Updates on the Cosmological Collider Physics

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References:

Quasi-Single Field Inflation and the Cosmological Collider:

Chen, YW 0909.0496, 0911.3380; Arkani-Hamed, Maldacena 1503.08043 Standard Model on the Cosmological Collider:

Chen, YW, Xianyu 1604.07841, 1610.06597, 1612.08122; Kumar, Sundrum 1711.03988 BSM on the Cosmological Collider:

Chen, YW, Xianyu 1805.02656; Kumar, Sundrum 1811.11200

High-Spin: Arkani-Hamed, Maldacena 1503.08043, Lee, Baumann, Pimentel 1607.03735

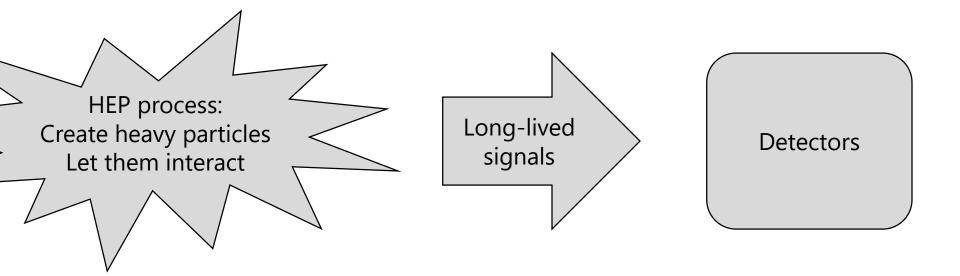
Parity and CP: Liu, Tong, YW, Xianyu, to appear

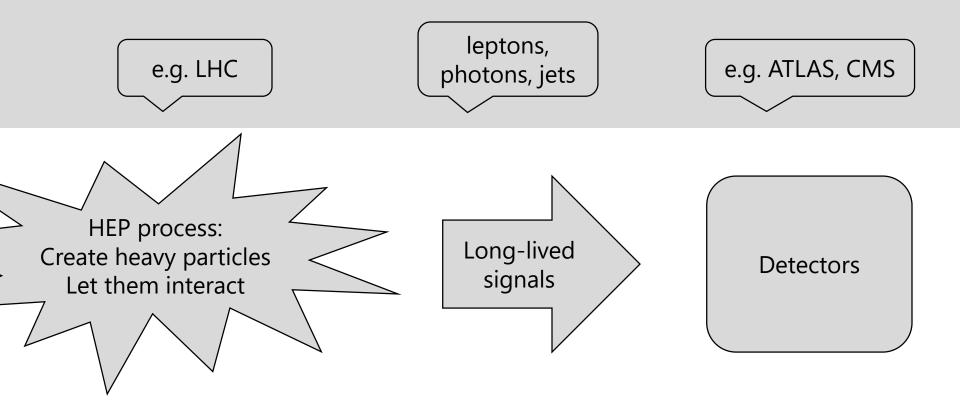
Isocurvature: Lu, YW, Xianyu, to appear

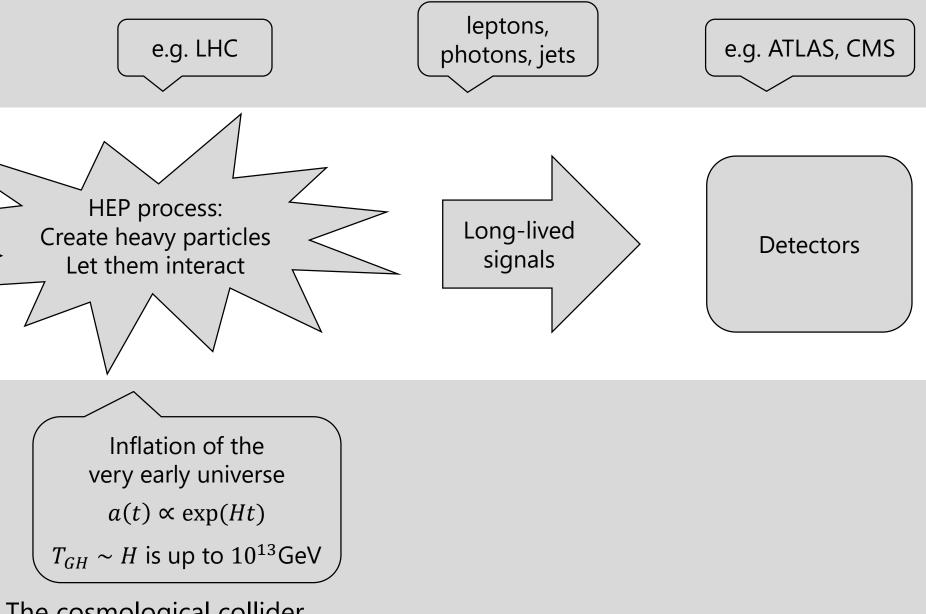
Quantum Primordial Standard Clocks: Chen, Namjoo, YW 1509.03930

