

Asymmetry Observables and the Origin of $R_{D^{(*)}}$ Anomalies

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Based on : 1810.06597, 1904.XXXXX

In collaboration with : Matthew Buckley, David Shih

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Outline

- $R_{D^{(*)}}$ Solutions
- Some Asymmetry Observables
- More On $F_{D^*}^L$

New Physics in the Flavor Experiments

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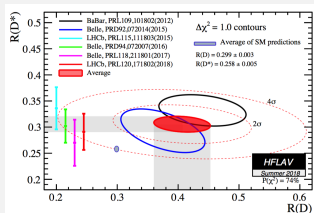
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$$R_{D^{(*)}} \equiv \frac{\Gamma(B \rightarrow D^{(*)} \tau \nu)}{\Gamma(B \rightarrow D^{(*)} l \nu)}, \quad l = e, \mu$$



$$R_D^{obs} = 0.407 \pm 0.046,$$

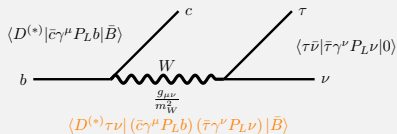
$$R_{D^*}^{obs} = 0.304 \pm 0.015,$$

$$R_D^{SM} = 0.299 \pm 0.003,$$

$$R_{D^*}^{SM} = 0.258 \pm 0.005.$$

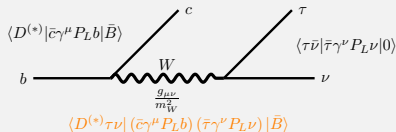
The Most General EFT

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- The most general dim-6 effective Hamiltonian:

$$\mathcal{H}_{\text{eff}} = \frac{4G_F V_{cb}}{\sqrt{2}} \sum_{\substack{X=S,V,T \\ M,N=L,R}} C_{MN}^X \mathcal{O}_{MN}^X,$$

$$\mathcal{O}_{MN}^S \equiv (\bar{c} P_M b) (\bar{\tau} P_N \nu),$$

$$\mathcal{O}_{MN}^V \equiv (\bar{c} \gamma^\mu P_M b) (\bar{\tau} \gamma_\mu P_N \nu),$$

$$\mathcal{O}_{MN}^T \equiv (\bar{c} \sigma^{\mu\nu} P_M b) (\bar{\tau} \sigma_{\mu\nu} P_N \nu),$$

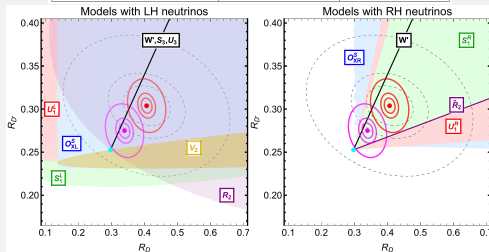
for $M, N = R$ or L (SM : $C_{LL}^V = 1$).

Minimal Models

Mediator	Operator Combination	Viability
Colorless Scalars	\mathcal{O}_{XL}^S	\times ($Br(B_c \rightarrow \tau\nu)$)
$W^{3\mu}$ (LH fermions)	\mathcal{O}_{LL}^V	\times (collider bounds)
S_1 LQ ($\bar{3}, 1, 1/3$) (LH fermions)	$\mathcal{O}_{LL}^S - x\mathcal{O}_{LL}^T, \mathcal{O}_{LL}^V$	✓
U_1^μ LQ ($3, 1, 2/3$) (LH fermions)	$\mathcal{O}_{RL}^S, \mathcal{O}_{LL}^V$	✓
R_2 LQ ($3, 2, 7/6$)	$\mathcal{O}_{LL}^S + x\mathcal{O}_{LL}^T$	✓
S_3 LQ ($\bar{3}, 3, 1/3$)	\mathcal{O}_{LL}^V	\times ($b \rightarrow s\nu\nu$)
U_3^μ LQ ($3, 3, 2/3$)	\mathcal{O}_{LL}^V	\times ($b \rightarrow s\nu\nu$)
V_2^μ LQ ($\bar{3}, 2, 5/6$)	\mathcal{O}_{RL}^S	\times ($R_{D^{(*)}}$ value)
Colorless Scalars	\mathcal{O}_{XR}^S	\times ($Br(B_c \rightarrow \tau\nu)$)
$W^{3\mu}$ (RH fermions)	\mathcal{O}_{RR}^V	✓
\tilde{R}_2 LQ ($3, 2, 1/6$)	$\mathcal{O}_{RR}^S + x\mathcal{O}_{RR}^T$	\times ($b \rightarrow s\nu\nu$)
S_1 LQ ($\bar{3}, 1, 1/3$) (RH fermions)	$\mathcal{O}_{RR}^V, \mathcal{O}_{RR}^S - x\mathcal{O}_{RR}^T$	✓
U_1^μ LQ ($3, 1, 2/3$) (RH fermions)	$\mathcal{O}_{LR}^S, \mathcal{O}_{RR}^V$	✓

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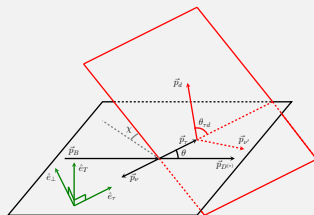
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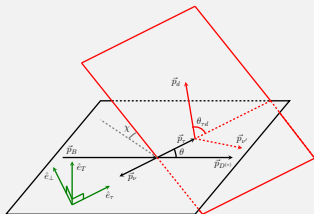
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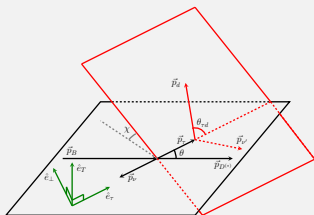
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$$\mathcal{P}_{\hat{e}}^{(*)} = \frac{\Gamma_{+\hat{e}}^{(*)} - \Gamma_{-\hat{e}}^{(*)}}{\Gamma_{+\hat{e}}^{(*)} + \Gamma_{-\hat{e}}^{(*)}}, \quad \mathcal{A}_{FB}^{(*)} = \frac{1}{\Gamma^{(*)}} \left(- \int_{\theta=0}^{\theta=\pi/2} + \int_{\theta=\pi/2}^{\theta=\pi} \right) d\theta \frac{d\Gamma^{(*)}}{d\theta}.$$

