

Is Vacuum Decay During Inflation Fatal?

Stephen Stopyra

Beyond General Relativity, Beyond Cosmological Standard Model

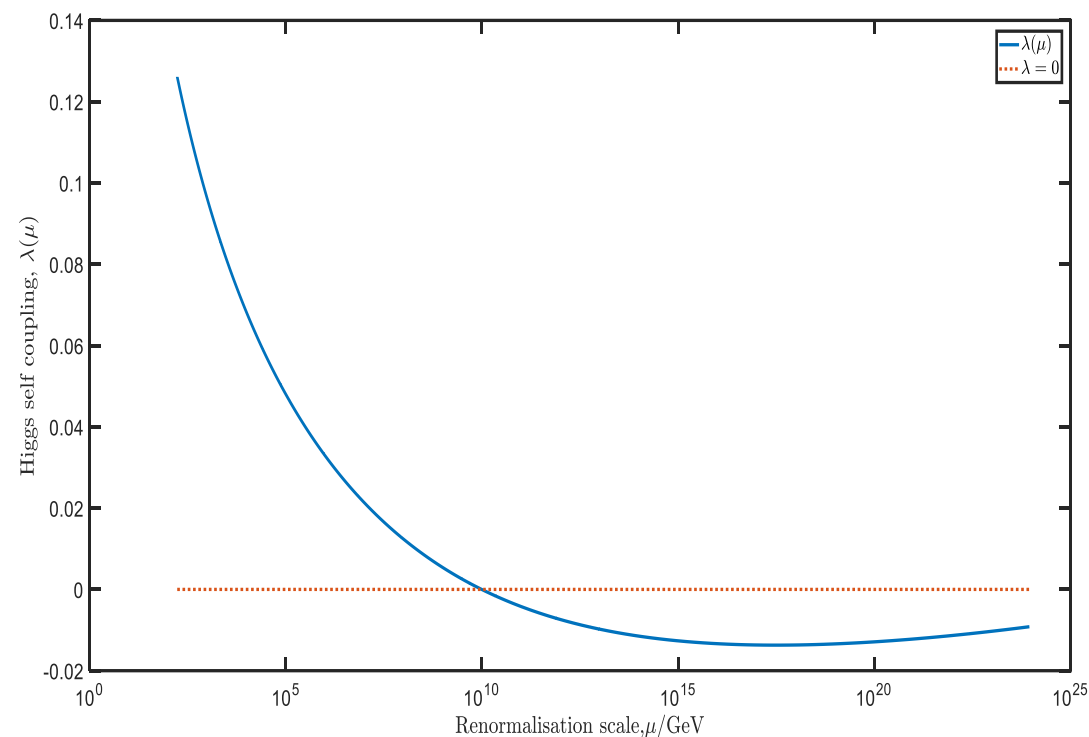
1st July 2019

Based on material in <https://spiral.imperial.ac.uk/handle/10044/1/66255>

See also: <https://doi.org/10.3389/fspas.2018.00040> (arXiv:1809.06923)

