Boson Stars from Repulsive Scalar Theory and the Gravitational Wave Signature
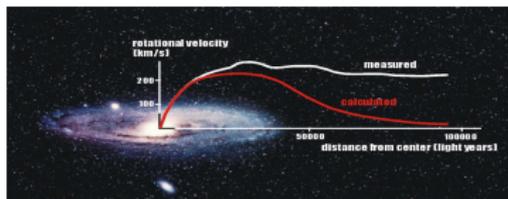
Chen Sun

Brown University/ CAS-ITP

Croon, Gleiser, Mohapatra, CS, 1802.08259
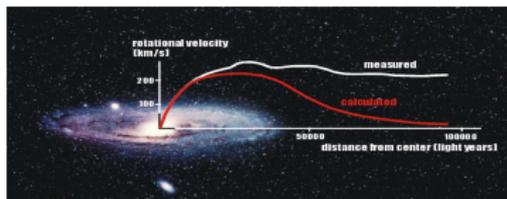*Croon, Fan, CS, 1810.01420*
Guo, Sinha, CS, 1904.07871

A Tale of Many Scales

- Large Scale:
    - Cosmic microwave background
    - Large scale structure
    - Supernovae, BAO, Bullet Cluster, etc.

- Small Scale:
    - Galactic rotation curves
    - Core-cusp problem
    - Missing satellites
    - Too-big-to-fail
    - Baryonic Tully-Fisher relation
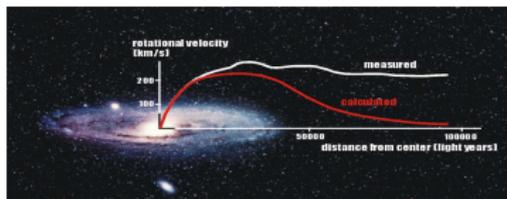
- ΛCDM

- Large Scale:
    - Cosmic microwave background
    - Large scale structure
    - Supernovae, BAO, Bullet Cluster, etc.

- Small Scale:
    - Galactic rotation curves
    - Core-cusp problem
    - Missing satellites
    - Too-big-to-fail
    - Baryonic Tully-Fisher relation

- ΛCDM $+...$?
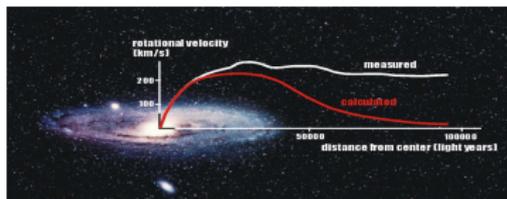
A Tale of Many Scales

- Large Scale:
  - Cosmic microwave background
  - Large scale structure
  - Supernovae, BAO, Bullet Cluster, etc.

- Small Scale:
  - Galactic rotation curves
  - Core-cusp problem
  - Missing satellites
  - Too-big-to-fail
  - Baryonic Tully-Fisher relation

- ΛCDM +... ?
  - "Everything can be fixed by baryon physics."
  - New type of dark matter is required. Particle nature of DM:
  - Gravity needs to be modified.

- Large Scale:
  - Cosmic microwave background
  - Large scale structure
  - Supernovae, BAO, Bullet Cluster, etc.

- Small Scale:
  - Galactic rotation curves
  - Core-cusp problem
  - Missing satellites
  - Too-big-to-fail
  - Baryonic Tully-Fisher relation

- ΛCDM +... ?
  - "Everything can be fixed by baryon physics."
  - New type of dark matter is required. Particle nature of DM:
    - SIDM (SIMP)
    - WDM
    - Superfluid DM
    - Axion like particles
    - BEC DM
  - Gravity needs to be modified.

Many body quantum system   Hartree-Fock approximation

$$\Phi(x) = \phi_1(x_1) \otimes \phi_2(x_2) \pm \phi_2(x_1) \otimes \phi_1(x_2)$$

Ground state of identical bosons   Gross-Pitaevskii Eqn
single wave function   aka nonlinear Schrödinger Eqn

$$\Phi(x) = \phi_1(x), \qquad s.t.$$
$$\mu\Phi = \left(-\frac{\nabla^2}{2m} + V(\mathbf{x}) + \lambda_{2n}|\Phi|^{2n-2} \right) \Phi$$

$$\mu\Phi = \left(-\frac{\nabla^2}{2m} + V(\mathbf{x}) + \lambda_{2n}|\Phi|^{2n-2}\right)\Phi$$

Self-gravitating BEC system:
- When BEC system is large enough that gravity is no longer negligible,
- $\lambda_4 = 0, +1, -1$
- $V(\mathbf{x}) = \frac{G_N}{2}\int\frac{|\Phi(\mathbf{x}')|^2}{|\mathbf{x}-\mathbf{x}'|}$

Soliton solutions:
- $V = 0$
- $\lambda_4 = \pm 1$
- Soliton solution.
  - $\lambda_4 > 0$, repulsive $\Phi \sim \tanh(x)$
  - $\lambda_4 < 0$, attractive $\Phi \sim \cosh(x)^{-1}$

Superconductor:
- Switching on magnetic field, $\nabla \rightarrow \nabla - ieA$,
- $n = 4$,
- $V = 0$,
- Ginsburg-Landau theory for type I superconductor.

