (new) identity

$$i\bar{n}\cdot\partial\left(S_{\bar{n}}^{\dagger}(x)O(x)S_{\bar{n}}(x)\right)=S_{\bar{n}}^{\dagger}(x)[i\bar{n}\cdot D,O(x)]S_{\bar{n}}(x)$$

- Apply this for k derivatives $\Rightarrow k$ commutators of $i\bar{n} \cdot D$
- $[iD^{\mu}, iD^{\nu}] = igG^{\mu\nu}$
- New result Moments over ω₁

$$\langle \omega_1^k h_{17} \rangle = (-1)^k \frac{1}{2M_B} \langle \bar{B} | \bar{h} \bar{n} (1+\gamma_5) \gamma_{\alpha}^{\perp} \underbrace{[i\bar{n} \cdot D, [i\bar{n} \cdot D, \cdots [i\bar{n} \cdot D]}_{k \text{ times}} [D^{\alpha}, i\bar{n} \cdot D] \cdots]] s^{\lambda} h | \bar{B} \rangle$$

AYESH GUNAWARDANA

Moments of the g_{17}

 Procedure to obtain these HQET matrix elements derived in [A. Gunawardana and G. Paz, JHEP 07(2017)137 [arXiv:1702.08904]]

$$\langle h_{17}
angle = 2\lambda_2 = 2\mu_G^2/3$$

 $\langle \omega_1^2 h_{17}
angle = \frac{2}{15} (5m_5 + 3m_6 - 2m_9)$ New result

m_i were extracted from data for the first time in 2016 [P. Gambino, K. J Healey, S. Turczyk PLB 763, 60 (2016)]

$$\mu_G^2 = 0.355 \pm 0.060 \text{ GeV}^2$$
 $m_5 = 0.072 \pm 0.045 \text{ GeV}^4$
 $m_6 = 0.060 \pm 0.164 \text{ GeV}^4$ $m_9 = -0.280 \pm 0.352 \text{ GeV}^4$

AYESH GUNAWARDANA

Reducing Uncertainties in $\bar{B} \rightarrow X_s \gamma$ Decay

• Relative errors are large:

• Relative errors are large: Numerical error is 17% for $\langle \omega_1^0 h_{17} \rangle$ Numerical error is 80% for $\langle \omega_1^2 h_{17} \rangle$

- Relative errors are large: Numerical error is 17% for $\langle \omega_1^0 h_{17} \rangle$ Numerical error is 80% for $\langle \omega_1^2 h_{17} \rangle$
- These moments still give useful information

- Relative errors are large: Numerical error is 17% for $\langle \omega_1^0 h_{17} \rangle$ Numerical error is 80% for $\langle \omega_1^2 h_{17} \rangle$
- These moments still give useful information
 - 2019 estimate $\left<\omega_1^2h_{17}
 ight>\in$ (0.03, 0.27) GeV 4
 - 2010 models provide $\langle \omega_1^2 h_{17} \rangle \in (-0.31, 0.49) \text{ GeV}^4$.
 - These older models were constructed before m_i were extracted
 - New estimate is significantly smaller than old estimate.

- Relative errors are large: Numerical error is 17% for $\langle \omega_1^0 h_{17} \rangle$ Numerical error is 80% for $\langle \omega_1^2 h_{17} \rangle$
- These moments still give useful information
 - 2019 estimate $\left<\omega_1^2h_{17}
 ight>\in$ (0.03, 0.27) GeV 4
 - 2010 models provide $\langle \omega_1^2 h_{17} \rangle \in (-0.31, 0.49) \text{ GeV}^4$.
 - These older models were constructed before m_i were extracted
 - New estimate is significantly smaller than old estimate.
- Expect in future
 - Further improvements on HQET matrix elements
 - Belle II or LQCD data \Rightarrow Better constrains on moments

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $ar{B}
ightarrow X_s \gamma$ Decay

Applications

• Properties of h_{17}

- Properties of h_{17}
 - Real and even function over ω_1
 - $\langle \omega_1^k h_{17}(\omega_1)
 angle = 0$ for $k = 1, 3, 5, \cdots$
 - h_{17} has a dimension of mass
 - Range of $\omega_1 \Rightarrow -\infty < \omega_1 < \infty$