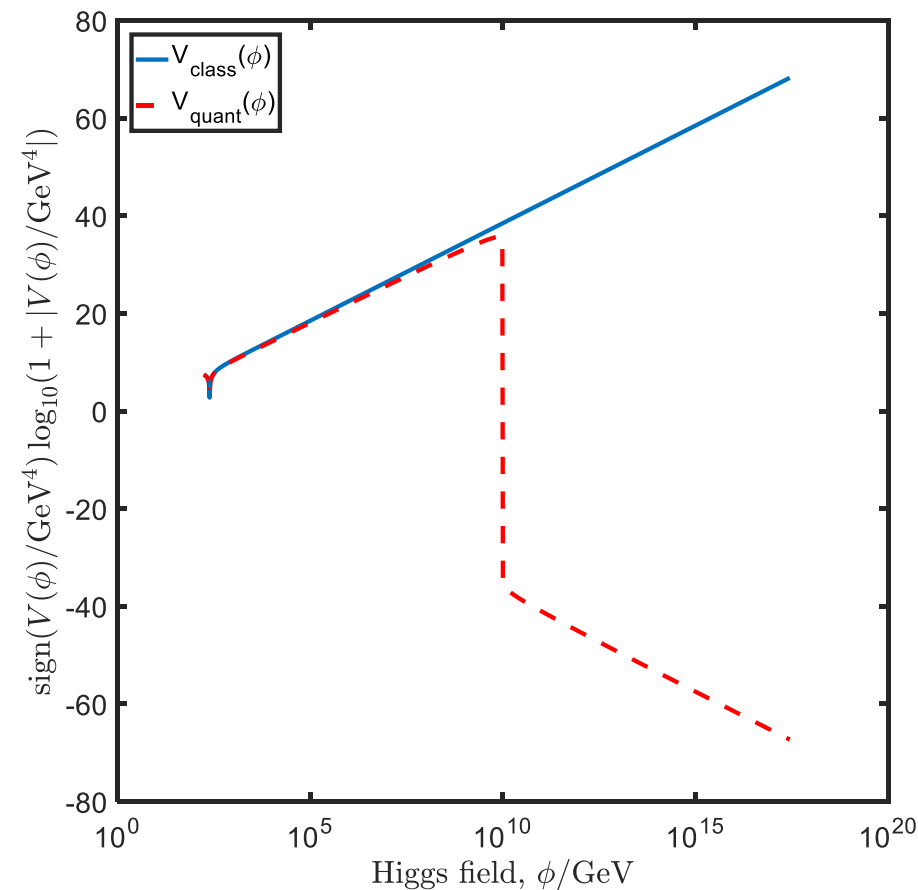
Introduction and Overview

- Current best estimates - $m_H = 125.10 \pm 0.14$ GeV, $m_T = 173.1 \pm 0.9$ GeV[1]
- Vacuum is metastable with these values, but long-lived.
- Flat space – bubbles expand.
- Negative energy density true vacuum – collapse to singularity.
- But what about curved space?
- Cosmological implications - see talk by Arttu Rajantie



Introduction and Overview

- Current best estimates - $m_H = 125.10 \pm 0.14$ GeV, $m_T = 173.1 \pm 0.9$ GeV[1]
- Vacuum is metastable with these values, but long-lived.
- Flat space – bubbles expand.
- Negative energy density true vacuum – collapse to singularity.
- But what about curved space?
- Cosmological implications - see talk by Arttu Rajantie



Introduction and Overview

- Situation less clear during inflation. Bubbles expand, but so does space.
- Some say anti-de-Sitter interior means bubble is inflated away.
- As we will see – not quite that simple.