Hydrogen atom:
- $\lambda = 0$
- $V = \frac{\alpha}{r}$

$$\mu\Phi = \left(-\frac{\nabla^2}{2m} + V(\mathbf{x}) + \lambda_{2n}|\Phi|^{2n-2}\right)\Phi$$

Self-gravitating BEC system:

- When BEC system is large enough that gravity is no longer negligible,
- $\lambda_4 = 0, +1, -1$
- $V(\mathbf{x}) = \frac{G_N}{2}\int\frac{|\Phi(\mathbf{x}')|^2}{|\mathbf{x}-\mathbf{x}'|}$

Soliton stars, oscillons (real), BEC dark matter, etc.

Superconductor:

- Switching on magnetic field, $\nabla \to \nabla - ieA$,
- $n = 4$,
- $V = 0$,
- Ginsburg-Landau theory for type I superconductor.

Hydrogen atom:

- $\lambda = 0$
- $V = \frac{\alpha}{r}$

$$\mu\Phi = \left(-\frac{\nabla^2}{2m} + V(\mathbf{x}) + \lambda_{2n}|\Phi|^{2n-2}\right)\Phi$$

Self-gravitating BEC system:

- When BEC system is large enough that gravity is no longer negligible,
- $\lambda_4 = 0, +1, -1$
- $V(\mathbf{x}) = \frac{G_N}{2}\int \frac{|\Phi(\mathbf{x'})|^2}{|\mathbf{x}-\mathbf{x'}|}$

Soliton stars, oscillons (real), BEC dark matter, etc.

BEC condition:

$$\frac{1}{mv} = \ell_d > \langle\ell\rangle = 1/n^{1/3}$$

$$T_c \sim mv^2 \sim n^{2/3}/m \sim \rho^{2/3}/m^{5/3}$$

In the context of cosmology,

Hydrogen atom:
- $\lambda = 0$
- $V = \frac{\alpha}{r}$

$$\left.\begin{array}{l}\rho \leq \Omega_{DM}\rho_c \\ v \sim 300 \text{ km/s}\end{array}\right\} \Rightarrow m < 1 \text{ eV}$$

$$\mu\Phi = \left(-\frac{\nabla^2}{2m} + V(\mathbf{x}) + \lambda_{2n}|\Phi|^{2n-2} \right) \Phi$$

Self-gravitating BEC system:

- When BEC system is large enough that gravity is no longer negligible,
- $\lambda_4 = 0, +1, -1$
- $V(\mathbf{x}) = \frac{G_N}{2} \int \frac{|\Phi(\mathbf{x}')|^2}{|\mathbf{x}-\mathbf{x}'|}$

Soliton stars, oscillons (real), BEC dark matter, etc.

BEC condition:

$$\frac{1}{mv} = \ell_d > \langle \ell \rangle = 1/n^{1/3}$$

$$T_c \sim mv^2 \sim n^{2/3}/m \sim \rho^{2/3}/m^{5/3}$$

In the context of cosmology,

$$\left. \begin{array}{l} \rho \le \Omega_{DM}\rho_c \\ v \sim 300 \text{ km/s} \end{array} \right\} \Rightarrow m < 1 \text{ eV}$$

- Nonlinear dispersion Φ^n
- solution does not scale $\Phi \to \Delta\Phi$
- $|\rho_\Phi|$ too small, gravity is negligible
- $|\rho_\Phi|$ too large, instability from
  - gravitational instability (GR)
  - non-linear dispersion (quantum)
  - space-time backreaction (GR + quantum)

$$\mu\Phi = \left(-\frac{\nabla^2}{2m} + V(\mathbf{x}) + \lambda_{2n}|\Phi|^{2n-2}\right)\Phi$$

Self-gravitating BEC system:

- When BEC system is large enough that gravity is no longer negligible,
- $\lambda_4 = 0, +1, -1$
- $V(\mathbf{x}) = \frac{G_N}{2}\int\frac{|\Phi(\mathbf{x}')|^2}{|\mathbf{x}-\mathbf{x}'|}$

Soliton stars, oscillons (real), BEC dark matter, etc.