- Properties of h_{17}
 - Real and even function over ω_1
 - $\langle \omega_1^k h_{17}(\omega_1)
 angle = 0$ for $k=1,3,5,\cdots$
 - h_{17} has a dimension of mass
 - Range of $\omega_1 \Rightarrow -\infty < \omega_1 < \infty$
- We use Hermite polynomials $H_n(x)$

- Properties of h_{17}
 - Real and even function over ω_1
 - $\langle \omega_1^k h_{17}(\omega_1)
 angle = 0$ for $k=1,3,5,\cdots$
 - h_{17} has a dimension of mass
 - Range of $\omega_1 \Rightarrow -\infty < \omega_1 < \infty$
- We use Hermite polynomials $H_n(x)$

• Our model:
$$\begin{array}{c} h_{17}(\omega_1) = \sum_n a_{2n} H_{2n}(\frac{\omega_1}{\sqrt{2}\sigma}) e^{\frac{-\omega_1^2}{2\sigma}} \\ & \text{- where} \\ a_0 = \frac{\langle \omega_1^0 h_{17} \rangle}{\sqrt{2\pi} |\sigma|}, \quad a_2 = \frac{\langle \omega_1^2 h_{17} \rangle - \sigma^2 \langle \omega_1^0 h_{17} \rangle}{4\sqrt{2\pi} |\sigma|^3}, \quad a_4 = \cdots \end{array}$$

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $ar{B}
ightarrow X_s \gamma$ Decay

- Properties of h_{17}
 - Real and even function over ω_1
 - $\langle \omega_1^k h_{17}(\omega_1)
 angle = 0$ for $k=1,3,5,\cdots$
 - h_{17} has a dimension of mass
 - Range of $\omega_1 \Rightarrow -\infty < \omega_1 < \infty$
- We use Hermite polynomials $H_n(x)$

• Our model:
$$\begin{array}{c} h_{17}(\omega_1) = \sum_n a_{2n} H_{2n}(\frac{\omega_1}{\sqrt{2}\sigma}) e^{\frac{-\omega_1^2}{2\sigma}} \\ & \text{- where} \\ a_0 = \frac{\langle \omega_1^0 h_{17} \rangle}{\sqrt{2\pi} |\sigma|}, \quad a_2 = \frac{\langle \omega_1^2 h_{17} \rangle - \sigma^2 \langle \omega_1^0 h_{17} \rangle}{4\sqrt{2\pi} |\sigma|^3}, \quad a_4 = \cdots \end{array}$$

• $|h_{17}| < 1$ GeV and no peaks beyond $\omega_1 = 1$ GeV AYESH GUNAWARDANA Wayne State University Reducing Uncertainties in $\bar{B} \rightarrow X_{s\gamma}$ Decay



Figure: 2019 model vs 2010 model for h_{17}

AYESH GUNAWARDANA

Wayne State University

Reducing Uncertainties in $\bar{B}
ightarrow X_s \gamma$ Decay



Figure: 2019 model vs 2010 model for h_{17}

• Orange dashed line: 2010 model $h_{17}(\omega_1,\mu) = \frac{2\lambda_2}{\sqrt{2\pi\sigma}} \frac{\omega_1^2 - \Lambda^2}{\sigma^2 - \Lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$



Figure: 2019 model vs 2010 model for h_{17}

• Orange dashed line: 2010 model $h_{17}(\omega_1,\mu) = \frac{2\lambda_2}{\sqrt{2\pi\sigma}} \frac{\omega_1^2 - \lambda^2}{\sigma^2 - \lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$

- $\sigma = 0.5 \text{ GeV}, \Lambda = 0.425 \text{ GeV}$ and $\Rightarrow \langle \omega_1^2 h_{17} \rangle = 0.49 \text{ GeV}^4$

AYESH GUNAWARDANA



Figure: 2019 model vs 2010 model for h_{17}

• Orange dashed line: 2010 model $h_{17}(\omega_1,\mu) = \frac{2\lambda_2}{\sqrt{2\pi\sigma}} \frac{\omega_1^2 - \lambda^2}{\sigma^2 - \lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$

- $\sigma = 0.5 \,\, {
m GeV}, \Lambda = 0.425 \,\, {
m GeV}$ and $\Rightarrow \langle \omega_1^2 h_{17}
angle = 0.49 \,\, {
m GeV}^4$