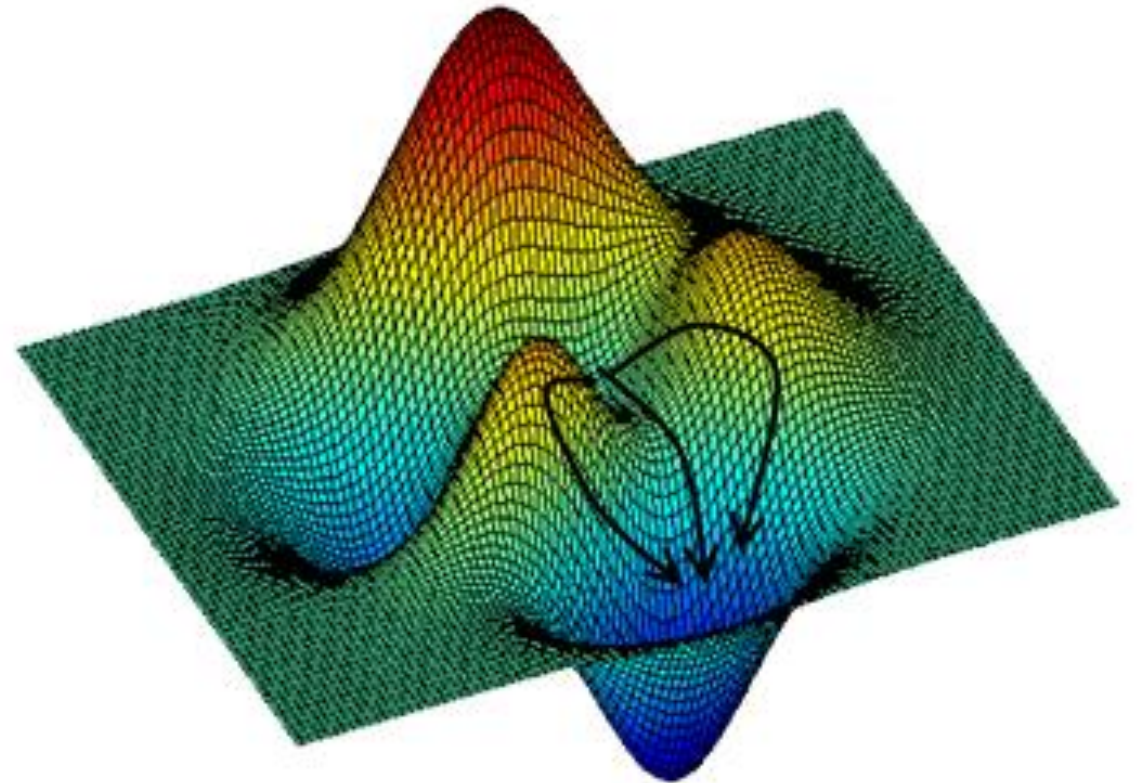


Introduction and Overview

- Review bubble expansion in flat space.
- Include gravity – formation of singularity.
- De Sitter Bubbles
- Conclusions