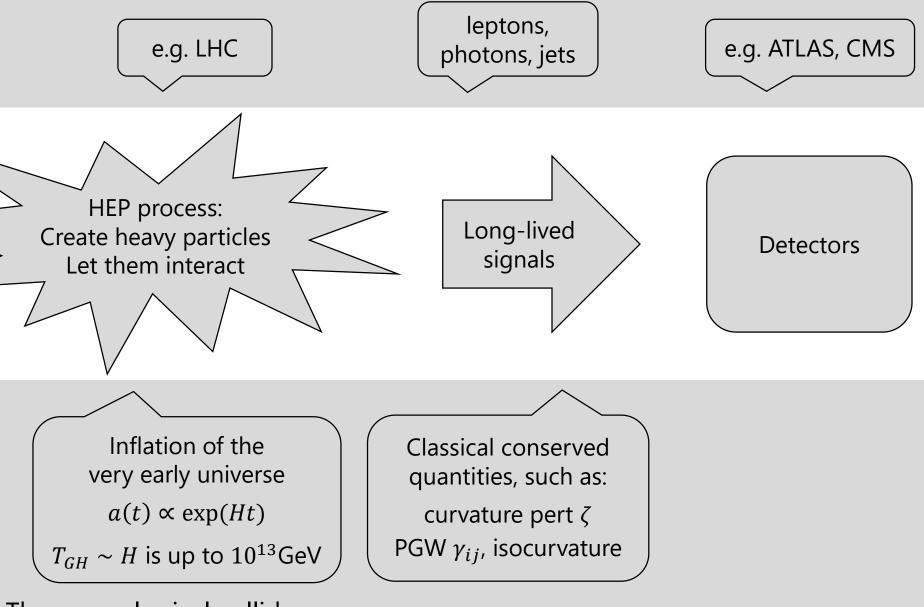
Why is inflation a cosmological "collider"?

What's needed as a "collider"?









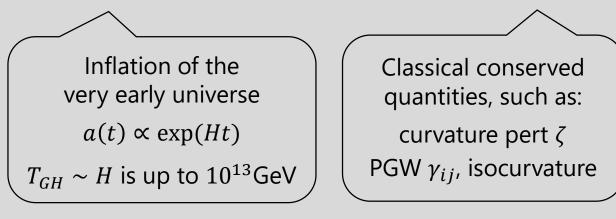
The curvature perturbation $\zeta(\mathbf{x}) \sim \delta N(\mathbf{x}) \sim \frac{H}{\dot{\phi}} \delta \phi \quad (\phi = \phi_0(t) + \delta \phi(\mathbf{x}, t))$

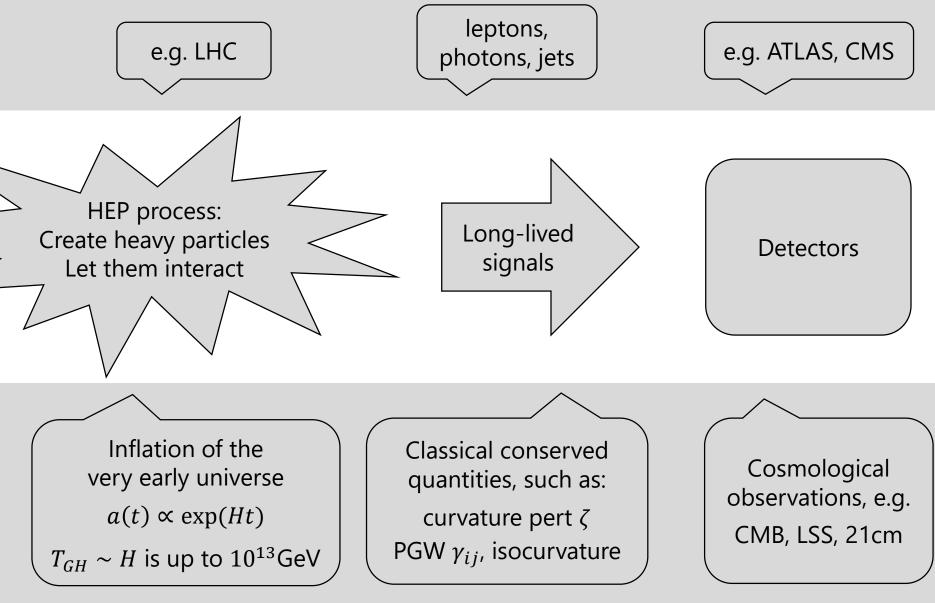
Intuitive (probably too rough) $T_{GH} \sim H \rightarrow \delta \phi \sim H$

Formalism: QFT in curved spacetime

$$\begin{split} S &= \int d^3x \, dt \, a^3(t) \left(\frac{\dot{\phi}^2}{2} + \cdots \right), \\ \langle \delta \phi^n(\mathbf{x}, t) \rangle &= \left\langle \left(\bar{T} e^{i \int^t dt \, H_I} \right) \delta \phi^n_{(I)} \left(T e^{-i \int^t dt \, H_I} \right) \right\rangle, \qquad \langle \delta \phi^2 \rangle \sim H^2 \,, \quad \langle \delta \phi^3 \rangle \cdots \end{split}$$