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Observable	\mathcal{A}_{FB}	\mathcal{A}_{FB}^*	\mathcal{P}_T	\mathcal{P}_T^*	\mathcal{P}_\perp	\mathcal{P}_\perp^*	\mathcal{P}_T	\mathcal{P}_T^*
SM value	-0.360	0.063	0.325	-0.497	-0.842	-0.499	0	0
Projected Precision	10%	-	3%	-	10%	-	-	-

Discerning Different Solutions at Belle II

- Let us assume we measure $R_{D^{(*)}}$ in Belle II and discover NP.

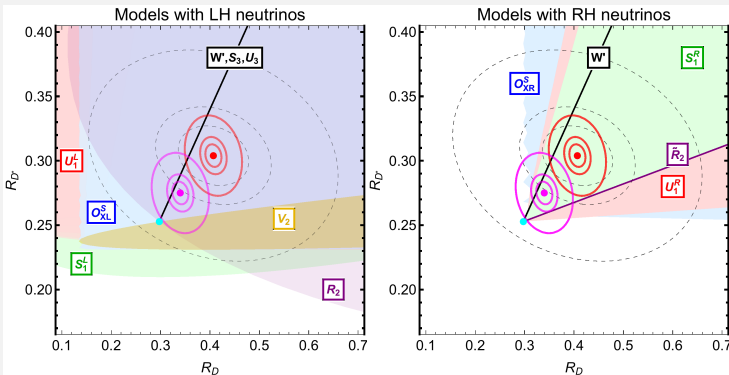
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- It highly depends on the measured $R_{D^{(*)}}$ value.

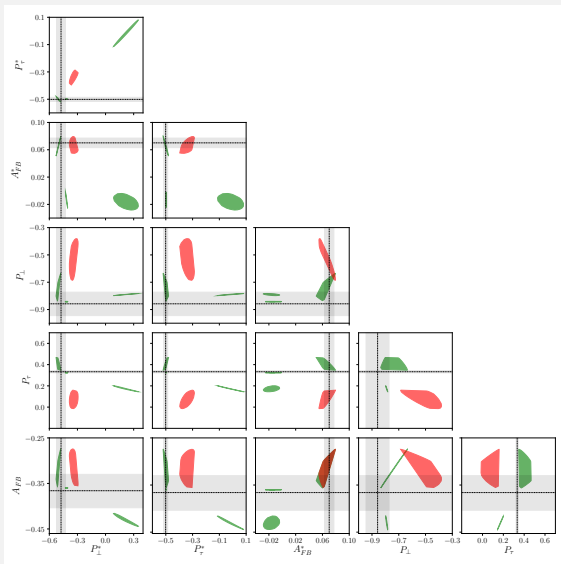
Two Extreme Outcomes for $R_{D^{(*)}}$



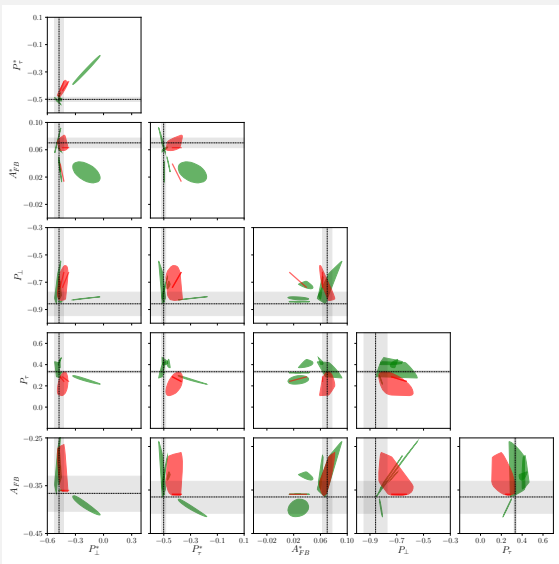
$$R_D = 0.407 \quad R_{D^*} = 0.304$$

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Distinguishing Various Minimal Models

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- We develop a simple χ^2 test to see how well each pair of models can be distinguished.
- Can tell all the models apart; we may need to resort the CP -odd observable $\mathcal{P}_T^{(*)}$ in the second scenario.

F_{D^L*}^L : Another Asymmetry Observable

$$F_{D^*}^L = \frac{\Gamma(\bar{B} \rightarrow D_L^* \tau \nu)}{\Gamma(\bar{B} \rightarrow D_L^* \tau \nu) + \Gamma(\bar{B} \rightarrow D_T^* \tau \nu)}.$$

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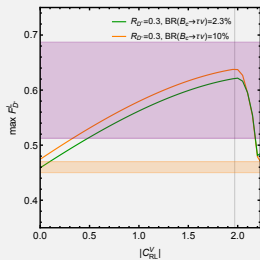
- None of the existing minimal models can accommodate this new observation.
- The observed value of $R_{D(*)}$ and $Br(B_c \rightarrow \tau \nu)$ are constraining.
- Is there any combination of the dim-6 operators that can explain the observed value?

Explaining the Observed $F_{D^*}^L$

- We look for the maximum of $F_{D^*}^L$ under certain constraints.
We show it can be achieved with all real WCs.