BEC condition:

$$\frac{1}{mv} = \ell_d > \langle\ell\rangle = 1/n^{1/3}$$

$$T_c \sim mv^2 \sim n^{2/3}/m \sim \rho^{2/3}/m^{5/3}$$

In the context of cosmology,

$$\left.\begin{array}{l}\rho \leq \Omega_{DM}\rho_c \\ v \sim 300 \text{ km/s}\end{array}\right\} \Rightarrow m < 1 \text{ eV}$$

- Nonlinear dispersion Φ^n
- solution does not scale $\Phi \to \Delta\Phi$
- $|\rho_\Phi|$ too small, gravity is negligible
- $|\rho_\Phi|$ too large, instability from
  - gravitational instability (GR)
  - non-linear dispersion (quantum)
  - space-time backreaction (GR + quantum)

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi^*)(\partial^\mu \phi) - \frac{1}{2}m^2|\phi|^2 - \frac{\lambda}{4!}|\phi|^4$$

- $G_N M_* < 1/m$
    - $M_* \gtrsim M_\odot \Rightarrow m \lesssim 10^{-10} \text{ eV}$ (can be relaxed for NR case)

Boson Star Mass Profile – Gravitational Bound

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi^*)(\partial^\mu \phi) - \frac{1}{2}m^2|\phi|^2 - \frac{\lambda}{4!}|\phi|^4$$
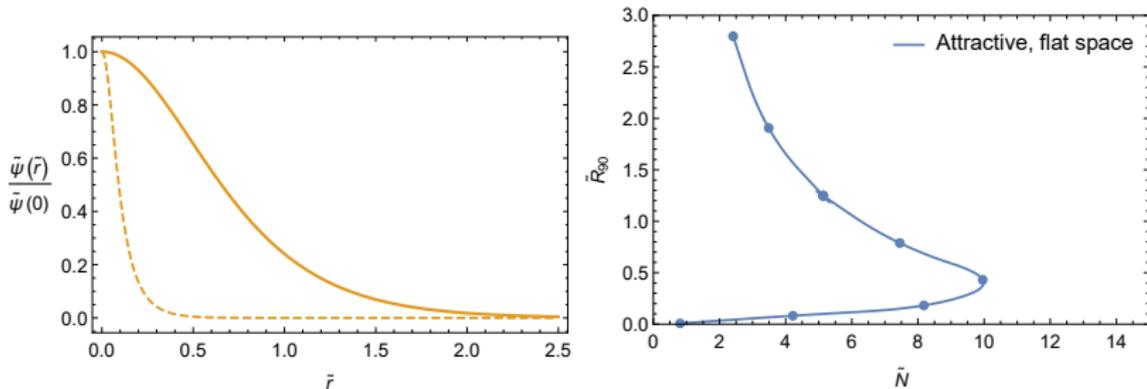
- $G_N M_* < 1/m$
  - $M_* \gtrsim M_\odot \Rightarrow m \lesssim 10^{-10} \text{ eV}$ (can be relaxed for NR case)

Macroscopic scale	Microscopic mechanism	Compactness $\frac{G_N M}{R}$	Mass
White dwarf	Pauli exclusion principle (degeneracy pressure of e^-)	$\sim \mathcal{O}(10^{-4})$	$\sim M_\odot$
Neutron star	Pauli exclusion principle (degeneracy pressure n)	~ 0.2	$\sim M_\odot$
Boson star	Heisenberg's uncertainty principle (kinetic energy of scalars)	?	$\sim \frac{M_{Pl}^2}{m}$

$$i\dot{\phi} = -\frac{1}{2m}\nabla^2\phi - G_N m^2 \phi \int d^3\mathbf{x}' \frac{\phi^*(\mathbf{x}')\phi(\mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|} + \frac{\lambda}{8}|\phi|^2\phi,$$
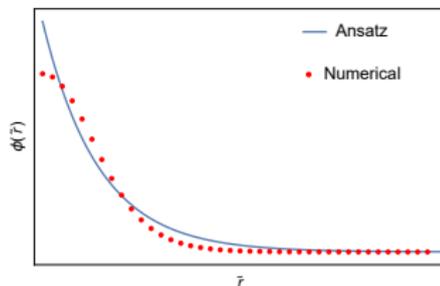
Numerical solution:

$$i\dot{\phi} = -\frac{1}{2m}\nabla^2\phi - G_N m^2 \phi \int d^3\mathbf{x}' \frac{\phi^*(\mathbf{x}')\phi(\mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|} + \frac{\lambda}{8}|\phi|^2\phi,$$

Analytical ansatz: $\phi(r) = \sqrt{\frac{N}{\pi R^3}} \mathrm{e}^{-r/R}$
approximates the wave function well except
for the center. Plug in the
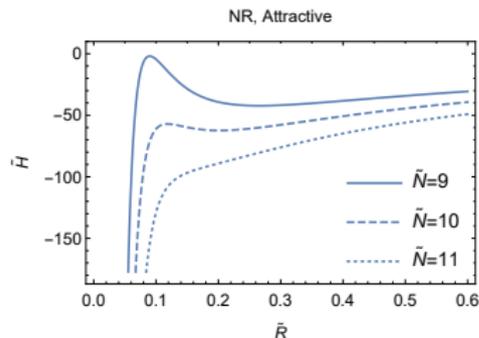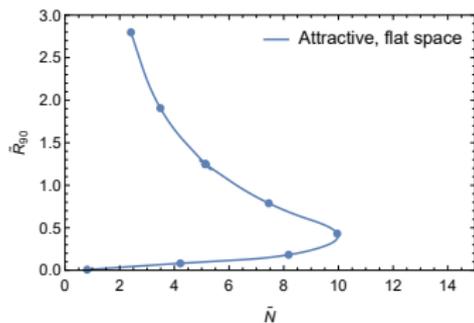Schrödinger-Newton system.