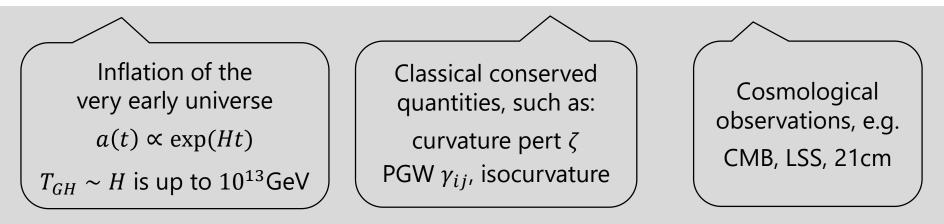
PGW & remaining isocurvature fluctuation (if any): similarly

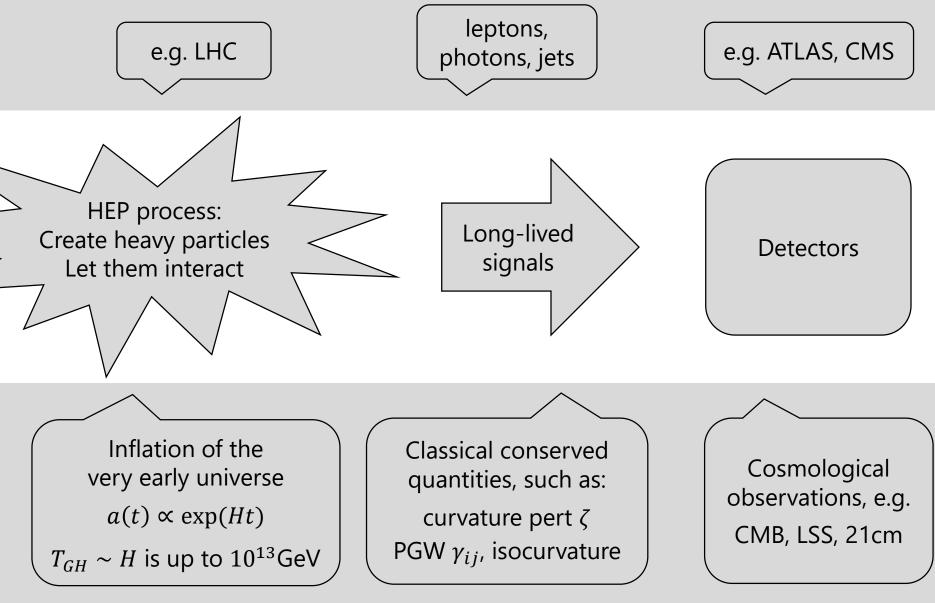




Observations: Correlation functions of

- Curvature perturbation ζ
 - From CMB $\Delta T/T$, LSS & 21cm $\delta \rho / \rho$
 - Status: 2pt well measured (COBE DMR)
 - 3pt, ... (non-Gaussianity) not yet observed. SphereX: 10X
- PGW: From CMB B-mode, not yet observed
- Isocurvature: From details of CMB/LSS, not yet observed