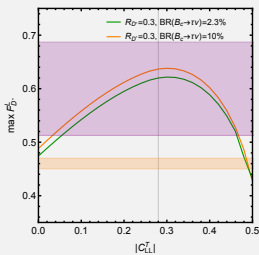
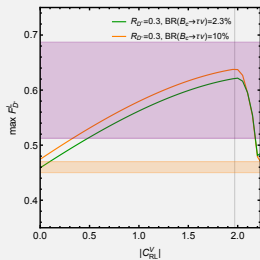
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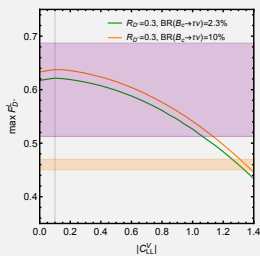
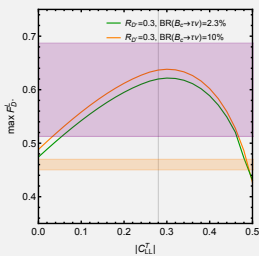
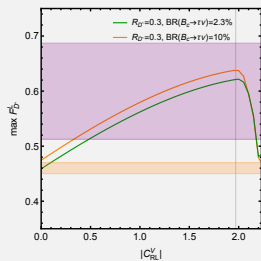
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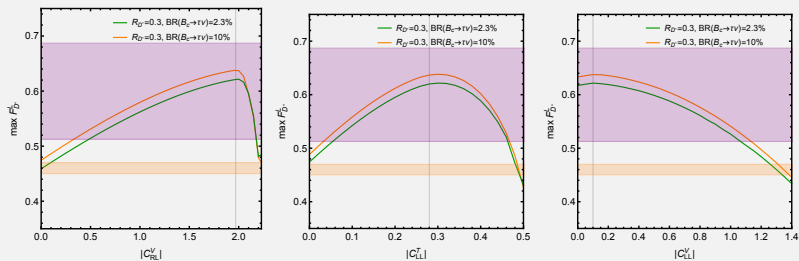
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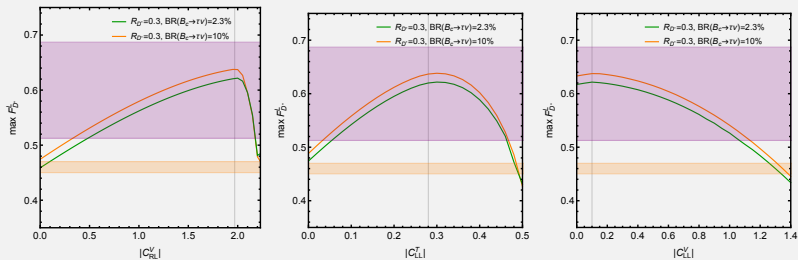
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- Relatively large C_{RL}^V , C_{LL}^T , and C_{LL}^V are required to explain the observed $F_{D^*}^L$.
- Not sensitive to individual operators. Need a combination of all.

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- We can resort to some asymmetry observables ($\mathcal{P}_T^{(*)}$, $\mathcal{A}_{FB}^{(*)}$, $\mathcal{P}_\perp^{(*)}$) to distinguish various models from one another.

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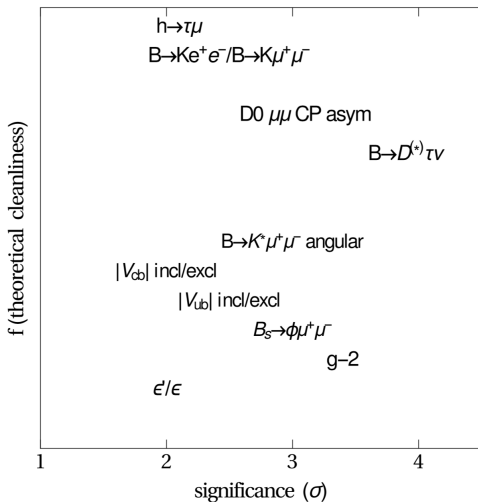
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THANK YOU!

BACK UP SLIDES

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Other Anomalies



Uncertainties

BaBar@Hadronic($\tau \rightarrow l$)

Source of uncertainty	(%)	
	$\mathcal{R}(D)$	$\mathcal{R}(D^*)$
Additive uncertainties		
PDFs		
MC statistics	4.4	2.0
$B \rightarrow D^{(*)}(\tau/\ell)\bar{\nu}$ FFs	0.2	0.2
$D^{**} \rightarrow D^{(*)}(\pi^0/\pi^\pm)$	0.7	0.5
$B(\bar{B} \rightarrow D^{**}\ell^-\bar{\nu}_\ell)$	0.8	0.3
$B(\bar{B} \rightarrow D^{**}\tau^-\bar{\nu}_\tau)$	1.8	1.7
$D^{**} \rightarrow D^{(*)}\pi\pi$	2.1	2.6
Cross-feed constraints		
MC statistics	2.4	1.5
$f_{D^{**}}$	5.0	2.0
Feed-up/feed-down	1.3	0.4
Isospin constraints	1.2	0.3
Fixed backgrounds		
MC statistics	3.1	1.5
Efficiency corrections	3.9	2.3
Multiplicative uncertainties		
MC statistics	1.8	1.2
$B \rightarrow D^{(*)}(\tau/\ell^+)\bar{\nu}$ FFs	1.6	0.4
Lepton PID	0.6	0.6
π^0/π^\pm from $D^* \rightarrow D\pi$	0.1	0.1
Detection/Reconstruction	0.7	0.7
$B(\tau^- \rightarrow \ell^-\bar{\nu}_\ell\nu_\tau)$	0.2	0.2
Total syst. uncertainty	9.6	5.5
Total stat. uncertainty	13.1	7.1
Total uncertainty	16.2	9.0