$$H_{NR} = \frac{N}{2mR^2} - \frac{N^2}{128\pi f^2 R^3} - \frac{5G_N m^2 N^2}{16R}$$

Boson Star Mass Profile – Non-linear Dispersion

Analytical ansatz: $\phi(r) = \sqrt{\frac{N}{\pi R^3}}\, e^{-r/R}$ approximates the wave function well except for the center. Plug in the Schrödinger-Newton system.
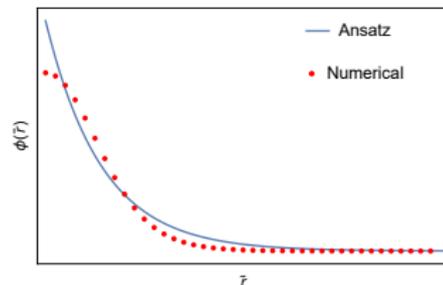
$$H_{NR} = \frac{N}{2mR^2} - \frac{N^2}{128\pi f^2 R^3} - \frac{5 G_N m^2 N^2}{16 R}$$

Self-gravitating BEC system:

- When BEC system is large enough that gravity is no longer negligible,
- $V(\mathbf{x}) = \frac{G_N}{2} \int \frac{|\Phi(\mathbf{x}')|^2}{|\mathbf{x}-\mathbf{x}'|}$
- $\lambda_4 = 0, +1, -1$

Soliton stars, boson stars, oscillons (real), BEC dark matter, etc.

BEC condition:

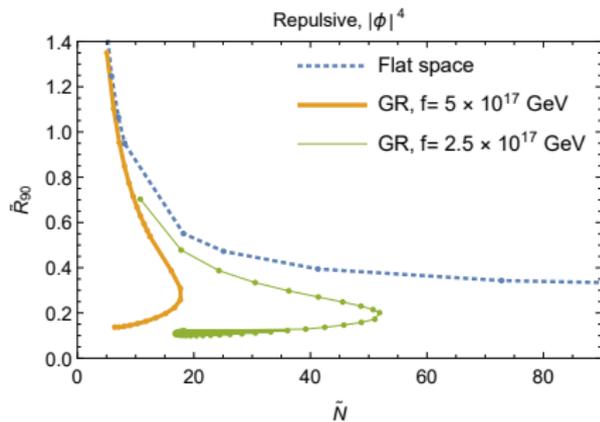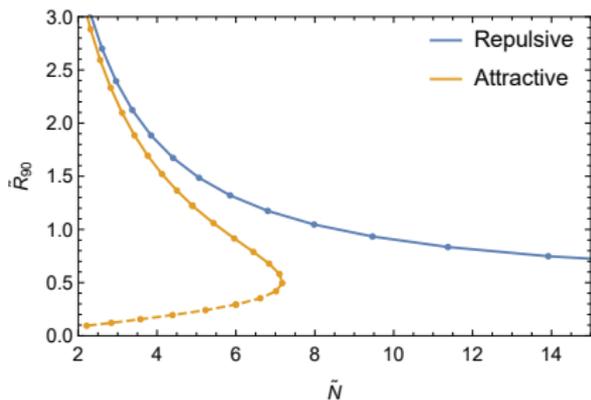$$\frac{1}{mv} = \ell_d > \langle \ell \rangle = 1/n^{1/3}$$

$$T_c \sim mv^2 \sim n^{2/3}/m \sim \rho^{2/3}/m^{5/3}$$

In the context of cosmology,

$$\left. \begin{array}{l} \rho \leq \Omega_{DM}\rho_c \\ v \sim 300 \text{ km/s} \end{array} \right\} \Rightarrow m < 1 \text{ eV}$$

- Nonlinear dispersion Φ^n
- solution does not scale $\Phi \to \Delta\Phi$
- $|\rho_\Phi|$ too small, gravity is negligible
- $|\rho_\Phi|$ too large, instability from
  - gravitational instability
  - non-linear dispersion
  - space-time backreaction

Boson Star Mass Profile – Space-time backreaction

Keep the exponential ansatz for the scalar wave function, but also account for the back reaction of the spacetime:

$$ds^2 = -(1 + 2V(r))dt^2 + (1 - 2V(r))dr^2 + r^2 d\theta^2 + r^2 \sin^2\theta d\phi^2,$$

$$\Phi(r) = \sqrt{\frac{N}{\pi m R^3}} \, e^{-r/R}, \qquad V(r) = -\frac{G_N M(r)}{r}.$$