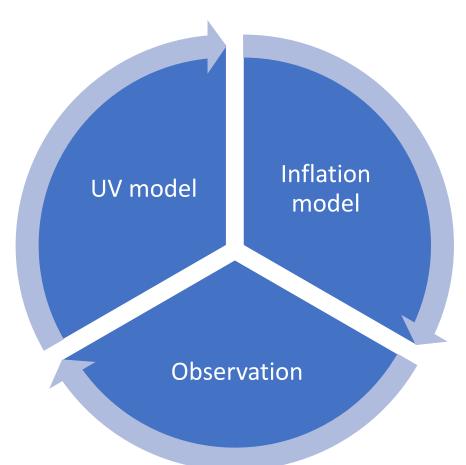




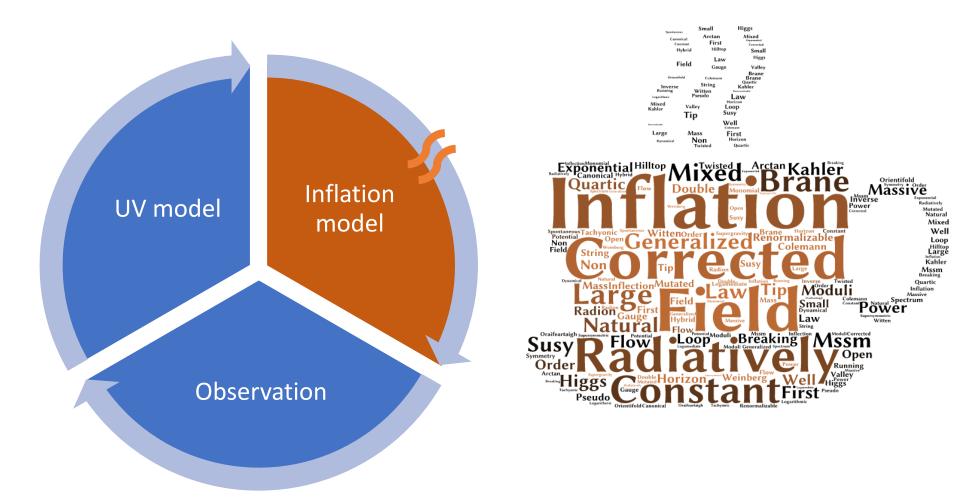


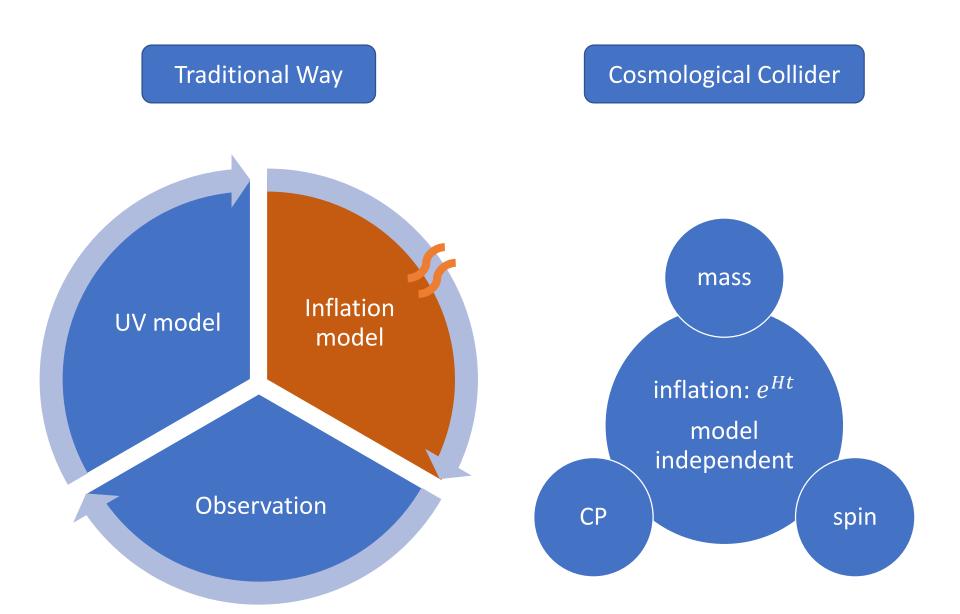
Ever since inflation was proposed, people use inflation to study HEP. What's new about the "cosmological collider"?





Traditional Way





Mass: from resonance

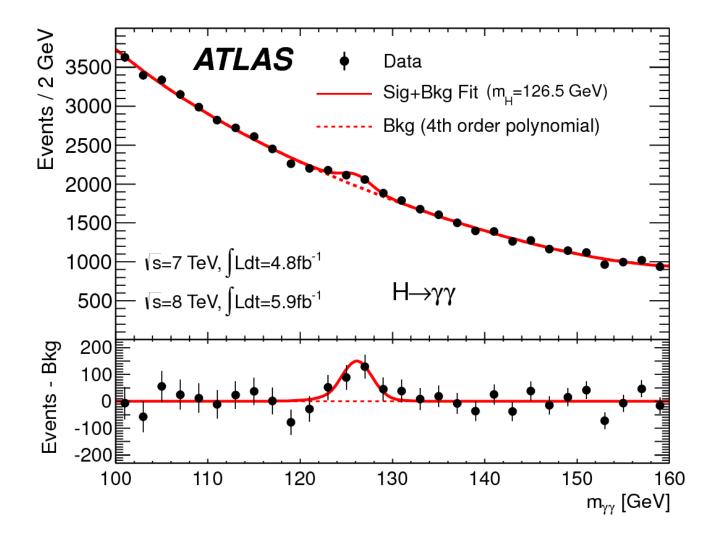


Image: ATLAS

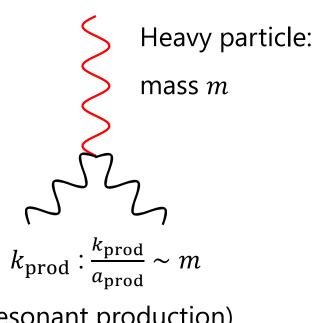
Dispersion relation for light and heavy particles during inflation

Light: $m \ll H$: $\omega = k/a$ (time dependent)