Belle@Semileptonic($\tau \rightarrow l$)

Sources	$\mathcal{R}(D^*)$ [%]
MC size for each PDF shape	2.2
PDF shape of the normalization in $\cos\theta_{B,D^*}\ell$	+1.1 -0.0
PDF shape of $B \rightarrow D^{**}\ell\nu_\ell$	+1.0 -1.7
PDF shape and yields of fake $D^{(*)}$	1.4
PDF shape and yields of $B \rightarrow X_c D^*$	1.1
Reconstruction efficiency ratio $\varepsilon_{\text{norm}}/\varepsilon_{\text{sig}}$	1.2
Modeling of semileptonic decay	0.2
$B(\tau^- \rightarrow \ell^-\bar{\nu}_\ell\nu_\tau)$	0.2
Total systematic uncertainty	+3.4 -3.5

Scales with MC statistics

Scales with DATA statistics

Theory/External

Irreducible
Requires additional studies

Belle@Hadronic($\tau \rightarrow h$)

Source	$R(D^*)$	P_τ
Hadronic B composition	+7.8% -6.9%	+0.14 -0.11
MC statistics for each PDF shape	+3.3% -2.8%	+0.13 -0.11
Fake D^* PDF shape	3.0%	0.010
Fake D^* yield	1.7%	0.016
$B \rightarrow D^{**}\ell^-\bar{\nu}_\ell$	2.1%	0.051
$\bar{B} \rightarrow D^{**}\tau^-\bar{\nu}_\tau$	1.1%	0.003
$\bar{B} \rightarrow D^*\ell^-\bar{\nu}_\ell$	2.4%	0.008
τ daughter and ℓ^- efficiency	2.1%	0.018
MC statistics for efficiency calculation	1.0%	0.018
EvtGen decay model	+0.8% -0.0%	+0.016 -0.000
Fit bias	0.3%	0.008
$B(\tau^- \rightarrow \pi^-\nu_\tau)$ and $B(\tau^- \rightarrow \rho^-\nu_\tau)$	0.3%	0.002
P_τ correction function	0.1%	0.018
Common sources		
Tagging efficiency correction	1.4%	0.014
D^* reconstruction	1.3%	0.007
D sub-decay branching fractions	0.7%	0.005
Number of $B\bar{B}$	0.4%	0.005
Total systematic uncertainty	+0.4% -9.5%	+0.20 -0.17

Individual Operator Effects

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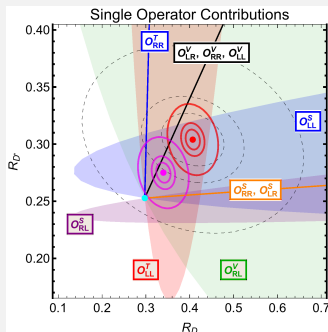
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All Operators

	Operator	Fierz identity	Allowed Current	$\delta\mathcal{L}_{\text{int}}$
\mathcal{O}_{V_L}	$(\bar{c}\gamma_\mu P_L b)(\bar{\tau}\gamma^\mu P_L\nu)$		$(\mathbf{1}, \mathbf{3})_0$	$(g_q\bar{q}_L\boldsymbol{\tau}\gamma^\mu q_L + g_\ell\bar{\ell}_L\boldsymbol{\tau}\gamma^\mu \ell_L)W'_\mu$
\mathcal{O}_{V_R}	$(\bar{c}\gamma_\mu P_R b)(\bar{\tau}\gamma^\mu P_L\nu)$		$\rangle(\mathbf{1}, \mathbf{2})_{1/2}$	$(\lambda_d\bar{q}_L d_R\phi + \lambda_u\bar{q}_L u_R i\tau_2\phi^\dagger + \lambda_\ell\bar{\ell}_L e_R\phi)$
\mathcal{O}_{S_R}	$(\bar{c}P_R b)(\bar{\tau}P_L\nu)$			
\mathcal{O}_{S_L}	$(\bar{c}P_L b)(\bar{\tau}P_L\nu)$			
\mathcal{O}_T	$(\bar{c}\sigma^{\mu\nu} P_L b)(\bar{\tau}\sigma_{\mu\nu} P_L\nu)$			
\mathcal{O}'_{V_L}	$(\bar{\tau}\gamma_\mu P_L b)(\bar{c}\gamma^\mu P_L\nu) \longleftrightarrow \mathcal{O}_{V_L}$			
\mathcal{O}'_{V_R}	$(\bar{\tau}\gamma_\mu P_R b)(\bar{c}\gamma^\mu P_L\nu) \longleftrightarrow -2\mathcal{O}_{S_R}$		$\rangle(\mathbf{3}, \mathbf{1})_{2/3}$	$(\lambda\bar{q}_L\gamma_\mu \ell_L + \tilde{\lambda}\bar{d}_R\gamma_\mu e_R)U^\mu$
\mathcal{O}'_{S_R}	$(\bar{\tau}P_R b)(\bar{c}P_L\nu) \longleftrightarrow -\frac{1}{2}\mathcal{O}_{V_R}$			
\mathcal{O}'_{S_L}	$(\bar{\tau}P_L b)(\bar{c}P_L\nu) \longleftrightarrow -\frac{1}{2}\mathcal{O}_{S_L} - \frac{1}{8}\mathcal{O}_T$		$(\mathbf{3}, \mathbf{2})_{7/6}$	$(\lambda\bar{u}_R\ell_L + \tilde{\lambda}\bar{q}_L i\tau_2 e_R)R$
\mathcal{O}'_T	$(\bar{\tau}\sigma^{\mu\nu} P_L b)(\bar{c}\sigma_{\mu\nu} P_L\nu) \longleftrightarrow -6\mathcal{O}_{S_L} + \frac{1}{2}\mathcal{O}_T$			
\mathcal{O}''_{V_L}	$(\bar{\tau}\gamma_\mu P_L c^c)(\bar{b}^c\gamma^\mu P_L\nu) \longleftrightarrow -\mathcal{O}_{V_R}$		$(\bar{\mathbf{3}}, \mathbf{2})_{5/3}$	$(\lambda\bar{d}_R^c\gamma_\mu \ell_L + \tilde{\lambda}\bar{q}_L^c\gamma_\mu e_R)V^\mu$
\mathcal{O}''_{V_R}	$(\bar{\tau}\gamma_\mu P_R c^c)(\bar{b}^c\gamma^\mu P_L\nu) \longleftrightarrow -2\mathcal{O}_{S_R}$			
\mathcal{O}''_{S_R}	$(\bar{\tau}P_R c^c)(\bar{b}^c P_L\nu) \longleftrightarrow \frac{1}{2}\mathcal{O}_{V_L}$		$(\bar{\mathbf{3}}, \mathbf{3})_{1/3}$	$\lambda\bar{q}_L^c i\tau_2 \boldsymbol{\tau} \ell_L \mathbf{S}$
\mathcal{O}''_{S_L}	$(\bar{\tau}P_L c^c)(\bar{b}^c P_L\nu) \longleftrightarrow -\frac{1}{2}\mathcal{O}_{S_L} + \frac{1}{8}\mathcal{O}_T$			
\mathcal{O}''_T	$(\bar{\tau}\sigma^{\mu\nu} P_L c^c)(\bar{b}^c\sigma_{\mu\nu} P_L\nu) \longleftrightarrow -6\mathcal{O}_{S_L} - \frac{1}{2}\mathcal{O}_T$		$\rangle(\bar{\mathbf{3}}, \mathbf{1})_{1/3}$	$(\lambda\bar{q}_L^c i\tau_2 \ell_L + \tilde{\lambda}\bar{u}_R^c e_R)S$