Integrate the energy density

$$H = \int_0^\infty dr \, 4\pi r^4 \, T_0^0$$

Keep the exponential ansatz for the scalar wave function, but also account for the back reaction of the spacetime:

$$ds^2 = -(1 + 2V(r))dt^2 + (1 - 2V(r))dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2,$$

$$\Phi(r) = \sqrt{\frac{N}{\pi m R^3}} \, e^{-r/R}, \qquad V(r) = -\frac{G_N M(r)}{r}.$$

Integrate the energy density by plugging the ansatz into:

$$H = \int_0^\infty dr \, 4\pi r^4 \left(\frac{\mu^2}{2(1 + 2V(r))} \Phi^2 + \frac{1}{2} m^2 \Phi^2 + \frac{1}{2(1 - 2V(r))} (\partial_r \Phi)^2 + \frac{1}{4!} \left(\frac{m^2}{f^2} \right) \Phi^4 \right)$$

Keep the exponential ansatz for the scalar wave function, but also account for the back reaction of the spacetime:

$$ds^2 = -(1 + 2V(r))dt^2 + (1 - 2V(r))dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2,$$

$$\Phi(r) = \sqrt{\frac{N}{\pi m R^3}} \, e^{-r/R}, \qquad V(r) = -\frac{G_N M(r)}{r}.$$

We recover all the NR limit, with an extra term:

$$H = \underbrace{mN}_{H_{mass}} + \underbrace{\frac{N}{2mR^2}}_{H_{kin}} - \underbrace{\frac{5G_N m^2 N^2}{16R}}_{H_{grav}} + \underbrace{\left(\frac{1}{128\pi f^2}\right)\frac{N^2}{R^3}}_{H_{int}} - \underbrace{\left(\frac{5G_N}{16}\right)\frac{N^2}{R^3}}_{H_{curv}}$$

In the limit of small self-interaction, we recover $C_{max} \sim 0.18$.

Keep the exponential ansatz for the scalar wave function, but also account for the back reaction of the spacetime:

$$ds^2 = -(1 + 2V(r))dt^2 + (1 - 2V(r))dr^2 + r^2 d\theta^2 + r^2 \sin^2\theta d\phi^2,$$

$$\Phi(r) = \sqrt{\frac{N}{\pi m R^3}}\, e^{-r/R}, \qquad V(r) = -\frac{G_N M(r)}{r}.$$
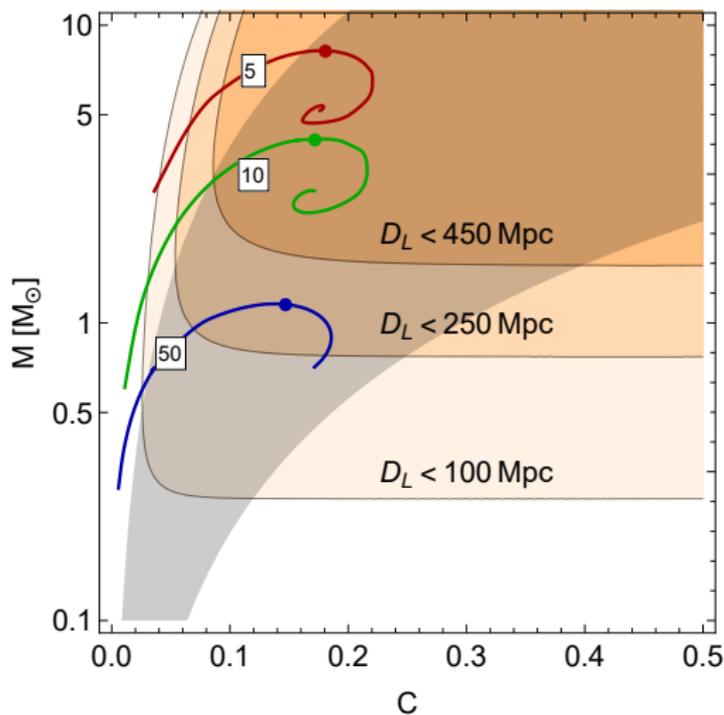
We recover all the NR limit, with an extra term:

$$H = \underbrace{mN}_{H_{mass}} + \underbrace{\frac{N}{2mR^2}}_{H_{kin}} - \underbrace{\frac{5G_N m^2 N^2}{16R}}_{H_{grav}} + \underbrace{\left(\frac{1}{128\pi f^2}\right)\frac{N^2}{R^3}}_{H_{int}} - \underbrace{\left(\frac{5G_N}{16}\right)\frac{N^2}{R^3}}_{H_{curv}}$$

$$T_0^0 = \frac{\mu^2}{2g_{00}}\Phi^2 + \frac{1}{2}m^2\Phi^2 + \frac{1}{2g_{rr}}(\partial_r\Phi)^2 + \frac{\lambda}{4!}\left(\frac{m^2}{f^2}\right)\Phi^4.$$

In the limit of small self-interaction, we recover $C_{max} \sim 0.18$.