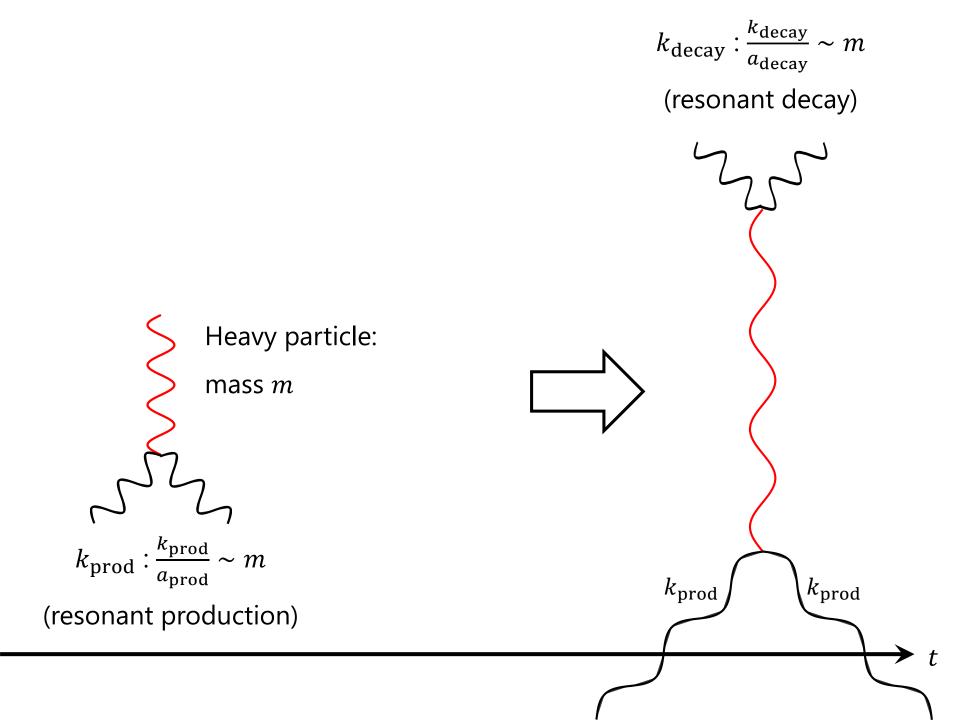
Heavy:
$$m \sim H$$
 or larger: $\omega = \sqrt{\left(\frac{k}{a}\right)^2 + m^2} \sim m$ (time independent)

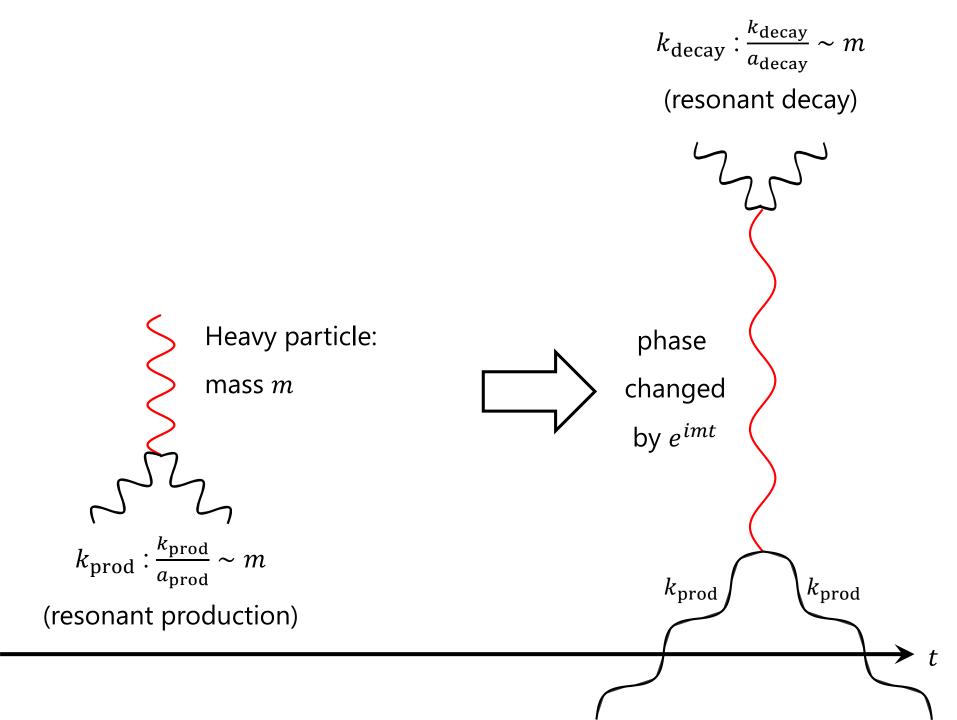
Thus can have a "resonant time" if these two coincide