Figure: [1506.08896]

Constrain I : $Br(B_c \rightarrow \tau\nu)$

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- Enhanced contribution from the scalar operators (same combination appearing in R_{D^*}).
- $Br(B_c \rightarrow \tau\nu) \leq 10\%$ from the $B_u \rightarrow \tau\nu$ at Z peak at LEP.

Constrain II : $b \rightarrow s\nu\nu$

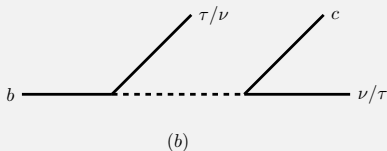
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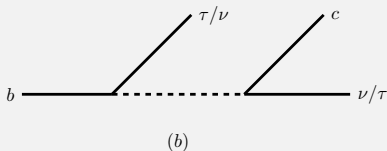
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These are neutral current constraints so will put severe bounds on the affected models.

Constrain II : $b \rightarrow s\nu\nu$

$$BR(B \rightarrow X_s \nu\nu) \leq 6.4 \times 10^{-4},$$

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$$\epsilon \equiv \frac{\sqrt{|C_L^\nu|^2 + |C_R^\nu|^2}}{|(C_L^\nu)^{SM}|}, \quad \eta \equiv -\frac{\text{Re}(C_L^\nu C_R^{\nu*})}{|C_L^\nu|^2 + |C_R^\nu|^2}.$$

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$$C_{LL}^V \leq 0.006,$$

$$C_{RR}^S \leq 0.01.$$

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On a W' coupled to the LH particles : The accompanying Z' is severely constrained. Ruled out unless Z' is a wide resonance.

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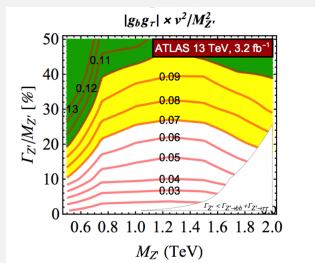


Figure: [1609.07138]

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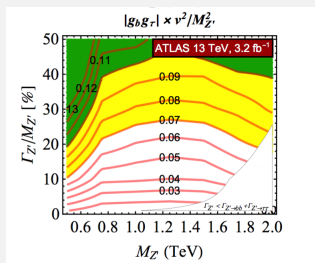


Figure: [1609.07138]

Things are better with RH neutrinos. But still severely tight from the LHC direct searches.

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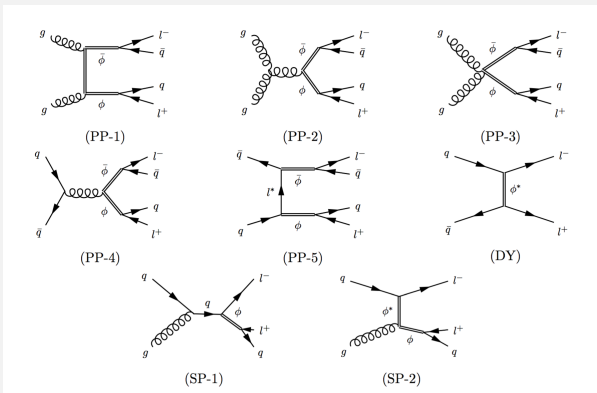
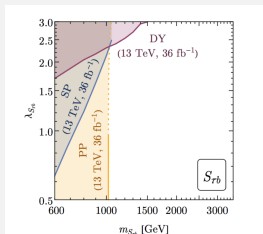
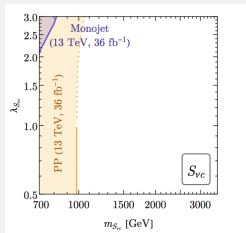
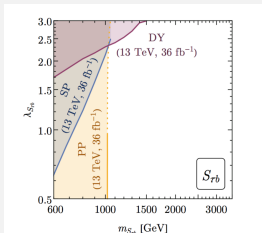