In the context of binary mergers,

From top to bottom, $f = 5 \times 10^{16}$ GeV, $f = 10^{17}$ GeV, $f = 5 \times 10^{17}$ GeV.

Self-gravitating BEC system:

- When BEC system is large enough that gravity is no longer negligible,
- $V(\mathbf{x}) = \frac{G_N}{2} \int \frac{|\Phi(\mathbf{x}')|^2}{|\mathbf{x} - \mathbf{x}'|}$
- $\lambda_4 = 0, +1, -1$

Soliton stars, boson stars, oscillons (real), BEC dark matter, etc.

BEC condition:

$$\frac{1}{mv} = \ell_d > \langle \ell \rangle = 1/n^{1/3}$$

$$T_c \sim mv^2 \sim n^{2/3}/m \sim \rho^{2/3}/m^{5/3}$$

In the context of cosmology,

$$\left. \begin{array}{l} \rho \leq \Omega_{DM}\rho_c \\ v \sim 300 \text{ km/s} \end{array} \right\} \Rightarrow m < 1 \text{ eV}$$

- Nonlinear dispersion Φ^n
- solution does not scale $\Phi \to \Delta\Phi$
- $|\rho_\Phi|$ too small, gravity is negligible
- $|\rho_\Phi|$ too large, instability from
  - gravitational instability
  - non-linear dispersion
  - space-time backreaction

Self-gravitating BEC system:

- When BEC system is large enough that gravity is no longer negligible,
- $V(\mathbf{x}) = \frac{G_N}{2} \int \frac{|\Phi(\mathbf{x}')|^2}{|\mathbf{x}-\mathbf{x}'|}$
- $\lambda_4 = 0, +1, -1$

Soliton stars, boson stars, oscillons (real), BEC dark matter, etc.

BEC condition:

$$\frac{1}{mv} = \ell_d > \langle \ell \rangle = 1/n^{1/3}$$

$$T_c \sim mv^2 \sim n^{2/3}/m \sim \rho^{2/3}/m^{5/3}$$

In the context of cosmology,

$$\left.\begin{array}{c} \rho \leq \Omega_{DM}\rho_c \\ v \sim 300 \text{ km/s} \end{array}\right\} \Rightarrow m < 1 \text{ eV}$$

- Nonlinear dispersion Φ^n
- solution does not scale $\Phi \rightarrow \Delta\Phi$
- $|\rho_\Phi|$ too small, gravity is negligible
- $|\rho_\Phi|$ too large, instability from
  - gravitational instability
  - non-linear dispersion
 - space-time backreaction
  - shape of potential

- Light scalars are very dangerous in EFT...
- ...unless it is a pseudo Nambu-Goldstone boson with an approximate shift symmetry

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi)(\partial^\mu \phi) \underbrace{-\frac{1}{2}m^2\phi^2 - \frac{\lambda_{eff}}{4!}\frac{m^2}{f^2}\phi^4 + \cdots}_{V(\phi)},$$

- It is known that QCD axions are attractive. That leads to small compactness, even if it is stable. $\Lambda(1 - \cos(\phi/f)) = m^2\phi^2 - \lambda\phi^4 + ...$
- Take the pheno model:

$$V(|\phi|) = \frac{1}{2}m^2|\phi|^2 + \frac{1}{2Q^2}m^2f^2\left[1 - \cos\left(\frac{Q|\phi|^2}{f^2}\right)\right],$$