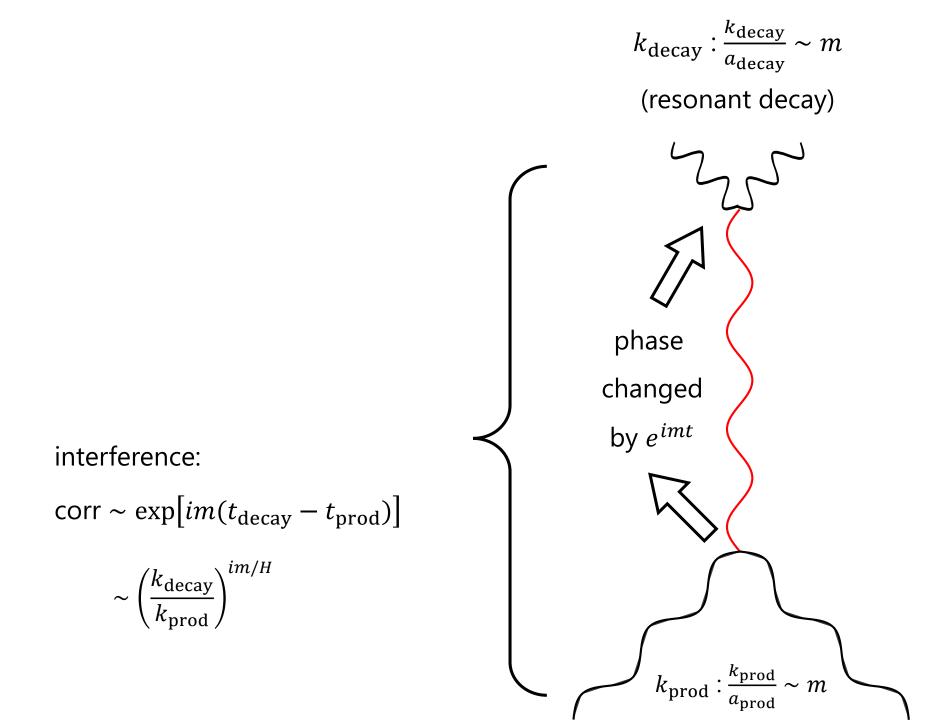
$$\int d\tau f(\tau) e^{-ik\tau} e^{imt}$$

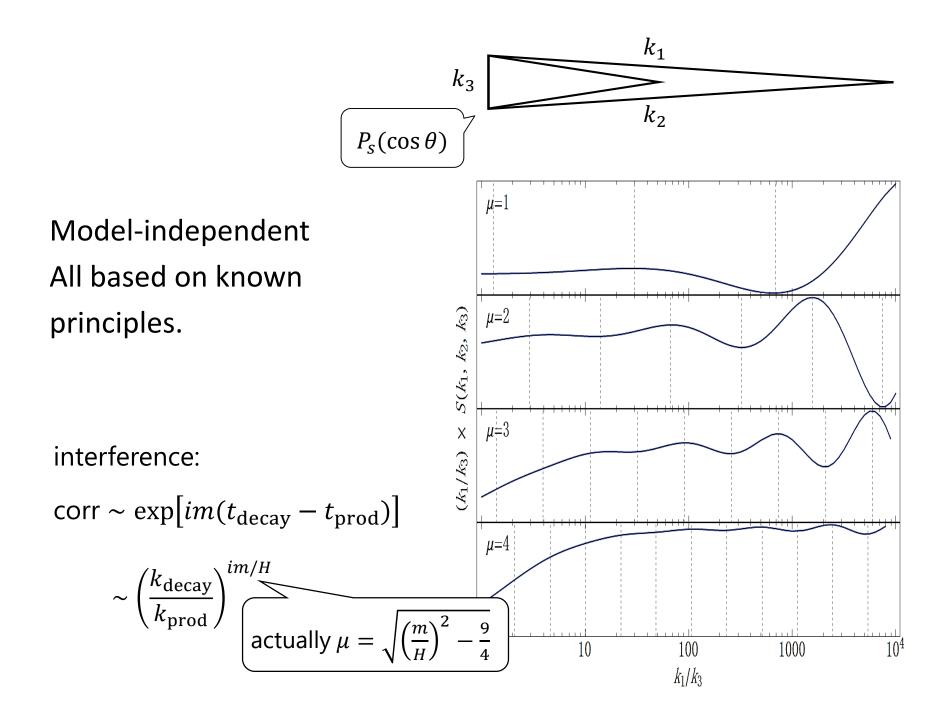


(resonant production)