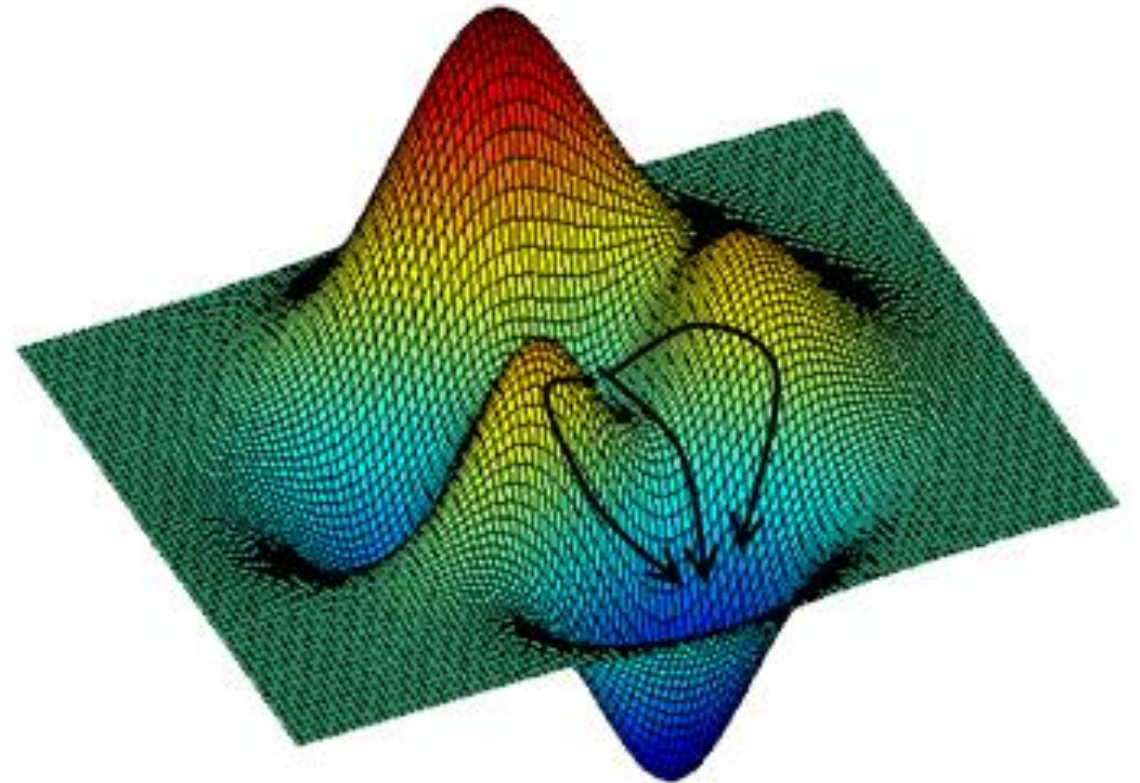
Review bubble expansion in flat space

- Compute shape and decay rate
 - quantum tunnelling with ∞ degrees of freedom.
- 1 d.o.f. \rightarrow WKB approx., $\Gamma \propto \exp\left(-2 \int_{x_1}^{x_2} dx \sqrt{2(V(x) - E)}\right)$
- ∞ d.o.f. – sum over many paths. Minimise *Euclidean* action:
 - $S_E[\phi] = \int d^4x \left[\frac{1}{2} (\nabla\phi)^2 + V(\phi) \right]$



Review bubble expansion in flat space

- Quantum Tunnelling in QFT \rightarrow PDE problem