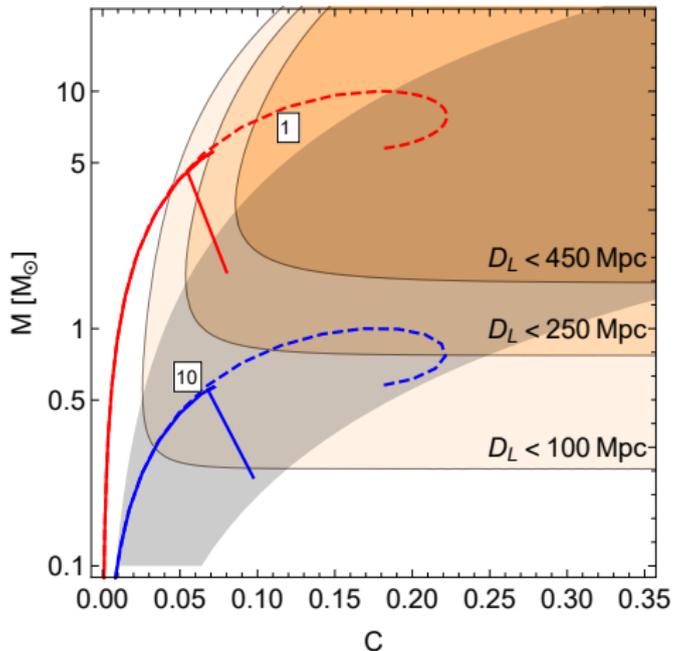The subleading terms can be significant
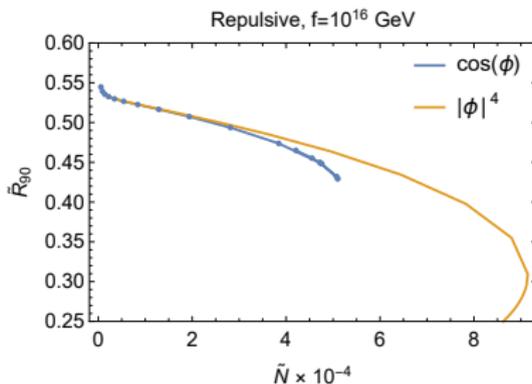
- Subleading ϕ^8 is attractive
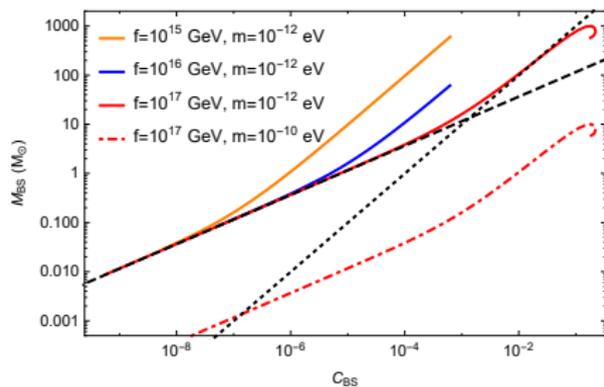- the potential is bounded from both below and above

The subleading terms can be significant:

Boson Star Linear VS Non-linear Regime

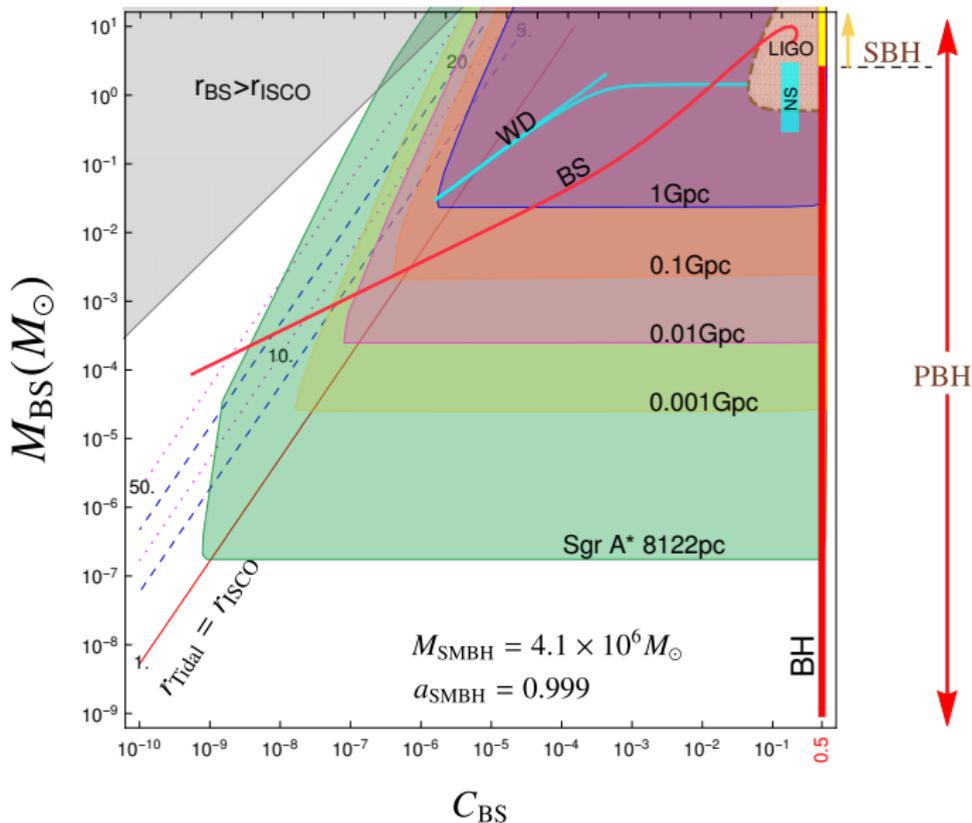$$H(N, R) \approx \frac{N}{2mR^2} + \frac{\lambda N^2}{32\pi f^2 R^3} - \frac{5 G_N m^2 N^2}{16R}$$