Figure: [1810.10017]

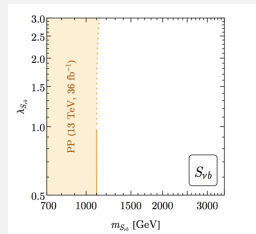
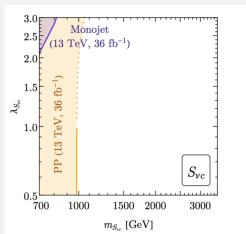
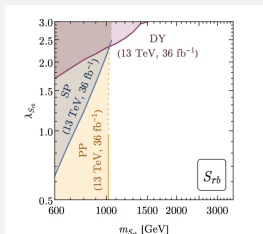
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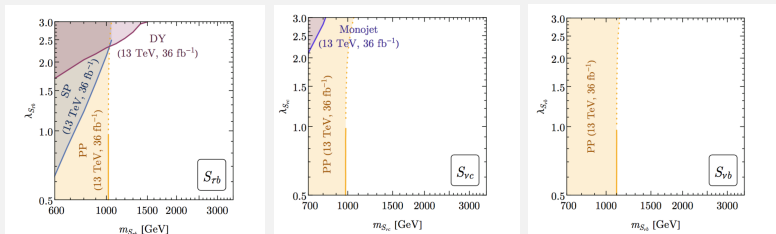


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- Not quite strong enough to kill any LQ yet.

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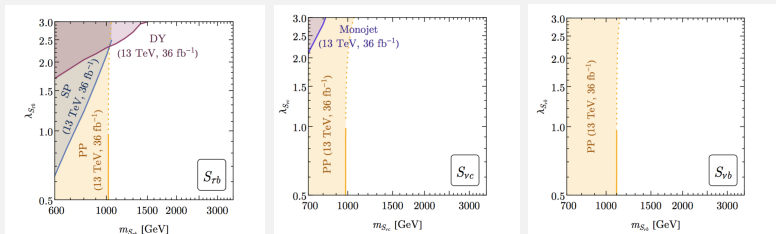


Figure: [1810.10017]

- Not quite strong enough to kill any LQ yet.
- Can always introduce a new decay channel that the direct searches are blind too. LHC is trying to close that gap.

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- Electroweak precision bounds : When introducing new gauge bosons or fermion mixings.

Constraining Hidden Channels

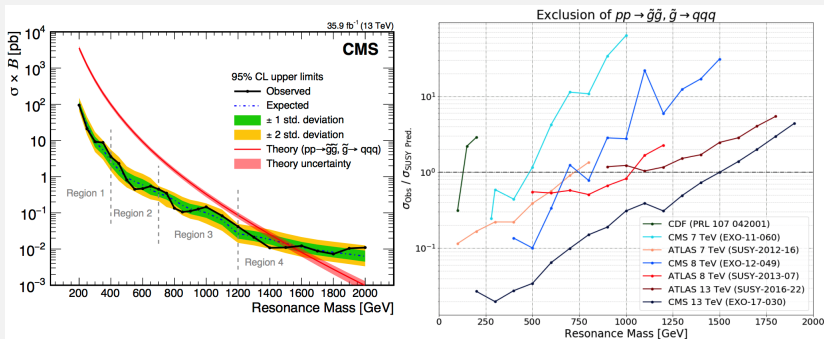
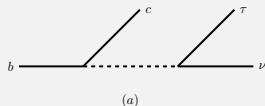


Figure: Talk by Abhijith Gandrakota

Calculation Steps



- Calculate the leptonic side matrix element.
- Use the available results (e.g. HQET or Lattice) for the Hadronic side.
- Integrate over various final state labels to get the numerical results.

LH \leftrightarrow RH

$$h_{\tau} \rightarrow -h_{\tau}, \quad C_{LL}^{S,T} \leftrightarrow \left(C_{RR}^{S,T}\right)^*, \quad C_{RL}^X \leftrightarrow \left(C_{LR}^X\right)^*,$$
$$1 + C_{LL}^V \leftrightarrow \left(C_{RR}^V\right)^*,$$

$$R_{D^{(*)}} \rightarrow R_{D^{(*)}}, \quad \mathcal{P}_X \rightarrow -\mathcal{P}_X, \quad \mathcal{A}_{FB} \rightarrow \mathcal{A}_{FB}.$$