- $O(4)$ symmetry[2]: PDE \rightarrow ODE
- Solve $\ddot{\phi}_B + \frac{3}{\lambda} \dot{\phi}_B - V'(\phi_B) = 0$,
with $\Gamma \propto e^{-S_E[\phi]}$
- 1 d.o.f. – emerge at single point on other side of barrier.
- ∞ d.o.f. – emerge at a field configuration, $\phi_B(r)$ given by solution by *analytic continuation*.



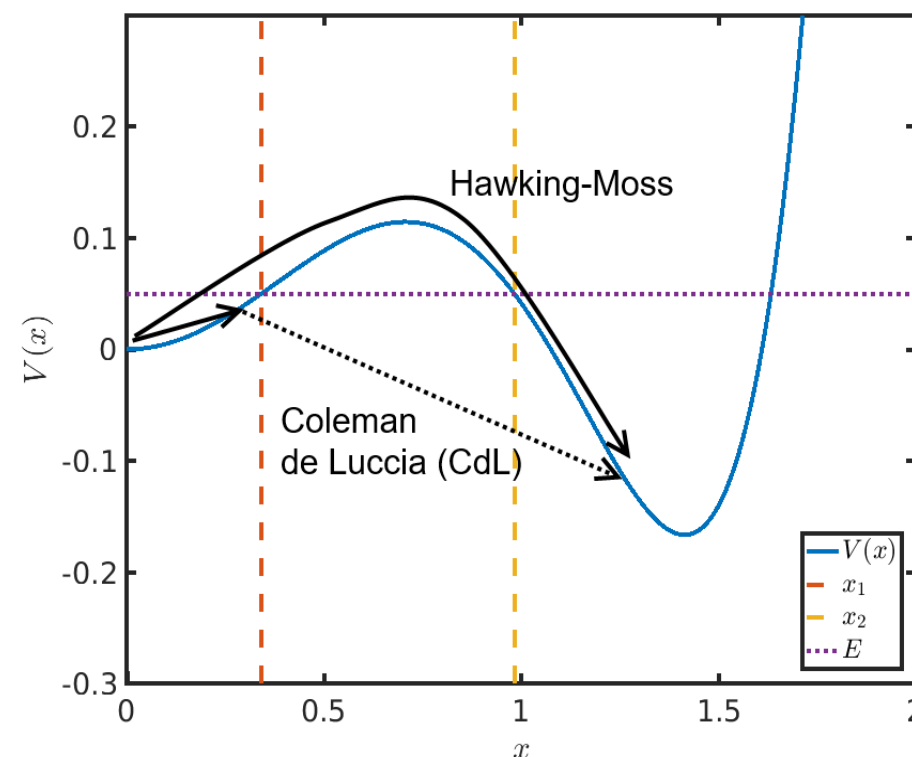
[2] S. Coleman, V. Glaser and A. Martin, Action minima among solutions to a class of euclidean scalar field equations, Communications in Mathematical Physics 58 (1978) 211.