• Blue line; 2019 model: $\sigma=0.5~{
m GeV}$ and $\langle \omega_1^2 h_{17}
angle=0.27~{
m GeV}^4$

AYESH GUNAWARDANA

Wayne State University



Figure: 2019 model vs 2010 model for h_{17}

• Orange dashed line: 2010 model $h_{17}(\omega_1,\mu) = \frac{2\lambda_2}{\sqrt{2\pi\sigma}} \frac{\omega_1^2 - \lambda^2}{\sigma^2 - \lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$

- $\sigma = 0.5 \,\, {
m GeV}, \Lambda = 0.425 \,\, {
m GeV}$ and $\Rightarrow \langle \omega_1^2 h_{17}
angle = 0.49 \,\, {
m GeV}^4$

- Blue line; 2019 model: $\sigma=0.5~{
 m GeV}$ and $\langle\omega_1^2h_{17}
 angle=0.27~{
 m GeV}^4$
- New function is 50% smaller than the 2010

AYESH GUNAWARDANA



Figure: 2019 model vs 2010 model for h_{17}

• Orange dashed line: 2010 model $h_{17}(\omega_1,\mu) = \frac{2\lambda_2}{\sqrt{2\pi\sigma}} \frac{\omega_1^2 - \lambda^2}{\sigma^2 - \lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$

- $\sigma=0.5~{
m GeV}, \Lambda=0.425~{
m GeV}$ and $\Rightarrow \langle \omega_1^2 h_{17}
angle=0.49~{
m GeV}^4$

- Blue line; 2019 model: $\sigma=0.5~{
 m GeV}$ and $\langle\omega_1^2h_{17}
 angle=0.27~{
 m GeV}^4$
- New function is 50% smaller than the 2010
 - New model give better constraints on $Q_1^c-Q_{7\gamma}$ contribution

AYESH GUNAWARDANA

Wayne State University



Figure: 2019 model vs 2010 model for h_{17}

• Orange dashed line: 2010 model $h_{17}(\omega_1,\mu) = \frac{2\lambda_2}{\sqrt{2\pi\sigma}} \frac{\omega_1^2 - \lambda^2}{\sigma^2 - \lambda^2} e^{-\frac{\omega_1^2}{2\sigma^2}}$

- $\sigma=0.5~{
m GeV}, \Lambda=0.425~{
m GeV}$ and $\Rightarrow \langle \omega_1^2 h_{17}
angle=0.49~{
m GeV}^4$

- Blue line; 2019 model: $\sigma=0.5~{
 m GeV}$ and $\langle\omega_1^2h_{17}
 angle=0.27~{
 m GeV}^4$
- New function is 50% smaller than the 2010
 - New model give better constraints on $Q_1^c Q_{7\gamma}$ contribution
- Consider also unknown higher moments, up to 6 Hermite polynomials

AYESH GUNAWARDANA

• Direct CP Asymmetry experimental bound: $A_{CP} = (1.5 \pm 2.0) \%$

[Y. Amhis et. al. EPJC 77, 895 (2017)]

• Direct CP Asymmetry experimental bound: $A_{CP} = (1.5 \pm 2.0) \%$

[Y. Amhis et. al. EPJC 77, 895 (2017)]

$$\mathcal{A}_{X_s\gamma}^{\rm SM} = \left(1.15 \times \frac{\tilde{\Lambda}_{17}^u - \tilde{\Lambda}_{17}^c}{300 \,\mathrm{MeV}} + 0.71\right) \% \text{ CP asymmetry}$$
$$\tilde{\Lambda}_{17}^u = \frac{2}{3} h_{17}(0)$$
$$\tilde{\Lambda}_{17}^c = \frac{2}{3} \int_{4m_c^2/m_b}^{\infty} \frac{d\omega_1}{\omega_1} \underbrace{f\left(\frac{m_c^2}{m_b\omega_1}\right)}_{\text{Perturbative}} \underbrace{h_{17}(\omega_1)}_{\text{Non-perturbative}}$$