Numerical Equations

$$\begin{aligned}
 A_{FB} &\approx \frac{1}{R_D} \left\{ -0.11 \left(|1 + C_{LL}^V + C_{RL}^V|^2 + |C_{RR}^V + C_{LR}^V|^2 \right) \right. \\
 &\quad - 0.35 \mathcal{R}e \left[(C_{LL}^S + C_{RL}^S)(C_{LL}^T)^* + (C_{RR}^S + C_{LR}^S)^*(C_{RR}^T) \right] \\
 &\quad - 0.24 \mathcal{R}e \left[(1 + C_{LL}^V + C_{RL}^V)(C_{LL}^T)^* + (C_{RR}^V + C_{LR}^V)^*(C_{RR}^T) \right] \\
 &\quad \left. - 0.15 \mathcal{R}e \left[(1 + C_{LL}^V + C_{RL}^V)(C_{LL}^S + C_{RL}^S)^* + (C_{RR}^V + C_{LR}^V)^*(C_{RR}^S + C_{LR}^S) \right] \right\} \\
 A_{FB}^* &\approx \frac{1}{R_{D^*}} \left\{ -0.813 \left(|C_{LL}^T|^2 + |C_{RR}^T|^2 \right) \right. \\
 &\quad + 0.016 \left(|1 + C_{LL}^V|^2 + |C_{RR}^V|^2 \right) - 0.082 \left(|C_{RL}^V|^2 + |C_{LR}^V|^2 \right) \\
 &\quad + 0.066 \mathcal{R}e \left[C_{RL}^V(1 + C_{LL}^V)^* + (C_{LR}^V)^* C_{RR}^V \right] \\
 &\quad + 0.095 \mathcal{R}e \left[(C_{RL}^S - C_{LL}^S)(C_{LL}^T)^* + (C_{LR}^S - C_{RR}^S)^* C_{RR}^T \right] \\
 &\quad + 0.395 \mathcal{R}e \left[(1 + C_{LL}^V - C_{RL}^V)(C_{LL}^T)^* + (C_{RR}^V - C_{LR}^V)^*(C_{RR}^T) \right] \\
 &\quad + 0.023 \mathcal{R}e \left[(C_{LL}^S - C_{RL}^S)(1 + C_{LL}^V - C_{RL}^V)^* + (C_{RR}^S - C_{LR}^S)^*(C_{RR}^V - C_{LR}^V) \right] \\
 &\quad \left. - 0.142 \mathcal{R}e \left[(C_{LL}^T)(1 + C_{LL}^V + C_{RL}^V)^* + (C_{RR}^T)^*(C_{RR}^V + C_{LR}^V) \right] \right\},
 \end{aligned}$$

Numerical Equations

$$\begin{aligned}
 \mathcal{P}_\tau &\approx \frac{1}{R_D} \left\{ 0.402 \left(|C_{LL}^S + C_{RL}^S|^2 - |C_{RR}^S + C_{LR}^S|^2 \right) \right. \\
 &+ 0.013 \left[|C_{LL}^T|^2 - |C_{RR}^T|^2 \right] + 0.097 \left[|1 + C_{LL}^V + C_{RL}^V|^2 - |C_{RR}^V + C_{LR}^V|^2 \right] \\
 &+ 0.512 \operatorname{Re} \left[(1 + C_{LL}^V + C_{RL}^V)(C_{LL}^S + C_{RL}^S)^* - (C_{RR}^V + C_{LR}^V)^*(C_{RR}^S + C_{LR}^S) \right] \\
 &- 0.099 \operatorname{Re} \left[(1 + C_{LL}^V + C_{RL}^V)(C_{LL}^T)^* - (C_{RR}^V + C_{LR}^V)^*(C_{RR}^T) \right] \left. \right\} \\
 \mathcal{P}_\tau^* &\approx \frac{1}{R_{D^*}} \left\{ -0.127 \left(|1 + C_{LL}^V|^2 + |C_{RL}^V|^2 - |C_{RR}^V|^2 - |C_{LR}^V|^2 \right) \right. \\
 &+ 0.011 \left(|C_{LL}^S - C_{RL}^S|^2 - |C_{RR}^S - C_{LR}^S|^2 \right) + 0.172 \left(|C_{LL}^T|^2 - |C_{RR}^T|^2 \right) \\
 &+ 0.031 \operatorname{Re} \left[(1 + C_{LL}^V - C_{RL}^V)(C_{RL}^S - C_{LL}^S)^* - (C_{RR}^V - C_{LR}^V)^*(C_{LR}^S - C_{RR}^S) \right] \\
 &+ 0.350 \operatorname{Re} \left[(1 + C_{LL}^V)(C_{LL}^T)^* - (C_{RR}^V)^*(C_{RR}^T) \right] \\
 &- 0.481 \operatorname{Re} \left[(C_{RL}^V)(C_{LL}^T)^* - (C_{LR}^V)^*(C_{RR}^T) \right] \\
 &+ 0.216 \operatorname{Re} \left[(1 + C_{LL}^V)(C_{RL}^V)^* - (C_{RR}^V)^*(C_{LR}^V) \right] \left. \right\}.
 \end{aligned}$$

Numerical Equations

$$\begin{aligned}
 \mathcal{P}_\perp &\approx \frac{1}{R_D} \text{Re} \left\{ -0.350 \left[(C_{LL}^T) (C_{LL}^S + C_{RL}^S)^* - (C_{RR}^T)^* (C_{RR}^S + C_{LR}^S) \right] \right. \\
 &- 0.357 \left[(1 + C_{LL}^V + C_{RL}^V) (C_{LL}^S + C_{RL}^S)^* - (C_{RR}^V + C_{LR}^V)^* (C_{RR}^S + C_{LR}^S) \right] \\
 &- 0.247 \left[(1 + C_{LL}^V + C_{RL}^V)^* (C_{LL}^T) - (C_{RR}^V + C_{LR}^V) (C_{RR}^T)^* \right] \\
 &\left. - 0.250 \left[|1 + C_{LL}^V + C_{RL}^V|^2 - |C_{RR}^V + C_{LR}^V|^2 \right] \right\}
 \end{aligned}$$