Spin

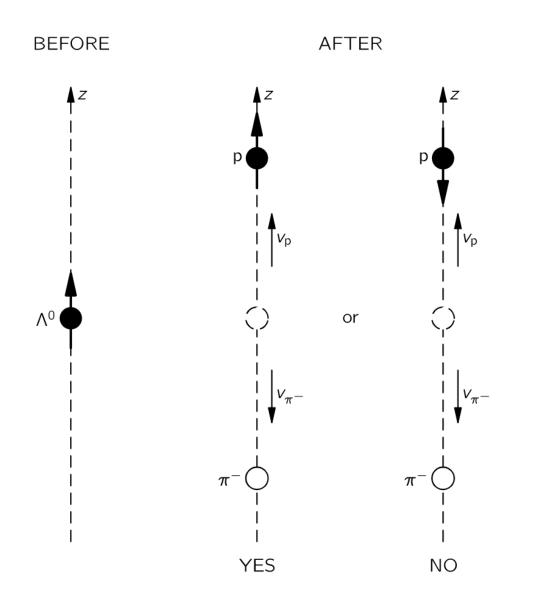
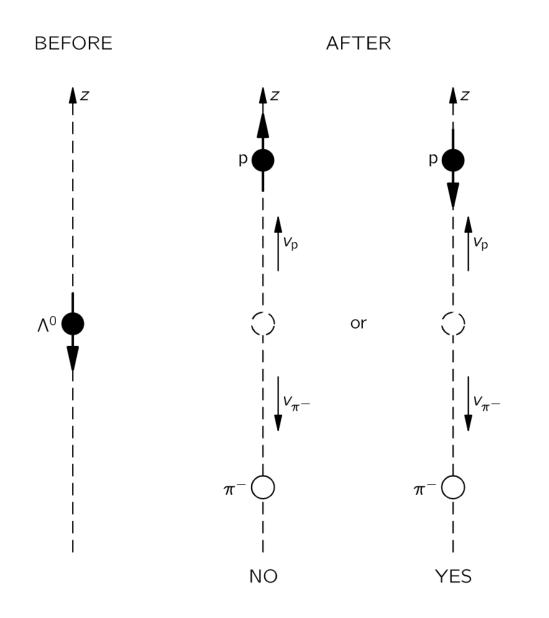
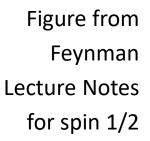
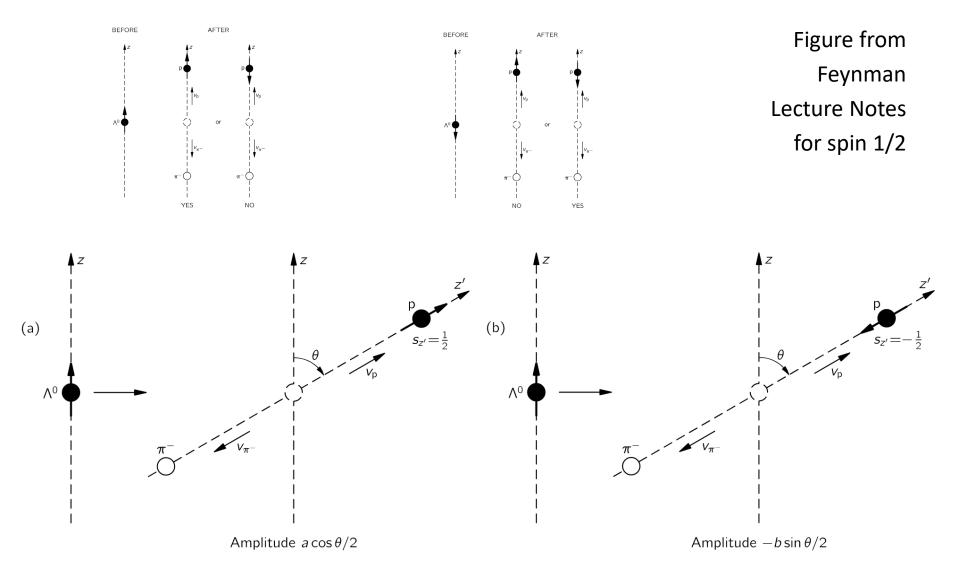


Figure from Feynman Lecture Notes for spin 1/2 Spin



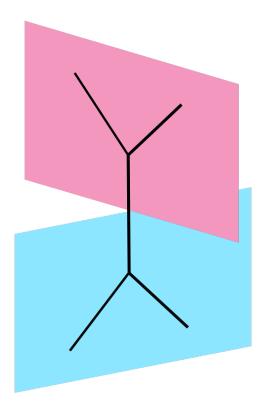


Spin: angular distribution



For integer spin: $P_s(\cos \theta)$ angular distribution

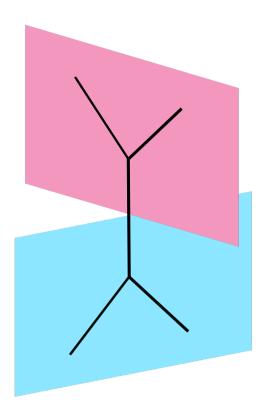
CP: decay plane correlation



CP arises from the plane correlation of the red and the blue

Liu, Tong, YW, Xianyu, to appear

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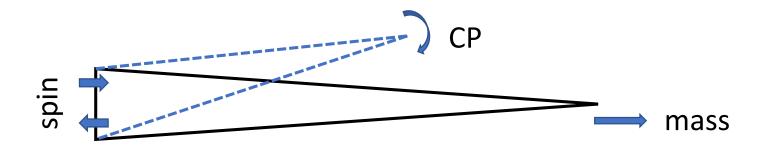
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Recap so far

Cosmological collider:

model-independently read off particle mass (resonance), spin (2D angle), CP (3D angle), ... from inflation



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model-independently read off particle mass (resonance), spin (2D angle), CP (3D angle), ... from inflation

Any target physics on the cosmological collider?

Accidentally near *H* ?