• Direct CP Asymmetry experimental bound: $A_{CP} = (1.5 \pm 2.0) \%$

[Y. Amhis et. al. EPJC 77, 895 (2017)]

$$\mathcal{A}_{X_s\gamma}^{\rm SM} = \left(1.15 \times \frac{\tilde{\Lambda}_{17}^u - \tilde{\Lambda}_{17}^c}{300 \,\mathrm{MeV}} + 0.71\right) \% \text{ CP asymmetry}$$
$$\tilde{\Lambda}_{17}^u = \frac{2}{3} h_{17}(0)$$
$$\tilde{\Lambda}_{17}^c = \frac{2}{3} \int_{4m_c^2/m_b}^{\infty} \frac{d\omega_1}{\omega_1} \underbrace{f\left(\frac{m_c^2}{m_b\omega_1}\right)}_{\text{Perturbative}} \underbrace{h_{17}(\omega_1)}_{\text{Non-perturbative}}$$

Previously known values:

• Direct CP Asymmetry experimental bound: $A_{CP} = (1.5 \pm 2.0) \,\%$

[Y. Amhis et. al. EPJC 77, 895 (2017)]

$$\mathcal{A}_{X_{s}\gamma}^{\mathrm{SM}} = \left(1.15 \times \frac{\tilde{\Lambda}_{17}^{u} - \tilde{\Lambda}_{17}^{c}}{300 \mathrm{MeV}} + 0.71\right) \% \text{ CP asymmetry}$$
$$\tilde{\Lambda}_{17}^{u} = \frac{2}{3} h_{17}(0)$$
$$\tilde{\Lambda}_{17}^{c} = \frac{2}{3} \int_{4m_{c}^{2}/m_{b}}^{\infty} \frac{d\omega_{1}}{\omega_{1}} \underbrace{f\left(\frac{m_{c}^{2}}{m_{b}\omega_{1}}\right)}_{\text{Perturbative}} \underbrace{h_{17}(\omega_{1})}_{\text{Non-perturbative}}$$

• Previously known values:

$$\begin{split} -330 \mathrm{MeV} < \tilde{\Lambda}^{\textit{u}}_{17} < +525 \mathrm{MeV} \\ -9 \mathrm{MeV} < \tilde{\Lambda}^{\textit{c}}_{17} < +11 \mathrm{MeV} \end{split}$$

[M. Benzke, S. J. Lee, M. Neubert and G. Paz PRL 106, 141801(2011)]

AYESH GUNAWARDANA

• Direct CP Asymmetry experimental bound: $A_{CP} = (1.5 \pm 2.0) \%$

[Y. Amhis et. al. EPJC 77, 895 (2017)]

$$\mathcal{A}_{X_{s}\gamma}^{\rm SM} = \left(1.15 \times \frac{\tilde{\Lambda}_{17}^{u} - \tilde{\Lambda}_{17}^{c}}{300 \,\mathrm{MeV}} + 0.71\right) \% \text{ CP asymmetry}$$
$$\tilde{\Lambda}_{17}^{u} = \frac{2}{3} h_{17}(0)$$
$$\tilde{\Lambda}_{17}^{c} = \frac{2}{3} \int_{4m_{c}^{2}/m_{b}}^{\infty} \frac{d\omega_{1}}{\omega_{1}} \underbrace{f\left(\frac{m_{c}^{2}}{m_{b}\omega_{1}}\right)}_{\text{Perturbative}} \underbrace{h_{17}(\omega_{1})}_{\text{Non-perturbative}}$$

Previously known values:

$$\begin{split} -330 \mathrm{MeV} < \tilde{\Lambda}^{\textit{u}}_{17} < +525 \mathrm{MeV} \\ -9 \mathrm{MeV} < \tilde{\Lambda}^{\textit{c}}_{17} < +11 \mathrm{MeV} \end{split}$$

[M. Benzke, S. J. Lee, M. Neubert and G. Paz PRL 106, 141801(2011)] • We plan to improve these estimates AYESH GUNAWARDANA Wayne State University Reducing Uncertainties in $\overline{B} \rightarrow X_{s}\gamma$ Decay

Conclusion

- $\bar{B}
 ightarrow X_s \gamma$ is a important New Physics probe
- Non perturbative error of the decay rate is 5%
- $Q_1^c Q_{7\gamma}$ is the largest contribution to the error
- Better estimates for $Q_1^c Q_{7\gamma}$ obtained from moments of h_{17}
- New estimates for CP asymmetry
- Reduce non-perturbative error on rate and CP asymmetry