Review bubble expansion in flat space

- Bounce $O(4)$ symmetric $\rightarrow O(3,1)$ symmetry in real space. Unique solution is $\phi(r, t) = \phi_B(\sqrt{r^2 - c^2 t^2})$.
- Position bubble with $\phi = \phi_0$ moves as $r^2 - c^2 t^2 = R_0^2$, $R_0 = \phi_B^{-1}(\phi_0)$.
- $v = \dot{r}(t) = \frac{c^2 t}{\sqrt{R_0^2 + c^2 t^2}} \rightarrow c$
- So bubble rapidly approaches speed of light.
- But what if $r < ct$? Corresponds to patch of spacetime *inside* the bubble. Need gravity to understand it.

Gravitational Collapse of Bubble

- Need to include gravity! See [3]
- $S_E = \int d^4x \sqrt{|g|} \left[\frac{1}{2} \nabla_\mu \phi \nabla^\mu \phi + V - \frac{M_P^2}{2} R \right],$
- $\Gamma \propto e^{-S_E}$
- $\ddot{\phi} + \frac{3\dot{a}}{a} \dot{\phi} - V'(\phi) = 0$
- $\dot{a}^2 = 1 - \frac{a^2}{3M_P^2} \left(-\frac{\dot{\phi}^2}{2} + V(\phi) \right)$
- $ds^2 = d\chi^2 + a^2(\chi) [d\psi^2 + \sin^2 \psi d\Omega_2^2]$
- How to continue this back?



[3] S. Coleman and F. De Luccia, Gravitational effects on and of vacuum decay, Phys. Rev. D 21 (1980) 3305.

Gravitational Collapse of Bubble

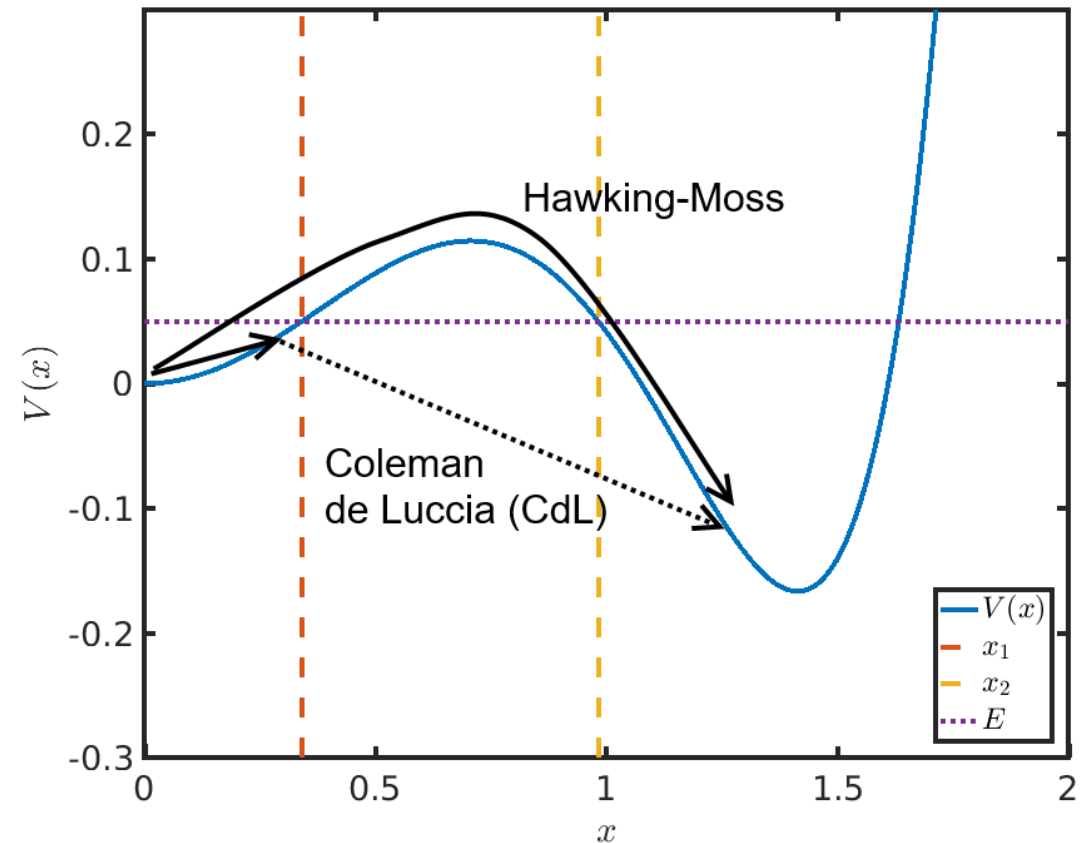
- Good explanation given by [4]. In flat space, simply do $t = -i\tau$.
- In curved space, define $f(\chi)$ by $f' = \frac{f}{a}$, $f(0) = 0$ and transform as:
 - $\left. \begin{array}{l} \tilde{r} = f(\chi) \sin \psi \\ \tilde{t} = f(\chi) \cos \psi \end{array} \right\} \Rightarrow ds^2 = \frac{a^2(\chi)}{f^2(\chi)} (d\tilde{t}^2 + d\tilde{r}^2 + \tilde{r}^2 d\Omega_2^2)$
 - Conformally flat, so now do $\tilde{t} = -i\tilde{t}$, (equivalently $\psi_+ = i(\psi - \pi/2)$)
 - $\left. \begin{array}{l} \tilde{r} = f(\chi) \cosh \psi_+ \\ \tilde{t} = f(\chi) \sinh \psi_+ \end{array} \right\} \Rightarrow ds^2 = \frac{a^2(\chi)}{f^2(\chi)} (-d\tilde{t}^2 + d\tilde{r}^2 + \tilde{r}^2 d\Omega_2^2)$