- Grand unification

Kumar, Sundrum 1811.11200

- Neutrino seesaw

Chen, YW & Xianyu, 1805.02656

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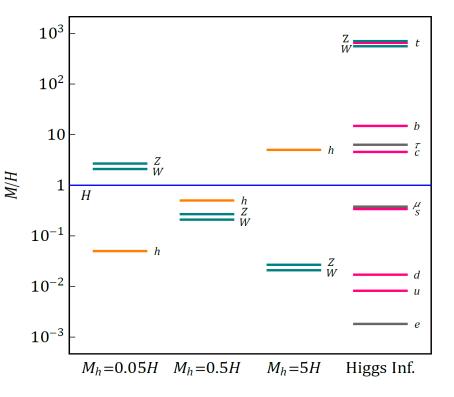
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Uplifted to *H* scale: - Standard Model $\langle h^2 \rangle \sim H^2$ $\lambda h^4 \supset \lambda \langle h^2 \rangle h^2 \sim m_{eff}^2 h^2$ also: possible $h^2 R \sim H^2 h^2$ Chen & YW, 0911.3380 Chen, YW & Xianyu, 1610.06597 Kumar & Sundrum, 1711.03988



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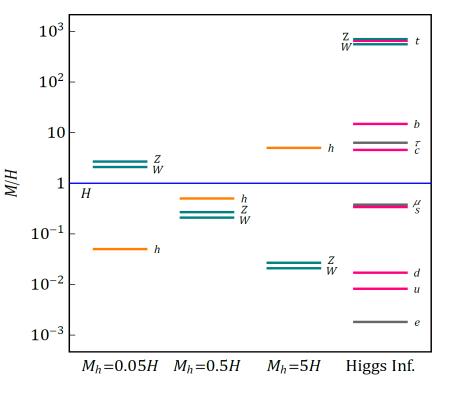
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Chen & YW, 0911.3380

Chen, YW & Xianyu, 1610.06597

Kumar & Sundrum, 1711.03988

- SUSY breaking

Baumann & Green, 1109.0292 Delacretaz, Gorbenko & Senatore 1610.04227 Beyond, beyond and beyond

Beyond $m \sim H$? Usually $e^{-\pi m/H}$ suppressed. However:

- Bumpier inflaton potential

Flauger, Mirbabayi, Senatore, Silverstein 1606.00513

- Non-minimal temperature Tong, YW, Zhou 1801.05688
- Chemical potential Chen, YW & Xianyu, 1805.02656

Beyond curvature perturbation:

- PGW Collider
- Cosmological Isocurvature Colliders Lu, YW, Xianyu, to appear
 e.g. Cosmological Higgs Collider with modulated reheating

Beyond cosmological collider:

- Quantum primordial standard clocks Chen, Namjoo, YW 1509.03930
- Quantum nature of fluctuation Maldacena 1508.01082; Liu, Sou, YW 1608.07909
- Hubble tension Adhikari, Huterer 1905.02278
- Impacts on inflation models Jiang, YW 1703.04477

Consider dim-5 operator $(\partial \phi)^2 \sigma / \Lambda$, $\mathcal{L}_2 \supset \frac{2\phi}{\Lambda} \delta \dot{\phi} \delta \sigma$

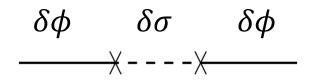
Note: $\dot{\phi} \sim 3600 H^2$. Even $\Lambda = M_p$, if $m \sim H$ for σ , still $\mathcal{L}_2 \sim \sqrt{\epsilon} H \times \delta \dot{\phi} \delta \sigma$

For $r \ge 10^{-3}$ (i.e. $\epsilon \ge 10^{-4}$)

Potentially observable change of

 $\Delta r/r$ and $\Delta (n_s - 1)/(n_s - 1)$

Observable M_p effect!



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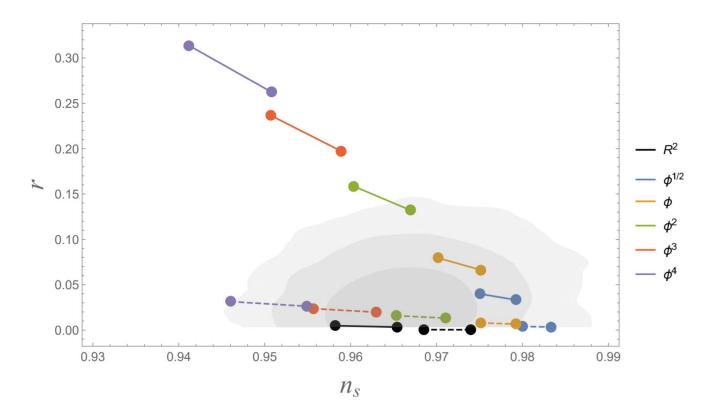
Observable M_p effect!

And many possible enhancement factors:

- 1. Larger *r*
- 2. Multi-field (all positive ΔP_{ζ})
- 3. IR growth if $m \preceq H$
- 4. If $M_{\text{string}} < M_p$
- 5. If $M_{\text{extra }D} < M_p$

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Tong, YW, Zhou 1708.01709, see also An, McAneny, Ridgway, Wise 1706.09971

Summary:

If we knew cosmological correlations infinitely precisely, we know mass and spin of all heavy fields during inflation.

Standard model and beyond, and beyond and beyond

Acknowledgment

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