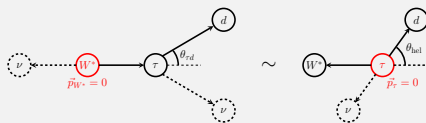
$$\begin{aligned}
 \mathcal{P}_\perp^* &\approx \frac{1}{R_{D^*}} \text{Re} \left\{ (C_{RR}^S - C_{LR}^S) \left[0.099 C_{RR}^T - 0.054 (C_{RR}^V - C_{LR}^V) \right]^* \right. \\
 &- (C_{LL}^S - C_{RL}^S)^* \left[0.099 C_{LL}^T - 0.054 (1 + C_{LL}^V - C_{RL}^V) \right] \\
 &+ (C_{RR}^T) \left[0.146 C_{RR}^V - 0.478 C_{LR}^V - 1.855 C_{RR}^T \right]^* \\
 &- (C_{LL}^T)^* \left[0.146 (1 + C_{LL}^V) - 0.478 C_{RL}^V - 1.855 C_{LL}^T \right] \\
 &+ (C_{LR}^V) \left[-0.081 C_{RR}^T + 0.025 C_{LR}^V - 0.075 C_{RR}^V \right]^* \\
 &- (C_{RL}^V)^* \left[-0.081 C_{LL}^T + 0.025 C_{RL}^V - 0.075 (1 + C_{LL}^V) \right] \\
 &\left. + (C_{RR}^V) \left[-0.071 C_{RR}^T - 0.075 C_{LR}^V + 0.126 C_{RR}^V \right]^* \right\}
 \end{aligned}$$

Numerical Equations

$$\begin{aligned}
 \mathcal{P}_T &\approx \frac{1}{R_D} \text{Im} \left\{ -0.350 \left[(C_{LL}^T) (C_{LL}^S + C_{RL}^S)^* - (C_{RR}^T)^* (C_{RR}^S + C_{LR}^S) \right] \right. \\
 &\quad - 0.357 \left[(1 + C_{LL}^V + C_{RL}^V) (C_{LL}^S + C_{RL}^S)^* - (C_{RR}^V + C_{LR}^V)^* (C_{RR}^S + C_{LR}^S) \right] \\
 &\quad \left. - 0.247 \left[(1 + C_{LL}^V + C_{RL}^V)^* (C_{LL}^T) - (C_{RR}^V + C_{LR}^V) (C_{RR}^T)^* \right] \right\} \\
 \mathcal{P}_T^* &\approx \frac{1}{R_{D^*}} \text{Im} \left\{ (C_{RR}^S - C_{LR}^S) \left[0.099 C_{RR}^T - 0.054 (C_{RR}^V - C_{LR}^V) \right]^* \right. \\
 &\quad - (C_{LL}^S - C_{RL}^S)^* \left[0.099 C_{LL}^T - 0.054 (1 + C_{LL}^V - C_{RL}^V) \right] \\
 &\quad + (C_{RR}^T) \left[0.146 C_{RR}^V - 0.478 C_{LR}^V \right]^* - (C_{LL}^T)^* \left[0.146 (1 + C_{LL}^V) - 0.478 C_{RL}^V \right] \\
 &\quad - (C_{LR}^V) \left[0.081 C_{RR}^T \right]^* + (C_{RL}^V)^* \left[0.081 C_{LL}^T \right] \\
 &\quad \left. - (C_{RR}^V) \left[0.071 C_{RR}^T \right]^* + (1 + C_{LL}^V)^* \left[0.071 C_{LL}^T \right] \right\}
 \end{aligned}$$

P_τ Measurement

$$\frac{1}{\Gamma} \frac{d\Gamma}{d\theta_{\text{hel}}} = \frac{1}{2} (1 + \alpha_d \mathcal{P}_\tau^* \cos \theta_{\text{hel}})$$

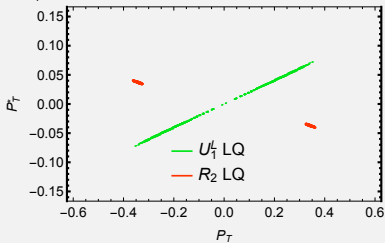
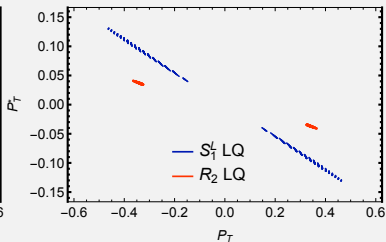
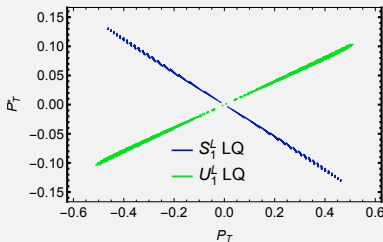


$$\cos \theta_{\tau d} = \frac{2E_\tau E_d - m_\tau^2 - m_d^2}{2|\vec{p}_\tau||\vec{p}_d|} \quad q^2 \text{ - frame}$$

$$|\vec{p}_\tau| = \frac{q^2 - m_\tau^2}{2\sqrt{q^2}} \quad q^2 \text{ - frame}$$

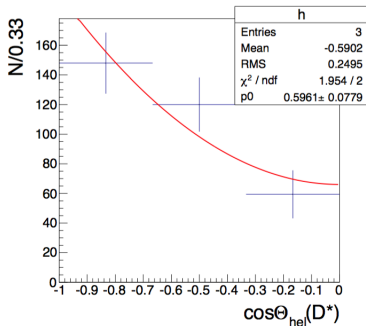
$$|\vec{p}_d^\tau| \cos \theta_{\text{hel}} = -\gamma \frac{|\vec{p}_\tau|}{E_\tau} E_d + \gamma |\vec{p}_d| \cos \theta_{\tau d} \quad \tau \text{ - frame}$$

$\mathcal{P}_T^{(*)}$



F_{D^{*}}^L Measurement

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_{\text{hel}}(D^*)} = \frac{3}{4} [2F_L^{D^*} \cos^2(\theta_{\text{hel}}(D^*)) + (1 - F_L^{D^*}) \sin^2(\theta_{\text{hel}}(D^*))]$$



Number of events in:

I bin: 151 ± 21

II bin: 125 ± 19

III bin: 55 ± 15

- signal yields corrected for acceptance variations

Dominant systematics:

- MC statistics (AR shape and peaking background)

= ± 0.03

Figure: Talk by Karol Adamczyk @ CKM 2018