Gravitational Collapse of Bubble

- $\Rightarrow \tilde{r}^2 - \tilde{t}^2 = f^2(\chi)$ (c.f. flat space, with $f(\chi) = \chi$)
- So solution is $\phi(\tilde{r}, \tilde{t}) = \phi_B \left(f^{-1} \left(\sqrt{\tilde{r}^2 - \tilde{t}^2} \right) \right)$ (c.f. flat space $\phi_B \left(\sqrt{\tilde{r}^2 - \tilde{t}^2} \right)$)
- But $\tilde{r}^2 - \tilde{t}^2 = f^2(\chi)$ clearly only works for $\tilde{r} > \tilde{t}$, ie, outside the lightcone of nucleated bubble.
- Burda et al. [3] showed you can get the $\tilde{r} < \tilde{t}$ portion with:
 - $\left. \begin{array}{l} \tilde{r} = f(\chi) \sinh \psi_- \\ \tilde{t} = f(\chi) \cosh \psi_- \end{array} \right\} \Rightarrow ds^2 = -d\chi^2 + a^2(\chi) [d\psi_-^2 + \sinh^2 \psi_- d\Omega_2^2]$

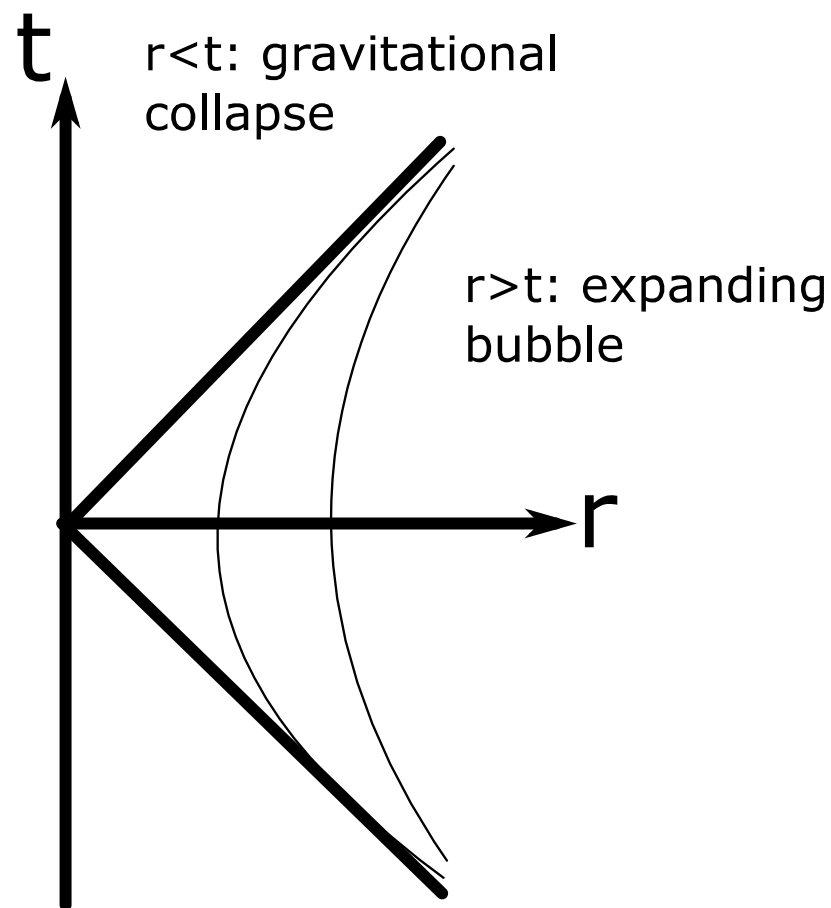
Gravitational Collapse of Bubble

- Same equation of motion as bounce, but *non-inverted* potential.
- $\ddot{a} = -\frac{a}{3M_P^2} \underbrace{(\dot{\phi}^2 - V(\phi))}_{>0 !!!}$
- $a = 0$, *big-crunch* singularity guaranteed in finite time. Same conclusion as [3].



Gravitational Collapse of Bubble

- So what happens next?
- $r > t$ describes the expanding bubble.
- $r < t$ describes the collapse.
- But expansion continues forever regardless of what is going on inside (information can't propagate from out of the $r < t$ region – peak of bubble wall moves at c)



De Sitter Bubbles

- So what happens to bubbles during inflation?
- For pure de Sitter tunnelling, get $a = \frac{1}{H} \sin H\chi$ for *Euclidean* (4-sphere) metric.

• $\Rightarrow f(\chi) = \tan \frac{H\chi}{2}$ and metric:

• $ds^2 = \frac{4}{H^2(1+\tilde{r}^2-\tilde{t}^2)^2} [-d\tilde{t}^2 + d\tilde{r}^2 + \tilde{r}^2 d\Omega_2^2]$

• Via
$$\left. \begin{aligned} t &= \frac{1}{2H} \log \left| \frac{1-\tilde{r}^2+2\tilde{t}+\tilde{t}^2}{1-\tilde{r}^2-2\tilde{t}+\tilde{t}^2} \right| \\ r &= \frac{2\tilde{r}}{H(1-\tilde{r}^2-\tilde{t}^2)} \end{aligned} \right\} \Rightarrow$$

• $ds^2 = -dt^2(1 - H^2r^2) + dr^2(1 - H^2r^2)^{-1} + r^2 d\Omega_2^2$

De Sitter Bubbles

- We map the (\tilde{r}, \tilde{t}) parabolas onto (r, t) to see how a de Sitter bubble expands.

- Point with $\phi = \phi_0, r_0 = f\left(\phi_B^{-1}(\phi_0)\right)$ moves as:

$$r(t) = \frac{2}{H(1+r_0^2)} \sqrt{r_0^2 + \frac{1}{4}(1-r_0^2)^2 \tanh^2 Ht} \rightarrow \frac{1}{H} \text{ as } t \rightarrow \infty$$

- Main observation – bubble expands to *comoving* radius $1/H$ and no further. Fills a full Hubble Volume.

Conclusions

- Anti-de-Sitter region at bubble centre indeed collapses.
- But observers outside the bubble never see this. Bubble continues to expand at the speed of light.
- Fills an entire Hubble volume.
- Once filled, this volume expands at the same rate as the rest of the universe, rather than shrinking away.
- Hubble rates of order $10^8 - 10^{10}$ GeV enough to be fatal. So new physics needed somewhere!