- Linear regime: kinetic energy vs gravity $N \sim 1/(m^3 R)$

$$C_{\text{BS}} \sim \frac{mN}{R} \sim m^2 M_{\text{BS}}^2$$

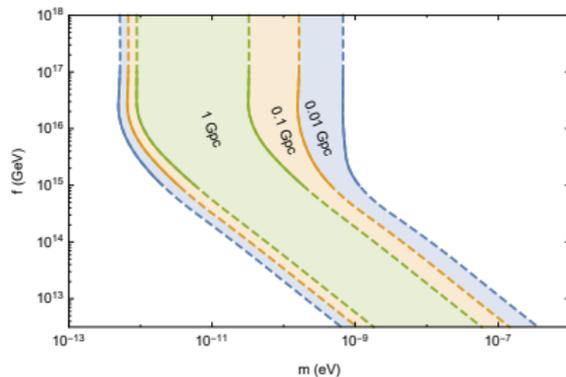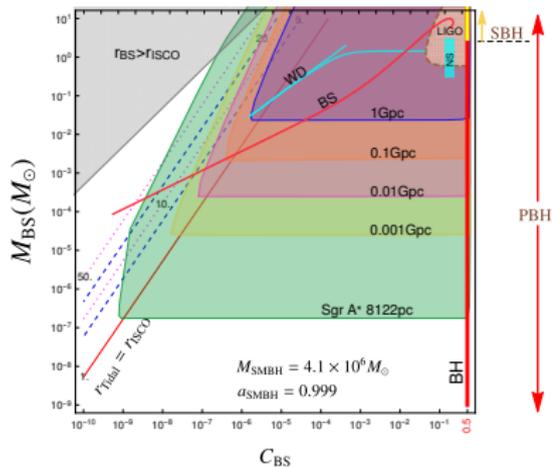- Nonlinear regime: repulsion vs gravity $R \sim 1/(mf)$

$$C_{\text{BS}} \sim \frac{M_{\text{BS}}}{R} \sim m \cdot f \cdot M_{\text{BS}}$$

c.f. Huai-Ke Guo's talk right after this one.

- Can probe the boson star in both linear and nonlinear regime
- Small boson star mass makes it distinguishable from all known sources
- Tidal force disruption further breaks the degeneracy with Primordial BH's.

Scalar Parameter Space

Outlook

- Cosmological scale
  DM self-interaction   ⟺ H_0, power spectrum ⟺   CMB, LSS, 21cm
  BEC of light scalars   ⟺ phase transition ⟺   CMB, BBN

- Galactic scale
  DM self-interaction   ⟺ shape of halo ⟺   lensing, survey $(\mathbf{x}, \mathbf{v})$
  dissipative DM   ⟺ DM substructure ⟺   lensing, survey $(\mathbf{x}, \mathbf{v})$
  BEC light scalars   ⟺ size of halo ⟺   lensing, survey $(\mathbf{x}, \mathbf{v})$

- Stellar scale
  axions   ⟺ gap of BH (m, a_*) ⟺   GW waveform
  ADM bound state   ⟺ dark star ⟺   gravitational wave
  ADM   ⟺ neutron star ⟺   gravitational wave
  BEC of light scalars   ⟺ boson star ⟺   gravitational wave

Outlook

- Cosmological scale
  DM self-interaction $\Leftrightarrow H_0$, power spectrum $\Leftrightarrow$ CMB, LSS, 21cm
  BEC of light scalars $\Leftrightarrow$ phase transition $\Leftrightarrow$ CMB, BBN

- Galactic scale
  DM self-interaction $\Leftrightarrow$ shape of halo $\Leftrightarrow$ lensing, survey $(\mathbf{x}, \mathbf{v})$
  dissipative DM $\Leftrightarrow$ DM substructure $\Leftrightarrow$ lensing, survey $(\mathbf{x}, \mathbf{v})$
  BEC light scalars $\Leftrightarrow$ size of halo $\Leftrightarrow$ lensing, survey $(\mathbf{x}, \mathbf{v})$

- Stellar scale
  axions $\Leftrightarrow$ gap of BH $(m, a_*) \Leftrightarrow$ GW waveform
  ADM bound state $\Leftrightarrow$ dark star $\Leftrightarrow$ gravitational wave
  ADM $\Leftrightarrow$ neutron star $\Leftrightarrow$ gravitational wave
  BEC of light scalars $\Leftrightarrow$ boson star $\Leftrightarrow$ gravitational wave

- Lifetime of the compact real scalar stars
- Excited states
- Extended potential structure, e.g. double well, nonminimal coupling, etc.
- Momentum corrections
- Merger process

Thank you!

(not to scale)

- When tidal force $\sim$ self-gravity, the star is severely deformed.

$$\ell = \left(\frac{C_{BH}}{C_*} \right) \left(\frac{M_*}{M_{BH}} \right)^{2/3} r_{BH}$$

- Want this condition met not too early before the boson star falls into the BH, *i.e.* $\ell \sim r_{BH}$.
- For $M_* \sim M_\odot$, this is true for dilute boson stars $C \sim 10^{-4}$.